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Prospects of Higgs pair production at the HL-LHC

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Outline of this talk

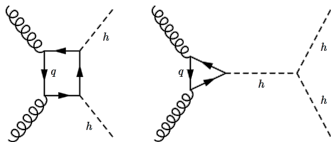
- HL-LHC prospects of non-resonant di-Higgs searches.
- A qualitative analysis of the sensitivity of di-Higgs search strategies on λ_{hhh} .
- Contamination to di-Higgs final states from BSM scenarios.

- The CMS and ATLAS collaborations have unambiguously confirmed the existence of a scalar boson at 125 GeV.
- Numerous studies performed to measure the Higgs coupling with SM particles.
- The Higgs self-coupling, λ_{hhh} still remains elusive.
- The only direct probe of λ_{hhh} in SM is offered by the **non-resonant Higgs pair production process**, where the final state Higgs eventually decays into SM particles.
- Precise measurement of λ_{hhh} is extremely challenging because of the **small** SM di-Higgs production cross-section.
- In the SM, the di-Higgs production proceeds through top-quark loop in the ggF mode.

Direct measurement is challenging :

Small SM production cross-section : $\sigma_{gghh_{SM}} = 39.56^{+7.32\%}_{-8.38\%}$ (at 14 TeV).^a
Other di-Higgs production modes furnish even smaller cross-sections ~ 1 fb

^aBorowka et. al., 2016



However, **phenomenologically rich final states can emerge**, offering a possibility to improve the discovery potential through highly optimized search strategies.

Current limits from non-resonant di-Higgs searches:

- ATLAS : 13.3 fb^{-1} : $4b$: ~ 29 times.
- ATLAS : 36.1 fb^{-1} : $4b$: ~ 13 times
- CMS : 35.9 fb^{-1} : $2b2\tau$: ~ 30 times.
- CMS : 35.9 fb^{-1} : $2b2\gamma$: ~ 19 times.

Strategy

- **Optimize the di-Higgs search strategy** using both, cut-based analysis and multivariate techniques, for a multitude of di-Higgs final states:

1. $b\bar{b}\gamma\gamma$

2. $b\bar{b}\tau_h\tau_h$

3. $b\bar{b}\tau_h\tau_l$

4. $b\bar{b}\tau_l\tau_l$

5. $b\bar{b}W^+W^- \rightarrow b\bar{b}l\nu jj$

6. $b\bar{b}W^+W^- \rightarrow b\bar{b}l\nu l\nu$

7. $W^+W^- \gamma\gamma \rightarrow l\nu jj\gamma\gamma$

8. $W^+W^- \gamma\gamma \rightarrow l\nu l\nu\gamma\gamma$

9. $4W \rightarrow ss2l$

10. $4W \rightarrow 3l$

11. $4W \rightarrow 4l$

The small signal rate is compensated by a very clean signature.

Backgrounds

- Dominant background : **QCD-QED $b\bar{b}\gamma\gamma$** .
Pure QED contribution is roughly $\sim O(1\%)$ of the QCD-QED counterpart.
- Associated production of Higgs with bottom and top quarks : $b\bar{b}h, t\bar{t}h$.
- Associated production of Higgs with Z-boson : Zh .
- Fake backgrounds : $c\bar{c}\gamma\gamma, jj\gamma\gamma, b\bar{b}jj, b\bar{b}j\gamma, c\bar{c}j\gamma$.

$$N^{c\bar{c}\gamma\gamma(jj\gamma\gamma)} = (N_{\text{ATLAS}}^{c\bar{c}\gamma\gamma(jj\gamma\gamma)} / N_{\text{ATLAS}}^{b\bar{b}\gamma\gamma}) \cdot N^{b\bar{b}\gamma\gamma}$$

$$N^{c\bar{c}j\gamma} = (N_{\text{ATLAS}}^{c\bar{c}j\gamma} / N_{\text{ATLAS}}^{b\bar{b}j\gamma}) \cdot N^{b\bar{b}j\gamma}$$

Event selection

- Exactly 2 b -tagged jets and exactly 2 γ .
- $p_T^{b_1} > 40$ GeV and $p_T^{b_2} > 30$ GeV.
- $|\eta_{b_1, b_2}| < 2.4$.
- $p_T^\gamma > 30$ GeV and $|\eta_\gamma| < 1.37, 1.52 < |\eta_\gamma| < 2.37$.
- Veto : Isolated leptons with $p_T > 25$ GeV and $|\eta| < 2.5$.

Selection cuts

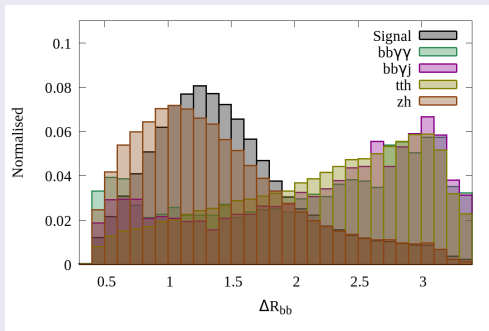
- $N_j < 6$
- $0.4 < \Delta R_{\gamma\gamma} < 2.0$, $0.4 < \Delta R_{bb} < 2.0$, $\Delta R_{\gamma b} > 0.4$
 - $p_{T,bb} > 80$ GeV, $p_{T,\gamma\gamma} > 80$ GeV
- 100 GeV $< m_{bb} < 150$ GeV, 122 GeV $< m_{\gamma\gamma} < 128$ GeV

Significantly reduces the $t\bar{t}h$ background, when either or both top quarks decay hadronically.

