

Universal Physics of Nuclear Few-body systems in Pionless Effective Theory

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Outline

□ Introduction & Motivation:

- Study of Few-Body Systems of light nuclei
- Search and Prediction of exotic bound states using Universality
↪ e.g., Exotic bound states of n -rich **Hypernuclei**

□ Universal physics in 3-body clusters:

- Two-body sector
- Three-body sector

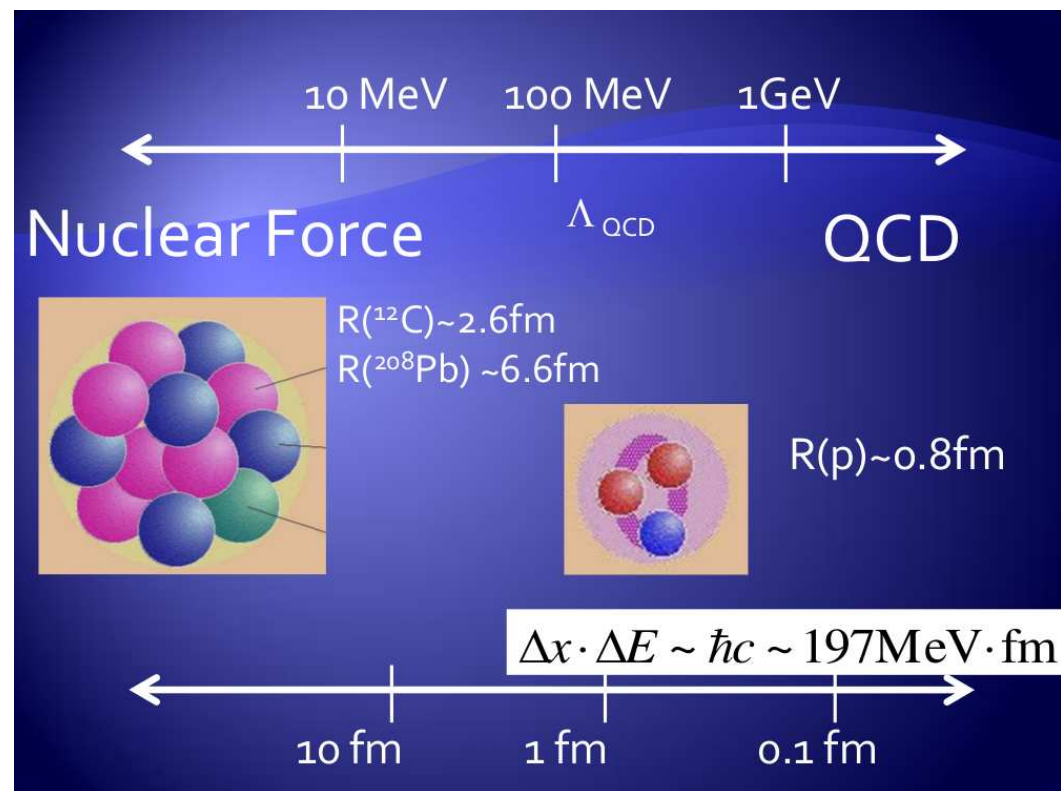
□ Framework: Low-energy Effective Field Theory

↪ **TOY MODEL:** System of 3 identical interacting bosons

□ Few Results: Feasibility of Efimov-like bound states in $nn\Lambda$

QCD and Strong Nuclear Force

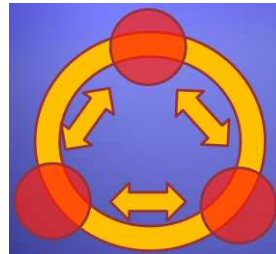
- ❑ **ASYMPTOTIC FREEDOM:** QCD becomes highly non-perturbative below $\lesssim 1 \text{ GeV}$
- ❑ How strong force binds nucleons and other baryons within nuclei?
- ❑ Focus on the Low-energy or long distance scales $\gtrsim 1 \text{ fm}$



Effective QCD & dynamics of light nuclei

👉 Analytical Approaches?

- **QCD** → True theory of underlying structure & dynamics of elementary hadrons
- Straightforward field theoretical extensions are not viable for systems with $A \geq 3$
QM Few-body Problem !
- **Classical 3-body Problem?** Very difficult to solve analytically
Classical ≥ 4 -body Problem? No general analytical solns. with **non-linear** ints.



- No general analytical solution of **non-linear** multi-hadron dynamics in QCD

👉 Approximate / Semi-analytic / Numerical Methodologies:

- **Mean-field & Variational approaches** → statistically large $A \rightarrow \infty$ (continuum matter)
- **Potential Models** → small A ; not very systematic, i.e., lack **universality**
- **Cluster EFTs** → $A \gtrsim 2,3,4, \dots$; simple yet powerful, systematic, model independent & **universal**

Exotics in Nuclear Spectroscopy

☞ Color-singlet hadrons:

- **Stable** hadronic bound states: *Mesons* ($q\bar{q}$) and *Baryons* (qqq)
- QCD does not restrict other **exotic metastable color-singlet** configurations

☞ Compact multi-quark / -gluon Exotics → Quark Model picture

Tetraquarks ($qq\bar{q}\bar{q}$), *Pentaquarks* ($qqqq\bar{q}$), *Dibaryons* ($qqqqqq$), ...

Quarkonia ($q\bar{q} - Q\bar{Q}$, $q\bar{Q} - Q\bar{q}$, $qQ - \bar{q}\bar{Q}$, ...) ; $Q = c, b, t$

Glueballs (gg, ggg, \dots), *Hybrids* ($Q\bar{Q}g, \dots$)

☞ Meson Exotics → "Molecular" nature to contrast with *tetraquark* structure

$f_0(500)$ or " σ " ($\pi\pi$), $K(800)$ or " κ " (πK), $f_0(980)$ & $a_0(980)$ ($K\bar{K}$), ...

