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SLAC NATIONAL
ACCELERATOR
LABORATORY

(arXiv:1808.10446)

(to appear)

**Gravity amplitudes
from the UV**

Amplitudes 2019 - Dublin
07/05/2019

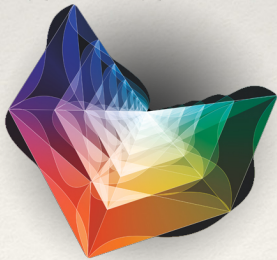
(0) Motivation

- UV behavior of (super-)gravities still mysterious

loop	(divergence) \times $(4\pi)^D (-i)^{L+1} \epsilon^{L+1} (i)^{L+1} \kappa^{2L}$
1-30	$-\frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] + \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] - \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $- 82 \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] - \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] + \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $- \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] - \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] - \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$
31-60	$\frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] + \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] + \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ 82 \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] + \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] - \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] + \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $- \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] + \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] - \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$
61-82	$-\frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] + \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] - \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $- 82 \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] - \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] + \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $- \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] - \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$ $+ \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right] - \frac{1}{2} \left[\frac{1}{(2\pi)^{2L}} \int d^D k \dots \right]$
sum	$1 = \frac{1}{2}(204 - 1)$

- “enhanced” cancellations, e.g. in
 - N=4 sugra, d=5 @ L=2; [Bern,Davies,Dennen,Huang: 1209.2472]
 - N=5 sugra, d=4 @ L=4 [Bern,Davies,Dennen: 1409.3089]
- N=8 sugra: 5-loop divergence in d=24/5 [Bern,Carrasco,Chen,Edison,Johansson,Parra-Martinez,Roiban,Zeng: 1804.09311]

- Beautiful picture of N=4 SYM, where nontrivial analytic properties are understood at integrand level [talk L. Ferro]



- Can we make similar progress in gravity?

Properties of amplitudes should be encoded before integration
 (integration = tracing out = forgetting details)

Outline

- ❖ i) Gravity integrands and the UV structure:
 - off-shell nonplanar integrands?
 - unitarity cuts, maximal cuts & power counting
 - multi-particle unitarity cuts and integrand enhanced cancellations
- ❖ ii) gravity integrand construction from UV constraints
 - unique construction of 2 & 3-loop gravity integrands
- ❖ iii) Conclusions & Outlook

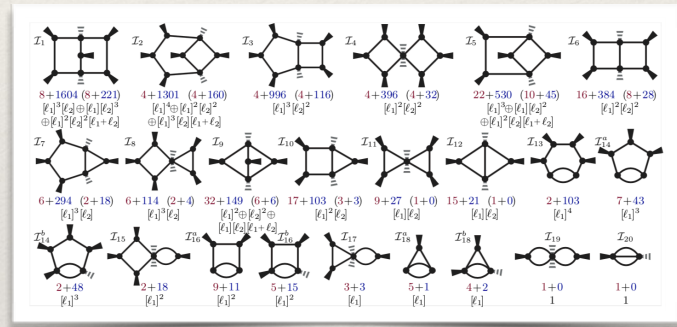
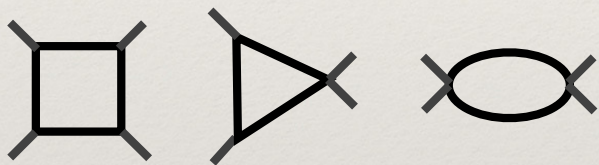
i-1) gravity integrands and the UV structure

- loop-amplitudes and integrands:

$$\mathcal{M}_n^{(L)} = \sum_k c_k \int \mathcal{I}_k(p_i, \ell_j) d^D \ell_1 \cdots d^D \ell_L$$

kinematic coefficients

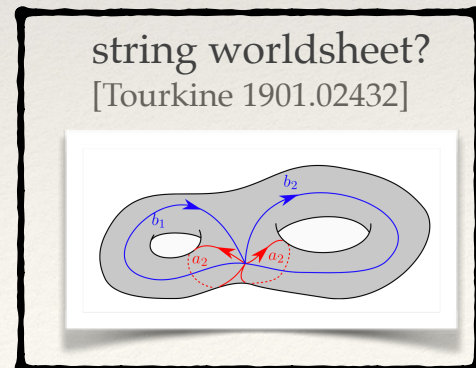
basis integrands [talk J. Bourjaily]



- line-1 problem: no *global* labels of an “integrand”

$$\mathcal{I}(p_i, \ell_j) \stackrel{?}{=} \sum_k c_k \mathcal{I}_k(p_i, \ell_j) \quad \times$$

