

Shear forces and tensor polarization



DIS-2019, Torino, April 9

Oleg Teryaev
JINR, Dubna



Main Topics

- Gravitational formfactors and (Extended) Equivalence Principle: spin $1/2$ and GPDs
- Spin 1 : INCLUSIVE processes with tensor polarizartion
- HERMES data and (Ex)EP
- New possibilities: DY and **hadronic** hard processes
- NICA@JINR

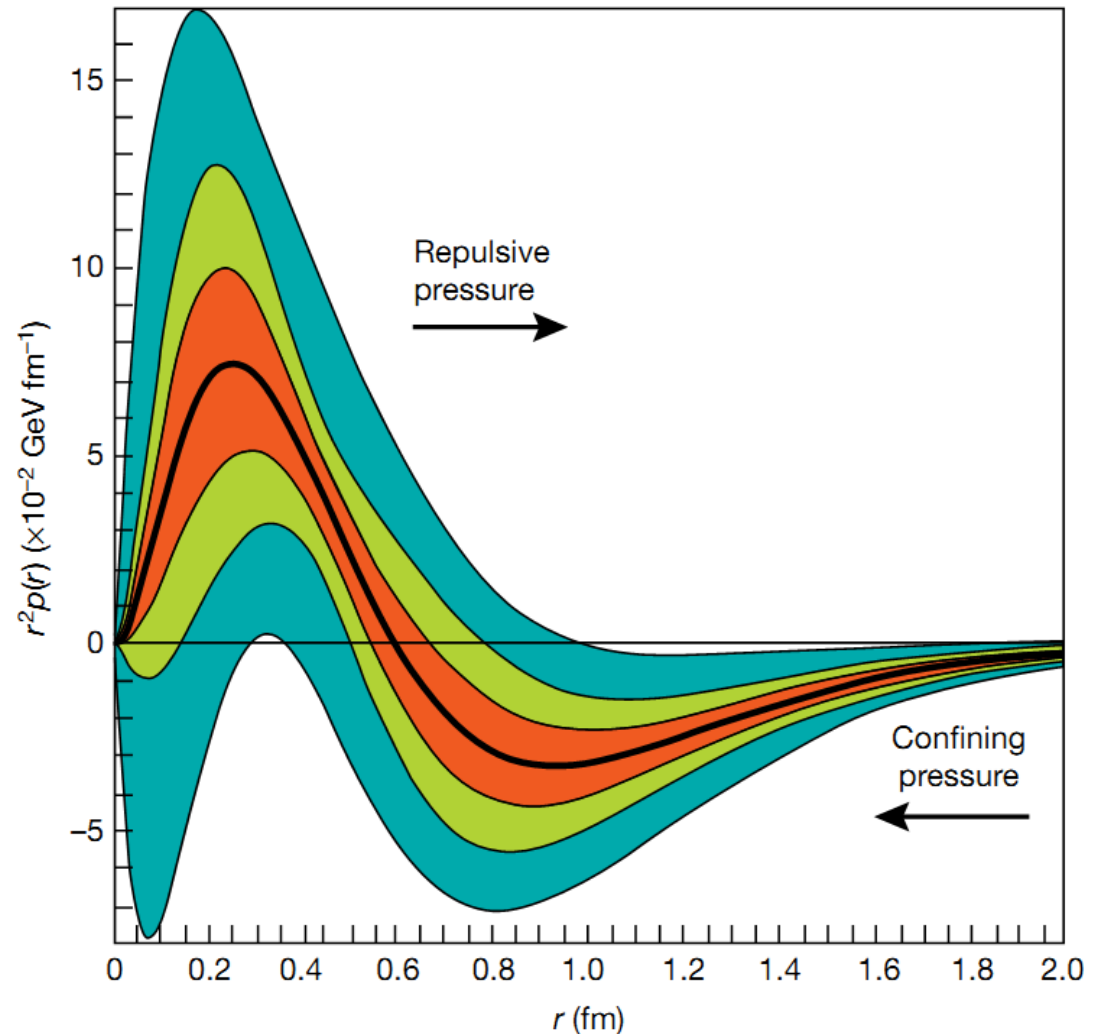


Recent development: pressure

- Published in Nature, already used for EoS (talk of A. Bacchetta)

The pressure distribution inside the proton

V. D. Burkert^{1*}, L. Elouadrhiri¹ & F. X. Girod¹



Gravitational Formfactors (spin 1/2)

$$\langle p' | T_{q,g}^{\mu\nu} | p \rangle = \bar{u}(p') \left[A_{q,g}(\Delta^2) \gamma^{(\mu} p^{\nu)} + B_{q,g}(\Delta^2) P^{(\mu} i \sigma^{\nu)\alpha} \Delta_\alpha / 2M \right] u(p)$$

- Conservation laws - zero Anomalous Gravitomagnetic Moment : $\mu_G = J$ (g=2)

$$P_{q,g} = A_{q,g}(0) \quad A_q(0) + A_g(0) = 1$$

$$J_{q,g} = \frac{1}{2} [A_{q,g}(0) + B_{q,g}(0)] \quad A_q(0) + B_q(0) + A_g(0) + B_g(0) = 1$$

- May be extracted from high-energy experiments/NPQCD calculations
- Describe the partition of angular momentum between quarks and gluons
- Describe interaction with both classical and TeV gravity

Electromagnetism vs Gravity (OT'99)

- Interaction – field vs metric deviation

$$M = \langle P' | J_q^\mu | P \rangle A_\mu(q)$$

$$M = \frac{1}{2} \sum_{q,G} \langle P' | T_{q,G}^{\mu\nu} | P \rangle h_{\mu\nu}(q)$$

