



**XXVII International Workshop on Deep-Inelastic
Scattering and Related Subjects
8-12 April 2019
Torino, Italy**

Some recent results related to low-x and forward physics

J. G. Contreras
Czech Technical University in Prague

April 8, 2019, Torino



Disclaimers

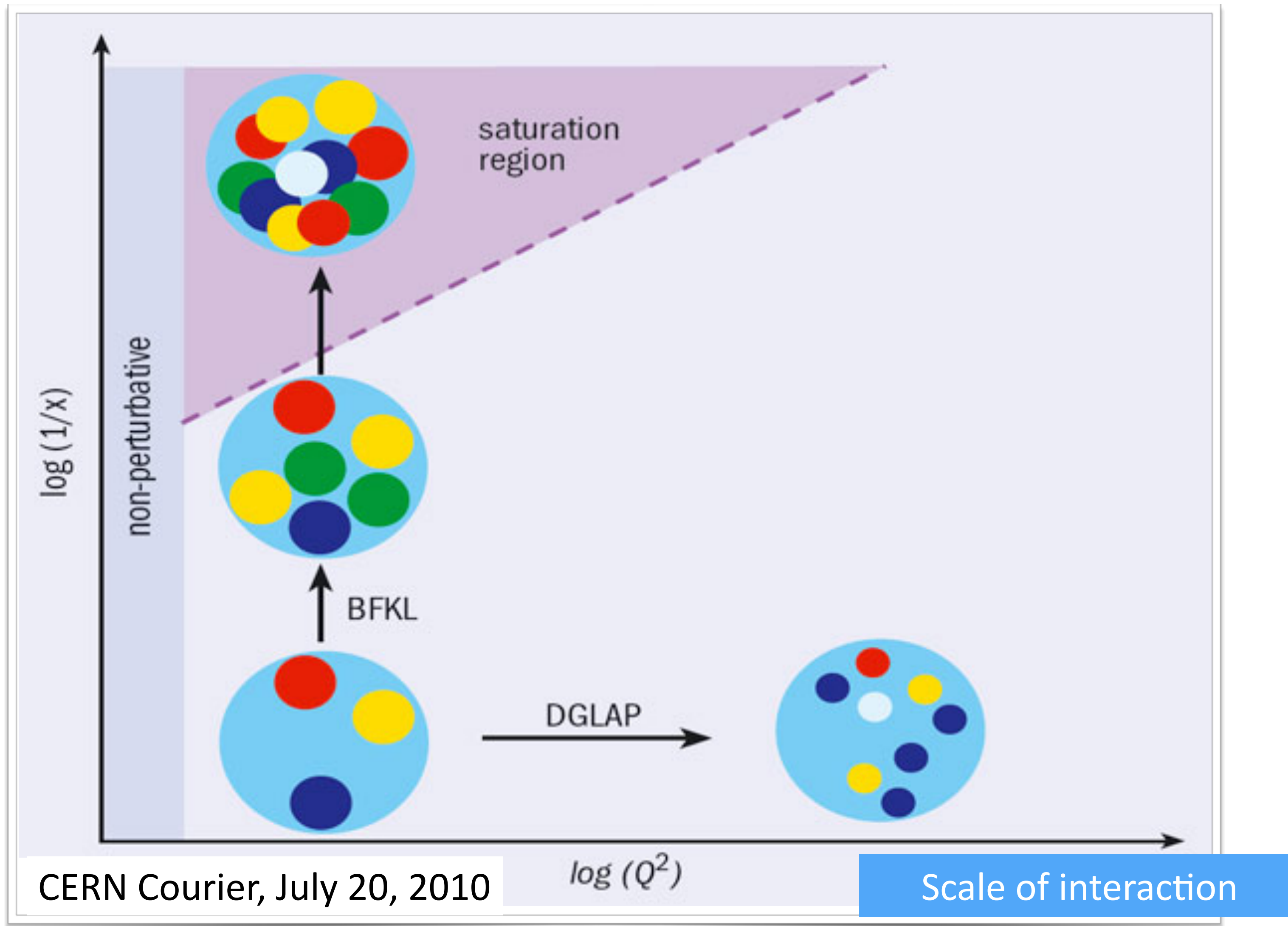
Selection of content based on:

- ◉ My ignorance of other topics.
- ◉ Not presented before in plenary talks in the few previous DIS meetings.
- ◉ Of potential interest to the DIS community.

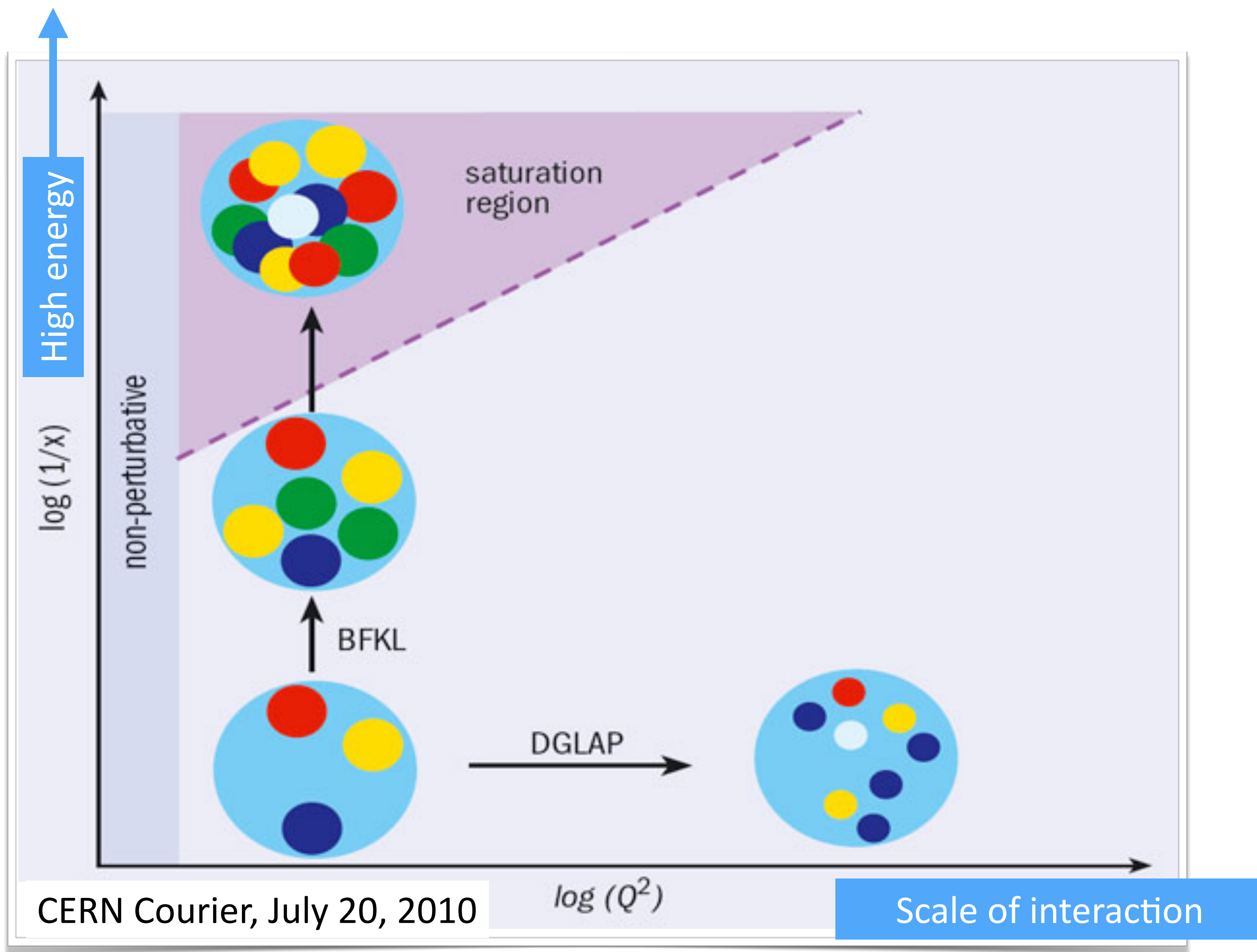
Apologies to the many nice results that I do not cover.

The structure of nuclei at low x

Low-x and the structure of hadrons

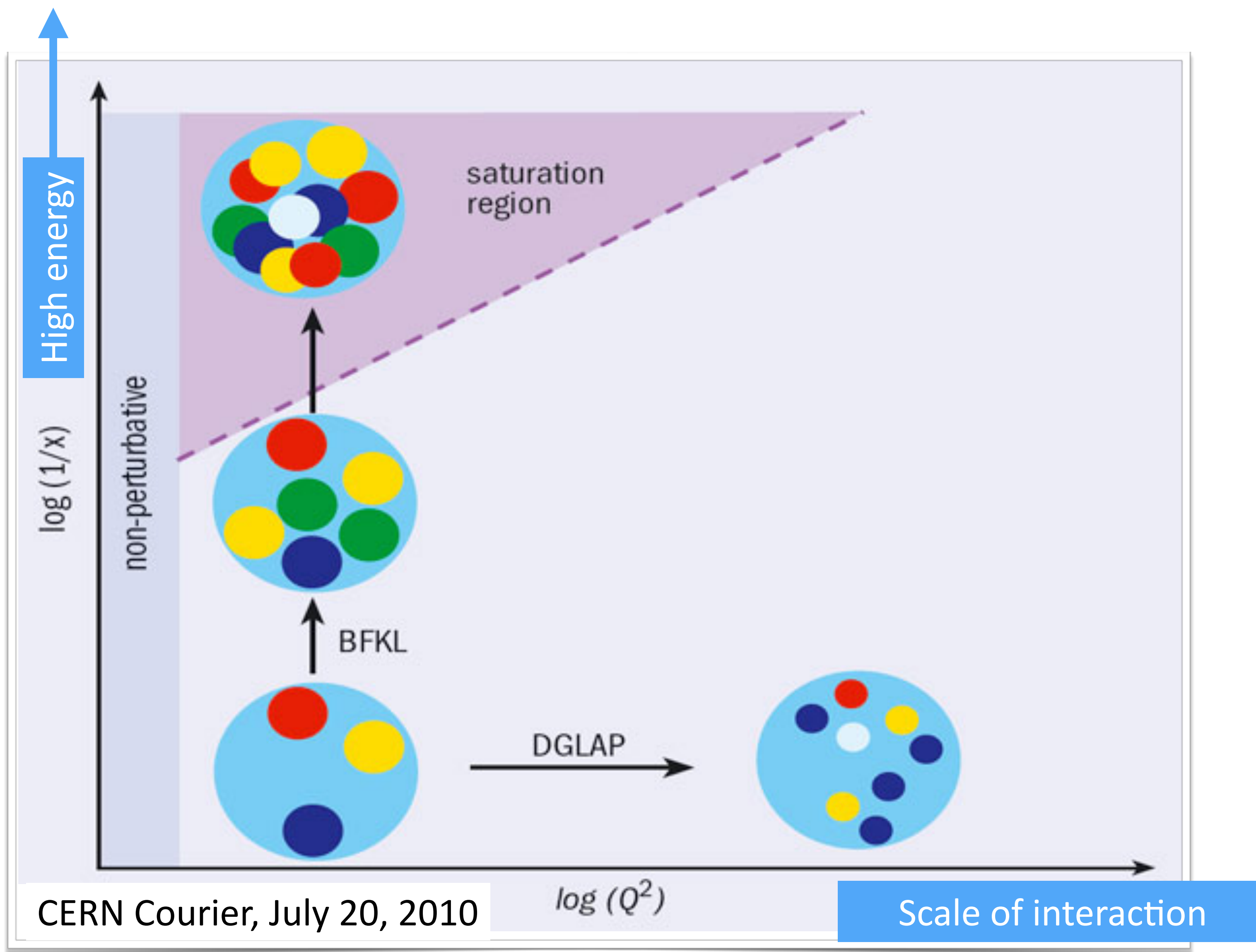


Low-x and the structure of hadrons



The high-energy limit of pQCD corresponds to the low x limit.

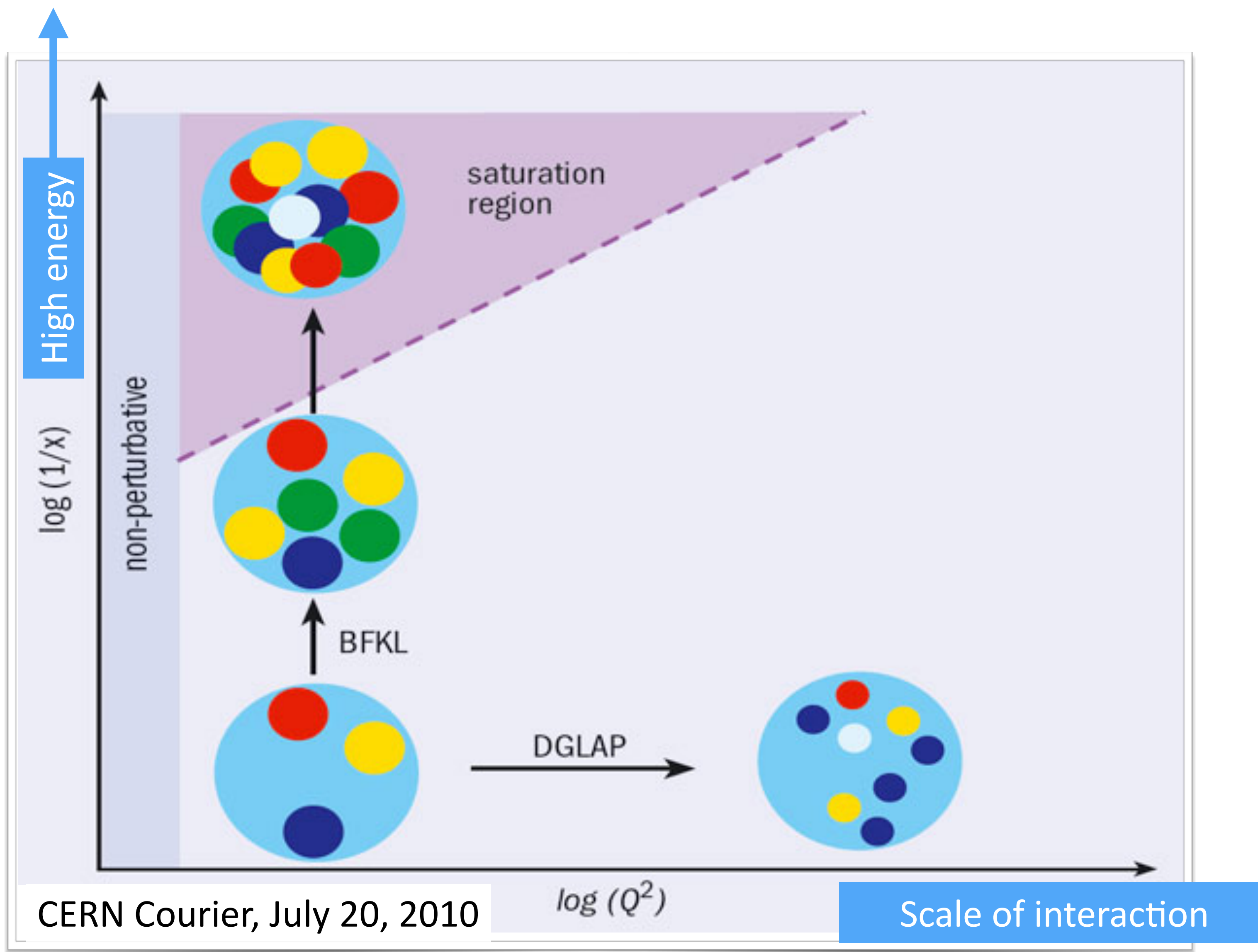
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In this limit, the structure of hadrons is expected to reach the saturation region.

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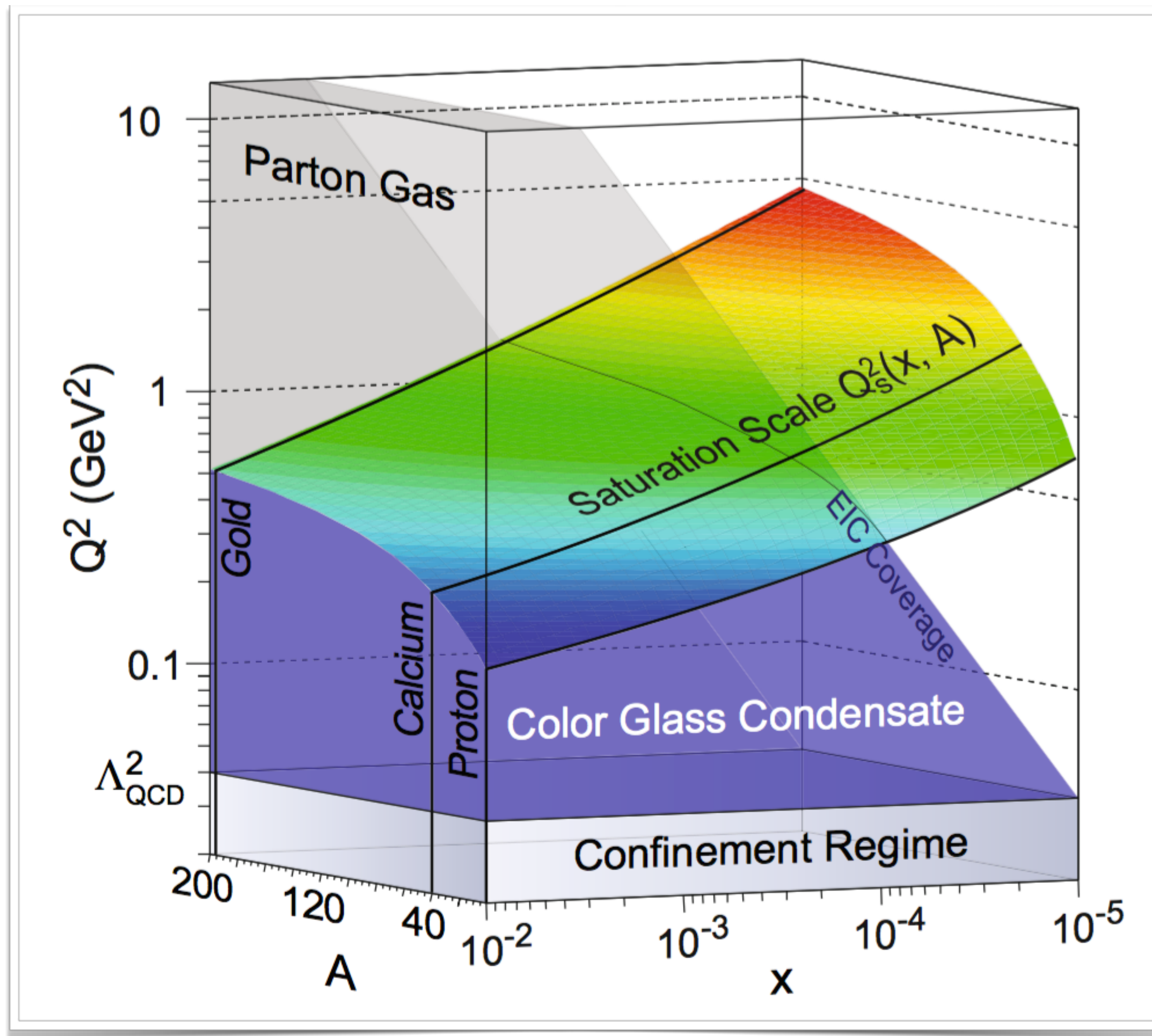
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One of the most interesting questions nowadays in pQCD is the precise location of the saturation region.

Low x and nuclear targets

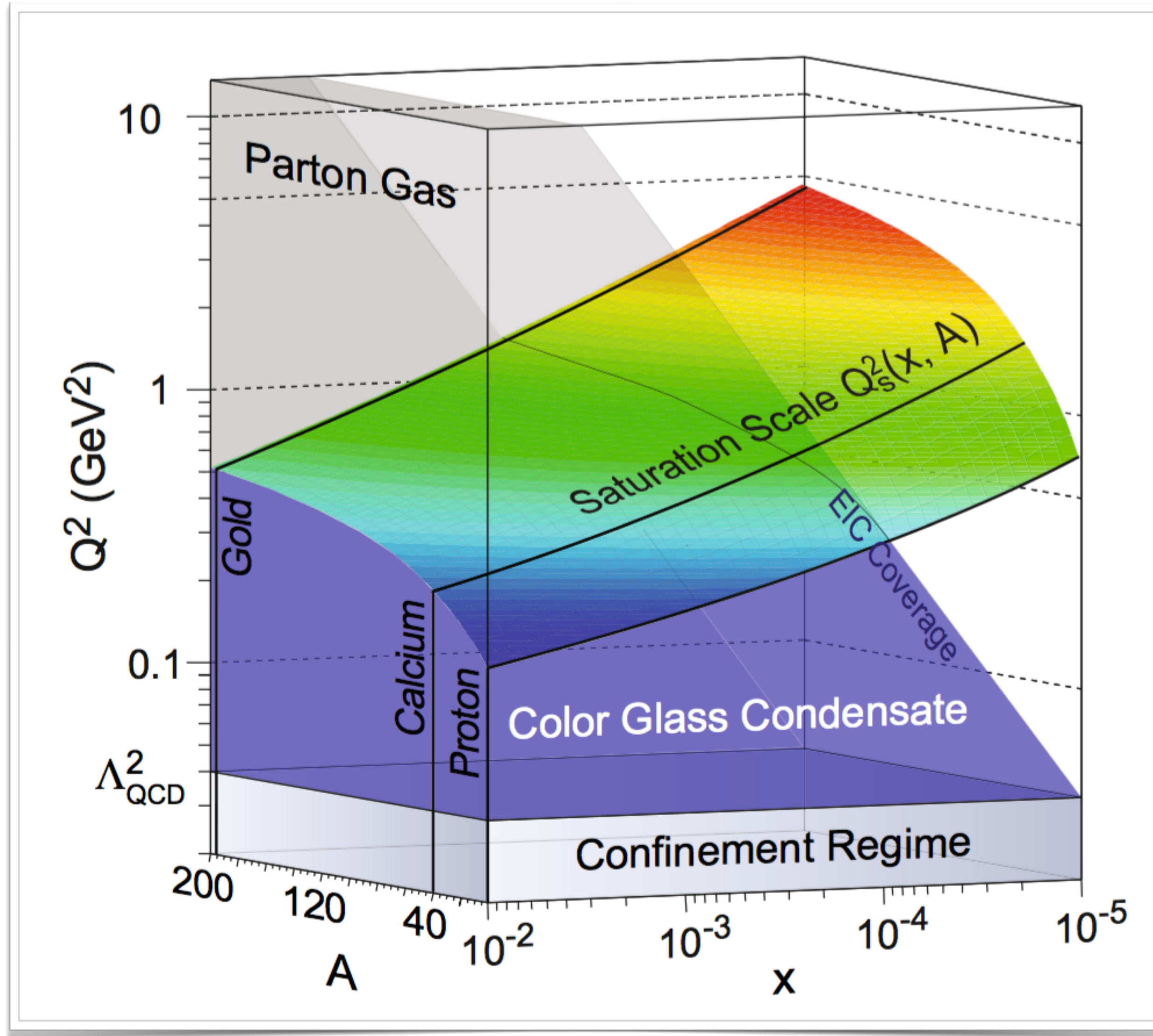
Accardi et al, EPJA 52 (2016) 268



The boundary between the dilute and saturated regimes is expected to appear at larger x in nuclei than in nucleons.

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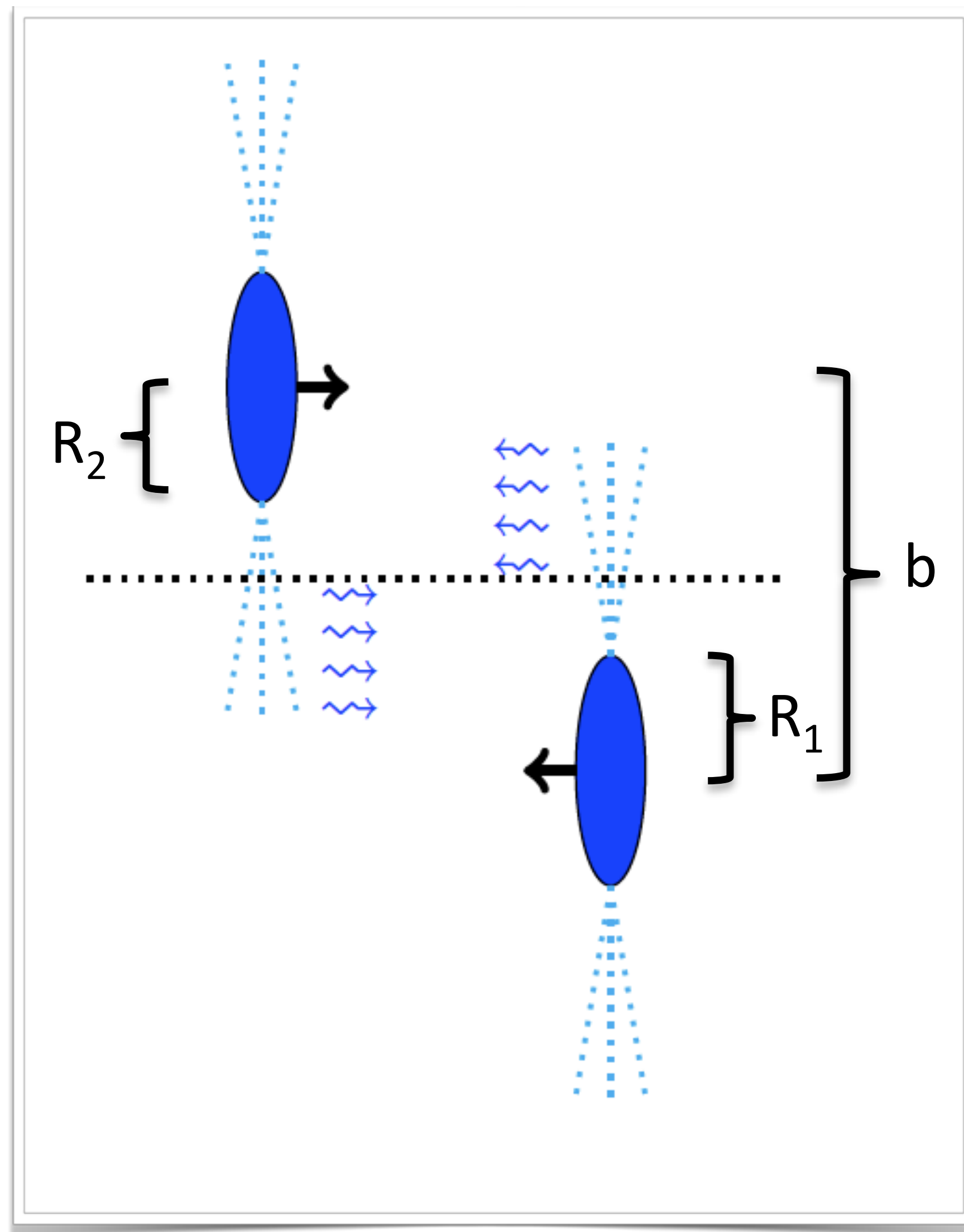
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How can we study the evolution of the pQCD structure of nuclear targets?
Today: RHIC and LHC.
Tomorrow: EIC.

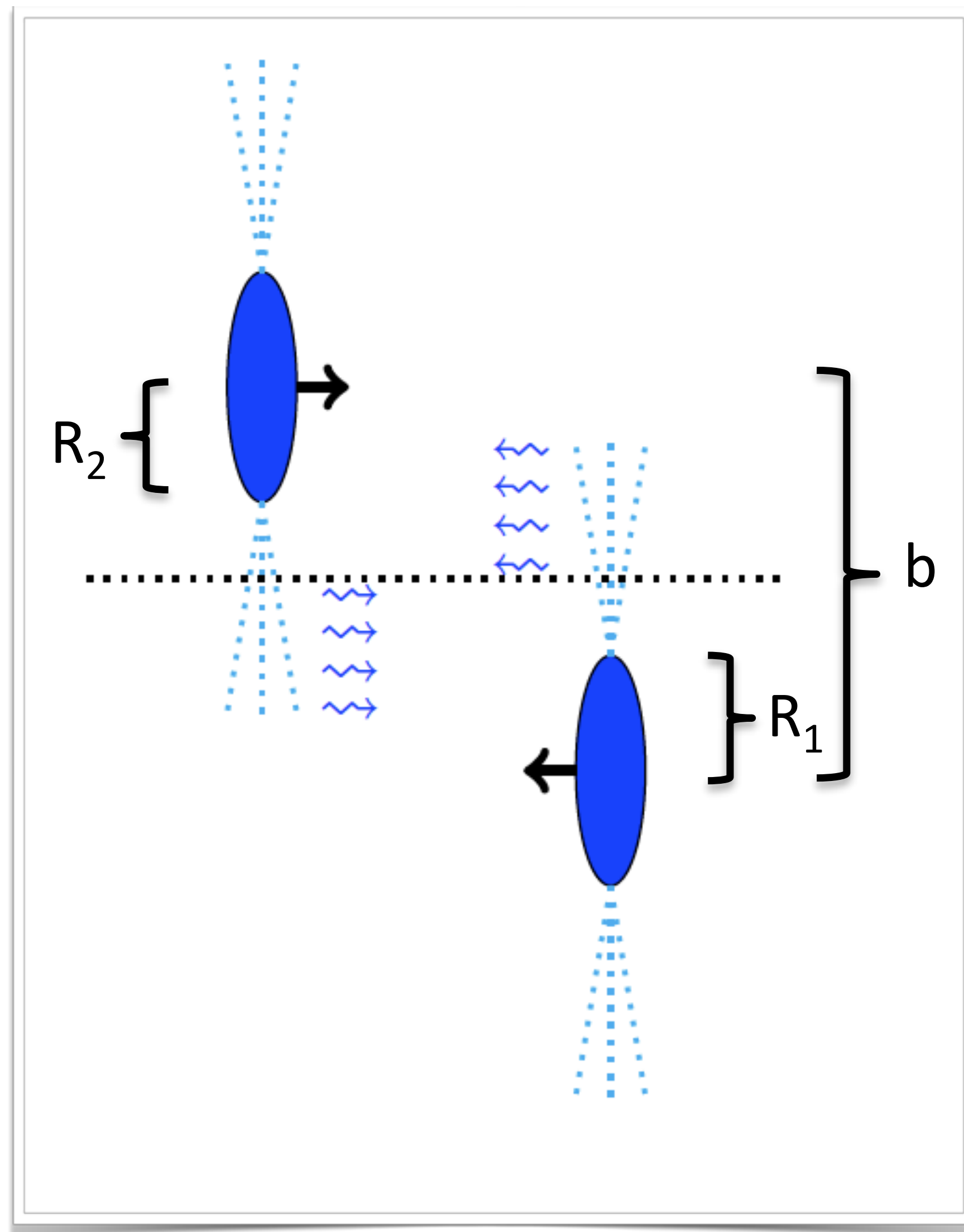
RHIC and LHC: our current photon-hadron colliders



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There are two potential sources, correspondingly two potential targets.

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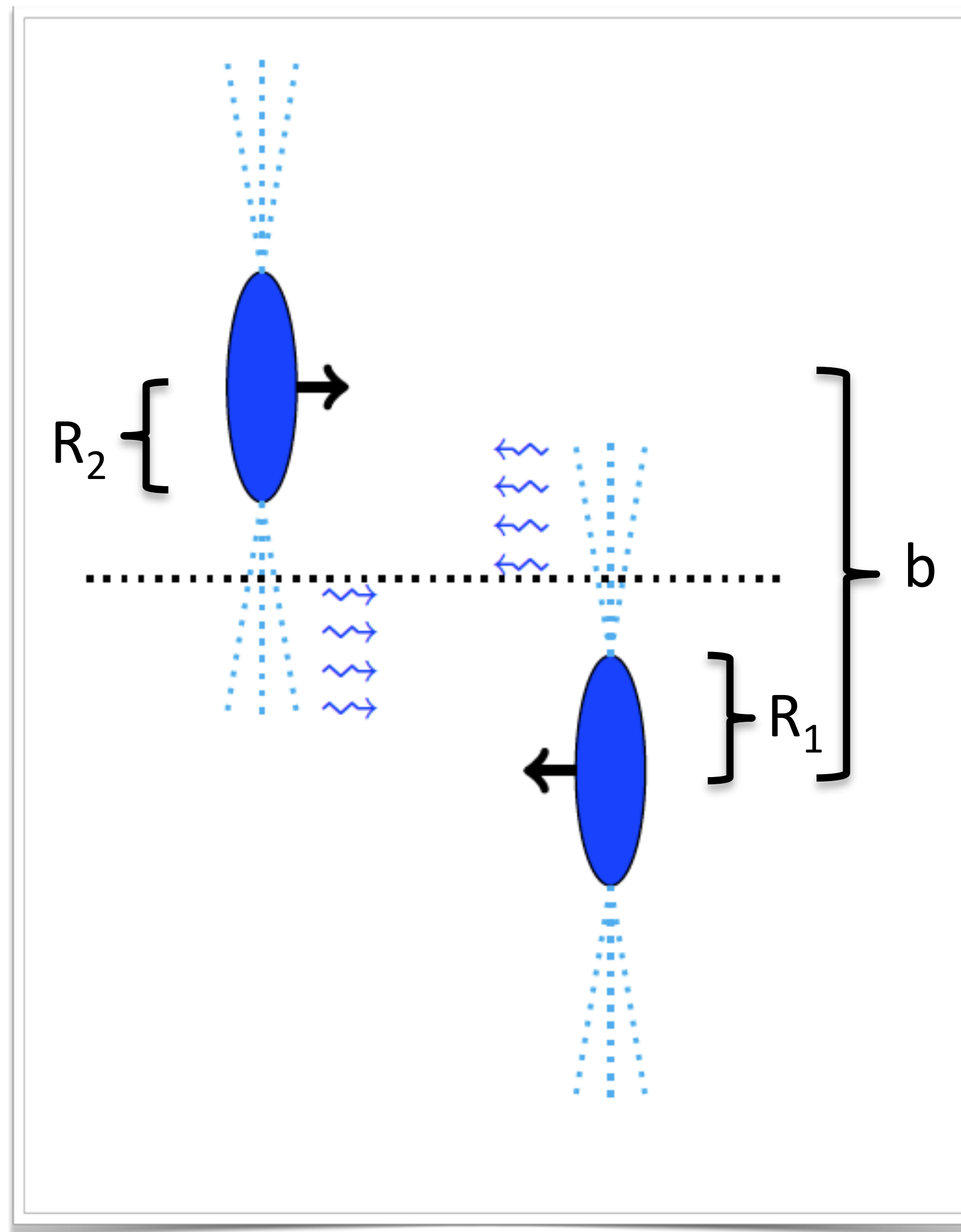
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The photons are coherently emitted by the source and its virtuality is restricted by the radius of the emitting particle:
Virtuality of photons from Pb: $Q^2 \approx (30 \text{ MeV})^2$.

The intensity of the photon beam is proportional to Z^2 .
Heavy nuclei are intense sources of quasi-real photons!

The max energy of the photons in the lab system is determined by the boost of the emitting particle:
Large photon energies are possible at RHIC and at LHC!

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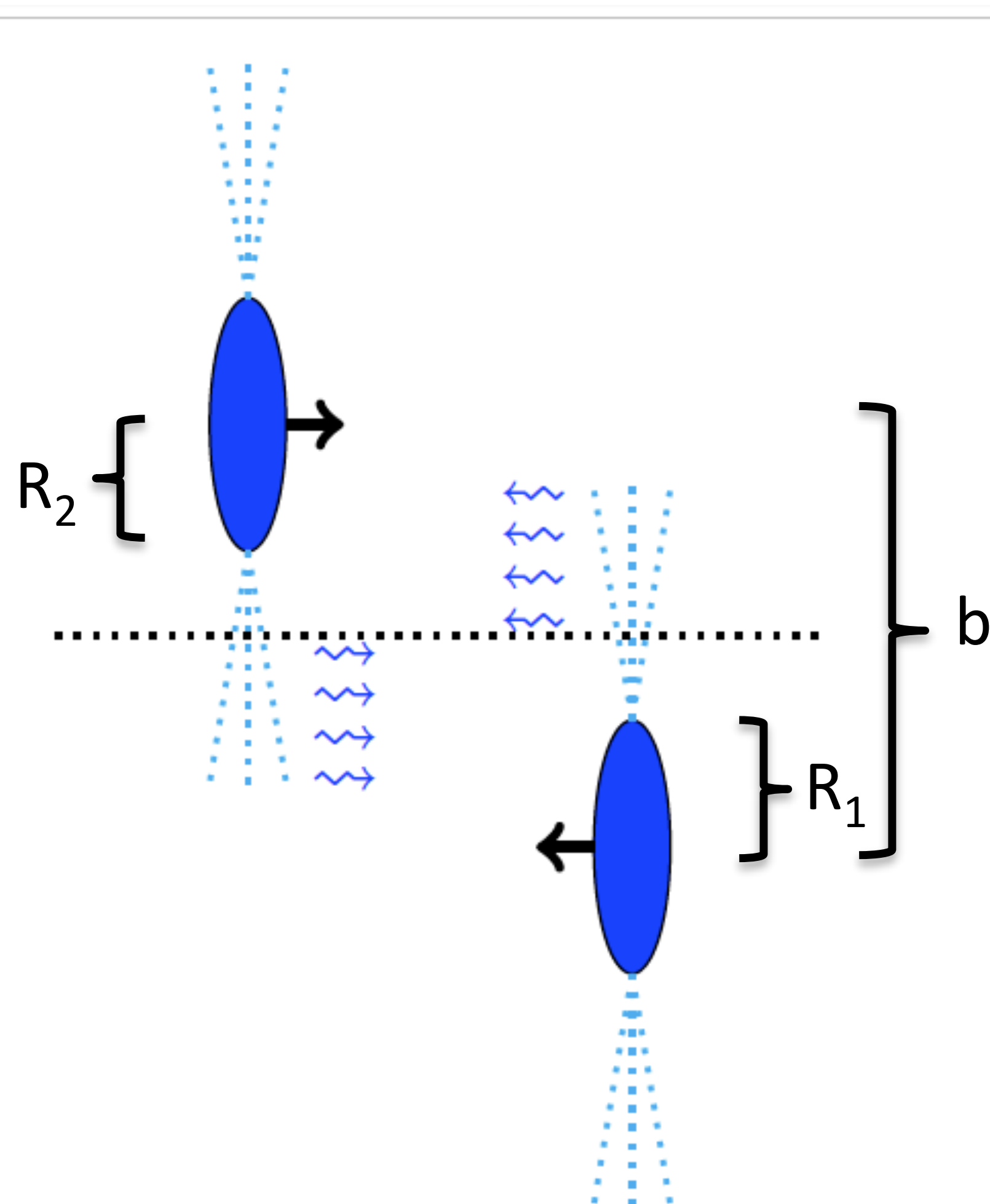
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RHIC and LHC: our current photon-hadron colliders



For $b > R_1 + R_2$ only photon induced processes remain.
These are called ultra-peripheral collisions (UPC).

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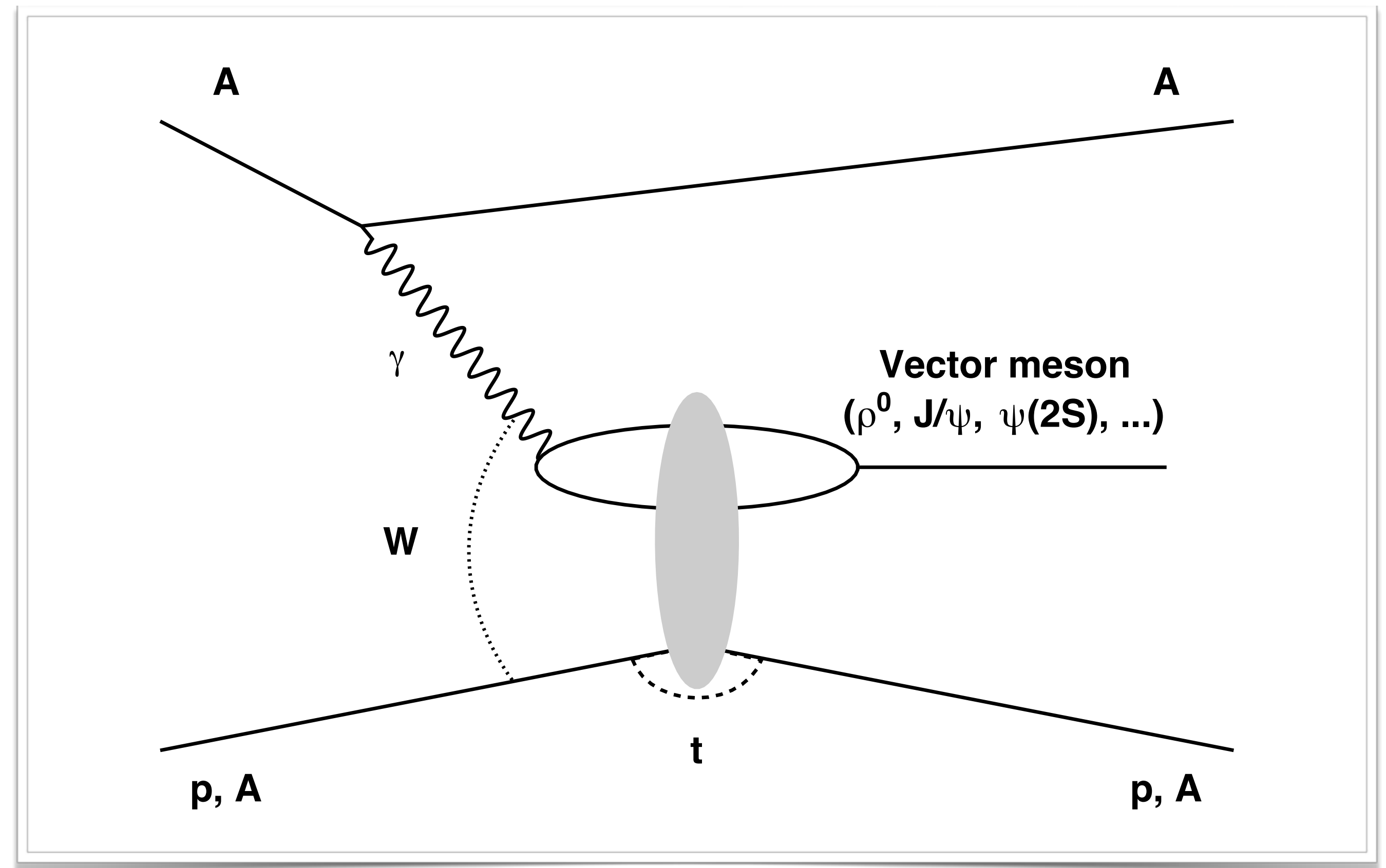
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Exclusive photoproduction of vector mesons

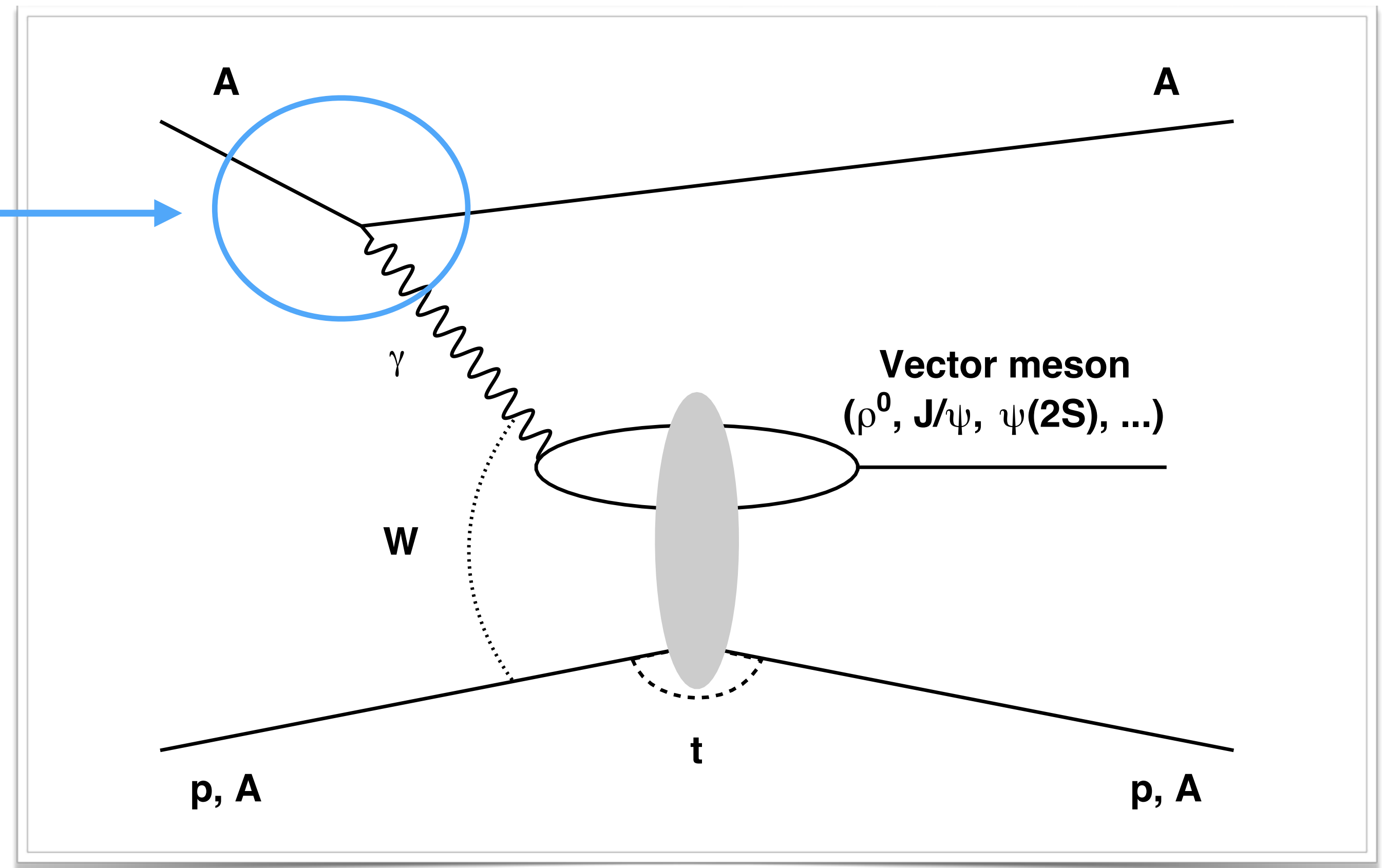


Exclusive photoproduction of vector mesons

The process can be factorised in two parts:

- Emission of the photon.

Only QED involved here

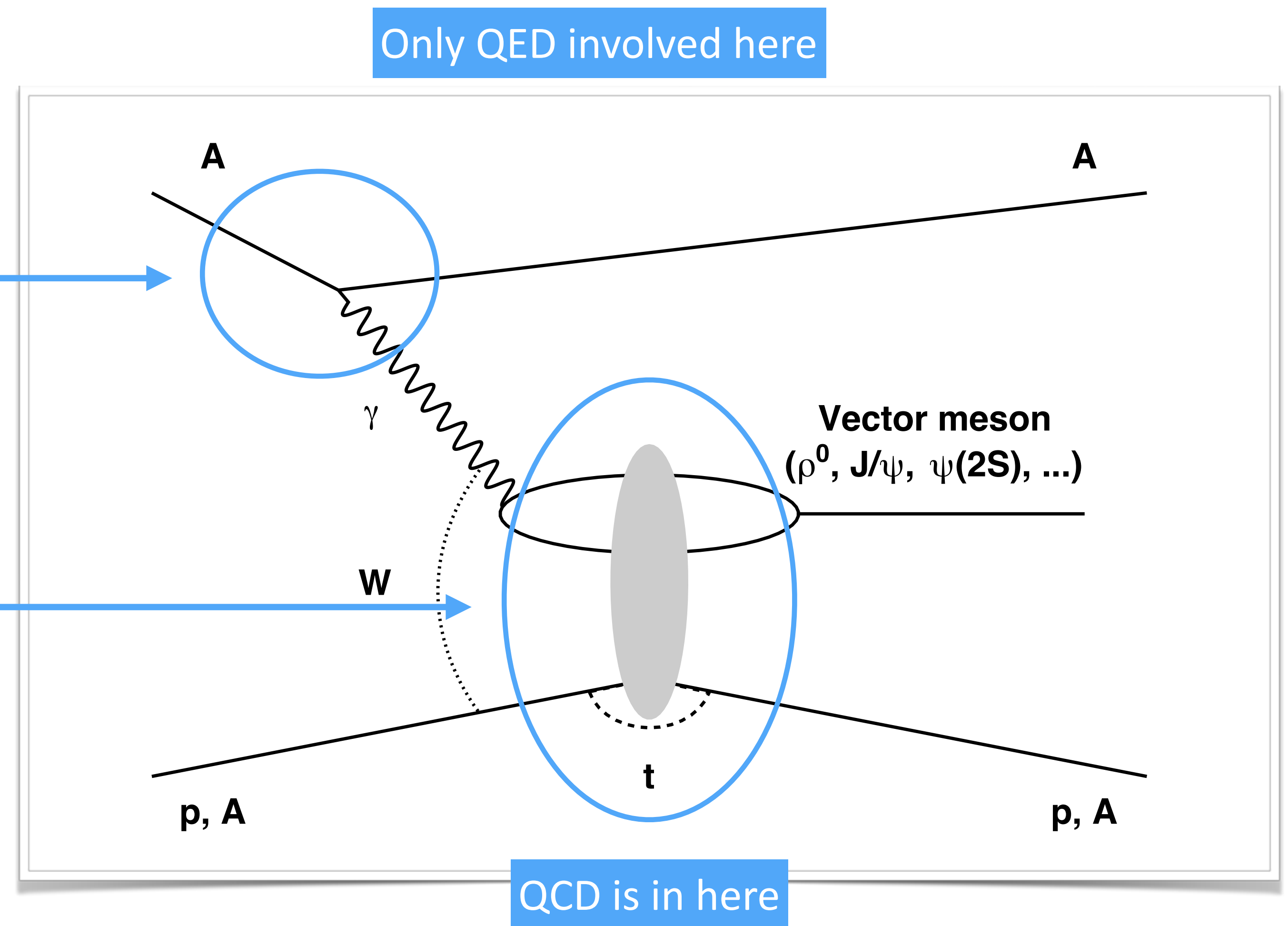


Exclusive photoproduction of vector mesons

The process can be factorised in two parts:

- Emission of the photon.

- Interaction of the photon with the target.



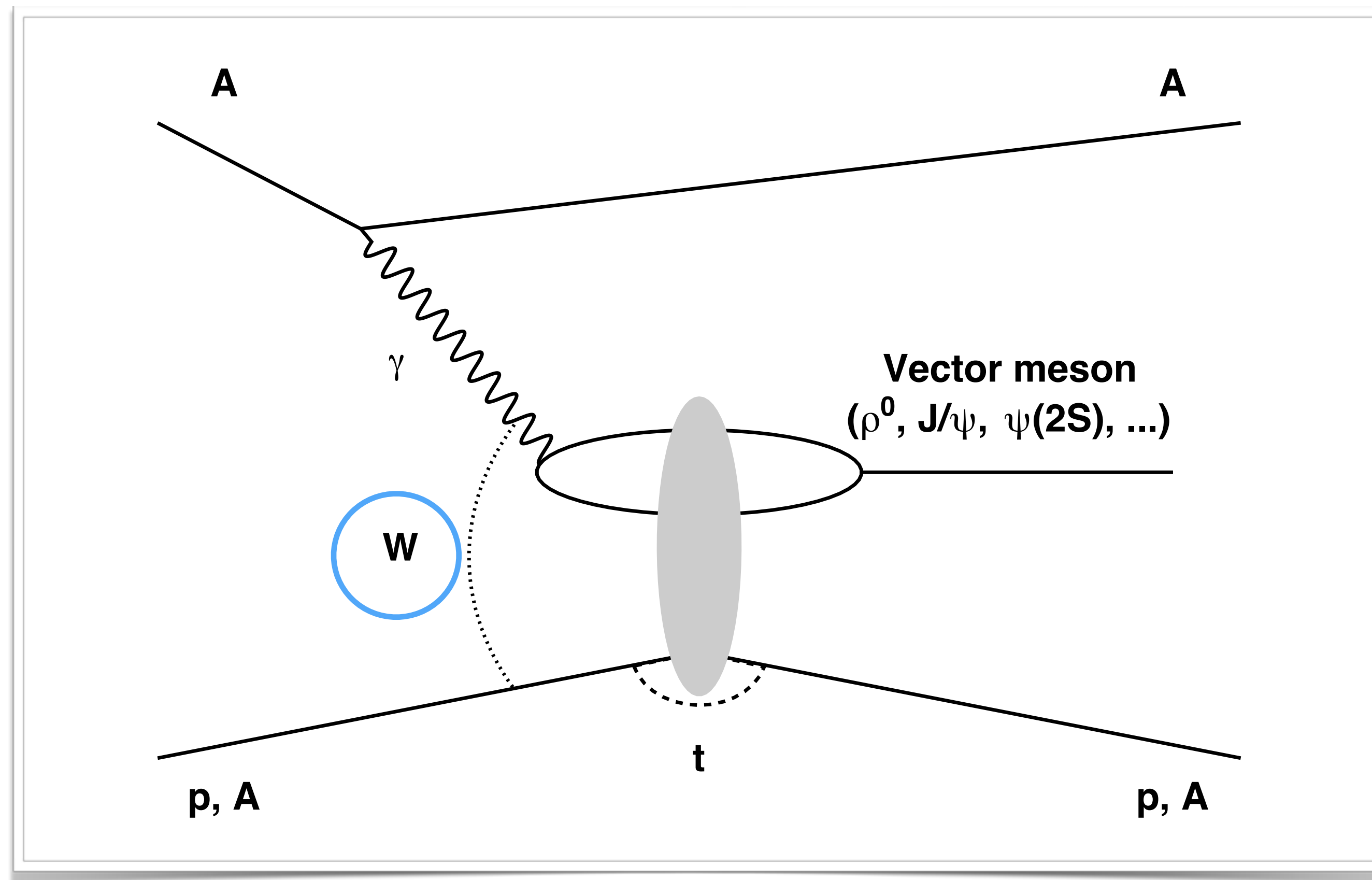
Kinematics: energy and rapidity

Centre-of-mass energy of the photon-target system:

$$W = \sqrt{2\omega 2E_{\text{beam}}}$$

where

- E_{beam} is the energy of the incoming beam and
- ω is the photon energy.



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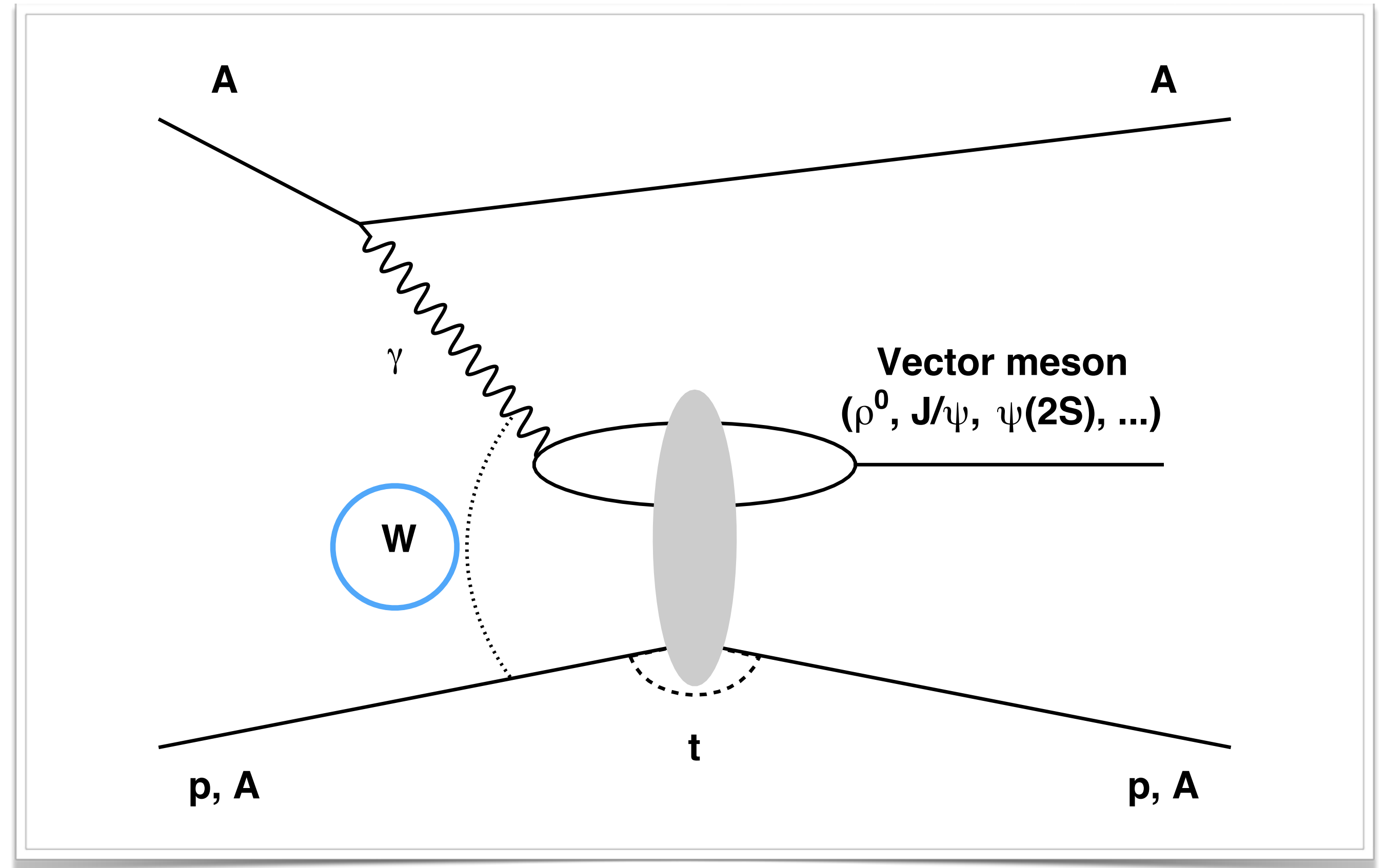
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$$2\omega = M \exp(-y)$$

where

- y is the rapidity of the vector meson measured in the laboratory frame **with respect to the direction of the target** and
- M is the mass of the vector meson.