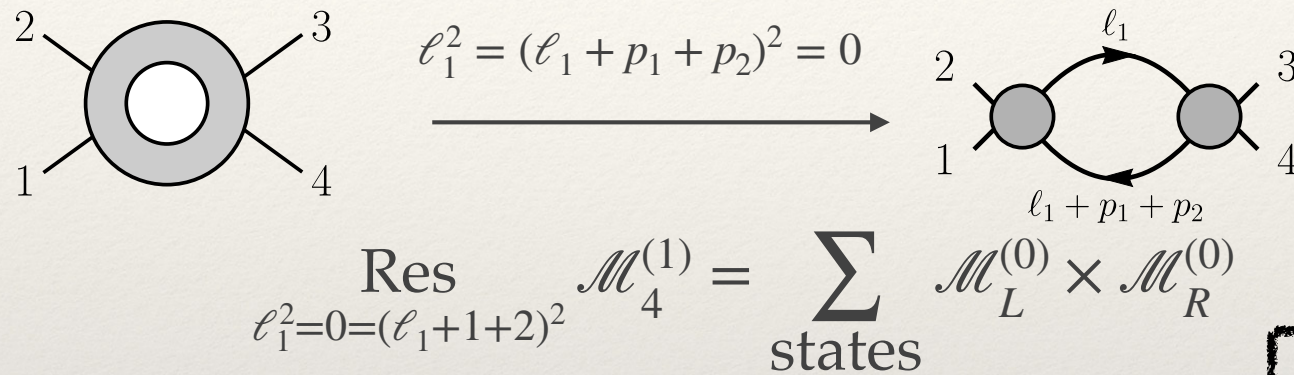
→ Study unitarity cuts of integrands!



string worldsheet?
[Tourkine 1901.02432]

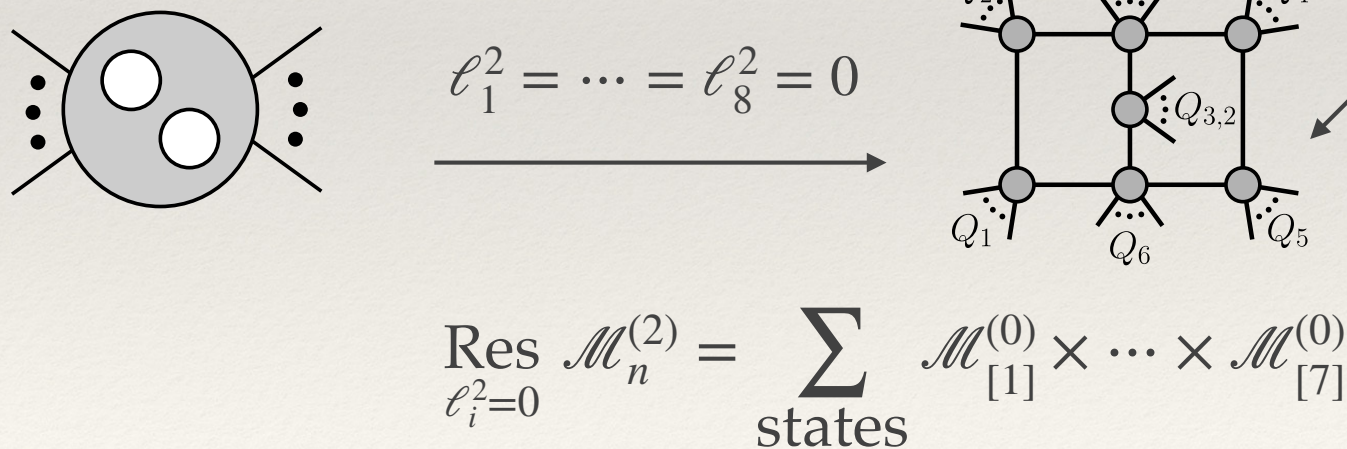
i-2) cuts of loop-integrands

❖ unitarity cut: [Cutkosky 1960]



