



PUCP



Breaking of CPT due to quantum decoherence tested at DUNE

Félix N. Díaz Desposorio

J.C. Carrasco, F.N. Díaz and A.M. Gago PhysRevD.99.075022

Pontificia Universidad Católica del Perú

XXVIIth International Conference on Supersymmetry and Unification of Fundamental Interactions
Corpus Christi - 2019



Outline

- 1. Introduction.**
- 2. Theoretical Approach.**
- 3. Simulation Details.**
- 4. Results.**
- 5. Conclusions.**

Introduction

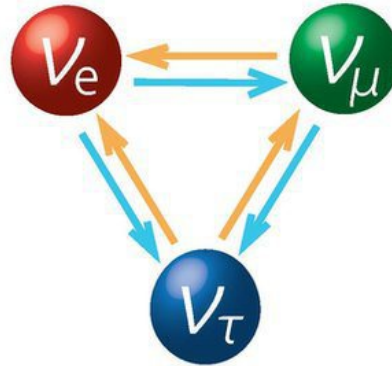
| | | | | | |
|----------------|---|---------------------------------------|--------------------------------------|-------------------------|-------------------------|
| mass → | ≈2.3 MeV/c ² | ≈1.275 GeV/c ² | ≈173.07 GeV/c ² | 0 | ≈126 GeV/c ² |
| charge → | 2/3 | 2/3 | 2/3 | 0 | 0 |
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| | u up | c charm | t top | g gluon | H Higgs boson |
| QUARKS | | | | | |
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$$\frac{\sigma(e^-e^-)}{\sigma(\nu_e e^-)} \approx \frac{10^{-33} \text{ cm}^2}{10^{-41} \text{ cm}^2} = 10^8$$

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|---|--|--|---|---|--|
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| QUARKS | mass → ≈4.8 MeV/c ² charge → -1/3 spin → 1/2 d down | mass → ≈95 MeV/c ² charge → -1/3 spin → 1/2 s strange | mass → ≈4.18 GeV/c ² charge → -1/3 spin → 1/2 b bottom | mass → 0 charge → 0 spin → 1 γ photon | |
| | LEPTONS | mass → 0.511 MeV/c ² charge → -1 spin → 1/2 e electron | mass → 105.7 MeV/c ² charge → -1 spin → 1/2 μ muon | mass → 1.777 GeV/c ² charge → -1 spin → 1/2 τ tau | mass → 91.2 GeV/c ² charge → 0 spin → 1 Z Z boson |
| mass → 2.2 eV/c ² charge → 0 spin → 1/2 ν _e electron neutrino | | mass → <0.17 MeV/c ² charge → 0 spin → 1/2 ν _μ muon neutrino | mass → <15.5 MeV/c ² charge → 0 spin → 1/2 ν _τ tau neutrino | mass → 80.4 GeV/c ² charge → ±1 spin → 1 W W boson | |
| | | | | | |

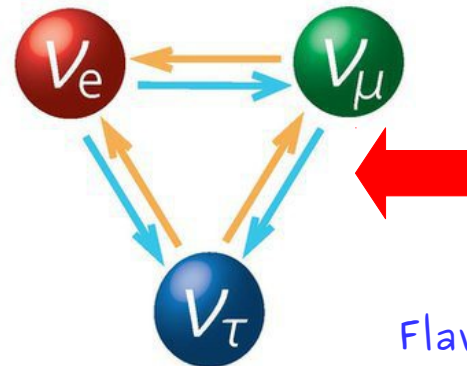
Oscillation



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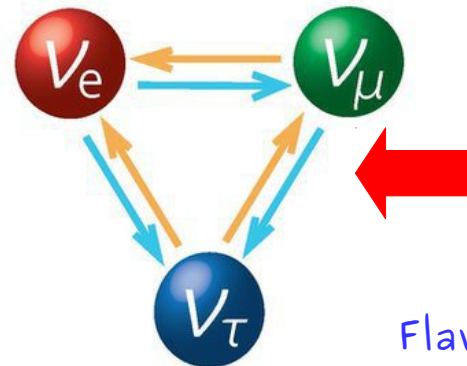
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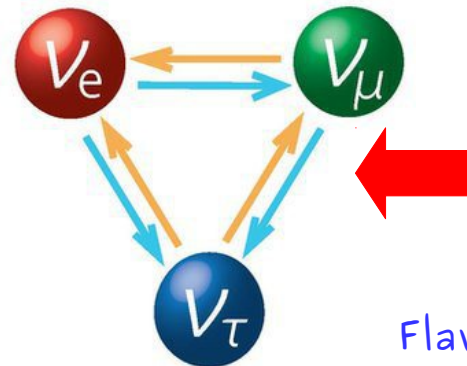
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Oscillation Parameters

$$U = U_{PMNS} = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\theta_{23} & \sin\theta_{23} \\ 0 & -\sin\theta_{23} & \cos\theta_{23} \end{pmatrix}}_{\text{atms \& acelerador}} \times \underbrace{\begin{pmatrix} \cos\theta_{13} & 0 & \sin\theta_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -\sin\theta_{13}e^{i\delta} & 0 & \cos\theta_{13} \end{pmatrix}}_{\text{atms \& reactor \& acelerador}} \times \underbrace{\begin{pmatrix} \cos\theta_{12} & \sin\theta_{12} & 0 \\ -\sin\theta_{12} & \cos\theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{solar \& reactor}}$$

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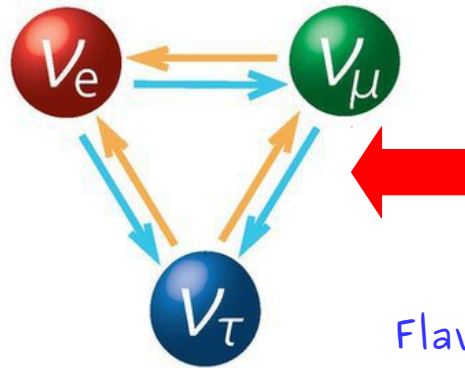
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- Mass Hierarchies
- Dirac Phase
- P. Degeneracy

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Mass Hierarchies

Dirac Phase

P. Degeneracy



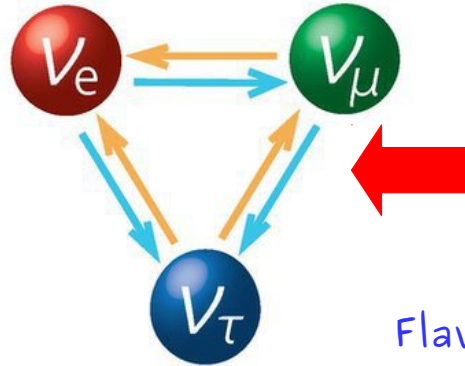
Neutrino Decay

Non Standard Interaction

Quantum Decoherence

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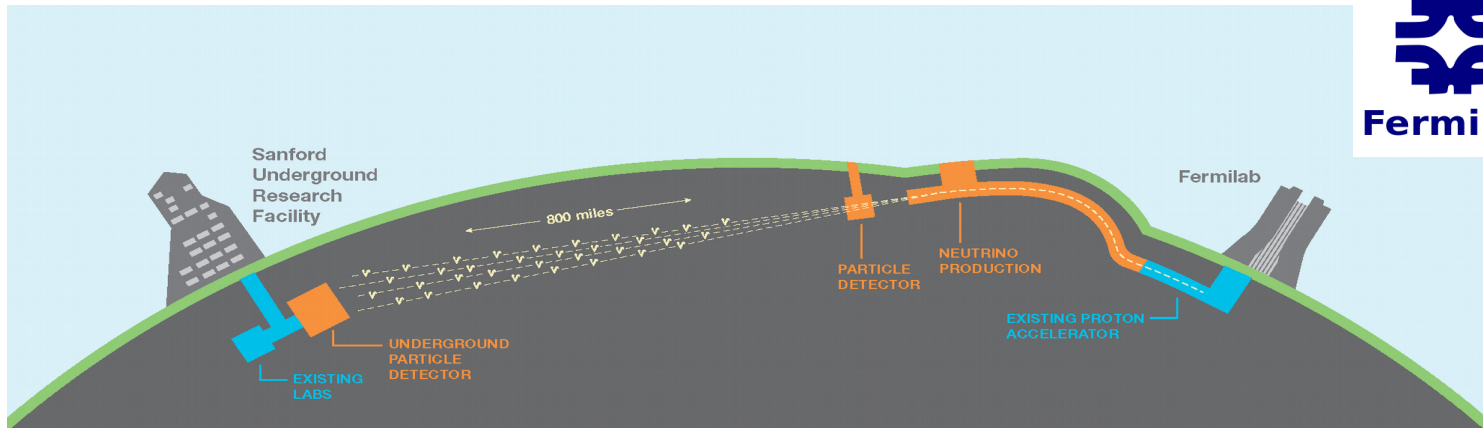