- Static limit

$$\langle P | J_q^\mu | P \rangle = 2e_q P^\mu$$

$$\sum_{q,G} \langle P | T_i^{\mu\nu} | P \rangle = 2P^\mu P^\nu$$
$$h_{00} = 2\phi(x)$$

$$M_0 = \langle P | J_q^\mu | P \rangle A_\mu = 2e_q M \phi(q)$$

$$M_0 = \frac{1}{2} \sum_{q,G} \langle P | T_i^{\mu\nu} | P \rangle h_{\mu\nu} = 2M \cdot M \phi(q)$$

- Mass as charge – equivalence principle



Gravitomagnetism

- Gravitomagnetic field (weak, except in gravity waves) – action on spin from $M = \frac{1}{2} \sum_{q,G} \langle P' | T_{q,G}^{\mu\nu} | P \rangle h_{\mu\nu}(q)$

$$\vec{H}_J = \frac{1}{2} \text{rot} \vec{g}; \quad \vec{g}_i \equiv g_{0i}$$

spin dragging twice
smaller than EM

- Lorentz force – similar to EM case: factor $1/2$ cancelled with 2 from frequency same as EM

$$h_{00} = 2\phi(x) \quad \text{Larmor}$$

$$\omega_J = \frac{\mu_G}{J} H_J = \frac{H_L}{2} = \omega_L \quad \vec{H}_L = \text{rot} \vec{g}$$

- Orbital and Spin momenta dragging – the same - Equivalence principle

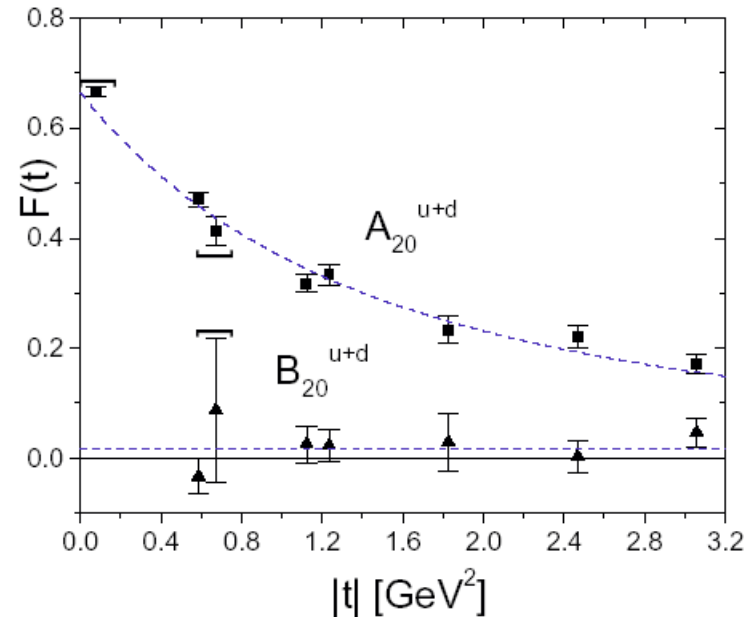
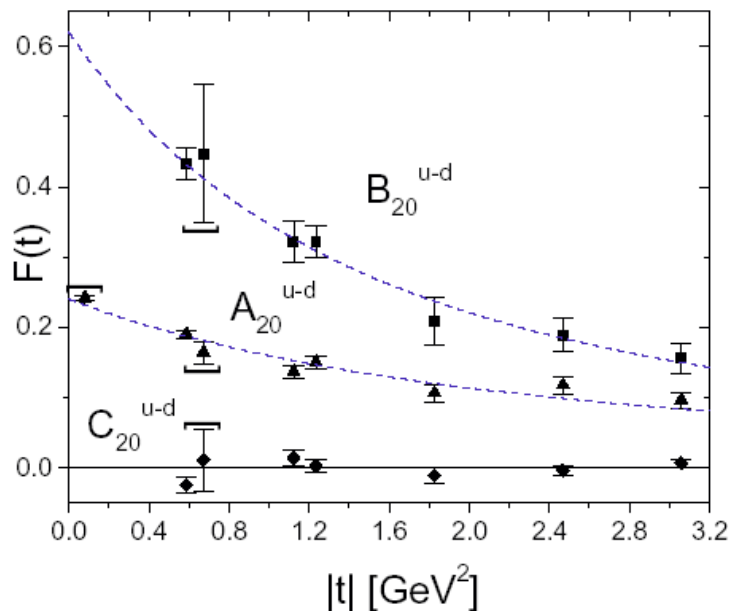


Equivalence principle

- Newtonian – “Falling elevator” – well known and checked (also for elementary particles)
- Post-Newtonian – gravity action on SPIN – known since 1962 (Kobzarev and Okun’; rederived from conservation laws - Kobzarev and V.I. Zakharov)
- Anomalous gravitomagnetic (and electric-CP-odd) moment is ZERO or
- Classical and QUANTUM rotators behave in the SAME way
- Gravitational analog of Ji’s SR $\int dx \times (\sum E_q + E_G) = 0!$

