

Parton distribution functions from lattice QCD

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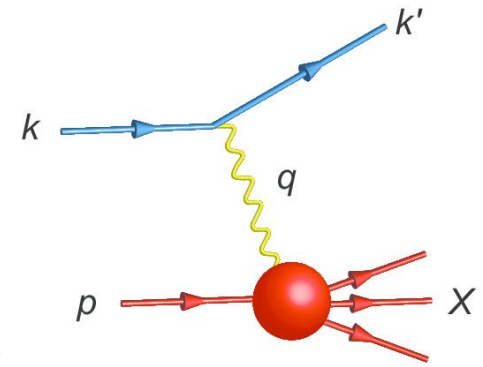
Introduction

- **Deep-inelastic lepton-nucleon scattering**
 - DIS X-sec is an **incoherent sum** of **virtual photon-quark partonic X-sec**, weighted by the **probability of finding each quark in the proton** with a given longitudinal mom. fraction

$$\frac{d^2\sigma}{dx dQ^2} \sim \sum_{q,\bar{q}} \int_x^1 \frac{dz}{z} f_q(z) \hat{\sigma}_{q\gamma^* \rightarrow X}(x/z)$$

Parton distribution function (PDF),
non-perturbative, universal

Perturbatively calculable,
asymptotic freedom of QCD



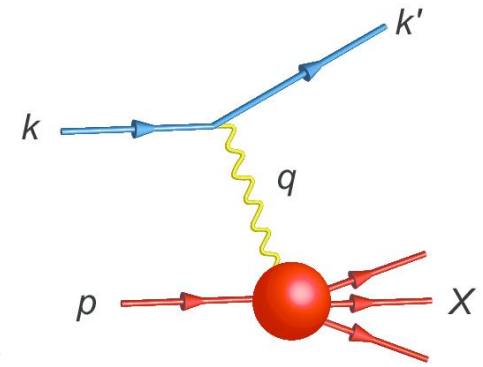
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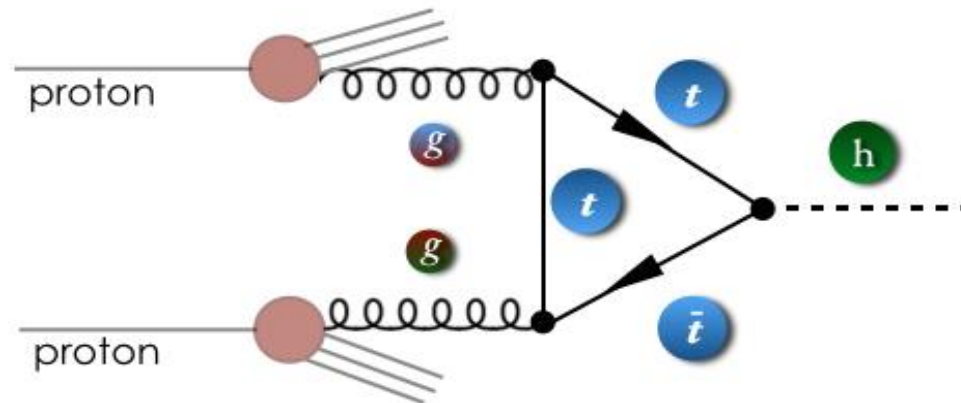
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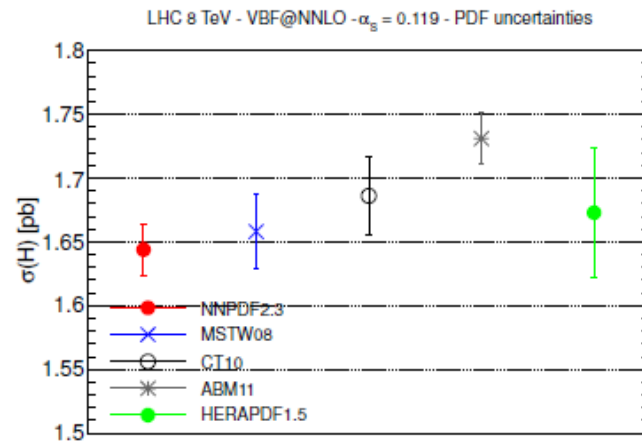
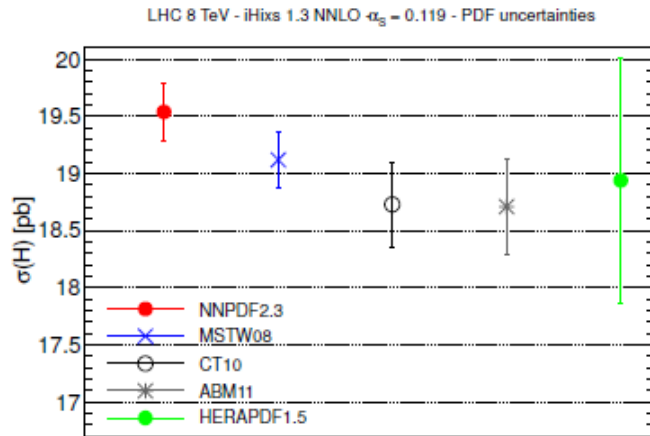