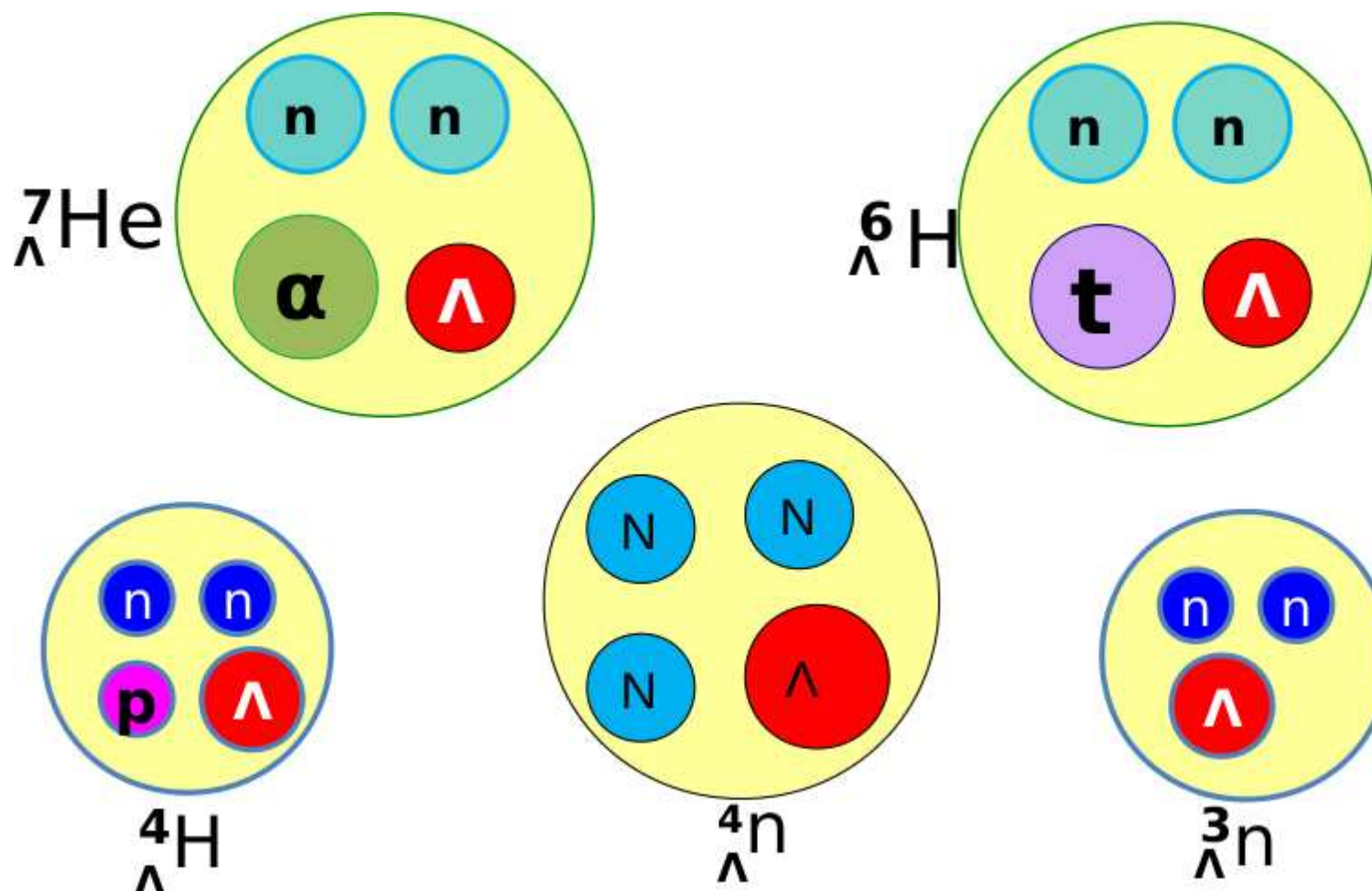
☞ Exotic Nuclei → "Halo" Clusters of "Molecular" nature

Neutron-rich states: *Tetra-neutron* (${}^4n \rightarrow nnnn$), ...

Hypernuclear states: *Hypertriton* (${}^3_{\Lambda}H \rightarrow \Lambda d$), ...

Recently suggested n -rich Λ -Hypernuclei

☞ **Single- Λ systems:** ${}^6_{\Lambda}\text{H}$, ${}^4_{\Lambda}\text{H}$, ${}^7_{\Lambda}\text{He}$, ${}^4_{\Lambda}\text{n}$, ${}^3_{\Lambda}\text{n}$, ...??

