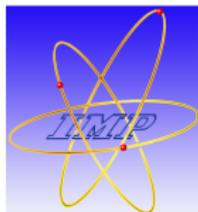


# Nucleon properties from basis light front quantization



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**DIS**

8-12 April

**2019**

**TORINO**

XXVII International Workshop on  
Deep Inelastic Scattering and  
Related Subjects

10<sup>th</sup> April, 2019

# Contents

- Methodology -Basis Light-Front Quantization (BLFQ)
- Numerical results
  - ✓ Form Factors
  - ✓ Parton distribution Functions (PDFs)
  - ✓ Generalized parton distributions (GPDs)
- Conclusions outlook

BLFQ: approach for solving quantum field theory



- **Nonperturbative:**  
for systems with strong interaction
- **First-principles:**  
effective Hamiltonian as input/ direct access to wavefunction of bound states
- **Light-front dynamics:**  
spectrum and light-front Fock-state wavefunctions are obtained from

$$H_{LF}|\psi\rangle = M^2|\psi\rangle$$

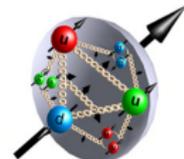
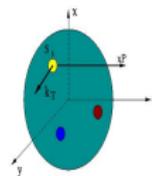
$$H_{LF} \equiv P_\mu P^\mu = P^+ P^- - \mathbf{P}_\perp^2$$

$$P^\pm = P^0 \pm P^3$$

LF wavefunctions

Proton 3D imaging

Proton spin



FF

GPD

TMD...

# General Procedure for BLFQ

- ✓ Derive/write the Light-Front Hamiltonian:  $P^- = H_{\text{eff}}$
- ✓ Construct the basis state:  $|\alpha\rangle$

- Calculate the Hamiltonian matrix elements:

$$H_{\text{eff}}^{\alpha'\alpha} = \langle\alpha'|H_{\text{eff}}|\alpha\rangle$$

$\langle\alpha'|$  &  $|\alpha\rangle$  are the basis state of BLFQ, such as  $|qqq\rangle$  .

- Diagonalize  $H_{\text{eff}}$  and obtain its eigen spectrum

$$H_{\text{eff}}|\beta\rangle = H_{\text{eff}}^{\beta}|\beta\rangle$$

- $|\beta\rangle$  is the physical state and eigenstate of Hamiltonian. In case of proton  $|\beta\rangle = |P_{\text{proton}}\rangle$ .

- ✓ Evaluate observables:

$$\mathcal{O} = \langle\beta|\hat{\mathcal{O}}|\beta\rangle$$

# Previous applications of BLFQ

## QCD systems

- **Heavy quarkonium**: decay constant, elastic form factor, radii, radiative transitions, distribution amplitude, GPDs

—Y Li, G Chen, X Zhao, P Maris, J Vary, L Adhikari, M Li, A El-Hady (2016 - 2019)

- **Light mesons**: spectrum, decay constant, elastic form factor, radii, distribution amplitude, PDFs and scale evolution

—S Jia, J Vary, J Lan, CM, X. Zhao (2018 - 2019)

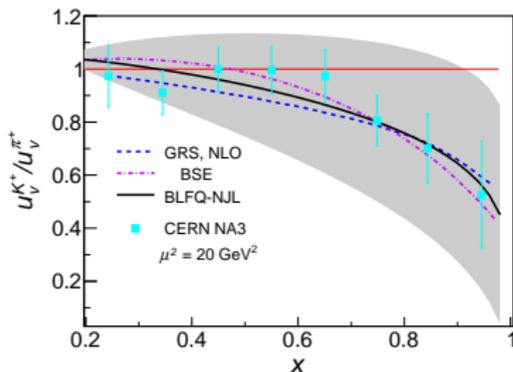
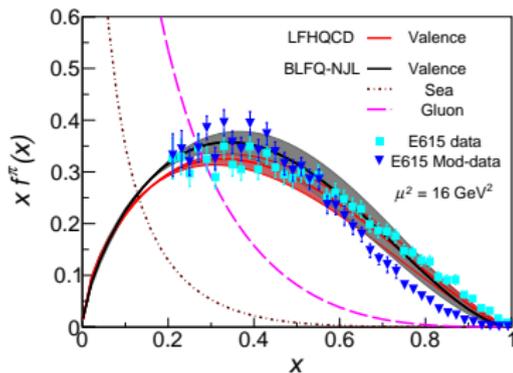
## QED systems