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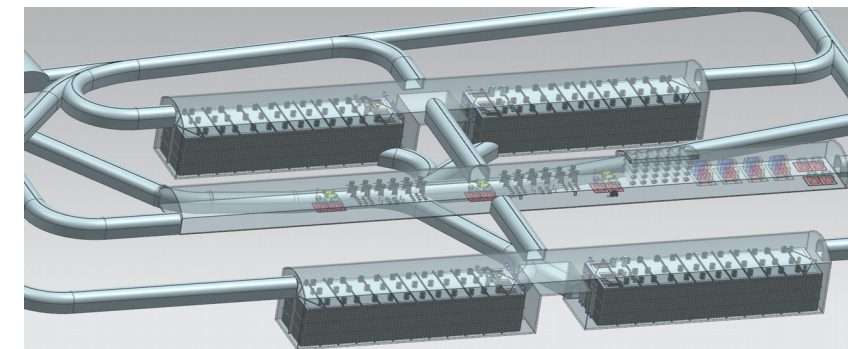
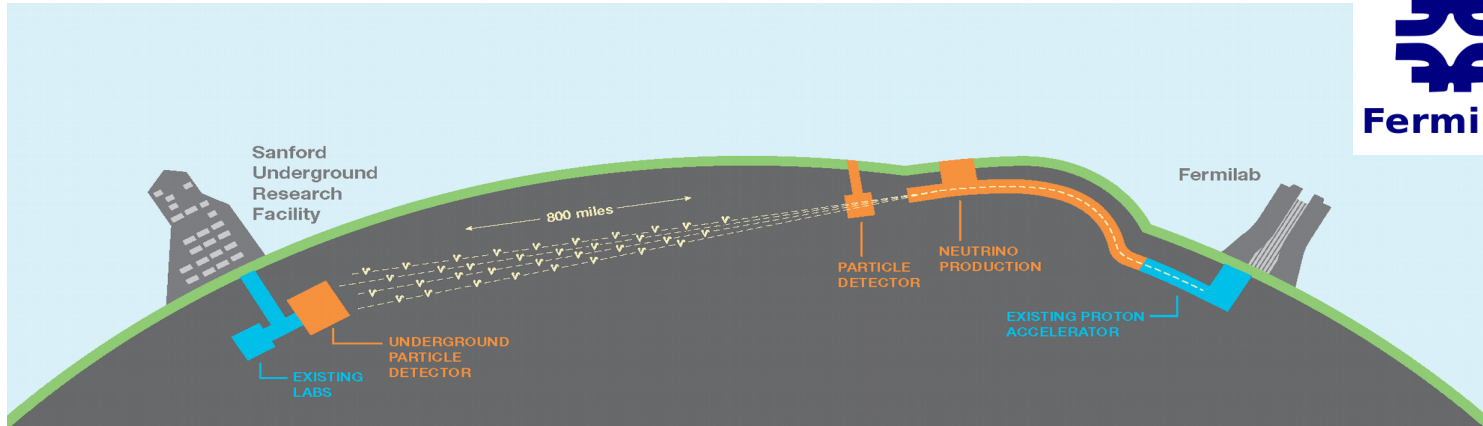
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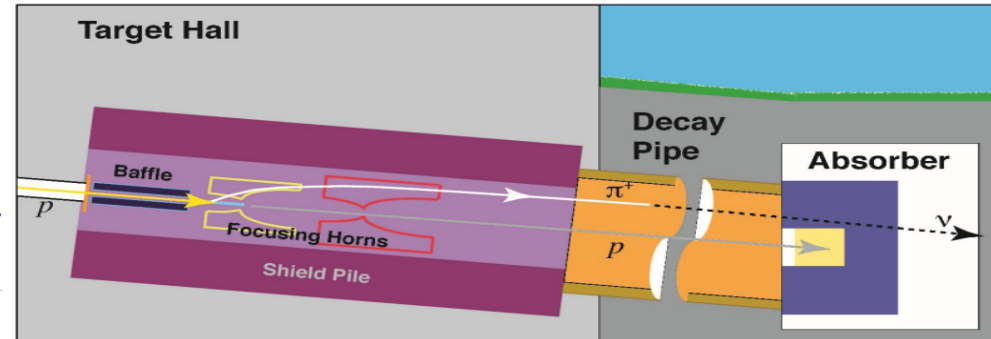


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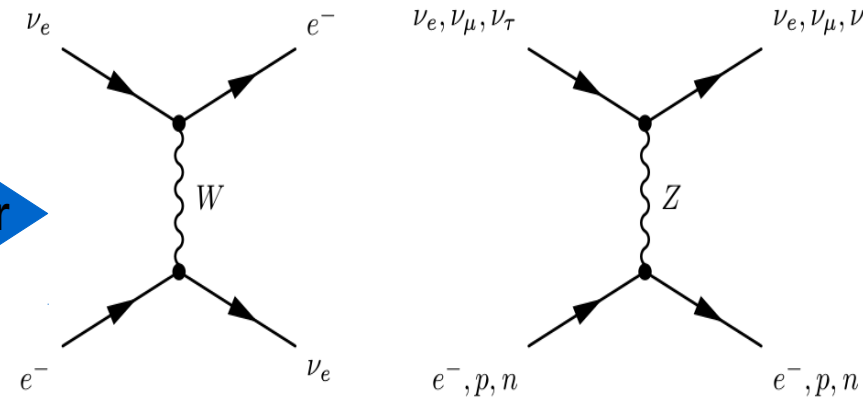


Liquid Argon

Beam



Interaction with matter



Ancillary files

Alion, T. et al - "Experiment Simulation Configurations Used in DUNE CDR"

Introduction

On the other hand...

Introduction

On the other hand...

PHYSICAL REVIEW

VOLUME 104, NUMBER 1

OCTOBER 1, 1956

Question of Parity Conservation in Weak Interactions*

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AND

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Parity?

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→ CPT

CPT symmetry and neutrino oscillation

$$P_{\nu_\alpha \rightarrow \nu_\beta} = P_{\bar{\nu}_\beta \rightarrow \bar{\nu}_\alpha}$$

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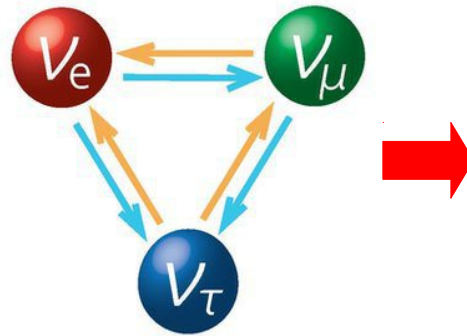
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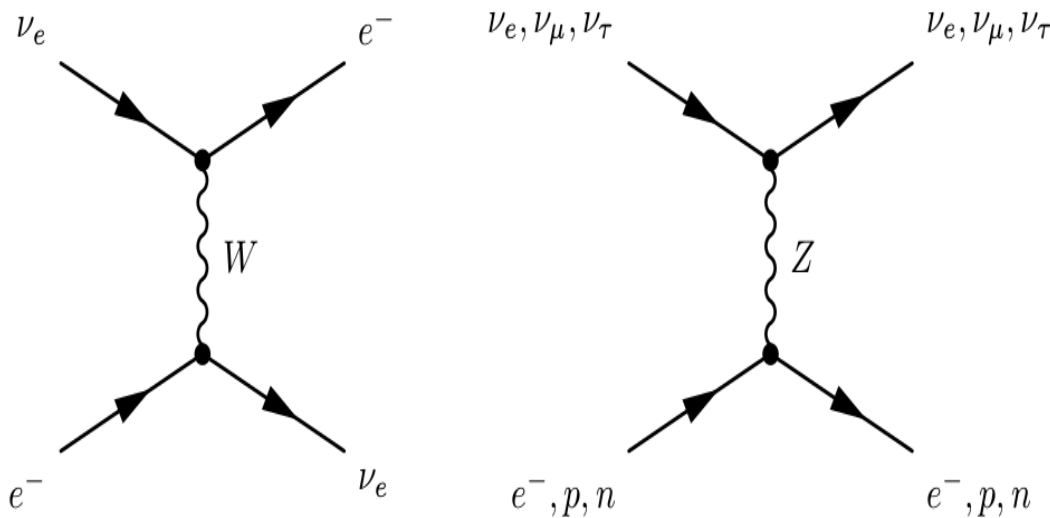
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- Neutrino Decay
- Non Standard Interaction
- Quantum Decoherence**

CPTV due to matter



$$P_{\nu_\alpha \rightarrow \nu_\beta} \neq P_{\bar{\nu}_\beta \rightarrow \bar{\nu}_\alpha} \quad \text{Matter interaction}$$

some scenarios

~~CPT~~

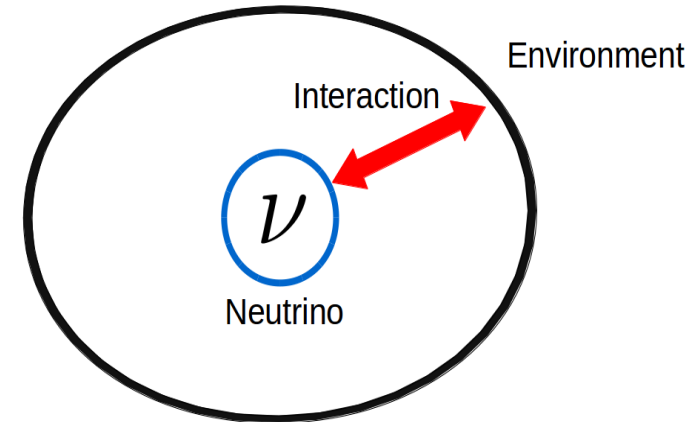
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Measurement



Theoretical Approach

Considering the neutrino like an open quantum system.



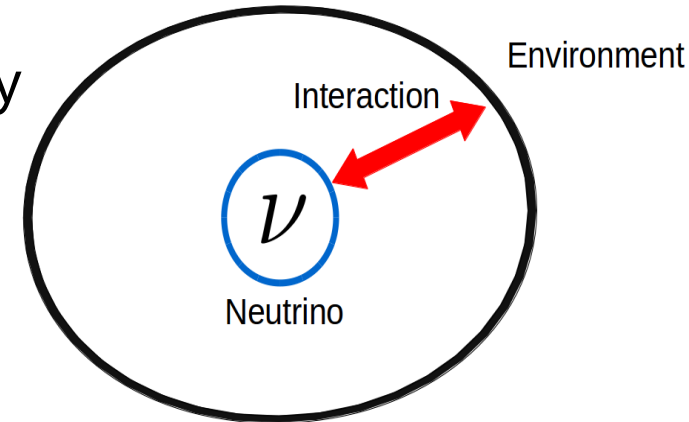
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The time evolution of our quantum system is given by

$$\frac{\partial \hat{\rho}(t)}{\partial t} = -i[\hat{H}, \hat{\rho}(t)] + \mathcal{D}[\hat{\rho}(t)]$$

Dissipative term



A. Gago et al arXiv:hep-ph/0208166

J.A. Carpio et al Phys. Rev. D 97, 115017

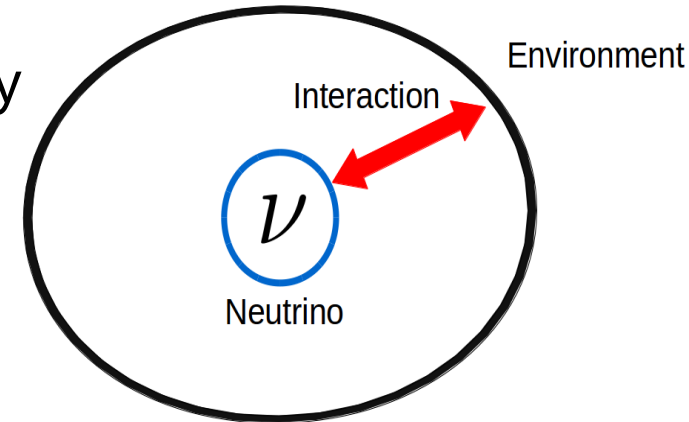
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Dissipative term