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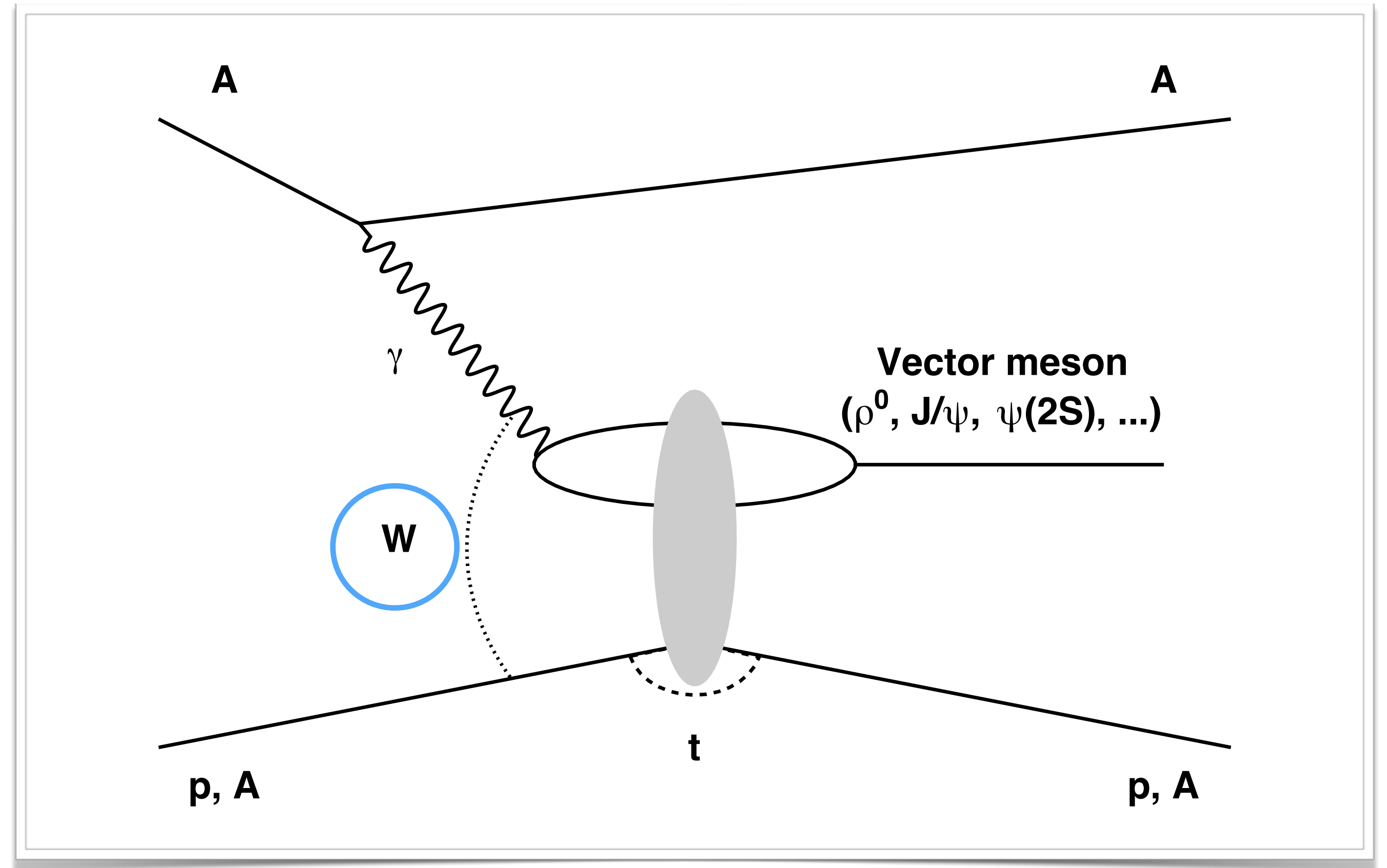
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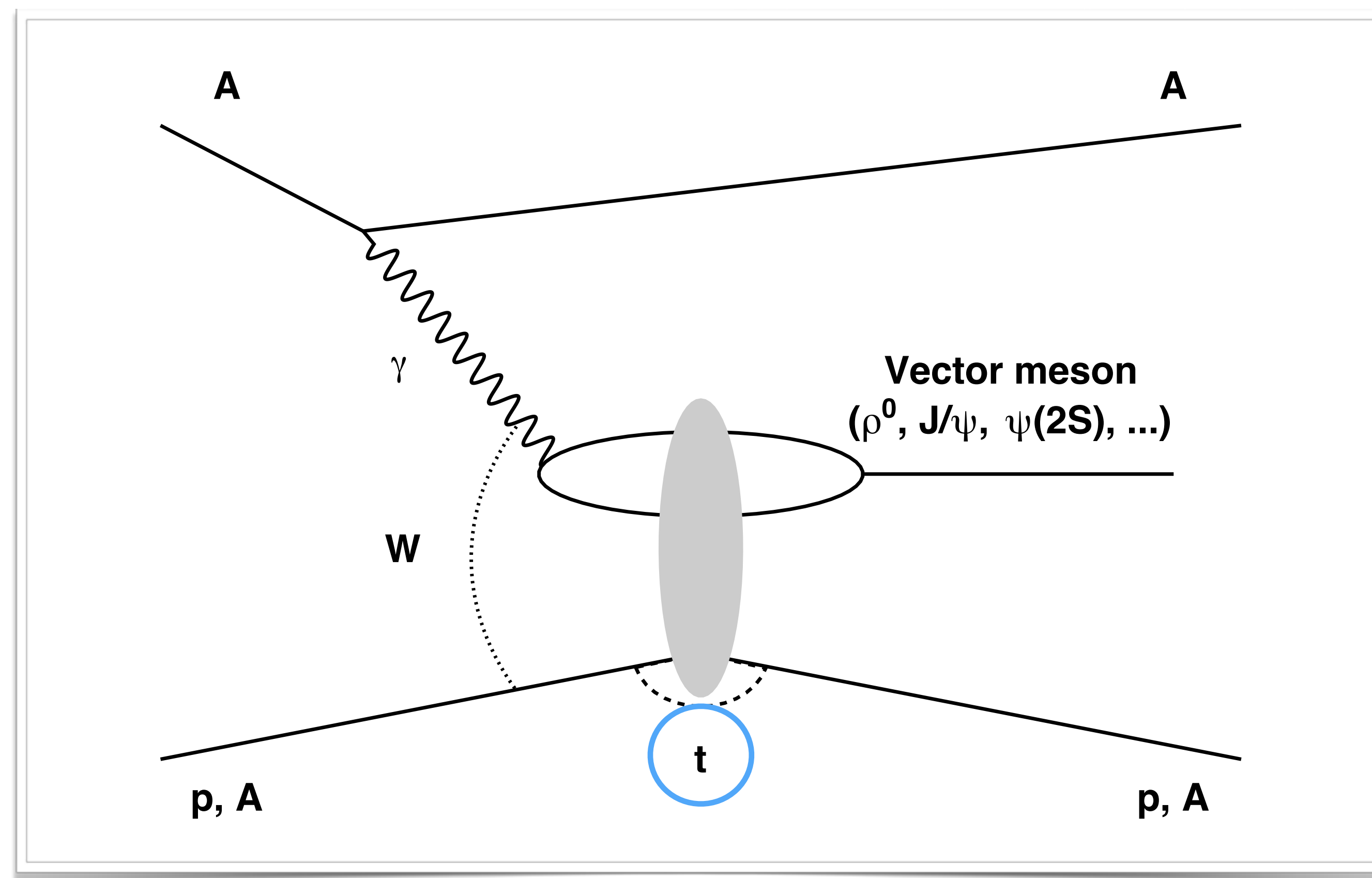


The rapidity of the vector meson determines the centre-of-mass energy of the photon-target interaction

Kinematics: transverse momentum

The square of the momentum transferred in the proton vertex $-t$ is related to the transverse momentum of the J/ψ :

$$-t \approx p_{\perp}^2$$

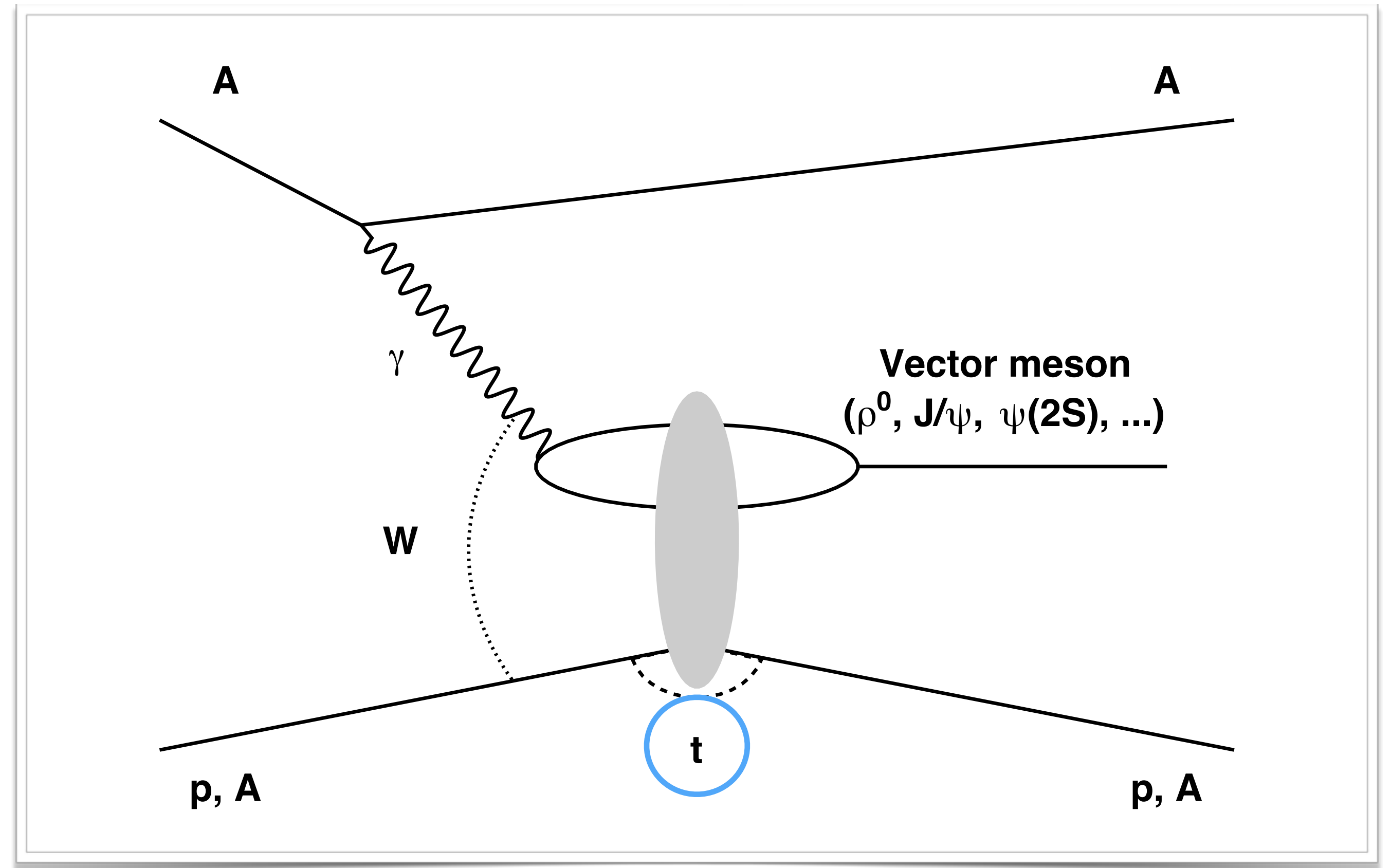


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The $-t$ distribution is determined by the **nuclear form factor**, which gives, through a Fourier transform, the distribution of gluons, in the transverse plane.

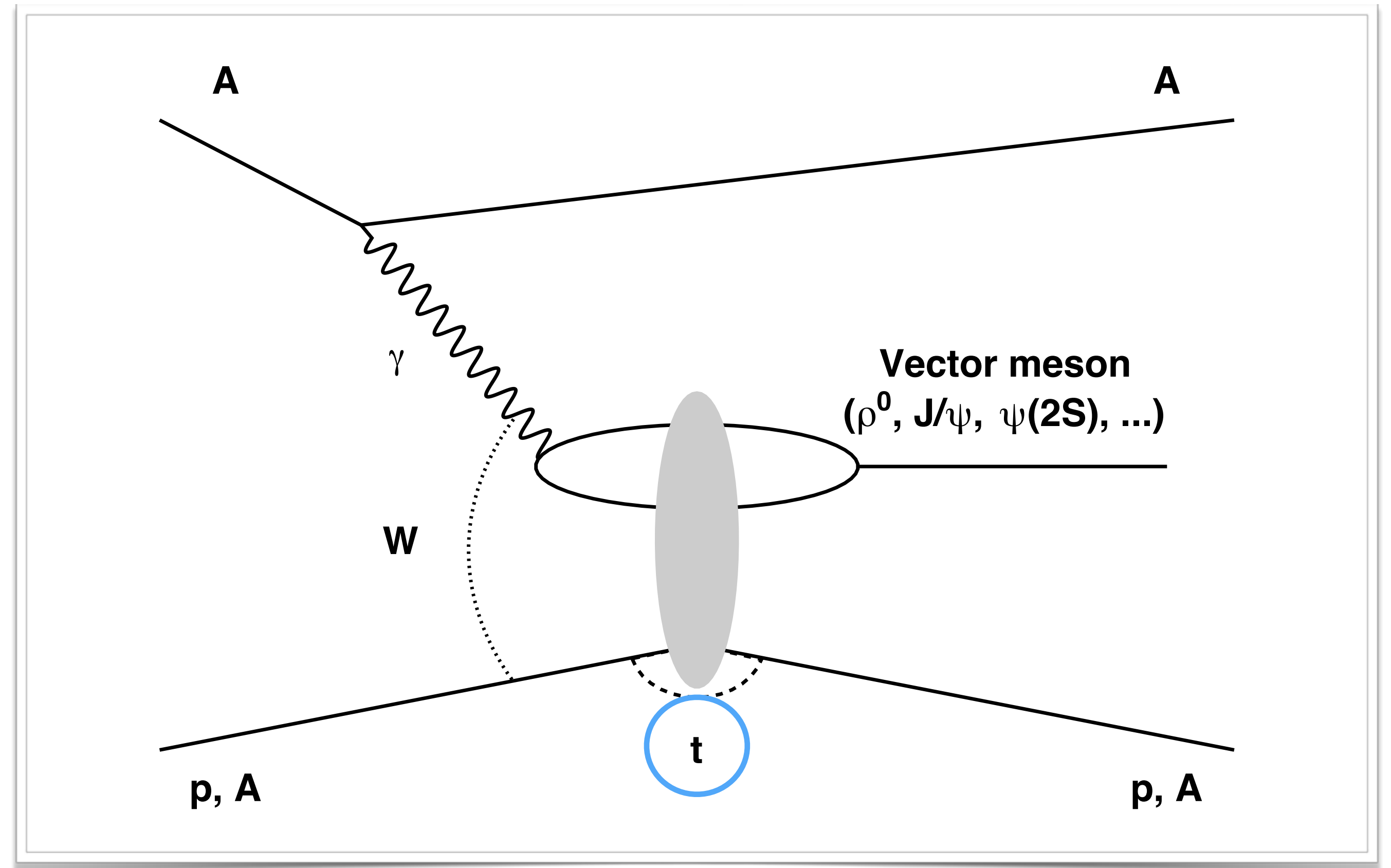


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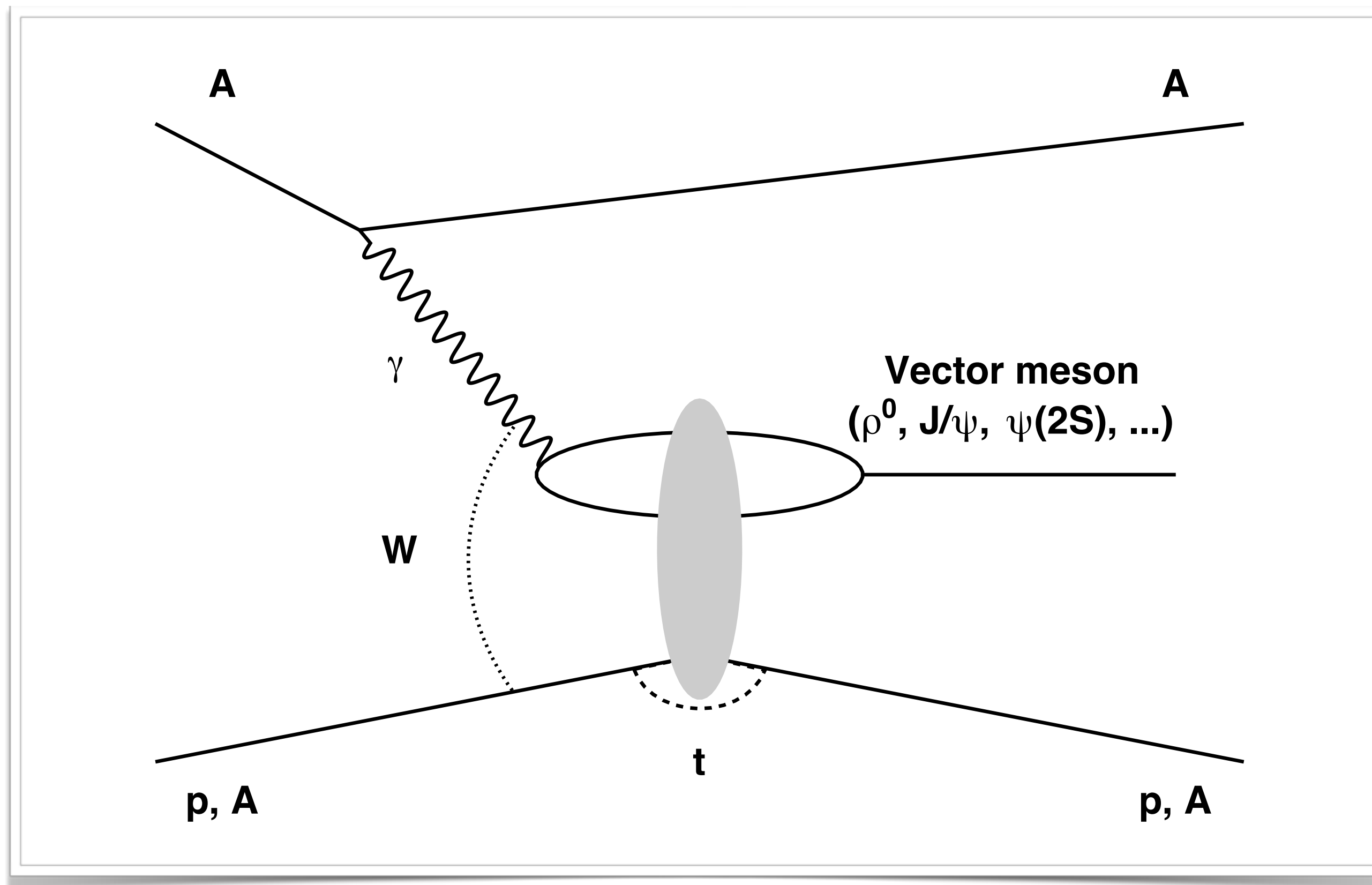
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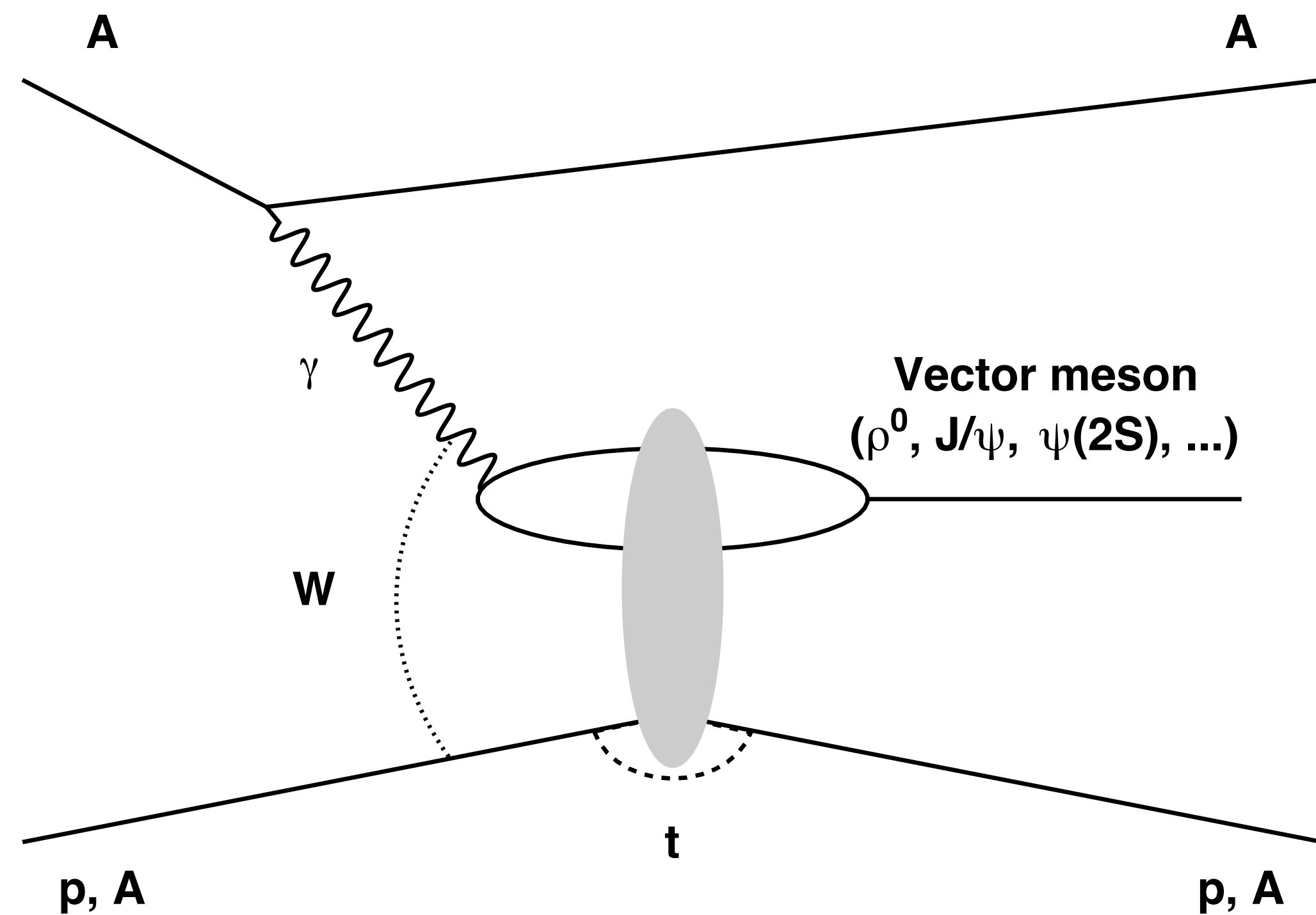


The $-t$ dependence of vector meson production determines the transverse distribution of matter in the target!

Exclusive photoproduction of vector mesons: experimental point of view



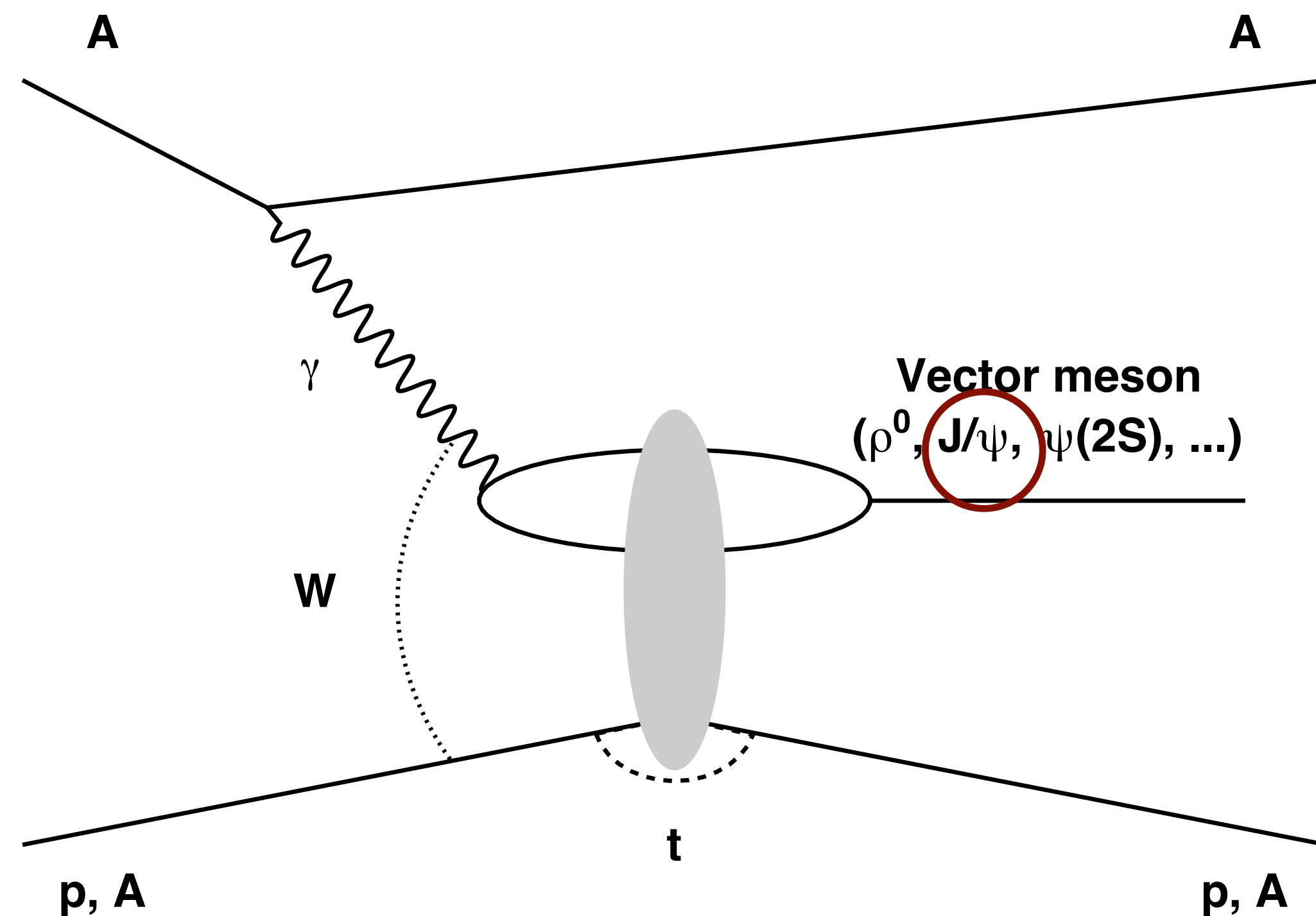
Exclusive photoproduction of vector mesons: experimental point of view



Very clean experimental signature:

- Vector mesons with low transverse momentum.
(Few tens of MeV/c for nuclear targets.)
- Nothing else in the detector.
(For AA collisions at RHIC/LHC there may be forward neutrons from independent photon exchanges.)

Exclusive photoproduction of vector mesons: experimental point of view



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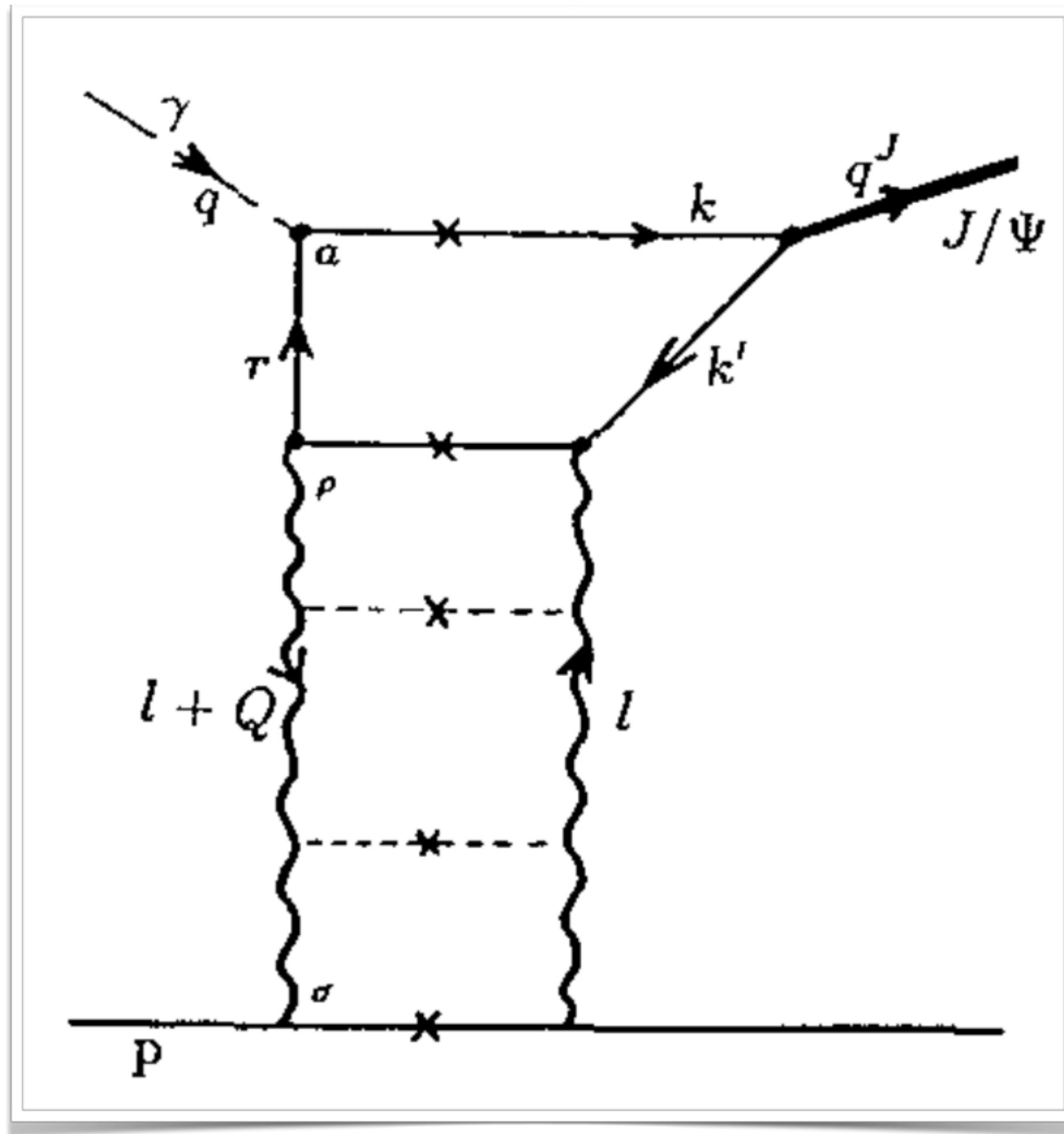
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J/ψ

- Small width and leptonic decays with a large BR.
- Large mass of the charm quark allows for pQCD.

Exclusive photoproduction of vector mesons: pQCD point of view

Ryskin: Z. Phys. C 57, 89-92 (1993)



$$\frac{d\sigma^T(\gamma p \rightarrow J/\Psi + p)}{dt} = \frac{|M|^2}{16\pi s^2}$$

$$= [F_N^{2G}(t)]^2 \frac{\alpha_s^2 \Gamma_{ee}^J m_J^3}{3\alpha_{em}} \pi^3$$

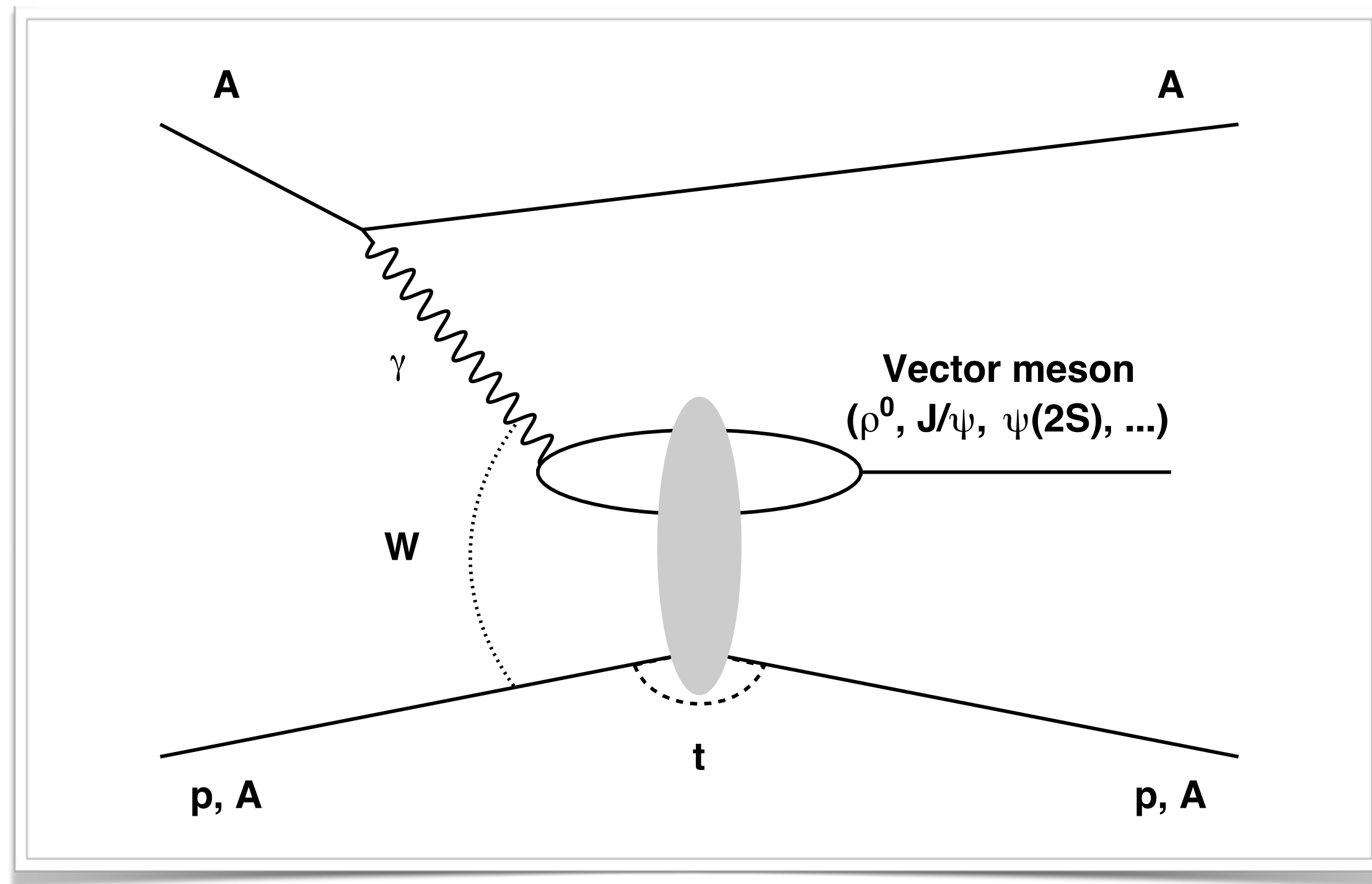
$$\times \left[\bar{x} G(\bar{x}, \bar{q}^2) \frac{2\bar{q}^2 - |q_t^J|^2}{(2\bar{q}^2)^3} \right]^2,$$

Cross section depends on the **square of the gluon distribution**
(Amplitude computed in the leading-log approximation of pQCD.)

The amplitude in the dipole picture

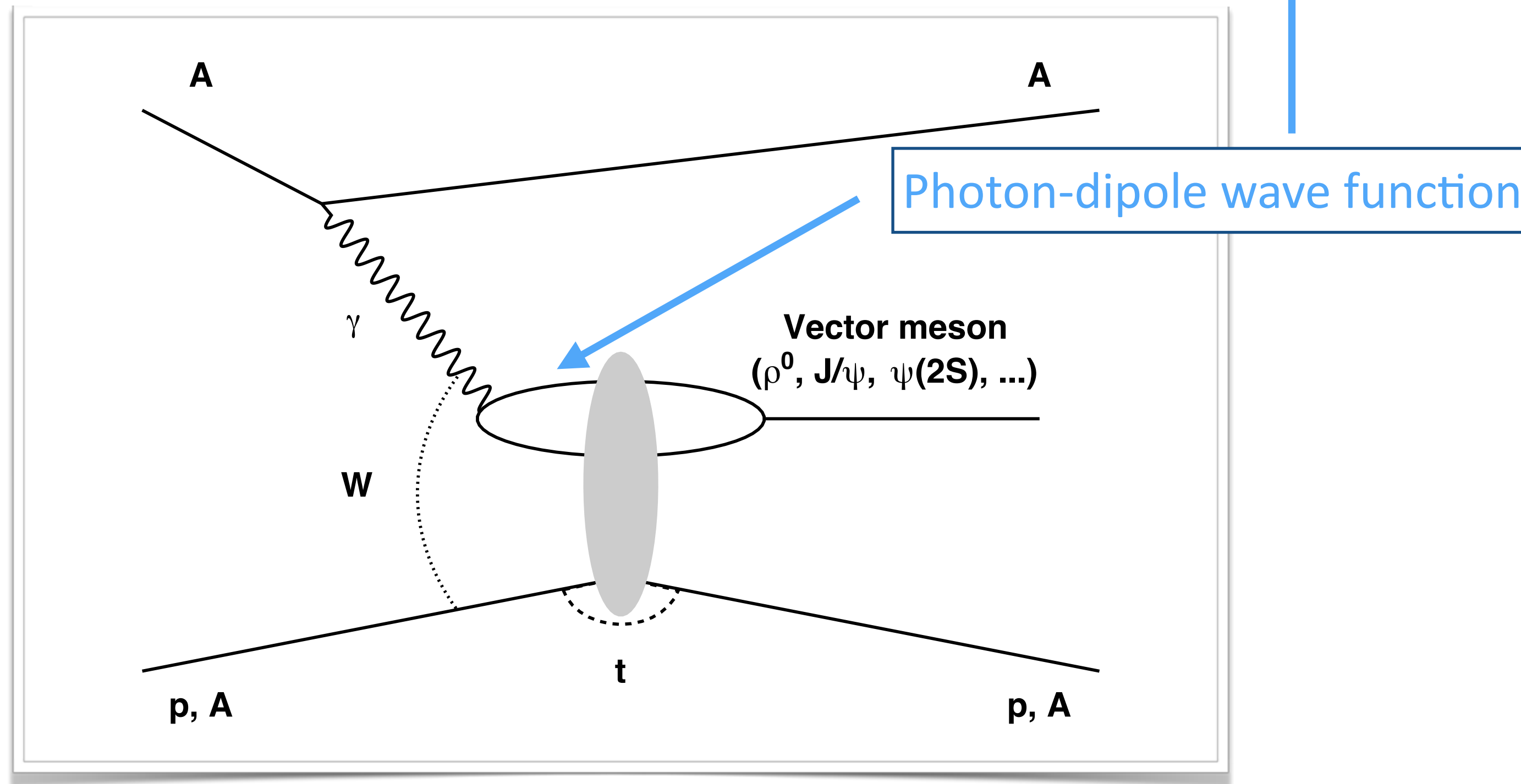
$$A(x, Q^2, \vec{\Delta})_{T,L} = i \int d\vec{r} \int_0^1 \frac{dz}{4\pi} (\Psi^* \Psi_V)_{T,L} \int d\vec{b} e^{-i(\vec{b} - (1-z)\vec{r}) \cdot \vec{\Delta}} \frac{d\sigma_{\text{dip}}}{d\vec{b}}$$

$\Delta^2 = -t$ Dipole size Quark energy fraction Impact parameter



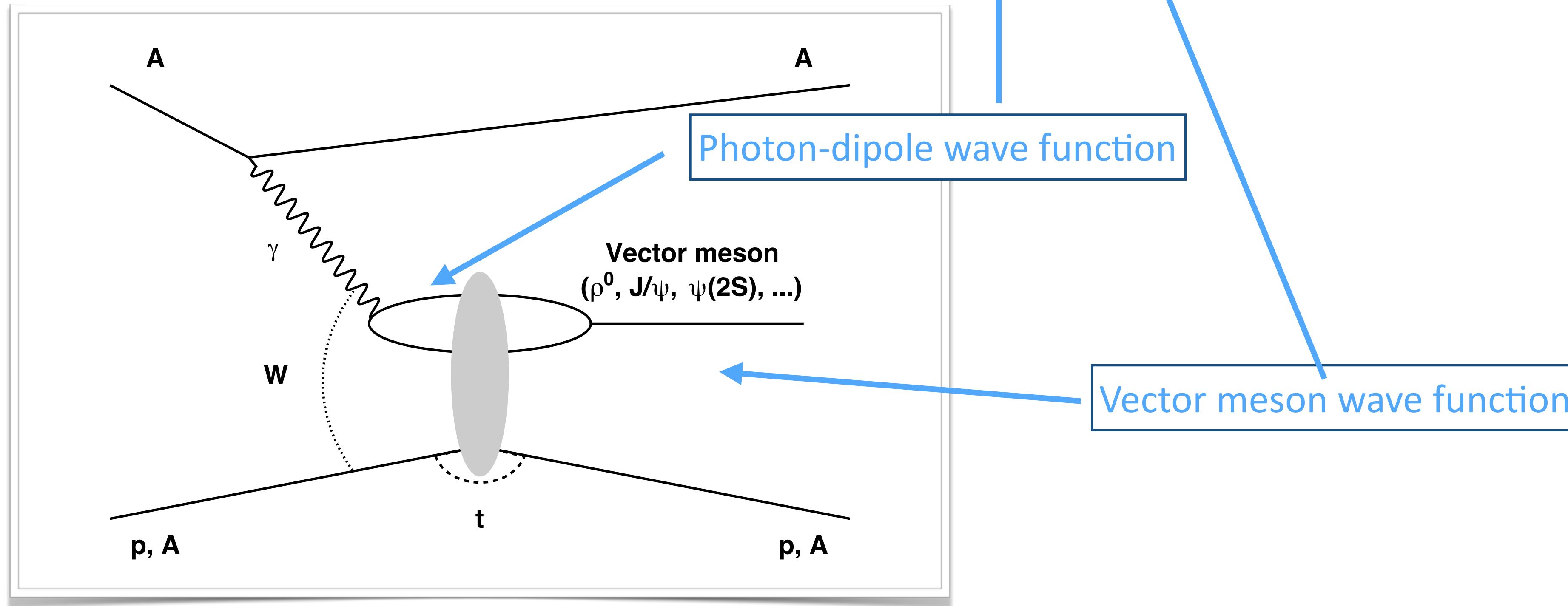
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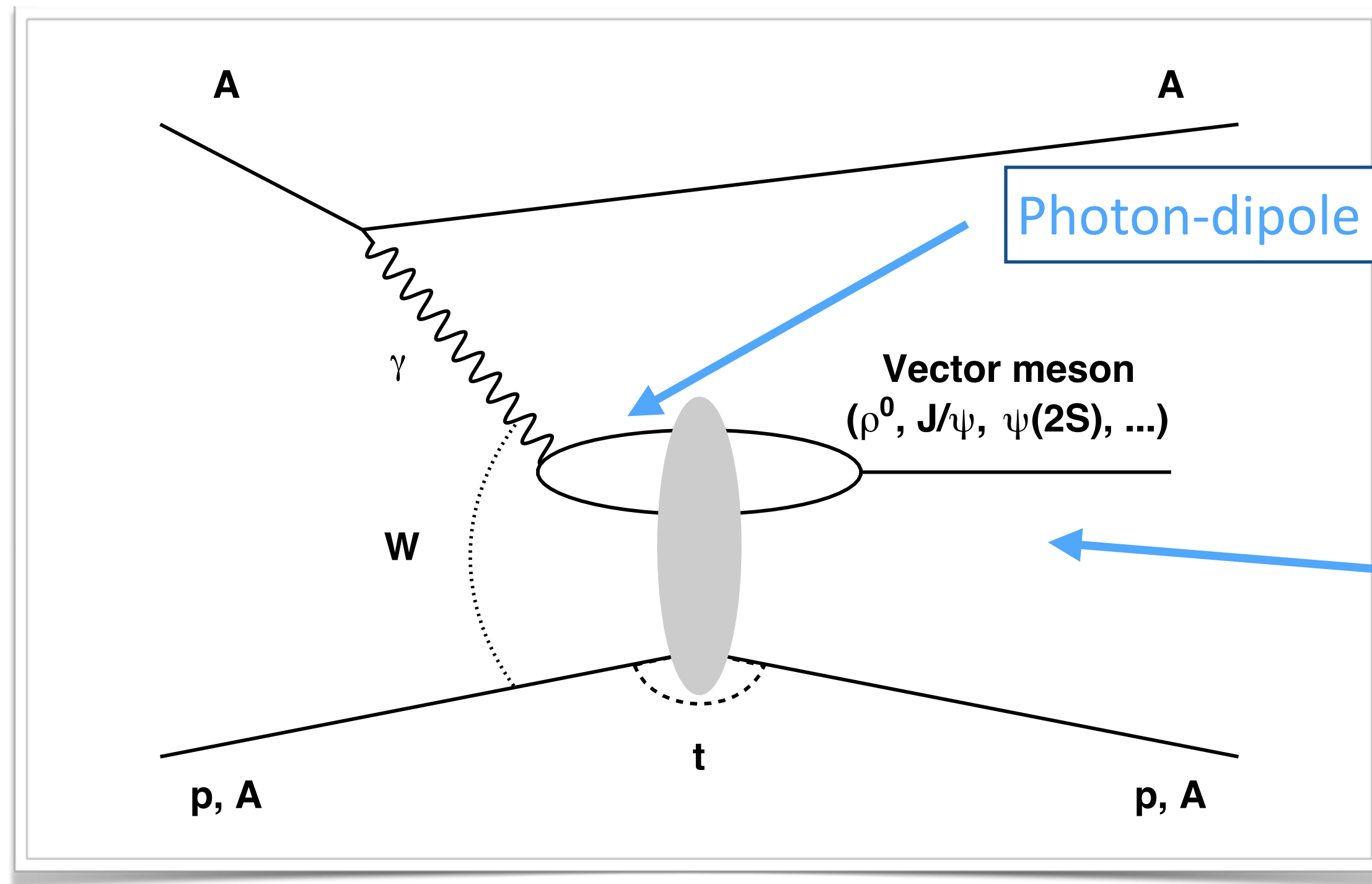


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Annotations for the equation:

- $\Delta^2 = -t$ points to $\vec{\Delta}$
- Dipole size points to $d\vec{r}$
- Quark energy fraction points to dz
- Impact parameter points to $d\vec{b}$
- $\frac{d\sigma_{\text{dip}}}{d\vec{b}}$ is highlighted in a box, with an arrow pointing to the final text block.

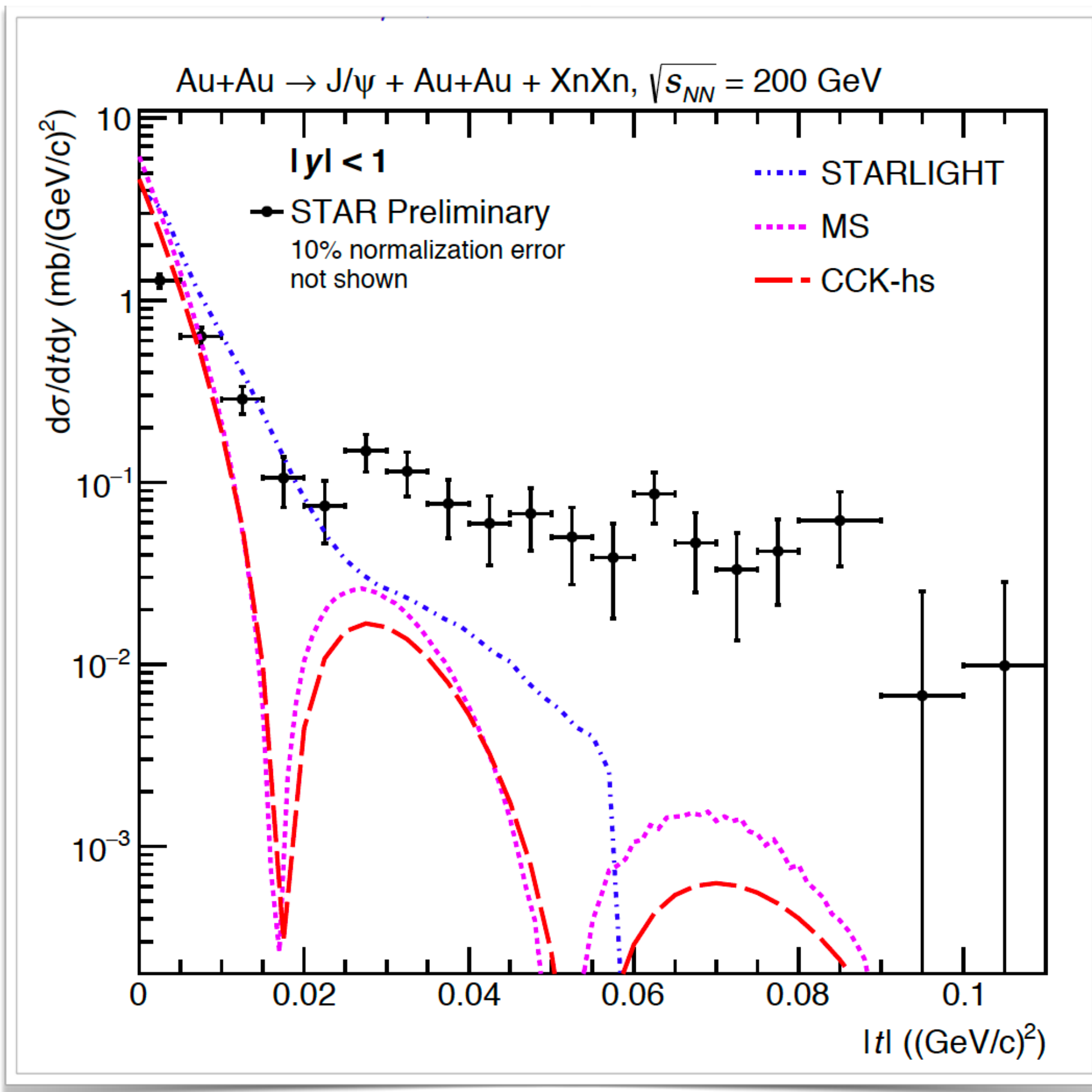


Photon-dipole wave function

Vector meson wave function

Dipole-Target cross section
pQCD physics gets here!

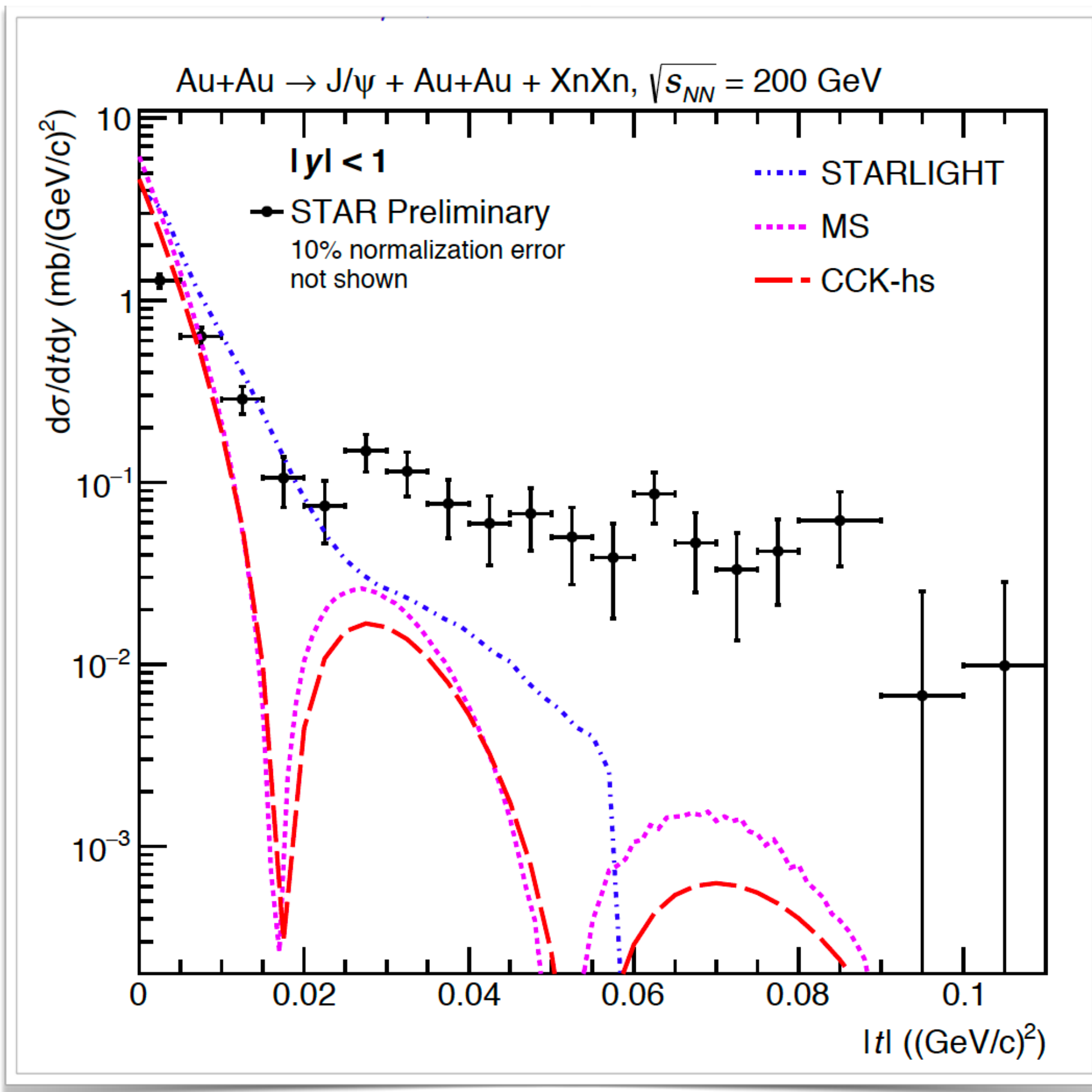
News from RHIC: coherent J/ψ photoproduction in UPC



Data from Au-Au UPC @ 200 GeV.

Measurement at mid-rapidity:
 $W_{\gamma Au} \approx 25$ GeV.

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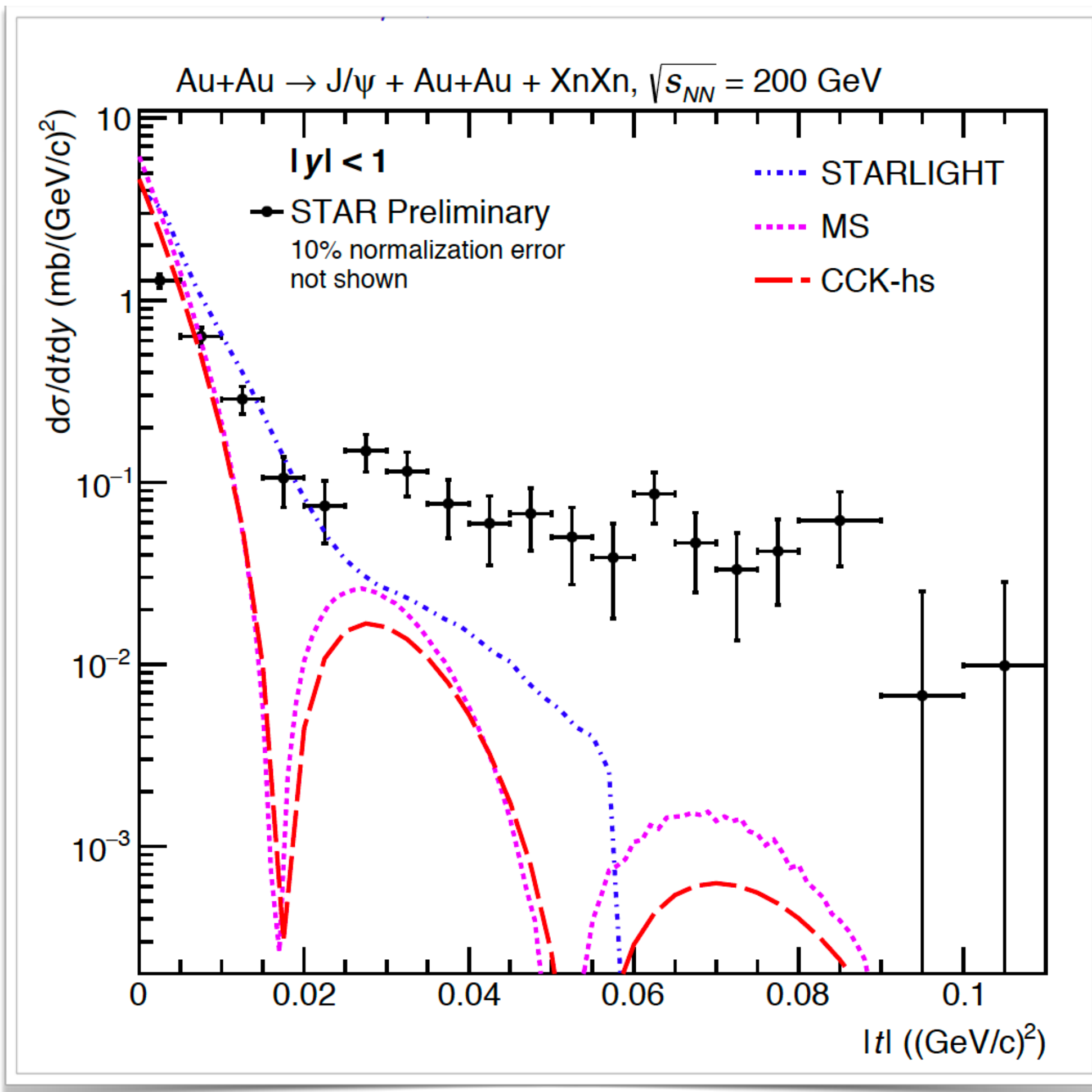
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Slope milder than dipole-based models
A hint of diffraction minimum.

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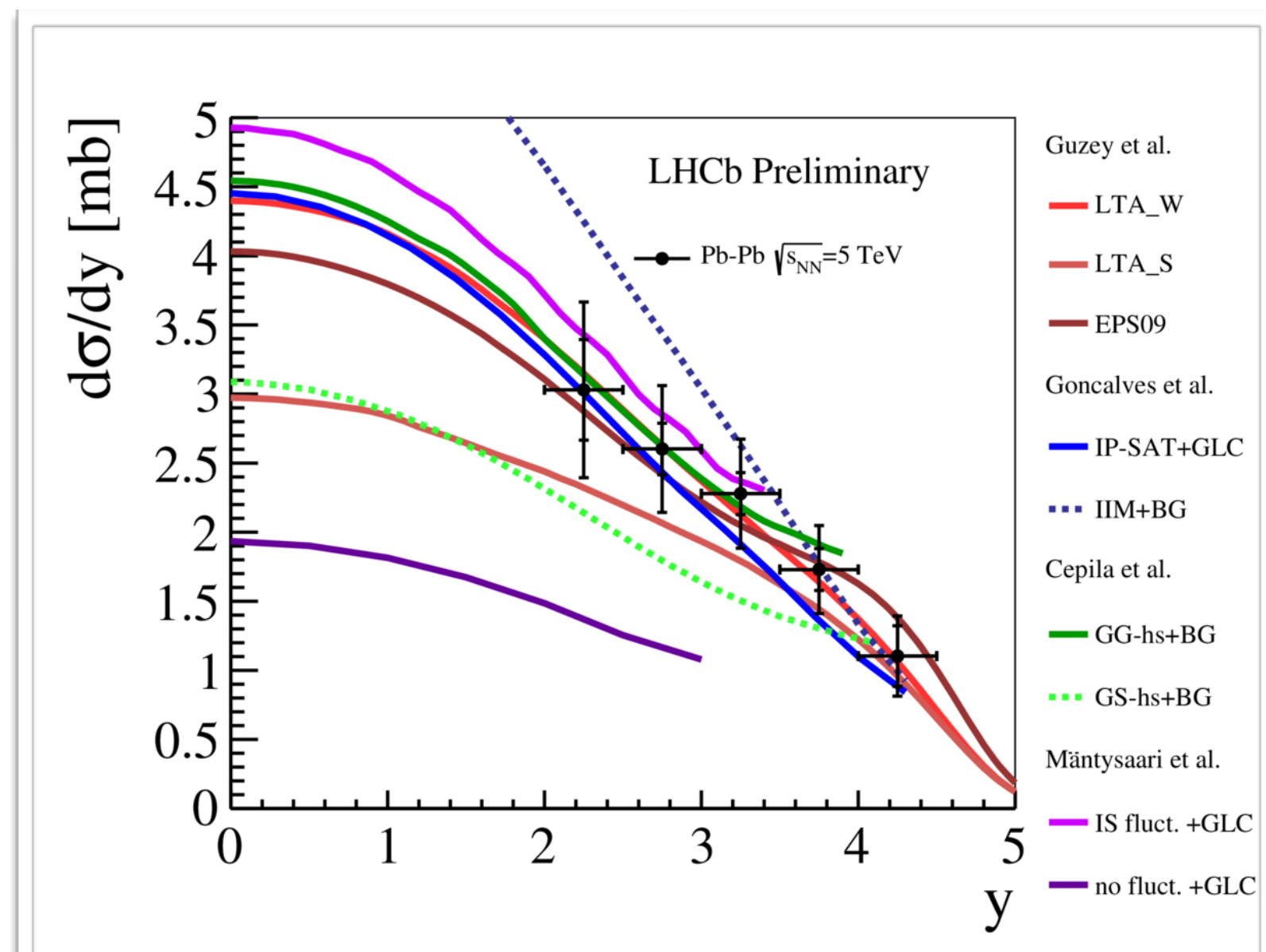
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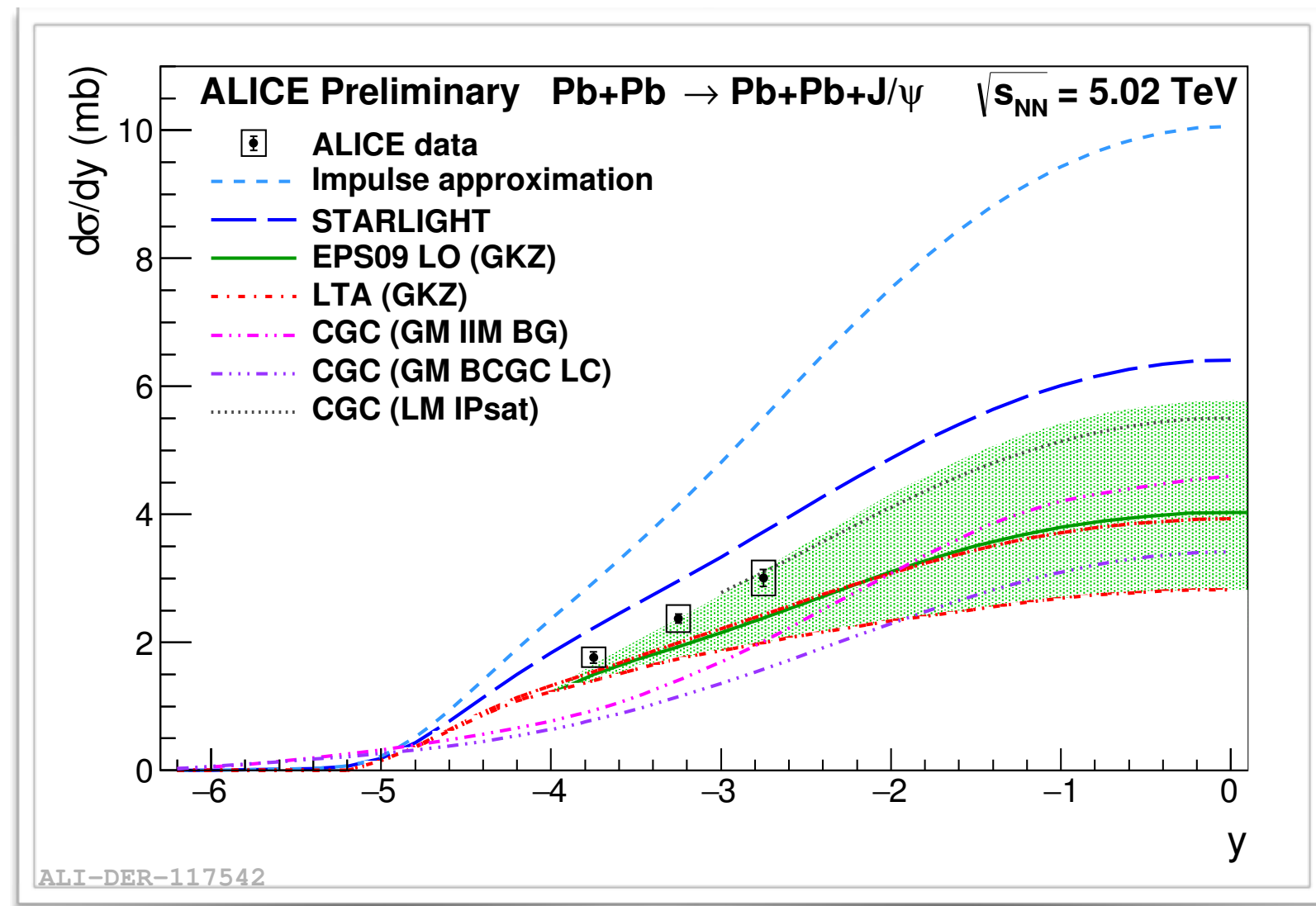
Wednesday in the WG2 session

Coherent J/ψ photoproduction in ultra-peripheral collisions at STAR, Jaroslav Adam.

