



# Calculation of Delayed-Photoneutron Production in Heavy-Water Reactors using Geant4

Evan Rand, Ph.D. - Canadian Nuclear Laboratories

ENSAR2 Workshop

April 24<sup>th</sup> - 26<sup>th</sup>, 2019

CW-01200-001-000



Canadian Nuclear Laboratories | Laboratoires Nucléaires Canadiens

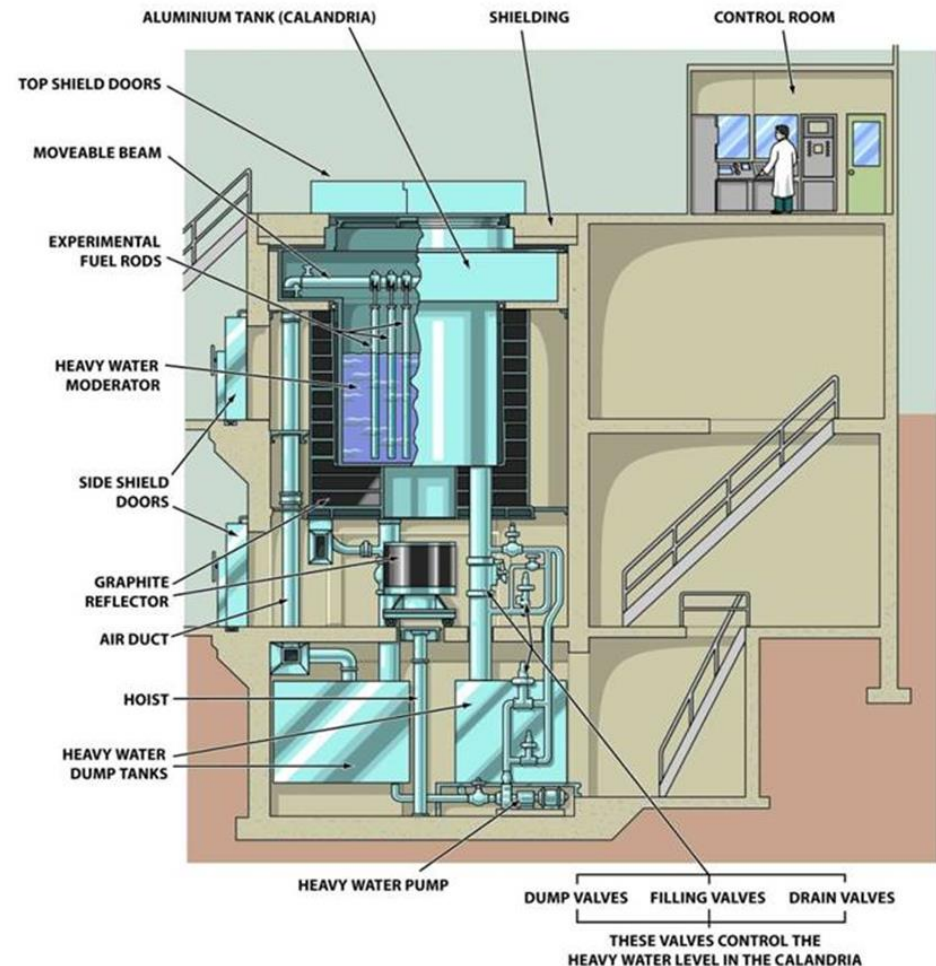
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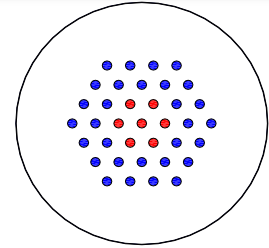
# ZED-2 - Zero Energy Deuterium

58 years and counting

- Successor to **ZEEP** – **Z**ero **E**nergy **E**xperimental **P**ile
- First went critical on September 7<sup>th</sup>, 1960
- Tank type: reactor control via moderator level
- 2521 cores built, 190 of them unique, and counting!



# Quick Facts



**Power:** up to  $\sim 200$  W (thermal)

**Peak Neutron Flux:**  $1 \times 10^9$  n/cm<sup>2</sup> s thermal,  $5 \times 10^8$  n/cm<sup>2</sup> s fast

**Calandria:** 3.36 m in diameter, 3.35 m in height

**Fuel:** Various types and assemblies

**Moderator:** Heavy water (99.8 to 97.5 weight% D<sub>2</sub>O), soluble poison capability, temperatures up to 90°C, and variable core height (criticality achieved by pumping moderator into calandria)

**Fuel Arrangement:** variable pitch square and hexagonal lattices

**Coolant:** Heavy water, light water, air, CO<sub>2</sub>, organics, Pb-Bi, etc. (not active). Temperatures up to 300°C in some channels.

**Flexibility:** We can operate with new fuels/coolants/materials as required



# Fuel

## Natural $\text{UO}_2$ Bundles:

- 7, 19, 28, 37, and 43-element

## Other Natural U flavours:

- Metal, Carbide, Silicide bundles
- ZEEP rods

## Mixed oxides:

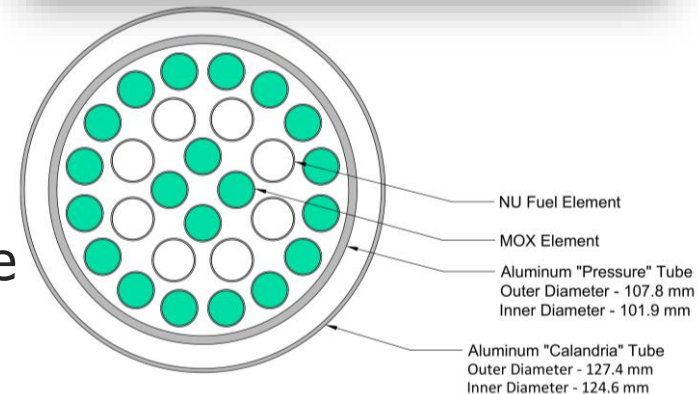
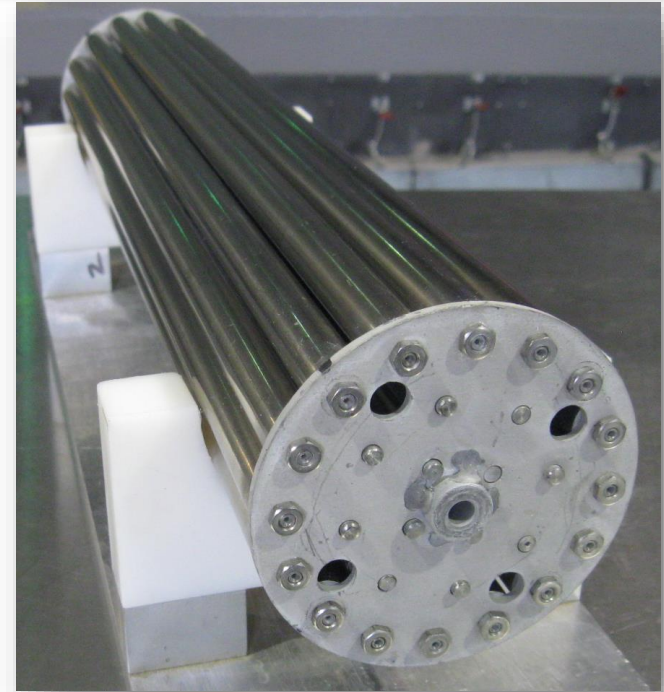
- Pu-U,  $^{233}\text{U}$ -Th, Pu-Th,  $^{235}\text{U}$ -Th

## Bundles with burnable absorber (Low Void Reactivity)

## Enriched or reprocessed $\text{UO}_2$ bundles:

- LEU, RU

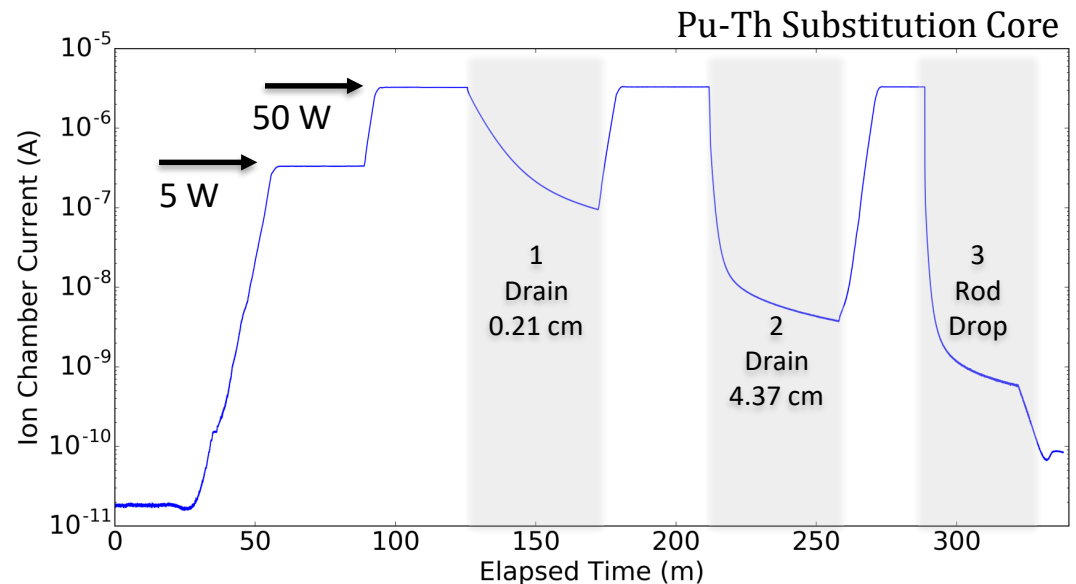
## Assembly geometry: bundles in Pressure Tube/Calandria Tube, clad rods, etc.



# Reactor Transient Experiments in ZED-2

Understanding and predicting the transient behaviour of reactor systems is paramount for reactor safety and operation. Transients occur during startup/shutdown, accident scenarios (coolant loss), etc. The behaviour depends on a number of factors such as fuel type(s) ( $^{nat}U$ , LEU, MOX), reactor configuration, etc.

*The study of these processes which control the time-dependent behavior of a nuclear reactor is known as **Reactor kinetics**.*



# Reactor Kinetics

The energy- and time-dependent diffusion equation:

$$\frac{1}{v} \frac{\partial \phi}{\partial t} = (F_p - M)\phi + S_d(r, E, t) + S(r, E, t)$$

Prompt-neutron production

$$F_p \phi = \chi_p(E) \int_0^\infty dE' v_p(E') \Sigma_f(r, E', t) \phi(r, E', t)$$

Neutron leakage/removal

$$M \phi = -\vec{\nabla} \cdot D(r, E) \vec{\nabla} \phi(r, E, t) + \Sigma(r, E, t) \phi(r, E, t) - \int_0^\infty dE' \Sigma_s(r, E' \rightarrow E, t) \phi(r, E', t)$$

Delayed neutrons

Neutron sources

*Almost impossible to solve in all practical scenarios, even when using today's most powerful computers. We are left with two approaches:*

***Deterministic (Point Kinetics) – An exact solution to an approximate problem.***

***Stochastic (Monte Carlo) – An approximate solution to the exact problem.***

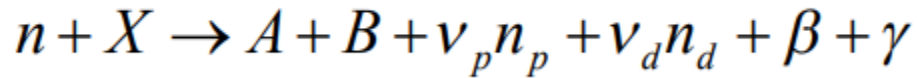


# Point Kinetics

...assumes flux is separable:

$$\phi(r, E, t) = n(t) \psi(r, E, t)$$

Nuclear fission:



Total neutron yield:

$$\nu = \nu_d + \nu_p$$

*Inverse Point Kinetics Solution:*

$$\frac{dn(t)}{dt} = \left( \frac{\rho(t) - \sum_{i=1}^{ngroups} \beta_i}{\Lambda} \right) n(t) + \sum_{i=1}^{ngroups} \lambda_i C_i$$

$$\frac{dC_i(t)}{dt} = -\lambda_i C_i(t) + \frac{\beta_i}{\Lambda} n(t)$$

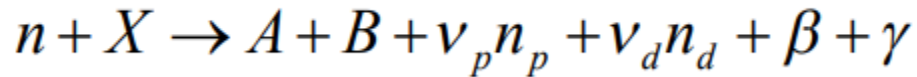


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*Reactivity*



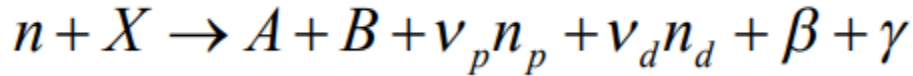


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*Effective multiplication factor*

$$\rho(t) = \frac{k_{eff}(t) - 1}{k_{eff}(t)}$$

*Reactivity*

critical  $k_{eff} = 1$   
 supercritical  $k_{eff} > 1$   
 subcritical  $k_{eff} < 1$

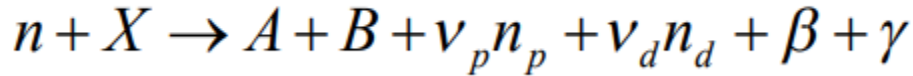


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*Prompt-neutron generation time*



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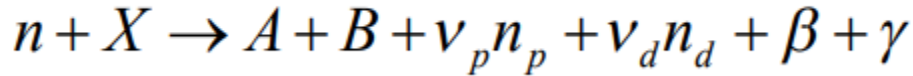


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Precursor concentrations      Effective multiplication factor

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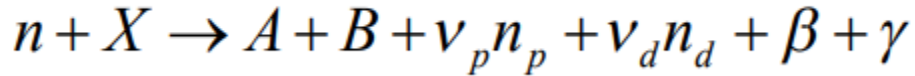


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Delayed-neutron fraction ( $\beta$ ):

$$\beta = \frac{\nu_d}{\nu} \longrightarrow \beta_i = \frac{\nu_{di}}{\nu}$$

*Delayed-neutron groups with differing time constants.*

*Inverse Point Kinetics Solution:*

$$\frac{dn(t)}{dt} = \left( \frac{\rho(t) - \sum_{i=1}^{ngroups} \beta_i}{\Lambda} \right) n(t) + \sum_{i=1}^{ngroups} \lambda_i C_i$$

