

# Effect of Freeze-out and Hadronization on Flow in ultrarelativistic HIC

Miklós Zétényi

with Yun Cheng, László P. Csernai, Volodymyr K. Magas, Daniel D. Strottman

Zimányi School '09

02 Dec 2009

# Introduction – the hydro description of ultrarelativistic HIC

3 stages of the evolution in Ultrarelativistic Heavy Ion Collisions:

initial nonequilibrium evolution  
& QGP formation

e.g. fire streak model

equilibrated QGP

numerical hydro

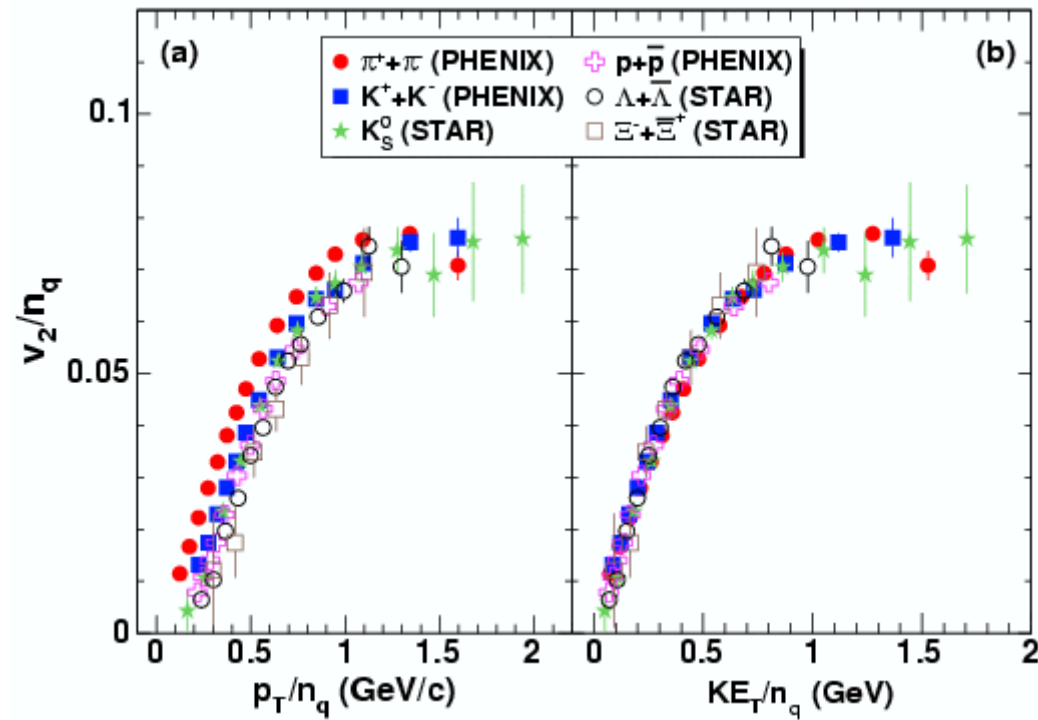
hadronic phase & Freeze-out

(hadronic transport, “afterburner”)

Hadronization & Freeze-out:

- a very fast, nonequilibrium process
- first approximation: **sudden** and **simultaneous** FO & Hadronization

# Constituent quark number scaling of $v_2(K E_T)$



Collective flow of hadrons can be described in terms of constituent quarks.

# Does hadronization accelerate the fluid?

Boundary conditions on the hadronization surface:  $[a] = a_{hadr} - a_{QGP}$

Baryon number:  $[N^\mu d\hat{\sigma}_\mu] = 0$        $N^\mu = n u^\mu$        $u^\mu = \gamma(1, \mathbf{v})$

Energy & momentum:  $[T^{\mu\nu} d\hat{\sigma}_\nu] = 0$        $T^{\mu\nu} = (e + P) u^\mu u^\nu - P g^{\mu\nu}$

