A critical look at the *B-*decay anomalies at the end of 2019

WHEPP XVI, IIT Guwahati December 1-10, 2019

A critical look at the *B-*decay anomalies at the end of 2019 *An experimental*

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Disclaimer

- Not a real expert, will try to summarize and introduce needed information for the anomalies.
- Hopeful that the material cover is sufficient for most of the people to follow the "discussions to follow" in the coming days.
- Apologies if some statement(s) is(are) wrong.
- Will introduce "*B-anomalies*" and their implications to be explained in detail by *Prof. R. Mohanta* "Combined explanation R_D/R_{D^*} and R_K/R_{K^*} " on Wednesday.

B decay anomalies

Two words are enough to wind up my presentation:

Lepton Flavor Universality Violation (LFUV) in *B* decays

$$
B \to D^{(*)} \tau \nu
$$

Not a rare decay, Branching fraction : 1-2 %

Belle, PRL99, 191807(2007)

 \triangleright Weak interaction can be tested precisely. \triangleright Allow one to test theory of low-energy QCD, which contribute to $\bar{B} \to D^*$ transition.

 $\bar B\to D^*\tau^-\bar\upsilon_\tau$ decays are sensitive to new scalar fields (New Physics at tree level)

BaBar, PRL100, 021801(2008)

e^+e^- colliders for B physics

How one measures $R_{D^{(*)}}$ at e^+e^- colliders

Hadronic tag

Beginning of "the R_{D^*} saga" BaBar

PRL 109, 101802(2012) PRD 88, 072012 (2013)

Results

BaBar PRL 109, 101802(2012) PRD 88, 072012 (2013)

How one should interpret the result BaBar PRD 88, 072012 (2013)

How one should interpret the result can BaBar PRD 88, 072012 (2013)

PS: Uncertainties in the extrapolation of type II 2HDM are not included

Hadronic tag Belle turn Belle turn Belle PRD92, 072014 (2015)

 $\tau \rightarrow \mu \nu \nu$ and $\tau \rightarrow e \nu \nu$ Four signal sample $D^0\ell$, $D^{*0}\ell$, $D^+\ell$ and $D^{*+}\ell$

We reconstruct D^+ mesons in the decays to $K^-\pi^+\pi^+$, $K_S^0 \pi^+$, $K_S^0 \pi^+ \pi^0$, and $K_S^0 \pi^+ \pi^+ \pi^-$; D^0 mesons to $K^- \pi^+$, $K^-\pi^+\pi^+\pi^-$, $K^-\pi^+\pi^0$, $K_S^0\pi^0$, and $K_S^0\pi^+\pi^-$; D^{*+} mesons to $D^0\pi^+$ and $D^+\pi^0$; and D^{*0} mesons to $D^0\pi^0$ and $D^0\gamma$.

Interpretation

Belle PRD92, 072014 (2015)

Belle result lies between SM expectation and the BaBar result (compatible with both).

Also compatible with 2HDM of type II around $tan\beta/m_{H^{\pm}}$ =0.50GeV⁻¹

Semileptonic *B* tag **Different Method** Belle PRD94, 072007 (2016)

Independent analysis of previous Belle measurement. Better efficiency in tagging *(somewhat double)* ν More background due to missing particle in tag side

2D fit to neural network output (O_{NB}) and E_{ECL}

 E_{ECL} : Sum of energy in ECL not associated with reconstruction

More distribution Belle PRD94, 072007 (2016)

Allow additional contribution from scalar and vector operators while disfavouring large additional contributions from a tensor operator with $+0.34 < C_T < +0.39$, and R_2 -type leptoquark model with $+0.34 < C_T < +0.38$, or an S_1 -type leptoquark model with $+ 0.22 < C_T < +0.28$

Other variable : $P_{\tau}(D^*)$

Hadronic Tag

 τ polarization :

 $P_{\tau}(D^*)_{SM} = -0.497 \pm 0.013$

M. Tanaka and R. Watanabe PRD87, 034028 (2013)

 τ polarization is sensitive to the NP contribution.

