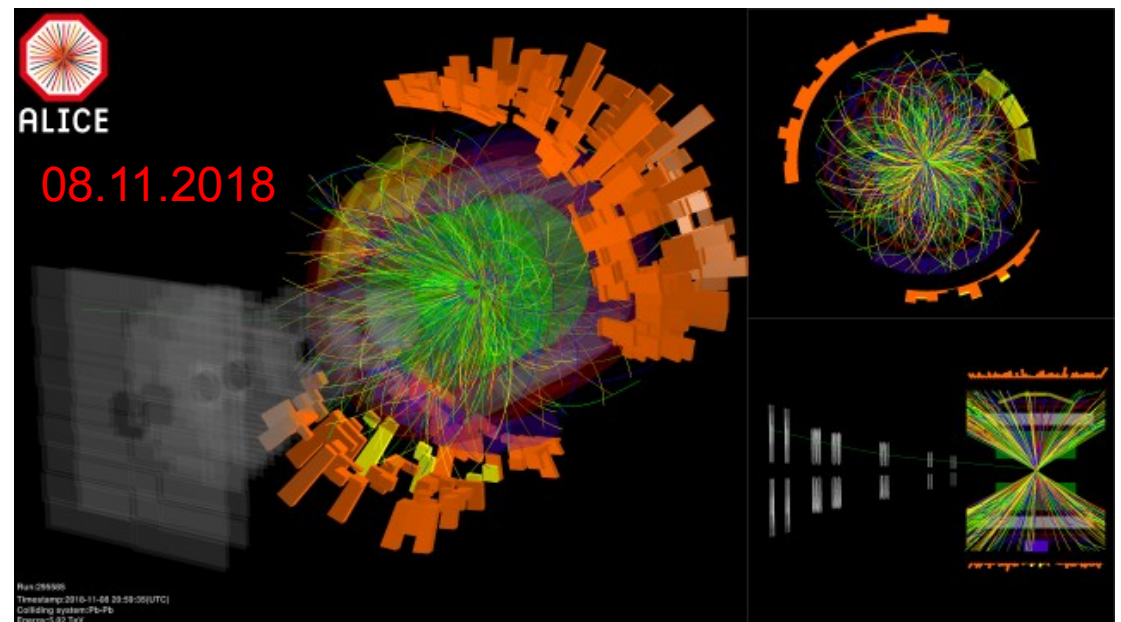


# Probing the Quark Gluon Plasma with anisotropic flow measurements in ALICE

A. Dobrin (CERN & ISS)  
for the ALICE Collaboration

ALICE papers

Phys. Lett. B 784 (2018) 82  
JHEP 07 (2018) 103  
JHEP 09 (2018) 006  
arXiv:1809.09371

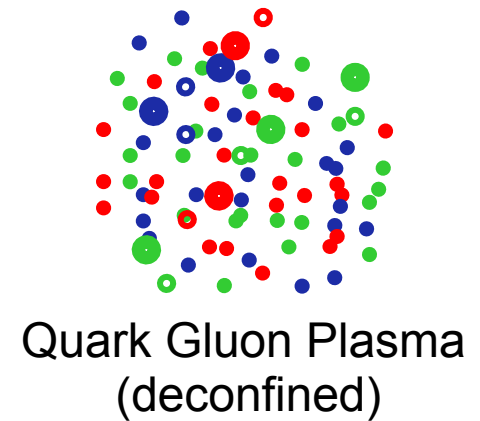
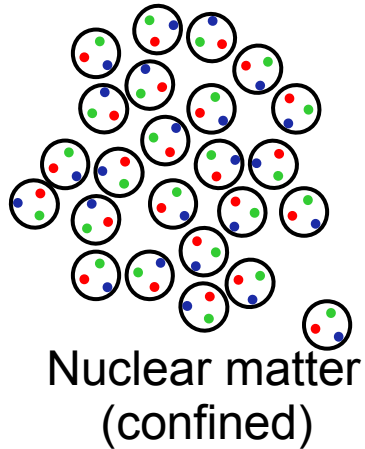
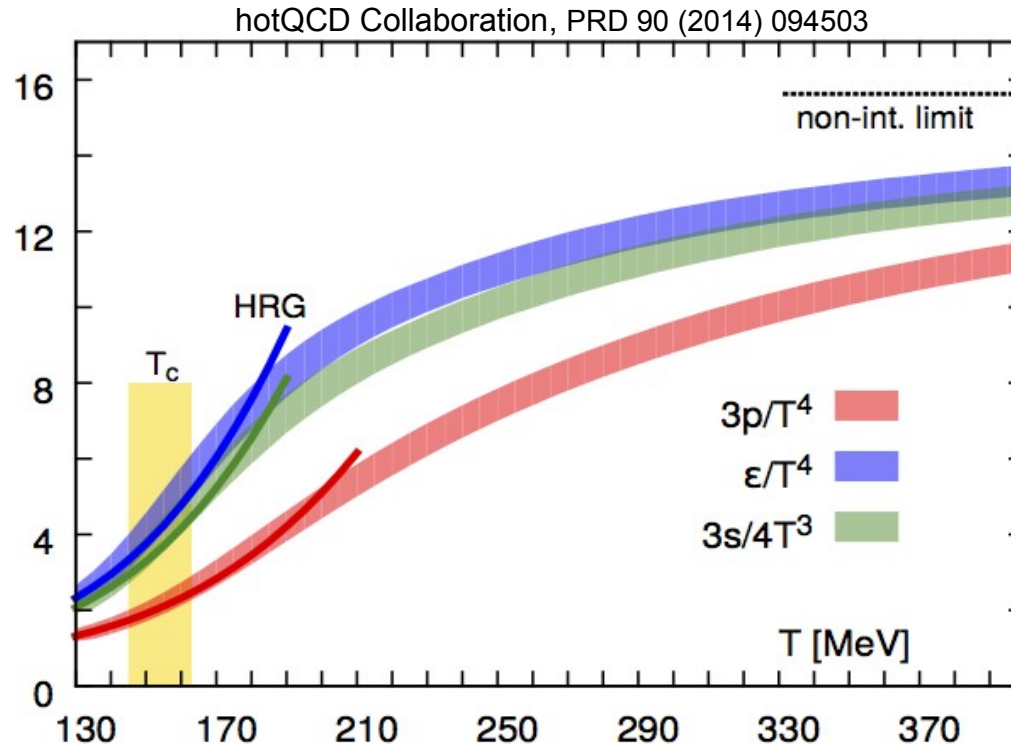




**ALICE**

# Introduction

# Lattice QCD



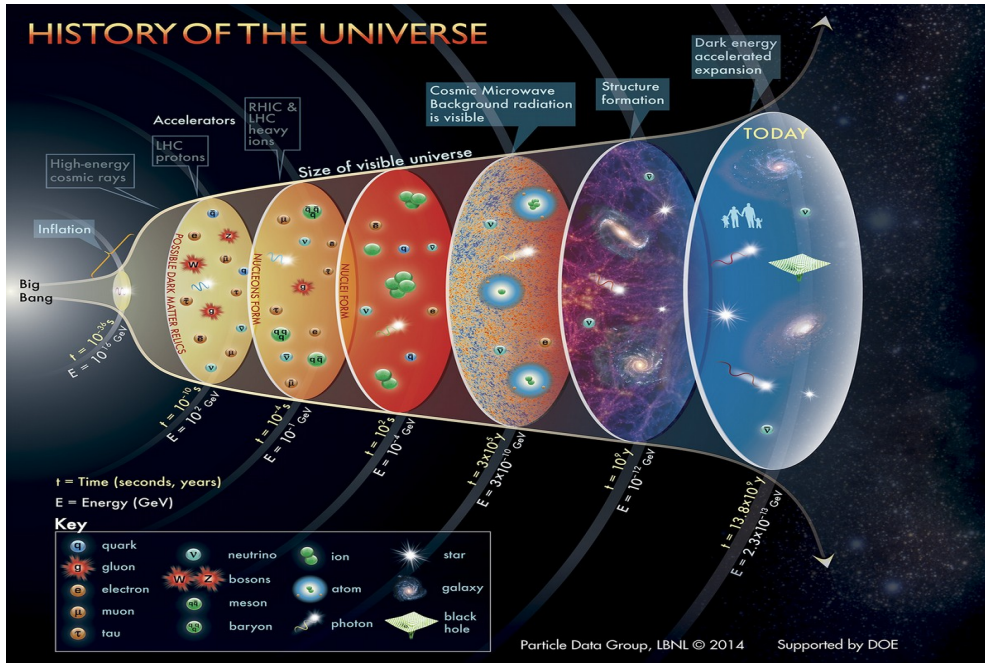
$$p = \frac{1}{3} \epsilon = g \frac{\pi^2}{90} T^4$$

$$\frac{\epsilon}{T^4} = g \frac{\pi^2}{30}$$

- Present understanding of QCD:

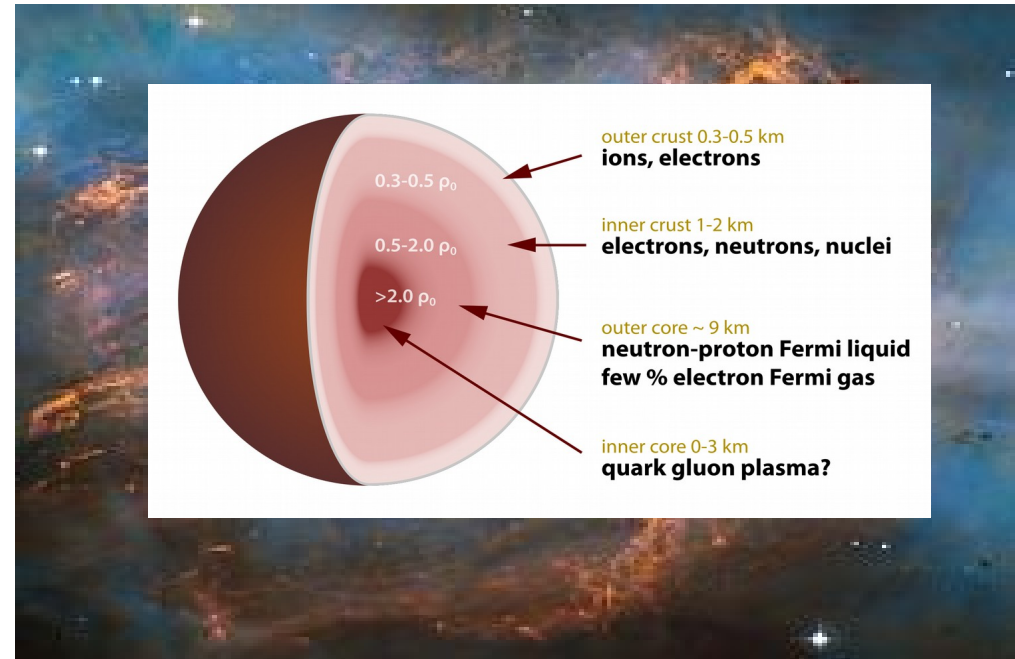
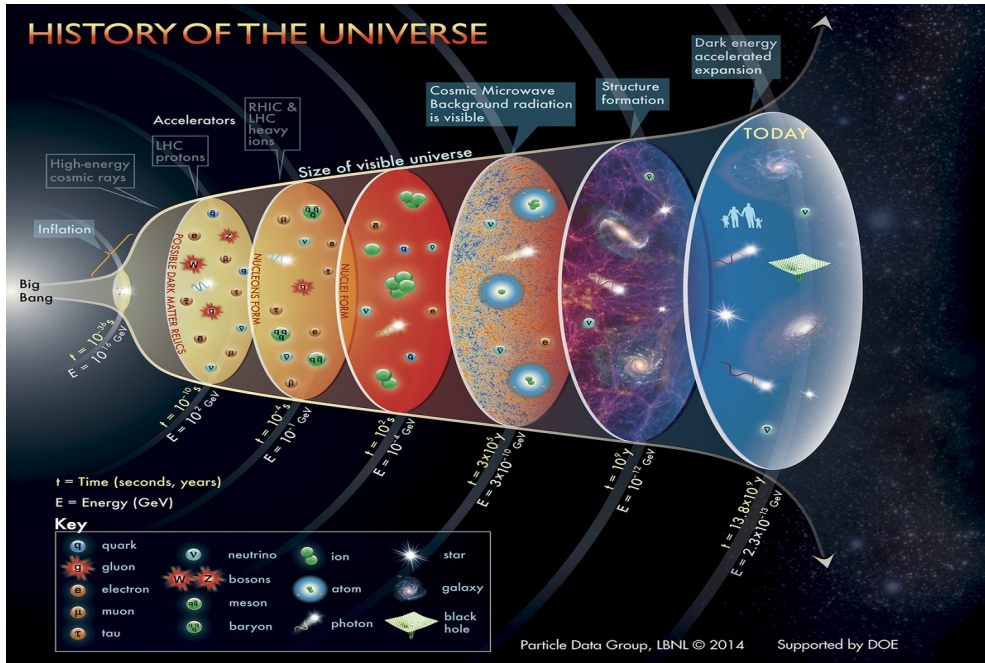
- Heating + compression → Quark Gluon Plasma (QGP): deconfined system of quarks and gluons
- Lattice QCD: transition expected to occur at critical energy density  $\epsilon_c \sim 0.5 \text{ GeV}/\text{fm}^3$  and temperature  $T_c \sim 156 \text{ MeV}$ 
  - Free quark gluon gas limit not reached → residual interactions

# Why study QGP?



QGP might have existed in the expanding Universe in the first  $\mu\text{s}$  after the Big Bang

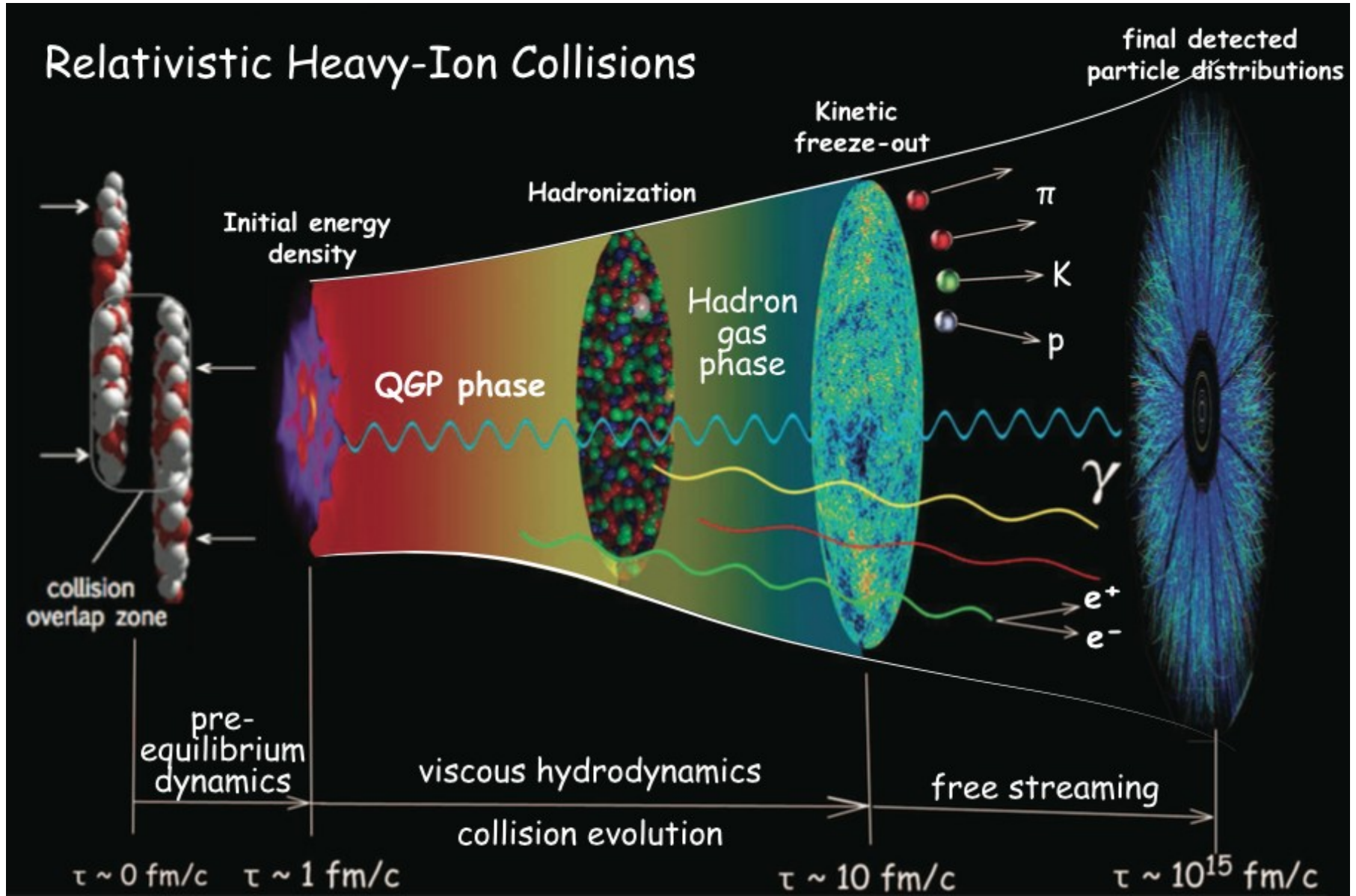
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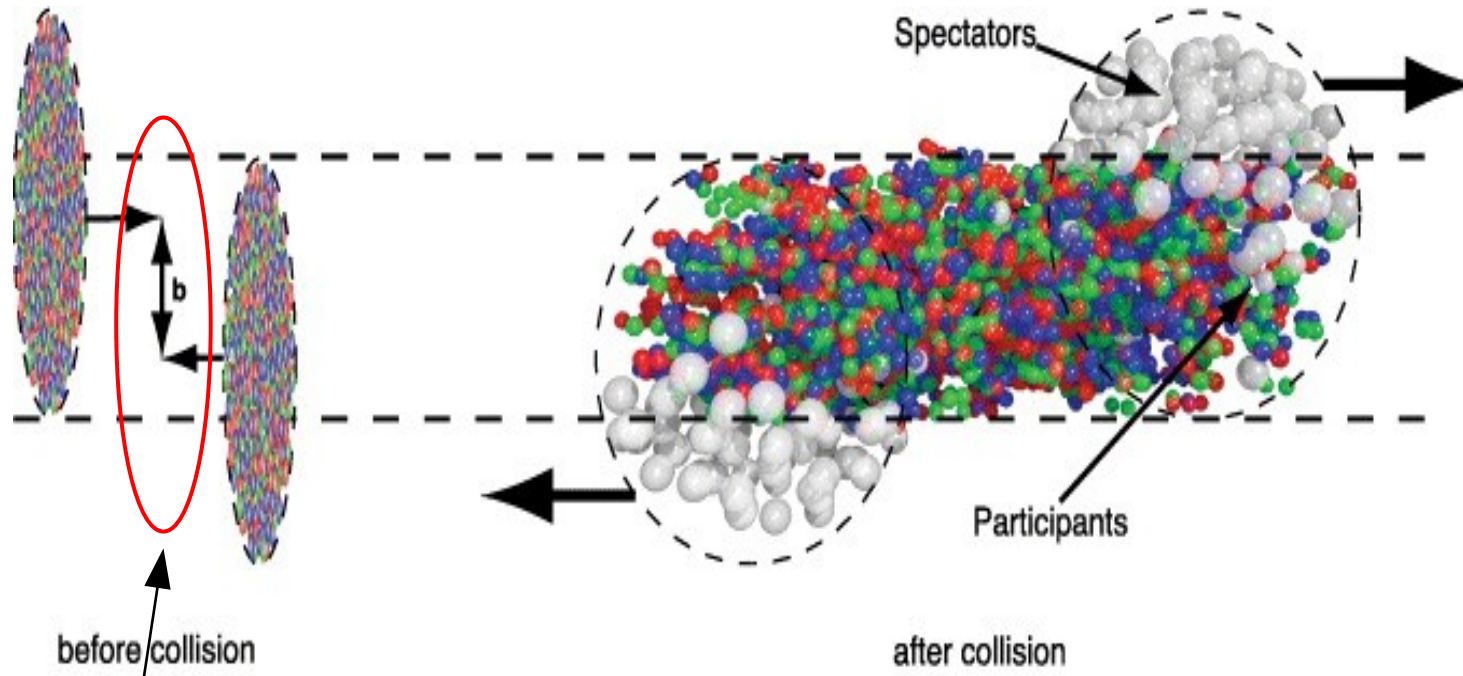
Neutron stars: a more likely place for QGP to exist  $\rightarrow$  mass controlled by the equation of state (EoS) of nuclear matter

# Time evolution of a heavy-ion collision



P. Sorensen / C. Shen

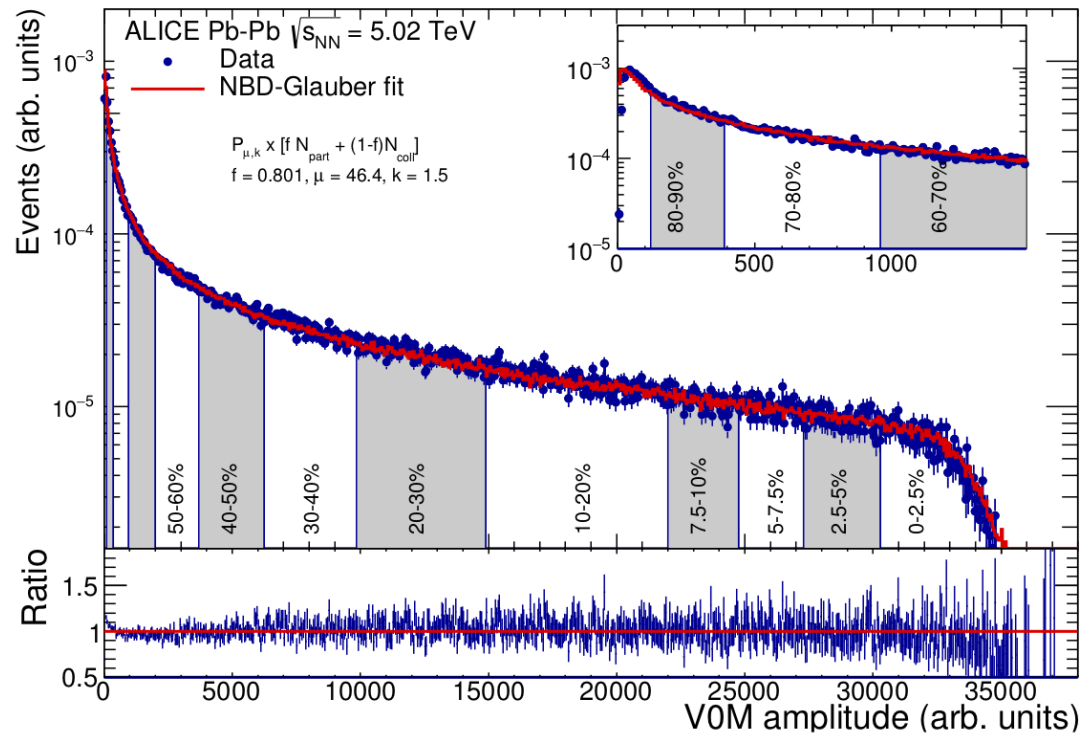
# Centrality: controlling volume and shape



## Impact parameter $b$

- Perpendicular to beam direction
- Connects centers of colliding nuclei
- Not measured directly → estimated by centrality

# Centrality: controlling volume and shape



## Impact parameter $b$

- Perpendicular to beam direction
- Connects centers of colliding nuclei
- Not measured directly → estimated by centrality
- Centrality determined from particle multiplicities