(Local) isochronous Freeze-out frame:  $d\hat{\sigma}^\mu = (1, \mathbf{0})$

$$[n\gamma] = 0$$

$$[(e + P)\gamma^2 - P] = 0$$

Can be solved numerically

$$[(e + P)\gamma^2 \mathbf{v}] = 0$$

## Acceleration, nonrelativistic limit

$v \ll 1$   $\gamma \approx 1$  in the local isochronous frame of the fluid

(good approximation if close to Bjorken flow)

$$[n]=0, \quad [e]=0, \quad [(e+P)\mathbf{v}]=0$$

$$\frac{v_{hadr}}{v_{QGP}} = \frac{e+P_{QGP}}{e+P_{hadr}}$$

acceleration if  $P_{hadr} < P_{QGP}$

## Role of temperature

thermal smearing  $\rightarrow$  anisotropy is reduced

# Equation of State

QGP: Bag EOS  $P = \frac{e}{3} - \frac{4}{3}B$        $e = e_0 + B$        $P = P_0 - B$

## Post Freeze-out/Hadronization:

NCQ scaling → description in terms of constituent quarks & antiquarks

(anti)quarks are in hadrons → No. of d.o.f. is different from a gas of free quarks and antiquarks → average degeneracy factor  $g_{av}$

relativistic Boltzmann gas of massive quarks + antiquarks

the main effects:

appearance of quark mass

disappearance of Bag constant

disappearance of gluons

extra degeneracy factor

## Equation of State II.

$$n(\mu, T) = \frac{g_{av}}{\pi^2} m^2 T \sinh\left(\frac{m\mu}{T}\right) K_2\left(\frac{m}{T}\right) \quad e(\mu, T) = 3P \left(1 + \frac{m}{3T} \frac{K_1}{K_2}\right)$$

$$P(\mu, T) = \frac{g_{av}}{\pi^2} m^2 T^2 \cosh\left(\frac{m\mu}{T}\right) K_2\left(\frac{m}{T}\right) \quad K_i = K_i\left(\frac{m}{T}\right)$$

$e, n, \mathbf{v}$ : obtained numerically from the boundary conditions

EOS can be solved for  $\mu$  and  $T$

thermal distribution for (anti)quarks (Juttner):

$$f_{q/\bar{q}}(x, p) = \frac{g_{av}}{(2\pi)^3} \exp\left(\frac{-p^\rho u_\rho \pm \mu}{T}\right)$$

Flow observables can be obtained using the Cooper-Frye formula.

$$E \frac{d^3 N}{d\mathbf{p}^3} = \int f(x, p) p^\mu d\sigma_\mu$$

# Entropy

Ideal gas:  $s/N = \text{const.}$   $\rightarrow$  entropy production = particle (pion) production

Coalescence:  $N$  is reduced by a factor of 2-3  
 $\rightarrow$  entropy is decreasing

Phase transition: EOS of state is different from ideal gas  
relevant degrees of freedom change

Bíró, Zimányi '07:

isentropically expanding fireball

with lattice QCD EOS

effective particle number:  $N_{\text{eff}} = V P(\mu, T)/T$

$N_{\text{eff}}$  is reduced by a factor of 1/3 as we go below  $T_c$



## Entropy II

Increasing entropy condition:  $[S^\mu d\hat{\sigma}_\mu] \geq 0$   $s^\mu = s u^\mu$