- **PP Collision (LHC)**

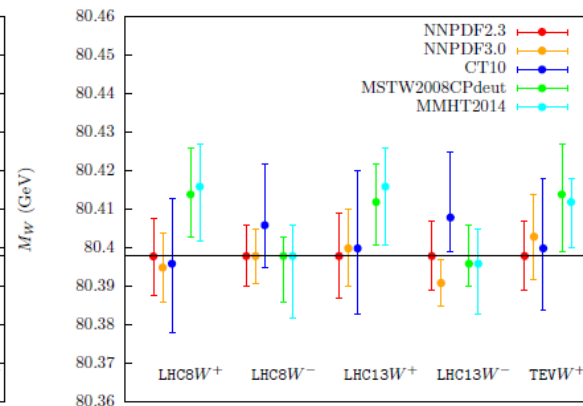
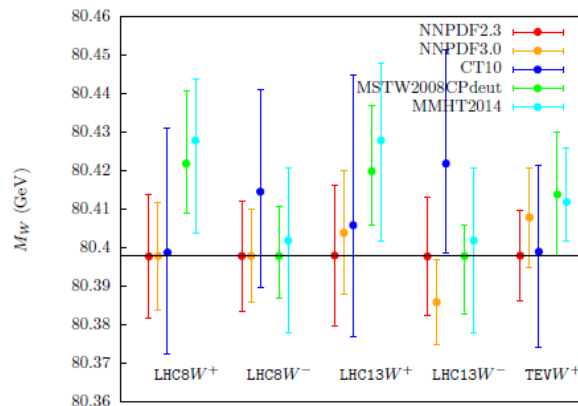


Introduction

- Precise knowledge of PDFs helps
 - Acquire a better understanding of the SM
 - Higgs production and self coupling



- W mass determination



- Disentangle new physics effects from SM

Introduction

- Precise knowledge of PDFs helps
 - Three-dim. imaging of the nucleon

Electron Ion Collider: The Next QCD Frontier

Imaging of the proton

*How are the **sea** quarks and gluons, and their spins, distributed in space and momentum inside the nucleon?*

EIC White Paper, 1212.1701



- **Ab initio** determination of PDFs from lattice QCD is desirable and provides complementary information to global analysis program

Introduction

- Parton picture arises in high-energy collisions where hadrons/probe move nearly at the speed of light, or with **infinite momentum**
- Parton physics usually formulated in terms of **light-cone quantization** [Dirac, Wilson]
 - light-cone coordinates $\xi^\pm = (t \pm z)/\sqrt{2}$
 - Example: [Collins and Soper, NPB 82']

$$q(x, \mu^2) = \int \frac{d\xi^-}{4\pi} e^{-ix\xi^- P^+} \langle P | \bar{\psi}(\xi^-) \gamma^+ \exp\left(-ig \int_0^{\xi^-} d\eta^- A^+(\eta^-)\right) \psi(0) | P \rangle$$

- Difficult to access from ***ab initio*** approaches like lattice QCD

$$t^2 - x^2 = 0 \rightarrow -\tau^2 - x^2 = 0$$

Introduction

- Lattice effort on calculating Mellin moments of PDFs

Quark density/unpolarized

$$\langle x^n \rangle_q = \int_{-1}^1 dx x^n q(x)$$

Helicity
longitudinally polarized

$$\langle x^n \rangle_{\Delta q} = \int_{-1}^1 dx x^n \Delta q(x)$$

Transversity
transversely polarized

$$\langle x^n \rangle_{\delta q} = \int_{-1}^1 dx x^n \delta q(x)$$

most well known

very poorly known

- PDFs can be fully reconstructed if all their moments are known

Introduction

- Lattice effort on calculating Mellin moments of PDFs

Full details of lattice-QCD calculations of higher moments of unpolarized and polarized PDFs.

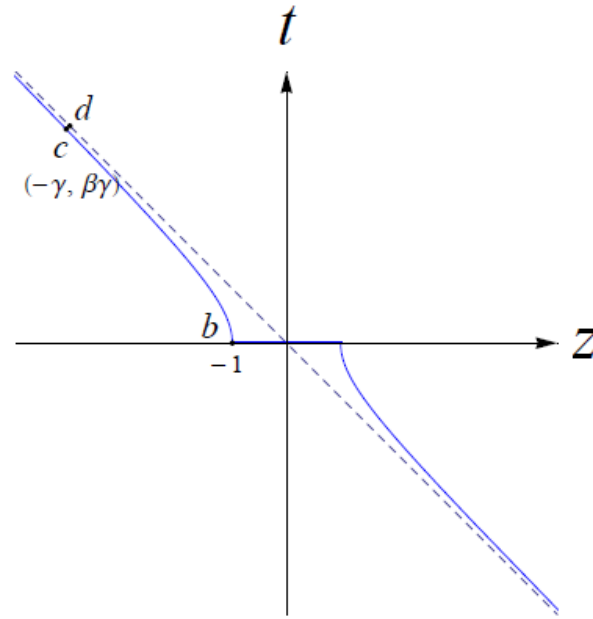
Ref.	Observables	Sea quarks	Valence quarks	Renormalization
LHPC '10 ^b [262]	$\langle x \rangle_{u^+ - d^+}, \langle x^2 \rangle_{u^- - d^-},$ $g_A, \langle x \rangle_{\Delta u^- - \Delta d^-}, \langle x^2 \rangle_{\Delta u^+ - \Delta d^+}$	2+1 staggered	domain wall	one-loop Z_O/Z_A
χ QCD '09 [280]	$\langle x \rangle_{u^+, d^+, s^+}$ (superseded by χ QCD '15), $\langle x^2 \rangle_{u^-, d^-, s^-}$	quenched	Wilson	one loop
LHPC '08 [295]	superseded by LHPC '10			
QCDSF '05c [313]	$\langle x^2 \rangle_{\Delta u^+ - \Delta d^+}$	2 clover	clover	Rome-Southampton
QCDSF '05b [93]	$\langle x \rangle_{u^+ - d^+}, \langle x^2 \rangle_{u^- - d^-}, \langle x^3 \rangle_{u^+ - d^+}$	quenched	clover	Rome-Southampton
QCDSF '05a ^a [314]	$\langle x \rangle_{u^+ - d^+}, \langle x^2 \rangle_{u^- - d^-}, \langle x^3 \rangle_{u^+ - d^+}$	2 clover	clover	one loop
LHPC and SESAM '02 [279]	$\langle x \rangle_{u^+ - d^+}, \langle x^2 \rangle_{u^- - d^-}, \langle x^3 \rangle_{u^+ - d^+},$ $g_A, \langle x \rangle_{\Delta u^- - \Delta d^-}, \langle x^2 \rangle_{\Delta u^+ - \Delta d^+}$	2 Wilson and quenched	Wilson	one loop
QCDSF '01 [315]	$\langle x^2 \rangle_{\Delta u^+ - \Delta d^+}$	quenched	clover	Rome-Southampton
QCDSF '96 [100]	$\langle x \rangle_{u^+ - d^+}, \langle x^2 \rangle_{u^- - d^-}, \langle x^3 \rangle_{u^+ - d^+},$ $g_A, \langle x^2 \rangle_{\Delta u^+ - \Delta d^+}$	quenched	Wilson	one loop