• Most central: 0-5% centrality

• Peripheral: 70-80% centrality

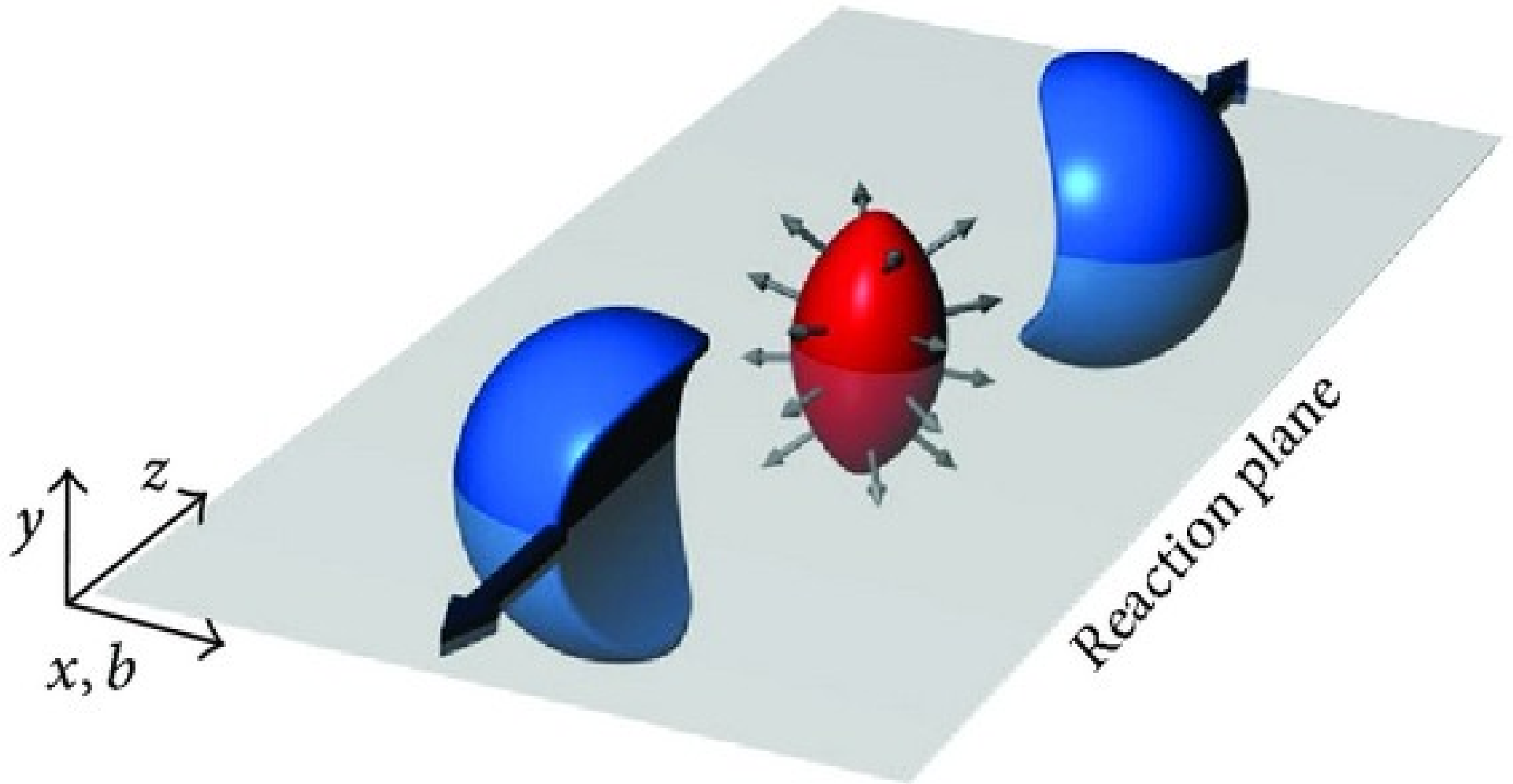




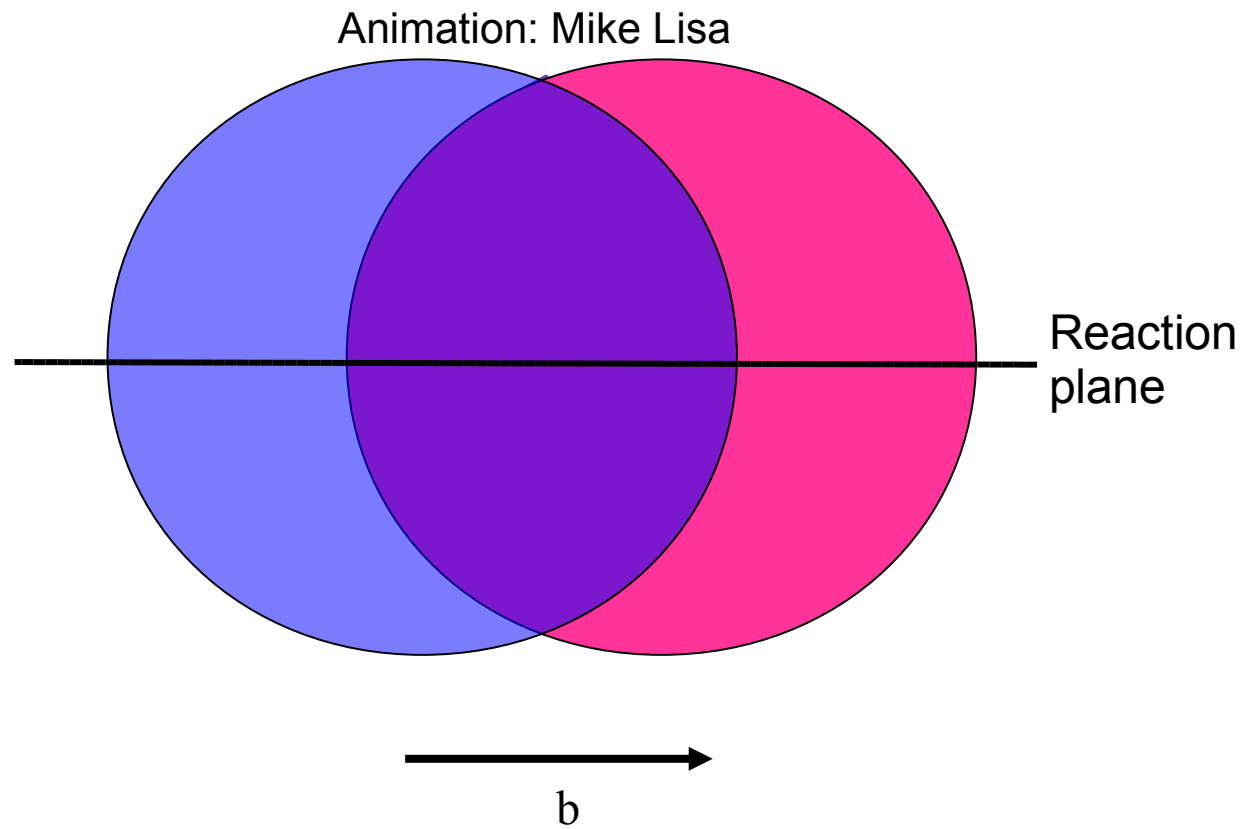
ALICE

# Anisotropic flow

# Elliptic flow



# Elliptic flow



# Elliptic flow

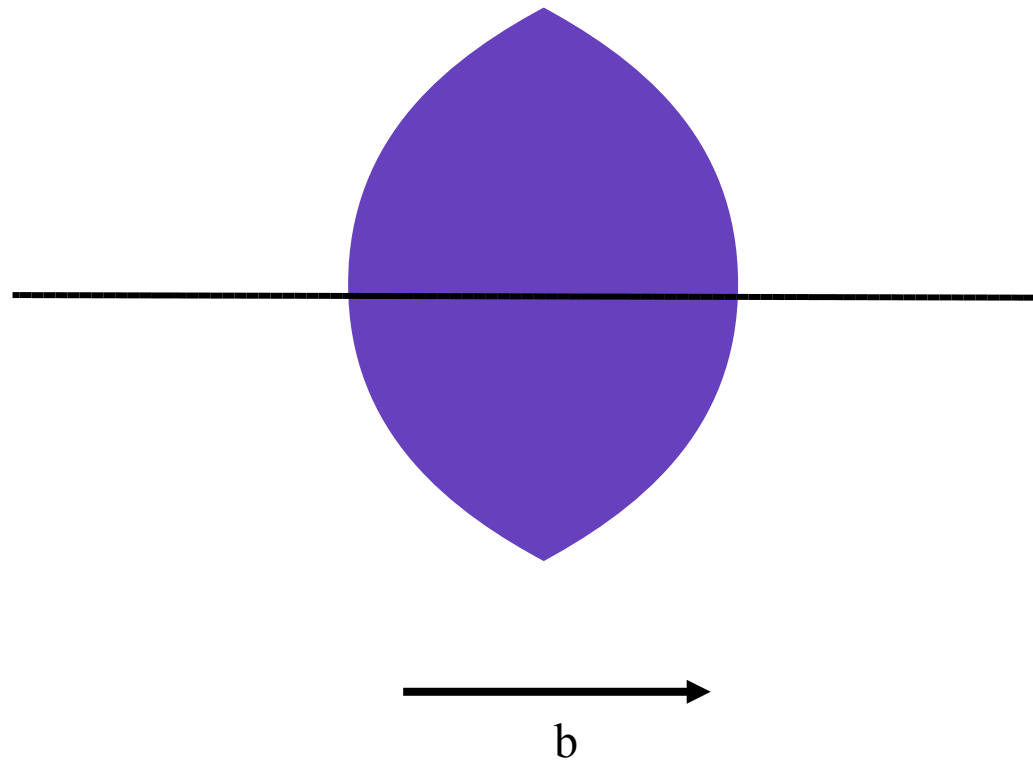


ALICE

## 1) Superposition of independent pp

$$\text{Eccentricity: } \varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

Animation: Mike Lisa



# Elliptic flow

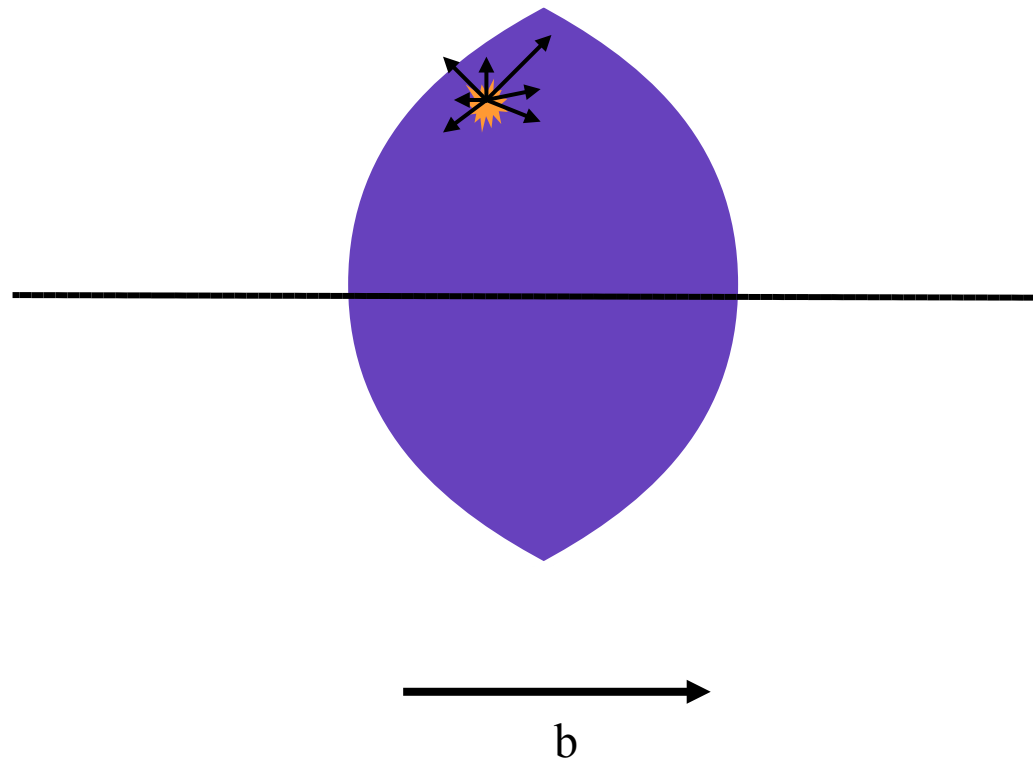


ALICE

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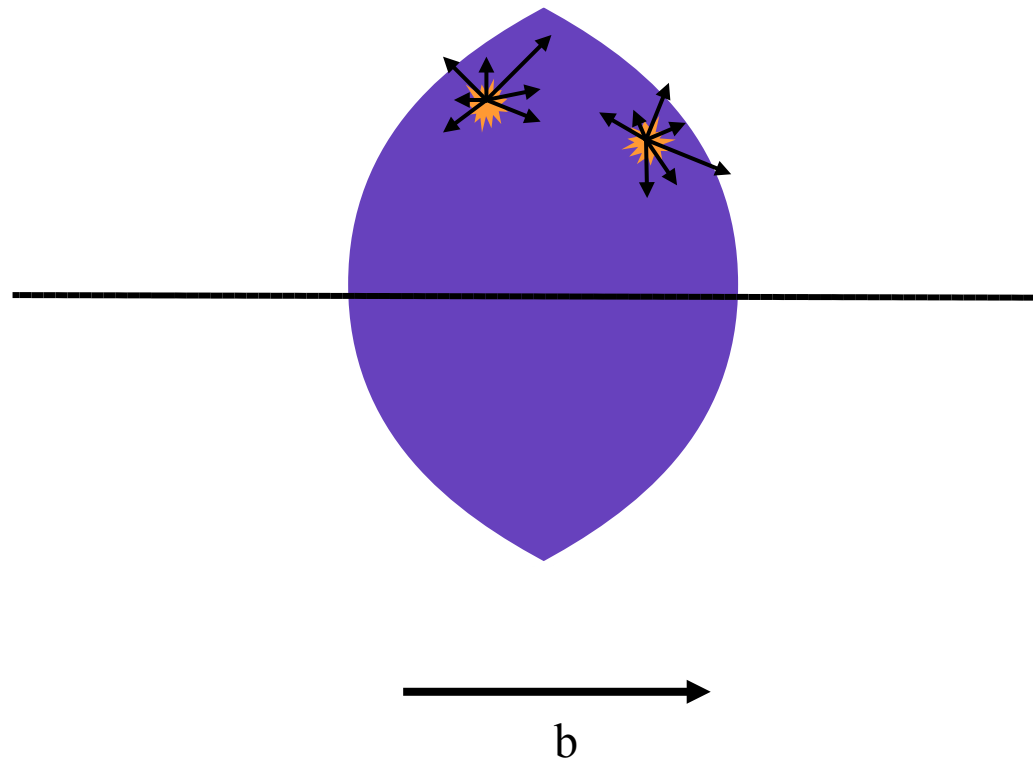


ALICE

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# Elliptic flow

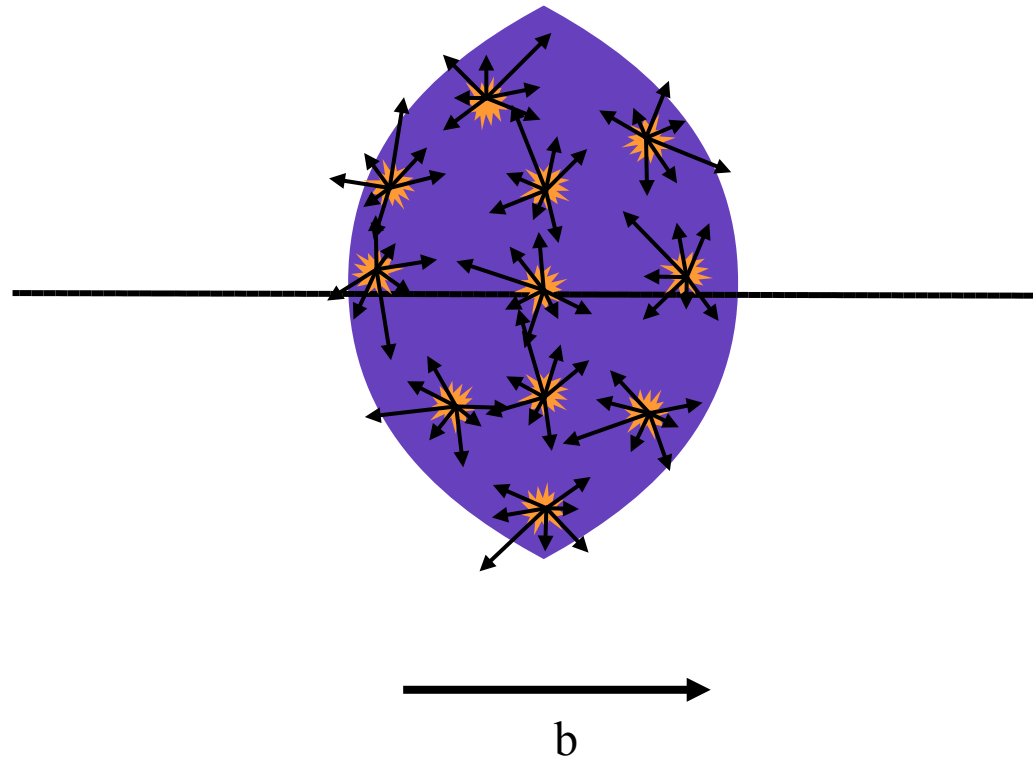


ALICE

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Animation: Mike Lisa



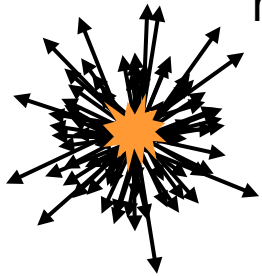
# Elliptic flow



ALICE

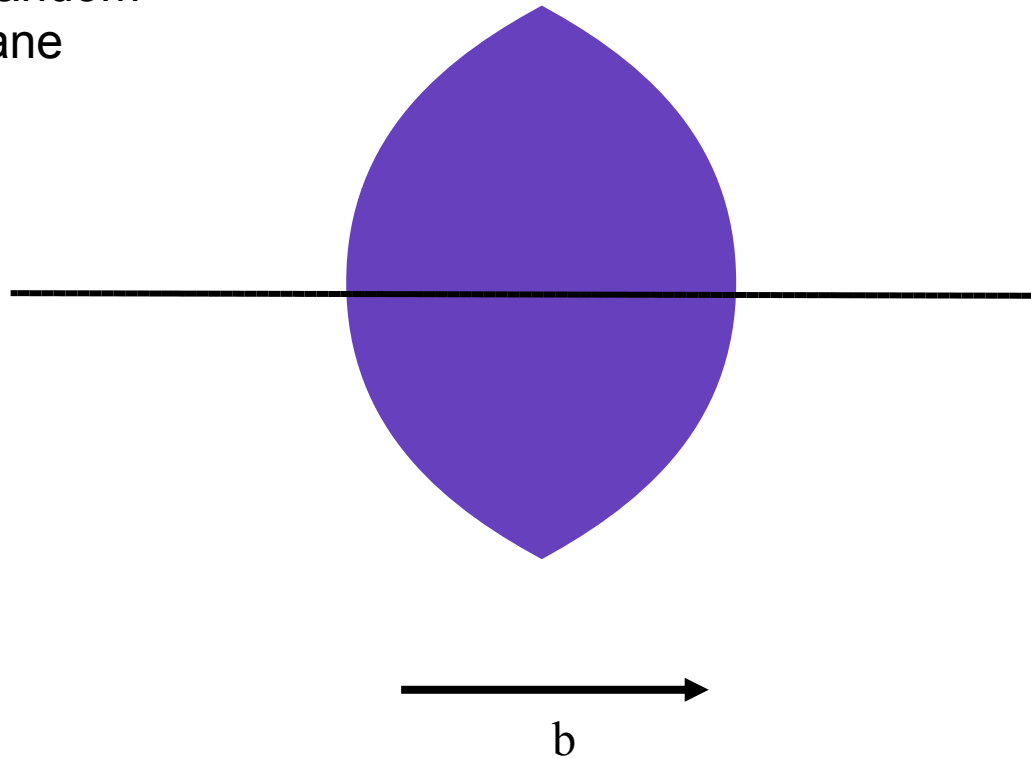
## 1) Superposition of independent pp

Momenta pointed at random  
relative to reaction plane



$$\text{Eccentricity: } \varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

Animation: Mike Lisa





# Elliptic flow

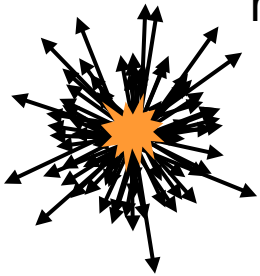


ALICE

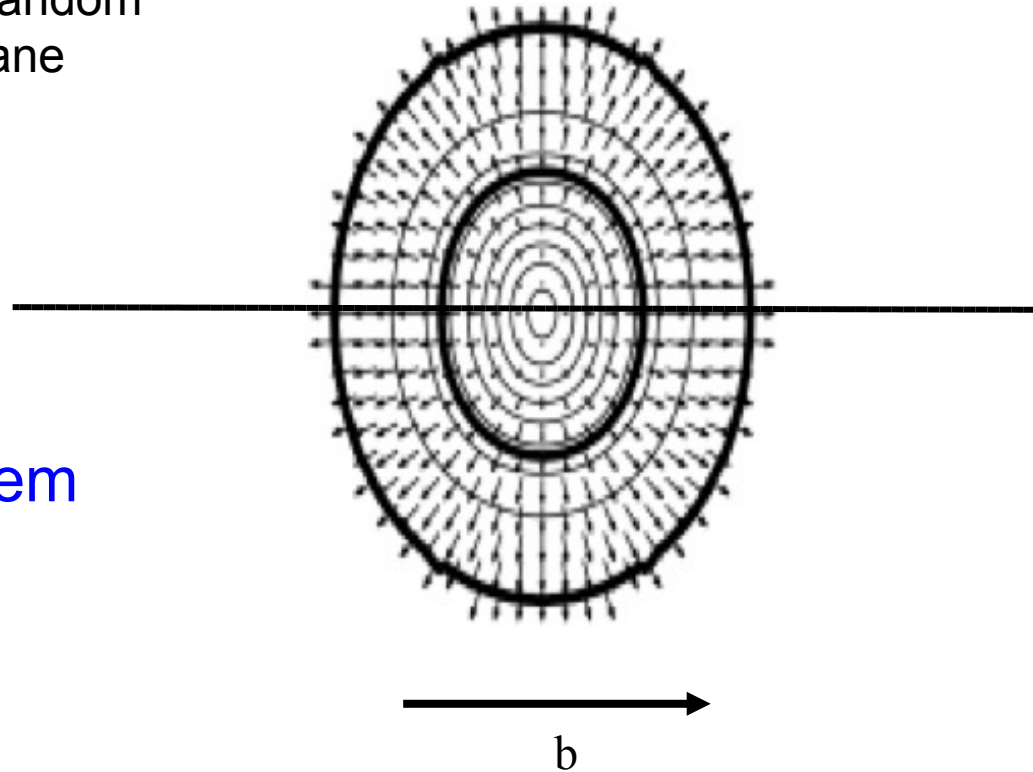
$$\text{Eccentricity: } \varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

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Animation: Mike Lisa



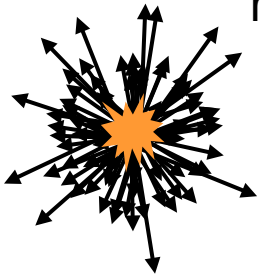
## 2) Evolution as a bulk system

# Elliptic flow

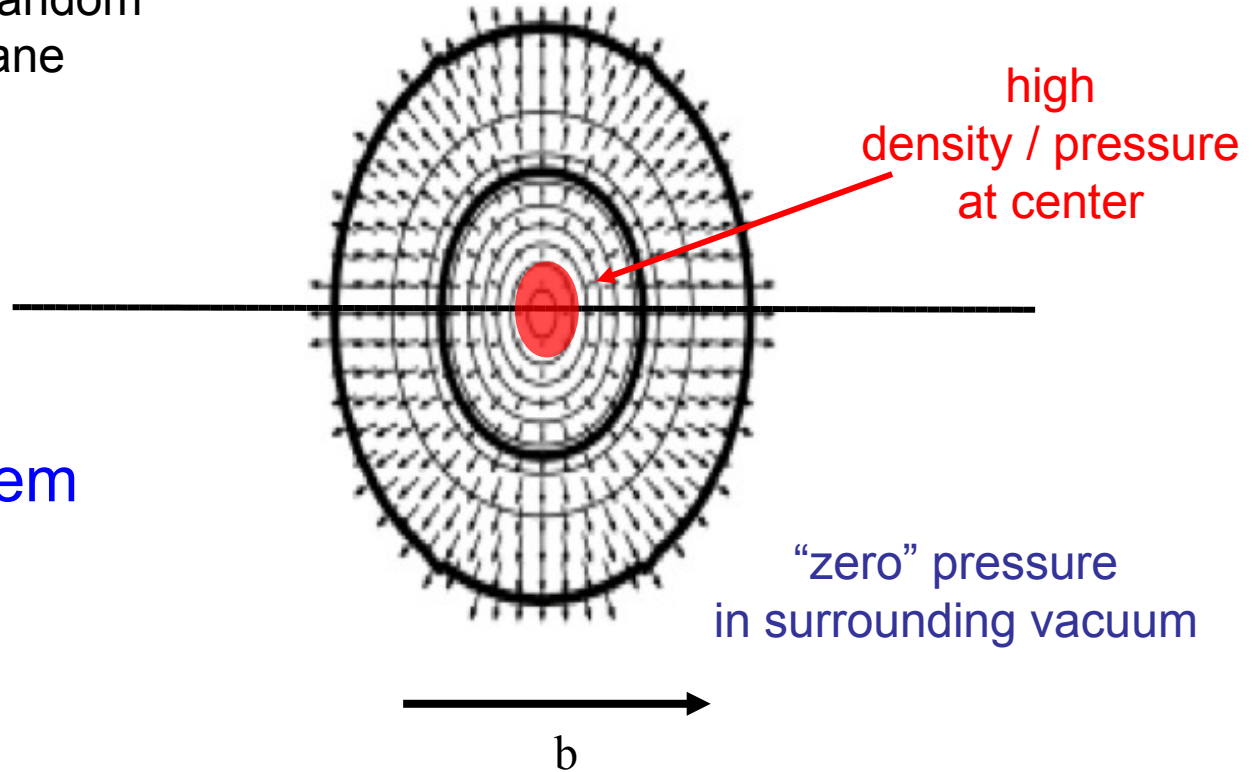
$$\text{Eccentricity: } \varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

## 1) Superposition of independent pp

Momenta pointed at random relative to reaction plane



Animation: Mike Lisa

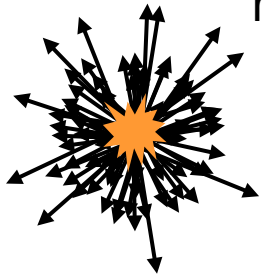


## 2) Evolution as a bulk system

# Elliptic flow

## 1) Superposition of independent pp

Momenta pointed at random relative to reaction plane

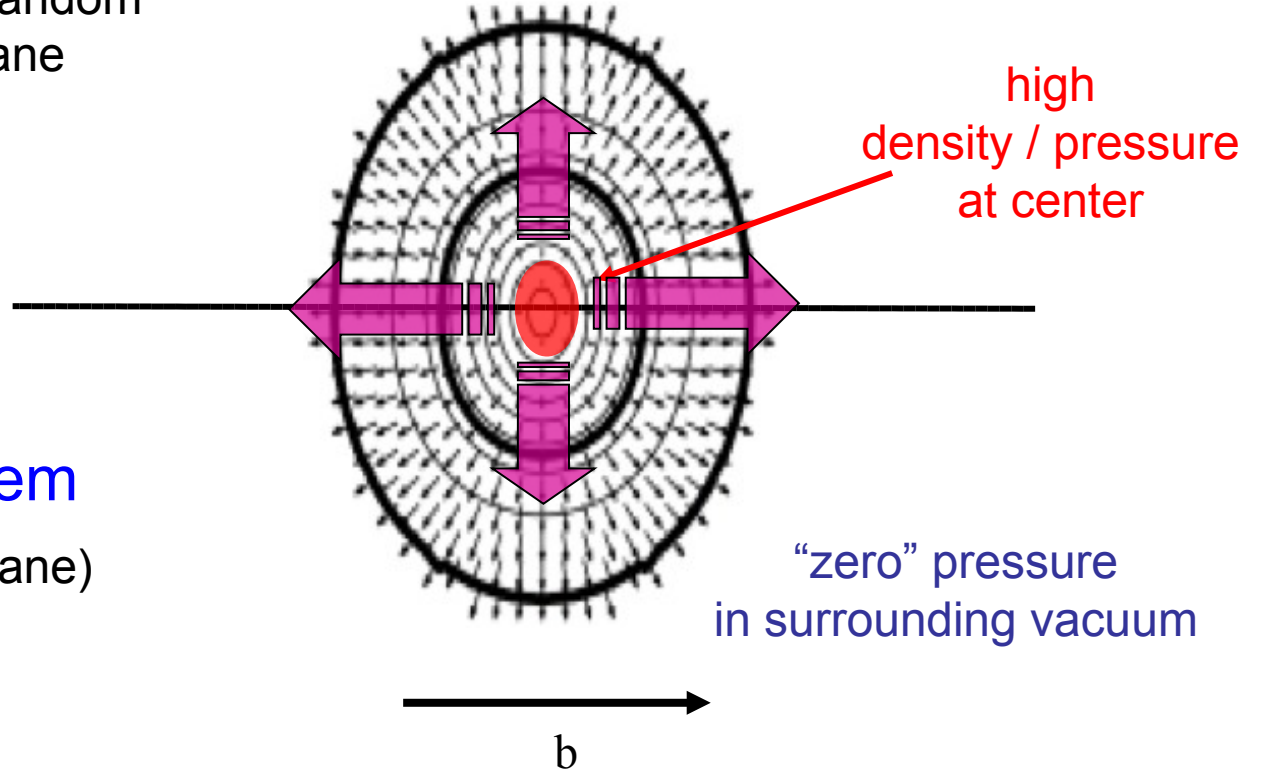


$$\text{Eccentricity: } \varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

Animation: Mike Lisa

## 2) Evolution as a bulk system

Pressure gradients (larger in-plane) push bulk out → “flow”

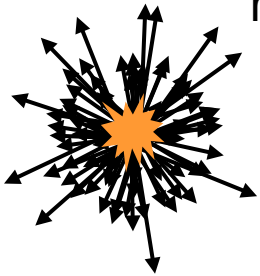


# Elliptic flow

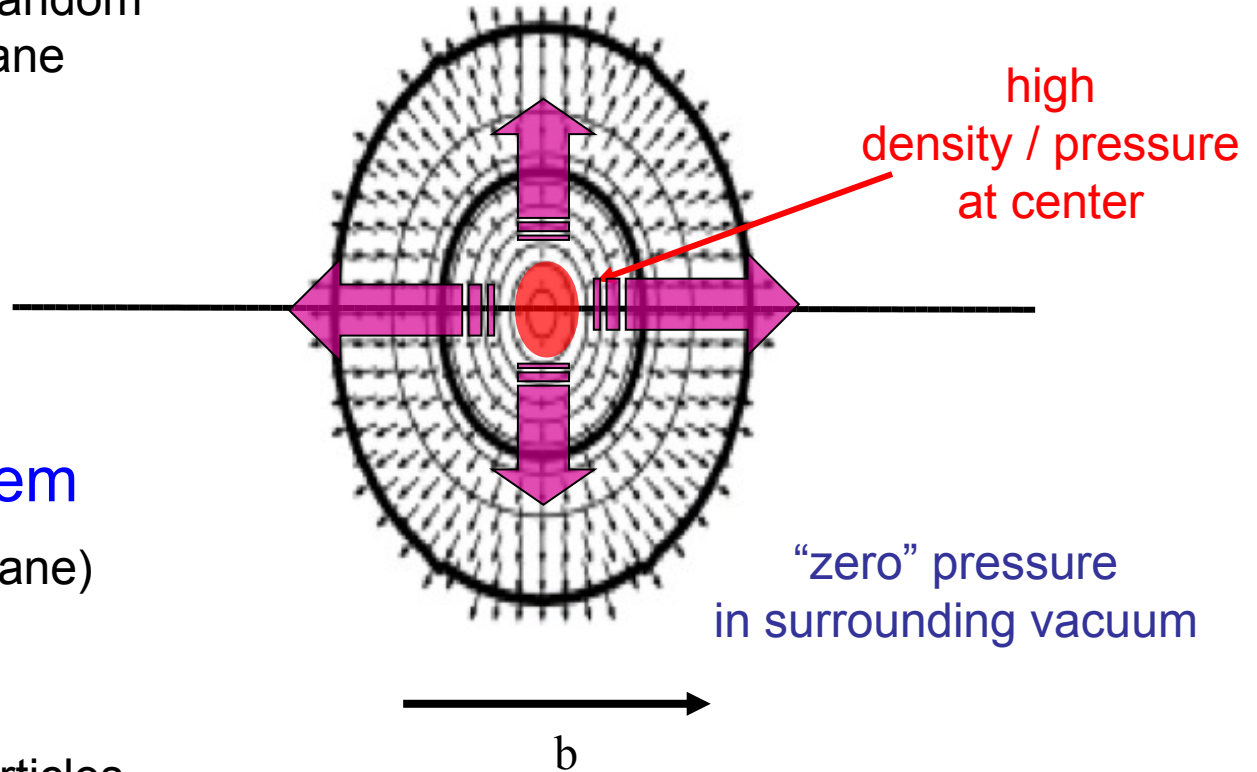
$$\text{Eccentricity: } \varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