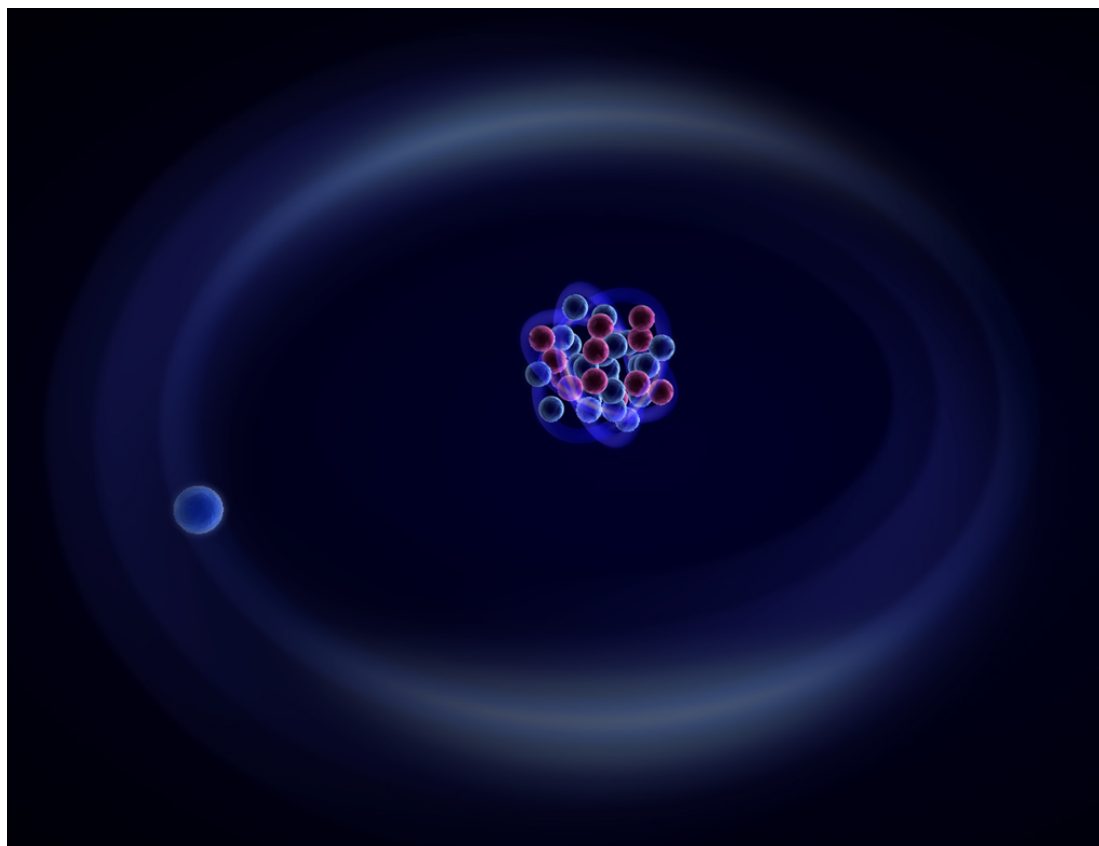


☞ **Double- Λ systems:** ${}^4_{\Lambda\Lambda}\text{H}$, ${}^6_{\Lambda\Lambda}\text{He}$, ${}^5_{\Lambda\Lambda}\text{H}$, ${}^5_{\Lambda\Lambda}\text{He}$, ${}^{10}_{\Lambda\Lambda}\text{Be}$, ${}^{13}_{\Lambda\Lambda}\text{B}$, ...??

What are Neutron-rich Halo Clusters?

- Loosely bound hadronic cluster with **tightly bound compact cores** with one to few **loosely bound neutrons** (near n -drip-lines) orbiting far away from the cores

e.g., famous ones \rightarrow ${}^6\text{He}$, ${}^{11}\text{Li}$



${}^{37}\text{Mg}$ one-neutron halo (credit: Ken-icheiro Yoneda, RIKEN, Japan)

Two-body Universality

☞ **Universality** of n -rich Halo/Cluster states:

E. Braaten, H.-W. Hammer, Phys. Rept. 428, 259 (2006)

P. Naidon, S. Endo, Rept. Prog. Phys. 80, 056001 (2017)

- ↪ Clusters of very different compositions behaving similarly at low-energies
- ↪ Similar long-range properties insensitive to short-distance details
- ↪ Suggests natural description in low-energy EFT framework

☞ \neq **EFT** → Model independent approach & distinct scale separation:

$$l_{\text{long}} \gg \frac{1}{\Lambda_{\text{Cut-off}} \approx m_{\pi}} \gtrsim l_{\text{natural/short}} \sim r_0 \} \text{ 2-body int. range}$$

↪ ideal for study of s-wave **threshold states** → $k_{\text{typical}} \rightsquigarrow 0$

☞ Existence of **shallow 2-body bound-states** whenever $a_0/r_0 \gg 1$

$$B_2^{\text{thr}} = -\frac{k_{\text{typical}}^2}{2\mu_{\text{reduced}}} = \lim_{a_0 \rightarrow \infty} \frac{-1}{2\mu_{\text{reduced}} a_0^2} \left[1 + \mathcal{O}\left(\frac{r_0}{a_0}\right) \right] \rightarrow 0$$

☞ Involves *unnatural scalings* & *fine-tuning effects* among low-energy parameters

↪ e.g., *unnaturally* large $|a_0| \gg r_0 \rightarrow l_{\text{natural/short}} = \{\lambda_\pi = m_\pi^{-1}, l_{vdW}, \dots\}$

☞ Short-distance 2-body interactions \rightarrow *Van der Waals* universality class

$$V_{L.Jones}(NN) \sim V_{vdW}(\text{atomic}) \stackrel{r \rightarrow \infty}{\sim} -\kappa_{\text{const.}}/r^6 ; l_{vdW} = \frac{1}{2}(2\mu_{\text{reduced}}\kappa_{\text{const.}})^{1/4}$$

Nuclear	$n - n$	$l_\pi \sim 1/m_\pi = 1.4 \text{ fm}$	$ a_0 = 23.7 \text{ fm}$	Fine-tuned
Atomic	${}^4\text{He} - {}^4\text{He}$	$l_{vdW} = 10.2 a_{\text{Bohr}}$	$ a_0 = 189 a_{\text{Bohr}}$	Fine-tuned
Atomic	${}^3\text{He} - {}^3\text{He}$	$l_{vdW} = 13.7 a_{\text{Bohr}}$	$ a_0 = 33 a_{\text{Bohr}}$	Natural

☞ **Fine-tuning \neq EFT Power Counting:** *Kaplan, Savage, Wise, U. van Kolck*

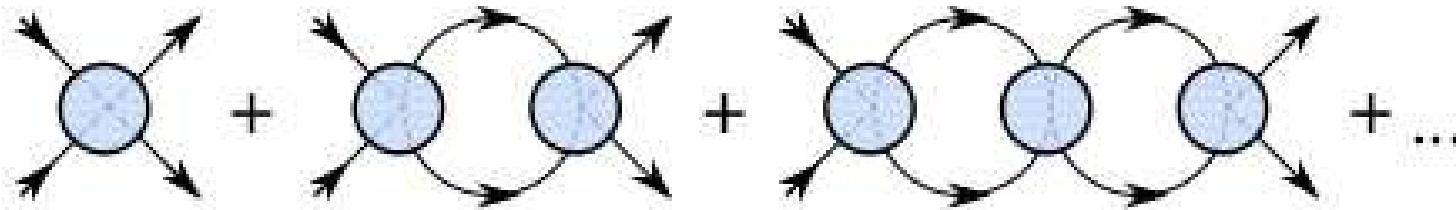
- non-perturbative resummation to all orders in $k_{\text{typical}}|a_0| \gg 1$
- perturbative fixed order expansion in $Q \equiv$ **Effective Range Expansion** *H. Bethe*

$$Q = k_{\text{typical}} r_0 \sim k_{\text{typical}}/\Lambda_{\text{Cut-off}} \ll 1$$

Toy Model: Identical Bosons

- ☞ Most general **Non-relativistic Effective** Lagrangian for two-body sector:
symmetries: P,C,T, Galilean Inv.