One can measure the τ polarization using its two-body decay

Hadronic Tag

Belle measures $P_{\tau}(D^*)$ Belle

PRL118,211801 (2017) PRD97, 012004 (2018)

New variable to search for NP

Provide independent study/confirmation of previous measurements

Angular distribution of τ decay

$$
\frac{1}{\Gamma} \frac{d\Gamma}{d\cos\theta_{\text{hel}}} = \frac{1}{2} \left[1 + \alpha P_{\tau}(D^*) \cos\theta_{\text{hel}} \right] \quad \alpha = \begin{cases} 1 \text{ for } \tau \to \pi\nu \\ 0.45 \text{ for } \tau \to \rho\nu \end{cases}
$$

 τ rest frame estimated

 $q = p_{e^+e^-} - p_{tag} - p_{D^*}$

$$
E_{\tau} = \frac{q^2 + m_{\tau}^2}{2\sqrt{q^2}} \qquad |\vec{p}_{\tau}| = \frac{q^2 - m_{\tau}^2}{2\sqrt{q^2}}
$$

$$
\cos \theta_{\tau d} = \frac{2E_{\tau}E_d - m_{\tau}^2 - m_d^2}{2|\vec{p}_{\tau}||\vec{p}_d|}
$$

$$
|\vec{p}_d^{\tau}| \cos \theta_{\text{hel}} = -\gamma |\vec{\beta}| E_d + \gamma |\vec{p}_d| \cos \theta_{\tau d}.
$$

Simultaneous fit to 8 samples $[B^0, B^+] \otimes [\pi^+, \rho^+] \otimes [\cos \theta_{hel} > 0, \cos \theta_{hel} < 0]$

 $D^0 \to K_S^0 \pi^0, \pi^+ \pi^-, K^- \pi^+, K^+ K^-,$ $K^-\pi^+\pi^0$, $K^0_S\pi^+\pi^-$, $K^0_S\pi^+\pi^-\pi^0$, $K^-\pi^+\pi^+\pi^-$, $D^+\to K^0_S\pi^+$, $K_S^0 K^+$, $K_S^0 \pi^+ \pi^0$, $K^- \pi^+ \pi^+$, $K^+ K^- \pi^+$, $K^- \pi^+ \pi^+ \pi^0$, and $K_S^0 \pi^+ \pi^+ \pi^-$

Result

 $R(D^*) = 0.270 \pm 0.035 \text{(stat)}_{-0.025}^{+0.028} \text{(syst)}$

 $P_{\tau}(D^*) = -0.38 \pm 0.51(\text{stat})_{-0.16}^{+0.21}(\text{syst}),$

 $SM = -0.497 \pm 0.013$

Measurement is consistent with the SM prediction Excludes $P_{\tau}(D^*)$ larger than 0.5 at 90% CL.

 $R(D^*) = 0.295 \pm 0.038 \pm 0.015$ *Previous from Hadron Tag*

 $R(D^*)_{SM} = 0.252 \pm 0.003$

Belle

Signal significance of 7.1 σ PRL118,211801 (2017) PRD97, 012004 (2018)

D^* polarization

 D^* polarization can give clue about the NP signature

$$
\frac{1}{\Gamma} \frac{d\Gamma}{d\cos\theta_{\text{hel}}} = \frac{3}{4} (2F_L^{D^*} \cos^2\theta_{\text{hel}} + (1 - F_L^{D^*}) \sin^2\theta_{\text{hel}})
$$

 θ_{hel} = angle in D^{*}- rest frame between D0 and B0 flight $F_L^{D^*}$ fraction of longitudinal polarization of D*

Evis and Xmiss $X_{miss} =$ $E_{miss} - |p_{D^*} - p_{d_{\tau}}|$ $E_{beam}^2 - m_{B^0}^2$

Peaking background from $B \to D^*(X) \ell \nu$

- All τ decays are useful (cross-feed no effect)
- Strong dependence on $cos\theta_{hel}$ and q^2 due to slow π from D^* [softer at $(cos \theta_{hel}) > 0$].

Measured F_L for $F_L^{D^*}(B \to D^{*-}e^+\nu) = 0.56 \pm 0.02$ SM : 0.54

First measurement of D^* polarization Belle, arXiv:1903.03102

 $F_L^{D^*} = 0.60 \pm 0.08 \pm 0.04$

$$
\left(F_L^{D^*}\right)_{SM} = 0.457 \pm 0.010
$$

Li et al PRD98, 095018(2018)

Agrees with SM within 1.6 σ

Update on semileptonic tagging by Belle

Updated result

Belle arXiv:1910.05864 *submitted to PRL*

 $R(D) = 0.307 \pm 0.037 \pm 0.016$ $R(D^*) = 0.282 \pm 0.018 \pm 0.014$

Most precise measurement to date !

 $R(D^*) = 0.302 \pm 0.030 \pm 0.011$ Previous semileptonic measurement $\mathcal{R}(D)_{SM} = 0.300 \pm 0.011$ $R(D^*)_{SM} = 0.252 \pm 0.003$

LHCb : beauty at the beast

LHCb is a single arm spectrometer optimized for beauty and charm physics at large η

Excellent vertex resolution: $20 \mu m$ resolution on impact parameter.

