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TECHNISCHE
UNIVERSITÄT
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ELECTROWEAK NLO CALCULATIONS

Rikkert Frederix

Technische Universität München

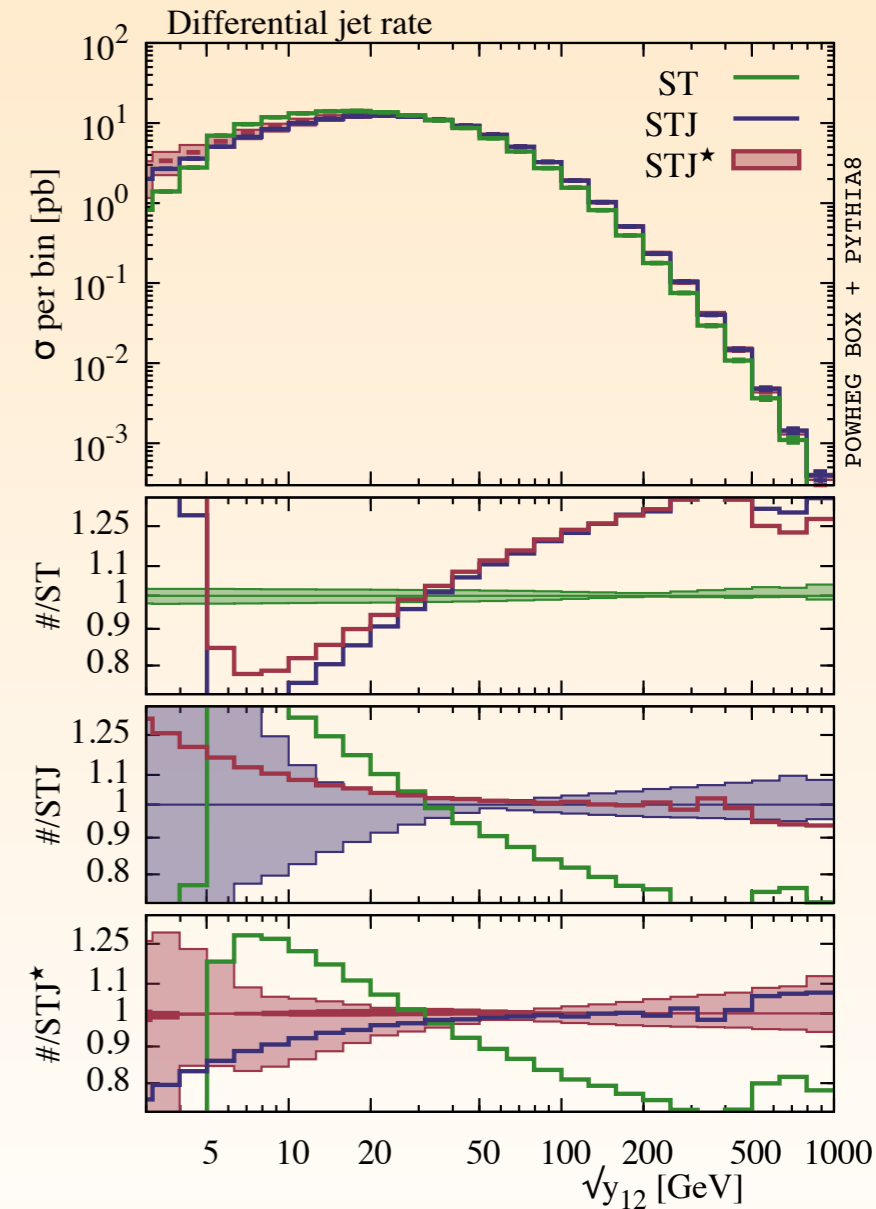
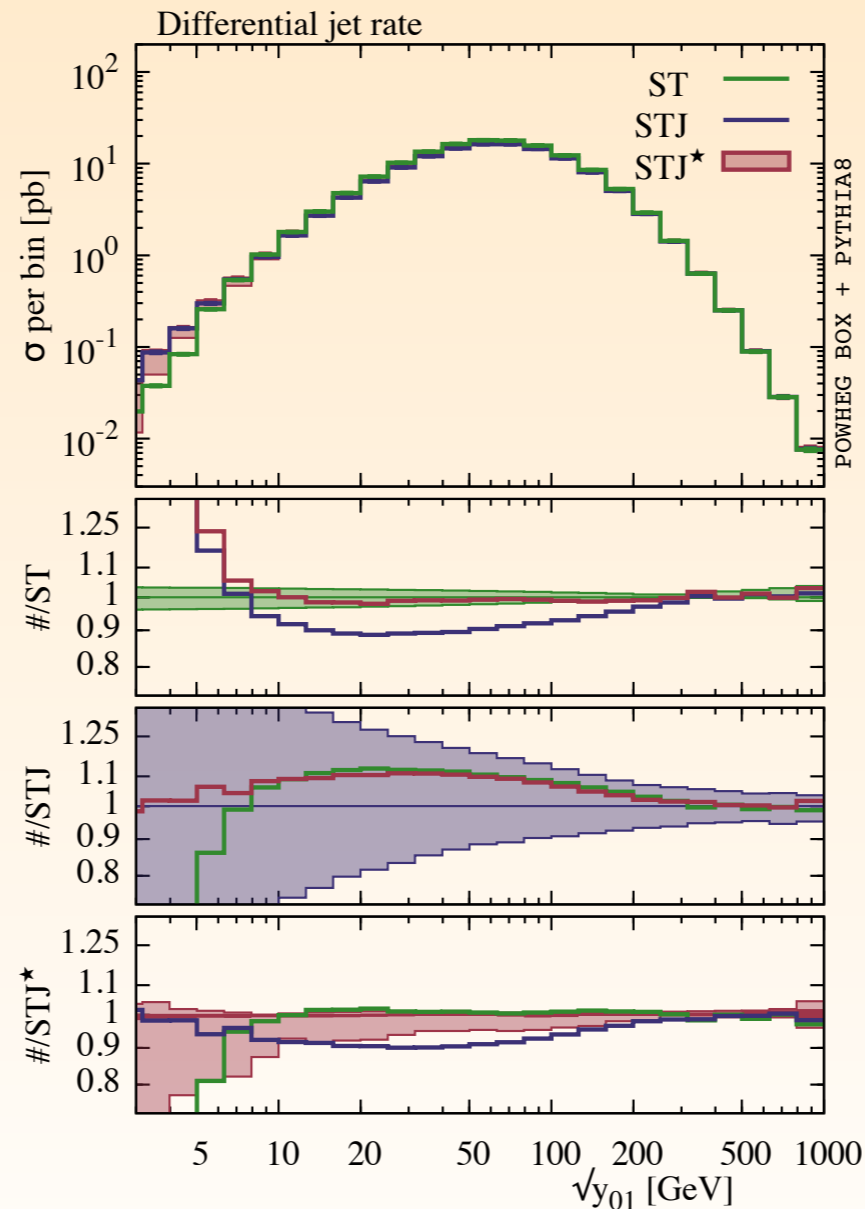
and on leave from Lund University