$$\mathcal{L}_{2\text{-body}} = \psi^\dagger \left[i\mathbf{v} \cdot \partial + \frac{(\mathbf{v} \cdot \partial)^2 - \partial^2}{2M} \right] \psi + \sum_{m+n=0,2,4,\dots} g_{m+n}^{(2)} \left[(\psi^\dagger \overleftrightarrow{\nabla}^m \psi) (\psi^\dagger \overleftrightarrow{\nabla}^n \psi) \right] + \dots$$



- ☞ **Unitarization:** Resum N^{th} order derivatively coupled tree vertices $\mathcal{V}_N^{\text{tree}} \sim \mathcal{O}(Q^N)$

$$Q\text{-Counting: } Q = \frac{k_{\text{typical}}}{\Lambda_{\text{Cut-off}} \sim m_\pi} \ll 1; \quad g_j^{(2)}(\mu) \sim Q^{-j-1}; \quad \mathcal{I}_{\text{loop}}(\mu) \sim \mu \sim Q$$

- ☞ **RG signature:**

↪ Existence of a **non-trivial UV fixed point** in $\{g_0^{(2)}, \dots, g_N^{(2)}\}$ space

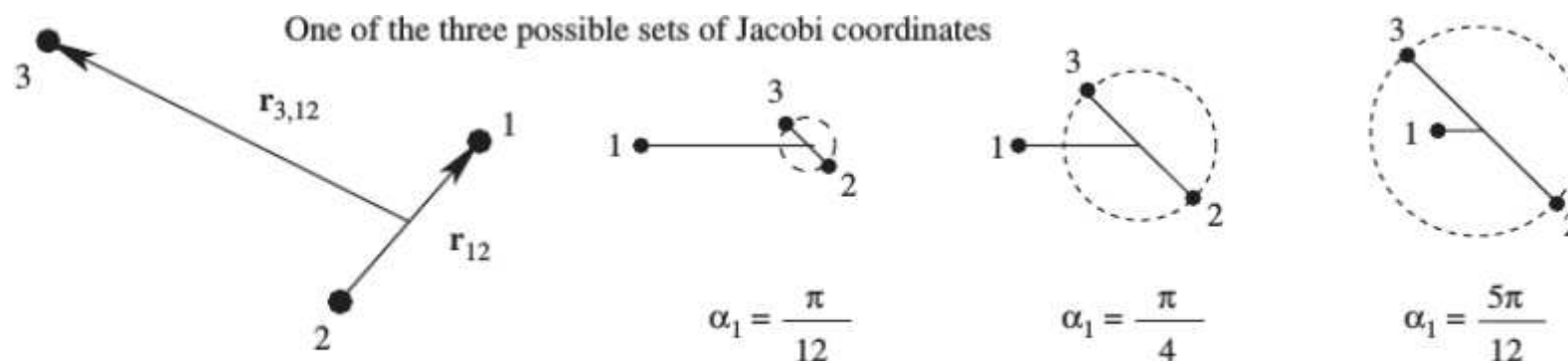
↪ **Unitary Limit:** Two-body interactions becoming resonant for $a_0 \rightarrow \pm\infty$

Three-body Universality: Interacting 3-boson system

☞ **Minimal Criterion:** at least two particle pairs with very large scattering lengths

↪ system becomes approximately scale invariant and resonant

☞ **Hyperspherical representation:** Hyper-radius R and Hyper-angles α_k



E. Braaten, H.-W. Hammer / Physics Reports 428 (2006) 259–390

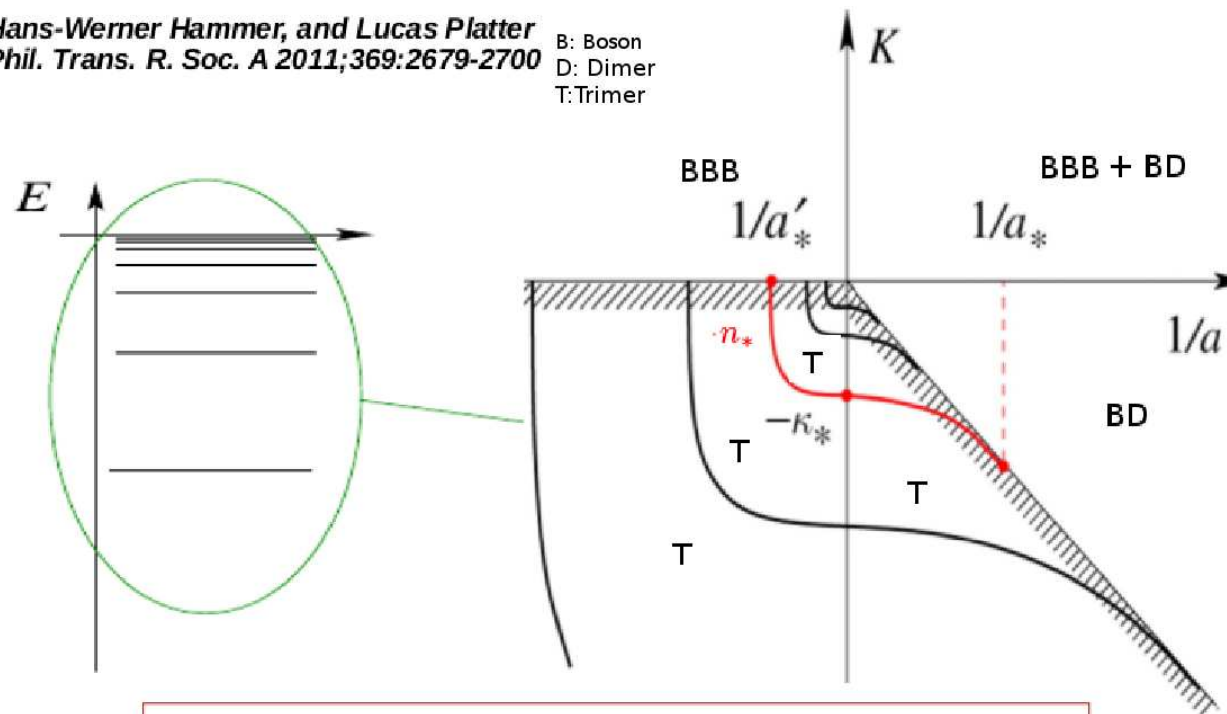
☞ **Radial Faddeev Eqns.** in co-ordinate representation with suitable asymptotic boundary conditions yields an *attractive* $1/R^2$ potential *V. Efimov, Phys. Lett. B33, 563 (1970)*

$$\frac{1}{2\mu_{\text{reduced}} a_0^2} \leq |B_3^{(n)}| \sim \left(e^{-2\pi/s_0} \right)^{n-n_*} \frac{\kappa_*^2}{2\mu_{\text{reduced}}} \leq \frac{1}{2\mu_{\text{reduced}} r_0^2} ; n = 0, \dots, n_{\text{max}}$$