on-shell functions

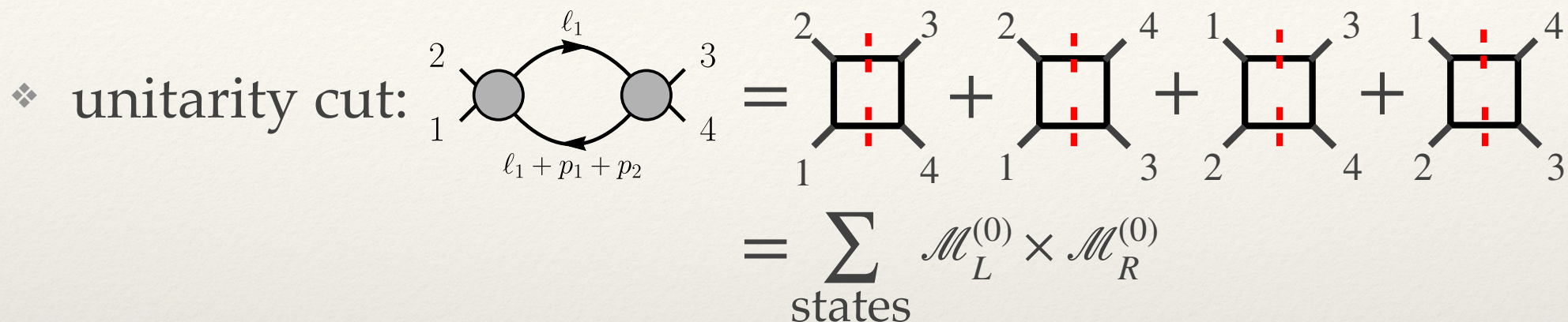
❖ generalized unitarity: [Bern,Dixon,Kosower:9708239,0404293; Britto,Cachazo,Feng:0412103]



→ well defined loop-variables on the cut

i-2) cuts of loop-integrands

❖ unitarity cut:

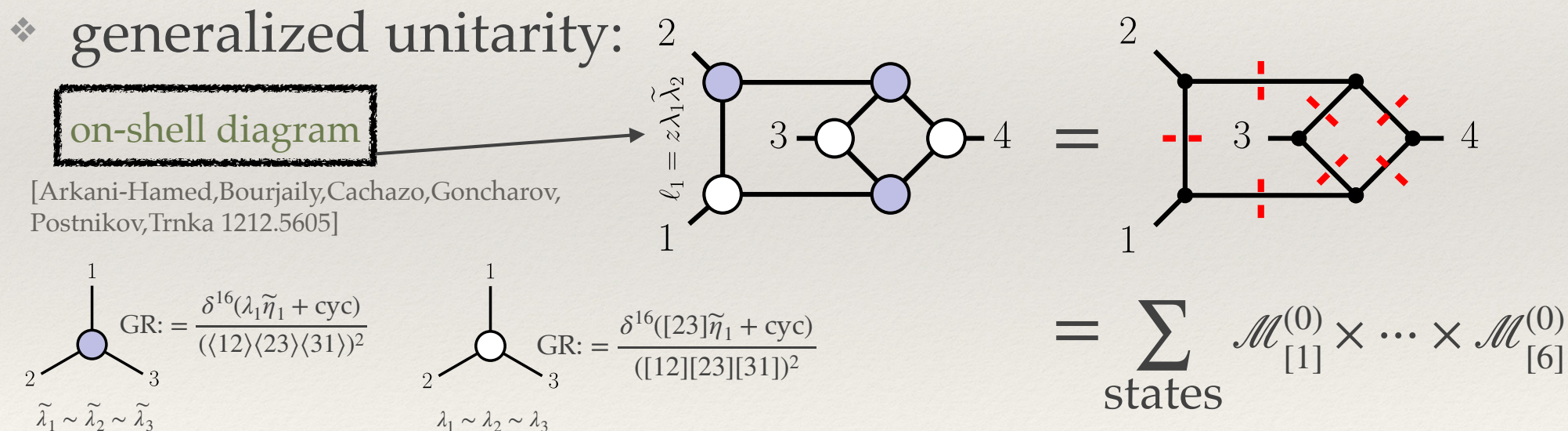


$$= \sum_{\text{states}} \mathcal{M}_L^{(0)} \times \mathcal{M}_R^{(0)}$$

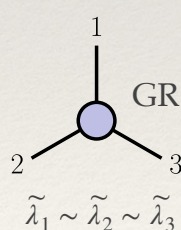
❖ generalized unitarity:

on-shell diagram

[Arkani-Hamed, Bourjaily, Cachazo, Goncharov, Postnikov, Trnka 1212.5605]

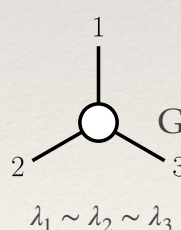


$$= \sum_{\text{states}} \mathcal{M}_{[1]}^{(0)} \times \dots \times \mathcal{M}_{[6]}^{(0)}$$



GR: $= \frac{\delta^{16}(\lambda_1 \tilde{\eta}_1 + \text{cyc})}{(\langle 12 \rangle \langle 23 \rangle \langle 31 \rangle)^2}$