Selection cuts

- $N_j < 6$
- $0.4 < \Delta R_{\gamma\gamma} < 2.0, 0.4 < \Delta R_{bb} < 2.0, \Delta R_{\gamma b} > 0.4$
- $p_{T,bb} > 80 \text{ GeV}, p_{T,\gamma\gamma} > 80 \text{ GeV}$
- $100 \text{ GeV} < m_{bb} < 150 \text{ GeV}, 122 \text{ GeV} < m_{\gamma\gamma} < 128 \text{ GeV}$

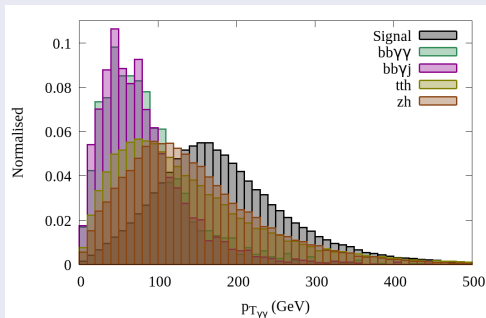
Highly effective in tackling the QED-QCD $b\bar{b}\gamma\gamma$ and fake backgrounds.



Selection cuts

- $N_j < 6$
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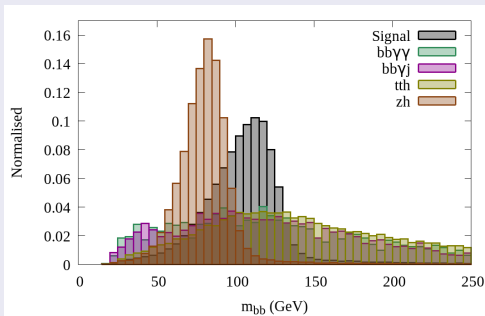
Effective in reducing the QED-QCD $b\bar{b}\gamma\gamma$, fake backgrounds and $t\bar{t}h$ backgrounds.



Selection cuts

- $N_j < 6$
- $0.4 < \Delta R_{\gamma\gamma} < 2.0$, $0.4 < \Delta R_{bb} < 2.0$, $\Delta R_{\gamma b} > 0.4$
 - $p_{T,bb} > 80$ GeV, $p_{T,\gamma\gamma} > 80$ GeV
- $100 \text{ GeV} < m_{bb} < 150 \text{ GeV}$, $122 \text{ GeV} < m_{\gamma\gamma} < 128 \text{ GeV}$

Eliminates a significant chunk of Zh and $b\bar{b}h$ background.



The $b\bar{b}\gamma\gamma$ channel

The cut flow table

Cut flow	Event rates with 3000 fb^{-1} of integrated luminosity							$\frac{S}{\sqrt{B}}$
	Signal $hh \rightarrow 2b2\gamma$	SM Backgrounds						
		$hb\bar{b}$	$t\bar{t}h$	Zh	$b\bar{b}\gamma\gamma^*$	Fake 1	Fake 2	
$2b + 2\gamma$	31.63	21.20	324.91	39.32	25890.31	1141.18	393.79	0.19
lepton veto	31.63	21.20	255.66	39.32	25889.94	1141.18	393.79	0.19
$N_j < 6$	31.04	21	192.05	39.23	25352.78	1064.64	167.32	0.19
ΔR cuts	22.19	7.75	38.71	23.48	4715.21	130.10	28.81	0.31
m_{bb}	12.71	1.53	13.80	1.09	862.37	22.11	6.88	0.42
$m_{\gamma\gamma}$	12.36	1.5	13.16	1.06	26.54	22.11	6.88	1.46
$p_{T,bb}, p_{T,\gamma\gamma}$	12.32	1.48	13.03	1.06	26.54	21.82	6.88	1.46

Signal significance : $\frac{S}{B} = 0.17$ and $\frac{S}{\sqrt{B}} = 1.46$.

⁰ $b\bar{b}\gamma\gamma^*$: $b\bar{b}\gamma\gamma + c\bar{c}\gamma\gamma + jj\gamma\gamma$

⁰ Fake 1: $b\bar{b}j\gamma + c\bar{c}j\gamma$, Fake 2: $b\bar{b}jj$.

The $b\bar{b}\gamma\gamma$ channel

Implications from \cancel{E}_T cut and modified range of $m_{b\bar{b}}$

- The signal lacks any \cancel{E}_T source.
- Leptonically decaying W -bosons in $t\bar{t}h$ can result in \cancel{E}_T .
- Jet energy corrections can result in a shift in the $m_{b\bar{b}}$ distributions.

