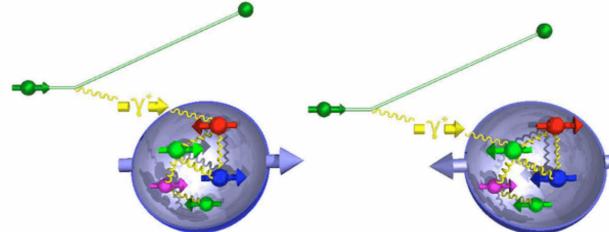
27<sup>th</sup> International Workshop on Deep-Inelastic Scattering and Related Subjects DIS 2019

8-12 April

Longitudinal double-spin asymmetries in semi-inclusive deep-inelastic scattering of electrons and positrons by protons and deuterons



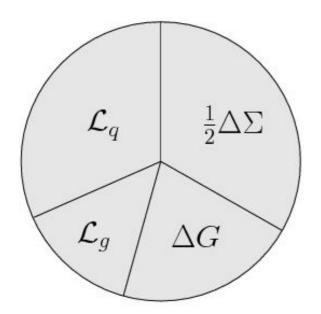
Gunar.Schnell @ DESY.de



Universidad del País Vasco

TORINC

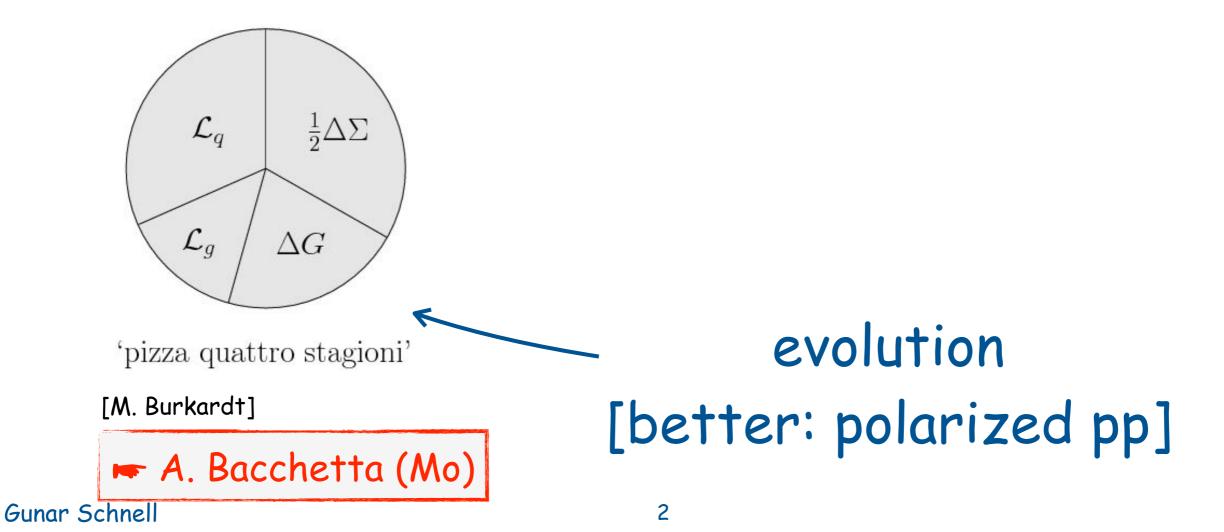
#### Access to angular momentum in (SI)DIS

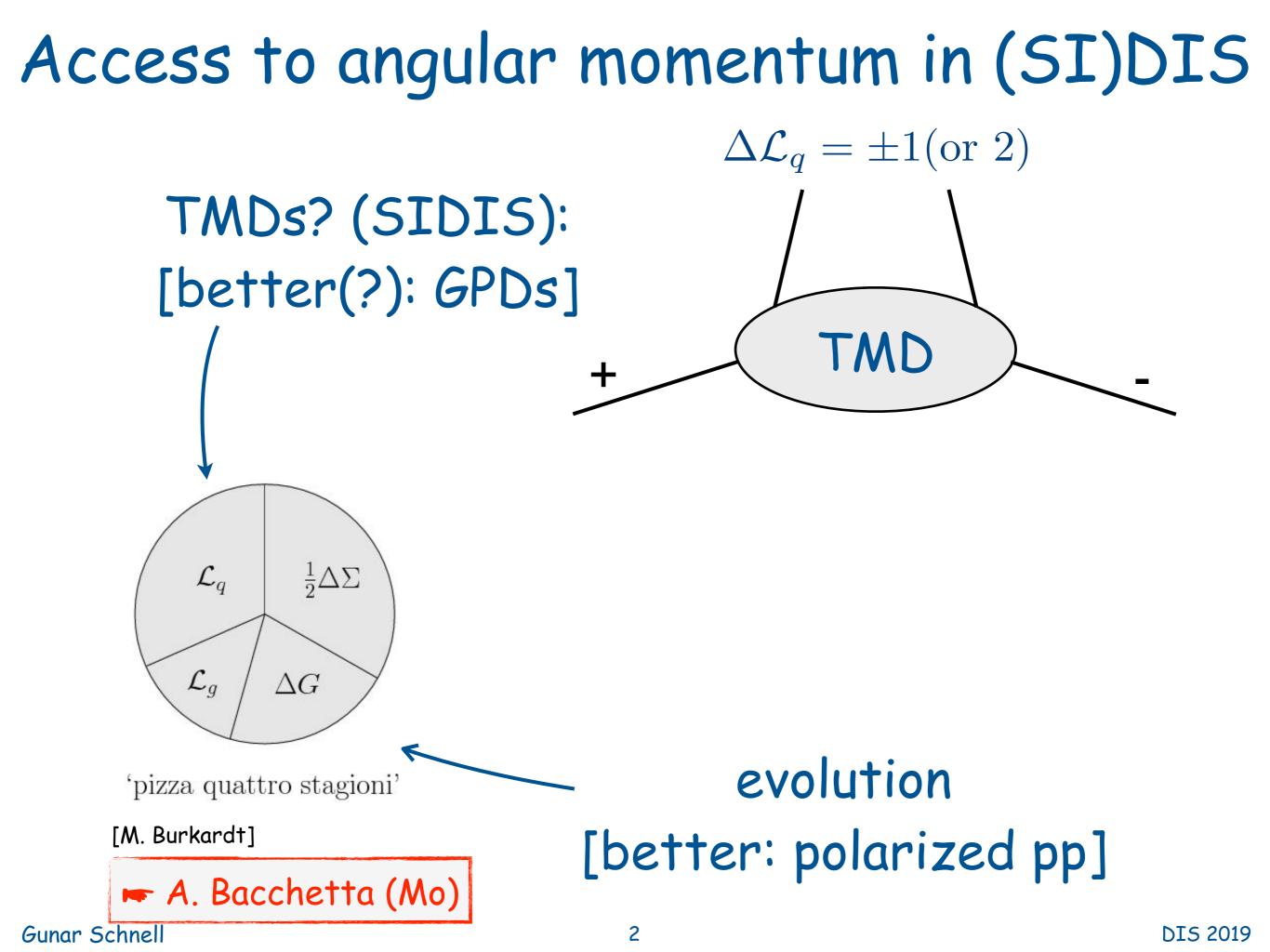


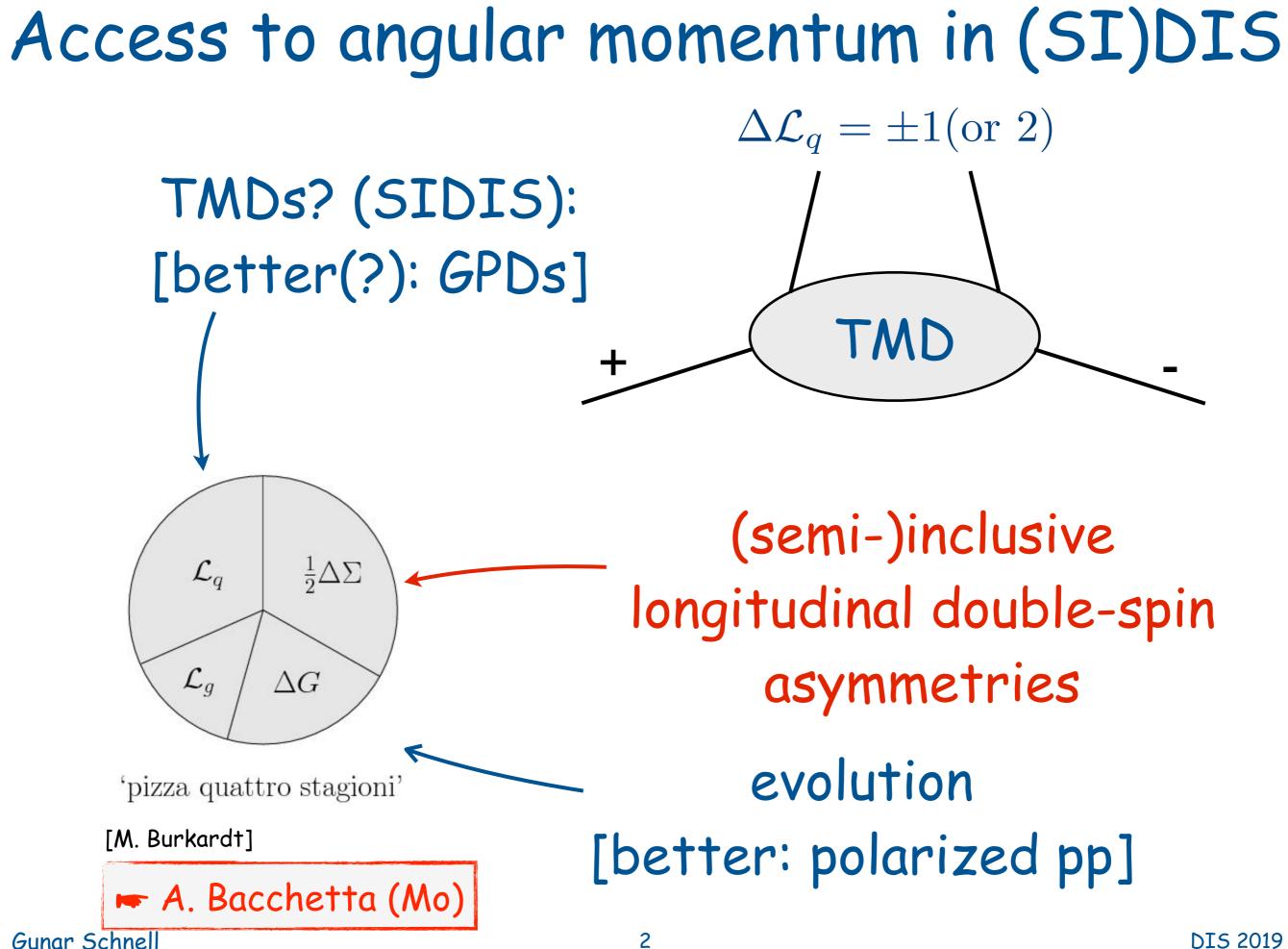
'pizza quattro stagioni'