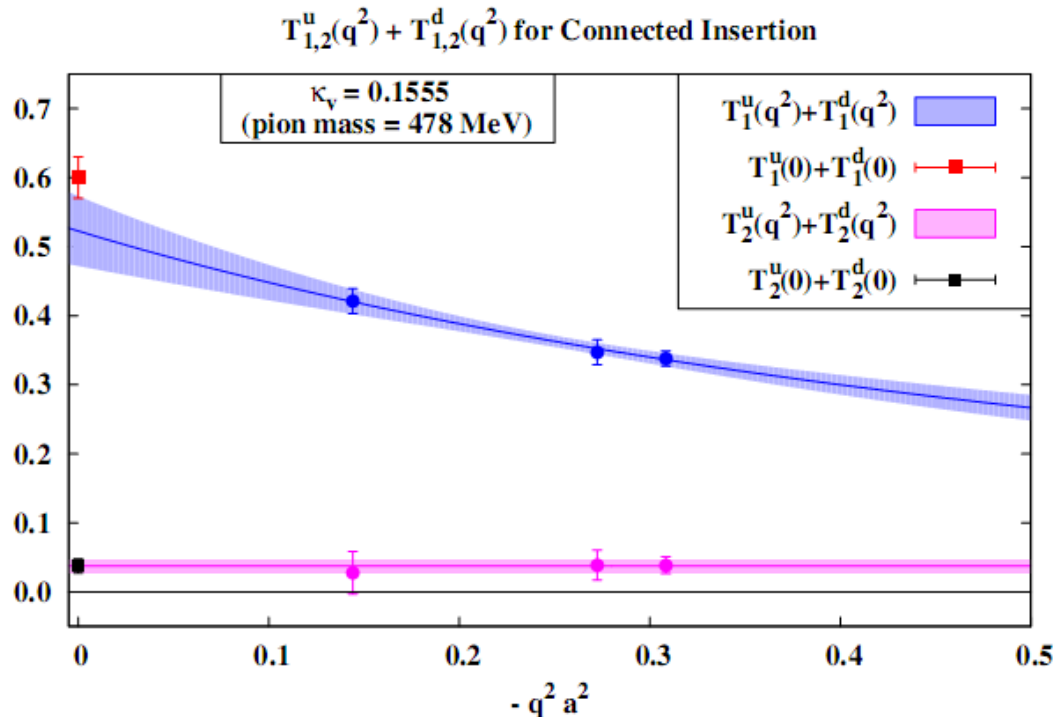
Generalization of Equivalence principle

- Various arguments: $AGM \approx 0$ separately for quarks and gluons – most clear from the lattice (LHPC/SESAM)



Recent lattice study (M. Deka et al. Phys.Rev. D91 (2015) no.1, 014505)

- Sum of u and d for Dirac (T1) and Pauli (T2) FFs



Extended Equivalence

Principle=Exact EquiPartition

- In pQCD – violated
- Reason – in the case of ExEP- no smooth transition for zero fermion mass limit (Milton, 73)
- Conjecture (O.T., 2001 – prior to lattice data) – valid in NP QCD – zero quark mass limit is safe due to chiral symmetry breaking
- Gravityproof confinement? Nucleons do not break even by black holes?
- Support by recent observation of smallness of C_{bar}



Spin-1 hadrons

- MANY new FFs!
- Recent extensive analysis
- Cosyn, Cotogno, Freese, Lorce:
1903.00408
- Polyakov, Sun: 1903.02738
- A lot of integral relations between GPDs
and FFs



Spin 1 EMT and **inclusive** processes

- Forward matrix element \rightarrow density matrix
- Contains **P-even** term: tensor polarization
- Symmetric and **traceless**: correspond to (average) **shear** forces
- For spin $1/2$: P-odd vector polarization requires another vector (q) to form vector product



Spin 1 in QCD

- Tensor polarization in QCD: Frankfurt, Strikman (81), Efremov, OT (81)
- Spin $\frac{1}{2}$: kinematically enhanced longitudinal polarization transverse-twist 3
- Spin 1: LL/TT related by tracelessness



SUM RULEs

- We (A.V. Efremov, OT'81) derived zero sum rules:
- 1st moment: also in parton model by Close and Kumano (90)
- 2nd moment (forward analog of Ji's SR)
- Average shear force (compensated between quarks and gluons)
- Gravity and (Ex)EP (zero average shear separately for quarks and gluons) – OT'09

Manifestation of post-Newtonian (Ex)EP for spin 1 hadrons

- Tensor polarization - coupling of EMT to spin in forward matrix elements - inclusive processes
- Second moments of tensor distributions should sum to zero

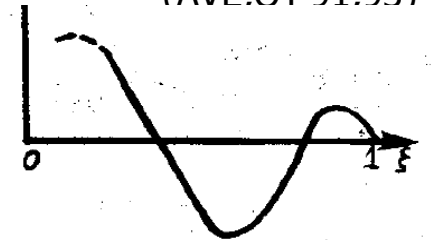
$$A_T = \frac{\sigma_+ + \sigma_- - 2\sigma_0}{3\bar{\sigma}}$$

$$\int_0^1 C_i^T(x) dx = 0$$

$$\langle P, S | \bar{\psi}(0) \gamma^\nu D^{\nu_1} \dots D^{\nu_n} \psi(0) | P, S \rangle_{\mu^2} = i^{-n} M^2 S^{\nu\nu_1} P^{\nu_2} \dots P^{\nu_n} \int_0^1 C_q^T(x) x^n dx \quad (\text{AVE.OT'91.93})$$

$$\sum_q \langle P, S | T_i^{\mu\nu} | P, S \rangle_{\mu^2} = 2P^\mu P^\nu (1 - \delta(\mu^2)) + 2M^2 S^{\mu\nu} \delta_1(\mu^2)$$

$$\langle P, S | T_g^{\mu\nu} | P, S \rangle_{\mu^2} = 2P^\mu P^\nu \delta(\mu^2) - 2M^2 S^{\mu\nu} \delta_1(\mu^2)$$