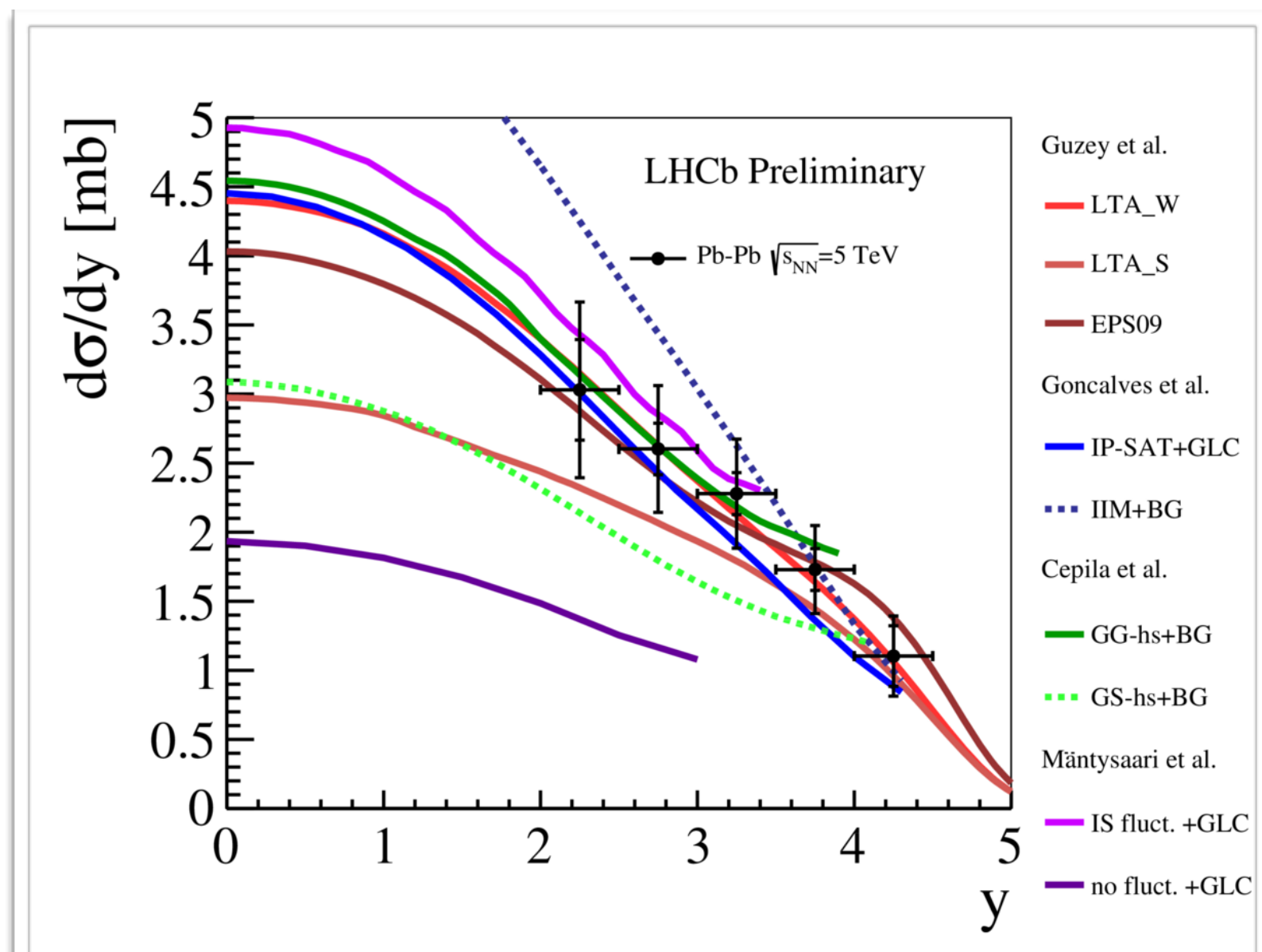
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Measurements by LHCb and ALICE.

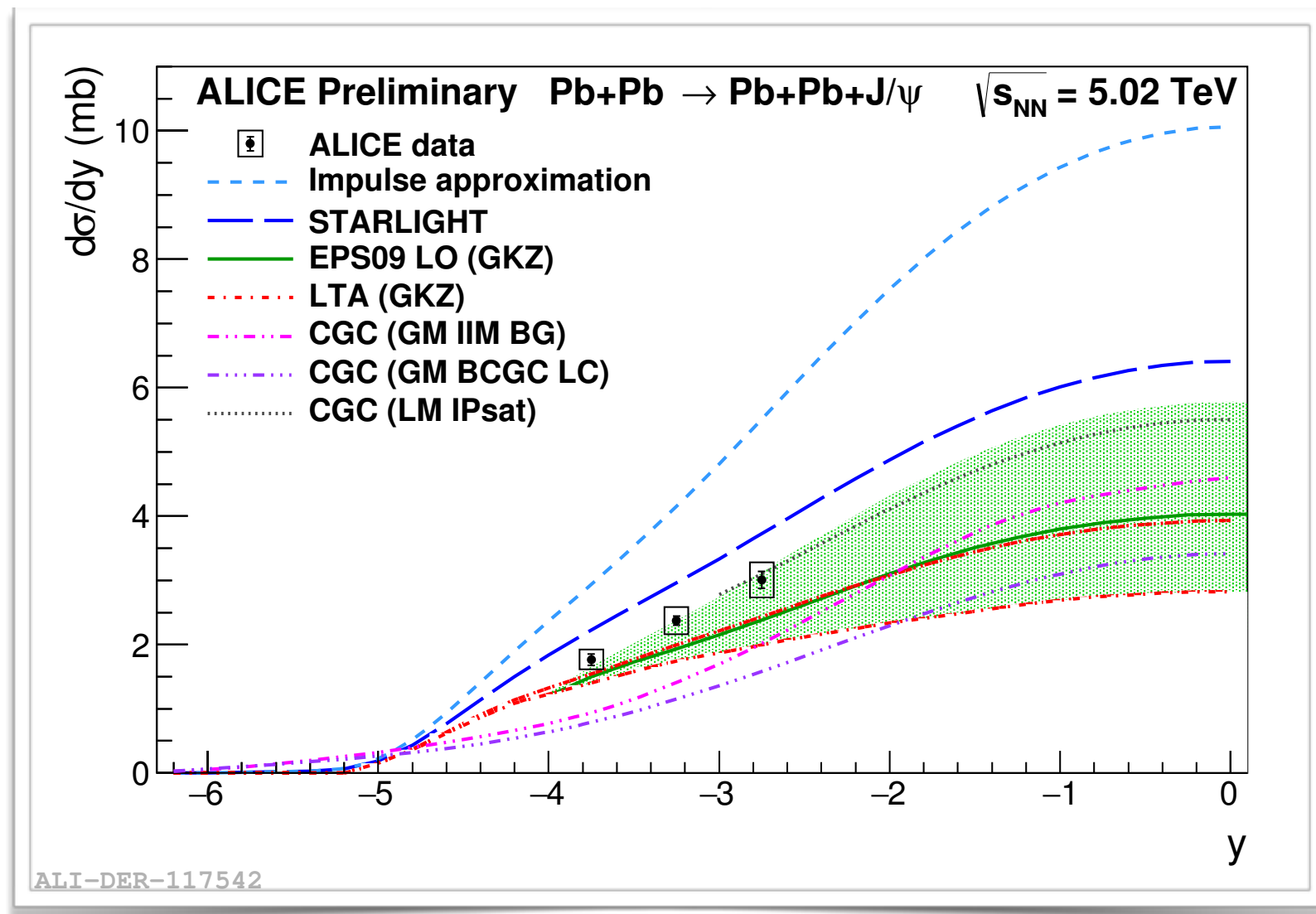


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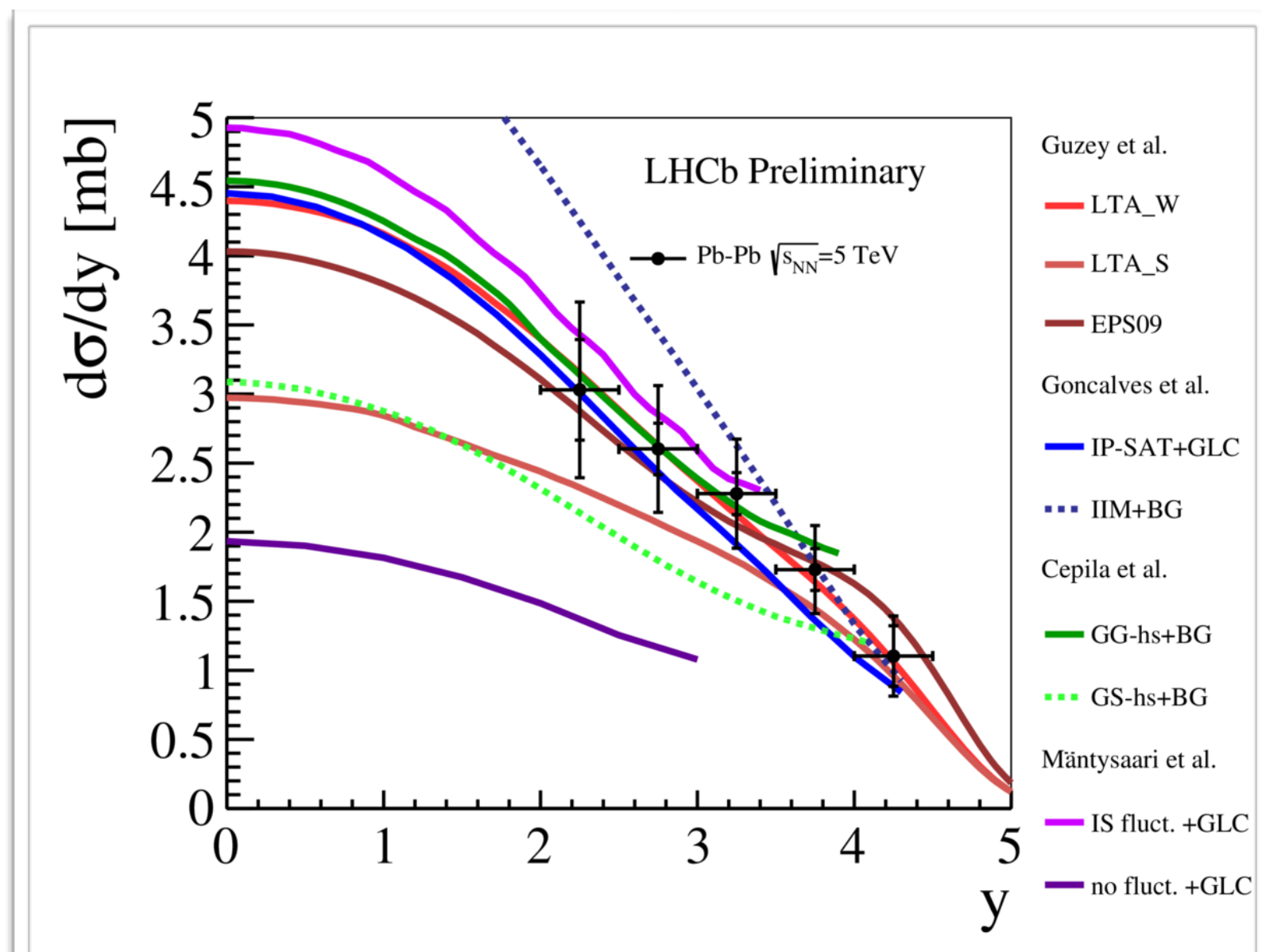


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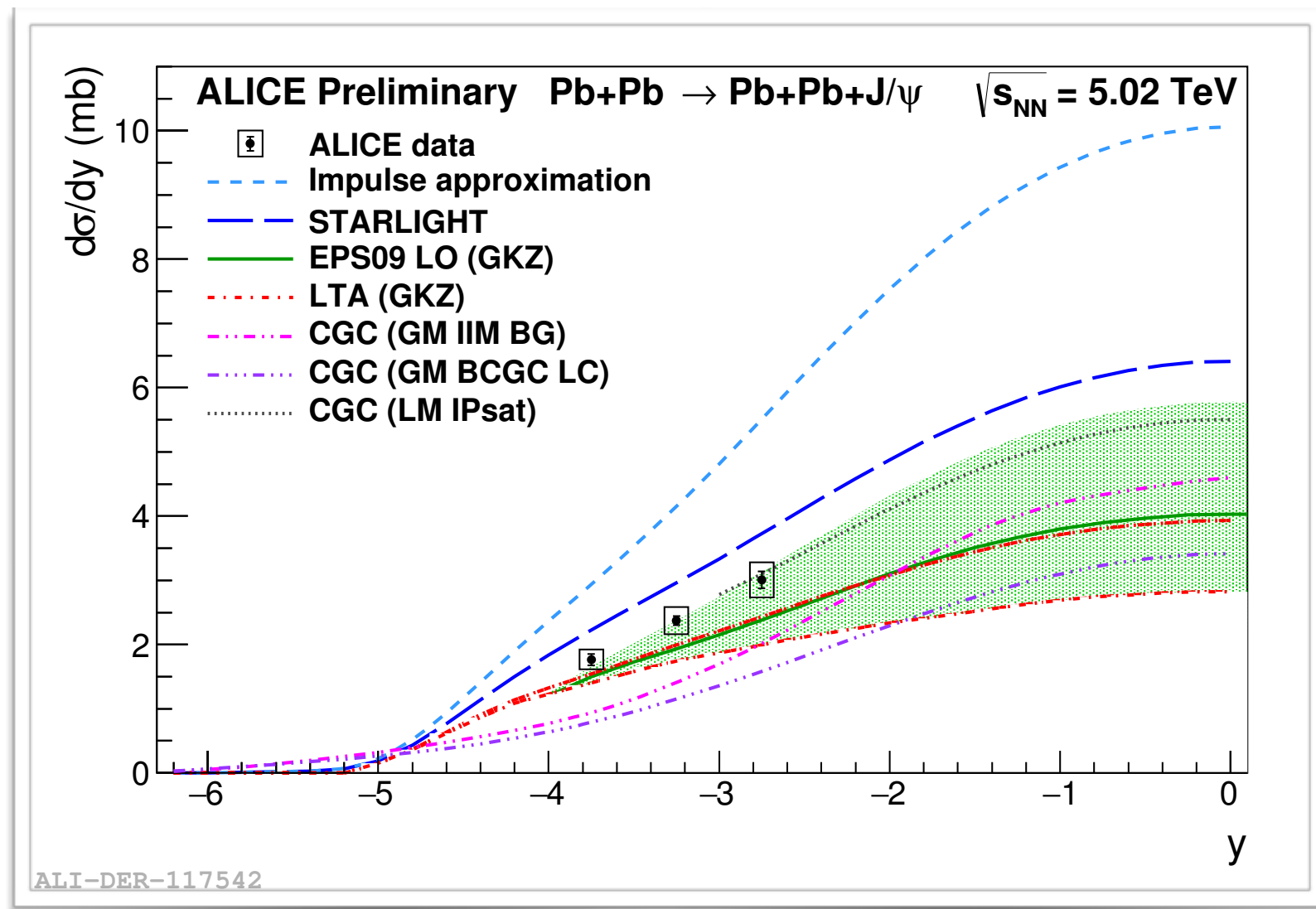


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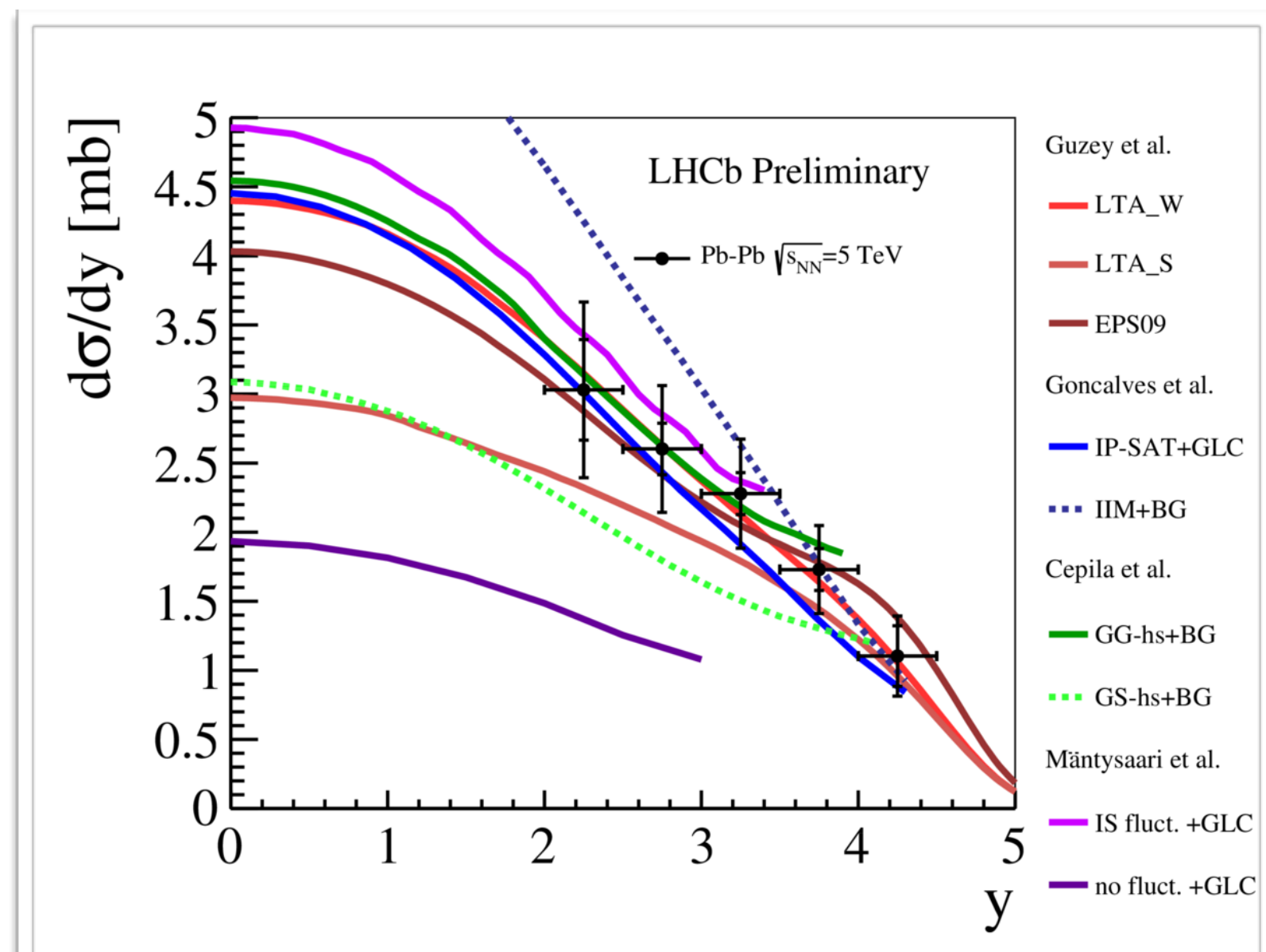
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Data is sensitive to variations in models.

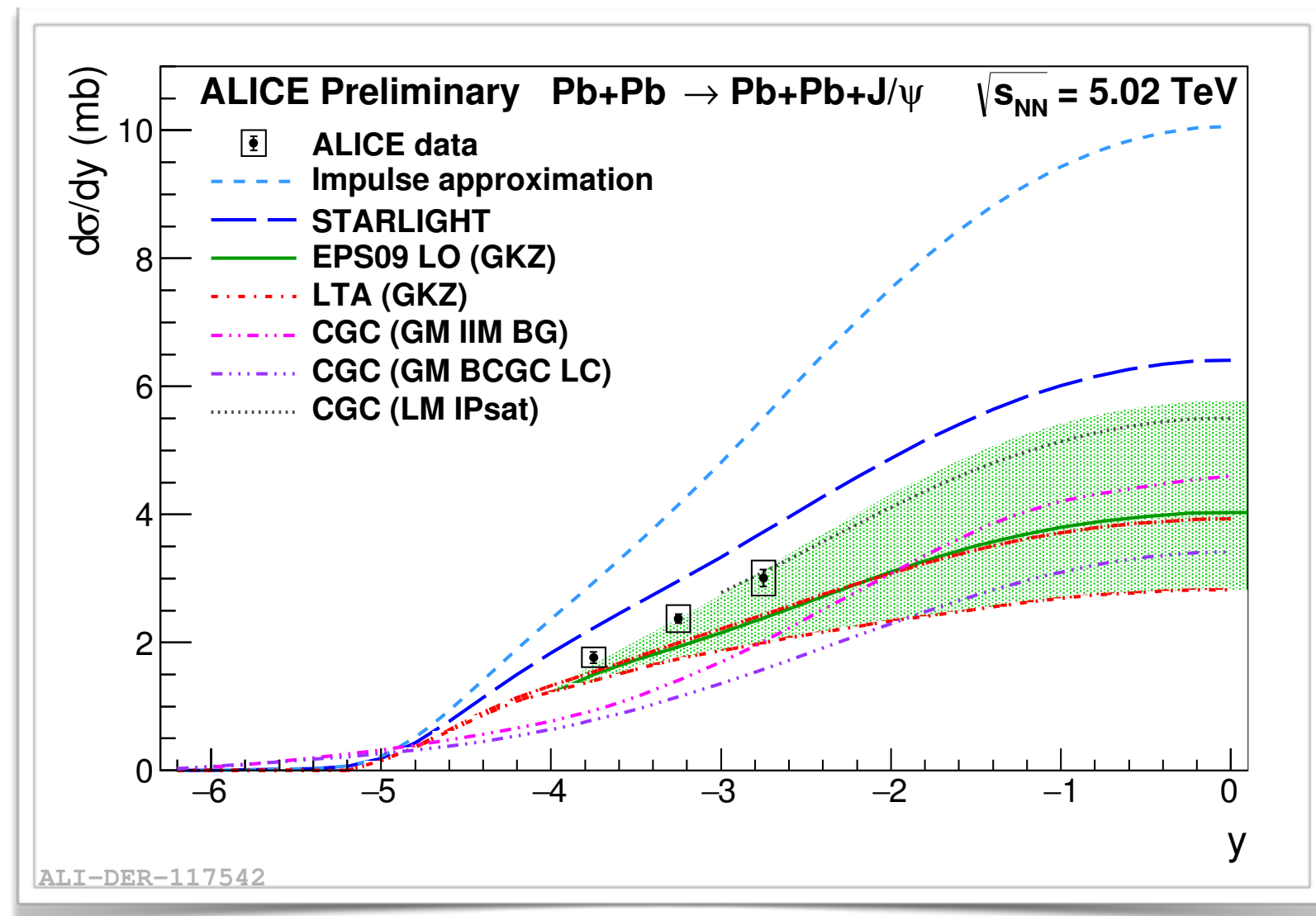
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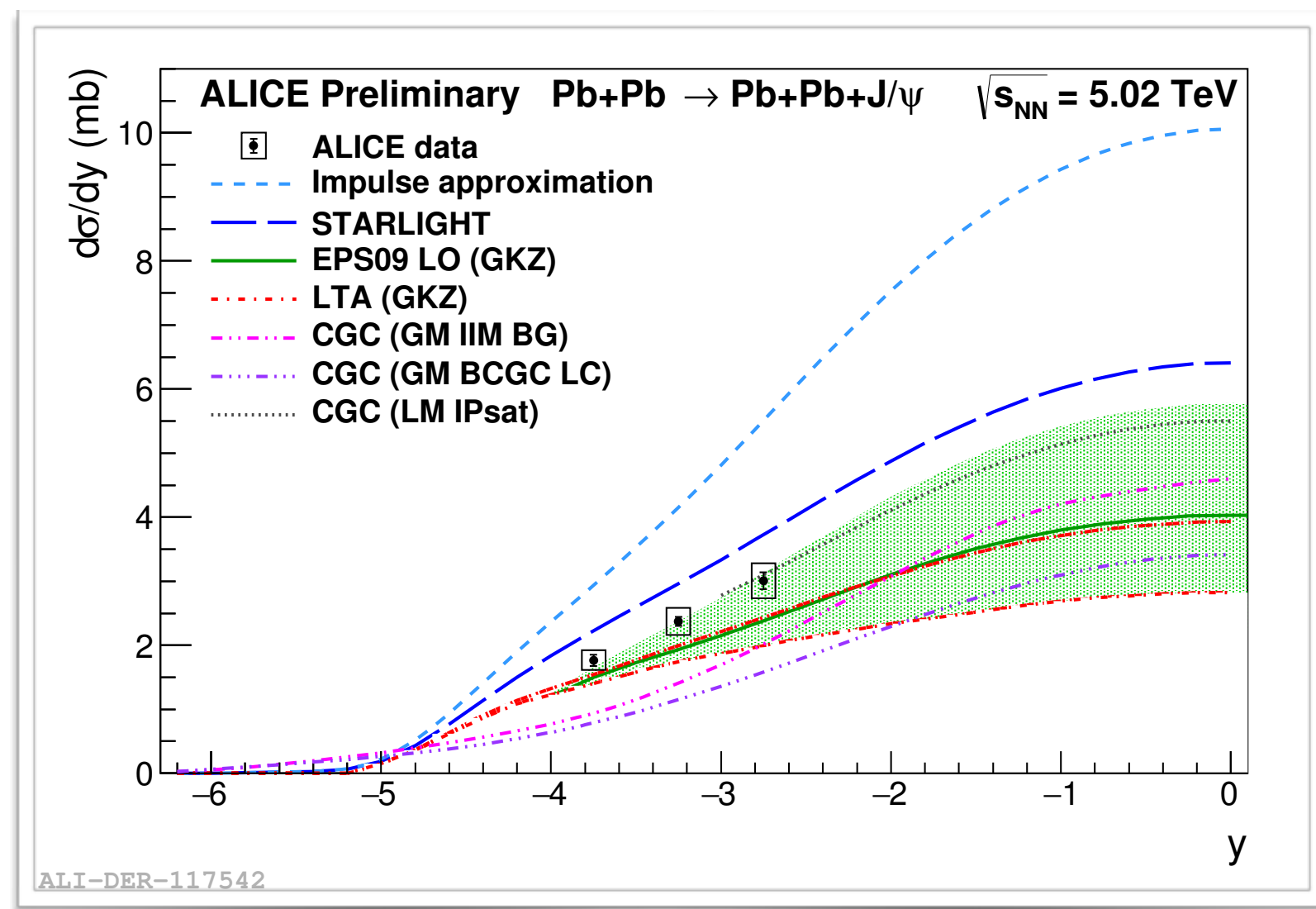
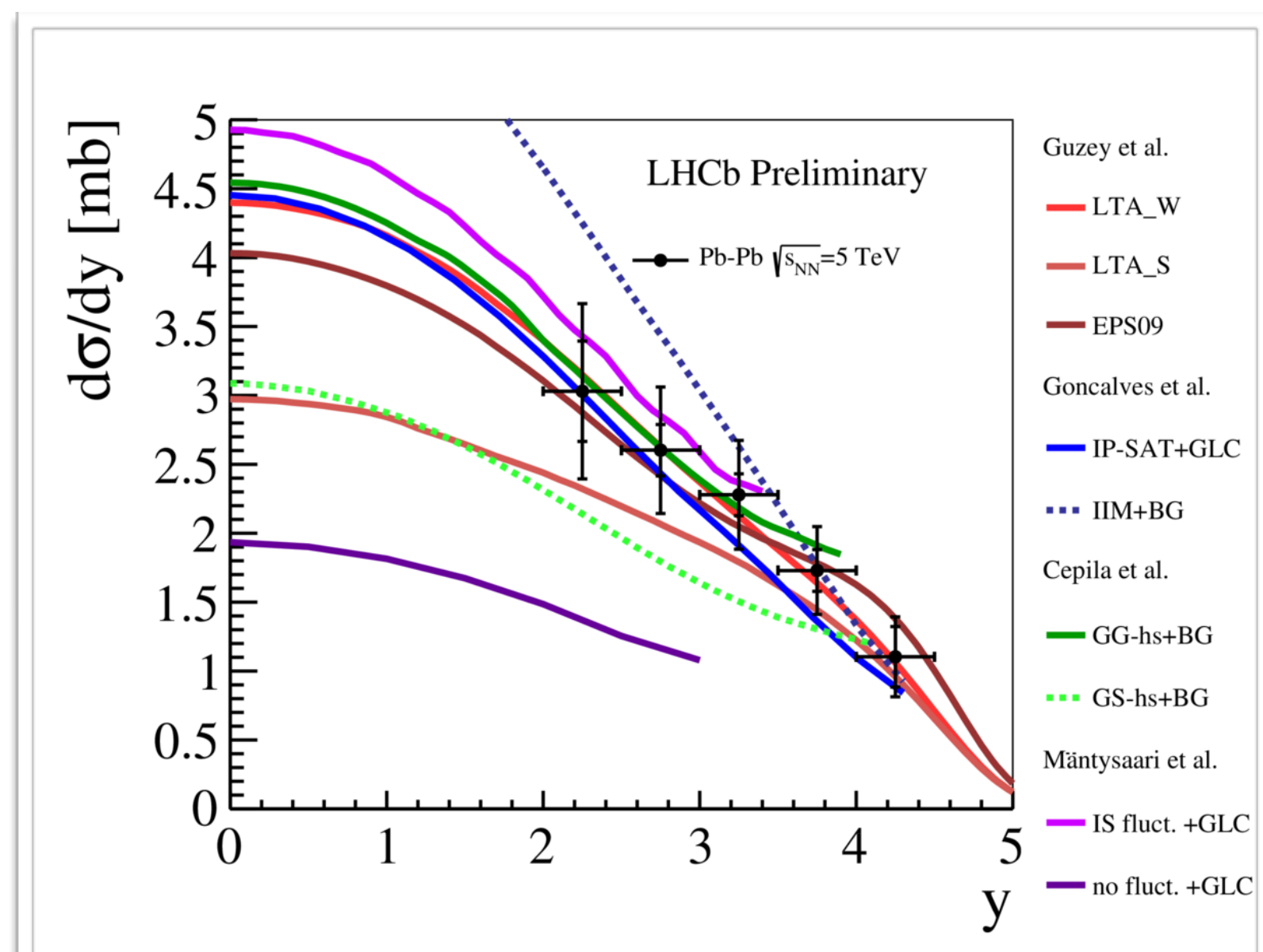
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Wednesday in the WG2 session

Recent ALICE results on coherent J/ψ photoproduction in ultra-peripheral Pb-Pb collisions, Evgeny Kryshen

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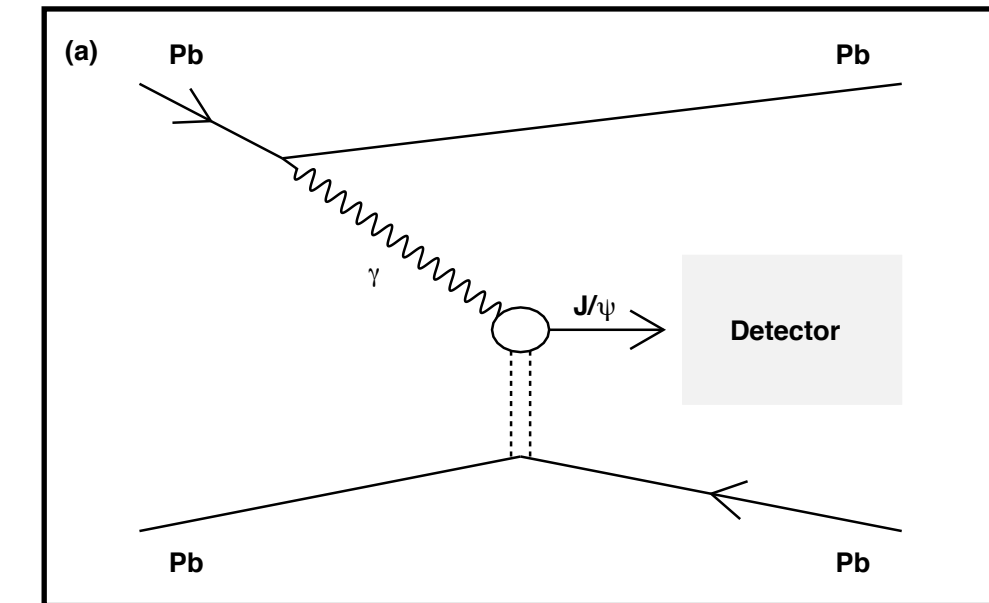
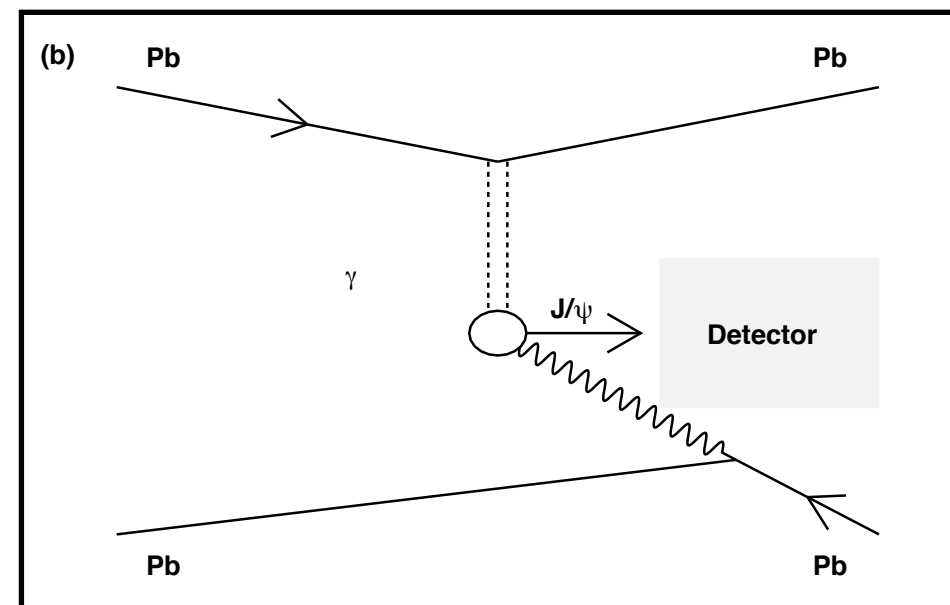
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What is the value of $W_{\gamma Pb}$?

Wednesday in the WG2 session

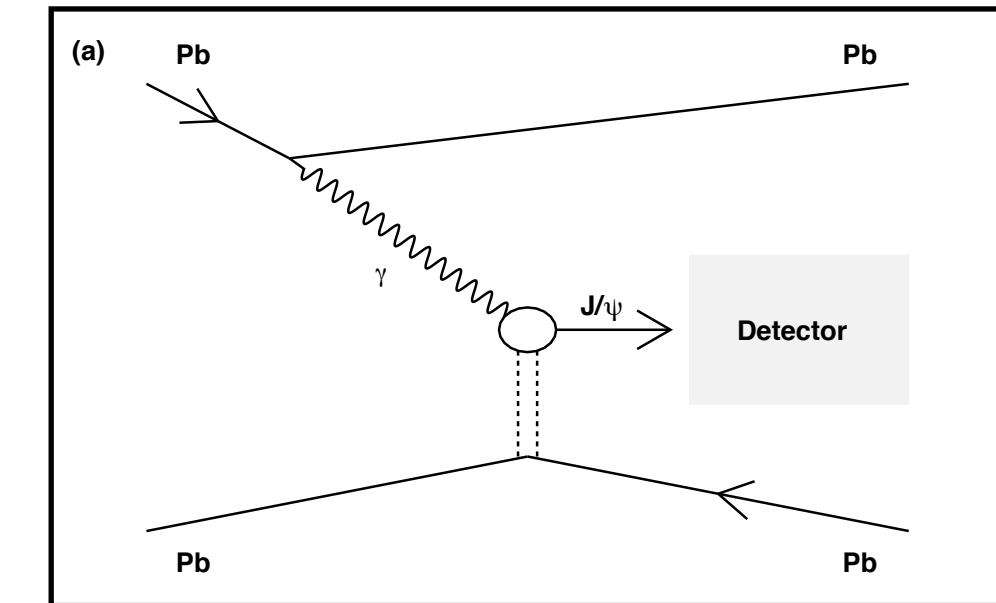
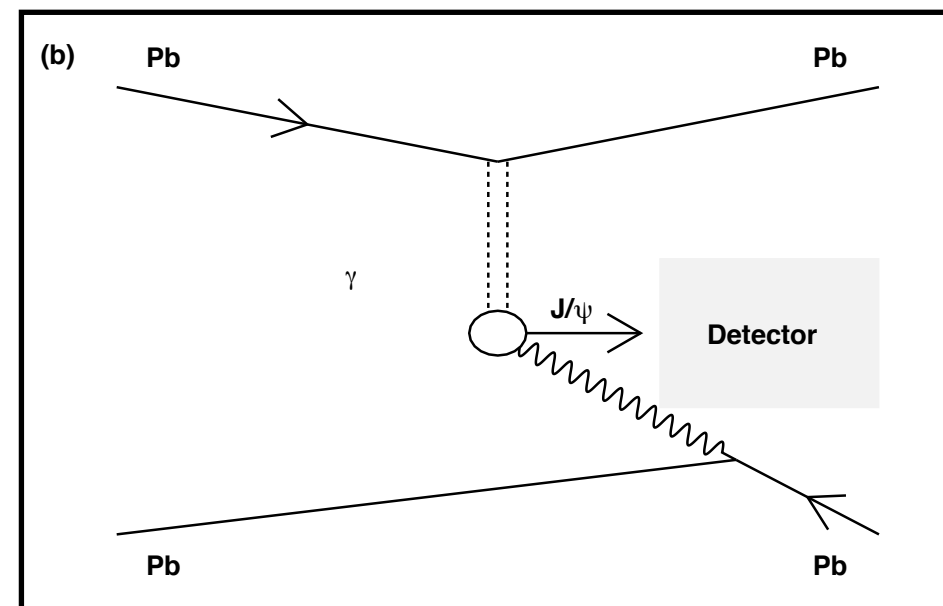
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The UPC and γ Pb cross section



$$\frac{d\sigma_{\text{PbPb}}}{dy} = n_{\gamma}(y; b_{1,2})\sigma_{\gamma\text{Pb}}(y) + n_{\gamma}(-y; b_{1,2})\sigma_{\gamma\text{Pb}}(-y)$$

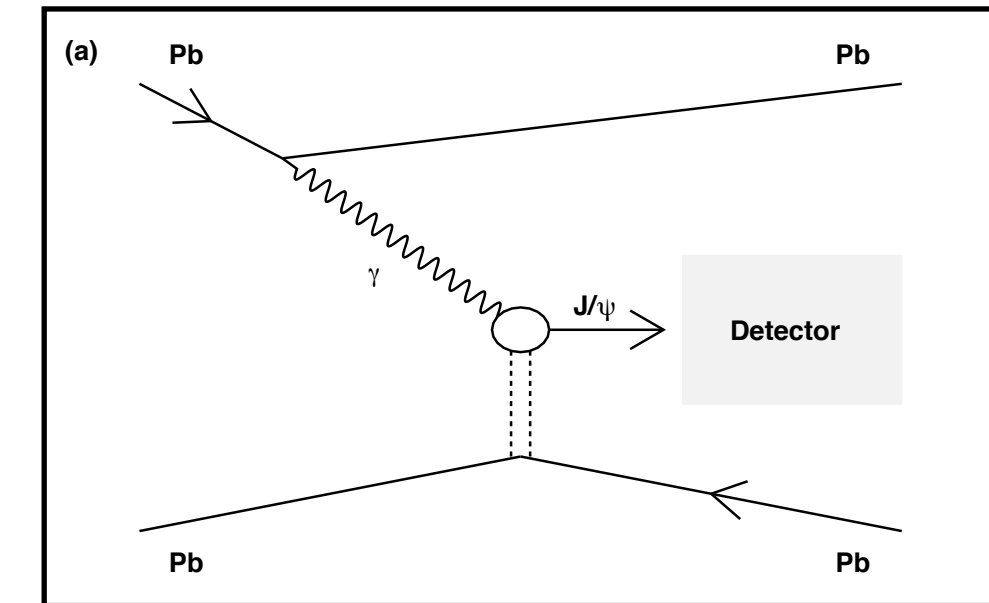
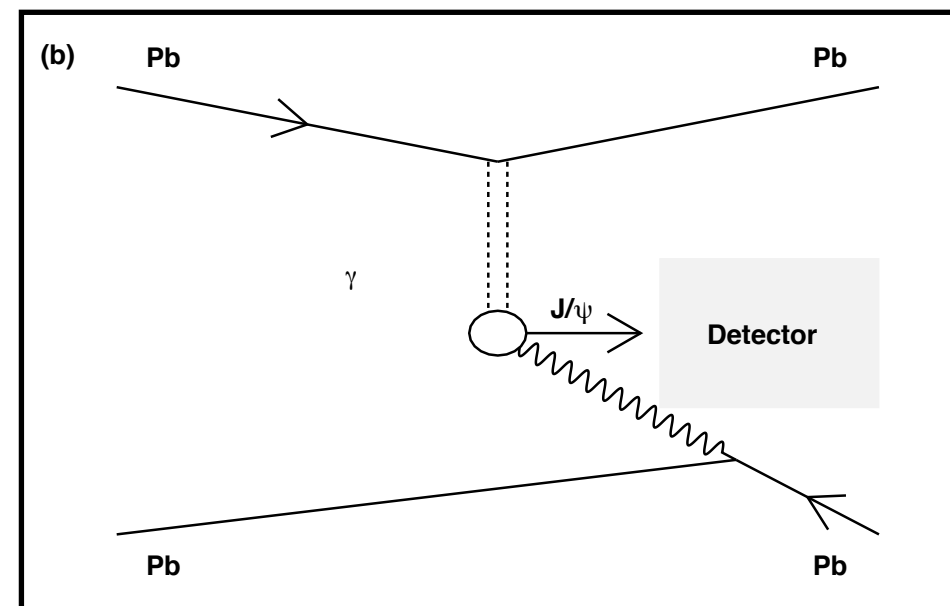
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What we measure.

The UPC and γ Pb cross section

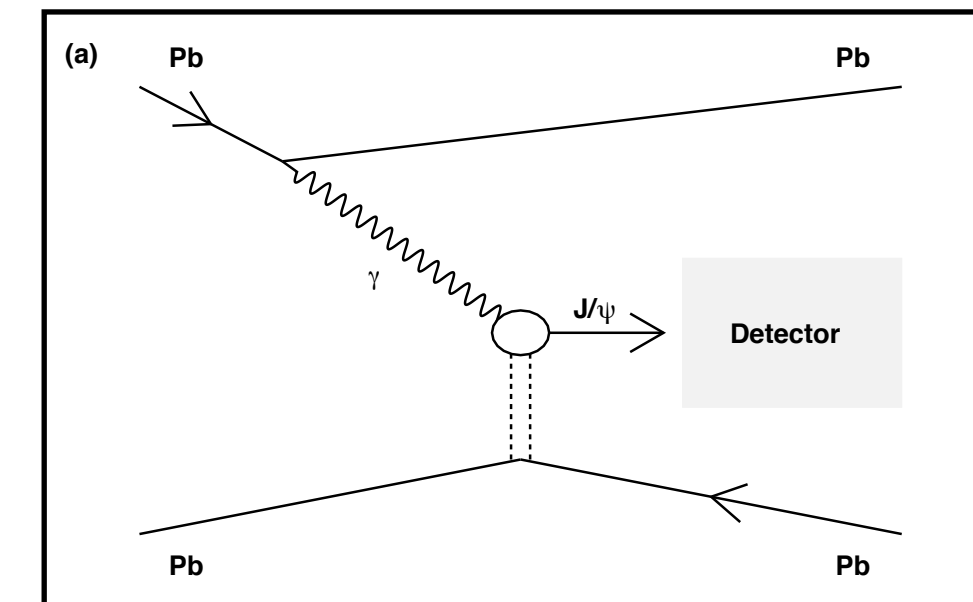
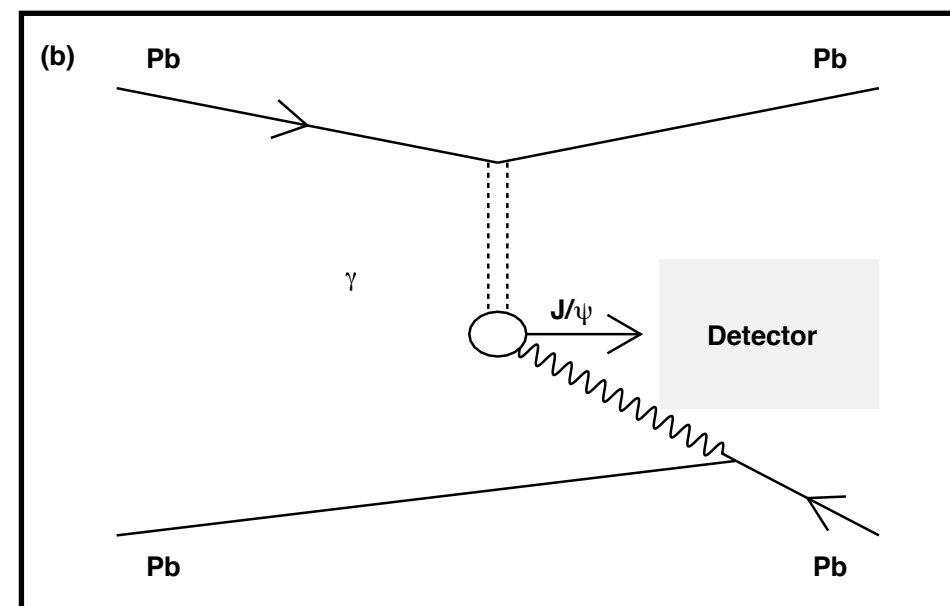


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Emission of the photon by the source.
We assume that we know it.

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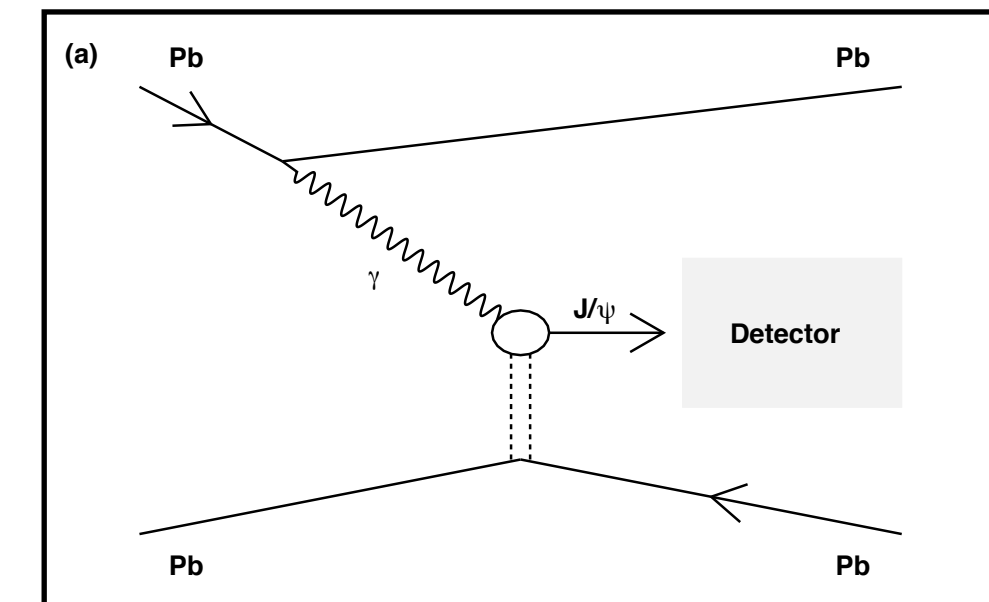
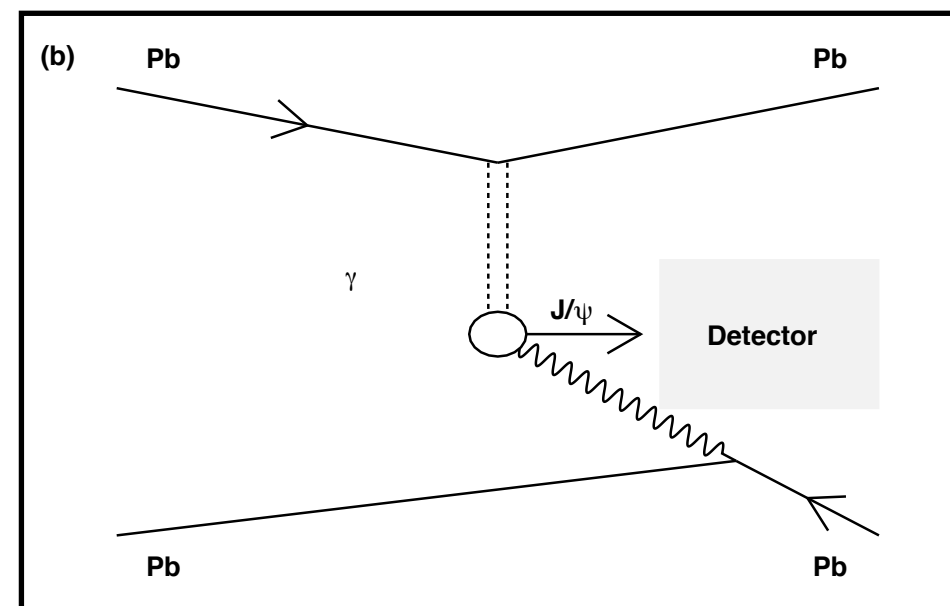
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What we want: Cross section for the
photon-Pb system at each rapidity.

The UPC and γ Pb cross section



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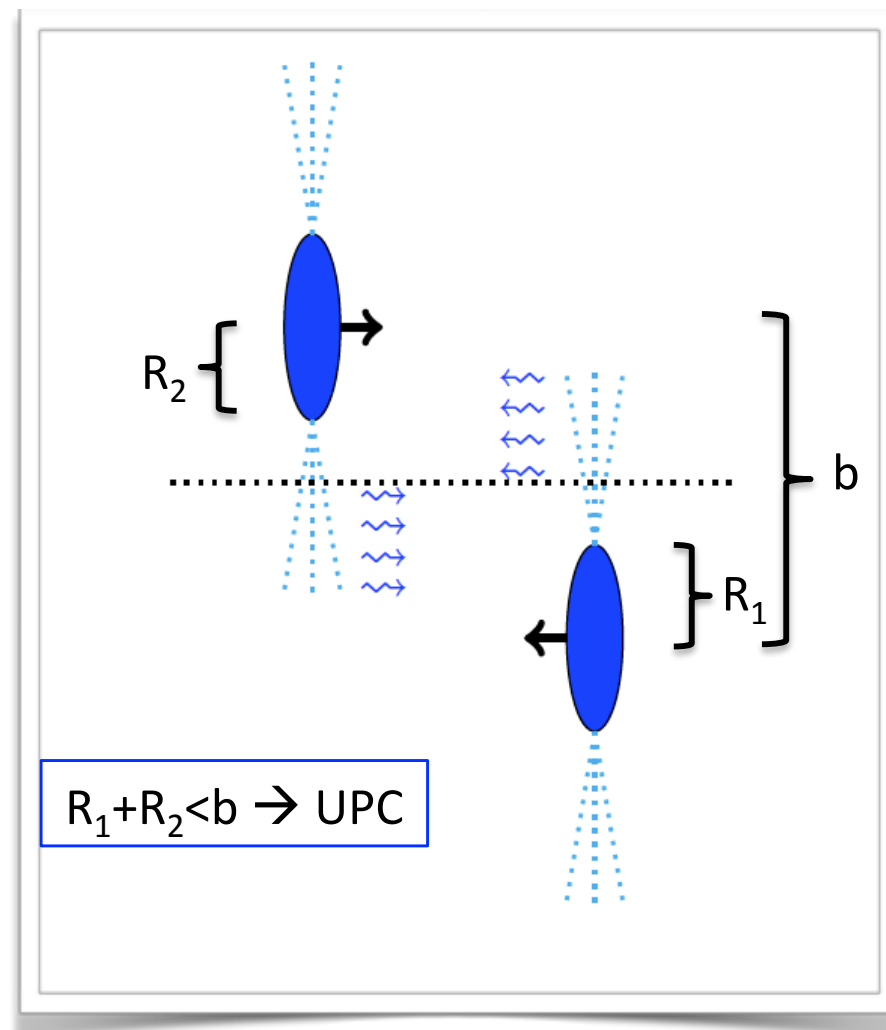
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What we want: Cross section for the
photon-Pb system at each rapidity.

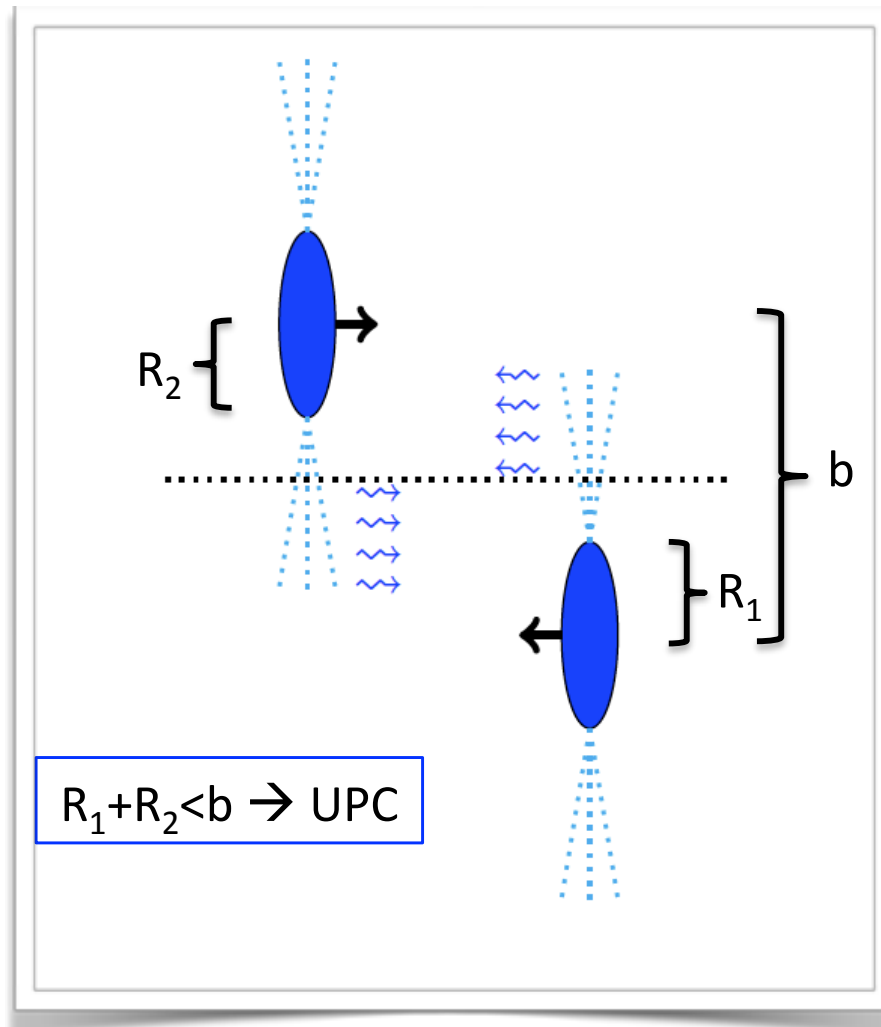
- At $y=0$ both UPC terms are equal.
- In pPb UPC one term is much larger than the other.
- In PbPb collisions we need to perform **the measurement at the same rapidity but at different impact parameters** in order to disentangle both contributions.

Coherent photoproduction in peripheral collisions @ 2.76 TeV



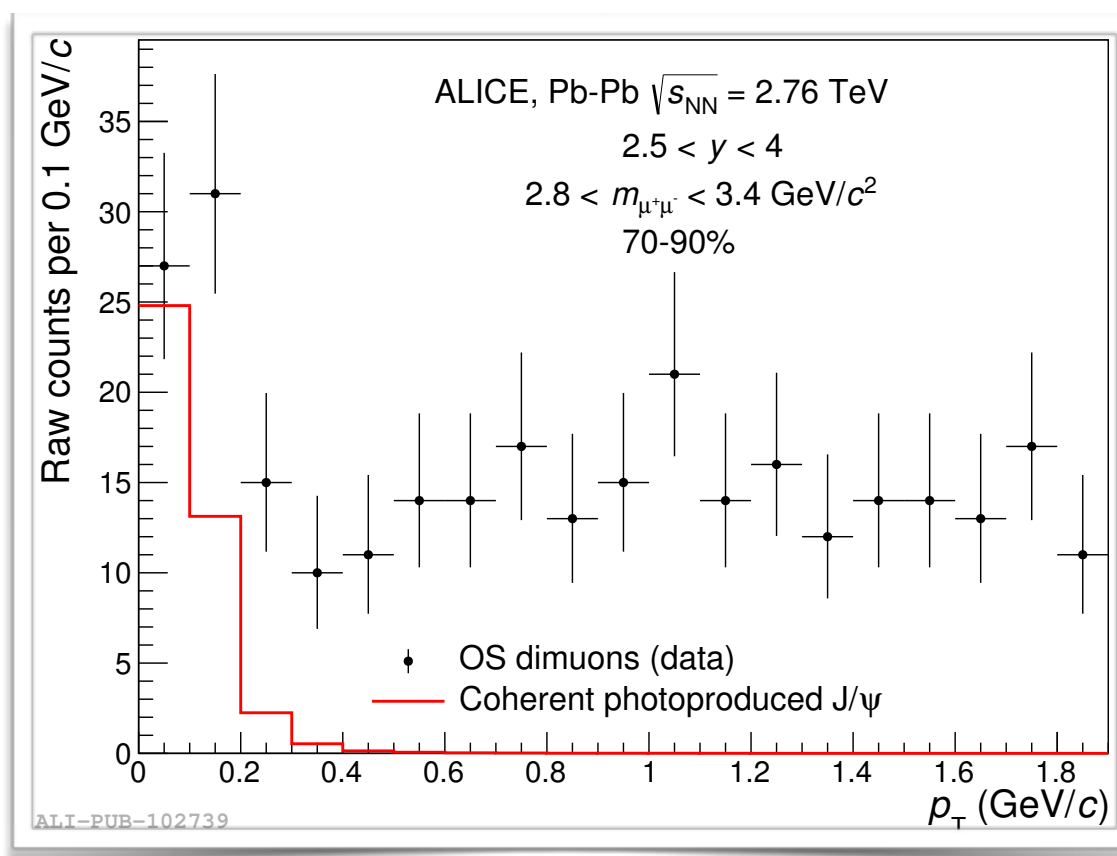
What happens if b is slightly smaller than $R_1 + R_2$?
The EM fields will still be there
... but the incoming nuclei will break.

