

# W+3 jet production

— *signal or background* —

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# W + 3 jets

- I. W + 3 jets measured at the Tevatron, but **LO varies by more than a factor 2** for reasonable changes in scales

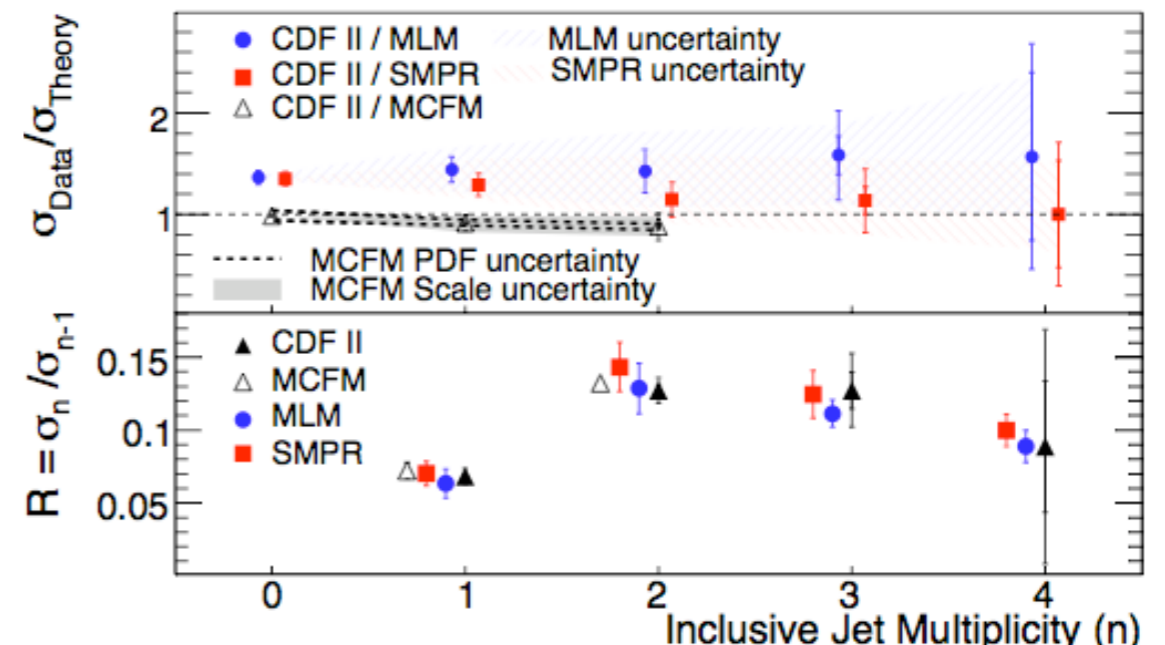
	$W^\pm$ , TeV	$W^+$ , LHC	$W^-$ , LHC
$\sigma$ [pb], $\mu = 40$ GeV	$74.0 \pm 0.2$	$783.1 \pm 2.7$	$481.6 \pm 1.4$
$\sigma$ [pb], $\mu = 80$ GeV	$45.5 \pm 0.1$	$515.1 \pm 1.1$	$316.7 \pm 0.7$
$\sigma$ [pb], $\mu = 160$ GeV	$29.5 \pm 0.1$	$353.5 \pm 0.8$	$217.5 \pm 0.5$

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II. CDF data for W + n jets with n=1,2 is described **exceptionally well by NLO QCD**  
 $\Rightarrow$  verify this for 3 and more jets



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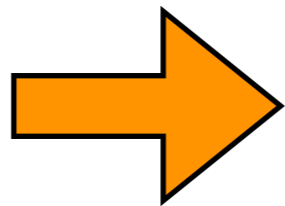
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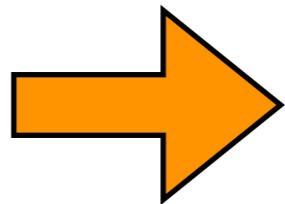
IV. Calculation highly non-trivial optimal testing ground

$$0 \rightarrow \bar{u} d g g g W^+$$



1203 + 104 Feynman diagrams

$$0 \rightarrow \bar{u} d \bar{Q} Q g W^+$$



258 + 18 Feynman diagrams

# Generalized unitarity

I will not explain the method.  
*I will concentrate on applications & recent results*

## References:

- Ellis, Giele, Kunszt '07 [Unitarity in  $D=4$ ]
- Giele, Kunszt, Melnikov '08 [Unitarity in  $D\neq 4$ ]
- Giele & GZ '08 [All one-loop  $N$ -gluon amplitudes]
- Ellis, Giele, Melnikov, Kunszt '08 [Massive fermions,  $ttggg$  amplitudes]
- Ellis, Giele, Melnikov, Kunszt, GZ '08 [ $W+5p$  one-loop amplitudes]
- Ellis, Melnikov, GZ '09, Melnikov & GZ '09 [ $W+3$  jets]

These papers heavily rely on previous work

- Bern, Dixon, Kosower '94 [Unitarity, oneloop from trees]
- Ossola, Pittau, Papadopoulos '06 [OPP]
- Britto, Cachazo, Feng '04 [Generalized cuts]
- [...]

# The F90 Rocket program

## Rocket science!

***Eruca sativa*** =Rocket=roquette=arugula=rucola  
Recursive unitarity calculation of one-loop amplitudes



### So far computed one-loop amplitudes:

- ✓ N-gluons
- ✓ qq + N-gluons
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NB: N is a parameter in Rocket  
In perspective, for gluons:

N = 6  $\Rightarrow$  10860 diags.

N = 7  $\Rightarrow$  168925 diags.

Successfully computed up to N=20



# Leading color adjustment

Define

$$\mathcal{R}_O = \frac{\int \mathcal{O}(p) d\sigma_{LO}^{\text{FC}}(\mu, p)}{\int \mathcal{O}(p) d\sigma_{LO}^{\text{LC}}(\mu, p)}$$

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$$\mathcal{O}^{\text{NLO}} = r \cdot \mathcal{O}^{\text{NLO,LC}}$$

*Leading color adjustment tested in  $W+1, W+2$  jets and  $W+3$  jets: always OK to 3 %*

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Other  $\mathcal{O}(1\%)$  effects neglected:

- CKM set to unity  $\Rightarrow \sim -1\%$
- W treated onshell  $\Rightarrow \sim +1\%$

# CDF cuts

$$p_{\perp,j} > 20\text{GeV} \quad p_{\perp,e} > 20\text{GeV} \quad E_{\perp,\text{miss}} > 30\text{GeV}$$

$$|\eta_e| < 1.1$$

$$M_{\perp,W} > 20\text{GeV}$$

$$\mu_0 = \sqrt{p_{\perp,W}^2 + M_W^2}$$

$$\mu = \mu_R = \mu_F = [\mu_0/2, 2\mu_0]$$

- PDFs: cteq6l1 and cteq6m
- CDF applies lepton-isolation cuts. This is a  $O(10\%)$  effect. Lepton-isolation has been corrected for (would not have been needed ...)  
**No lepton isolation applied**
- CDF uses JETCLU with  $R = 0.4$ , but this is **not infrared safe**, use a different jet-algorithm

# Jet-algorithms

- CDF uses JETCLU which is not infrared safe
- NLO calculation with JETCLU not possible
- use e.g. SIScone and anti-kt algorithm which are IR safe
- can compare Leading order results for these algorithm (even if meaning of LO for JETCLU is questionable ...)

## Leading order:

Algorithm	$R$	$E_{\perp}^{\text{jet}} > 20 \text{ GeV}$	$E_{\perp}^{\text{3rdjet}} > 25 \text{ GeV}$
JETCLU	0.4	$1.845(2)^{+1.101(3)}_{-0.634(2)}$	$1.008(1)^{+0.614(2)}_{-0.352(1)}$
SIScone	0.4	$1.470(1)^{+0.765(1)}_{-0.560(1)}$	$0.805(1)^{+0.493(1)}_{-0.281(1)}$
anti- $k_{\perp}$	0.4	$1.850(1)^{+1.105(1)}_{-0.638(1)}$	$1.010(1)^{+0.619(1)}_{-0.351(1)}$

SIScone: Salam & Soyez '07;  
anti-kt: Cacciari, Salam, Soyez '08

At LO anti-kt  $R = 0.4$  is closer to JETCLU

## Moral:

*precision comparison with theory require that experiments use IR-safe algorithms*

# Cross-section at the Tevatron

$$\sigma_{W+3j}(p_{\perp,j} > 25 \text{ GeV}) = (0.84 \pm 0.24) \text{ pb}$$

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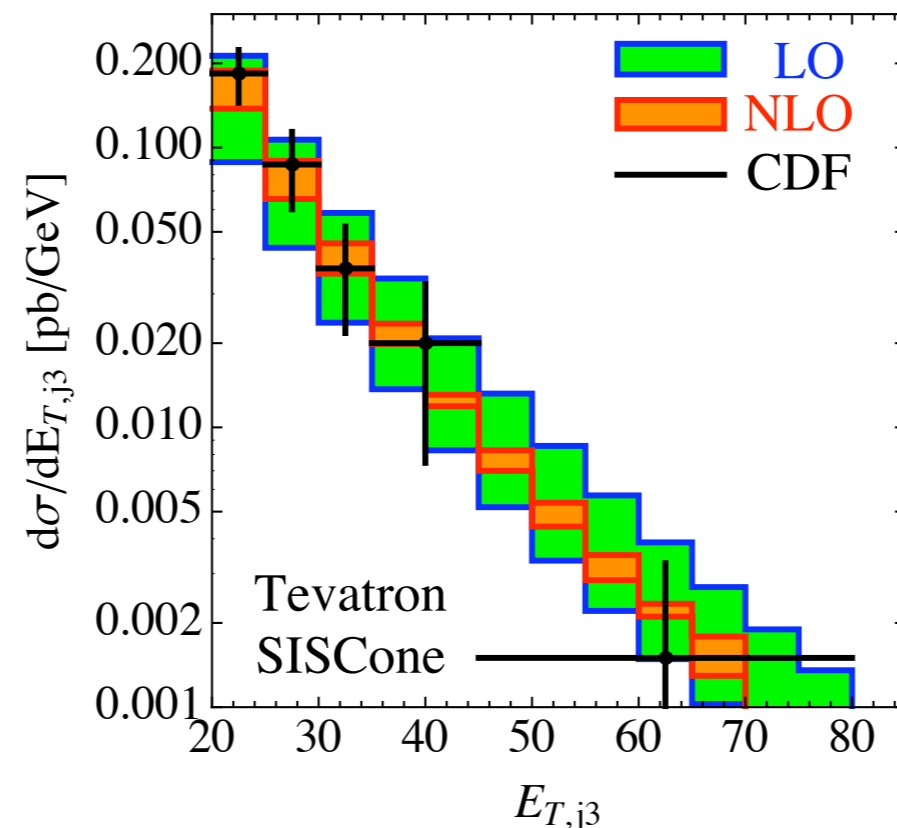
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⇒ important (10% or more) differences due to different jet-algorithms.

