

Inelastic Boosted Dark Matter Searches at ICARUS – Gran Sasso

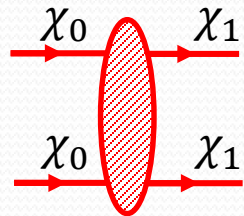


Doojin Kim

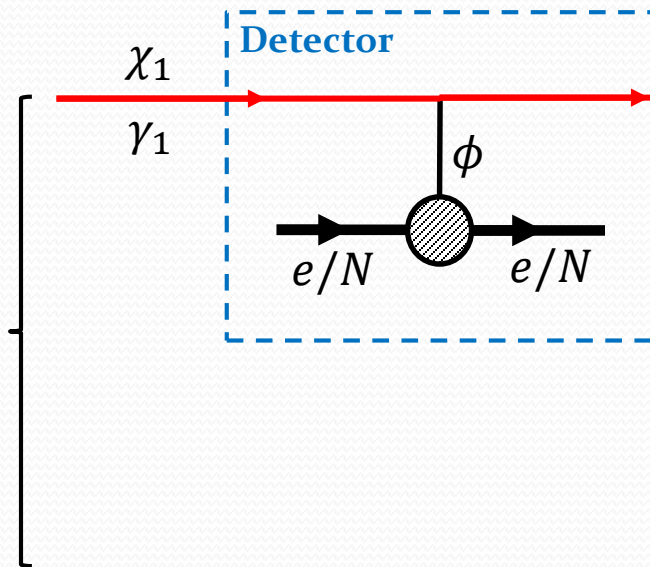
Generic BDM Signal Processes

(a) Elastic scattering (eBDM) (cf. eBDM at DUNE [Necib, Moon, Wongjirad, Conrad (2016); Alhazmi, Kong, Mohlabeng, Park (2016)])

Our focus here is
 $m_0 = E_1 = \sim 30 \text{ MeV} \sim 2 \text{ GeV}$



Galactic Center

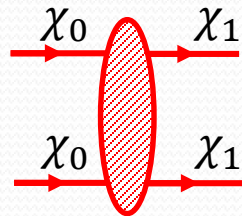


Signal **featureless**, hence
 challenging to separate from
 potential BGs

- χ_0 : heavier DM
- χ_1 : lighter DM
- γ_1 : boost factor of χ_1
- χ_2 : massive unstable dark-sector state
- ϕ : mediator/portal particle

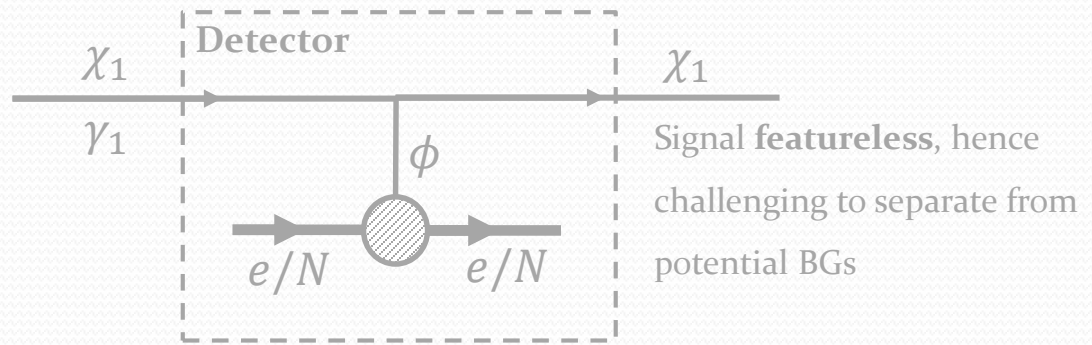
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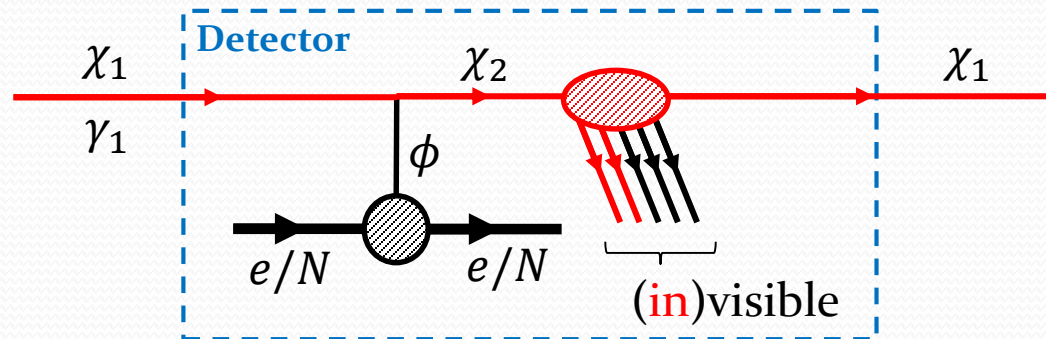