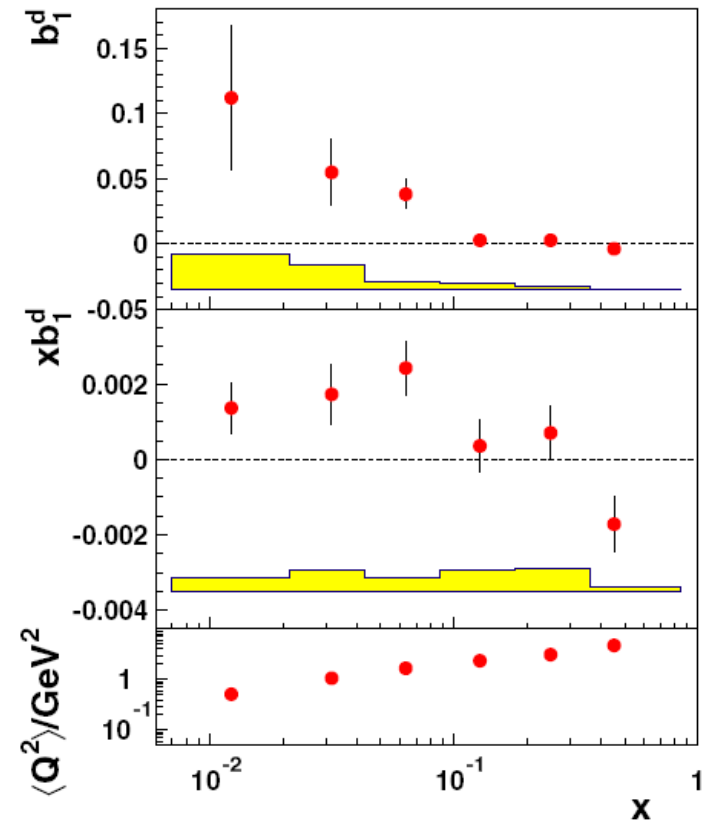


$$\sum_q \int_0^1 C_i^T(x) x dx = \delta_1(\mu^2) = 0 \quad \text{for ExEP}$$

HERMES – data on tensor spin structure function

PRL 95, 242001 (2005)

- Isoscalar target – proportional to the sum of u and d quarks – combination required by (Ex)EP
- Second moments – compatible to zero better than the first one (collective glue \ll sea)





Where else to test?

- COMPASS
- EIC
- DY@J-PARC:
(Song, Kumano: 1902.04712)
- However: ET'81-**any** hard process

- $f_{AI} \sim b_1$

$$\frac{P_{xx} - 2P_{yy} - 2P_{zz}}{3 \int_0^1 d\xi f(\xi) \text{Sp}[\hat{P}E(\xi, P)]} = \frac{2F_{AI}(x_1, x_2)}{3F(x_1, x_2)}$$

- Suggestion: **hadronic** tensor SSA



Vector vs Tensor SSA

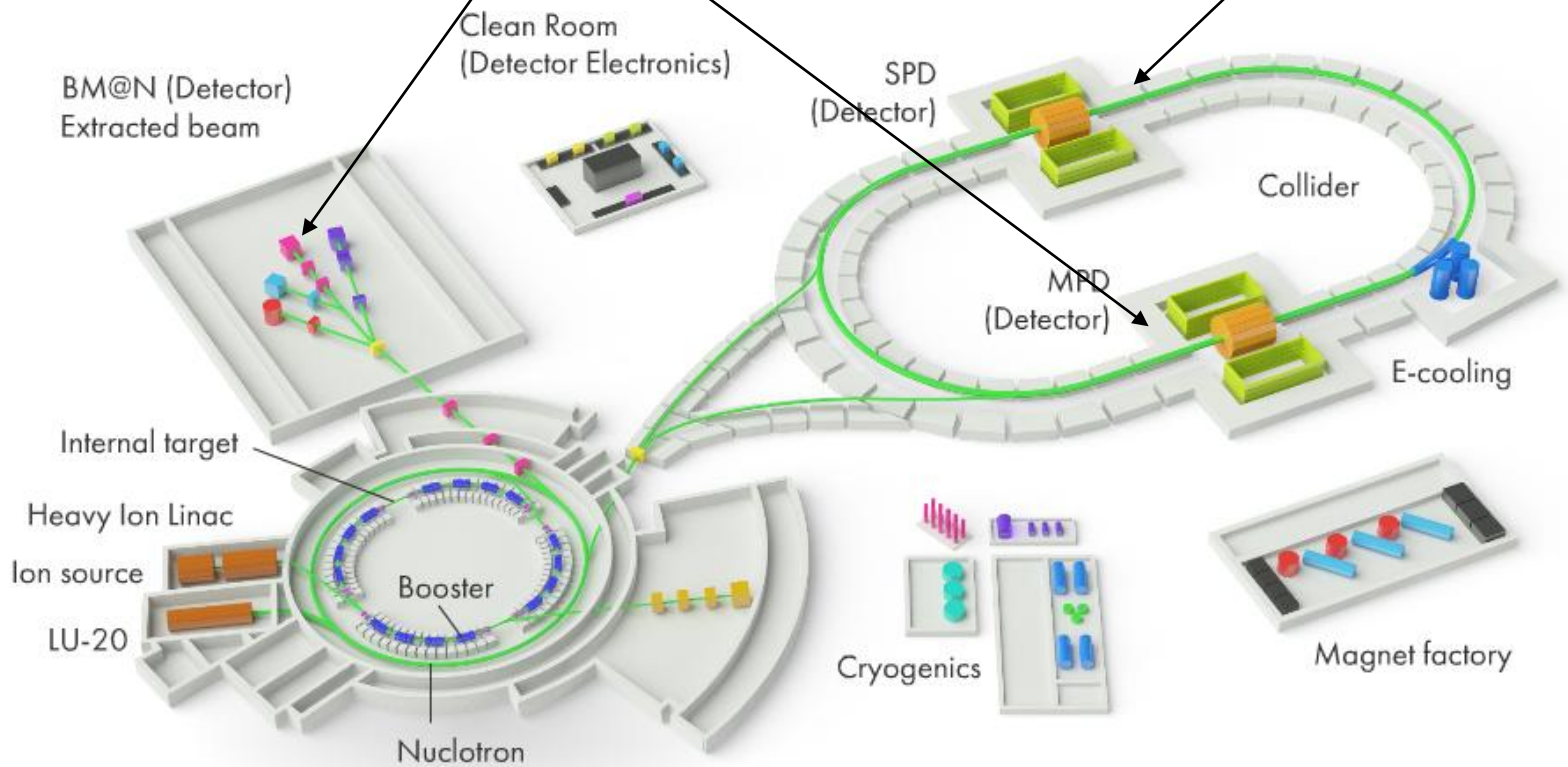
- Vector: $A = (\sigma(+)-\sigma(-))/(\sigma(+)+\sigma(-))$
- Tensor:
 $A = (\sigma(+)+\sigma(-))/(\sigma(+)+\sigma(-)+\sigma(0))$
- Inclusive pion production: (T-odd) vector SSA may be also excluded by summing $\sigma(L)+\sigma(R)$