$\tilde{\lambda}_1 \sim \tilde{\lambda}_2 \sim \tilde{\lambda}_3$

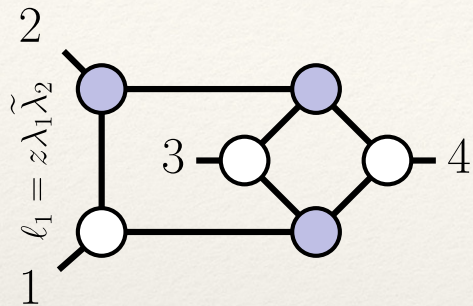


GR: $= \frac{\delta^{16}([23] \tilde{\eta}_1 + \text{cyc})}{([12][23][31])^2}$

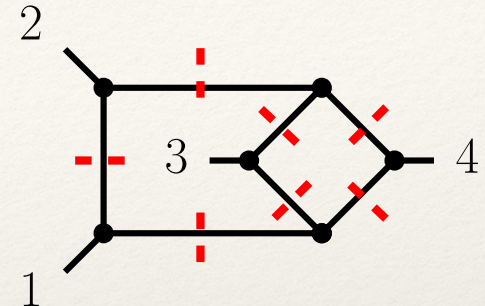
$\lambda_1 \sim \lambda_2 \sim \lambda_3$

➔ maximal cut: isolates single diagram

i-2) maximal cuts and power counting



=

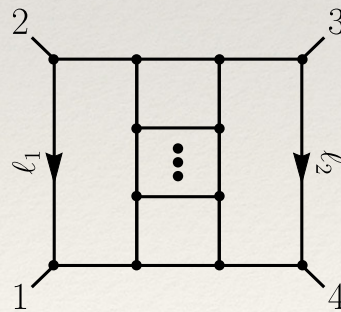
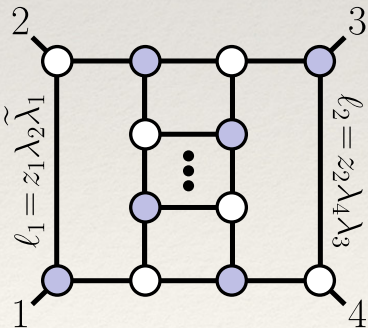


$$\sum_{\text{states}} \mathcal{M}_{[1]}^{(0)} \times \dots \times \mathcal{M}_{[6]}^{(0)} = f(z; p_i) = \frac{[12][34]^2}{z \langle 12 \rangle \langle 13 \rangle (z \langle 13 \rangle + \langle 23 \rangle) (z \langle 14 \rangle + \langle 24 \rangle) \langle 41 \rangle}$$

$$= \text{Res}_{\ell_i^2=0} \frac{\text{Num}(p_i, \ell_j)}{\prod \ell_i^2} = \frac{\text{Num}(p_i, \ell_j^*)}{\mathcal{F}}$$

$$\mathcal{F} = z s_{12} \langle 14 \rangle (z[24] - [14]) \langle 13 \rangle (z[23] - [13])$$

→ Num(p_i, ℓ_j^*) loop-independent



$$N(\ell_i, p_j) \sim \begin{cases} f_{\text{YM}}(p_j) \times (\ell_1 \cdot \ell_2)^{L-3}, & \mathcal{N} = 4 \text{ SYM} \\ f_{\text{GR}}(p_j) \times (\ell_1 \cdot \ell_2)^{2(L-3)}, & \mathcal{N} = 8 \text{ sugra} \end{cases}$$

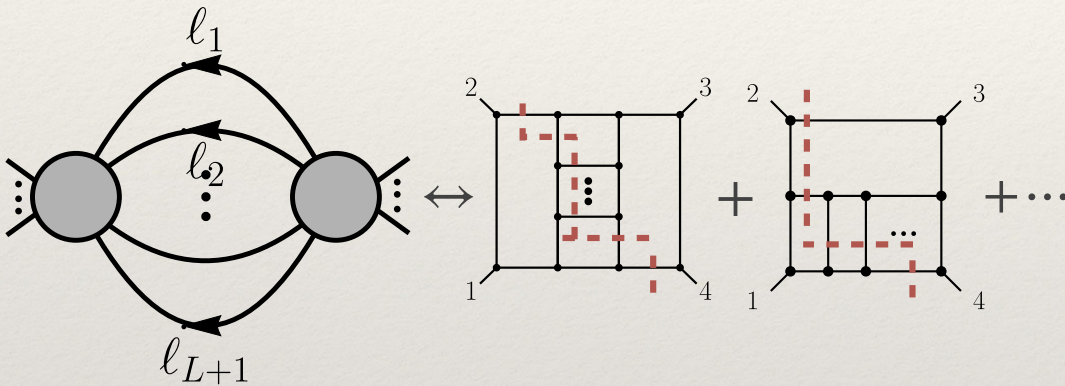
e.g. in [Bern,Davies,Dennen 1409.3089]

z-scaling of maximal cut dictates power counting of numerator!