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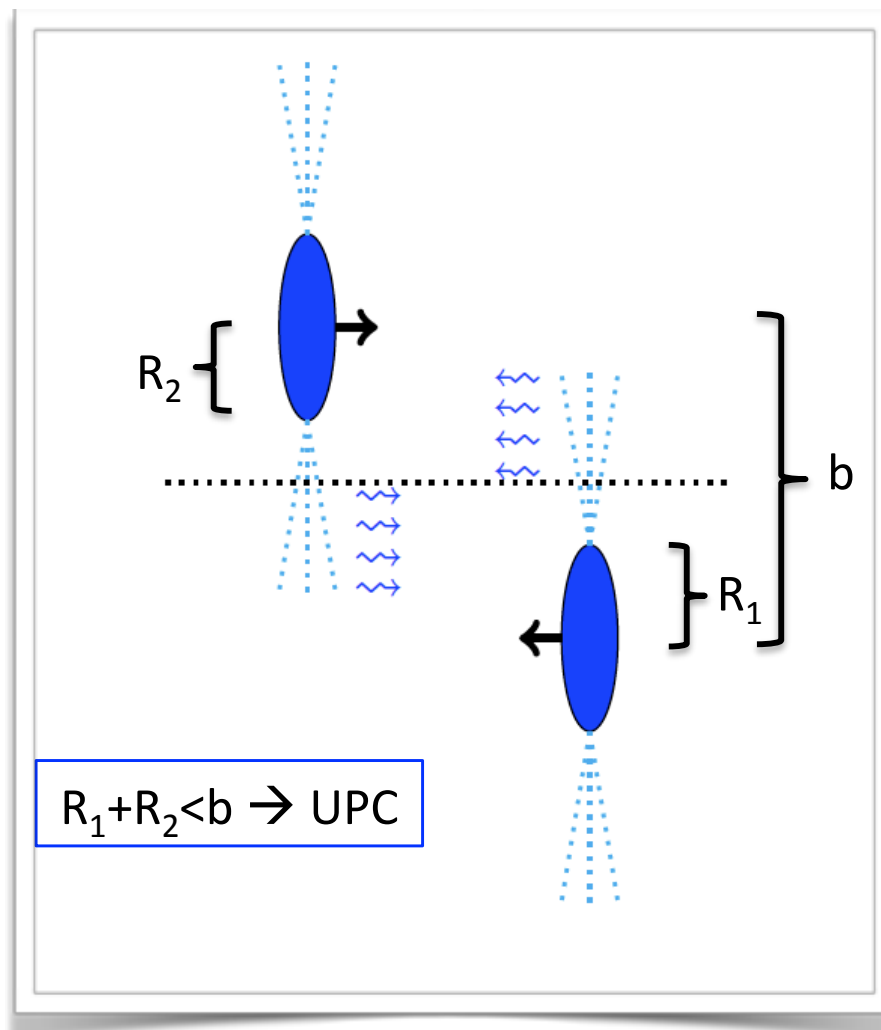
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ALICE: Phys.Rev.Lett. 116 (2016) 222301



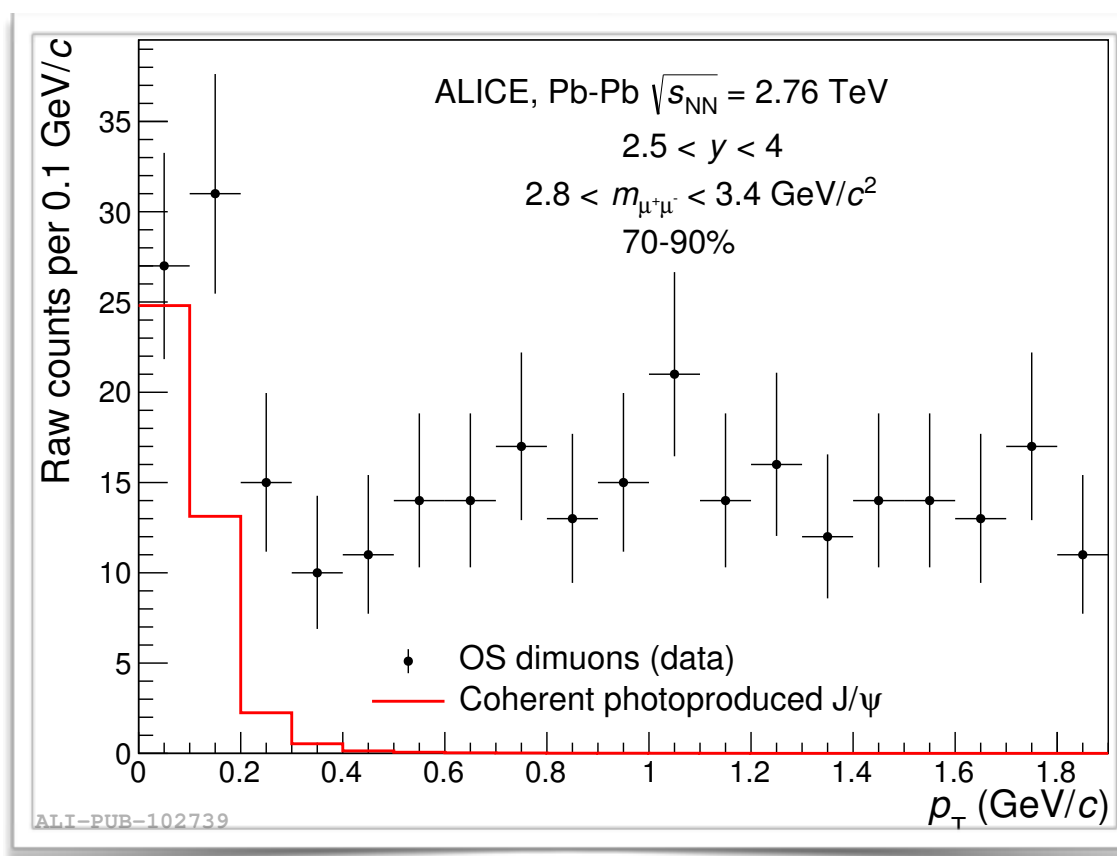
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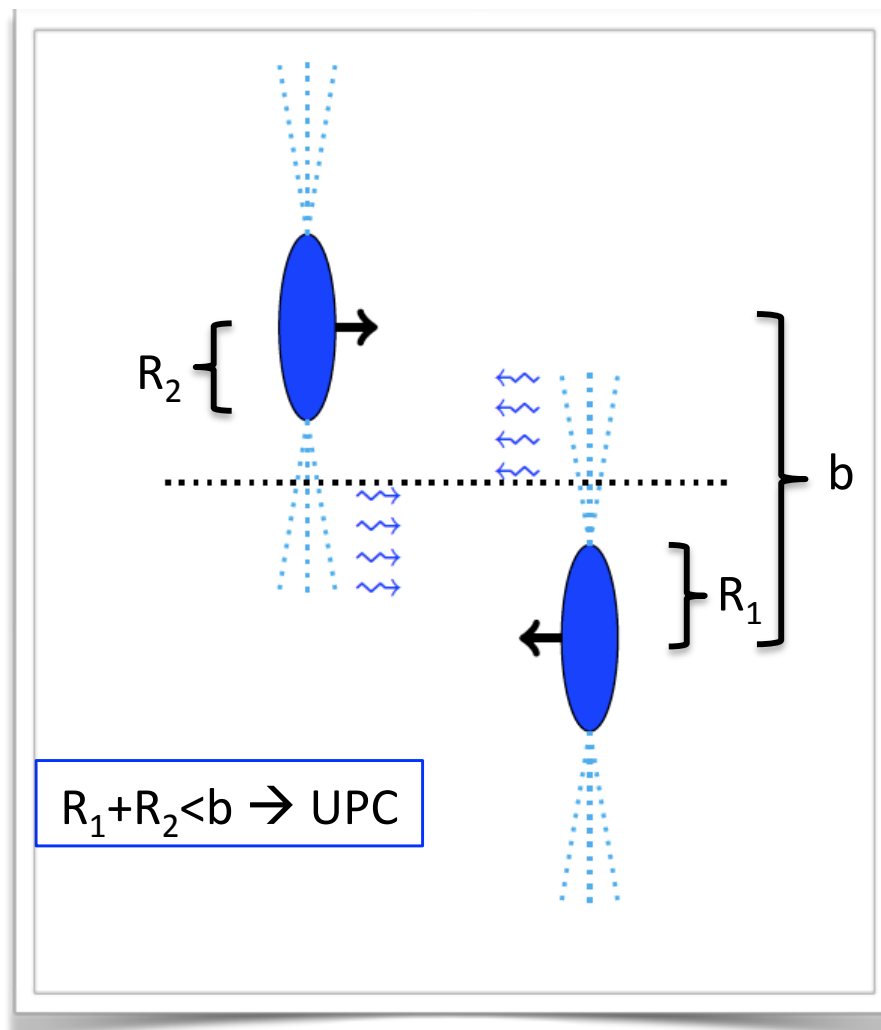
ALICE: Phys.Rev.Lett. 116 (2016) 222301



Cent. (%)	$d\sigma_{J/\psi}^{\text{coh}}/dy$ (μb)
0–10	< 318
10–30	< 290
30–50	$73 \pm 44^{+26}_{-27} \pm 10$
50–70	$58 \pm 16^{+8}_{-10} \pm 8$
70–90	$59 \pm 11^{+7}_{-10} \pm 8$

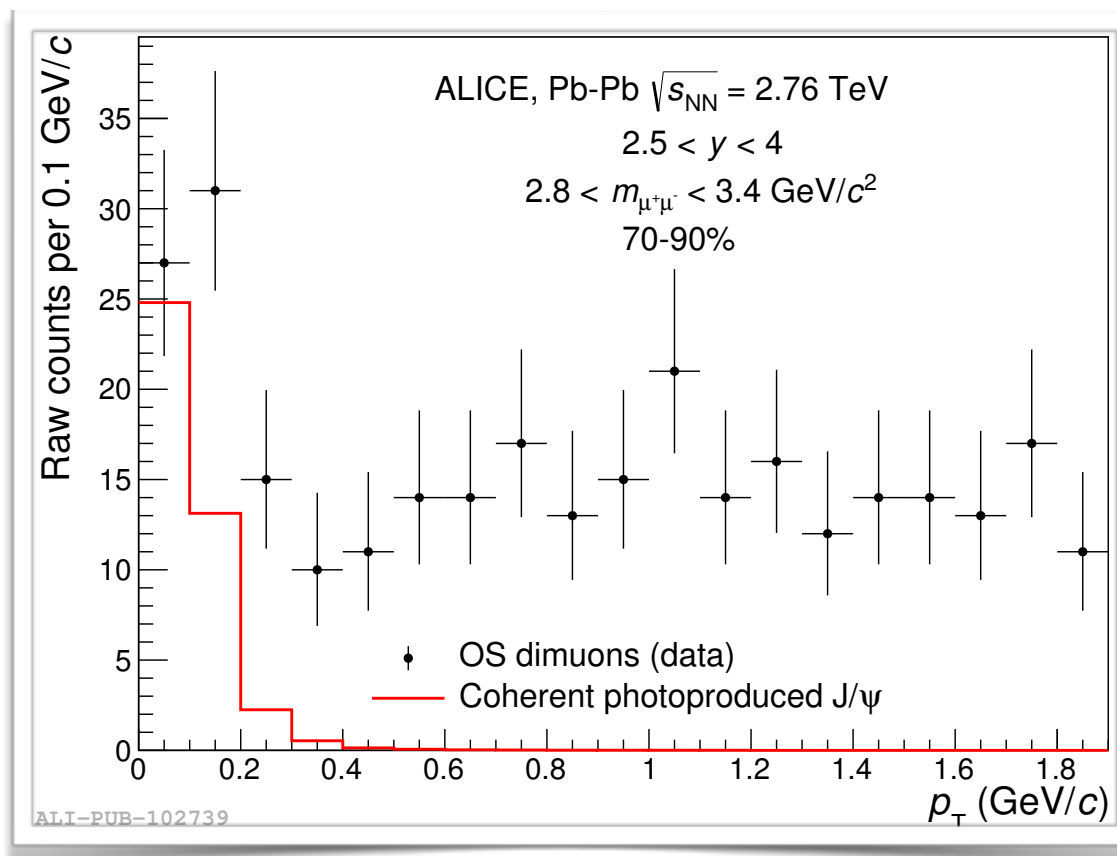
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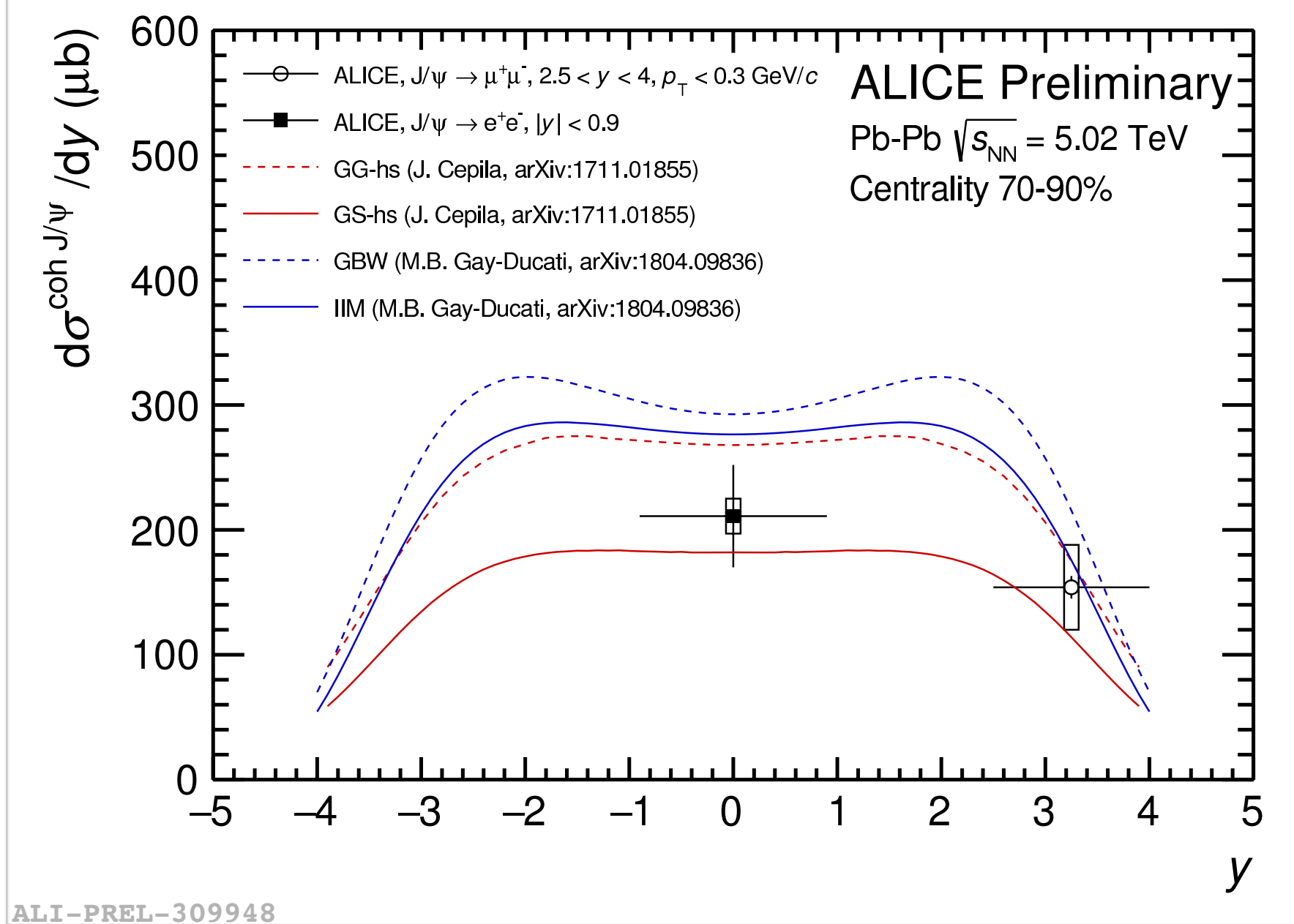


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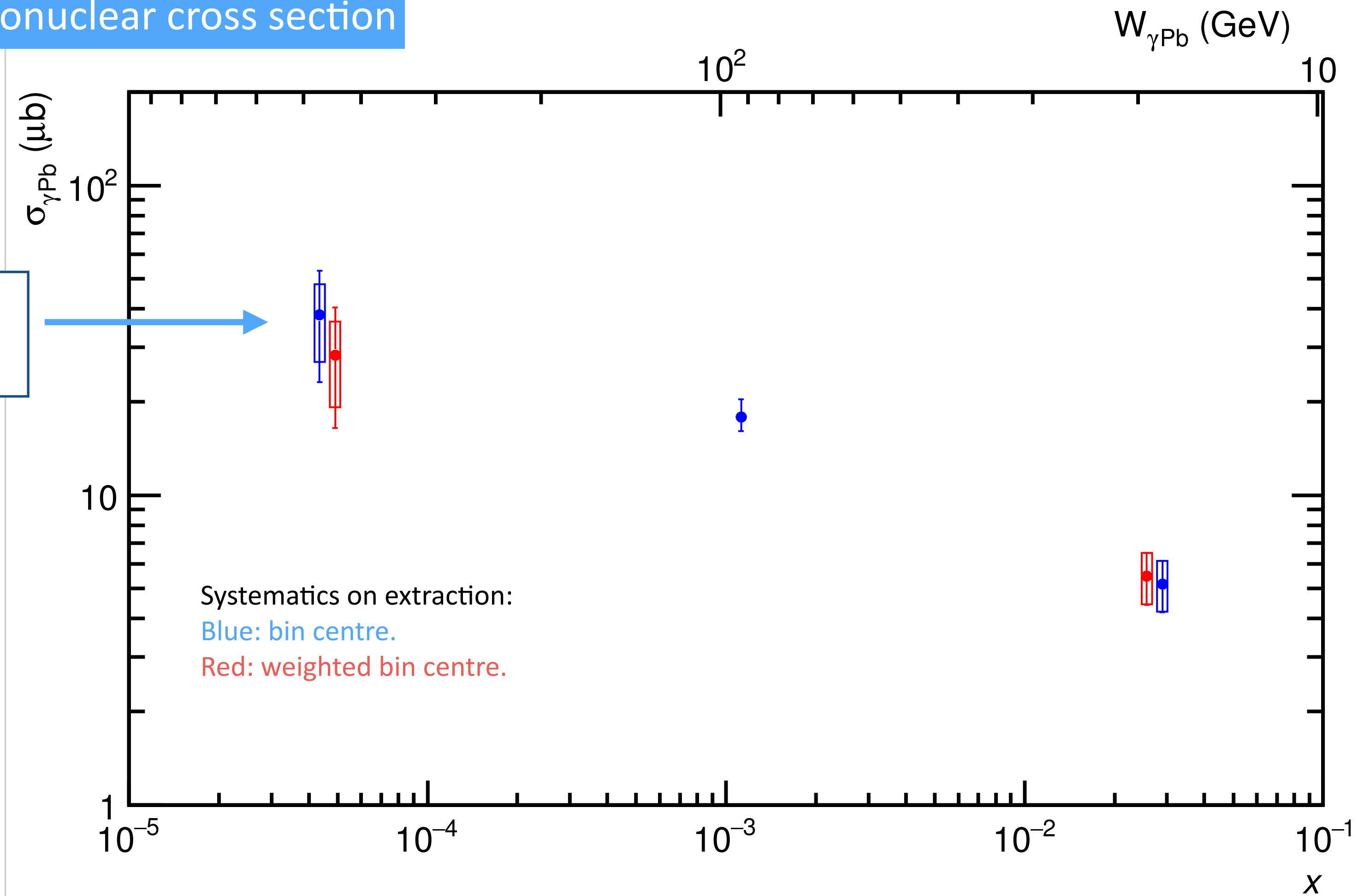


Coherent photoproduction in peripheral collisions!

Photonuclear cross section extracted from ALICE data

Photonuclear cross section

$W_{\gamma\text{Pb}}$ up to 470 GeV,
x down to 4×10^{-5} !

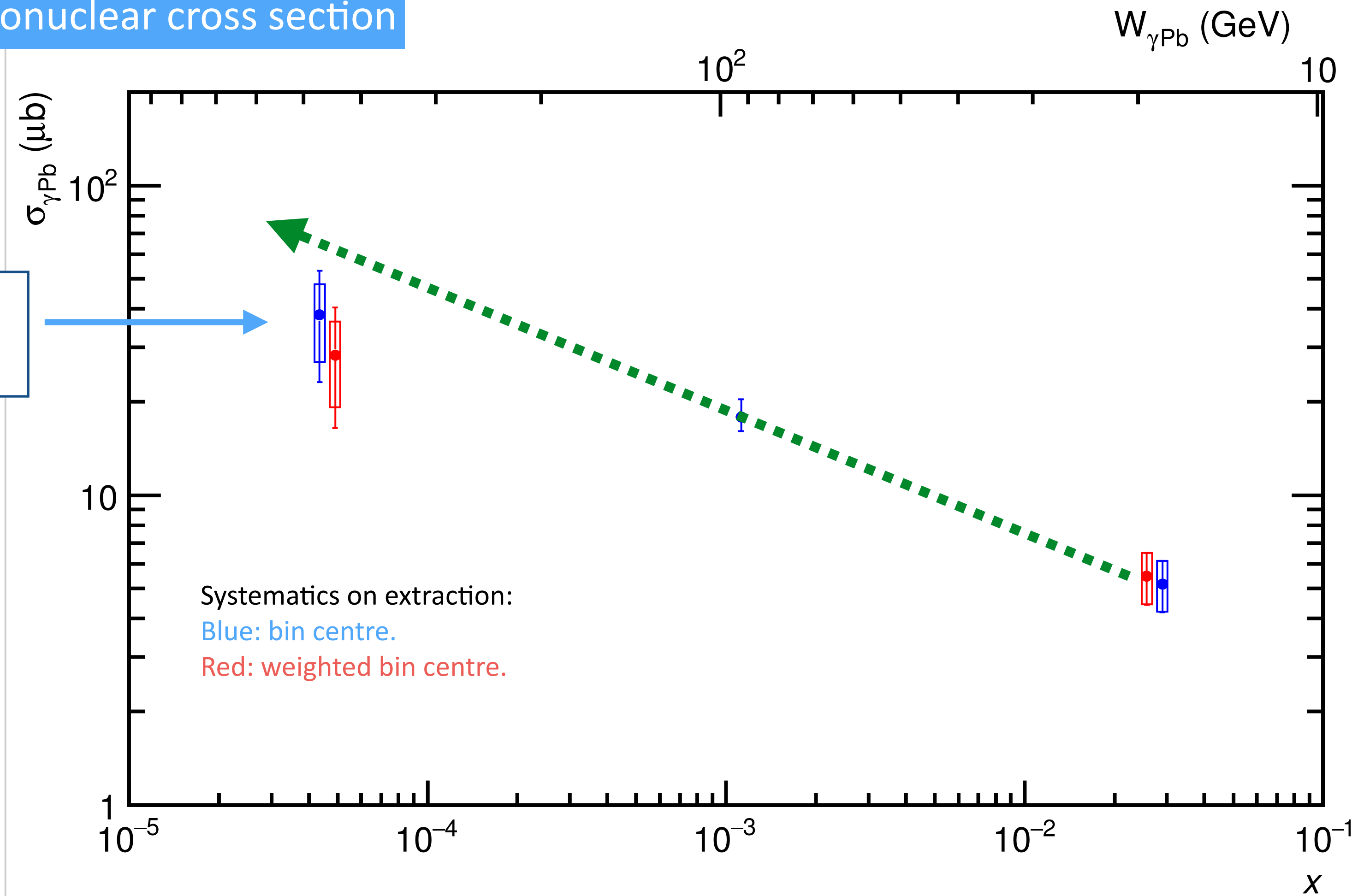


$$x = \left(\frac{M}{W} \right)^2$$

Photonuclear cross section extracted from ALICE data

Photonuclear cross section

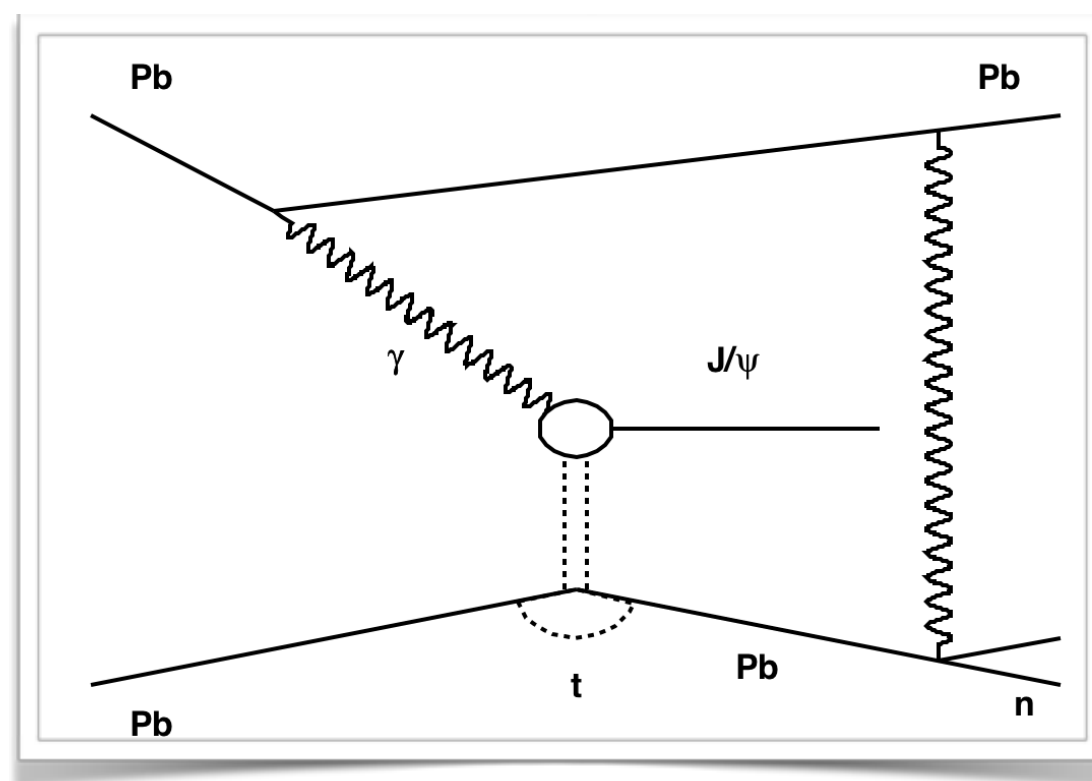
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Second option: use forward neutrons!

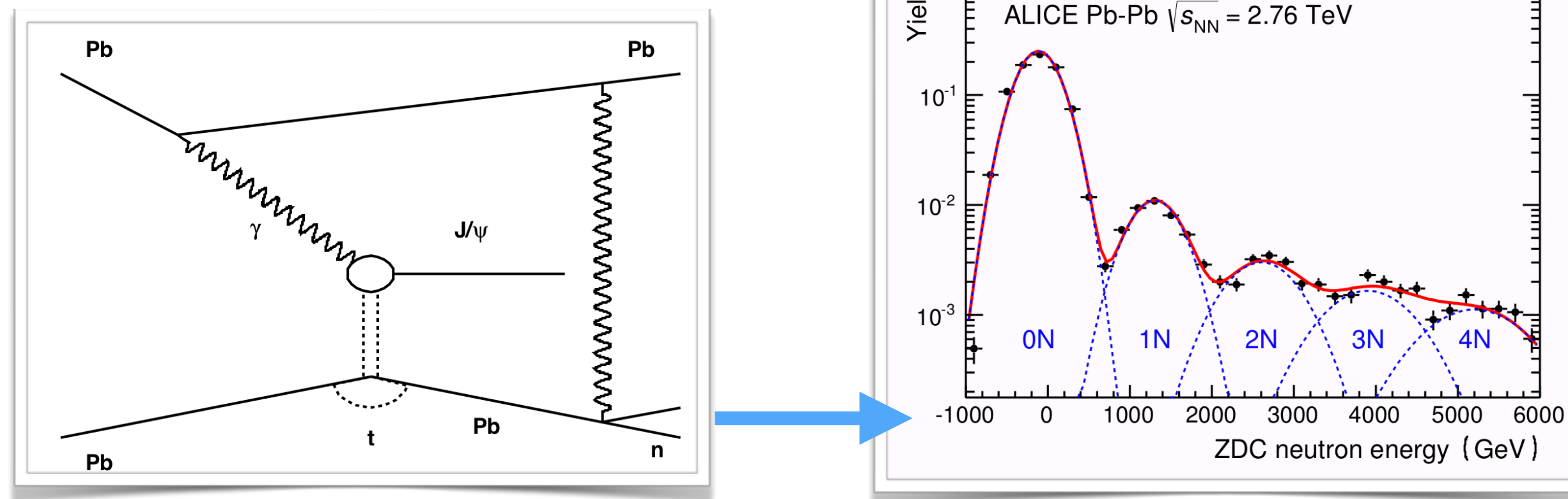
Independent soft electromagnetic interactions between the nuclei excite them; upon de-excitation they **emit neutrons** in the forward direction



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ALICE: JHEP 09 (2015) 095

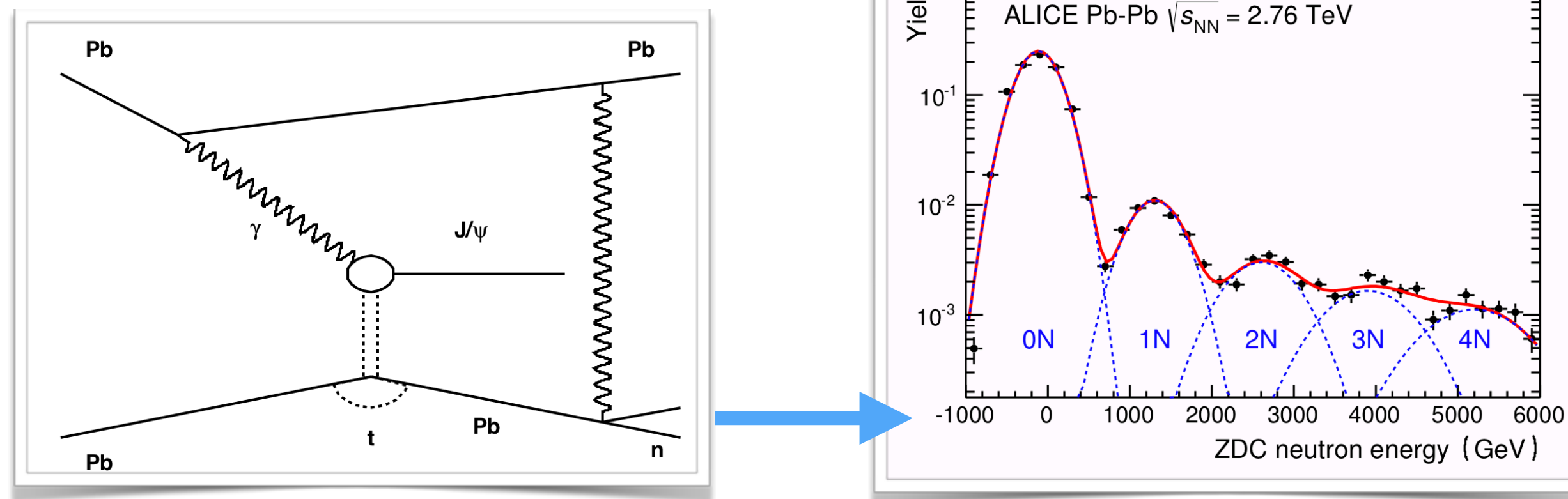


LHC/RHIC experiments can measure these neutrons!

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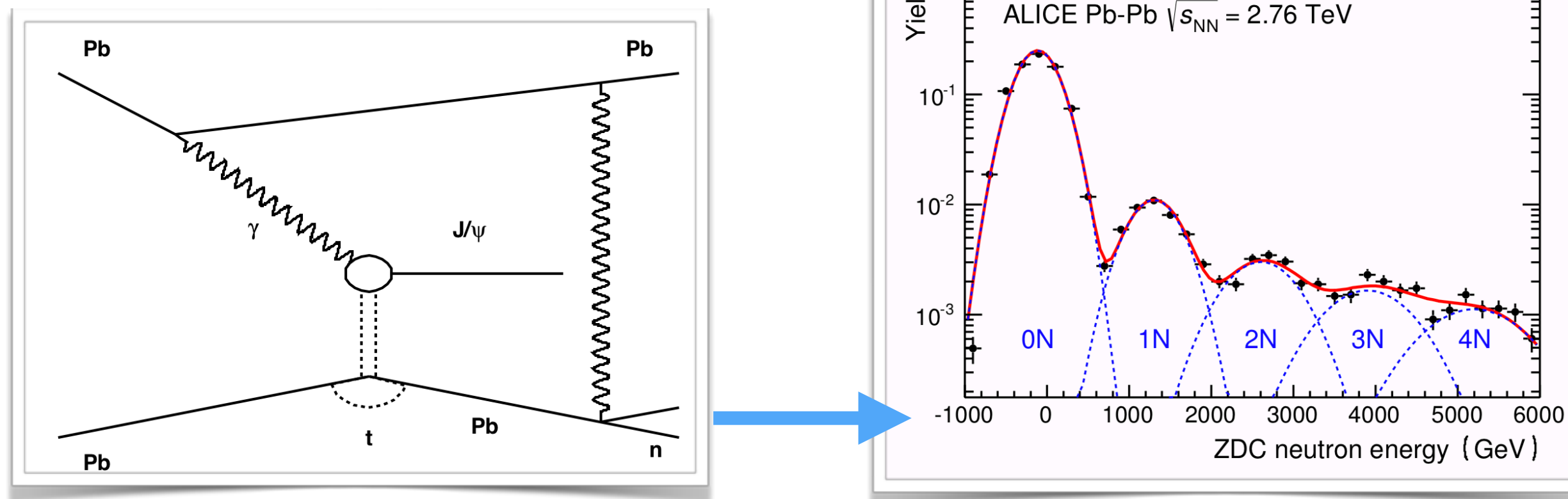
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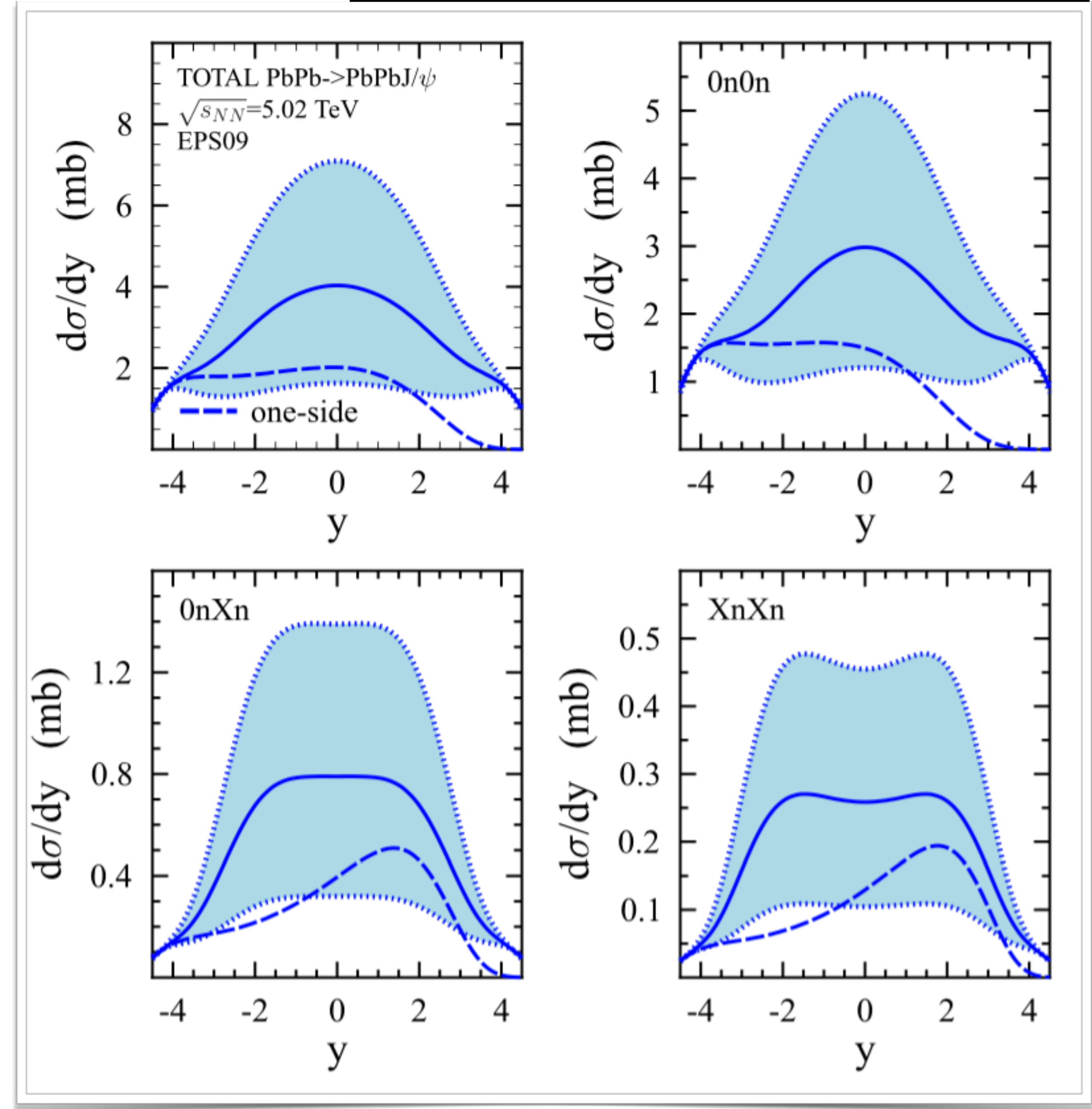
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Guzay, Kryshen, Zhalov PRC 93 (2016)055206



(a parenthesis)

Photon induced processes as a potential probe of QGP

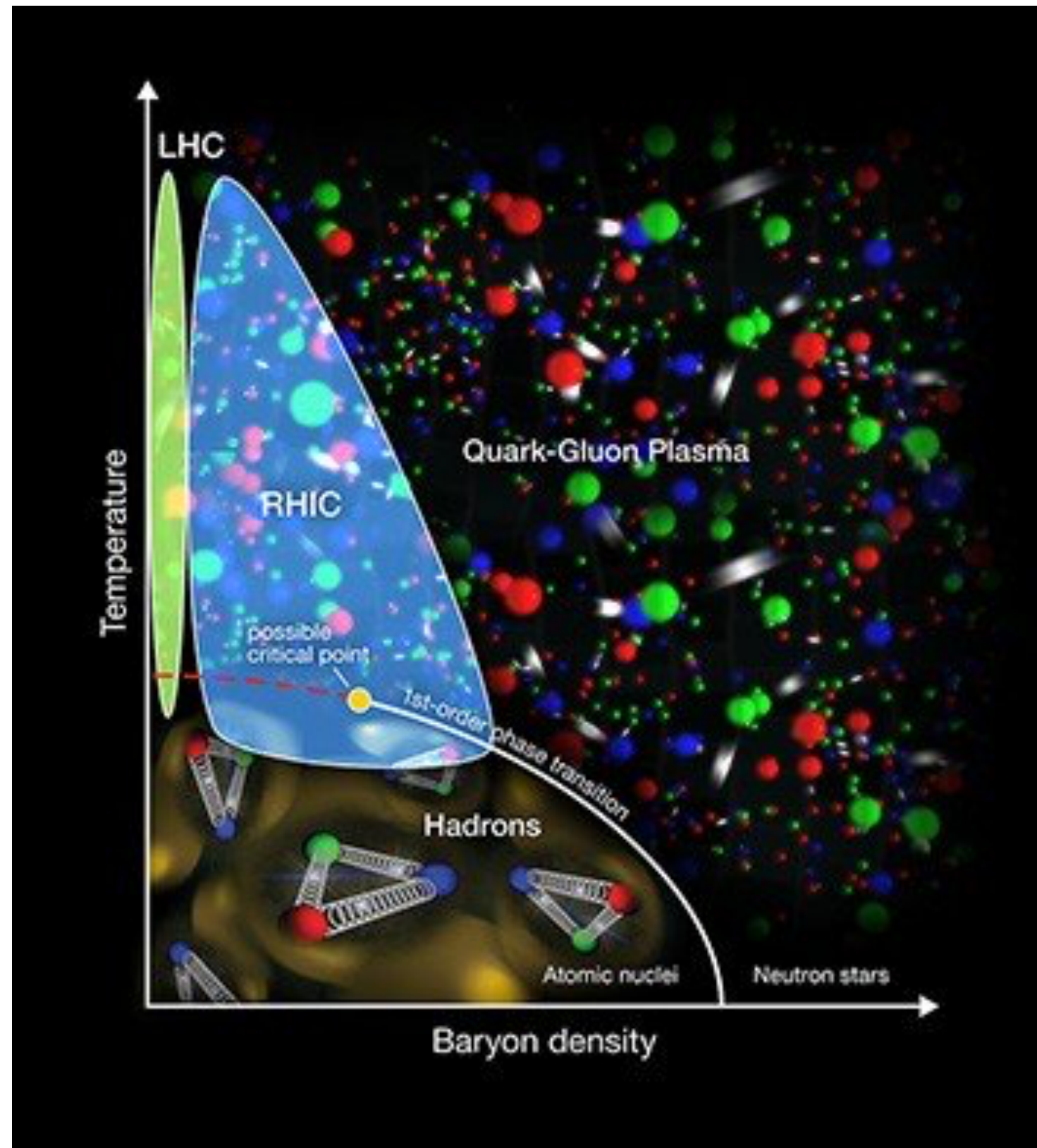
The quark-gluon plasma and where to find it

QCD is richer than just hadrons.
It has a rich phase space with different states of matter.

The quark-gluon plasma and where to find it

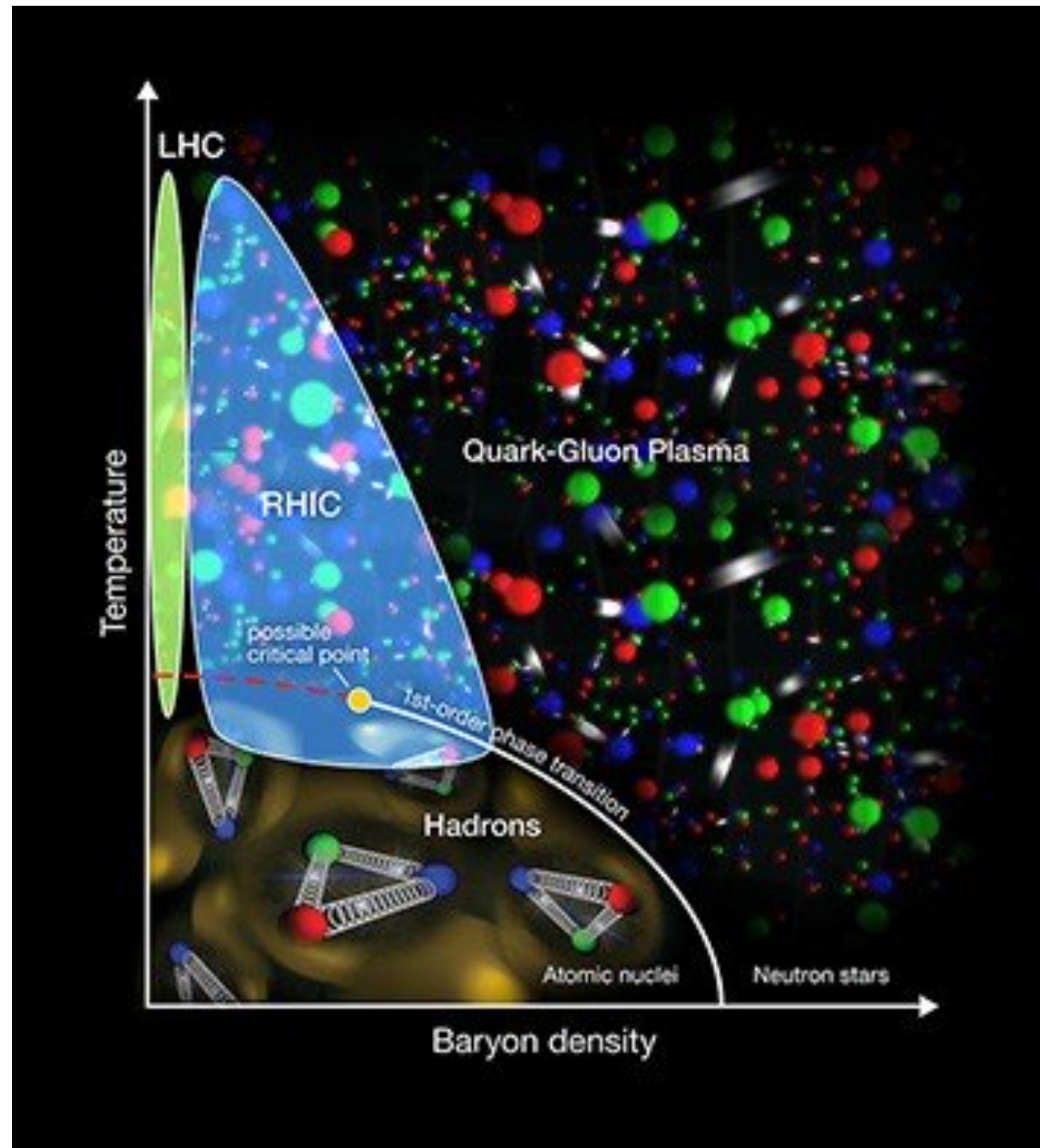
QCD is richer than just hadrons.
It has a rich phase space with different states of matter.

The study of the quark-gluon plasma (QGP) and its properties is one of the main topics of research at RHIC and LHC where such a state of matter is created in head-on (so called 'central') collisions of heavy nuclei.



Credit: Brookhaven National Laboratory

The quark-gluon plasma and where to find it



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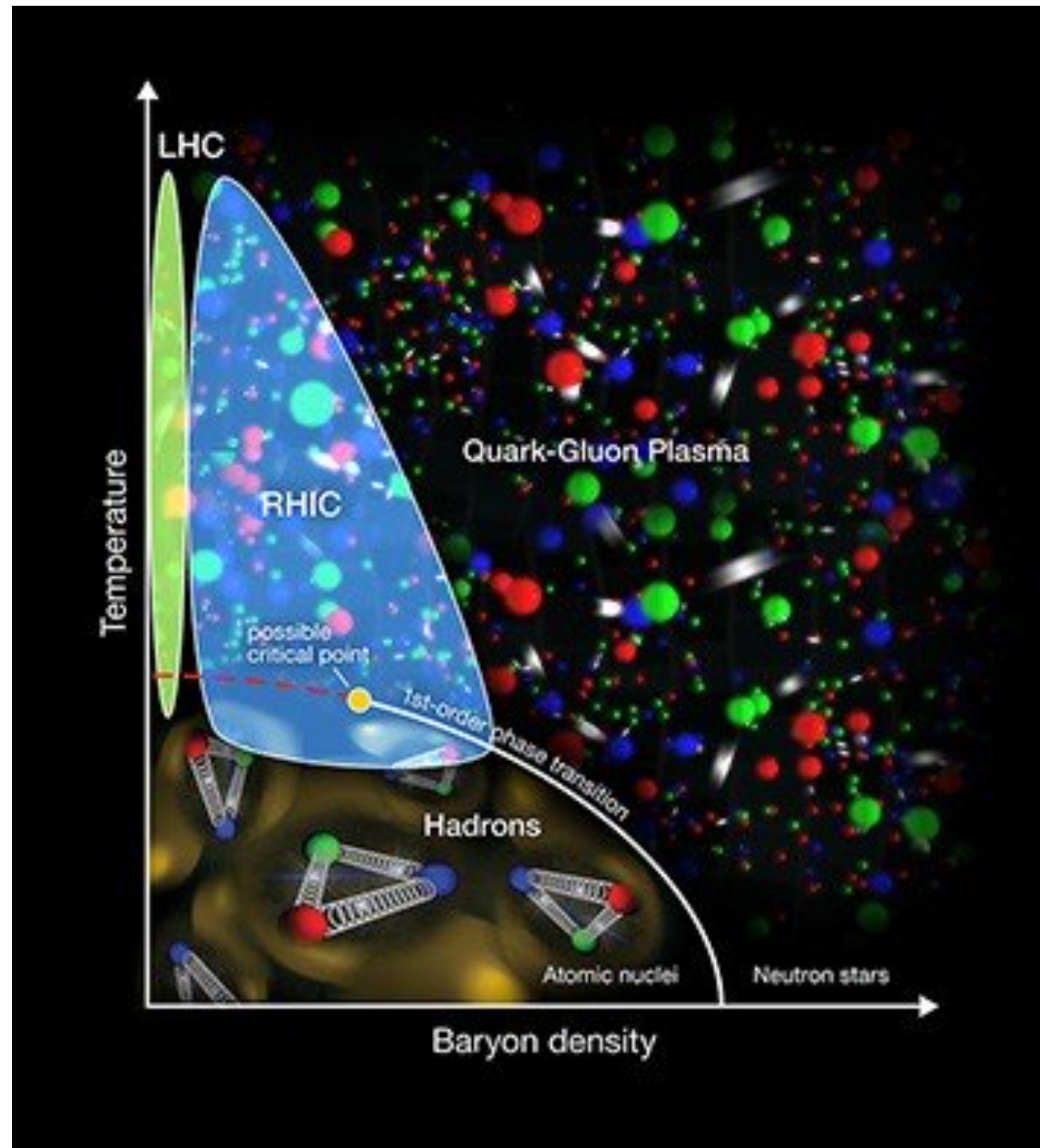
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The quark-gluon plasma and where to find it



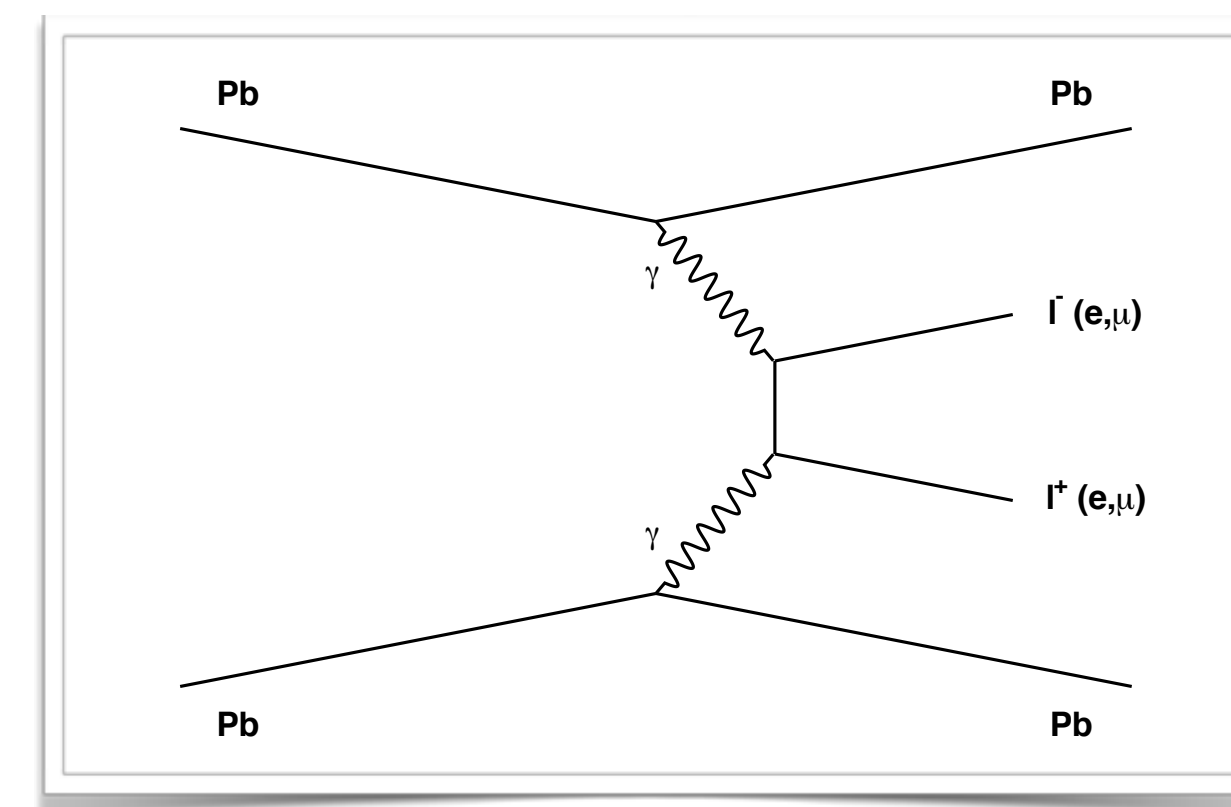
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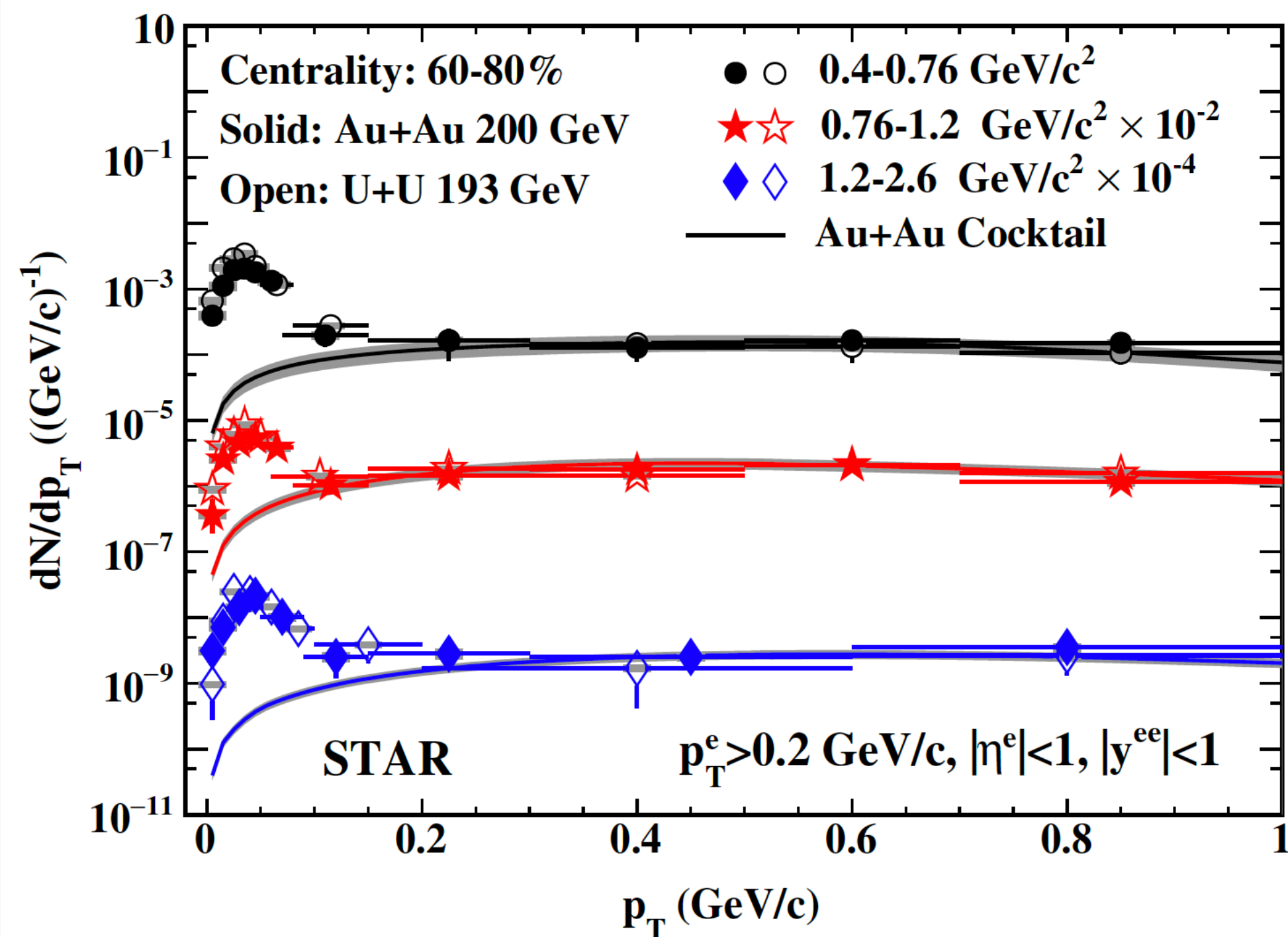
It is expected that the QGP creates strong electromagnetic fields.

Can we use photon-induced processes to study the QGP?



News from RHIC: electron-positron pair production in Au-Au and U-U collisions

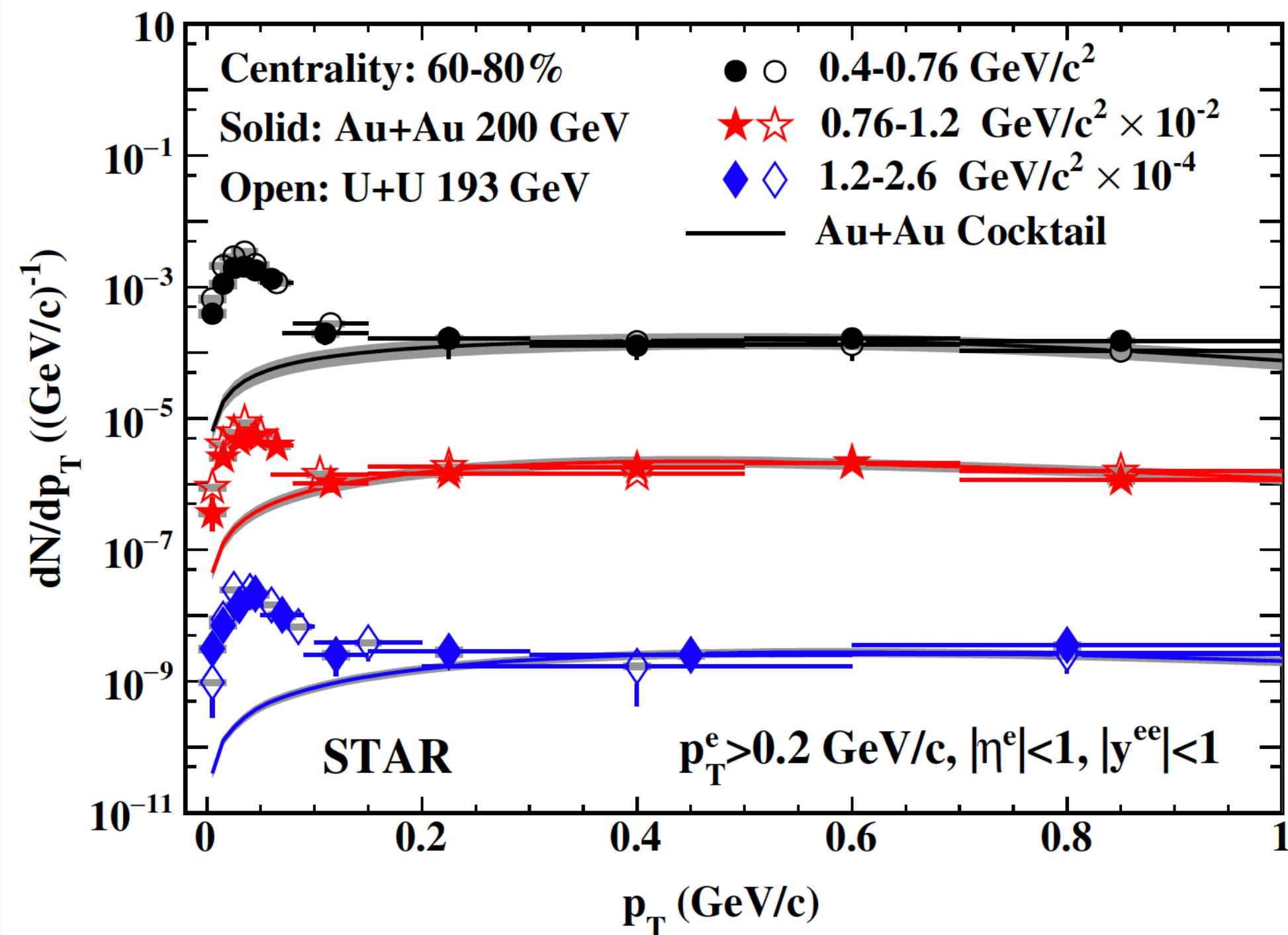
STAR, PRL 121 (2018) 132301



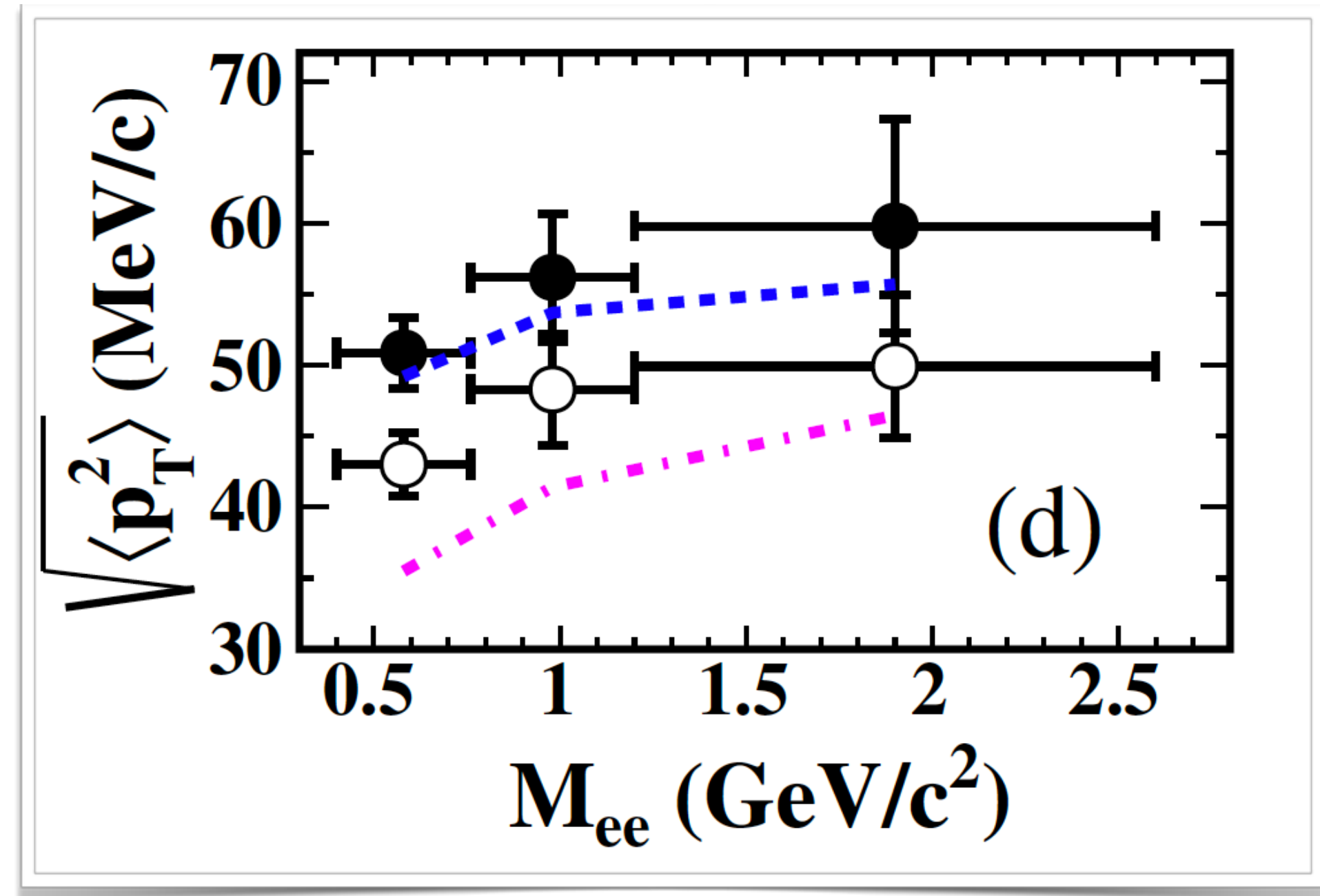
Different colors: different masses.
Line: expectations from hadronic collisions.
Clear excess at low transverse momentum.