## 1) Superposition of independent pp

Momenta pointed at random relative to reaction plane

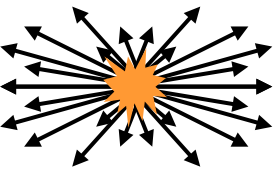


Animation: Mike Lisa



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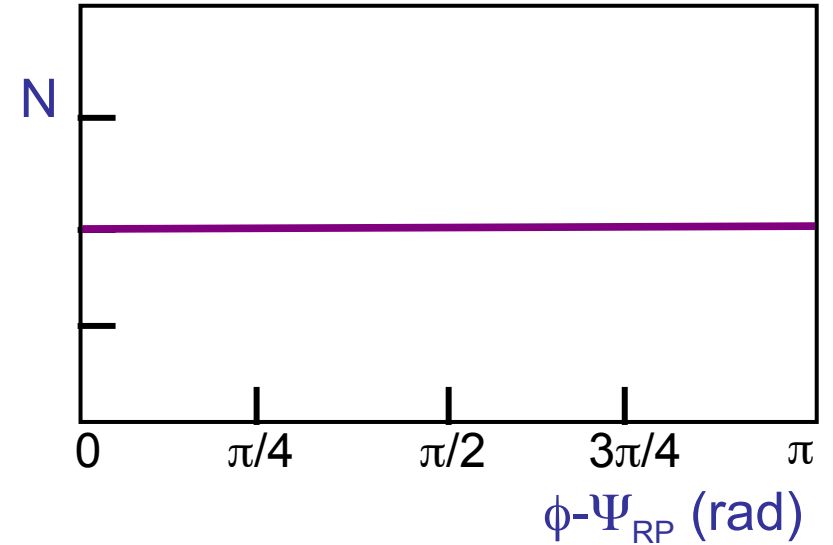
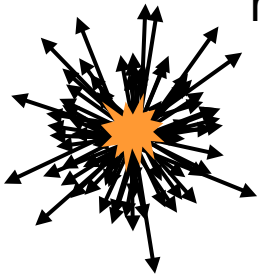
More, faster particles seen in-plane

# Elliptic flow



## 1) Superposition of independent pp

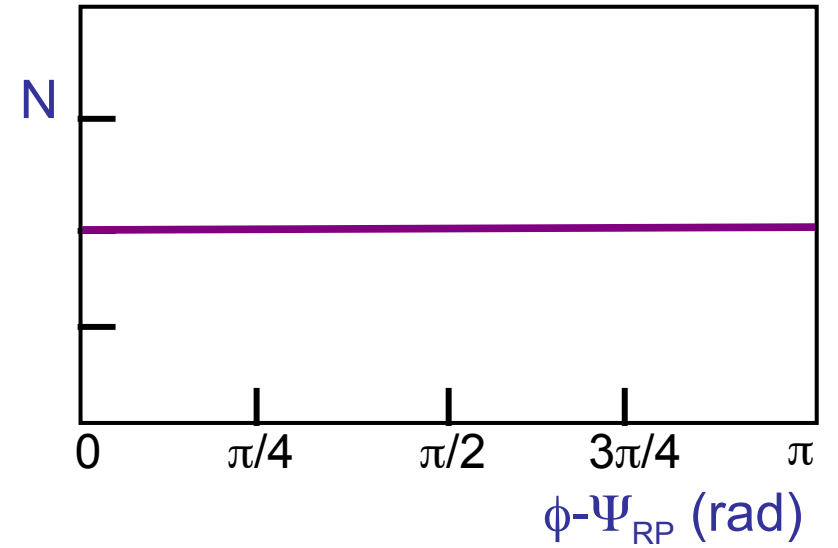
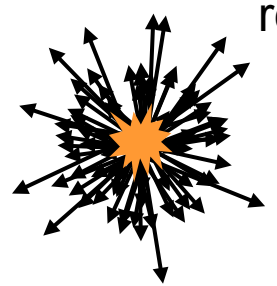
Momenta pointed at random  
relative to reaction plane



# Elliptic flow

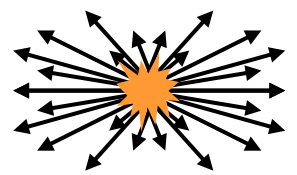
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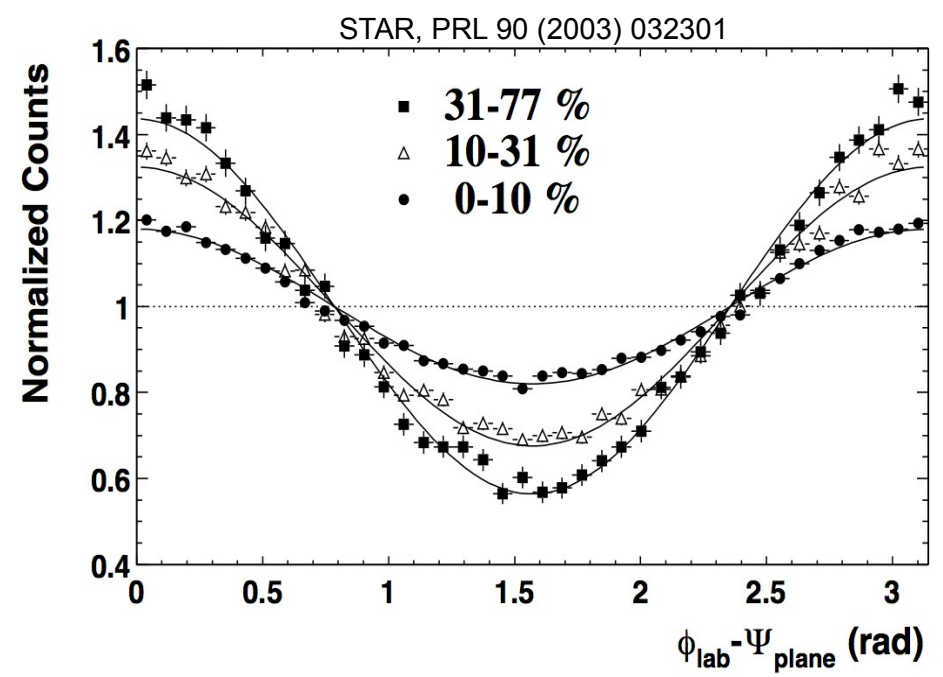


## 2) Evolution as a bulk system

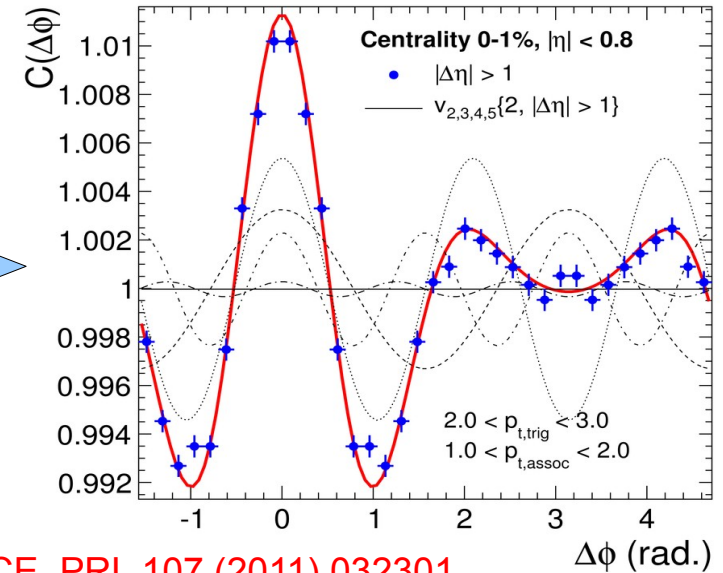
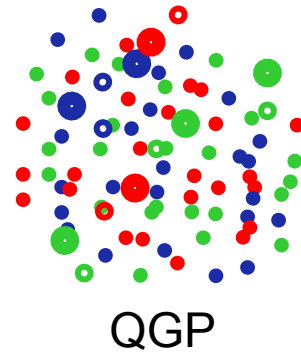
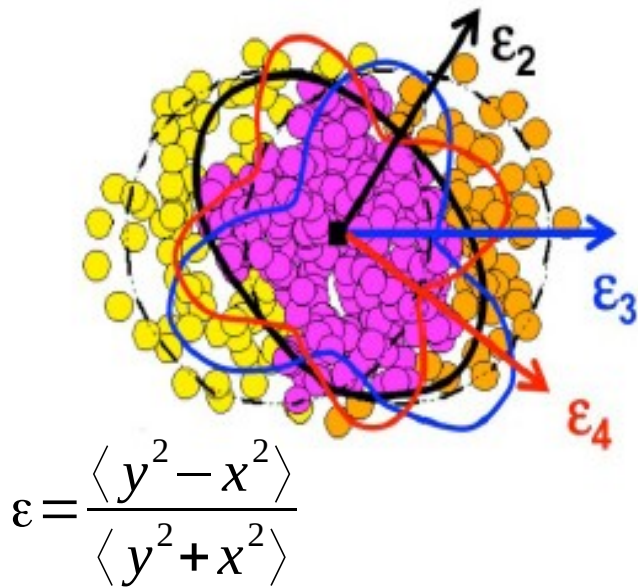
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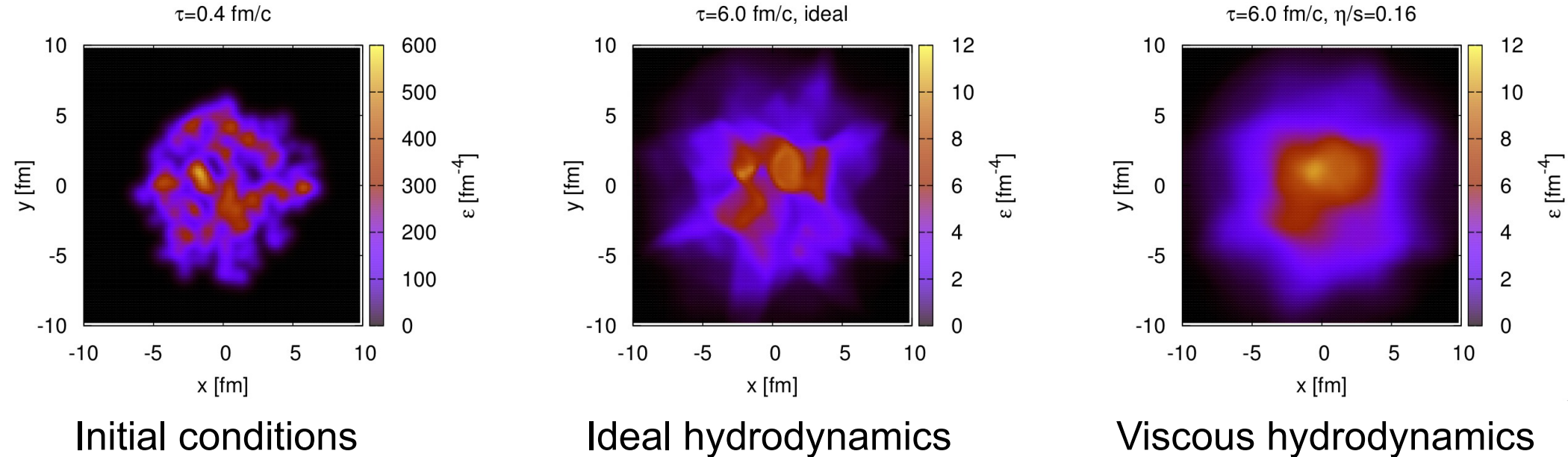
# Anisotropic flow



ALICE, PRL 107 (2011) 032301

- Most central collision: fluctuations of participating nucleons
  - Higher (odd) harmonics since  $\Psi_{\text{RP}} \rightarrow \Psi_n$  (n-th order symmetry plane)
- Anisotropic flow: the transfer of initial spatial anisotropy into the final anisotropy in momentum space via collective interactions

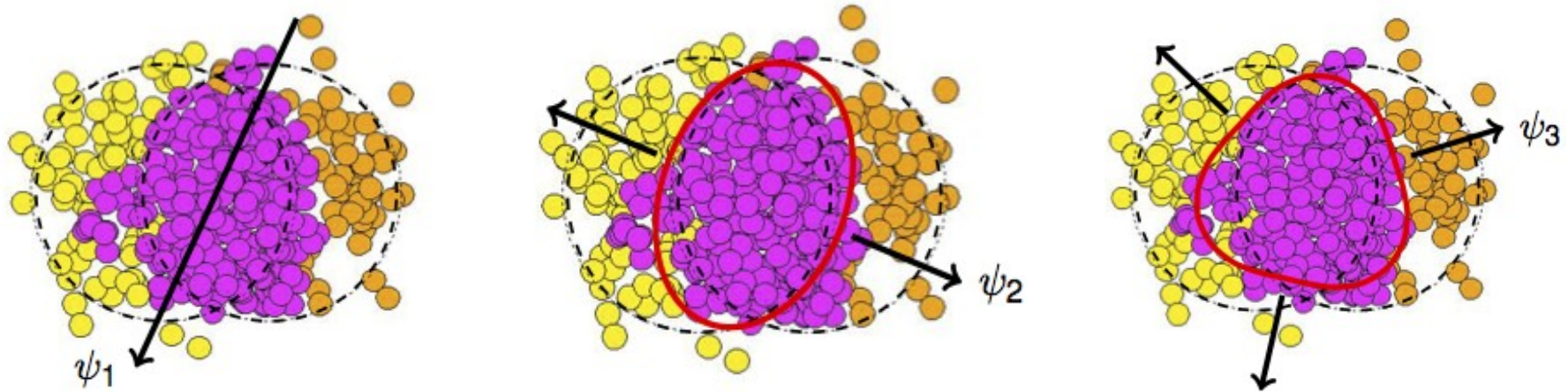
B. Schenke et al., PRL 106 (2011) 042301



- Most central collision: fluctuations of participating nucleons
  - Higher (odd) harmonics since  $\Psi_{RP} \rightarrow \Psi_n$  (n-th order symmetry plane)
- Anisotropic flow: the transfer of initial spatial anisotropy into the final anisotropy in momentum space via collective interactions
  - Sensitive to the system evolution
    - Constrain initial conditions, equation-of-state (EOS), transport properties



# Quantifying anisotropic flow



M. Luzum, J. Phys. G: Nucl. Part. Phys. 38 (2011) 124026

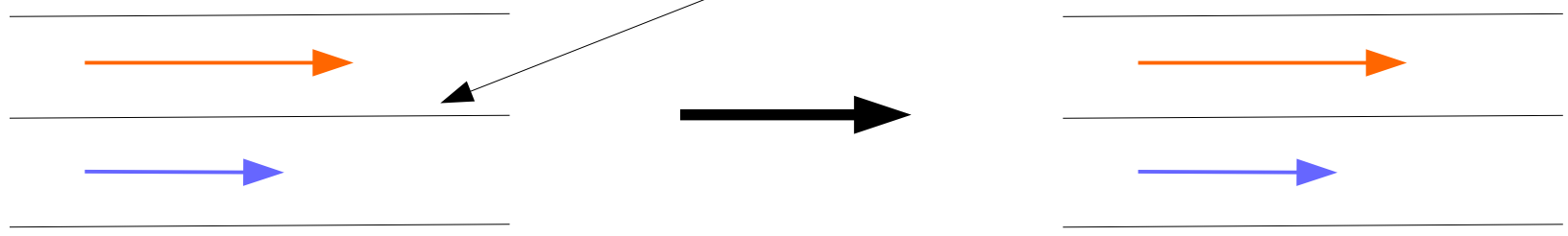
$$E \frac{d^3 N}{d^3 p} = \frac{1}{2\pi} \frac{d^2 N}{p_T dp_T dy} \left( 1 + \sum_{n=1}^{\infty} 2 v_n \cos(n(\varphi - \Psi_n)) \right)$$

$$v_n = \langle \cos(n(\varphi - \Psi_n)) \rangle$$

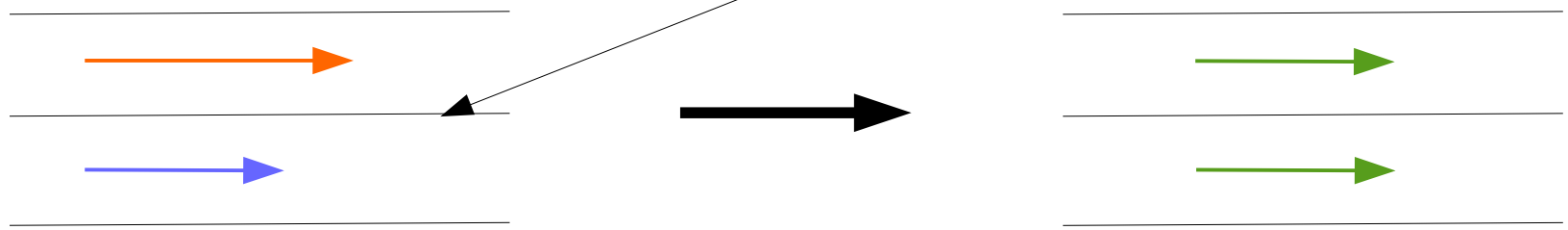
- Particle azimuthal distribution measured with respect to the symmetry plane is not isotropic → Fourier series S. Voloshin and Y. Zhang, Z. Phys. C 70, (1996) 665
- $v_n$  quantify the event anisotropy
  - $v_1$  directed flow,  $v_2$  elliptic flow,  $v_3$  triangular flow, ...

# Shear viscosity over entropy density ( $\eta/s$ )

No friction,  $\eta/s = 0$

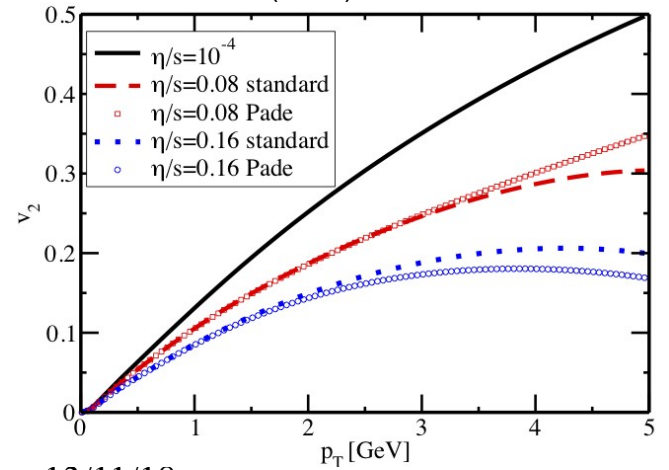


Friction between layers,  $\eta/s > 0$



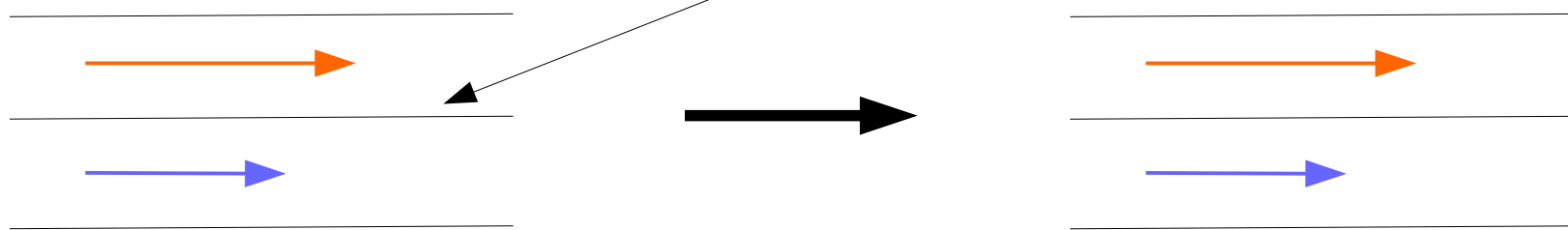
- Shear viscosity will make the velocities  $u_1$ ,  $u_2$ ,  $u_3$  equal and destroy the elliptic flow

M. Luzum and P. Romatschke, PRC 78 (2008) 034915  
 Erratum: PRC 79 (2009) 039903

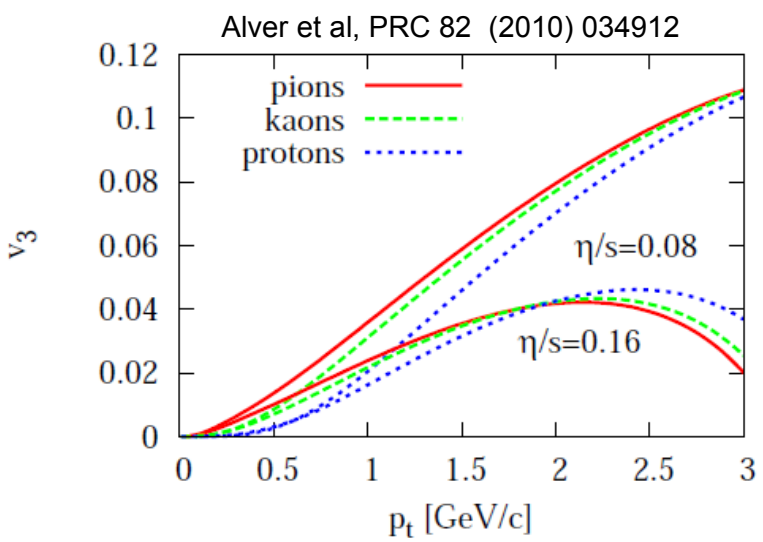
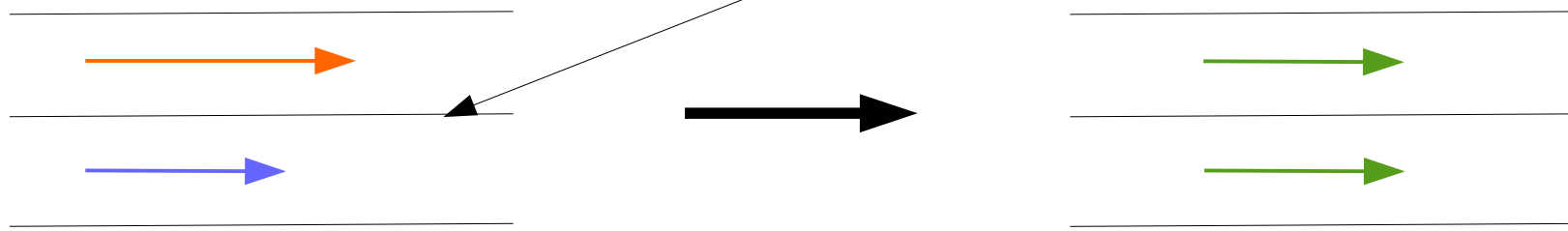


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Friction between layers,  $\eta/s > 0$



- Shear viscosity will make the velocities  $u_1$ ,  $u_2$ ,  $u_3$  equal and destroy the elliptic flow
- Higher harmonics get destroyed more easily due to small differences between velocities  $\rightarrow$  sensitive probes to  $\eta/s$

# Flow fluctuations

- 2- and 4-particle azimuthal correlations for an event

$$\langle 2 \rangle \equiv \langle \cos(n(\varphi_i - \varphi_j)) \rangle, i \neq j$$

$$\langle 4 \rangle \equiv \langle \cos(n(\varphi_i + \varphi_j - \varphi_k - \varphi_l)) \rangle, i \neq j \neq k \neq l$$

- Averaging over all events, the 2<sup>nd</sup> and 4<sup>th</sup> order cumulants

$$c_n\{2\} = \langle \langle 2 \rangle \rangle = v_n^2$$

$$c_n\{4\} = \langle \langle 4 \rangle \rangle - 2 \langle \langle 2 \rangle \rangle^2 = -v_n^4$$

- Flow fluctuations: affect methods differently

$$v_n\{2\} \approx \langle v_n \rangle + \sigma_{v_n}^2 / (2 v_n)$$

$$v_n\{4\} \approx \langle v_n \rangle - \sigma_{v_n}^2 / (2 v_n)$$

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- For Bessel-Gaussian fluctuations