Applying, $\cancel{E}_T < 50$ GeV

	Process	Events
Background	$hb\bar{b}$	1.31
	$t\bar{t}h$	7.87
	Zh	1.03
	$b\bar{b}\gamma\gamma^*$	23.18
	Fake 1	20.69
	Fake 2	6.52
	Total	60.60
Signal ($hh \rightarrow 2b2\gamma$)		11.75
Significance (S/\sqrt{B})		1.51

Applying, $m_{bb} : 90\text{GeV} < m_{bb} < 130$ GeV

	Process	Events
Background	$hb\bar{b}$	1.55
	$t\bar{t}h$	11.91
	Zh	4.43
	$b\bar{b}\gamma\gamma^*$	28.41
	Fake 1	22.39
	Fake 2	7.25
	Total	75.94
Signal ($hh \rightarrow 2b2\gamma$)		14.27
Significance (S/\sqrt{B})		1.64

The $b\bar{b}\gamma\gamma$ channel

A multivariate analysis is also performed by utilising the Boosted Decision Tree (BDT) algorithm.

Variables used

$$m_{bb}, p_{T,\gamma\gamma}, \Delta R_{\gamma\gamma}, p_{T,bb}, \Delta R_{b_1\gamma_1}, p_{T,\gamma_1}, \Delta R_{bb}, \\ p_{T,\gamma_2}, \Delta R_{b_2\gamma_1}, \Delta R_{b_2\gamma_2}, p_{T,b_1}, \Delta R_{b_1\gamma_2}, p_{T,b_2}, \cancel{p_{T,T}}$$

The most effective variables in isolating the signal and backgrounds:

$$m_{bb}, p_{T,\gamma\gamma}, \Delta R_{b_1\gamma_1}, \Delta R_{bb}$$

Features a $\sim 20\%$ improvement in signal significance.

Sl. No.	Process	Events
Background	$hb\bar{b}$	2.75
	$t\bar{t}h$	14.85
	Zh	12.28
	$b\bar{b}\gamma\gamma^*$	34.46
	Fake 1	14.25
	Fake 2	8.46
	Total	87.05
	Signal ($hh \rightarrow 2b2\gamma$)	16.46
	Significance (S/\sqrt{B})	1.76

The $b\bar{b}\tau^+\tau^-$ channel

The τ leptons can decay either leptonically (34%) and hadronically.

Three final states : $b\bar{b}\tau_l\tau_l$, $b\bar{b}\tau_l\tau_h$, $b\bar{b}\tau_h\tau_h$.

Background

- Dominant background : $t\bar{t}$

Contributions arise from hadronic, leptonic and semi-leptonic decay modes of $t\bar{t}$.

- $b\bar{b}Z^*/\gamma^* \rightarrow b\bar{b}\tau^+\tau^-$.

- $b\bar{b}h$.

- Zh .

- Associated production of bosons with $t\bar{t}$ pair : $t\bar{t}h$, $t\bar{t}W$, $t\bar{t}Z$

- Fake backgrounds : $b\bar{b}jj$.

- Other potential sources can be $W(\rightarrow \tau\nu) + \text{jets}$, Wh , WZ , $h \rightarrow ZZ^*$, however, are neglected because of small rates.

The $b\bar{b}\tau^+\tau^-$ channel

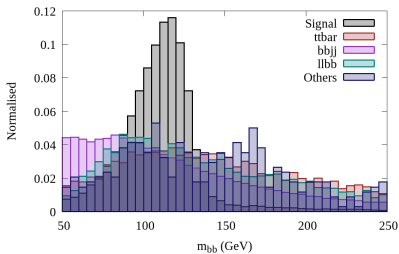
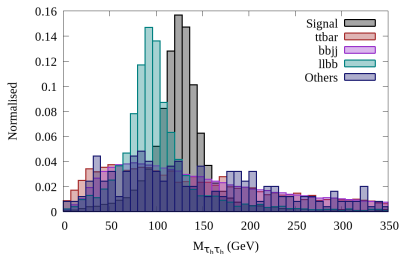
- The signal is marred by an extremely large $t\bar{t}$ background.
- Cut-based analysis, motivated from a similar search by ATLAS, results in a signal significance of ~ 0.44 , 0.26 and 0.044 , for $b\bar{b}\tau_h\tau_h$, $b\bar{b}\tau_l\tau_h$ and $b\bar{b}\tau_l\tau_l$, respectively.
- Imposing an additional cut on $M_{\tau\tau} < 200$ GeV improves S/\sqrt{B} to 0.65 , 0.44 and 0.07 .
- Multivariate analysis performed using the variables :

$\tau_h\tau_h$: $p_{T,bb}$, m_{bb} , ΔR_{bb} , $M_{\tau_h\tau_h}$, m_{T2} , $\Delta\phi_{\tau_{h1}}\cancel{E}_T$, m_{hh}^{vis} , $p_{T,hh}^{\text{vis}}$, $\Delta R_{hh}^{\text{vis}}$,

$\tau_h\tau_l$: $p_{T,bb}$, m_{bb} , ΔR_{bb} , $M_{\tau_h\tau_l}$, m_{T2} , $\Delta\phi_{\tau_h}\cancel{E}_T$, $\Delta\phi_{\tau_l}\cancel{E}_T$, m_{hh}^{vis} , $\Delta R_{hh}^{\text{vis}}$

$\tau_l\tau_l$: $p_{T,bb}$, m_{bb} , ΔR_{bb} , $M_{\tau_l\tau_l}$, m_{T2} , $\Delta\phi_{\tau_{l1}}\cancel{E}_T$, $\Delta\phi_{\tau_{l2}}\cancel{E}_T$, m_{hh}^{vis}

results in a signal significance of 0.74 , 0.49 and 0.077 , respectively.