For 3 generations, the probability is given by

$$P_{\nu_\alpha \rightarrow \nu_\beta} = \frac{1}{3} + \frac{1}{2} \left(\sum_{i,j} \rho_i^\beta \rho_j^\alpha [e^{Mt}]_{ij} \right)$$

Where $M = H + D$

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Theoretical Approach

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Theoretical Approach

$$P_{\nu_\alpha \rightarrow \nu_\beta} = \frac{1}{3} + \frac{1}{2} \left(\sum_{i,j} \rho_i^\beta \rho_j^\alpha [e^{Mt}]_{ij} \right)$$

Where $M = H + D$

$$\rho_0^\alpha = \sqrt{2/3}$$

$$\rho_1^\alpha = 2 \operatorname{Re} (U_{\alpha 1}^* U_{\alpha 2})$$

$$\rho_2^\alpha = -2 \operatorname{Im} (U_{\alpha 1}^* U_{\alpha 2})$$

$$\rho_3^\alpha = |U_{\alpha 1}|^2 - |U_{\alpha 2}|^2$$

$$\rho_4^\alpha = 2 \operatorname{Re} (U_{\alpha 1}^* U_{\alpha 3})$$

$$\rho_5^\alpha = -2 \operatorname{Im} (U_{\alpha 1}^* U_{\alpha 3})$$

$$\rho_6^\alpha = 2 \operatorname{Re} (U_{\alpha 1}^* U_{\alpha 3})$$

$$\rho_7^\alpha = -2 \operatorname{Im} (U_{\alpha 2}^* U_{\alpha 3})$$

$$\rho_8^\alpha = \frac{1}{\sqrt{3}} (|U_{\alpha 1}|^2 + |U_{\alpha 2}|^2 - 2|U_{\alpha 3}|^2)$$

A. Gago et al arXiv:hep-ph/0208166

In a CPTV case one coefficient of the factor $\rho_i \rho_j$ must be ρ_2 or ρ_5 or ρ_7 and the other should be anyone of the others $\rho_1, \rho_3, \rho_4, \rho_6, \rho_8$ having in total fifteen cases.

Theoretical Approach

- **CPT Asymmetry**

We use the difference of survival probability to refer the CPT violation

$$\Delta P_{\text{CPT}} = P_{\nu_{\alpha} \rightarrow \nu_{\alpha}} - P_{\bar{\nu}_{\alpha} \rightarrow \bar{\nu}_{\alpha}}$$

Theoretical Approach

• CPT Asymmetry

We use the difference of survival probability to refer the CPT violation

$$\Delta P_{CPT} = P_{\nu_\alpha \rightarrow \nu_\alpha} - P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\alpha}$$

Diagonal elements equal to Γ

Same for neutrinos and antineutrinos

One β_{ij} at a time.



$$\mathbf{D} = \begin{pmatrix} \Gamma & \beta_{12} & 0 & 0 & 0 & 0 & 0 & 0 \\ \beta_{12} & \Gamma & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & \Gamma & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \Gamma & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & \Gamma & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \Gamma & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & \Gamma & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \Gamma \end{pmatrix}$$

Theoretical Approach

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$$\Delta P_{\text{CPT}} = \beta_{ij} \frac{(e^{\Omega_{\beta_{ij}} t} - e^{-\Omega_{\beta_{ij}} t})}{\Omega_{\beta_{ij}}} \rho_i^\alpha \rho_j^\alpha e^{-\Gamma t}$$

Where: $\Omega_{\beta_{ij}} = \sqrt{\Delta_{\beta_{ij}}^2 - \beta_{ij}^2}$ $\Delta_{\beta_{ij}} = \frac{\Delta m_{ij}^2}{2p}, \quad i, j = 1, 2, 3$

Theoretical Approach

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We use the difference of survival probability to refer the CPT violation

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One β_{ij} at a time.

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$$\Delta P_{\text{CPT}} = \beta_{ij} \frac{(e^{\Omega_{\beta_{ij}} t} - e^{-\Omega_{\beta_{ij}} t})}{\Omega_{\beta_{ij}}} \rho_i^\alpha \rho_j^\alpha e^{-\Gamma t} \rightarrow \text{Damping}$$

Where: $\Omega_{\beta_{ij}} = \sqrt{\Delta_{\beta_{ij}}^2 - \beta_{ij}^2}$ $\Delta_{\beta_{ij}} = \frac{\Delta m_{ij}^2}{2p}$, $i, j = 1, 2, 3$

Theoretical Approach

Considering the muon disappearance channel and the DUNE baseline

$$\Delta P_{\cancel{CPT}} = P_{\nu_{\mu} \rightarrow \nu_{\mu}} - P_{\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{\mu}}$$

Theoretical Approach

Considering the muon disappearance channel and the DUNE baseline

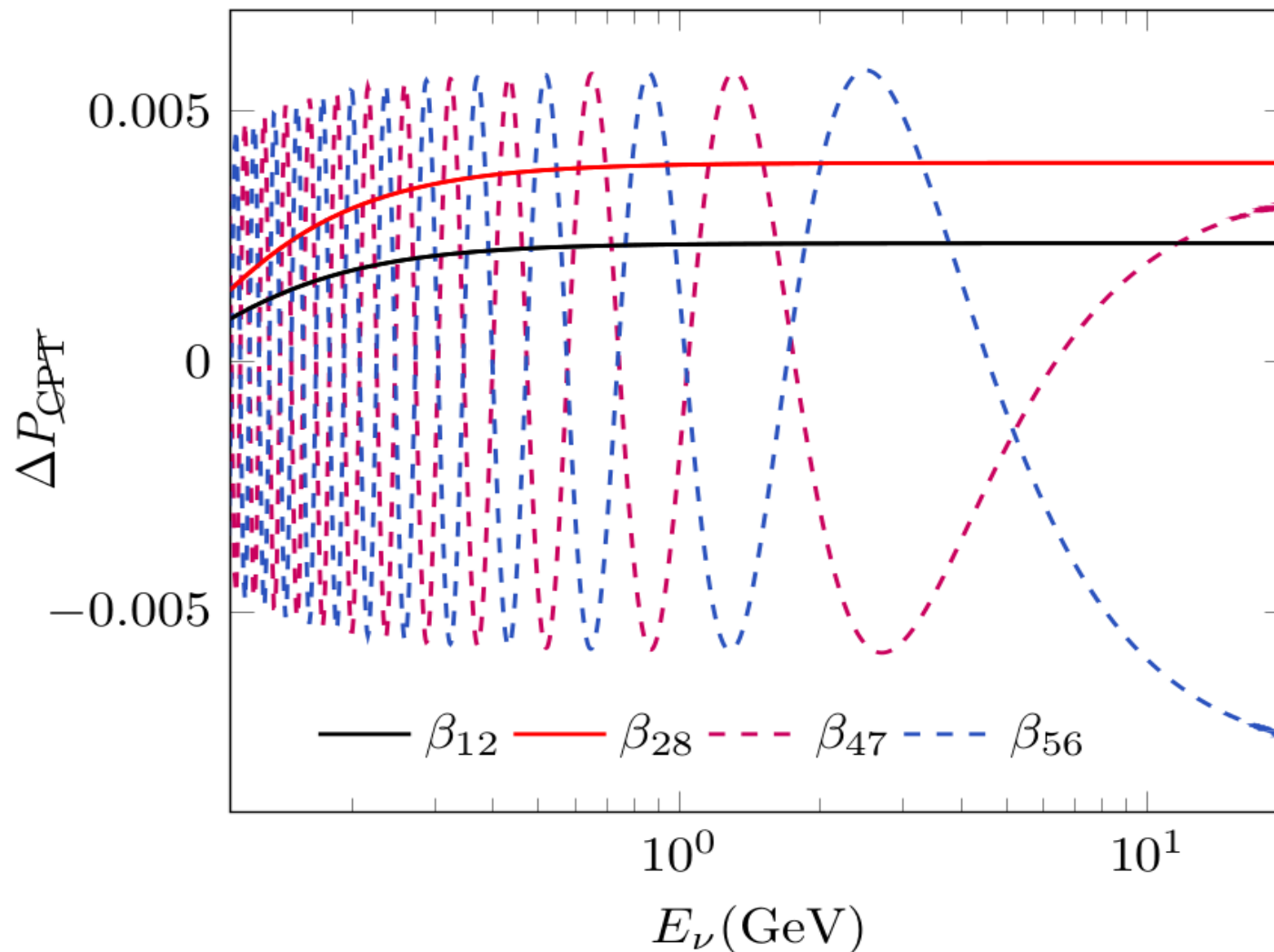
$$\Delta P_{\text{CPT}} = P_{\nu_{\mu} \rightarrow \nu_{\mu}} - P_{\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{\mu}}$$

$$\Gamma = 10^{-23} \text{ GeV}$$

$$\delta_{\text{CP}} = 3\pi/2$$

$$\beta_{ij} = \Gamma/\sqrt{3}$$

$$L = 1300 \text{ Km}$$



$$\Delta P_{\text{CPT}} = P_{\nu_\mu \rightarrow \nu_\mu} - P_{\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu}$$