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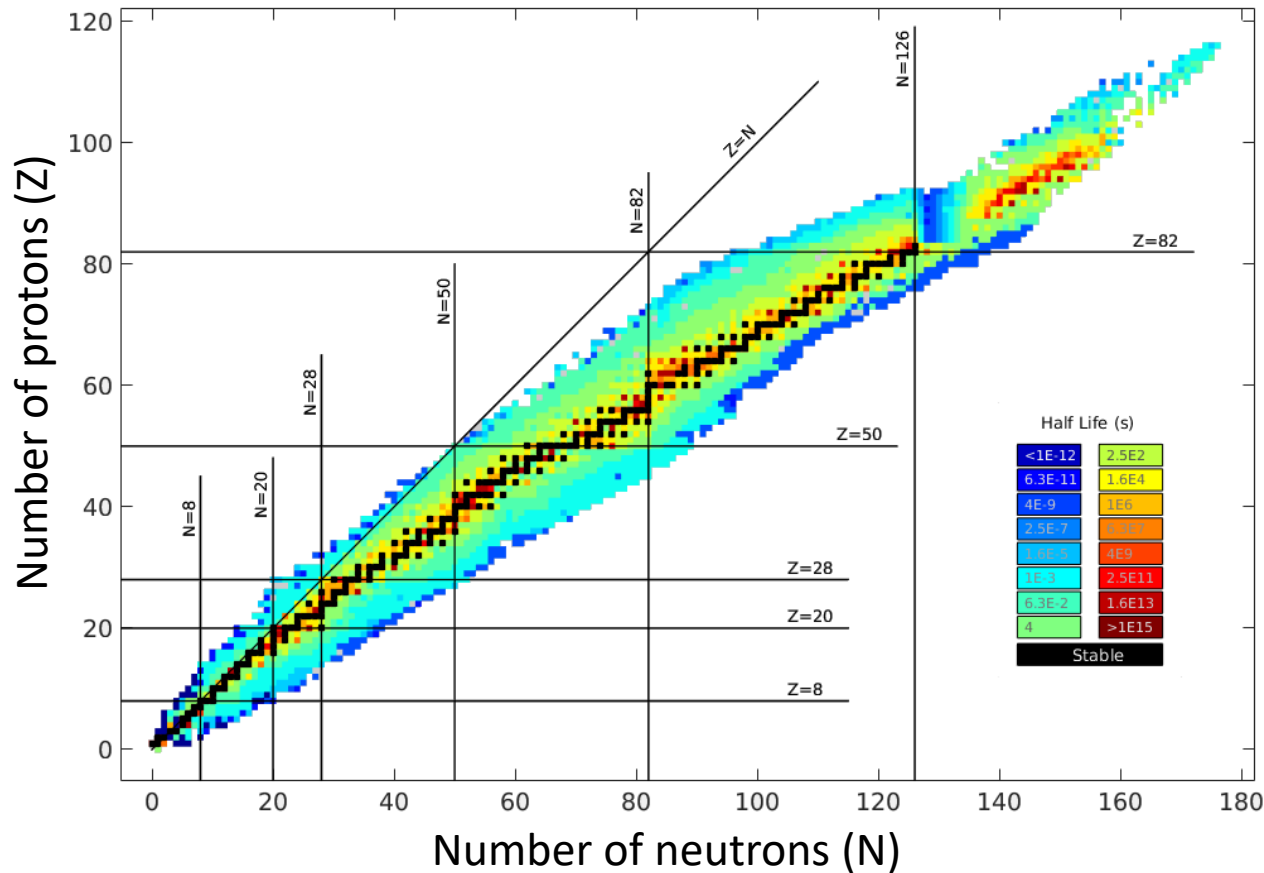
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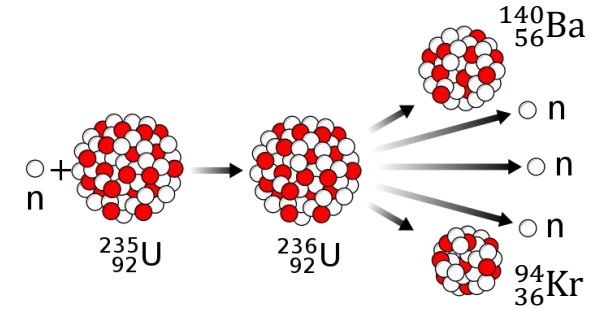
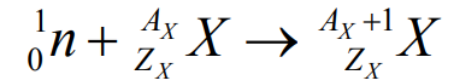
*Prompt-neutron generation time*



# Nuclear Landscape



## Nuclear Fission



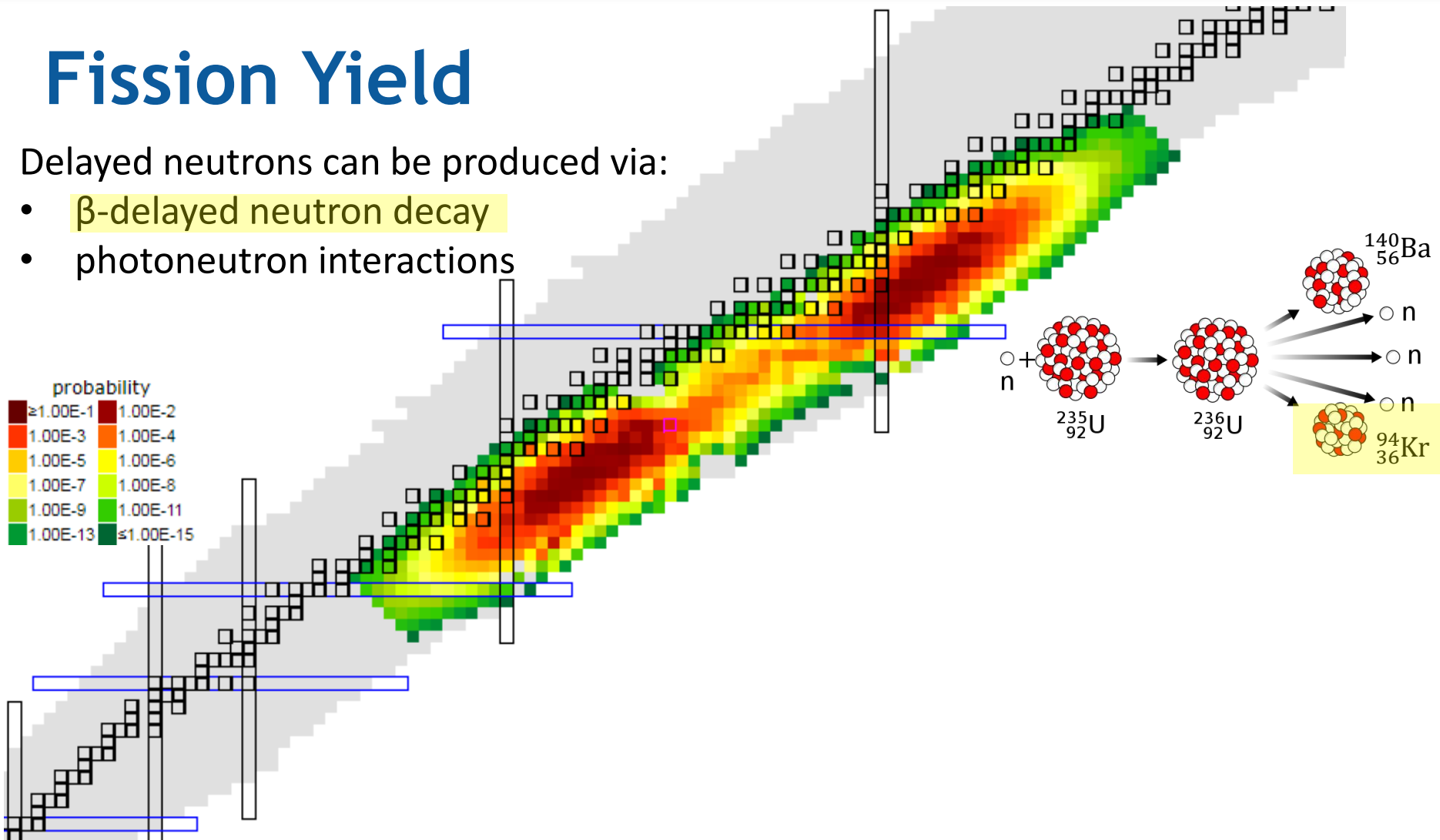
*The presence of delayed neutrons is perhaps most important aspect of the fission process from the viewpoint of reactor control!*



# Fission Yield

Delayed neutrons can be produced via:

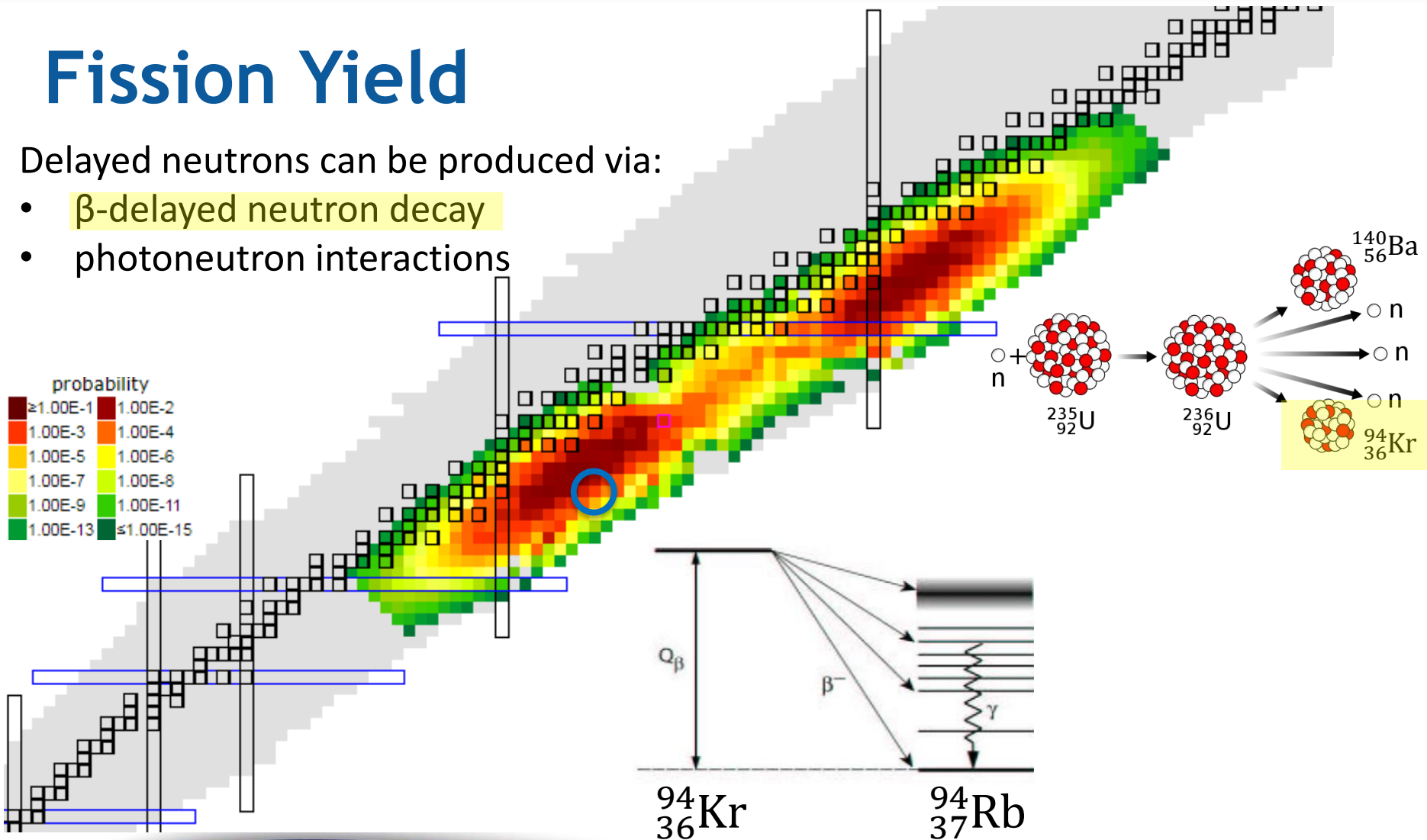
- $\beta$ -delayed neutron decay
- photoneutron interactions



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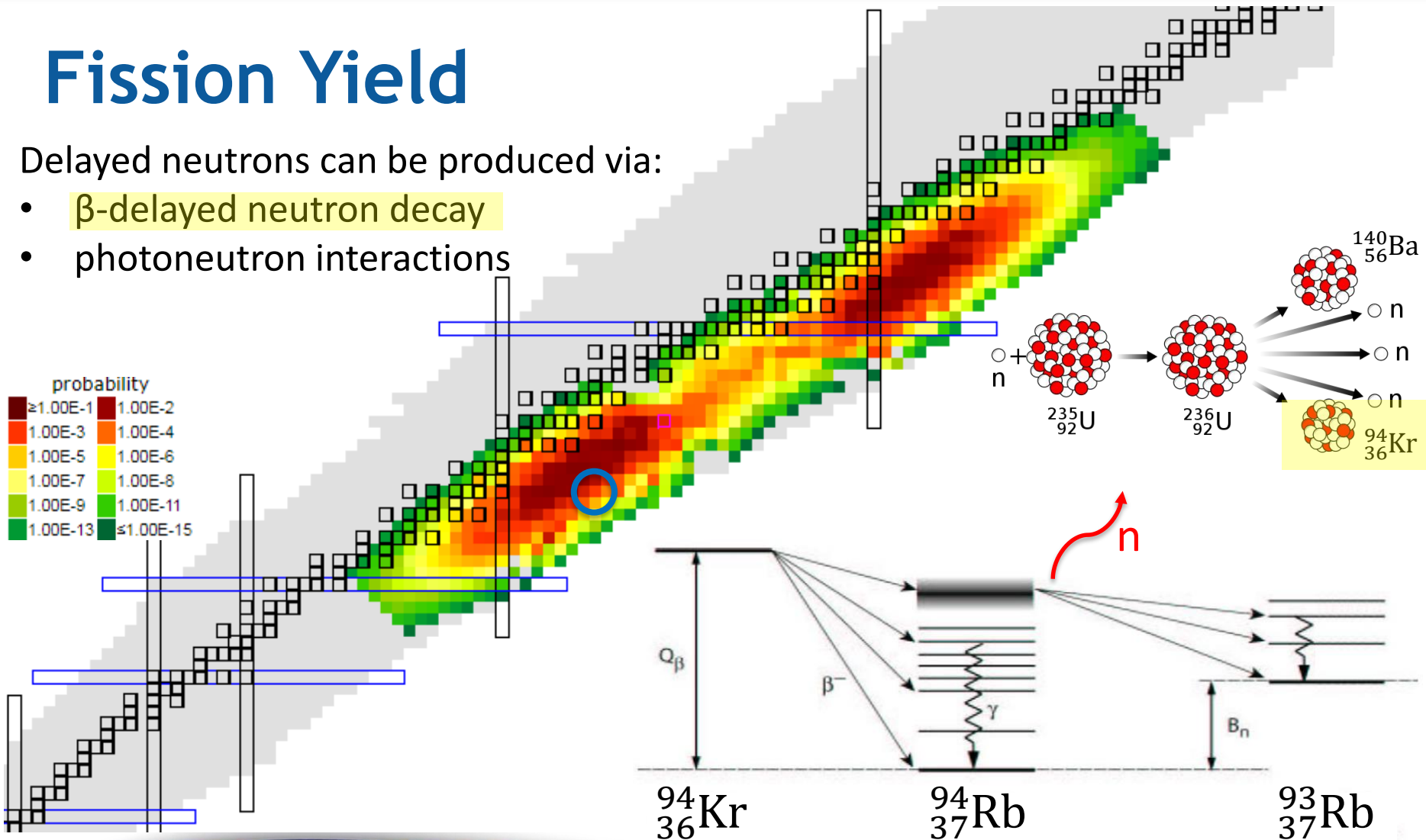
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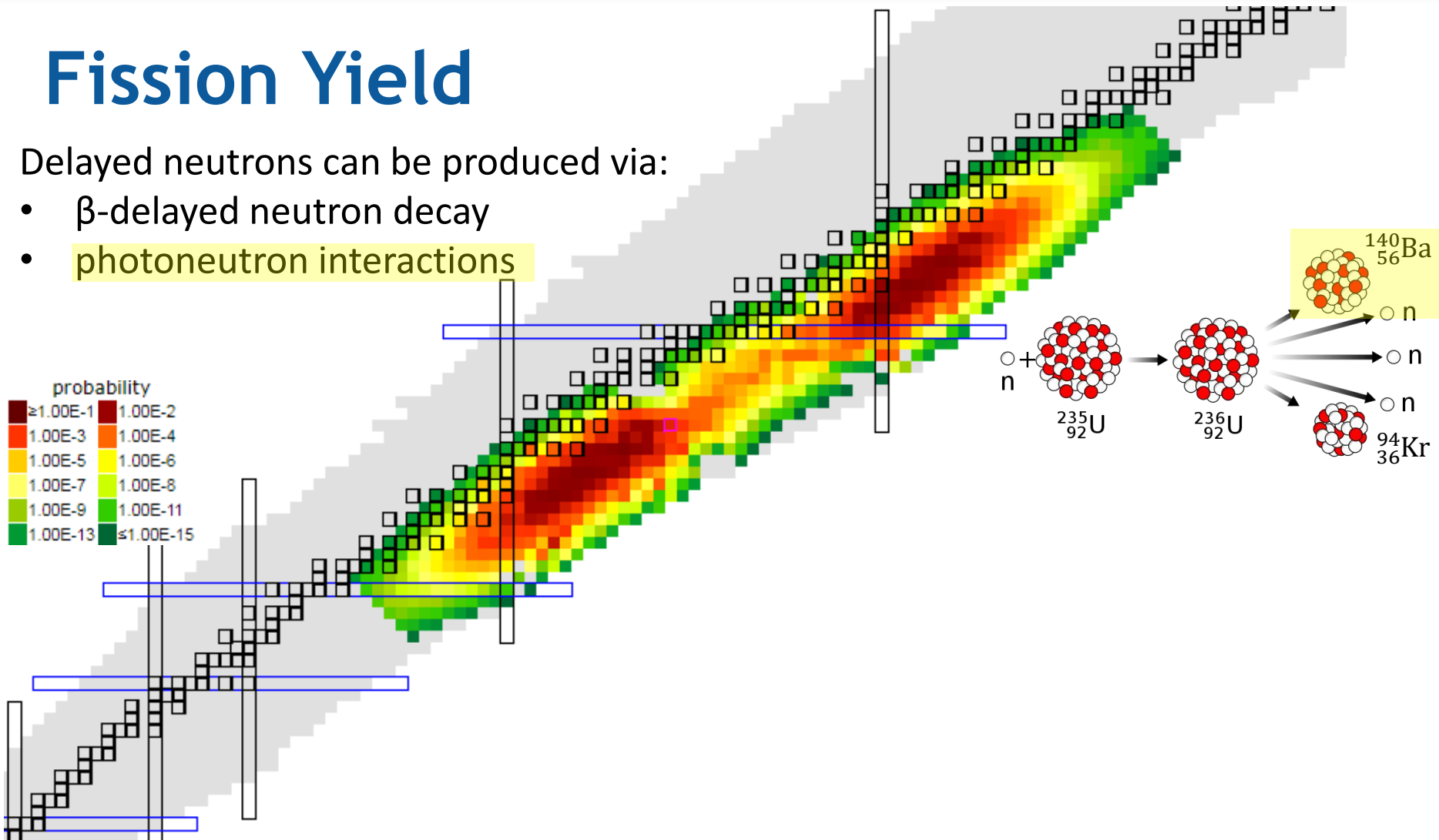




