

Enhanced production of Λ_c in proton-proton collisions at the LHC

Rafał Maciuła

Institute of Nuclear Physics PAN, Kraków, Poland

in collaboration with A. Szczurek
based on Phys.Rev. D98, 014016 (2018)

The 27th Workshop on Deep-Inelastic Scattering and Related Subjects - DIS2019,
8-12 April 2019, Turin, Italy



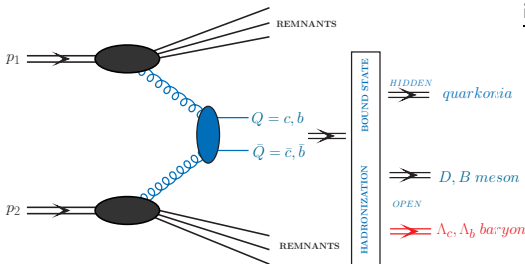
Heavy Flavours at high energies

Heavy-flavours (HF) \Rightarrow open or hidden charm and beauty \Rightarrow study of QCD

- among the most important tools in high-energy pp , pA and AA collisions

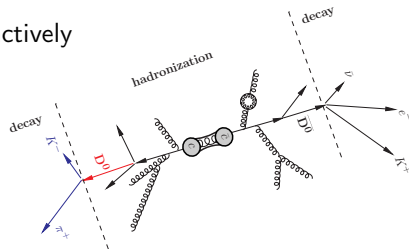
in proton-proton interactions:

- important tests of our understanding of various aspects of QCD quarkonium bound state \Rightarrow non-perturbative aspects of QCD calculations
- heavy quark mass \Rightarrow perturbative QCD
- B and D mesons and/or Λ -baryons \Rightarrow non-perturbative aspects of hadronization in QCD



Production of HFs is one of the most actively studied topics at the LHC

- indirect method** \Rightarrow leptons from semileptonic decays of heavy flavours
- direct method** \Rightarrow open heavy mesons
- Λ_c -baryon \Rightarrow an interesting extension of the standard studies



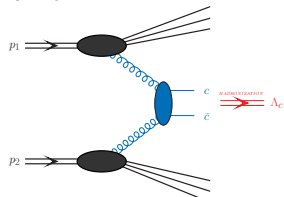
Studies of charm production at the LHC

our last years activity in Heavy Flavours

- model of calculations of charm cross section at high energies:
 k_T -factorization approach \Rightarrow off-shell initial state partons and unintegrated PDFs
 + independent parton fragmentation picture
 - inclusive D meson (transverse momentum and rapidity distributions)
 - $D\bar{D}$ -pair production (more exclusive correlation studies)
 - double and triple charm production (multi-parton interactions)
 - charm meson associated with jets
 - inclusive J/ψ -meson (color evaporation model)
 - nonphotonic leptons (semileptonic decays of charm meson)
 - prompt atmospheric neutrino flux (charm in the Earth's atmosphere)

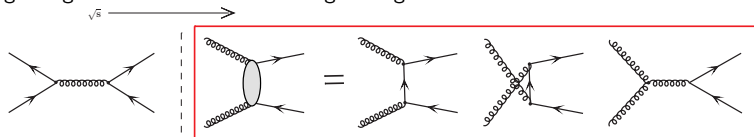
very recently: **First measurements of inclusive Λ_c production**
 in proton-proton scattering at $\sqrt{s} = 7$ TeV

- LHCb data: $2 < y < 4.5$ and $2 < p_T < 8$ GeV
- ALICE data: $|y| < 0.5$ and $1 < p_T < 8$ GeV
- the data sets not fully understood
- reported enhanced production of Λ_c at ALICE and LHCb



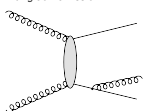
Charm cross section at high energies

- The leading-order (LO) partonic processes for $Q\bar{Q}$ production \Rightarrow gluon-gluon fusion dominant at high energies

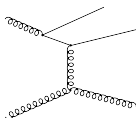


- Main classes of the next-to-leading order (NLO) diagrams:

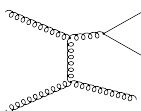
pair creation
with gluon emission



flavour excitation



gluon splitting



the NLO corrections of
a special importance
for charm production!

the observables of the interest \Rightarrow single-particle transverse momentum (inclusive spectra)
correlation observables (less inclusive $\Delta\varphi$, M_{inv} , pair p_T)

collinear approach:

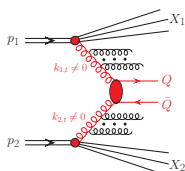
- stat of the art for single particle spectra at NLO (FONLL, GM-VFNS)
- MC@NLO for correlations
- NNLO not available for charm/bottom

k_T -factorization:

- exact kinematics from the very beginning
- correlation observables directly calculable
- some contributions beyond the NLO available (also differentially)



k_T -factorization (high-energy factorization) approach



off-shell initial state partons \Rightarrow

initial transverse momenta explicitly included $k_{1,t}, k_{2,t} \neq 0$

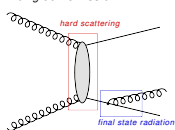
- additional hard dynamics coming from transverse momenta of incident partons (virtualities taken into account)
- very efficient for less inclusive studies of kinematical correlations
- more exclusive observables, e.g. pair transverse momentum or azimuthal angle very sensitive to the incident transverse momenta

multi-differential cross section:

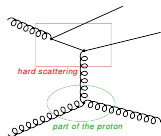
$$\frac{d\sigma}{dy_1 dy_2 d^2 p_{1,t} d^2 p_{2,t}} = \int \frac{d^2 k_{1,t}}{\pi} \frac{d^2 k_{2,t}}{\pi} \frac{1}{16\pi^2 (x_1 x_2 s)^2} \overline{|\mathcal{M}_{g^*g^* \rightarrow Q\bar{Q}}|^2} \times \delta^2(\vec{k}_{1,t} + \vec{k}_{2,t} - \vec{p}_{1,t} - \vec{p}_{2,t}) \mathcal{F}_g(x_1, k_{1,t}^2, \mu) \mathcal{F}_g(x_2, k_{2,t}^2, \mu)$$

- $\mathcal{F}_g(x, k_t^2, \mu)$ - (unintegrated) transverse momentum dependent PDFs
- the LO off-shell matrix elements $\overline{|\mathcal{M}_{g^*g^* \rightarrow Q\bar{Q}}|^2}$ available (analytic form)
- the higher-order matrix elements only at tree-level (KaTie Monte Carlo generator)
- part of **higher-order (real) corrections effectively included in uPDF**

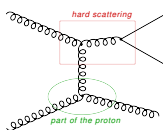
pair creation
with gluon emission



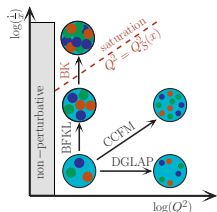
flavour excitation



gluon splitting



Unintegrated parton distribution functions (uPDFs)

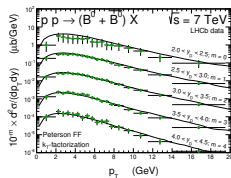
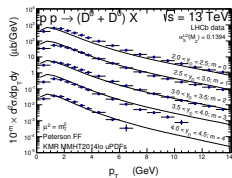


Most popular models:

- H. Jung et al. (CCFM, broad range of x)
- H. Jung et al. (DGLAP + Parton-Branching, broad range of x)
- Kimber-Martin-Ryskin (DGLAP-BFKL, broad range of x)
- Kwieciński-Martin-Staśto (BFKL-DGLAP, rather small x -values)
- Kutak-Staśto, Kutak-Sapeta (BK+saturation, only small x -values)

As a default set: Kimber-Martin-Ryskin (KMR) approach:

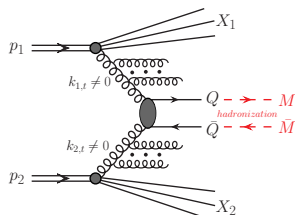
- calculated from collinear PDFs (most up-to-date PDF sets can be used)
- the unique feature: $k_T > \mu$ included as a consequence of angular ordering
 \Rightarrow **hard emissions from the uPDF**
- in fact, only the KMR model with angular ordering effectively includes real higher-order contributions that correspond, e.g. for charm, to associated production with one (NLO) and two (NNLO) jets (or minijets)



k_T -factorization + KMR uPDF works very well for inclusive open charm and bottom mesons at the LHC (including correlation observables)



Heavy quark to open meson and/or baryon transition



- **independent parton fragmentation** picture → phenomenological fragmentation functions from e^+e^- data
- **often used for heavy flavours (scale-independent):** Peterson et al., Braaten et al., Kartvelishvili et al.
- **more up-to-date (scale-dependent):** Kneesch et al. (KKKS08) + DGLAP evolution

as a default: We assume the same fragmentation scenario for Λ_c -baryon as for D -meson