Galactic Center

(a) Elastic scattering (eBDM) (cf. eBDM at DUNE [Necib, Moon, Wongjirad, Conrad (2016); Alhazmi, Kong, Mohlabeng, Park (2016)])



(b) Inelastic scattering (iBDM) (cf. iBDM at DUNE [DK, Park, Shin (2016)])



Many signal **features**, helping veto BGs

- χ_0 : heavier DM
- χ_1 : lighter DM
- γ_1 : boost factor of χ_1
- χ_2 : massive unstable dark-sector state
- ϕ : mediator/portal particle

Benchmark Model: Building Blocks

$$\mathcal{L}_{\text{int}} \ni -\frac{\epsilon}{2} F_{\mu\nu} X^{\mu\nu} + g_{11} \bar{\chi}_1 \gamma^\mu \chi_1 X_\mu + g_{12} \bar{\chi}_2 \gamma^\mu \chi_1 X_\mu + \text{h. c.} + (\text{others})$$

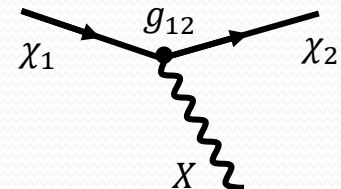
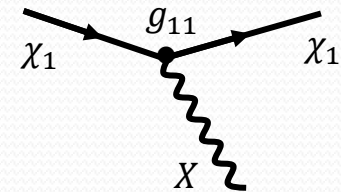
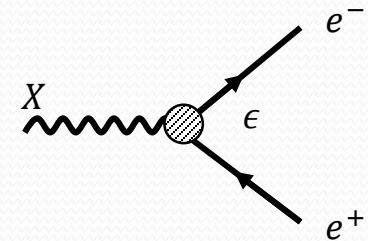
□ **Vector portal** (e.g., dark gauge boson scenario) [Holdom (1986)]

□ Fermionic DM

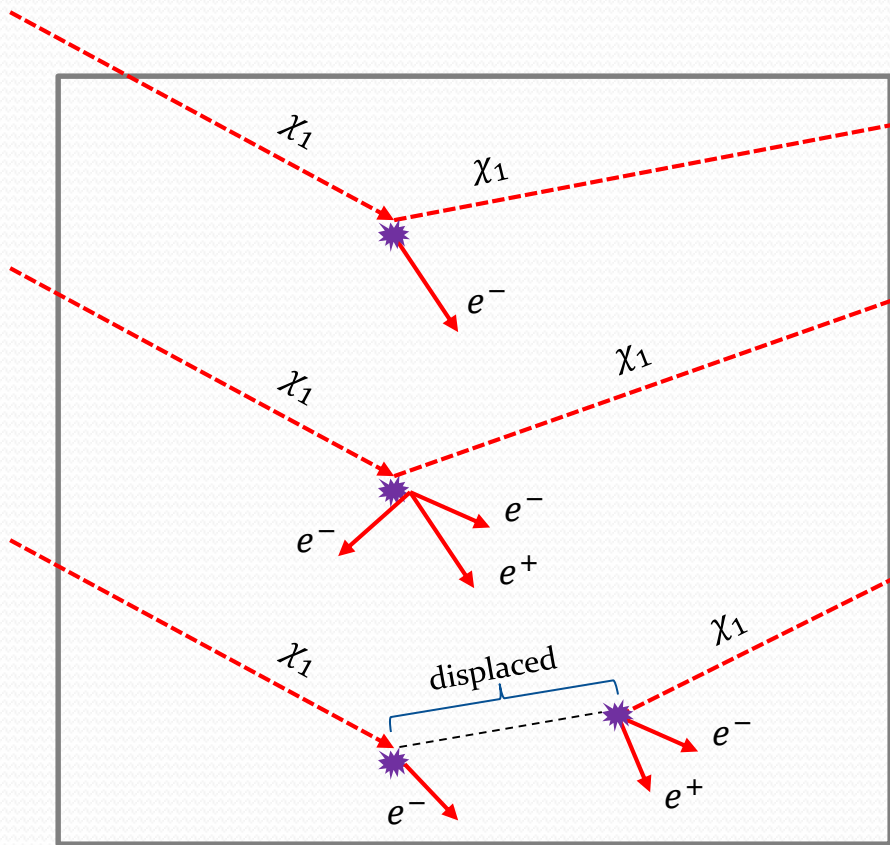
❖ χ_2 : a heavier (unstable) dark-sector state