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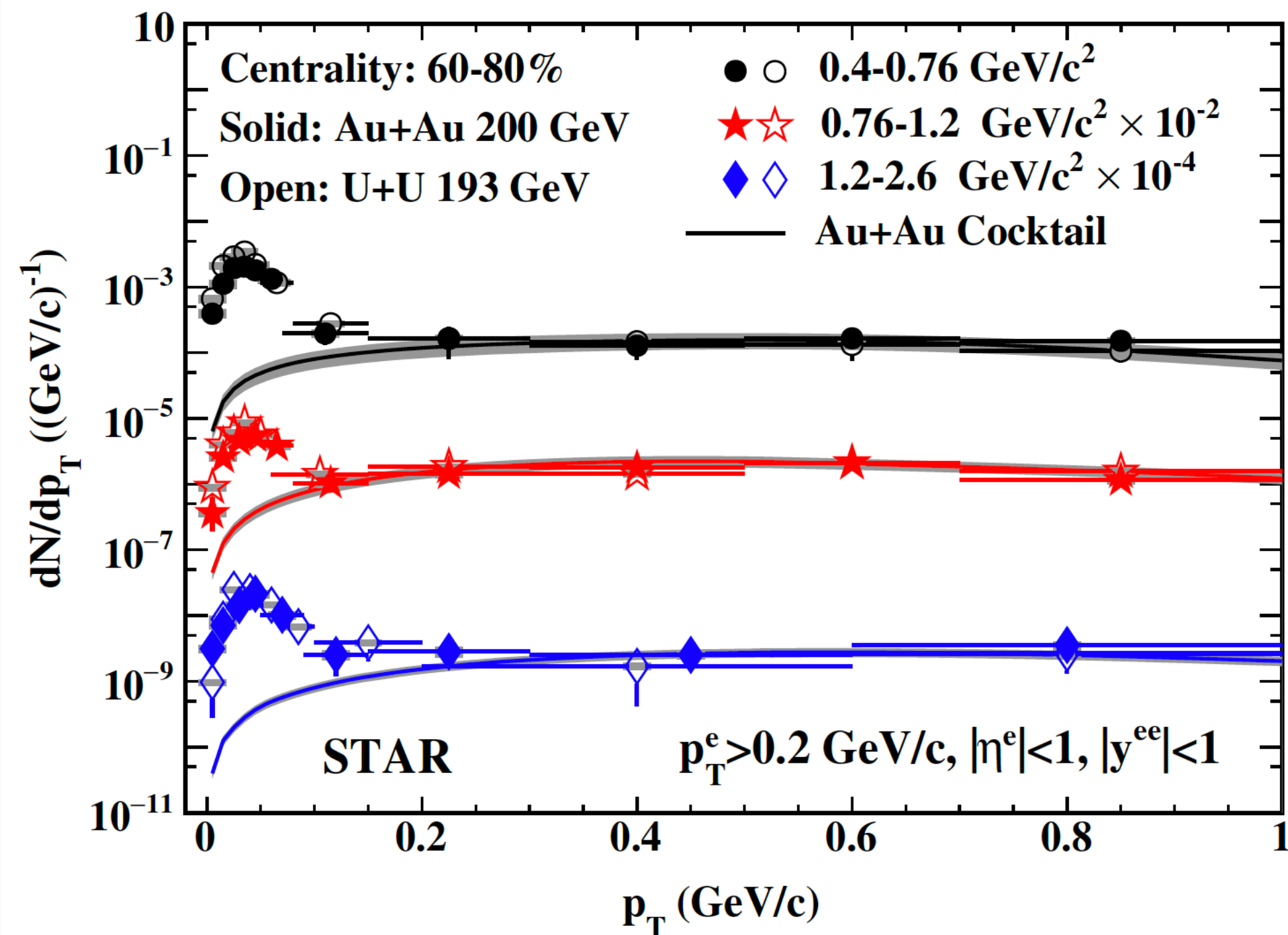


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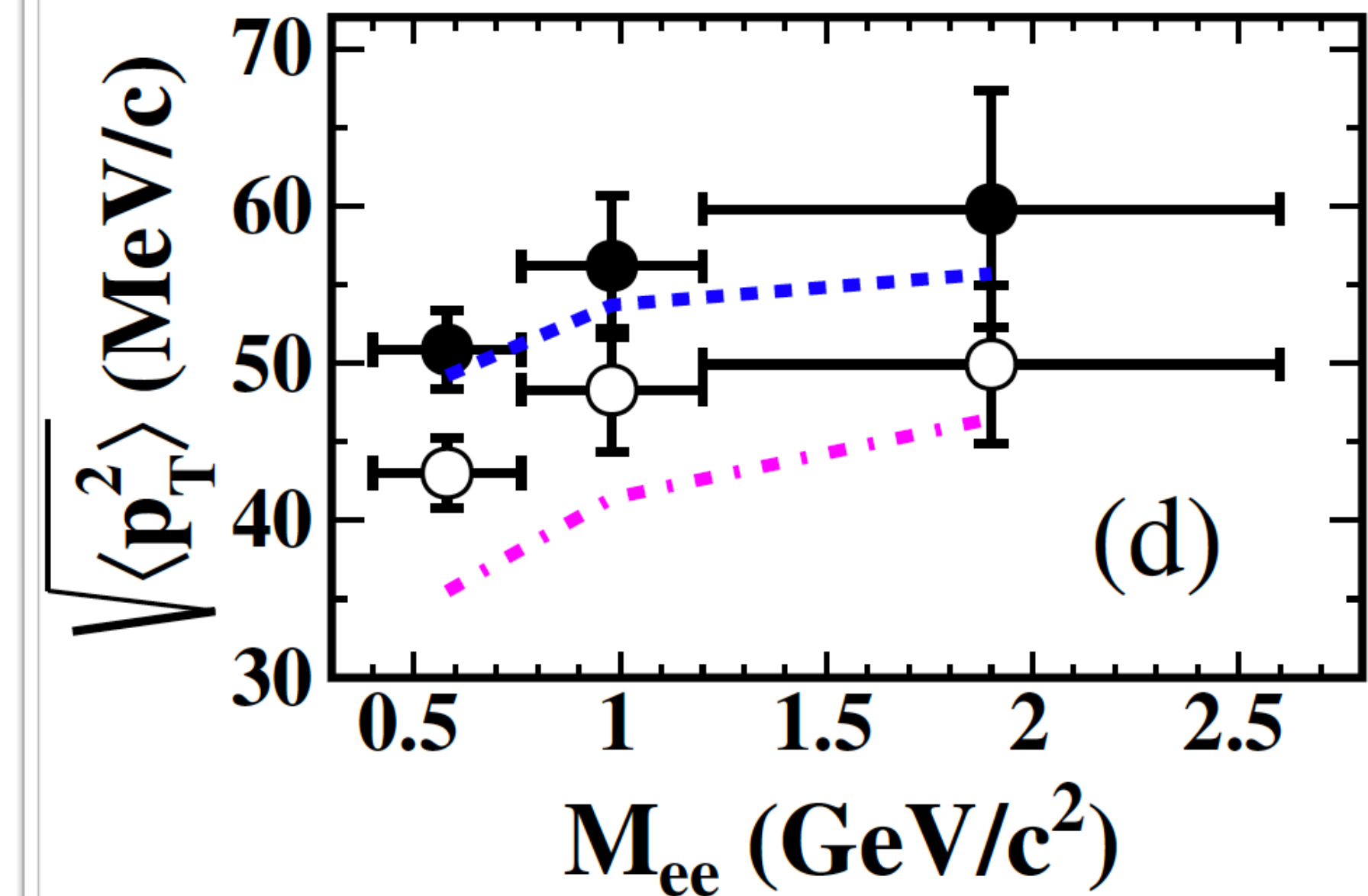


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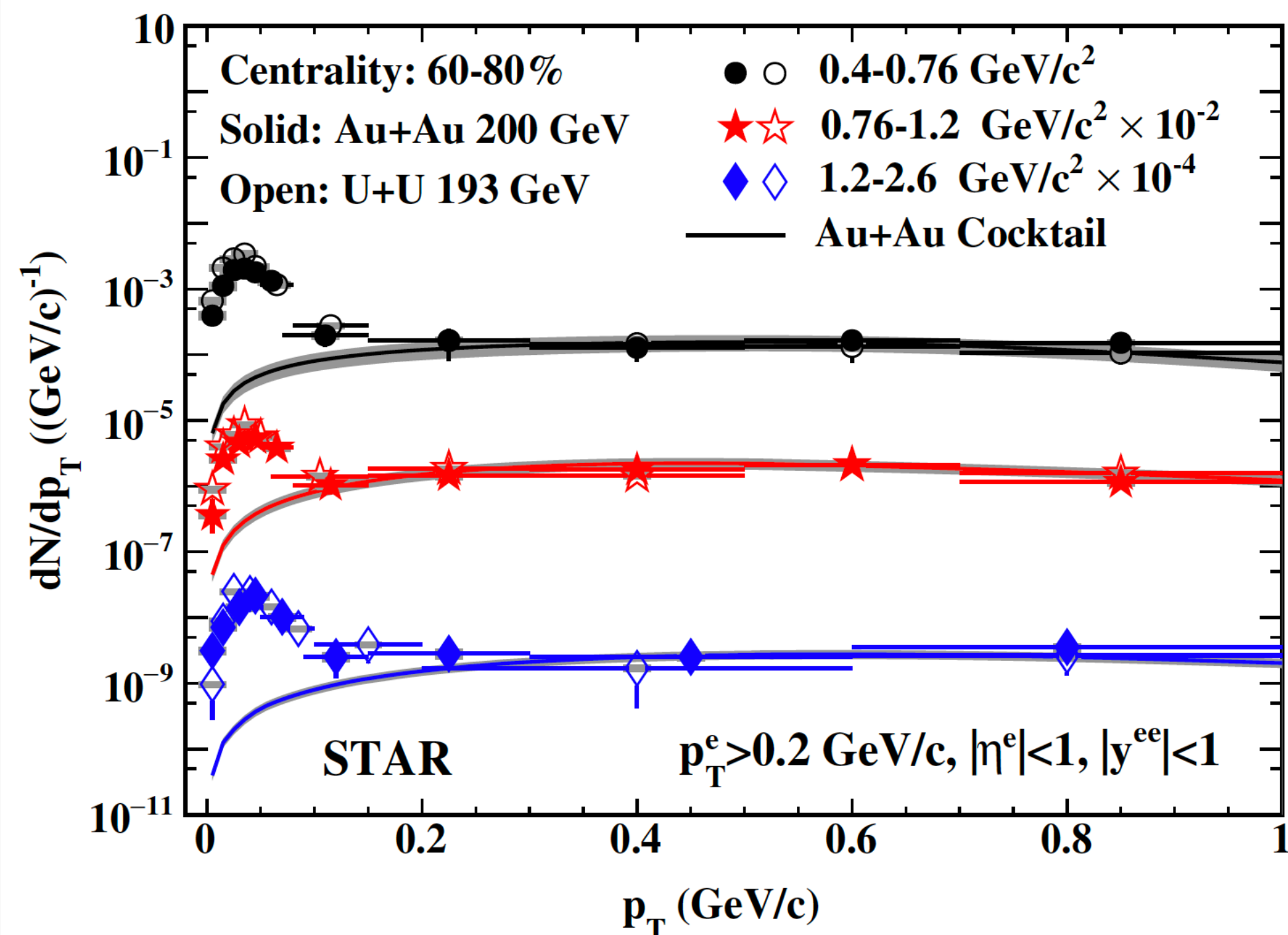
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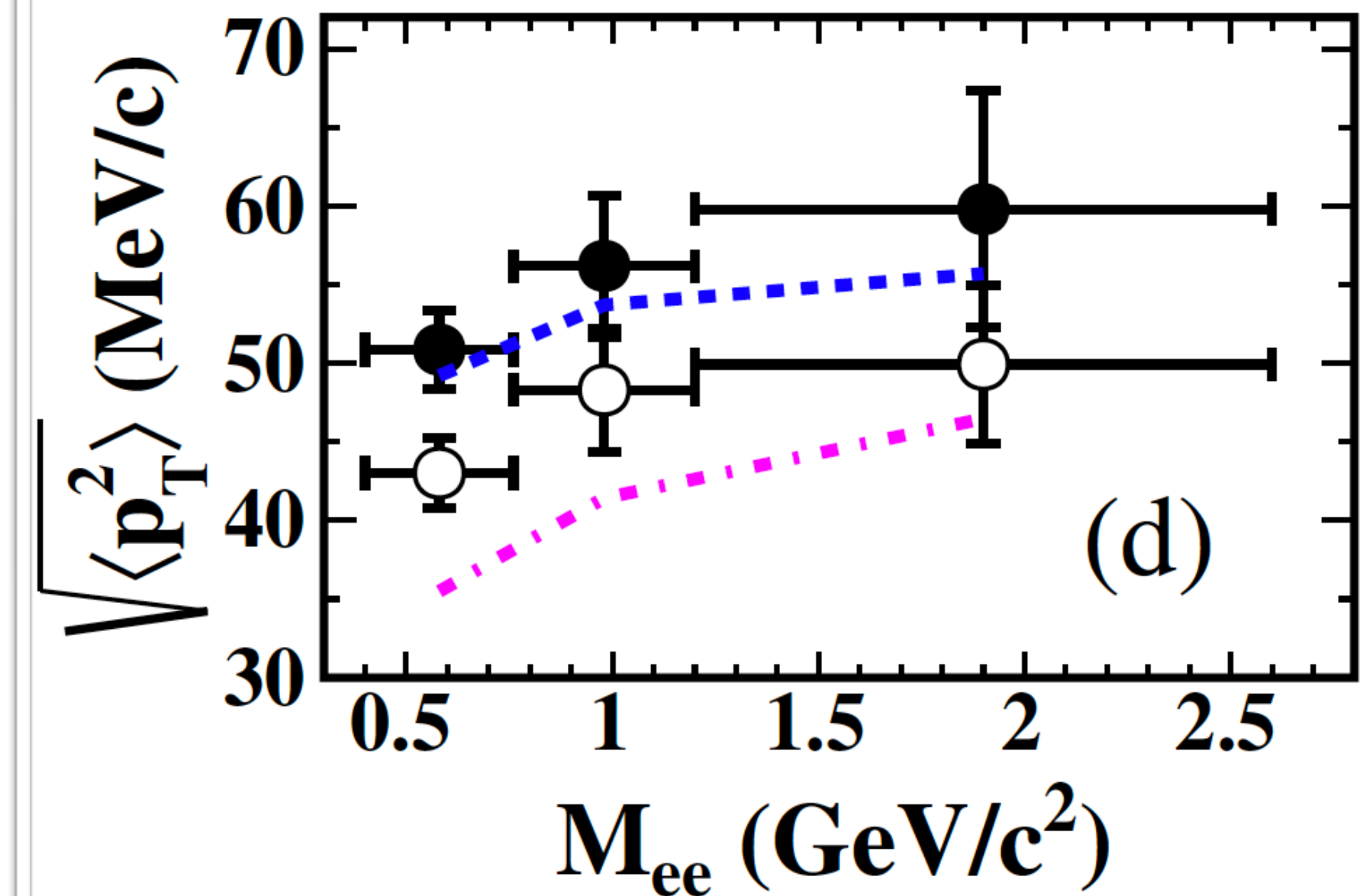
Magenta: 'standard' $\gamma\gamma$ contributions.
Blue: Magnetic field of 10^{14} T !
(Zha et al, PLB 781 (2018) 182)

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STAR, PRL 121 (2018) 132301



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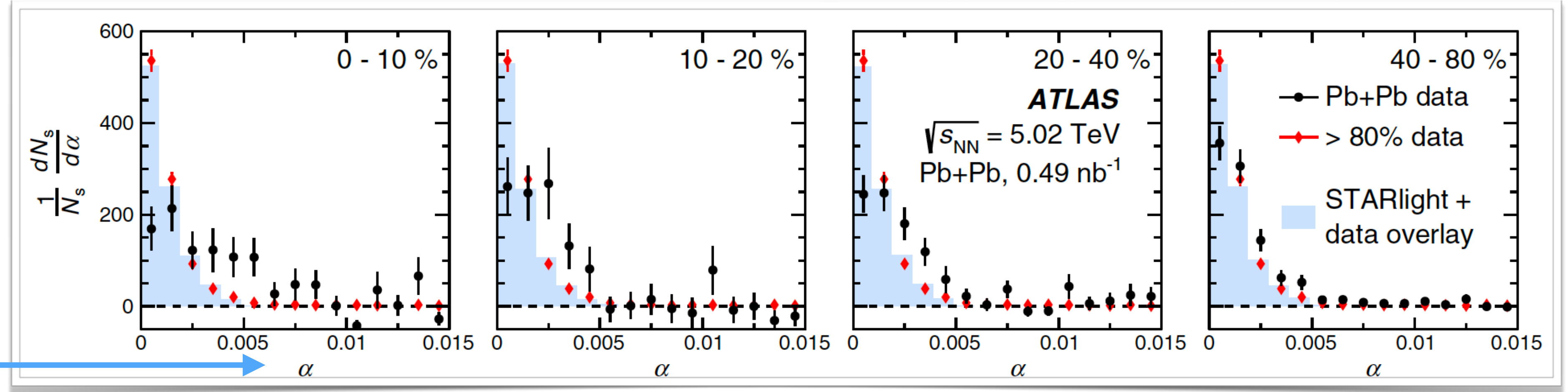
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(Zha et al, PLB 781 (2018) 182)

But other explanations are available:
Kłusek-Gawenda et al, PLB790 (2019) 339.
Zha et al, arXiv 1812.02820.

News from the LHC: muon-pair production

ATLAS, PRL 121 (2018) 212301

$$\alpha \equiv 1 - \frac{|\phi^+ - \phi^-|}{\pi}$$

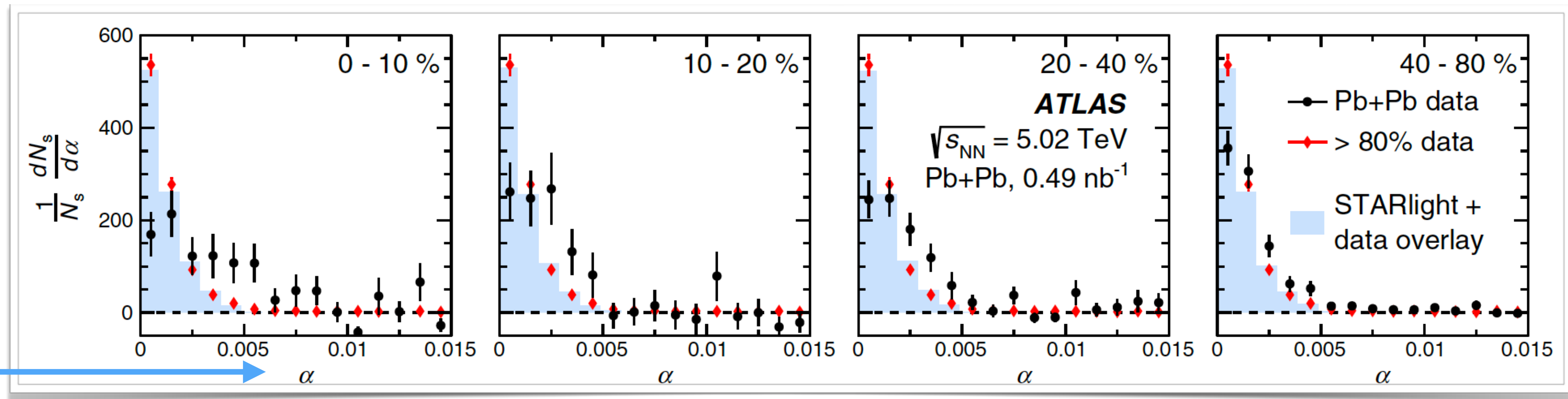


Peripheral collisions consistent with $\gamma\gamma \rightarrow \mu\mu$.
Shape changes with centrality.

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ATLAS, PRL 121 (2018) 212301

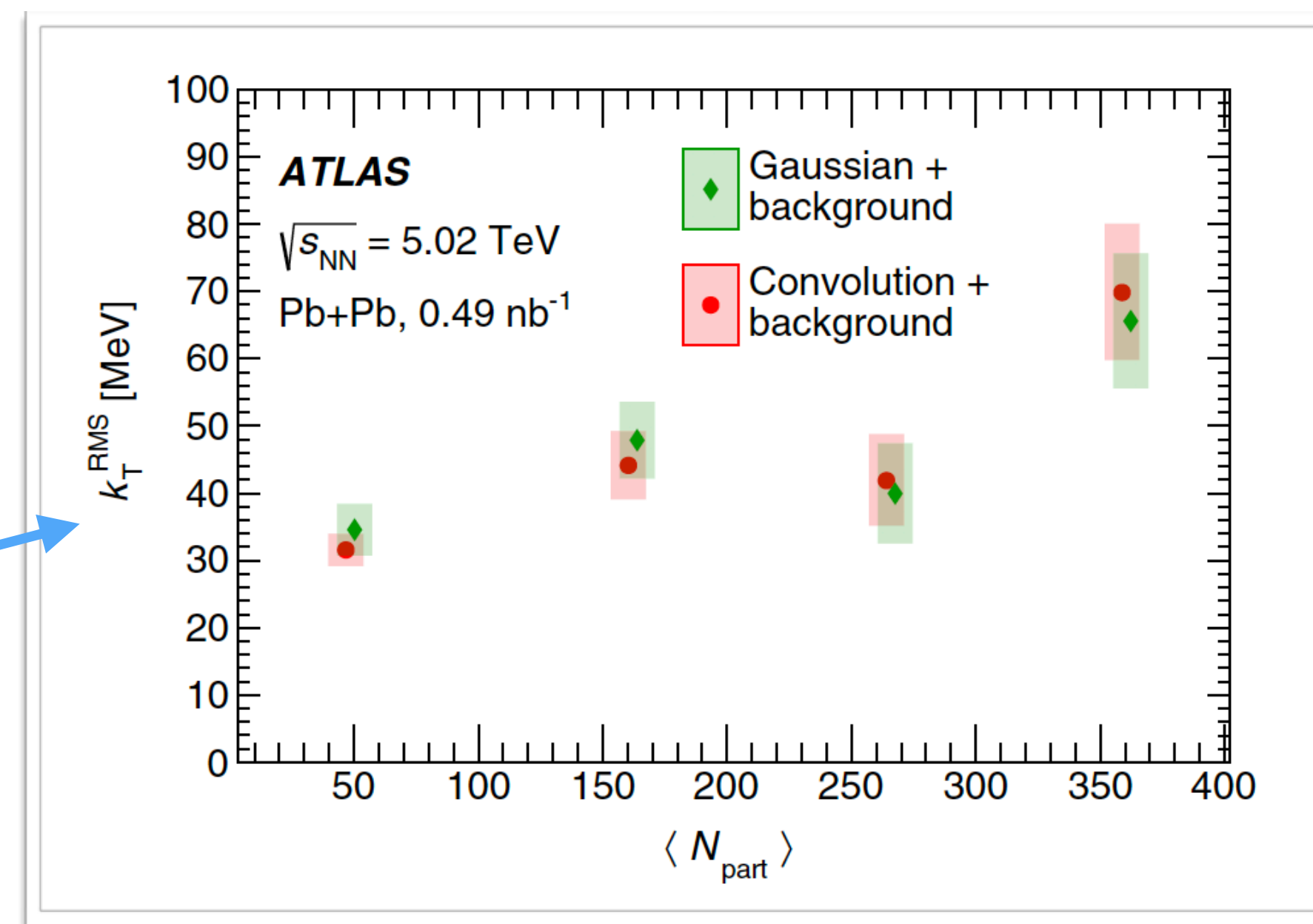
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Shape changes with centrality.

Assuming that the broadening of the α distributions results from a physical process that transfers a *small* amount of transverse momentum, $|\vec{k}_T| \ll p_T^\pm$, to each muon then the variance of the α distribution can be approximated as

$$\langle \alpha^2 \rangle = \langle \alpha^2 \rangle_0 + \frac{1}{\pi^2} \frac{\langle \vec{k}_T^2 \rangle}{\langle p_{T\text{avg}}^2 \rangle}, \quad (1)$$

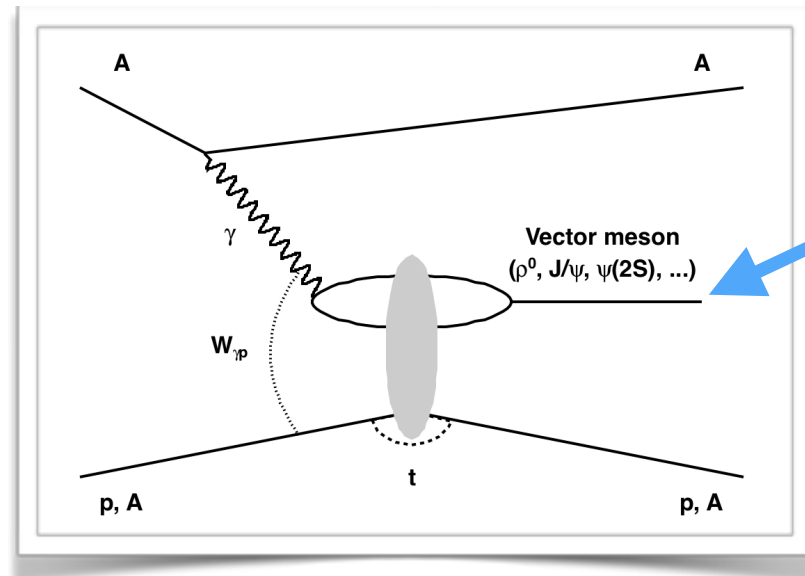


Increase in k_T due to re-scattering in QGP?

The structure of protons at small x

News about wave functions

$$A(x, Q^2, \vec{\Delta})_{T,L} = i \int d\vec{r} \int_0^1 \frac{dz}{4\pi} (\Psi^* \Psi_V)_{T,L} \int d\vec{b} e^{-i(\vec{b} - (1-z)\vec{r}) \cdot \vec{\Delta}} \frac{d\sigma_{\text{dip}}}{d\vec{b}}$$

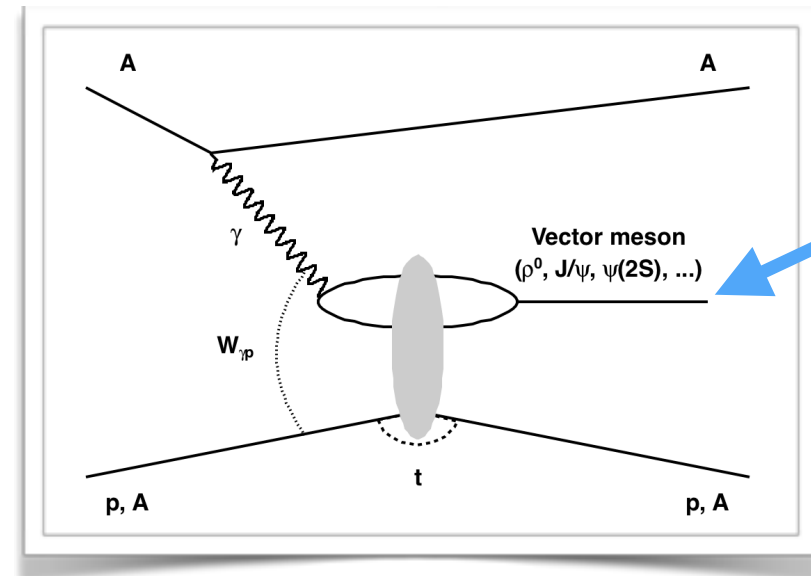


Vector meson
wave function

News about wave functions

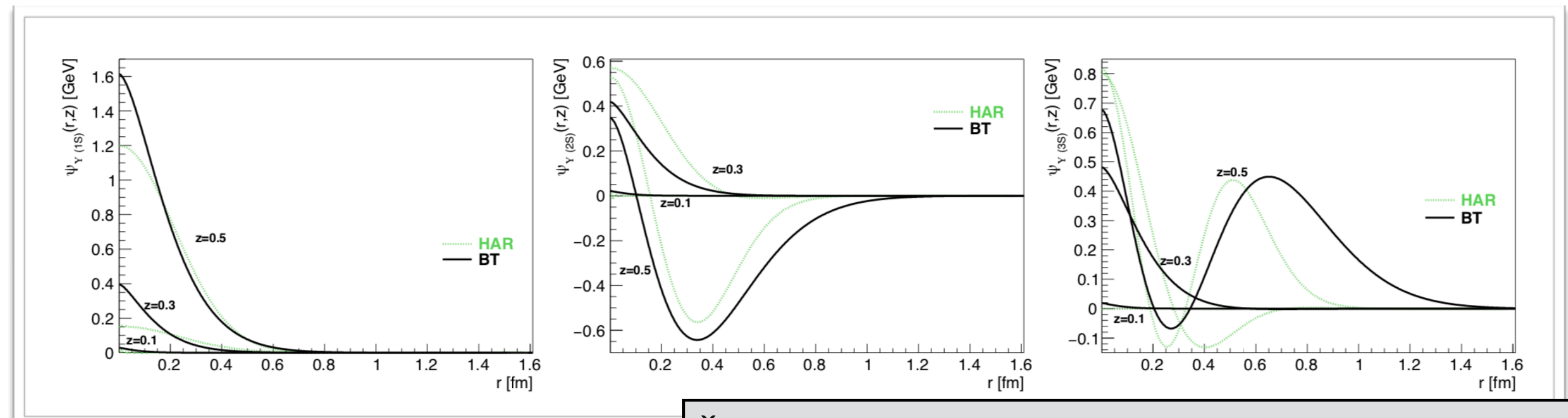
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Wave functions with a spacial and a spin part



Vector meson
wave function

Spacial part from solution of Schrödinger equation with different potentials



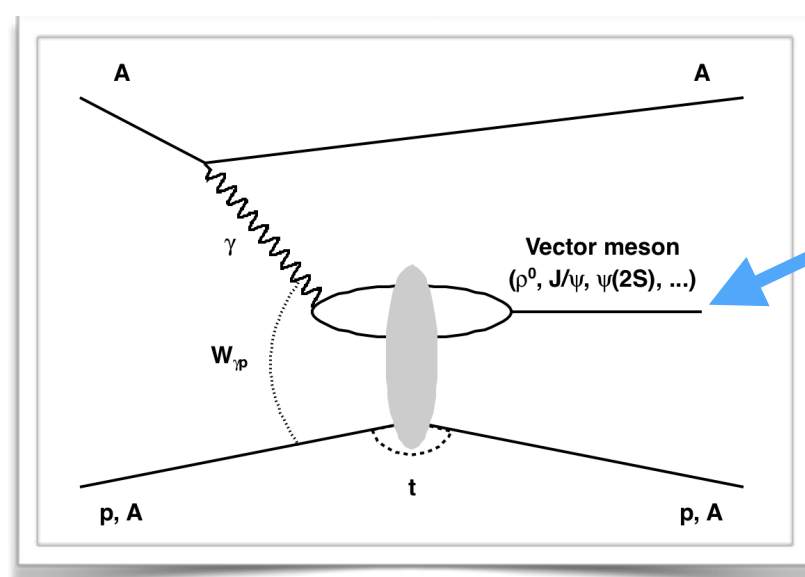
Čepila, Křelina, Němčík, Pasechnik, arXiv 1901.02664

News about wave functions

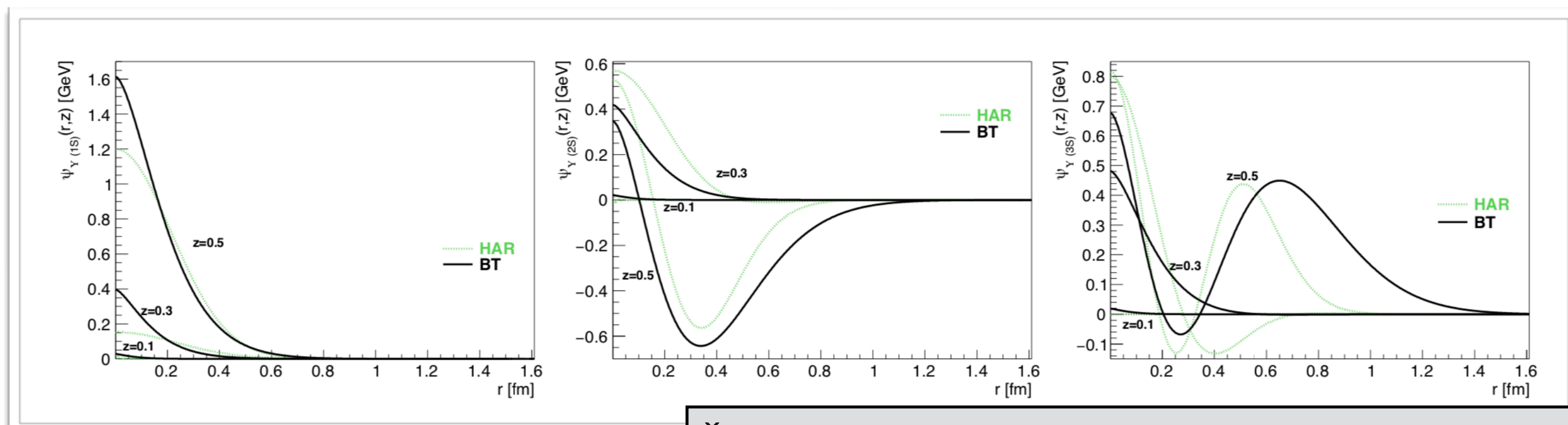
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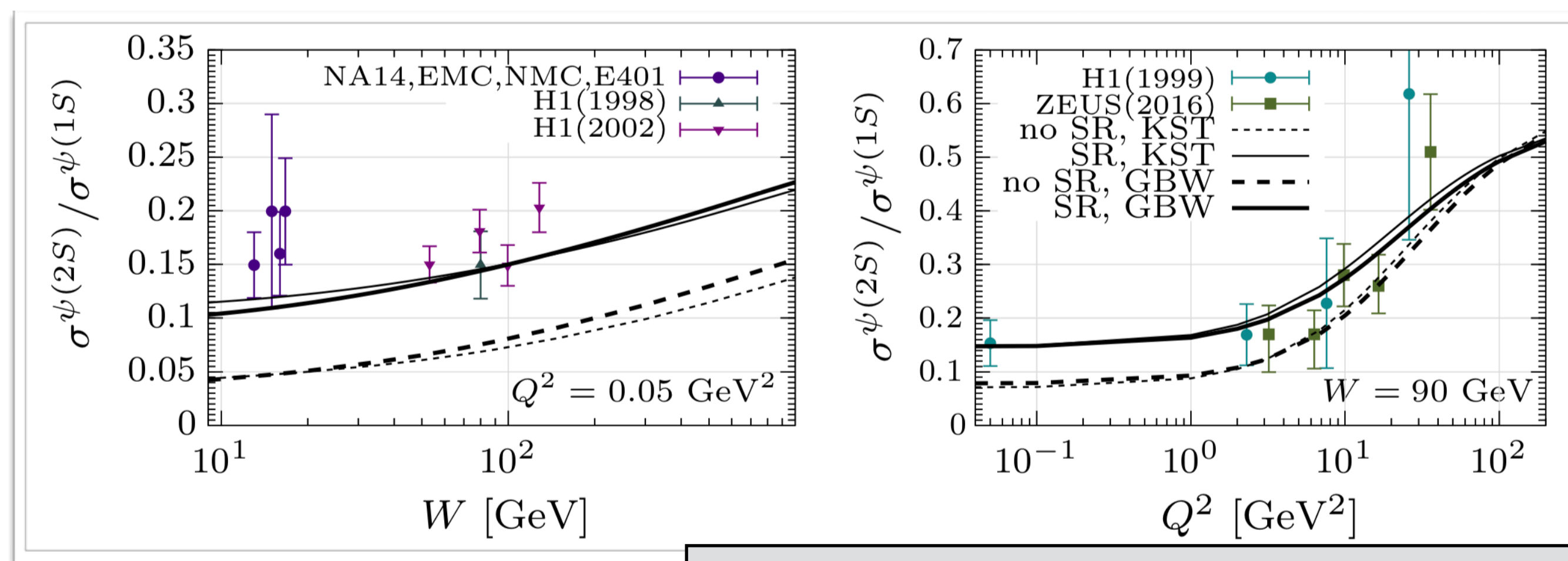


Vector meson
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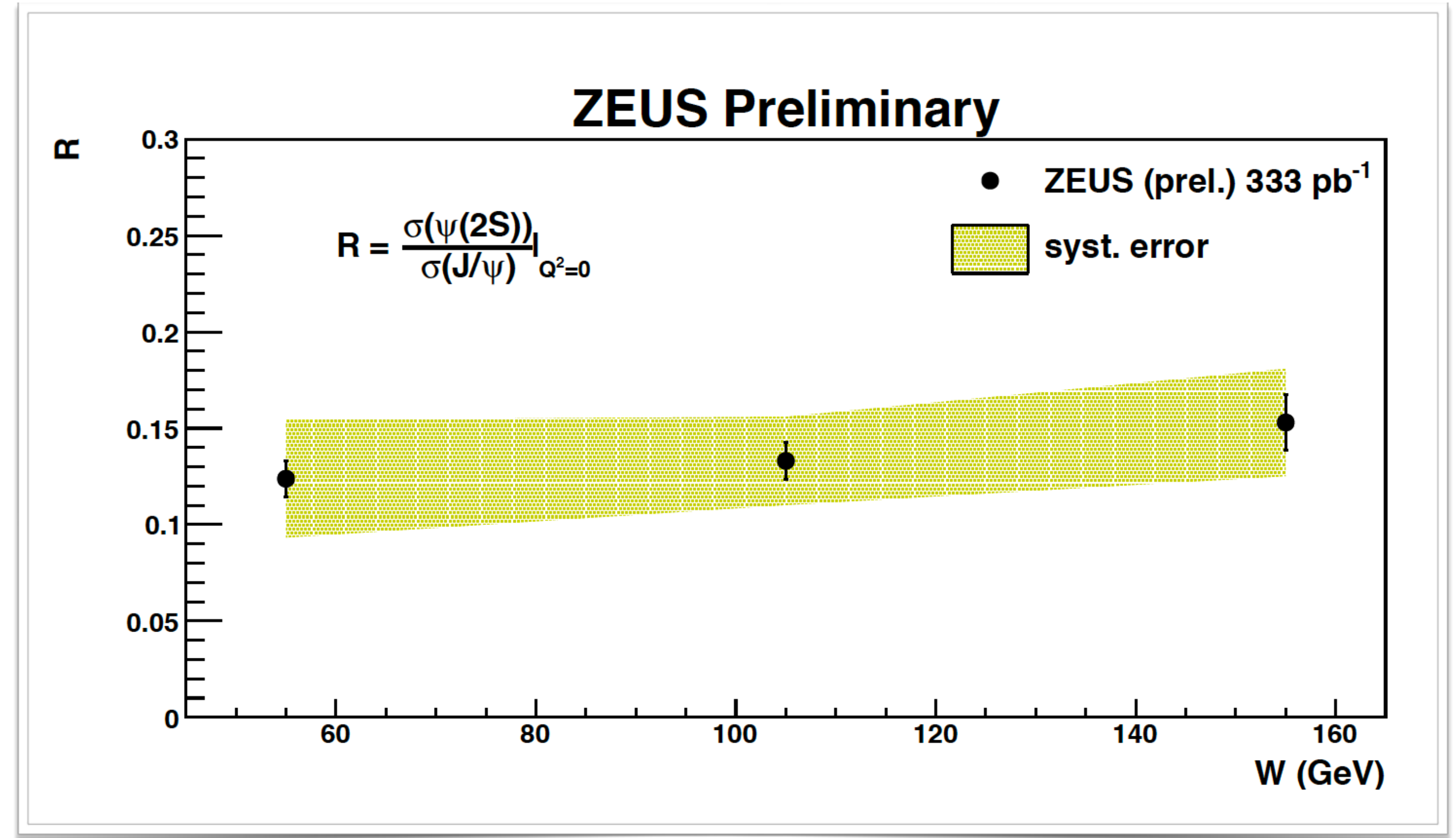
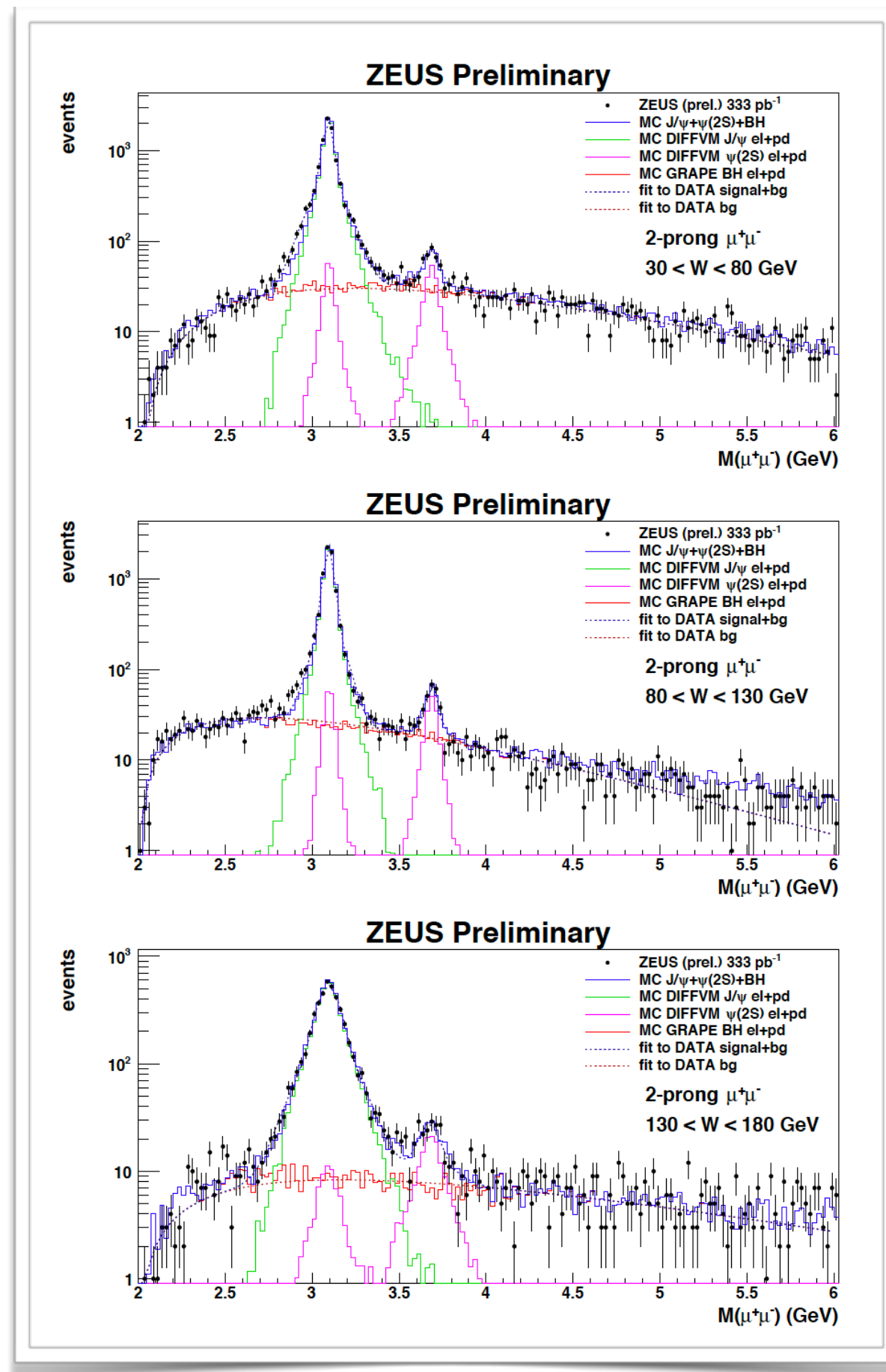
Čepila, Křelina, Němčík, Pasechnik, arXiv 1901.02664

Inclusion of spin rotation effects have a large effect on 2S states due to a nodal structure of the corresponding wave functions.



Čepila, Křelina, Němčík, Pasechnik, EPJC79 (2019) 154

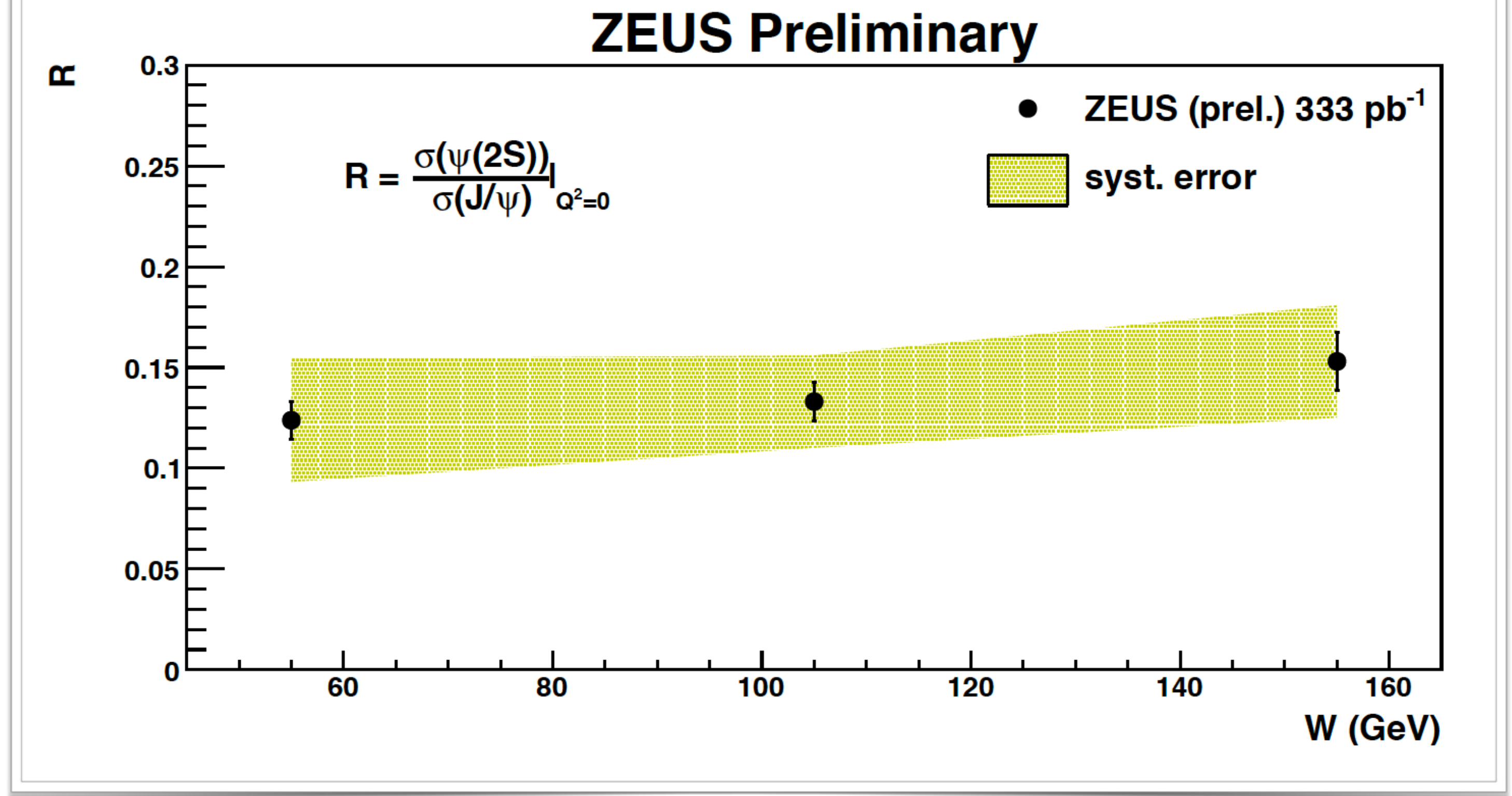
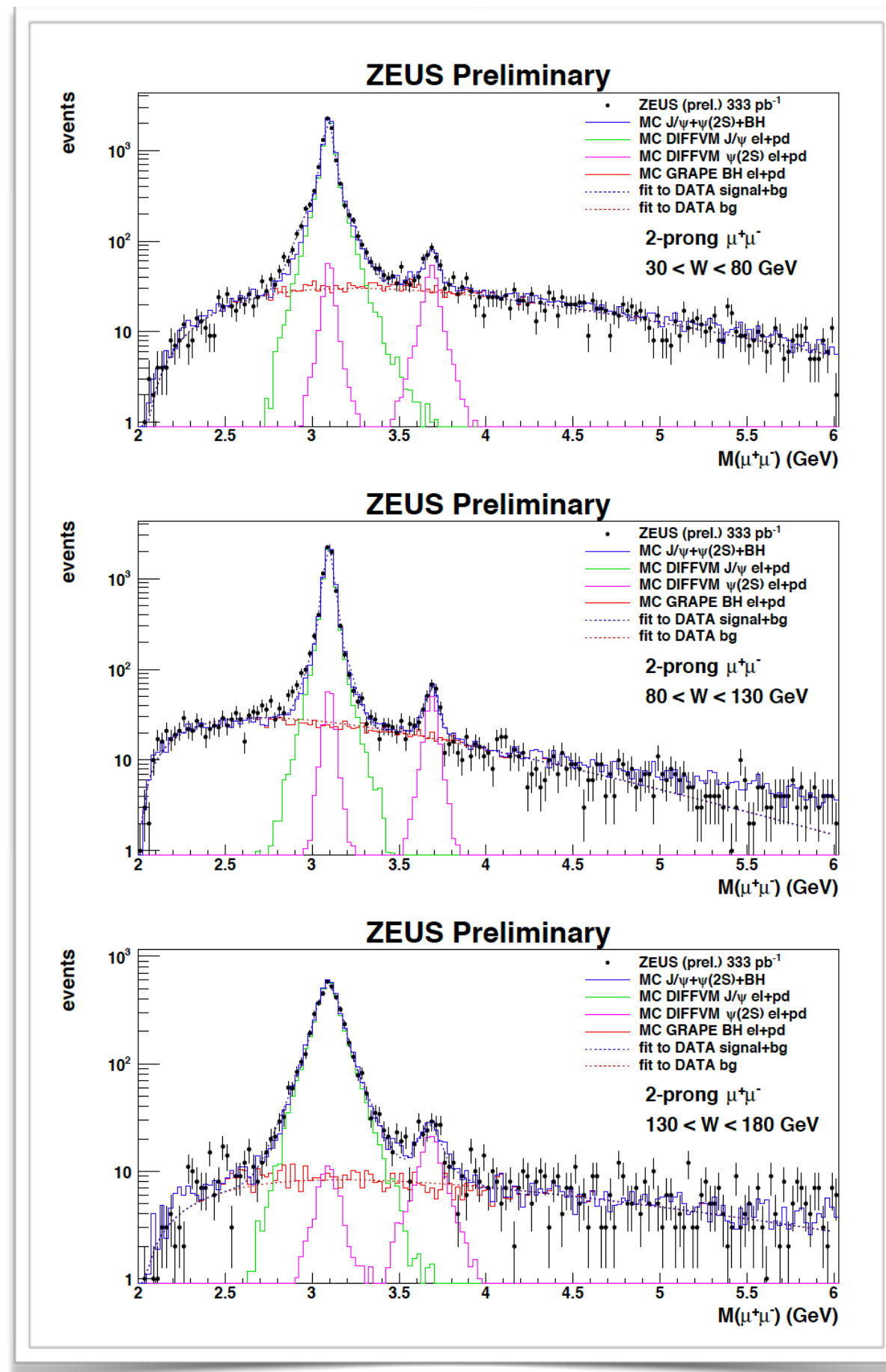
News from HERA: ψ' to J/ψ ratio in photoproduction



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Wednesday in the WG2 session

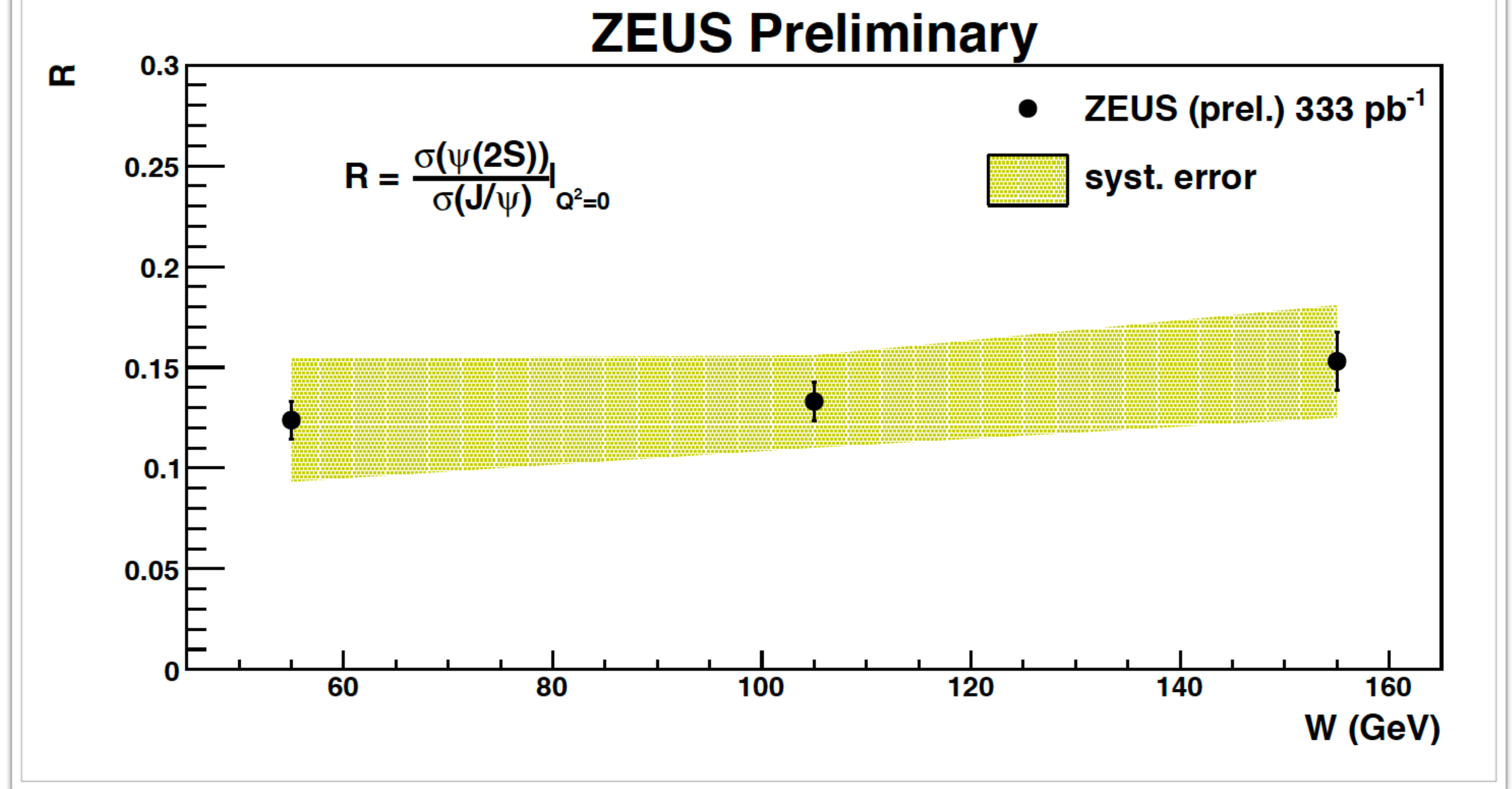
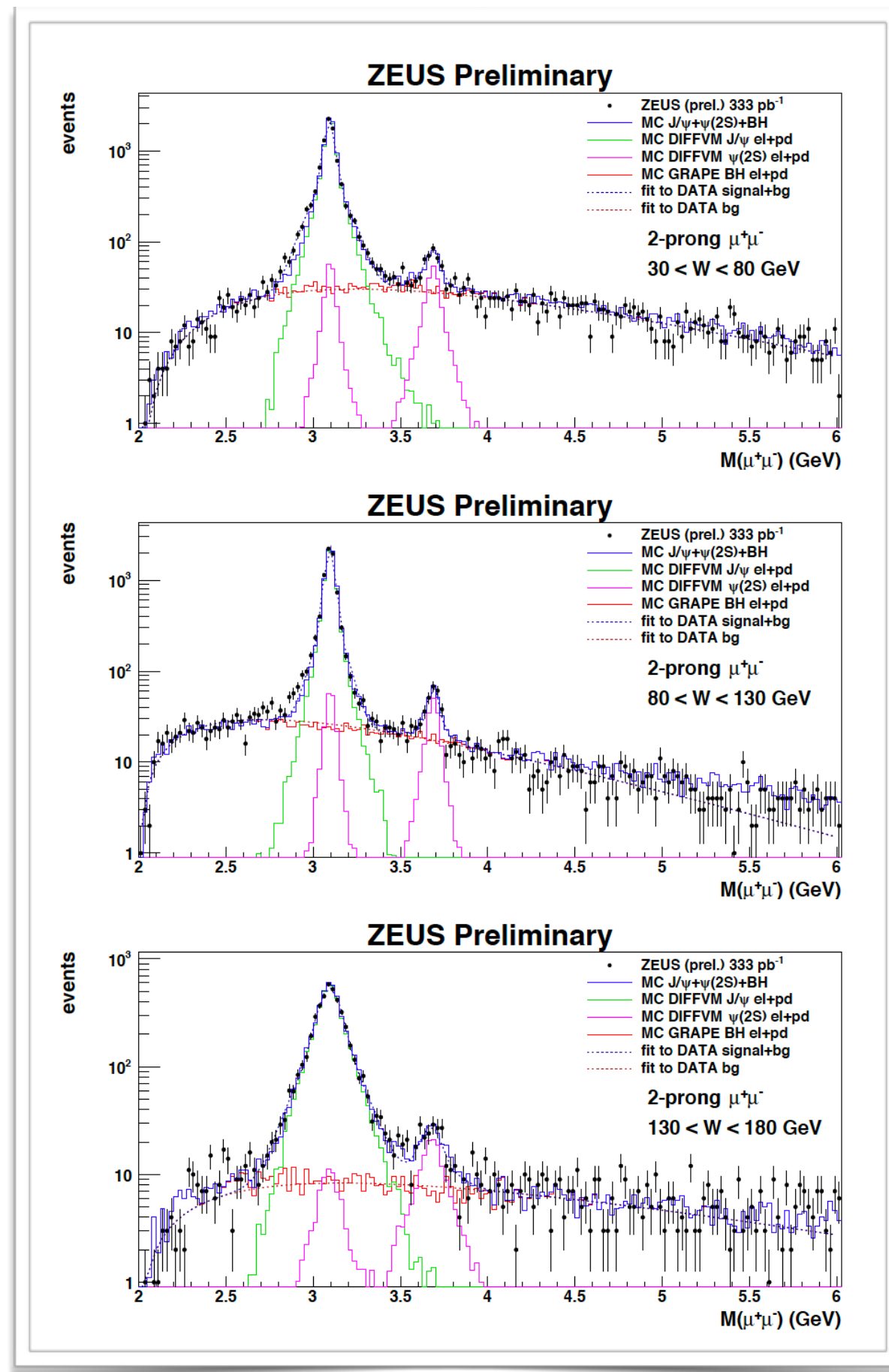
Measurement of the $\Psi(2S)$ to J/Ψ cross section ratio in photoproduction with the ZEUS detector at HERA, Alessia Bruni