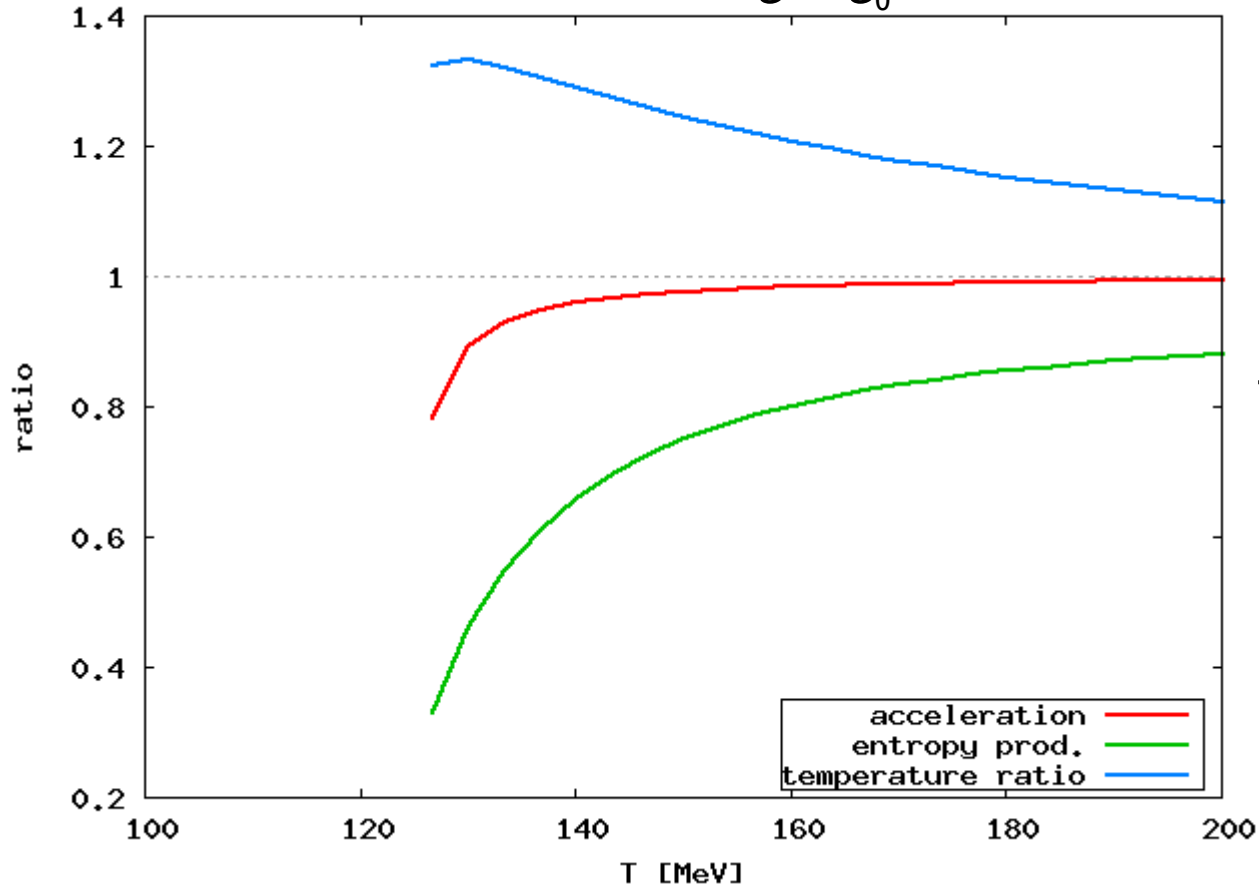
Isochronous Freeze-out frame:  $[s\gamma] = 0$

In the constituent quark + antiquark model:

$$S = \left( \frac{\partial P}{\partial T} \right)_\mu = \frac{g_{av}}{\pi^2} m^2 \left[ \cosh\left(\frac{\mu}{T}\right) (4T K_2 + m K_1) - \mu \sinh\left(\frac{\mu}{T}\right) K_2 \right]$$