i-3) multi-particle unitarity cuts and “enhanced cancellations”

[different cuts analyzed in Bern,Dixon,Roiban: 0611086 based on no-triangle hypothesize for 1-loop amplitudes in N=8 sugra]

maximal cut isolates single diagram \rightarrow cuts with minimum # of on-shell props
 No cancellations cancellations between diagrams possible



- ❖ (L+1) legs cut
- ❖ almost all diagrams contribute
- ❖ flexibility to “approach” infinity

D-dimensional parametrization of on-shell loop momenta:

$$\begin{aligned}
 \ell_i &\rightarrow \ell_i + zq_i & i &= 1 \dots L \\
 q_i \cdot q_j &= 0 & \ell_i \cdot q_i &= 0 \\
 \Rightarrow \ell_i \cdot x &\rightarrow \ell_i \cdot x + zq_i \cdot x
 \end{aligned}$$

$\xrightarrow{D=4 \text{ limit}}$

$$\ell_i = \lambda_{\ell_i} \left(\tilde{\lambda}_{\ell_i} + zc_i \tilde{\chi} \right)$$

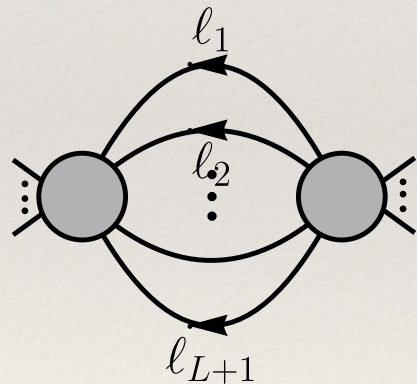
gravity cut scaling compared to a) SYM b) individual diagrams ?

i-3) multi-unitarity cut UV scaling

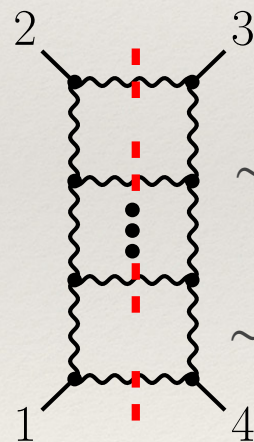
❖ Naive scaling of Feynman diagrams

$$(\ell_i \cdot \ell_j) \sim z, \quad \ell^2 \sim z, \quad \epsilon \cdot \ell \sim z$$

$2L+2$ vertices $3L+1$ props - $(L+1)$ cuts = $2L$ uncut props



(super-) Yang-Mills

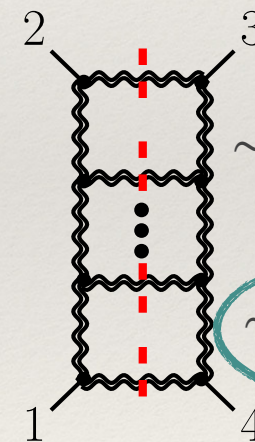


$$\sim \frac{(\epsilon \cdot \ell)^4 [\ell_i \cdot \ell_j]^{\frac{2L+2-4}{2}}}{[\ell^2]^{2L}}$$

$$\sim \frac{z^{L+3}}{z^{2L}}$$

$$= z^{-L+3} \xrightarrow{\text{susy}} z^{-L-1}$$

(super-) gravity

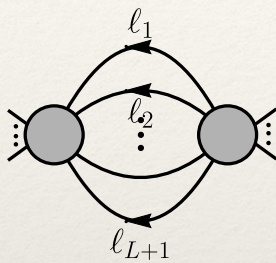
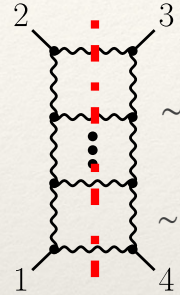
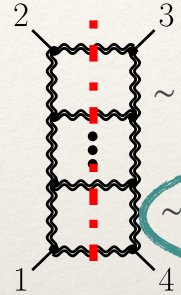


$$\sim \frac{(\epsilon \cdot \ell)^8 [\ell_i \cdot \ell_j]^{2L+2-4}}{[\ell^2]^{2L}}$$