- **Electron**: anomalous magnetic moments
- **positronium** wave function, spectroscopy
- GPDs of the electron and positronium

—X. Zhao, P. Wiecki, Y. Li, H. Honkanen, D. Chakrabarti, P. Maris, J. P. Vary, S. J. Brodsky (2013 - 2018)

# Example: light meson PDFs

Lan, CM, Jia, Zhao, Vary: to appear in PRL (2019)



Light front effective Hamiltonian,  $H_{\text{eff}}$ : ( $\mu_{0\pi}^2 = 0.240 \pm 0.024 \text{ GeV}^2$ )

$$\underbrace{\frac{\vec{k}_\perp^2 + m_q^2}{x} + \frac{\vec{k}_\perp^2 + m_{\bar{q}}^2}{1-x}}_{\text{LF Kinetic energy}} + \underbrace{\kappa^4 x(1-x) \vec{r}_\perp^2}_{\text{Transverse}} - \underbrace{\frac{\kappa^4}{(m_q + m_{\bar{q}})^2} \partial_x (x(1-x) \partial_x)}_{\text{Longitudinal}} + H_{\text{NJL}}^{\text{eff}},$$

- Light front wavefunctions ► eigenvectors of this Hamiltonian.  
PDFs evolution ► based on the NNLO DGLAP equations.
- The robustness of our results motivates the application of analogous effective Hamiltonians to the baryons.

# Effective Light-front Hamiltonian for nucleon

## Light-Front Hamiltonian

$$P_{baryon}^- = H_{K.E.} + H_{trans} + H_{longi} + H_{OGE}$$

$$H_{K.E.} = \sum_i \frac{p_i^2 + m_q^2}{p_i^+} \quad m_u = 0.6 \text{ GeV}, m_d = 0.57 \text{ GeV}$$

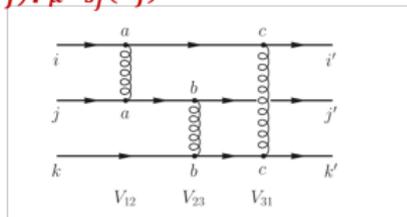
$$H_{trans} \sim \kappa_T^4 b^4 \zeta^2 \quad \text{-- Brodsky, Teramond arXiv: 1203.4025}$$

$$H_{longi} = - \sum_{ij} \kappa_L^4 b^4 \partial_{x_i} (x_i x_j \partial_{x_j}) \quad \text{--Y Li, X Zhao, P Maris, J Vary PLB 758(2016)}$$

$$H_{OGE} = - \frac{C_F 4\pi\alpha_s(Q^2)}{Q^2} \sum_{i,j(i<j)} \bar{u}_{s'_i}(k'_i) \gamma^\mu u_{s_i}(k_i) \bar{u}_{s'_j}(k'_j) \gamma_\mu u_{s_j}(k_j)$$

$$\text{Infrared cut off : } m_g = 0.05 \text{ GeV}, C_F = -\frac{2}{3}$$

$$|P_{baryon}\rangle = |qqq\rangle + |qqqg\rangle + |qqq q\bar{q}\rangle + \dots$$



Although we truncate to the **leading Fock sector**, we can solve the **baryon system** with **multi-particle** (at least three particle). In future, we can include **higher Fock sector**, and study the **sea quark** and **gluon** distribution function.