$$L = 1300 \text{ Km}$$

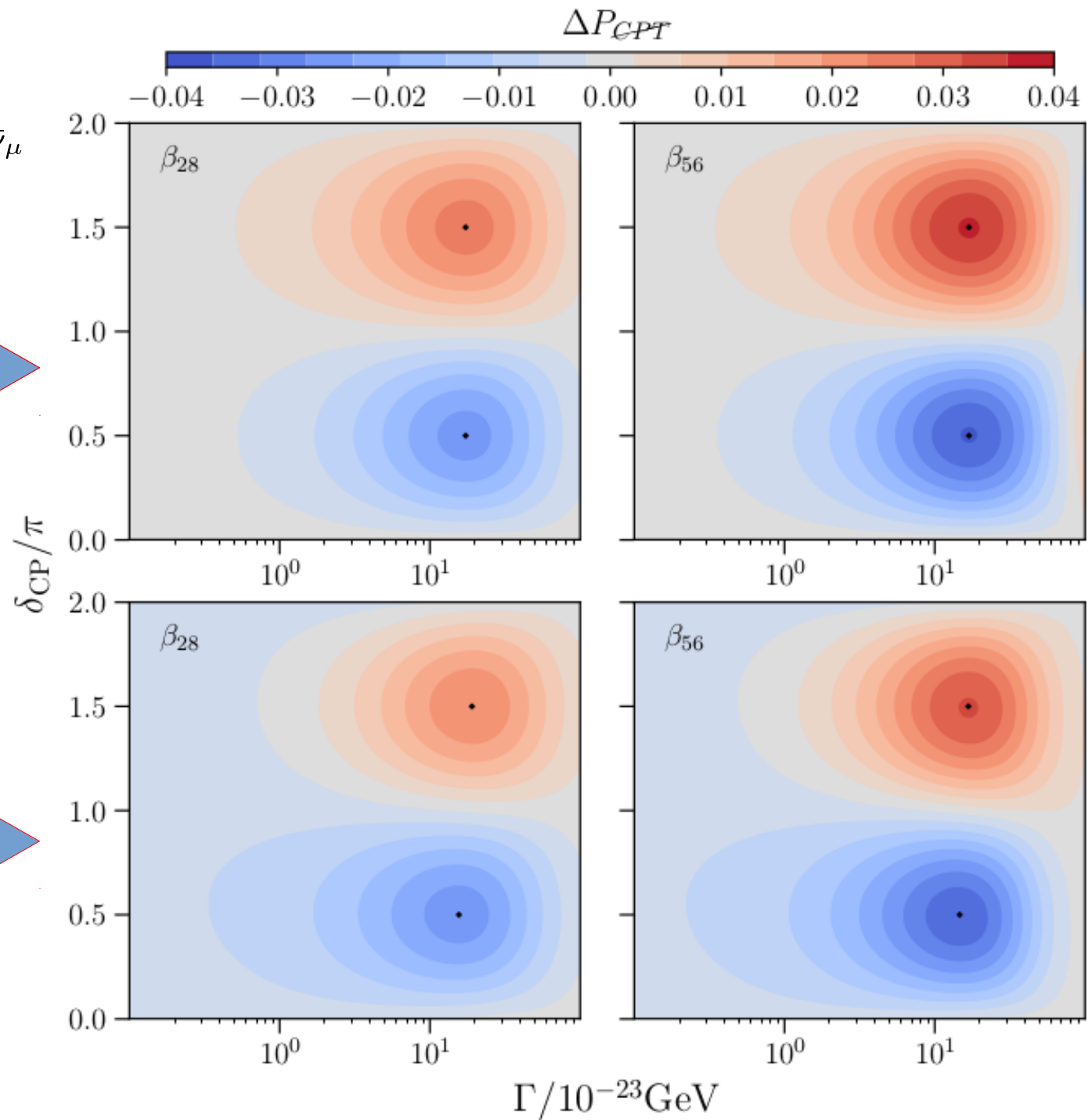
Vacuum

$$\beta_{28} = \Gamma / \sqrt{3}$$

$$\beta_{56} = \Gamma / \sqrt{3}$$

$$E_\nu = 2.4 \text{ GeV}$$

Matter



Theoretical Approach

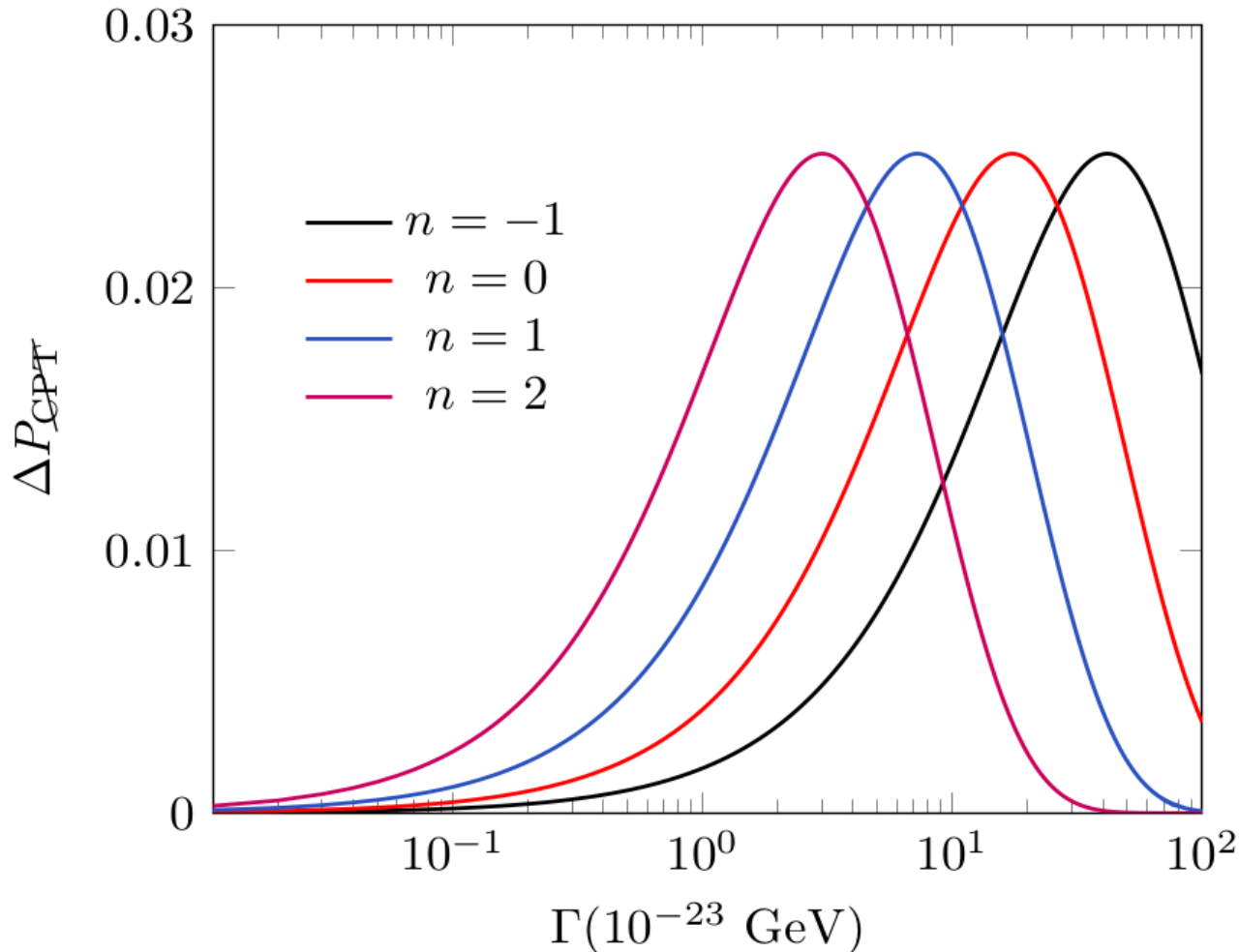
- Energy Dependence

$$\Gamma_{E_\nu} = \Gamma \left(\frac{E}{\text{GeV}} \right)^n$$

Theoretical Approach

- Energy Dependence

$$\Gamma_{E\nu} = \Gamma \left(\frac{E}{\text{GeV}} \right)^n$$

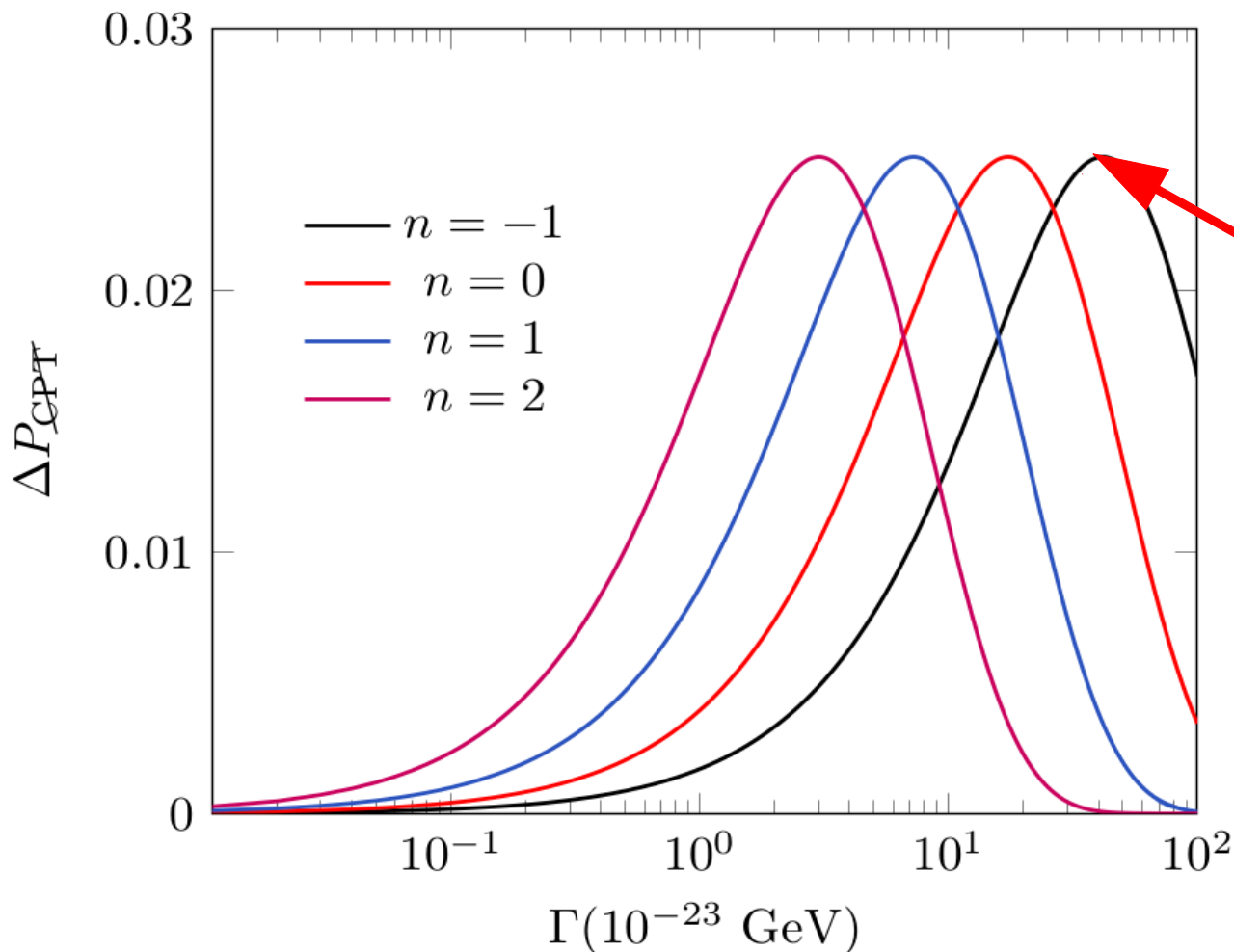


$$E = 2.4 \text{ GeV}, \beta_{28} = \Gamma / \sqrt{3}$$
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Theoretical Approach