[M. Burkardt]

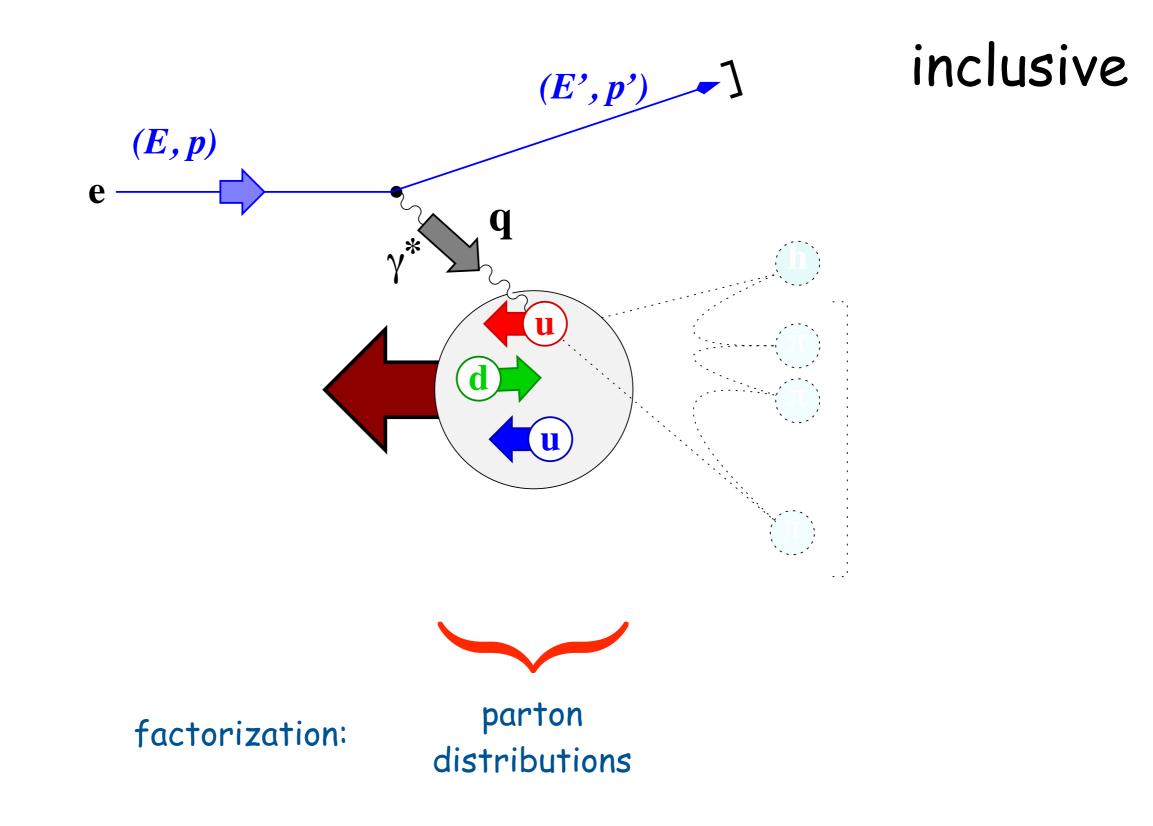
#### Access to angular momentum in (SI)DIS



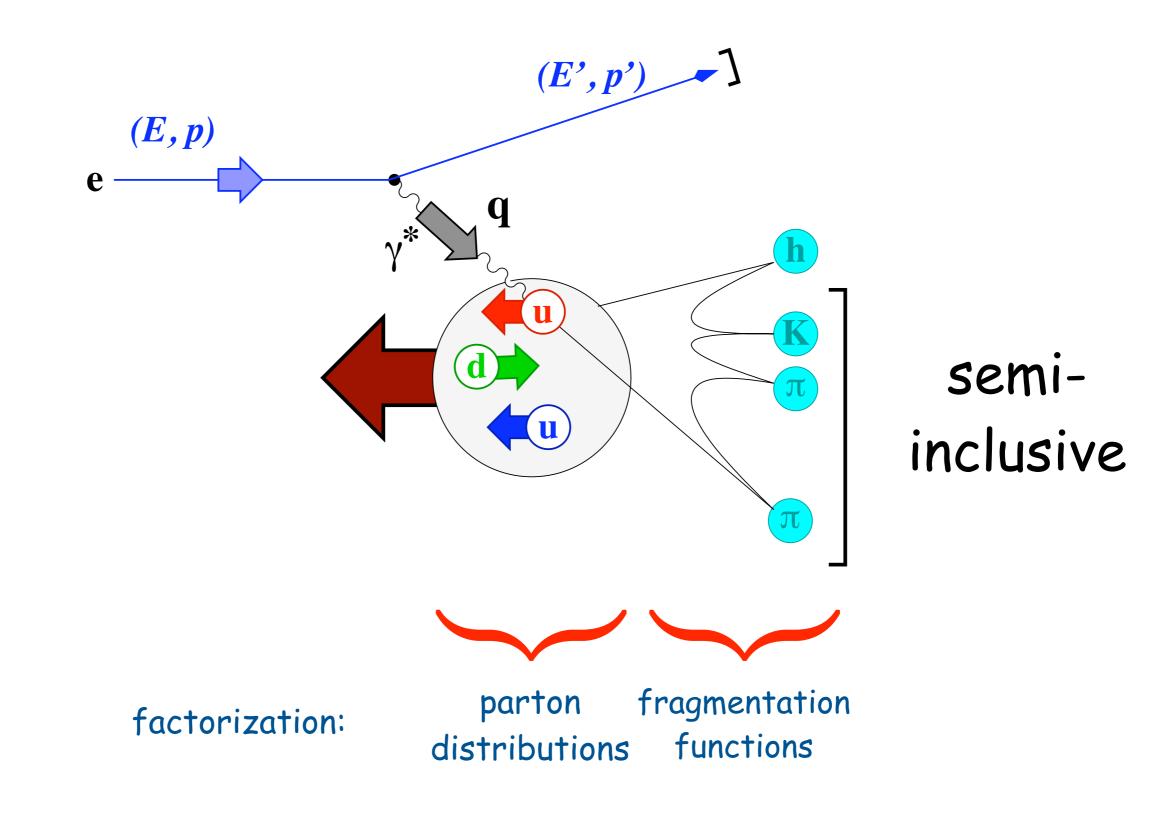




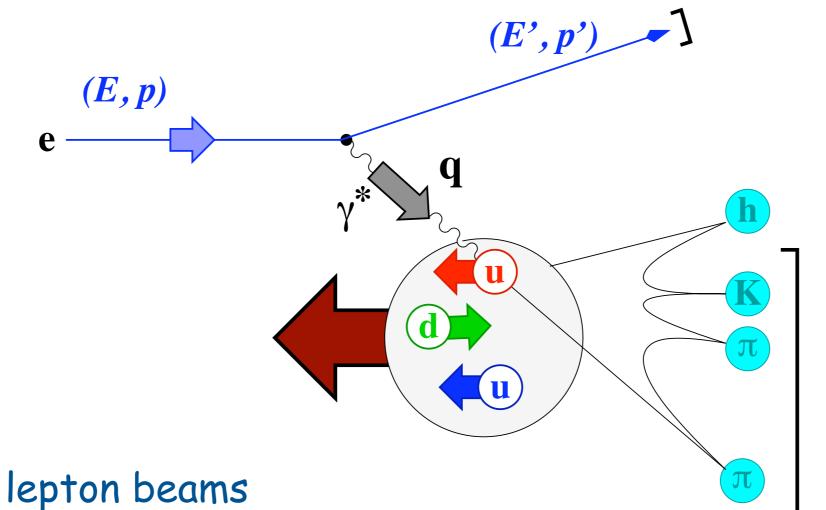
#### deep-inelastic scattering



### deep-inelastic scattering



### deep-inelastic scattering



semiinclusive

- polarized lepton beams
- polarized targets
- Iarge-acceptance spectrometer
- good particle identification (PID)