*High precision comparison impossible if using different algorithms*

# Tevatron: sample distribution: $E_{T,j3}$

***NB:** CDF  $\Rightarrow$  JetCLU VERSUS NLO Theory  $\Rightarrow$  SISCone*



Ellis et al '09

- ☺ agreement with CDF data (within currently large errors)
- ☺ small  $K=1.0-1.1$ , reduced uncertainty: 50% (LO)  $\rightarrow$  10% (NLO)
- ☺ first applications of new techniques to  $2 \rightarrow 4$  LHC processes

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## Standard procedure

- study a given process with **signal cuts**  $\Rightarrow$  refine theoretical tools
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*How reliable is this procedure ?*

Purpose of background cuts: push into corners of phase-space the SM process, therefore the robustness of the procedure is not assured.  
**NLO QCD predictions for non-trivial processes can shed light on this.**

# $W^+ + 3$ jets at the LHC

In the following: use highly non-trivial NLO calculation of  $W^+ + 3$  jets to illustrate/study this issue

Signal-cut setup (inspired by CMS studies):

$$E_{\text{CM}} = 10 \text{ TeV}$$

$$E_{\perp, \text{jet}} = 30 \text{ GeV}$$

$$E_{\perp, e} = 20 \text{ GeV}$$

$$E_{\perp, \text{miss}} = 15 \text{ GeV}$$

$$M_{\perp, W} = 30 \text{ GeV}$$

$$|\eta_e| < 2.4$$

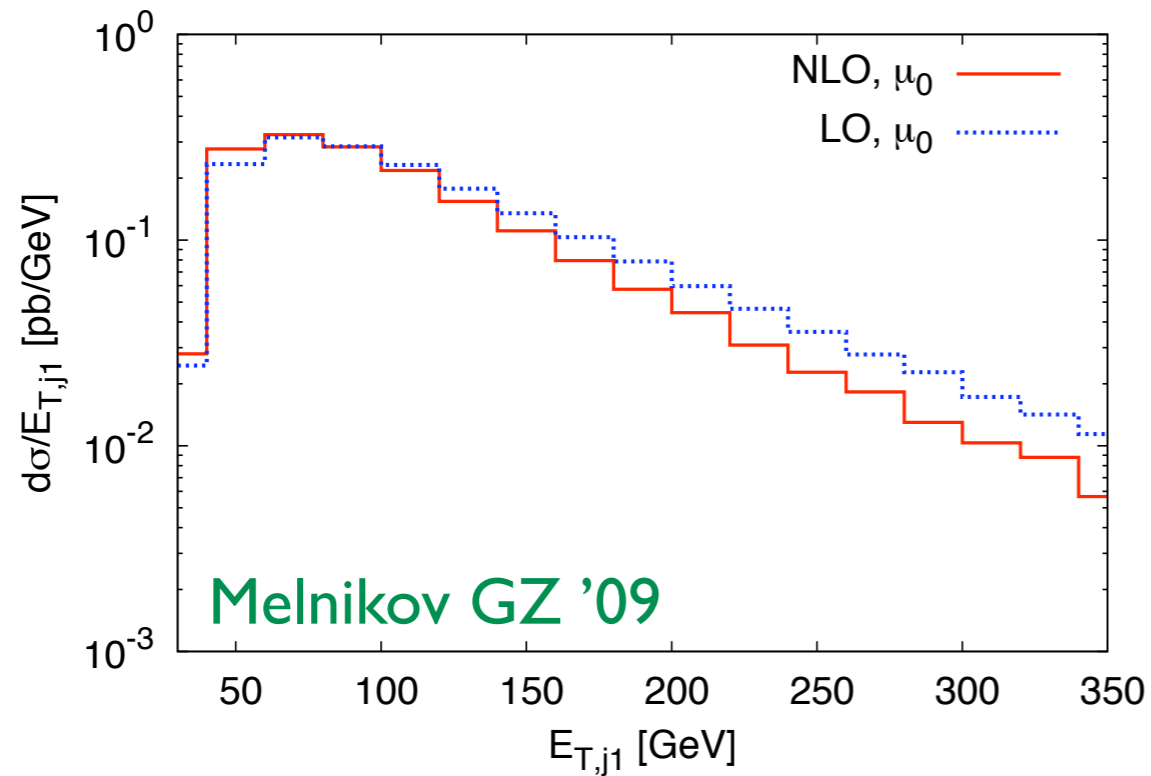
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Jets: SISCone with  $R = 0.5$ ; PDFs: cteq6l1/cteq6m

# Sample transverse energy distribution



Renormalization and factorization scale set to

$$\mu_0 = \sqrt{p_{T,W}^2 + m_W^2}$$

- with scale  $\mu_0$ : considerable change in shape between LO and NLO (extrapolation of LO from low  $p_t$  to high  $p_t$  would fail badly)
- but origin of the change in shape well understood: at high  $E_T$ ,  $\mu_0$  is smaller than typical scales of the QCD branching  $\Rightarrow$  LO overshoots the result

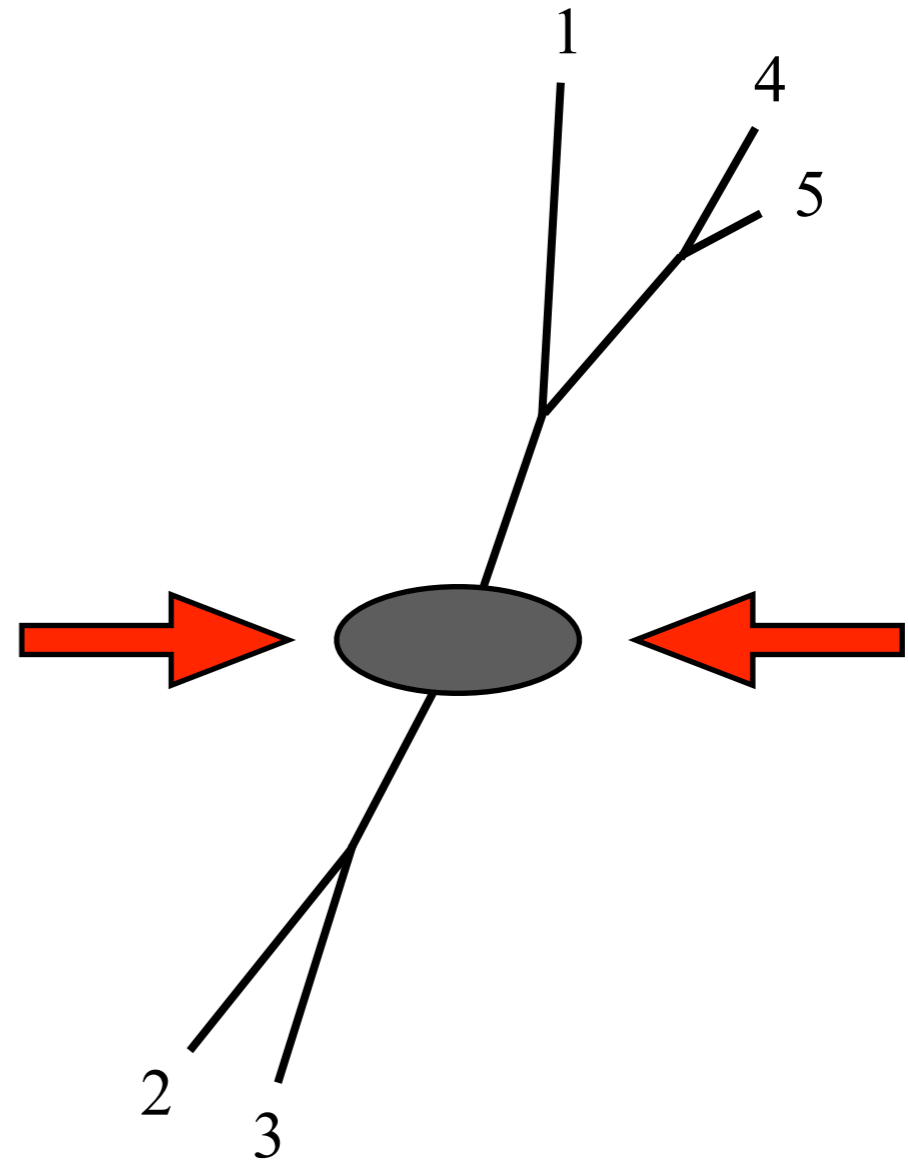
Can one do a more sophisticated LO calculation?



# Local (CKKW) scale

## Local scale choice (CKKW):

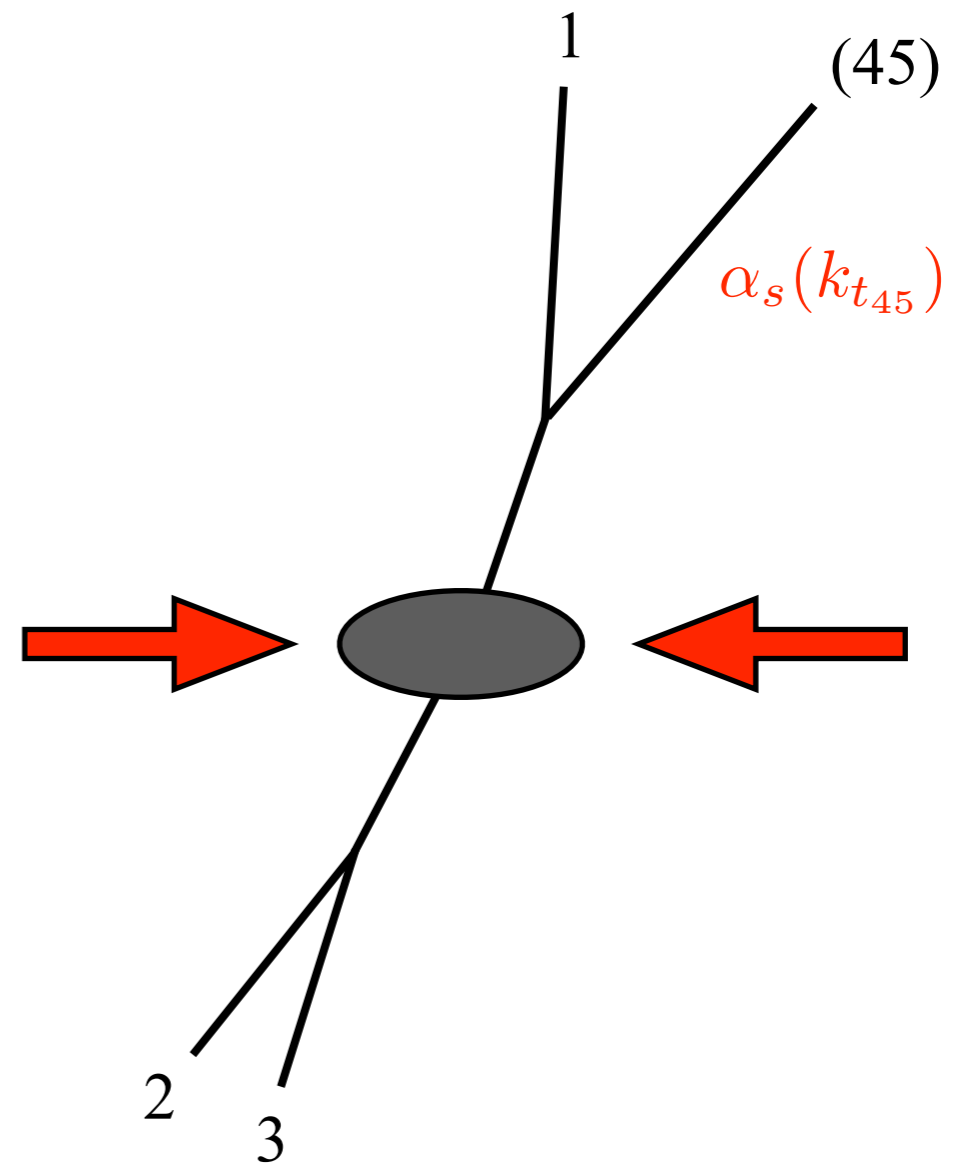
- given a partonic event reconstruct a branching history: cluster partons into jets using  $k_t$ -algorithm
- at each branching the scale in the coupling to set to the relative  $k_t$  of the daughter partons
- local scale = CKKW scale choice, but no Sudakov reweighting, no parton shower