- Difficulties in going beyond the first few moments
 - Power divergent mixing between high and low moments operators
 - Increasing stochastic noise for operators with many derivatives

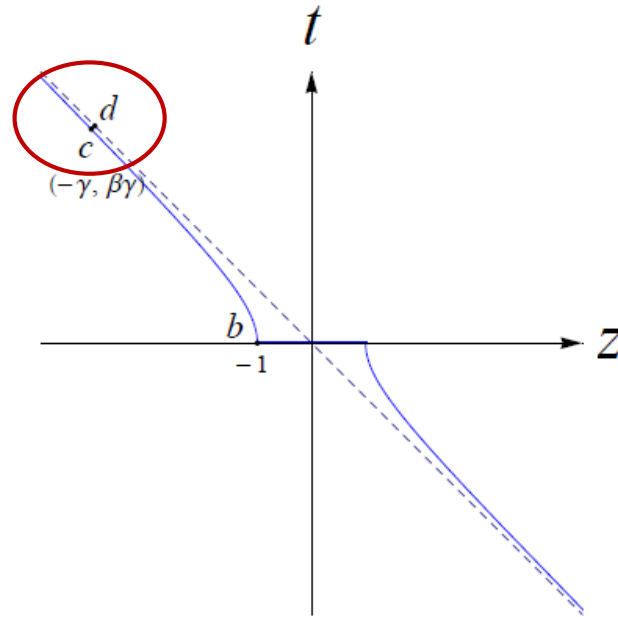
Introduction

- Obstacle to access PDFs on the lattice
 - They are defined on the light-cone
- However, they were originally introduced by Feynman as the **infinite momentum limit** of **frame-dependent** quantities
$$q(x) = \lim_{P_Z \rightarrow \infty} \tilde{q}(x, P_Z)$$
- Boost to infinite momentum leads to light-cone correlations
- If we can construct a $\tilde{q}(x, P_Z)$ such that it is calculable on the lattice, and all P_Z -dependence can be systematically removed
- Then we can calculate $q(x)$!

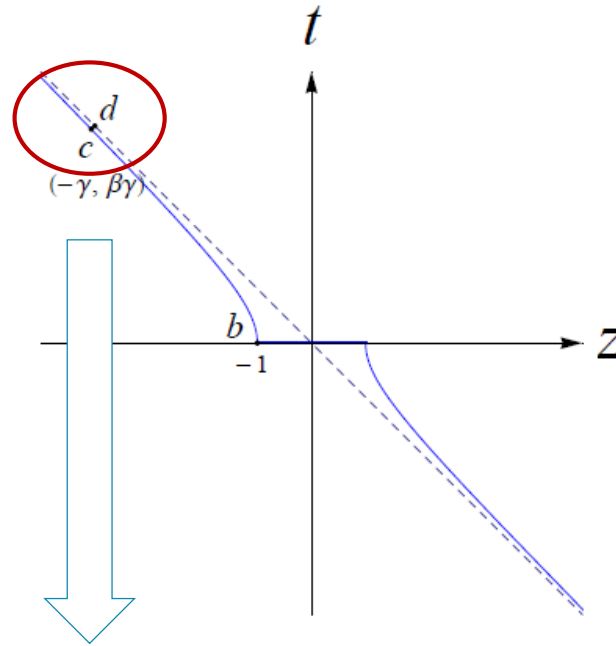
Introduction



Introduction

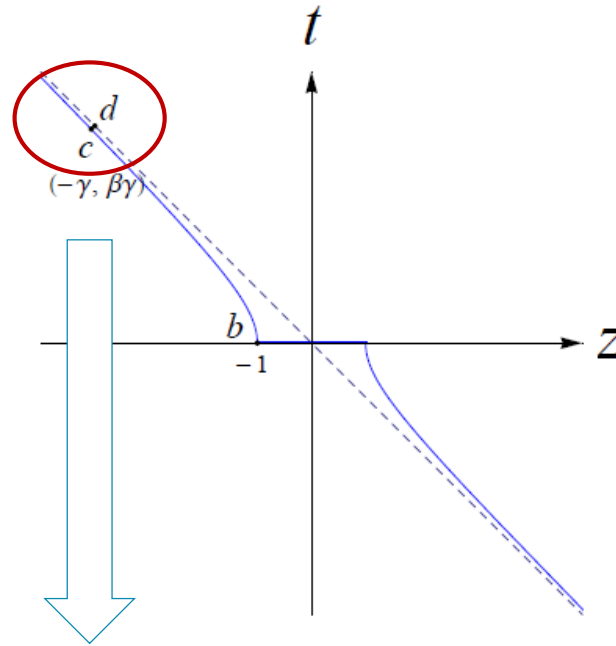


Introduction



- Systematic connection through **large momentum effective theory (LaMET)** [Ji, PRL 13', Sci. China Phys. Mech. Astron. 14']

