

Spectrally novel FF sum rules and dynamical generation of mass

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*Based on: Accardi, Signori, arXiv:1903.04458
Accardi, Bacchetta, PLB 773 (2017) 632
+ in progress w/ Bacchetta, Radici, Signori*

Overview

□ “Inclusive jet” correlator

- Quarks are not asymptotic states

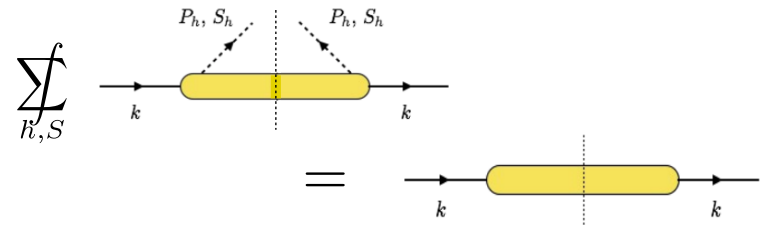


□ Gauge invariant spectral representation

- jet/dressed quark mass

□ New FF sum rules

- Jet/ dressed quark mass is experimentally observable!



□ New phenomenology

□ Conclusions

Inclusive jet correlator

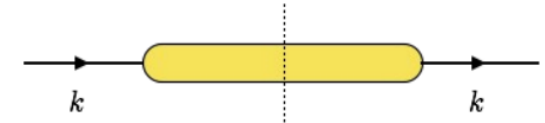
AA, Signori, 1903.04458
Sterman, NPB 281 ('87)

Inclusive $q \rightarrow X$ “inclusive jet” correlator

$$\Xi_{ij}(k; n_+) = \text{Disc} \int \frac{d^4\xi}{(2\pi)^4} e^{ik \cdot \xi} \frac{\text{Tr}_c}{N_c} \langle \Omega | \mathcal{T} W_{(\infty, \xi)}^{n_+} \psi_i(\xi) \bar{\psi}_j(0) W_{(0, \infty)}^{n_+} | \Omega \rangle$$

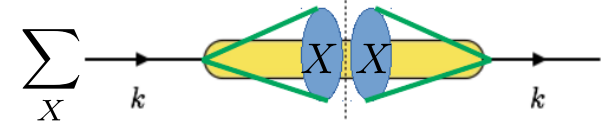
□ Partonic picture: **gauge invariant dressed quark correlator**

- Quarks are not asymptotic states
- Note color averaging



□ Hadronic picture: **“inclusive jet” correlator**

- Hadronization products pass the cut
- Interpret as a gauge invariant quark-to-jet amplitude squared
- No measured hadrons \rightarrow no jet cone / energy



□ Can study **fragmentation w/o fragments**

- In particular, dynamical mass generation & χ -symmetry breaking

Gauge invariant spectral representation

AA, Signori, 1903.04458

□ First: convolution representation

$$\Xi_{ij}(k) = \text{Disc} \int d^4p \frac{\text{Tr}_c}{N_c} \langle \Omega | \tilde{S}_{ij}(p) \widetilde{W}(k-p) | \Omega \rangle ,$$

where

$$\tilde{S}_{ij}(p) = \int \frac{d^4\xi}{(2\pi)^4} e^{i\xi \cdot p} \mathcal{T} \psi_i(\xi) \bar{\psi}_j(0) ,$$

$$\widetilde{W}(k-p) = \int \frac{d^4\xi}{(2\pi)^4} e^{i\xi \cdot (k-p)} \mathcal{T} W(0, \xi) .$$

□ Invariant decomposition of quark's bilinear operator:

$$\tilde{S}_{ij}(p) = \hat{s}_3(p^2) \not{p}_{ij} + \sqrt{p^2} \hat{s}_1(p^2) \mathbb{I}_{ij} + g \cdot f \cdot \not{n}$$

“Spectral operators”

gauge fixing term
(for axial gauges)