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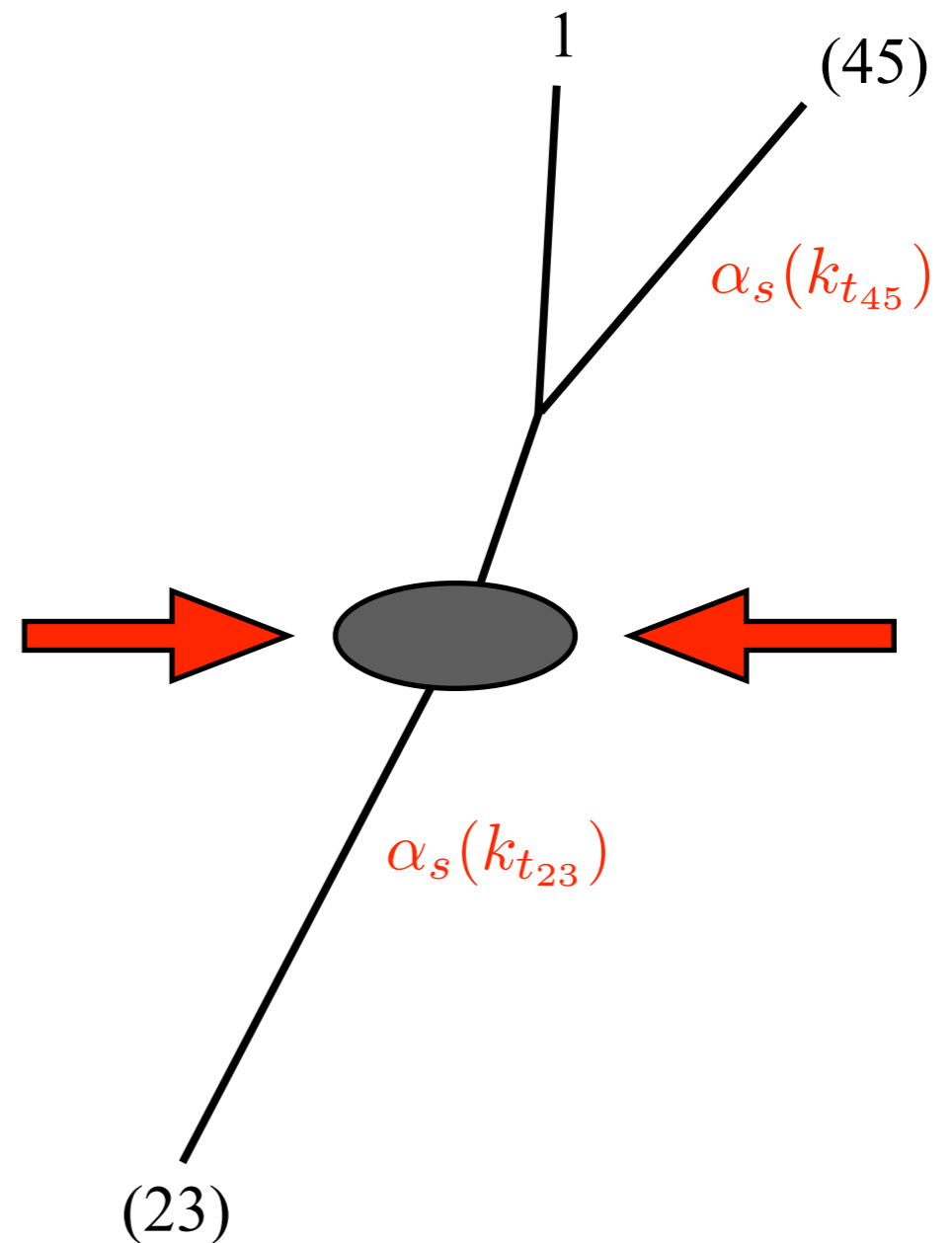
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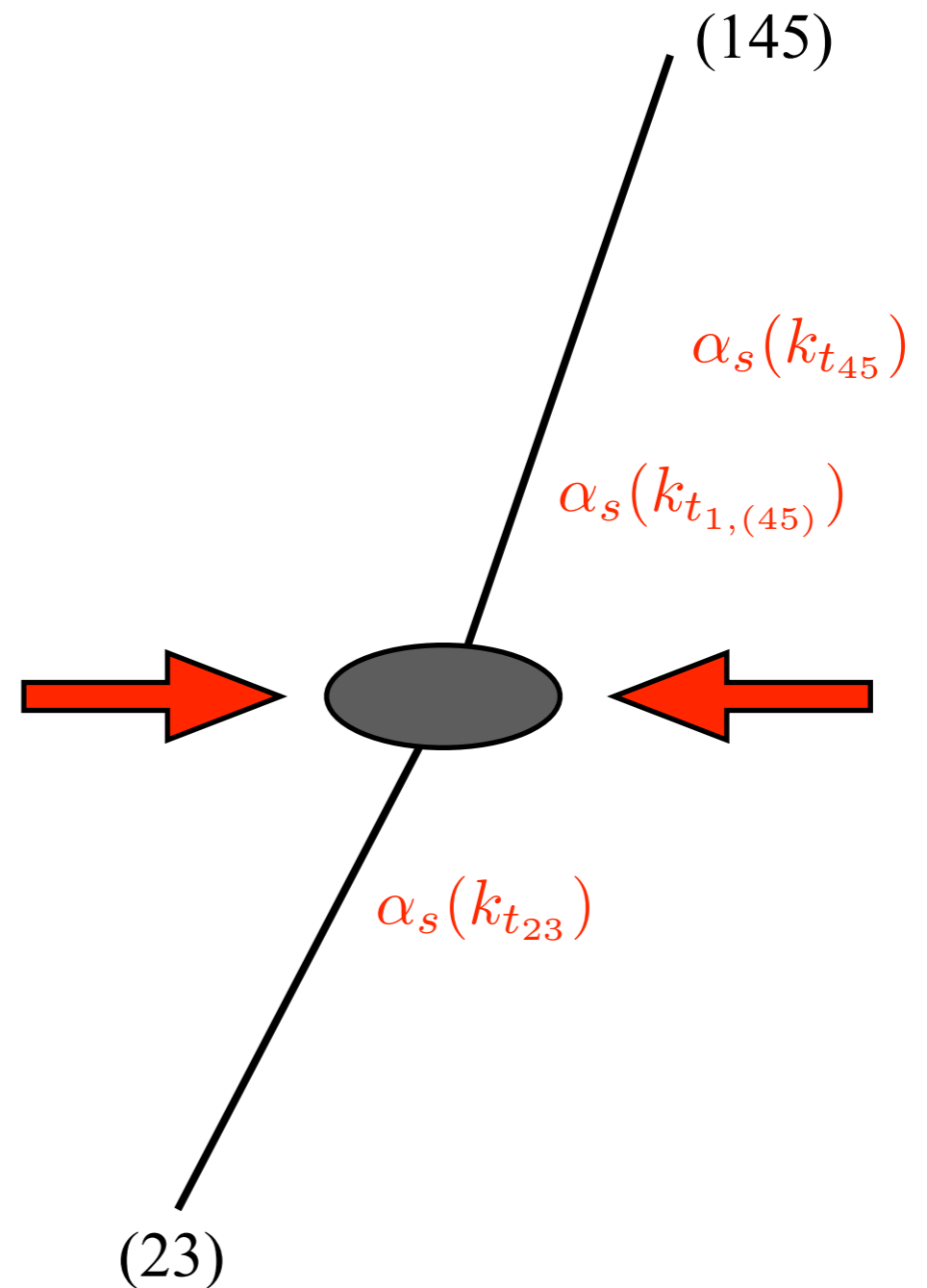
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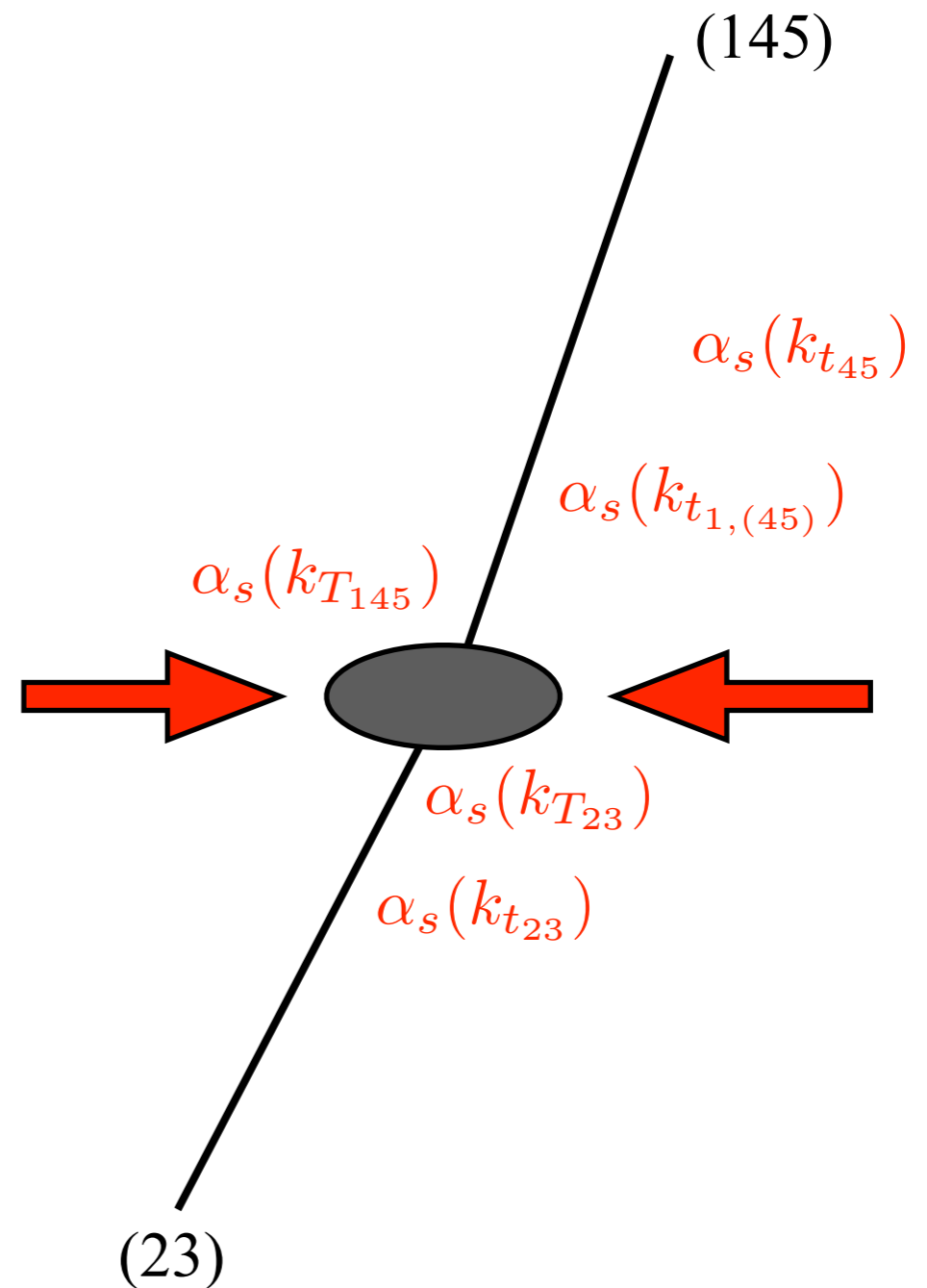
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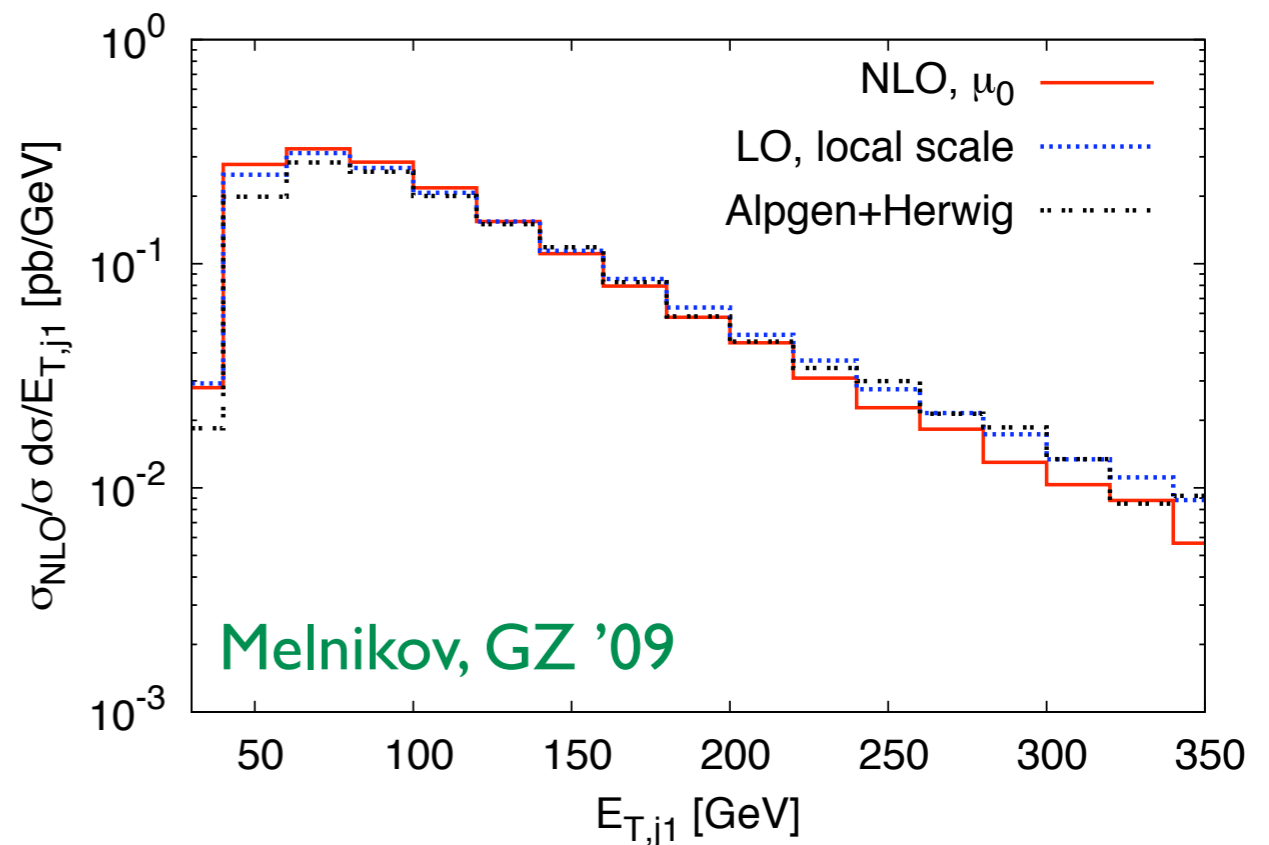
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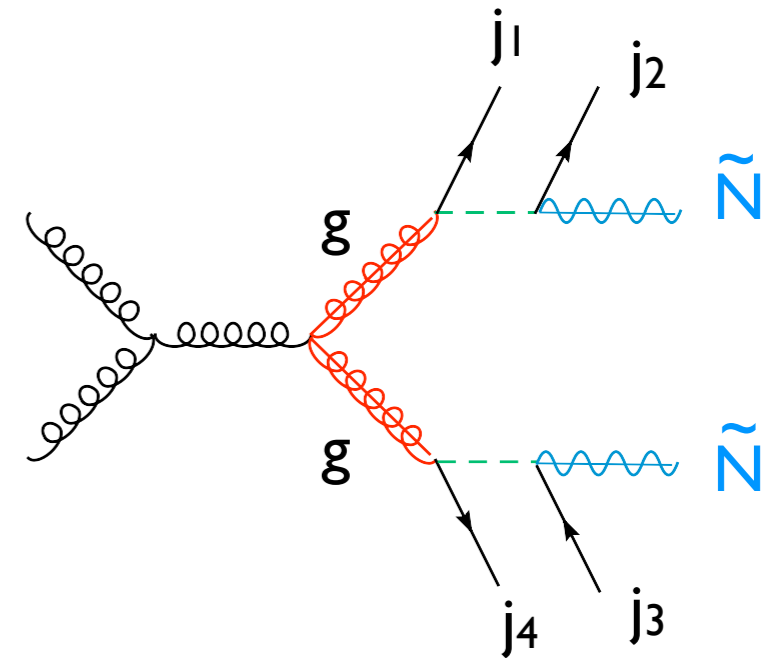
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- ➡ local scale choice reproduces the shape of the NLO distribution well
- ➡ the difference between LO with local scale and full Alpgen+Herwig indicative of the importance of the parton shower