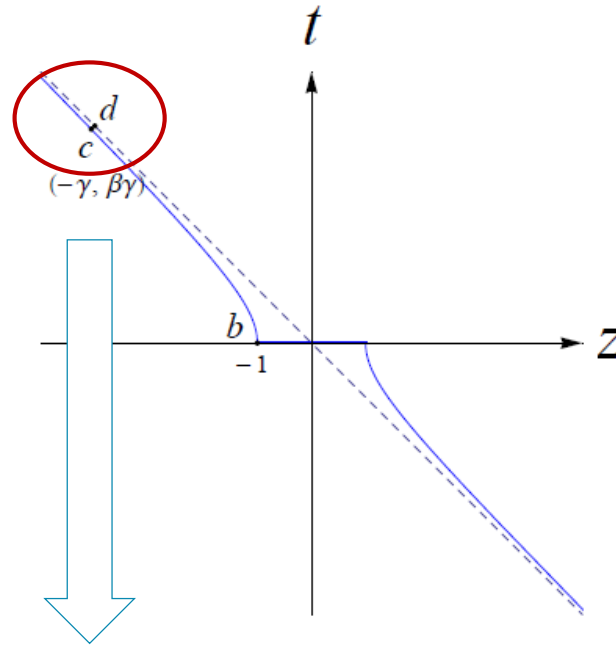
Introduction



- Systematic connection through **large momentum effective theory (LaMET)** [Ji, PRL 13', Sci. China Phys. Mech. Astron. 14']
 - Appropriately chosen $\tilde{q}(x, P_z)$ can be calculated on the Euclidean lattice, e.g.

$$\tilde{q}(x, \Lambda, P^z) = \int_{-\infty}^{\infty} \frac{dz}{4\pi} e^{izkz} \langle P | \bar{\psi}(0, 0_{\perp}, z) \gamma^z \exp \left(-ig \int_0^z dz' A^z(0, 0_{\perp}, z') \right) \psi(0) | P \rangle$$

Introduction



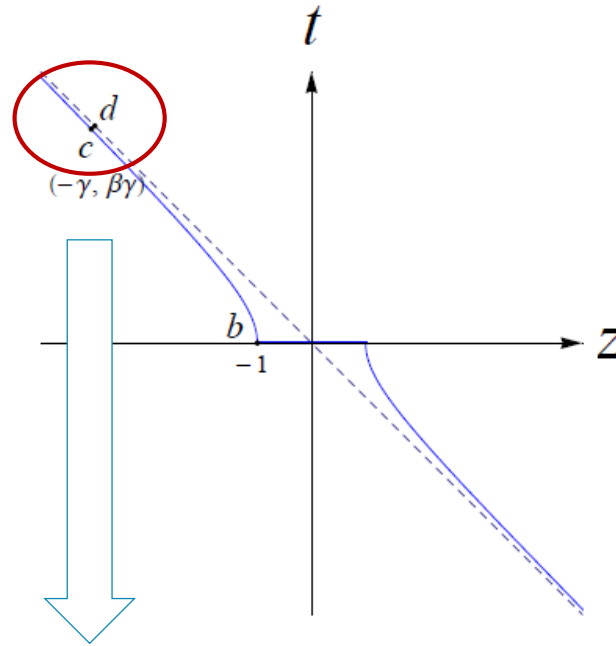
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- A finite but large P_z already offers a good approximation, where **(leading) frame-dependence can be removed through a factorization formula**

$$\tilde{q}(x, P_z, p_z^R, \mu_R) = \int_{-1}^1 \frac{dy}{|y|} C \left(\frac{x}{y}, r, \frac{yP_z}{\mu}, \frac{yP_z}{p_z^R} \right) q(y, \mu) + \mathcal{O} \left(\frac{M^2}{P_z^2}, \frac{\Lambda_{\text{QCD}}^2}{P_z^2} \right)$$

Introduction



- Systematic connection through **large momentum effective theory (LaMET)** [Ji, PRL 13', Sci. China Phys. Mech. Astron. 14']
- **Parton model is an effective theory for the nucleon moving at large momentum**

Related proposals

- **They all share the same property of computing correlations at spacelike separations**
- Current-current correlation functions
 - [Liu and Dong, PRL 94']
 - [Detmold and Lin, PRD 06']
 - [Braun and Müller, EPJC 08']
 - [Davoudi and Savage, PRD 12']
 - [Chambers et al., PRL 17']
- Lattice cross sections
 - [Ma and Qiu, 14' & PRL 17']
- Ioffe-time /pseudo-distribution
 - [Radyushkin, PRD 17']

PDFs from LaMET

Bare lattice
matrix element

Non-pert. Renorm.

renormalized
matrix element

Fourier transform

Quasi-PDF

Factorization

PDF

Ji, JHZ, Zhao, PRL 18'
Ishikawa et al, PRD 17'
Green et al, PRL 18'
Stewart, Zhao, PRD 18'
Chen, JHZ et al, PRD 18'
Alexandrou et al, NPB 17'
Monahan, Orginos, JHEP 17'
JHZ et al, 18' & Wang, JHZ et al, 19'

Cont. limit,

Ji, PRL 13'
Xiong, Ji, JHZ, Zhao, PRD 14'
Chen, JHZ et al, PRD 18'
Stewart, Zhao, PRD 18'
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$$\tilde{h}(z, P_z, a^{-1}) = \frac{1}{2P^0} \langle P | O_{\gamma^t}(z) | P \rangle$$

$$O_{\Gamma}(z) = \bar{\psi}(z) \Gamma U(z, 0) \psi(0)$$

$$U(z, 0) = P \exp \left(-ig \int_0^z dz' A_z(z') \right)$$

$$\tilde{h}_R(z, P_z, p_z^R, \mu_R)$$

$$= Z^{-1}(z, p_z^R, a^{-1}, \mu_R) \tilde{h}(z, P_z, a^{-1}) \Big|_{a \rightarrow 0}$$

$$Z(z, p_z^R, a^{-1}, \mu_R) = \frac{\sum_s \langle p, s | O_{\gamma^t}(z) | p, s \rangle}{\sum_s \langle p, s | O_{\gamma^t}(z) | p, s \rangle_{\text{tree}}} \Big|_{\substack{p^2 = -\mu_R^2 \\ p_z = p_z^R}}$$