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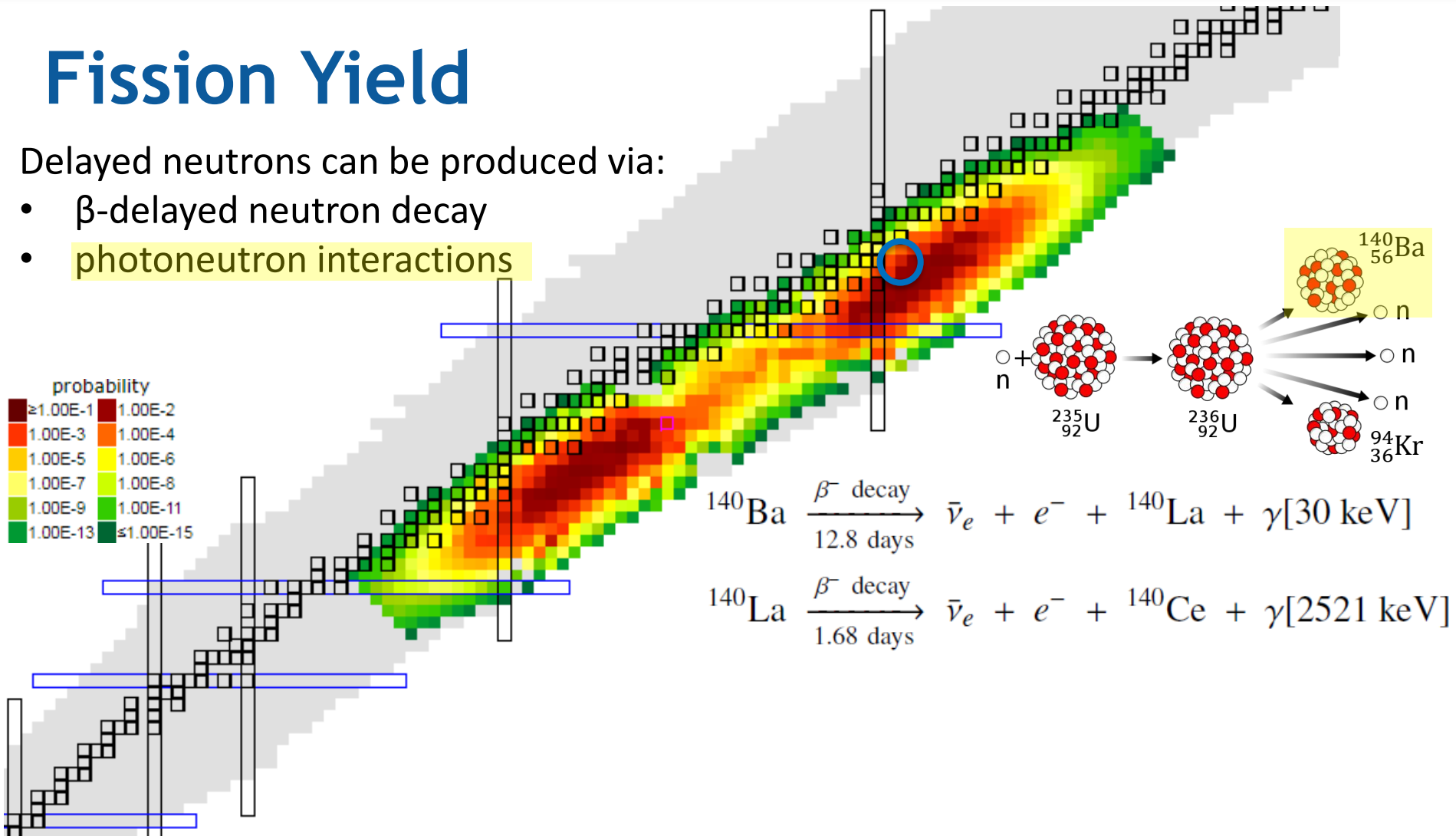
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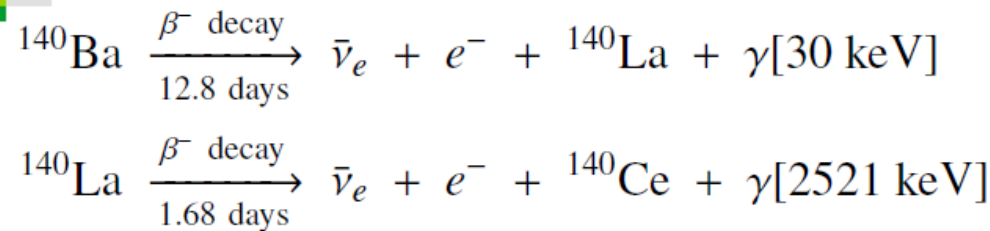
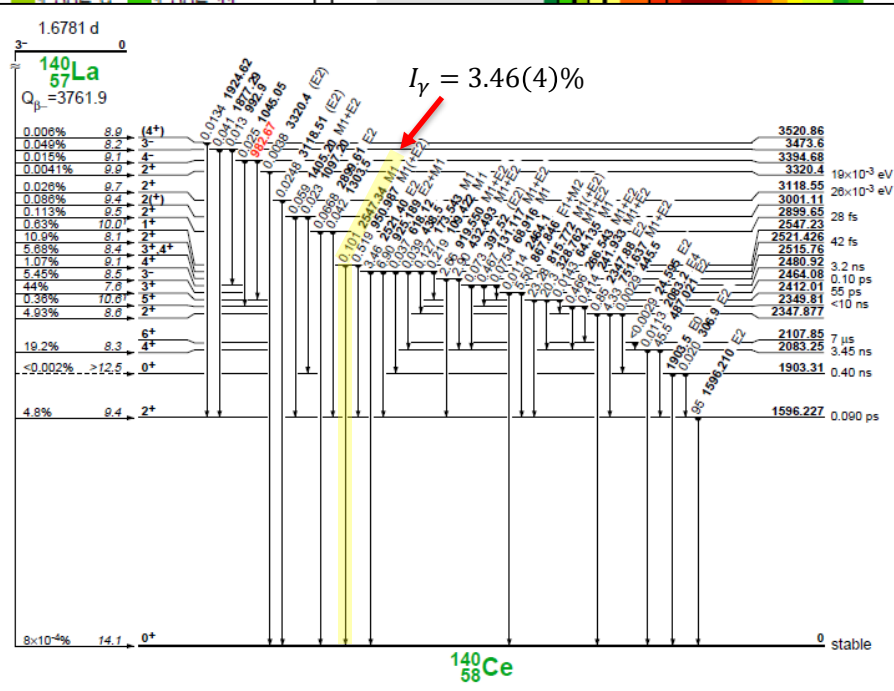
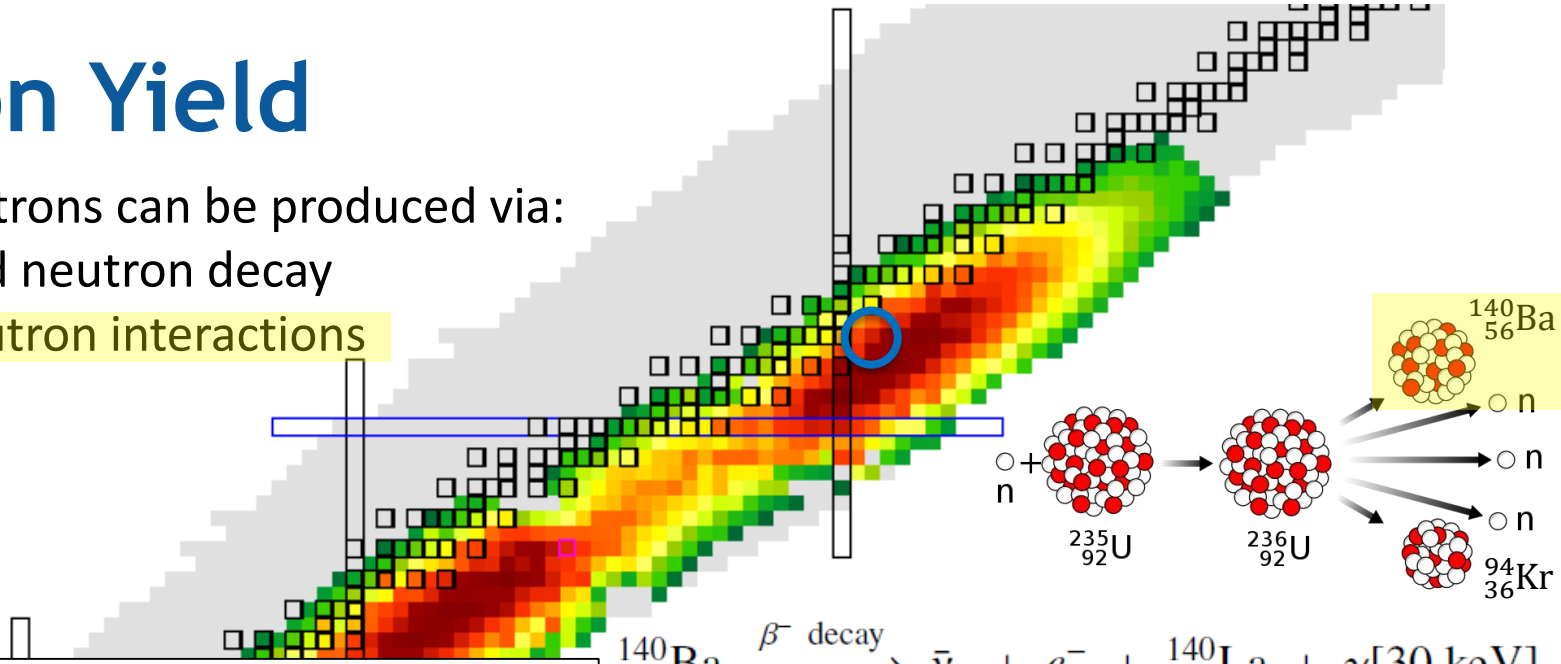
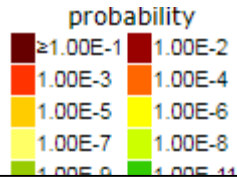
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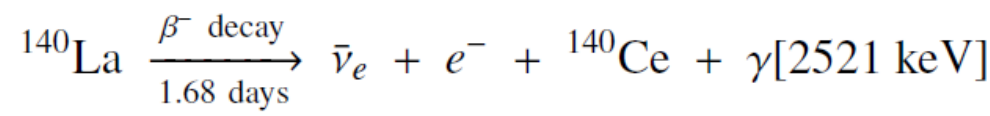
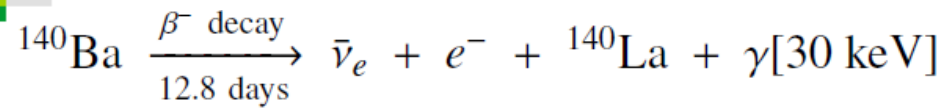
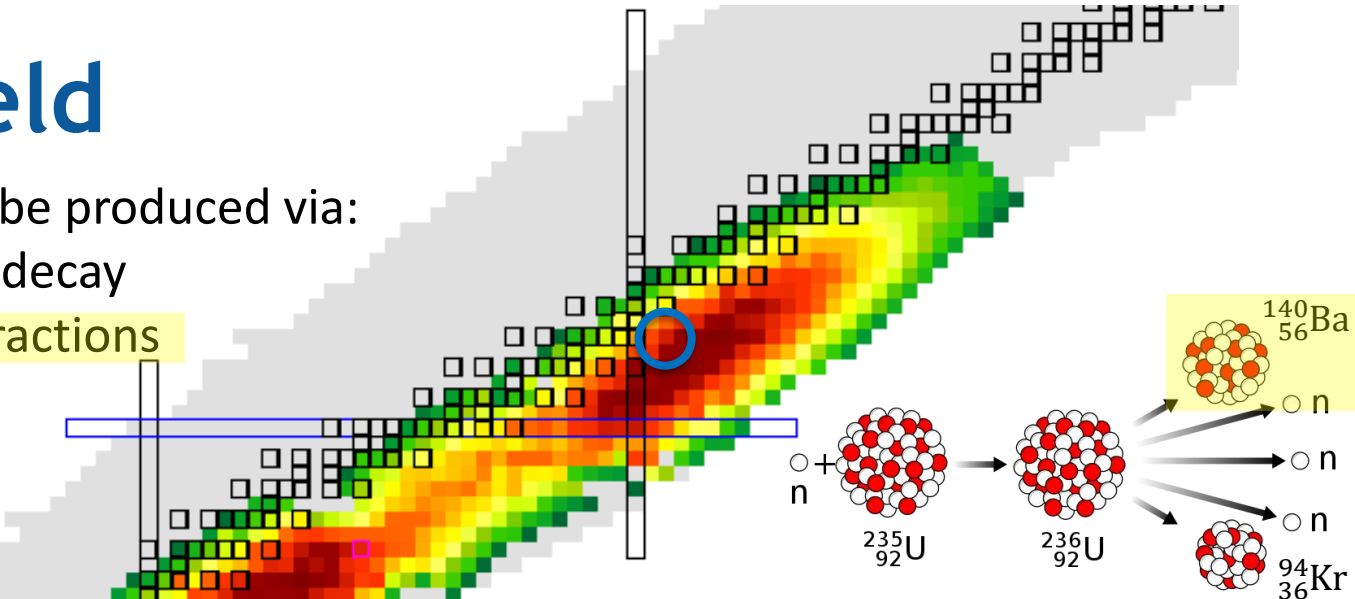
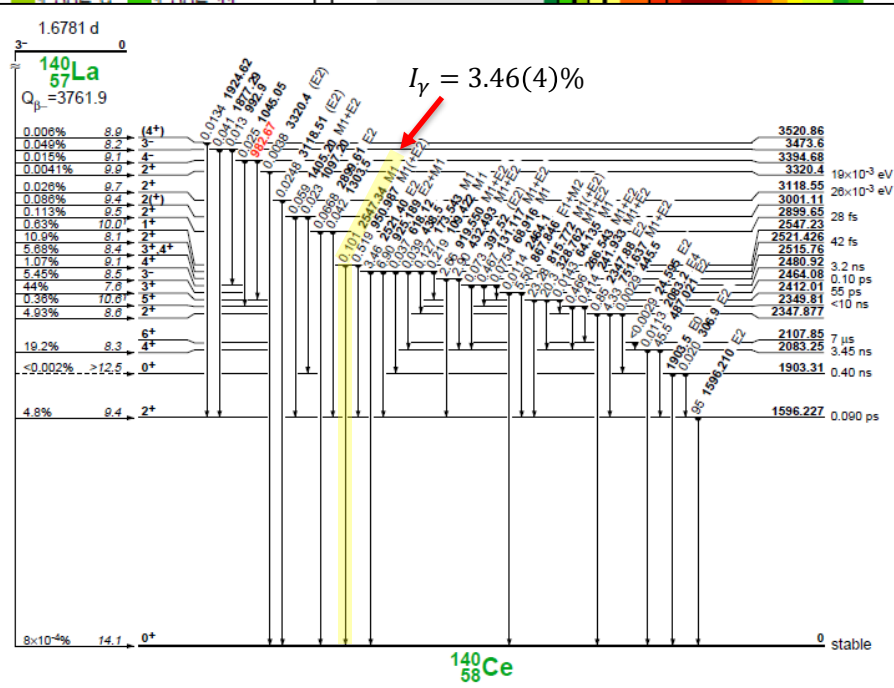
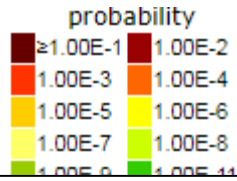
- $\beta$ -delayed neutron decay
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# Fission Yield

Delayed neutrons can be produced via:

- $\beta$ -delayed neutron decay
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Deuterium breakup reaction:

$$BE\left(\begin{smallmatrix} A \\ Z \end{smallmatrix} X\right) = Z \cdot M_p c^2 + N \cdot M_n c^2 - M'\left(\begin{smallmatrix} A \\ Z \end{smallmatrix} X\right) c^2$$

$$BE\left(^2\text{H}\right) = 2.225 \text{ MeV}$$

$$\gamma[2521 \text{ keV}] + ^2\text{H} \rightarrow p + \textcircled{n}$$

# Delayed Neutrons

$t_{1/2}$ (s)	1	2	3	4	5	6	7	8	9	10	11	← group number			
<b>Photo-neutron</b>	1,107,360	190,800	15,840	5,924	1,620	462.1	144.1	55.7	22.7	-	6.22	2.3	-	-	-
<b>Direct-delayed</b>	-	-	-	-	-	-	-	55.6	24.5	16.3	5.21	2.37	1.04	0.424	0.195
								1	2	3	4	5	6	7	8

*$t_{1/2} \sim 13$  days* (arrow pointing to group 1)

*$t_{1/2} \sim 1$  min* (arrow pointing to group 8)

A set of parameters covering the all photoneutron timescales was recommended for **CANDU** analyses\*. Interestingly, the data was stitched together from three different sources (Bernstein et al., Ergen, and Baumann), using experimental  $^{235}\text{U}$  data from as far back as the 1940s.

