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### Testing the weak gravity conjecture using type I strings with broken supersymmetry

Quentin Bonnefoy

Centre de Physique Théorique - École Polytechnique

based on arXiv:1811.11199 and work in progress, in collaboration with E. Dudas and S. Lüst

> SUSY 2019 May 23rd, 2019

The swampland and the WGC  $_{\odot OOOO}$ 

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# The swampland and the weak gravity conjecture

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#### Consistent effective field theories: anomalies, unitarity, etc

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#### Consistent effective field theories: anomalies, unitarity, etc

When coupled to quantum gravity: a landscape and a swampland Vafa '05

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Swampland conjectures characterize the landscape



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#### Consistent effective field theories: anomalies, unitarity, etc

When coupled to quantum gravity: a landscape and a swampland Vafa '05

Swampland conjectures characterize the landscape , ex:

•	no exact global symmetries	see Banks, Seiberg '10
•	completeness of the charge la	ttice Polchinski '03
•	distance conjecture	Ooguri, Vafa '06
•	weak gravity conjecture	Arkani-Hamed, Motl, Nicolis, Vafa '06
•	no stable non-SUSY AdS	Ooguri, Vafa '16
•	de Sitter conjecture	Obied, Ooguri, Spodyneiko, Vafa '18

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Weak gravity conjecture:

Arkani-Hamed, Motl, Nicolis, Vafa '06

For every gauge field, there must exist a charged state such that  $gQM_P \ge M$ 



Motivation:

- enable extremal black holes to decay
- account for the absence of global symmetries in quantum gravity

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Motivation:

- enable extremal black holes to decay
- account for the absence of global symmetries in quantum gravity

One of the best motivated conjectures: arguments from BHs, holography, string theory...

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Can be extended (multiple gauge fields, higher dimensions, scalar fields, p-forms...)



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Can be extended (multiple gauge fields, higher dimensions, scalar fields, *p*-forms...)

For p-forms in dimension d with a dilatonic gravity:

Arkani-Hamed, Motl, Nicolis, Vafa '06, Heidenreich, Reece, Rudelius '15

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Can be extended (multiple gauge fields, higher dimensions, scalar fields, *p*-forms...)

For p-forms in dimension d with a dilatonic gravity:

Arkani-Hamed, Motl, Nicolis, Vafa '06, Heidenreich, Reece, Rudelius '15

For every *p*-form field, there must exist a charged (extended) object such that  $e^2Q^2 \ge 8\pi G\left(\frac{\alpha^2}{2} + \frac{p(d-p-2)}{d-2}\right)T^2$ 

The swampland and the WGC  $0000 \bullet$ 

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# String theory has p-forms. For R-R forms, the charged objects are D-branes.

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String theory has p-forms. For R-R forms, the charged objects are D-branes.

In the superstring, they are BPS-states and saturate the WGC.

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How does supersymmetry breaking affect this statement?

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String theory has p-forms. For R-R forms, the charged objects are D-branes.

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# How does supersymmetry breaking affect this statement?

In what follows:

a test of the WGC for the R-R 2-form in type I string theory with broken supersymmetry

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## The string theory setup

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#### Type I string theory: unoriented type IIB

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#### **Type I string theory**: unoriented type IIB

Massless spectrum: 10D  $\mathcal{N} = 1$  supergravity with SO(32) gauge group



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Massless spectrum: 10D  $\mathcal{N} = 1$  supergravity with SO(32) gauge group

Unoriented worldsheets, ex: one-loop closed amplitude in 10-dimensional spacetime

$$\frac{1}{2}\mathcal{T} = \frac{1}{2}\# \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2^6} \left| \frac{V_8 - S_8}{\eta^8} \right|^2(\tau) \\ \left[ V_8 = \frac{\theta_3^4 - \theta_4^4}{2\eta^4} , \ S_8 = \frac{\theta_2^4 + \theta_1^4}{2\eta^4} \right]$$

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#### Supersymmetry breaking: Scherk-Schwarz mechanism

Scherk, Schwarz '79



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Generally: mass shifts via twisted boundary conditions for a field theory defined on a compact manifold Setup 00●00 D1-D1 interactions and WGC 00000