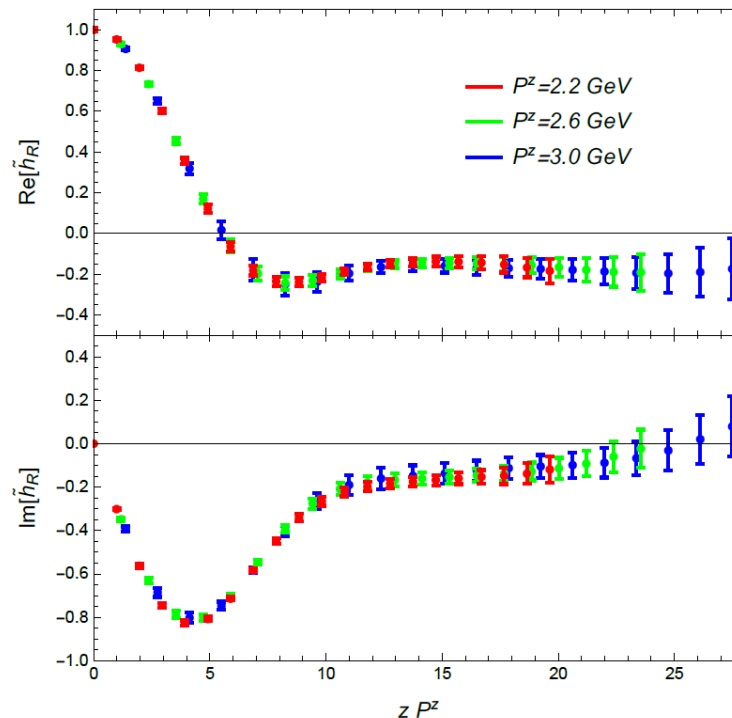
PDFs from LaMET

Bare lattice
matrix element

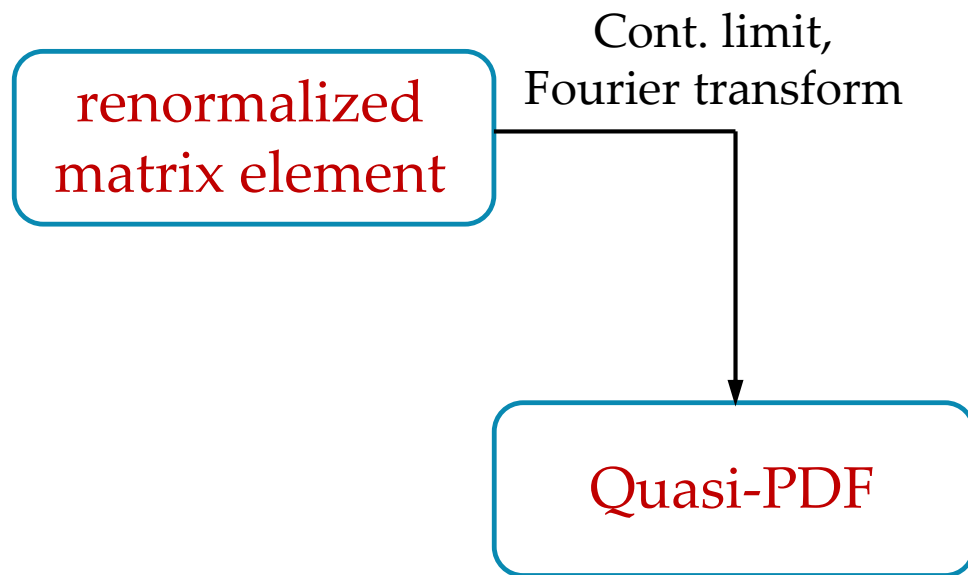
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PDFs from LaMET



$$\tilde{q}_R(x, P_z, p_z^R, \mu_R) = P_z \int \frac{dz}{2\pi} e^{ixP_z z} \tilde{h}_R(z, P_z, p_z^R, \mu_R)$$

PDFs from LaMET

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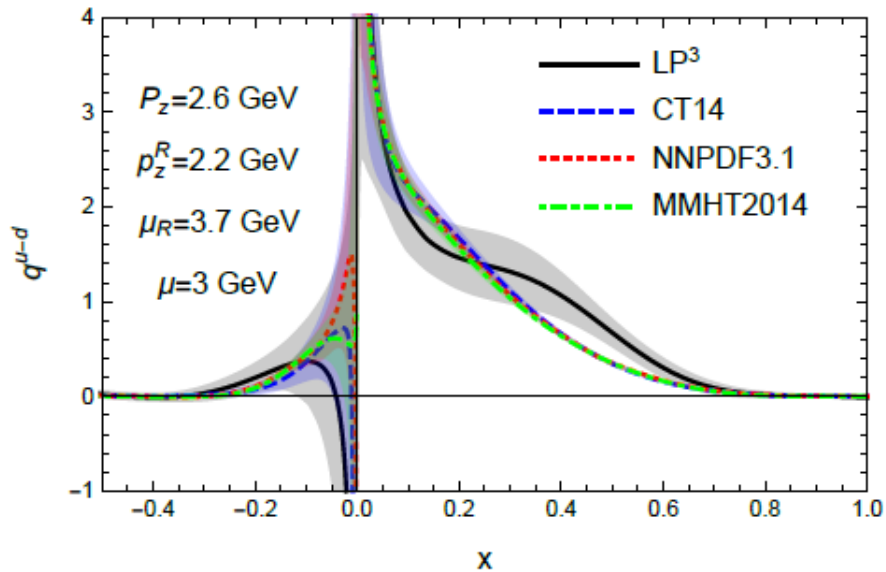
Quasi-PDF