S. Voloshin et al., PLB 659 (2008) 537

$$v_n\{2\}^2 = v_{n,0}^2 + 2 \sigma_{vn}^2$$

$$v_n\{m\} = v_{n,0}, m \geq 4$$

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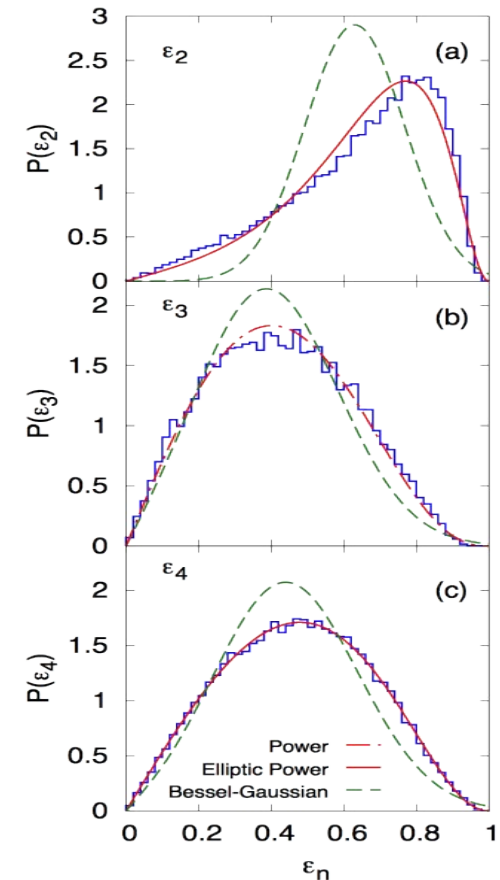
S. Voloshin et al., PLB 659 (2008) 537

$$v_n\{2\}^2 = v_{n,0}^2 + 2 \sigma_{v_n}^2$$

$$v_n\{m\} = v_{n,0}, m \geq 4$$

- Other parametrizations available

L. Yan and J.Y. Ollitrault, PRL 112 (2014) 082301  
 L. Yan et al., PRC 90 (2014) 024903

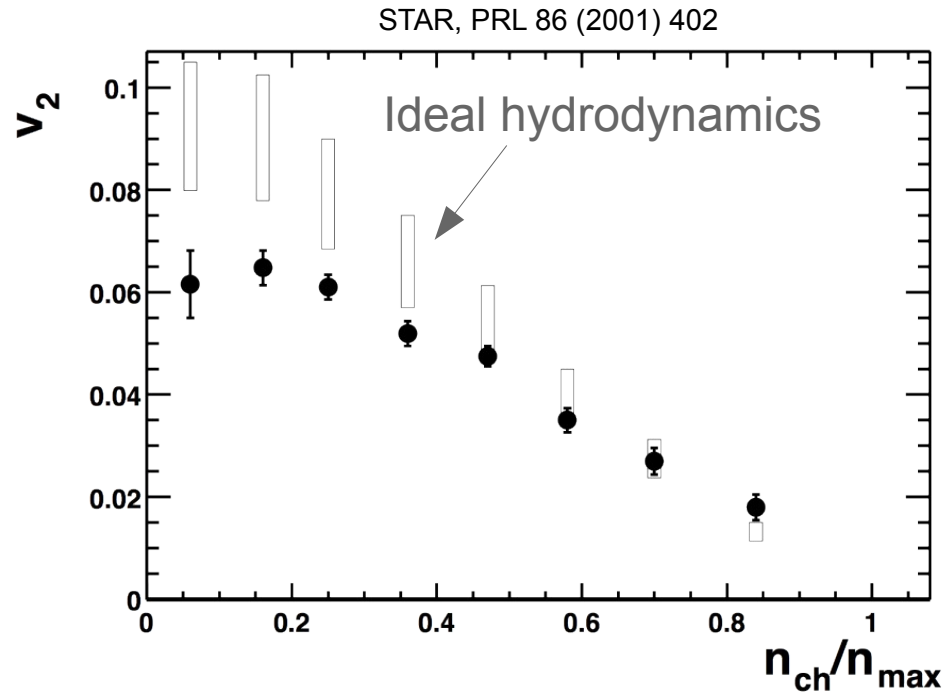




ALICE

# $v_n$ measurements

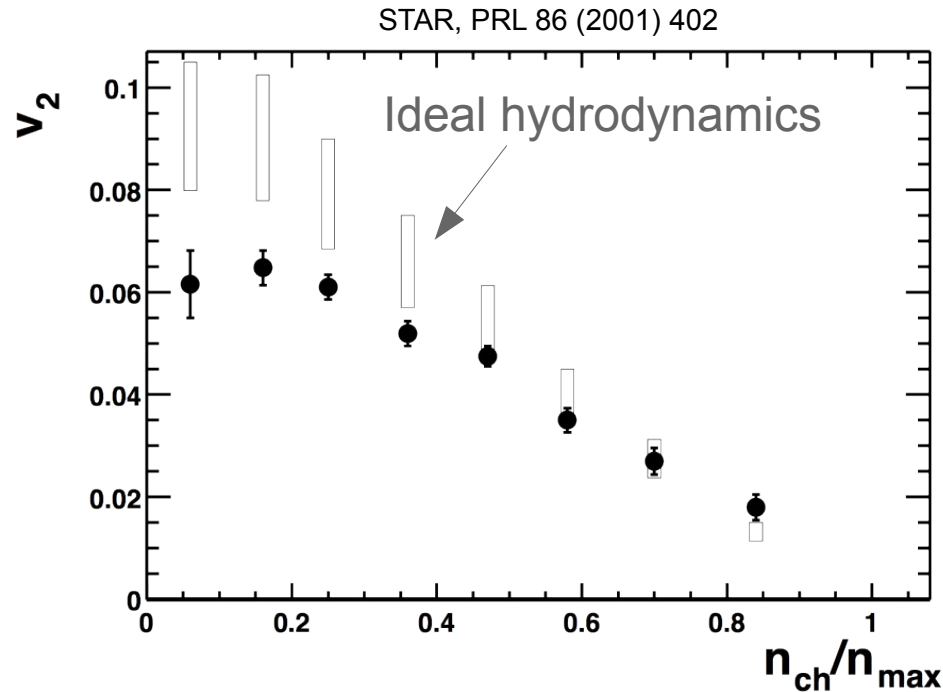
# First $v_2$ @ RHIC



- First  $v_2$  measurement @ RHIC
  - Good agreement with hydrodynamic predictions with  $\eta/s=0$



# First $v_2$ @ RHIC



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## RHIC Scientists Serve Up 'Perfect' Liquid

New state of matter more remarkable than predicted — raising many new questions

April 18, 2005

TAMPA, FL — The four detector groups conducting research at the Relativistic Heavy Ion Collider (RHIC) — a giant atom “smasher” located at the U.S. Department of Energy’s Brookhaven National Laboratory — say they’ve created a new state of hot, dense matter out of the quarks and gluons that are the basic particles of atomic nuclei, but it is a state quite different and even more remarkable than had been predicted. In peer-reviewed papers summarizing the first three years of RHIC findings, the scientists say that instead of behaving like a gas of free quarks and gluons, as was expected, the matter created in RHIC’s heavy ion collisions appears to be more like a *liquid*.

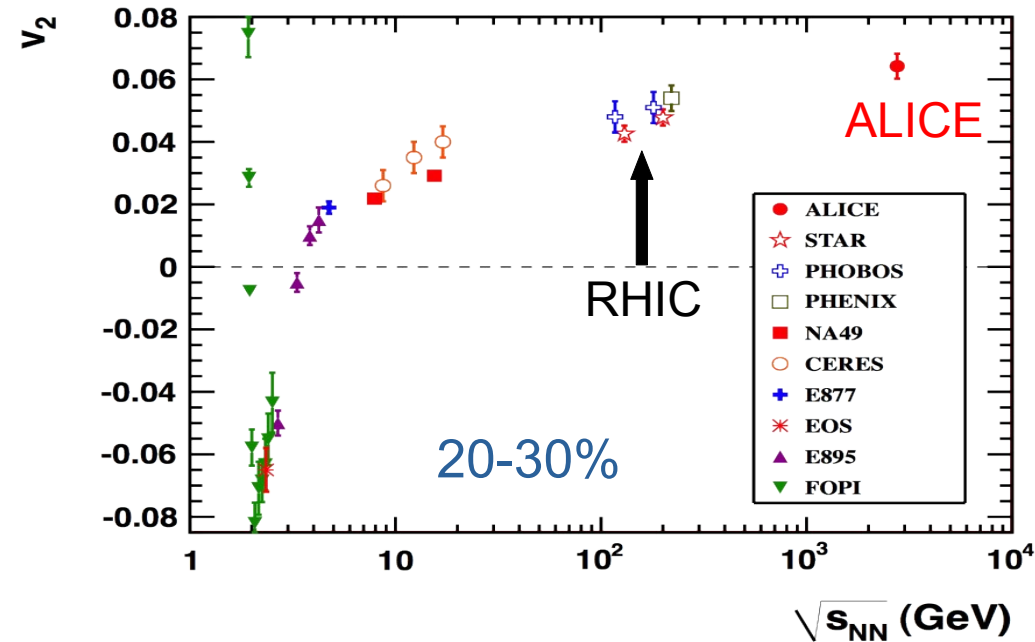
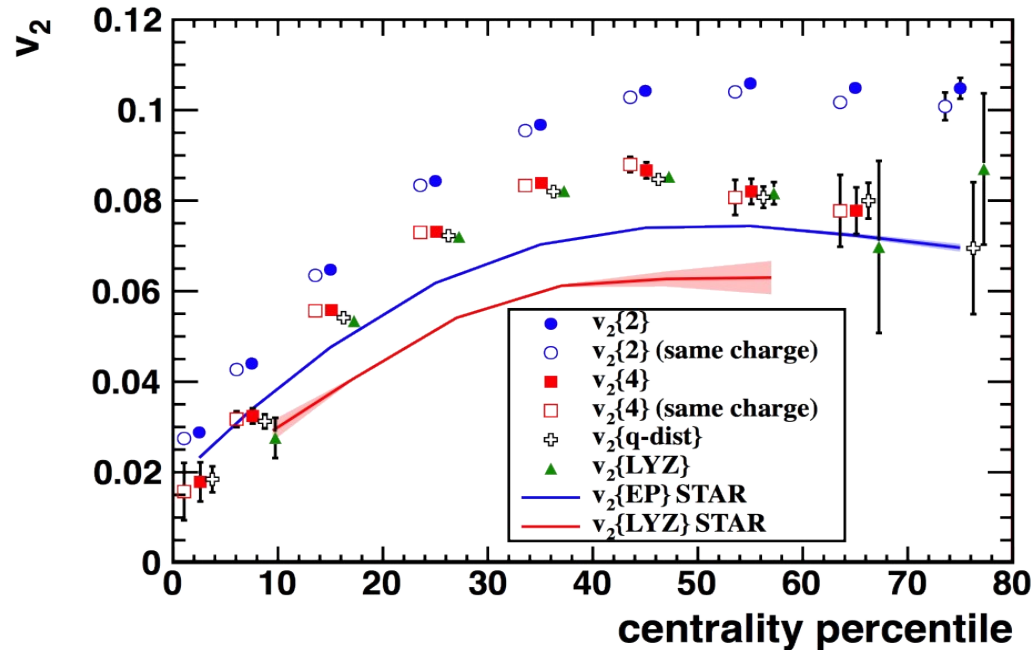
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- First  $v_2$  measurement @ RHIC
  - Good agreement with hydrodynamic predictions with  $\eta/s=0$ 
    - “Perfect” liquid (almost zero friction)

ALICE, PRL 105 (2010) 252302

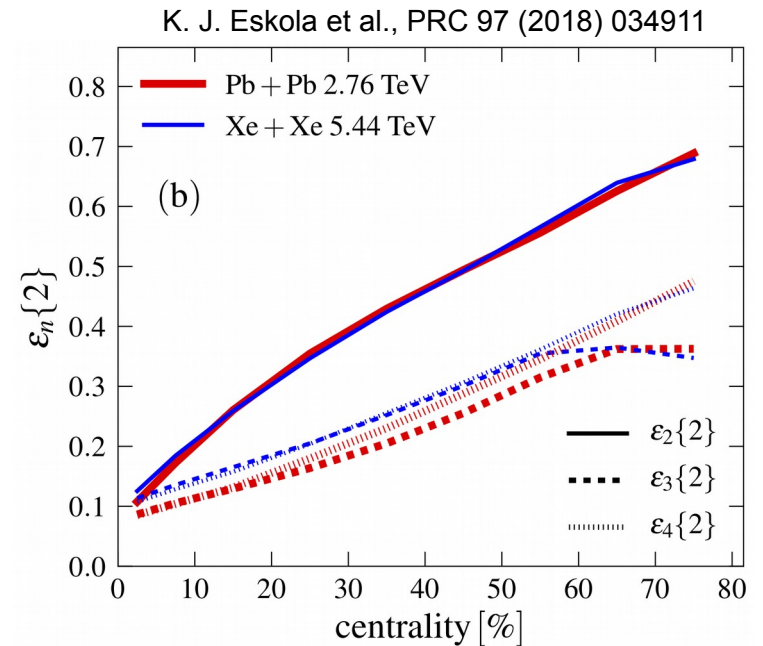
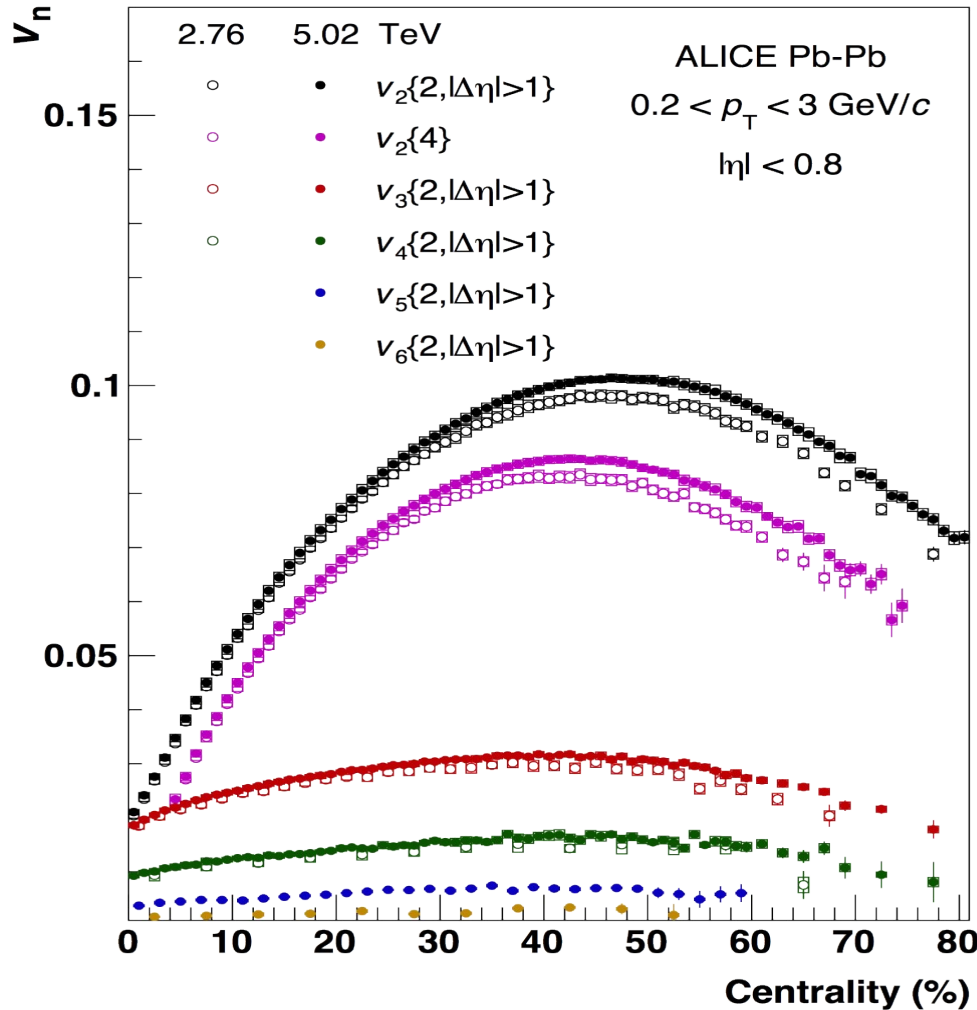


- First  $v_2$  measurement @ LHC

- Elliptic flow increases by  $\sim 30\%$  when compared to RHIC energies
  - The system created at the LHC behaves like a “perfect” liquid

# Integrated $v_n$

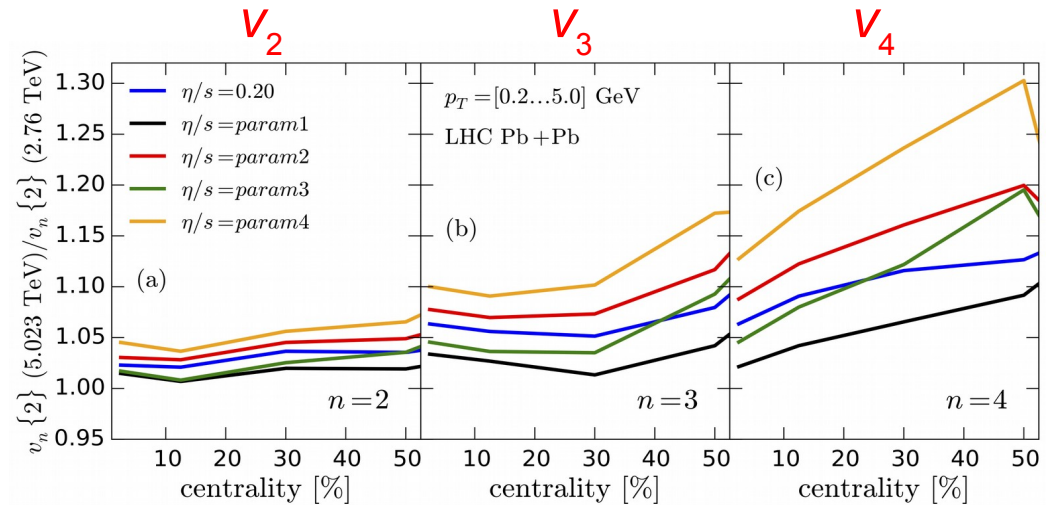
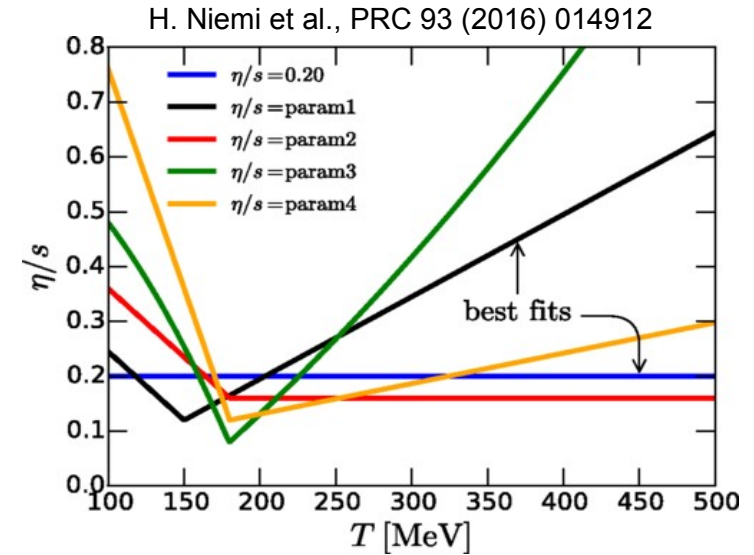
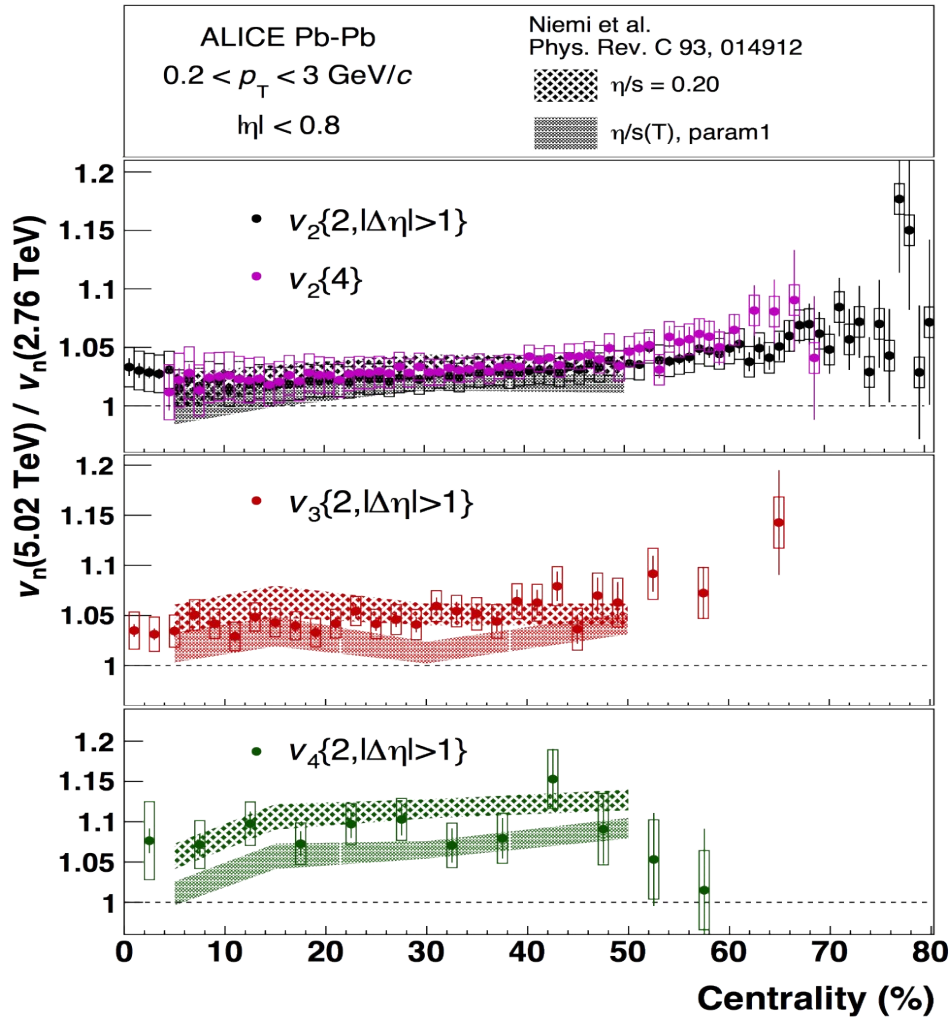
ALICE, JHEP 07 (2018) 103



- Integrated  $v_n$  measured up to  $v_6$  using cumulants
  - Increase of  $\langle p_T \rangle$  responsible for differences between  $\sqrt{s_{NN}} = 2.76 \text{ TeV}$  and  $5.02 \text{ TeV}$

# Integrated $v_n$

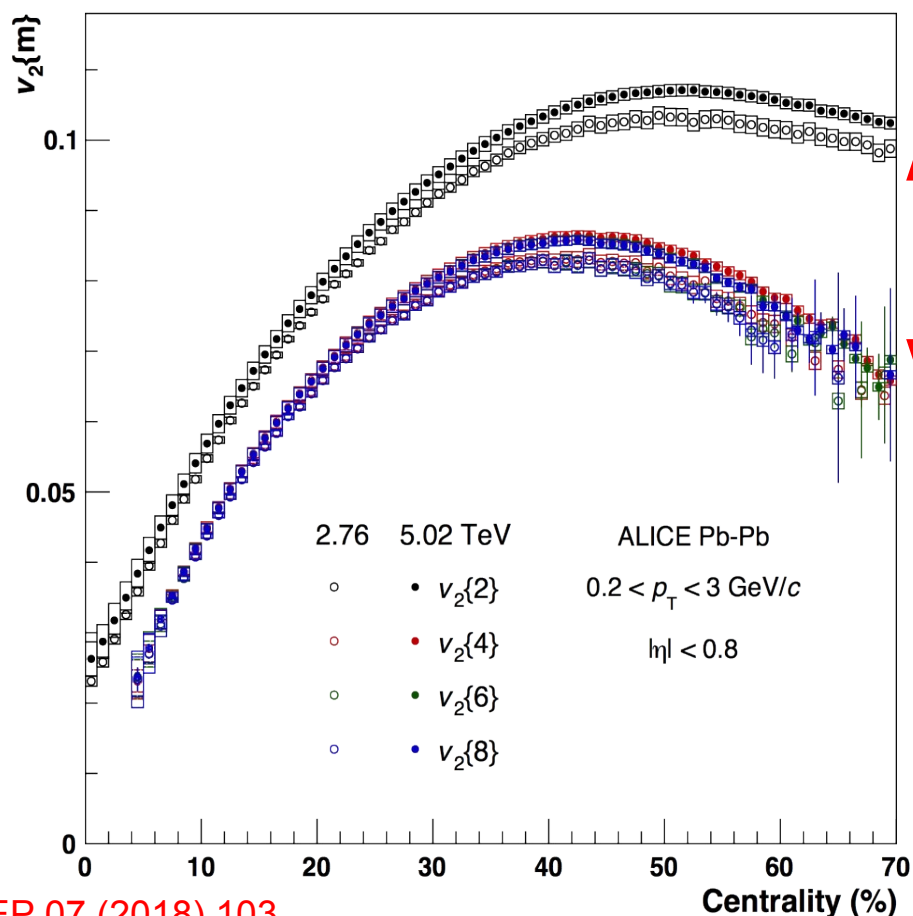
ALICE, JHEP 07 (2018) 103



- Integrated  $v_n$  measured up to  $v_6$  using cumulants

- Increase of  $\langle p_T \rangle$  responsible for differences between  $\sqrt{s_{NN}} = 2.76$  TeV and 5.02 TeV
- Ratios of  $v_n$  at different energies constrain initial conditions and  $\eta/s(T)$

# Integrated $v_2$ : cumulants



ALICE, JHEP 07 (2018) 103

$$v_2\{2\} \approx \langle v_2 \rangle + \sigma_{v_2}^2 / (2 v_2)$$

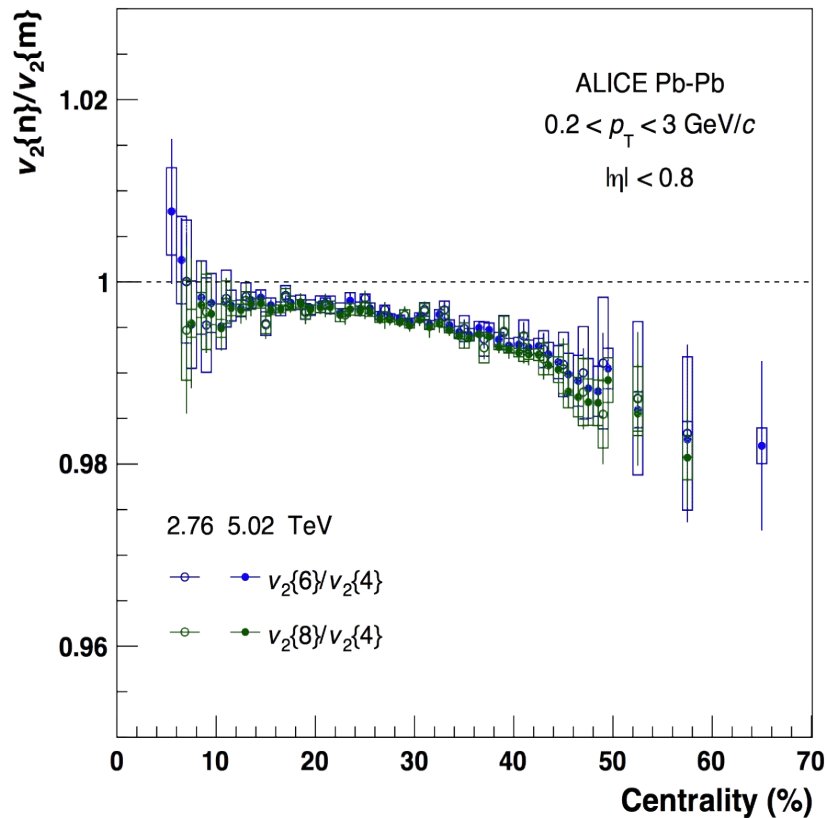
$$v_2\{4\} \approx \langle v_2 \rangle - \sigma_{v_2}^2 / (2 v_2)$$

- Integrated  $v_2$  measured with 2-, 4-, 6-, 8-particle cumulants

- Different sensitivities to flow fluctuations
- Allow to extract flow probability distribution function (p.d.f.)
- Increase of  $\langle p_T \rangle$  responsible for differences between  $\sqrt{s_{NN}} = 2.76$  TeV and 5.02 TeV