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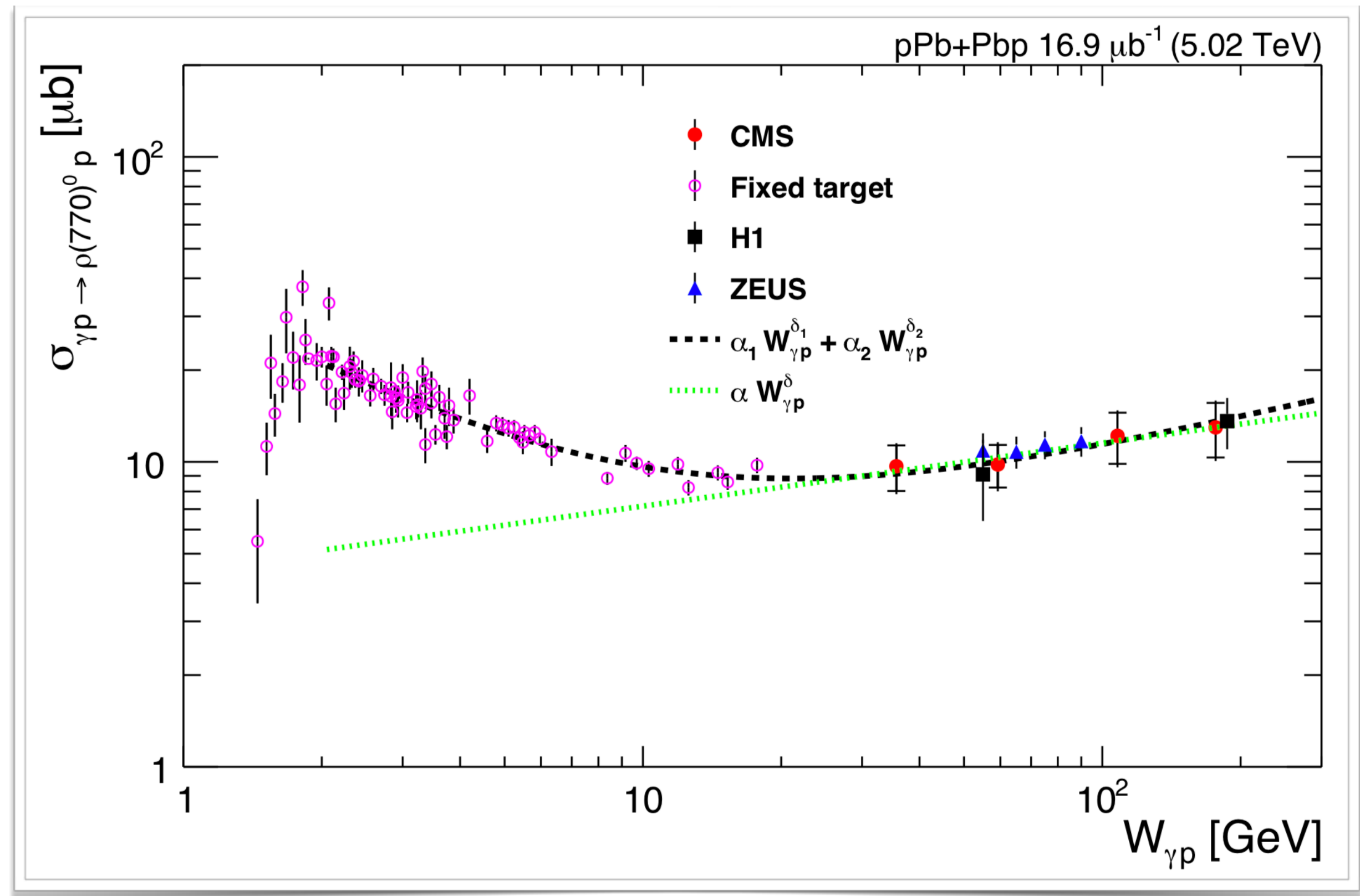
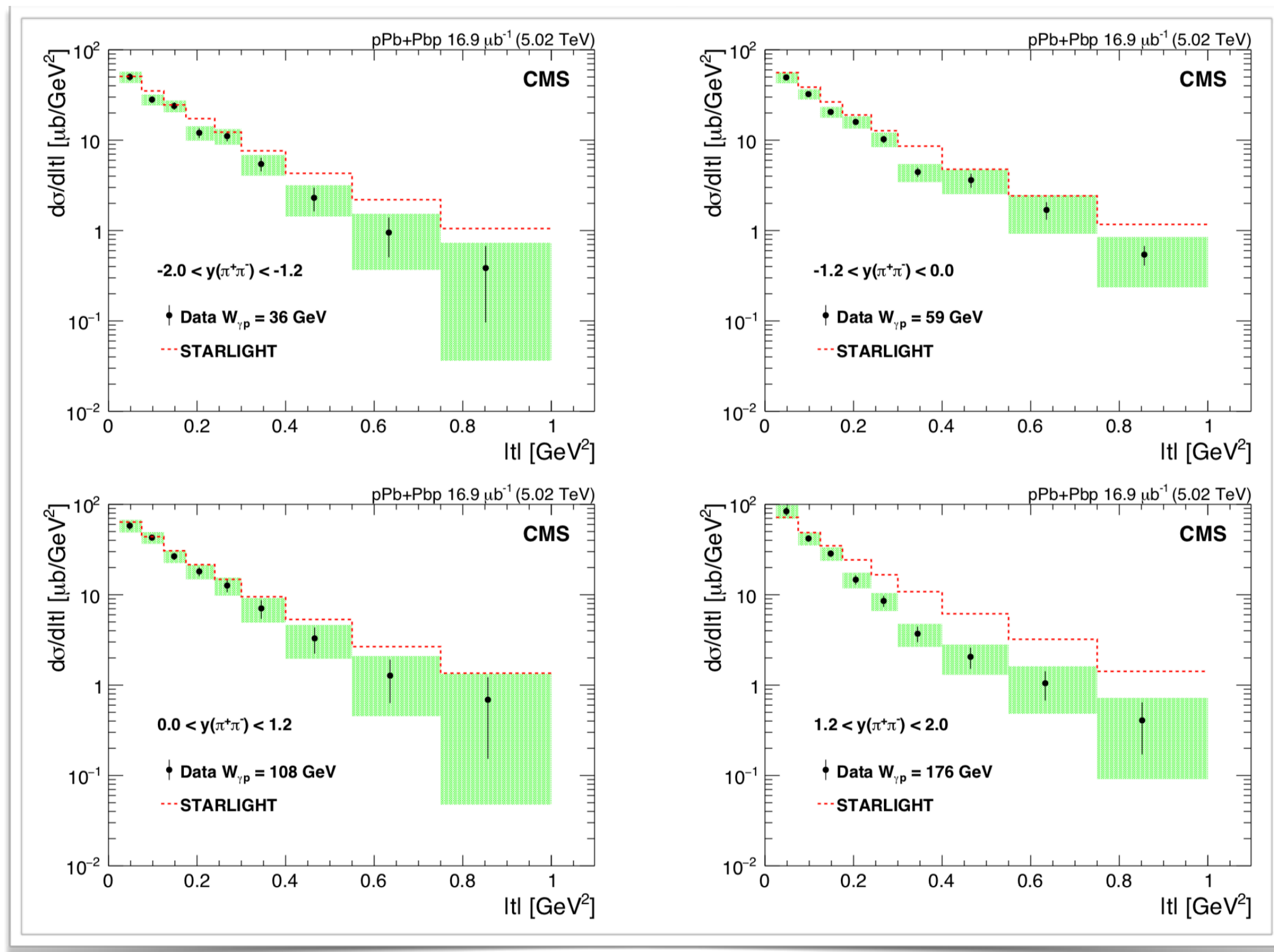


Wednesday in the WG2 session

Soft QCD and Central Exclusive Production at LHCb,
Kucharczyk, Marcin

News from the LHC: ρ photoproduction off protons

CMS arXiv:1902.01339



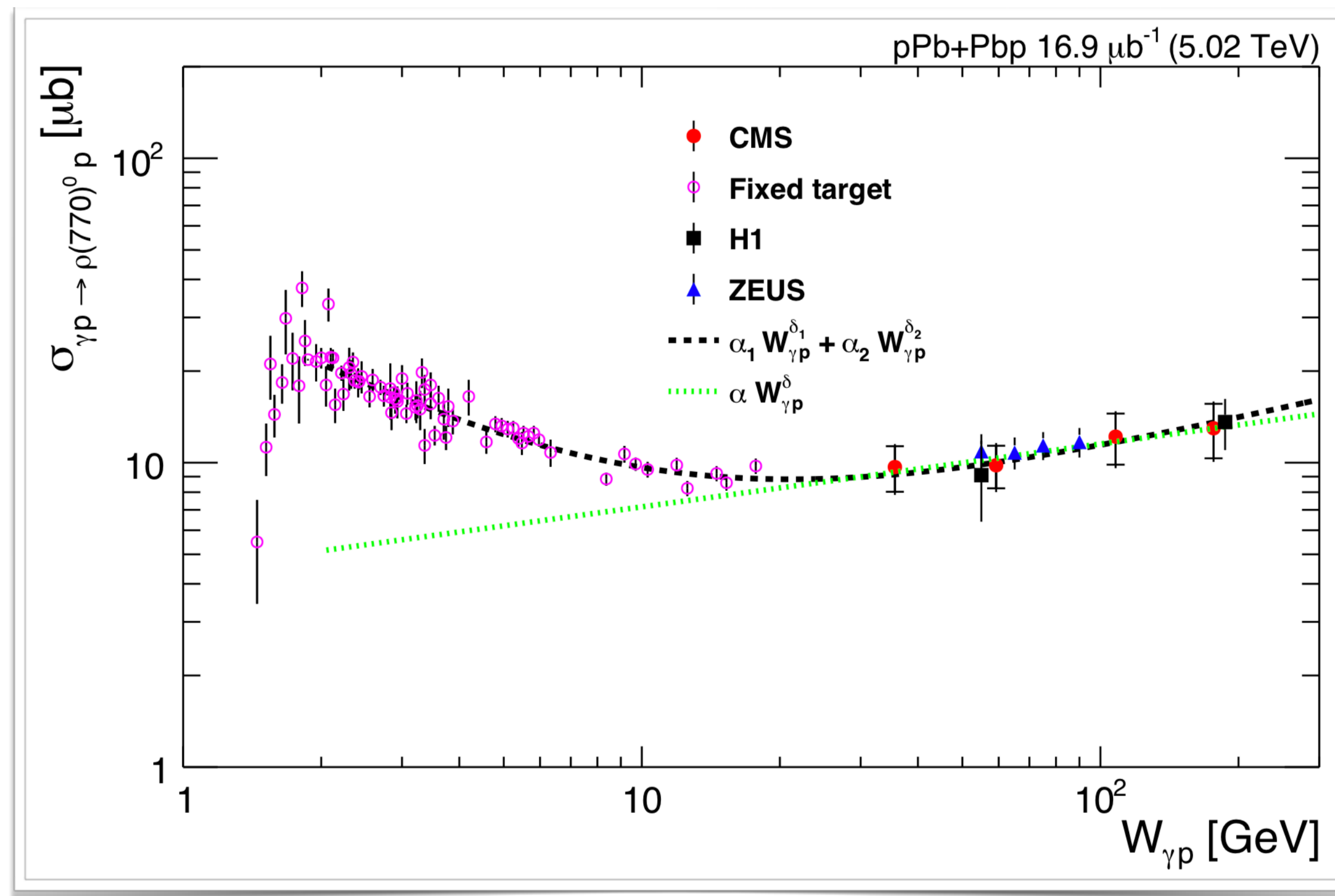
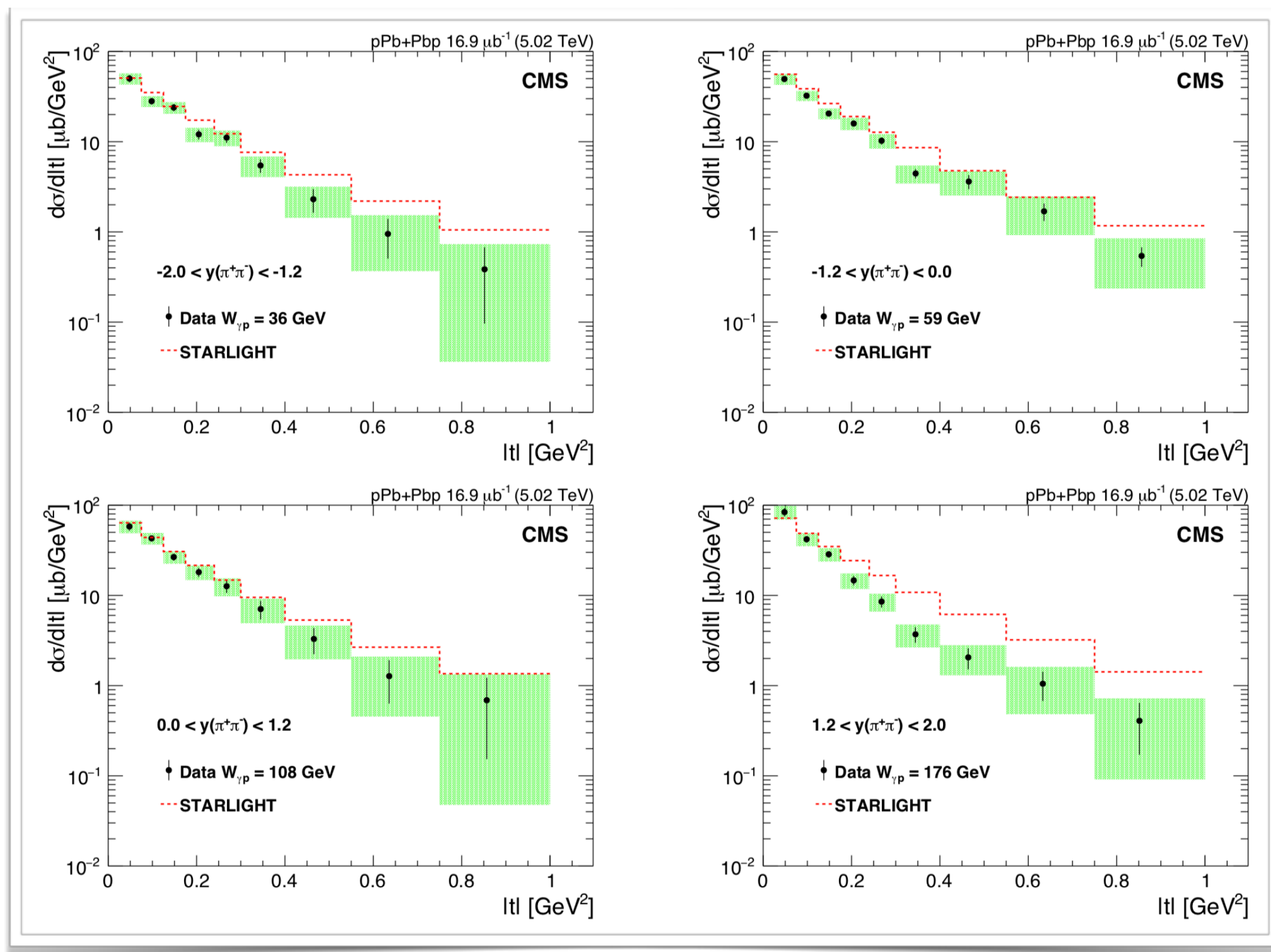
Measurements of the $|t|$ and energy dependence of this process!

News from the LHC: ρ photoproduction off protons

Wednesday in the WG2 session

Recent CMS results on exclusive processes Bylinkin Aleksandr

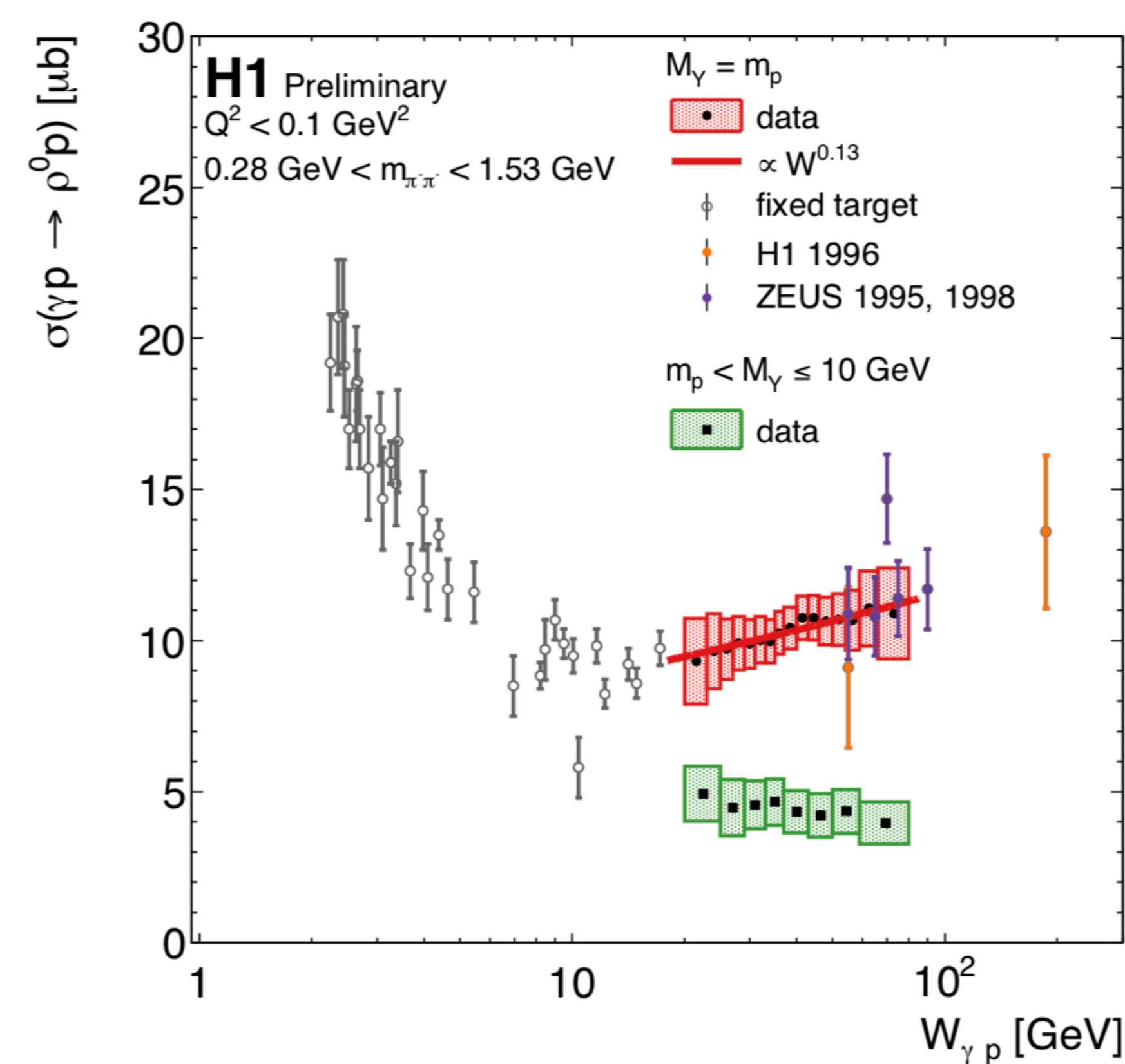
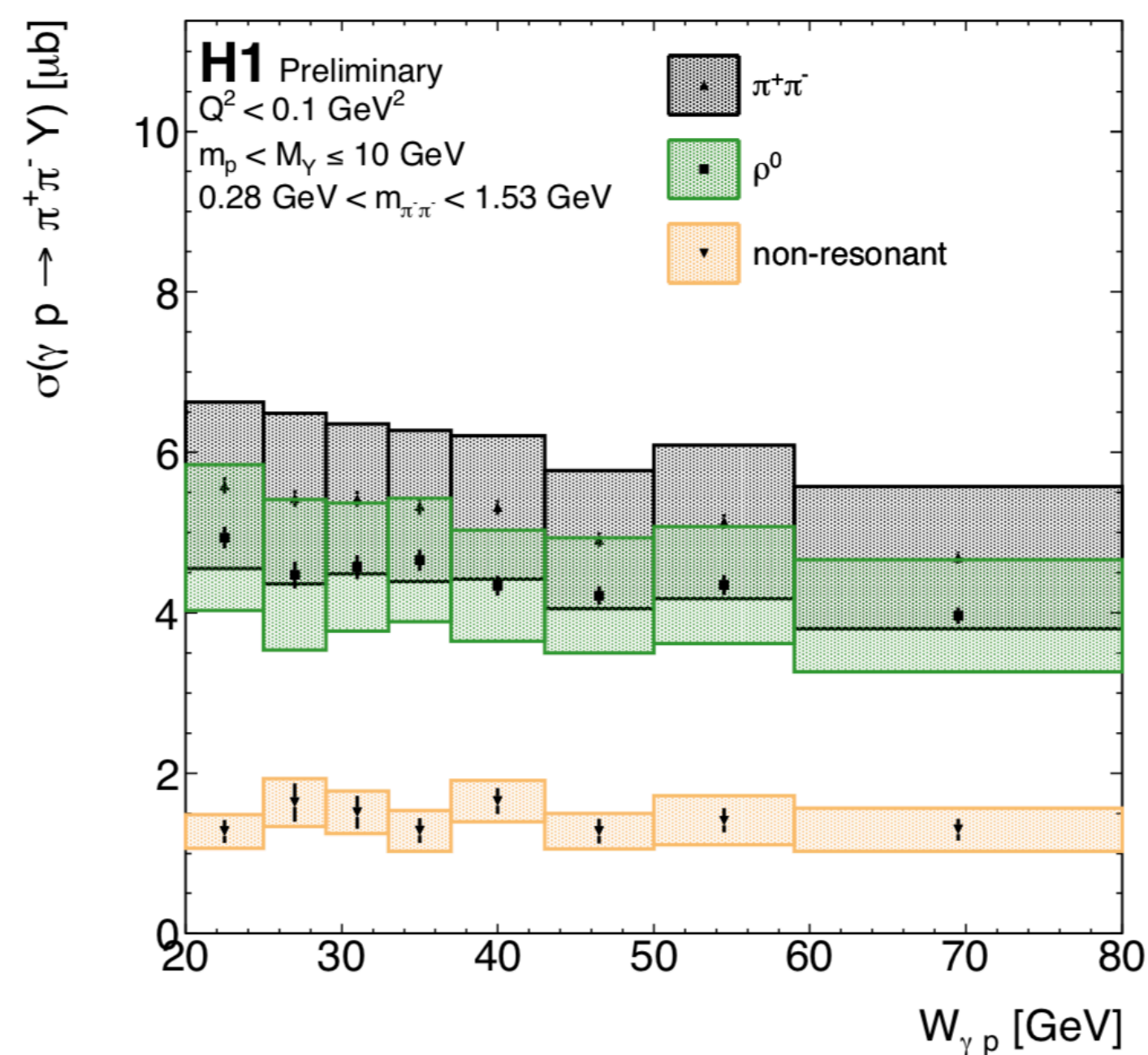
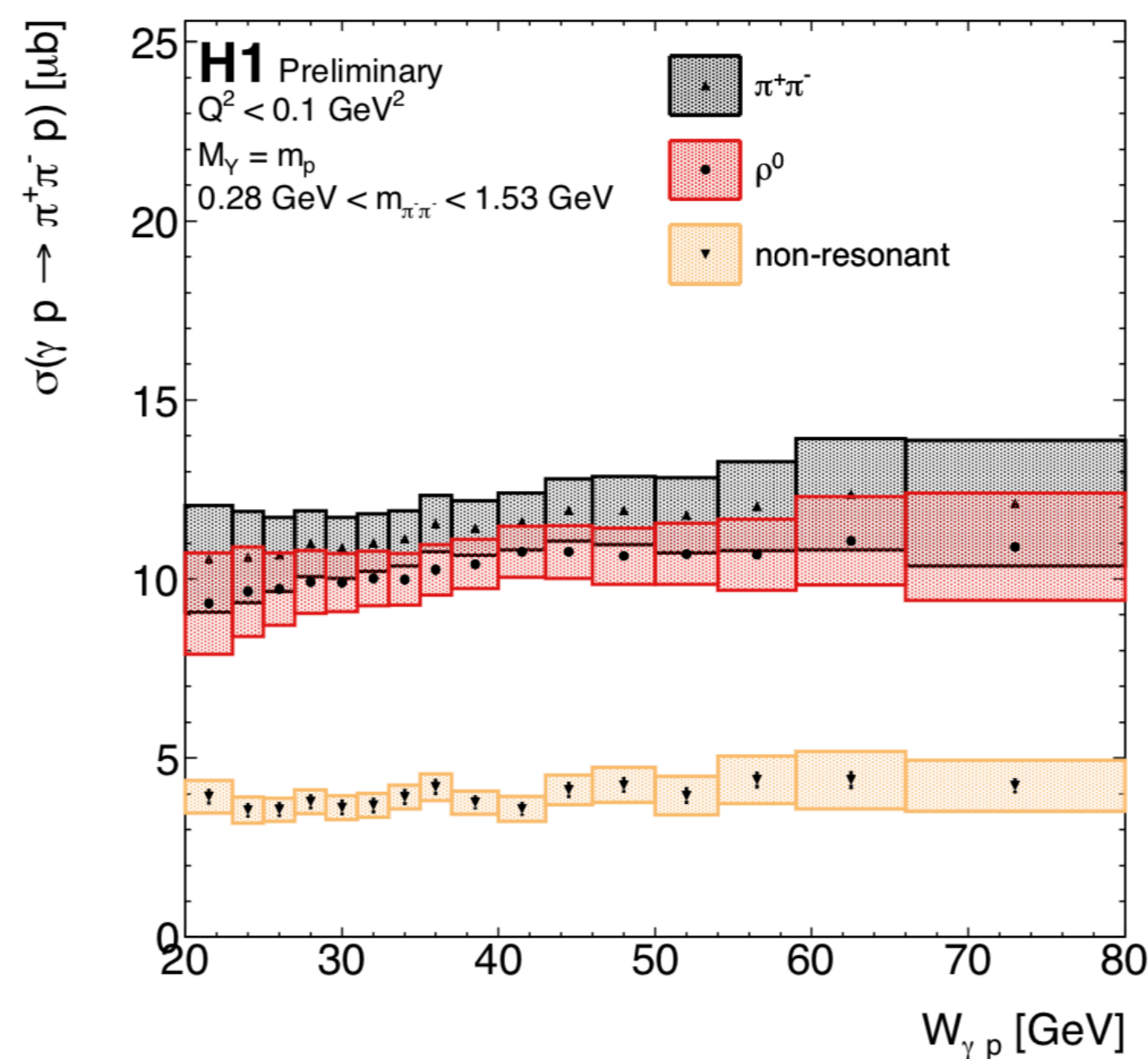
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News from HERA: ρ photoproduction off protons

<http://www-h1.desy.de/publications/htmlsplit/H1prelim-18-012.long.html>



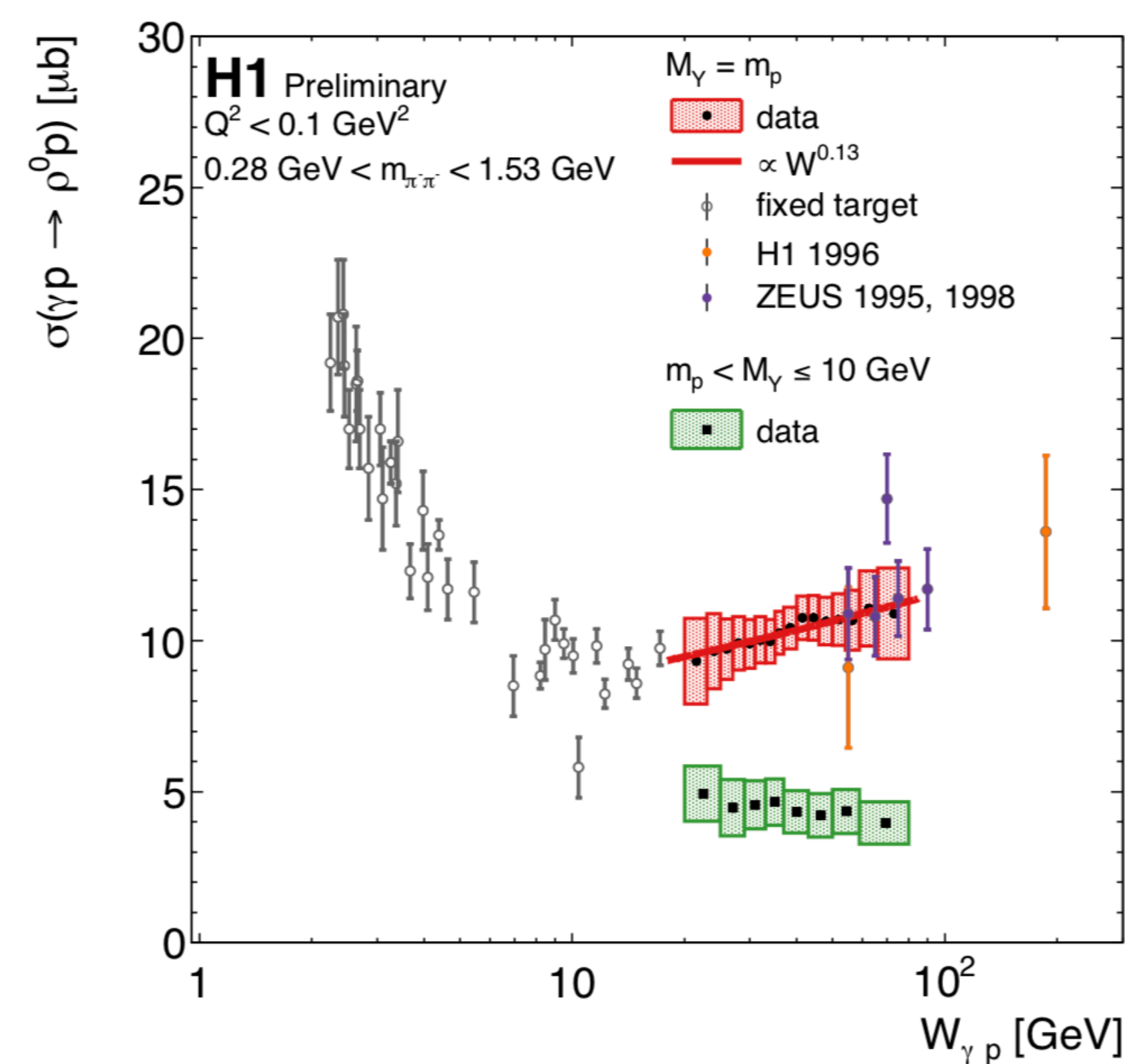
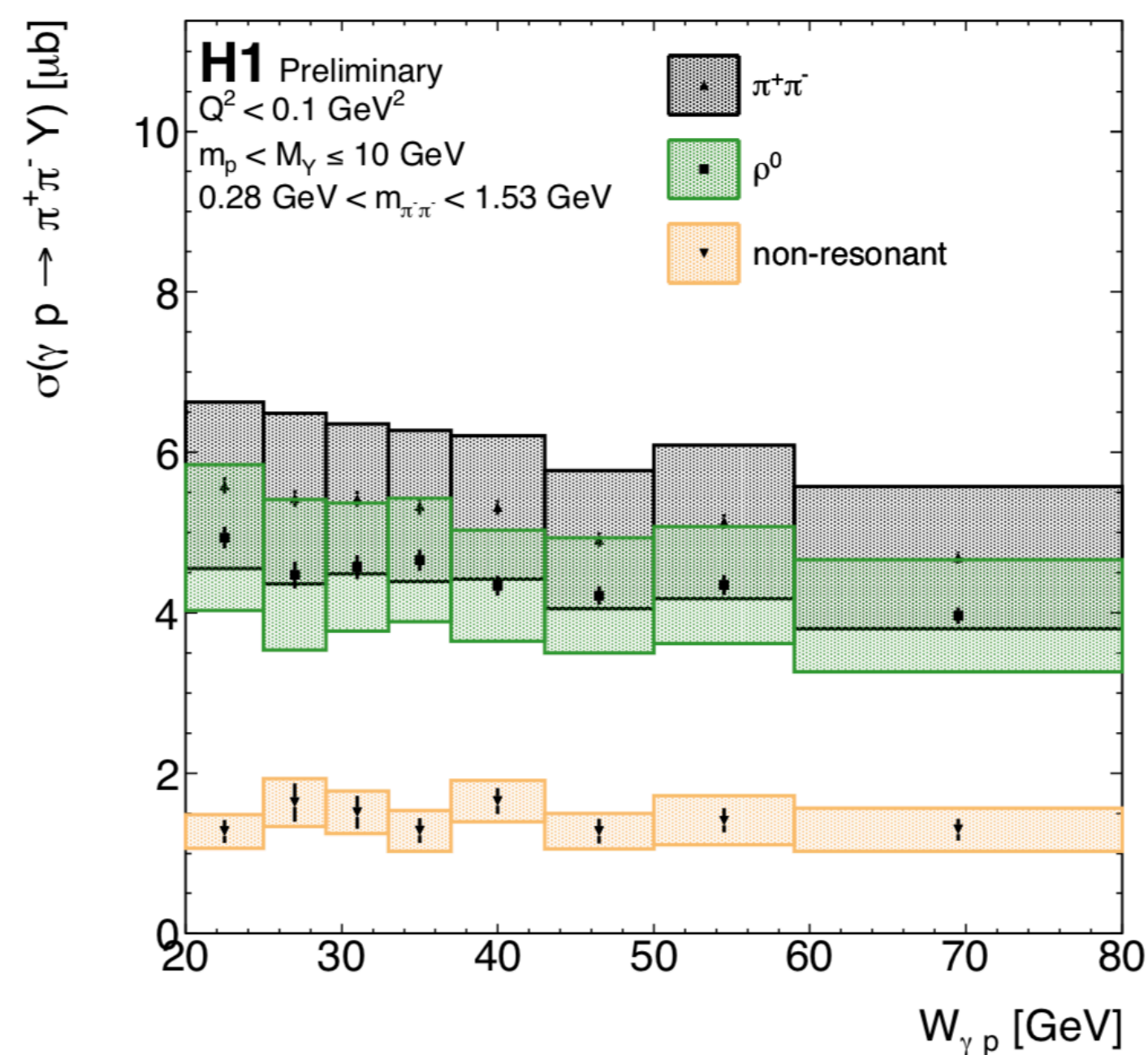
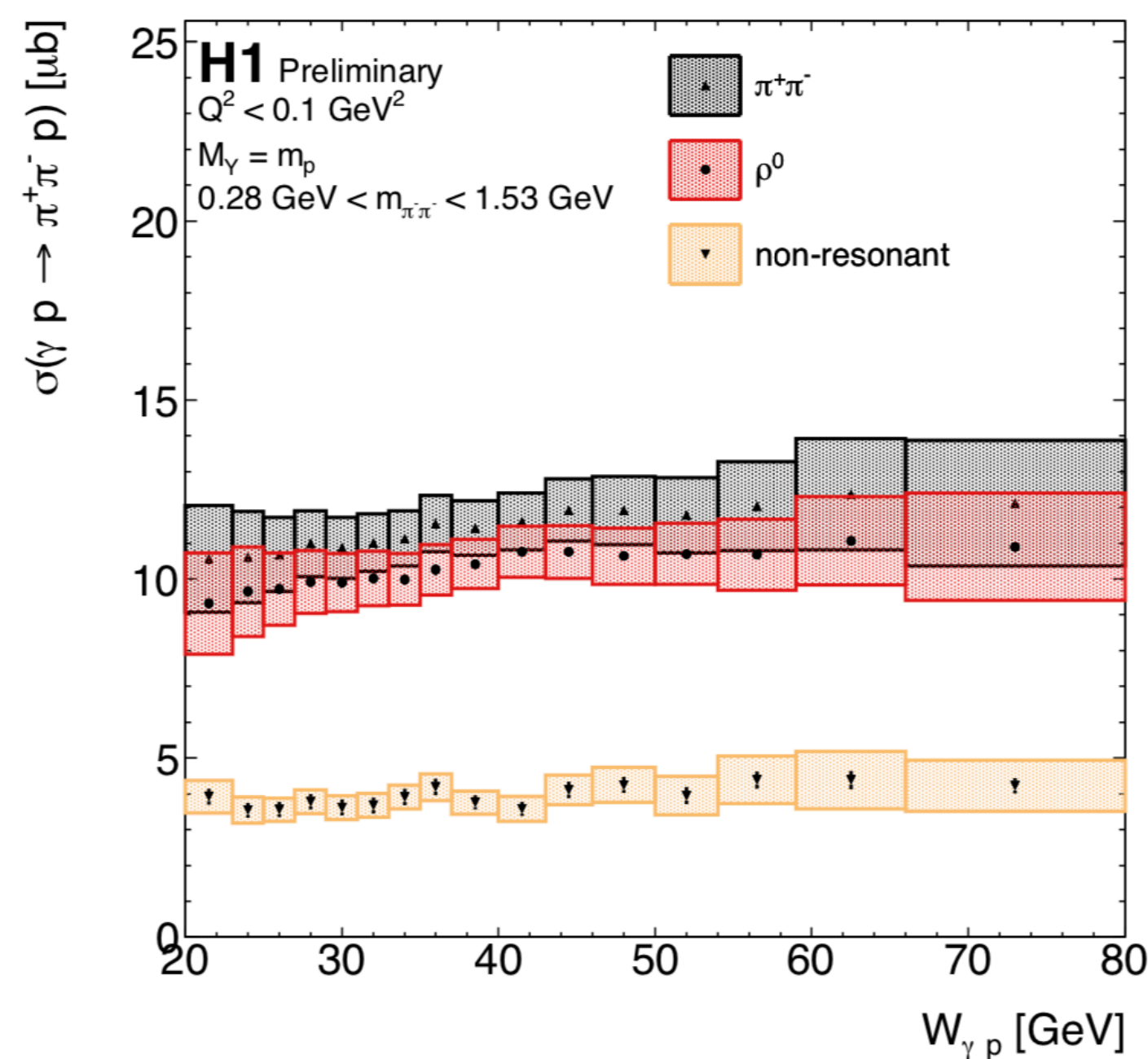
Measurement of exclusive and **dissociative** cross sections!

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Wednesday in the WG2 session

Exclusive $\rho(770)$ photoproduction at HERA, Arthur Bolz.

<http://www-h1.desy.de/publications/htmlsplit/H1prelim-18-012.long.html>



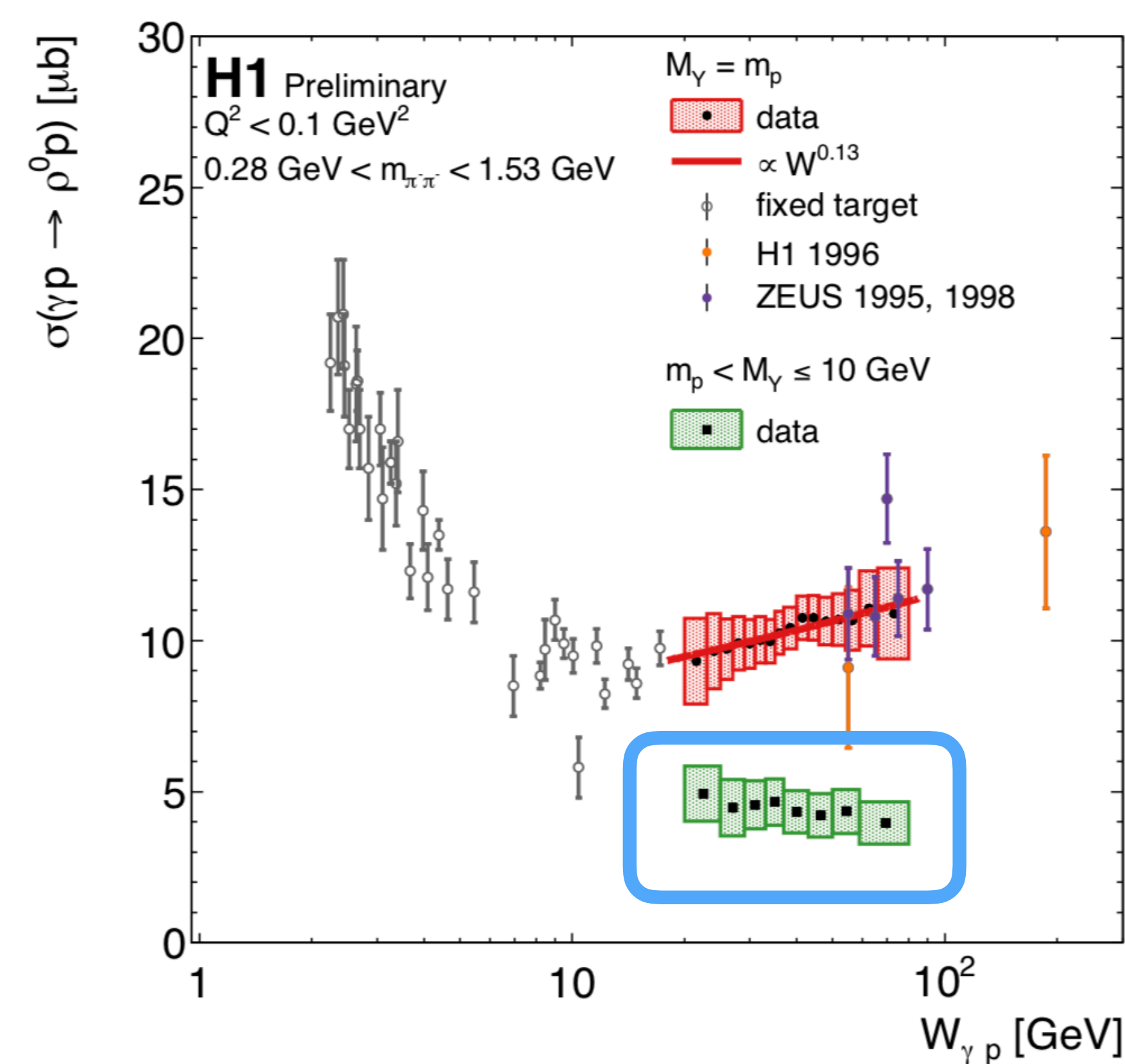
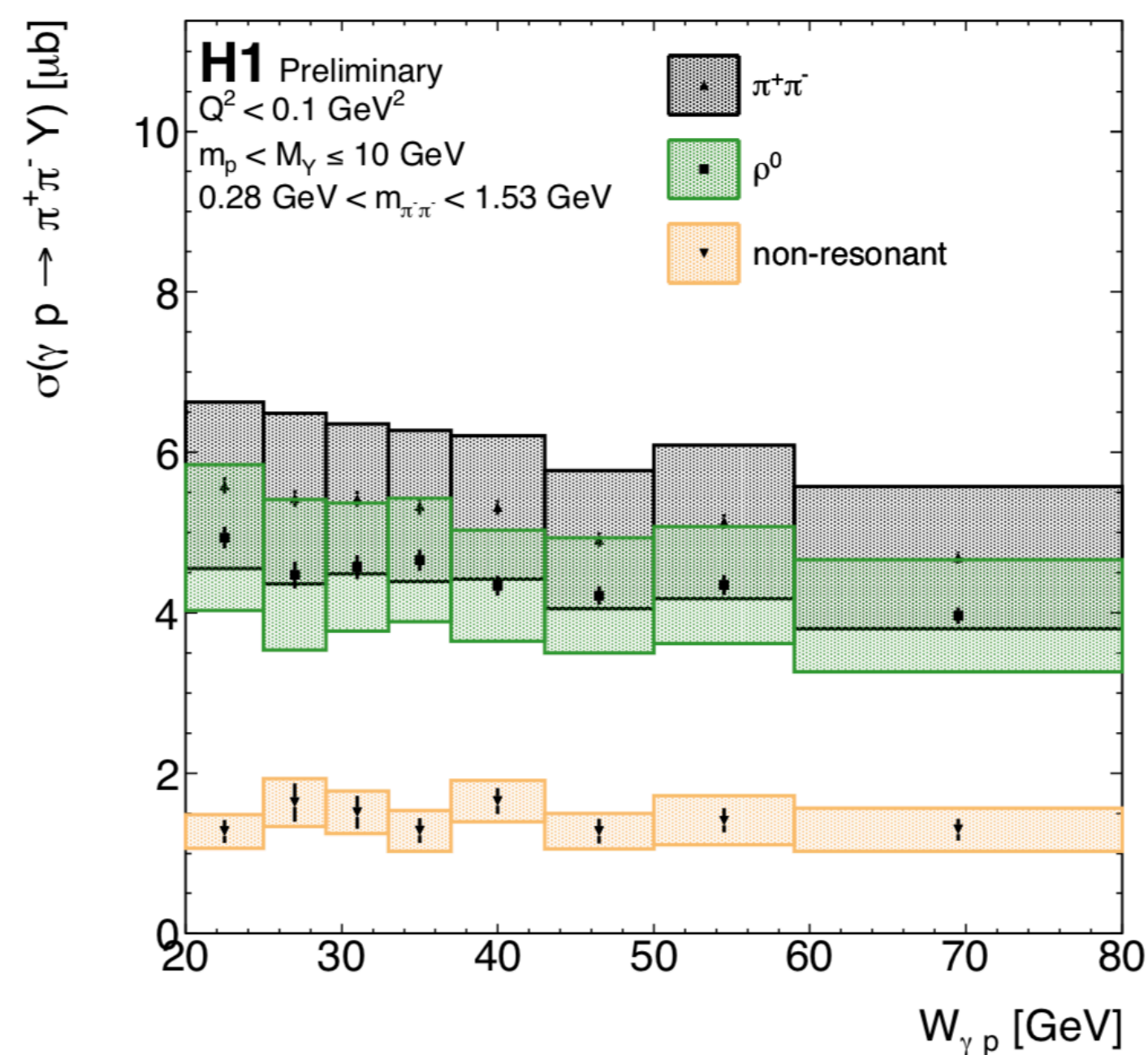
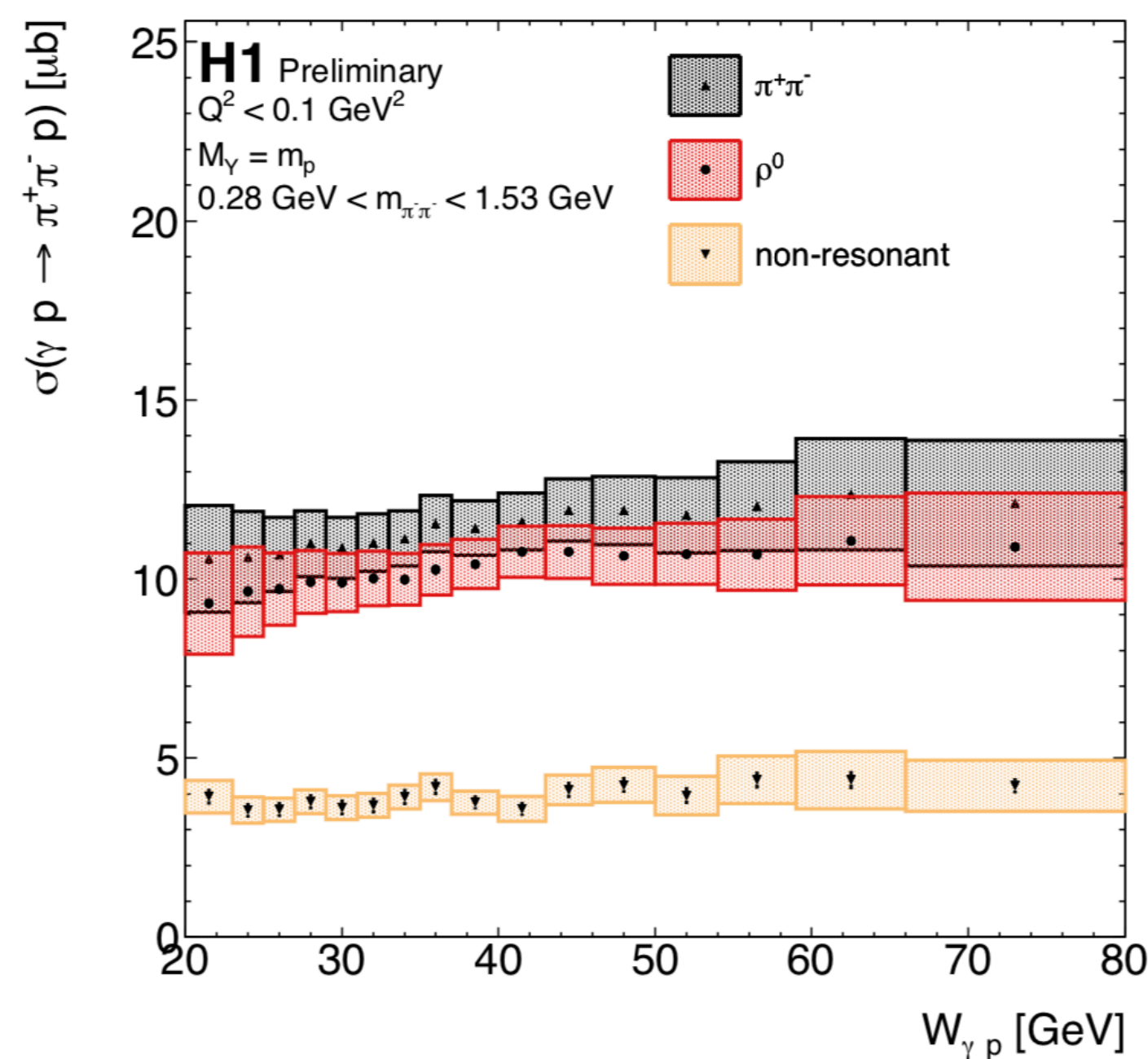
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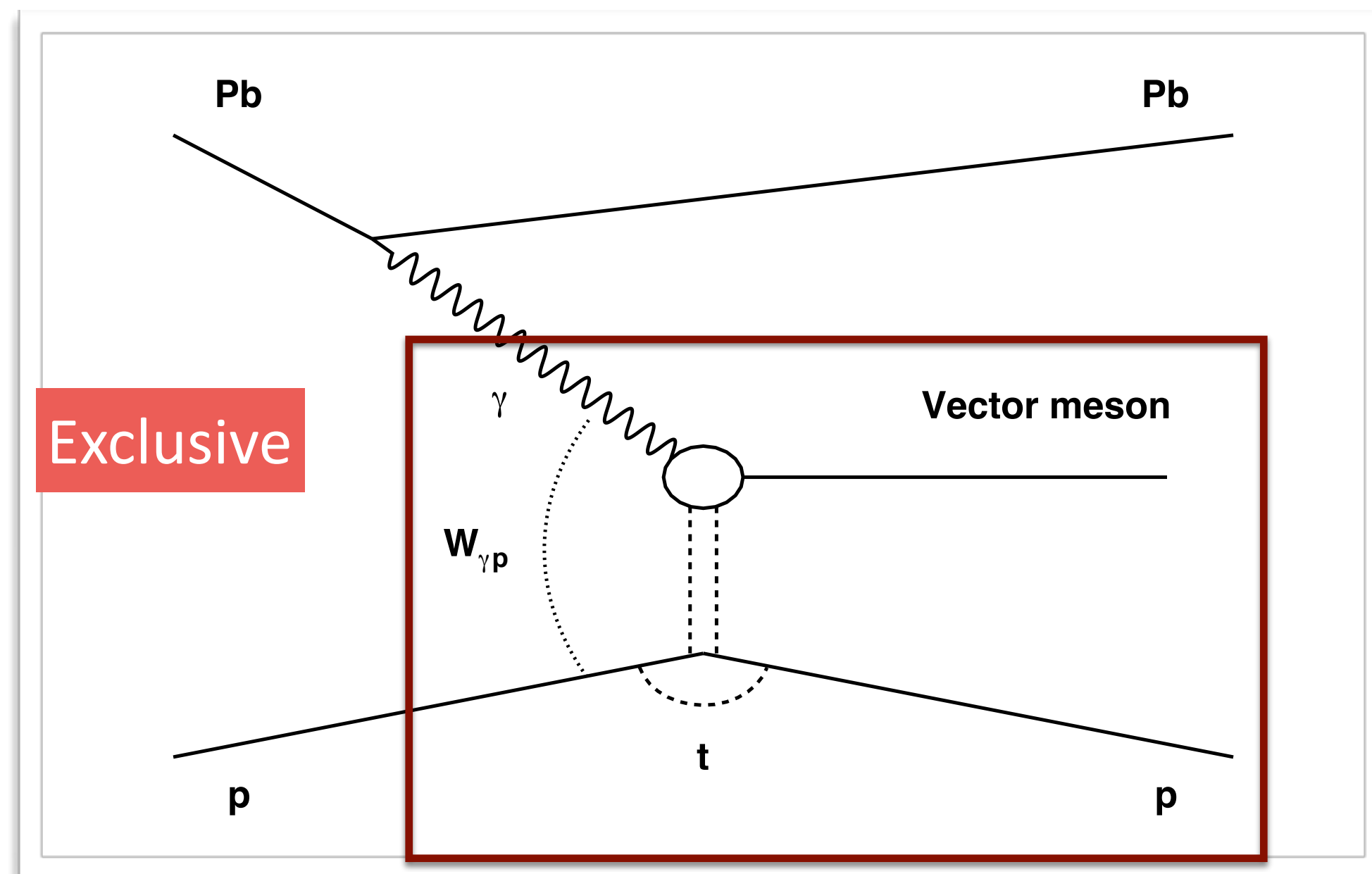
<http://www-h1.desy.de/publications/htmlsplit/H1prelim-18-012.long.html>



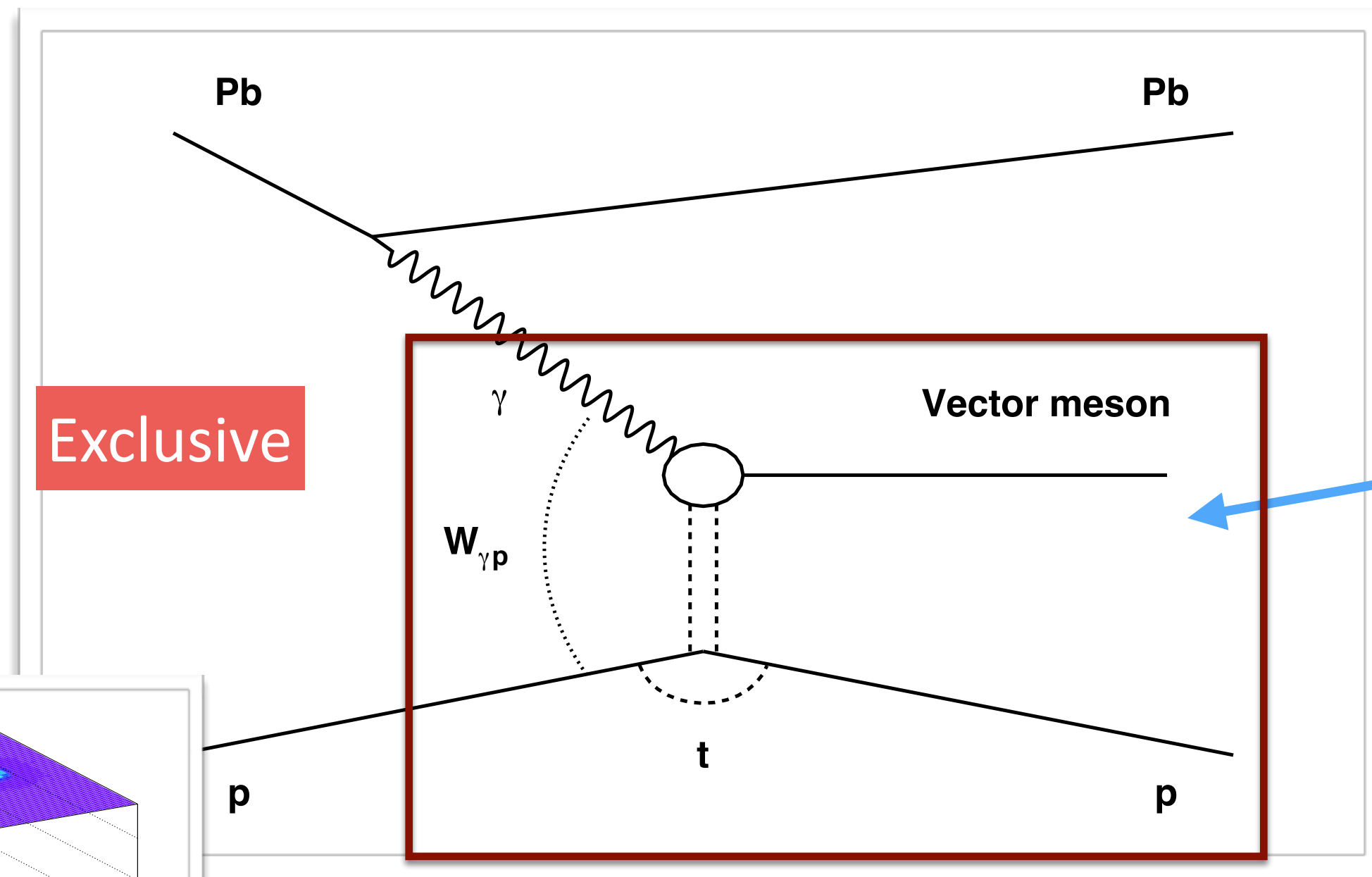
Measurement of exclusive and **dissociative** cross sections!

Dissociation: a new window to saturation at future EIC?