# Numerical results for acceleration, entropy prod.

$v = 0.99 c, \quad g = g_0$



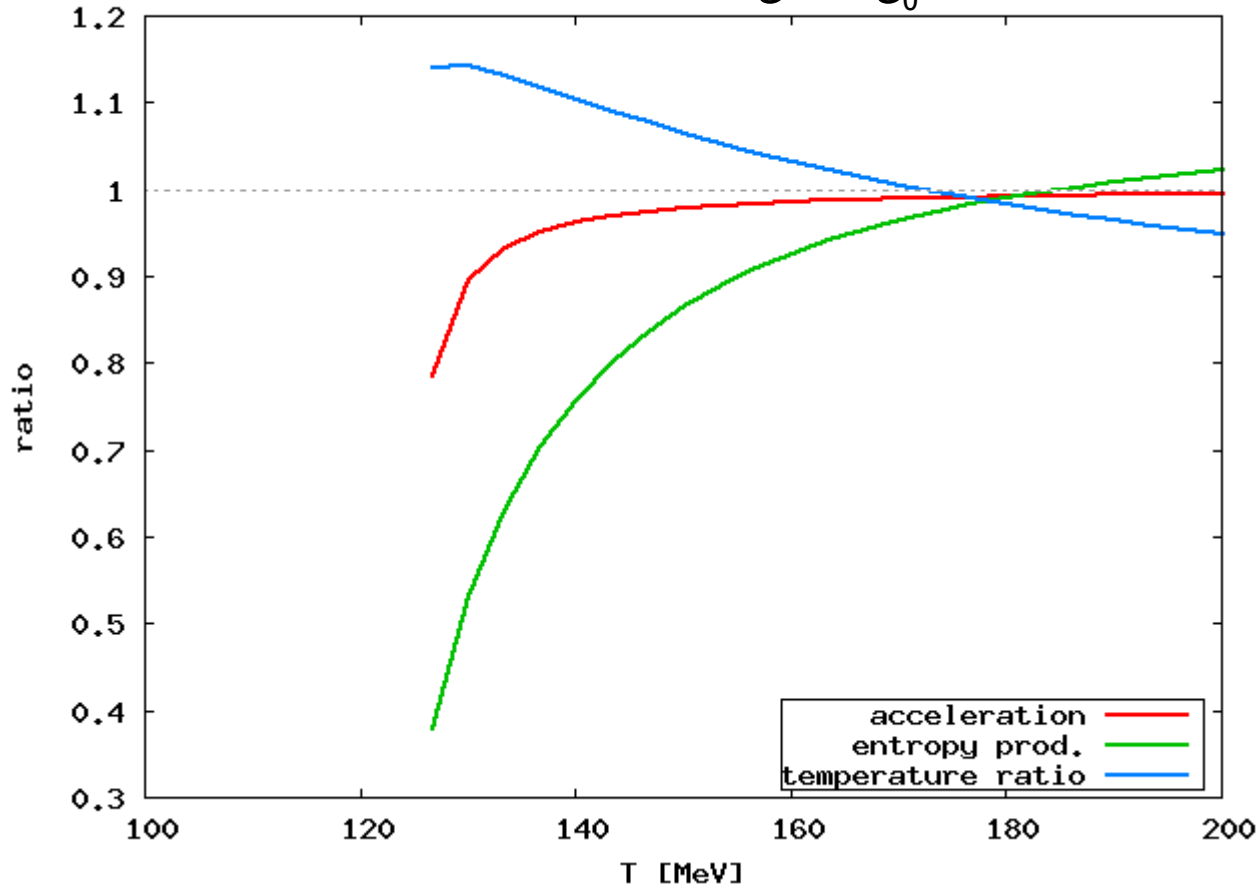
$$\text{acceleration} = v_{\text{hadr}}/v_{\text{QGP}}$$