Factorization

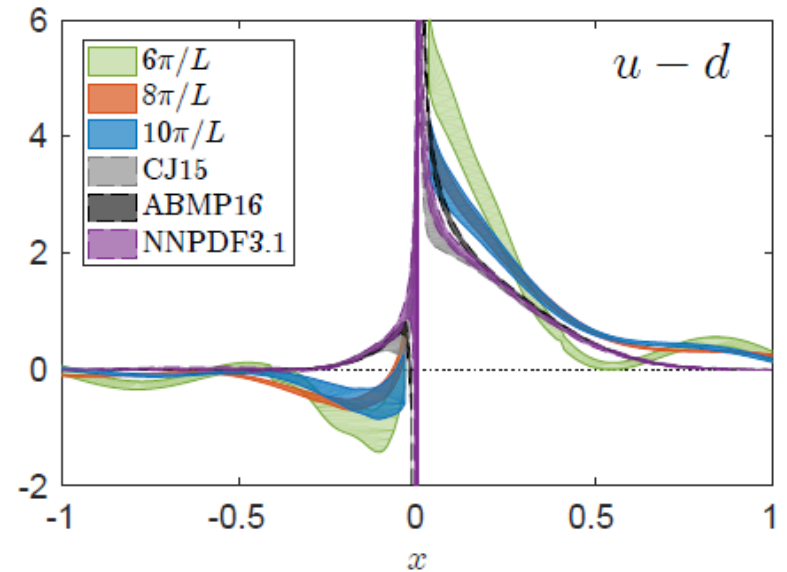
Ji, PRL 13'
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Chen, JHZ et al, PRD 18'
Stewart, Zhao, PRD 18'
Izubuchi et al, PRD 18'
Wang, JHZ et al, 19'

PDF

Isovector quark unpolarized PDF



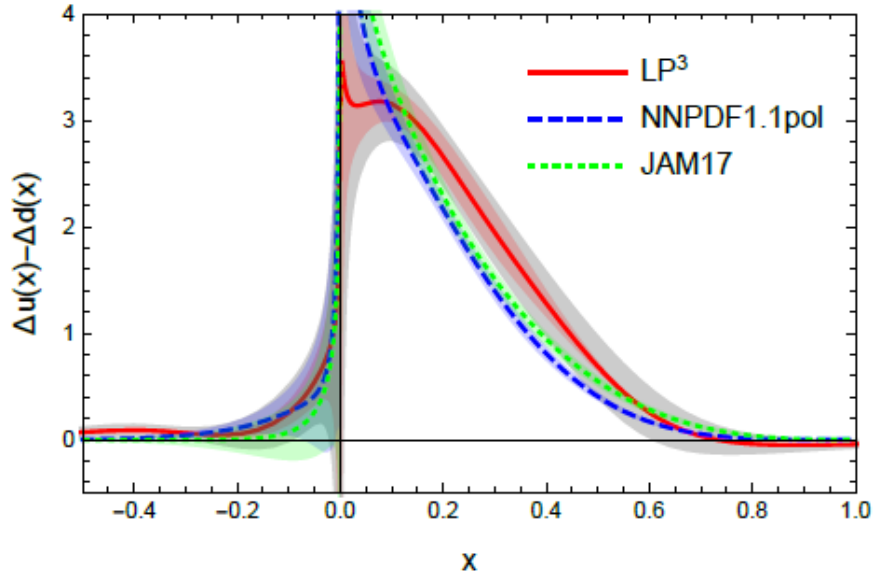
Liu, JHZ et al, 1803.04393 & 1807.06566,
 $m_\pi \approx 135 \text{ MeV}, a = 0.09 \text{ fm}, L \approx 5.8 \text{ fm}$



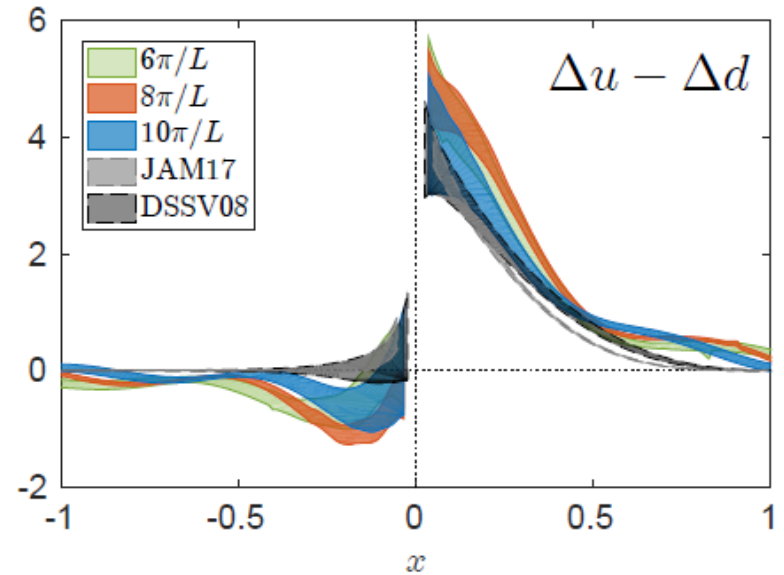
Alexandrou et al, 1803.02685, $m_\pi \approx$
130 MeV, $a = 0.094 \text{ fm}, L \approx 4.5 \text{ fm}$

- Systematics
 - Higher-order + higher-twist corrections
 - Continuum extrapolation
 - Discretization errors
 - Finite volume effects