### HERMES (†2007) @ DESY

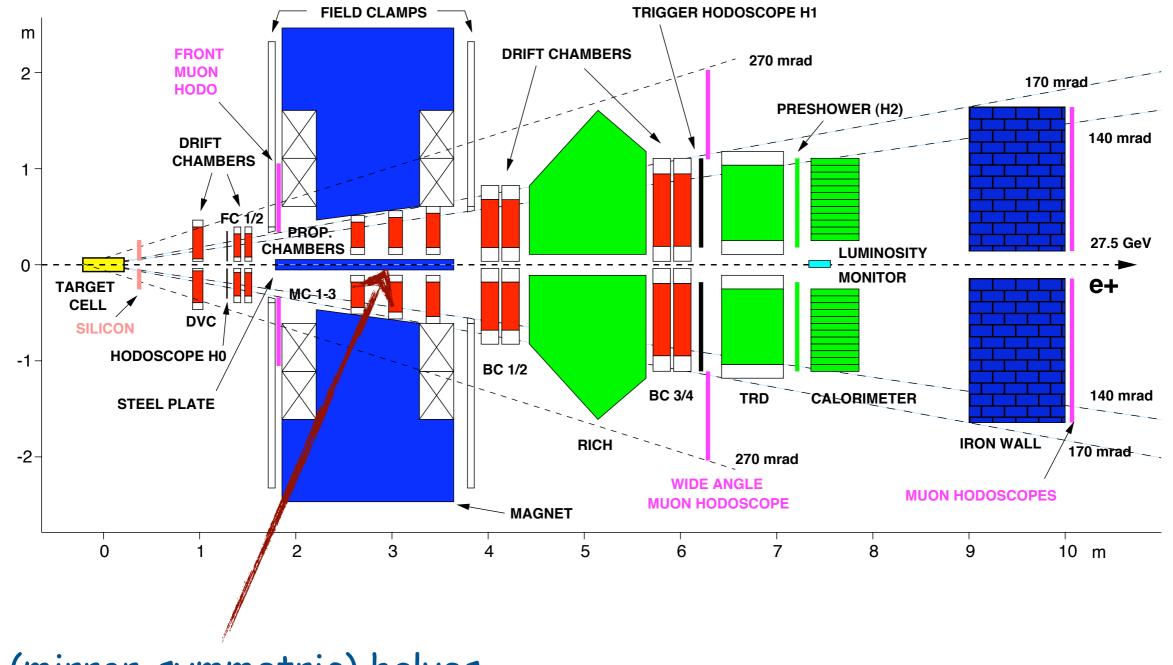
# 27.6 GeV **polarized** e<sup>+</sup>/e<sup>-</sup> **beam** scattered off ...



 unpolarized (H, D, He,..., Xe)
 as well as transversely (H) and longitudinally polarized (pure)
 H, D & <sup>3</sup>He gas targets

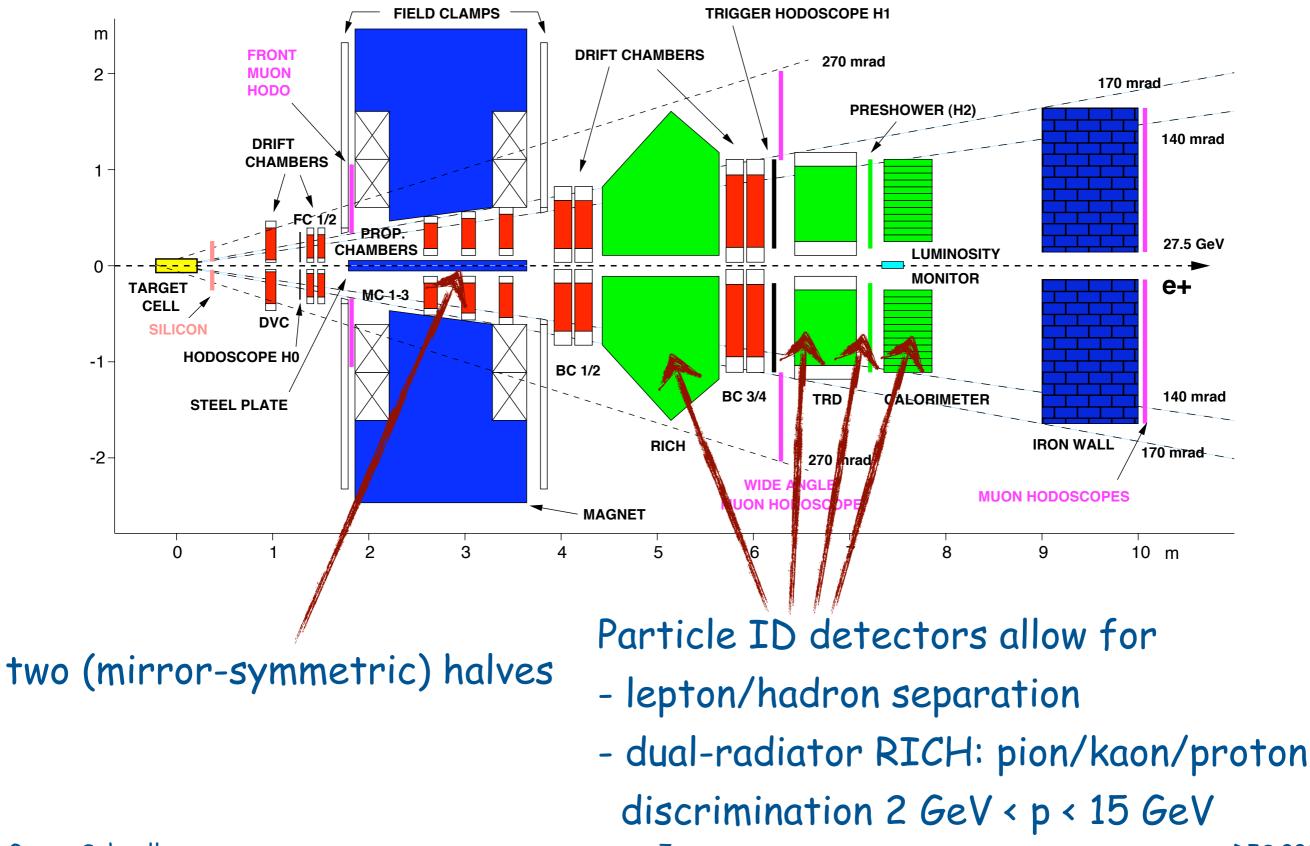


### HERMES (1998-2005) schematically



two (mirror-symmetric) halves

### HERMES (1998-2005) schematically



### semi-inclusive DIS

excluding transverse polarization:

$$\begin{aligned} \frac{\mathrm{d}\sigma^{h}}{\mathrm{d}x\,\mathrm{d}y\,\mathrm{d}z\,\mathrm{d}P_{h\perp}^{2}\,\mathrm{d}\phi} &= \frac{2\pi\alpha^{2}}{xyQ^{2}}\frac{y^{2}}{2(1-\epsilon)}\left(1+\frac{\gamma^{2}}{2x}\right)\\ \left\{F_{UU,T}^{h} + \epsilon F_{UU,L}^{h} + \lambda\Lambda\sqrt{1-\epsilon^{2}}F_{LL}^{h} \right.\\ &+ \sqrt{2\epsilon}\left[\lambda\sqrt{1-\epsilon}\,F_{LU}^{h,\sin\phi} + \Lambda\sqrt{1+\epsilon}\,F_{UL}^{h,\sin\phi}\right]\sin\phi\\ &+ \sqrt{2\epsilon}\left[\lambda\Lambda\sqrt{1-\epsilon}\,F_{LL}^{h,\cos\phi} + \sqrt{1+\epsilon}\,F_{UU}^{h,\cos\phi}\right]\cos\phi\\ &+ \Lambda\epsilon\,F_{UL}^{h,\sin2\phi}\sin2\phi + \epsilon\,F_{UU}^{h,\cos2\phi}\cos2\phi\right\}\\ \end{aligned}$$

Beam  $(\lambda)$  / Target  $(\Lambda)$ helicities

1′

qi

S

 $\mathbf{P}_{h\perp}$ 

 $\mathbf{S}_{\perp}$ 

 $\mathbf{P}_h$ 

φ

1

### semi-inclusive DIS

excluding transverse polarization:

$$\begin{aligned} \frac{\mathrm{d}\sigma^{h}}{\mathrm{d}x\,\mathrm{d}y\,\mathrm{d}z\,\mathrm{d}P_{h\perp}^{2}\,\mathrm{d}\phi} &= \frac{2\pi\alpha^{2}}{xyQ^{2}}\frac{y^{2}}{2(1-\epsilon)}\left(1+\frac{\gamma^{2}}{2x}\right)\\ F_{UU,T}^{h} + \epsilon F_{UU,L}^{h} + \lambda\Lambda\sqrt{1-\epsilon^{2}}F_{LL}^{h} \\ &+\sqrt{2\epsilon}\left[\lambda\sqrt{1-\epsilon}\,F_{LU}^{h,\sin\phi} + \Lambda\sqrt{1+\epsilon}\,F_{UL}^{h,\sin\phi}\right]\sin\phi\\ &+\sqrt{2\epsilon}\left[\lambda\Lambda\sqrt{1-\epsilon}\,F_{LL}^{h,\cos\phi} + \sqrt{1+\epsilon}\,F_{UU}^{h,\cos\phi}\right]\cos\phi\\ &+\Lambda\epsilon\,F_{UL}^{h,\sin2\phi}\sin2\phi + \epsilon\,F_{UU}^{h,\cos2\phi}\cos2\phi \end{aligned}$$

• double-spin asymmetry 
$$A_{LL}^h \equiv \frac{\sigma_{++}^h - \sigma_{+-}^h + \sigma_{--}^h - \sigma_{-+}^h}{\sigma_{++}^h + \sigma_{+-}^h + \sigma_{--}^h + \sigma_{-+}^h}$$