$$\text{entropy prod.} = (\gamma s)_{\text{hadr}}/(\gamma s)_{\text{QGP}}$$

$$\text{temperature ratio} = T_{\text{hadr}}/T_{\text{QGP}}$$

# Numerical results for acceleration, entropy prod.

$v = 0.99 c, \quad g = 2g_0$



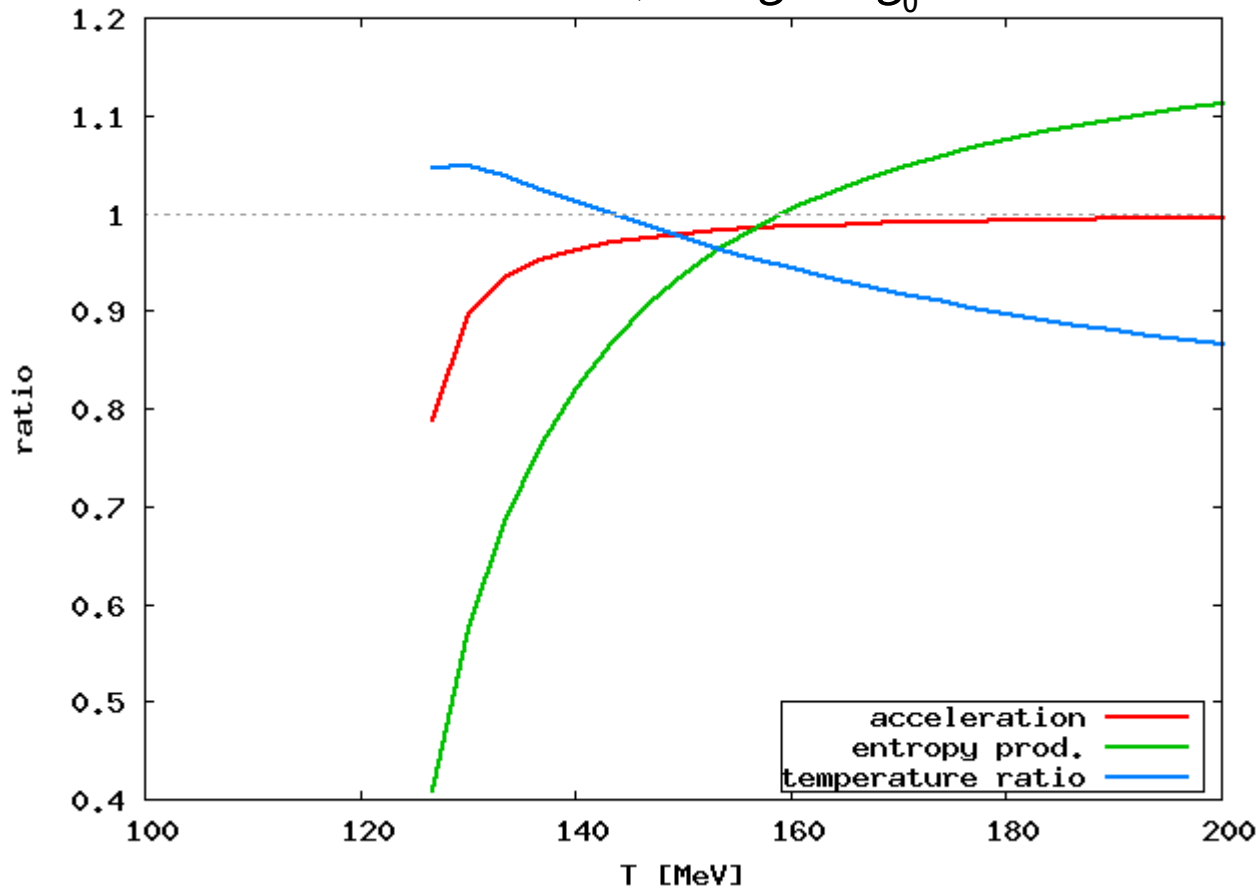
$$\text{acceleration} = v_{\text{hadr}}/v_{\text{QGP}}$$

$$\text{entropy prod} = (\gamma s)_{\text{hadr}}/(\gamma s)_{\text{QGP}}$$

$$\text{temperature ratio} = T_{\text{hadr}}/T_{\text{QGP}}$$

# Numerical results for acceleration, entropy prod.

$v = 0.99 c,$       $g = 3g_0$



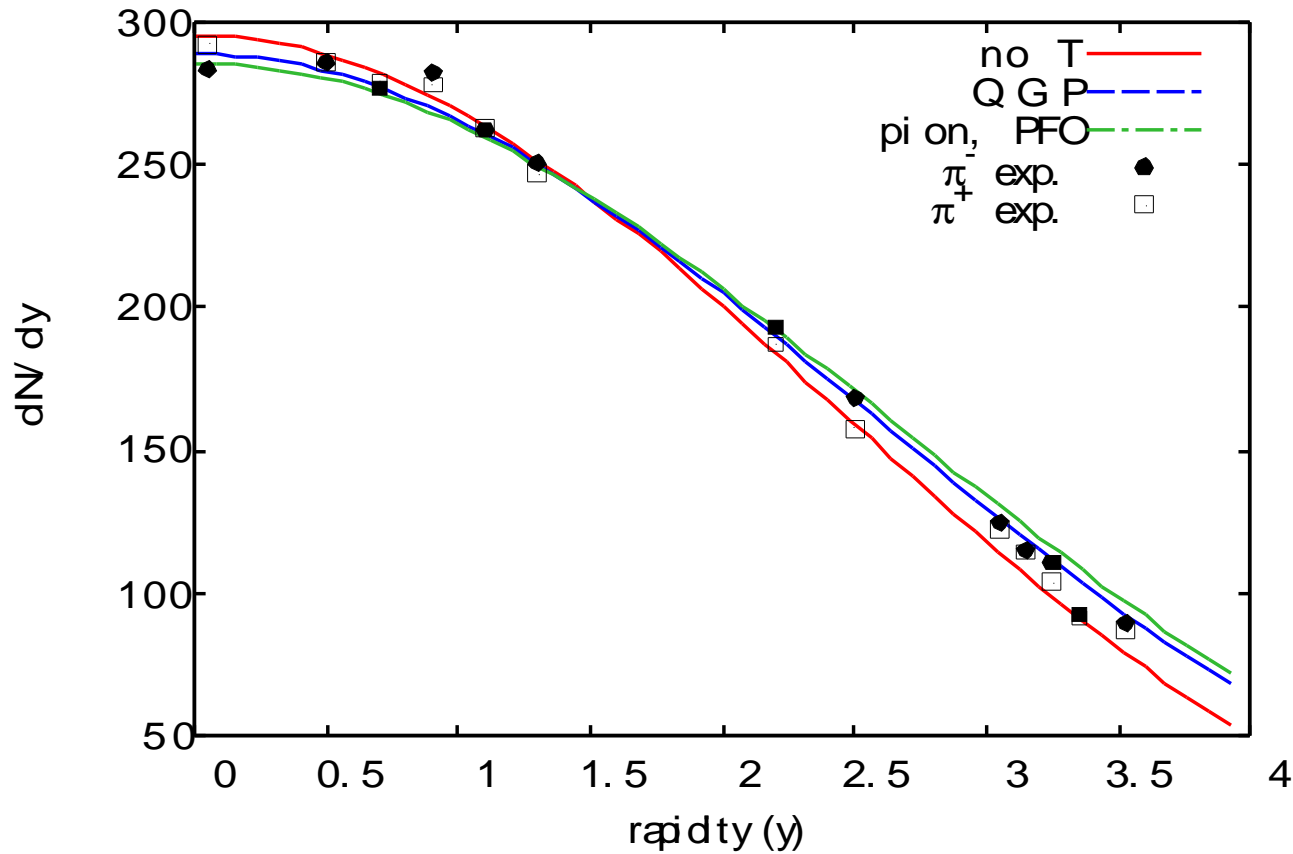
$$\text{acceleration} = v_{\text{hadr}}/v_{\text{QGP}}$$