Efimov Effect: Energy Spectrum of 3-boson system

Hans-Werner Hammer, and Lucas Platter
Phil. Trans. R. Soc. A 2011;369:2679-2700

B: Boson
 D: Dimer
 T: Trimer



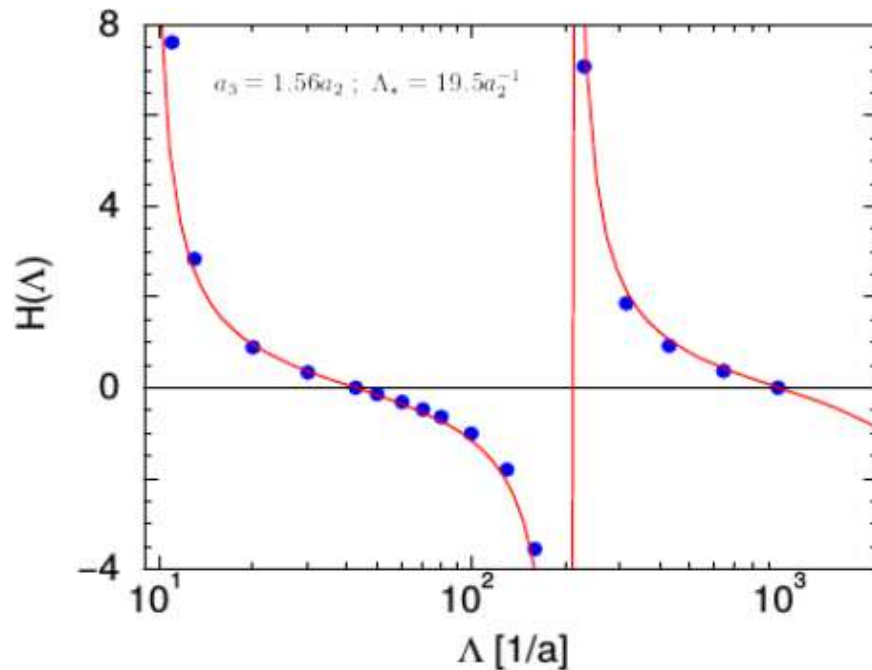
$$B_3^{(n)} / B_3^{(n+1)} \xrightarrow{1/a \rightarrow 0} 515.035\dots$$

- ☞ **Unitary Limit:** Infinitely many arbitrarily shallow 3-body bound states emerging from zero-energy threshold
- ☞ **Universality:** Only depends on Ratio of masses & Total Spin and Isospin

RG Limit Cycle: Interacting 3-boson system

- ☞ **RG signature:** Onset of a **IR Limit Cycle** behavior in the 3-body coupling $H(\Lambda)$
- ☞ Breakdown of *continuous scale invariance* \rightsquigarrow *discrete scaling symmetry*
- ☞ Momentum regulator (Λ) dependence in three-body coupling H :

$$\mathcal{L}_{3\text{-body}} \supset \frac{H(\Lambda)}{\Lambda^4} (\psi^\dagger \psi)^3 + \dots$$



$$\frac{\text{doll}^{(n)}}{\text{doll}^{(n+1)}} \approx e^{2\pi/s_0} \approx 515.035\dots$$

with $s_0 = 1.00624$

P.F. Bedaque et al. / Nuclear Physics A 646 (1999) 444–466

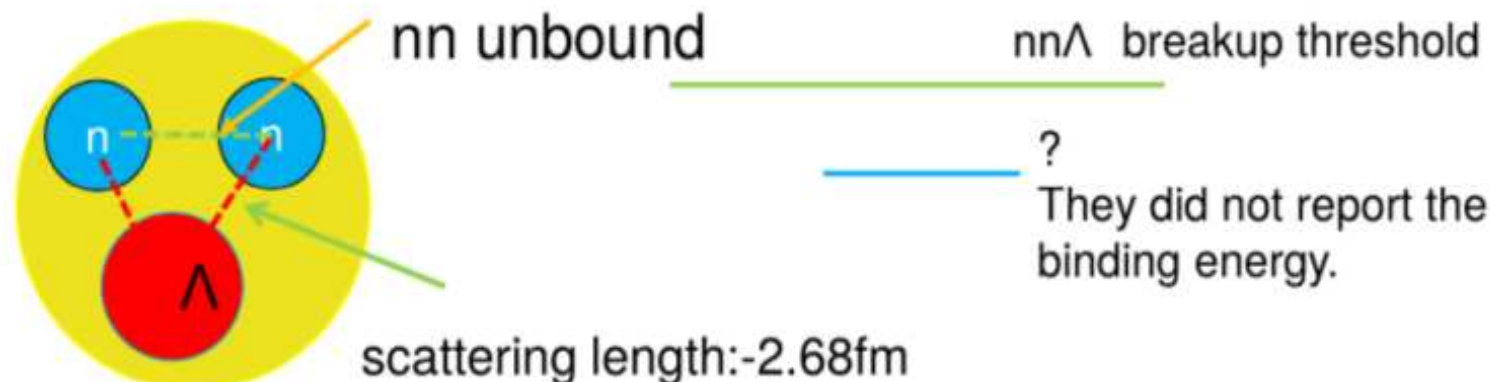
- ☞ **RG Limit Cycle:** Periodic self-similar repetition as in *Matryoshka dolls*

$J = 1/2, I = 1$ & $S = -1$ $nn\Lambda$ Hypernucleus??