MINLO' T-CHANNEL SINGLE-TOP

- ◆ **ST**: default POWHEG NLO t-channel single-top predictions
- STJ**: new t-channel single-top+jet NLO in POWHEG
- STJ***: Minlo' merged **ST+STJ** (without merging scale through enforcing unitarity)

- ◆ Small differences between **ST/STJ** and **STJ*** at small scales, but this is deep in the Sudakov region, where higher accurate resummation is needed (and non-perturbative effects play an important role as well)

- ◆ Uncertainty band for **ST** y_{12} is too small -> artefact of POWHEG methodology



- ◆ **The preferred (i.e. most-accurate) predictions for t-channel single-top production**

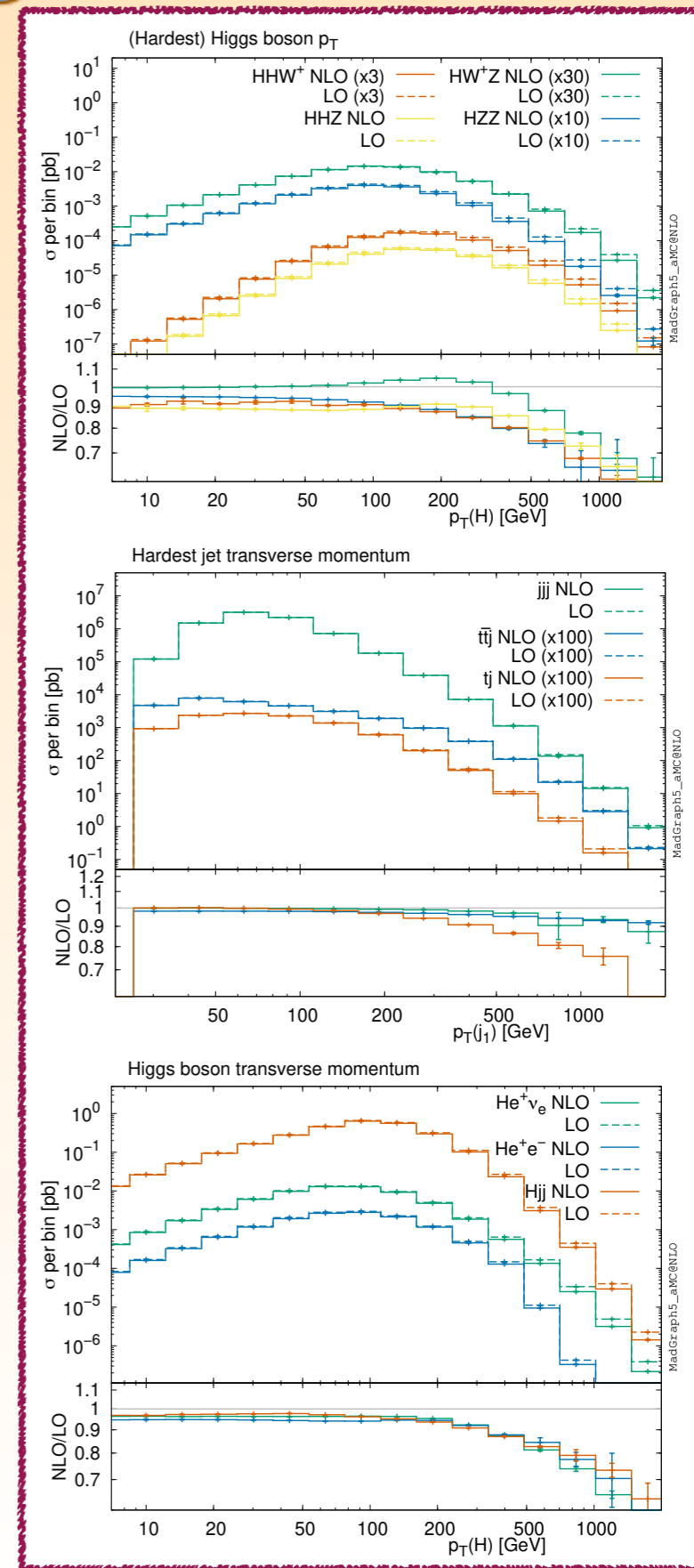
NLO EW CORRECTIONS

EW CORRECTIONS

- ◆ Just as one can have a perturbative series in the strong coupling, one can also include higher order corrections in the electroweak (EW) coupling
- ◆ By comparing the strength of the strong to the EW coupling, one **expects that NNLO QCD corrections of similar importance to NLO EW corrections**
 - On top of that, **EW corrections can be enhanced in certain kinematical regions**, where they can result in several tens of percents:
 - ◆ Close to EW resonances, radiation from decay products results in sizeable changes
 - ◆ When photon luminosity is important
 - ◆ Large transverse) momenta or invariants result in large EW corrections
 - ❖ Important in BSM searches, **particularly when understanding shapes of backgrounds is a must**

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EW VS STRONG CORRECTIONS

- ◆ When including higher order corrections in the strong coupling, renormalisation (and factorisation) scale dependence is reduced in the predictions
- ◆ This is not the case for EW corrections: scale dependence is effectively the same in LO and NLO EW computations
 - Instead, scheme dependence is reduced
 - Note that scheme dependence is typically not considered to be an uncertainty: it is quite obvious which scheme is preferred

EW SCHEME CHOICES

- ◆ The EW sector of the SM has 3 independent parameters for the gauge interactions. Historically taken to be α , M_W and M_Z (with α measured in Thomson scattering, and M_W and M_Z the on-shell weak boson masses)
 - Other EW parameters are then predictions: v_{ev} , G_F , $\sin(\theta_W)$, λ , ρ , ...
- ◆ Alternatively, by using other input parameters, and updating the renormalisation conditions accordingly, one resums some important higher order contributions
 - At LO, scheme dependence is only through the numerical value of the input parameters (which effectively means the value of α)

COMMON EW SCHEMES: OVERVIEW

$\{\alpha(0), M_W, M_Z\} \rightarrow \alpha(0)$ scheme

$$\delta Z_e = -\frac{1}{2}\delta Z_{AA} - \frac{s_W}{c_W} \frac{1}{2}\delta Z_{ZA}$$