1′

a

S

 $\mathbf{P}_{h\perp}$ 

 $\mathbf{S}_{\perp}$ 

 $\mathbf{P}_h$ 

1

Gunar Schnell

φ

### semi-inclusive DIS

- in experiment extract instead A<sub>||</sub> which differs from A<sub>LL</sub> in the way the polarization is measured:
  - ALL: along virtual-photon direction
  - A<sub>II</sub>: along beam direction (results in small admixture of transverse target polarization and thus contributions from A<sub>LT</sub>)

1

 $\mathbf{S}$ 

 $\mathbf{P}_{h\perp}$ 

 $\mathbf{S}_{\perp}$ 

 $\mathbf{P}_h$ 

 $D = \frac{1 - (1 - y)\epsilon}{1 + \epsilon R}$ 

 $\eta = \frac{\epsilon \gamma y}{1 - (1 - u) \epsilon}$ 

DIS 2019

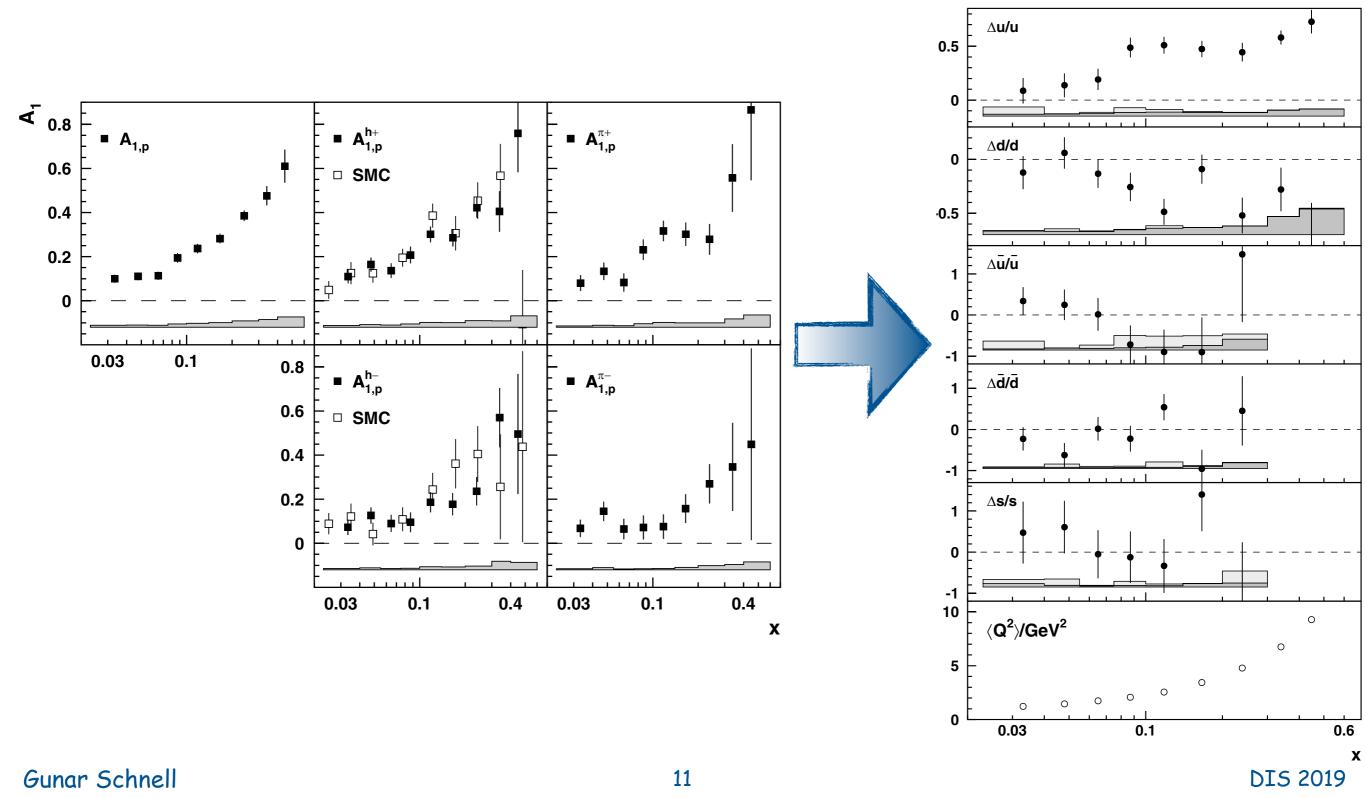
•  $A_{||}$  related to virtual-photon-nucleon asymmetry  $A_1$ 

$$A_1^h = \frac{1}{D(1+\eta\gamma)} A_{\parallel}^h$$

### previous HERMES analysis

(semi-) inclusive asymmetries used for LO extraction of helicity PDFs

PHYSICAL REVIEW D 71, 012003 (2005)



### re-analysis of double-spin asymmetries

- revisited [PRD 71 (2005) 012003] A1 analysis at HERMES in order to
  - exploit slightly larger data set (less restrictive momentum range)
  - provide  $A_{\parallel}$  in addition to  $A_1$

$$A_1^h = \frac{1}{D(1+\eta\gamma)} A_{\parallel}^h \qquad D = \frac{1-(1-y)\epsilon}{1+\epsilon R}$$

R (ratio of longitudinal-to-transverse cross-sec'n) still to be measured! [only available for inclusive DIS data, e.g., used in  $g_1$  SF measurements]

- correct for D-state admixture (deuteron case) on asymmetry level
- correct better for azimuthal asymmetries coupling to acceptance
- look at multi-dimensional (x, z,  $P_{h\perp}$ ) dependences
- extract twist-3 cosine modulations

#### azimuthal-asymmetry corrections

measured

"polarized Cahn" effect etc.~

 $\tilde{A}^{h}_{\parallel}(x,Q^{2},z,P_{h\perp}) = \frac{\int \mathrm{d}\phi \ \sigma^{h}_{\parallel}\left(x,Q^{2},z,P_{h\perp},\phi\right) \ \xi(\phi)}{\int \mathrm{d}\phi \ \sigma^{h}_{IIII}\left(x,Q^{2},z,P_{h\perp},\phi\right) \ \xi(\phi)}$ 

Boer-Mulders and Cahn effects etc.

- both numerator and in particular denominator  $\phi$  dependent
  - in theory integrated out
  - in praxis, detector acceptance also  $\phi$  dependent
  - convolution of physics & acceptance leads to bias in normalization of asymmetries

azimuthal

acceptance

### azimuthal-asymmetry corrections

measured

"polarized Cahn" effect etc.~

 $\tilde{A}^{h}_{\parallel}(x,Q^{2},z,P_{h\perp}) = \frac{\int \mathrm{d}\phi \ \sigma^{h}_{\parallel}\left(x,Q^{2},z,P_{h\perp},\phi\right) \ \xi(\phi)}{\int \mathrm{d}\phi \ \sigma^{h}_{UU}\left(x,Q^{2},z,P_{h\perp},\phi\right) \ \xi(\phi)}$ 

Boer-Mulders and Cahn effects etc.

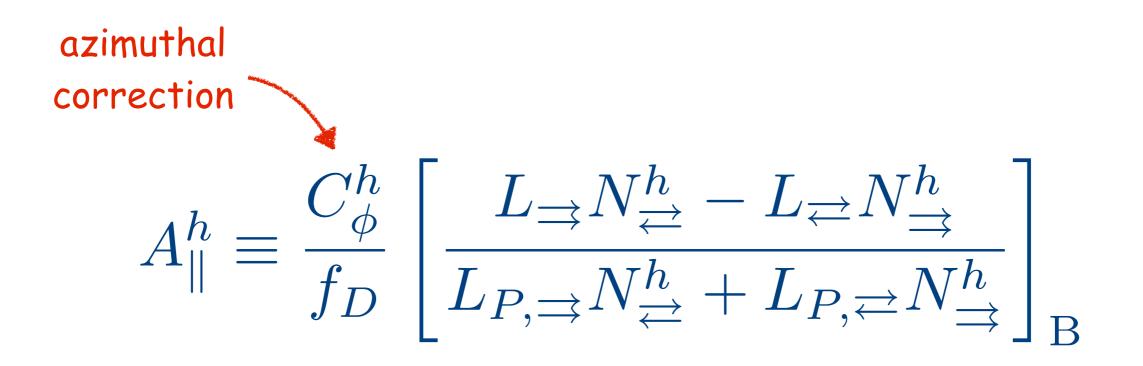
- both numerator and in particular denominator  $\phi$  dependent
  - in theory integrated out
  - in praxis, detector acceptance also  $\phi$  dependent
  - convolution of physics & acceptance leads to bias in normalization of asymmetries
- Implement data-driven model for azimuthal modulations [PRD 87 (2013) 012010] into MC rectract correction factor & apply to data