PHYSICAL REVIEW C 88, 041001(R) (2013)

Search for evidence of ${}^3_{\Lambda}n$ by observing $d + \pi^-$ and $t + \pi^-$ final states in the reaction of ${}^6\text{Li} + {}^{12}\text{C}$ at 2A GeV

C. Rappold,^{1,2,*} E. Kim,^{1,3} T. R. Saito,^{1,4,5,†} O. Bertini,^{1,4} S. Bianchin,¹ V. Bozkurt,^{1,6} M. Kavatsyuk,⁷ Y. Ma,^{1,4} F. Maas,^{1,4,5} S. Minami,¹ D. Nakajima,^{1,8} B. Özel-Tashenov,¹ K. Yoshida,^{1,5,9} P. Achenbach,⁴ S. Ajimura,¹⁰ T. Aumann,^{1,11} C. Ayerbe Gayoso,⁴ H. C. Bhang,³ C. Caesar,^{1,11} S. Erturk,⁶ T. Fukuda,¹² B. Göküzüm,^{1,6} E. Guliev,⁷ J. Hoffmann,¹ G. Ickert,¹ Z. S. Ketenci,⁶ D. Khanef, ^{1,4} M. Kim,³ S. Kim,³ K. Koch,¹ N. Kurz,¹ A. Le Fèvre,^{1,13} Y. Mizoi,¹² L. Nungesser,⁴ W. Ott,¹ J. Pochodzalla,⁴ A. Sakaguchi,⁹ C. J. Schmidt,¹ M. Sekimoto,¹⁴ H. Simon,¹ T. Takahashi,¹⁴ G. J. Tambave,⁷ H. Tamura,¹⁵ W. Trautmann,¹ S. Voltz,¹ and C. J. Yoon³
(HypHI Collaboration)



Formation recently suggested by the experimental studies by the **HypHI Collaboration (2013)**

$nn\Lambda$ Searches and Predictions

☞ **Currently still a controversial topic:**

- *No concrete* experimental data other than rough bounds by **HypHI**

- **Supported in Lattice QCD:** *Beane et. al. NPLQCD Collaboration (2013)*

- **Inconsistent with Potential Models:** *Barnea, Gall, Contessi, Garcilazo, et al....*

Models with $\Lambda N - \Sigma N$ mixing fail to reproduce $n\Lambda$ scattering length, ${}^3_{\Lambda}\text{H}$ binding energy, and the ground and first excited states energies of ${}^4_{\Lambda}\text{H}$ and ${}^4_{\Lambda}\text{He}$.

- Expect more definitive exp. results in future from **ALICE, FAIR,..**

- ☞ First **EFT** analysis to investigate possible formation of Efimov-like three-body bound states *Ando et al., PRC 92 (2015) 024325*

$nn\Lambda$ @ LO in $\not\chi$ EFT

- ☞ **RG Limit cycle signature:** Qualitative investigation of Efimov-like bound states
- ☞ $\not\chi$ **EFT:** Pions are integrated out, $\Lambda_H \sim m_\pi$
- ☞ **Avoid coupled-channel analysis:** $\Sigma^{0,\pm}$ -hyperons not explicitly incorporated

$$\delta M = M_\Sigma - M_\Lambda \simeq 80 \text{ MeV} \sim \Lambda_{\text{Hard}} = m_\pi \gg B_{nn\Lambda} \lesssim 1 \text{ MeV}$$

☞ **2-body input Parameters in EFT:**

- nn -Singlet Scat. length: $a_{nn} = -18.5 \text{ fm}$ (exp.) *Teramond & Gabioud*
- $n\Lambda$ -Singlet Scat. length: $a_{s(n\Lambda)} \xrightarrow{\text{isospin } 1} a_{p\Lambda} = -2.90 \text{ fm}$ χ EFT *Polinder et. al.*
- $n\Lambda$ -Triplet Scat. length: $a_{t(n\Lambda)} \xrightarrow{\text{isospin } 3} a_{p\Lambda} = -1.60 \text{ fm}$ χ EFT *Polinder et. al.*

☞ **3-body input:** (no datum currently) Rough estimate given from **HypHI Coll.**

$$\hookrightarrow B_{nn\Lambda} \sim 0.01 - 1.1 \text{ MeV}$$

1-Body Sector Lagrangian

☞ Explicit degrees of freedom:

- Neutron field N
- Λ -Hyperon field Λ

$$\mathcal{L}_n = N^\dagger \left[i v \cdot \partial + \frac{(v \cdot \partial)^2 - \partial^2}{2M_n} \right] N + \text{higher derivatives}$$

$$\mathcal{L}_\Lambda = \Lambda^\dagger \left[i v \cdot \partial + \frac{(v \cdot \partial)^2 - \partial^2}{2M_\Lambda} \right] \Lambda + \text{higher derivatives}$$

☞ 4-velocity vector which is chosen as $v^\mu = (1, \mathbf{0})$

Quasiparticle Formalism: 2-Body Sector

Beane, Badaque, Hammer, Ubi. vanKolck, Birse, et al

☞ **Cluster fields** or **Dibaryons** (non-dynamical auxiliary d.o.f.s):

- Spin-Singlet nn -Dibaryon field $\rightarrow s_{(nn)}$
- Spin-Singlet $n\Lambda$ -Dibaryon field $\rightarrow s_{(n\Lambda)}$
- Spin-Triplet $n\Lambda$ -Dibaryon field $\rightarrow t_{(n\Lambda)}$

☞ **Unitarization:** Improves convergence in the vicinity of bound states

$$\mathcal{L}_{s(nn)} = -s_{(nn)}^\dagger \left[iv \cdot \partial + \frac{(v \cdot \partial)^2 - \partial^2}{4M_n} \right] s_{(nn)} - y_{s(nn)} \left[s_{(nn)}^\dagger \left(N^T P_{(nn)}^{(1S_0)} N \right) + \text{h.c.} \right] + \dots$$

$$\mathcal{L}_{s(n\Lambda)} = -s_{(n\Lambda)}^\dagger \left[iv \cdot \partial + \frac{(v \cdot \partial)^2 - \partial^2}{2(M_n + M_\Lambda)} \right] s_{(n\Lambda)} - y_{s(n\Lambda)} \left[s_{(n\Lambda)}^\dagger \left(N^T P_{(n\Lambda)}^{(1S_0)} \Lambda \right) + \text{h.c.} \right] + \dots$$

$$\mathcal{L}_{t(n\Lambda)} = -\sum_{j=1}^3 t_{(n\Lambda)j}^\dagger \left[iv \cdot \partial + \frac{(v \cdot \partial)^2 - \partial^2}{2(M_n + M_\Lambda)} \right] t_{(n\Lambda)j} - y_{t(n\Lambda)} \left[\sum_{j=1}^3 t_{(n\Lambda)j}^\dagger \left(N^T P_{(n\Lambda)j}^{(3S_1)} \Lambda \right) + \text{h.c.} \right] + \dots$$