Gauge invariant spectral representation

AA, Signori, 1903.04458

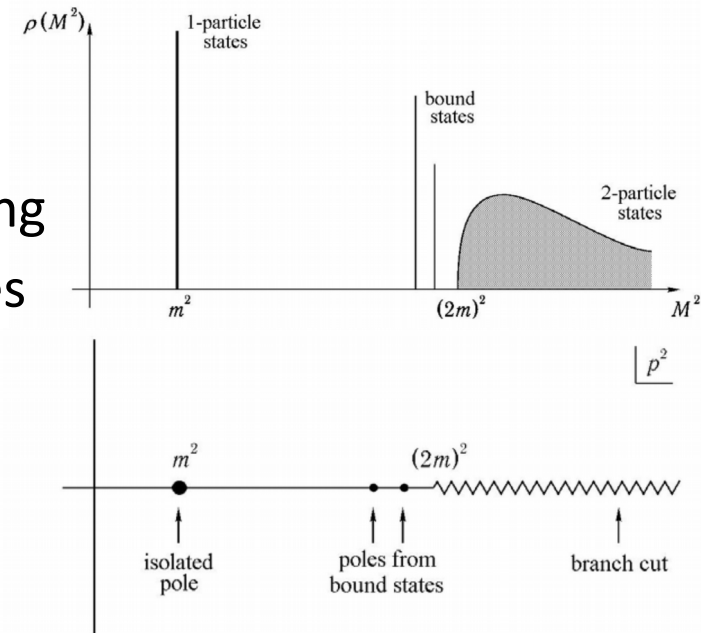
□ Kallen-Lehman representation for Feynman propagator

$$\frac{\text{Tr}_c}{N_c} \langle \Omega | \tilde{S}(p) | \Omega \rangle = \int_{-\infty}^{+\infty} \frac{d\mu^2}{(2\pi)^4} \frac{\mathbf{i}}{p^2 - \mu^2 + i\epsilon} \left\{ \not{p} \rho_3(\mu^2) + \sqrt{\mu^2} \rho_1(\mu^2) \right\} \theta(\mu^2)$$

$\rho_{1,3}$ are spectral functions:

→ strength of quark-to-multihadron coupling

→ color averaging: only colorless final states



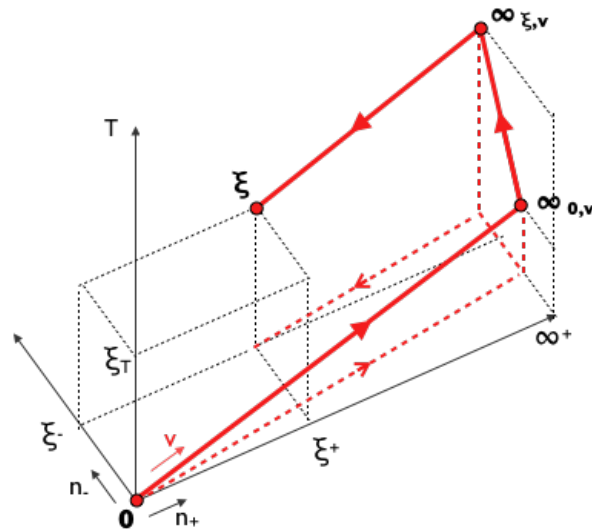
□ In terms of **spectral propagators**:

$$(2\pi)^3 \text{Disc} \frac{\text{Tr}_c}{N_c} \langle \Omega | \hat{s}_{1,3}(p^2) | \Omega \rangle = \rho_{1,3}(p^2) \theta(p^2) \theta(p^-)$$