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azimuthal

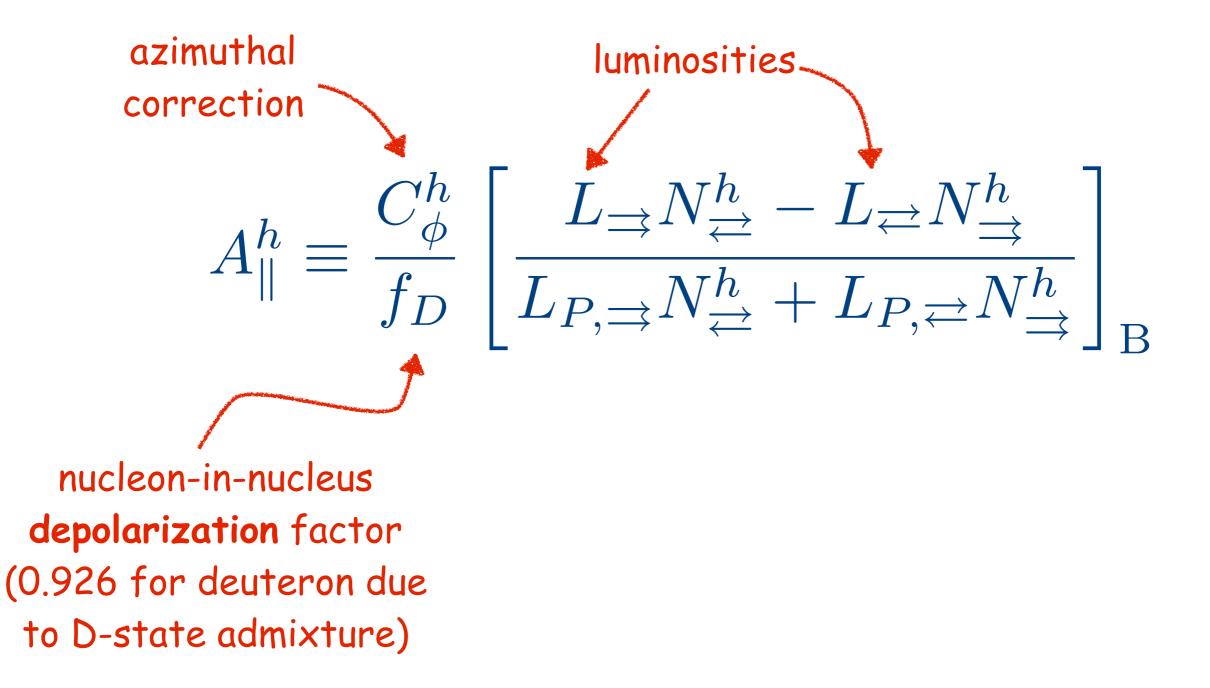
acceptance

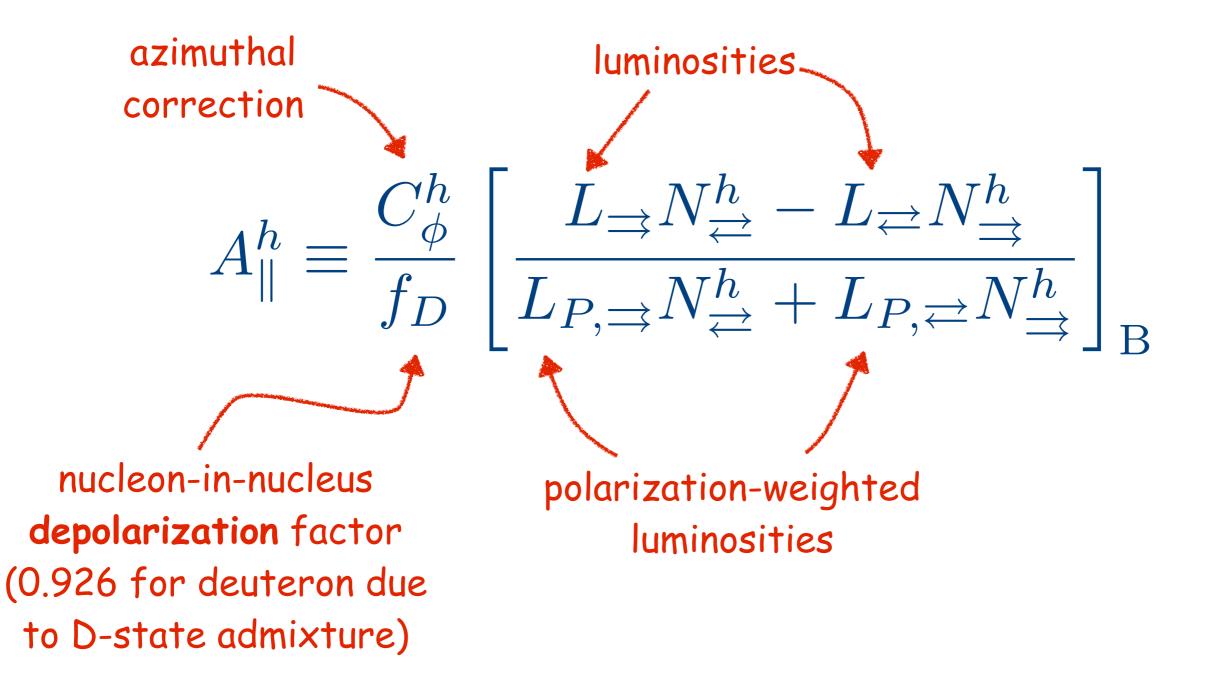
$$A^{h}_{\parallel} \equiv \frac{C^{h}_{\phi}}{f_{D}} \left[ \frac{L_{\rightrightarrows} N^{h}_{\rightleftharpoons} - L_{\rightleftharpoons} N^{h}_{\rightrightarrows}}{L_{P, \rightrightarrows} N^{h}_{\rightleftharpoons} + L_{P, \rightleftharpoons} N^{h}_{\rightrightarrows}} \right]_{B}$$

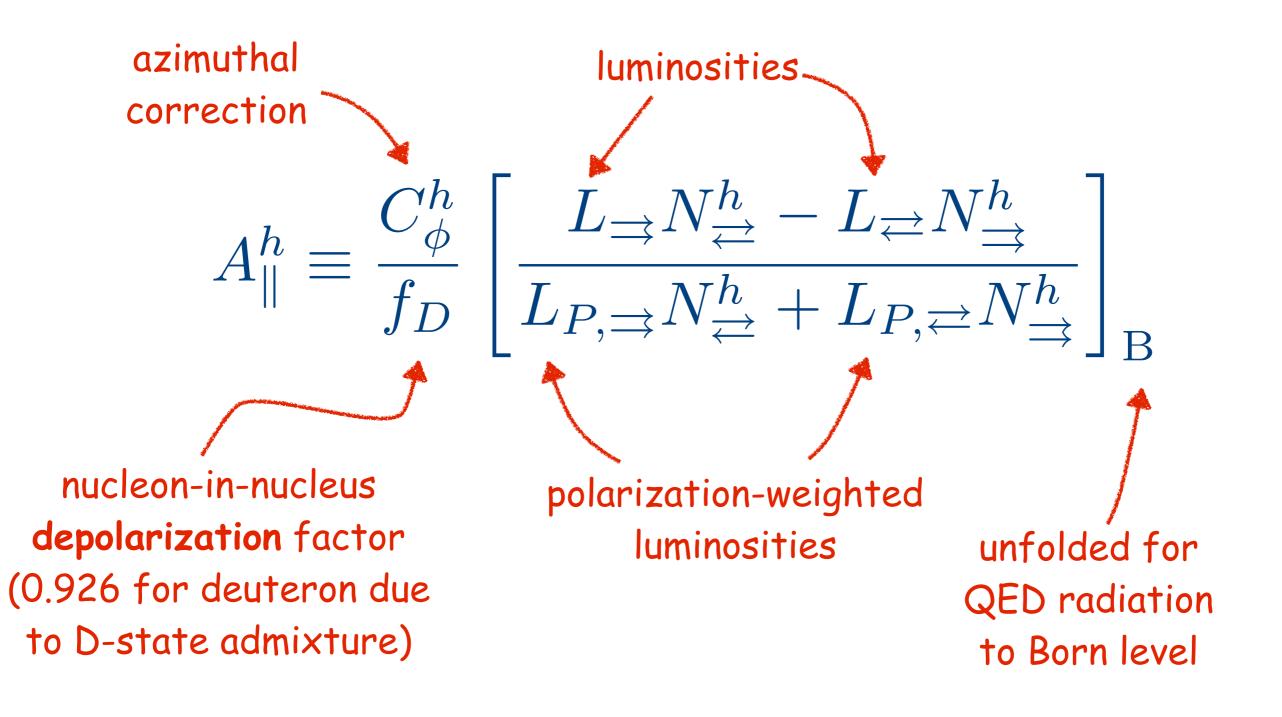


azimuthal  
correction
$$A_{\parallel}^{h} \equiv \frac{C_{\phi}^{h}}{f_{D}} \begin{bmatrix} L_{\Rightarrow} N_{\Rightarrow}^{h} - L_{\Rightarrow} N_{\Rightarrow}^{h} \\ L_{P,\Rightarrow} N_{\Rightarrow}^{h} + L_{P,\Rightarrow} N_{\Rightarrow}^{h} \end{bmatrix}_{B}$$
nucleon-in-nucleus
depolarization factor
(0.926 for deuteron due

to D-state admixture)







$$A^{h}_{\parallel} \equiv \frac{C^{h}_{\phi}}{f_{D}} \left[ \frac{L_{\rightrightarrows} N^{h}_{\rightleftharpoons} - L_{\rightleftharpoons} N^{h}_{\rightrightarrows}}{L_{P, \rightrightarrows} N^{h}_{\rightleftharpoons} + L_{P, \rightleftharpoons} N^{h}_{\rightrightarrows}} \right]_{B}$$

$$A^{h}_{\parallel} \equiv \frac{C^{h}_{\phi}}{f_{D}} \left[ \frac{L_{\rightrightarrows} N^{h}_{\rightleftharpoons} - L_{\rightleftharpoons} N^{h}_{\rightrightarrows}}{L_{P, \rightrightarrows} N^{h}_{\rightleftharpoons} + L_{P, \rightleftharpoons} N^{h}_{\rightrightarrows}} \right]_{B}$$