$\{\alpha(M_Z), M_W, M_Z\} \rightarrow \alpha(M_Z)$ scheme

$$\delta Z_e|_{\alpha(Q^2)} \equiv \delta Z_e|_{\alpha(0)} - \frac{1}{2}\Delta\alpha(Q^2)$$

$\{G_\mu, M_W, M_Z\} \rightarrow G_\mu$ scheme

$$\delta Z_e|_{G_\mu} \equiv \delta Z_e|_{\alpha(0)} - \frac{1}{2}\Delta r =$$

$$\delta Z_e|_{\alpha(m_Z^2)} - \frac{1}{2} \left(-\frac{c_W^2}{s_W^2} \Delta\rho + \Delta r_{\text{rem}} \right)$$

$$\alpha(0) \sim 1/137 \quad \alpha(M_Z) \sim 1/128 \quad \alpha|_{G_\mu} \sim 1/132$$

As a rule of thumb, for a generic process at the LHC, the G_μ scheme is superior and has to be preferred. However, if a photon is present in the Born final-state, $\alpha(0)$ and the corresponding renormalisation should be used for the associated QED vertex.

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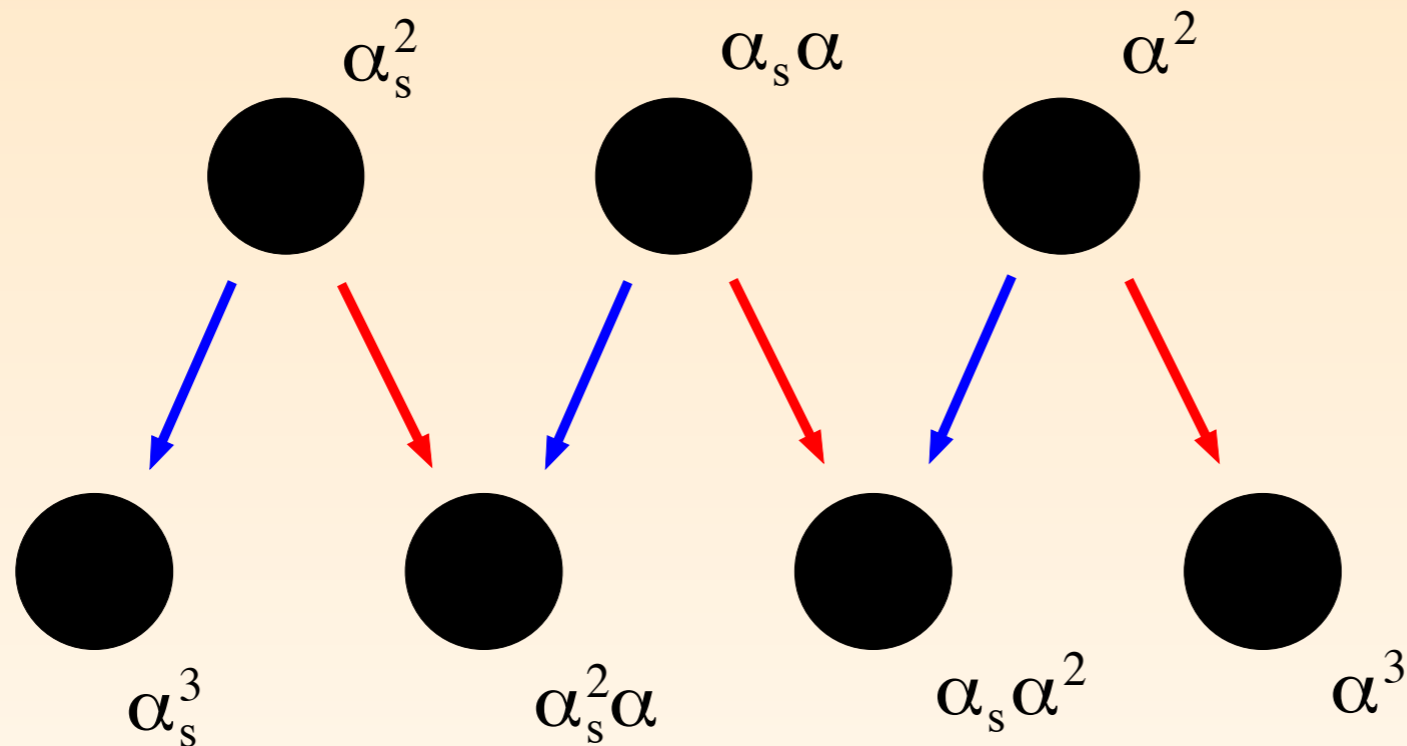
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NLO DISSECTION

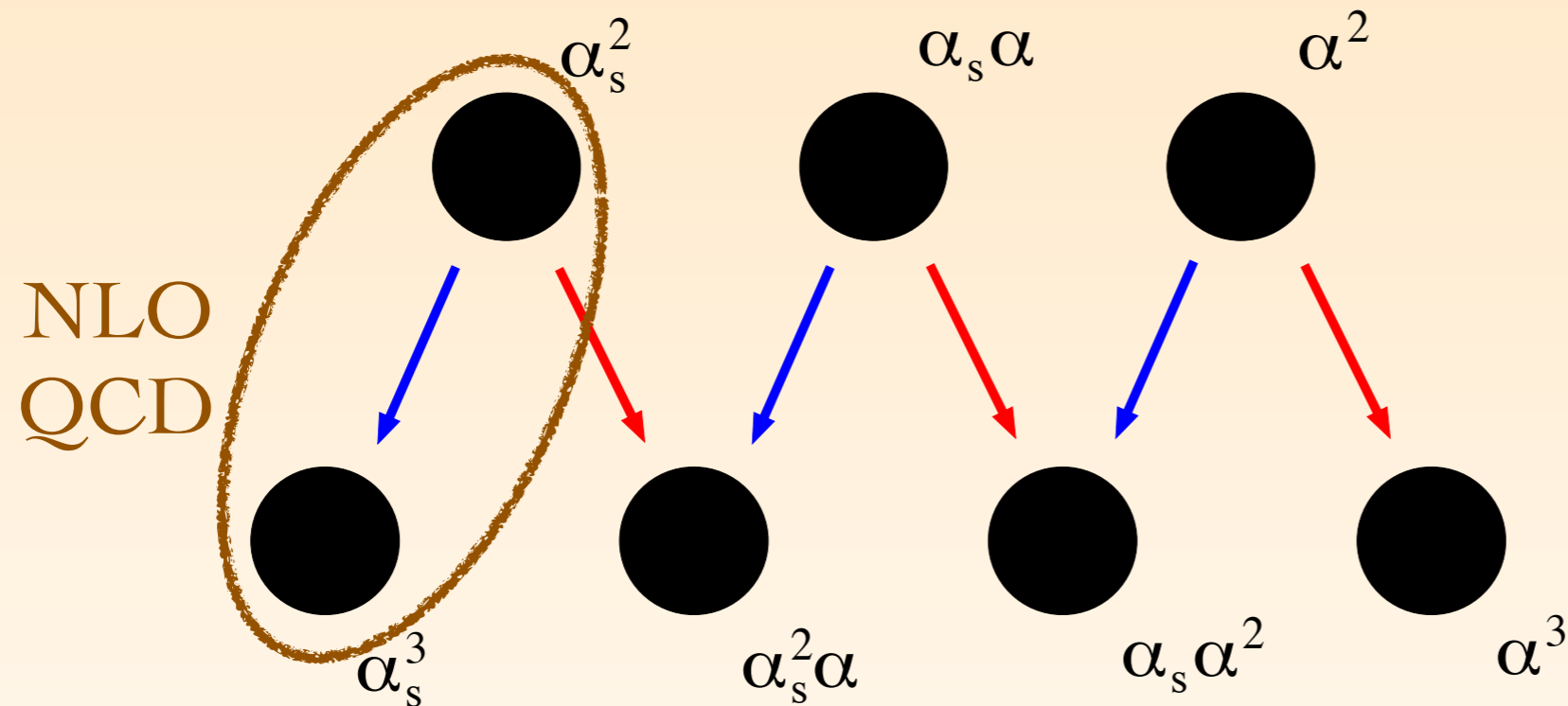
- ◆ For example: consider di-jet production