$$\text{entropy prod} = (\gamma s)_{\text{hadr}}/(\gamma s)_{\text{QGP}}$$

$$\text{temperature ratio} = T_{\text{hadr}}/T_{\text{QGP}}$$

# Effect of thermal smearing – Landau model

$\pi^\pm$  rapidity spectra, Au+Au, 200 GeV, BRAHMS



Inclusion of thermal effects slightly improves the agreement with data.

# Problems to be faced in a more realistic treatment

The distribution on the PFO side is not necessarily thermal.

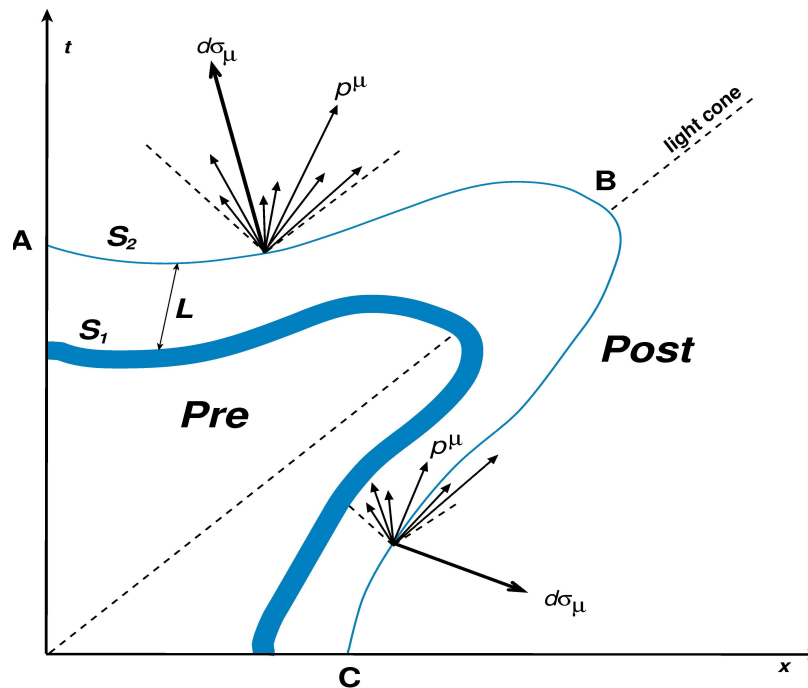
e.g.: hypersurface with spacelike normal