Tensor polarized beams

- Opportunity: NICA@JINR with polarized **hadronic** beams
- Polarized deuterons is easier to accelerate: no depolarizing resonances
- DY, J/ψ (+**hadronic** SSA)

NICA: heavy ions and hadrons





Conclusions

- Tensor polarization: way to test of gravitational coupling of quarks and gluons in inclusive processes
- HERMES data are compatible with validity of EP separately for quarks and gluons
- Tests in DY and **hadronic** TSSA at NICA are possible



BACKUP



One more gravitational formfactor

- Quadrupole

$$\langle P + q/2 | T^{\mu\nu} | P - q/2 \rangle = C(q^2)(g^{\mu\nu} q^2 - q^\mu q^\nu) + \dots$$

- Cf vacuum matrix element – cosmological constant (vacuum pressure)

$$\langle 0 | T^{\mu\nu} | 0 \rangle = \Lambda g^{\mu\nu}$$

$$\Lambda = C(q^2) q^2$$

- Inflation \sim annihilation ($q^2 > 0$) OT'15
- How to measure experimentally – Deeply Virtual Compton Scattering



D-term interpretation: Inflation and annihilation

- Quadrupole gravitational FF

$$\langle P + q/2 | T^{\mu\nu} | P - q/2 \rangle = C(q^2)(g^{\mu\nu} q^2 - q^\mu q^\nu) + \dots$$

- Moment of D-term – positive
- Vacuum – Cosmological Constant $\langle 0 | T^{\mu\nu} | 0 \rangle = \Lambda g^{\mu\nu}$

- 2D effective CC – negative in scattering, positive in annihilation

$$\Lambda = C(q^2)q^2$$

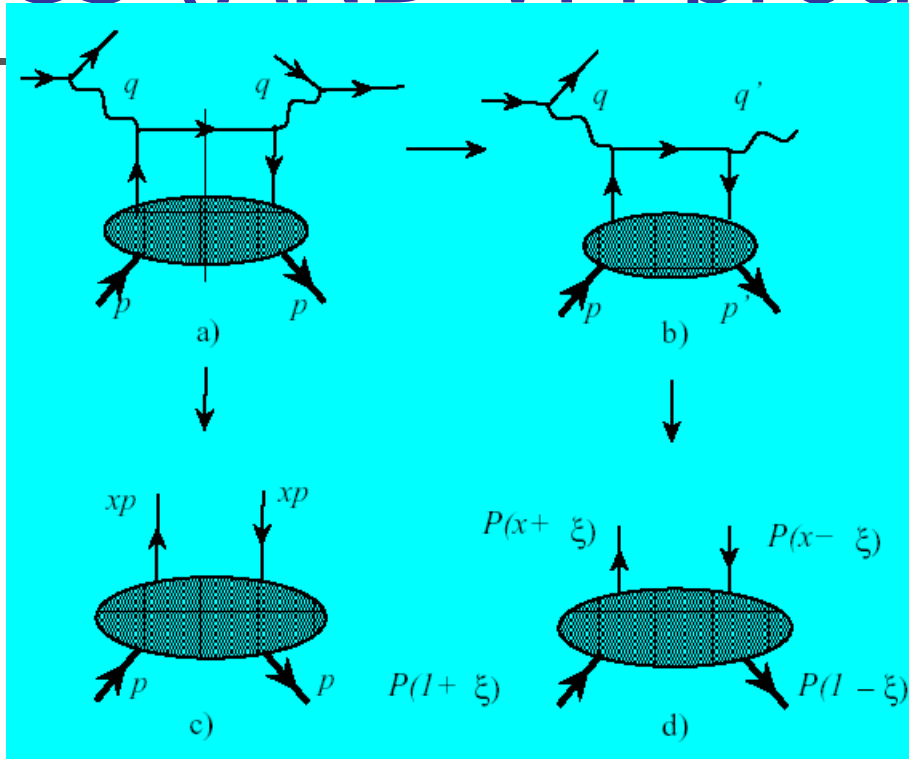
- Similarity of inflation and Schwinger pair production – Starobinsky, Zel'dovich
- Was OUR Big Bang resulting from one graviton annihilation at extra dimensions??! Version of "ekpyrotic" ("pyrotechnic") universe



C vs Cbar

- Cancellations of Cbars – negative pressure (cf Chaplygin gas)
- Cancellation in vacuum; Pauli (divergent), Zel'dovich (finite)
- Flavour structure of pressure: DVMP!