We performed detailed “microscopic” Geant4 calculations based on nuclear data to extract the photoneutron group structure, which had never been performed in the past.

Calculations performed on  $^{235}\text{U}$ ,  $^{233}\text{U}$ , and  $^{239}\text{Pu}$ .



\*P. J. Laughton, Revised delayed photoneutron data for use in CANDU-reactor analysis, in: Proc. Ann. Conf. Canadian Nucl. Soc., no. 2, 1998, pp. 376–388

# Simulation

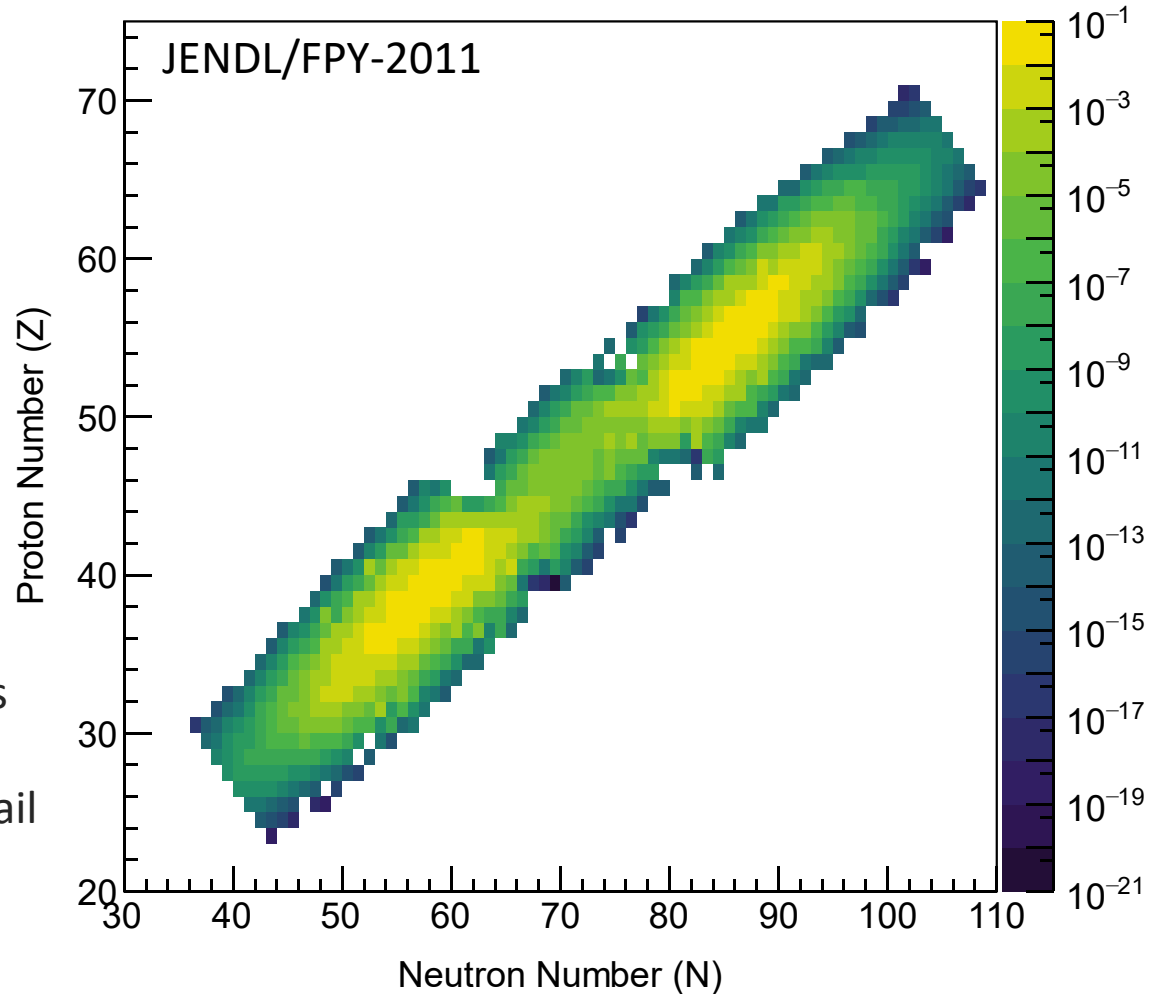
The simulation was built using Geant4 (10.03.p02), and applied radioactive decay and photon emission physics:

**G4RadioactiveDecay5.1.1**

**G4PhotonEvaporation4.3.2**

**G4ENSDFSTATE2.1.**

No geometry was simulated, rather, the emission of  $\beta$  particles and  $\gamma$ -rays from the radioactive decay process was tracked in detail from input isotope fission yields.



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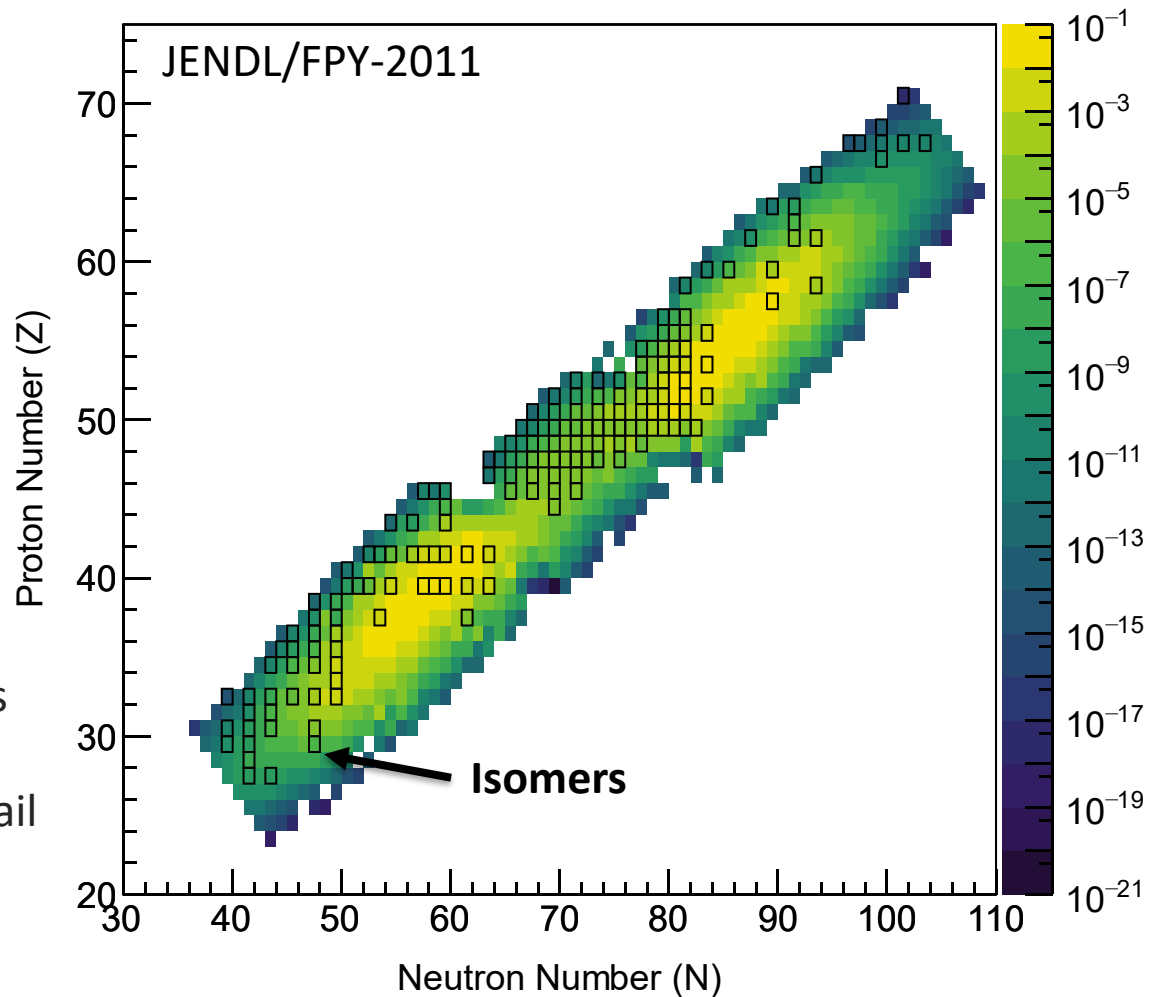
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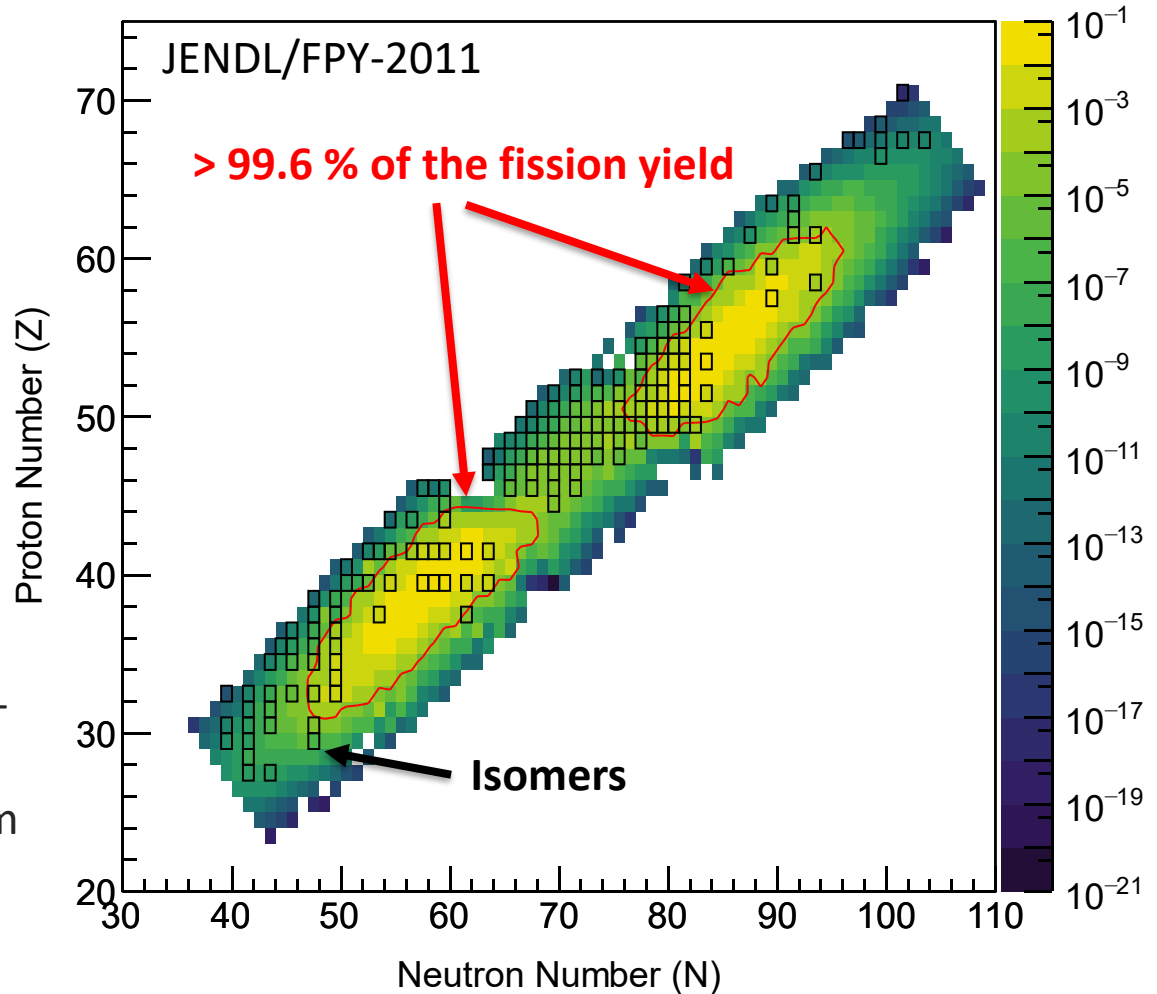
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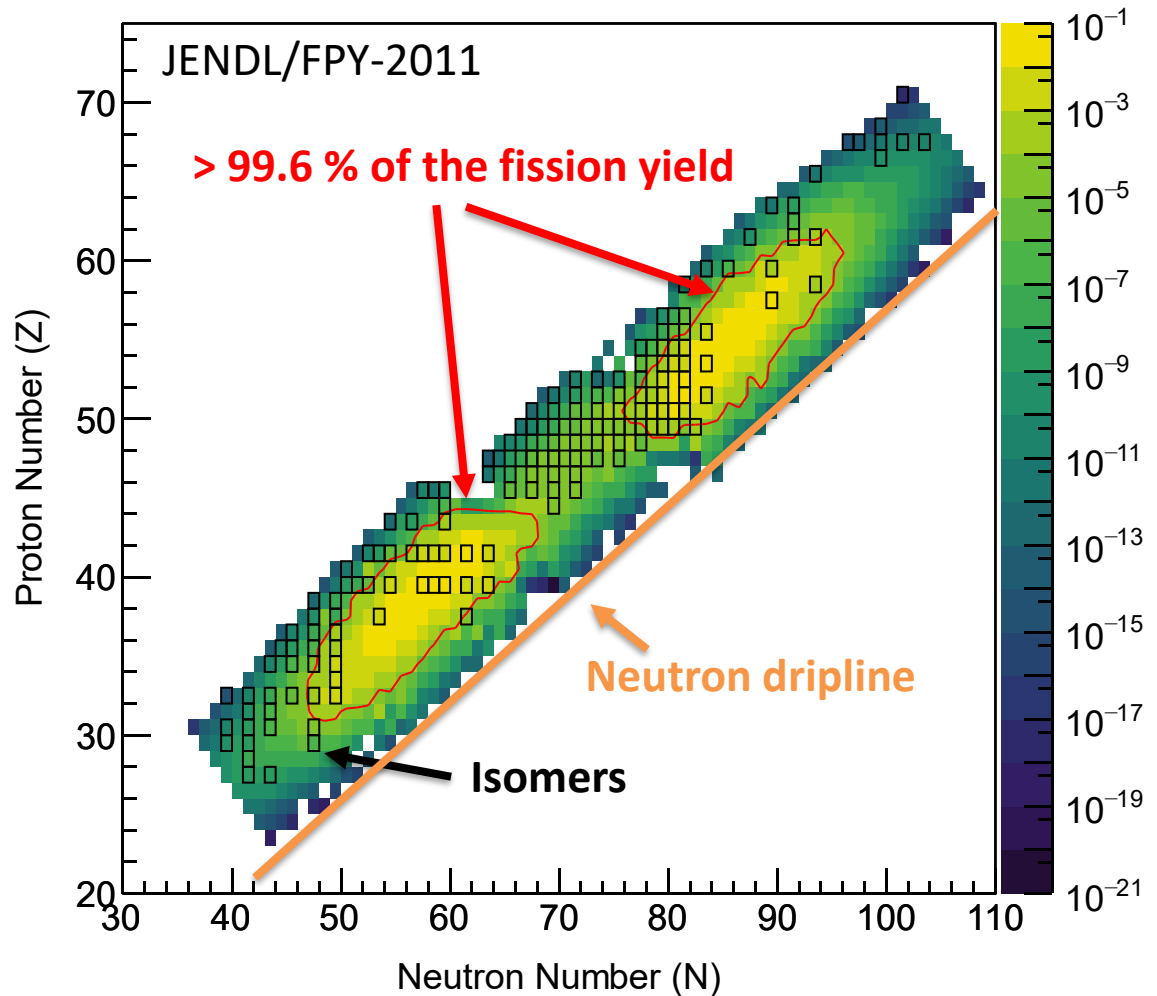
# Simulation

## Approximations

1. Nuclei near the neutron dripline with no known nuclear data information were input into the simulation using theoretical  $\beta$ -decay half-lives\*. These calculations used a proton-neutron quasiparticle random-phase approximation (pn-QRPA) and applied mass formula predictions from Hilf et al.† which are thought to be the most reliable far from stability.