# Basis construction

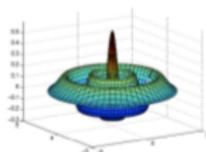
- Fock's space expansion:

$$|N\rangle_{\text{baryon}} = a|qqq\rangle + b|qqqg\rangle + c|qqqq\bar{q}\rangle + \dots$$

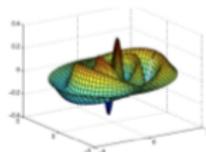
- For each Fock particle:

- **Transverse**: 2D harmonic oscillator basis:  $\Phi_{n,m}^b(\vec{p}_\perp)$   
labeled by radial (angular) quantum number  $n$  ( $m$ ); scale parameter  $b$

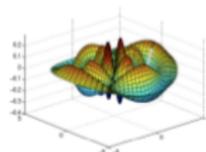
e.g.,  $n=4$



$m=0$



$m=1$



$m=2$

- **Longitudinal**: plane-wave basis, labeled by  $k$
- **Helicity**: labeled by  $\lambda$

- For the first Fock sector:

$$|qqq\rangle = |n_{q_1}, m_{q_1}, k_{q_1}, \lambda_{q_1}\rangle \otimes |n_{q_2}, m_{q_2}, k_{q_2}, \lambda_{q_2}\rangle \otimes |n_{q_3}, m_{q_3}, k_{q_3}, \lambda_{q_3}\rangle.$$

# Basis Truncation Scheme

- Symmetries of Hamiltonian:

- Net fermion number:

$$\sum_i n_i^f = N^f$$

- Total angular momentum projection:

$$\sum_i (m_i + s_i) = J_z$$

- Longitudinal momentum:

$$\sum_i k_i = K$$

- Further truncation:

- Fock sector truncation

- Discretization in longitudinal direction

$$k_i = \begin{cases} 1, 2, 3, \dots & \text{bosons} \\ 0.5, 1.5, 2.5, \dots & \text{fermions} \end{cases}$$

- " $N_{\max}$ " truncation in transverse directions

$$\sum_i [2n_i + |m_i| + 1] \leq N_{\max}$$

UV cutoff  $\Lambda \sim b\sqrt{N_{\max}}$ ; IR cutoff  $\lambda \sim b/\sqrt{N_{\max}}$

# Flavor Form Factor in BLFQ



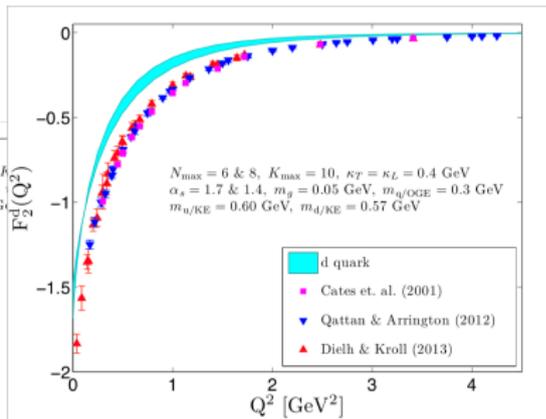
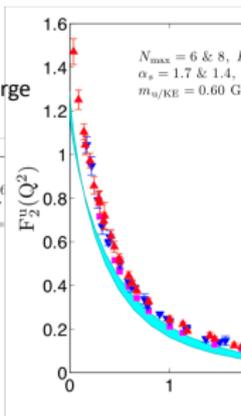
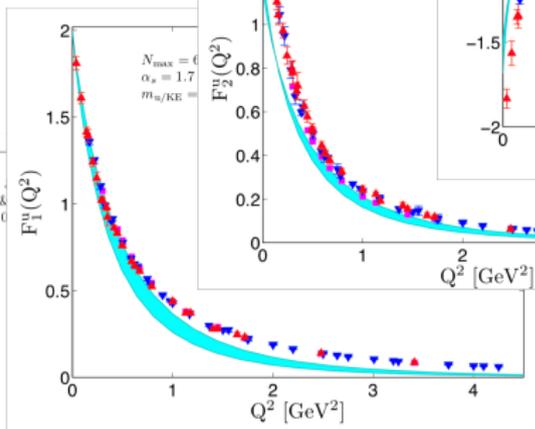
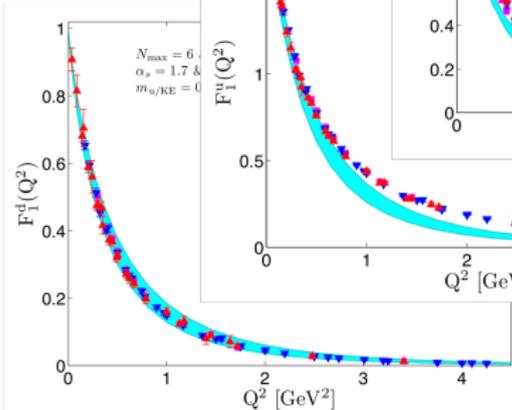
[Work in progress, C. Mondal, Siqi Xu, et. al.]