- the standard approach:

$$\frac{d\sigma(y, p_t^M)}{dyd^2p_t^M} \approx \int \frac{D_{Q \rightarrow M}(z)}{z^2} \cdot \frac{d\sigma(y, p_t^Q)}{dyd^2p_t^Q} dz$$

where: $p_t^Q = \frac{p_t^M}{z}$ and $z \in (0, 1)$

approximation: rapidity unchanged in the fragmentation process → $y_Q = y_M$

⇒ rescaling of the transverse momentum at a constant rapidity (angle)

- Peterson FF parameter: $\varepsilon_c^{\Lambda_c} = \varepsilon_c^D = 0.05$ (the same shapes of the FFs)
- fragmentation probabilities from different experiments:
 $f_{c \rightarrow \Lambda_c} = 0.0526 \pm 0.0031$ (e^+e^-), 0.0540 ± 0.0195 (DIS), 0.0639 ± 0.0122 (pp)

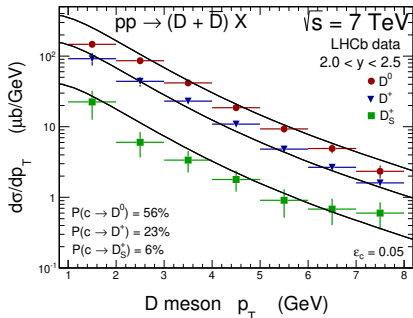
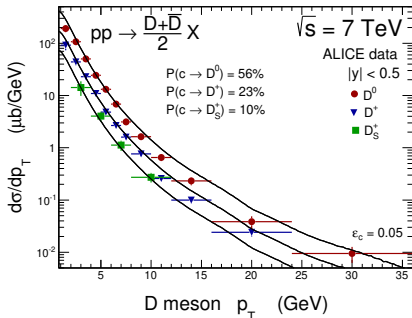


D-meson

ALICE and LHCb data

Theoretical computations:

k_T -factorization: $g^* g^* \rightarrow c\bar{c} + \text{KMR uPDF} + \text{Peterson FF for } c \rightarrow D \text{ transition}$



- very good description of the data for the D^0 and D^+ meson with the $f_{c \rightarrow D}$ known from the combined analysis of charm-quark fragmentation fraction measurements in the former e^+e^- , ep and pp experiments
- some excess for the D_s -meson
 \Rightarrow we need larger $f_{c \rightarrow D_s}$ for the ALICE data than in the case of the LHCb
- we cannot describe both sets of data with the same $f_{c \rightarrow D_s}$, but both are consistent with those extracted from other experiments (large uncertainties)

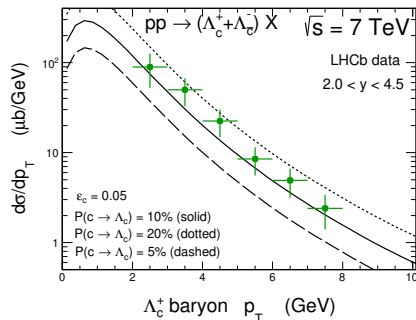
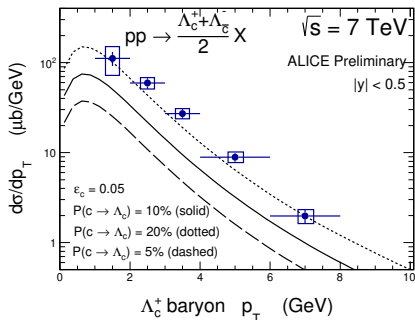


Λ_c -baryon

ALICE and LHCb data

Theoretical computations:

k_T -factorization: $g^* g^* \rightarrow c\bar{c} + \text{KMR uPDF} + \text{Peterson FF for } c \rightarrow \Lambda \text{ transition}$



- the $f_{c \rightarrow \Lambda_c} = 0.05$ (typical pre-LHC value) clearly underpredicts both data sets
- the agreement between data and predictions with the increased $f_{c \rightarrow \Lambda_c}$ is better
- a visible difference appears in the observed agreement for the midrapidity ALICE and forward LHCb regimes
- the two types of the enhancement: the LHC to the pre-LHC results
the ALICE to the LHCb results