QCD Factorization for DIS and DVCS (AND VM production)



- Manifestly spectral

$$\mathcal{H}(x_B) = \int_{-1}^1 dx \frac{H(x)}{x - x_B + i\epsilon}$$

- Extra dependence on ξ

$$\mathcal{H}(\xi) = \int_{-1}^1 dx \frac{H(x, \xi)}{x - \xi + i\epsilon}$$



Unphysical regions

- DIS : Analytical function – polynomial in $1/x_B$ if $1 \leq |X_B|$

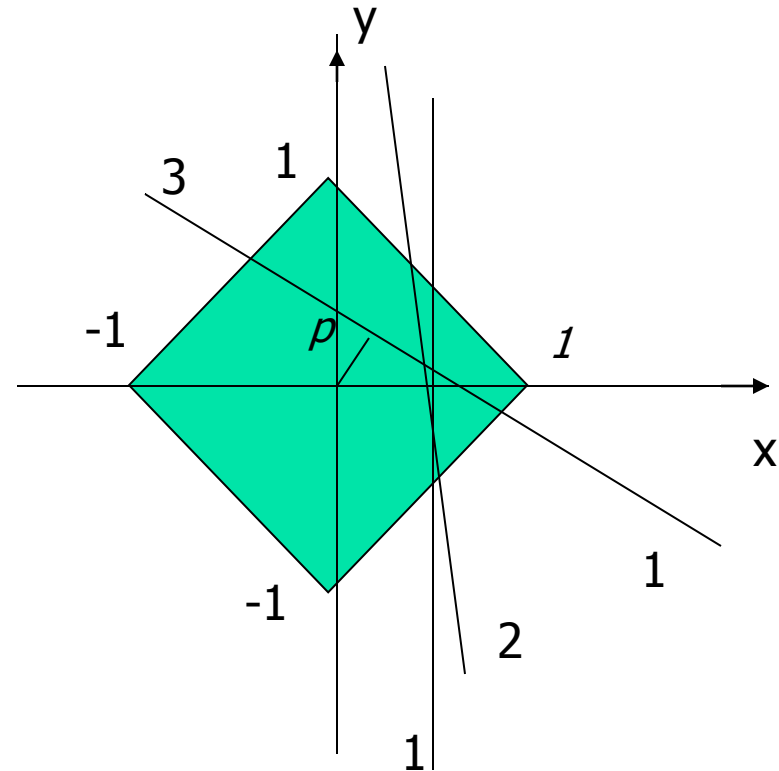
$$H(x_B) = -\int_{-1}^1 dx \sum_{n=0}^{\infty} H(x) \frac{x^n}{x_B^{n+1}}$$

- DVCS – additional problem of analytical continuation of $H(x, \xi)$
- Solved by using of Double Distributions Radon transform

$$H(z, \xi) = \int_{-1}^1 dx \int_{|x|-1}^{1-|x|} dy (F(x, y) + \xi G(x, y)) \delta(z - x - \xi y)$$

Double distributions and their integration

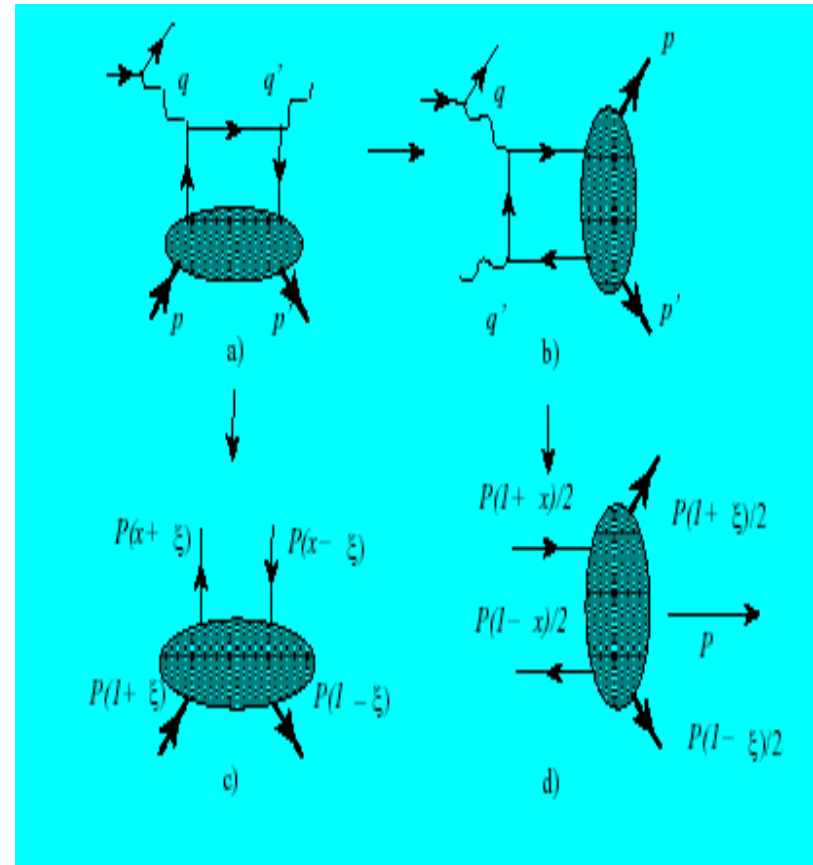
- Slope of the integration line-skewness
- Kinematics of DIS: $\xi = 0$
("forward") - vertical line (1)
- Kinematics of DVCS: $\xi < 1$
- line 2
- Line 3: $\xi > 1$ unphysical region - required to restore DD by inverse Radon transform: tomography



$$\begin{aligned}
 f(x, y) &= -\frac{1}{2\pi^2} \int_0^\infty \frac{dp}{p^2} \int_0^{2\pi} d\phi |\cos\phi| (H(p/\cos\phi + x + ytg\phi, t g\phi) - H(x + ytg\phi, t g\phi)) = \\
 &= -\frac{1}{2\pi^2} \int_{-\infty}^\infty \frac{dz}{z^2} \int_{-\infty}^\infty d\xi (H(z + x + y\xi, \xi) - H(x + y\xi, \xi))
 \end{aligned}$$

Crossing for DVCS and GPD

- DVCS \rightarrow hadron pair production in the collisions of real and virtual photons
- GPD \rightarrow Generalized Distribution Amplitudes
- Duality between s and t channels
(Polyakov, Shuvaev, Guzey, Vanderhaeghen)



GDA -> back to unphysical regions for DIS and DVCS

- Recall DIS

$$H(x_B) = - \int_{-1}^1 dx \sum_{n=0}^{\infty} H(x) \frac{x^n}{x_B^{n+1}}$$

- Non-positive powers of x_B

- DVCS

$$H(\xi) = - \int_{-1}^1 dx \sum_{n=0}^{\infty} H(x, \xi) \frac{x^n}{\xi^{n+1}}$$