Isovector quark helicity PDF



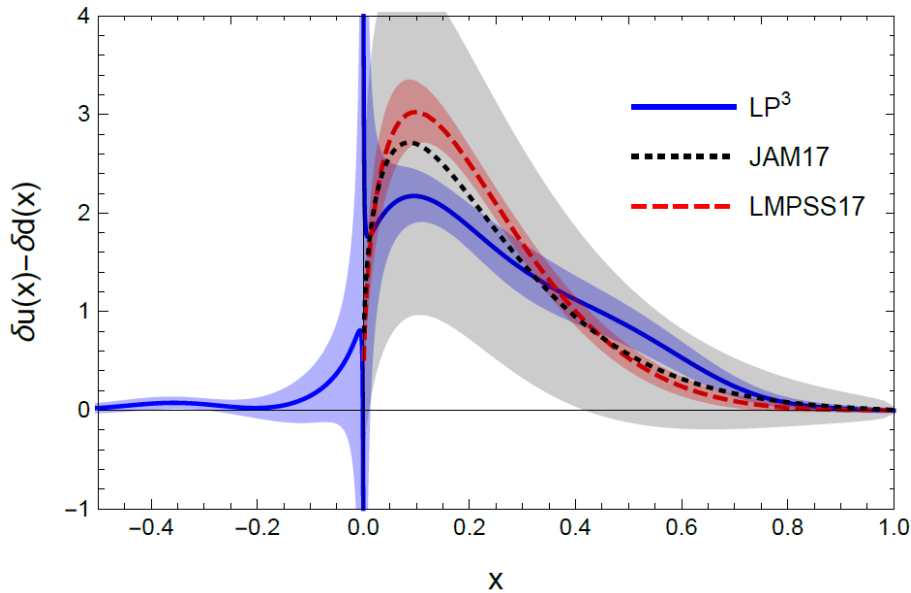
Lin, JHZ et al, 1807.07431, $m_\pi \approx 135 \text{ MeV}$, $a = 0.09 \text{ fm}$, $L \approx 5.8 \text{ fm}$



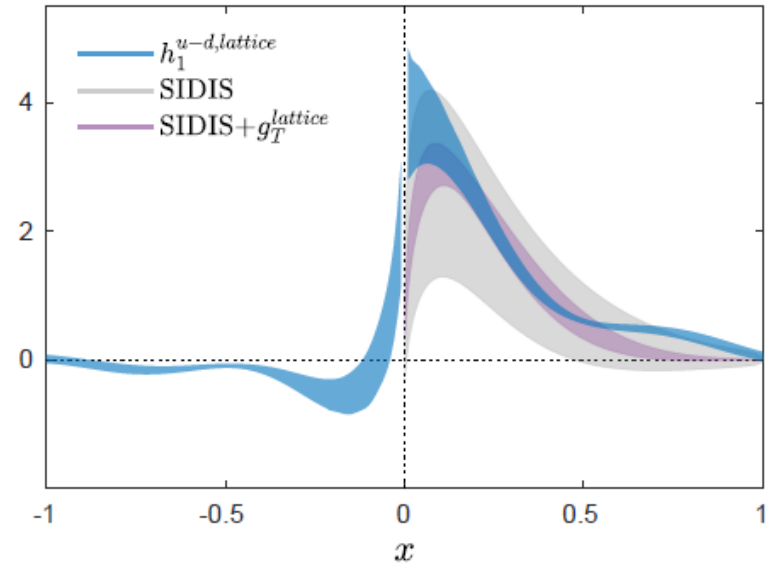
Alexandrou et al, 1803.02685, $m_\pi \approx 130 \text{ MeV}$, $a = 0.094 \text{ fm}$, $L \approx 4.5 \text{ fm}$

- Systematics
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 - Continuum extrapolation
 - Discretization errors
 - Finite volume effects

Isovector quark transversity PDF



LP3, $m_\pi \approx 135 \text{ MeV}$, $a = 0.09 \text{ fm}$, $L \approx 5.8 \text{ fm}$



Alexandrou et al, 1807.00232, $m_\pi \approx 130 \text{ MeV}$, $a = 0.094 \text{ fm}$, $L \approx 4.5 \text{ fm}$

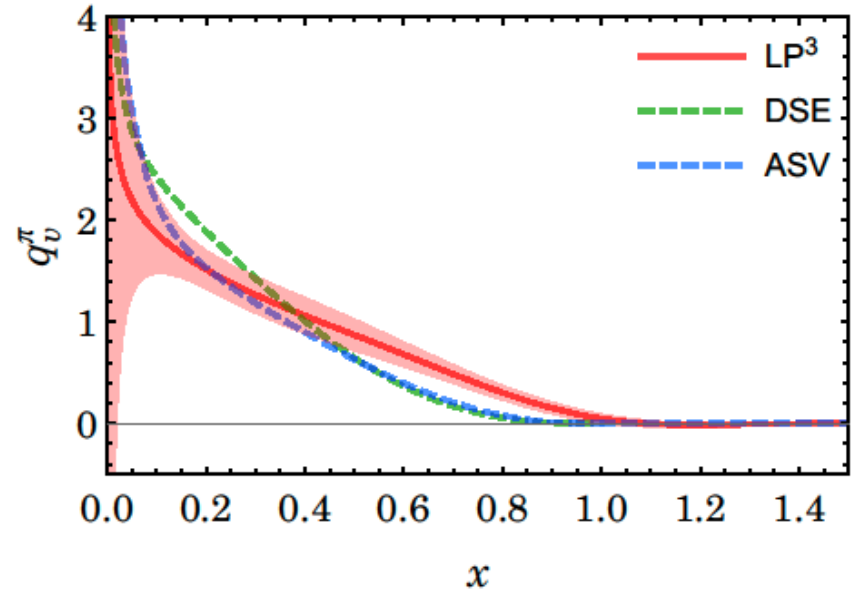
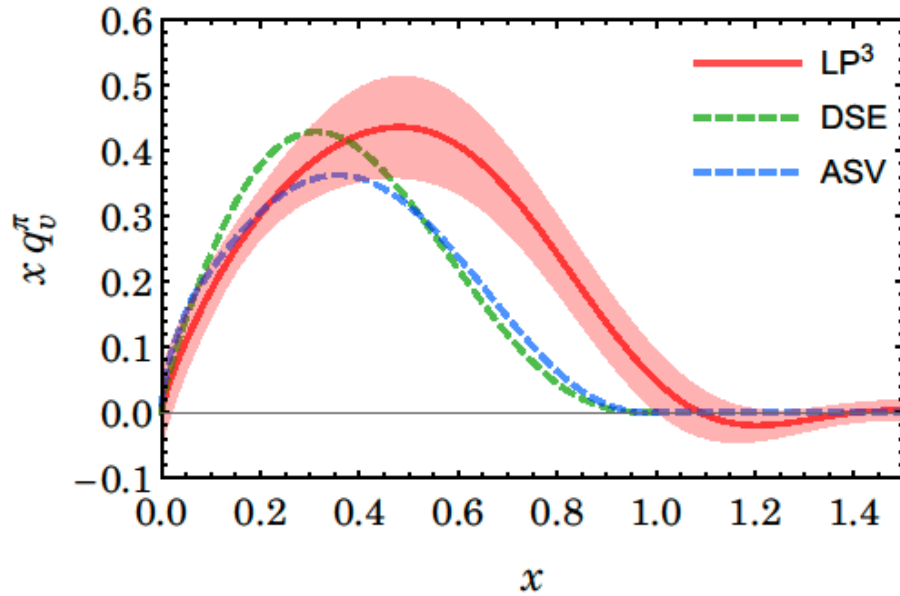
- Systematics
 - Higher-order + higher-twist corrections
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Summary and outlook