❖ **Flavor-conserving neutral current** \Rightarrow elastic scattering

❖ **Flavor-changing neutral current** \Rightarrow inelastic scattering



Expected Signatures with Electron Recoil



- Ordinary elastic scattering: electron recoil (ER) only, i.e., single track
- “Prompt” inelastic scattering: ER + e^+e^- pair (from the decay of on-shell X), i.e., **three tracks**
- “Displaced” inelastic scattering: ER + e^+e^- pair (typically from a three-body decay of χ_2), i.e., again **three tracks**
- Note that **tracks will pop up inside the fiducial volume.**
- Straightforwardly applicable to proton recoil (up to form factor, DIS etc.)

Model-independent Reach

□ **Non-trivial** to find appropriate parameterizations for providing **model-independent reaches** due to many parameters involved in the model

□ Number of signal events N_{sig} is

$$N_{\text{sig}} = \sigma_{\epsilon} \mathcal{F} A t_{\text{exp}} N_e, \quad (3)$$

- σ_{ϵ} : scattering cross section between χ_1 and (target) electron
 - \mathcal{F} : flux of incoming (boosted) χ_1
 - A : acceptance
 - t_{exp} : exposure time
 - N_e : total # of target electrons
- } **Controllable!** (once a detector is determined)

Here determined by **distance between the primary (ER) and the secondary vertices**, other factors such as **cuts, energy threshold, etc are absorbed into σ_{ϵ}** . Depending on analyses, some factors can be reabsorbed into A .

Model-independent Reach: Comprehensive

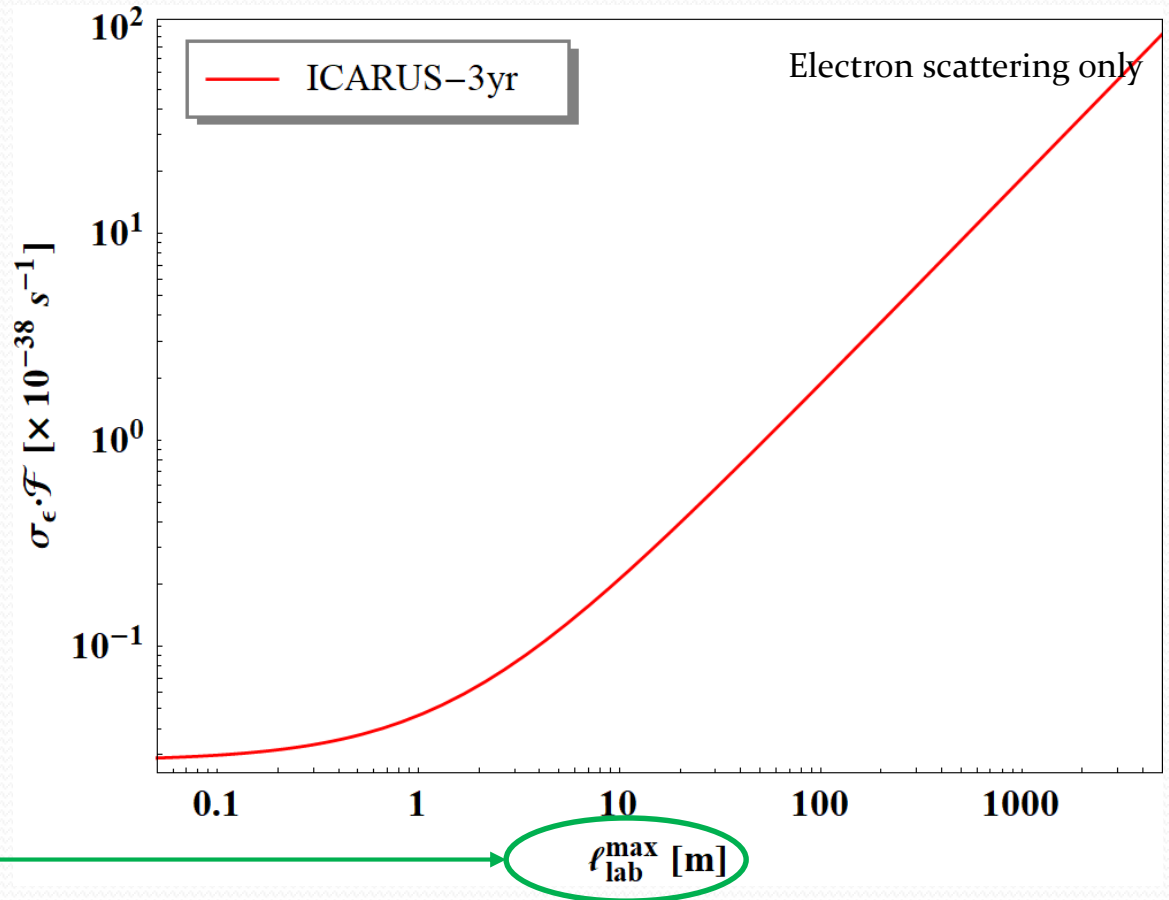
Two completely separated modules are assumed.