- ◆ "NLO EW" is a bit of a misnomer:
 NLO_2 and NLO_3 part of a "mixed" expansion
- ◆ "Complete-NLO" takes all the LO and NLO contributions in the mixed coupling expansion into account

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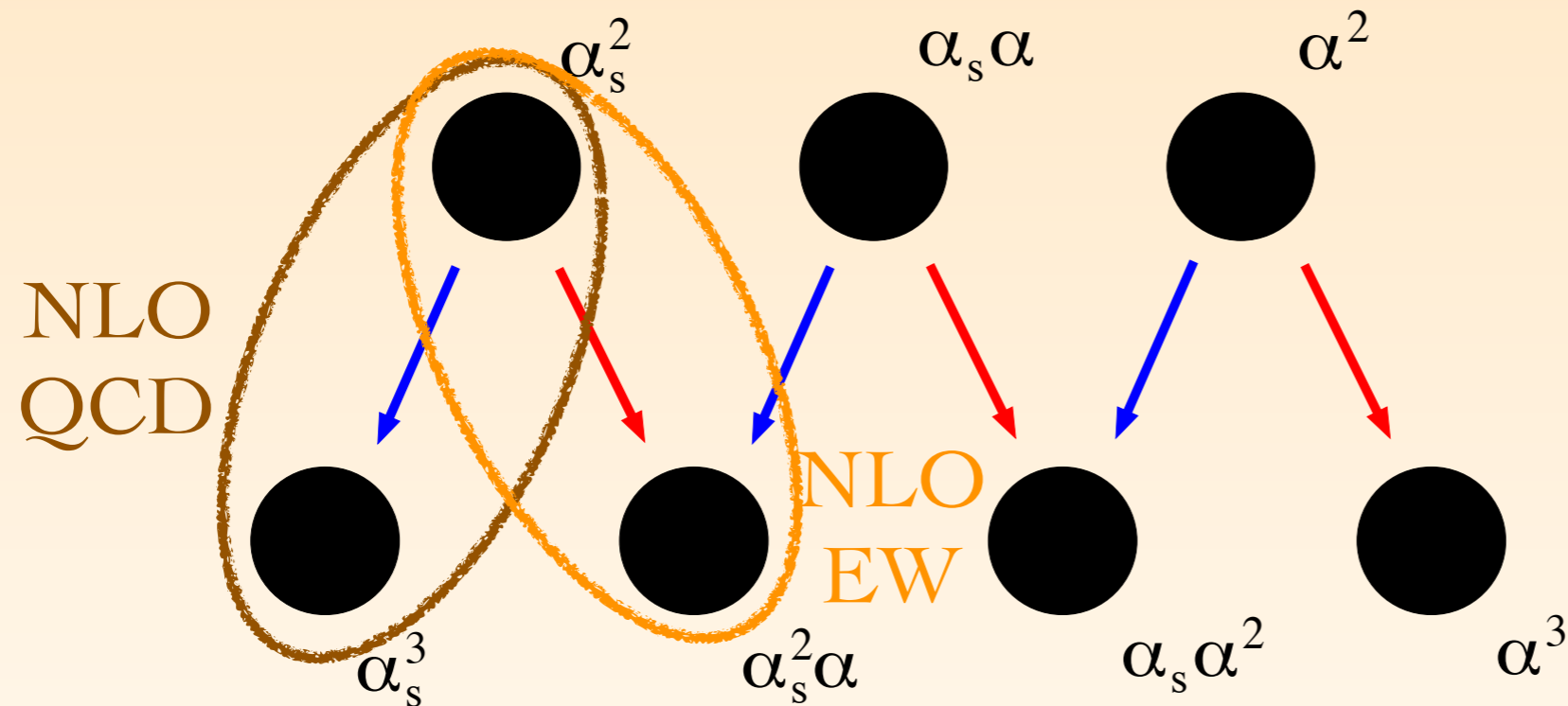