# SUSY signature

SUSY with R-parity: e.g. gluino pair production,  
each decays into 2 jets and neutralino

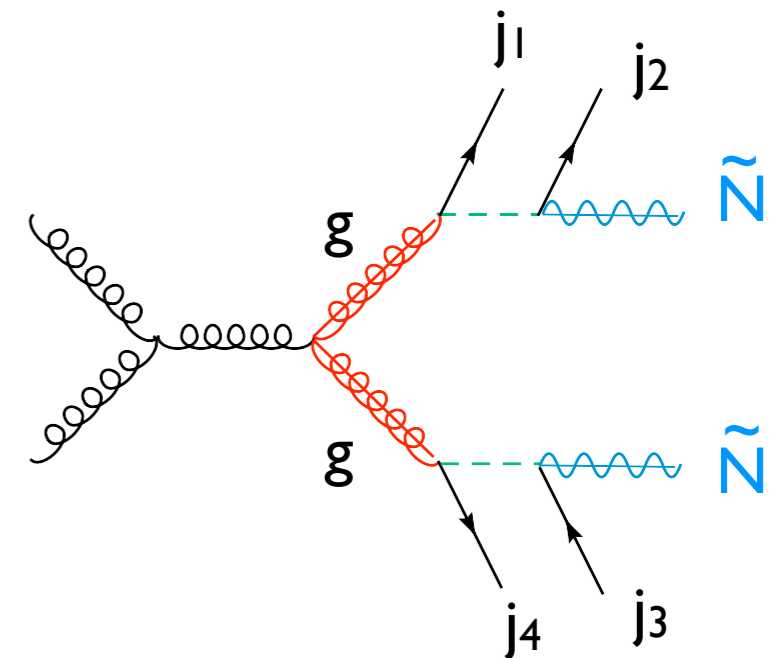
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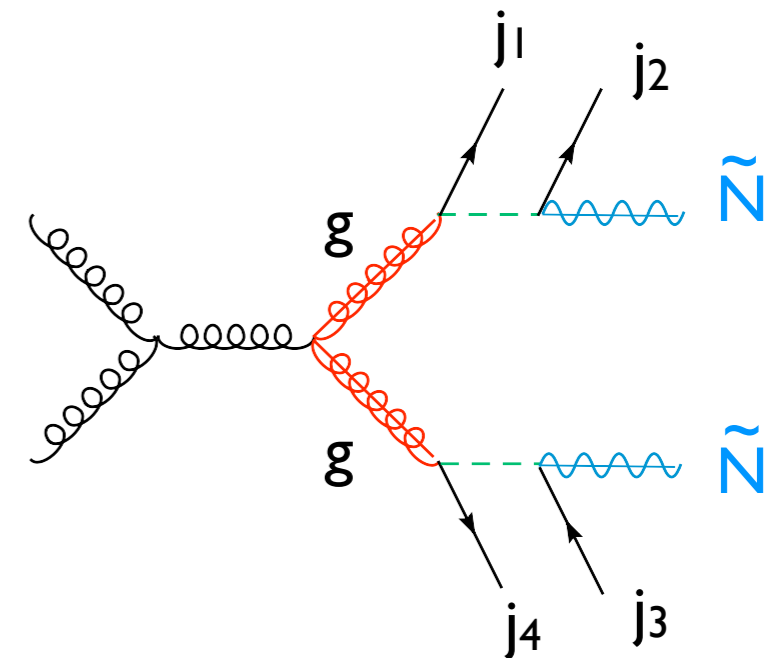
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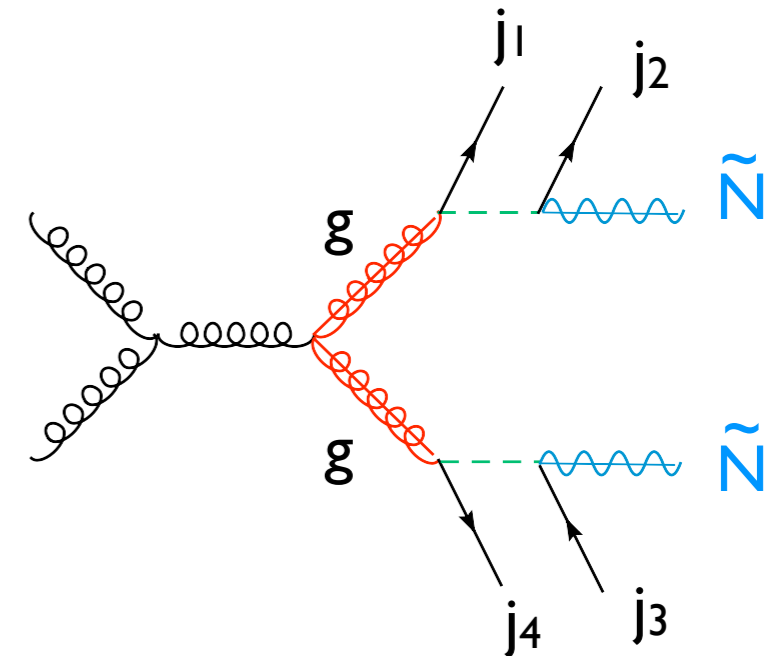
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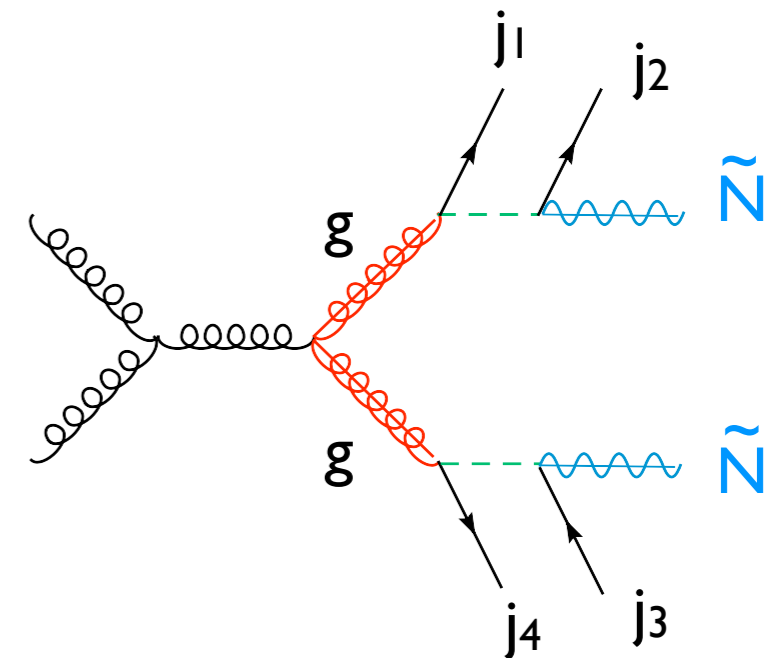
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$\Rightarrow$  important to consider this source of background as well