- Energy Dependence

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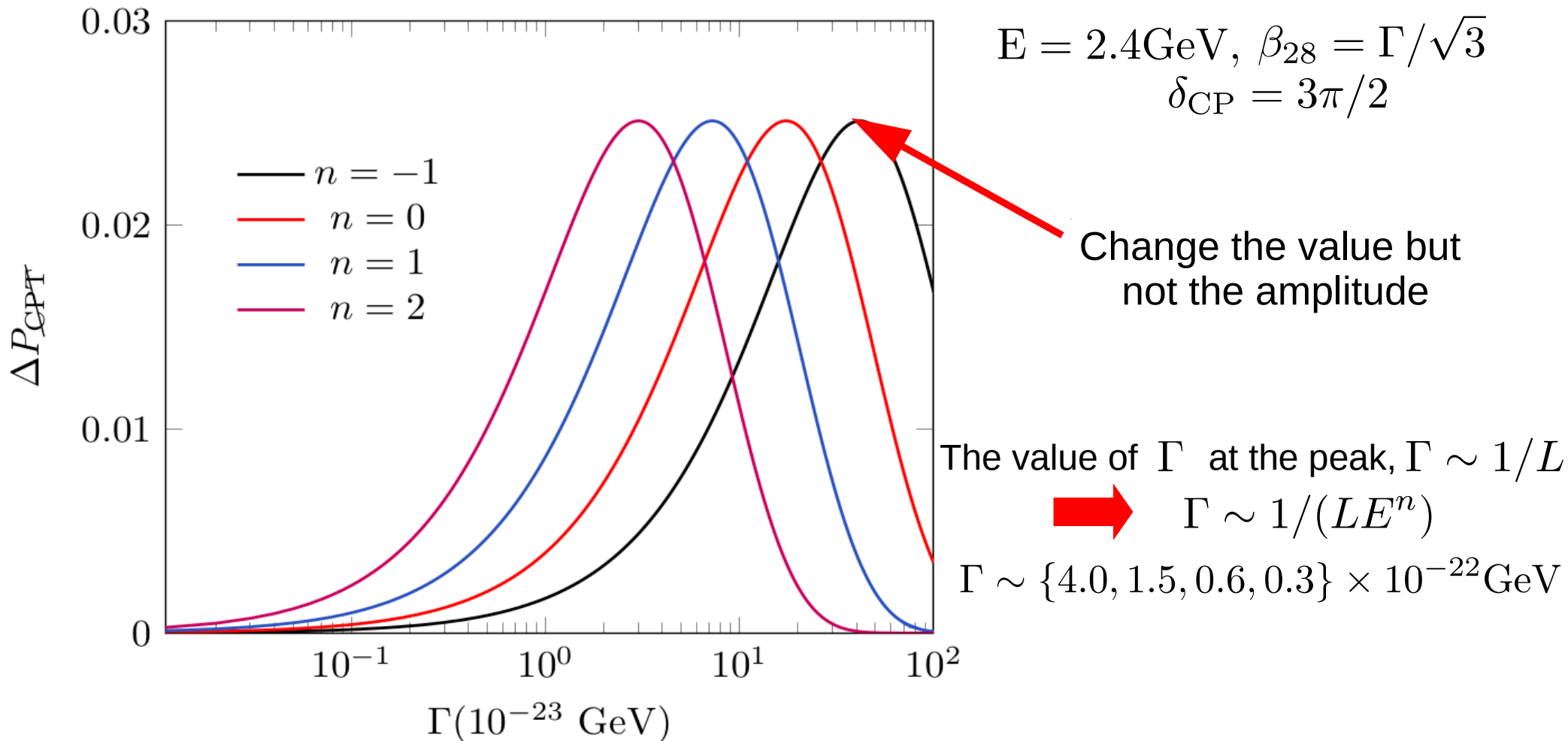
$$E = 2.4\text{GeV}, \beta_{28} = \Gamma/\sqrt{3}$$
$$\delta_{\text{CP}} = 3\pi/2$$

Change the value but not the amplitude

Theoretical Approach

- Energy Dependence

$$\Gamma_{E\nu} = \Gamma \left(\frac{E}{\text{GeV}} \right)^n$$



Theoretical Approach

- Optimal case

$$\beta_{28} = \Gamma/\sqrt{3}$$

$$\beta_{12} = (\sqrt{2/3})\Gamma/3$$

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Theoretical Approach

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$$\beta_{28} = \Gamma/\sqrt{3} \quad \beta_{12} = (\sqrt{2/3})\Gamma/3 \quad \beta_{56} = -\beta_{47} = \Gamma/3$$

- CPT Asymmetry in matter

CPTV in Standard Oscillation (SO)

$$P_{\nu_\alpha \rightarrow \nu_\beta}^{\text{SO}} \neq P_{\bar{\nu}_\beta \rightarrow \bar{\nu}_\alpha}^{\text{SO}} \quad \text{Matter interaction}$$

$$\rightarrow \Delta P_{\text{CPT}}^{\text{SO}} \neq 0$$

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Standard Oscillation and Diagonal
Decoherence Matrix

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Standard Oscillation

$$\Delta P_{\text{CPT}}^{\text{SO} \oplus \text{DDM}} = e^{-\Gamma t} \Delta P_{\text{CPT}}^{\text{SO}}$$

Simulation Details

In order to show tangible results, we define the observable of CPT asymmetry depending of the number of events of neutrinos and antineutrinos.

$$\Delta N = N(\nu_{\mu}) - N(\bar{\nu}_{\mu})$$

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$$\Delta N = N(\nu_{\mu}) - N(\bar{\nu}_{\mu})$$

To study and differentiate the CPTV due to the effect of quantum decoherence from the CPTV due to the matter effect, we define the ratio

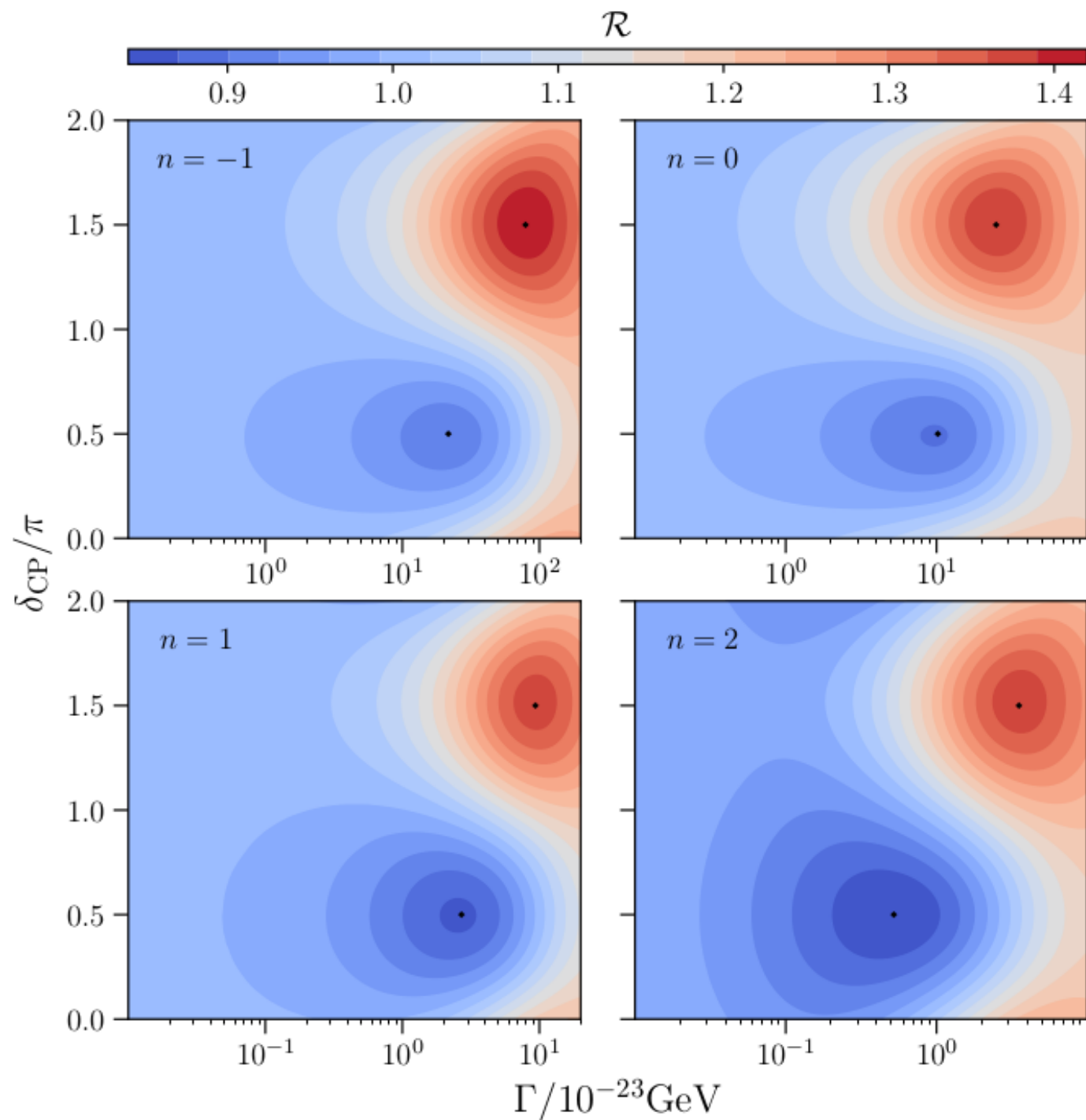
$$\mathcal{R} = \frac{\Delta N^{\text{SO} \oplus \text{DEC}}}{\Delta N^{\text{SO}}}$$

The uncertainty for the event rate are considered as \sqrt{N} .