Excellent particle identification

Dipole magnet polarity periodically flipped to change the sign of many reconstruction asymmetries

Semileptonic decays : High trigger efficiency

 momentum direction is determined from unit vector to B decay vertex from the associated PV. Component of B momentum along the B momentum along the beam axis is approximated using $(p_B)_z = ({}^{m_B} /_{m_{reco}}) (p_{reco})_z$

Deviate from SM LHCb, PRL 115, 111803 (2015)

 $3fb^{-1}$ 2011-2012 data

$$
B → D*τν
$$

\n
$$
B → D*Hc(→ ivX')X
$$

\n
$$
B → D*tν
$$

\n
$$
B → D*tν
$$

\n
$$
B → D*tν
$$

\nCombinatorial
\nMisidentified μ

$$
\sim 16400 \text{ events for } \overline{B^0} \to D^{*+} \tau^- \overline{\nu_\tau}
$$

 $R(D^*) = 0.336 \pm 0.027 \pm 0.030$ $(2.1\sigma$ from SM)

 $R(D^*)_{SM} = 0.252 \pm 0.003$

- Absence of charged lepton avoids background from semileptonic $b \rightarrow c$ decays
- Three prong helps in getting precise τ decay vertex.
- Only one ν emitted at the τ vertex.
- Background $B \to D^{*-}DX$ leads to nice mass peaks and not the signal. Provide handle to control various background.

$$
R_{D^*}
$$
 using 3-prong $\tau^- \to \pi^- \pi^+ \pi^- (\pi^0) \nu$ decays

Signal candidates are required to be well isolated.

Events with extra charged particles pointing to B and/or τ vertices are vetoed.

$$
\mathcal{K}(D^{*-}) \equiv \frac{\mathcal{B}(B^0 \to D^{*-}\tau^+\nu_\tau)}{\mathcal{B}(B^0 \to D^{*-}3\pi)} = \frac{N_{\text{sig}}}{N_{\text{norm}}} \frac{\varepsilon_{\text{norm}}}{\varepsilon_{\text{sig}}} \frac{1}{\mathcal{B}(\tau^+ \to 3\pi\bar{\nu}_\tau) + \mathcal{B}(\tau^+ \to 3\pi\pi^0\bar{\nu}_\tau)}
$$

Most of the systematic cancel in this ratio.

$$
\mathcal{R}(D^*) = \mathcal{K}(D^{*-}) \times \frac{\mathcal{B}(B^0 \to D^{*-} \pi^+ \pi^- \pi^+)}{\mathcal{B}(B^0 \to D^{*-} \mu^+ \nu_{\mu})}
$$

 N_{sig} from a 3D fit to q^2 (8 bins), 3π decay time (8 bins) and BDT (4 bins) $N_{\rm sig}$ from a fit to $M(D^{*-}\pi^+\pi^-\pi^+)$

LHCb,PRD 97, 072013 (2018)

$$
R_{D^*}
$$
 using 3-prong $\tau^- \to \pi^- \pi^+ \pi^- (\pi^0) \nu$ decays

BDT trained to suppress main background from $D^{*-}D_s^+X$ events

Input variables: 3π dynamics, $D^*3\pi$ dynamics, neutrals isolation

Result by 3 prong

LHCb,PRD 97, 072013 (2018)

 $3fb^{-1}$ data

 $R(D^*) = 0.291 \pm 0.019 \pm 0.026 \pm 0.013$ $(1.1\sigma$ from SM)

 $R(D^*) = 0.336 \pm 0.027 \pm 0.030$ $(2.1\sigma$ from SM) Previous one prong result

LHCb, PRL 115, 111803 (2015)

 $R(D^*)_{SM} = 0.252 \pm 0.003$

Current scenario

Skip

Anomaly $\mathcal{R}_{I/\psi}$

LHCb,PRL120, 121801 (2018)

 $= 0.71 \pm 0.17$ (stat) ± 0.18 (syst).

SM predicts within range: 0.25-0.28

$h \rightarrow s \ell \ell$

Rare decays, FCNC

- Small branching fraction : $\mathcal{O}(10^{-6})$
- \triangleright Decays are sensitive to NP
- \triangleright Modify the decay rate and the angular distribution of final state

- \triangleright Give strong constraints on many BSM by probing energy scales higher than direct searches
- \triangleright Experimentally: full reconstruction but background dominated

Processes which are suppressed or even forbidden in the SM, one expect the NP effect to be relatively large.