Faddeev-type Integral Eqns. in momentum space

- These **STM** (*Skornyakov, Ter-Martirosyan*) **Eqns.** must be **numerically solved** to obtain possible energy eigenvalues
- If Efimov effect is manifest then the integral equations are ill-defined in the UV limit
- Loops need to be regulated with cut-off Λ dependent **3-body counter-terms** from additional 3-body Lagrangian $\mathcal{L}_{3\text{-body}}$ (counter-terms not shown)

$$\begin{array}{c}
 t(n\Lambda) \\
 \begin{array}{c} \text{---} \\ \text{---} \end{array} \begin{array}{c} n \\ t(n\Lambda) \end{array} \\
 \mathcal{T}_a \text{ (hatched)} \\
 n
 \end{array}
 =
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 +
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 \mathcal{T}_a \text{ (hatched)}
 +
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 \mathcal{T}_b \text{ (vertical stripes)}
 +
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 \mathcal{T}_c \text{ (grid)}$$

$$\begin{array}{c}
 t(n\Lambda) \\
 \begin{array}{c} \text{---} \\ \text{---} \end{array} \begin{array}{c} n \\ s(n\Lambda) \end{array} \\
 \mathcal{T}_b \text{ (vertical stripes)} \\
 n
 \end{array}
 =
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 +
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 \mathcal{T}_a \text{ (hatched)}
 +
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 \mathcal{T}_b \text{ (vertical stripes)}
 +
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 \mathcal{T}_c \text{ (grid)}$$

$$\begin{array}{c}
 nn \\
 \begin{array}{c} \text{---} \\ \text{---} \end{array} \begin{array}{c} n \\ t(n\Lambda) \end{array} \\
 \mathcal{T}_c \text{ (grid)} \\
 \Lambda
 \end{array}
 =
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 +
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 \mathcal{T}_a \text{ (hatched)}
 +
 \begin{array}{c} \text{---} \\ \text{---} \end{array}
 \mathcal{T}_b \text{ (vertical stripes)}$$

3-body Lagrangian (Part)

- ☞ **EFT TENET**: Introduce non-derivatively coupled contact terms with 3-body coupling $g(\Lambda)$ in the elastic channel, e.g.,

$$n + t(n\Lambda) \rightarrow n + t(n\Lambda)$$

Promoted to Leading Order in EFT

$$\mathcal{L}_{3\text{-body}} \supset \overbrace{-\frac{1}{6} M_\Lambda y_{t(n\Lambda)}^2 \frac{g(\Lambda)}{\Lambda^2} t_{(n\Lambda)i}^\dagger N^\dagger \sigma_i \sigma_j N t_{(n\Lambda)j}} + \dots$$

+ neglect sub-leading derivative operators in all channels

- ☞ **3-body input**: $B_{nn\Lambda} \lesssim 1 \text{ MeV}$ can be used to fix the coupling $g(\Lambda)$.

Asymptotic Analysis

- Investigate asymptotic behavior of integral equations to determine possible **Efimov-effect in unitary limit**

$$p \sim \Lambda \gg B_{nn\Lambda}, k \sim 1/a_{nn}, 1/a_{t(n\Lambda)}, 1/a_{s(n\Lambda)}$$

- Channel amplitudes obey simple power-law behavior: $\mathcal{T}_{a,b,c}(p) \sim p^{s-1}$

→ integral equations decouple

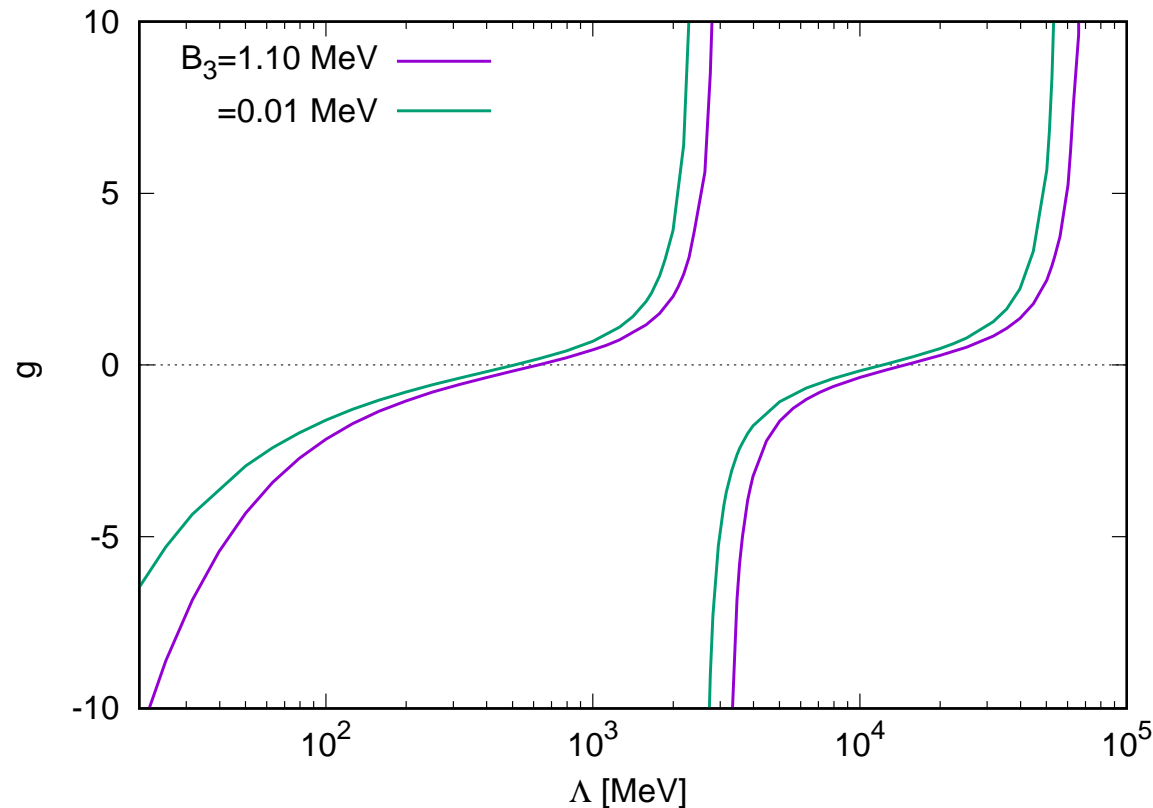
- Asymptotic expansion followed by Mellin transformations:**