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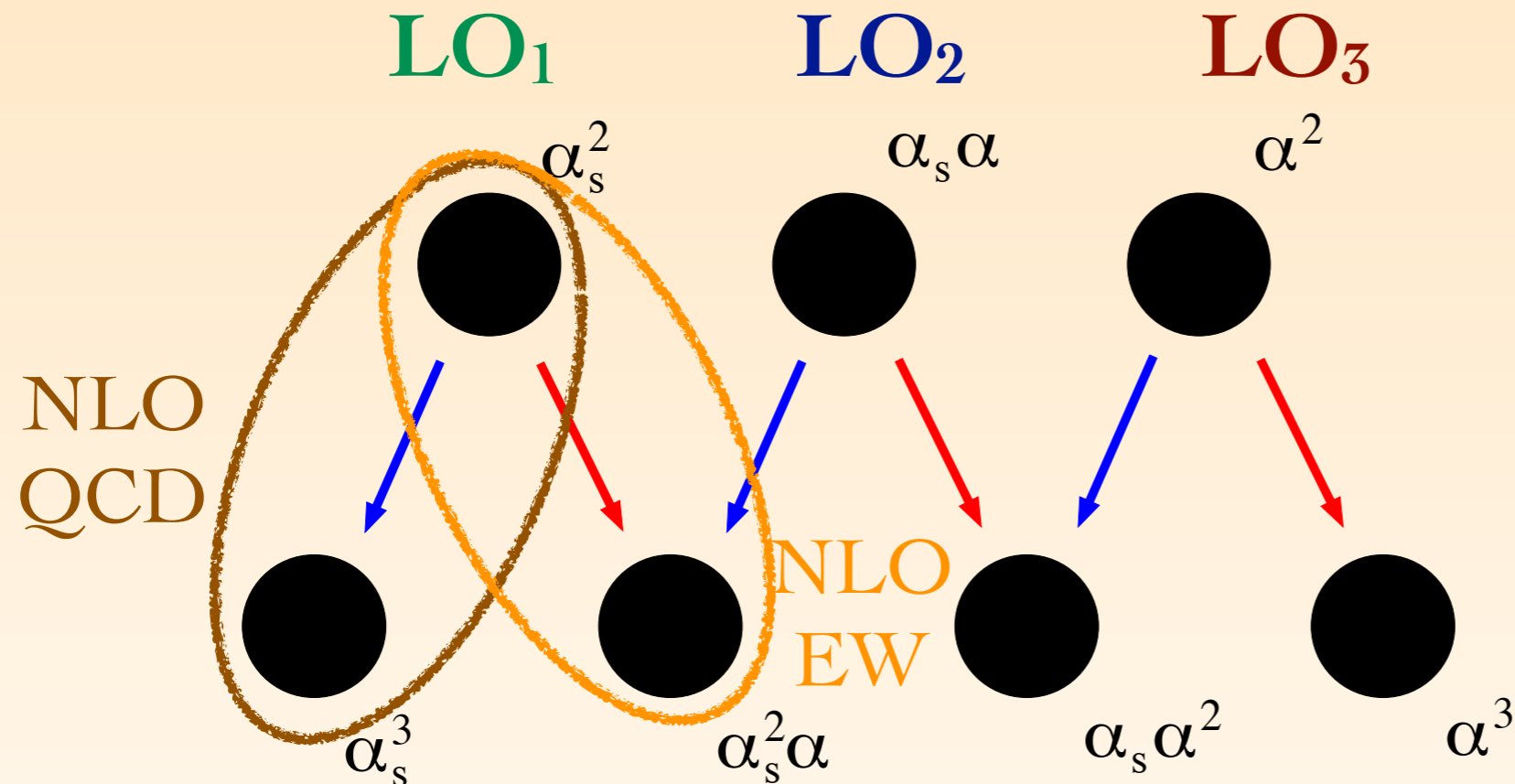
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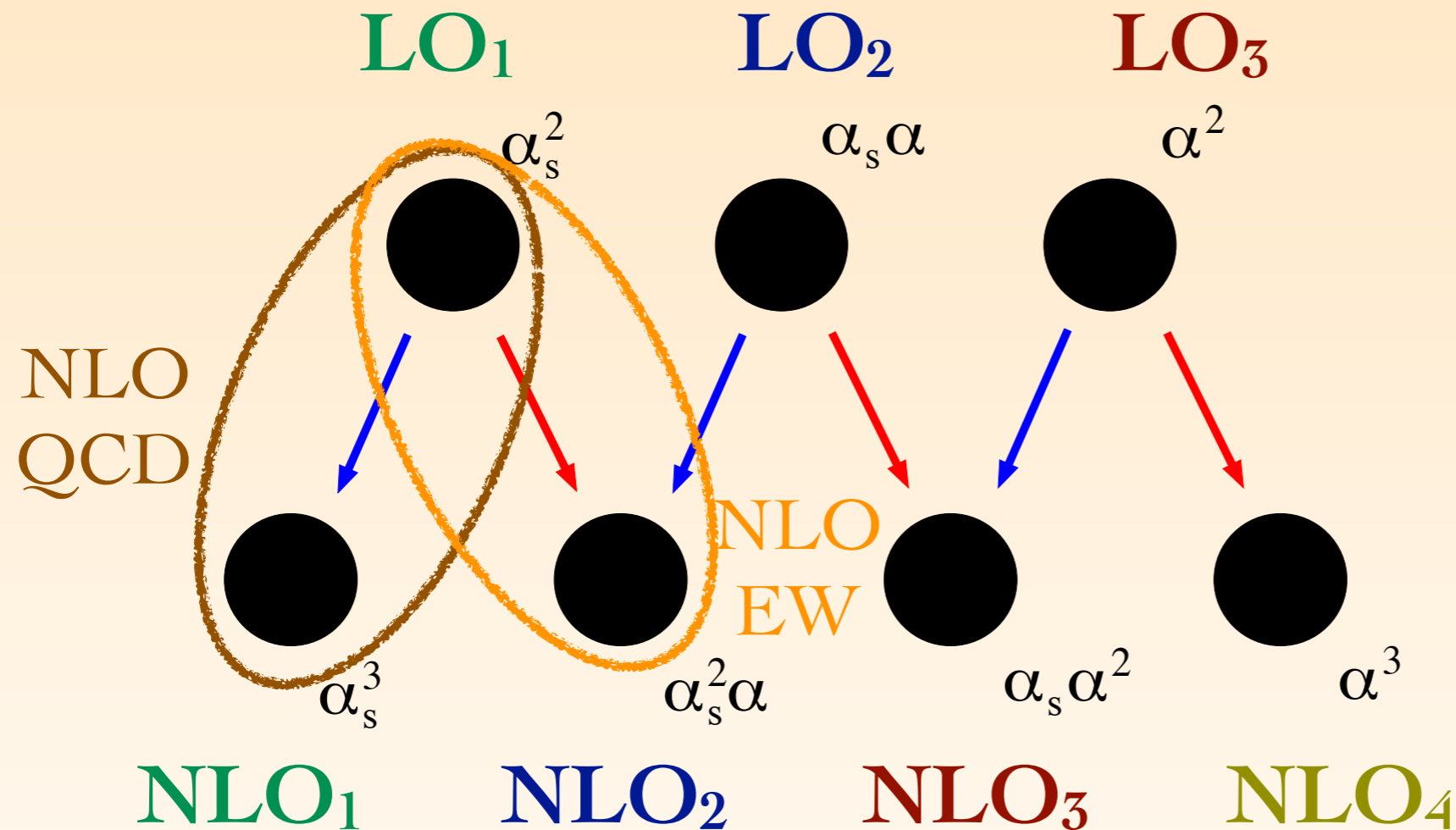
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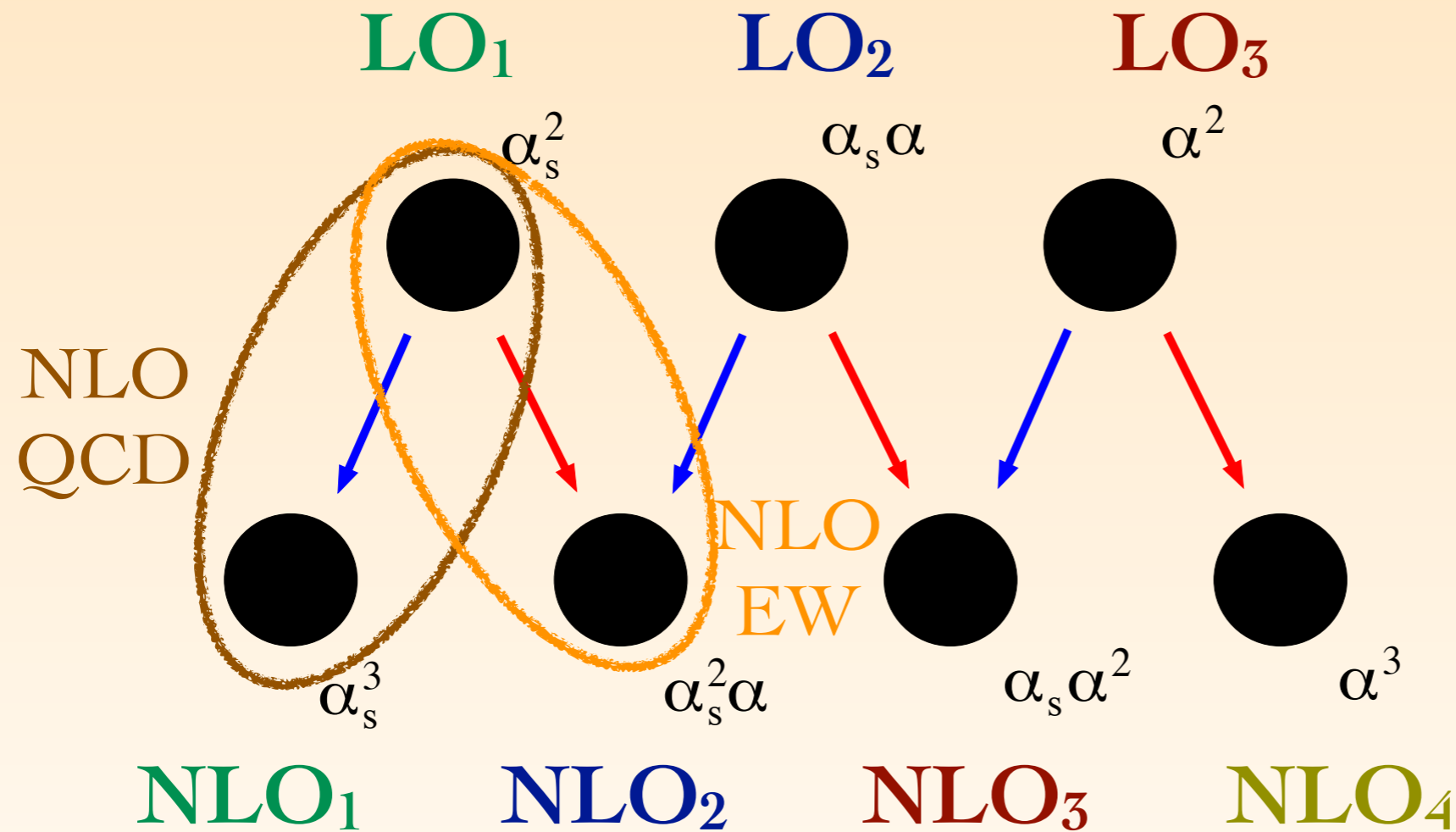