# Atlas setup

Cuts designed by ATLAS to suppress W+3j background

$$p_{T,j} > 50 \text{ GeV} \quad p_{T,j1} > 100 \text{ GeV} \quad p_{t1} < 20 \text{ GeV}$$

$$E_{T,\text{miss}} > \max(100 \text{ GeV}, 0.2 H_T) \quad H_T = \sum_j p_{T,j} + E_{T,\text{miss}}$$

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Yamazaki [ATLAS and CMS Col.] 0805.3883

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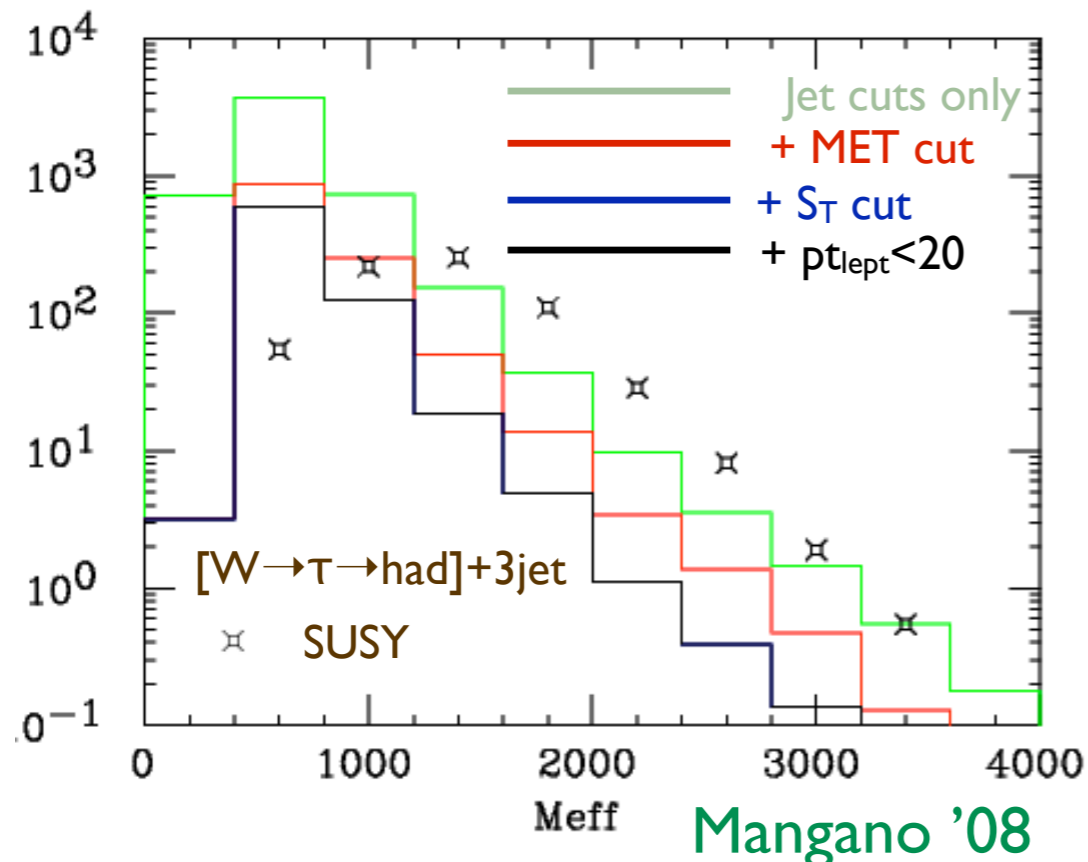
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- each cut suppresses background by factor  $\sim 3$  without modifying the shape
- cut on collinear unsafe sphericity  $S_T$  not applied in the following study

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Our calculation includes only the leptonic decay of the  $W$  (in  $e, \mu$  or  $\tau$ ) but not the hadronic subsequent decay of  $\tau$ . However

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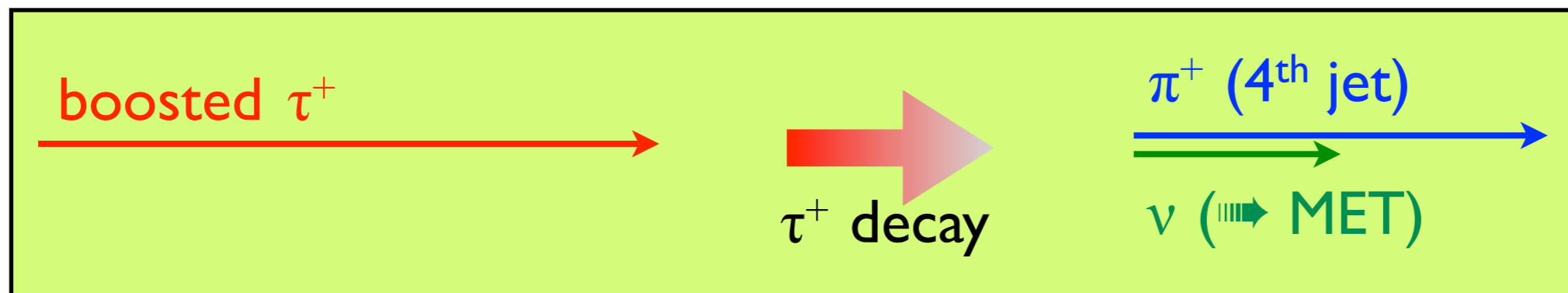
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Theoretical robust approximation:

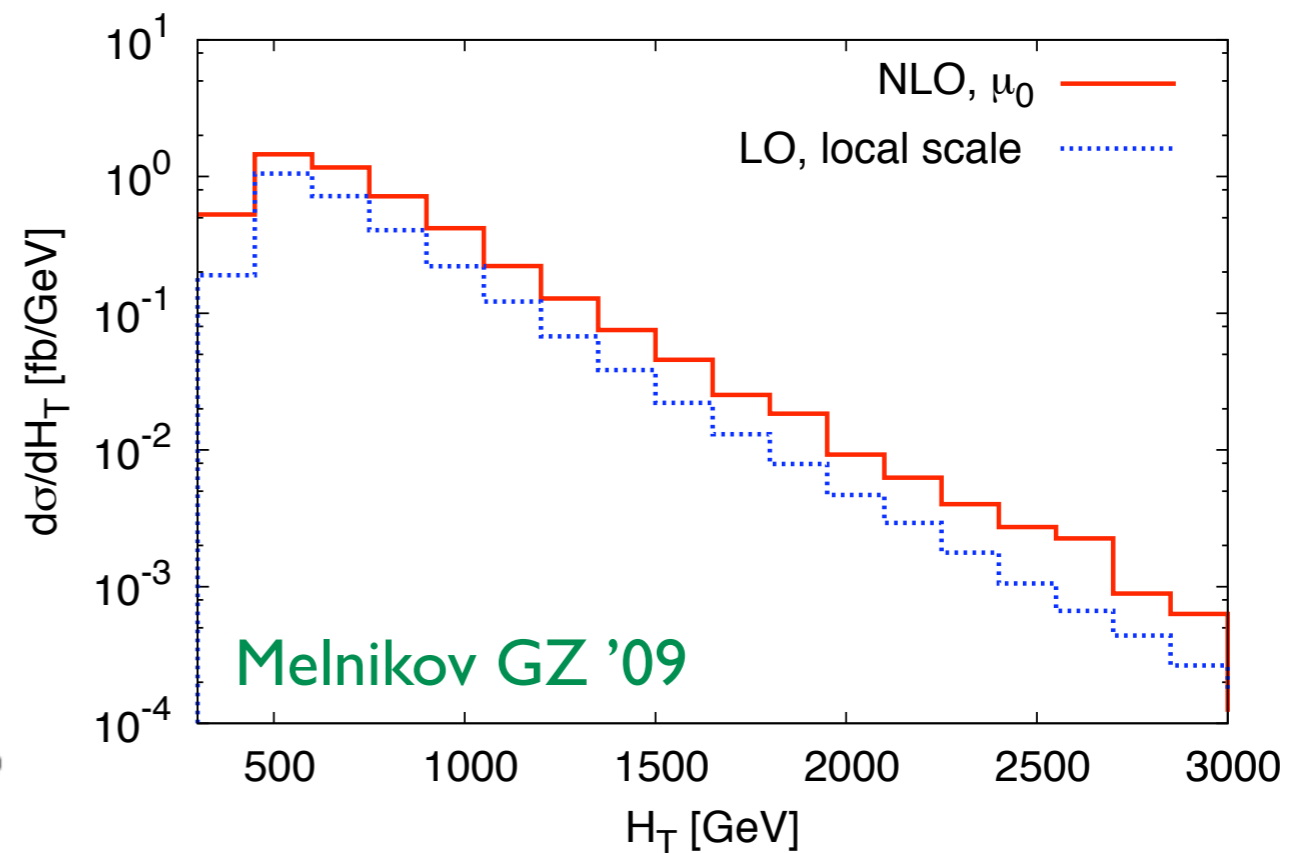
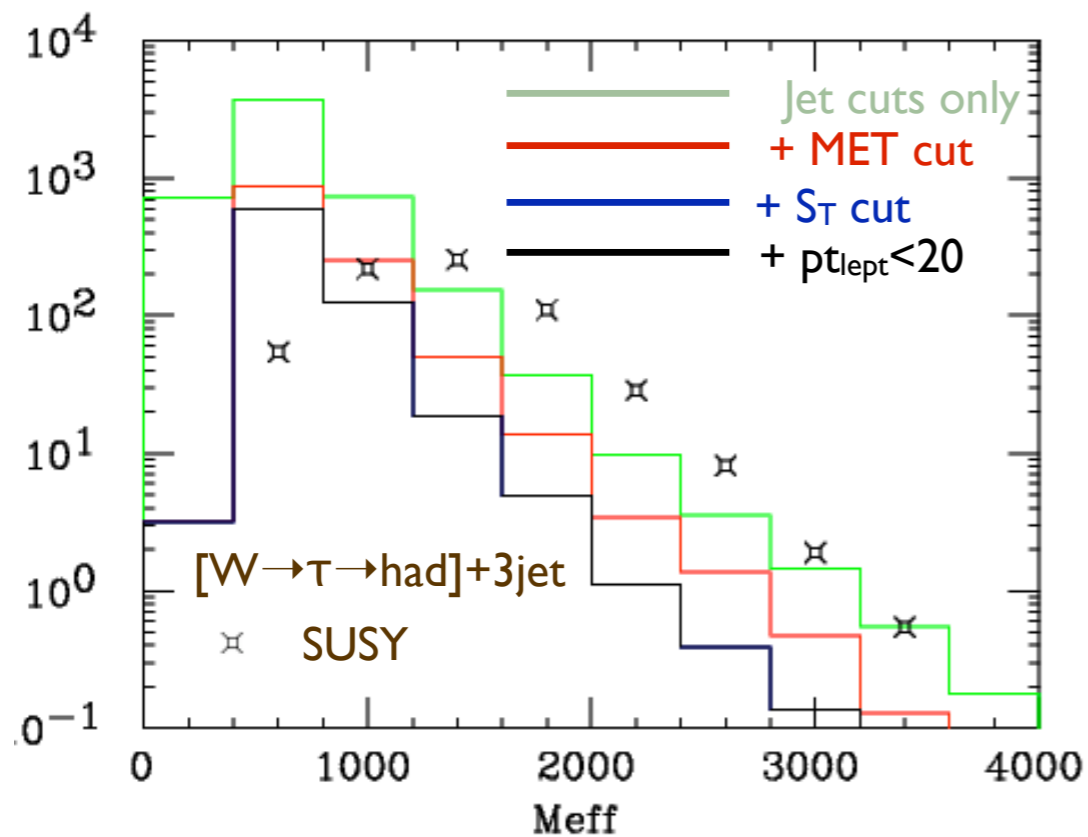
simulate the W decay as a perfect collinear branching with momentum fractions 2/3 ( $\pi^+$ ) and 1/3 ( $\nu$ )