\*A. Staudt et al., Atomic Data and Nuclear Data Tables 44 (1) (1990) 79-132.

†E. Hilf et al., Gross theory of nuclear masses and radii, in: the 3rd International Conference on Nuclei far from Stability (1976)



# Simulation

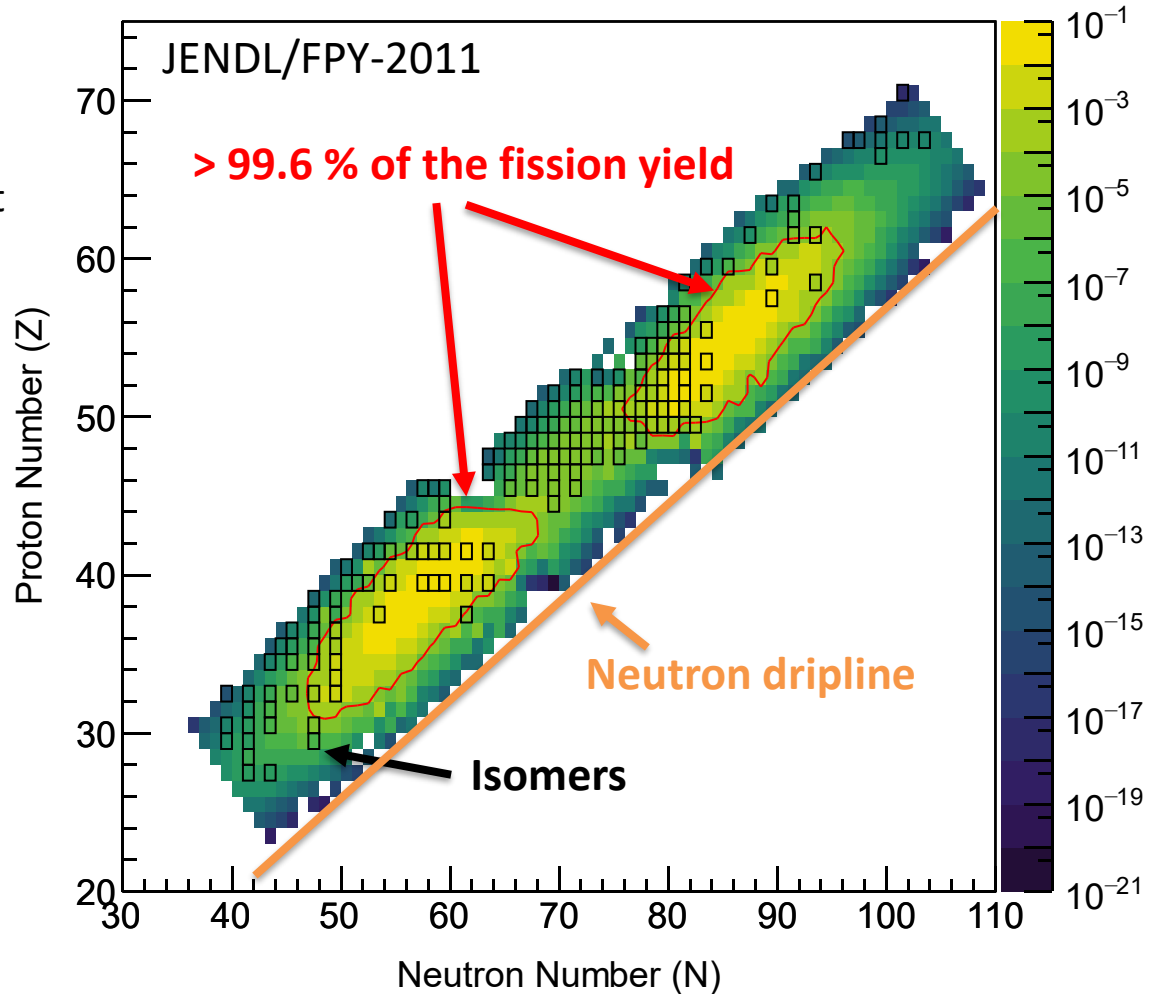
## Approximations

2. Geant4 currently does not support  $\beta$ -delayed neutron decay.

**2019 Planned Feature in Geant4!**

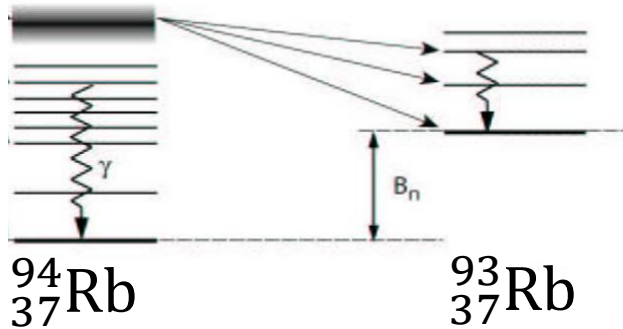
Fortunately, we calculated  $< 1\%$  of the total fission yield directly populates  $\beta$ -delayed neutron branches in  $^{235}\text{U}$ , determined from measured branching ratios and predictions using an updated empirical Kratz Herrmann formula (KHF)\* where data was lacking.

\* B. Pfeiffer et al., Progress in Nuclear Energy 41 (1) (2002) 39-69.

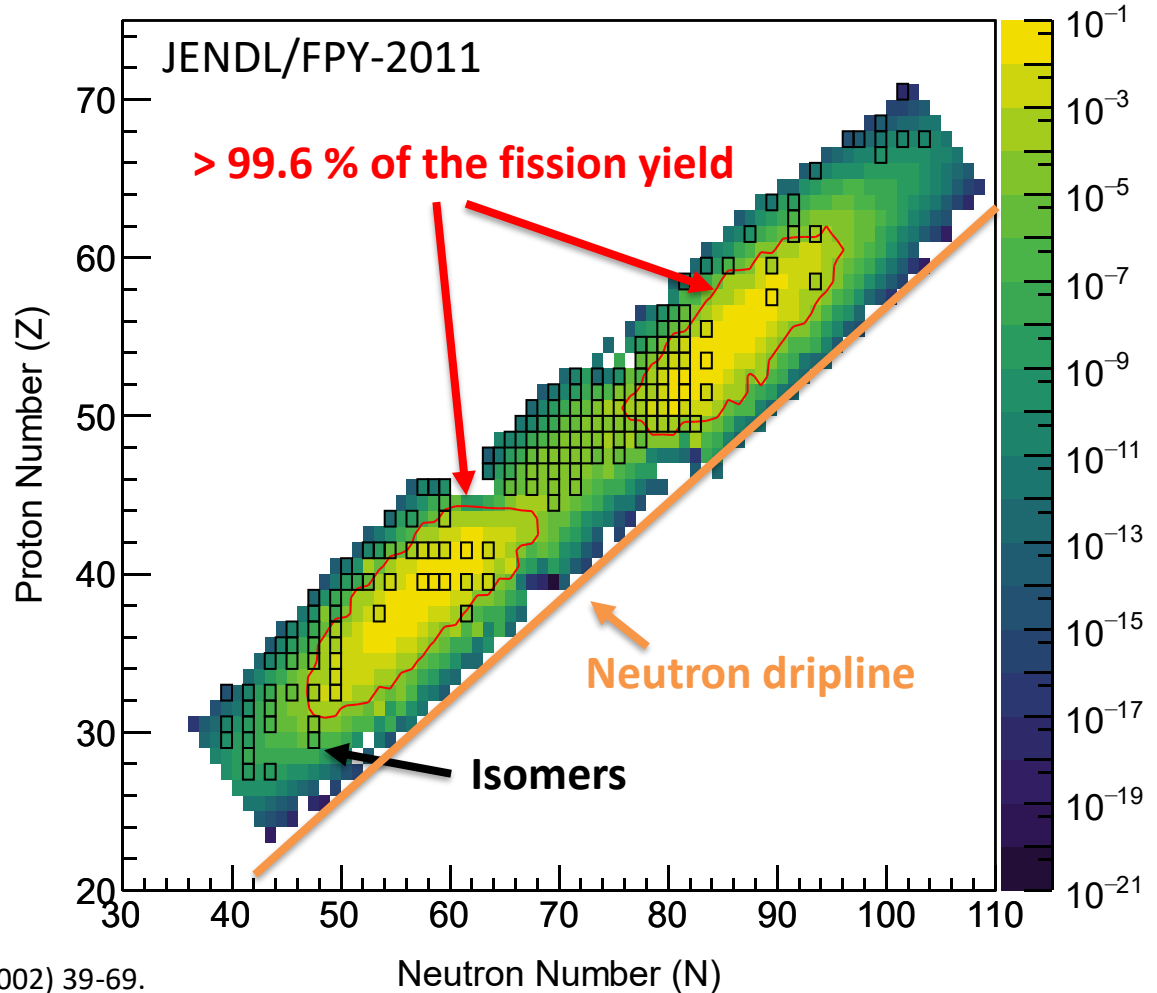


# Simulation

## Approximations



Correlations between  $\beta$ -delayed neutron emission and high-energy  $\gamma$  decay was also found to be small. Determined by conservatively assuming  $N_\gamma$  to be reduced by  $P_n$  for each isotope. In total, the number of high-energy  $\gamma$ -rays emitted per fission was only reduced by  $< 1.5\%$ .



\* B. Pfeiffer et al., Progress in Nuclear Energy 41 (1) (2002) 39-69.

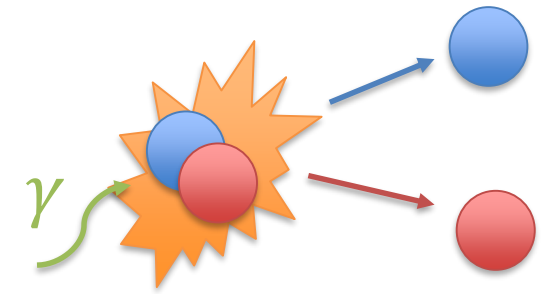


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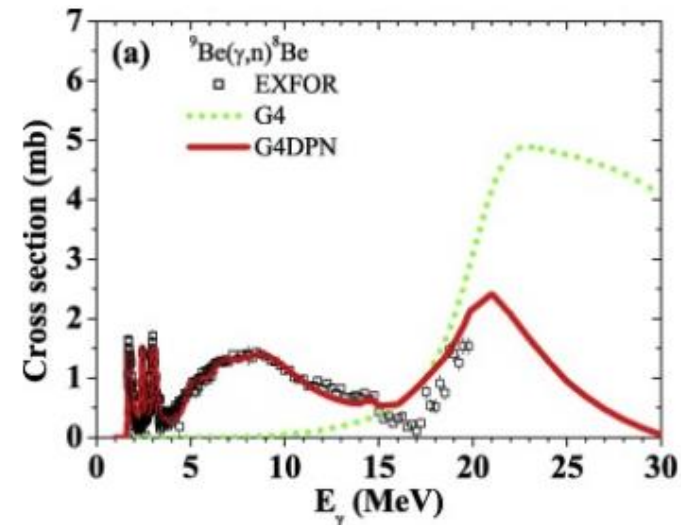
## Modifications to Geant4 Physics

Photonuclear interactions with nuclei are handled by the **G4PhotoNuclearProcess** class in Geant4 (using a Bertini-style cascade). The model has only been validated at much higher energies (> **60 MeV**) and fails to produce neutrons from  ${}^2\text{H}$  in our energy regime.

Large discrepancies for other isotopes observed by others, *J. W. Shin, Nucl Instrum Methods Phys Res B 358 (2015)*. In this work a data-based (ENDF/B-VII.1) photonuclear reaction model for Geant4 was developed for a more accurate simulation of photon-induced neutron production < 30 MeV.

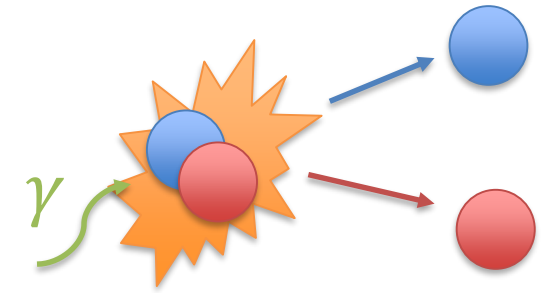
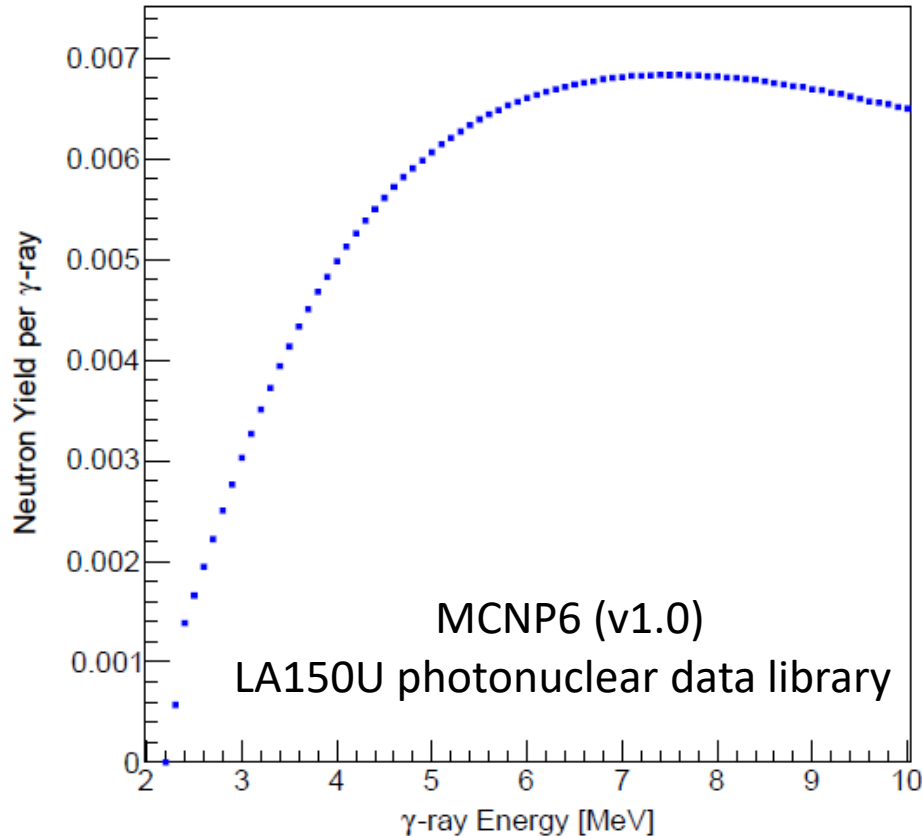


$$\sigma_{\text{Geant4}} = 0$$
$${}^2\text{H}(\gamma, n) {}^1\text{H}$$



# Simulation

## Modifications to Geant4 Physics



MCNP6 photonuclear physics\* for deuterium was hardcoded into our Geant4 simulation.