$$\left(\frac{M_\Lambda}{2\pi\mu_{n\Lambda}C_1} \right) \left[\frac{2\pi \sin \left[s \sin^{-1} \frac{\mu_{n\Lambda}}{M_\Lambda} \right]}{s \cos[\pi s/2]} \right] + \left(\frac{M_n}{\pi^2\mu_{n\Lambda}C_1C_2} \right) \left[\frac{2\pi \sin \left[s \cot^{-1} \sqrt{\frac{2M_n - \mu_{n\Lambda}}{\mu_{n\Lambda}}} \right]}{s \cos[\pi s/2]} \right]^2 = 1$$

$$C_1 = \sqrt{\frac{\mu_{n\Lambda}}{\mu_{n(n\Lambda)}}}, \quad C_2 = \sqrt{\frac{M_n}{\mu_{\Lambda(nn)}}}$$

$$\mu_{n\Lambda} = \frac{M_n M_\Lambda}{M_n + M_\Lambda}, \quad \mu_{n(n\Lambda)} = \frac{M_n(M_n + M_\Lambda)}{2M_n + M_\Lambda}, \quad \mu_{\Lambda(nn)} = \frac{2M_n M_\Lambda}{2M_n + M_\Lambda}$$

Limit Cycle RG of 3-Body Coupling for $nn\Lambda$



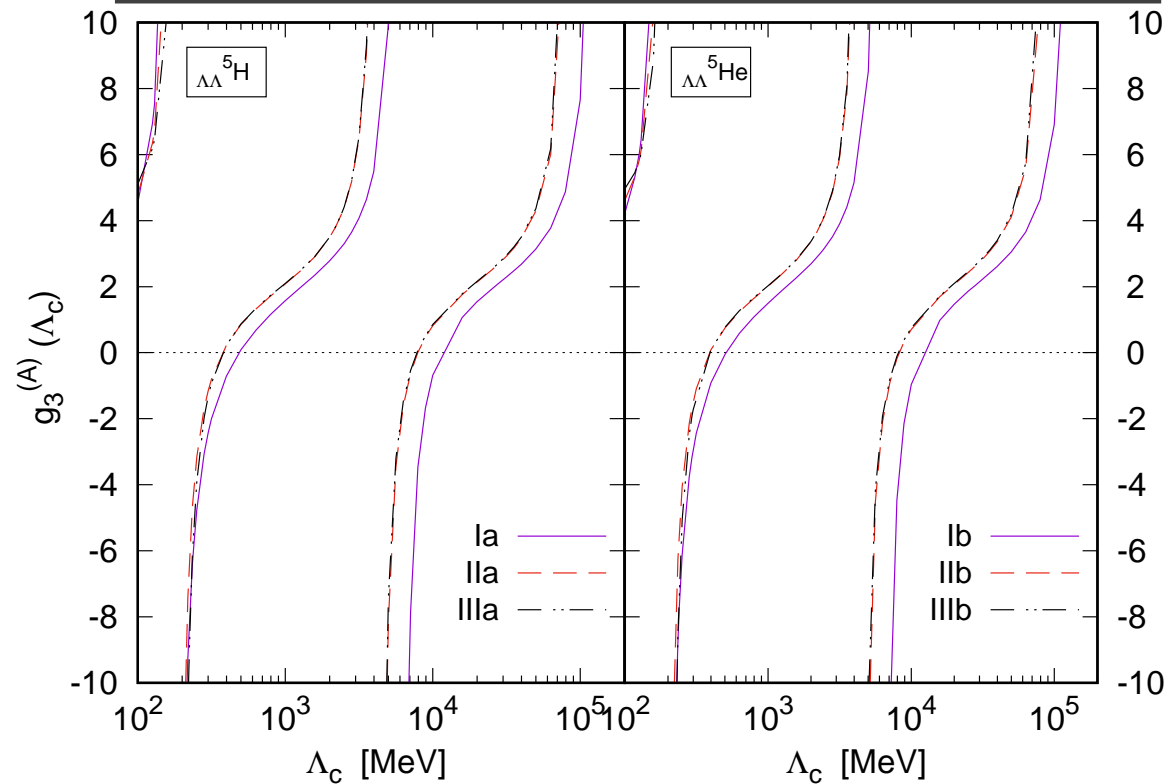
☞ **Efimov Effect:** Solving for s yields imaginary values \rightarrow **Limit Cycle**

$$s = \pm i s_0 \quad ; \quad s_0 = 1.00760\dots$$

☞ **Cyclic Periodicity:** $g(\Lambda_n) = g(\Lambda_1 \sim 500 - 600 \text{ MeV})$ with

$$\Lambda_n = \Lambda_1 e^{n\pi/s_0}$$

Limit Cycle RG in (${}^5_{\Lambda\Lambda}\text{H}$, ${}^5_{\Lambda\Lambda}\text{He}$) *G. Meher et al*



☞ **Efimov Effect:** Solving for s yields imaginary values \rightarrow **Limit Cycle**

$$s = \pm i s_0 \quad ; \quad s_0 \approx 1.0351\dots$$

☞ **Cyclic Periodicity:** $g(\Lambda_n) = g(\Lambda_1 \sim 300 - 400 \text{ MeV})$ with

$$\Lambda_n = \Lambda_1 e^{n\pi/s_0}$$