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Primary observable is  $H_T$  (previously called  $M_{\text{eff}}$ ) which 'measures' the SUSY scale:

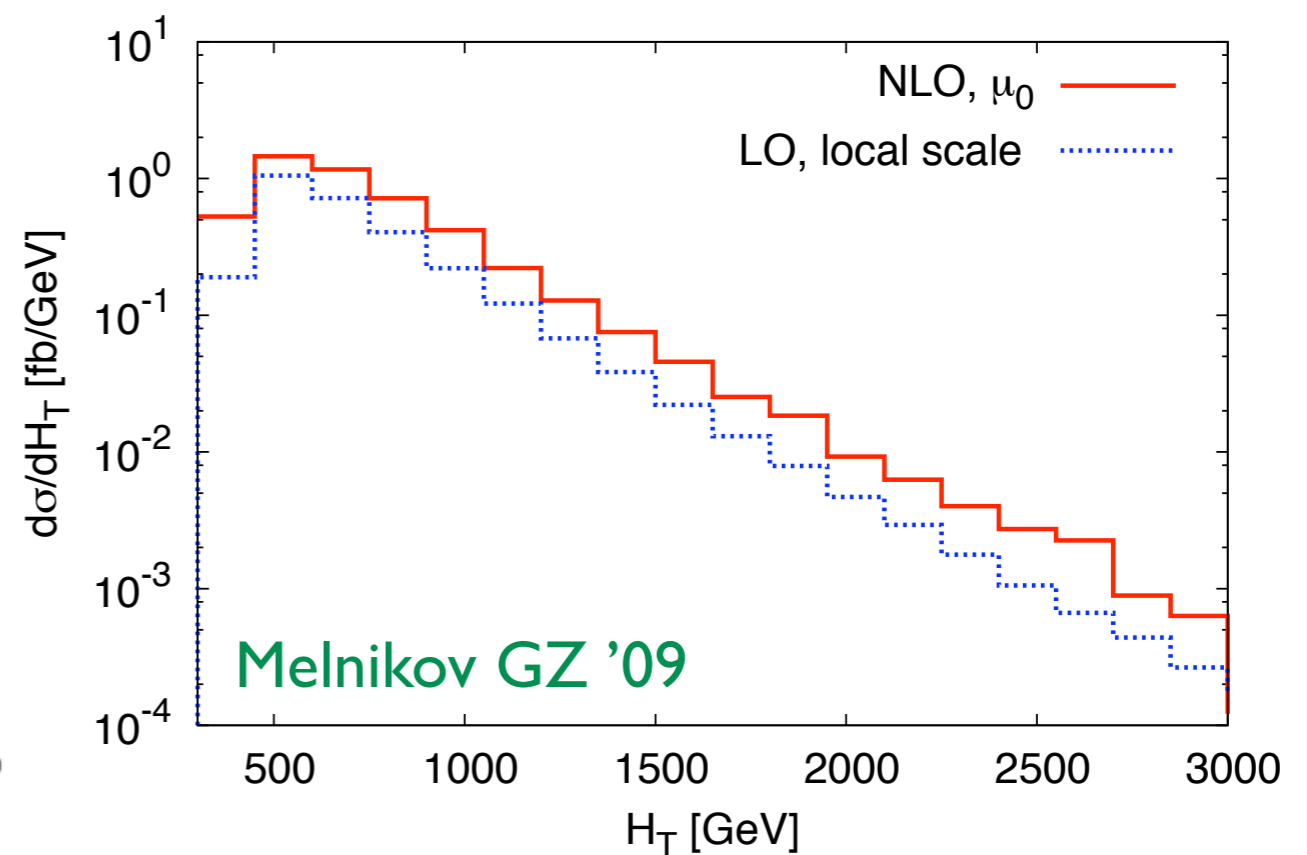
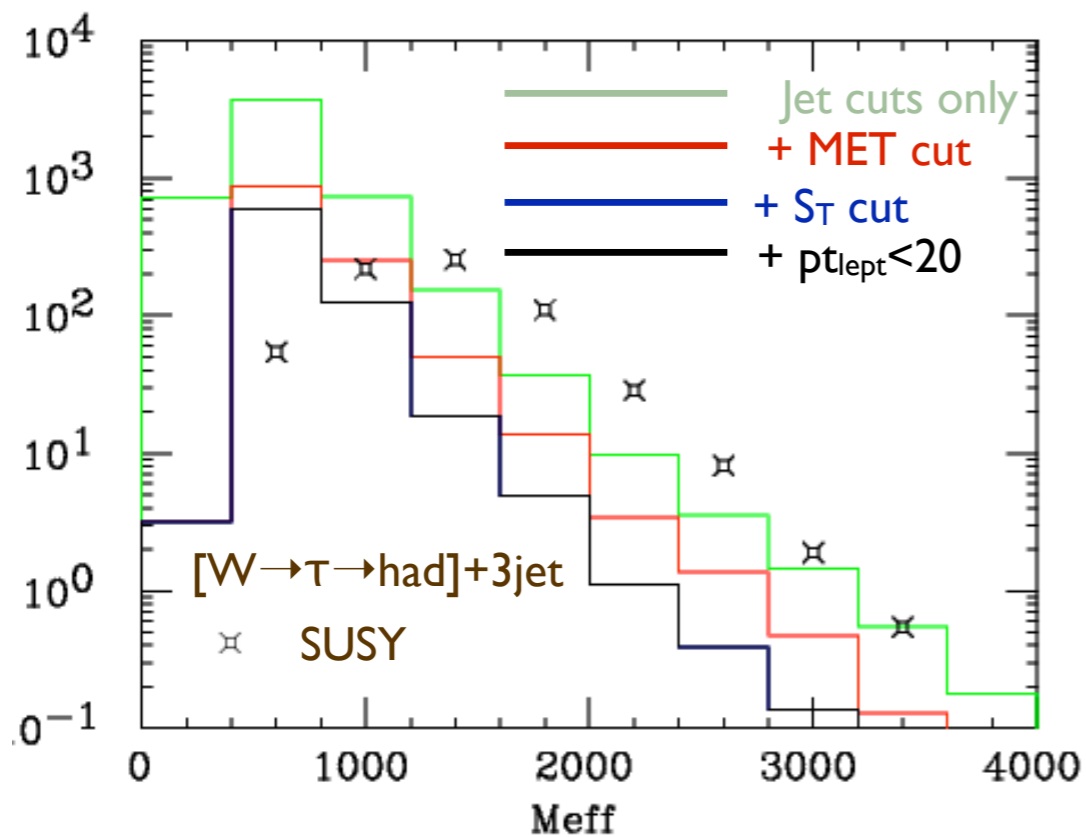
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$$H_T = \sum_j p_{T,j} + E_{T,\text{miss}}$$



- ☛ universal enhancement (K-factor  $\sim 3$ ) of LO without distorting the shape  
NB: *same observable* with cuts as shown before had K-factor  $\sim 1$
- ☛ NLO effect similar to that of cuts but *works in opposite direction*

# CMS style indirect lepton veto cut

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Indirect lepton veto = no explicit lepton veto, but other cuts force contribution from  $W$ +jets to become naturally small

$$p_{T,j} > 30\text{GeV} \quad p_{T,j1} > 180\text{GeV} \quad p_{T,j2} > 110\text{GeV} \quad E_{T,\text{miss}} > 200\text{GeV}$$