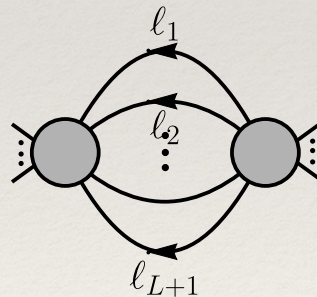
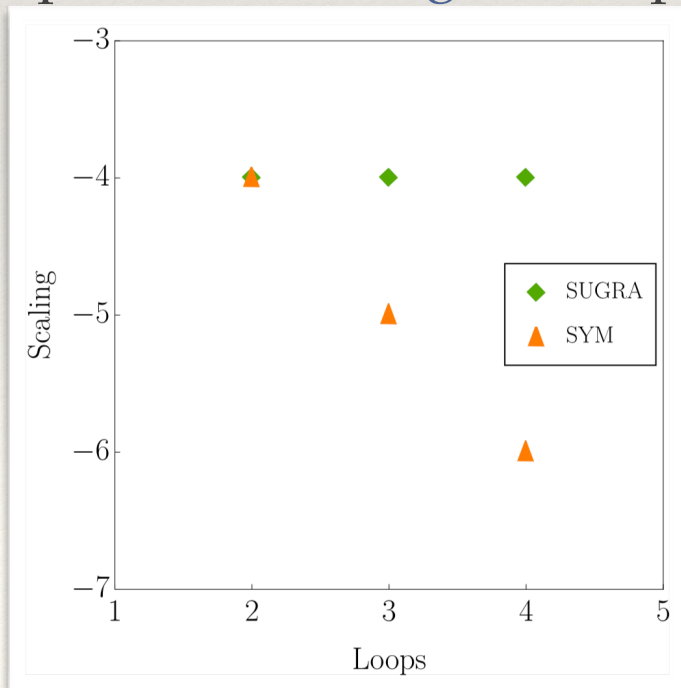
$$\sim z^6 \xrightarrow{\text{susy}} z^{-2}$$

Compare diagram scaling to full cut?

i-3) cut vs. diagram UV scaling

$(\ell_i \cdot \ell_j) \sim z, \ell^2 \sim z, \epsilon \cdot \ell \sim z$	sYM	suGRa
	 $\sim \frac{(\epsilon \cdot \ell)^4 [\ell_i \cdot \ell_j]^{\frac{2L+2-4}{2}}}{[\ell^2]^{2L}}$ $\sim \frac{z^{L+3}}{z^{2L}} = z^{-L+3} \xrightarrow{\text{susy}} z^{-L-1}$	 $\sim \frac{(\epsilon \cdot \ell)^8 [\ell_i \cdot \ell_j]^{2L+2-4}}{[\ell^2]^{2L}}$ $\sim z^6 \xrightarrow{\text{susy}} z^{-2}$