Wilson line structure

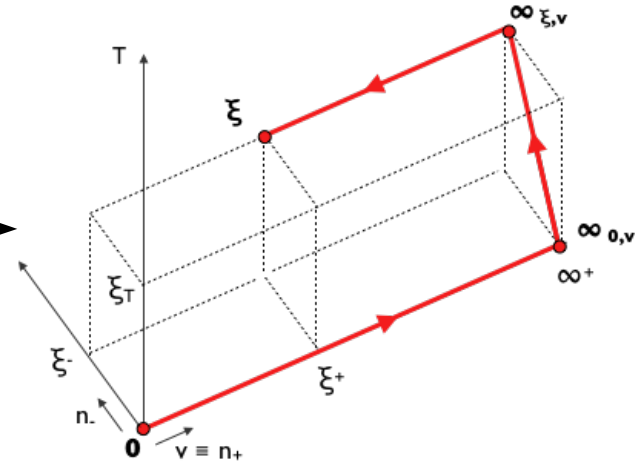
AA, Signori, 1903.04458

- Focus on (l.c.) staple-like Wilson lines
 - But spectral convolution method is general



Generic v

Light-cone $v = n_+$



TMD jet correlator

AA, Signori, 1903.04458

- Boost the quark at large light-cone momentum
 - (e.g., as it happens in large- Q DIS)

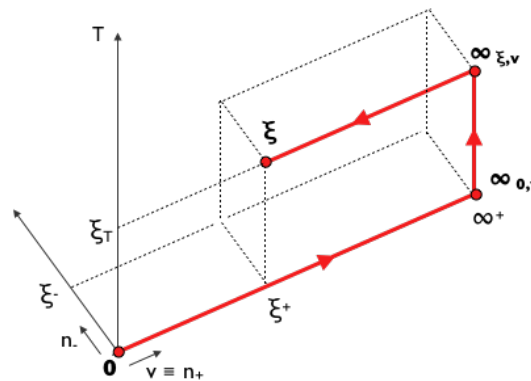
$$k^- \gg |\mathbf{k}_T| \gg k^+$$

- Integrate out the small momentum component:

TMD Inclusive jet correlator

$$J_{ij}(k^-, \mathbf{k}_T) \equiv \frac{1}{2} \int dk^+ \Xi_{ij}(k),$$

obtain standard staple,



TMD jet correlator in full glory

AA, Signori, 1903.04458
+ work in progress

- Expand in Dirac structures, order in powers of $1/k^-$

$$J(k^-, k_T) = \left\{ \gamma^+ + \frac{M_j}{k^-} + \frac{\not{k}_T}{k^-} + \frac{(K_j^2 + T_j^2 + g.f.) + k_T^2}{2(k^-)^2} \not{n}_- \right\} \theta(k^-)$$

- where, using spectral convolution in light-cone gauge,

$$M_j = \int_0^\infty d\mu^2 \mu \rho_1(\mu^2)$$

Jet “mass” \sim dressed quark mass
 $\sim O(100 \text{ MeV})$

$$K_j^2 = \int_0^\infty d\mu^2 \mu^2 \rho_3(\mu^2)$$

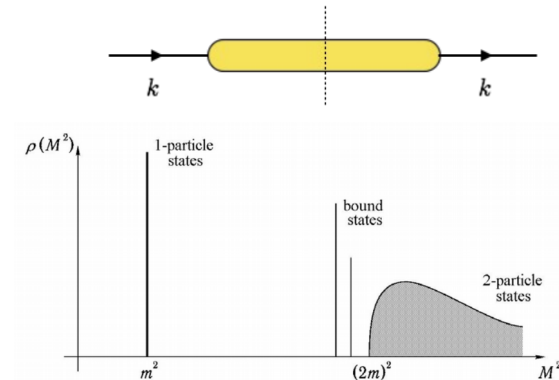
Jet’s “virtuality”

$$T_j^2 \sim \langle \langle \mu_T^2 \rangle \rangle$$

Jet’s “transverse size”

NOTE:

- Average jet shape (dynamics of hadronization) encoded in TMD jet correlator !!
- Explicit g.f. contributions pushed to twist-4 in light-cone gauge