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Supersymmetry breaking: Scherk-Schwarz mechanism Scherk, Schwarz '79

Generally: mass shifts via twisted boundary conditions for a field theory defined on a compact manifold

On a circle of radius R: field $(x + 2\pi R)$  = symmetry × field(x)

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$$\implies \phi(x) = e^{\frac{ixQ}{2R}} \sum_m e^{\frac{imx}{R}} \phi_n \implies M_{\text{KK}} = \left| \frac{m}{R} + \frac{Q}{2R} \right|$$

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Always understood as a **spontaneous breaking** 

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#### Supersymmetry breaking: Scherk-Schwarz mechanism Scherk, Schwarz '79

For our type I string, use  $(-)^F \delta_{x \to x + \pi R}$ .



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#### Supersymmetry breaking: Scherk-Schwarz mechanism Scherk, Schwarz '79

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$$\mathcal{T}_{\text{SUSY}} = \int_{\mathcal{F}} \frac{d^2 \tau}{\tau_2^{11/2}} \left| \frac{V_8 - S_8}{\eta^8} \right|^2 \Lambda_{m,n}$$

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SUSY generates exponential quintessence-like runaway potentials:

$$V = -\left(\frac{\mathcal{T}}{2} + \mathcal{K} + \mathcal{A} + \mathcal{M}\right) \sim \Lambda^4 e^{-c\Phi}$$

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Favored by the **de Sitter conjecture**:

Obied, Ooguri, Spodyneiko, Vafa '18

Every scalar potential V should verify  $|\nabla V| \ge cV$ , with c > 0 and  $c = \mathcal{O}(1)$ 

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SUSY enables tests of the mutual compatibility of swampland conjectures

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### D1-D1 interactions and WGC

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#### We compute one-loop interactions between branes

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#### We compute one-loop interactions between branes

Closed string exchange  $\iff$  open-string cylinder calculation (with Dirichlet-Dirichlet boundary conditions):



#### figure from JHEP 0305 (2003) 055

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We focus on D1-D1 systems:

$$\mathcal{A}_{11} = \frac{1}{\pi\sqrt{\alpha'}} \int_0^\infty \frac{d\tau_2}{\tau_2^{3/2}} e^{-\frac{\tau_2 r^2}{4\pi\alpha'}} \left[ P_m - P_{m+1/2} \right] \frac{\theta_2^4}{2\eta^{12}} \left( \frac{i\tau_2}{2} \right)$$

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One-loop attraction

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One-loop attraction

#### Does this violate the WGC?

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• "charge=tension" at massless level: higher-order amplitudes important at large distance

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- "charge=tension" at massless level: higher-order amplitudes important at large distance
- decrease of the tension of a brane: source for those higher-order effects

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- "charge=tension" at massless level: higher-order amplitudes important at large distance
- **decrease of the tension** of a brane: source for those higher-order effects

$$V_{11} \sim \frac{1}{M_P^2} \left[ \frac{\frac{4}{3}Q^2 - M^2 - \frac{1}{3}M^2 e^{-m_{\phi}r}}{r} - \frac{Q^2}{6} \frac{e^{-r\sqrt{\frac{R^2}{\alpha'^2} - \frac{2}{\alpha'}}}}{r} \right]$$

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**SUSY breaking in string theory** generates the necessary ingredients for **non-trivial and compatibility tests of the swampland conjectures**: brane interactions, runaway potentials...

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Our calculation supports the weak gravity conjecture for the R-R 2-form of type I string theory with broken SUSY

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Need additional effort (higher-loop calculations/corrections to the effective action, application to actual black holes, more realistic quintessence models)

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SUSY breaking in string theory generates the necessary ingredients for non-trivial and compatibility tests of the swampland conjectures: brane interactions, runaway potentials...

#### Our calculation supports the weak gravity conjecture for the R-R 2-form of type I string theory with broken SUSY

Need additional effort (higher-loop calculations/corrections to the effective action, application to actual black holes, more realistic quintessence models)

Future directions: other SUSY and non-SUSY tests (ex: more gauge fields)

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# Thank you!

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