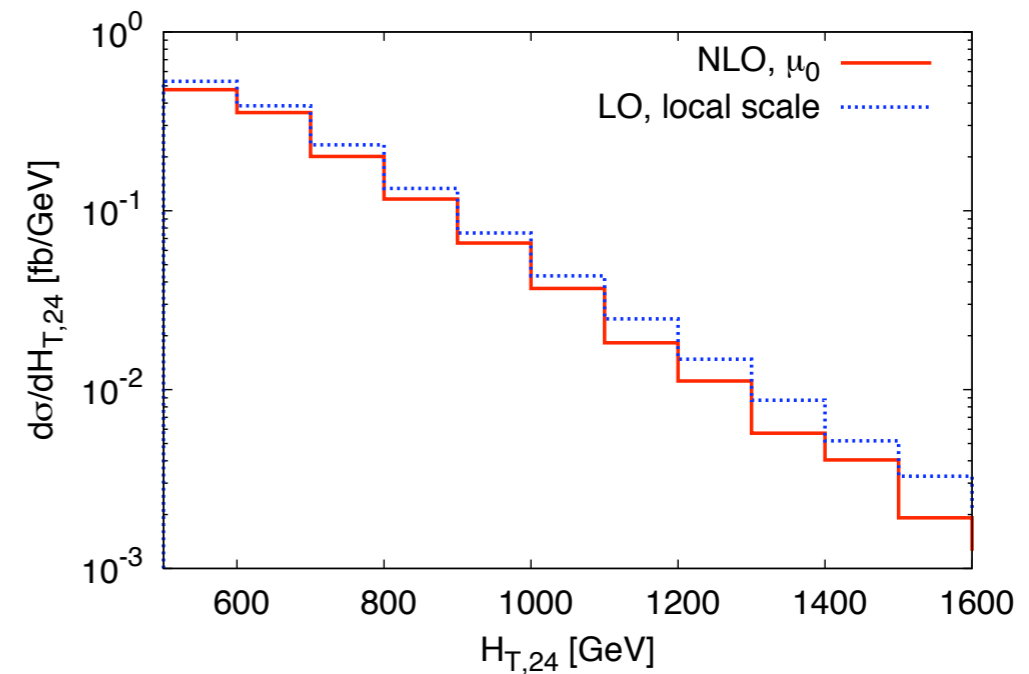
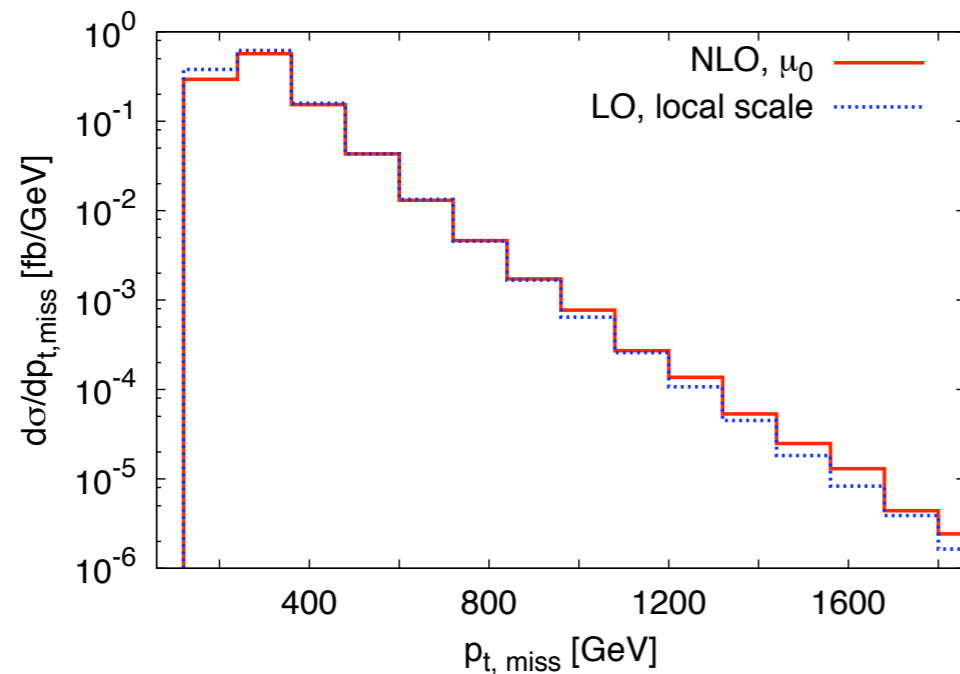
$$|\eta_{\text{lead jet}}| < 1.7 \quad |\eta_{\text{other jets}}| < 3 \quad H_{T,24} = \sum_{j=2}^4 p_{T,j} + E_{T,\text{miss}} > 500\text{GeV}$$

CMS Collaboration Journal Phys. G: Nucl. Part. Phys. 34 (2007) 995

# CMS style indirect lepton veto cut

## Primary search observables

distribution in transverse missing energy and total effective mass  $H_{T,24}$



- NLO correction to cross-section small, K-factor  $\sim 1$
- shapes of LO mostly OK, but moderate shape distortion at high  $H_{T,24}$

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- ☑ all this emphasizes the need to extend NLO corrections to other processes ( $Z+3j, W+4j \dots$ )

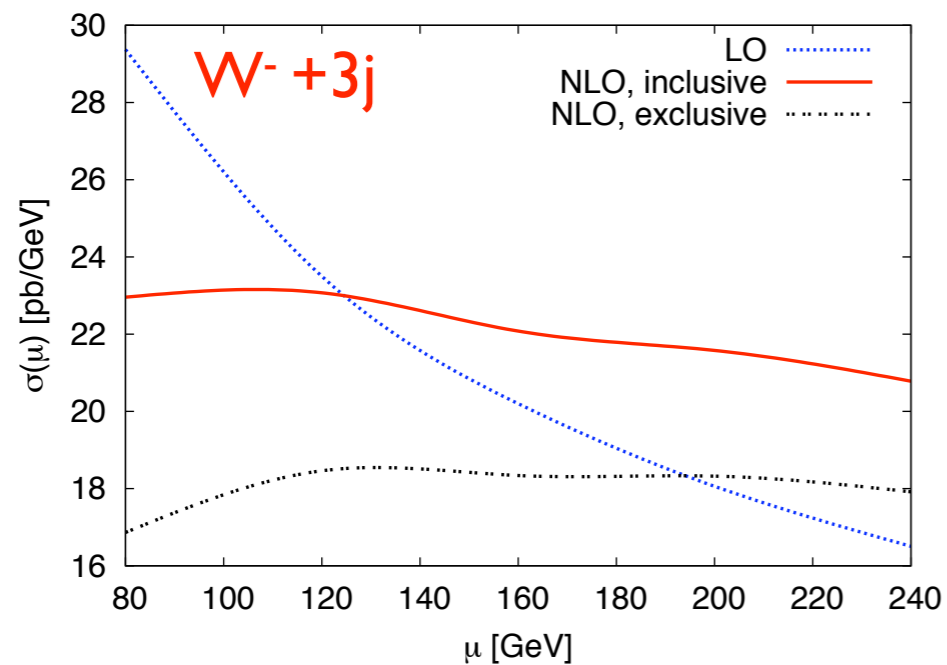
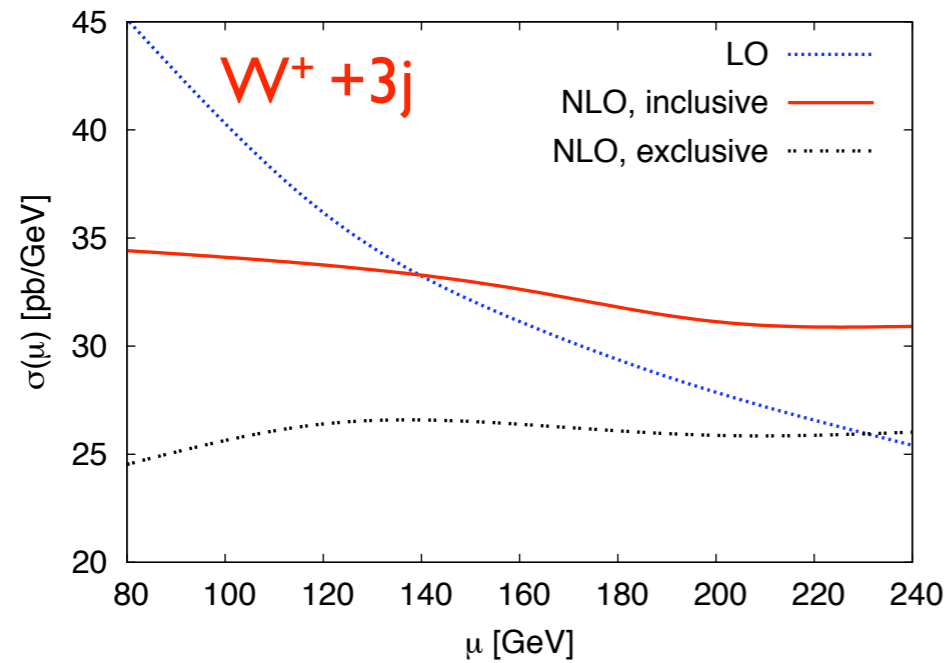
*Extra slides*

# Cross-section calculation

- Consider the NLO **leading color approximation**, keep  $n_f$  dependence exact (important for beta function) but neglect  $1/N_c^2$  terms
- Real radiation part:
  - leading color tree level  **$W+6$  parton amplitudes computed recursively**
  - we use **Catani-Seymour subtraction** terms modified to deal with the minimal set of color structures needed at leading color
- Real + virtual implemented in the **MCFM parton level integrator**