The jet/quark mass

AA, Signori, 1903.04458

Mass associated with **chiral-odd component** of jet amplitude squared:

$$M_{jet} \sim \frac{\text{Tr}_c}{N_c} \int dk^+ \text{Tr}_D \left(\begin{array}{c} + \\ \longrightarrow \\ k \end{array} \text{---} \text{---} \text{---} \begin{array}{c} \text{---} \text{---} \text{---} \\ | \\ \text{---} \text{---} \text{---} \end{array} \begin{array}{c} \text{---} \text{---} \text{---} \\ | \\ \text{---} \text{---} \text{---} \end{array} \begin{array}{c} \longrightarrow \\ - \\ k \end{array} \right)$$

- In light cone gauge, it is interpreted as average mass of the color-neutral QCD d.o.f (“hadrons”) through cut

$$M_{jet} = \int_0^\infty d\mu^2 \mu \rho_1(\mu^2)$$

- Provides definition for **mass of a colored-screened dressed quark**, which is:

- Gauge-invariant
- Renormalization-scale dependent (grows with jet energy scale k^-)
- Calculable theoretically (through spectral functions)
- **Most importantly, measurable via a new momentum FF sum rule**

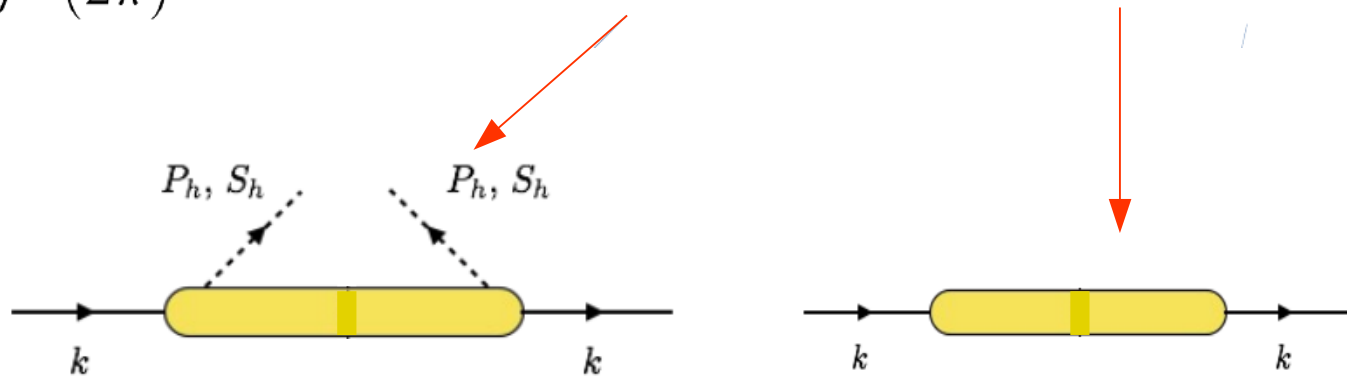
Momentum sum rule - operator level

AA, Signori, 1903.04458

+ work in progress

- Extend field-theoretical technique of *Meissner, Metz, Pitonyak, PLB 2010* :

$$\sum_{h, S_h} \int \frac{d^4 P_h}{(2\pi)^3} \delta(P_h^2 - M_h^2) P_h^\mu \Delta^h(k, P_h, S_h) = k^\mu \Xi^{\text{uncut}}(k)$$

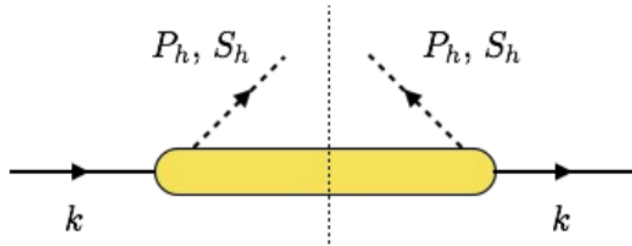


- **Dressed quark propagator as “average” on-shell four momentum produced by hadronization**
- **Dirac projections give momentum sum rules for TMD FFs !**