90% C.L. with zero BG

$$\sigma_\epsilon \mathcal{F} > \frac{2.3}{A(\ell_{\text{lab}}) t_{\text{exp}} N_e}, \quad (4)$$

Evaluated under the assumption of cumulatively isotropic χ_1 flux

ℓ_{lab} different event-by-event, so taking $\ell_{\text{lab}}^{\text{max}}$ for more conservative limit



Model-independent Reach: More Familiar Form

- More familiar parameterization possible with the below modification!

$$\sigma_{\epsilon} \geq \frac{2.3}{\mathcal{F} \cdot A \cdot t_{\text{exp}} \cdot N_e}$$

$$\mathcal{F} = 1.6 \times 10^{-4} \text{ cm}^{-2} \text{ s}^{-1} \times \left(\frac{\langle \sigma v \rangle_{0 \rightarrow 1}}{5 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}} \right) \times \left(\frac{\text{GeV}}{m_0} \right)^2, \quad (1)$$

set to be $5 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$

- Then having

$$\sigma_{\epsilon} \text{ vs. } m_0 (= E_1 = \gamma_1 m_1)$$

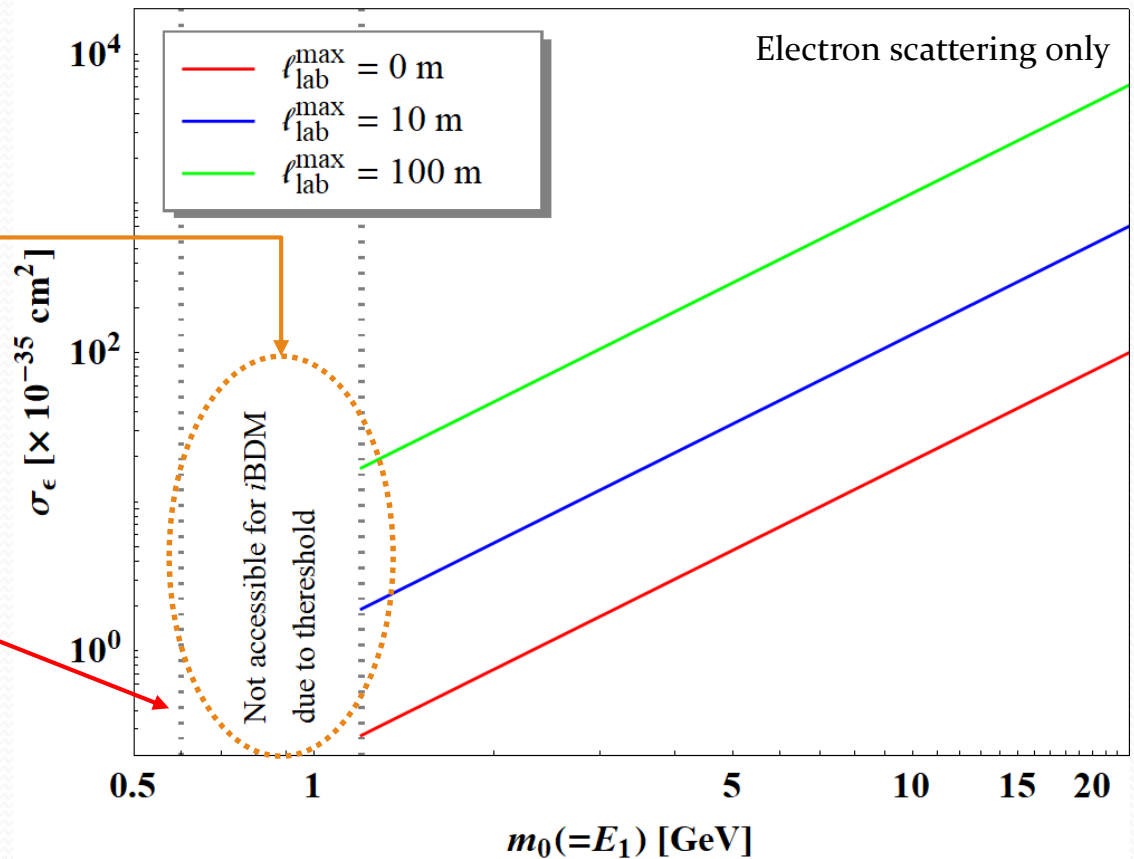
just like σ vs. m_{DM} in conventional WIMP searches

Model-independent Reach: More Familiar Form

3-year data collection assumed.

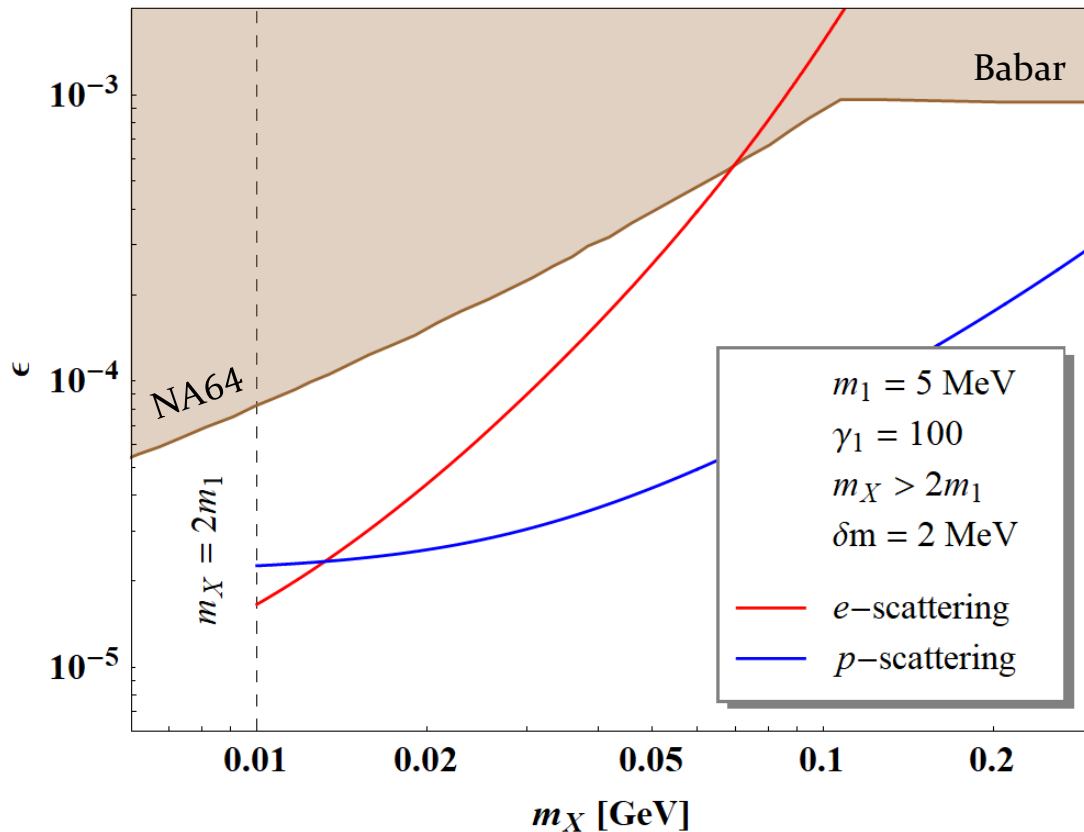
Absolute lower bound for visible tri-track events due to the threshold energy of 400 MeV. (The actual lower bound may involve minor model-dependence.)