# Integrated $v_2$ : ratios of cumulants

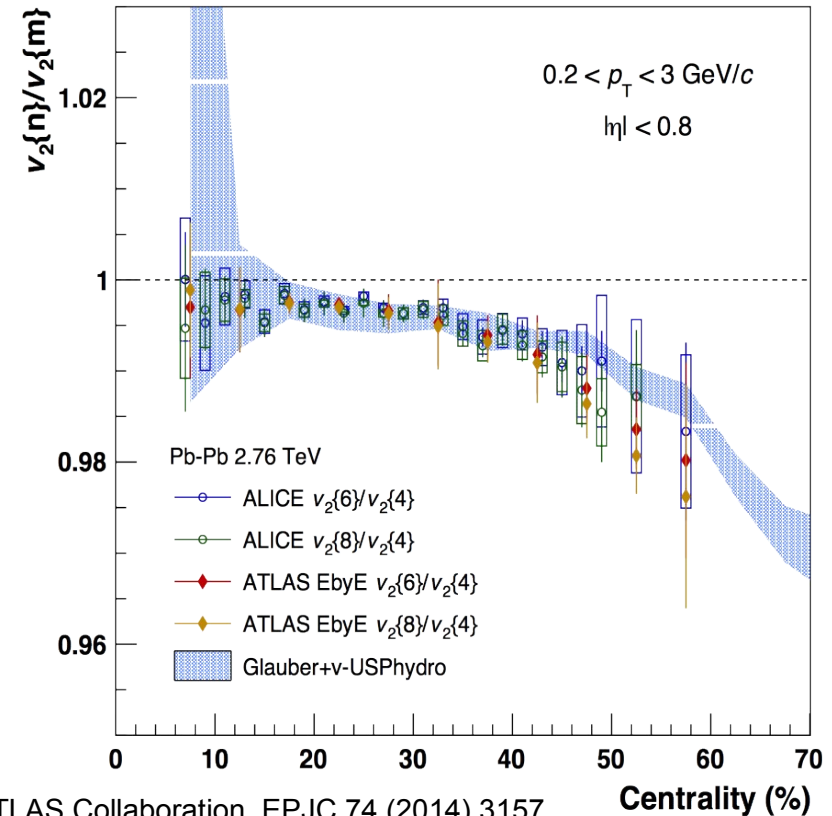
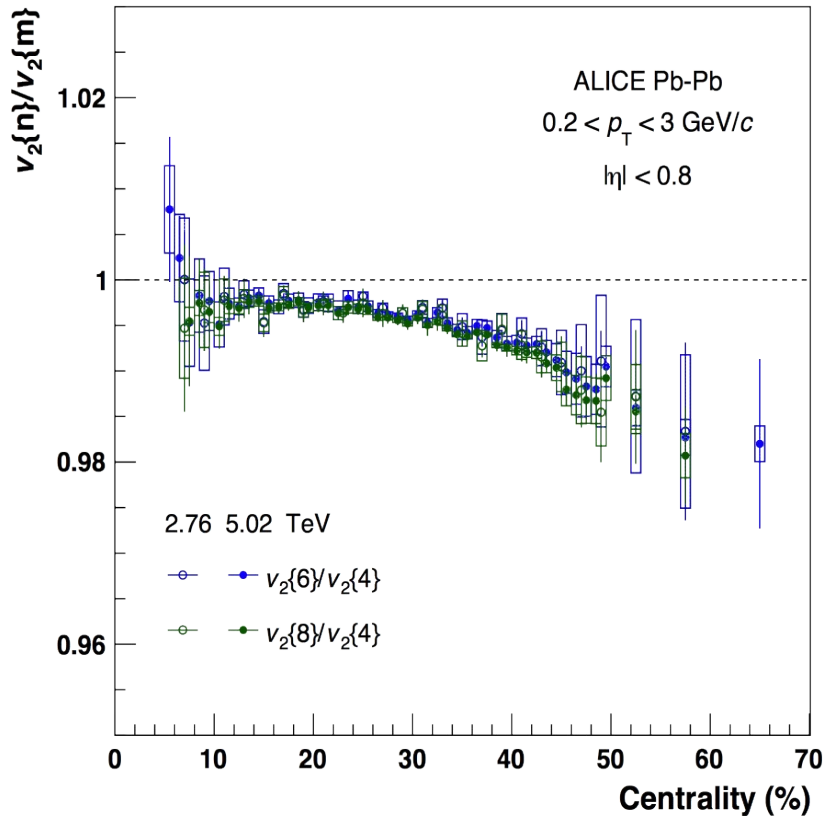
ALICE, JHEP 07 (2018) 103



- Ratios  $v_2\{6\}/v_2\{4\}$  and  $v_2\{8\}/v_2\{4\}$  below unity  $\rightarrow$  non-Gaussian fluctuations
  - Small centrality dependence consistent between  $\sqrt{s_{NN}} = 2.76 \text{ TeV}$  and  $5.02 \text{ TeV}$

# Integrated $v_2$ : ratios of cumulants

ALICE, JHEP 07 (2018) 103

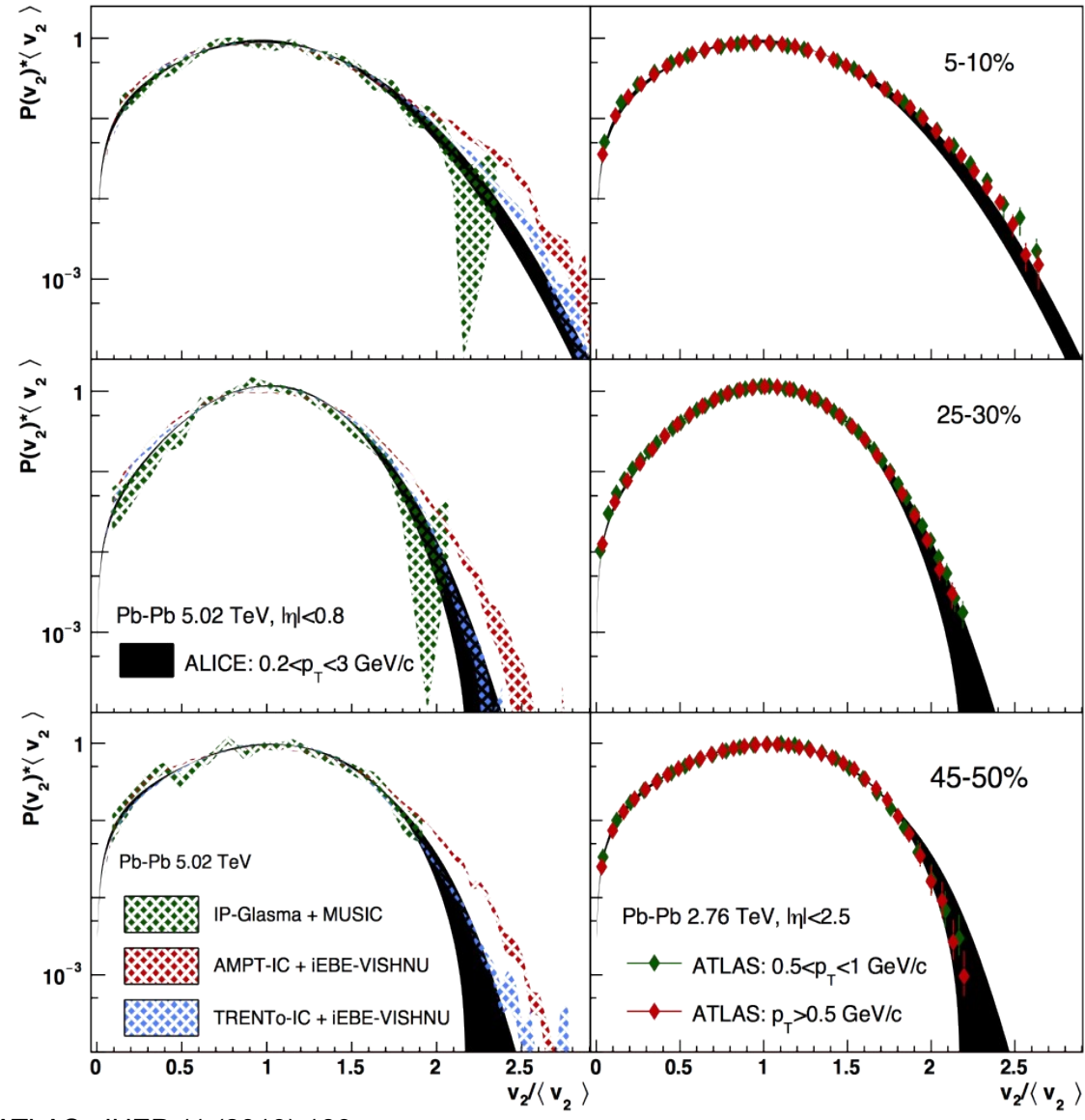


ATLAS Collaboration, EPJC 74 (2014) 3157  
 G. Giacalone et al., PRC 95 (2017) 014913

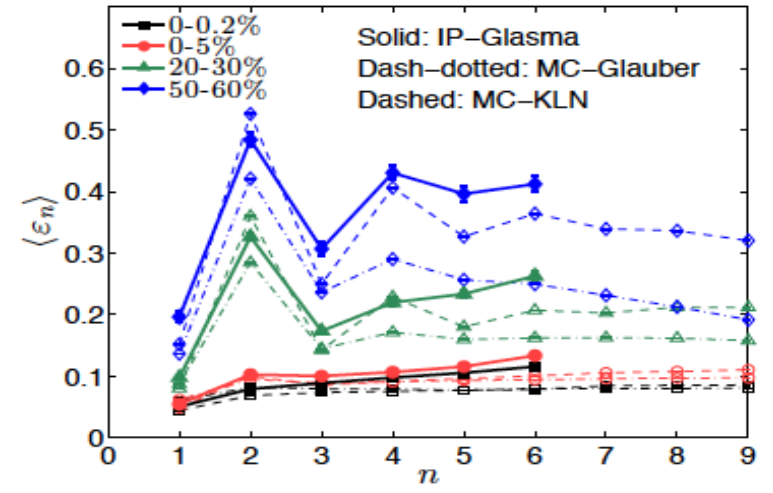
- Ratios  $v_2\{6\}/v_2\{4\}$  and  $v_2\{8\}/v_2\{4\}$  below unity  $\rightarrow$  non-Gaussian fluctuations
  - Small centrality dependence consistent between  $\sqrt{s_{NN}} = 2.76 \text{ TeV}$  and  $5.02 \text{ TeV}$
  - Good agreement with ATLAS results and hydrodynamic calculations

# $v_2$ p.d.f.

ALICE, JHEP 07 (2018) 103



- Use  $v_2\{m\}$  to determine  $v_2$  p.d.f  $P(v_2)$
- $P(v_2)$  scaled by  $\langle v_2 \rangle$  agrees with ATLAS results at  $\sqrt{s_{NN}} = 2.76$  TeV
  - Flow fluctuations at low  $p_T$  depend weakly on  $p_T$  and collision energy
- Good agreement with hydrodynamic calculations with TRENTo and IP-Glasma initial conditions
  - Constrain initial state models



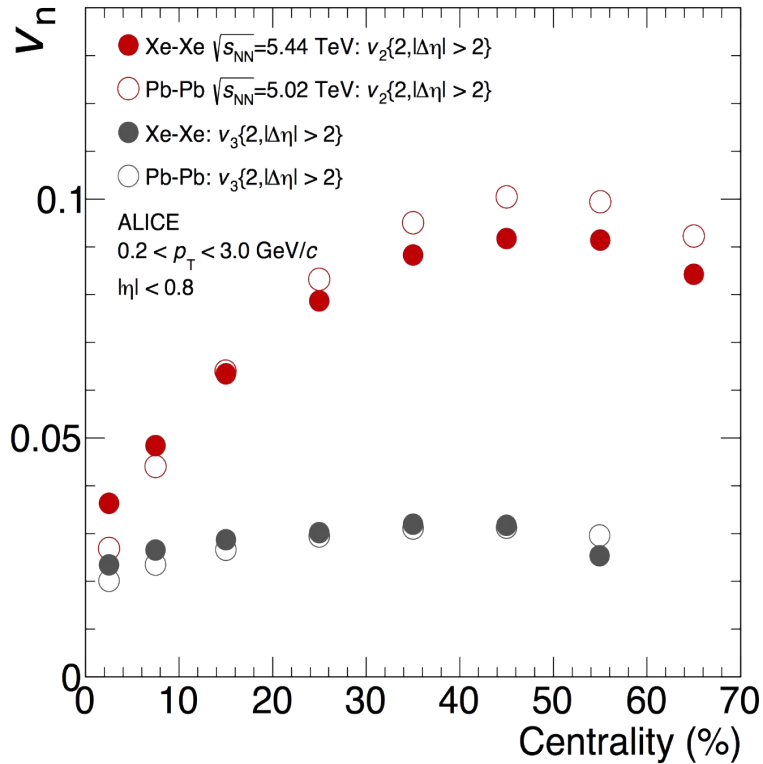
ATLAS, JHEP 11 (2013) 183  
 S. McDonald et al., PRC 95 (2017) 064913  
 W. Zhao et al., EPJC 77 (2017) 645

U. Heinz and R. Snellings,  
 Ann.Rev.Nucl.Part.Sci. 63 (2013) 123



# $v_n$ : Pb-Pb vs Xe-Xe

ALICE, PLB 784 (2018) 82

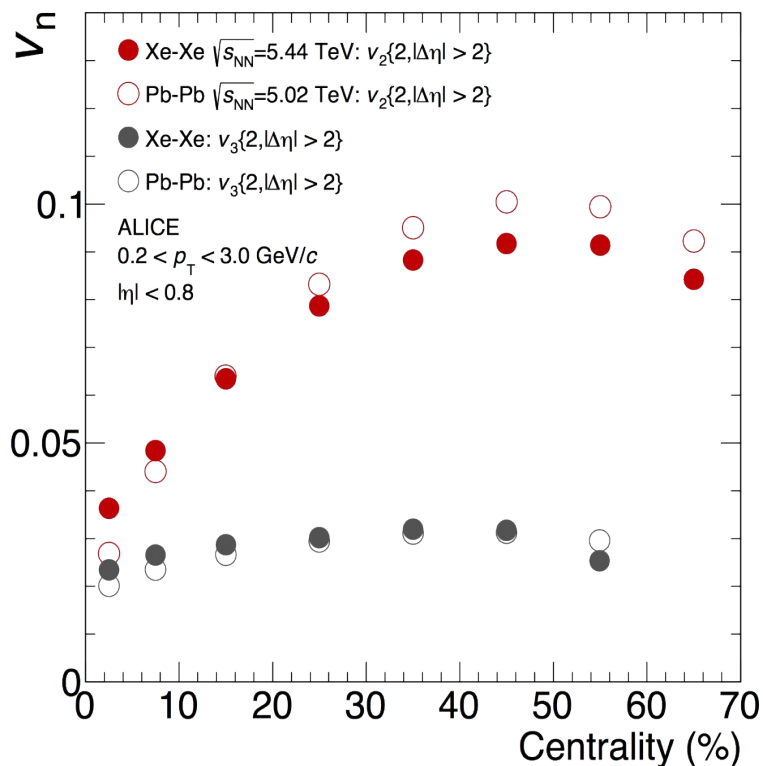


ALI-PUB-150777

- $v_2$ : differences within 10% except in 0-5% centrality interval where reached 35% (Xe deformation)

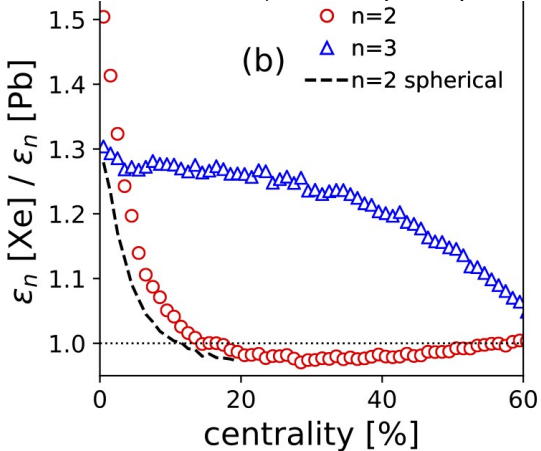
# $v_n$ : Pb-Pb vs Xe-Xe

ALICE, PLB 784 (2018) 82



ALI-PUB-150777

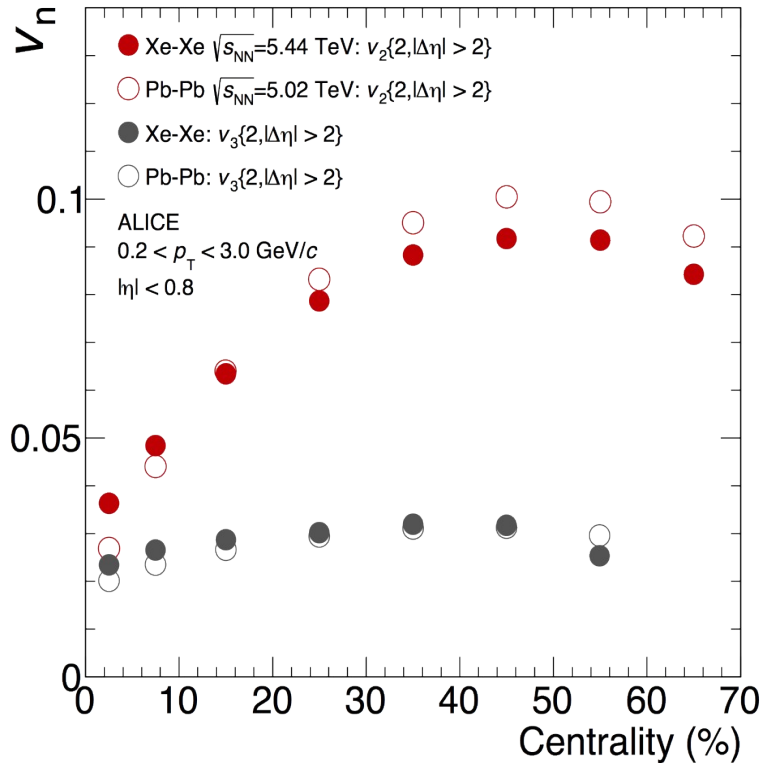
G. Giacalone et al., PRC 97 (2018) 034904



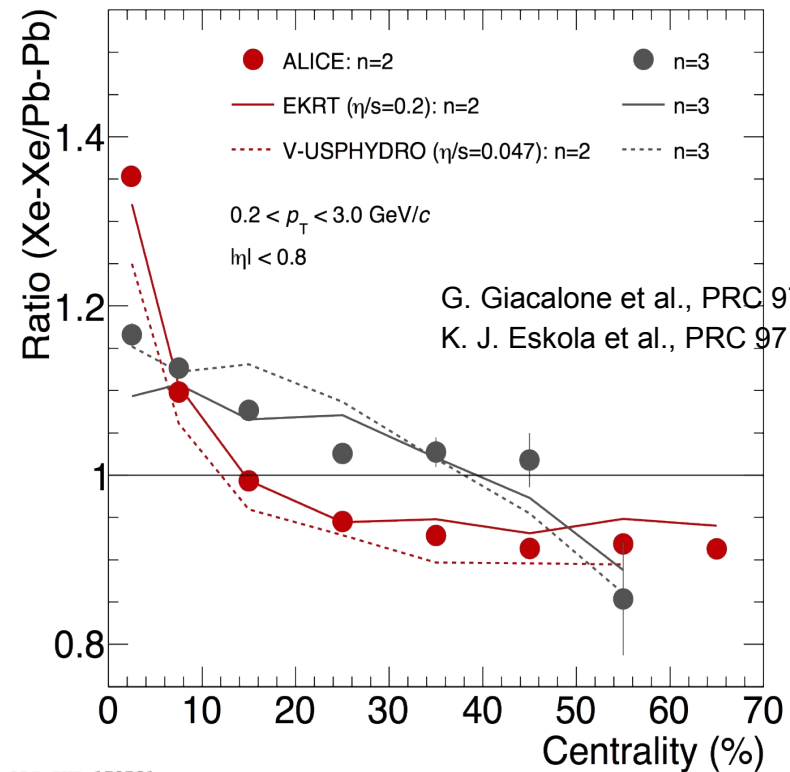
- $v_2$ : differences within 10% except in 0-5% centrality interval where reached 35% (Xe deformation)
- $v_3$ : larger in Xe-Xe due to larger initial state fluctuations

# $v_n$ : Pb-Pb vs Xe-Xe

ALICE, PLB 784 (2018) 82

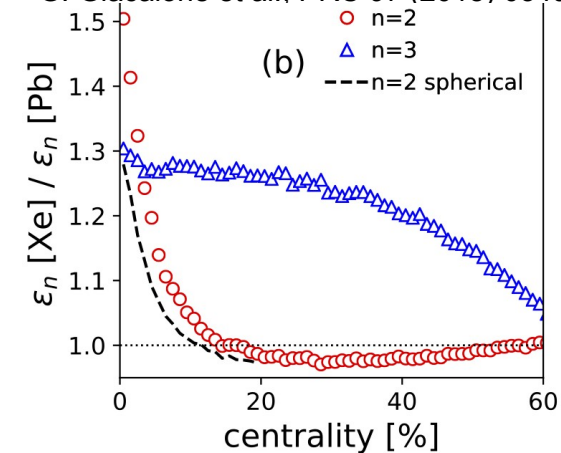


ALI-PUB-150777



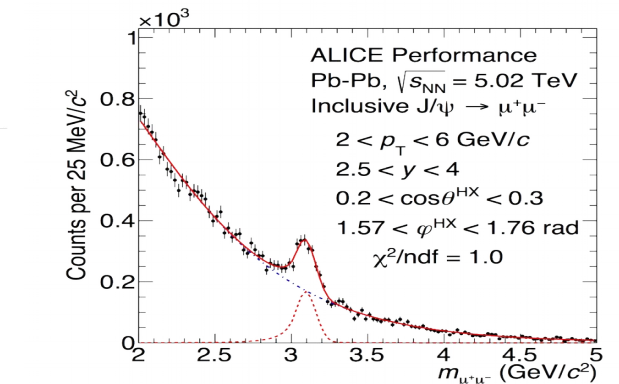
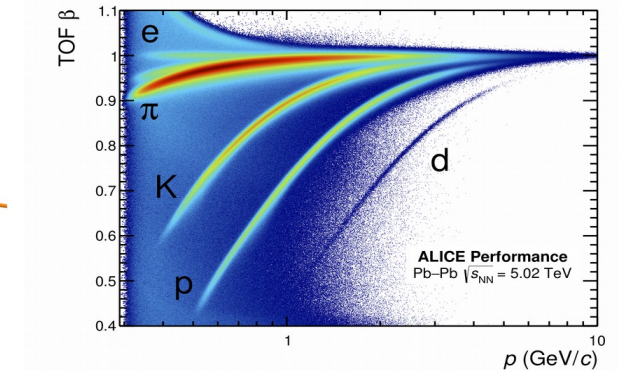
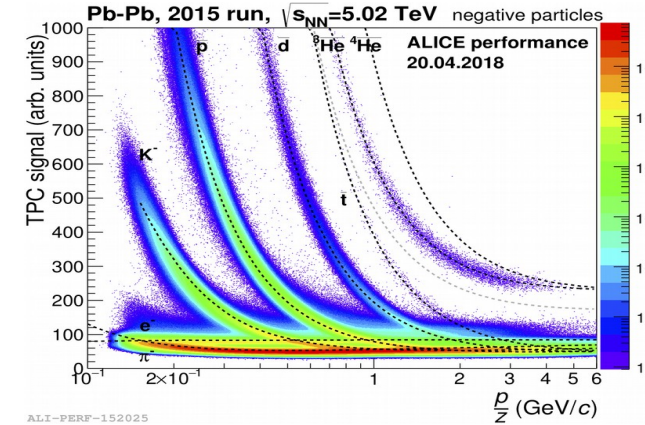
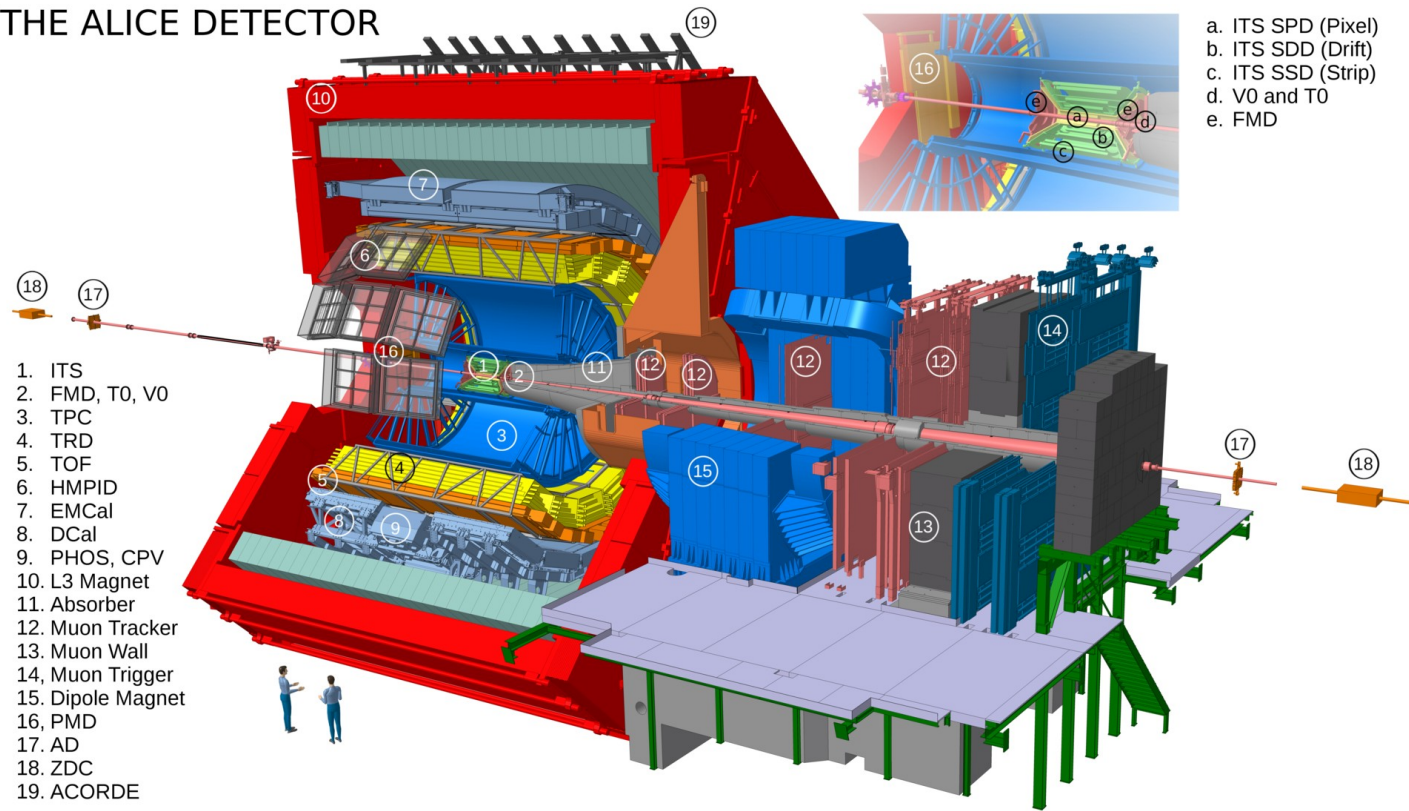
ALI-PUB-150781

G. Giacalone et al., PRC 97 (2018) 034904



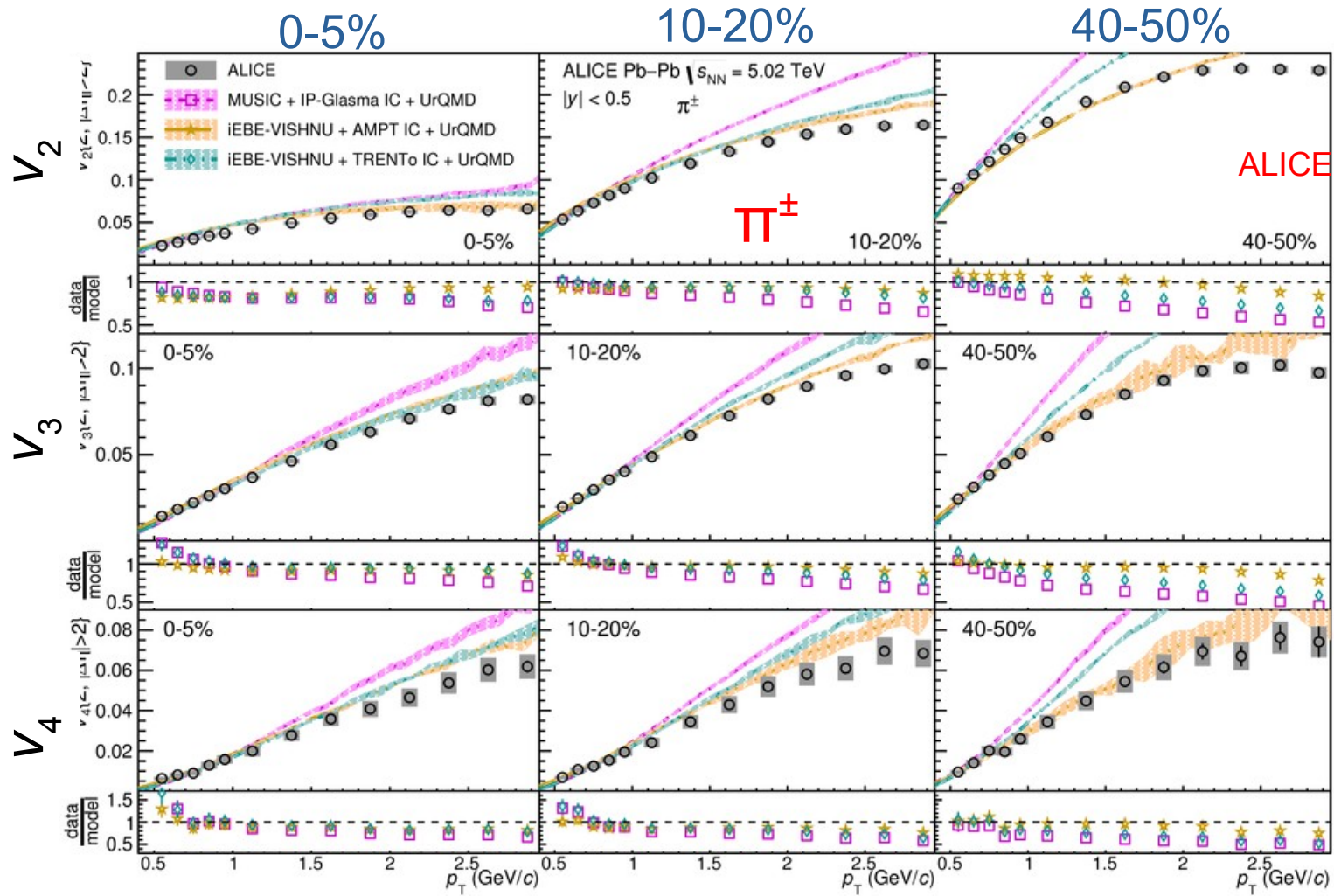
- $v_2$ : differences within 10% except in 0-5% centrality interval where reached 35% (Xe deformation)
- $v_3$ : larger in Xe-Xe due to larger initial state fluctuations
- Results described quantitatively by hydrodynamic models

## THE ALICE DETECTOR



- $\pi^\pm$ ,  $K^\pm$ , p identified using TPC and TOF
- Topological reconstruction for  $K_S^0$ ,  $\Lambda$ ,  $J/\Psi$  and D-mesons

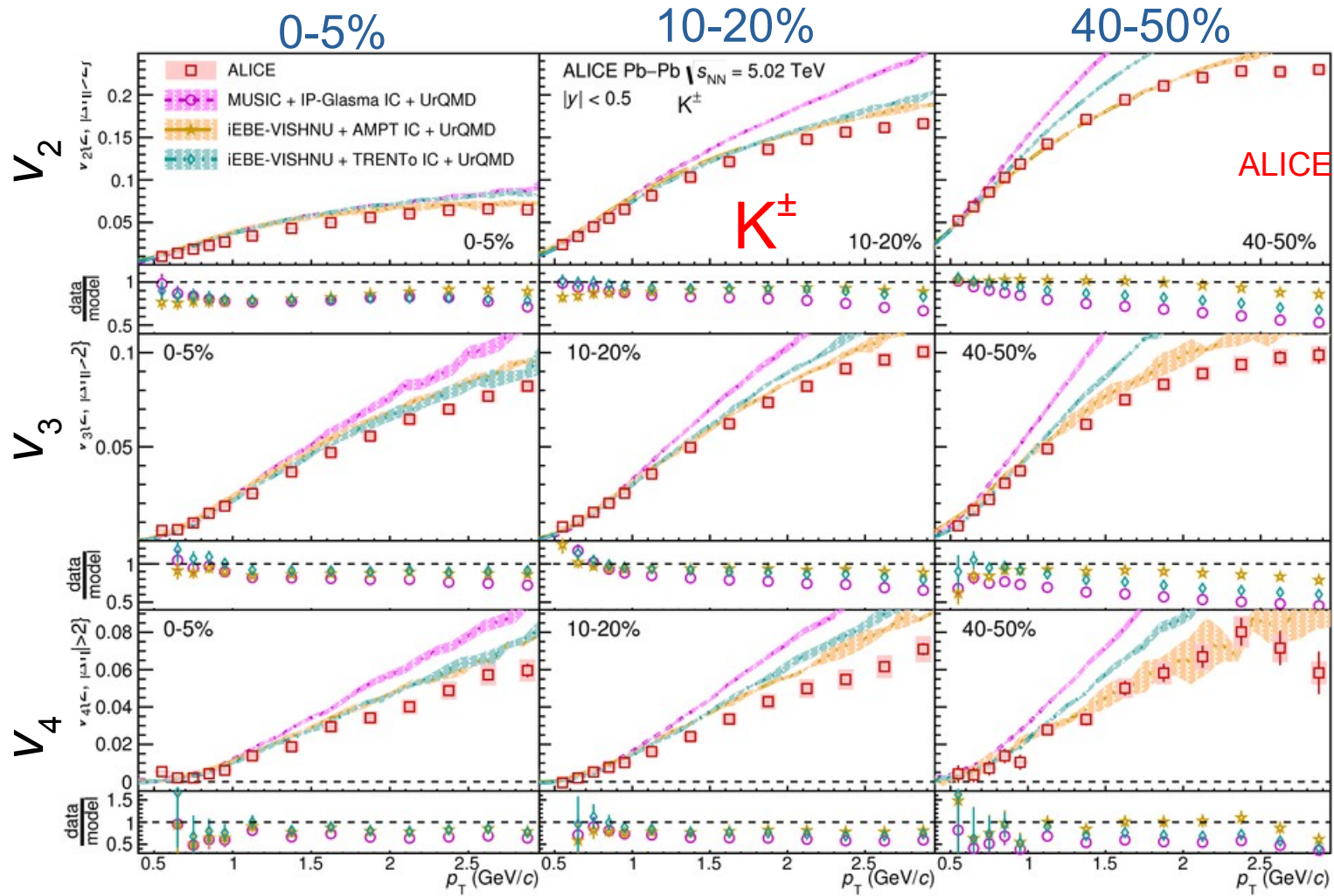
# $\pi^\pm v_n$ vs hydrodynamic calculations



• Hydrodynamic calculations coupled to a hadronic cascade model (UrQMD) describe data at low  $p_T$

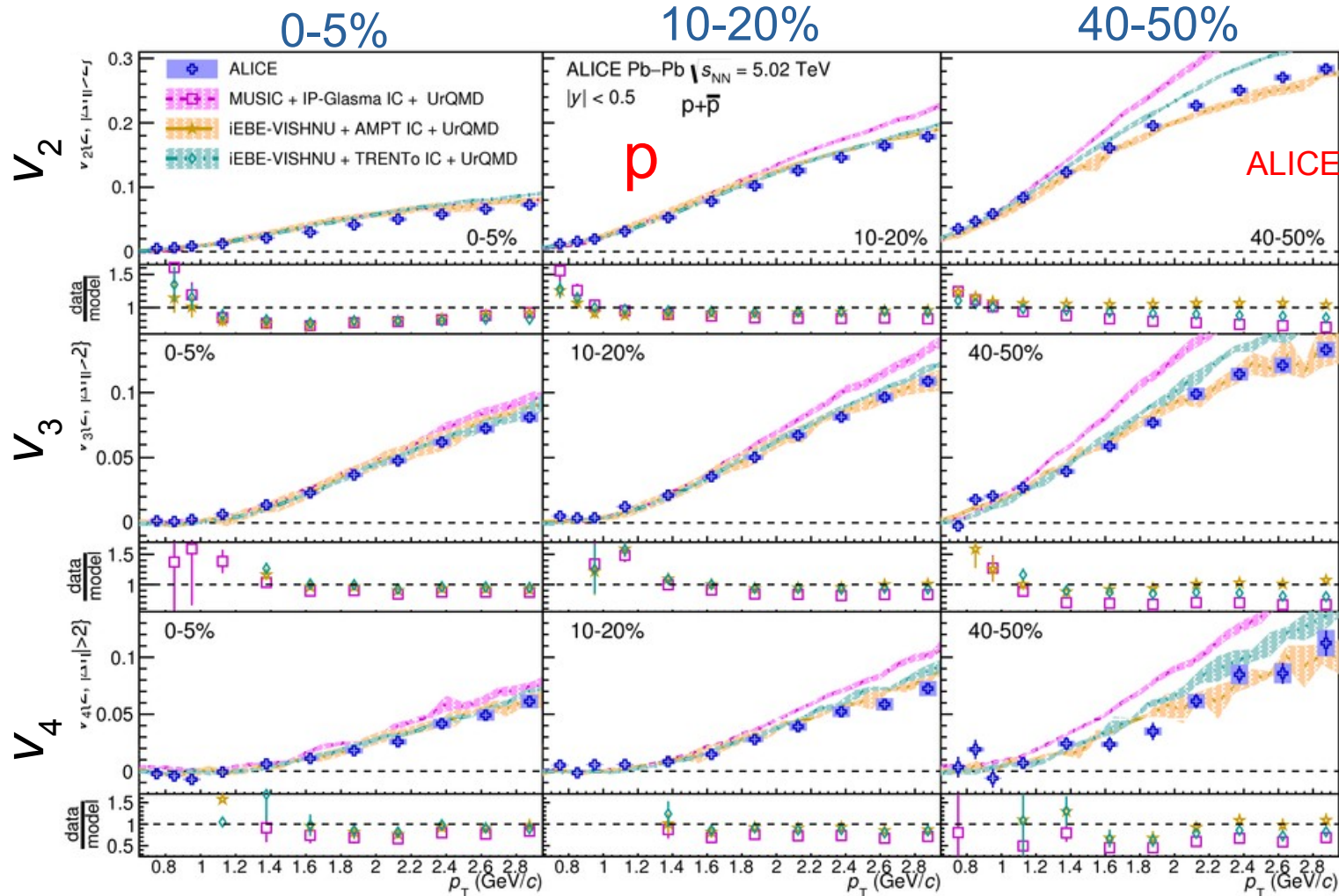
- MUSIC with IP-Glasma IC,  $\zeta/s(T)$ ,  $\eta/s=0.095$  for  $p_T < 1$  GeV/c S. McDonald et al., PRC 95 (2017) 064913
- iEBE-VISHNU with AMPT IC,  $\zeta/s=0$ ,  $\eta/s=0.08$  for  $p_T < 2$  GeV/c W. Zhao et al., EPJC 77 (2017) 645
- iEBE-VISHNU with TRENTo IC,  $\zeta/s(T)$ ,  $\eta/s(T)$  for  $p_T < 1-2$  GeV/c

# $K^\pm v_n$ vs hydrodynamic calculations



- Hydrodynamic calculations coupled to a hadronic cascade model (UrQMD) describe data at low  $p_T$ 
  - MUSIC with IP-Glasma IC,  $\zeta/s(T)$ ,  $\eta/s=0.095$  for  $p_T < 1$  GeV/c S. McDonald et al., PRC 95 (2017) 064913
  - iEBE-VISHNU with AMPT IC,  $\zeta/s=0$ ,  $\eta/s=0.08$  for  $p_T < 2$  GeV/c W. Zhao et al., EPJC 77 (2017) 645
  - iEBE-VISHNU with TRENTo IC,  $\zeta/s(T)$ ,  $\eta/s(T)$  for  $p_T < 1-2$  GeV/c

# $p v_n$ vs hydrodynamic calculations



• Hydrodynamic calculations coupled to a hadronic cascade model (UrQMD) describe data at low  $p_T$

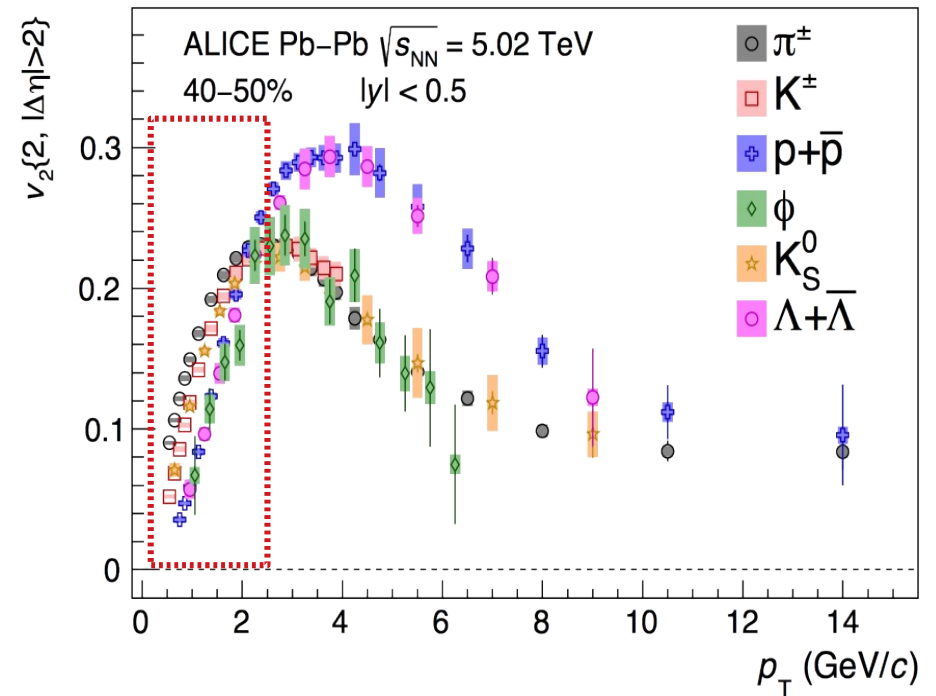
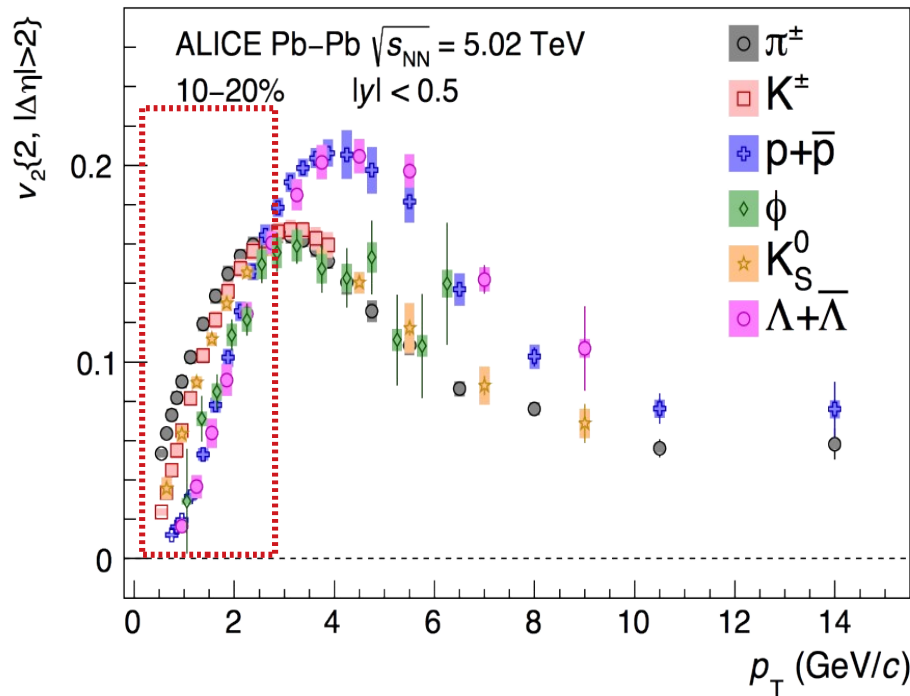
- MUSIC with IP-Glasma IC,  $\zeta/s(T)$ ,  $\eta/s=0.095$  for  $p_T < 1$  GeV/c S. McDonald et al., PRC 95 (2017) 064913
- iEBE-VISHNU with AMPT IC,  $\zeta/s=0$ ,  $\eta/s=0.08$  for  $p_T < 3$  GeV/c W. Zhao et al., EPJC 77 (2017) 645
- iEBE-VISHNU with TRENTo IC,  $\zeta/s(T)$ ,  $\eta/s(T)$  for  $p_T < 2-3$  GeV/c

# PID $v_2(p_T)$

10-20%

ALICE, JHEP 09 (2018) 006

40-50%



- For  $p_T < 2$  GeV/c:  $v_2$  of lighter particles is larger than heavier ones  $\rightarrow$  mass ordering
  - Interplay between elliptic and radial flow
    - Radial flow (isotropic expansion) pushes particles to higher  $p_T$

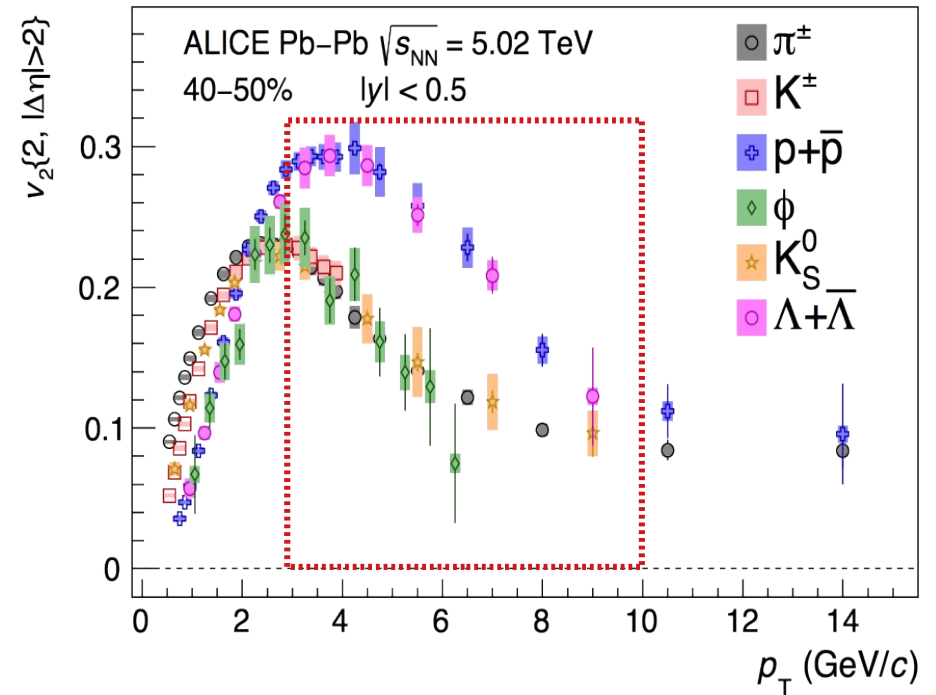
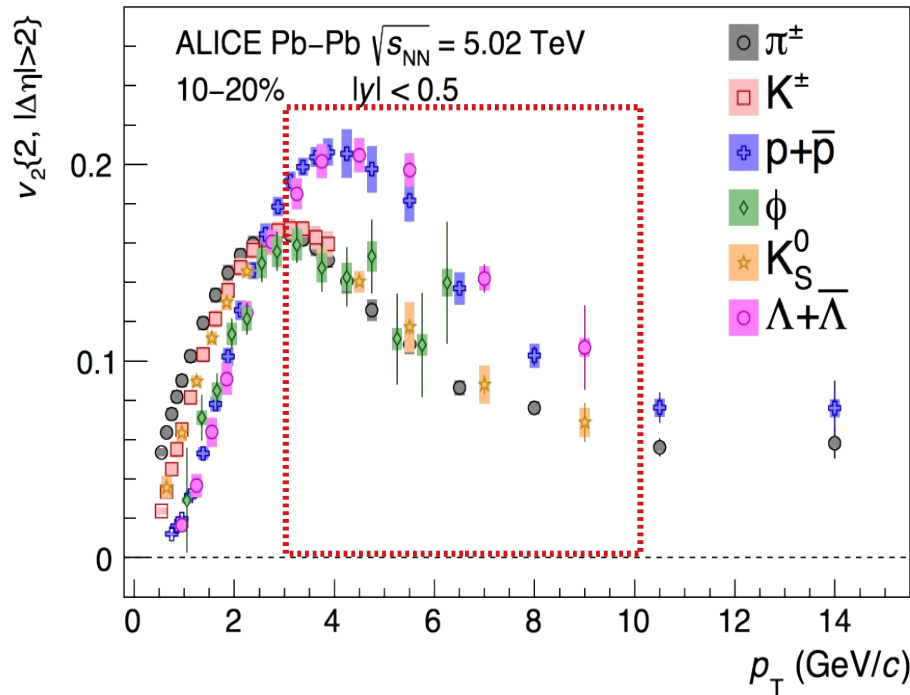


# PID $v_2(p_T)$

10-20%

ALICE, JHEP 09 (2018) 006

40-50%



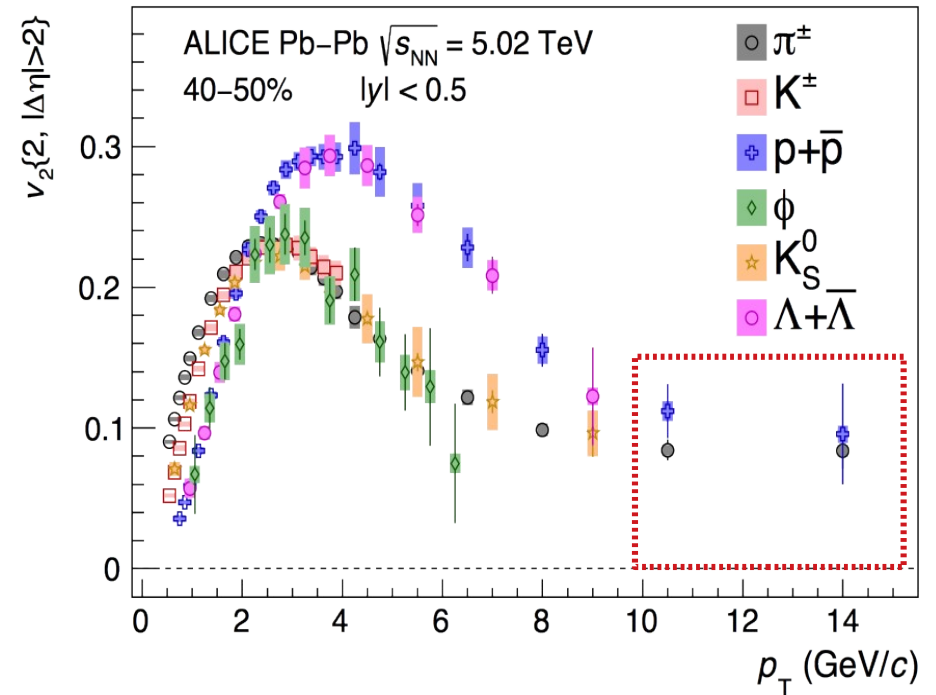
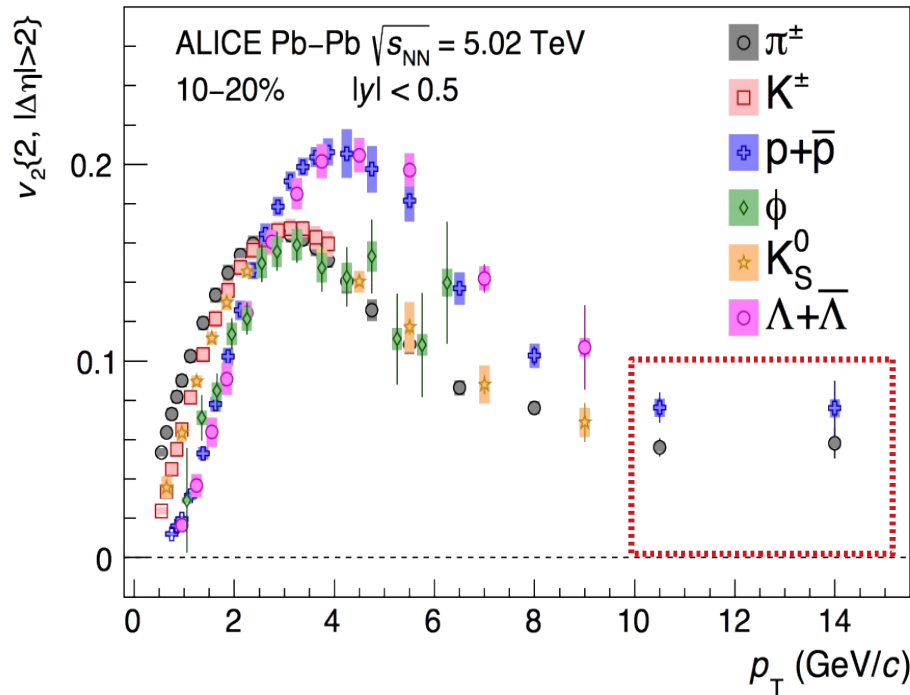
- For  $p_T < 2$  GeV/c:  $v_2$  of lighter particles is larger than heavier ones → mass ordering
  - Interplay between elliptic and radial flow
    - Radial flow (isotropic expansion) pushes particles to higher  $p_T$
- For  $3 < p_T < 10$  GeV/c: particles tend to group into mesons and baryons
- $\phi$ -meson ( $m \sim 1$  GeV/c<sup>2</sup>)  $v_2$  tests both mass ordering and particle type scaling

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10-20%

ALICE, JHEP 09 (2018) 006

40-50%



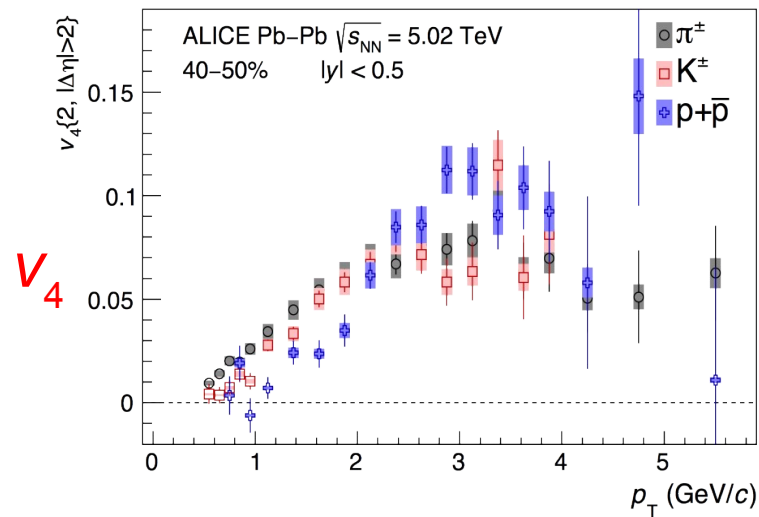
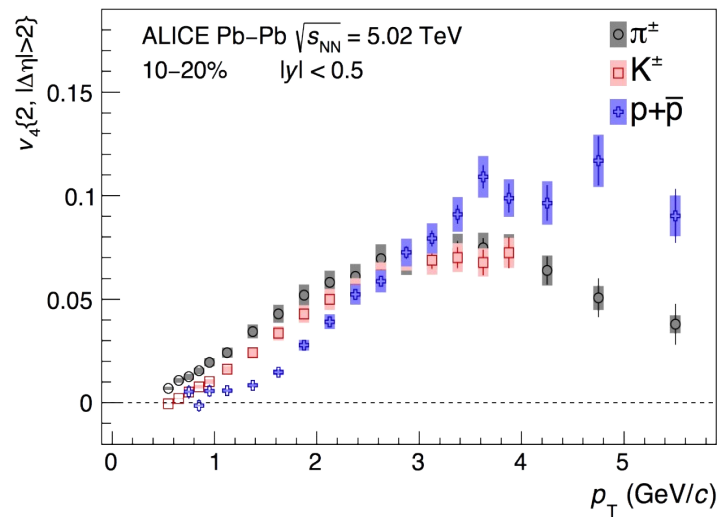
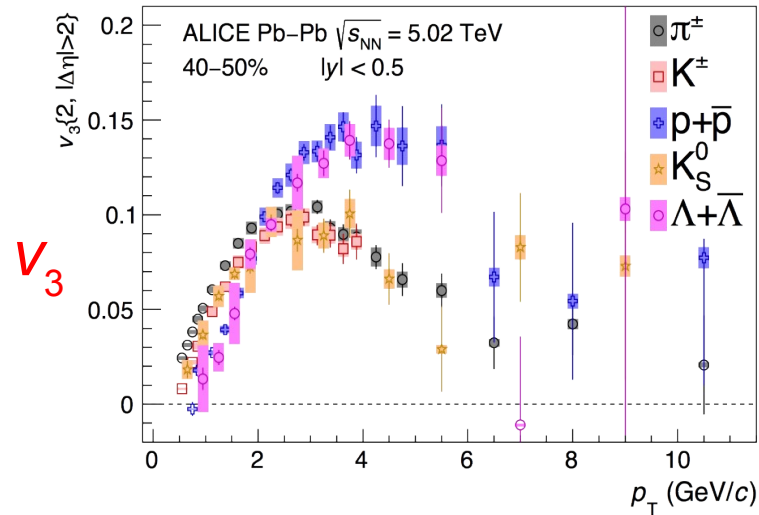
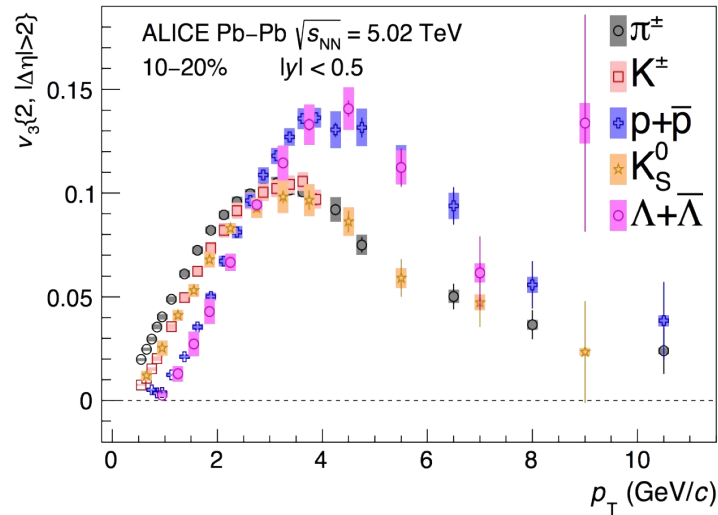
- For  $p_T < 2$  GeV/c:  $v_2$  of lighter particles is larger than heavier ones  $\rightarrow$  mass ordering
  - Interplay between elliptic and radial flow
    - Radial flow (isotropic expansion) pushes particles to higher  $p_T$
- For  $3 < p_T < 10$  GeV/c: particles tend to group into mesons and baryons
- $\phi$ -meson ( $m \sim 1$  GeV/c<sup>2</sup>)  $v_2$  tests both mass ordering and particle type scaling
- For  $p_T > 10$  GeV/c: no particle type dependence within uncertainties
  - Parton energy loss as hadronization mechanism

# PID $v_n(p_T)$

10-20%

ALICE, JHEP 09 (2018) 006

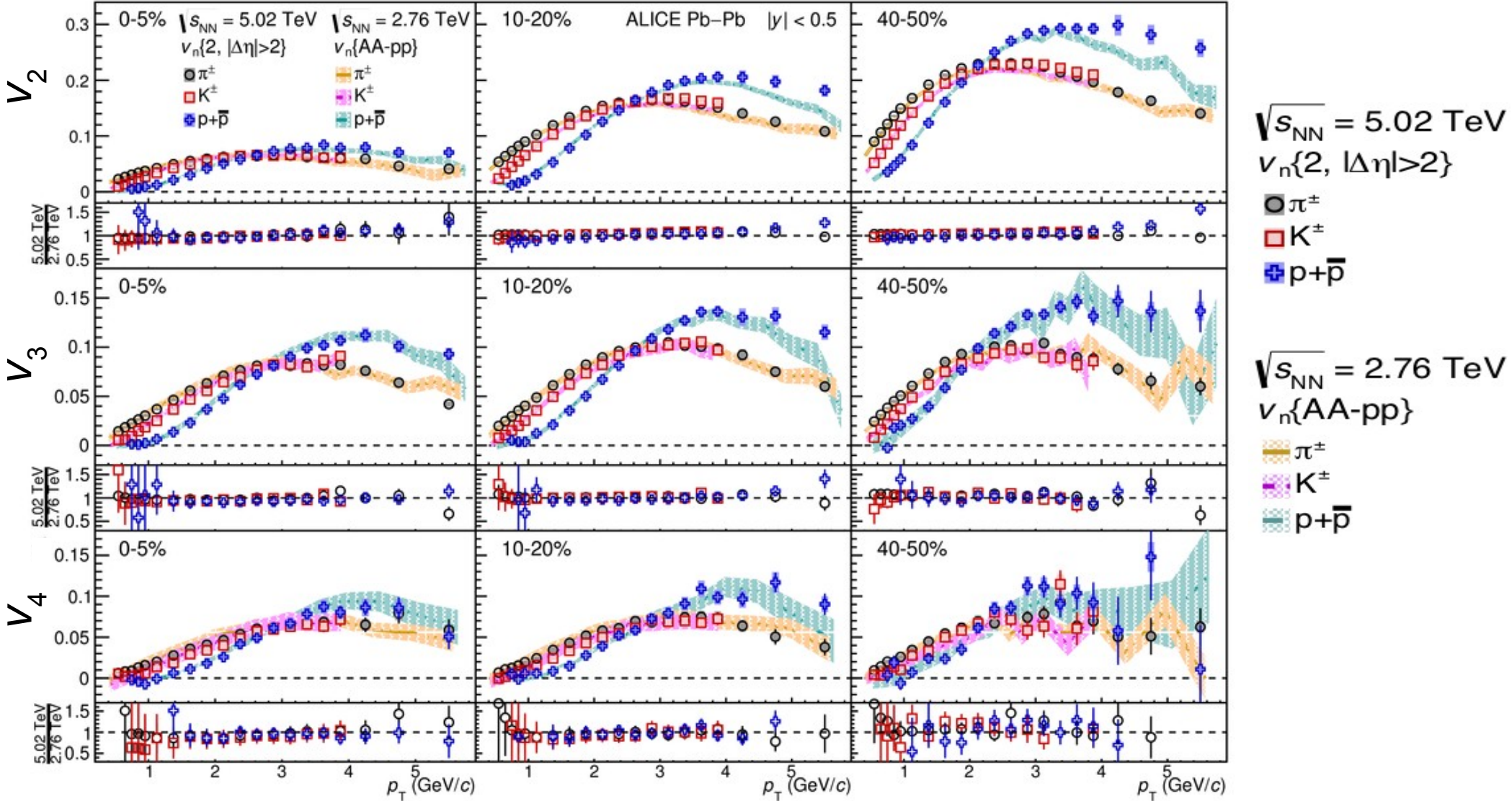
40-50%



- Analogous to the trend of  $v_2$

# PID $v_n(p_T)$ : energy dependence

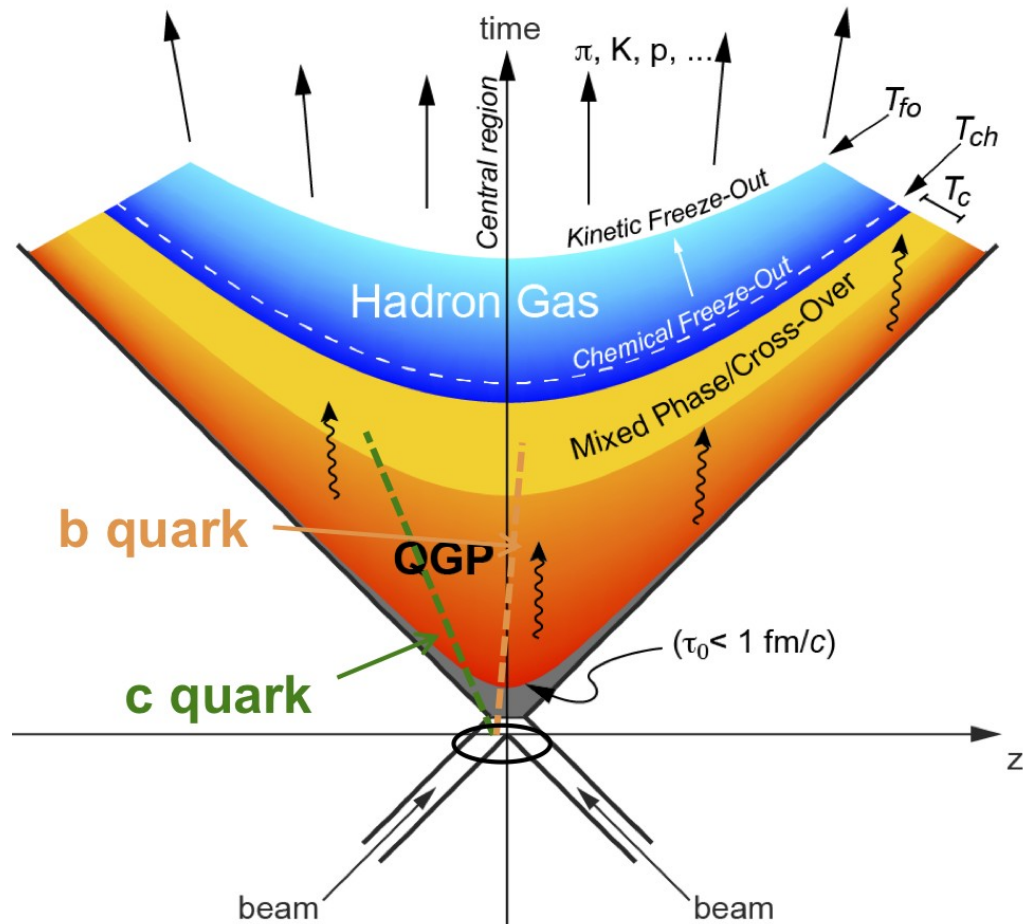
0-5%                      10-20%                      40-50%



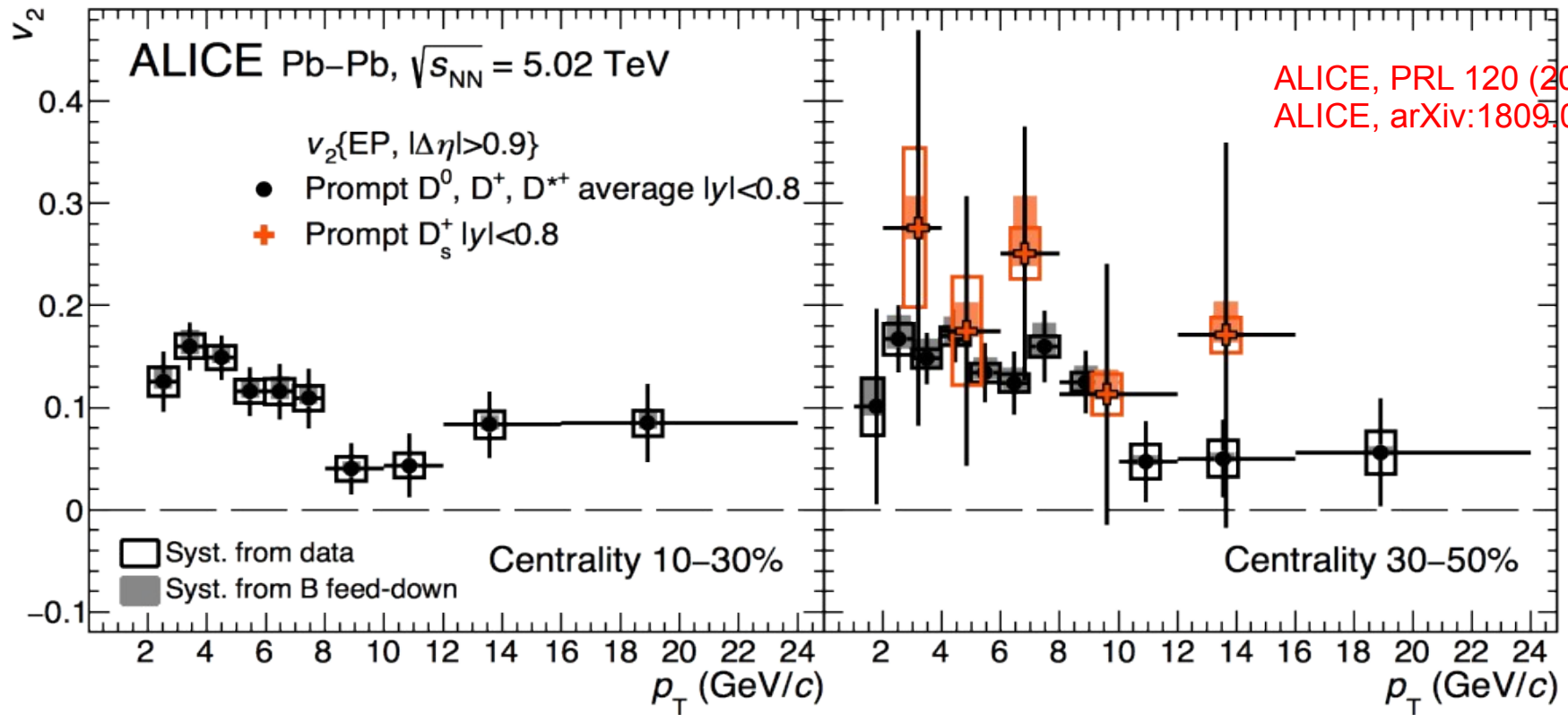
- Measurements are compatible within uncertainties

ALICE, JHEP 09 (2018) 006  
 ALICE, JHEP 09 (2016) 164

# What about heavy quarks?

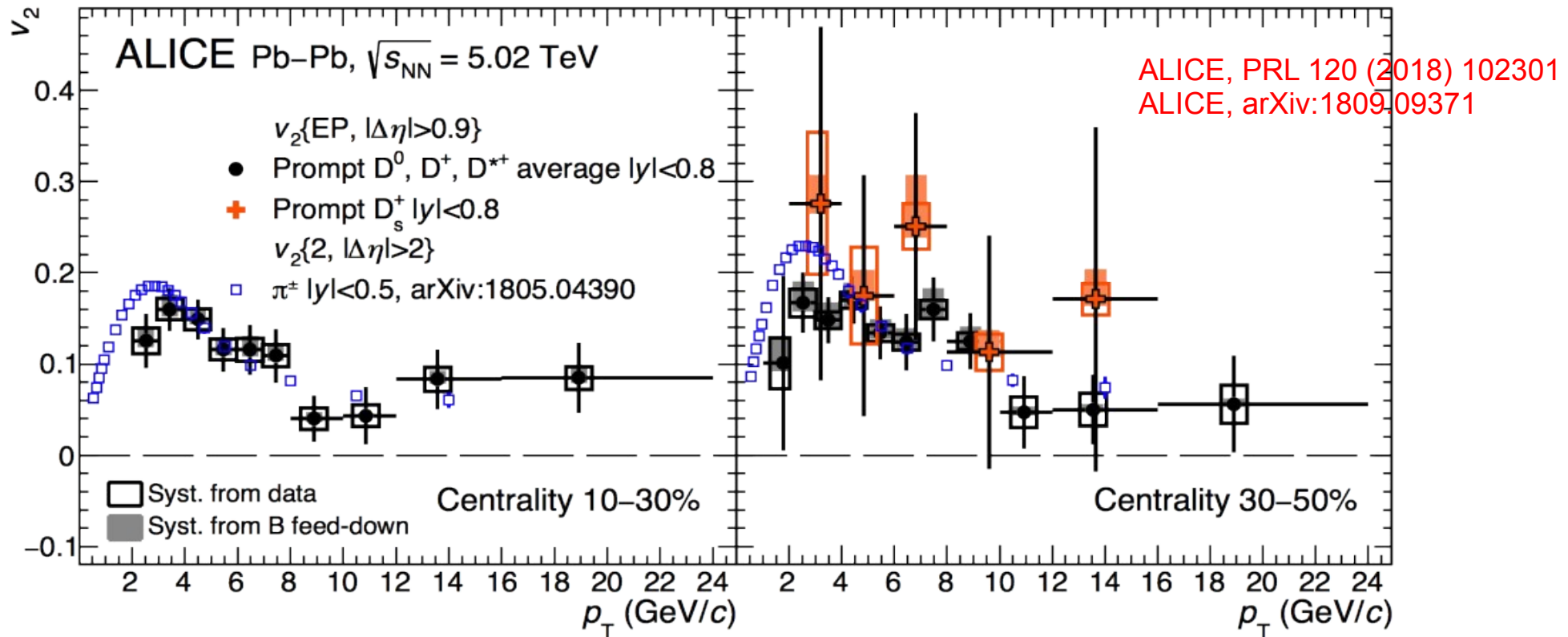


- Heavy quarks produced early  $\rightarrow$  calculable with perturbative QCD
- Large mass  $\rightarrow$  short formation time  $\rightarrow$  probe the evolution of the QGP
  - $1/2m_c$  ( $\sim 0.07 \text{ fm}/c$ )  $<$  QGP formation time ( $\sim 0.1-1 \text{ fm}/c$ )
- How do they interact with the perfect liquid?
  - $\rightarrow v_2$  provides a measure of the heavy-quark diffusion



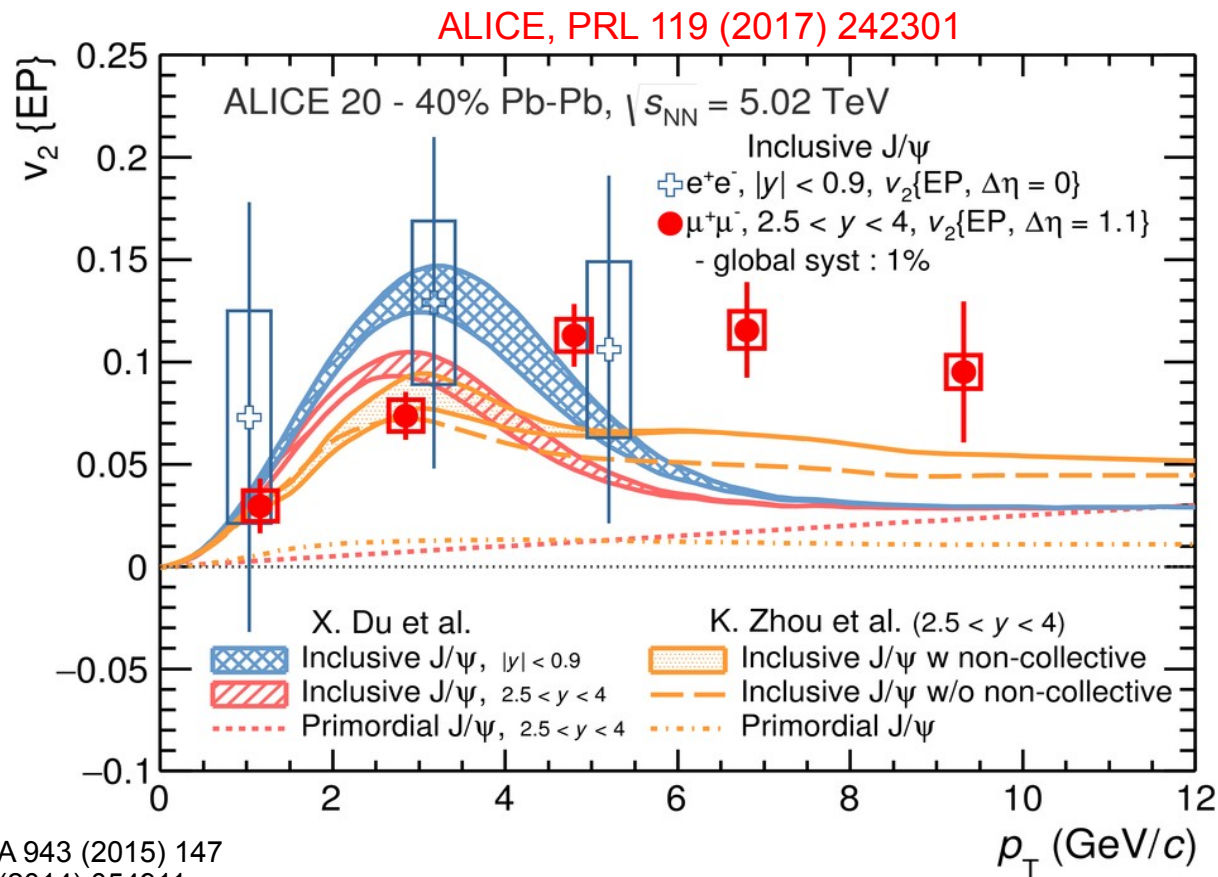
ALICE, PRL 120 (2018) 102301  
ALICE, arXiv:1809.09371

- D-meson  $v_2$  larger than 0 in  $2 < p_T < 10$  GeV/c
  - Indication of strong coupling of c-quark to the medium
- First measurement of  $D_s v_2$  at LHC



- D-meson  $v_2$  larger than 0 in  $2 < p_T < 10$  GeV/c
  - Indication of strong coupling of c-quark to the medium
- First measurement of  $D_s v_2$  at LHC
- D-meson  $v_2$  similar to that of  $\pi^\pm$
- Is light-quark  $v_2$  responsible for D-meson  $v_2$  (via interactions)?

# J/ψ $v_2$

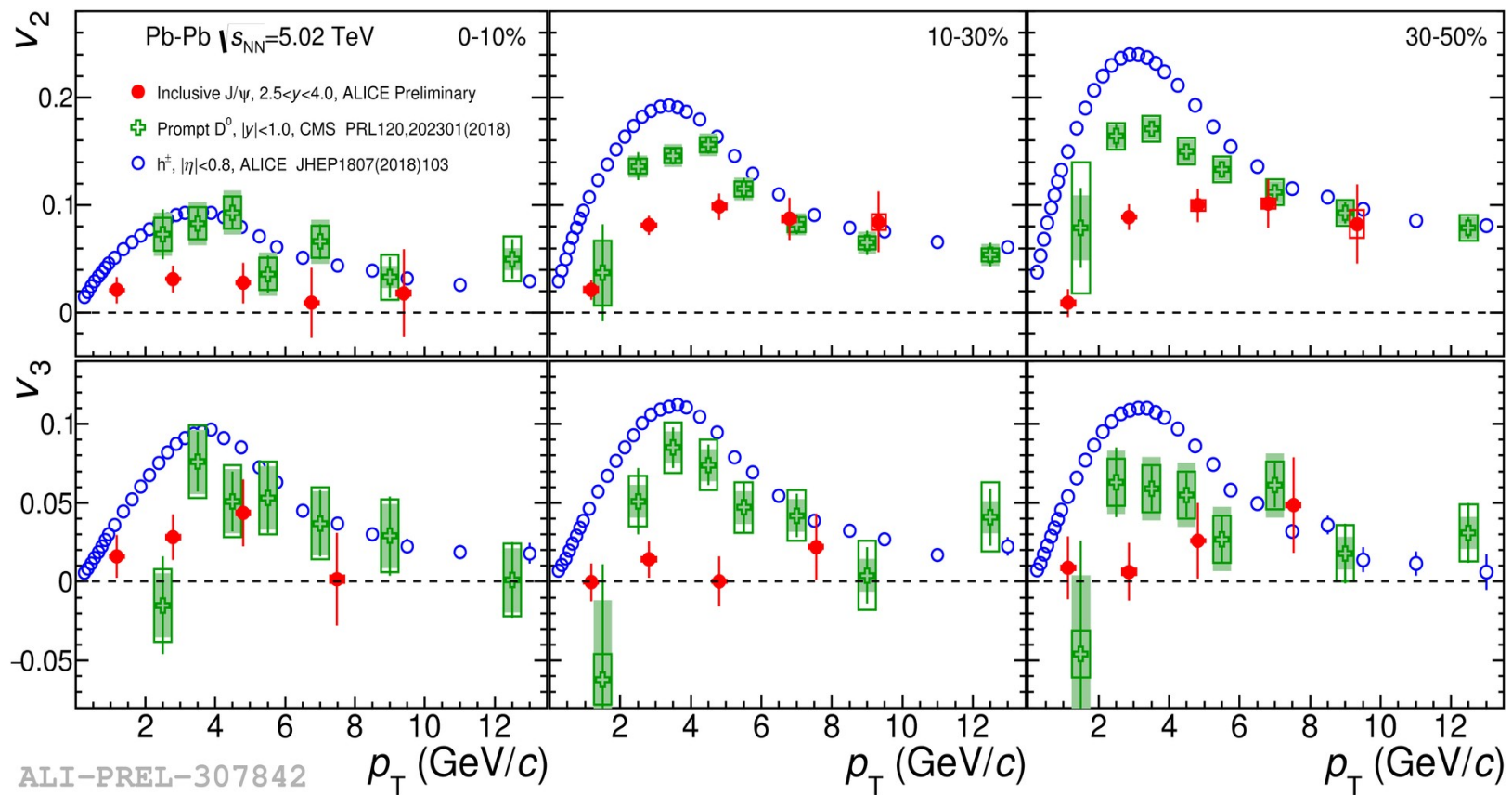


X. Du and R. Rapp, NPA 943 (2015) 147  
 K. Zhou et al., PRC 89 (2014) 054911

- Significant  $v_2$  is observed in different  $p_T$  ranges
- Comparison to transport model calculations
  - Indication of strong coupling of c-quark to the medium

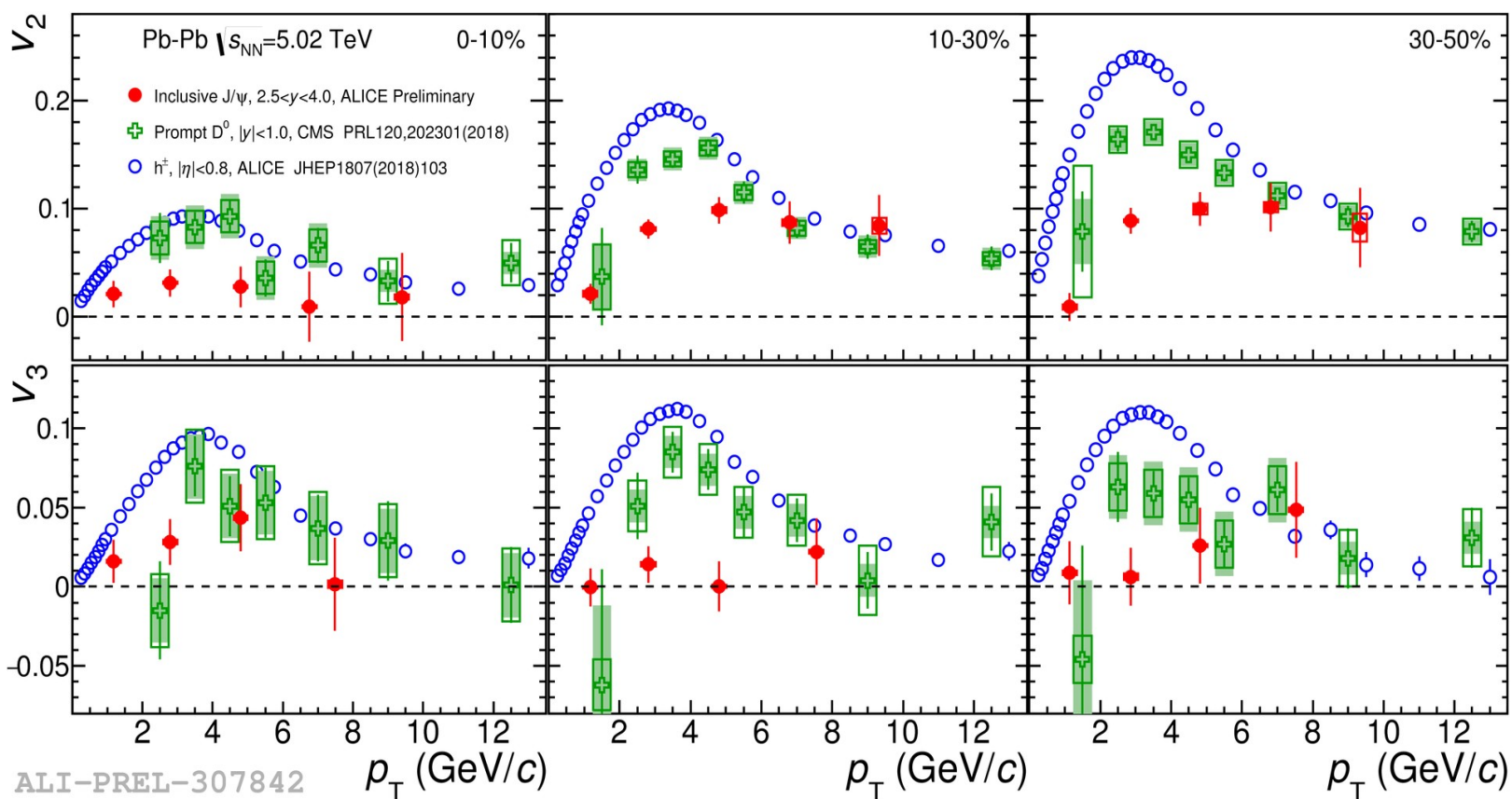


# J/ψ $v_n$

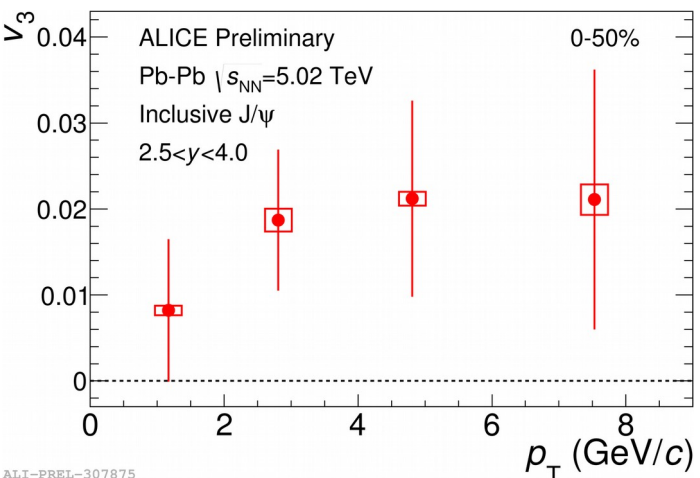


- Comparison between different flavors
  - Clear ordering for  $p_T < 6$  GeV/c:  $v_n(\text{J}/\Psi) < v_n(\text{D}^0) < v_n(\text{h}^\pm)$
  - Convergence for  $p_T > 6$  GeV/c:  $v_n(\text{J}/\Psi) \approx v_n(\text{D}^0) \approx v_n(\text{h}^\pm)$

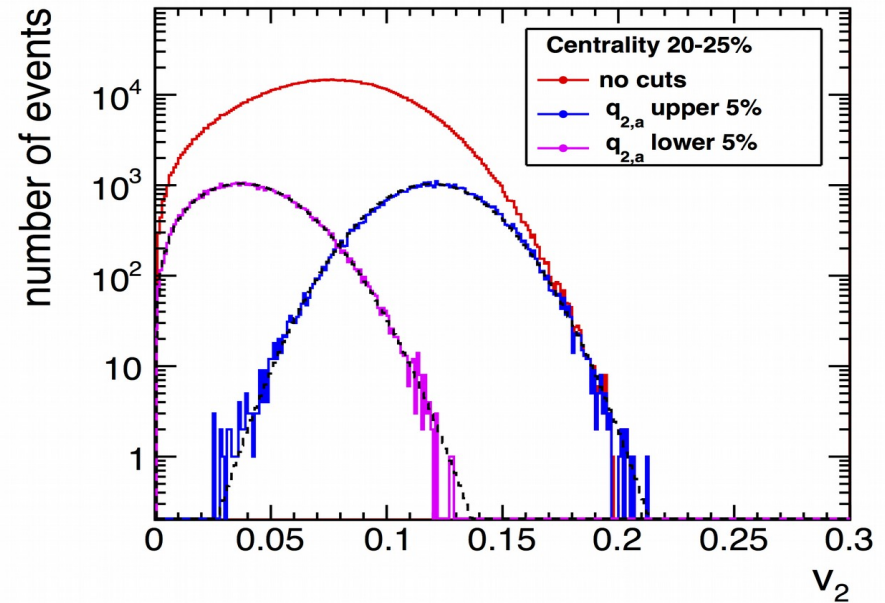
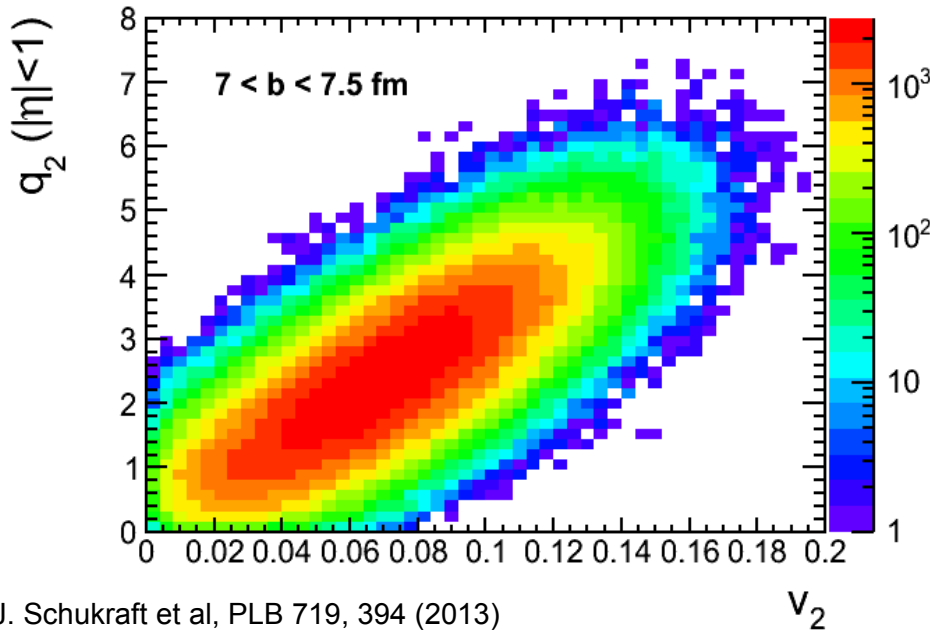
# J/ψ $v_n$



ALI-PREL-307842



- Comparison between different flavors
  - Clear ordering for  $p_T < 6$  GeV/c:  $v_n(J/\Psi) < v_n(D^0) < v_n(h^\pm)$
  - Convergence for  $p_T > 6$  GeV/c:  $v_n(J/\Psi) \approx v_n(D^0) \approx v_n(h^\pm)$
- First evidence for  $v_3(J/\Psi) > 0$  ( $3.7\sigma$  significance)



J. Schukraft et al, PLB 719, 394 (2013)  
 H. Petersen et al, PRC 88, 044918 (2013)  
 P. Huo et al, PRC 90, 024910 (2014)

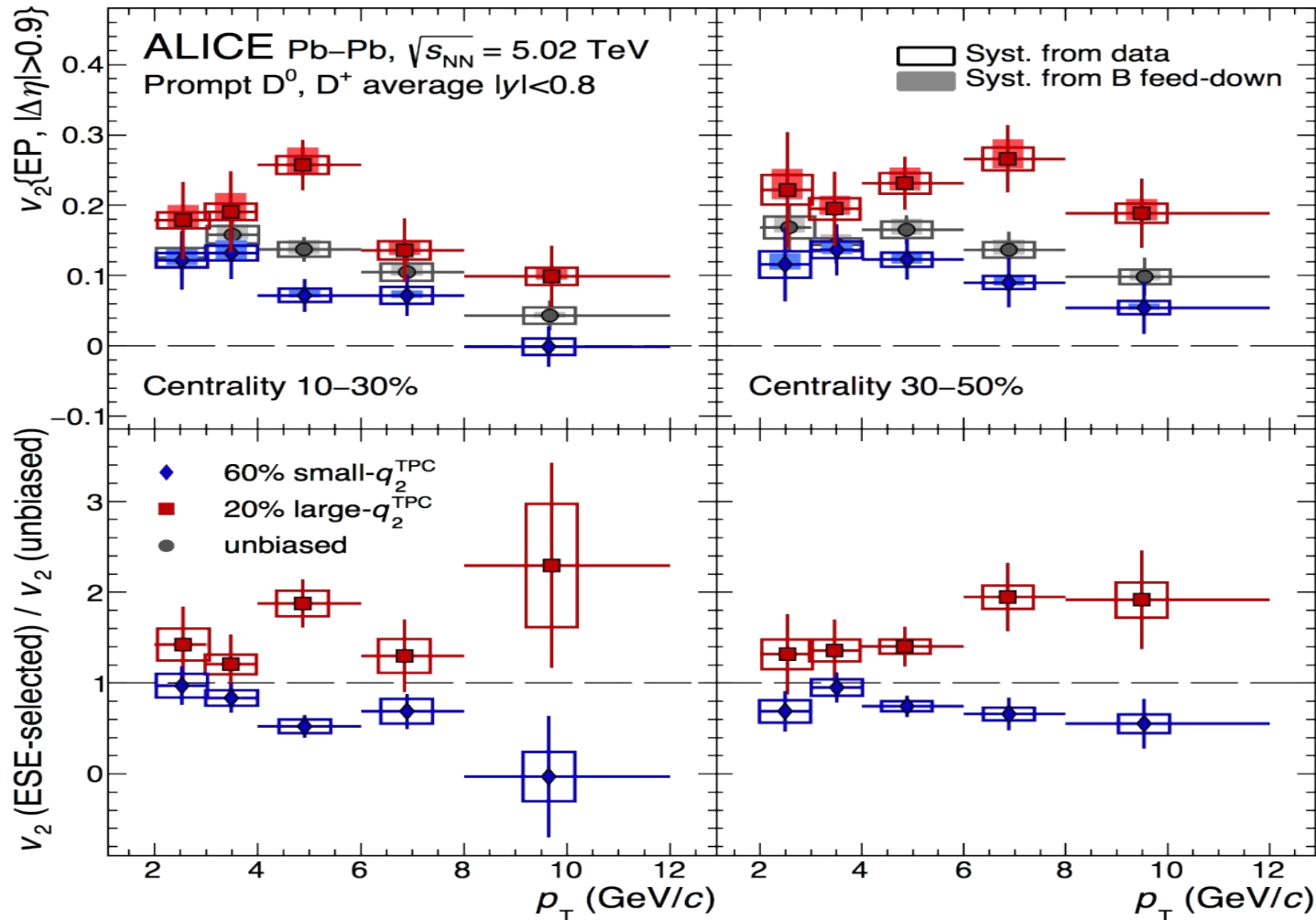
Select events with similar centralities (volume) and different shapes based on the event-by-event flow/eccentricity fluctuations

Flow vector  $\rightarrow$  q-distributions

$$\begin{aligned}
 Q_{n,x} &= \sum_i \cos(n\varphi_i) \\
 Q_{n,y} &= \sum_i \sin(n\varphi_i)
 \end{aligned}
 \rightarrow
 \begin{aligned}
 Q_n &= \{Q_{n,x}, iQ_{n,y}\} \\
 q_n &= |Q_n| / \sqrt{M}
 \end{aligned}$$

# D-meson $v_2$ with ESE

ALICE, arXiv:1809.09371



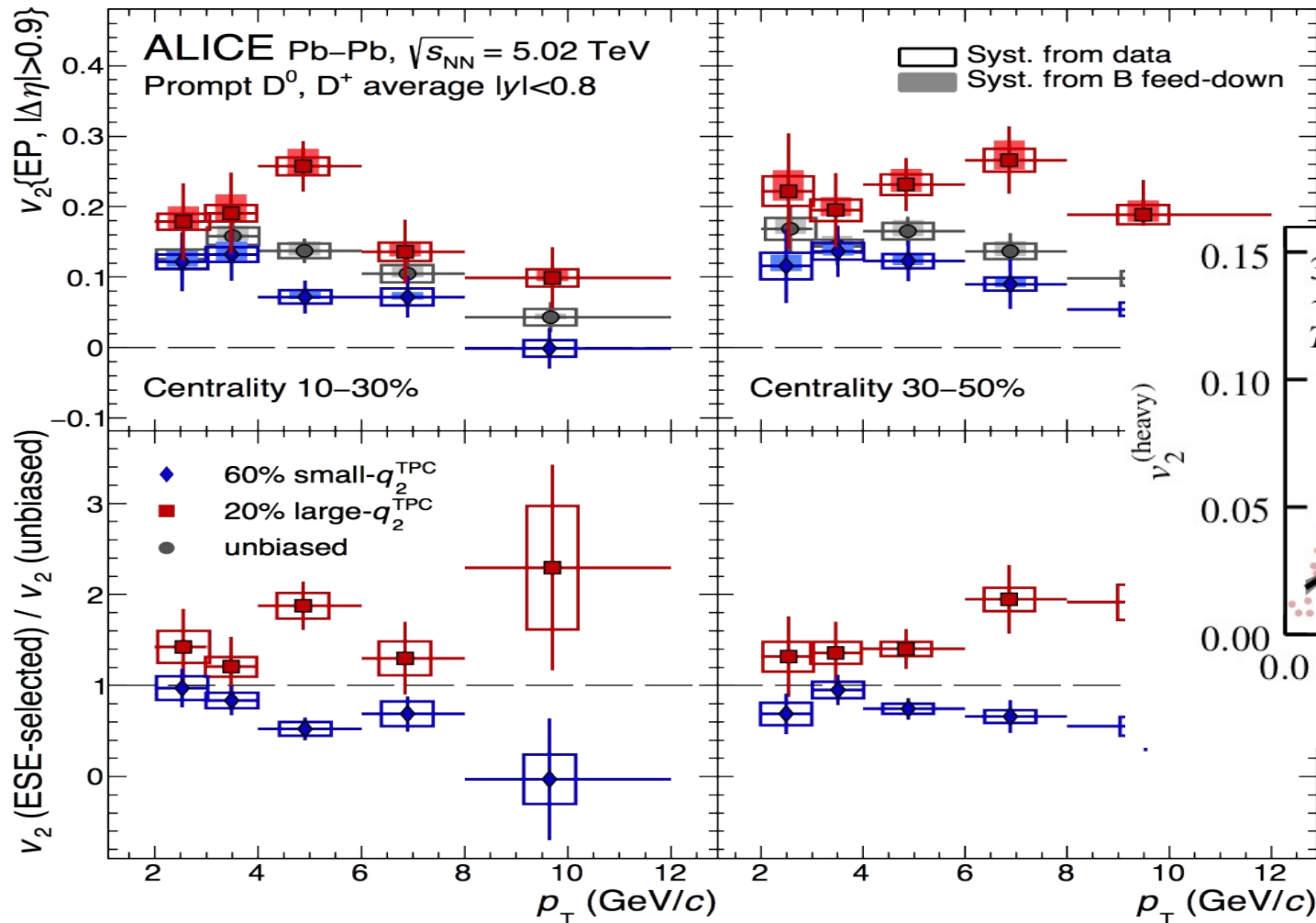
Large- $q_2$ : 20% high

Small- $q_2$ : 60% low

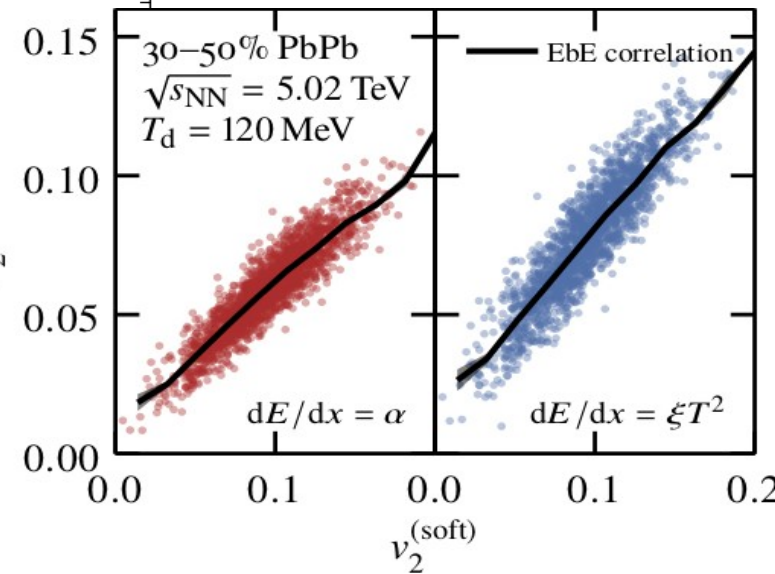
- Correlation between bulk (light charged particles used for ESE) and D-meson  $v_2$ 
  - Charm sensitive to bulk  $v_2$  and initial state fluctuations

# D-meson $v_2$ with ESE

ALICE, arXiv:1809.09371

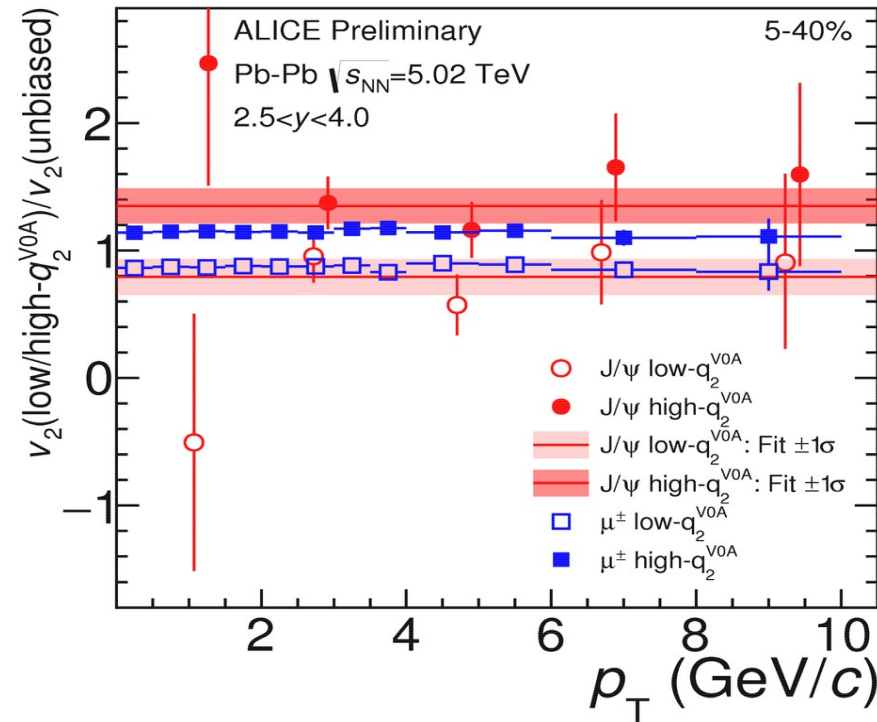
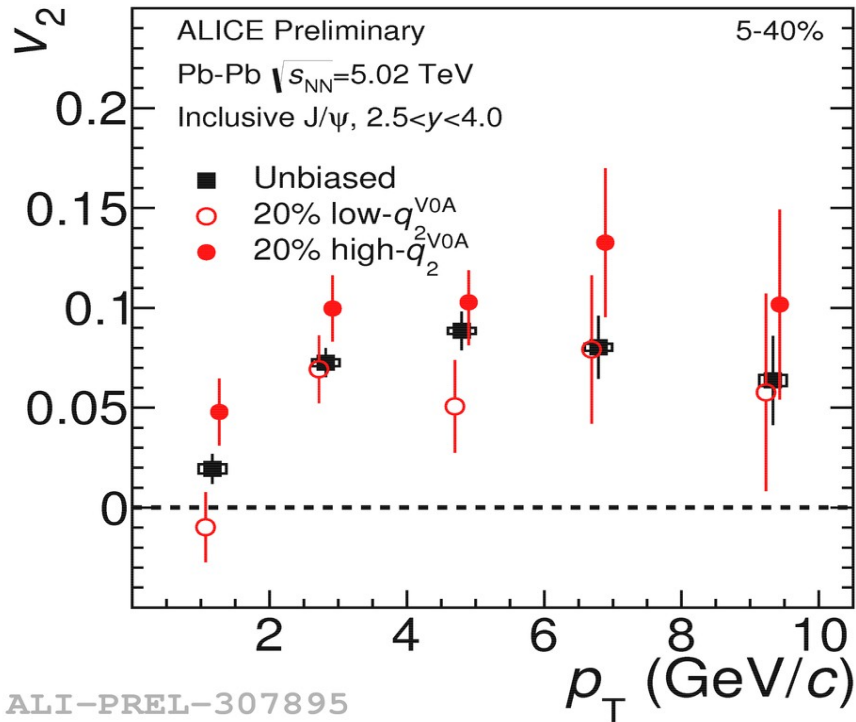


C. Prado et al., PRC 96 (2017) 064903



- Correlation between bulk (light charged particles used for ESE) and D-meson  $v_2$ 
  - Charm sensitive to bulk  $v_2$  and initial-state fluctuations
- Further constraints on the theory

# J/Ψ $v_2$ with ESE

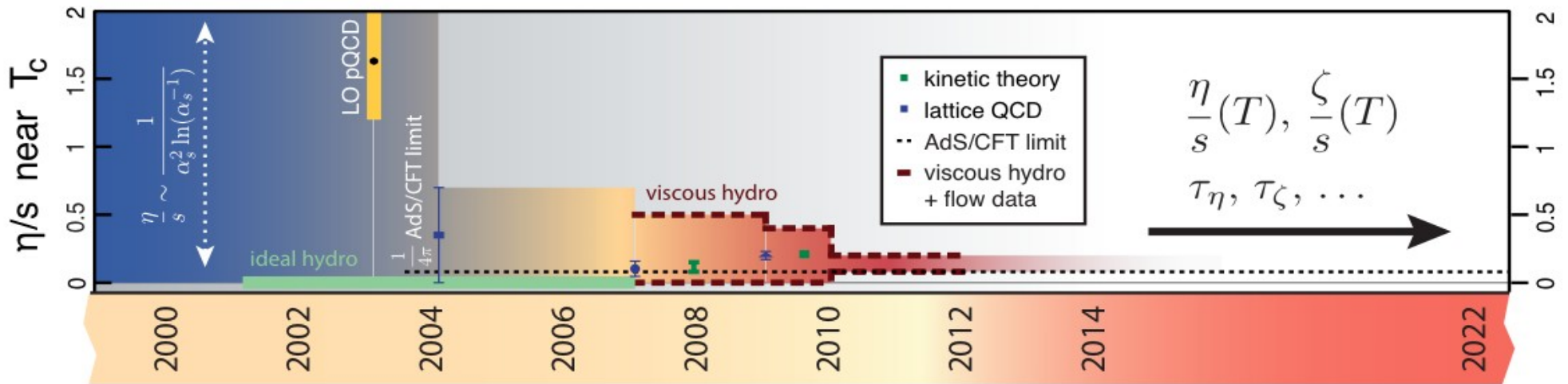


Large- $q_2$ :  
20% high

Small- $q_2$ :  
20% low

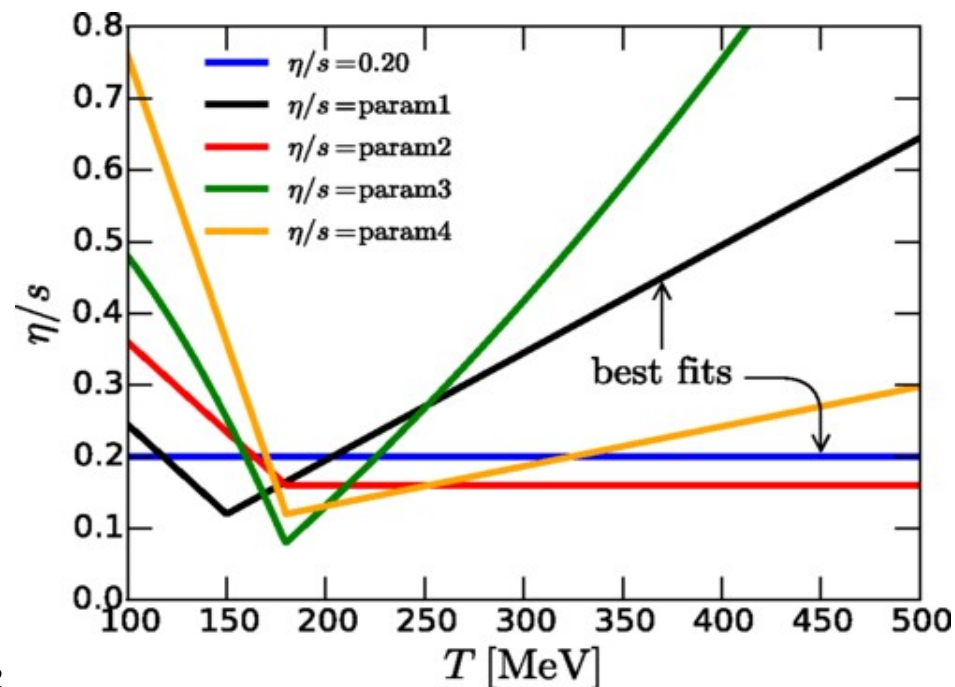
- J/Ψ  $v_2$  is larger or smaller than the average with ESE
- Ratios (ESE/unbiased) of J/Ψ  $v_2$  consistent with those of single muons within uncertainties
  - J/Ψ  $v_2$  compatible with the expected variations of the eccentricity

- The properties of the QGP and parts of the QCD phase diagram are understood much better



C. Gale et al, Int. J. Mod. Phys. A 28 (2013) 1340011

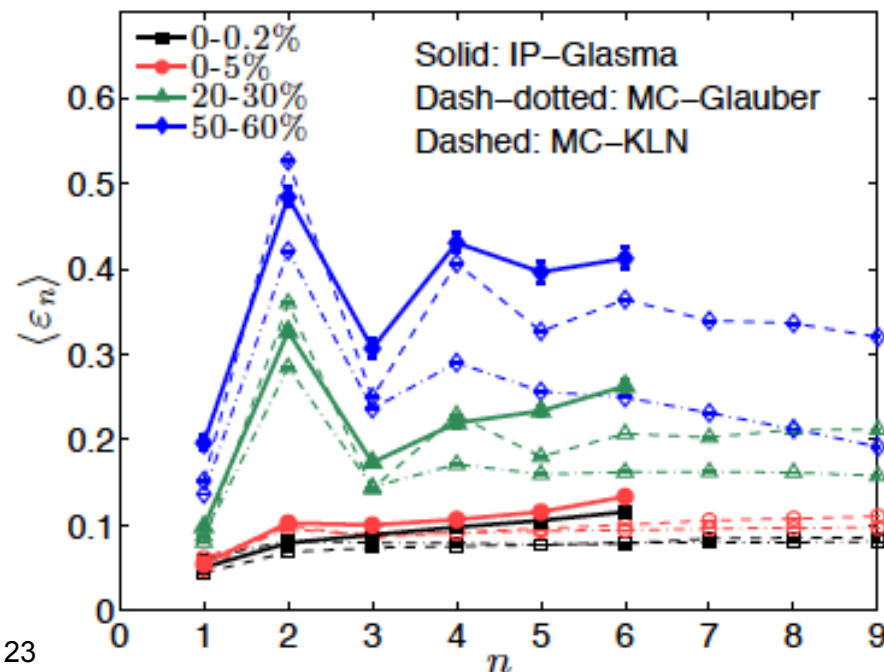
- The properties of the QGP and parts of the QCD phase diagram are understood much better
  - Start constraining using  $v_n$  for inclusive and identified particle and their energy dependence
    - Temperature dependence of shear viscosity





# Summary

- The properties of the QGP and parts of the QCD phase diagram are understood much better
  - Start constraining using  $v_n$  for inclusive and identified particle and their energy dependence
    - Temperature dependence of shear viscosity
    - Initial conditions



U. Heinz and R. Snellings,  
Ann.Rev.Nucl.Part.Sci. 63 (2013) 123

- Further constrain temperature dependence of shear viscosity and equation-of-state using Bayesian statistics
- Collectivity in small systems