*different quark masses*

in kinetic and one gluon exchange term to  
*minimize the effect of higher Fock Sector*

**Dirac form factor**

**d quark:** good agreement with data;  
**u quark:** deviates from the data at large  $Q^2$



**Pauli Form factor**

**d quark:**  $F_2(0) = -1.73$  (Exp. :  $-2.03$ )  
**u quark:**  $F_2(0) = 1.27$  (Exp. :  $1.67$ )

With *increasing the basis size*,  
our result *approach to the experiment data*.

# Nucleon Form Factor

$$G_E(Q^2) = \sum_q e_q F_1^q(Q^2) - \frac{Q^2}{4M^2} \sum_q e_q F_2^q(Q^2),$$

$$G_M(Q^2) = \sum_q e_q F_1^q(Q^2) + \sum_q e_q F_2^q(Q^2).$$

## Charge form factor

**Neutron** : agrees with experiment data

**Proton** : also agrees with the data

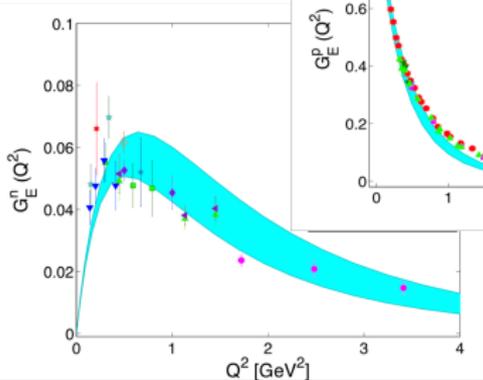
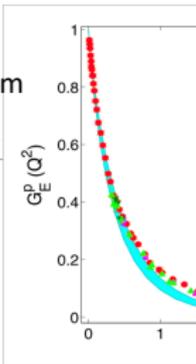
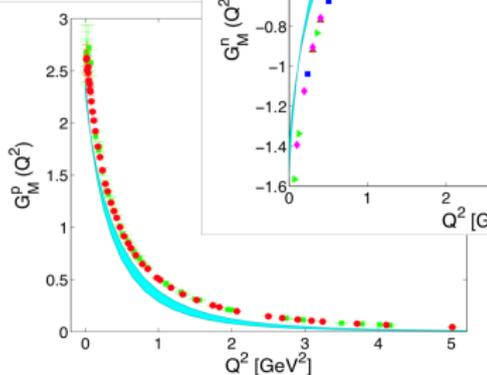
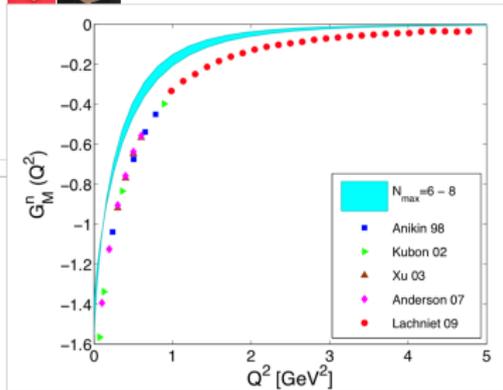
## Magnetic form factor

**Proton** : more or less agrees with data

**Neutron** : deviates from the data



[Work in progress, C. Mondal, Siqi Xu, et.al]



## Anomalous magnetic moment

Proton :  $\mu_N = 2.42$  (Exp. : 2.79)

Neutron :  $\mu_N = -1.57$  (Exp. : -1.91)

With *increasing the basis size*, the magnetic form factor *approach to the experiment data*

$$\langle r_E^2 \rangle^N = -6 \frac{dG_E^N(Q^2)}{dQ^2} \Big|_{Q^2=0},$$

$$\langle r_M^2 \rangle^N = -\frac{6}{G_M^N(0)} \frac{dG_M^N(Q^2)}{dQ^2} \Big|_{Q^2=0}$$