The $b\bar{b}W^+W^-$ channel

Three possible final states :

- Fully leptonic : $b\bar{b}ll + \cancel{E}_T$. **Cleaner signature and lesser background.**
- Semi leptonic : $b\bar{b}l + jets + \cancel{E}_T$.
- Hadronic : $b\bar{b} + jets..$

Backgrounds

- Dominant background : $t\bar{t}$
Fully leptonic $t\bar{t}$ contributes as dominant background to $b\bar{b}ll + \cancel{E}_T$.
Leptonic and semi-leptonic contributes as dominant background to $b\bar{b}l + jets + \cancel{E}_T$.
- Second most dominant background for the semi-leptonic case : $Wb\bar{b} + jets$
- Other contributing backgrounds : $llb\bar{b}, b\bar{b}h, t\bar{t}h, t\bar{t}V, Vh, Vb\bar{b}$ and VVV

The $b\bar{b}W^+W^-$ channel

The $2b2l + \cancel{E}_T$ channel :

Multivariate analysis is performed using BDT algorithm.

Variables used :

$$p_{T,\ell_{1/2}}, \cancel{E}_T, m_{\ell\ell}, m_{bb}, \Delta R_{\ell\ell}, \Delta R_{bb}, p_{T,bb}, p_{T,\ell\ell}, \Delta\phi_{bb\ell\ell},$$

- The BDT algorithm is trained by using the signal sample and the $t\bar{t}$ background only.
- Best discriminatory variables : m_{bb} , $m_{\ell\ell}$, $p_{T,bb}$ and $p_{T,\ell\ell}$

Sl. No.	Process	Events
Background	$t\bar{t}$ lep	2080.52
	$t\bar{t}h$	131.66
	$t\bar{t}Z$	106.31
	$t\bar{t}W$	35.97
	$hb\bar{b}$	~ 0
	$\ell\ell b\bar{b}$	842.72
	Total	3197.18
	Signal ($hh \rightarrow b\bar{b}WW \rightarrow b\bar{b}\ell\ell + \cancel{E}_T$)	35.20
	Significance (S/\sqrt{B})	0.62

The $b\bar{b}W^+W^-$ channel

The $1/2j2b + \cancel{E}_T$ channel :

- Dominant background : Semi-leptonic and leptonic decays of $t\bar{t}$
- Input variables used :

$$p_{T,\ell}, \cancel{E}_T, m_{jj}, m_{bb}, \Delta R_{jj}, \Delta R_{bb}, p_{T,bb}, p_{T,\ell jj}, \Delta\phi_{bb\ell jj}, \Delta R_{\ell jj},$$

where $p_{T,\ell jj}$ refer to the p_T of the ℓjj system (Eg., for the signal, ensuing from the $h \rightarrow WW^* \rightarrow \ell\nu jj$ decay).

- Background training is performed using $t\bar{t}$.
- The best discriminatory observables : m_{bb} , p_{T,ℓ_1} , $p_{T,bb}$ and \cancel{E}_T .

Sl. No.	Process	Events
Background	$t\bar{t}$ semi-lep	866990.56
	$t\bar{t}$ lep	96147.82
	$t\bar{t}h$	4508.25
	$t\bar{t}Z$	5192.52
	$t\bar{t}W$	2949.65
	$Wb\bar{b}$ + jets [LO]	121313.52
	$\ell b\bar{b}$	5780.47
	Total	1102882.79
	Signal ($hh \rightarrow 2b2W$)	134.34
	Significance (S/\sqrt{B})	0.13

The $W^+W^-\gamma\gamma$ channel

We analyse two different final states : • Fully leptonic : $l^+l^-\gamma\gamma + \cancel{E}_T$. • Semi-leptonic : $lj\gamma\gamma + \cancel{E}_T$.

Features smaller signal yield, however, has lesser background.

The pure leptonic decay :

Backgrounds : • $t\bar{t}h$, • Zh , • $ll\gamma\gamma$.