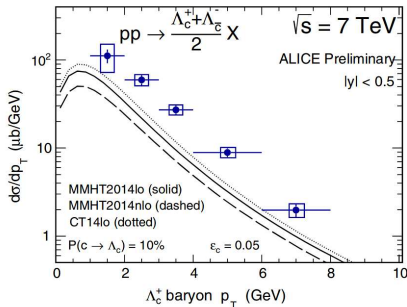
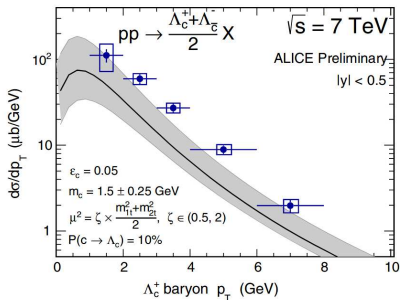


Λ_c -baryon

ALICE and LHCb data

Theoretical computations:

k_T -factorization: $g^* g^* \rightarrow c\bar{c} + \text{KMR uPDF} + \text{Peterson FF for } c \rightarrow \Lambda$ transition



- the uncertainties related to the choice of the perturbative calculation parameters and due to the choice of the gluon collinear PDFs in the calculation of the KMR uPDF
- the upper limit of the uncertainties almost agrees with the ALICE experimental data
- however, this would not be the case for the LHCb data, where already the central value is consistent with the data
- the set of parameters which describes the ALICE data for Λ_c production would also lead to an overestimation of D -meson data
- still, at least the $f_{c \rightarrow \Lambda_c} = 0.1$ is needed

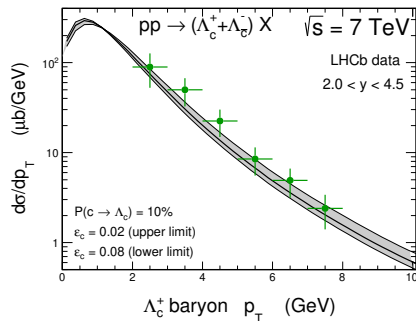
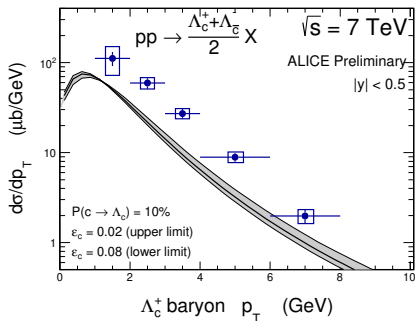


Λ_c -baryon

ALICE and LHCb data

Theoretical computations:

k_T -factorization: $g^* g^* \rightarrow c\bar{c} + \text{KMR uPDF} + \text{Peterson FF for } c \rightarrow \Lambda \text{ transition}$



- Can the situation somehow change when the shape of the $c \rightarrow \Lambda_c$ fragmentation function is different?
- we modify the ϵ parameter for the $c \rightarrow \Lambda_c$ case
- some uncertainties appear but rather small and do not change the overall picture

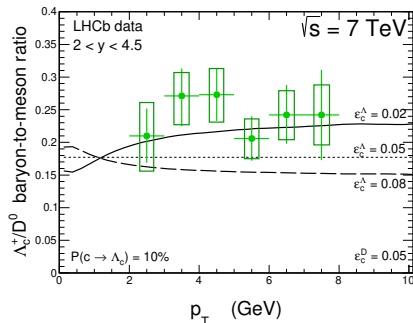
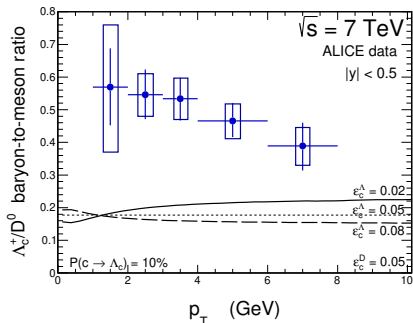


Λ_c/D^0 baryon-to-meson ratio

ALICE and LHCb data

Theoretical computations:

k_T -factorization: $g^*g^* \rightarrow c\bar{c} + \text{KMR uPDF} + \text{Peterson FF for } c \rightarrow \Lambda$ transition



- transverse momentum dependence of the Λ_c/D^0 baryon-to-meson ratio experimentally observed
- from our default set \Rightarrow naturally flat distribution
- some dependence may appear when different fragmentation functions are used for the $c \rightarrow \Lambda_c$ and $c \rightarrow D$ (different ϵ -parameters)
- softer (peaked at intermediate- z) FF for the $c \rightarrow \Lambda_c$ than for the $c \rightarrow D$ suggested by the data