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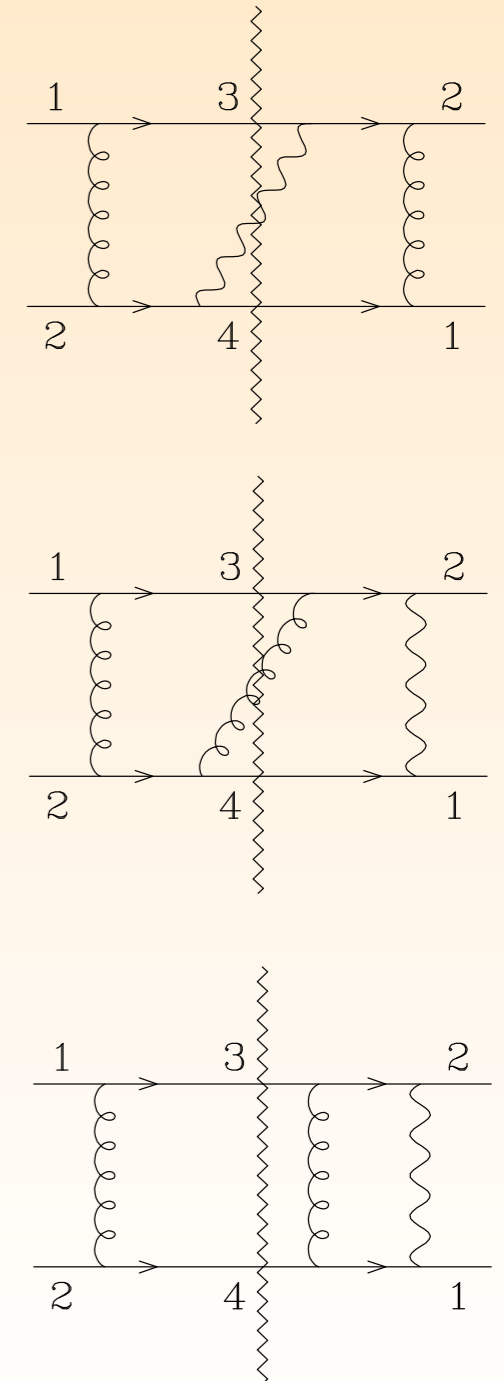
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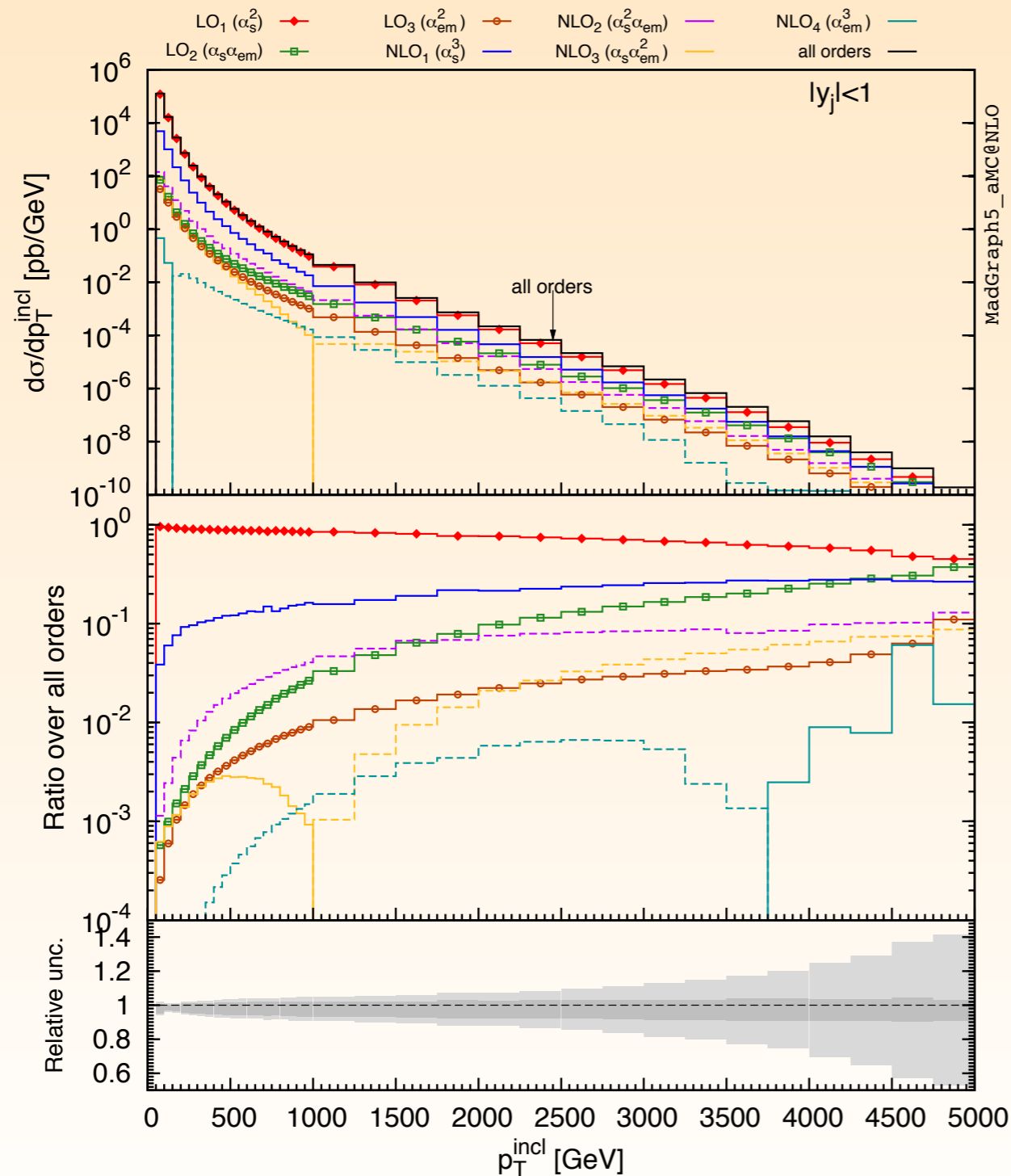
For example: NLO_2 has both "QCD" and "EW" contributions