- Variables for multivariate analysis :

$$p_{T,\ell(1,2)}, \cancel{E}_T, m_{\ell\ell}, m_{\gamma\gamma}, \Delta R_{\gamma\gamma(\ell\ell)}, p_{T,\ell\ell}, p_{T,\gamma\gamma}, \Delta\phi_{\ell\ell\gamma\gamma}, \quad (1)$$

- Best discriminating variables : $m_{\ell\ell}, \cancel{E}_T, p_{T,\gamma\gamma}, m_{\gamma\gamma}$

Sl. No.	Process	Events
Background	$t\bar{t}h$	0.89
	$Zh + \text{jets}$	0.20
	$ll\gamma\gamma + \text{jets}$	0.33
Total		1.42
Signal		0.57

- Features one of the best $\frac{S}{B}$ values ~ 0.40 .
Has the potential to become an important channel for HE-LHC owing to higher $\frac{S}{B}$.

The $W^+W^-\gamma\gamma$ channel

The semi-leptonic case :

Backgrounds

- Irreducible background : $l\nu\gamma\gamma$ ($\sigma = 3.28$ fb)
- Mis-identified background : $ll\gamma\gamma$ ($\sigma = 1.05$ fb)
 - Both, generated with $120 \text{ GeV} < m_{\gamma\gamma} < 130 \text{ GeV}$
- Other important backgrounds: $t\bar{t}h$, Zh +jets, Wh

Variables used : p_{T,ℓ_1} , $E_{T,\nu}$, $m_{\gamma\gamma}$, $\Delta R_{\gamma\gamma}$, $p_{T,\gamma\gamma}$, p_{T,ℓ_j} , $\Delta\phi_{\ell_j\gamma\gamma}$, ΔR_{ℓ_j} , m_T

Most discriminating variables :

ΔR_{ℓ_j} , $p_{T,\gamma\gamma}$, $m_{\gamma\gamma}$ and m_T

Sl. No.	Process	Events
Background	$t\bar{t}h$	6.49
	Zh + jets	1.71
	Wh + jets	5.13
	$l\nu\gamma\gamma$ + jets	2.57
	$ll\gamma\gamma$ + jets	1.07
	Total	16.97
Signal		1.85

The 4W channel

We choose these final states :

- Same-sign di-leptons (SS2l): $e^\pm l^\pm + 4j + \cancel{E}_T$, • Tri-leptons (3l) : $3l + 2j + \cancel{E}_T$,
- Four leptons (4l) : $4l + \cancel{E}_T$.

Backgrounds

- SS2l :
 - Dominant backgrounds : $WZ(W \rightarrow l\nu, Z \rightarrow ll)$ and same-sign W -pair.
 - Other important ones : $Vh, t\bar{t}X$ ($X = W^\pm, Z, h$), • Fakes : $t\bar{t}$.
- 3l :
 - Dominant backgrounds : Wh, WZ , • Fake backgrounds : $t\bar{t}$
 - Others : Zh ($Z \rightarrow ll, h \rightarrow W^+W^-$), $t\bar{t}X$ ($X = W^\pm, Z, h$) and ZZ
- 4l :
 - Fake backgrounds : $t\bar{t}$ • Other important backgrounds : $Wh, t\bar{t}h, t\bar{t}, ZZ$ and Zh

Variables used for multivariate analysis :

- SS2l: $m_{e^\pm l^\pm}, \Delta R_{l_{ijk}}, m_{jj}$.
- 3l: $m_{e_l l_j}, \Delta R_{l_i l_j}, m_{lll}, m_{\text{eff}}, \cancel{E}_T, p_{T,l_i}, n_{\text{jet}}$.

The 4W channel

We choose these final states :

- Same-sign di-leptons (SS2l): $l^\pm l^\pm + 4j + \cancel{E}_T$, • Tri-leptons (3l) : $3l + 2j + \cancel{E}_T$,
- Four leptons (4l) : $4l + \cancel{E}_T$.

Backgrounds

- SS2l :
 - Dominant backgrounds : $WZ(W \rightarrow \nu\bar{\nu}, Z \rightarrow \ell\bar{\ell})$ and same-sign W -pair.
 - Other important ones : $Vh, t\bar{t}X$ ($X = W^\pm, Z, h$), • Fakes : $t\bar{t}$.
- 3l :
 - Dominant backgrounds : Wh, WZ , • Fake backgrounds : $t\bar{t}$
 - Others : Zh ($Z \rightarrow \ell\bar{\ell}, h \rightarrow W^+W^-$), $t\bar{t}X$ ($X = W^\pm, Z, h$) and ZZ
- 4l :
 - Fake backgrounds : $t\bar{t}$ • Other important backgrounds : $Wh, t\bar{t}h, t\bar{t}, ZZ$ and Zh

SS2l :

		Events
Background	Total	12366.53
Signal		11.96
Significance (S/\sqrt{B})		0.11

3l :

		Events
Background	Total	5374.45
Signal		15.01
Significance (S/\sqrt{B})		0.20

4l :

		\cancel{E}_T cut > 50 GeV
Background	5876.22	131.51
Signal	2.02	1.42

- Combination of results from the previous channels yield a combined significance of 2.1σ .
- Small di-Higgs production rate makes the prospect of discovering di-Higgs at the HL-LHC quite weak.
- The $b\bar{b}\gamma\gamma$ channel offers the highest signal significance.
- The $WW\gamma\gamma$ channel has the potential to play an important role in di-Higgs searches at higher luminosities or HE-LHC, owing to its excellent S/B value.

