Theory Summary For Working Group 6: Spin and 3D Structure

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DIS 2019, Turin Italy

• ≈ 32 theory presentations

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 - PDFs 🖌
 - TMD PDFs (transverse momentum dependent distributions)
 - GPDs (generalized parton distributions)

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- Theoretical Refinements
- Extraction and Pheno ✓

Semi-Inclusive DIS

• Description of large q_T :



- Higher orders? Refine fragmentation function fits?

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Semi-Inclusive DIS



regions of b_T is positive, i.e. $\exp[S_{\text{pert}}(b_*)] > 1$ allowing for an unphysical Sudakov enhancement. In particular in COMPASS-like kinematics, this enhancement dominates over

Drell-Yan

• TMD PDF fits with large amount of data:



• New tools:

repository: https://github.com/VladimirovAlexey/artemide-public

• Many tools now for TMD physics, with high orders in all parts.

Fragmentation Functions

New ways of constraining fragmentation and hadronization dynamics:



Sum rules for quarks into unpolarized hadrons, up to twist-3
 (only thing missing for twist-4: full FF-TMD analysis)
 dependent DiFFs D^{h₁,h₂}/_{q,s'}

$$\begin{split} &\sum_{h \ S_h} \int dz z D_1^h(z) = 1 \\ &\text{Collins-Soper} \end{split} \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h E^h(z) = M_j \\ &\text{NEW} \quad \sum_{h \ S_h} \int dz M_h E^h(z) = M_j \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h H^h(z) = 0 \\ &\sum_{h \ S_h} \int dz M_h \tilde{H}^h(z) = 0 \\ &\sum_{h \ S_h} \int dz M_h \tilde{H}^h(z) = 0 \\ \text{Schaefer-Teryaev} \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 D^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 D^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 G^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 G^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 G^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 G^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\ \text{NEW} \quad &\sum_{h \ S_h} \int dz M_h^2 \tilde{G}^{\perp (1) \ h}(z) = 0 \\$$

1 *q*,s

A. Accardi

Hadronization

 Progress in incorporating polarization in Monte Carlo simulations



Comparison between Pythia + ${}^{3}P_{0}$ and stand alone ${}^{3}P_{0}$:

A different option for the final state in PYTHIA + ${}^{3}P_{0}$



A. Kerbizi

Proton Tensor Charge

• Tension?





Progress in Large TM Sivers

• Polarization dependent results important for fully global TMD pheno program:

$$\begin{split} f_{1T;q\leftarrow h;\mathrm{DY}}^{\perp}(x,\vec{b};\mu,\zeta) &= \pi T(-x,0,x) + \pi a_s(\mu) \Big\{ -2\mathbf{L}_{\mu}P \otimes T + C_F \left(-\mathbf{L}_{\mu}^2 + 2\mathbf{l}_{\zeta}\mathbf{L}_{\mu} + 3\mathbf{L}_{\mu} - \frac{\pi^2}{6} \right) T(-x,0,x) \\ &+ \int d\xi \int_0^1 dy \delta(x-y\xi) \Big[\left(C_F - \frac{C_A}{2} \right) 2\bar{y} \ T(-\xi,0,\xi) + \frac{3y\bar{y}}{2} \frac{G_+(-\xi,0,\xi) + G_-(-\xi,0,\xi)}{\xi} \Big] \Big\} \\ &+ O(a_s^2) + O(\vec{b}^2) \end{split}$$

$$\mu^{2} \frac{d}{d\mu^{2}} T(-x,0,x) = 2a_{s}(\mu)P \otimes T = 2a_{s} \int d\xi \int_{0}^{1} dy \delta(x-y\xi) \left\{ \left(C_{F} - \frac{C_{A}}{2} \right) \left[\left(\frac{1+y^{2}}{1-y} \right)_{+} T(-\xi,0,\xi) + (2y-1)_{+} T(-x,\xi,x-\xi) - \Delta T(-x,\xi,x-\xi) \right] \right. \\ \left. + \frac{C_{A}}{2} \left[\left(\frac{1+y}{1-y} \right)_{+} T(-x,x-\xi,\xi) + \Delta T(-x,x-\xi,\xi) \right] \right. \\ \left. + \frac{1-2y\bar{y}}{4} \frac{G_{+}(-\xi,0,\xi) + Y_{+}(-\xi,0,\xi) + G_{-}(-\xi,0,\xi) + Y_{-}(-\xi,0,\xi)}{\xi} \right\},$$

Tests of Process Dependence

• Gluon Sivers function at RHIC



Exclusive Processes

Nucleon tomography:

Results



H. Moutarde, P. S., J. Wagner "*Border and skewness functions from a leading order fit to DVCS data*" Eur. Phys. J. C78 (2018) 11, 890

Goal: global extraction of Compton Form Factors (CFFs) from DVCS data using LO/LT formalism

Analysis done within PARTONS framework



Model studies are useful to get insight on complex physics phenomena

Model-dependent relations between distributions can be useful

But they should not be extrapolated to different models

S. Rodini

Summary of Summary

• Interesting results and progress with phenomenology.

• Interesting sources of tension.

• Apologies for all talks that I missed!