Amplitude of a hadron decay process is described as

 $\mathcal{A}(I \to F) = \langle F | \mathcal{H}_{eff} | I \rangle =$ G_F 2 \sum i $V_{CKM}^{i}C_{i}(\mu)\langle F|O_{i}(\mu)|M$ **CKM** couplings coefficients **Wilson** Hadronic Matrix Elements

At μ scale

Wilson Coefficients C_i = Perturbative short distance effects Operators O_i = non-perturbative long distance effects $i = 7$: Photon penguin

i =9,10 : Electroweak penguin

NP modify the SM operator contribution (C_i) and /or enter through new operators

In the SM, couplings of the gauge bosons to leptons are independent of lepton flavour Any sign of lepton non-universal interaction would be a direct sign of new physics

What we observe is effect of the particle involved affecting the B

$$
m_t = 173.0 \pm 0.4 \text{ GeV}/c^2
$$

Provide unique way to look for the (new) physics

Forward-backward Asymmetry in $B \to K^* \ell \ell$

A.Ali et al PLB 273, 505 (1991)

Interference between γ and weak coupling in $b \to s \ell^+ \ell^-$ production give rise to a forwardbackward asymmetry

$$
\frac{N_F - N_B}{N_F + N_B} = \frac{A_{FB}(B \rightarrow K^* \ell^+ \ell^-) = -C_{10} \xi(q^2) \left[Re(C_9)F_1 + \frac{1}{q^2} C_7 F_2 \right] \qquad \text{from factors}
$$
\n
$$
Z \text{ contribution}
$$
\nA. Ali et al PRD66,034002 (2002)\n
$$
= \frac{1}{2 \cdot \text{flipped-sign } C_9 \text{ model}}
$$
\n
$$
= 0.5 \qquad \frac{1}{3 \cdot \text{flipped-sign } C_9 \text{ model}}
$$
\n
$$
= 0.5 \qquad \frac{1}{3}
$$
\n $$

SM predictions are not sensitive to QCD corrections

 $\ell^+\ell^-$ rest frame

Measurement of \mathcal{A}_{FB}

$B \to K^{*0} \ell^+ \ell^-, K^{*+} \ell^+ \ell^-, K^+ \ell^+ \ell^ 357fb^{-1}$ $K^{*0} \to K^+\pi^-$, A_{FB} (bkg-sub)
 \circ
 \circ
 \circ
 \circ $K^{*+} \to K_S^0 \pi^+, K^+ \pi^0$ $A_7 > 0$ SM (solid line) $A_7 = -0.330$; $A_9 = 4.069$; $A_7 = -4.213$; -0.5 $A_{10} > 0$ $A_{10} > 0 A_7 > 0$ $12²$ 2 10 14 16 18 20 q^2 GeV²/c² $\mathcal{A}_{FB}(B \to K^* \ell^+ \ell^-) = 0.50 \pm 0.15 \pm 0.02$ $\mathcal{A}_{FB}(B \to K^+ \ell^+ \ell^-) = 0.10 \pm 0.14 \pm 0.01$

Belle, PRL96, 251801(2006)

It seems $C_7 = -C_7^{SM}$ preferred the \mathcal{A}_{FB} data

Flipping the sign of photonic penguin, would turn on the large rate of $B \to X_s \ell \ell$ (then what seen in experiments) Cite19

Addition to measuring A_{FB} one can also measure K* longitudinal polarization F_L

between I^+ and $B^0(B^+)$ in dilepton rest frame.

$$
\frac{d\Gamma}{d\cos\theta_{B\ell}} = \frac{3}{4}F_L(1 - \cos_{\theta_{B\ell}}^2) + \frac{3}{8}(1 - F_L)\left(1 + \cos_{\theta_{K^*}}^2\right) + \mathcal{A}_{FB}\cos_{\theta_{B\ell}}
$$

TABLE III. Results for the fits to the $K\ell^+\ell^-$ and $K^*\ell^+\ell^$ samples. N_s is the number of signal events in the m_{ES} fit. The quoted errors are statistical only.

| Decay | q^2 | N_S | F_L | $\mathcal{A}_{\rm FB}$ |
|-------------------|-------|----------------|------------------------|-------------------------|
| $K\ell^+\ell^-$ | low | 26.0 ± 5.7 | | $+0.04_{-0.24}^{+0.16}$ |
| | high | 26.5 ± 6.7 | | $+0.20_{-0.22}^{+0.14}$ |
| $K^*\ell^+\ell^-$ | low | 27.2 ± 6.3 | 0.35 ± 0.16 | $+0.24_{-0.23}^{+0.18}$ |
| | high | 36.6 ± 9.6 | $0.71_{-0.22}^{+0.20}$ | $+0.76^{+0.52}_{-0.32}$ |

 $A_i \simeq C_i$ A_i are q^2 independent real term of C_i $C_7C_9C_{10}$ are real up to higher order corrections.

One need to use Wilson coefficient complex (not only for NP, but also for SM) Wou, Hovhannisyan, Mahajan PRD 77, 014016 (2008)

BaBar, PRD79, 031102 (R) (2009).