- The HL-LHC study by ATLAS^a predicts a sensitivity of $-0.8 < \lambda_{hhh}/\lambda_{SM} < 7.7$.
- Change in λ_{hhh} can alter the di-Higgs production rate as well as the kinematics.

^aATL-PHYS-PUB-2017-001

5 different values of $\lambda_{hhh}/\lambda_{SM}$ is considered : $-1, 1, 2, 5, 7$, and we do the following:

Signal production is done with the new λ_{hhh} values, and

- Case A : Passed through the cut-based analysis optimized for SM case ($\lambda_{hhh}/\lambda_{SM} = 1$).
- Case B : Passed through the BDT analysis optimized for SM case ($\lambda_{hhh}/\lambda_{SM} = 1$).
- Case C : The BDT framework is trained with an alternative $\lambda_{hhh}/\lambda_{SM} = 5$
- Case D : Training is performed for each specific value of λ_{hhh} .

Ramifications of varying λ_{hhh}

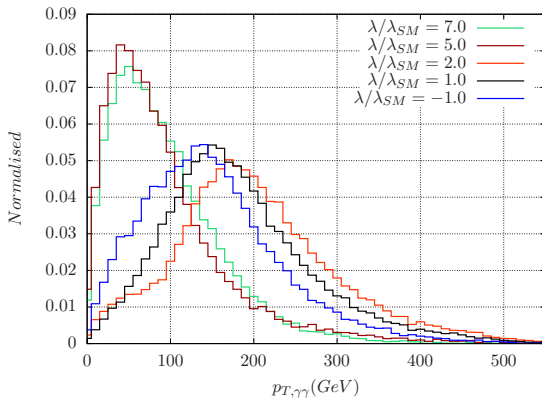
Case A : Passed through the cut-based analysis optimized for SM case ($\lambda_{hhh}/\lambda_{SM} = 1$)

Cut Based (optimised for $\lambda_{hhh}/\lambda_{SM} = 1$)					
λ/λ_{SM}	Sigma xs (fb)	Efficiency	Signal yield	Background yield	S/\sqrt{B}
-1	0.40	0.027	32.40	70.81	3.85
1	0.105	0.039	12.28		1.46
2	0.05	0.046	6.90		0.82
5	0.26	0.008	6.24		0.74
7	0.70	0.010	21.00		2.49

Case B : Passed through the BDT analysis optimized for SM case ($\lambda_{hhh}/\lambda_{SM} = 1$).

BDT (optimised for $\lambda_{hhh}/\lambda_{SM} = 1$)					
λ/λ_{SM}	Signal xs (fb)	Efficiency	Signal yield	Background yield	S/\sqrt{B}
-1	0.40	0.035	41.76	87.05	4.48
1	0.105	0.052	16.46		1.76
2	0.05	0.063	9.42		1.01
5	0.26	0.010	7.84		0.84
7	0.70	0.011	23.10		2.48

Difference in kinematic distributions cropping from variation in λ_{hhh} :



Case D

Training the BDT with corresponding λ_{hhh} samples \rightarrow

- the training becomes more **tuned to the modified kinematic distributions**.

Marks an improvement in the signal significance.

BDT (optimised for each λ_{hhh})					
λ/λ_{SM}	Signal xs (fb)	Efficiency	Signal yield	Background yield	S/\sqrt{B}
-1	0.40	0.049	58.80	166.13	4.55
1	0.105	0.052	16.46	87.05	1.76
2	0.05	0.068	10.20	85.54	1.10
5	0.26	0.046	35.88	455.51	1.69
7	0.70	0.049	102.90	466.97	4.76

- Changes accounting from variations in λ_{hhh} requires a careful analysis.
- Using the kinematic distributions to segregate the modifications emerging from changes in λ_{hhh} would be essential in searches for new physics.**

Contaminations to non-resonant di-Higgs pair production.

- The SM signal is susceptible to contamination or distortion from small new physics effects.
- Statistically significant deviations may be considered as New Physics signatures.
- Deviations can also be solely attributed to **modifications in λ_{hhh}** .

The BDT was trained to maximize the SM di-Higgs yield → **Do the NP signatures mimic the SM signal and render the BDT optimisation insensitive?**

Contaminations can arise if :

- the discriminatory kinematics variables used to train the BDT, have a **significant overlap with the SM counterparts**.
- **large new physics cross-section**, with a smaller kinematic overlap.

We study some such imposters ensuing from various well-motivated new physics scenarios through benchmark studies.

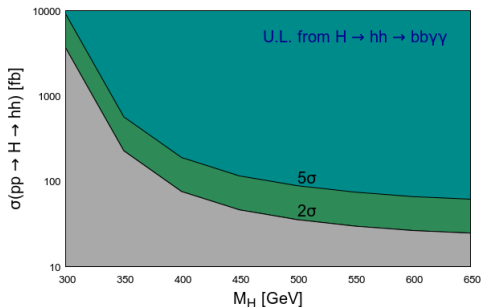
A typical example would be extensions of SM, with an extended Higgs sector. Resonant heavy Higgs production followed by decay into SM-like Higgs pair can enhance the di-Higgs cross-section.