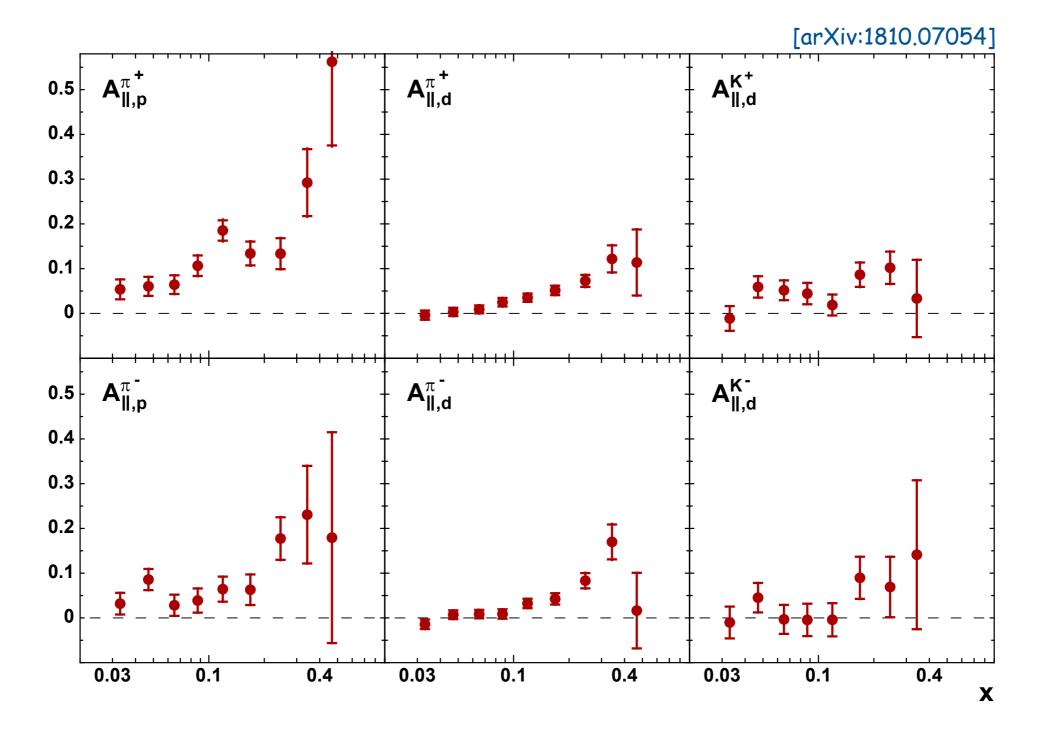
• dominated by statistical uncertainties

$$A^{h}_{\parallel} \equiv \frac{C^{h}_{\phi}}{f_{D}} \left[ \frac{L_{\Rightarrow} N^{h}_{\rightleftharpoons} - L_{\rightleftharpoons} N^{h}_{\Rightarrow}}{L_{P,\Rightarrow} N^{h}_{\rightleftharpoons} + L_{P,\rightleftharpoons} N^{h}_{\Rightarrow}} \right]_{B}$$

- dominated by statistical uncertainties
- main systematics arise from
  - polarization measurements [6.6% for hydrogen, 5.7% for deuterium)
  - azimuthal correction [O(few %)]

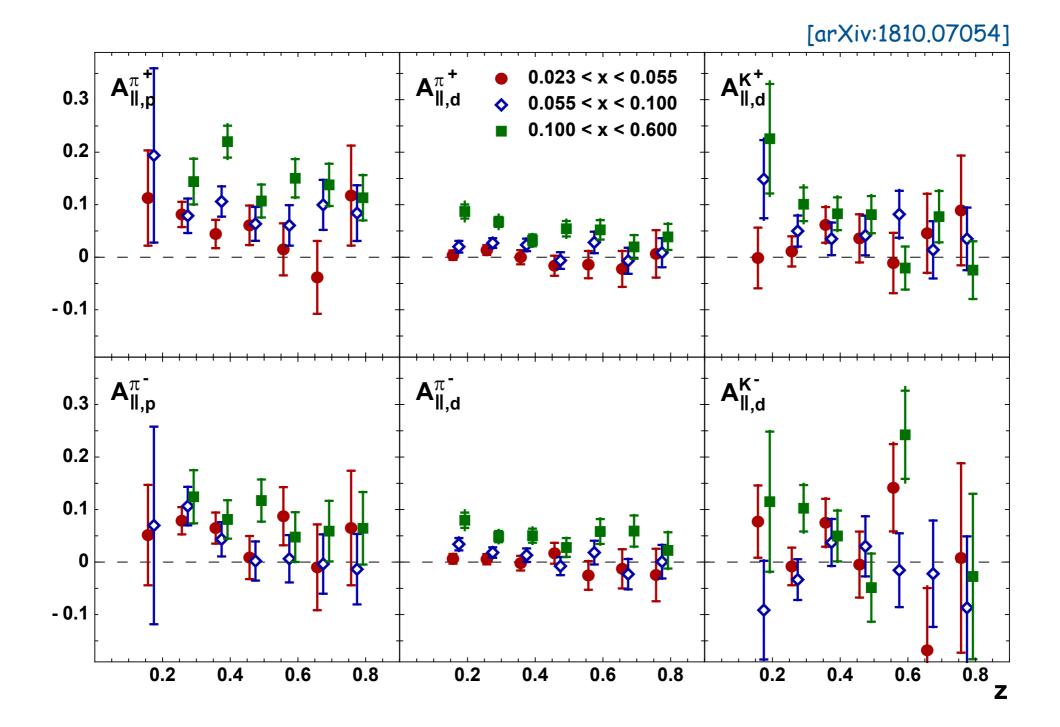
#### x dependence of A<sub>||</sub>

consistent with previous HERMES publication [PRD 71 (2005) 012003]



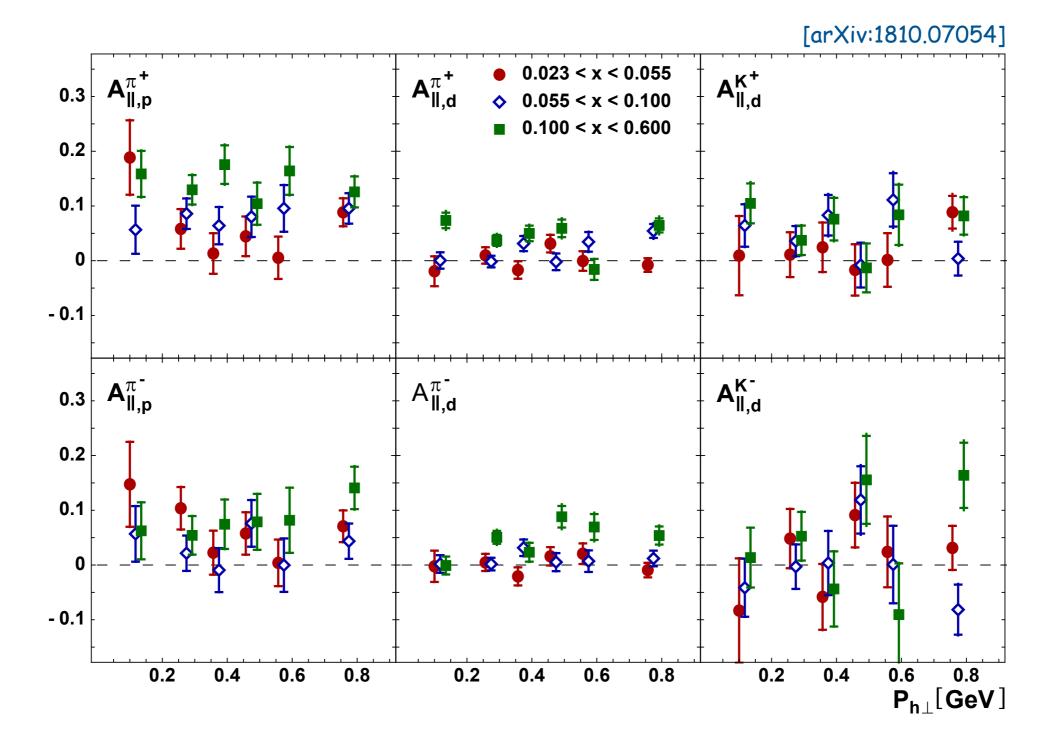
## z dependence of $A_{||}$ (three x ranges)