Smaller thresholds allow to probe smaller cosmological dark matter mass.



Dark Photon Parameter Space: Invisible X Decay

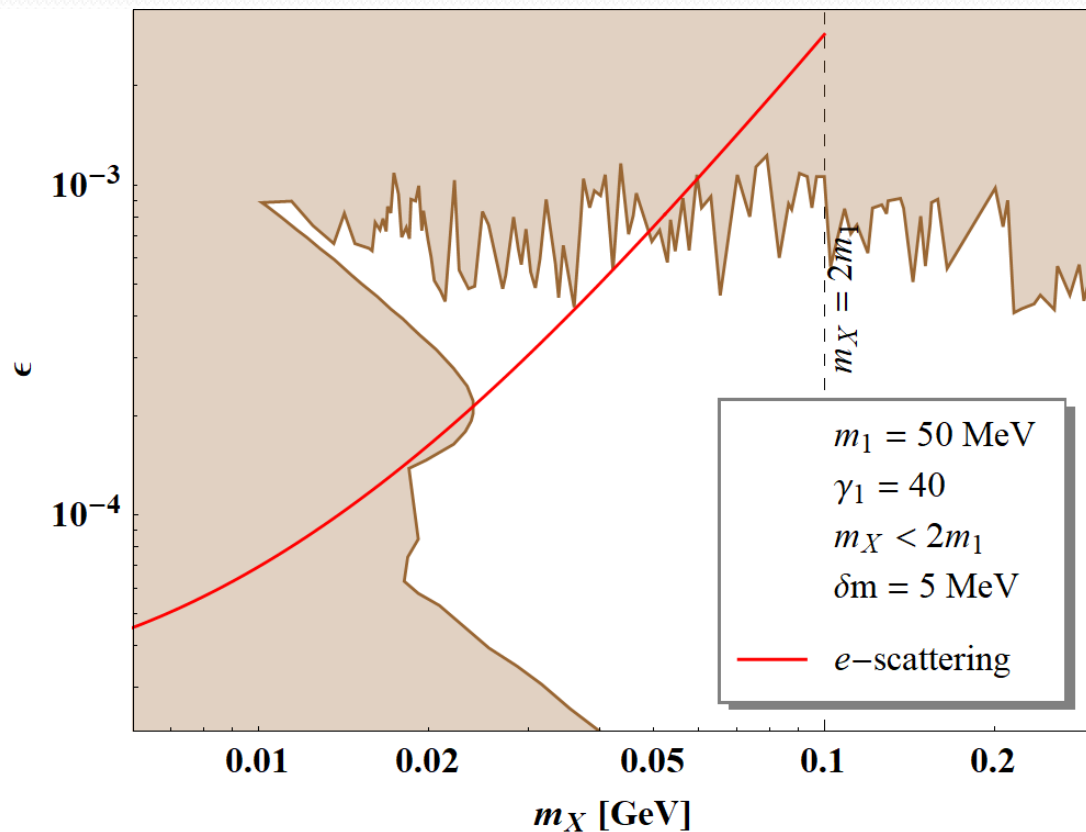
□ Case study 1: mass spectra for which dark photon decays into DM pairs, i.e., $m_X > 2m_1$



- 3-year data collection is assumed.
- 400 MeV threshold is assumed.
- ICARUS can probe the uncovered parameter region by an order of magnitude in the ϵ axis.
- p -scattering is preferred for heavier dark photon masses [DK, Machado, Park, Shin, in progress]

Dark Photon Parameter Space: Visible X decay

- Case study 2: mass spectra for which dark photon decays into lepton pairs, i.e., $m_X < 2m_1$



- 3-year data collection is assumed.
- 400 MeV threshold is assumed.
- ICARUS can probe the uncovered parameter region by half order of magnitude in the ϵ axis.