Dirac structures

AA, Signori, 1903.04458

□ TMD Fragmentation Functions



$$\Delta^h(z, P_{h\perp}) = \frac{1}{2}\gamma^+ D_1^h + \frac{M_h}{2P_h^-} E^h \mathbb{I} + \frac{\not{P}_{h\perp}}{2zP_h^-} D^{\perp h} + \text{quark polarized terms}$$

□ For the inclusive jet correlator



$$J(k^-, k_T) = \frac{1}{2}\gamma^+ + \frac{M_j}{2k^-} \mathbb{I} + \frac{\not{k}_T}{2k^-} + \text{higher-twist terms}$$

Mass sum rule

AA, Signori, 1903.04458

- Projecting the sum rule onto the identity matrix,

$$M_j = \sum_{h, S_h} \int dz M_h E^h(z)$$

jet/quark mass as
**average of produced hadron masses
weighted by chiral-odd FFs**

- **Dynamical mass component:**

EOM relations:

$$E = \tilde{E} + z \frac{m_q}{M_h} D_1$$

q-g-q correlations
neglecting

WW approx.

$$M_j = m_q$$



full QCD

**Dynamical
mass!**

$$M_j \equiv m_q + m^{corr}$$



$$m^{corr} = \sum_{h, S_h} \int dz M_h \tilde{E}^h(z)$$

Expect non-zero in χ -limit \rightarrow **observable χ -symmetry order parameter!**

Full set of sum rules

AA, Signori, PoS(DIS2018)

+ in progress

□ Sum rules for quarks into unpolarized hadrons, up to twist-3

- (only thing missing for twist-4: full FF-TMD analysis)

$$\sum_h \int dz z D_1^h(z) = 1$$

Collins-Soper

NEW

$$\sum_h \int dz M_h E^h(z) = M_j$$

NEW

$$\sum_h \int dz M_h \tilde{E}^h(z) = M_j - m_{q0} = m_q^{corr}$$

NEW

$$\sum_h \int dz M_h H^h(z) = 0$$

$$\sum_h \int dz M_h \tilde{H}^h(z) = 0$$

Diehl-Sapeta

fully dynamical
quantities

$$\sum_h \int dz z M_h H_1^{\perp(1)h}(z) = 0$$

Schaefer-Teryaev

NEW

$$\sum_h \int dz M_h^2 D^{\perp(1)h}(z) = 0$$

NEW

$$\sum_h \int dz M_h^2 \tilde{D}^{\perp(1)h}(z) = -\frac{1}{2} \langle P_{\perp}^2 / z \rangle$$

NEW

$$\sum_h \int dz M_h^2 G^{\perp(1)h}(z) = 0$$

NEW

$$\sum_h \int dz M_h^2 \tilde{G}^{\perp(1)h}(z) = 0 .$$

Conclusions

- We can quantitatively connect quark fragmentation to the dynamical generation of mass
 - **Gauge invariant definition for dressed quark mass, M_j**
 - The dynamical component $m^{corr} = M_j - m_q$ is recognized as an **observable order parameter for χ -simmetry breaking**

$$m^{corr} = \sum_{h, S_h} \int dz M_h \tilde{E}^h(z)$$

Practical exp. recipe:

- measure \tilde{E} , obtain m^{corr}
- flavor decomposition, too!