\mathcal{A}_{FB} in $B \to K^* \ell^+ \ell^-$

Belle updated the study and found the opposite sign again

 $C_7 = -C_7^{SM}$ seems to be favourable. If then data favoured 4th generation case.

Should be visible else where (sin2 ϕ_{B_S} negative)

LHCb, JHEP08, 131 (2013)

\mathcal{A}_{FB} in Sum of Exclusive

Inclusive measurement is theoretically cleaner than exclusive, but experimentally challenging

P'_5 anomaly Continuing with $B \to K^* \ell^+ \ell^-$

Angular analysis of $B \to \ell^+ \ell^- K^*(892)(\to K^-\pi^+).$

One can simply measure the angles between the direction of flight of all the different particles as function of q^2

Differential Angular distribution can be written as

$$
\frac{1}{d\Gamma/dq^2} \frac{d^4\Gamma}{d\cos\theta_\ell d\cos\theta_K d\phi dq^2} = \frac{9}{32\pi} \left[\frac{3}{4} (1 - F_L) \sin^2\theta_K + F_L \cos^2\theta_K + \frac{1}{4} (1 - F_L) \sin^2\theta_K \cos 2\theta_\ell - F_L \cos^2\theta_K \cos 2\theta_\ell + S_3 \sin^2\theta_K \sin^2\theta_\ell \cos 2\phi + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi + S_5 \sin 2\theta_K \sin \theta_\ell \cos \theta_\ell + S_7 \sin 2\theta_K \sin \phi + S_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + S_9 \sin^2\theta_K \sin^2\theta_\ell \sin 2\phi \right],
$$

 F_{L} is the longitudinal polarization fraction One has additional angular observables $S_n(n =$ 3,4,5,7,8,9) from the decay amplitude, which are functions of the Wilson coefficients and form factors Suggested by Decoste-Genon etal JHEP 2013, 137 (2013)

$$
P'_{4,5,6,8} = \frac{S_{4,5,6,8}}{\sqrt{F_L(1 - F_L)}}
$$

These observables are largely free from form factor uncertainties (especially at low q^2)

Each transformation preserves the first five terms and corresponding S_i term. Resulting angular distribution depend only on F_L , S_3 and one of the observables $S_{4,5,7,8}$

$$
P'_{4}, S_{4}: \begin{cases} \phi \to -\phi & \text{for } \phi < 0\\ \phi \to \pi - \phi & \text{for } \theta_{\ell} > \pi/2\\ \theta_{\ell} \to \pi - \theta_{\ell} & \text{for } \theta_{\ell} > \pi/2, \end{cases}
$$

$$
P'_{5}, S_{5}: \begin{cases} \phi \to -\phi & \text{for } \phi < 0\\ \theta_{\ell} \to \pi - \theta_{\ell} & \text{for } \theta_{\ell} > \pi/2, \end{cases}
$$

$$
P'_6, S_7: \begin{cases} \phi \to \pi - \phi & \text{for } \phi > \pi/2 \\ \phi \to -\pi - \phi & \text{for } \phi < -\pi/2 \\ \theta_\ell \to \pi - \theta_\ell & \text{for } \theta_\ell > \pi/2, \end{cases}
$$

$$
P'_{8}, S_{8}: \begin{cases} \phi \to \pi - \phi & \text{for } \phi > \pi/2 \\ \phi \to -\pi - \phi & \text{for } \phi < -\pi/2 \\ \theta_{K} \to \pi - \theta_{K} & \text{for } \theta_{\ell} > \pi/2 \\ \theta_{\ell} \to \pi - \theta_{\ell} & \text{for } \theta_{\ell} > \pi/2. \end{cases}
$$

First anomaly observed in P'_5

LHCb,PRL111, 191801 (2013)

$$
B \to K^*(892)\mu^+\mu^-
$$

$$
K^*(892) \to K^-\pi^+
$$

Local discrepancy of 3.7 σ is observed in the interval $4.30 < q^2 < 8.68$ GeV² for P'_5

Integrating over 1-6, is 2.5 σ

Belle data support similar trend Bernat *et al* JHEP 10,075 (2016)

 $B \to K^*(892)\ell^+\ell^ K^*(892) \to K^-\pi^+, K_s\pi^+, K^+\pi^0$

 ℓ for e and μ observable to test LFU, $Q_i = P_i^{\prime \mu} - P_i^{\prime e}$ Deviation from zero test of SM.