❖ explicit local integrand representations and check **d-dim** large z-scaling



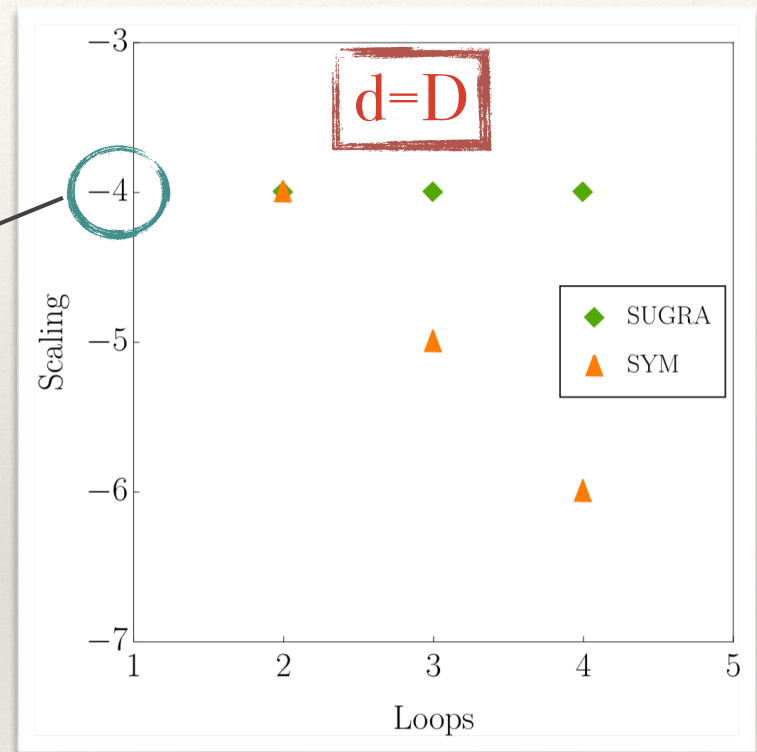
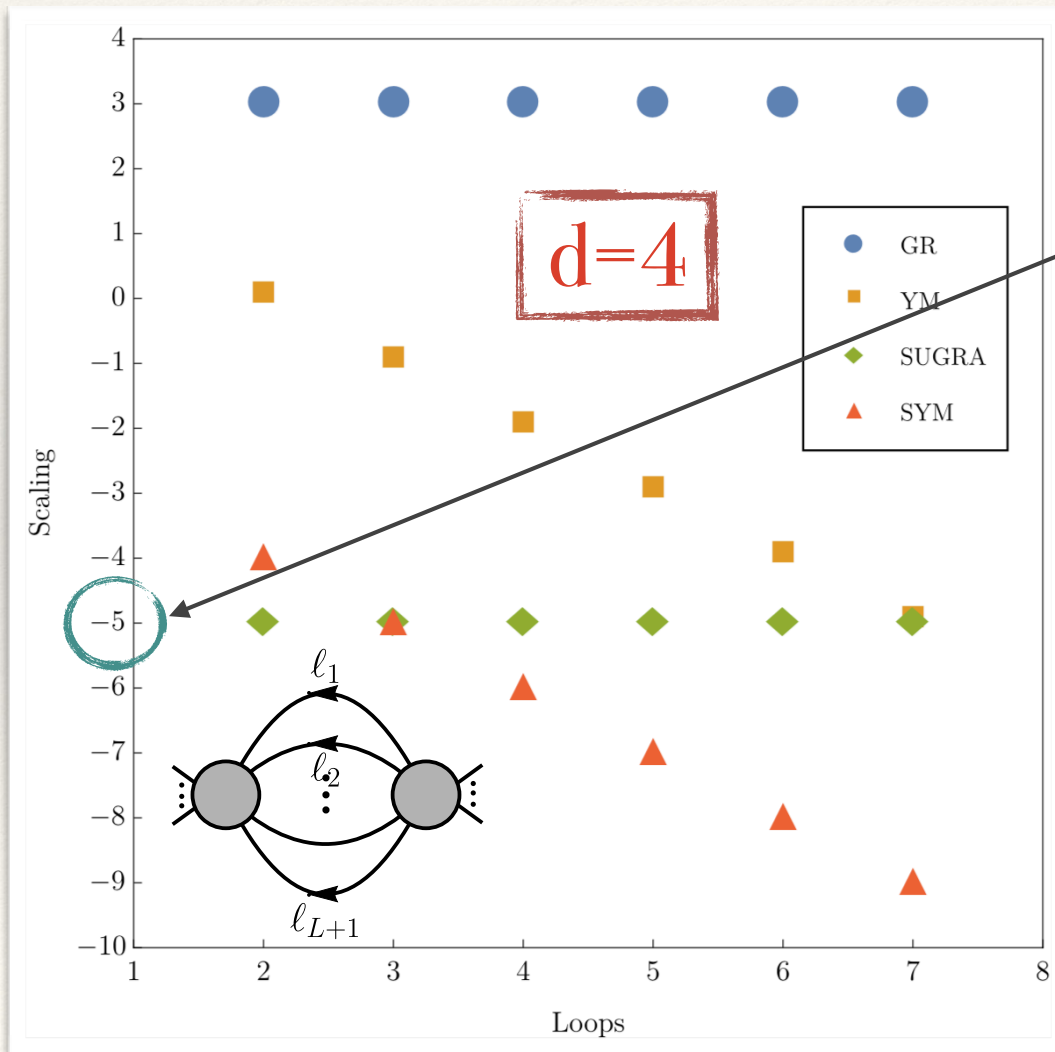
SYM : z^{-L-2} vs. naive z^{-L-1}

sugra : z^{-4} vs. naive z^{-2}

Is there anything special about $d=4$?

i-3) $d=4$ UV scaling of cut

- use KLT and BCFW for N=4 SYM amps to check 4-dim large z-scaling



In $d=4$, SUGRA drops 1 power in large z scaling!