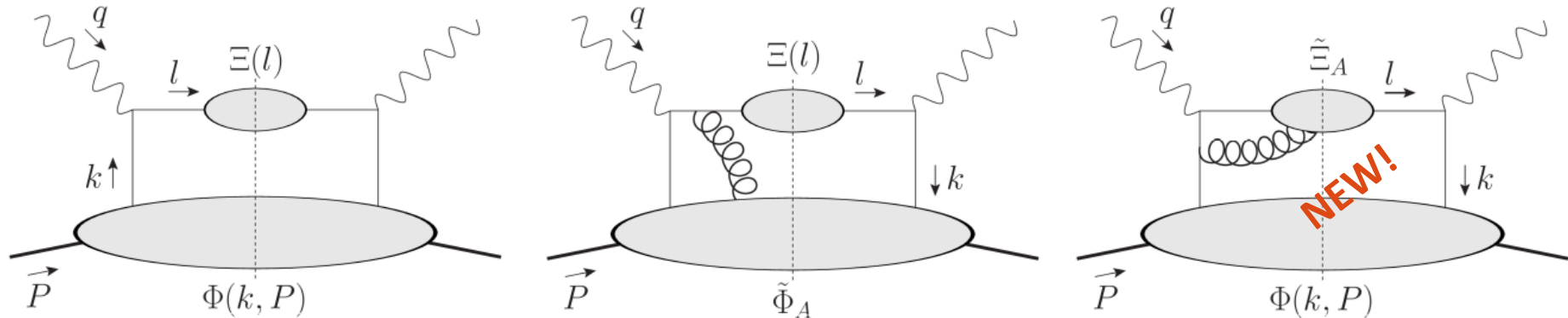
- **Novel phenomenology:**
 - Transversity in g2, same side di-hadrons, ...
- **New sum rules:** guidance for future fits

- **New spectral convolution technique** for treating Wilson lines

Backup

Inclusive DIS with jet correlators

AA, Bacchetta, PLB 773 ('17) 632



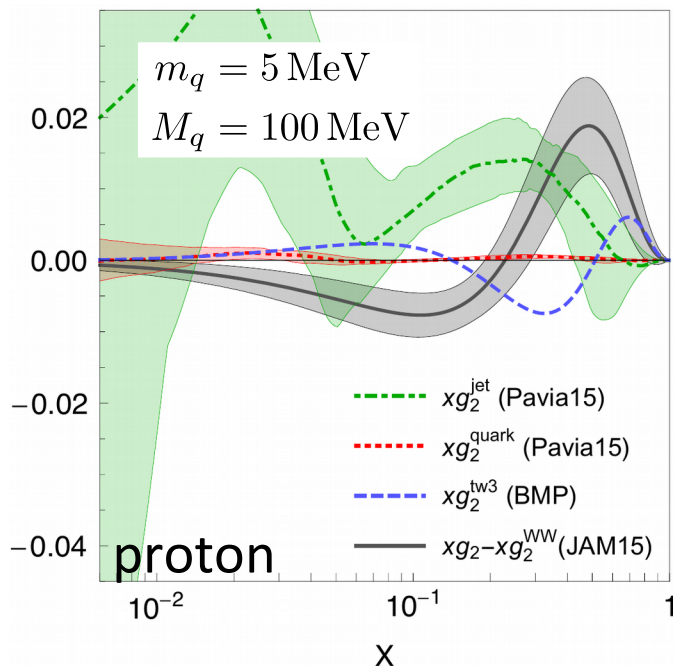
Jet correlators: \rightarrow non-asymptotic quark states / dressed quarks

$$\begin{aligned} & \text{Diagram 1: } l \rightarrow \text{shaded oval} \rightarrow \Xi_{ij}(l, n_+) = F.T. \langle \Omega | W_{(+\infty, \xi)}^{n_+} \psi_i(\xi) \bar{\psi}_j(0) W_{(0, +\infty)}^{n_+} | \Omega \rangle \\ & \text{Diagram 2: } l \rightarrow \text{shaded oval} \rightarrow (\Xi_A^\mu)_{ij} = F.T. \langle \Omega | W_{(+\infty, \xi)}^{n_+} g A^\mu(\xi) \psi_i(\xi) \bar{\psi}_j(0) W_{(0, +\infty)}^{n_+} | \Omega \rangle \end{aligned}$$

g2 structure function revisited

□ Integrating SIDIS, and using EOM, Lorentz Invariance Relations:

$$\begin{aligned}
 g_2(x_B) - g_2^{WW}(x_B) &\stackrel{\text{red}}{=} g_2^{\text{quark}} \\
 &= \frac{1}{2} \sum_a e_a^2 \left(g_2^{q,\text{tw}3}(x_B) + \frac{m_q}{M} \left(\frac{h_1^q}{x} \right)^* (x_B) + \boxed{\frac{M_j - m_q}{M} \frac{h_1^q(x_B)}{x_B}} \right) \stackrel{\text{red}}{=} g_2^{\text{jet}}
 \end{aligned}$$