- Polynomiality (general property of Radon transforms): moments - integrals in x weighted with x^n - are polynomials in $1/\xi$ of power $n+1$
- As a result, analyticity is preserved: only non-positive powers of ξ appear



Holographic property (OT'05)

Factorization
Formula

->

- Analyticity ->
Imaginary part ->
Dispersion relation:

$$\mathcal{H}(\xi) = \int_{-1}^1 dx \frac{H(x, \xi)}{x - \xi + i\epsilon}$$

$$\mathcal{H}(\xi) = \int_{-1}^1 dx \frac{H(x, x)}{x - \xi + i\epsilon}$$

$$\Delta\mathcal{H}(\xi) \equiv \int_{-1}^1 dx \frac{H(x, x) - H(x, \xi)}{x - \xi + i\epsilon}$$

- “Holographic”
equation (DVCS **AND**
VM)

$$= \sum_{n=1}^{\infty} \frac{1}{n!} \frac{\partial^n}{\partial \xi^n} \int_{-1}^1 H(x, \xi) dx (x - \xi)^{n-1} = \text{const}$$



Holographic property - II

- Directly follows from double distributions

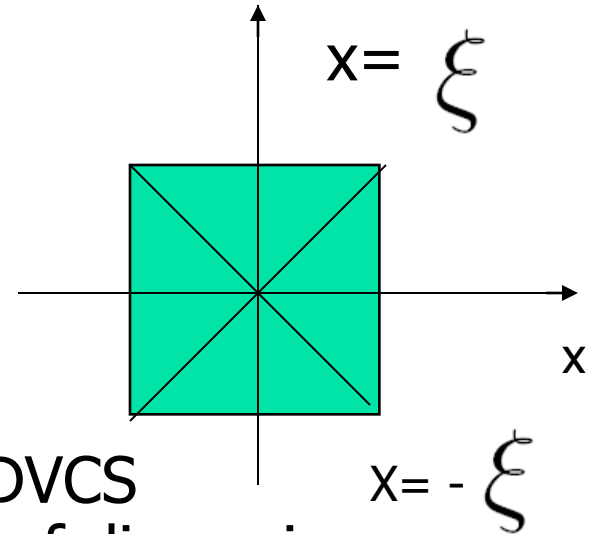
$$H(z, \xi) = \int_{-1}^1 dx \int_{|x|-1}^{1-|x|} dy (F(x, y) + \xi G(x, y)) \delta(z - x - \xi y)$$

- Constant is the SUBTRACTION one - due to the (generalized) Polyakov-Weiss term $G(x, y)$

$$\begin{aligned} \Delta \mathcal{H}(\xi) &= \int_{-1}^1 dx \int_{|x|-1}^{1-|x|} dy \frac{G(x, y)}{1 - y} \\ &= \int_{-\xi}^{\xi} dx \frac{D(x/\xi)}{x - \xi + i\epsilon} = \int_{-1}^1 dz \frac{D(z)}{z - 1} = \text{const} \end{aligned}$$

Holographic property - III

- 2-dimensional space \rightarrow 1-dimensional section!
- Momentum space: any relation to holography in coordinate space ?!
- Strategy (now adopted) of GPD's studies: start at diagonals (through SSA due to imaginary part of DVCS amplitude) and restore by making use of dispersion relations + subtraction constants





Holography vs NLO

- Depends on factorization scheme
- Special role of scheme preserving the coefficient function
- Nucleon as (scheme dependent) black hole – 3D information encoded in 2D

Pressure in hadron pairs production

- Back to GDA region
- -> moments of $H(x,x)$ - define the coefficients of powers of cosine! - $1/\xi$
- Higher powers of cosine ξ in t-channel - threshold in s-channel
- Larger for pion than for nucleon pairs because of less fast decrease at $x \rightarrow 1$
- Stability defines the sign of GDA and (via soft pion theorem) DA: work in progress

$$\begin{aligned} \mathcal{H}(\xi) &= - \int_{-1/\xi}^{1/\xi} dx \sum_{n=0}^{\infty} H(x, \xi) \frac{x^n}{\xi^{n+1}} \\ &= - \int_{-1/\xi}^{1/\xi} dx \sum_{n=0}^{\infty} H(x, x) \frac{x^n}{\xi^{n+1}} + \Delta \mathcal{H}. \end{aligned}$$

Analyticity of Compton amplitudes in energy plane (Anikin, OT'07)

- Finite subtraction implied

$$\operatorname{Re}\mathcal{A}(\nu, Q^2) = \frac{\nu^2}{\pi} \mathcal{P} \int_{\nu_0}^{\infty} \frac{d\nu'^2}{\nu'^2} \frac{\operatorname{Im}\mathcal{A}(\nu', Q^2)}{(\nu'^2 - \nu^2)} + \Delta \quad \Delta = 2 \int_{-1}^1 d\beta \frac{D(\beta)}{\beta - 1}$$

$$\Delta_{\text{CQM}}^p(2) = \Delta_{\text{CQM}}^n(2) \approx 4.4, \quad \Delta_{\text{latt}}^p \approx \Delta_{\text{latt}}^n \approx 1.1$$

- Numerically close to Thomson term for real proton (but NOT neutron) Compton Scattering!
- Duality (sum of squares vs square of sum; proton: $4/9+4/9+1/9=1$)?!