Results

Our observable \mathcal{R}
and its dependence on
 δ_{CP} and Γ

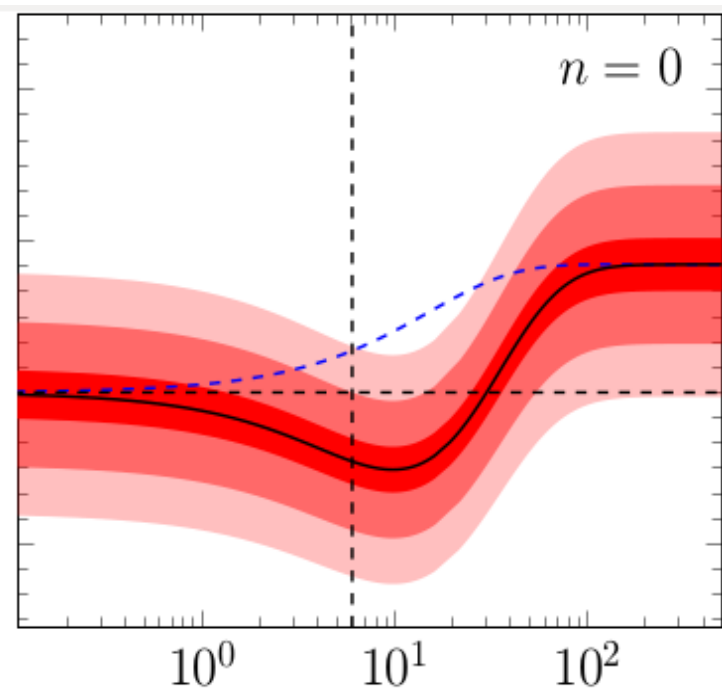
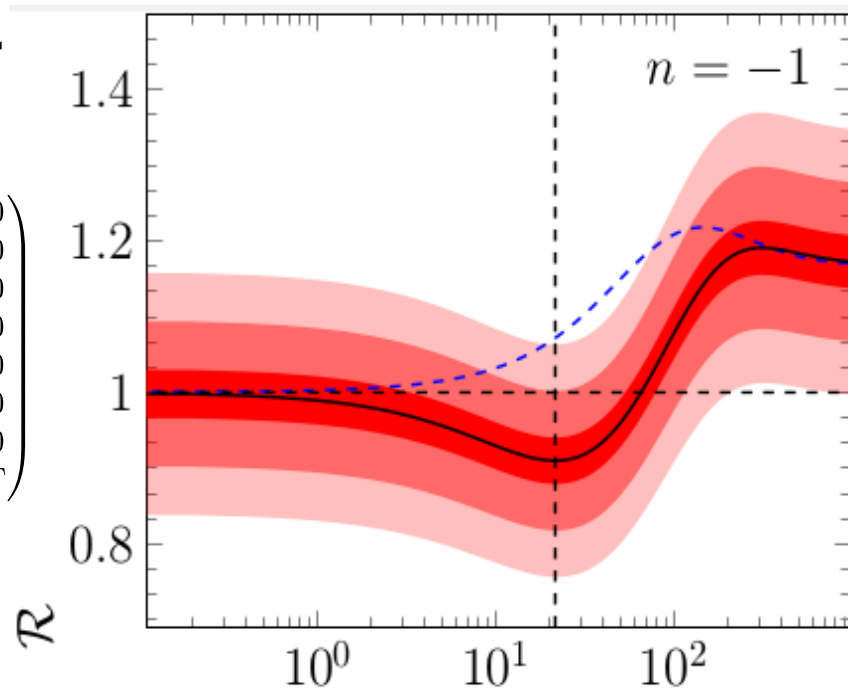
Optimal Case
Energy range
0.5 – 20 GeV
5 years FHC
5 years RHC
Normal hierarchy



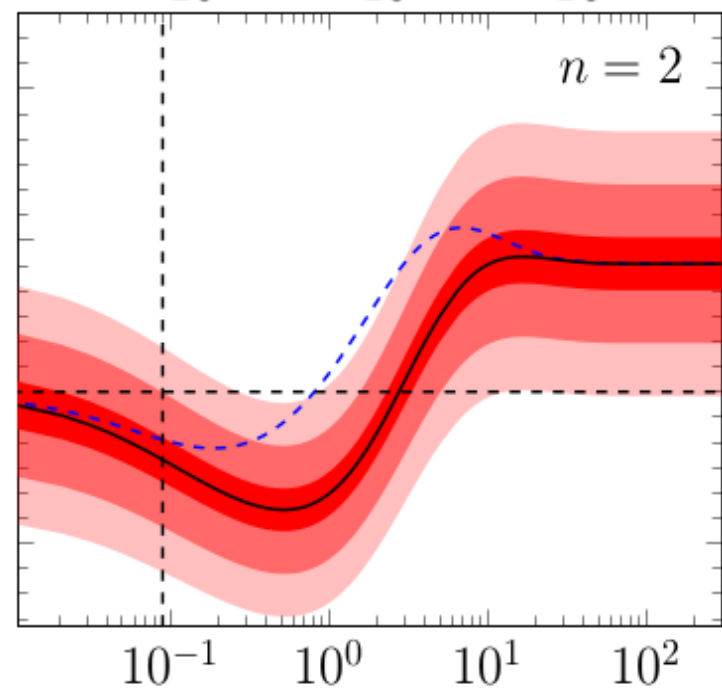
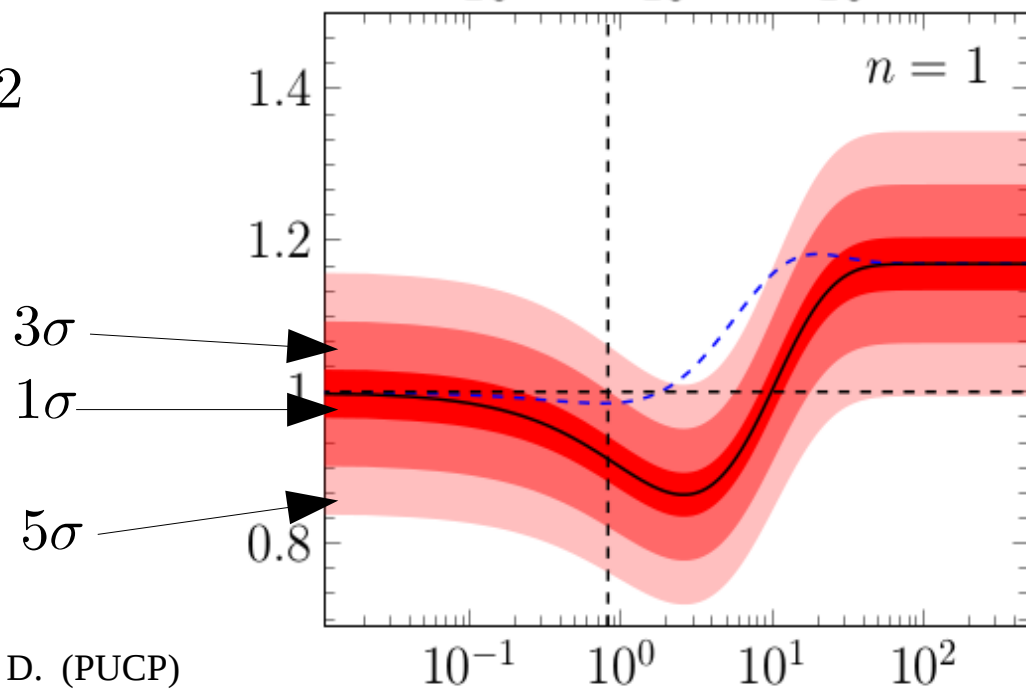
Results

Optimal case ———
 DDM - - - - -

$$D = \begin{pmatrix} \Gamma & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \Gamma & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & \Gamma & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \Gamma & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & \Gamma & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \Gamma & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & \Gamma & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \Gamma \end{pmatrix}$$