Three possible scenarios which can contaminate :

- Double Higgs production : $pp \rightarrow hh(+X)$
- Single Higgs production : $pp \rightarrow h + X$
- Null Higgs scenario : $pp \rightarrow X$

We derive upper limits on $\sigma(pp \rightarrow H \rightarrow hh)$ as a function of M_H , in a **model independent fashion**, using the inequality $S_{NP}^{UL}/\sqrt{B_{SM}} \geq N\sigma$.

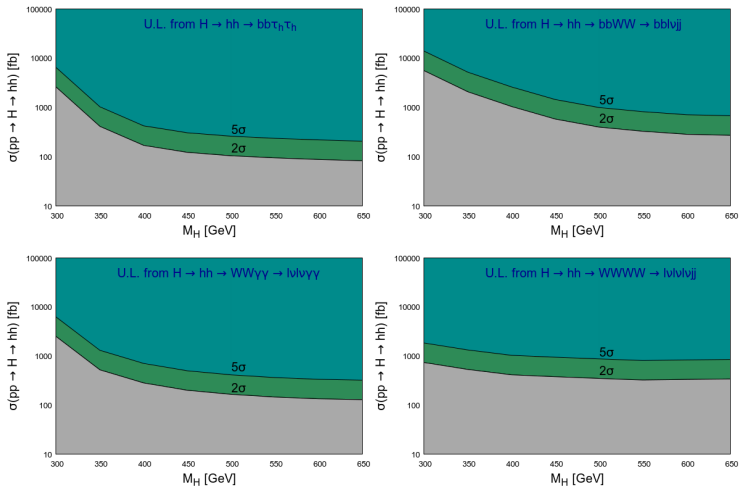
- B_{SM} includes the SM backgrounds as well as the SM di-Higgs contribution.
- The BDT optimization derived for non-resonant SM di-Higgs has been used to quantify the amount of contamination.



- The grey region do not contaminate the SM expectation.
- The blue and green region contaminate at 5 σ and 2 σ .

The $hh(+X)$ channel

Upper limits corresponding to all other non-resonant di-Higgs final states considered earlier, were obtained.



The $hh(+X)$ channel

Generic features :

- **Strongest upper limits are obtained for $b\bar{b}\gamma\gamma$.** Around ~ 25 fb at $M_H = 650$ GeV.
- $b\bar{b}\tau\tau$ offers the second strongest bound. Cross-section varies between 170 – 80 fb.
- $b\bar{b}WW^*$, varying between 228 fb and 40 fb for m_H varying between 450 GeV and 650 GeV.
- The 2σ upper limit plateaus between 129 fb and 282 fb for the fully leptonic case $WW\gamma\gamma$.
- Bounds from the $4W$ modes are much weaker.

The $h + X$ channels

Resonant production of pseudoscalar Higgs boson, followed by its decay into $A \rightarrow Zh$, can imitate a majority of di-Higgs final states.

Within MSSM, $A \rightarrow Zh$ can be appreciable below $t\bar{t}$ threshold.

We derive upper limits on $\sigma(pp \rightarrow A \rightarrow Zh)$ as a function of M_A .

- **Weaker limits** than those from resonant scalar production.
- **$b\bar{b}\tau\tau$ yields the strongest upper limit** : 292 fb to 186 fb.
- Upper limits from $b\bar{b}\gamma\gamma$ vary in between : ~ 330 fb to ~ 197 fb (at 650 GeV).
- Upper limits from di-leptonic $b\bar{b}WW$ strengthens from 1236 fb to 110 fb (at 650 GeV).
- $m_{b\bar{b}}$ serves as an important discriminator between the Zh and hh signals, thereby, weakening the contamination from $A \rightarrow Zh$.
- Hence, more cross-section required to contaminate the SM signal to the same degree as in the $H \rightarrow hh$ channel.

The $h + X$ channels

SUSY contamination to $h + X$

- Bounds on squarks and gluino have already surpassed a TeV.
- The electroweakino sector still remains as the plausible search area.
- We choose a benchmark point marginally outside the projected exclusion obtained by ATLAS for HL-LHC.
- Pair production of wino dominated $\chi_2^0 \chi_1^\pm$ is considered (cross-section = 420 fb).