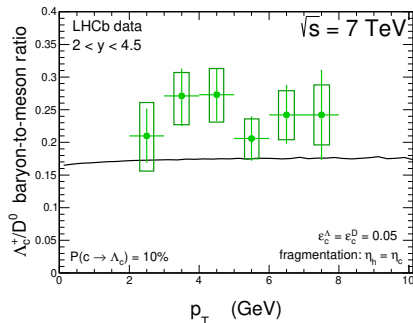
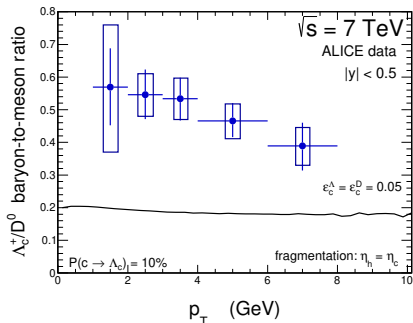


Λ_c/D^0 baryon-to-meson ratio

ALICE and LHCb data

Theoretical computations:

k_T -factorization: $g^*g^* \rightarrow c\bar{c} + \text{KMR uPDF} + \text{Peterson FF for } c \rightarrow \Lambda \text{ transition}$

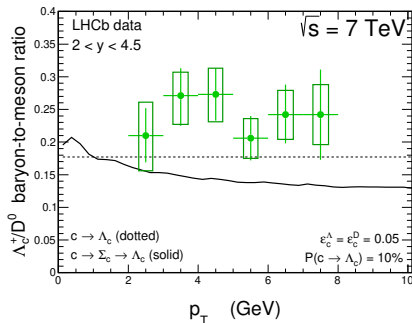
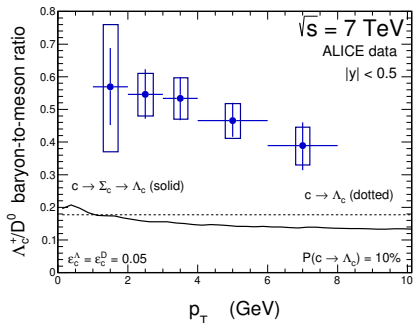


- we follow the prescription: $\eta_h = \eta_c$, i.e., we assume that the charmed hadron is emitted in the same direction (in the pp cms frame) as the quark/antiquark
- here the energies are rescaled and the masses of the incident parton and final hadron are implicitly taken into account
- some p_T -dependence obtained but almost negligible
- works in the right direction (small- p_T enhanced for ALICE, and lowered for LHCb) but the effect is definitely too small



Λ_c/D^0 baryon-to-meson ratio

ALICE and LHCb data

Theoretical computations: k_T -factorization: $g^* g^* \rightarrow c\bar{c} + \text{KMR uPDF} + \text{Peterson FF for } c \rightarrow \Lambda \text{ transition}$ 

- in principle, Λ_c -baryons do not need to be exclusively directly produced
- they may come from a feed-down mechanism from excited baryonic states (examples of interest are Σ_c baryons both for $J = 1/2$ ($\Sigma_c(2455)$) and $J = 3/2$ ($\Sigma_c(2520)$))
- the feed-down mechanism could modify transverse momentum distributions.
- this mechanism leads to a small enhancement of the Λ_c/D^0 ratio at small transverse momenta, but it also causes it to decrease at larger p_T 's



Conclusions

- we found that the fragmentation fraction $f_{c \rightarrow \Lambda_c} = 0.1 - 0.15$ describes the recent data from the LHCb but fails to describe the new ALICE data.
- even for LHCb, this number is slightly bigger than the values from the compilation of world results obtained from experimental data on e^+e^- , ep , pp and B meson decays
- although we could agree with the ALICE data using the standard estimation of model uncertainties related to the factorization/renormalization scale, quark mass, and PDF, we were not able to describe simultaneously the ALICE and LHCb Λ_c -baryon data as well as data on D -meson production with the same set of parameters.
- The new data from the ALICE suggests a much bigger $f_{c \rightarrow \Lambda_c}$ hadronization fraction than those obtained in other processes and in the recent LHCb measurements.
- The interpretation of the increased fragmentation fraction $f_{c \rightarrow \Lambda_c}$ is at present not clear and requires further studies, both on the theoretical and experimental sides

Thank You for attention!