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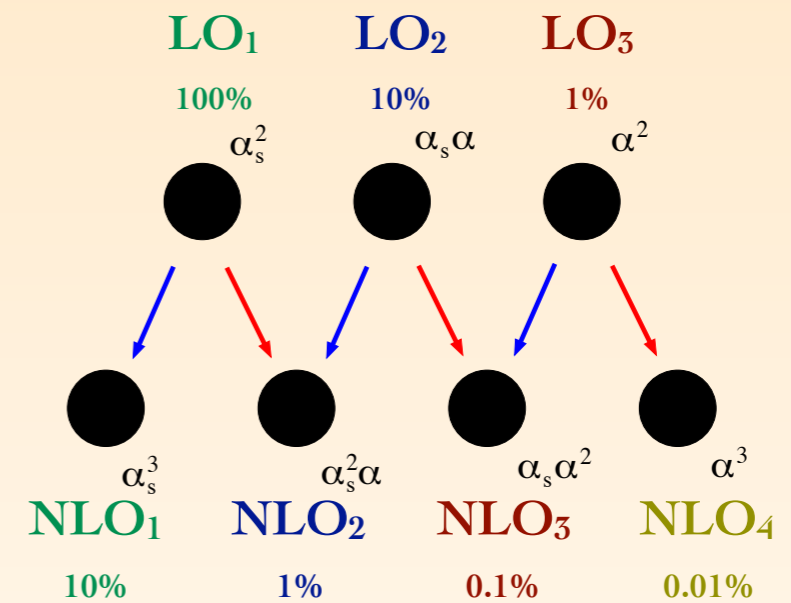
COMPLETE-NLO TO INCLUSIVE JET-PT