Full-color NLO calculation done by Berger et al. '09

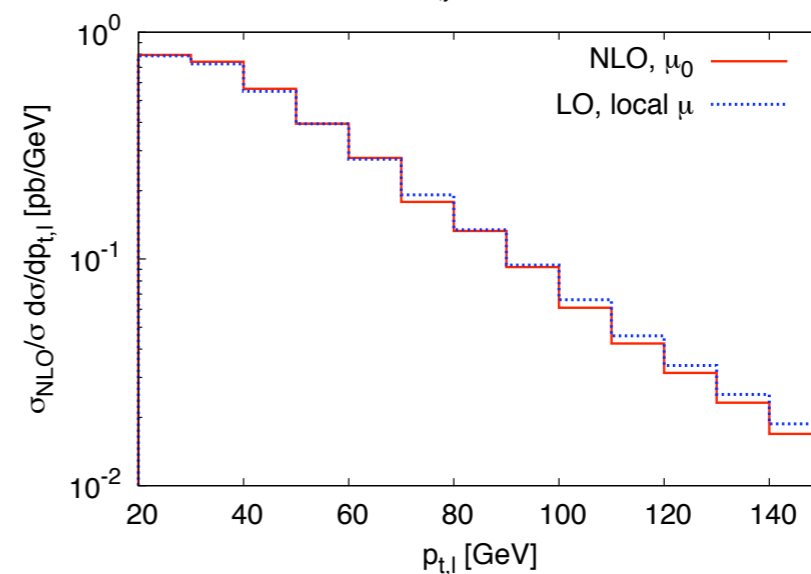
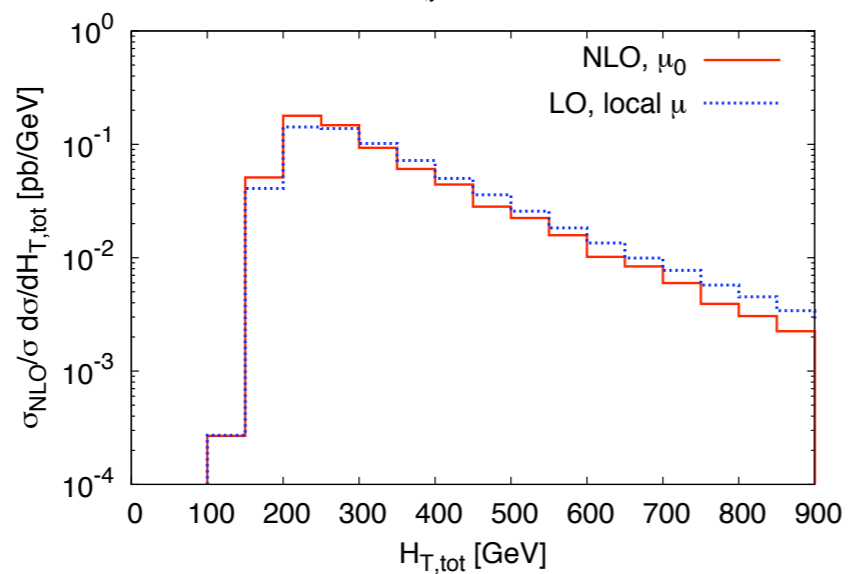
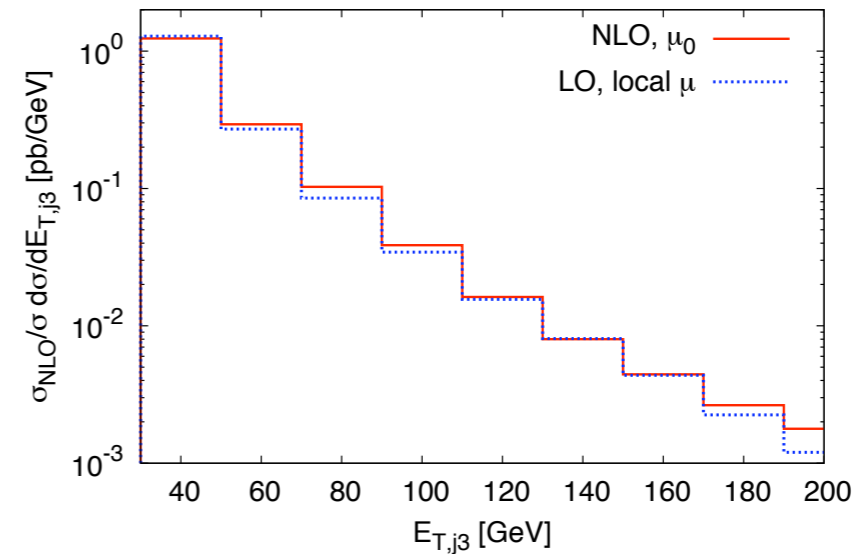
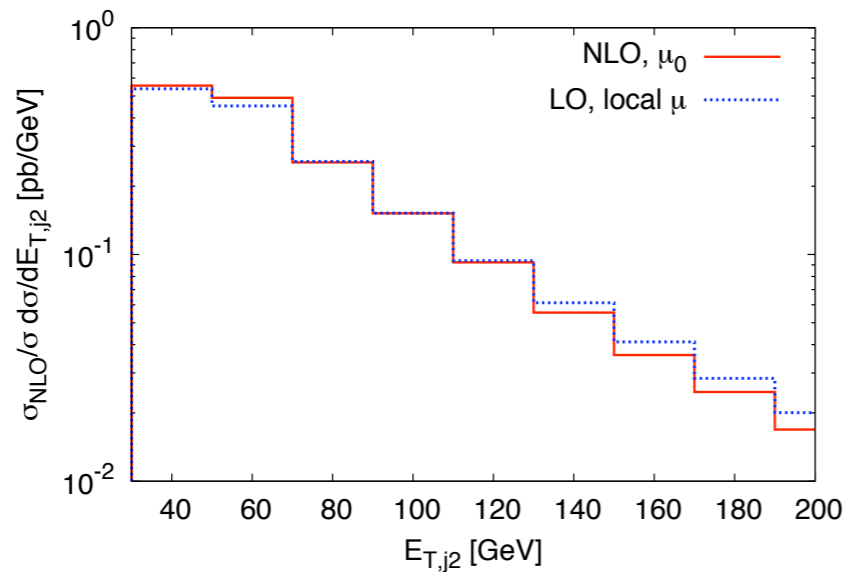
# Scale dependence



Melnikov & GZ '09

- scale dependence considerably reduced at NLO (both inclusive and exclusive)
- NLO tends to reduce cross-section
- because of very large scale dependence of LO, quoting a K-factor not very meaningful

# Other hadronic distributions



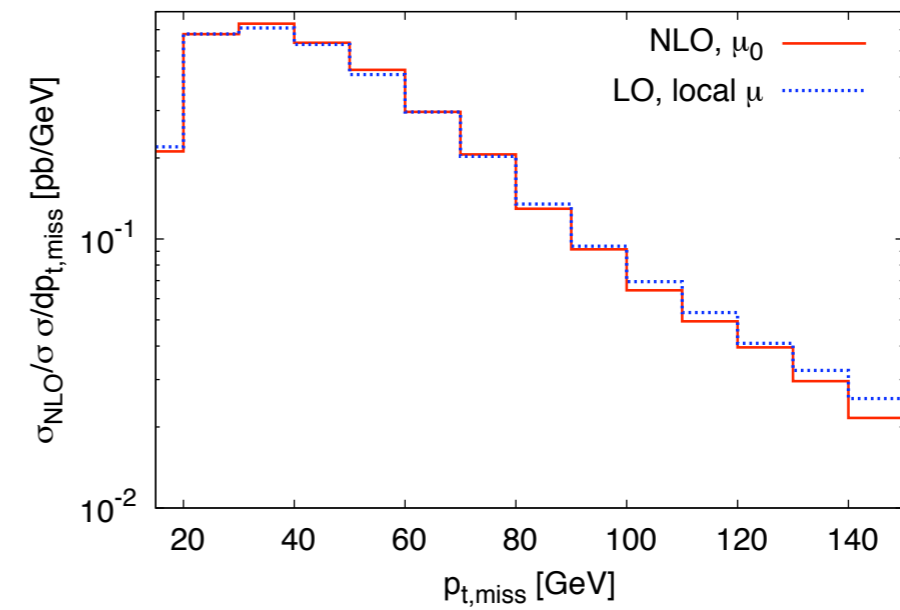
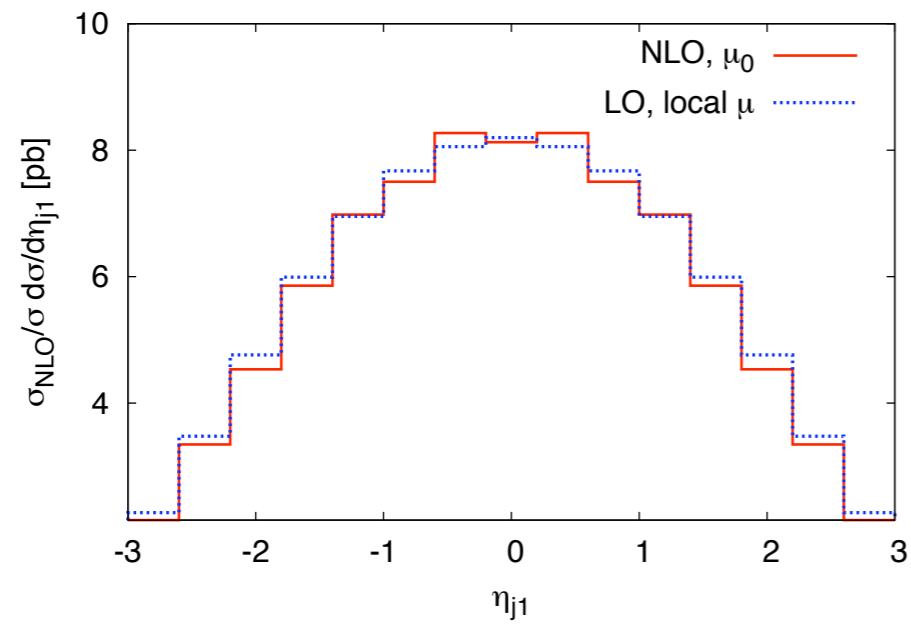
Melnikov GZ '09

👉 LO with local scale does a very reasonable job in reproducing shapes

**NB:**  
normalization of LO remains out of control. LO is normalized to NLO in above plots



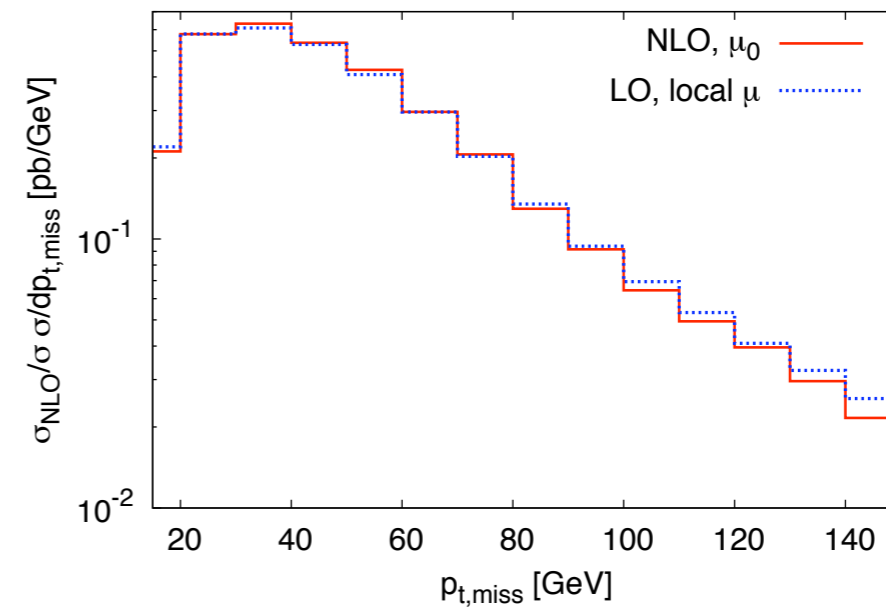
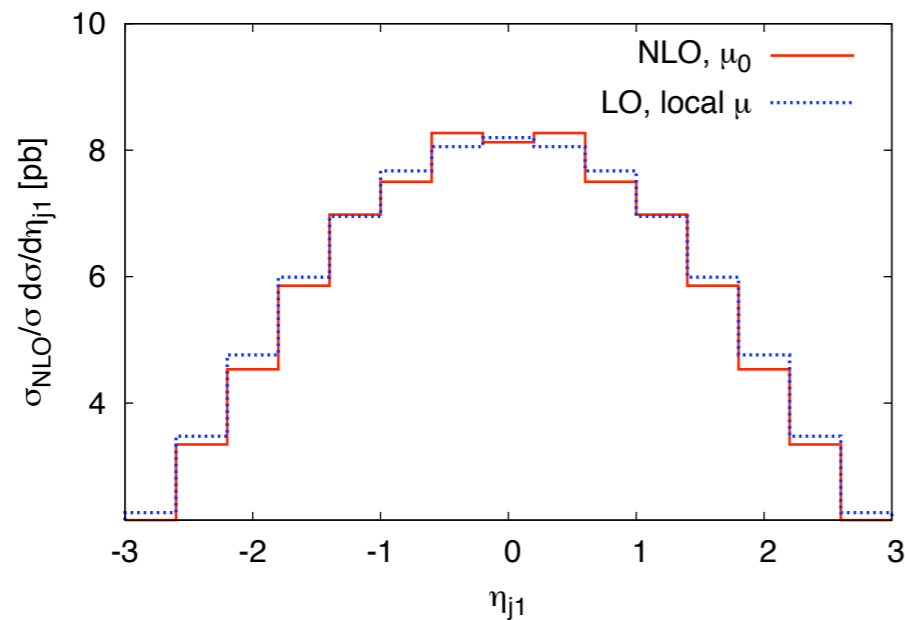
# Leptonic distributions



Melnikov GZ '09

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Melnikov GZ '09

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*How solid (cut-independent) is this statement ?*

*See what happens with different cuts.*

*Consider two sets of cuts where  $W+3\text{jet}$  plays the role of unwanted background*