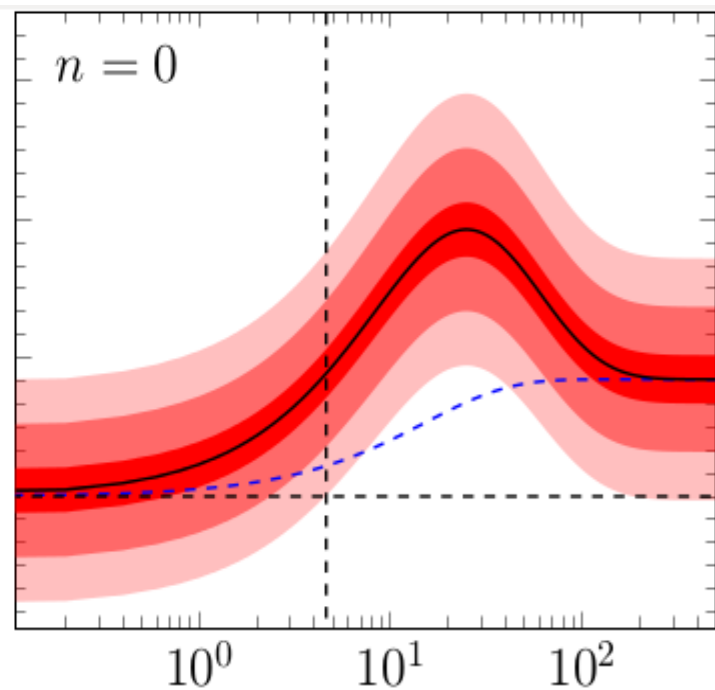
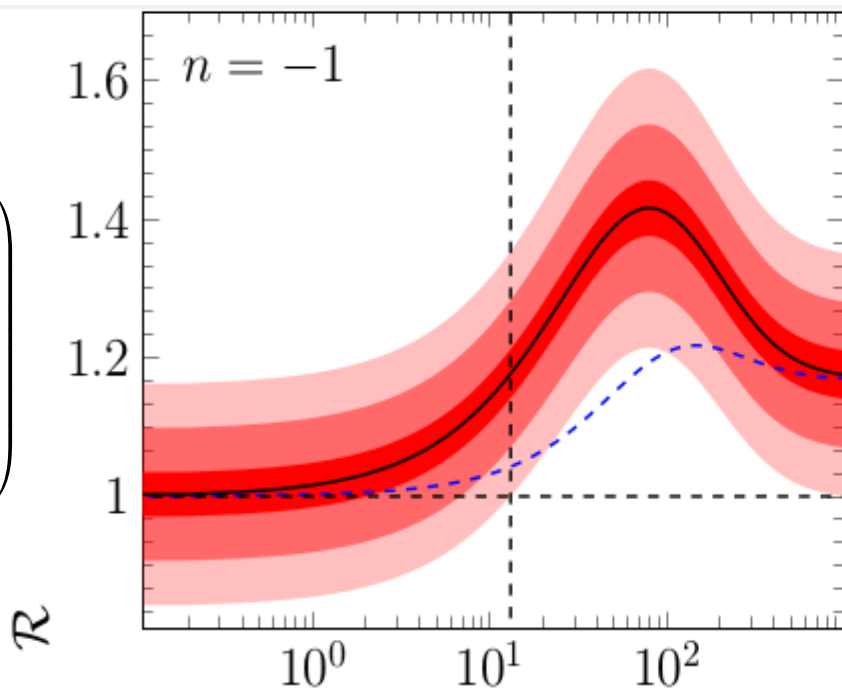
$\delta_{CP} = \pi/2$



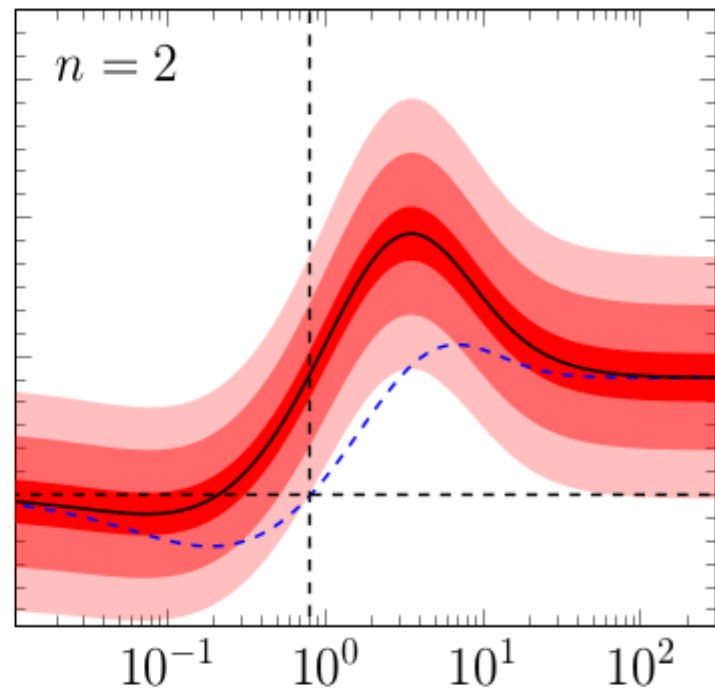
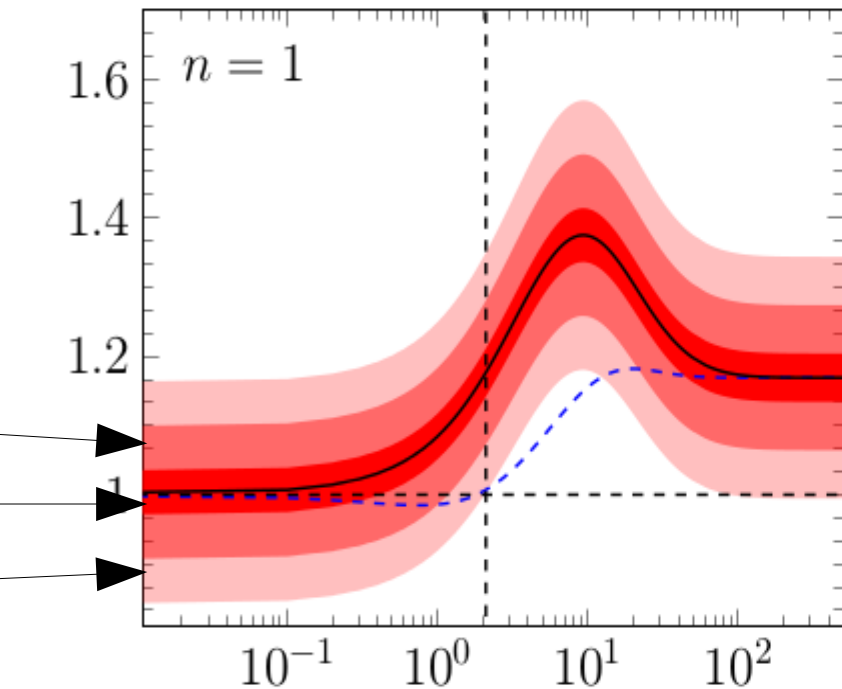
Results

Optimal case ———
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$\delta_{CP} = 3\pi/2$



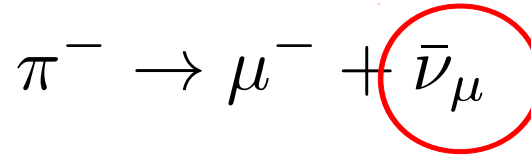
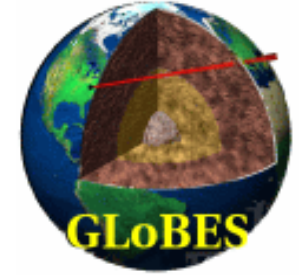
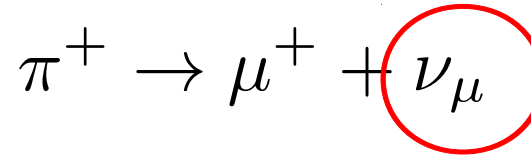
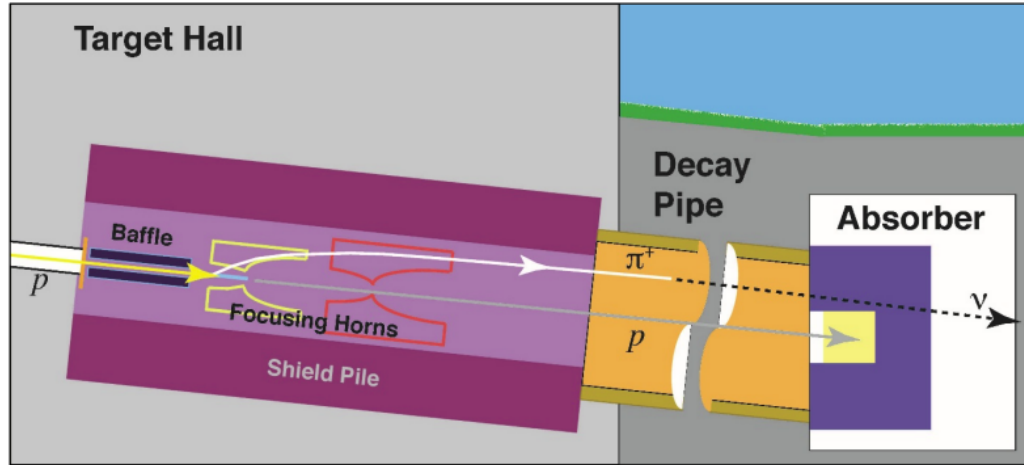
Conclusions

- We have shown that a breakdown of the CPT symmetry can take place when the neutrino system is affected by the environment. we have quantified a possible measurement of this CPTV using the disappearance channels $\nu_\mu \rightarrow \nu_\mu$ and $\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$, with their corresponding backgrounds, and an observable \mathcal{R} . All in the context of the DUNE experiment.
- The simulated measurements of \mathcal{R} have been performed considering four hypothesis of energy dependence on the decoherence parameters. For $\delta_{CP} = 3\pi/2$ and a NDM, we achieve a 5σ for \mathcal{R} with respect to its expectation value at the SO case, $\mathcal{R} = 1$, For $\delta_{CP} = \pi/2$, we reach discrepancies of the order of 3σ in our best case.
- The observations of CPTV appear when the neutrino system is treated as an open system. The latter means that it is likely that if we had access to the information of the environment, the overall CPT symmetry is conserved. For this reason, it deserves a more profound discussion to ascertain if this CPTV is a breaking at the fundamental level or it is only an apparent one, because of our lack of information from the environment.