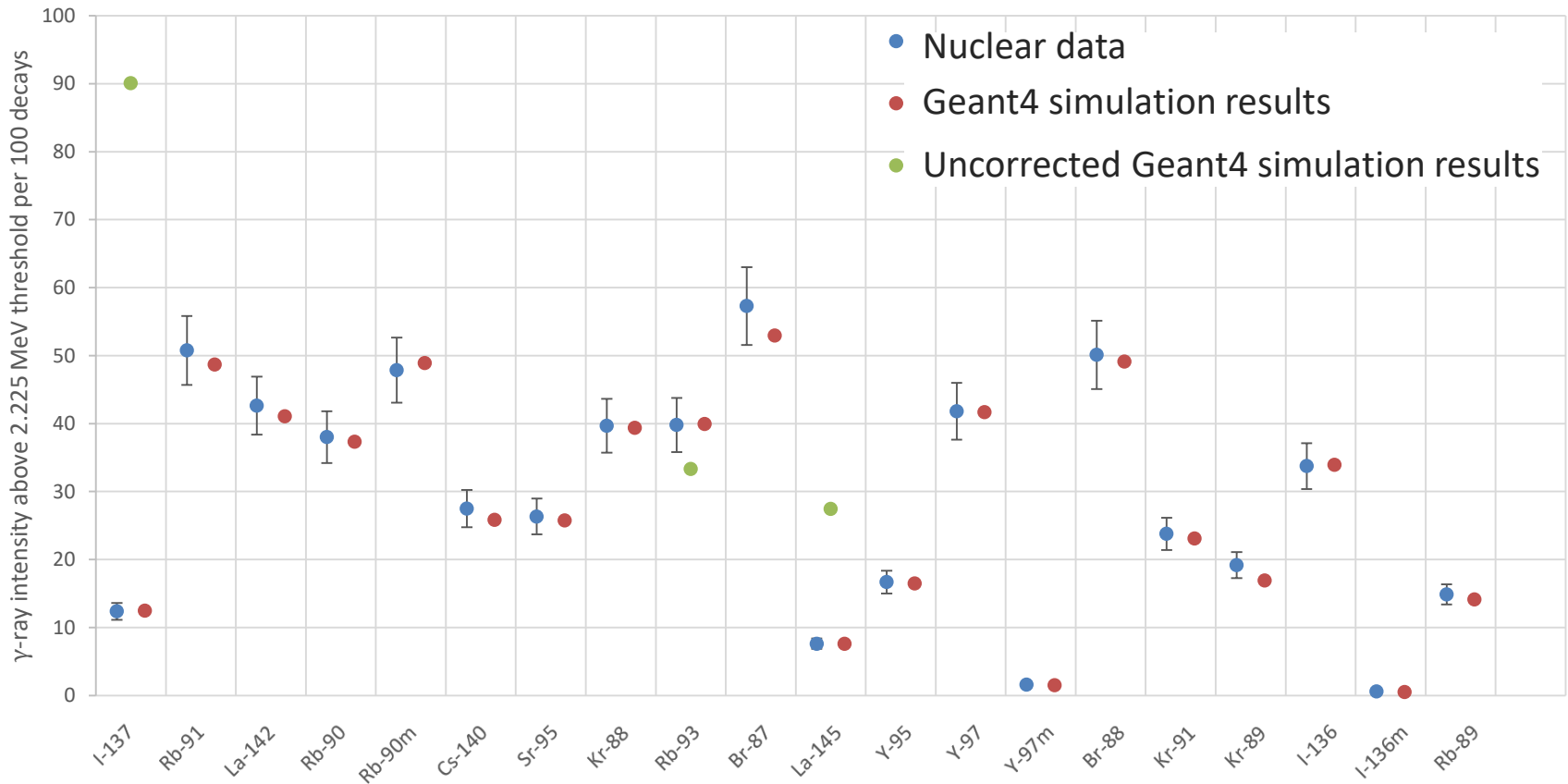
MCNP6 photoneutron production simulation from a point-source of monoenergetic  $\gamma$ -rays within a 10 m radius homogeneous sphere of heavy water, which effectively mimics an infinite heavy water bath for  $\gamma$ -ray energies < 10 MeV.

\*M. C. White, Release of the LA150U photonuclear data library, Tech. Rep. memo X-5:MCW-00-87(U), Los Alamos National Lab (2000)



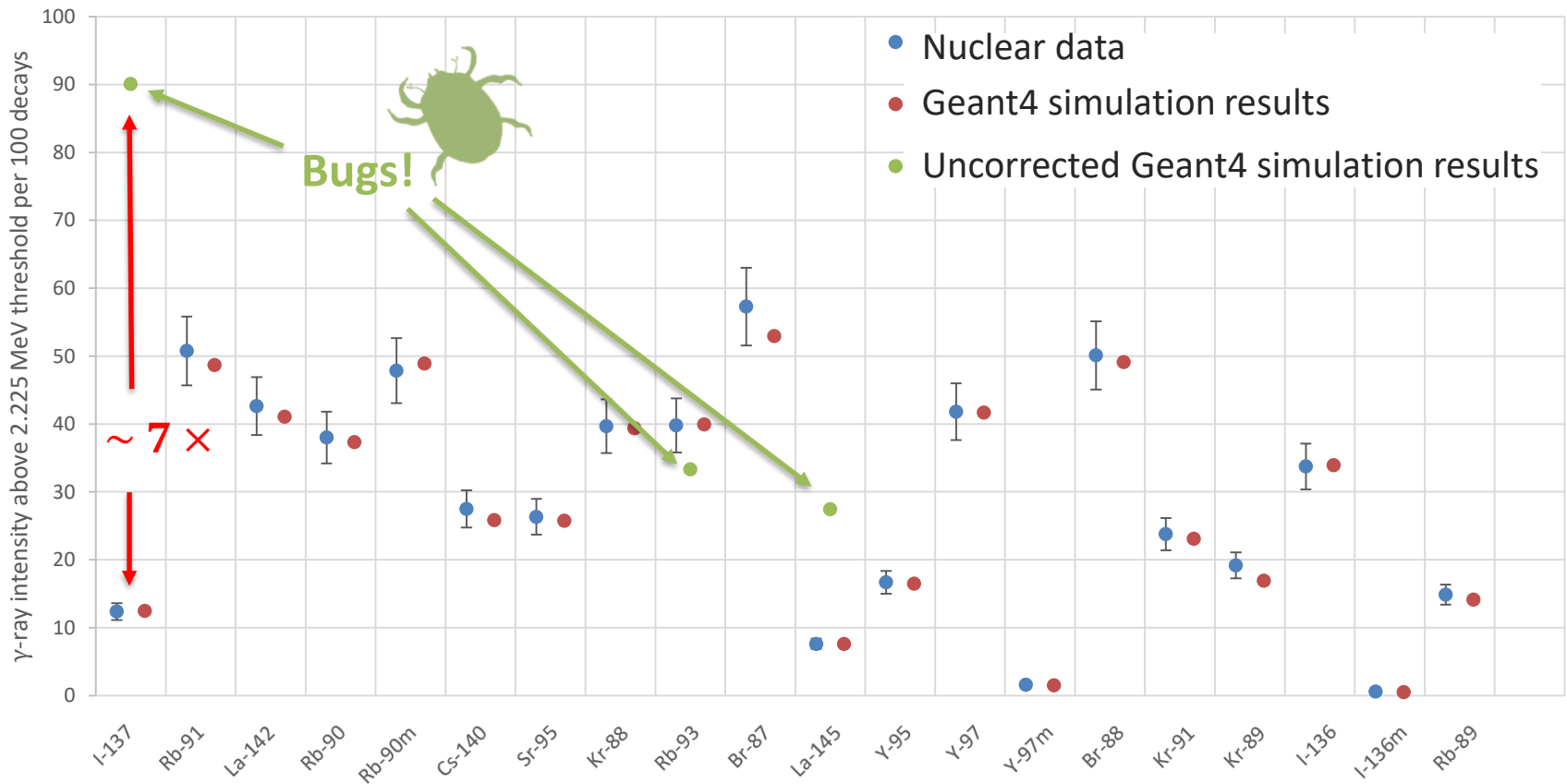
# Validation Work

Total  $\gamma$ -ray yield above 2.225 MeV

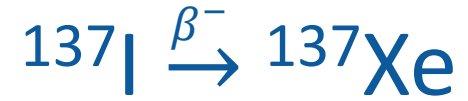


# Validation Work

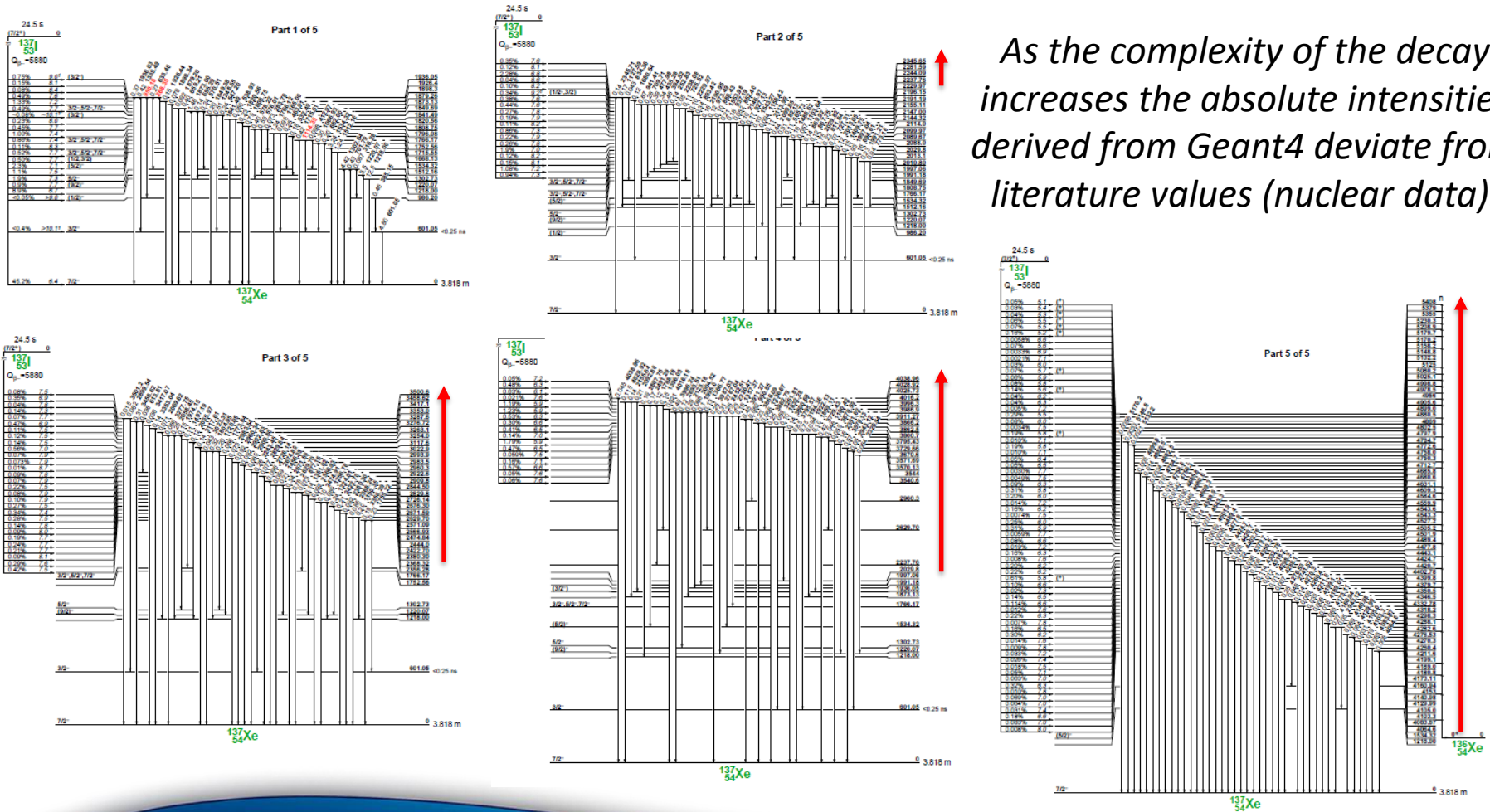
Total  $\gamma$ -ray yield above 2.225 MeV



# Validation Work



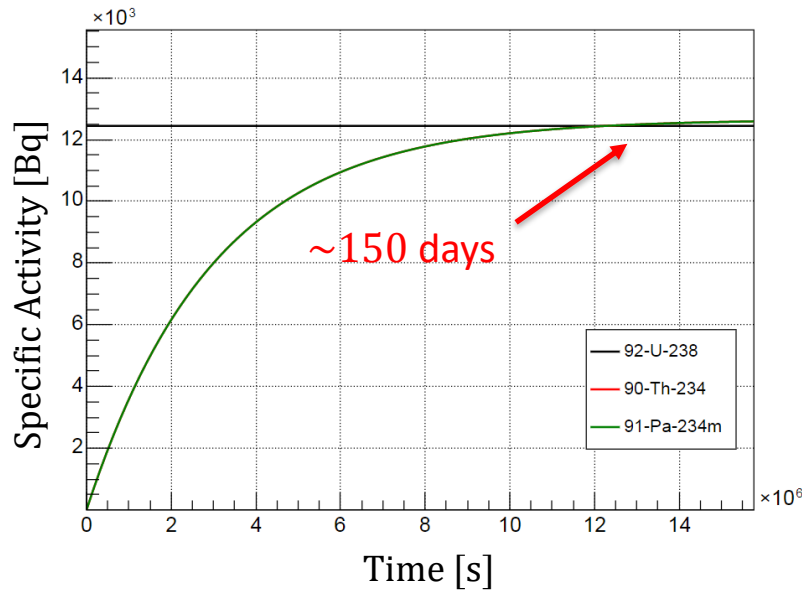
As the complexity of the decay increases the absolute intensities derived from Geant4 deviate from literature values (nuclear data).