Exclusive and dissociative vector meson production in a Good-Walker approach

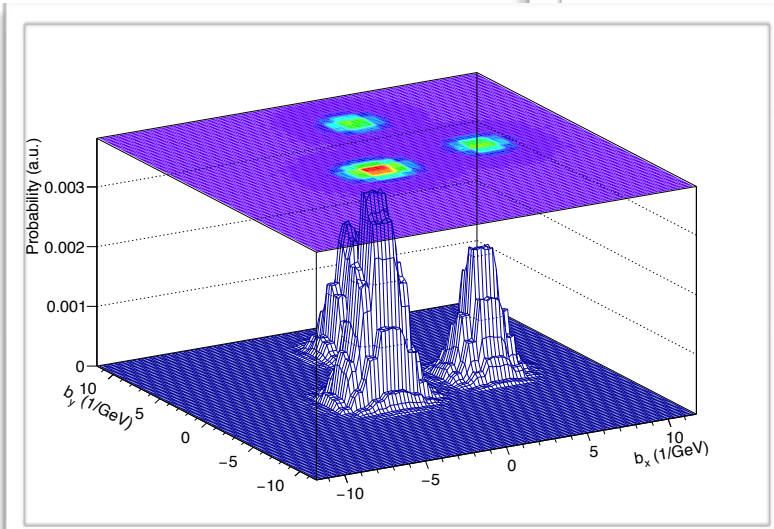


Exclusive and dissociative vector meson production in a Good-Walker approach



$$\frac{d\sigma(\gamma p \rightarrow J/\psi p)}{dt} = \frac{R_g^2}{16\pi} \left| \langle A(x, Q^2, \vec{\Delta}) \rangle \right|^2$$

Average over configurations



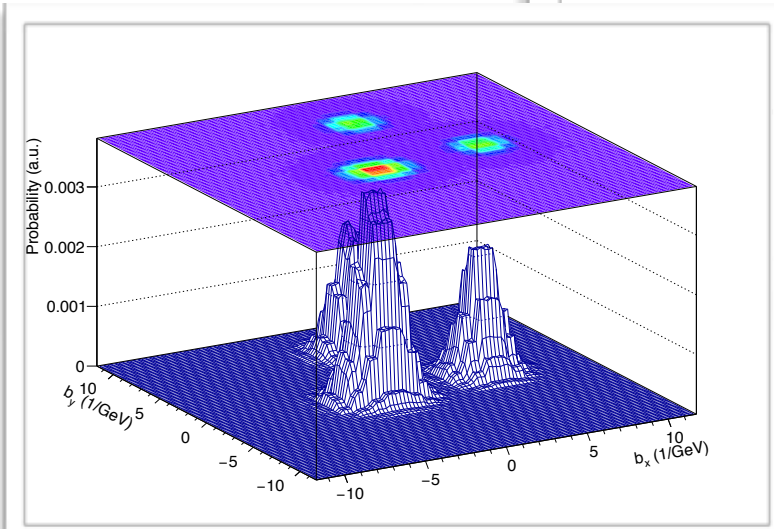
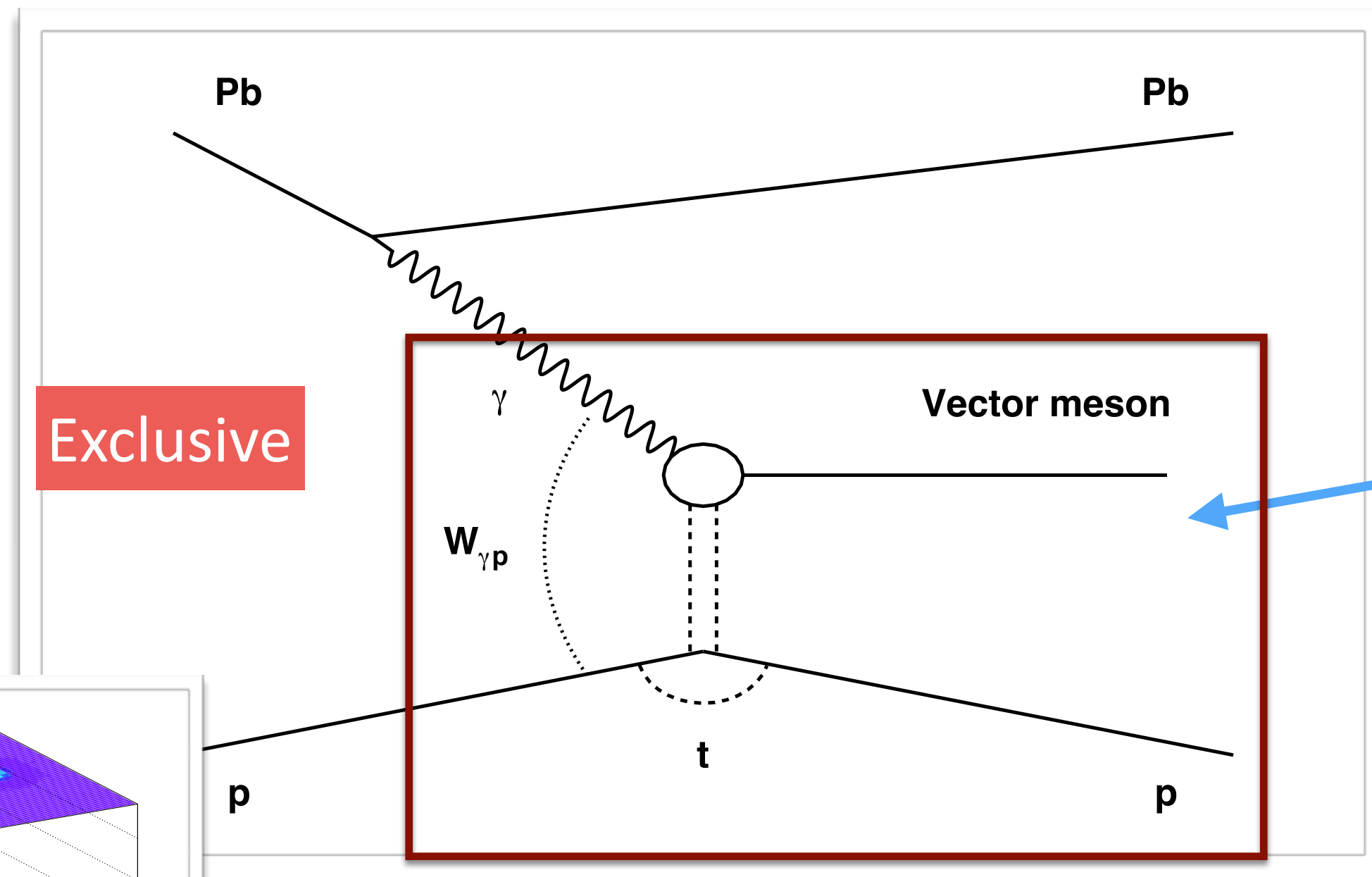
Good, Walker, PR 120 (1960) 1857

Miettinen, Pumplin, PRD18 (1978) 1696

Mantysaari, Schenke, PRL 117 (2016) 052301

Exclusive and dissociative vector meson production in a Good-Walker approach

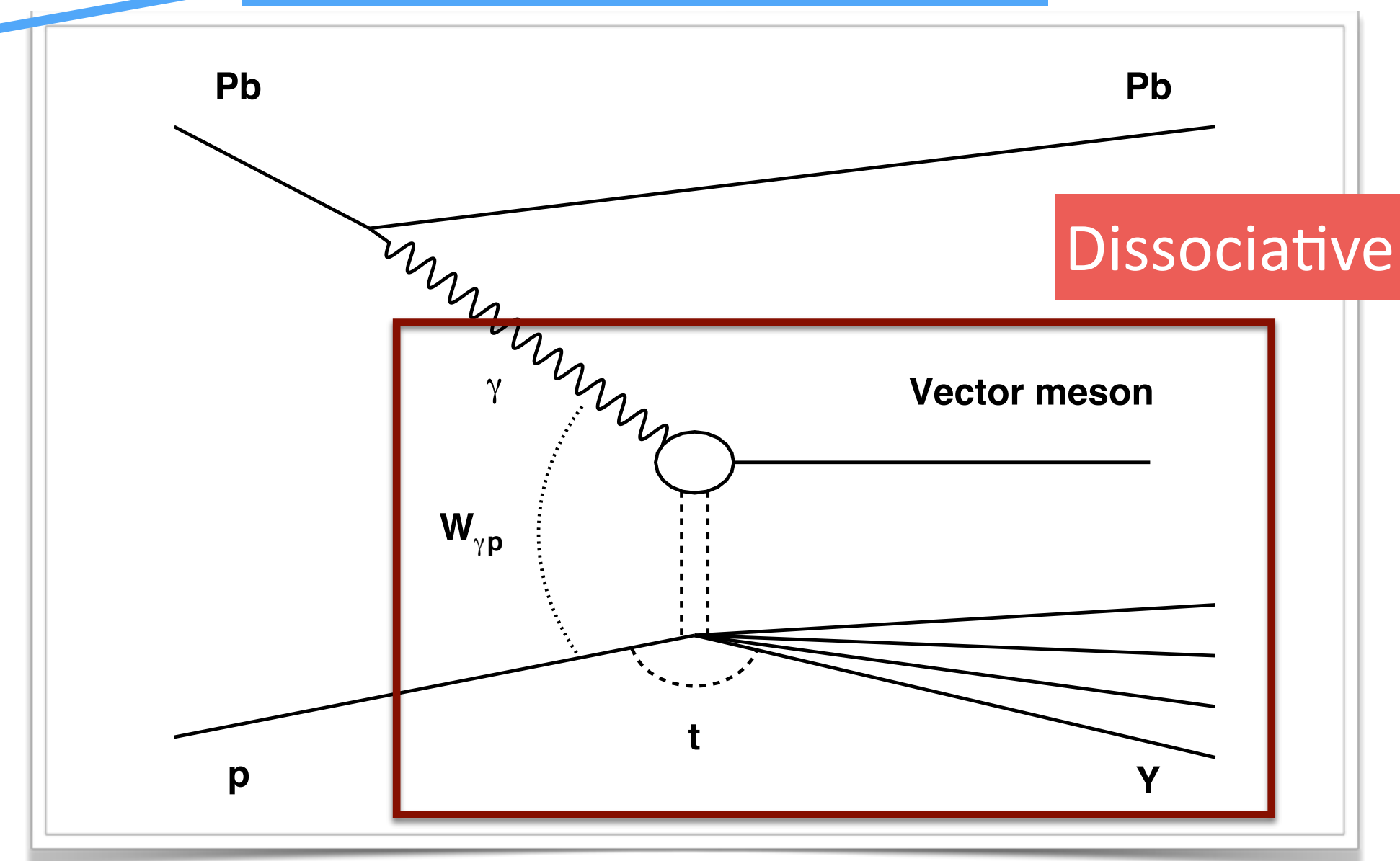
Exclusive



$$\frac{d\sigma(\gamma p \rightarrow J/\psi p)}{dt} = \frac{R_g^2}{16\pi} \left| \langle A(x, Q^2, \vec{\Delta}) \rangle \right|^2$$

Average over configurations

Dissociative



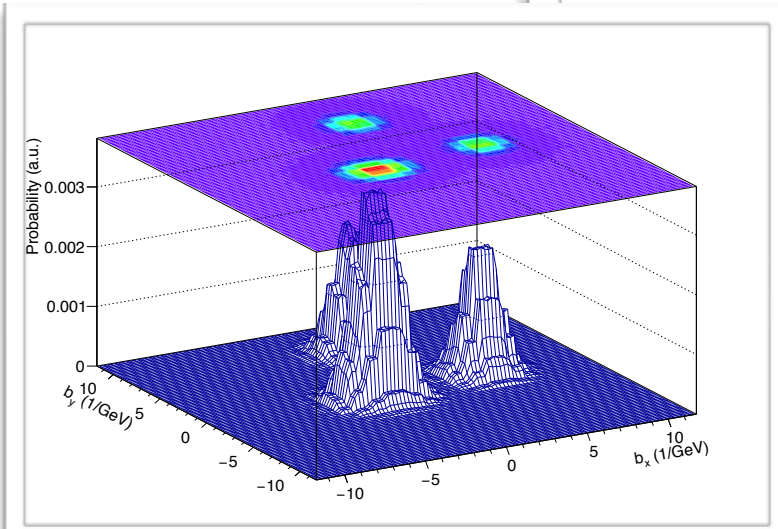
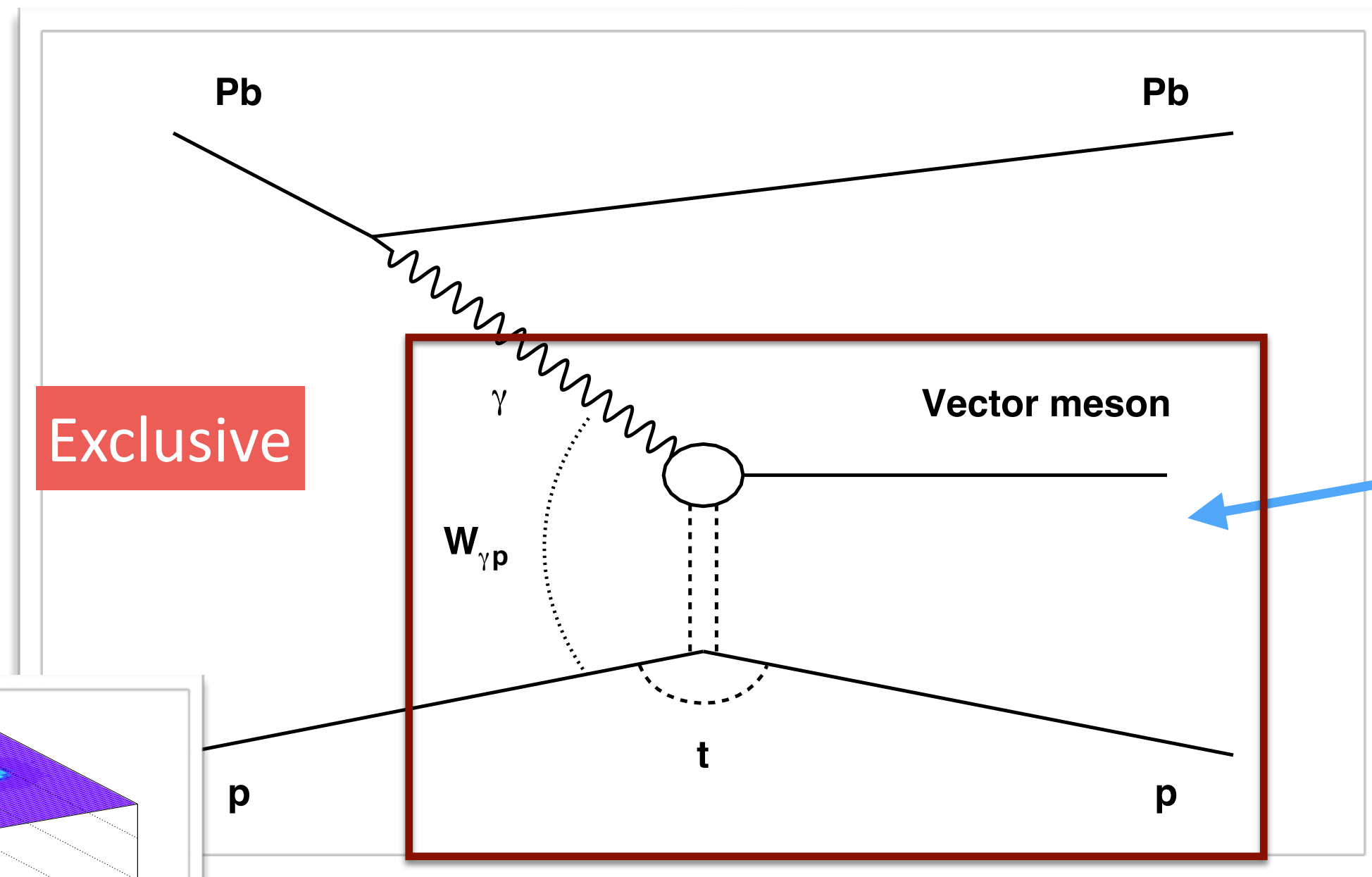
Good, Walker, PR 120 (1960) 1857

Miettinen, Pumplin, PRD18 (1978) 1696

Mantysaari, Schenke, PRL 117 (2016) 052301

Exclusive and dissociative vector meson production in a Good-Walker approach

Exclusive



Good, Walker, PR 120 (1960) 1857

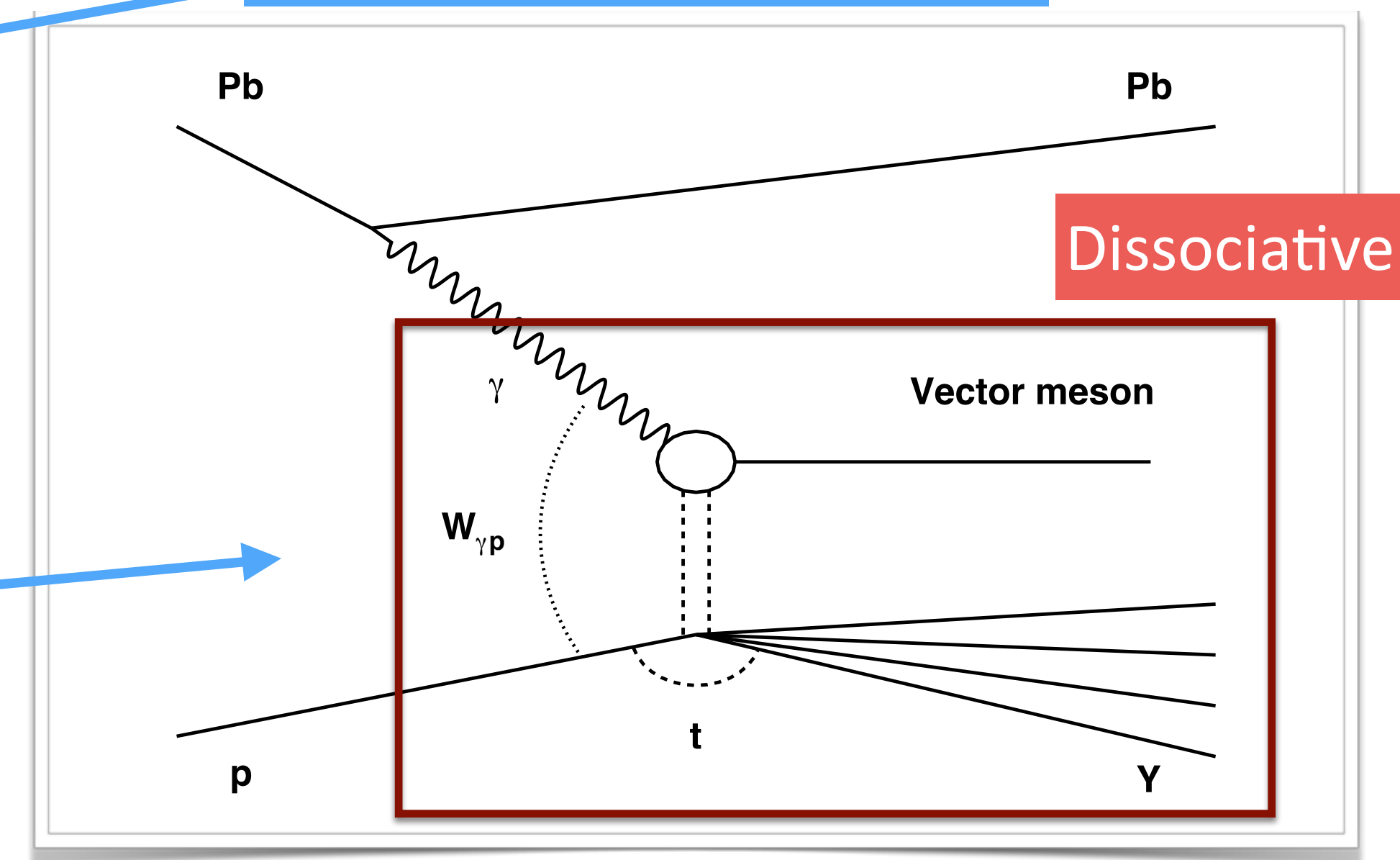
Miettinen, Pumplin, PRD18 (1978) 1696

Mantysaari, Schenke, PRL 117 (2016) 052301

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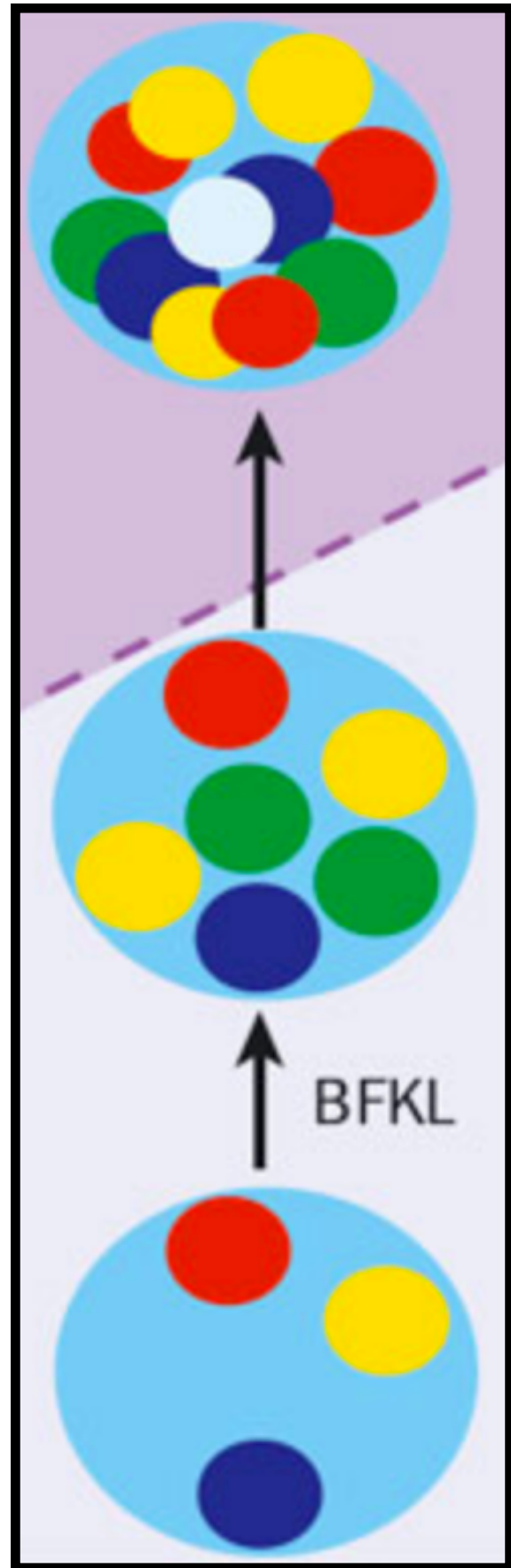
Dissociative



Variance over configurations

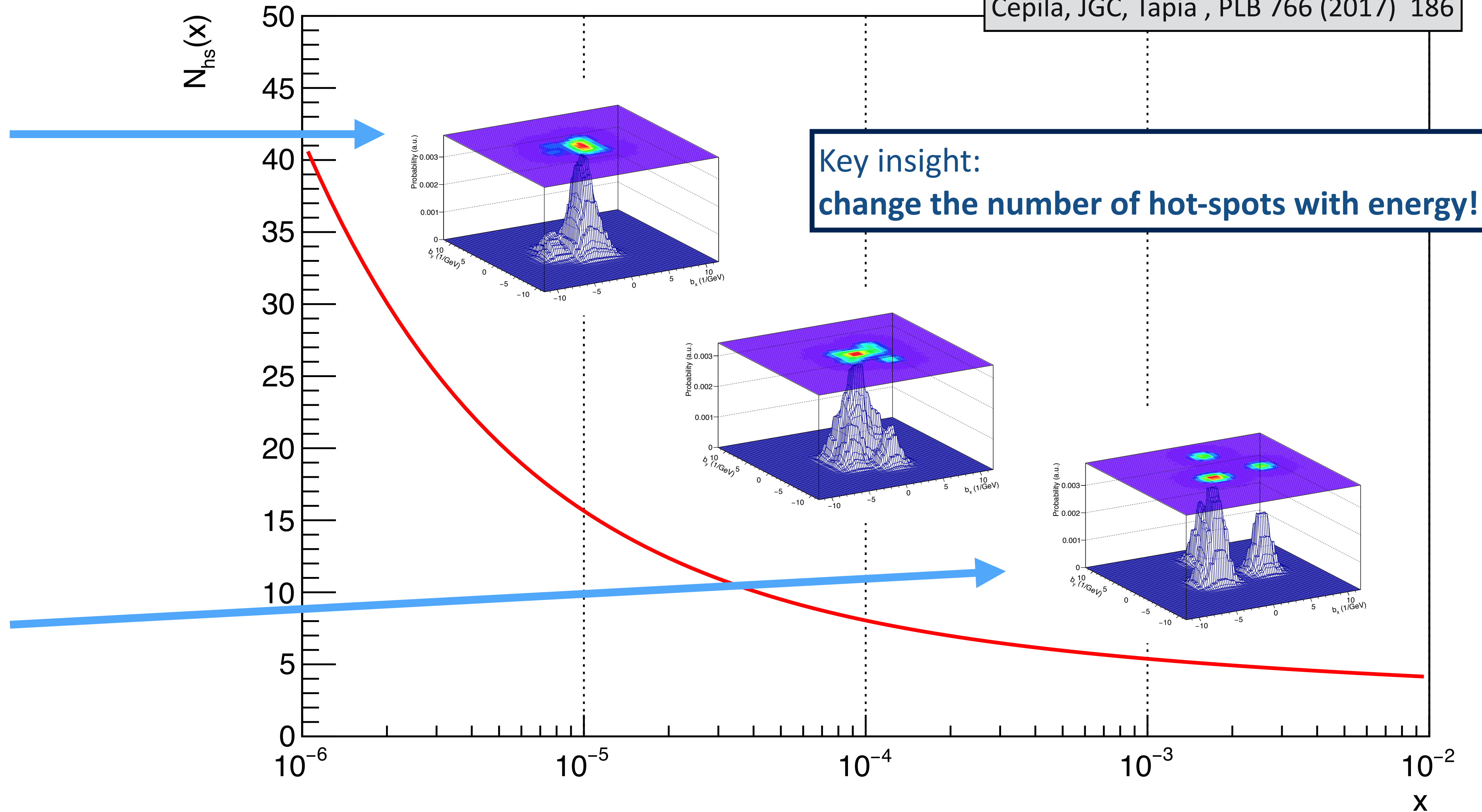
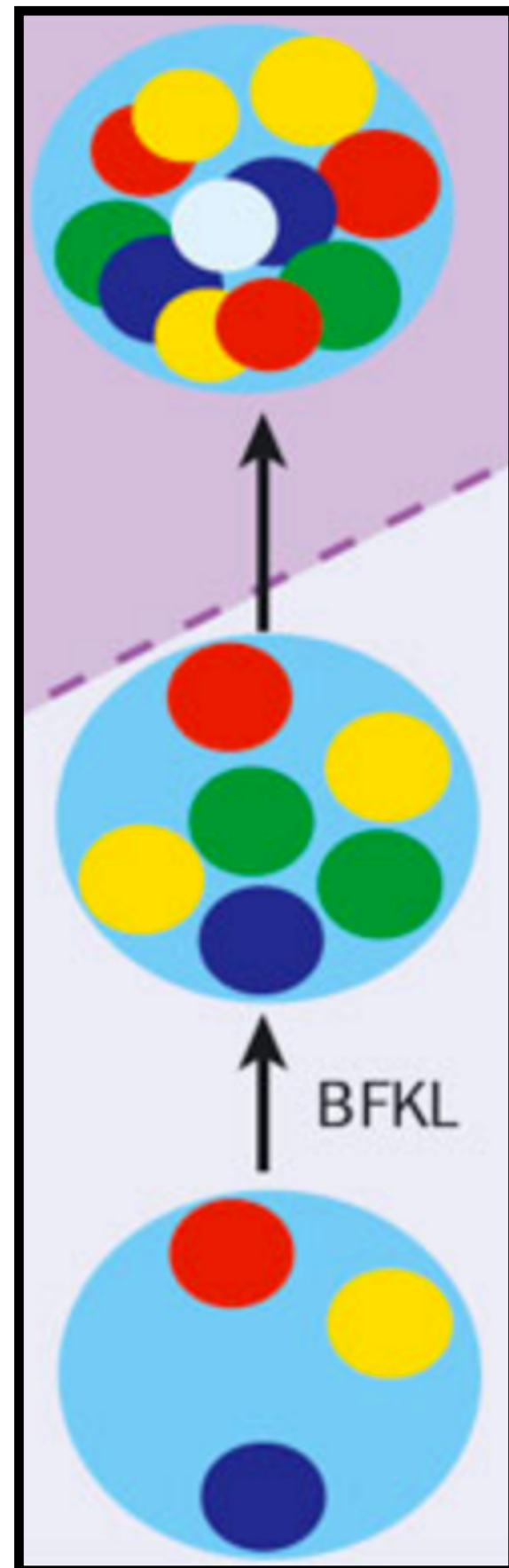
$$\frac{d\sigma(\gamma p \rightarrow J/\psi Y)}{dt} = \frac{R_g^2}{16\pi} \left(\langle |A(x, Q^2, \vec{\Delta})|^2 \rangle - \left| \langle A(x, Q^2, \vec{\Delta}) \rangle \right|^2 \right)$$

Energy-dependent hot spot model



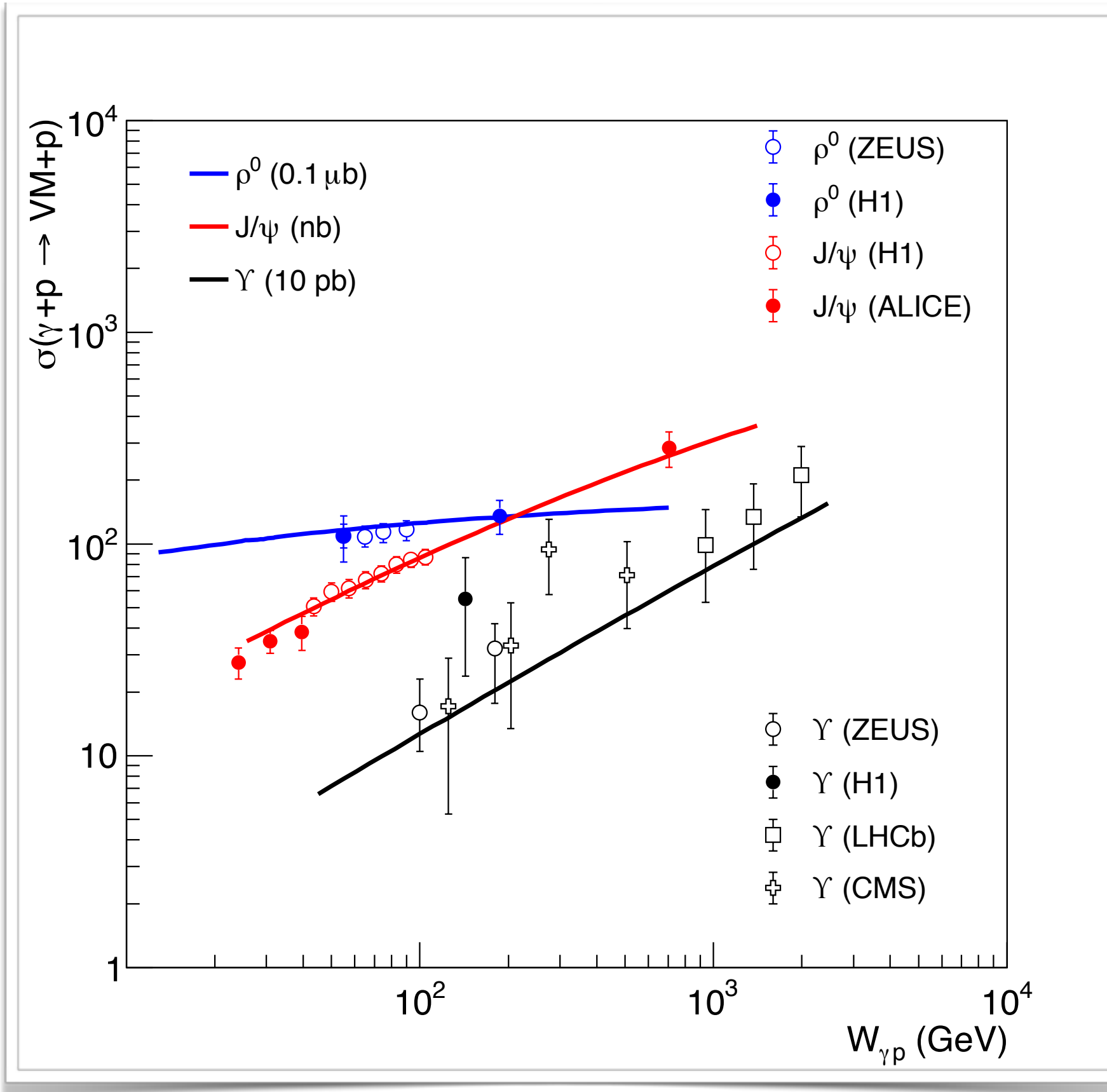
Energy-dependent hot spot model

Čepila, JGC, Tapia , PLB 766 (2017) 186

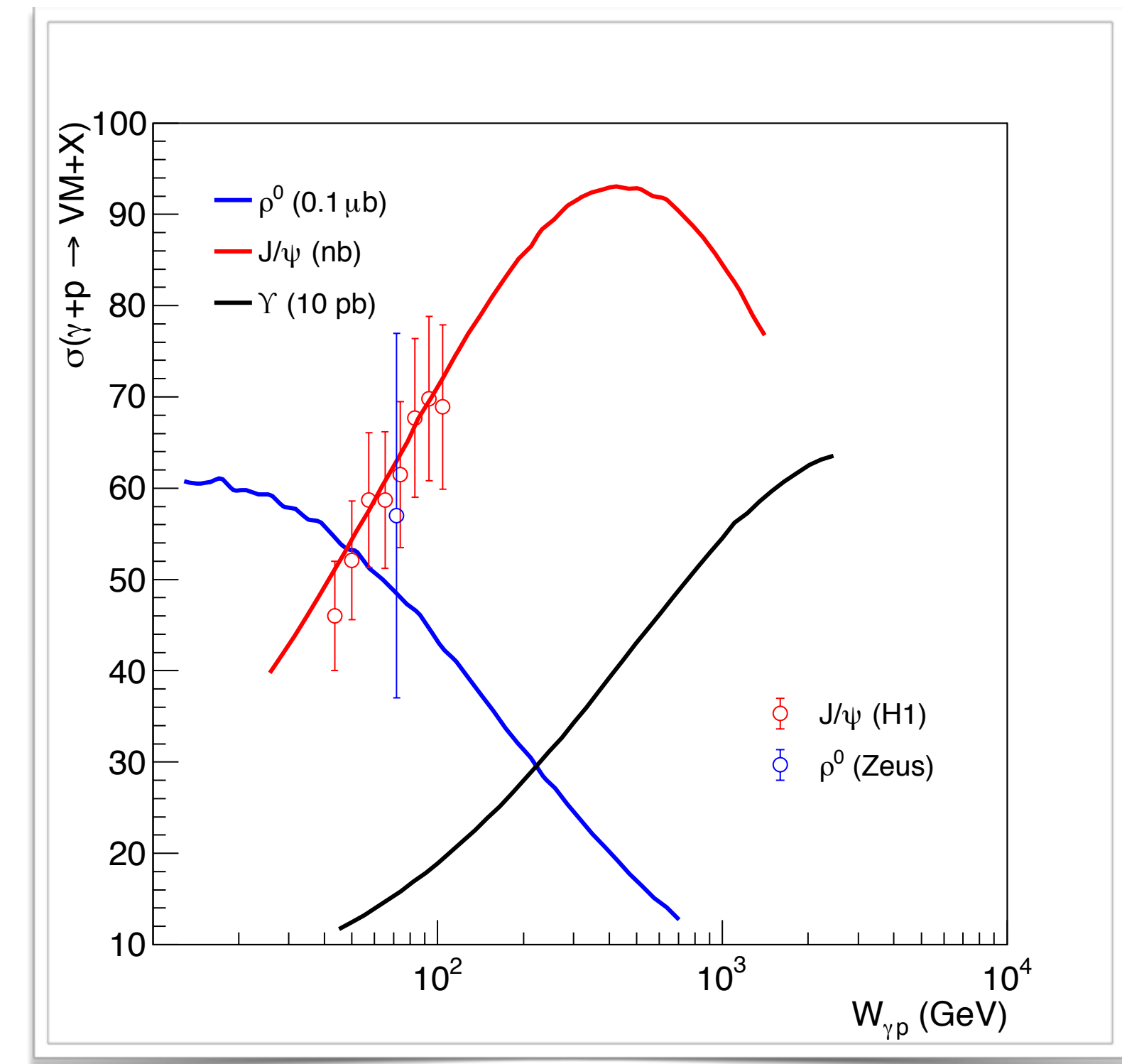


Prediction for the energy dependence of dissociative production

Čepila, JGC, Krelina, Tapia NPB934 (2018) 330

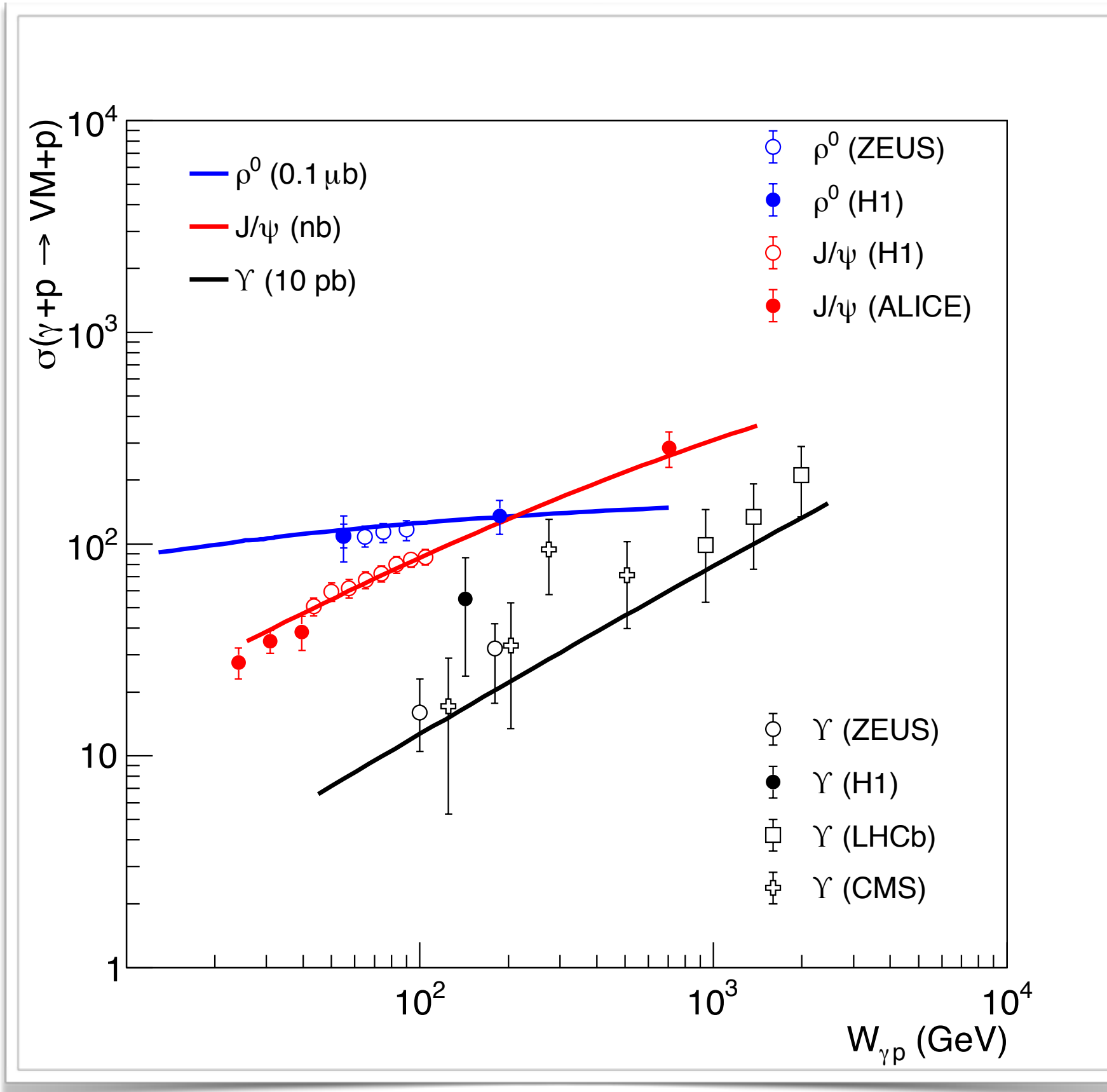


Correct description of available data.

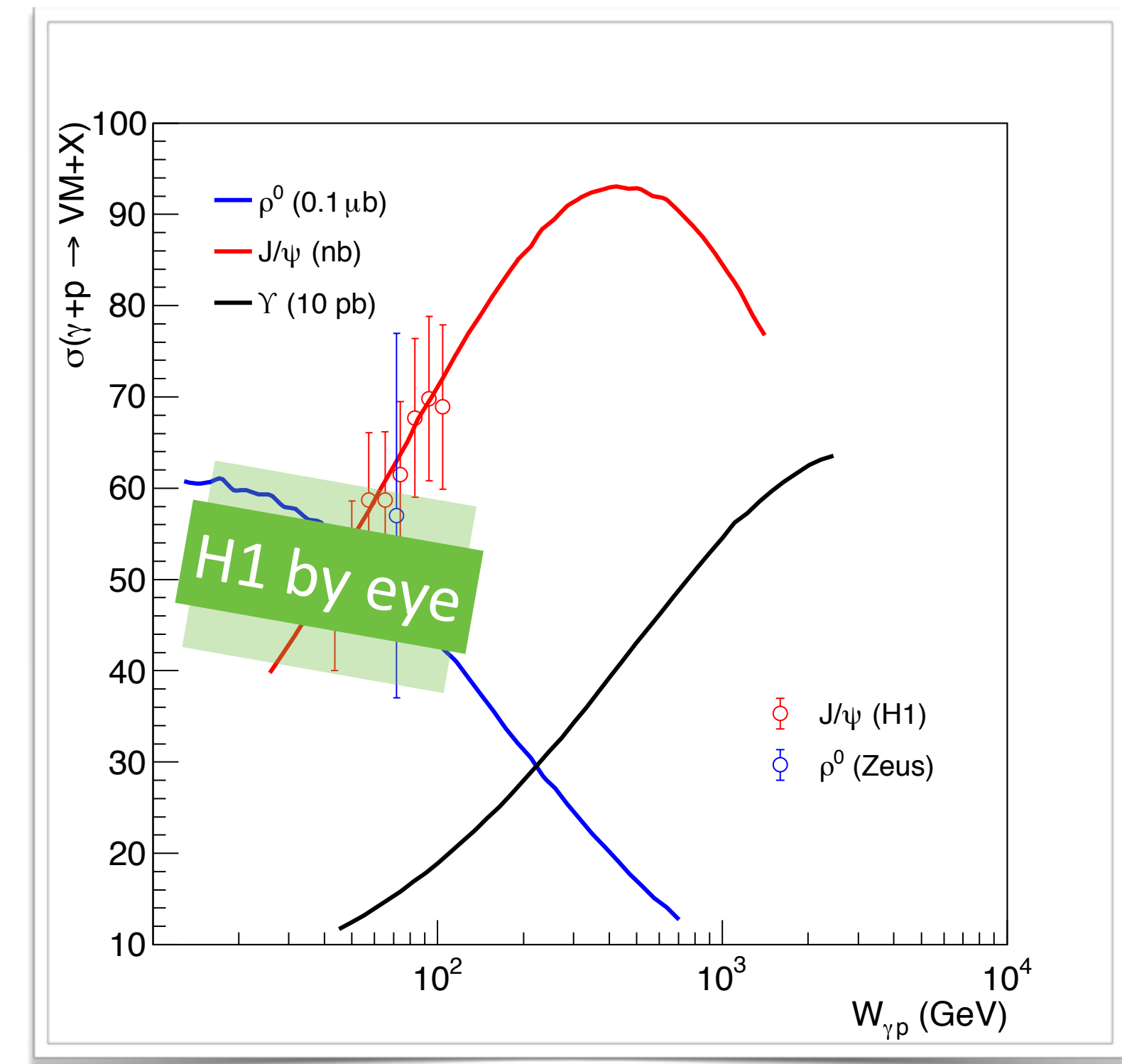


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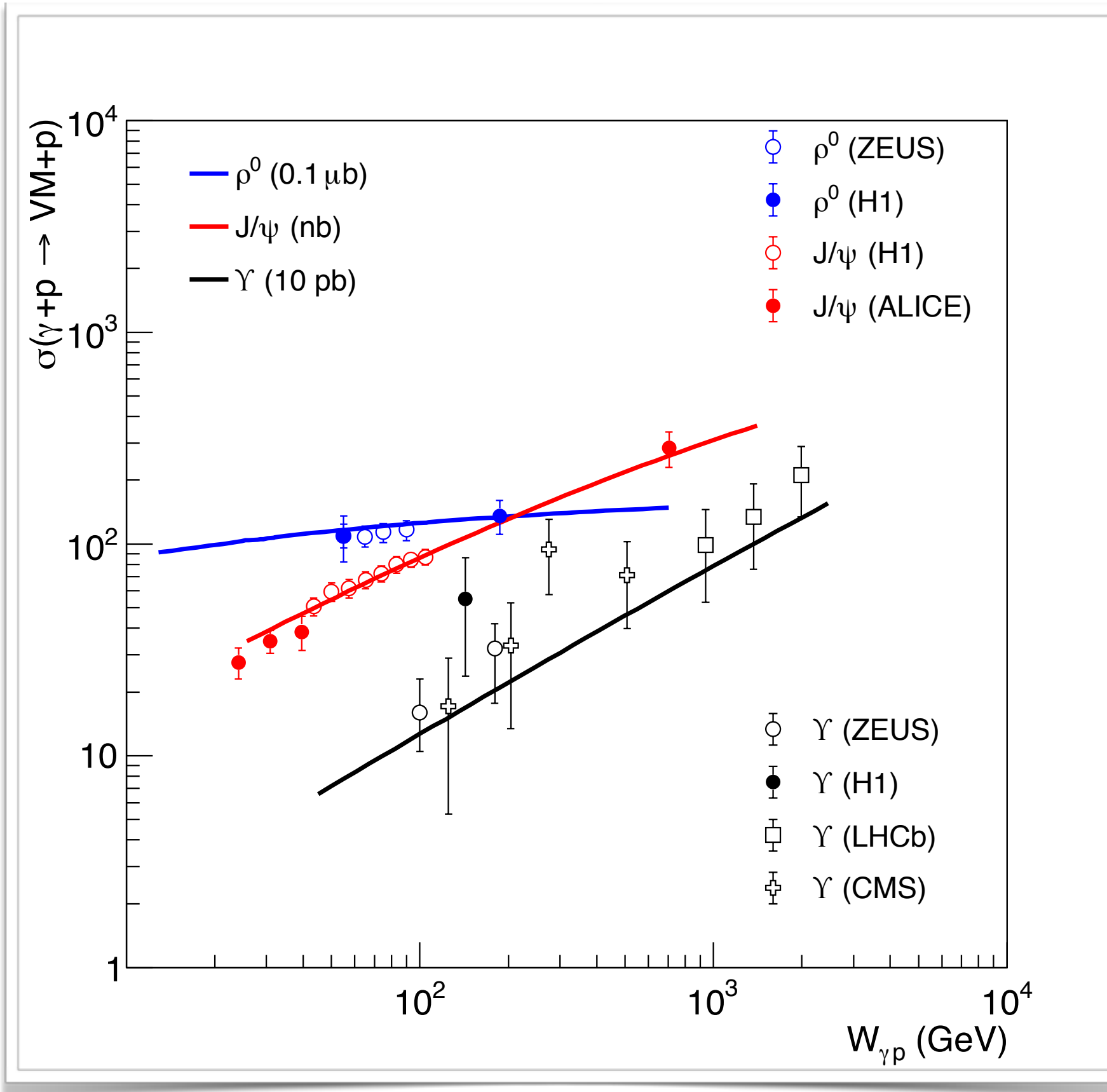


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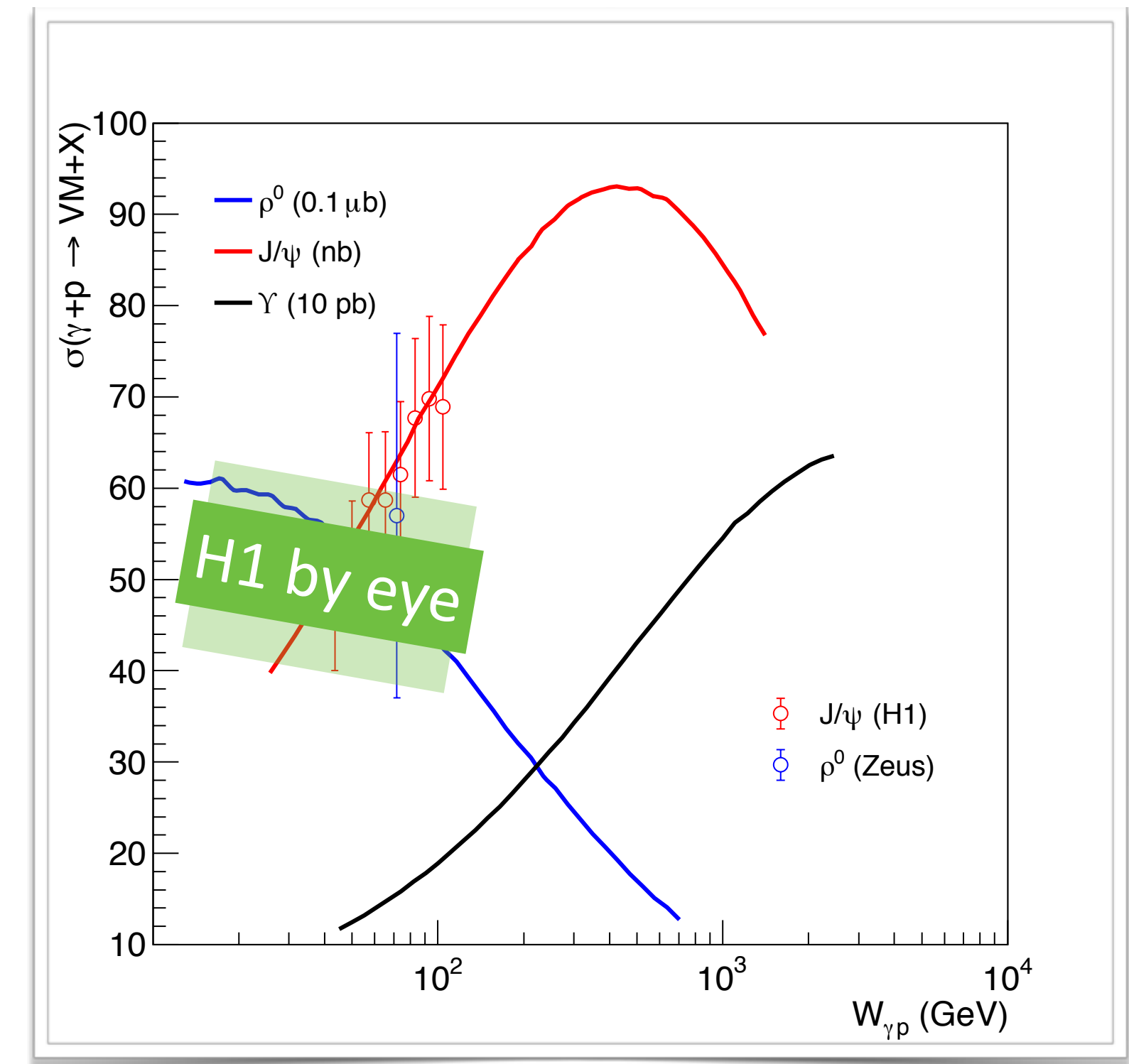


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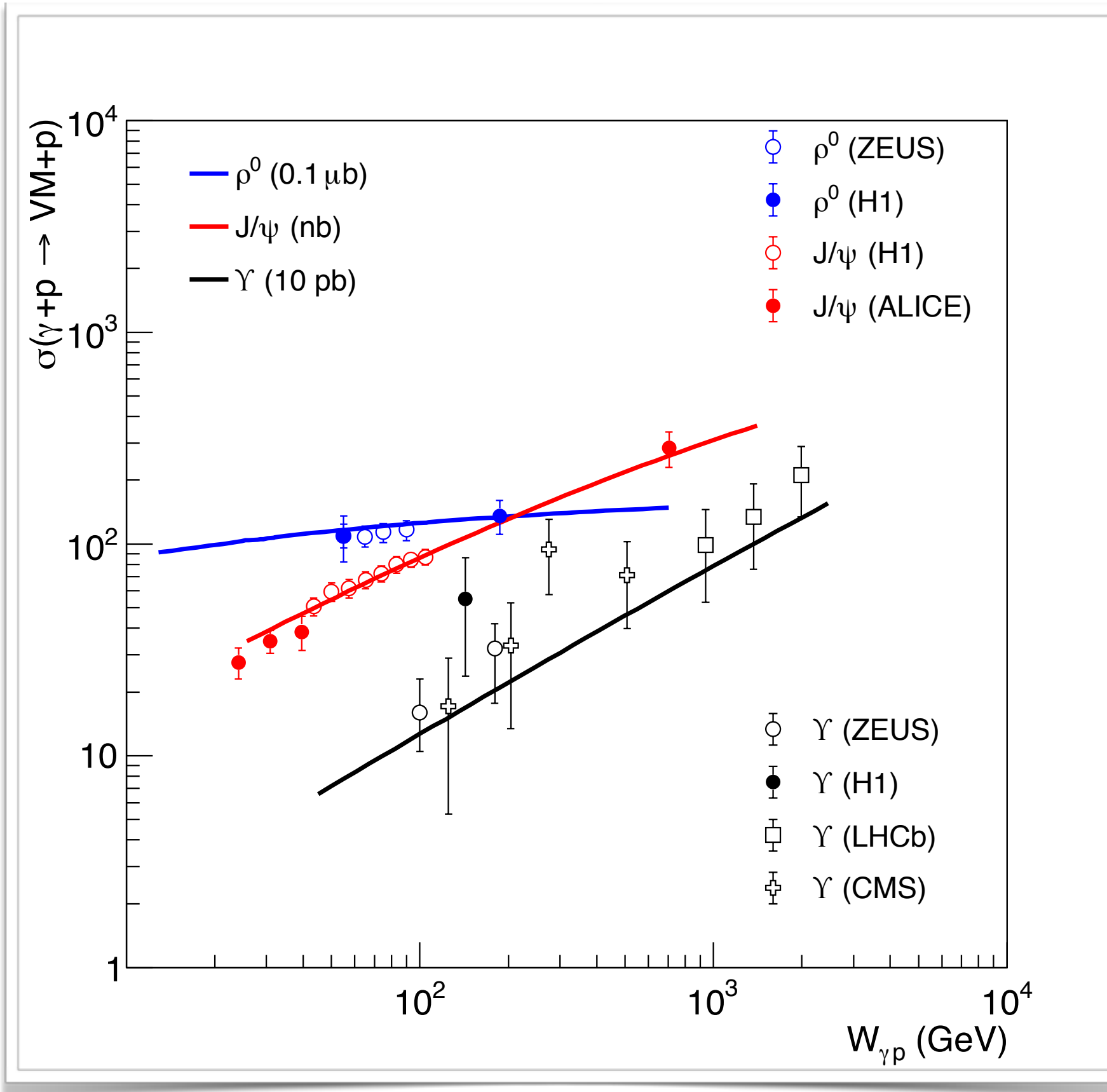
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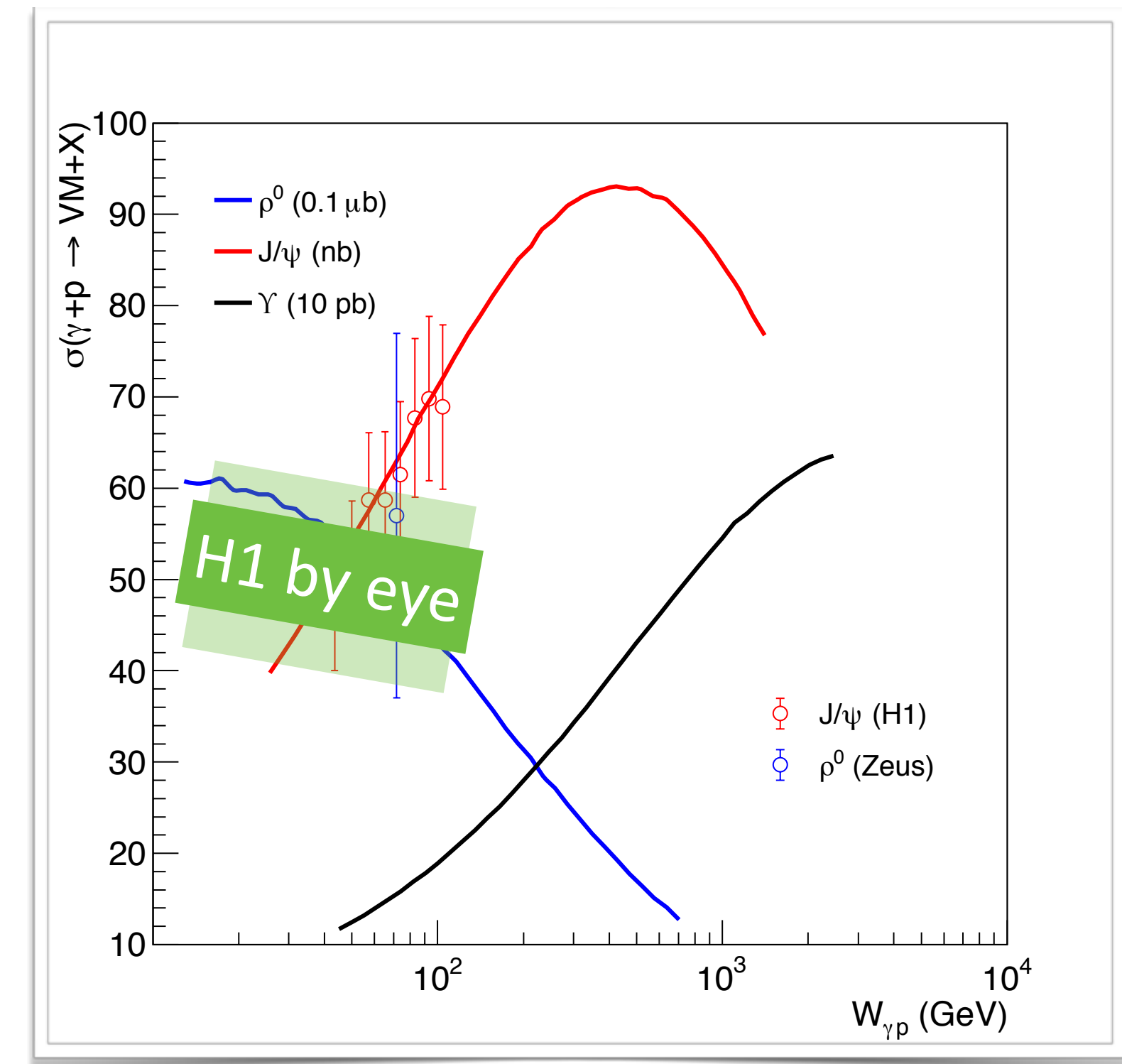
Behaviour of dissociative cross section depends on the **mass** of the vector meson.