- Recent years have witnessed rapid progress on direct computations of x -dependence of hadron structure from lattice QCD
- **Large momentum effective theory** & related proposals
- Applications to nucleon PDFs (isovector quark), as well as meson PDFs & DAs have yielded encouraging results
- Flavor-singlet quark PDF and Gluon PDF
 - **Renormalization and factorization**
 - [JHZ et al, 1808.10824] (see also [Li, Ma, Qiu, 1809.01836])
 - [Wang, JHZ et al, 1904.00978]
- GPDs, TMDs ...

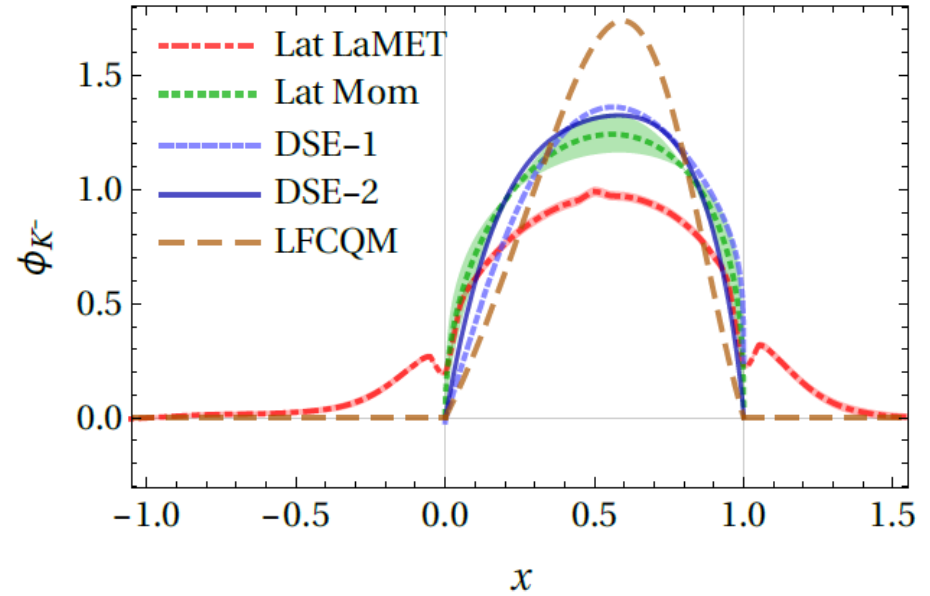
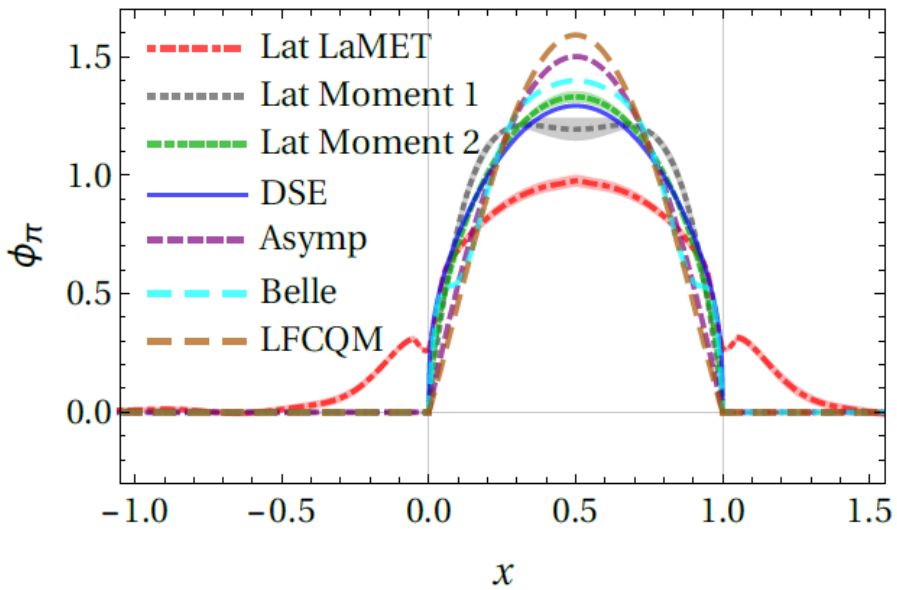
Backup Slides

Pion valence quark PDF



JHZ et al, 1804.01483, $m_\pi \approx 310 \text{ MeV}$, $a = 0.12 \text{ fm}$, $L \approx 3 \text{ fm}$

Meson distribution amplitudes



Chen, JHZ et al, 1712.10025, $m_\pi \approx 310 \text{ MeV}$, $a = 0.12 \text{ fm}$, $L \approx 3 \text{ fm}$