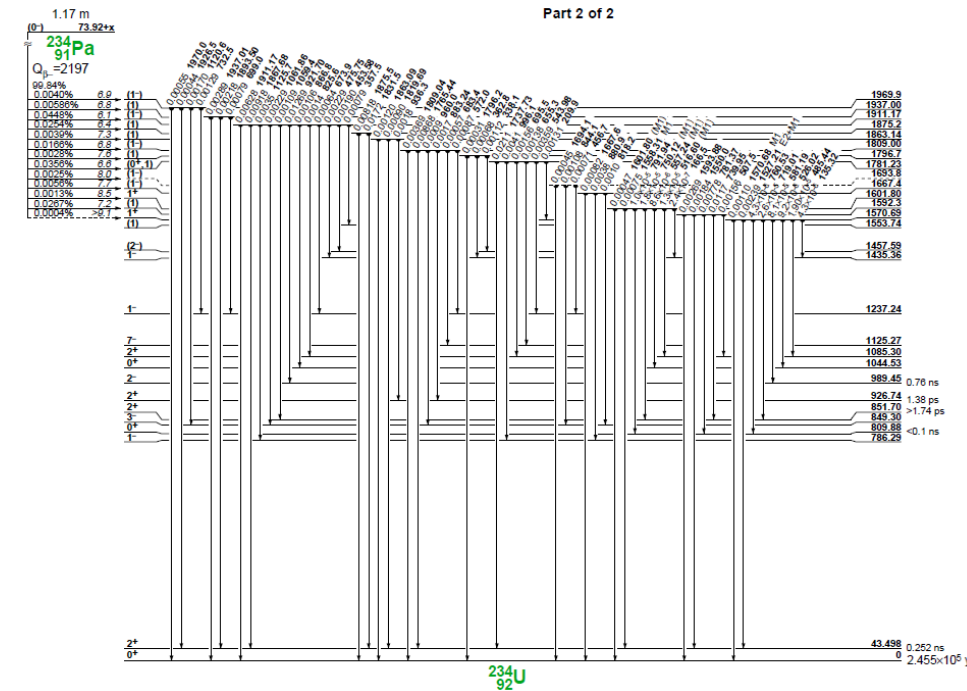
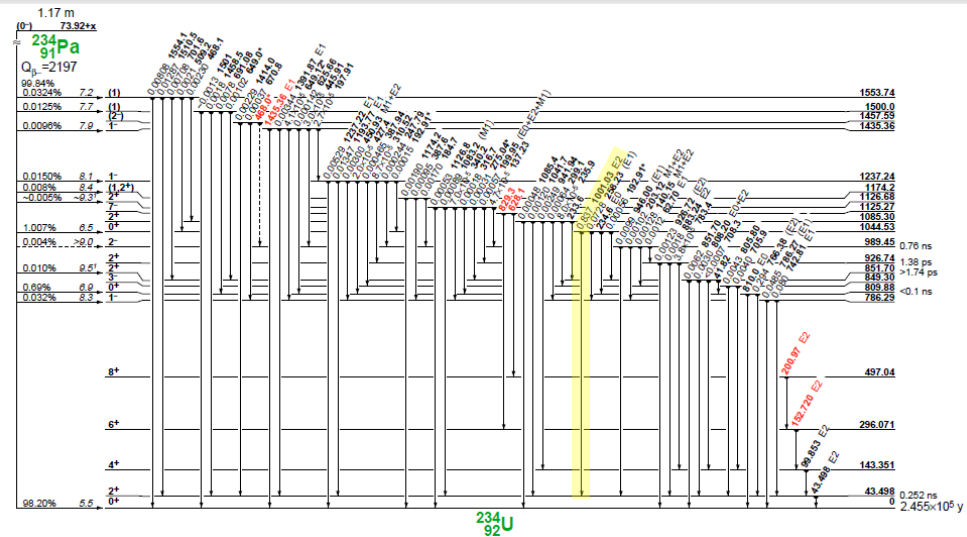


# Validation Work

Another important example in the detection of  $^{238}\text{U}$



The absolute intensity of the 1 MeV  $\gamma$ -ray from  $^{234\text{m}}\text{Pa}$  decay was found to be 0.9125% using the Geant4 data libraries, which is approximately 10% larger than what is found in literature.



# Results

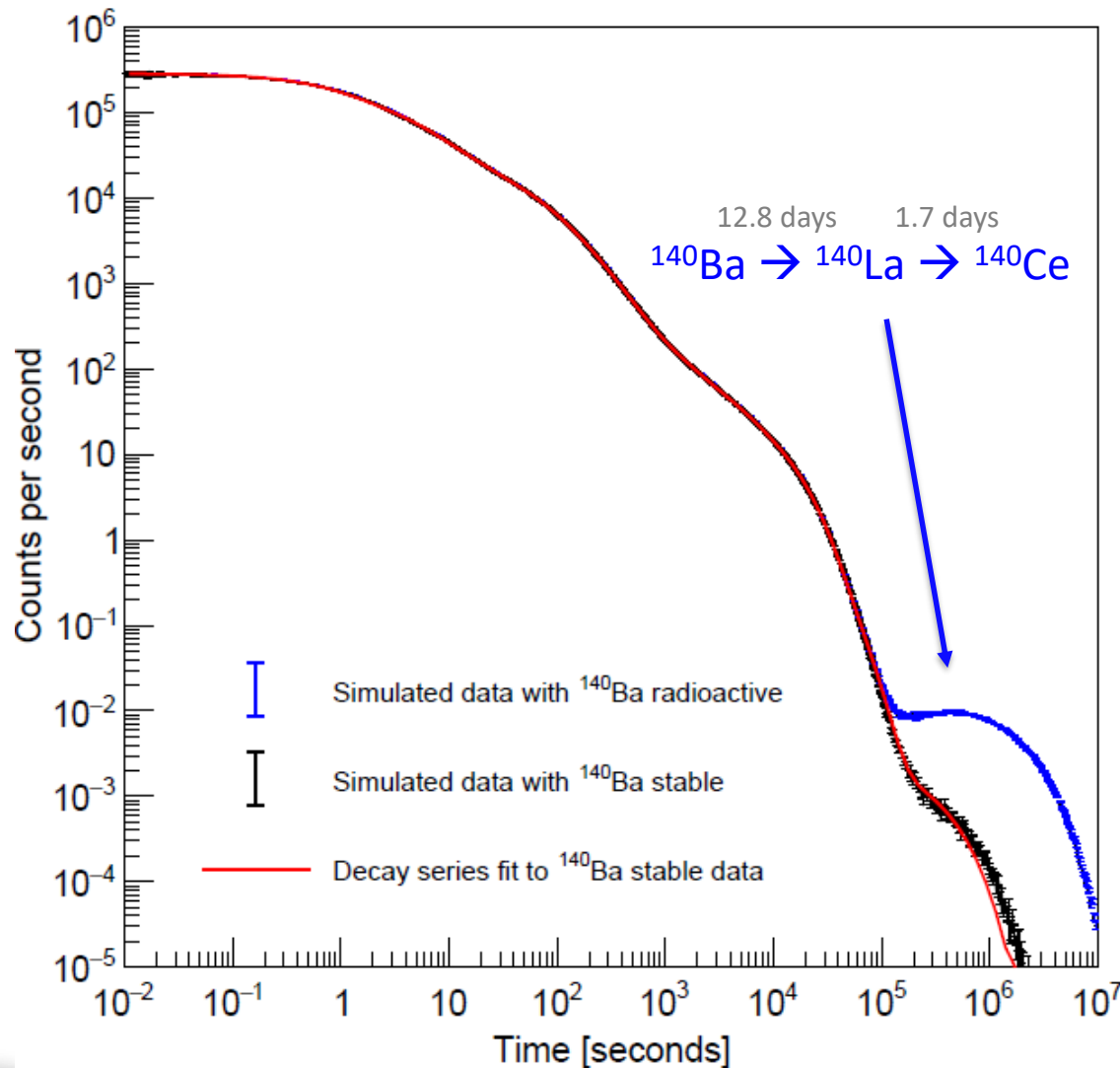
One billion  $^{233}\text{U}$ ,  $^{235}\text{U}$  and  $^{239}\text{Pu}$  fission fragment events were simulated. For **computational efficiency** the pn prod. curve was increased 145-fold. Thus, each simulation represented approximately **72.5 billion effective fission events**.

Similar to past experimental work, data was fit to a series of exponential terms,

$$f(t) = \sum_{i=1}^{12} A_i e^{-\lambda_i t},$$

using ROOT and the MINUIT2  $\chi^2$  minimization toolkit. Relative pn yields for each group was extracted:

$$\frac{\beta_{\gamma_i}}{\beta_{\gamma}} = \frac{A_i / \lambda_i}{\sum_{i=1}^{12} A_i / \lambda_i}.$$



# Results

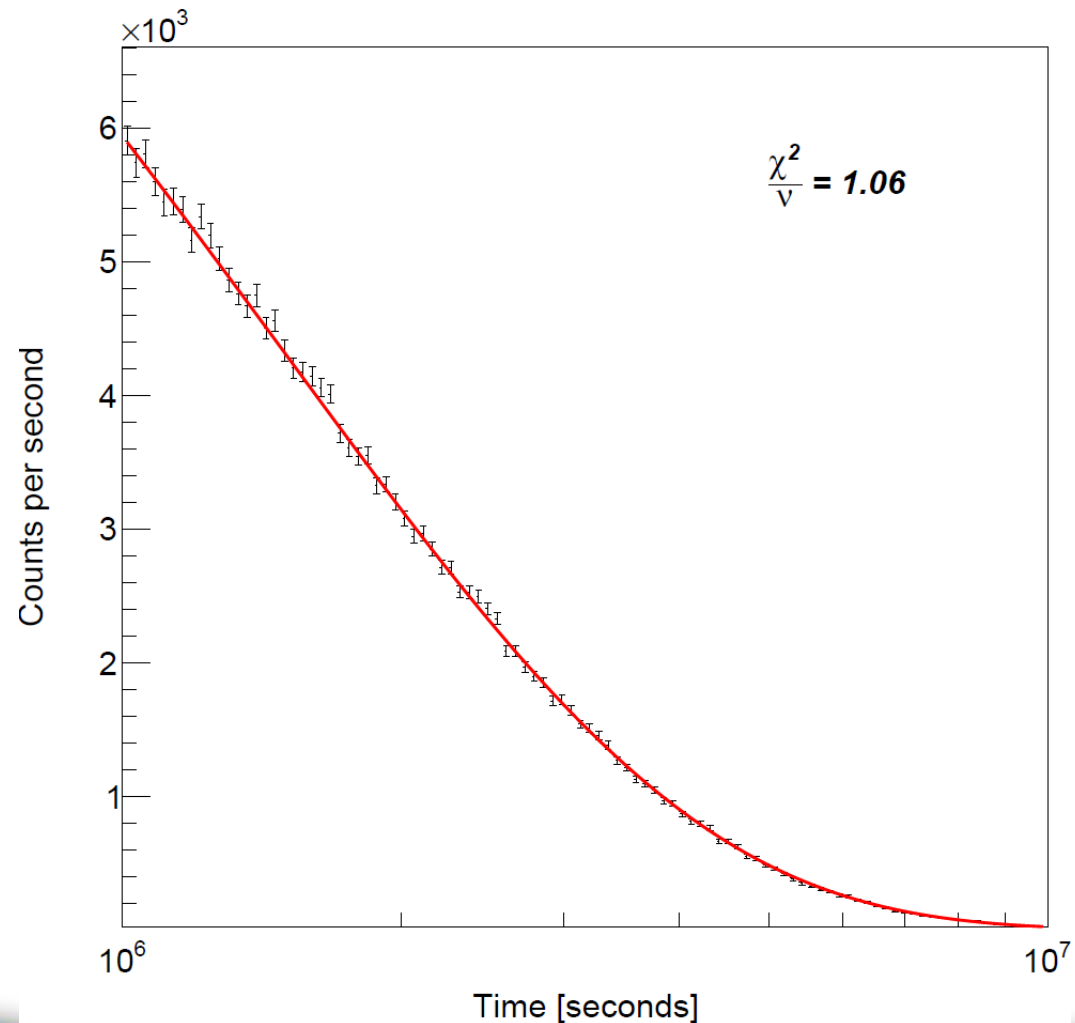
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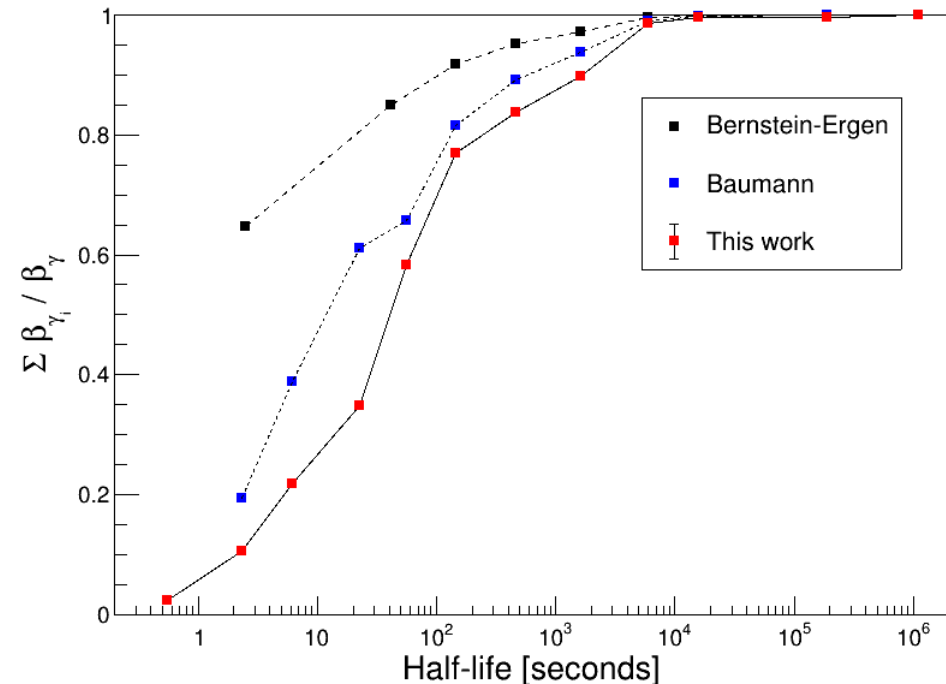
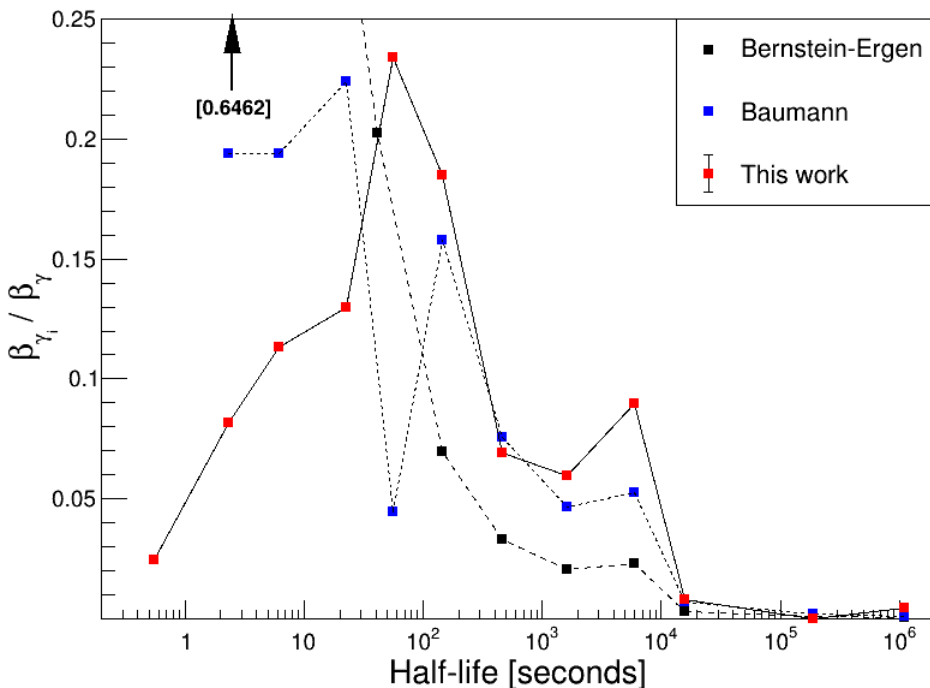
# Results

Relative photoneutron group yields for  $^{233}\text{U}$ ,  $^{235}\text{U}$  and  $^{239}\text{Pu}$  derived from simulation. The errors are purely statistical resulting from the  $\chi^2$  minimization fitting and do not include uncertainties persisting in the nuclear data.