Belle,PRL 118, 111801 (2017)

First lepton-flavour-dependent measurement for P'_5 is reported

TABLE I. Fit results for P'_4 and P'_5 for all decay channels and separately, for the electron and muon modes. The first uncertainties are statistical and the second systematic.

| q^2 in GeV ² / c^2 | | | | | P^{μ}_{\leq} |
|-----------------------------------|--|----------|---|--|---------------------------------|
| [1.00, 6.00] | | | | $-0.45_{-0.22}^{+0.23} \pm 0.09 \ -0.72_{-0.39}^{+0.40} \pm 0.06 \ -0.22_{-0.34}^{+0.35} \pm 0.15 \ -0.23_{-0.22}^{+0.21} \pm 0.07 \ -0.22_{-0.41}^{+0.39} \pm 0.03 \ -0.43_{-0.28}^{+0.26} \pm 0.10$ | |
| [0.10, 4.00] | | | | $0.11_{-0.31}^{+0.32} \pm 0.05 \qquad 0.34_{-0.45}^{+0.41} \pm 0.11 \quad -0.38_{-0.48}^{+0.50} \pm 0.12 \qquad 0.47_{-0.28}^{+0.27} \pm 0.05 \qquad 0.51_{-0.46}^{+0.39} \pm 0.09 \qquad 0.42_{-0.39}^{+0.39} \pm 0.14$ | |
| [4.00, 8.00] | | | | $-0.34_{-0.17}^{+0.18} \pm 0.05$ $-0.52_{-0.22}^{+0.24} \pm 0.03$ $-0.07_{-0.31}^{+0.32} \pm 0.07$ $-0.30_{-0.19}^{+0.19} \pm 0.09$ $-0.52_{-0.26}^{+0.28} \pm 0.03$ $-0.03_{-0.30}^{+0.31} \pm 0.09$ | |
| [10.09, 12.90] | $-0.18_{-0.27}^{+0.28} \pm 0.06$ | \cdots | $-0.40^{+0.33}_{-0.29} \pm 0.09$ $-0.17^{+0.25}_{-0.25} \pm 0.01$ | . | $0.09_{-0.29}^{+0.29} \pm 0.02$ |
| | $[14.18, 19.00]$ $-0.14^{+0.26}_{-0.26} \pm 0.05$ $-0.15^{+0.41}_{-0.40} \pm 0.04$ $-0.10^{+0.39}_{-0.39} \pm 0.07$ $-0.51^{+0.24}_{-0.22} \pm 0.01$ $-0.91^{+0.36}_{-0.30} \pm 0.03$ $-0.13^{+0.39}_{-0.35} \pm 0.06$ | | | | |

TABLE II. Results for the lepton-flavor-universality-violating observables Q_4 and Q_5 . The first uncertainty is statistical and the second systematic.

| q^2 in GeV ² / c^2 | Q_4 | $\varrho_{\scriptscriptstyle{5}}$ |
|-----------------------------------|------------------------------|-----------------------------------|
| [1.00, 6.00] | $0.498 \pm 0.527 \pm 0.166$ | $0.656 \pm 0.485 \pm 0.103$ |
| [0.10, 4.00] | $-0.723 \pm 0.676 \pm 0.163$ | $-0.097 \pm 0.601 \pm 0.164$ |
| [4.00, 8.00] | $0.448 \pm 0.392 \pm 0.076$ | $0.498 \pm 0.410 \pm 0.095$ |
| [14.18, 19.00] | $0.041 \pm 0.565 \pm 0.082$ | $0.778 \pm 0.502 \pm 0.065$ |

First time

Results are compatible with SM within statistical uncertainty. Largest discrepancy is 2.6 σ in P'_5 for the muon channel

Results from ATLAS and CMS on P'_5

Differential branching fraction in $B \to K\mu\mu$, $B_s \to \phi\mu\mu$, $B \to K^*\mu\mu$, $\Lambda_b \to \Lambda\mu\mu$ consistently lower than SM prediction

JHEP06,133(2014), JHEP09,179(2015), JHEP06,115(2015)

$R_{K^*}R_K$ anomaly

Lepton flavour universality (LFU): electroweak couplings treat all flavours of leptons same *(observed difference are due to their mass differences)*

LFU can be tested very precisely

$$
R_K = \frac{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\Gamma[B^+ \to K^+ \mu^+ \mu^-]}{dq^2} dq^2}{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\Gamma[B^+ \to K^+ e^+ e^-]}{dq^2} dq^2} dq^2
$$

In ratio, hadronic uncertainties cancel and SM prediction is near unity.

$$
R_K^{SM} = \frac{\mathcal{B}(B^+ \to K^+ \mu^+ \mu^-)}{\mathcal{B}(B^+ \to K^+ e^+ e^-)} = 1 \pm \mathcal{O}(10^{-4})
$$

C. Bobeth et al, JHEP 07, 040 (2004) HPQCD, PRL111, 162002 (2013)

 R_H , $H = K$, K^* , X_s provides constraint to the New Physics G. Hiller, F. Kruger PRD69. 074020 (2004)

0.001 typical error

In 2014, LHCb measured and saw 2.6 hint of deviation from the SM.