From D-term to pressure

- *Inverse -> 1st moment (model)*
- *Kinematical factor – moment of pressure $C \sim \langle p r^4 \rangle$ ($\langle p r^2 \rangle = 0$)*

M. Polyakov '03

$$T_{\mu\nu}^Q(\vec{r}, \vec{s}) = \frac{1}{2E} \int \frac{d^3\Delta}{(2\pi)^3} e^{i\vec{r}\cdot\vec{\Delta}} \langle p', S' | \hat{T}_{\mu\nu}^Q(0) | p, S \rangle$$

$$T_{ij}(\vec{r}) = s(r) \left(\frac{r_i r_j}{r^2} - \frac{1}{3} \delta_{ij} \right) + p(r) \delta_{ij}$$

- *Stable equilibrium $C > 0$:*



Loss of stability?

- $D=0$ -> extra node required (cf tensor distribution - Efremov, OT- mechanical analogy – c.m. and c.i.)
- Smooth decrease – two extra nodes
- + + + + -----
- + + + + + + + + ----- + + + + + -----
- $J=2$ (Talk of Barbara Pasquini, comment by Maxim – zeros of Bessel functs?!)



CONCLUSIONS/OUTLOOK

- “Macroscopical” aspects of GPDs
- Pressure of quark flavours/gluons – DVMP
- Pressure from TMDs (TMD/GPD relations)?
- Comparison to QCD matter (HIC)



BACKUP



Is D-term independent?

- Fast enough decrease at large energy -

$$> \quad \text{Re } \mathcal{A}(\nu) = \frac{\mathcal{P}}{\pi} \int_{\nu_0}^{\infty} d\nu'^2 \frac{\text{Im } \mathcal{A}(\nu')}{\nu'^2 - \nu^2} + \mathbf{C}_0.$$

$$\begin{aligned} \mathbf{C}_0 &= \Delta - \frac{\mathcal{P}}{\pi} \int_{\nu_0}^{\infty} d\nu'^2 \frac{\text{Im } \mathcal{A}(\nu')}{\nu'^2} \\ &= \Delta + \mathcal{P} \int_{-1}^1 dx \frac{H^{(+)}(x, x)}{x}. \end{aligned}$$

- FORWARD limit of Holographic equation

$$\begin{aligned} \Delta &= \mathcal{P} \int_{-1}^1 dx \frac{H^{(+)}(x, 0) - H^{(+)}(x, x)}{x} & \mathbf{C}_0(t) &= 2\mathcal{P} \int_{-1}^1 dx \frac{H(x, 0, t)}{x} \\ &= 2\mathcal{P} \int_{-1}^1 dx \frac{H(x, 0) - H(x, x)}{x}, \end{aligned}$$



“D – term” 30 years before...

- Cf Brodsky, Close, Gunion'72 (**seagull** \sim **pressure**) – but NOT DVMP
- D-term – a sort of renormalization constant
- May be calculated in effective theory if we know fundamental one
- OR
- Recover through special regularization procedure (D. Mueller)?



Vector mesons and EEP

- $J=1/2 \rightarrow J=1$. QCD SR calculation of Rho's AMM gives g close to 2.
- Maybe because of similarity of moments
- $g-2 = \langle E(x) \rangle$; $B = \langle xE(x) \rangle$
- Directly for charged Rho (combinations like $p+n$ for nucleons unnecessary!). Not reduced to non-extended EP:



EEP and AdS/QCD

- Recent development – calculation of Rho formfactors in Holographic QCD (Grigoryan, Radyushkin)
- Provides $g=2$ identically!
- Experimental test at time –like region possible



EEP and Siverson function

- Siverson function – process dependent (effective) one
- T-odd effect in T-conserving theory- phase
- FSI – Brodsky-Hwang-Schmidt model
- Unsuppressed by M/Q twist 3
- Process dependence- colour factors
- After Extraction of phase – relation to universal (T-even) matrix elements



EEP and Sivers function -II

- Qualitatively similar to OAM and Anomalous Magnetic Moment (talk of S. Brodsky)
- Quantification : weighted TM moment of Sivers PROPORTIONAL to GPD E (OT'07, **hep-ph/0612205**) : $x f_T(x) : xE(x)$
- Burkardt SR for Sivers functions is then related to Ji's SR for E and, in turn, to Equivalence Principle

$$\sum_{q,G} \int dx x f_T(x) = \sum_{q,G} \int dx x E(x) = 0$$



EEP and Sivers function for deuteron

- EEP - smallness of deuteron Sivers function
- Cancellation of Sivers functions – separately for quarks (before inclusion gluons)
- Equipartition + small gluon spin – large longitudinal orbital momenta (BUT small transverse ones –Brodsky, Gardner)

Another relation of Gravitational FF and NP QCD (first reported at 1992: **hep-ph/9303228**)

- BELINFANTE (relocalization) invariance :

decreasing in coordinate –

$$M^{\mu,\nu\rho} = \frac{1}{2} \epsilon^{\mu\nu\rho\sigma} J_{S\sigma}^5 + x^\nu T^{\mu\rho} - x^\rho T^{\mu\nu}$$

smoothness in momentum space

$$M^{\mu,\nu\rho} = x^\nu T_B^{\mu\rho} - x^\rho T_B^{\mu\nu}$$

- Leads to absence of massless pole in singlet channel – U_A(1)

$$\epsilon_{\mu\nu\rho\alpha} M^{\mu,\nu\rho} = 0.$$

- Delicate effect of NP QCD

$$(g_{\rho\nu} g_{\alpha\mu} - g_{\rho\mu} g_{\alpha\nu}) \partial^\rho (J_{5S}^\alpha x^\nu) = 0$$

- Equipartition – deeply related to relocalization

$$q^2 \frac{\partial}{\partial q^\alpha} \langle P | J_{5S}^\alpha | P + q \rangle = (q^\beta \frac{\partial}{\partial q^\beta} - 1) q_\gamma \langle P | J_{5S}^\gamma | P + q \rangle$$

$$\langle P, S | J_\mu^5(0) | P + q, S \rangle = 2MS_\mu G_1 + q_\mu (Sq) G_2, \\ q^2 G_2|_0 = 0$$

invariance by QCD evolution