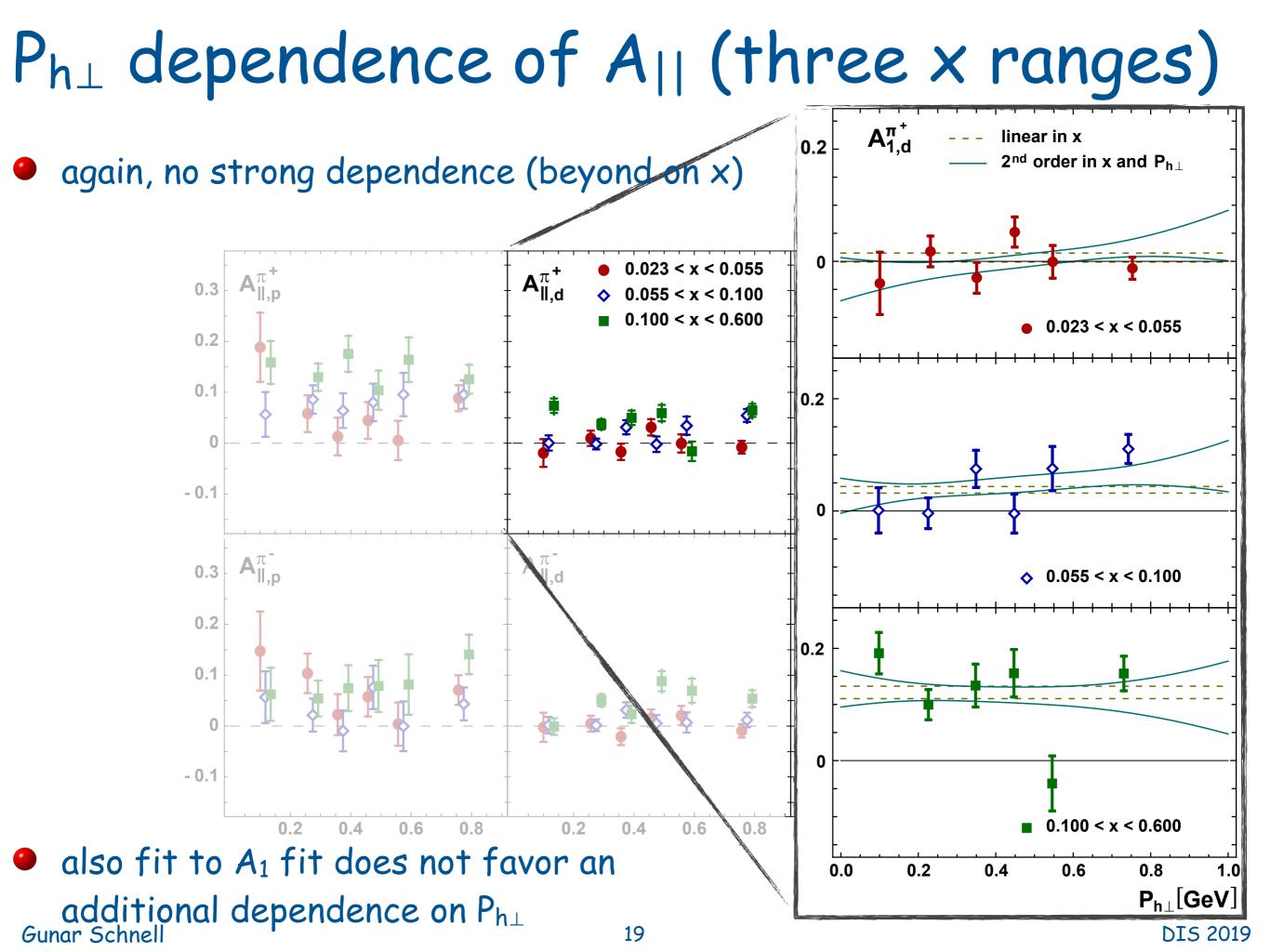
in general, no strong z-dependence visible



## $P_{h\perp}$ dependence of $A_{||}$ (three x ranges)

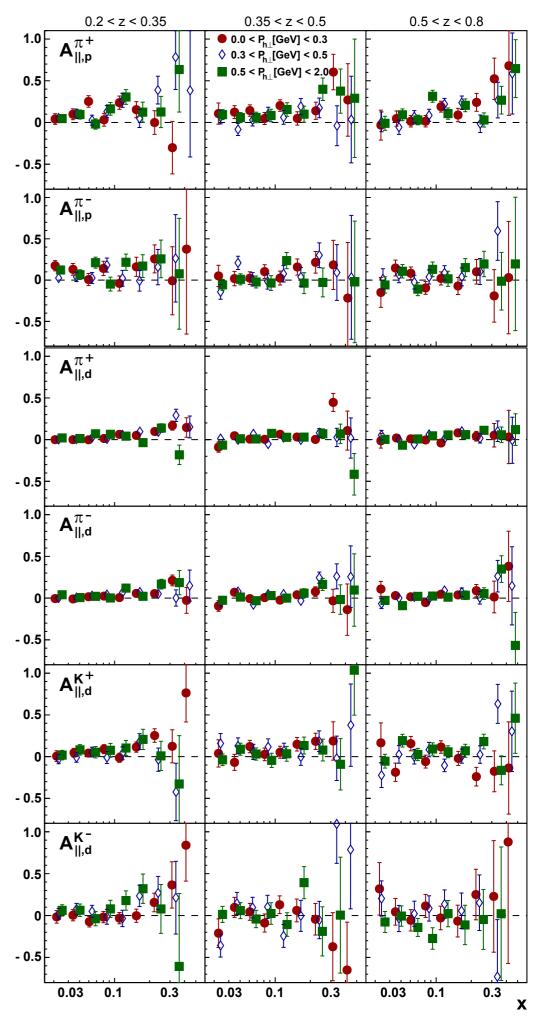
again, no strong dependence (beyond on x)





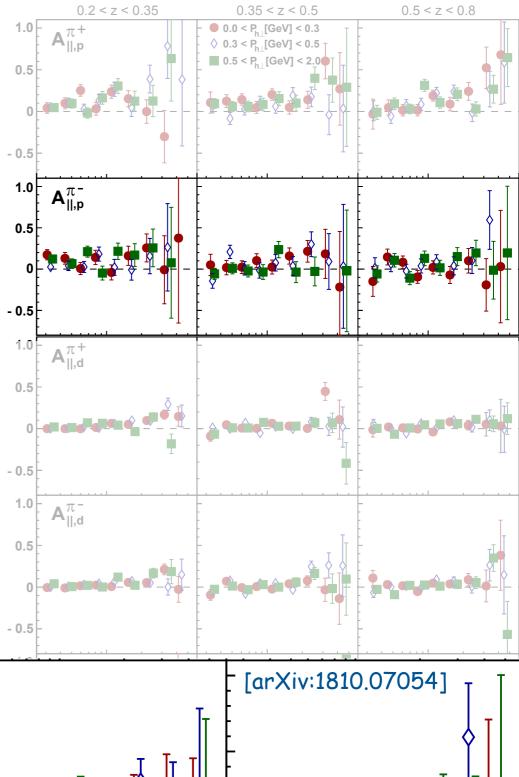
## 3-dimensional binning

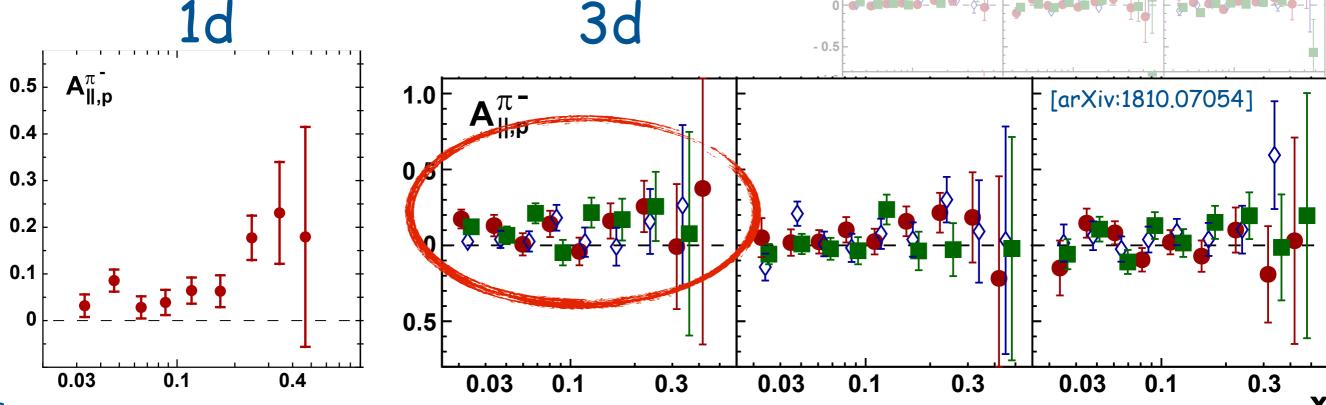
 3d dependences provides transversemomentum dependence



### 3-dimensional binning

- 3d dependences provides transversemomentum dependence
- but also extra flavor sensitivity, e.g.,
  - π<sup>-</sup> asymmetries mainly coming from low-z region where disfavored fragmentation large and thus sensitivity to the large positive up-quark polarization





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$$A_1^{h^+ - h^-}(x) \equiv \frac{\left(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}\right) - \left(\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-}\right)}{\left(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}\right) + \left(\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-}\right)}$$

$$A_1^{h^+ - h^-}(x) \equiv \frac{\left(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}\right) - \left(\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-}\right)}{\left(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}\right) + \left(\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-}\right)}$$

• at leading-order and leading-twist, assuming charge conjugation symmetry for fragmentation functions:

$$A_{1,d}^{h^+ - h^-} \stackrel{\text{LO LT}}{=} \frac{g_1^{u_v} + g_1^{d_v}}{f_1^{u_v} + f_1^{d_v}}$$

$$A_1^{h^+ - h^-}(x) \equiv \frac{\left(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}\right) - \left(\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-}\right)}{\left(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}\right) + \left(\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-}\right)}$$

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$$A_{1,d}^{h^+ - h^-} \stackrel{\text{lolt}}{=} \frac{g_1^{u_v} + g_1^{d_v}}{f_1^{u_v} + f_1^{d_v}}$$

assuming also isospin symmetry in fragmentation:

$$A_{1,p}^{h^+ - h^-} \stackrel{\text{lo lt}}{=} \frac{4g_1^{u_v} - g_1^{d_v}}{4f_1^{u_v} - f_1^{d_v}}$$

$$A_1^{h^+ - h^-}(x) \equiv \frac{\left(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}\right) - \left(\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-}\right)}{\left(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}\right) + \left(\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-}\right)}$$