**THANK YOU VERY MUCH FOR
YOUR ATTENTION**

BACK UP

Simulation Details



nuSQuIDS

Probability

Efficiency

Fluxes

Cross Section

Smearing



Event Rate N



5 years neutrino mode (FHC)

5 years antineutrino mode (RHC)

Optimal case

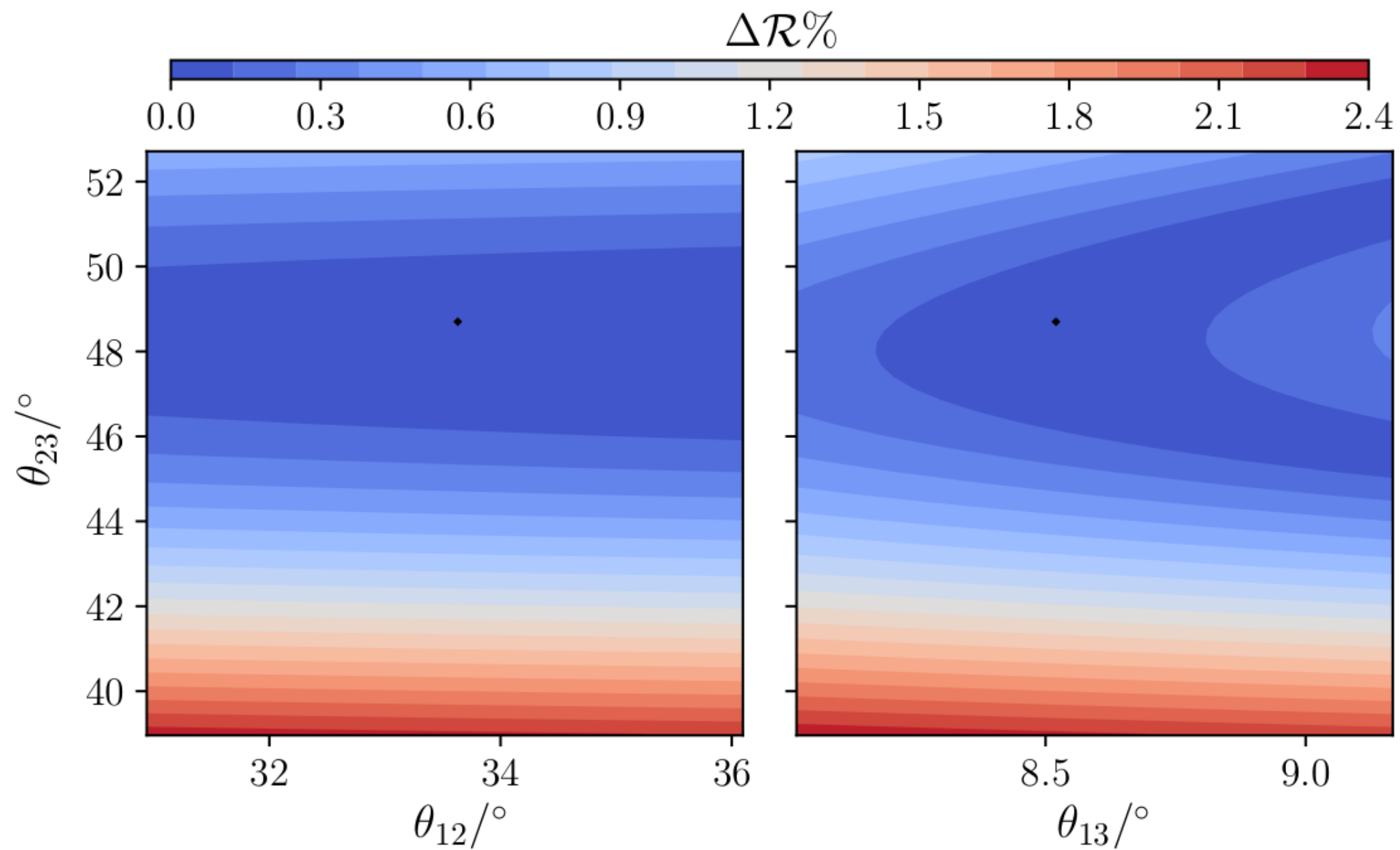
Energy range 0.5 – 20 GeV

$\delta_{CP} = 3\pi/2$

Ancillary files

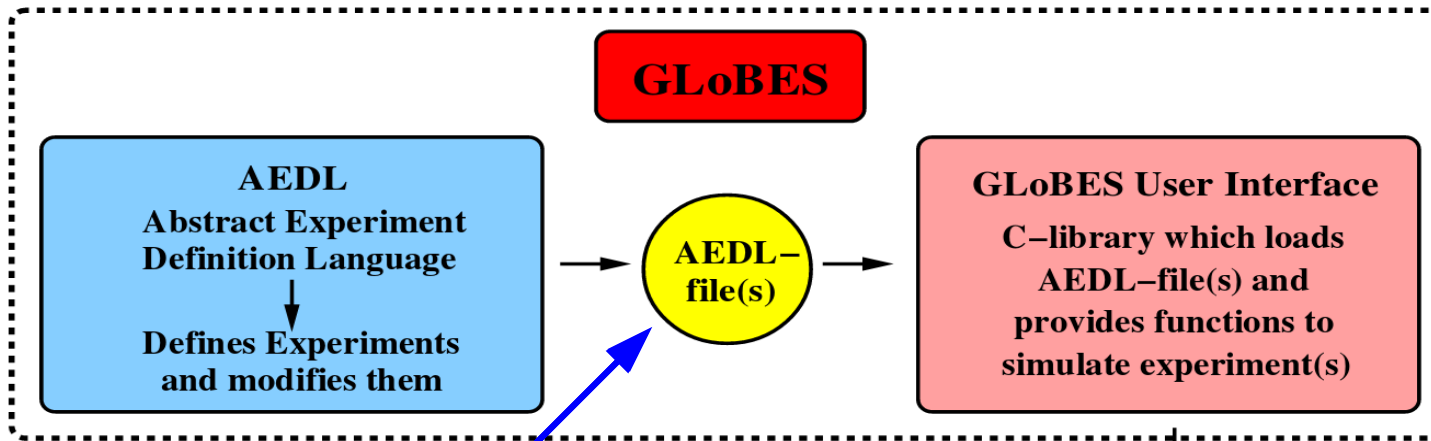
Alion, T. et al - arXiv:1606.09550

| $\Gamma = 10^{-23}$ GeV | Std | $n = -1$ | $n = 0$ | $n = 1$ | $n = 2$ |
|---|-------|----------|---------|---------|---------|
| Neutrino mode | | | | | |
| $\nu_\mu + \bar{\nu}_\mu$ CC signal | 11749 | 11841 | 11965 | 11573 | 11932 |
| NC background | 109 | 109 | 109 | 109 | 109 |
| $\nu_\tau + \bar{\nu}_\tau$ CC background | 43 | 43 | 46 | 74 | 87 |
| Antineutrino mode | | | | | |
| $\bar{\nu}_\mu + \nu_\mu$ CC signal | 5903 | 5897 | 5846 | 5237 | 4816 |
| NC background | 58 | 58 | 58 | 58 | 58 |
| $\nu_\tau + \bar{\nu}_\tau$ CC background | 27 | 27 | 29 | 50 | 60 |

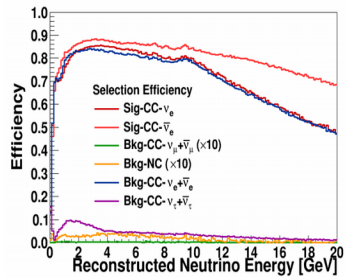


General Long Baseline Experiment Simulator

GLOBES

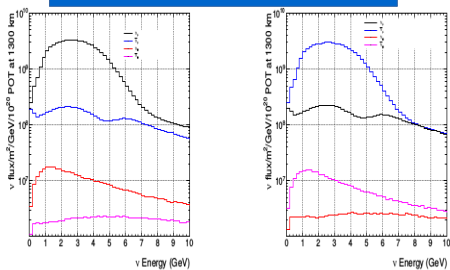


Efficiency

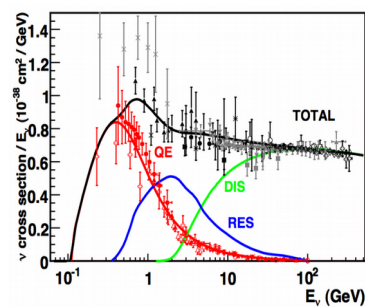


Inputs

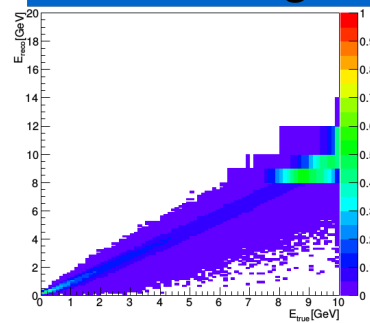
Flux



Cross Section

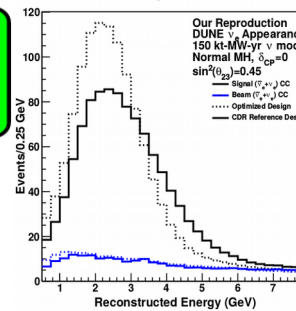


Smearing

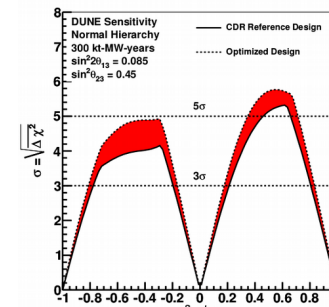


Application software to compute high-level sensitivities, precision etc.

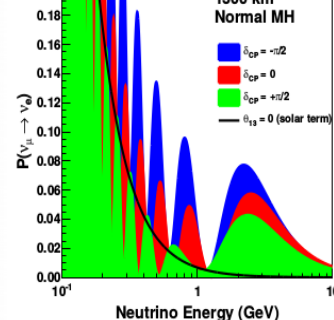
DUNE νe Appearance



CP Violation Sensitivity



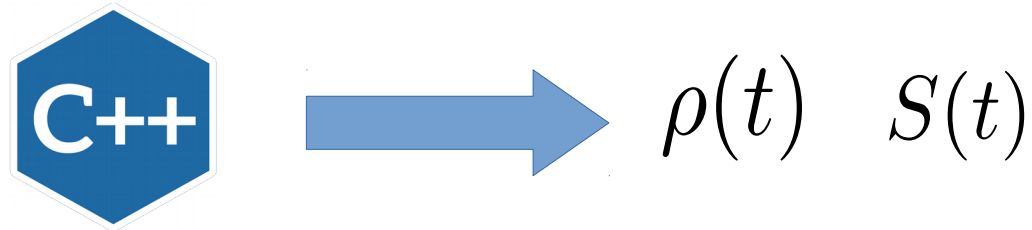
1300 km Normal MH



nuSQuIDS



A Simple Quantum Integro-Differential Solver: SQuIDS



How it works?, Expresses all operators on the basis $SU(N)$

It allows us to add new terms to the right of the Lindblad equation.

$$\frac{\partial \rho(t)}{\partial t} = -i[H_{eff}, \rho(t)] + \mathcal{D}[\rho(t)]$$

Matter effect



Decoherence effect