Both BaBar and LHCb seems to favour a value below one.

Belle result didn't contradict q^2 dependent.

$$
R_K[1.0 < q^2 < 6.0 \ GeV^2] = 0.745^{+0.090}_{-0.074} \pm 0.036
$$

LHCb, PRL113, 151601(2014) BaBar, PRD86, 032012(2012) Belle, PRL103, 171801(2009)

 R_{K^*} anomaly

Trend similar to R_K

Belle result on R_{K^*}

 K^* from $K^+\pi^-$, $K^+\pi^0$, and $K^0_S\pi^+$

 $B^0 \to K^{*0} \mu^+ \mu^-$, $B^+ \to K^{*+} \mu^+ \mu^-$, $B^0 \to K^{*0} e^+ e^-$ and $B^+ \to K^{*+} e^+ e^-$

LHCb, PRL122, 191801 (2019)

 $B \to K^+ \ell \ell$

2019 LHCb R_K update

 $R_{K, \text{Run 1}}^{\text{new}} = 0.717_{-0.071}^{+0.083} \text{ (stat)}_{-0.016}^{+0.017} \text{ (syst)} ,$ $R_{K, \text{Run 2}} = 0.928_{-0.076}^{+0.089} \text{ (stat)} \pm_{-0.017}^{+0.020} \text{ (syst)}$.

 R_K [1.1 < q^2 < 6.0 GeV²] = 0.846^{+0.060+0.016}

Most precise measurement to data Consistent with SM at the level of 2.5 σ deviation

 $R_K[1.0 < q^2 < 6.0 \; GeV^2] = 0.745^{+0.090}_{-0.074} \pm 0.036$ LHCb, PRL113, 151601(2014)

 $B \to K \ell \ell$ R_K result from Belle Belle, arXiv:190801848

Result is consistent with LHCb as well as SM

More precision at Belle II

R_K result from Belle

⁺ to B^0

Sospin asymmetry
$$
A_{I} = \frac{\left(\frac{7}{(\tau_{B} + /\tau_{B^0})}\beta(B^0 \rightarrow K^0 \ell^+ \ell^-) - \beta(B^+ \rightarrow K^+ \ell^+ \ell^-)\right)}{(\tau_{B} + /\tau_{B^0})\beta(B^0 \rightarrow K^0 \ell^+ \ell^-) + \beta(B^+ \rightarrow K^+ \ell^+ \ell^-)}
$$

 A_I for all bins have negative asymmetry. For bin $1 < q^2 < 6$ GeV²/ c^4 deviates from zero by 2.7 σ for muon final state.

First time provided by Belle

Anomalies leading SM

or

https://www.mycustomer.com/experience/engagement/how-can-leaders-create-conditions-for-all-employees-to-deliver-outstanding-cx

SM leading anomalies

https://www.mckinsey.com/featured-insights/leadership/lead-at-your-best

https://images-na.ssl-images-amazon.com/images/I/51fvLhfOxVL.jpg

A**hulu** DOCUMENTARY **THE NEXT STEP**

AS TOLD BY MORGAN FREEMAN

Thank you

Given the above formula and input for $b_n^{0,+}$, the SM predicts $R_{\pi}^{\text{SM}} = 0.641 \pm 0.016$, whereas the experimental data suggests $R_{\pi}^{\text{exp.}} \simeq 1.05 \pm 0.51$ by using $\mathcal{B}(B \to \pi \ell \bar{\nu}_{\ell}) = (1.45 \pm 0.05) \times$ 10^{-4} [77]. Thus, at present the experimental result is consistent with the SM prediction $R_{\pi}^{5\,\text{ab}^{-1}} = 0.64 \pm 0.23$, $R_{\pi}^{50\,\mathrm{ab}^{-1}} = 0.64 \pm 0.09$.

B2TiP report

$$
R(D^*) = 0.270 \pm 0.035 \text{(stat)}^{+0.028}_{-0.025} \text{(syst)},
$$

$$
P_{\tau}(D^*) = -0.38 \pm 0.51 \text{(stat)}^{+0.21}_{-0.16} \text{(syst)},
$$

Prospects for data & physics harvesting

LHCb Ungrade II Physics Case

[LHCB-PUB-2018-009] arXiV:1808.08865

Table 10.1: Summary of prospects for future measurements of selected flavour observables for LHCb, Belle II and Phase-II ATLAS and CMS. The projected LHCb sensitivities take no account of potential detector improvements, apart from in the trigger. The Belle-II sensitivities are taken from Ref. [608].