→ cut-Juttner, cancelling Juttner distribution

Freeze-out in finite time: hypersurface → finite layer

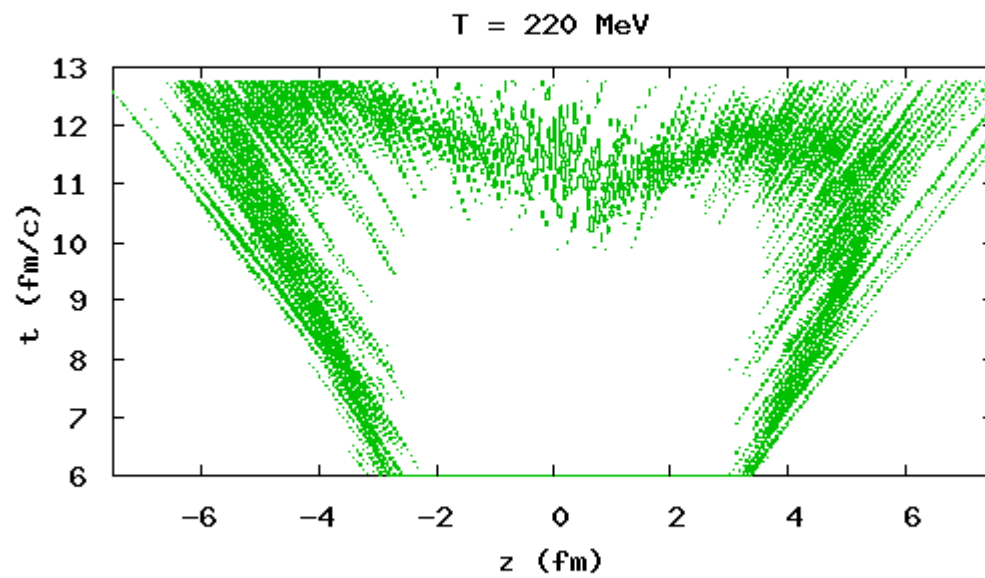
(approximate solution of Boltzmann eq.)

Molnár, Csernai, Magas,... '03



# Freeze-out hypersurface in numerical hydro

e.g. isotherm ( $T = 220$  MeV):



# Conclusions

NCQ scaling → flow can be studied in terms of constituent quarks

FO & Hadronization can affect flow observables in two ways:  
increase/decrease the collective flow velocity  
change in temperature – thermal smearing

The effect is not very large

It depends very much on the specific hadronization model



Backup slides

# Hadronization via recombination

Momentum distribution of mesons in the recombination model:

$$\frac{d^3 N}{dp^3} \propto \int \prod_{i=1}^2 d^3 x_i d^3 p_i f_q(x_1, p_1) f_q(x_2, p_2) W_M(p, p_1, p_2, x_1, x_2)$$

meson Wigner function:

momentum conservation

$$W_M(p, p_1, p_2, x_1, x_2) = \Phi_M(x_1 - x_2, p_1 - p_2) \delta(p_T - p_{T1} - p_{T2})$$

comoving quark and antiquark:  $\Phi_M \propto \delta^3(x_1 - x_2) \delta^3(p_1 - p_2)$

For the momentum distribution of mesons we get:

$$\frac{d^3 N_M}{p_T dp_T dy d\phi} \propto \int d^3 x f_q(x, p_t/2)^2$$

similar expression  
for baryons,  $2 \rightarrow 3$

# Hadronization via recombination II.

calculation of flow moments:

$$v_n(p_T) = \frac{\int dy d\phi \cos n\phi \frac{d^3 N}{p_T d p_T dy d\phi}}{\int dy d\phi \frac{d^3 N}{p_T d p_T dy d\phi}}$$

elliptic flow of mesons:

$$v_{2,M}(p_T) = \frac{2v_{2,q}(p_T/2)}{1+2v_{2,q}^2(p_T/2)} \quad \frac{v_{2,M}(p_T)}{2} = v_{2,q}(p_T/2)$$

baryons:

$$v_{2,B}(p_T) = \frac{3v_{2,q}(p_T/3)+3v_{2,q}^3(p_T/3)}{1+6v_{2,q}^2(p_T/3)} \quad \frac{v_{2,B}(p_T)}{3} = v_{2,q}(p_T/3)$$

in a medium  $p_T$  is not necessarily conserved,  $KE_T = m_T - m$  might be conserved

→ scaling in the variable  $KE_T$  J. Jia & C. Zhang '07