The Sachs form factors are defined as

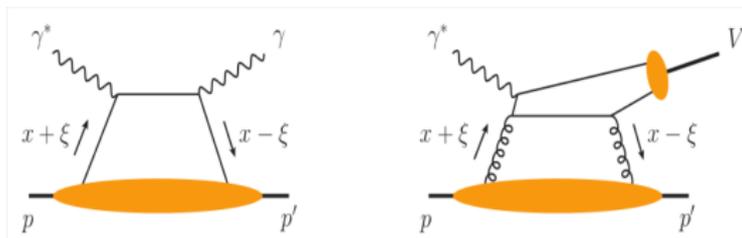
$$G_E^N(Q^2) = F_1^N(Q^2) - \frac{Q^2}{4M_N^2} F_2^N(Q^2),$$

$$G_M^N(Q^2) = F_1^N(Q^2) + F_2^N(Q^2).$$

Quantity	$N_{\max} = 6$	$N_{\max} = 8$	Data from PDG
$r_E^p$ (fm)	0.966	0.877	$0.877 \pm 0.005$
$r_M^p$ (fm)	1.094	1.008	$0.777 \pm 0.016$
$\langle r_E^2 \rangle^n$ (fm <sup>2</sup> )	-0.1175	-0.1398	$-0.1161 \pm 0.0022$
$r_M^n$ (fm)	1.278	1.2045	$0.862^{+0.009}_{-0.008}$

# Generalized Parton Distribution functions (GPDs)

➤ **Deeply Virtual Compton Scattering (DVCS)** / vector meson productions **experiment:**



Encode the information about  
three dimensional spatial structure  
the spin and orbital angular momentum

➤ **GPDs** appear in **DVCS** processes.

• GPDs are functions of three variables :

▪ **Longitudinal momentum fraction**  $x = \frac{k^+}{p^+}$

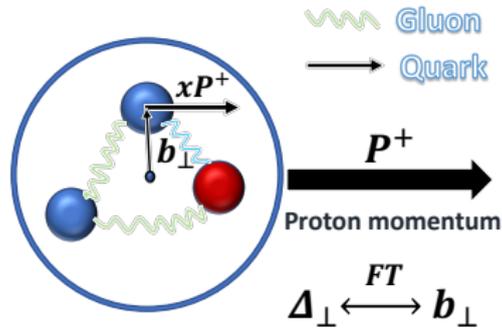
$$\frac{k^+}{p^+}$$

▪ **Longitudinal momentum transfer**  $\rightarrow$

$$\text{skewness } \xi = \frac{\Delta^+}{p^+} = 0$$

▪ **Square of total mom transfer**

$$t = \Delta^2 = (\mathbf{P}' - \mathbf{P})^2$$

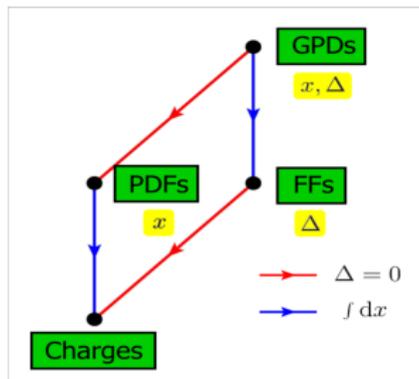


where the  $\mathbf{b}_{\perp}$  is transverse position of parton

# Generalized Parton Distribution Functions (GPD)

**GPD** encode the information about three dimensional spatial structure of nucleon, as well as the spin and orbital angular momentum of the constituents

With **increasing momentum transfer ( $t$ )**, the **peaks of distributions** shift to **larger  $x$** ;  
 At **large  $x$** , **BLFQ results** follow the **quark-diquark model**

