### **Critical Opalescence: A Smoking Gun Signature for a Critical Point**

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**based on arXiv:0903.0669 [nucl-th]**

# **4 steps 4 definitive CEP measurements**



#### What about

- random fields?
- experimentally measurable order pars ?
- 1st order PT: speed of sound, latent heat?

#### 1. Identify:

What type of transition chiral ? deconfinement ? quarkionic ? liquid-gas?

2. Locate: Where is  $({\sf T}_{_{E^{\prime}}}\,\mu_{_{E}})?$ At what centrality,  $\sqrt{s_{NN}}$  ? critical opalescence onset of 1<sup>st</sup> order PT

> 3. Characterize: measure order parameters, critical exponents, universality classes.

4. Controll: Cross-checks for consistency, significance, quality.

### **Correlations for VARIOUS Quark Matters**

#### Transition to hadron gas may be:

1<sup>st</sup> order (strong) 2<sup>nd</sup> order (Critical Point, CP) Cross-over Non-equilibrium, e.g. from a supercooled state (scQGP)

#### Type of phase transition: its correlation signature:

Strong 1<sup>st</sup> order QCD phase transition: (Pratt, Bertsch, Rischke, Gyulassy)  $R_{out} >> R_{side}$ 

2<sup>nd</sup> order QCD phase transition: (T. Cs, S. Hegyi, T. Novák, W.A. Zajc) non-Gaussian shape

Cross-over quark matter-hadron gas: (lattice QCD, Buda-Lund hydro fits) hadrons appear from

Supercooled QGP (scQGP) -> hadrons: (T. Cs, L.P. Csernai) example to pion flash  $(R_{out} \sim R_{side})$ 

α(Lévy) decreases to 0.5

a region with  $T > T_c$ 

same freeze-out for all particles strangeness enhancement no mass-shift of  $\phi$ 

### **Excitation of 3d Gaussian fit parameters**

#### STAR, Phys.Rev.C71:044906,2005



#### These data indicate

$$
\mathsf{R}_{\text{out}} \sim \mathsf{R}_{\text{side}}
$$

hence exclude: Strong, equilibrium 1st order phase transit. > 50 hydro models

For a second order PT:

check excitation function of non-Gaussian parameter  $α$ 

New analysis / new data are needed

HBT Radii independent of energy perhaps initial volume ? subtle mt dependencies?

### **Excitation of 3d Gaussian fit parameters**



PHENIX, Phys. Rev. Lett. 93, 152302 (2004)

### **Critical Opalescence**

#### **Critical Opalescence: a laboratory method to observe a 2nd order PT**

**correlation length diverges, clusters on all scales appear incl. the wavelength of the penetrating (laser) probe**

**side view:** http://www.msm.cam.ac.uk/doitpoms/tlplib/solidsolutions/videos/laser1.mov



**front view: matter becomes opaque at the critical point (CP)**



### **Opical opacity** κ **and attenuation length** λ

$$
I = I_0 \exp(-\kappa x) = I_0 \exp(-x/\lambda)
$$
  
\n
$$
\kappa = \frac{I(\text{generated}) - I(\text{transmitted})}{I(\text{generated})\Delta x}
$$
  
\n
$$
R_{AA} = \frac{I(\text{transmitted})}{I(\text{generated})} = \frac{I(\text{measured})}{I(\text{expected})}
$$
  
\n
$$
I(\text{measured}) = \frac{1}{N_{event}^{AA}} \frac{d^2 N_{AA}}{dy d p_t}
$$
  
\n
$$
I(\text{expected}) = \frac{\langle N_{coll} \rangle}{\sigma_{inel}^{NN}} \frac{d^2 \sigma_{NN}}{dy d p_t}
$$

$$
\frac{\partial I}{\partial x} = -\kappa I
$$

$$
\kappa = -\frac{\ln(R_{AA})}{R_{HBT}}
$$

### **Opical opacity** κ **and attenuation length** λ

$$
\kappa = -\frac{\ln(R_{AA})}{R_{HBT}} \qquad \lambda = \frac{1}{\kappa} = -\frac{R_{HBT}}{\ln(R_{AA})}
$$



**Table 1:** Examples of opacities  $\kappa$  and attenuation lengths  $\lambda$  in  $\sqrt{s_{NN}}$  = 200 GeV Au+Au reactions, evaluated from nuclear modification factors measured at  $p_t = 4.75$  GeV/c in ref. [36] and using HBT radii of ref. [37], averaged over both directions and charge combinations, at the same centrality class as  $R_{AA}$ .

No max in opacity or min in attenuation length is seen wrt centrality Contrast: for a 5 GeV  $\gamma$  on lead,  $\lambda = 6.2$  mm **RHIC perfect fluid is more opaque (to jets), than lead (to** γ**) - by 12 orders of magnitude** for details: see nucl-th/0903.0669v3

Critical Opalescence: a smoking gun signature of a 2<sup>nd</sup> order PT New experimental definition of opacity / attenuation length: A combination of attenuation  $(R_{AA})$  and lengthscale (e.g.  $R_{HBT}$ ) is needed



$$
I = I_0 \exp(-\kappa x) = I_0 \exp(-x/\lambda)
$$

$$
\kappa = -\frac{\ln(R_{AA})}{R_{HBT}}
$$

CPOD@BNL, 10/06/09 T. Csörgő Use optical opacity and look for maximal opalescence! Alternative lengthscale measurement:  $R(HBT)$  =  $p_{_{0}} + p_{_{1}}N_{_{part}}$ 1/3 Estimate Npart and take  $\bm{{\mathsf{p}}}_{_{\!0}}$  and  $\bm{{\mathsf{p}}}_{_{\!1}}$  from HBT measurements Possible: azimuthally sensitive  $\mathsf{R}_{_{\mathsf{A}\mathsf{A}}}$  and  $\mathsf{R}_{_{\mathsf{H}\mathsf{B}\mathsf{T}}}$  : azimuthally sensitive opacity Further refinements: beyond Gaussian approximation, e.g. use R(Lévy)

# **Characterizing critical phenomena**

Theoretical order parameter of QCD - quark condensate:<br>Experimental order parameter is needed:

- Experimental order is needed: The formulation of in-medium mass-shift
	- for deconfinement, signal of quark degrees of freedom

Understandable in laymen's terms: quark scaling of particle ratios

$$
\frac{\overline{\Lambda}|\overline{\Sigma}}{\Lambda|\Sigma} = \frac{\overline{p}}{p} \left[\frac{K}{\overline{K}}\right],
$$
\n
$$
\frac{\overline{\Xi}}{\Xi} = \frac{\overline{p}}{p} \left[\frac{K}{\overline{K}}\right]^2,
$$
\n
$$
\frac{\overline{\Lambda}|\overline{\Sigma}}{\Xi} = \frac{\overline{p}}{p} \left[\frac{K}{\overline{K}}\right]^3,
$$
\n
$$
\frac{\overline{\Lambda}|\overline{\Sigma}}{\Lambda|\Sigma} = \left[\frac{\overline{Q}}{Q}\right]^2 \left[\frac{\overline{Q}}{Q}\right]^2
$$
\n
$$
\frac{\overline{\Xi}}{\Xi} = \left[\frac{\overline{Q}}{Q}\right]^2 \left[\frac{\overline{Q}}{Q}\right]^2
$$
\n
$$
\frac{\overline{Q}}{\Xi} = \left[\frac{\overline{Q}}{Q}\right] \left[\frac{\overline{Q}}{Q}\right]^2
$$

$$
\frac{\overline{p}}{p} = \left[\frac{\overline{Q}}{Q}\right]^3,
$$
  

$$
\frac{\overline{\Lambda}|\overline{\Sigma}}{\Lambda|\Sigma} = \left[\frac{\overline{Q}}{Q}\right]^2 \left[\frac{\overline{S}}{S}\right],
$$
  

$$
\frac{\overline{\Xi}}{\overline{\Xi}} = \left[\frac{\overline{Q}}{Q}\right] \left[\frac{\overline{S}}{S}\right]^2,
$$
  

$$
\frac{\overline{\Omega}}{\overline{\Omega}} = \left[\frac{\overline{S}}{S}\right]^3,
$$
  

$$
\frac{\overline{K}}{K} = \frac{\overline{Q}S}{Q\overline{S}}.
$$

J. Zimányi, T. Biró, T.Cs, P. Lévai, Phys.Lett.B472:243-246<br>A. Bialas, Phys.Lett.B442:449-452,1998 A. Bialas, Phys.Lett.B442:449-452,1998

#### **Order parameter for chiral symmetry from HBT, using**  $\lambda$ **(m<sub>t</sub>, √s<sub>NN</sub>)**



#### **Order parameter for deconfinement from identified particle v<sub>2</sub>**



Idea: look for the break down of the quark number scaling If scaling: quark degrees of freedom are active (exp. view) Measure:  $v_2/n_q$  as a function of  $KE_\tau/n_q$  of identified particles needs high statistics PID measurement at low  $\sqrt{s}_{_{\sf NN}}$ 

# **Critical Exponents at 2nd order PT**

Relevant and important quantities: critical exponents, universality classes Reduced temperature:  $t = (T - T_c)/T_c$ Exponent of specific heat:  $|C(T) \sim |t|^{-\alpha} + \text{less singular.}$  $\epsilon$  and  $\epsilon$ Exponent of order  $\langle |\phi| \rangle \sim |t|^{\beta}$  for  $t < 0$ parameter: Exponent of correlation  $\xi \sim |t|^{-\nu}$ length : Exponent of susceptibility:

 $\int d^3x \; G_{\alpha\beta}(x) \sim t^{-\gamma}$ 

# **Critical Exponents (2)**

Exponent of the Fourier-transformed correlation function:

Exponent of order parameter in external field:

$$
G_{\alpha\beta}(k \to 0) \sim k^{-2+\eta}
$$

$$
\langle |\phi| \rangle (t=0, H \rightarrow 0) \sim H^{1/\delta}
$$

 There are thus 6 critical exponents, α, β, γ, δ, ν, η but only 2 are independent:

Exponents <-> universality class!

$$
\alpha = 2 - d\nu
$$
  

$$
\beta = \frac{\nu}{2}(d - 2 + \eta)
$$
  

$$
\gamma = (2 - \eta)\nu
$$
  

$$
\delta = \frac{d + 2 - \eta}{d - 2 + \eta}.
$$

# **Two particle Interferometry**

#### for non-interacting identical bosons



$$
A_{12} = \frac{1}{\sqrt{2}} [e^{ip_1 \cdot (r_1 - x)} e^{ip_2 \cdot (r_2 - y)} + e^{ip_1 \cdot (r_1 - y)} e^{ip_2 \cdot (r_2 - x)}]
$$
  
so that

$$
\mathcal{P}_{12} = \int d^4\pmb{x} \, d^4\pmb{y} \, |A_{12}|^2 \rho(\pmb{x}) \rho(\pmb{y}) = 1 + |\tilde{\rho}(q)|^2 \equiv C_2(q)
$$

#### emission function

$$
C(p_1, p_2) = 1 + \frac{\left| \int d^4x \cdot S(x, K) \cdot e^{iq \cdot x} \right|^2}{\left| \int d^4x \cdot S(x, K) \right|^2}
$$

$$
q = p_1 - p_2
$$

$$
q = p_1 - p_2
$$
  $K = \frac{1}{2}(p_1 + p_2)$ 

## **Correlations for various collisions**



Correlations have more information (3d shape analysis) Use advanced techniques & extract it ( $\sim$  medical imaging)

# **CEP: Scale invariant (Lévy) sources**

Fluctuations appear on many scales,

final position is a sum of many random shifts:

$$
x = \sum_{i=1}^{n} x_i, \qquad f(x) = \int \Pi_{i=1}^{n} dx_i \Pi_{j=1}^{n} f_j(x_j) \delta(x - \sum_{k=1}^{n} x_k).
$$

correlation function measures a Fourier-transform,

that of an n-fold convolution:

$$
\tilde{f}(q) = \int dx \exp(iqx) f(x), \quad \tilde{f}(q) = \prod_{i=1}^{n} \tilde{f}_i(q),
$$

Lévy: generalized central limit theorems adding one more step in the convolution does not change the shape

$$
\tilde{f}_i(q) = \exp(iq\delta_i - |\gamma_i q|^\alpha), \qquad \prod_{i=1}^n \tilde{f}_i(q) = \exp(iq\delta - |\gamma q|^\alpha)
$$

$$
\gamma^\alpha = \sum_{i=1}^n \gamma_i^\alpha, \qquad \delta = \sum_{i=1}^n \delta_i.
$$

# **Correlation functions for Lévy sources**

Correlation funct of stable sources:

$$
C(q;\alpha)=1+\lambda\exp{(-|qR|^\alpha)}
$$

R: scale parameter

- $\alpha$ : shape parameter or Lévy index of stability
- $\alpha = 2$  Gaussian,  $\alpha = 1$  Lorentzian sources

Further details: T. Cs, S. Hegyi and W. A. Zajc, EPJ C36 (2004) 67



# **Correlation signal of the CEP**

If the source distribution at CEP is a Lévy, it decays as:

$$
\rho(R) \propto R^{-(1+\alpha)}\Big|
$$

at CEP, the tail decreases as:

$$
\rho(R) \propto R^{-(d-2+\eta)}
$$

Levy index of stability 2.5 2 1.5 ರ  $0.5$  $\theta$  $-0.5$  $0.5$  $\Omega$  $-1$ T

#### hence:  $\alpha$  as a function of  $\tau = \mid T - T_c \mid / T_c$

$$
\alpha(L\acute{e}vy)=\eta(3d\;Ising)=0.50\pm0.05
$$

T. Cs, S. Hegyi, T. Novák, W.A.Zajc,

Acta Phys. Pol. B36 (2005) 329-337

#### **Critical exponents, universality class**



#### **Lévy fits to PHENIX prel. Au+Au @ 200 GeV**



### **Summary**

**4 steps for a definitive result on CEP:** 

- **identify type of phase transition**
- **locate**
- **characterize**
- **cross-check**

**Concept of optical opacity:** 

- **both attenuation measure, R**<sub>AA</sub>
- <u>a sa sababaran sa sa</u> <u>- and lengthscale measure R<sub>HBT</sub> needed</u>



**perfect fluid at RHIC: 12 orders of magn more opaque than Pb for** γ

**Critical Opalescence: Smoking gun signature to locate CEP**

**Levy stable Bose-Einstein/HBT correlations**

**critical exponent** η

# **Backup slides**

**Signal of first order phase transitions from HBT Rout/Rside(m<sup>t</sup> , √sNN)** 



Idea: look first order phase transition where  $R_{out}/R_{side}$  >>1 Measure: Gaussian HBT radii for pions (if possible kaons too)

# **Lattice QCD: EoS of QCD Matter**



At the Critical End Point, the chiral phase transition is of 2nd order. Stepanov, Rajagopal, Shuryak: (Static) universality class of QCD = 3d Ising model PRL 81 (1998) 4816

## **Measure by two-particle correlations**

Single particle spectrum: averages over space-time information

$$
E\frac{dN}{d\mathbf{p}} = \int dx^4 S(x, \mathbf{p})
$$

 sensitivity to space-time information Correlations:

$$
C_2(\mathbf{q}) = \frac{dN_2/d\mathbf{p}_1 d\mathbf{p}_2}{(dN_1/d\mathbf{p}_1)(dN_1/d\mathbf{p}_2)} \approx \int d\mathbf{r} \left[ \Phi(\mathbf{r}, \mathbf{q}) \right]^2 S(\mathbf{r}, \mathbf{q})
$$
  
FSI Solve function

Intensity interferometry, HBT technique, femtoscopy ….

### **Search for a 1st order QCD PT**

**QGP has more degrees of freedom than pion gas**

**Entropy should be conserved during fireball evolution**

 **Look in** *hadronic* **phase Hence: for signs of: Large size, Large lifetime, Softest point of EOS**

**possible signal of a 1st order phase trans.**



#### **But are the correlation data Gaussian?**



#### **Non-Gaussians, 2d E802 Si+Au data**



E802 Si+Au data,  $sqrt(s_{NN}) = 5.4$  GeV Z. Fodor, S.D. Katz:  $T_F = 162 \pm 2$  MeV,  $\mu_{\text{\tiny E}} = 360$ ±40 MeV l (hep-lat/0402006 systematics: this meeting

Best Gaussian: bad shape

#### **Non-Gaussian structures, 2d, UA1 data**



### **Femptoscopy signal of sudden hadronization**

Buda-Lund hydro: RHIC data follow the predicted (1994-96) scaling of HBT radii

 $\textsf{hep-ph}/9406365$   $\textcolor{red}{\blacksquare}$   $\textcolor{red}{\blacksquare}$ T. Cs, L.P. Csernai T. Cs, B. Lörstad hep-ph/9509213

Hadrons with T>T<sub>c</sub>: 1st order PT excluded hint of a cross-over M. Csanád, T. Cs, B. Lörstad and A. Ster, nucl-th/0403074



# **Backup slides (2)**

# **1 st milestone: new phenomena**



**Suppression of high p<sup>t</sup> particle production in Au+Au collisions at RHIC**

### **2 nd milestone: new form of matter**



# **3 rd milestone: Top Physics Story 2005**



**http://arxiv.org/abs/nucl-ex/0410003**

**PHENIX White Paper: second most cited in nucl-ex during 2006** 

#### **4 th Milestone: A fluid of quarks**



#### **Strange and even charm quarks participate in the flow**

**Milestone # 5: Perfection at limit! All "realistic" hydrodynamic calculations for RHIC fluids to date have assumed zero viscosity**

- η **= 0** →**"perfect fluid"**
- **But there is a (conjectured) quantum limit:**

$$
\eta \ge \frac{\hbar}{4\pi} (Entropy Density) = \frac{\hbar}{4\pi} s
$$

**"A Viscosity Bound Conjecture", P. Kovtun, D.T. Son, A.O. Starinets, hep-th/0405231 Where do "ordinary" fluids sit wrt this limit?** 200 **(4** π**)** η**/s > 10 ! Example 19** Helium 0.1 MPa Nitrogen 10 MPa

**RHIC's perfect fluid (4** π**)** η**/s ~1 on this scale: The hottest (T > 2 Terakelvin) and the most perfect filuid ever made…**<br>**fluid ever made…** 



# **World Context: Search for the CEP**