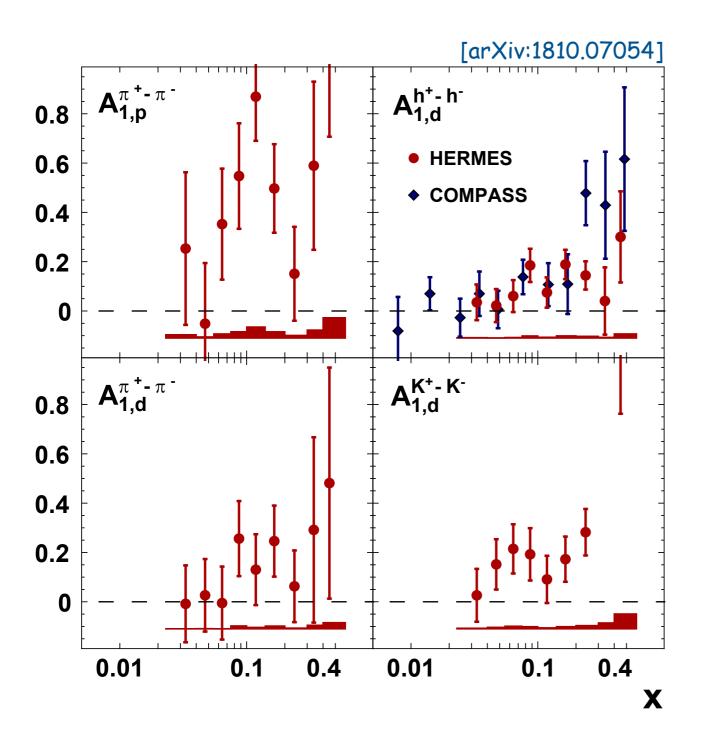
• at leading-order and leading-twist, assuming charge conjugation symmetry for fragmentation functions:

$$A_{1,d}^{h^+ - h^-} \stackrel{\text{lo lt}}{=} \frac{g_1^{u_v} + g_1^{d_v}}{f_1^{u_v} + f_1^{d_v}}$$

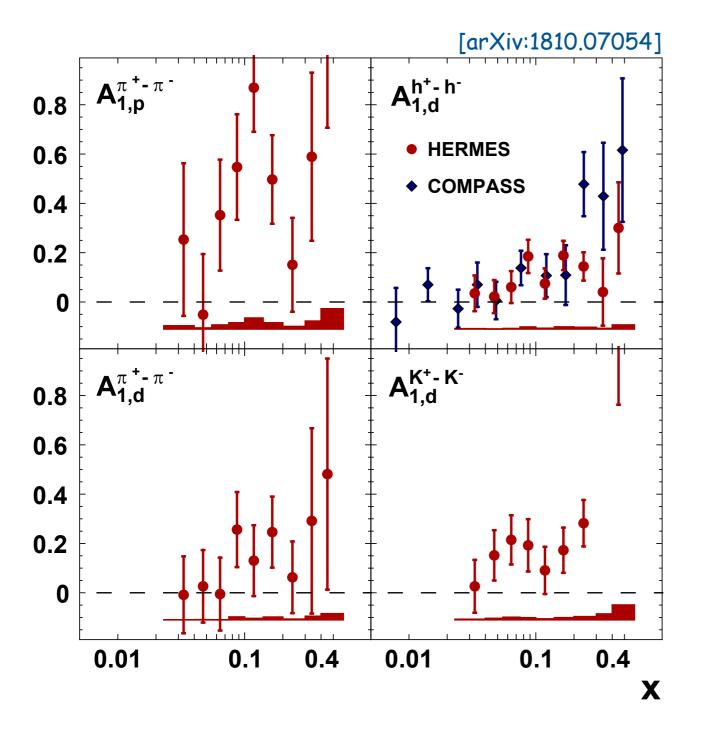
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$$A_{1,p}^{h^+ - h^-} \stackrel{\text{LO LT}}{=} \frac{4g_1^{u_v} - g_1^{d_v}}{4f_1^{u_v} - f_1^{d_v}}$$

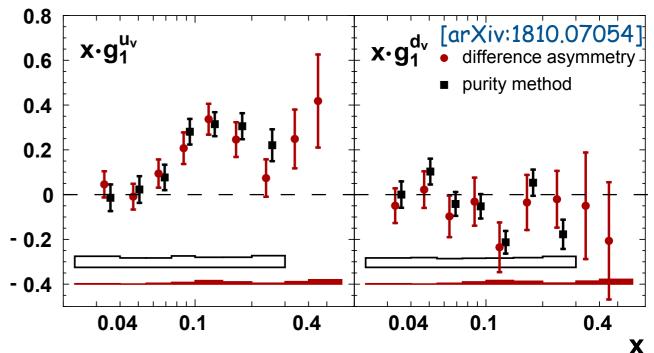
can be used to extract valence helicity distributions



- no significant hadron-type dependence for deuterons
- deuteron results (unidentified hadrons) consistent with COMPASS



- no significant hadron-type dependence for deuterons
- deuteron results (unidentified hadrons) consistent with COMPASS
- valence distributions consistent with JETSETbased extraction:



### azimuthal modulations

- twist-3 ratious contributions
- most prominent: "polarized Cahn effect"

$$xg_L^{\perp}D_1 \oplus \frac{M_h}{Mz}g_1\tilde{D}^{\perp} \oplus xe_LH_1^{\perp} \oplus \frac{M_h}{Mz}h_{1L}^{\perp}\tilde{E}$$

only one surviving WW-type approximations

🖛 K. Tezgin (We, WG-6)

l'

S

 $\mathbf{P}_{h\perp}$ 

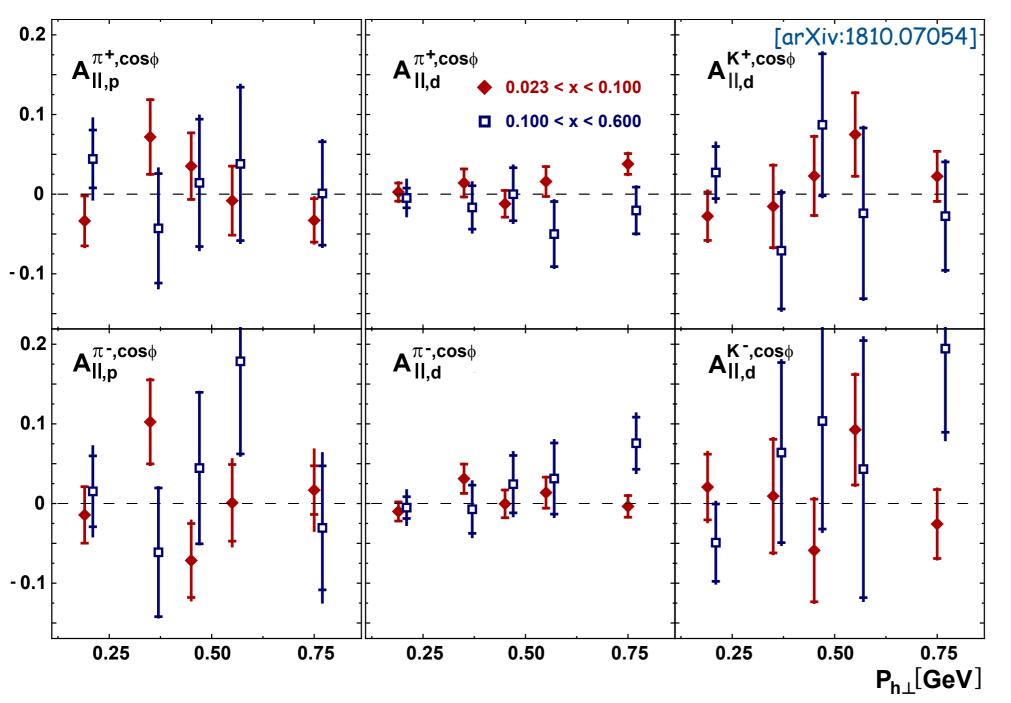
 $\mathbf{S}_{\perp}$ 

 $\mathbf{P}_h$ 

1

### azimuthal modulations

- twist-3 ratious contributions
- most prominent: "polarized Cahn effect"
- cosine modulations largely consistent with zero



 $\mathbf{P}_{h\perp}$ 

 $\mathbf{P}_h$ 

 $\mathbf{S}$ 

#### summary

- several longitudinal double-spin asymmetries in SIDIS have been presented that
  - extend the analysis of previous HERMES publications to include also transverse-momentum dependence and for the first time also a 3d binning
  - provide  $A_{||}$  in addition to  $A_1$
- within precision of the measurements, the virtual-photon-nucleon asymmetries display no significant dependence on z and  $P_{h\perp}$
- hadron-charge difference asymmetries in agreement with COMPASS
  - used for LO, leading-twist extraction of valence helicity PDFs
- $\cos \phi$  moments of semi-inclusive double-spin asymmetry compatible with zero

backup

#### data sets

	Beam	Target	Hadron	Hadron Momentum
Year	Type	Gas	Type	$P_h$
1996	$e^+$	Η	$\pi^{\pm}$	$4-13.8  {\rm GeV}$
1997	$\mathrm{e}^+$	Η	$\pi^{\pm}$	$4 - 13.8  {\rm GeV}$
1998	$e^-$	D	$\pi^{\pm}, K^{\pm}$	$215~\mathrm{GeV}$
1999	$\mathrm{e}^+$	D	$\pi^{\pm}, K^{\pm}$	$215~\mathrm{GeV}$
2000	$e^+$	D	$\pi^{\pm}, K^{\pm}$	$2-15 \mathrm{GeV}$