Parameters (GeV)	Mass (GeV)	Processes	Branching Fraction
$M_1 = 150, M_2 = 300$ $\mu = 1000, m_{\tilde{t}_R} = 3000$	$m_{\chi_2^0} = 296.7$ $m_{\chi_1^\pm} = 296.7$ $m_{\chi_1^0} = 149.3$ $m_h = 125.0$ $m_{H^\pm} = 1003.0$ $m_H = 1000.0$	$\chi_1^\pm \rightarrow \chi_1^0 W^\pm$ $\chi_2^0 \rightarrow \chi_1^0 h$	100% 93.5%

The $h + X$ channels

The $Wh + \cancel{E}_T$ final state has the potential to contaminate some of the di-Higgs search channels :

$$b\bar{b}WW^* \rightarrow b\bar{b}ljj + \cancel{E}_T, \gamma\gamma WW^* \rightarrow \gamma\gamma ljj + \cancel{E}_T, 4W \rightarrow l^\pm l^\pm jjjj + \cancel{E}_T, 3ljj + \cancel{E}_T$$

Channel	SM background	SM hh production	BP2 contamination
$b\bar{b}ljj + \cancel{E}_T$	1103017.13	134.34	382.88
$SS2ljj + \cancel{E}_T$	12378.49	11.96	270.31
$3ljj + \cancel{E}_T$	5389.46	15.01	291.91

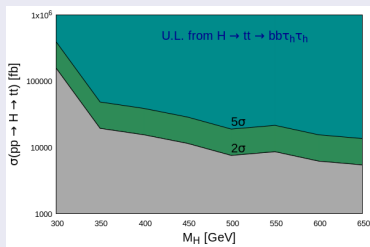
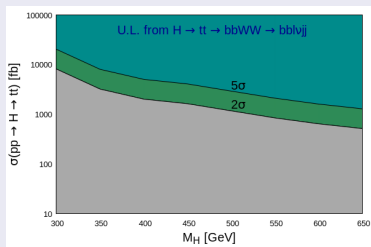
- SM di-Higgs expectations for these channels are relatively small as compared to the SUSY contamination.
- Observation of significant number of events in these channels over and above SM backgrounds can be a potential signature of New Physics.

The null Higgs channels

This category of contaminant has no Higgs in the production or decay mode.

$pp \rightarrow H(A) \rightarrow t\bar{t}$ can potentially contaminate the $b\bar{b}\tau\tau$ and $b\bar{b}WW$ final states.

- We derive upper limits on $\sigma(pp \rightarrow H(A) \rightarrow t\bar{t})$ for both final states from $b\bar{b}WW$, and $b\bar{b}\tau\tau$.
- Relatively weaker final states are obtained $\rightarrow \sigma \times Br$ values reaching to pb values.



- Lower contamination is ensured by the BDT variable $m_{b\bar{b}}$, owing to difference in kinematics of the final state $b\bar{b}$ pair.
- Large production cross-section of heavier Higgs will be required in order to contaminate the SM di-Higgs signal.

Null Higgs channel

SUSY contamination to null Higgs channel.

We consider stop pair production : $pp \rightarrow \tilde{t}_1 \tilde{t}_1^*$, using the following benchmark :

Benchmark Points	Parameters (GeV)	Mass (GeV)	Processes	Branching Fraction
3rd benchmark BP3 $pp \rightarrow \tilde{t}_1 \tilde{t}_1^*$ (Cross-section: 200 fb)	$M_1 = 500, M_2 = 1000$ $\mu = 1000, m_{\tilde{t}_R} = 625$	$m_{\tilde{t}_1} = 609.3$ $m_{\chi_1^0} = 498.1$ $m_h = 125.0$ $m_{H^\pm} = 1003.0$ $m_H = 1000.0$	$\tilde{t}_1 \rightarrow \chi_1^0 b W^+$	99.9%

- The benchmark point is chosen such that mass difference between \tilde{t}_1 and χ_1^0 is less than the top mass, ensuring the stop decays as $\tilde{t}_1 \rightarrow Wb\chi_1^0$.
- It can potentially contaminate $pp \rightarrow b\bar{b}WW$ and $pp \rightarrow b\bar{b}\tau\tau$ final states.
Contamination from stop pair production in the $b\bar{b}WW \rightarrow b\bar{b}lj\bar{j} + \cancel{E}_T$ final state comes out to be of same order as the SM signal.

SM background	SM hh production	BP3 contamination
1103017.13	134.34	101.83

- A clear understanding of the Higgs self-coupling will require a precise analysis of the di-Higgs channel.
- Combination of results from multiple search channels yield a signal significance of 2.1σ .
- The effects of varying λ_{hhh} was analyzed, and it was observed that varying λ_{hhh} can significantly alter our conclusions.
- New Physics scenarios can contaminate the SM di-higgs signals. Contaminations can creep in from $H \rightarrow hh$, $A \rightarrow Zh$, $H \rightarrow t\bar{t}$, $H^\pm \rightarrow t\bar{b}$, $\tau\nu_\tau$.
- Contaminations from benchmark scenarios of squark pair production, electroweakino pair production and stop pair production were quantified. In most cases, the contaminations can pose a serious threat to the search analyses.

- A detailed study of the less explored di-higgs final states in the context of HE-LHC. We are currently looking into the future prospect of $WW\gamma\gamma$ channel at HE-LHC.
- Moreover, a global combination of various di-Higgs signals at HE-LHC could provide interesting results concerning the projected capability of HE-LHC in probing λ .
- A detailed qualitative understanding of the alterations in the distribution of various kinematic variables emerging from different values of Higgs tri-linear coupling and aim to improve upon the robustness of the existing search strategies.
- Compilation of various allowed BSM benchmark scenarios which could possibly contaminate the di-Higgs final states.

Thank you.