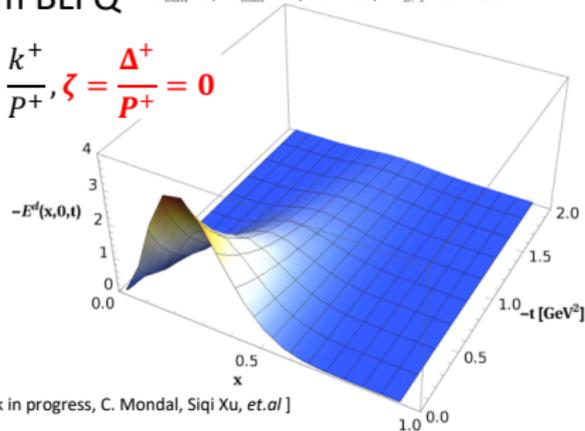
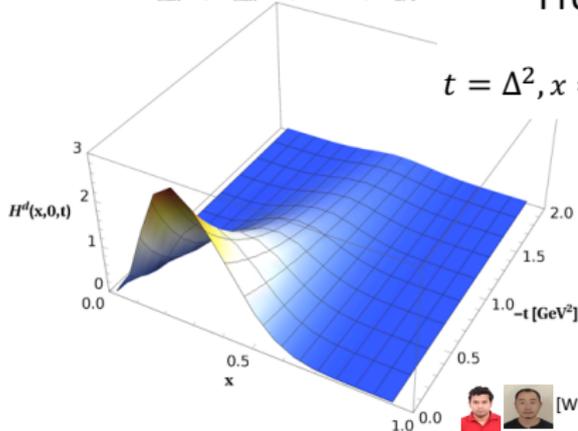


$N_{\max}=8, K_{\max}=16, \alpha=0.8, \kappa_{L/T}=0.4 \text{ GeV}$

From BLFQ

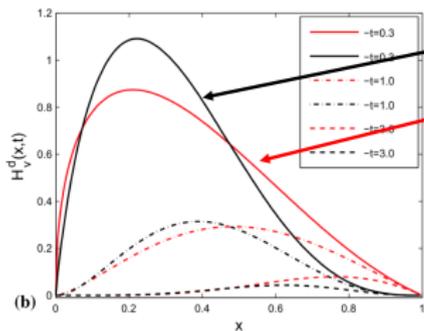
$N_{\max}=8, K_{\max}=16, \alpha=0.8, \kappa_{L/T}=0.4 \text{ GeV}$

$$t = \Delta^2, x = \frac{k^+}{p^+}, \zeta = \frac{\Delta^+}{p^+} = 0$$



[Work in progress, C. Mondal, Siqi Xu, et.al]