@ L=2: drop due to Gram determinant that vanishes in $d=4$

ii) gravity integrand construction from UV constraints

- ❖ Turn things upside down

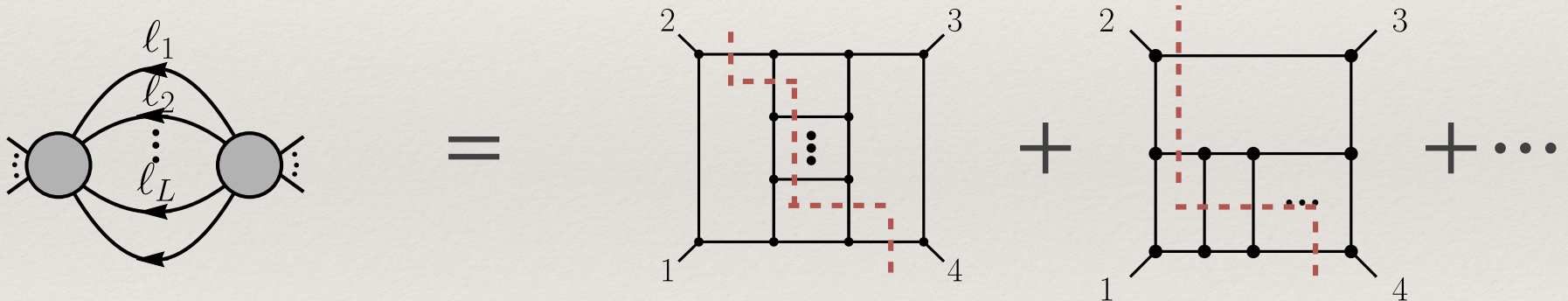


Impose the good large- z scaling on multi-particle unitarity cuts
as constraints on an ansatz?

Is there a **unique** object that satisfies the large- z scaling?

ii) gravity integrand construction from UV constraints

- ❖ write numerator ansatz for each local diagram
 - triangle-power counting
 - numerator respects diagram symmetries
- ❖ **multi-unitarity cut!** $L+1$ props on-shell



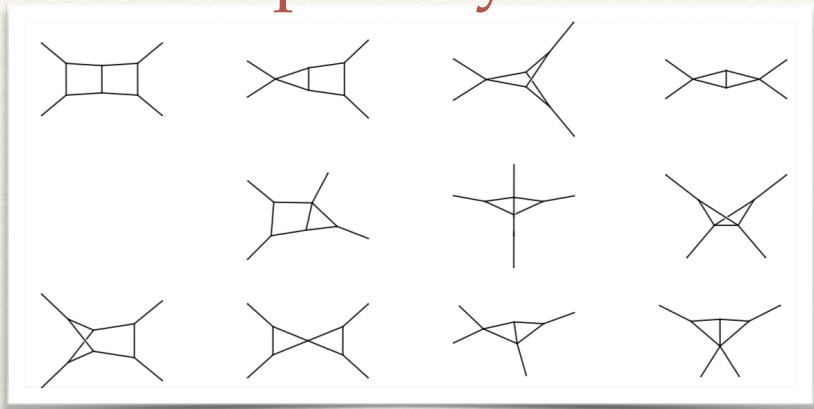
$$\mathcal{M}(z) |_{\text{cut}} \sim z^a = z^{b_1} + z^{b_2} + z^{b_3} + \dots$$

- ❖ demand: $\forall b_i > a$, set coefficients to zero as $\ell_i(z) \xrightarrow{z \rightarrow \infty} \infty$

Can impose the scaling in all directions at infinity!

ii) gravity integrand construction from UV constraints

❖ 2-loop analysis:



$$\left. \begin{aligned} &\sim \frac{1}{z^5} \\ &\sim \frac{1}{z^4} \end{aligned} \right\} \begin{aligned} N_{db} &= N_{npdb} = a s^2 \\ N_{\text{other}} &= 0 \end{aligned}$$

❖ 3-loop analysis:

~ 2700 parameters, 83 diagrams

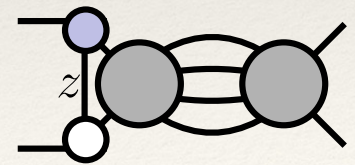
$$\begin{aligned} &\sim \frac{1}{z^5} & & \sim \frac{1}{z^5} \\ &\sim \frac{1}{z^5} & & \sim \frac{1}{z^4} \end{aligned}$$

50 unfixed dof

problem:

$$= a_1 s^4 + a_2 s^3 t + a_3 s^2 t^2 \sim \frac{1}{z^6}$$

distinguish via BCFW-shift
of external momenta



We find a **unique** solution after shifting external kinematics as well

iii) Conclusions & Outlook

- ❖ analyzed cuts of (super-)gravity integrands through 7-loops!
 - ❖ improvement of d-dimensional cut scaling
 - ❖ further drop in UV scaling in d=4
 - status of 7-loop divergence?
 - ❖ how much of the UV scaling can be understood from trees?
 - analyze more general deformations of gravity tree-amplitudes
- ❖ imposing UV constraints on ansatz fixed N=8 sugra amplitude uniquely to L=3, n=4
 - ❖ How does this extend to higher loops, higher points? [Bourjaily,EH,Trnka 1812.11185]
 - ❖ Including deformation of external momenta, do we have the full list of homogeneous constraints that “define” gravity?
 - ❖ Can we “geometrize” these properties?

THANK YOU FOR YOUR ATTENTION