Consequences:

- h1 accessible in inclusive DIS
↔ Potentially large signal
- Burkardt-Cottingham sum rule broken

$$\int_0^1 g_2(x) = (M_j - m_q) \int_0^1 dx \frac{h_1(x)}{x}$$

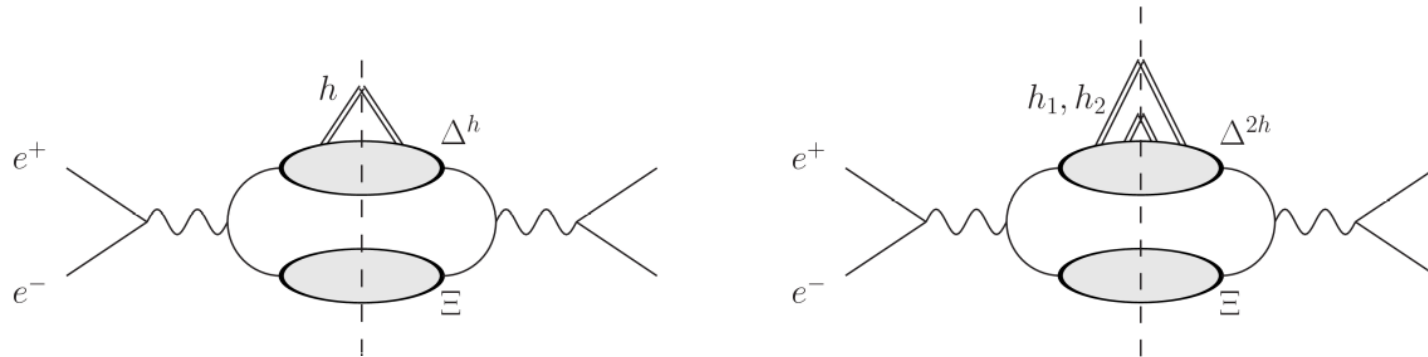
- ETL: novel way to measure tensor charge

$$\int_0^1 x g_2^{q-\bar{q}}(x) = 2 (M_j - m_q) \int_0^1 dx h_1^{q-\bar{q}}(x)$$

Measuring the jet correlator

Accardi, Bacchetta, Signori, Radici, in progress

□ Jet mass M_{jet} can be measured in polarized $e^+ + e^-$:



– Needs **LT asymmetry** in semi-inclusive **Lambda** production

$$\frac{d\sigma^L(e^+e^- \rightarrow \text{jet } h X)}{d\Omega dz} = \frac{3\alpha^2}{Q^2} \lambda_e \sum_a e_a^2 \left\{ \frac{C(y)}{2} \lambda_h G_1 + D(y) |S_T| \cos(\phi_S) \frac{2M_h}{Q} \left(\frac{G_T}{z} + \frac{M_q - m_q}{M_h} H_1 \right) \right\}$$

– Similarly a **LU asymmetry** in unpolarized **dihadron** production

Where can we measure jet correlators?

❑ Can we get a (polarized) $e^+ e^-$ collider at JLab / BNL?

– At JLab12 ? EIC + positron beam ?

❑ Are **existing facilities** enough?

	BEPC	super KEKB	ILC	JLab/BNL
E beam [GeV]	1.9	4 (e^-) 7 (e^-)	250	?
\sqrt{s} [GeV]	3 – 5	10	500	?
polarization	?	maybe	80% e^- 60% e^+	YES!

❑ What else?

A new “universal” fits

□ Chiral-odd collinear sector across processes:

