

# Correlations in the High Temperature QCD Dirac Spectrum

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spectrum of  $D(A)$   $\iff$   $\langle \bar{\psi}\psi \rangle$  (chiral order parameter)

$$\begin{aligned}\langle \bar{\psi}\psi \rangle &= \lim_{m \rightarrow 0} \lim_{V \rightarrow \infty} \frac{1}{V} \langle \sum_x \text{Tr}(D+m)^{-1} \rangle \\ &= \lim_{m \rightarrow 0} \int_0^\infty d\lambda \rho(\lambda) \frac{1}{\lambda^2 + m^2} \\ &= \pi \rho(0),\end{aligned}$$

where  $\rho(\lambda)$  is the spectral density per unit volume of  $D$

# Chirally symmetric and broken phase

$$\rho(0) \neq 0$$

$\chi$ -symmetry broken

- $T < T_c$
- $T > T_c$  and  $\text{Re}(P) < 0$   
("unphysical" P-loop sector)

$$\rho(\lambda) \approx \Sigma$$

Random matrix theory (RMT)  
 $\Rightarrow$  statistical description of low eigenvalues.

$$\rho(0) = 0$$

$\chi$ -symmetry intact

- $T > T_c$  and  $\text{Re}P > 0$   
("physical" Polyakov sector)

$$\rho(\lambda) \approx C \cdot \lambda^\alpha$$
$$\rho(\lambda) \approx C \cdot (\lambda - G)^\alpha$$

No simple statistical description of low eigenvalues.

- $SU(2)_c$  quenched.
- Overlap Dirac operator.
- $T = 2.6 \cdot T_c$
- $N_t = 4$ ;  $L = 12 - 32$ , 8 different spatial volumes.
- $N_t = 6$ ;  $L = 18 - 32$ , 8 different spatial volumes.
- Only  $Q = 0$  topological sector used.  
(at high  $T$  few  $Q \neq 0$  confs.)
- Antiperiodic b.c. for the fermions in the time direction,  
 $P$ -loop  $> 0$  sector  $\Rightarrow \rho(0) = 0$   
 $P$ -loop  $< 0$  sector  $\Rightarrow \rho(0) \neq 0$

# Is there any difference between the eigenmodes?

## Localization

- How to measure spatial “extension” of (normalized) eigenmodes?



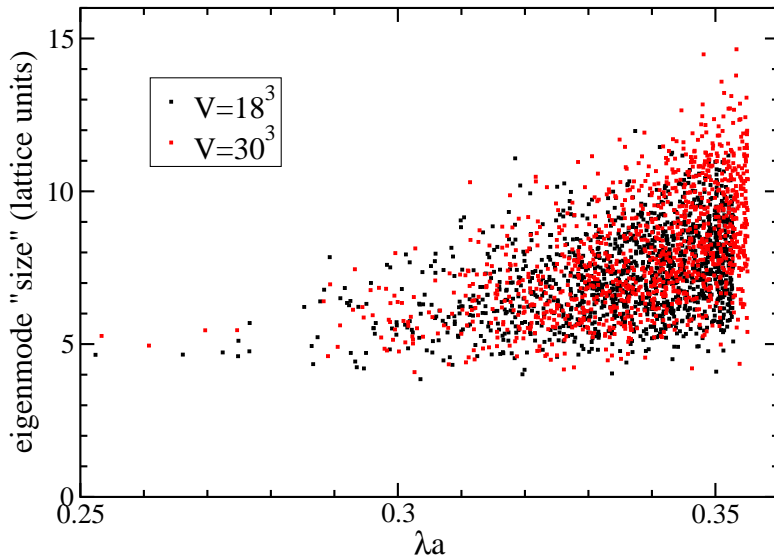
$$\mathcal{V} = \left[ \sum_{\mathbf{x}} (\psi^\dagger \psi(\mathbf{x}))^2 \right]^{-1}.$$

- If  $\psi$  is constant in volume  $V$  and zero elsewhere  $\Rightarrow \mathcal{V} = V$ .
- $\mathcal{V}$  measures “volume” occupied by eigenmode.
- $\Rightarrow$  Define 3-dimensional “size”:

$$d = \left[ \frac{\mathcal{V}}{N_t} \right]^{1/3}.$$

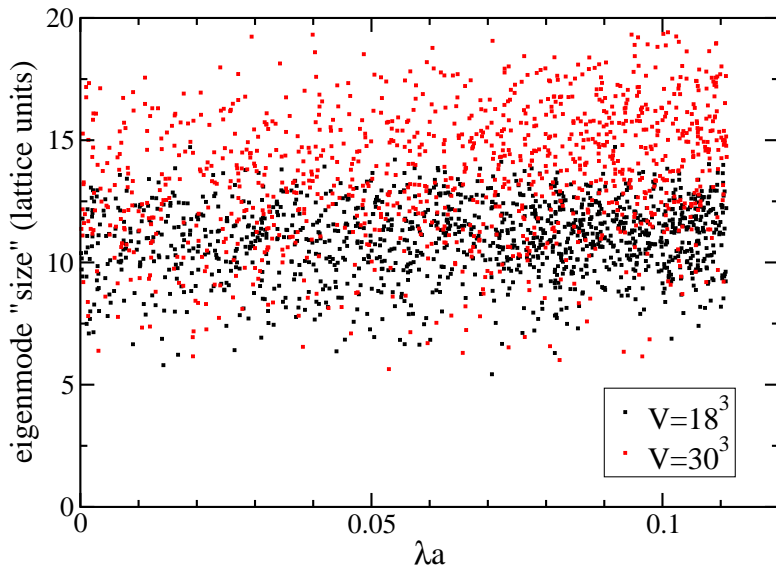
# Localization of low eigenmodes of $D$ :

$\rho(0) = 0$  sector (Polyakov loop in physical sector)



# Localization of low eigenmodes of $D$ :

$\rho(0) \neq 0$  sector (Polyakov loop in un-physical sector)



# Differences in eigenmodes

$$\rho(0) \neq 0$$

Eigenmodes extended

Extension grows with box size

$$\rho(0) = 0$$

Eigenmodes localized

Extension independent of box size

$\rho(0) = 0$  simple picture

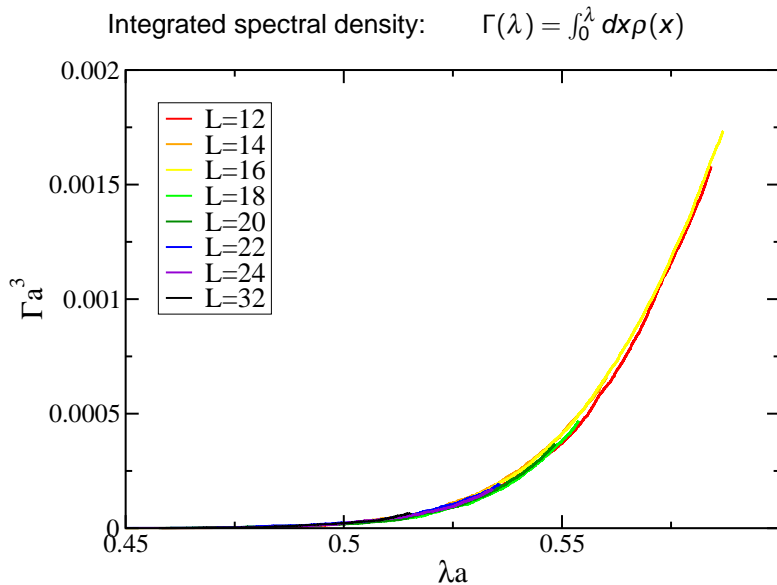
- Low eigenvalues produced independently in many small subvolumes of  $\approx d^3$ .
- Avg. number of “small” eigenvalues in volume  $d^3$ :  $\ll 1$ .
- $\Rightarrow$  Small eigenvalues occur independently.



# Assumptions for low eigenvalues in the $\rho(0) = 0$ phase

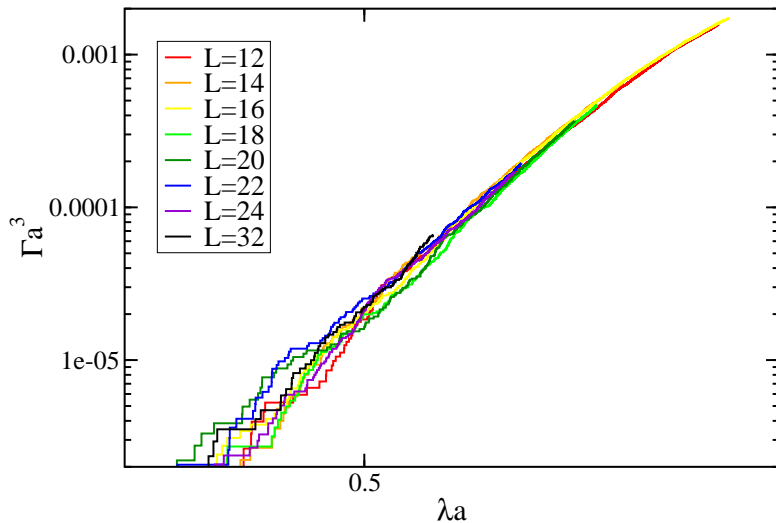
- For any given temperature  $\rho(\lambda)$  is well-defined, i.e. the spectral density scales with the 3-volume.
- $\rho(\lambda) \approx C \cdot \lambda^\alpha$ .
- The number of eigenvalues in any two disjoint intervals are independent random variables.  
 $\Rightarrow$  The spectral density encodes all statistical properties of the spectrum.

# Scaling of the spectral density with the 3-volume



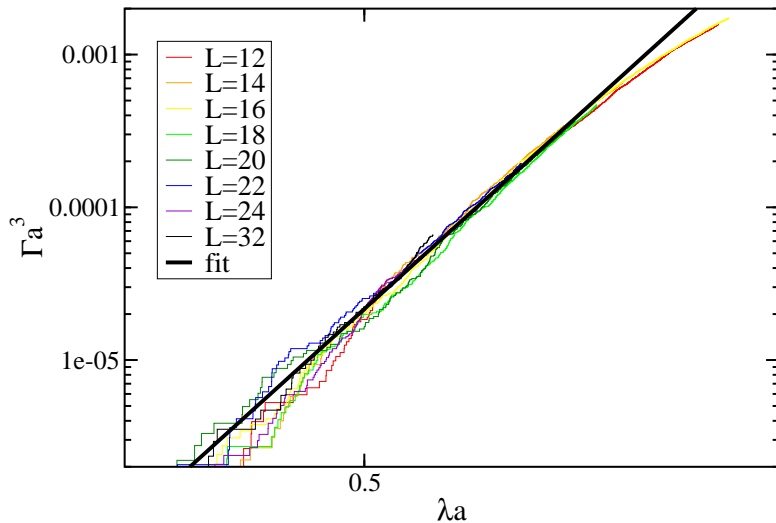
# Scaling of the spectral density with the 3-volume

$$\Gamma(\lambda) = \int_0^\lambda dx \rho(x) \quad \text{log-log scale}$$



# Spectral density around zero

$$\Gamma(\lambda) = \frac{C}{\alpha+1} \cdot \lambda^{\alpha+1} \quad \text{fit}$$



# Correlations in the spectrum

Assume that different regions in the spectrum are un-correlated.

- Number of eigenvalues in any two disjoint intervals are independent random variables.
- It is not the correlations among different eigenvalues! (Usually computed in RMT.)
- Hard to test directly.
- Assumption  $\Rightarrow$  Fully determines:
  - Statistical properties of low eigenvalues.
  - Volume dependence of statistics.
- All statistical information encoded in spectral density.
- $\Rightarrow$  Easily verifiable analytic predictions.

# Distribution of smallest eigenvalue

Let's first compute  $\mathcal{P}_0(\lambda_1, \lambda_2)$ , the probability that there is no eigenvalue in the  $[\lambda_1, \lambda_2]$  interval:

Average number of eigenvalues in a small  $\Delta x (\rightarrow 0)$  interval:

$$V\rho(x)\Delta x$$

- Probability of  $\geq 2$  eigenvalue:  $\rightarrow 0$
- Probability of 0 eigenvalue:  $1 - V\rho(x)\Delta x$
- Probability of 1 eigenvalue:  $V\rho(x)\Delta x$

No correlation  $\Rightarrow$  Probabilities factorize:

$$\mathcal{P}_0(\lambda, x + \Delta x) = \mathcal{P}_0(\lambda, x) \cdot [1 - V\rho(x)\Delta x]$$

$$\Rightarrow \frac{d\mathcal{P}_0(\lambda, x)}{dx} = -V\rho(x) \cdot \mathcal{P}_0(\lambda, x) = -VCx^\alpha \cdot \mathcal{P}_0(\lambda, x)$$

# Distribution of the smallest eigenvalue contd.

$\mathcal{P}_0(\lambda_1, \lambda_2)$  contd.:

Solution with initial value

$$\mathcal{P}_0(\lambda, \lambda) = 1:$$

$$\mathcal{P}_0(\lambda, x) = \exp \left[ -\frac{VC}{\alpha+1} \left( x^{\alpha+1} - \lambda^{\alpha+1} \right) \right]$$

Probability density of smallest eigenvalue,  $\rho_1(\lambda)$ :

Probability of  $\begin{cases} \text{no eigenvalue in } [0, \lambda] : & \mathcal{P}_0(0, \lambda) \\ \text{one eigenvalue in } [\lambda, \lambda + d\lambda] : & V\rho(\lambda) d\lambda \end{cases}$

$$\mathcal{P}_0(0, \lambda) \cdot V\rho(\lambda) d\lambda = \rho_1(\lambda) d\lambda$$

$$\Rightarrow \rho_1(\lambda) = \exp \left( -\frac{CV}{\alpha+1} \lambda^{\alpha+1} \right) CV\lambda^\alpha$$

# Further analytic predictions

Distribution of the second smallest eigenvalue,  $p_2(\lambda)$ :

$$\text{Probability of } \begin{cases} \text{Smallest eigenvalue in } [x, x + \Delta x]: & p_1(x)\Delta x \\ \text{No eigenvalue in } [x + \Delta x, \lambda]: & \mathcal{P}_0(x, \lambda) \\ \text{One eigenvalue in } [\lambda, \lambda + d\lambda]: & V\rho(\lambda)d\lambda \end{cases}$$

$$\begin{aligned} \Rightarrow p_2(\lambda) &= \int_0^\lambda dx p_1(x) \cdot \mathcal{P}_0(x, \lambda) \cdot V\rho(\lambda) \\ &= \frac{C^2 V^2}{\alpha + 1} \exp\left(-\frac{CV}{\alpha + 1} \lambda^{\alpha+1}\right) \lambda^{2\alpha+1}. \end{aligned}$$

Average smallest eigenvalues

$$\langle \lambda_1 \rangle = (CV\mu)^{-\mu} \cdot \Gamma(1 + \mu)$$

$$\langle \lambda_2 \rangle = (CV\mu)^{-\mu} \cdot \Gamma(2 + \mu)$$

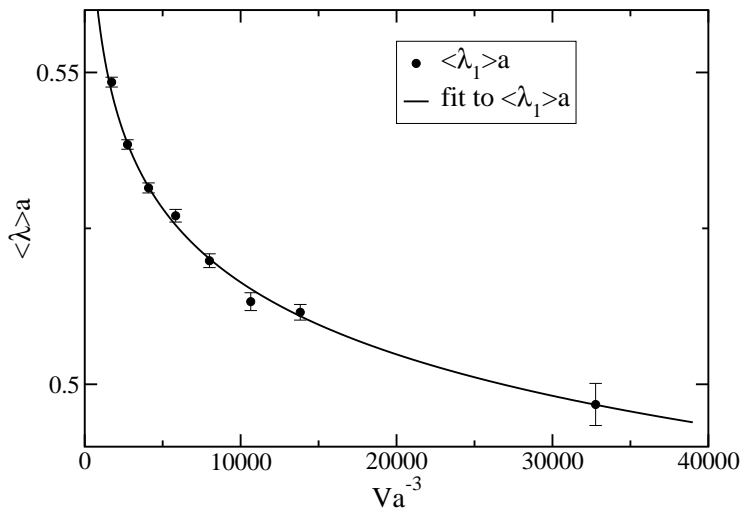
$$\mu = \frac{1}{\alpha+1}$$



# Numerical tests: volume dependence of $\langle \lambda_1 \rangle$ ( $N_T = 4$ )

Fit:  $\langle \lambda_1 \rangle = (CV\mu)^{-\mu} \cdot \Gamma(1 + \mu)$

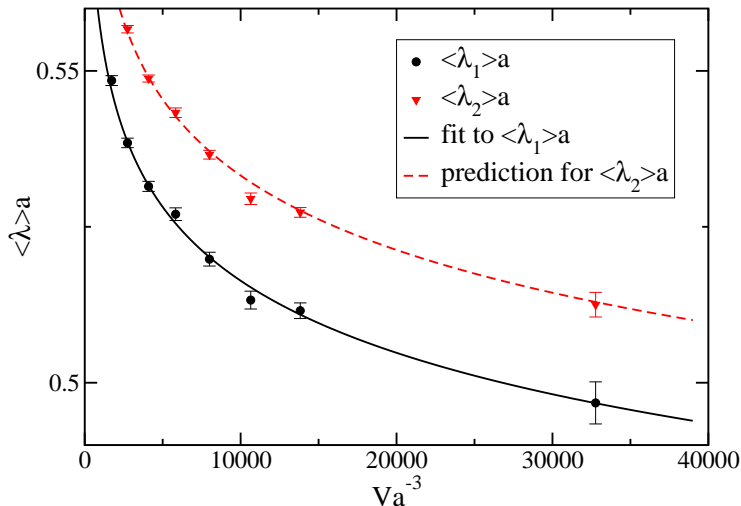
Fit parameters:  $\mu, C$



# Volume dependence of $\langle \lambda_2 \rangle$ ( $N_T = 4$ )

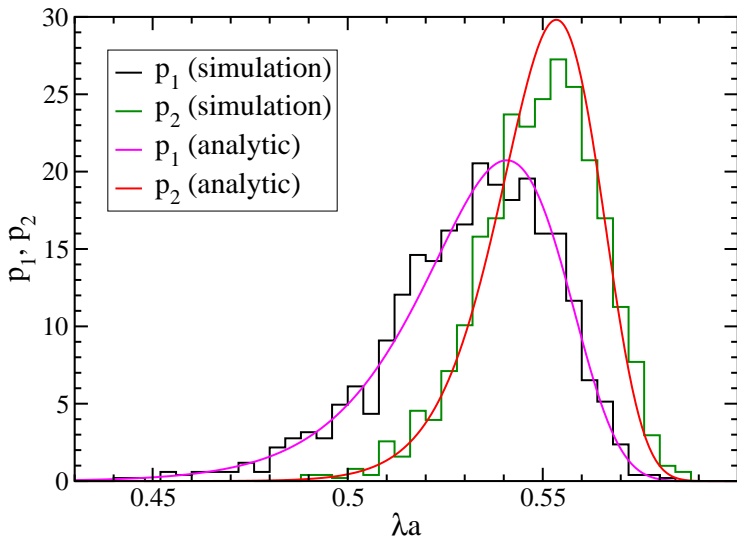
$$\langle \lambda_2 \rangle = (CV\mu)^{-\mu} \cdot \Gamma(2 + \mu)$$

No free parameter!



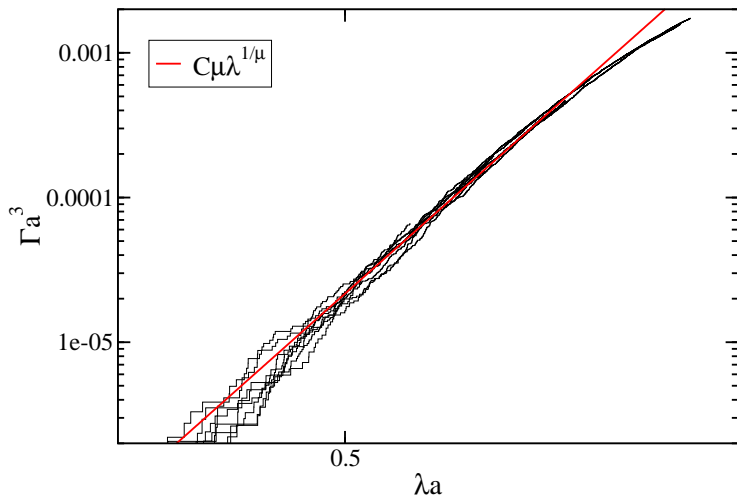
# Distribution of the smallest two eigenvalues ( $V = 16^3$ )

Not a fit,  $p_1(\lambda), p_2(\lambda)$  with previously fitted parameters  $\mu, C$ .



# Integrated spectral density, $\Gamma(\lambda)$

Not a fit,  $\Gamma(\lambda) = \mu C \lambda^{\frac{1}{\mu}}$  with previously fitted parameters  $\mu, C$ .



# Is it really incompatible with RMT?

- Isn't it possible that this works also for  $\rho(0) \neq 0$ ?
- RMT analytic predictions are different.  
E.g. distribution of smallest eigenvalue:

After a

$$z = \frac{CV}{\alpha + 1} \cdot \lambda^{\alpha+1}$$

transformation (in RMT:  $\alpha = 0, C = \Sigma$ ):

$$\text{RMT: } \rho_1(z) = \frac{2+z}{4} \exp\left(-\frac{z}{2} - \frac{z^2}{8}\right)$$

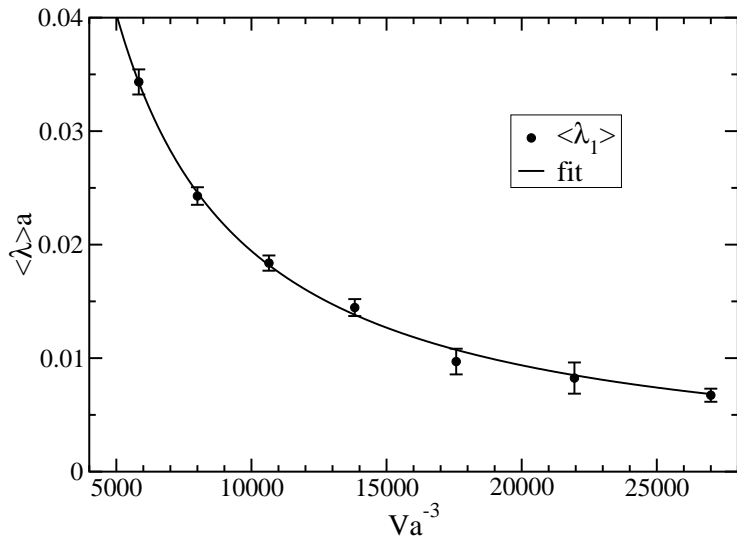
$$\text{Here: } \rho_1(z) = \exp(-z)$$

- Can we see the difference numerically?
- Same analysis for the  $\rho(0) \neq 0$  case:  
in un-physical P-loop sector.

# Volume dependence of $\langle \lambda_1 \rangle$ ( $N_T = 6, P < 0$ )

Fit:  $\langle \lambda_1 \rangle = (CV\mu)^{-\mu} \cdot \Gamma(1 + \mu)$

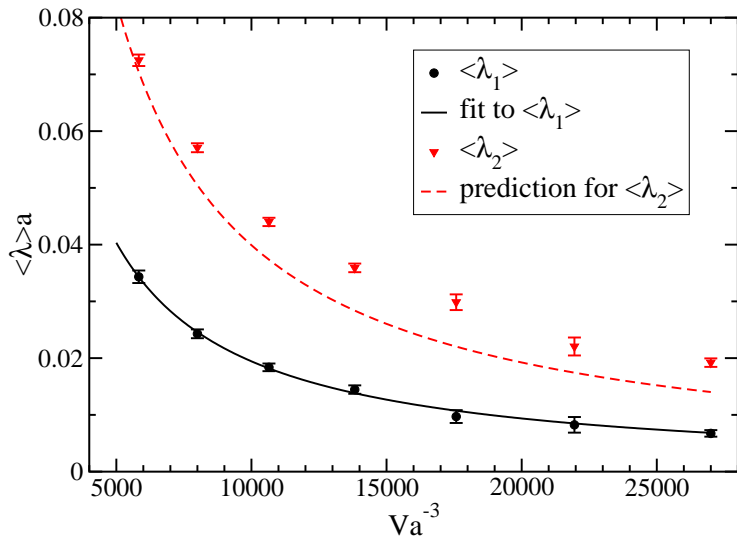
Best fit parameter:  $\mu = 1.05(3)$



# Volume dependence of $\langle \lambda_2 \rangle$

$$\langle \lambda_2 \rangle = (CV\mu)^{-\mu} \cdot \Gamma(2 + \mu)$$

No free parameter!



# Summary

- Main assumption:  
number of eigenvalues in disjoint intervals are uncorrelated
- Chirally symmetric phase ( $\rho(0) = 0$ ):  
good description of the low end of the Dirac spectrum
- Chirally broken phase ( $\rho(0) \neq 0$ ):  
random matrix theory is the proper description.

## Further questions:

- Are these results generally true if  $\rho(0) = 0$ ?  
( $SU(3)_c$ , dynamical fermions, other fermion representations...?)
- Possible connection to Anderson localization?  
[A.M. Garcia-Garcia and J.C. Osborn, Phys. Rev. D75 \(2007\) 034503.](#)
- Possible connection to “generalized” random matrix models?  
[Jackson and Verbaarschot, Phys. Rev. D53 \(1996\) 7223.](#)