# Generalized Parton Distribution Functions



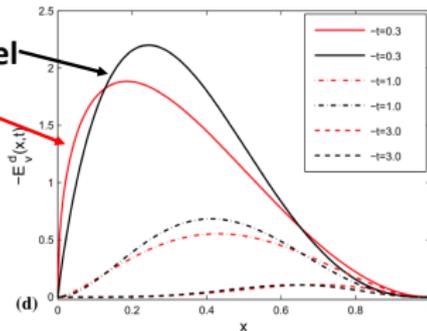
quark-diquark model

AdS/QCD

$$t = \Delta^2$$

$$x = \frac{k^+}{p^+}, \zeta = \frac{\Delta^+}{p^+} = 0$$

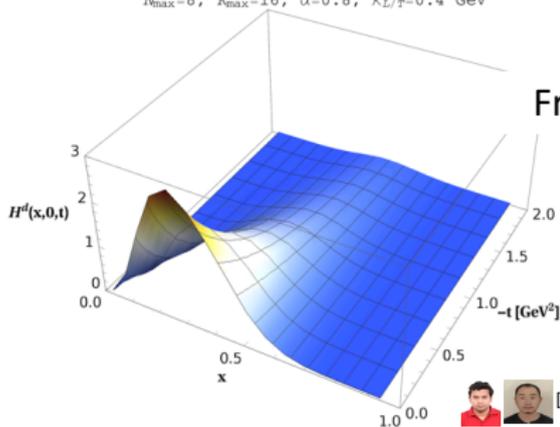
For d quark



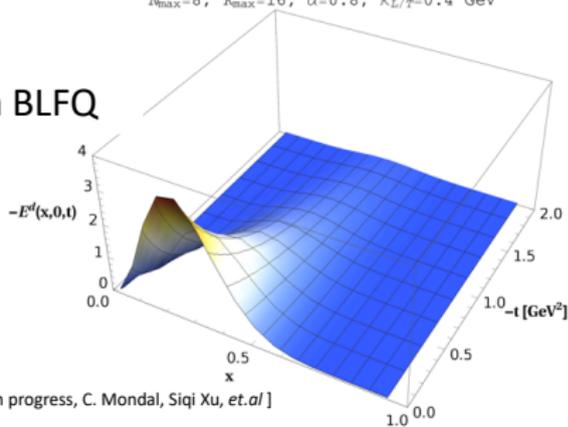
With increasing momentum transfer ( $t$ ), the peaks of distributions shift to larger  $x$ ;  
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$N_{\max}=8, K_{\max}=16, \alpha=0.8, \kappa_{L/T}=0.4$  GeV

$N_{\max}=8, K_{\max}=16, \alpha=0.8, \kappa_{L/T}=0.4$  GeV



From BLFQ



[Work in progress, C. Mondal, Siqi Xu, et al.]

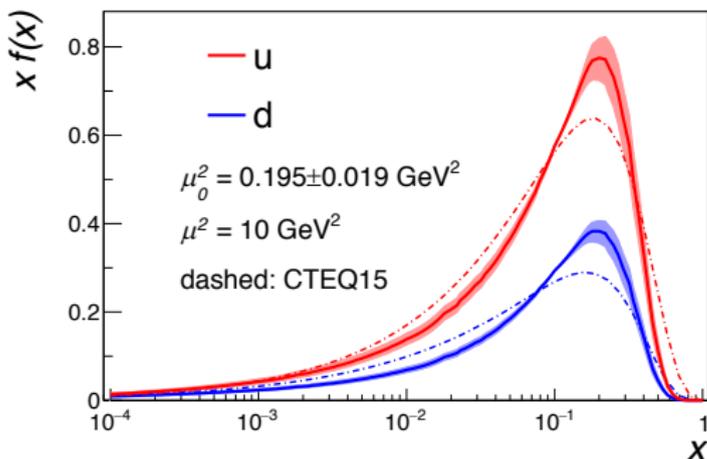
# Parton distribution functions (PDF)

## Parton distribution functions in BLFQ:

$$\langle x_q \rangle = \int dx x f_1^q(x)$$

Average momentum is carried by the quark.

Here, we use the NNLO DGLAP equation to evolve the PDF. Qualitative behavior of PDF is more or less same with the global fit CTEQ 15 PDF.



### U Quark

**Low  $x$  region:** agrees with global fit

**High  $x$  region:** deviates from the fit

### D Quark

**Low  $x$  region:** deviates from the fit

**High  $x$  region:** more or less agrees with global fit

We need the contribution from **higher Fock sector** such as  $|qqqq\rangle, |qqq q\bar{q}\rangle$  to obtain **better agreement**, and to study the **sea quark** and **gluon** distributions.



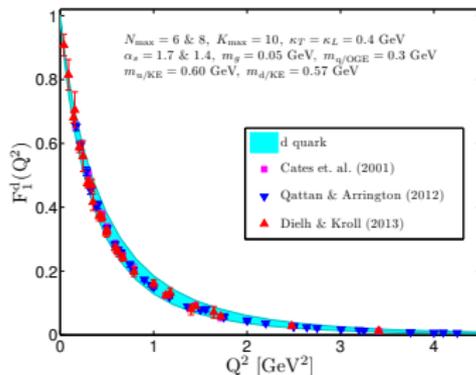
[Work in progress, C. Mondal, Siqi Xu, Jiangshan Lan, et al.]

## Conclusions

- We have discussed the preliminary results of nucleon form factors, PDFs & GPDs in BLFQ approach.
- In the effective Hamiltonian, we have the kinetic energy & the confining potential in both the transverse and longitudinal direction and one gluon exchange with fixed coupling. Here, we consider only the leading Fock sector.
- BLFQ formalism provides promising results in order to understand the nucleon structure.

## Outlook:

- Increase basis size
- Include the higher Fock component  $|qqqg\rangle$ .
- Investigate other nucleon properties..
- Investigate the structure of other baryons.



Thank You