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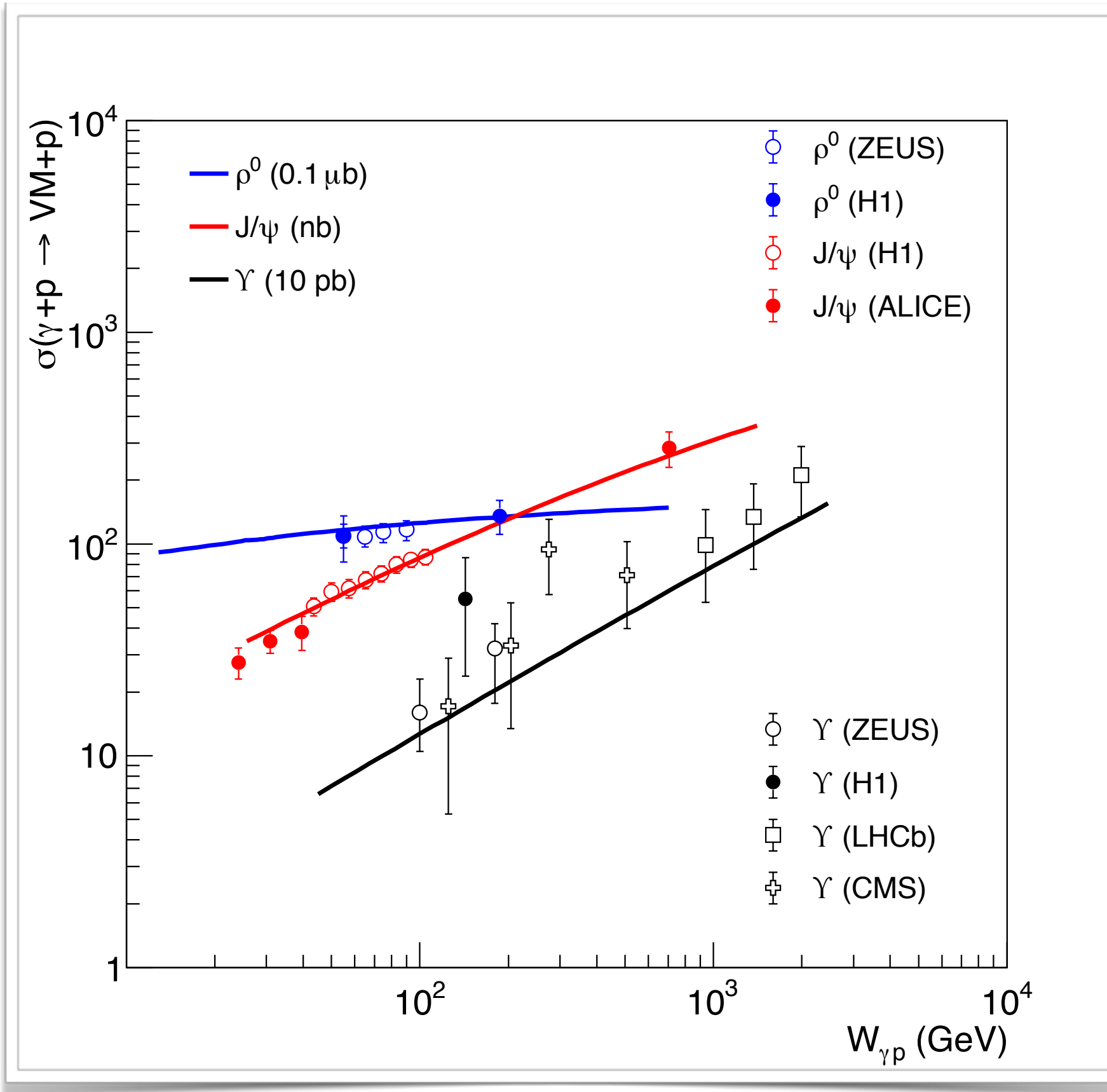


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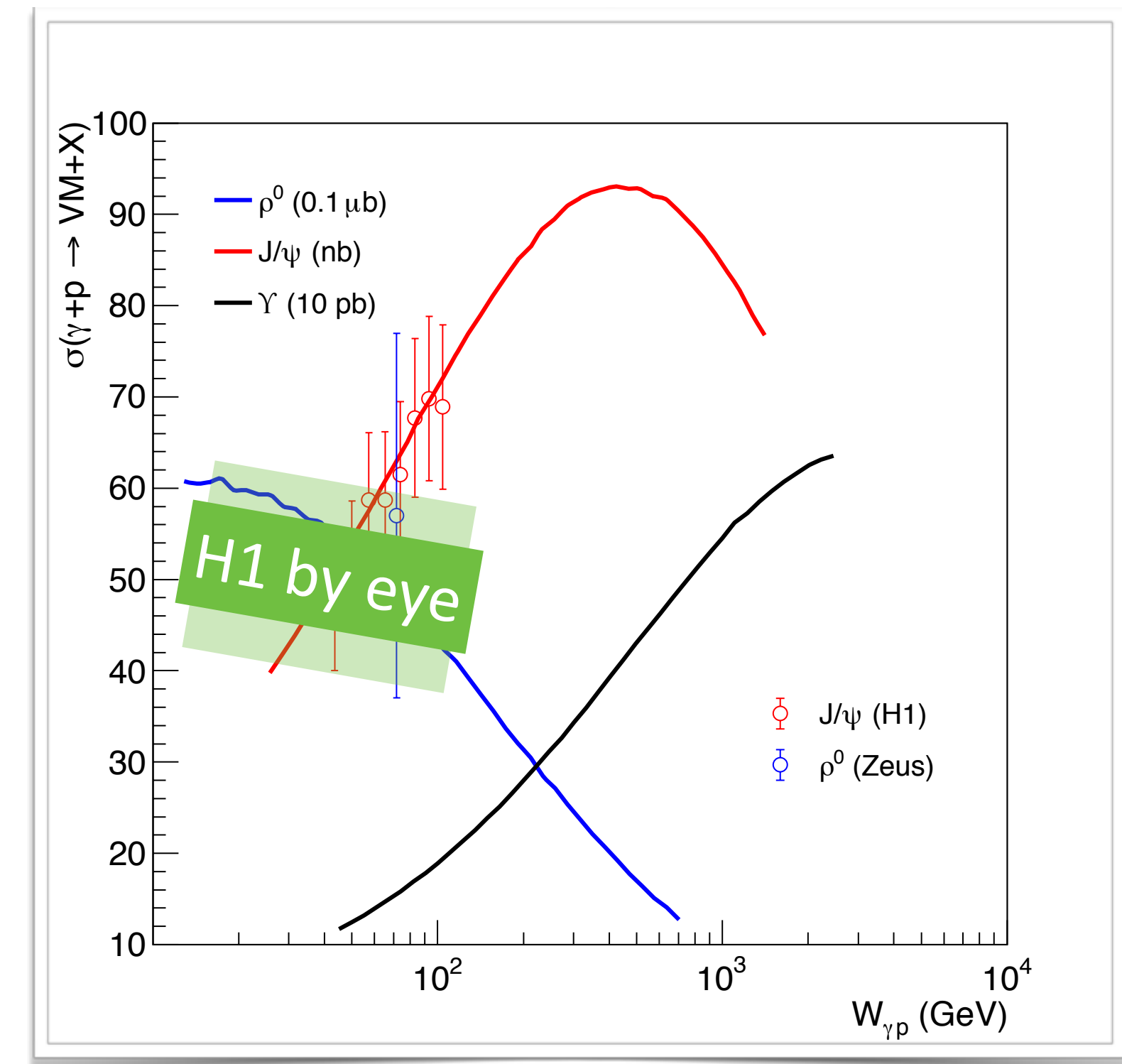
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All targets look the same: **saturation in transverse plane.**

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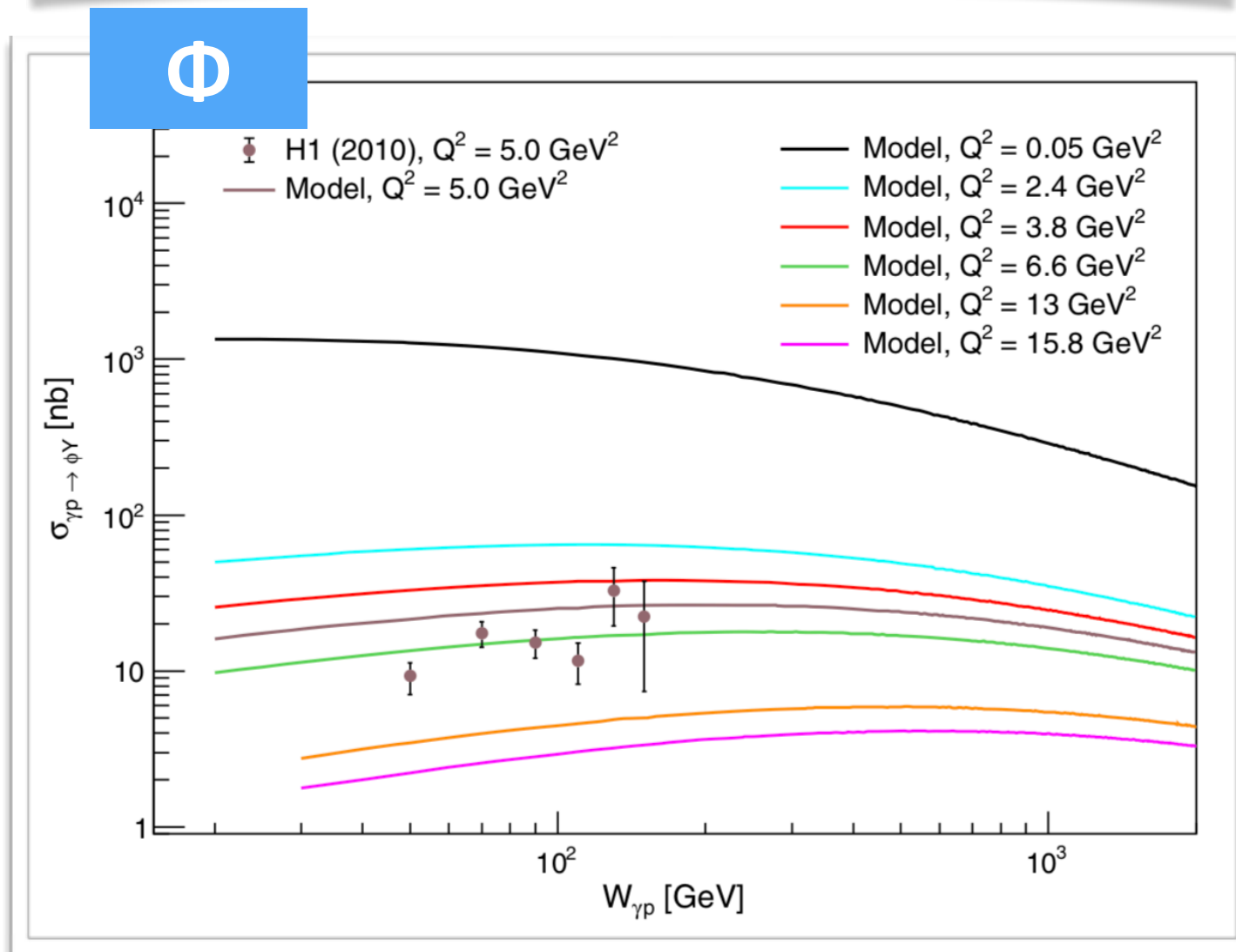
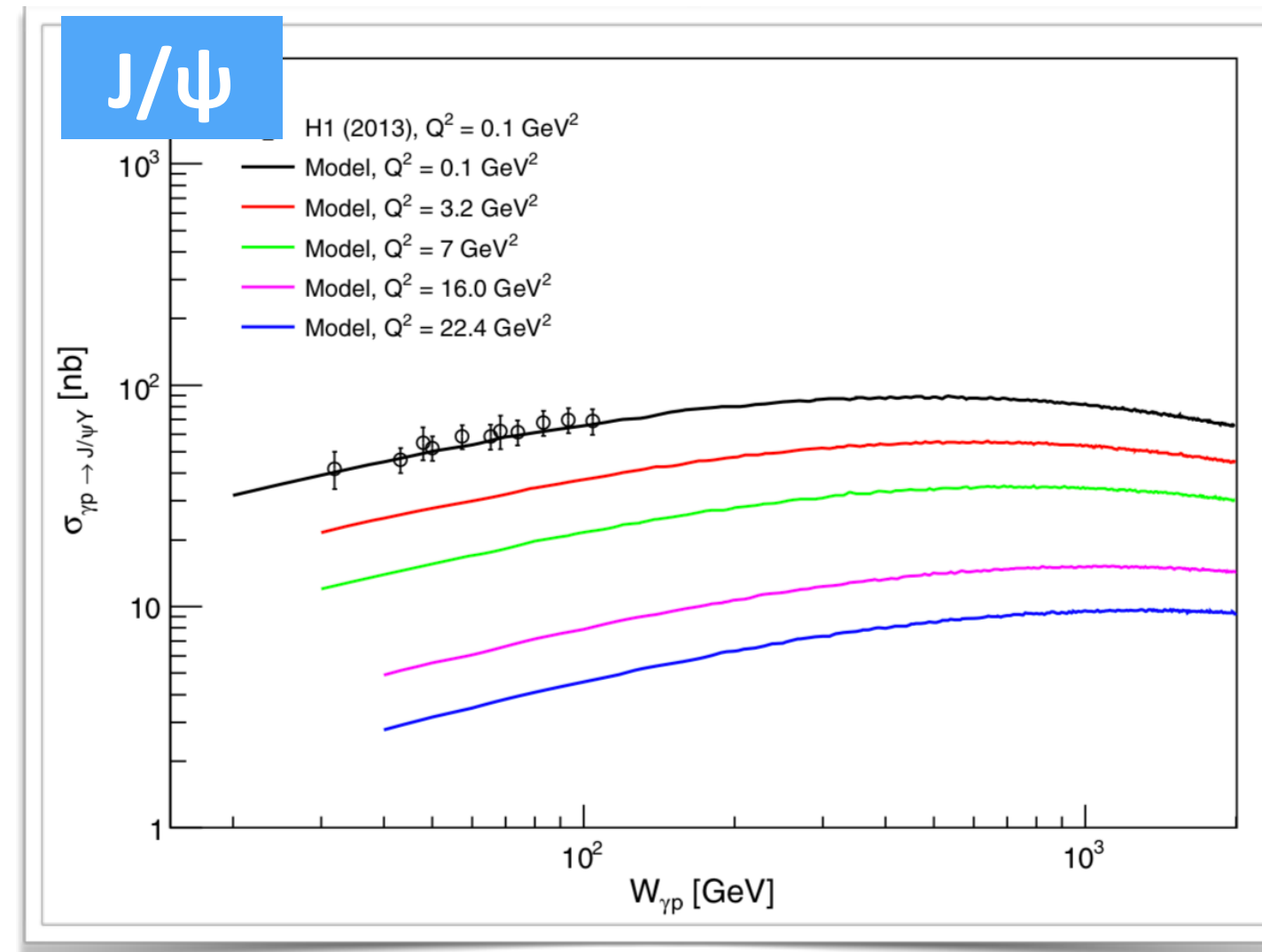
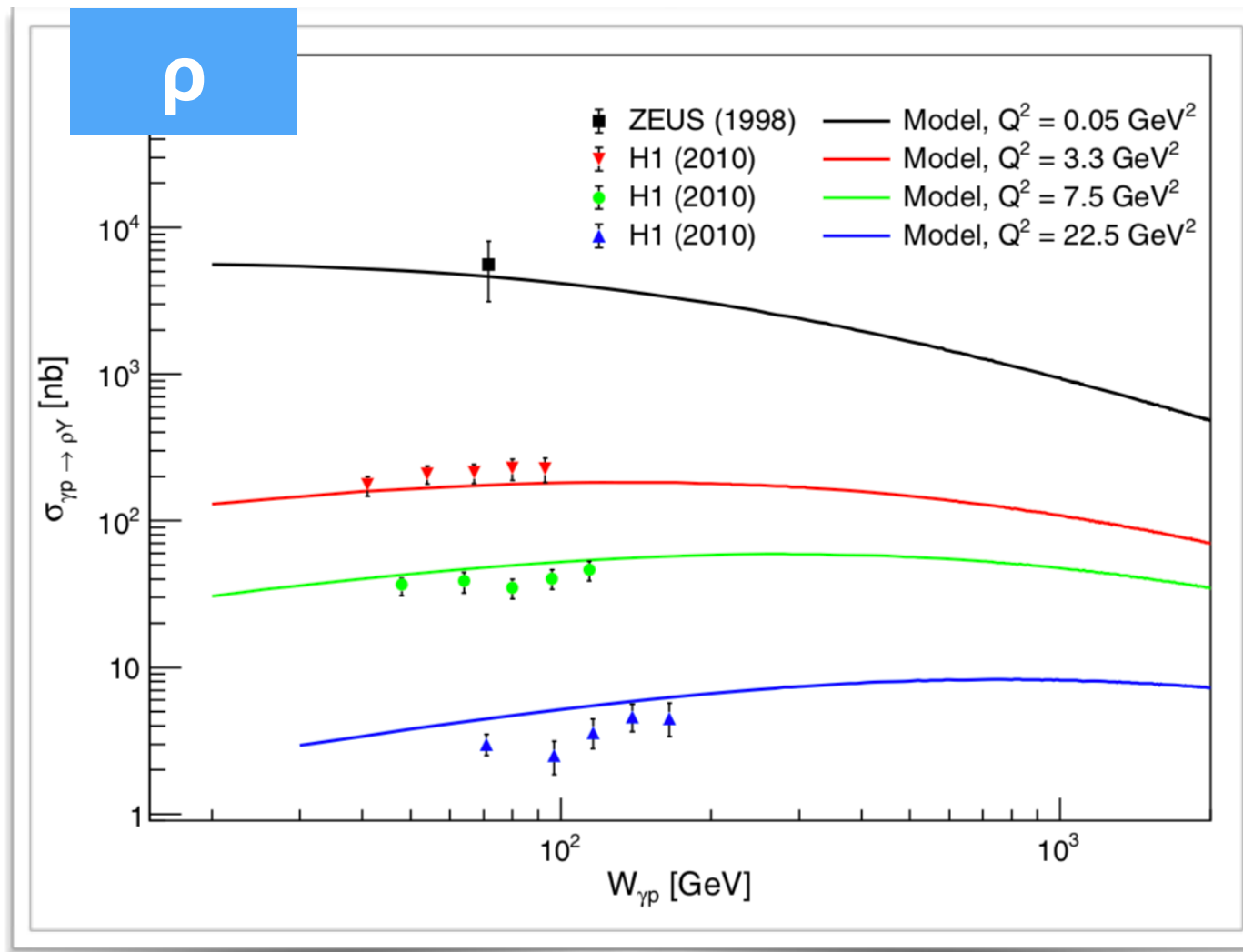
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Saturation effects kick in at the position of the maximum

Dissociative electroproduction of vector mesons: new window to saturation?

Bendová, Čepila, JGC, PRD99 (2019) 034025

Different colours, different Q^2 values.

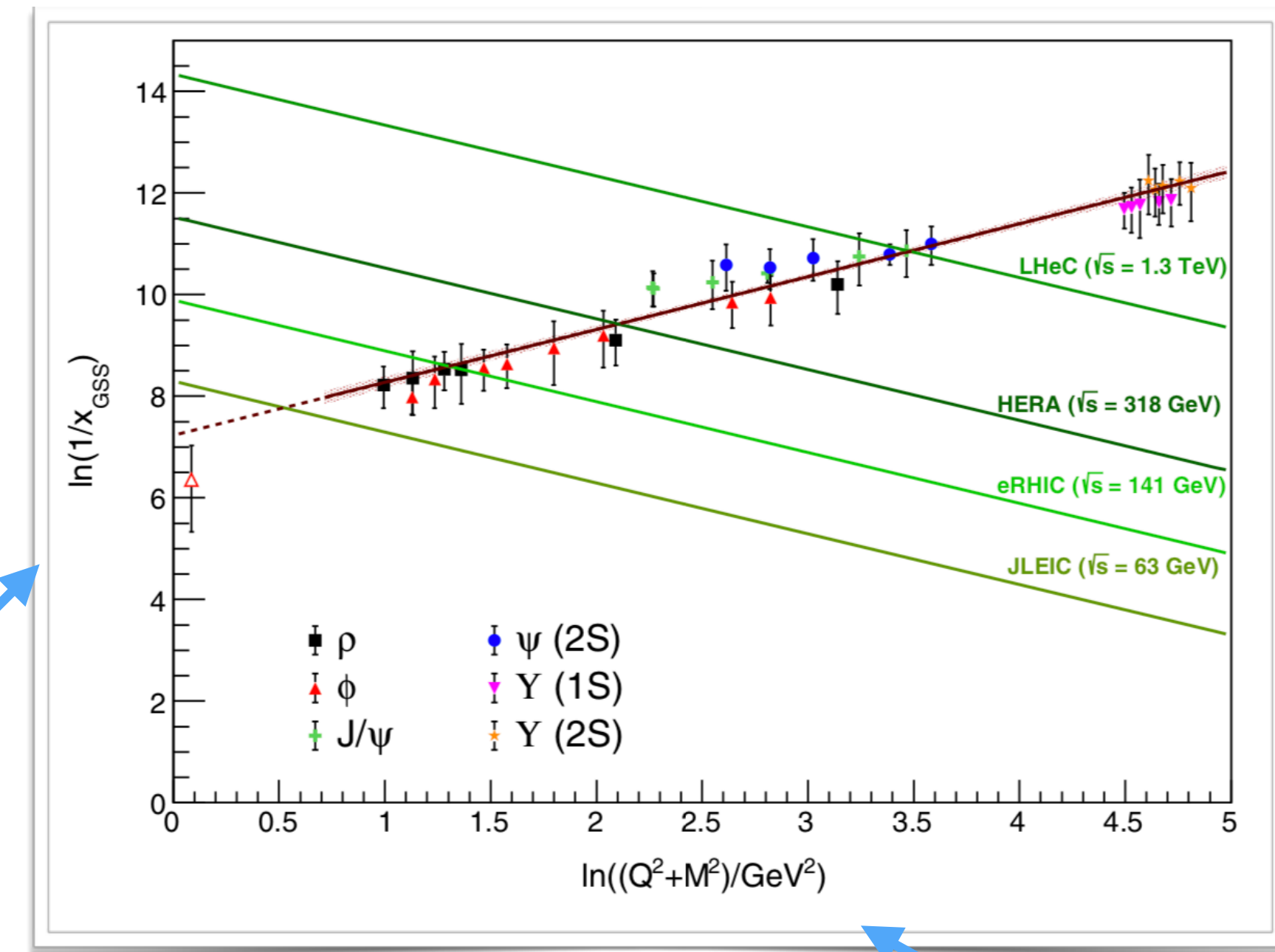
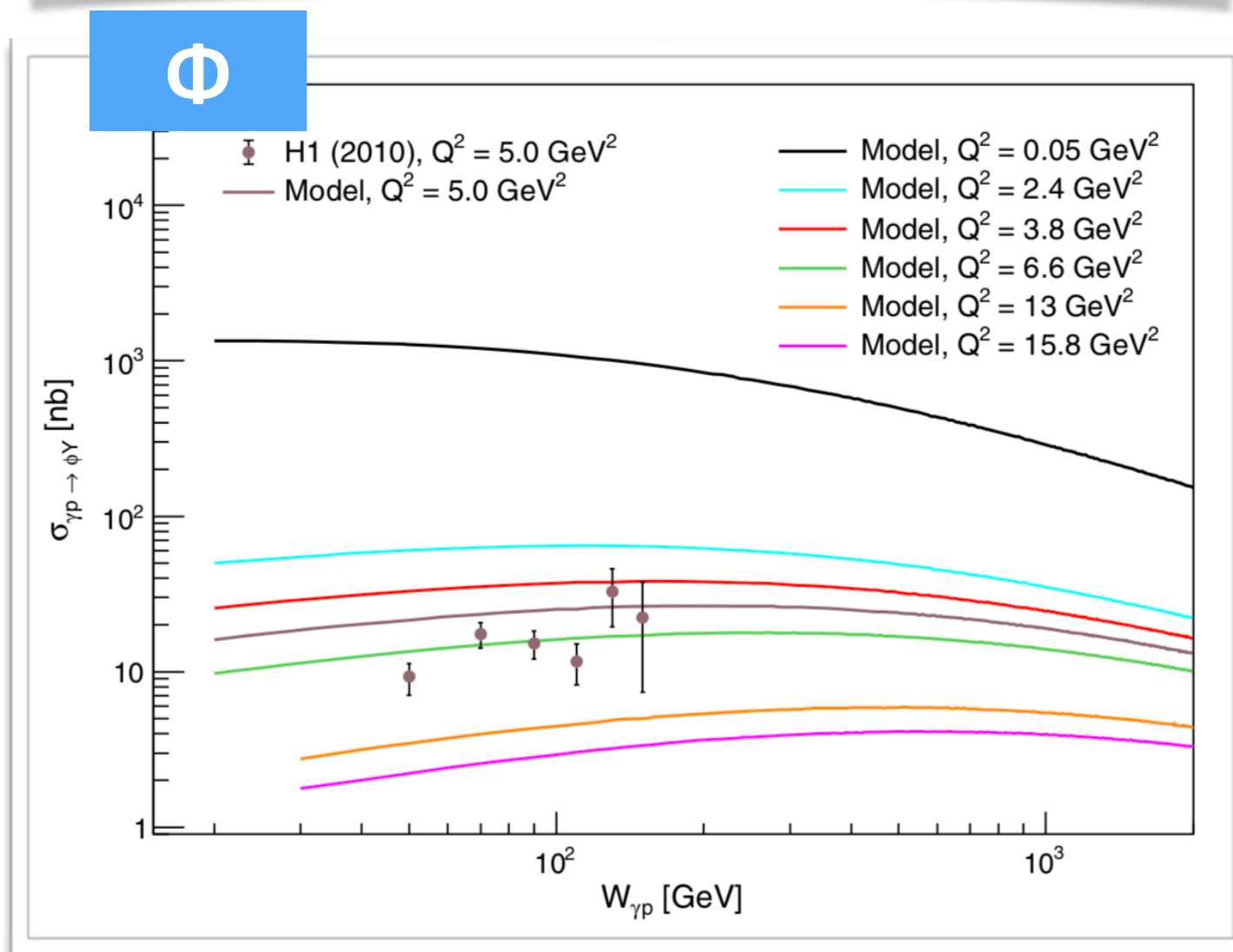
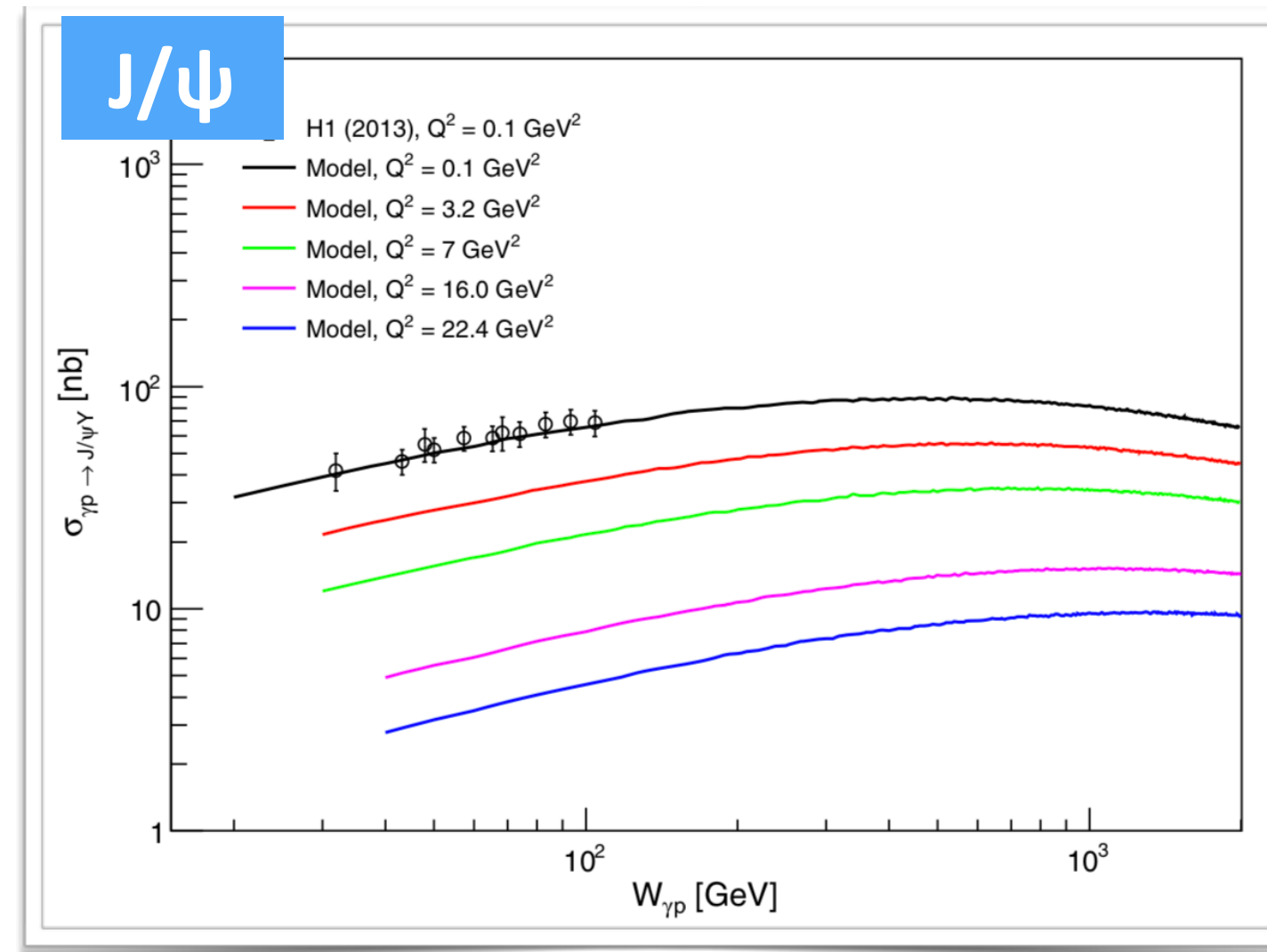
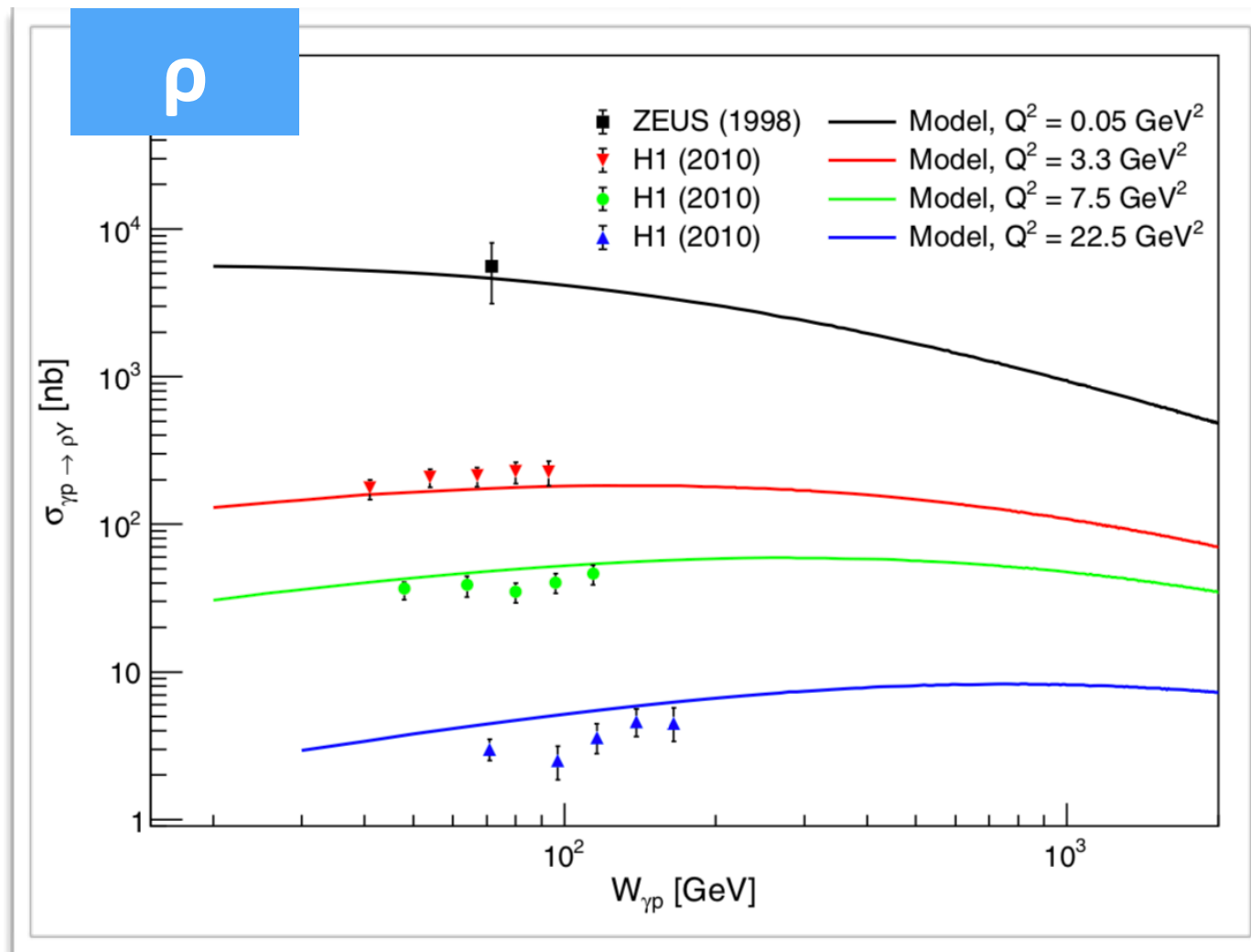


Position of the maximum of the dissociative cross section depends on the scale.

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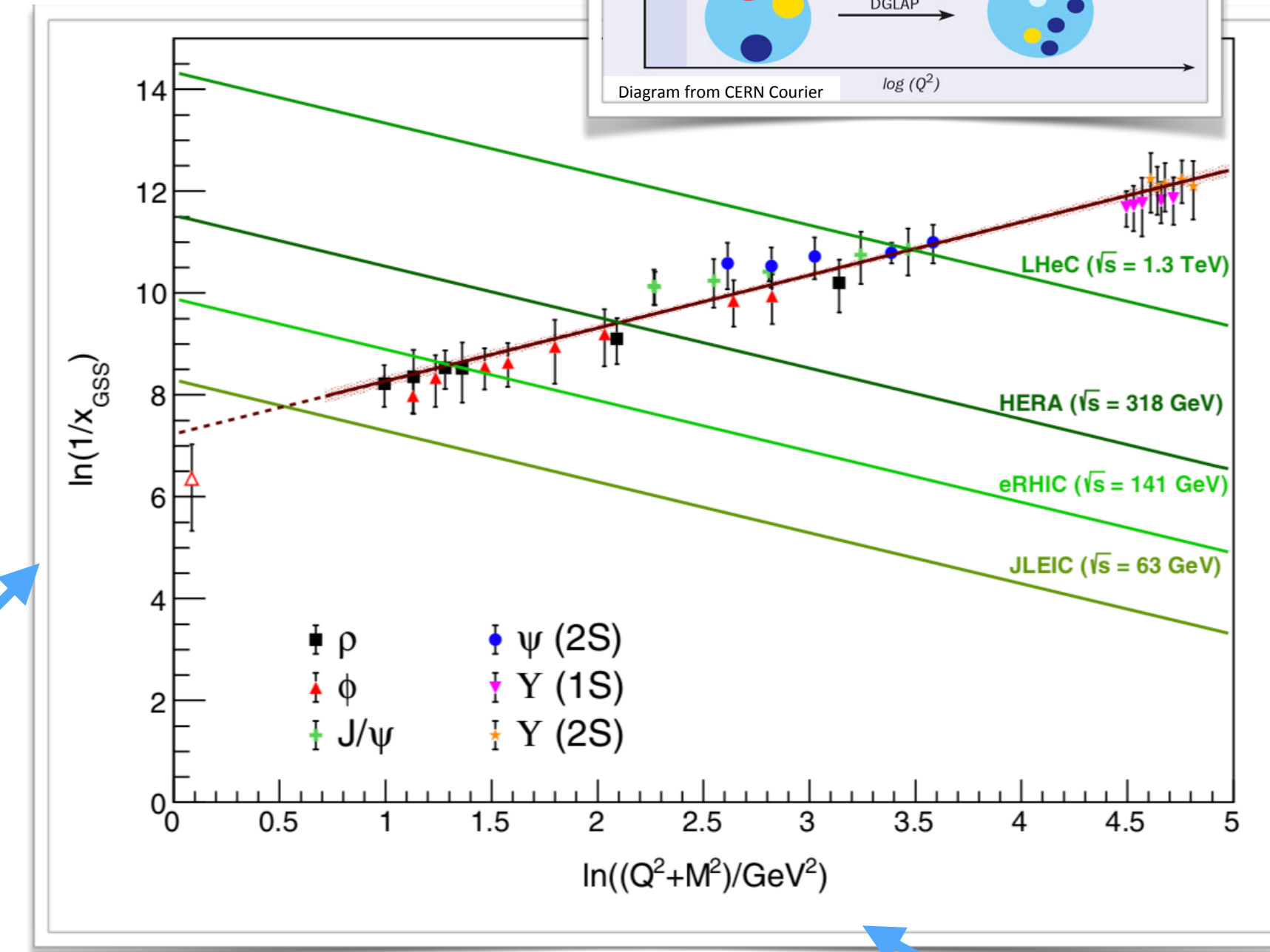
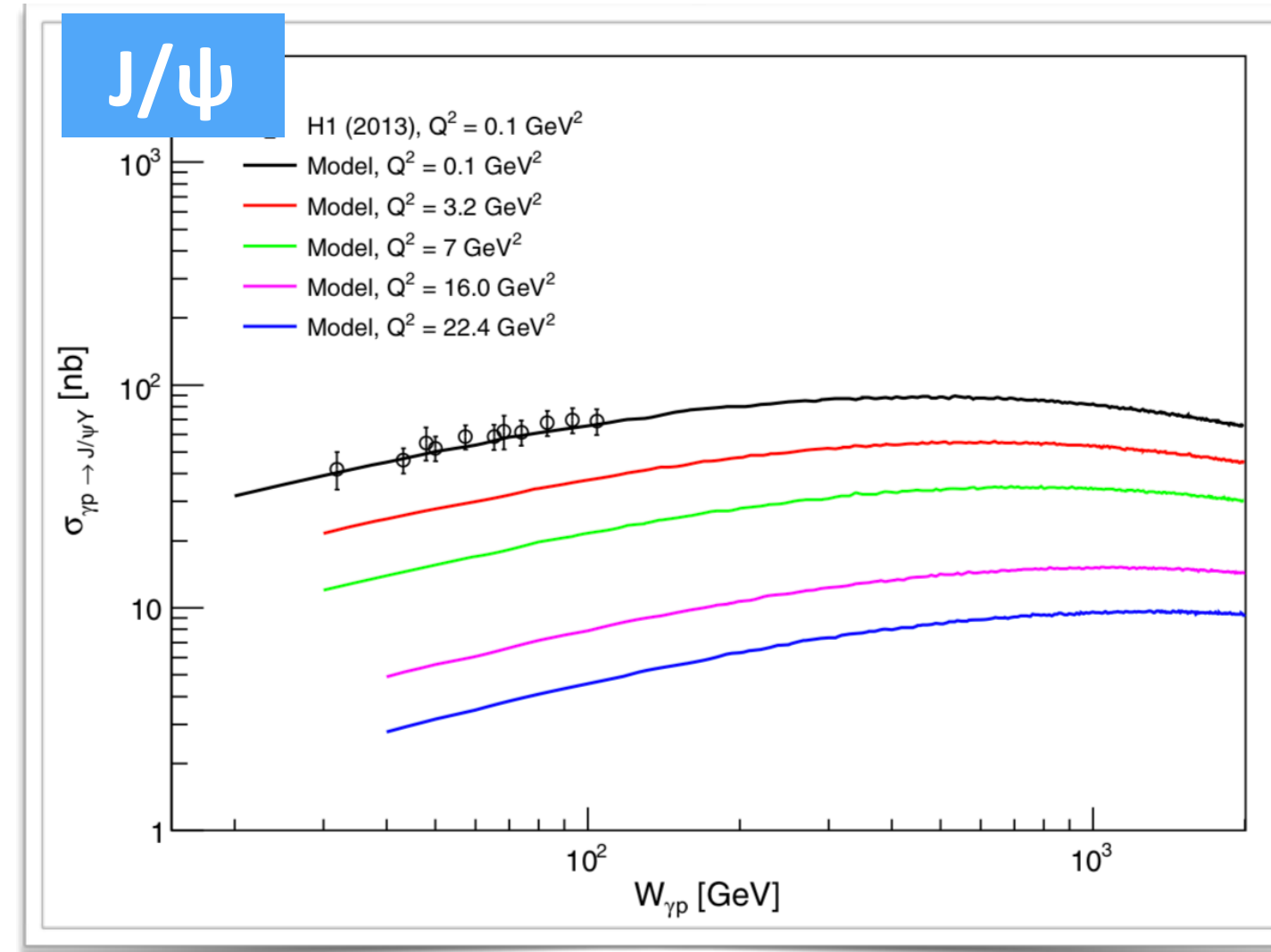
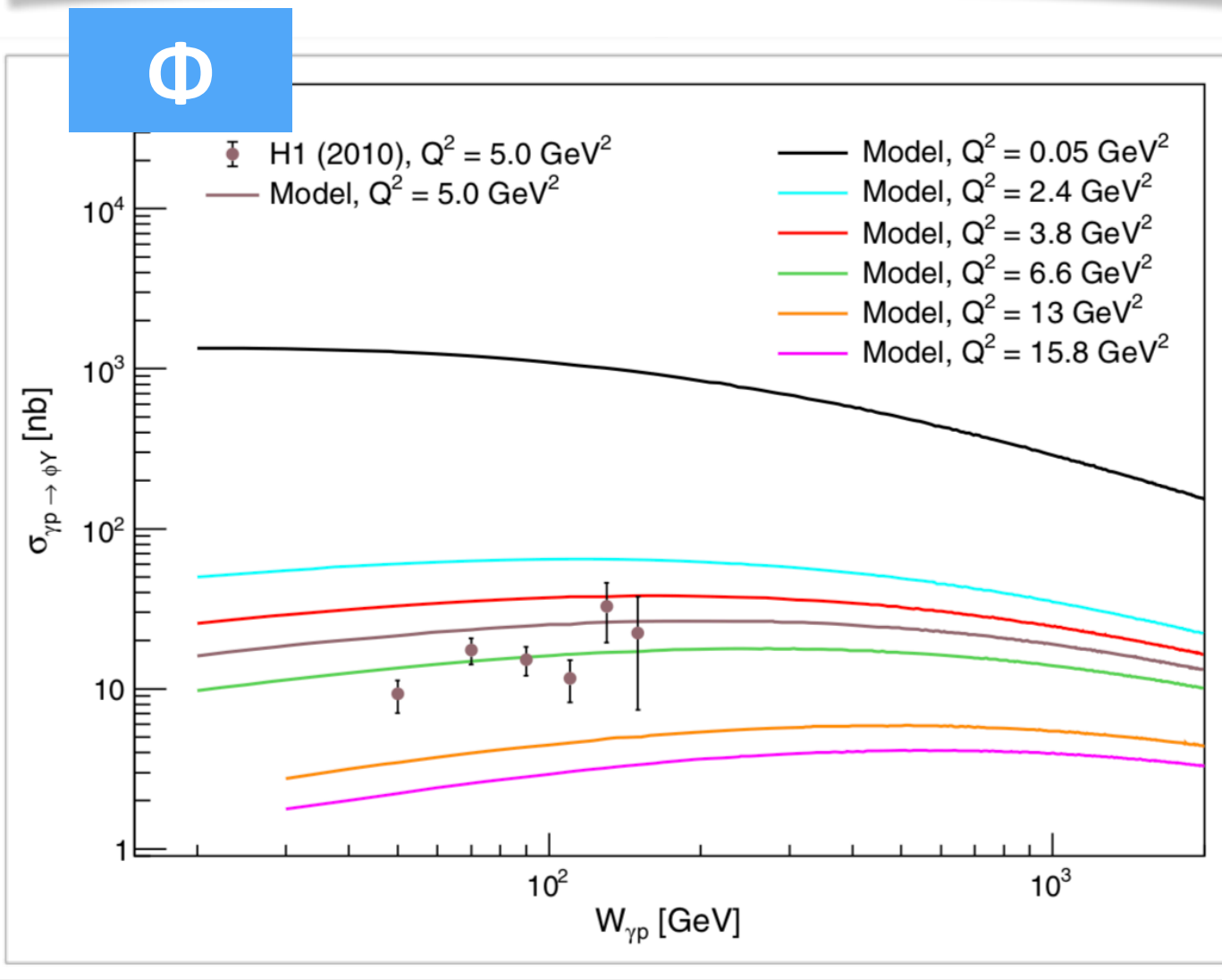
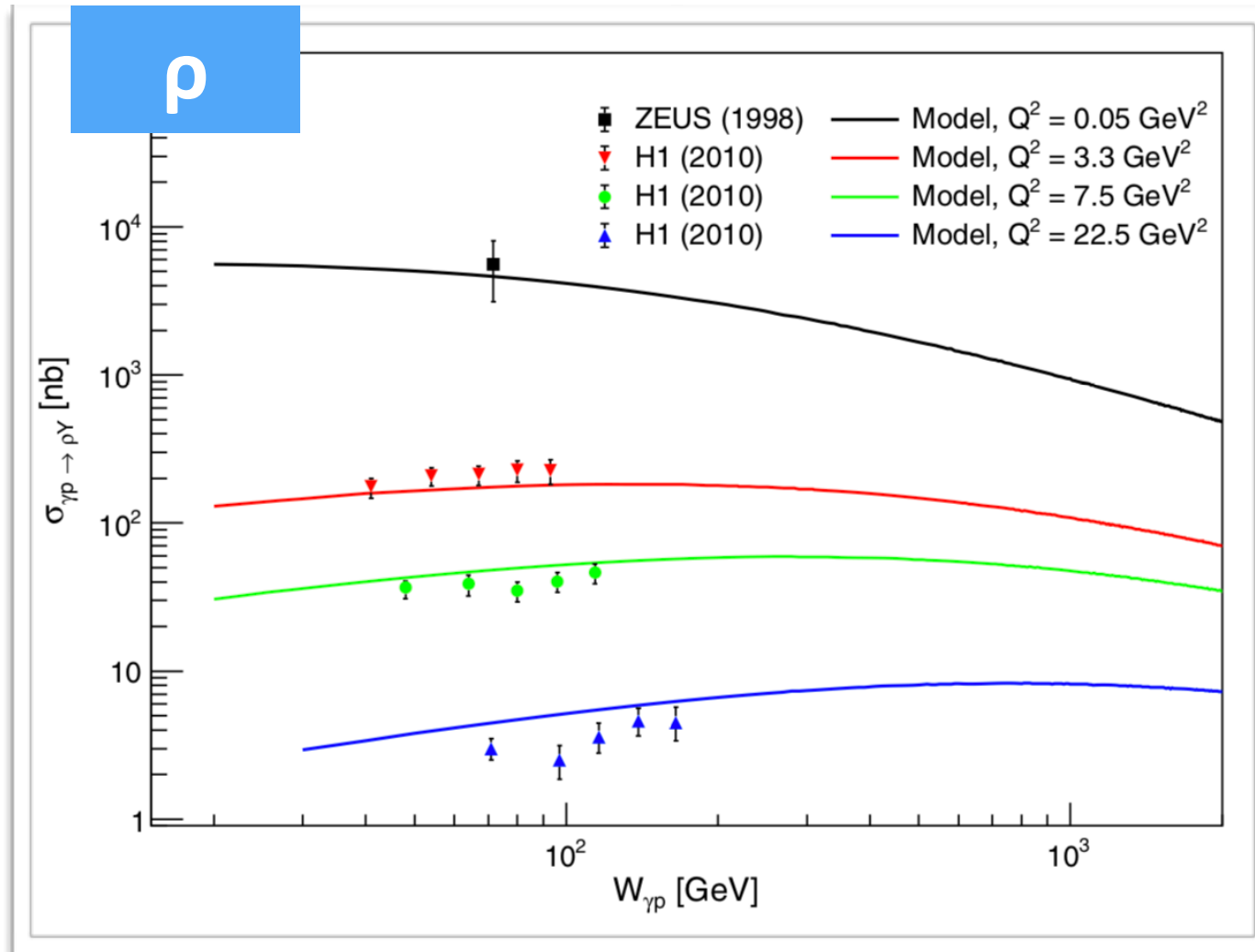


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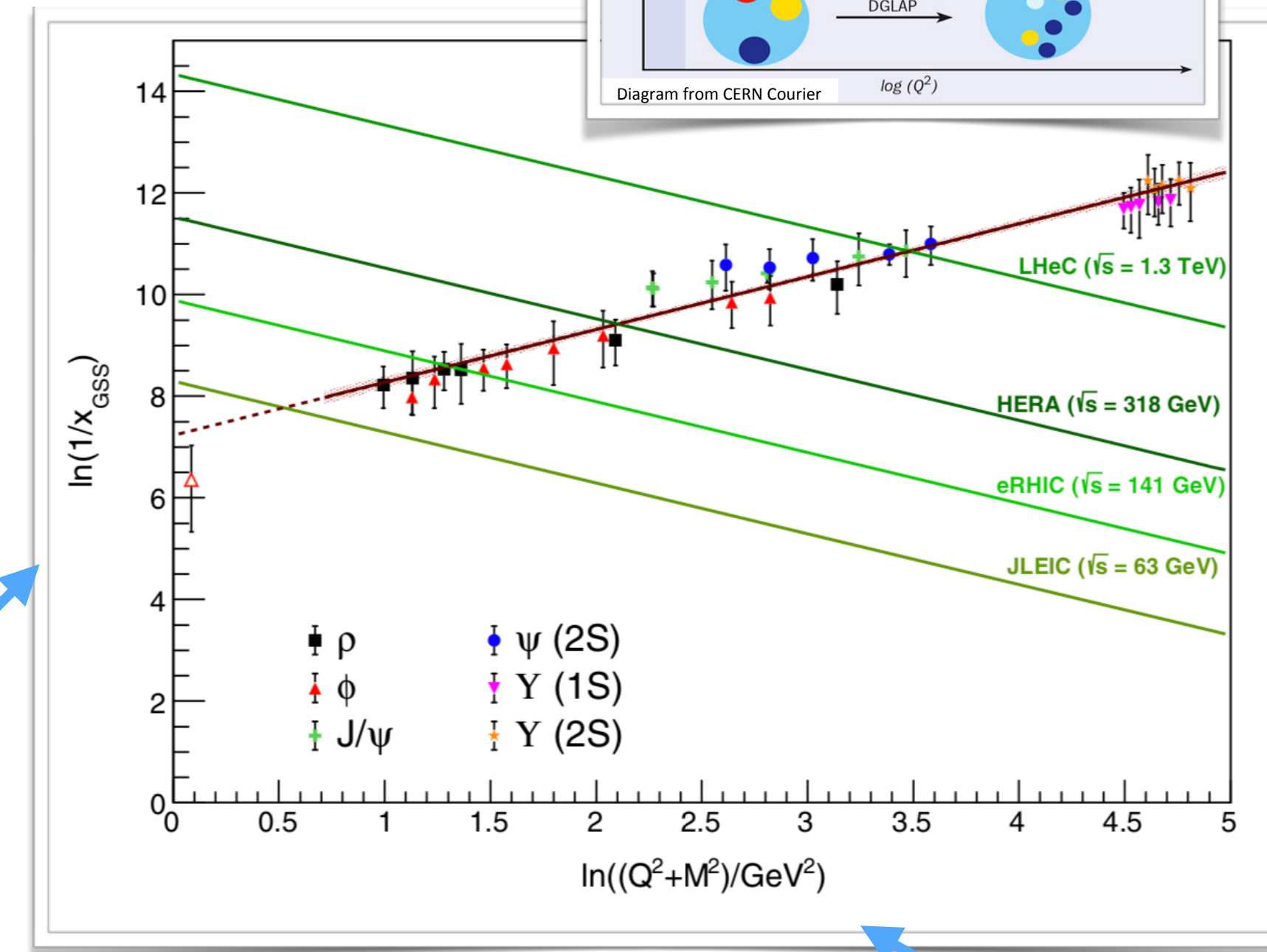
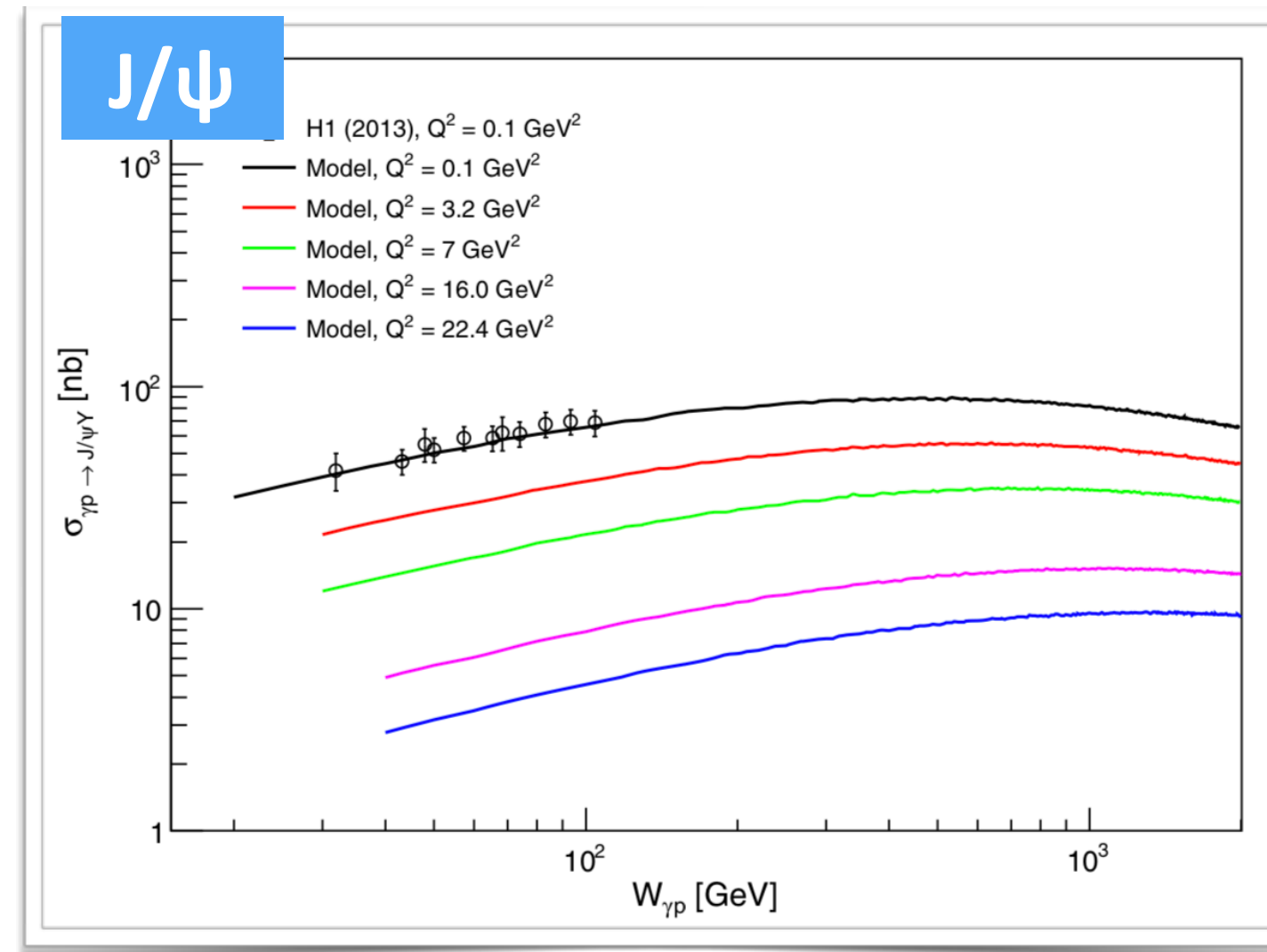
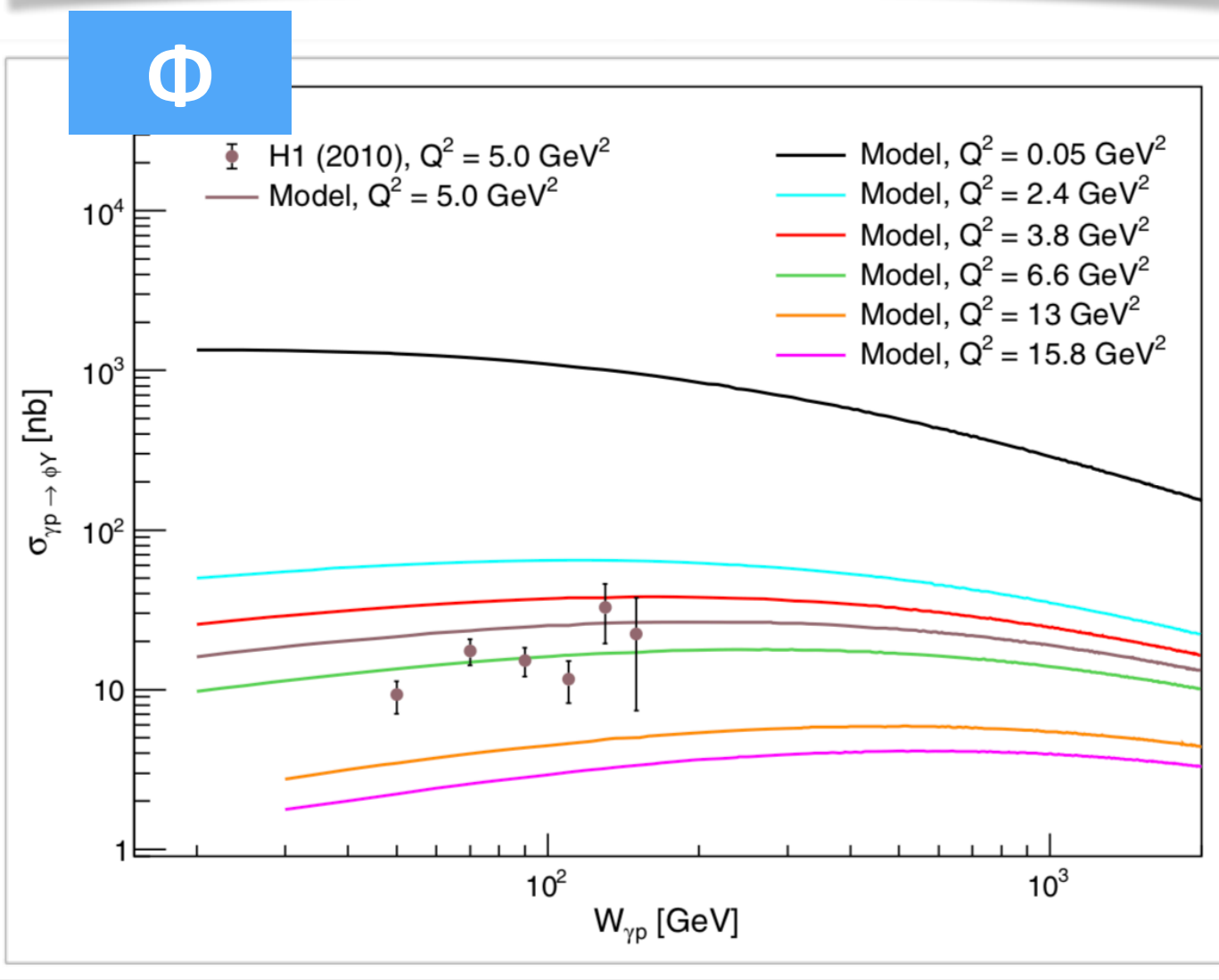
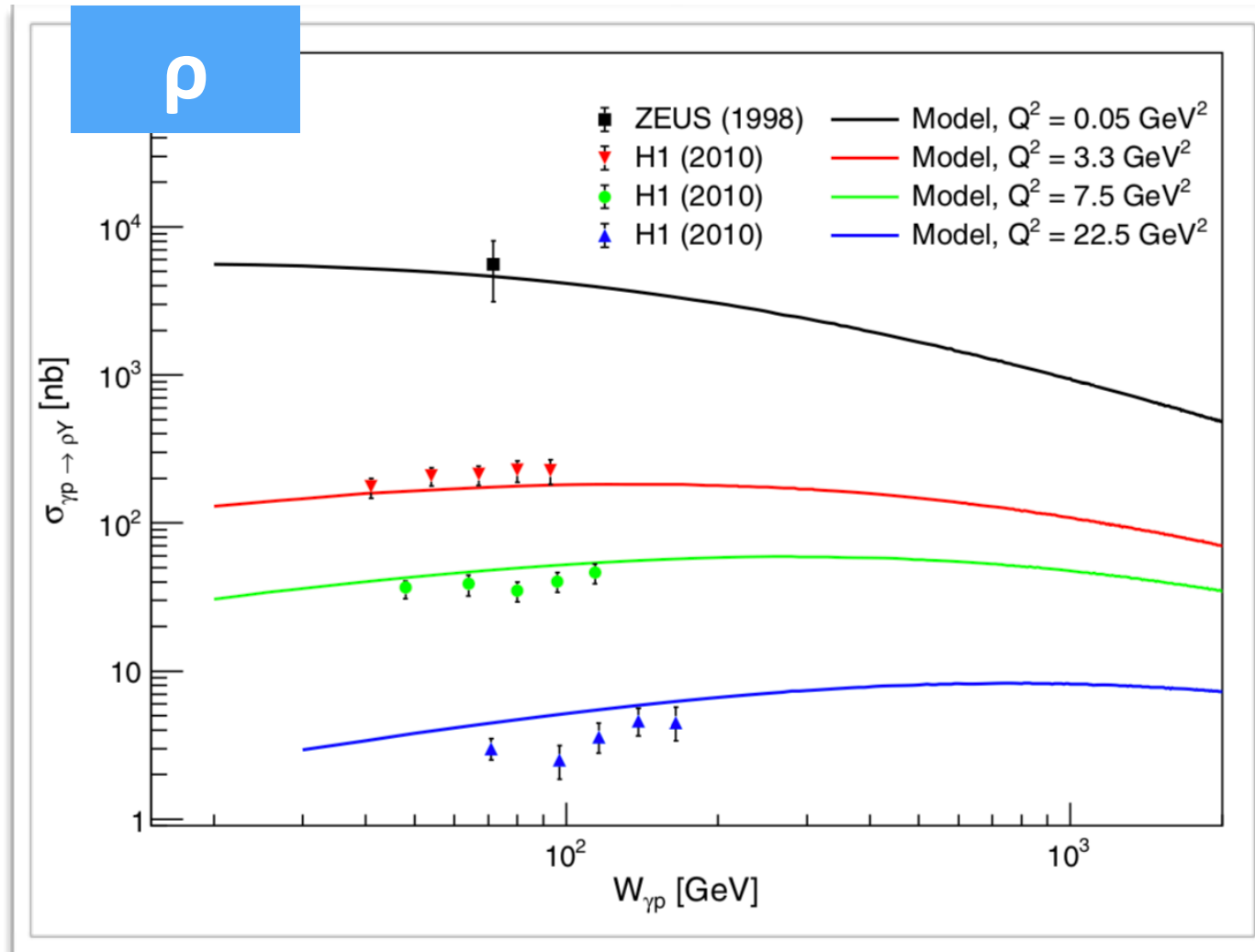


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It will be possible to open this window to saturation at the EIC!

Summary, conclusions and outlook

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Enjoy DIS2019!