B2TIP report | arXiv:1808.10567

Integrated Luminosity (ab⁻¹)

 $R(D^*)$

Fig. 72: Expected Belle II constraints on the R_D vs R_{D^*} plane (left) and the R_{D^*} vs $P_{\tau}(D^*)$ plane (right) compared to existing experimental constraints from Belle. The SM predictions are indicated by the black points with theoretical error bars. In the right panel, the NP scenarios "Scalar", "Vector" and "Tensor" assume contributions from the operators \mathcal{O}_{S_1} , \mathcal{O}_{V_1} and O_T , respectively.

 $B \to D^{(*)} \tau \bar{\nu}$:

 $B\to \pi\tau\bar\nu$:

 ${\rm Leptonic}$ B ${\rm Decays:}$

Semileptonic $b\to s$ Penguin Decays:

Semileptonic $b \to d$ Penguin Decays:

TABLE II. The systematic uncertainties in $R(D^*)$ and $P_{\tau}(D^*)$, where the values for $R(D^*)$ are relative errors. The group "common sources" identifies the common systematic uncertainty sources in the signal and the normalization modes, which cancel to a good extent in the ratio of these samples. The reason for the incomplete cancellation is described in the text.

| Source | $R(D^*)$ | $P_\tau(D^*)$ |
|--|---------------------|----------------------|
| Hadronic B composition | $+7.7%$ $-6.9%$ | $+0.134$ -0.103 |
| MC statistics for PDF shape | $+4.0%$ $-2.8%$ | $+0.146$ -0.108 |
| Fake D^* | 3.4% | 0.018 |
| $\bar{B} \to D^{**} \ell^- \bar{\nu}_{\ell}$ | 2.4% | 0.048 |
| $\bar{B} \to D^{**} \tau^- \bar{\nu}_\tau$ | 1.1% | 0.001 |
| $\bar{B} \to D^* \ell^- \bar{\nu}_e$ | 2.3% | 0.007 |
| τ daughter and ℓ^- efficiency | 1.9% | 0.019 |
| MC statistics for efficiency estimation | 1.0% | 0.019 |
| $\mathcal{B}(\tau^- \to \pi^- \nu_\tau, \rho^- \nu_\tau)$ | 0.3% | 0.002 |
| $P_{\tau}(D^*)$ correction function | 0.0% | 0.010 |
| Common sources | | |
| Tagging efficiency correction | 1.6% | 0.018 |
| D^* reconstruction | 1.4% | 0.006 |
| Branching fractions of the <i>D</i> meson | 0.8% | 0.007 |
| Number of $B\bar{B}$ and $\mathcal{B}(\Upsilon(4S) \rightarrow B^+B^-$ or $B^0\bar{B}^0$) | 0.5% | 0.006 |
| Total systematic uncertainty | $+10.4%$ $-9.4%$ | $+0.21$ -0.16 |

R(D)* systematics

LHCb one prong 2015

LHCb three prong 2017

Belle 2019 semileptonic

TABLE I. Systematic uncertainties contributing to the $\mathcal{R}(D^{(*)})$ results, together with their correlation.

Belle 2016 semileptonic

TABLE I. Summary of the systematic uncertainties on $\mathcal{R}(D^*)$ for electron and muon modes combined and separated. The uncertainties are relative and are given in percent.

| | $\mathcal{R}(D^*)$ (%) | | |
|---|---|------------------|------------------|
| Sources | $\ell^{sig} = e, \mu \ell^{sig} = e \ell^{sig} = \mu$ | | |
| MC size for each PDF shape | 2.2 | 2.5 | 3.9 |
| PDF shape of the normalization in $\cos \theta_{B-D^* \ell}$ | $+1.1$ -0.0 | $+2.1$ -0.0 | $+2.8$ -0.0 |
| PDF shape of $B \to D^{**} \ell \nu_{\ell}$ | $+1.0$ -1.7 | $+0.7$ -1.3 | $+2.2$ -3.3 |
| PDF shape and yields of fake $D^{(*)}$ | 1.4 | 1.6 | 1.6 |
| PDF shape and yields of | 1.1 | 1.2 | 1.1 |
| $B \to X_c D^*$ | | | |
| Reconstruction efficiency ratio $\varepsilon_{\text{norm}}/\varepsilon_{\text{sig}}$ | 1.2 | 1.5 | 1.9 |
| Modeling of semileptonic decay | 0.2 | 0.2 | 0.3 |
| $\mathcal{B}(\tau^- \to \ell^- \bar{\nu}_e \nu_\tau)$ | 0.2 | 0.2 | 0.2 |
| Total systematic uncertainty | $+3.4$ -3.5 | $+4.1$ -3.7 | $+5.9$ -5.8 |