Group (i)	Half-life (s)	$\beta_{\gamma_i}/\beta_{\gamma} (\times 10^3)$				
		$^{235}\text{U}$	$^{233}\text{U}$	$^{239}\text{Pu}$	Baumann	Bernstein-Ergen
1	1107360	$4.49 \pm 0.03$	$5.27 \pm 0.03$	$6.83 \pm 0.04$	1.1	0.5
2	190800	$0.205 \pm 0.002$	$0.261 \pm 0.003$	$0.568 \pm 0.005$	2.3	1.0
3	15840	$8.04 \pm 0.03$	$12.75 \pm 0.04$	$6.77 \pm 0.04$	7.3	3.2
4	5924	$89.48 \pm 0.08$	$128.2 \pm 0.1$	$109.6 \pm 0.1$	52.7	23.2
5	1620	$59.6 \pm 0.1$	$65.7 \pm 0.1$	$99.3 \pm 0.2$	46.6	20.5
6	462.1	$69.0 \pm 0.2$	$94.3 \pm 0.2$	$80.5 \pm 0.2$	75.7	33.4
7	144.1	$185.5 \pm 0.3$	$202.7 \pm 0.3$	$98.9 \pm 0.4$	157.6	69.5
8	55.7	$234.4 \pm 0.4$	$245.9 \pm 0.4$	$263.0 \pm 0.5$	44.8	-
-	41	-	-	-	-	202.5
9	22.7	$129.6 \pm 0.3$	$115.9 \pm 0.4$	$172.4 \pm 0.5$	223.9	-
10	6.22	$113.3 \pm 0.3$	$75.9 \pm 0.3$	$61.2 \pm 0.3$	194.0	-
-	2.5	-	-	-	-	646.2
11	2.3	$81.6 \pm 0.2$	$44.6 \pm 0.2$	$78.5 \pm 0.2$	194.0	-
12	0.548	$24.81 \pm 0.06$	$8.41 \pm 0.05$	$22.48 \pm 0.08$	0	-



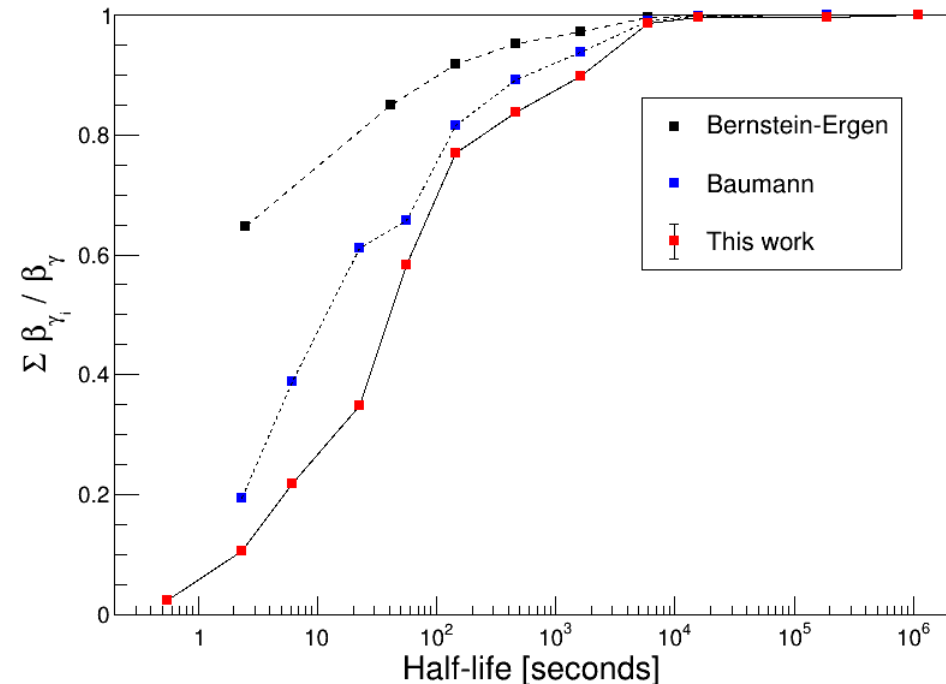
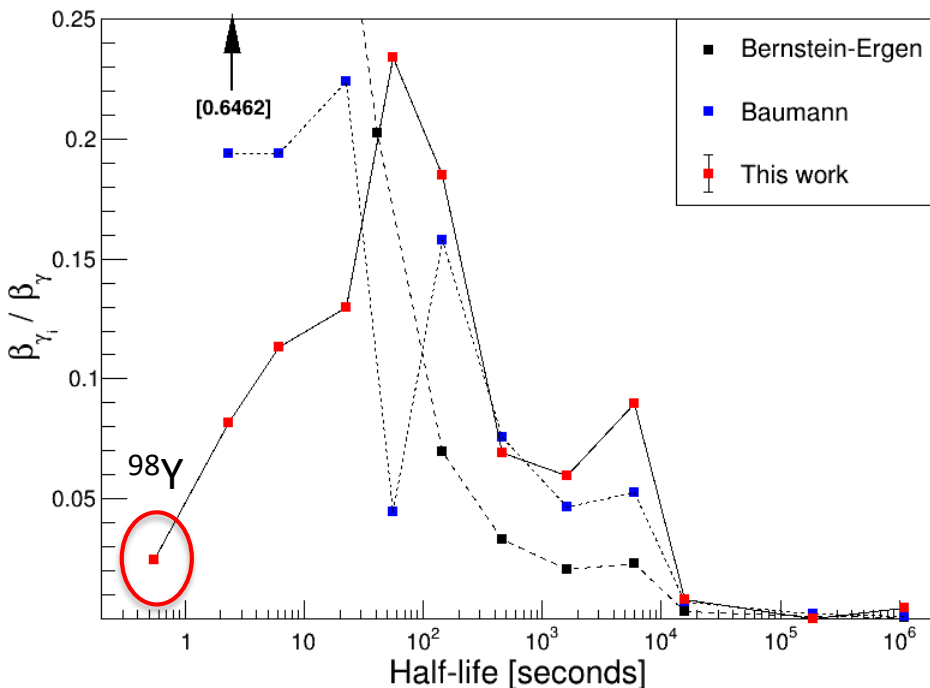
# Results



Relative photoneutron group yields (a) and running total (b) for the thermal fission of  $^{235}\text{U}$  calculated using Geant4. The current recommended values based on the  $^{235}\text{U}$  experimental work of Baumann and others are presented under the label “Baumann”.



# Results



Relative photoneutron group yields (a) and running total (b) for the thermal fission of  $^{235}\text{U}$  calculated using Geant4. The current recommended values based on the  $^{235}\text{U}$  experimental work of Baumann and others are presented under the label “Baumann”.



# Conclusions

- Large discrepancies between the simulation results and the experimentally derived data were found, most significantly for the shortest and longest half-life groups.
- These results are **inconsistent** with the (“stitched together”) experimentally derived group data.

*For example, significantly more yield was found in simulation for the 55.7 sec half-life group. We can perform a sanity check!*

## Major contributor of this group: $^{91}\text{Rb}$ , using the

- Fission yield [JENDL-FPY 2011]
- Photoneutron production yield [MCNP6]
- $\gamma$ -ray energies from nuclear data [Nuclear Data Sheets 114 (10) (2013)]
- Total neutron yield of  $\bar{\nu} = 2.437 \pm 0.017$  [ENDF/B-VII.1]

We calculate a nuclide specific photoneutron yield of  $\beta_{\gamma}^{91\text{Rb}} = (4.0 \pm 0.2) \times 10^{-5}$

Baumann reports  $3 \times 10^{-5}$  for the entire group, under-predicting its importance.

Other contributors with half-lives near 55.7 seconds:  $^{86,87}\text{Br}$ ,  $^{139}\text{Xe}$ ,  $^{140}\text{Cs}$ , and  $^{144}\text{La}$

E.T. Rand, J.E. Atfield, L.R. Yaraskavitch, "Microscopic calculation of delayed-photoneutron production in  $D_2O$  using Geant4", Annals of Nuclear Energy **129**, 390-398 (2019)



# Limitations and Future Improvement

- 1.  $\beta$ -delayed neutron decay physics in Geant4.**
  - Current status? Nuclear data driven or model driven?
- 2. Large discrepancies are observed in  $(\gamma, n)$  interactions with nuclei  $< 30$  MeV. Cross section on deuterium was found to be zero, thus we ported MCNP6 physics into our simulation.**
  - A nuclear data driven photonuclear reaction model for Geant4 was developed by others (J. W. Shin, Nucl Instrum Methods Phys Res B 358, 2015).
  - Possibility to validate and incorporate this reaction model?
- 3. Absolute  $\gamma$ -ray yields from the Geant4 radioactive process deviate from literature values for many isotopes.**
  - There is no “easy fix”, left to the users to verify and check.
  - Using the per-decay sampling approach has its advantages (physically correct, event particle coincidences, angular correlations, etc.) However, there exists a community of users who only need to reproduce absolute  $\gamma$ -ray intensities. Currently users can hack radioactive decay and photonevaporation data files, write their own generator, or use GPS.
  - A supported and simplistic  $\gamma$ -decay model would meet the needs of these users and help identify issues in the radioactive decay module.







# Thank you!

## Questions?

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Applied Physics Branch  
Canadian Nuclear Laboratories  
[evan.rand@cnl.ca](mailto:evan.rand@cnl.ca)



Canadian Nuclear Laboratories | Laboratoires Nucléaires Canadiens

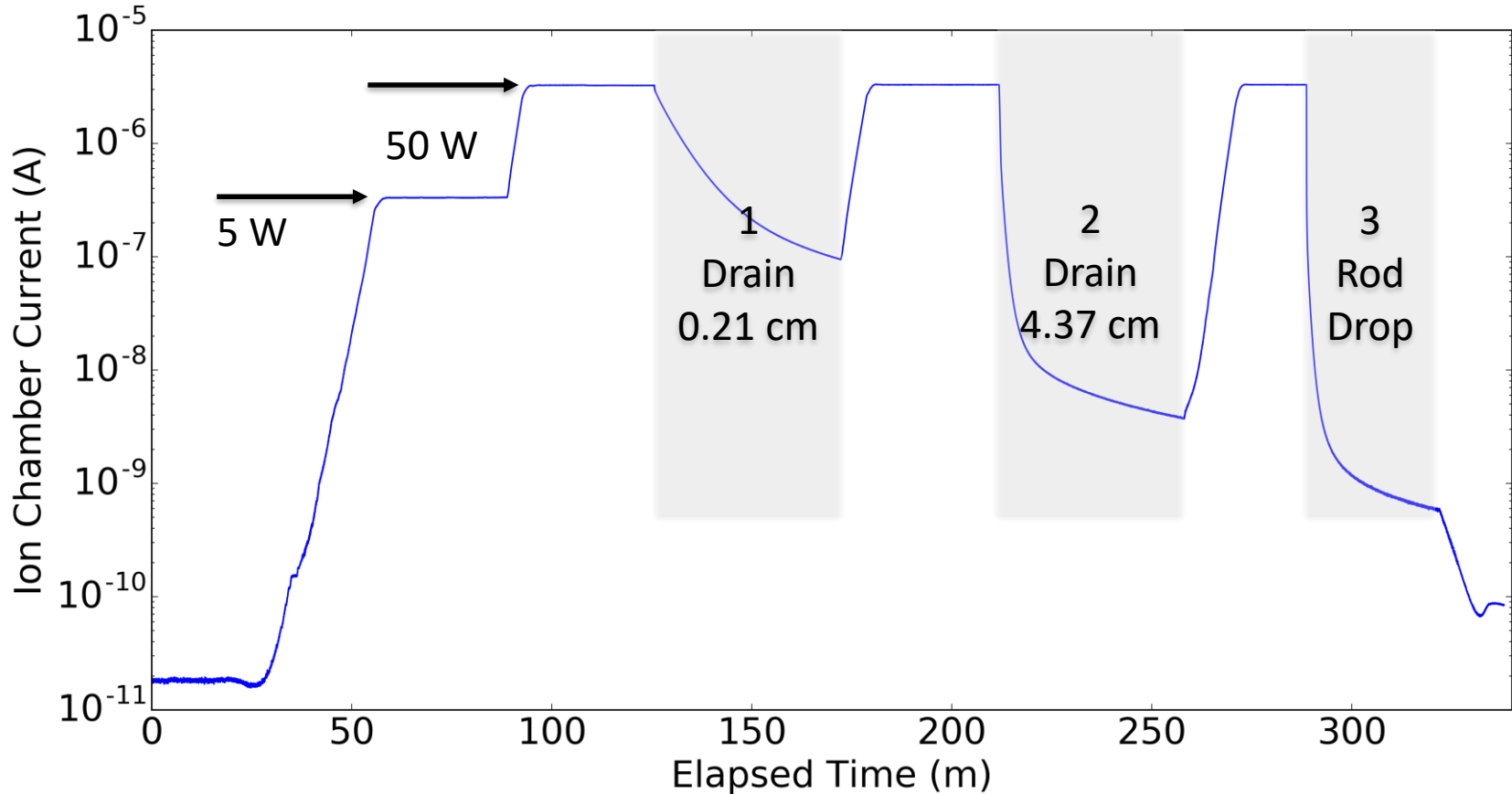
A grayscale, slightly blurred image of a microscope's objective lenses and eyepiece, serving as a background for the slide title.

# Backup Slides



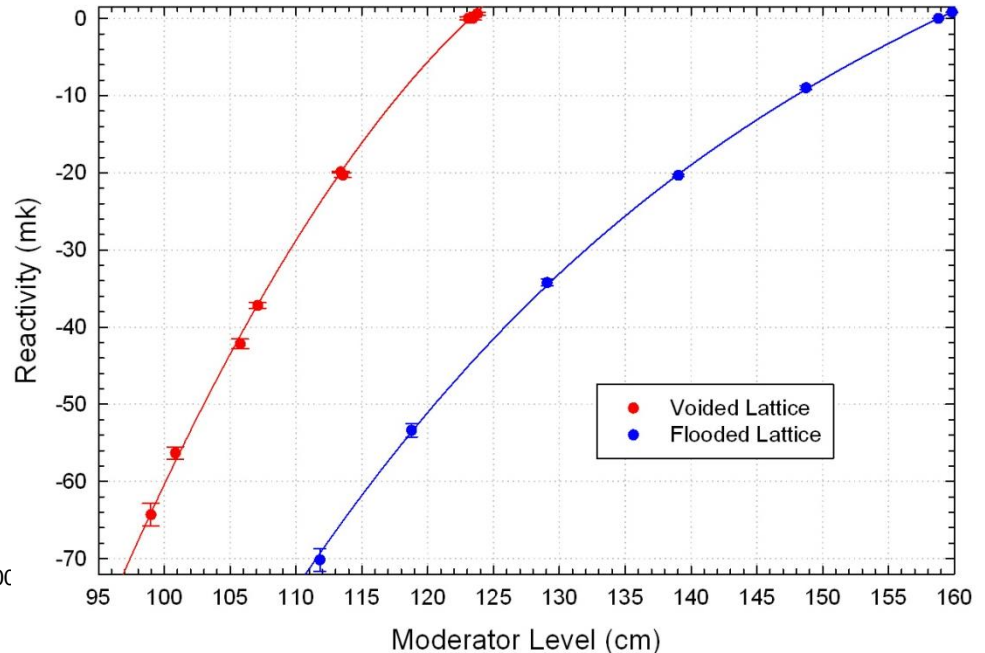
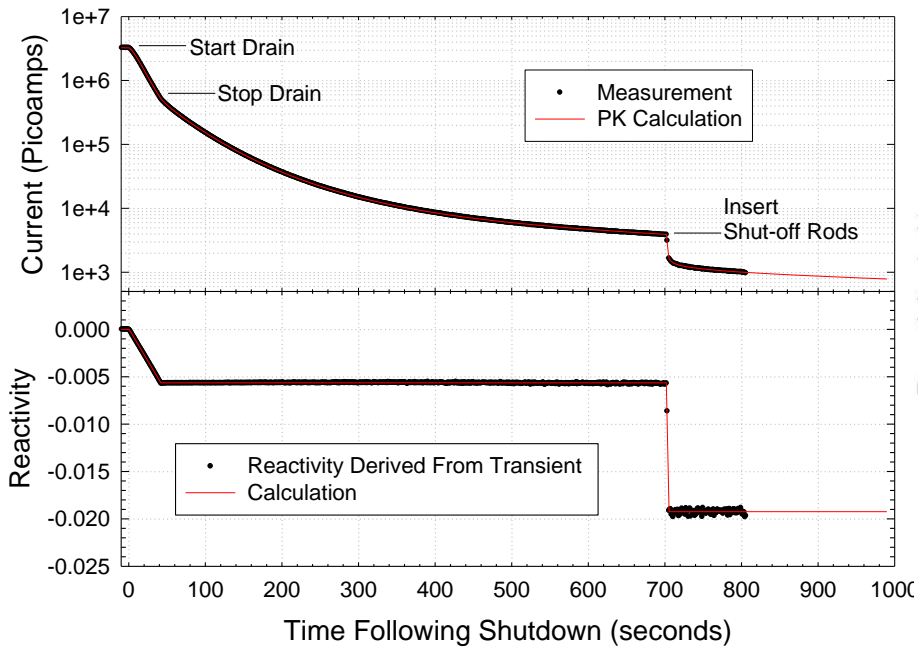
# Pu-Th Substitution Core

- Pu-Th Test Core



# Existing Capabilities

- Direct Measurement of Reactivity Effect on Voiding



Zeller and Atfield, "Transient experiments in ZED-2 to investigate the impact of leakage on reactor physics phenomena", AECL Nuclear Review, 4 (2015) 67-73.

**Other recent work:** Atfield, Yaraskavitch, and Rand, "Kinetics experiments in ZED-2 using heterogeneous cores of advanced nuclear fuels", Annals of Nuclear Energy, 121 (2018) 36-49.