[RF, Frixione, Hirschi, Pagani, Shao, Zaro, 2017]



◆ Inclusive jet- p_T

◆ Expectation (assume $\alpha_s=0.1$, $\alpha=0.01$):

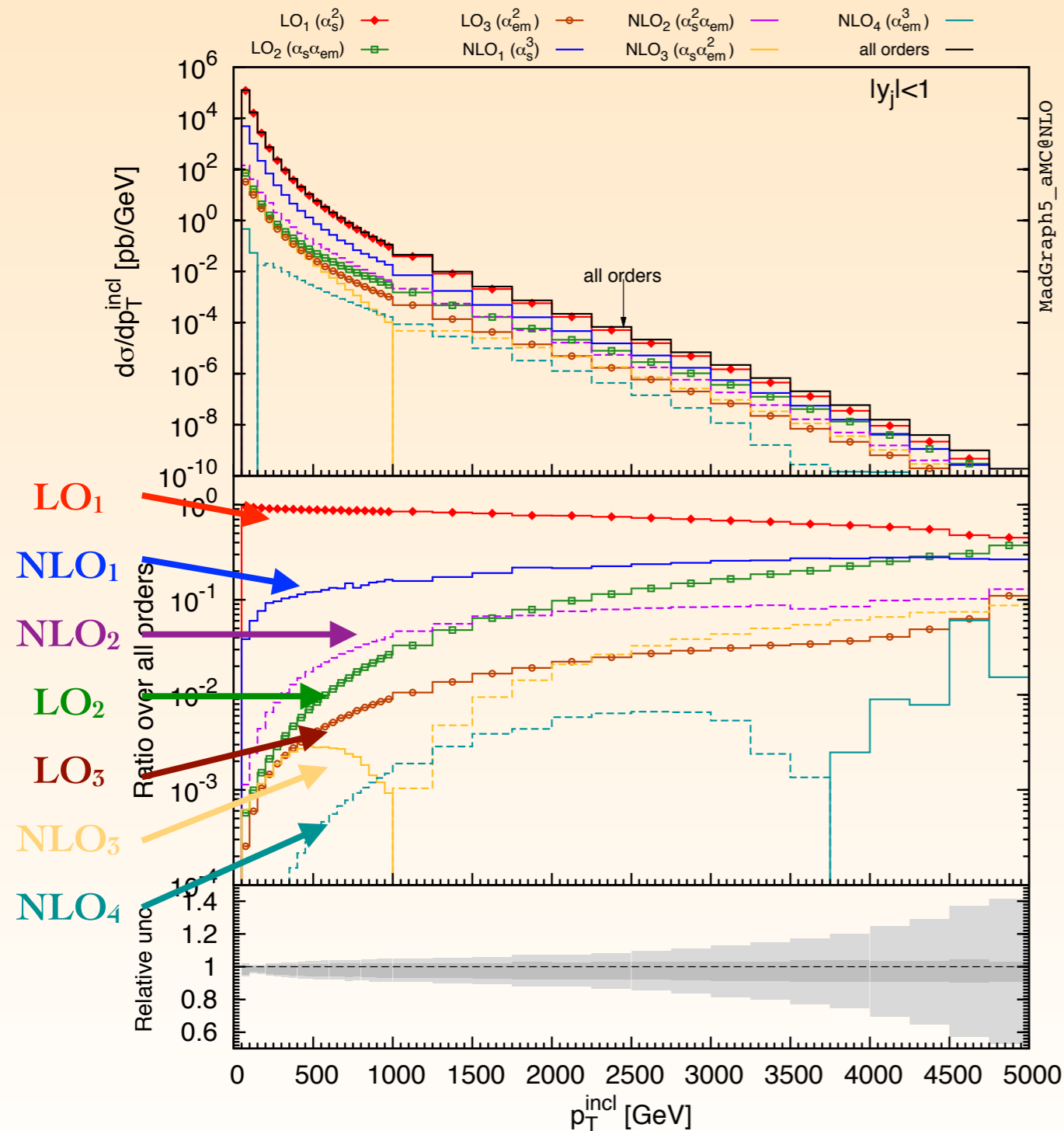


◆ Size of corrections mostly follows what one expects from the coupling combinations

○ Apart from the very far tail where NLO₃ is slightly larger than one would expect

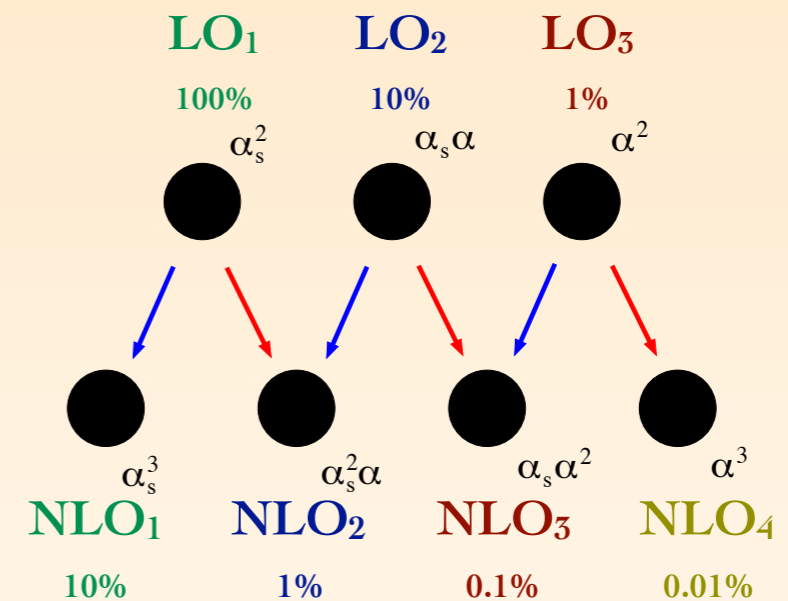
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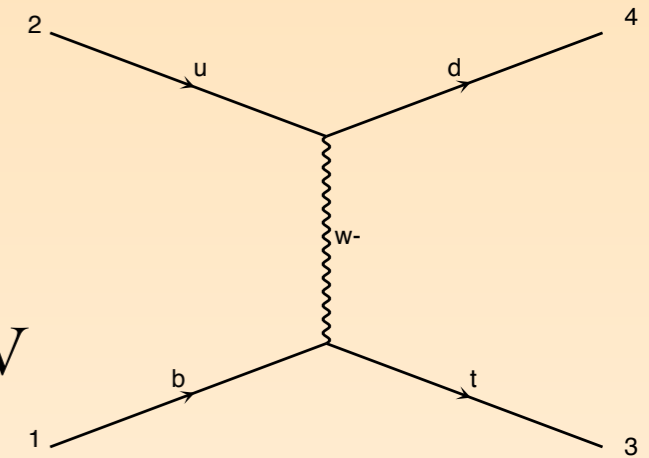


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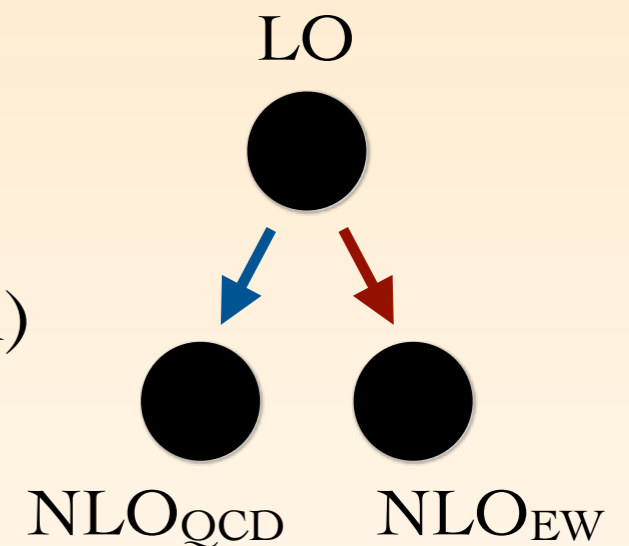
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SINGLE-TOP PRODUCTION

T- AND S-CHANNEL

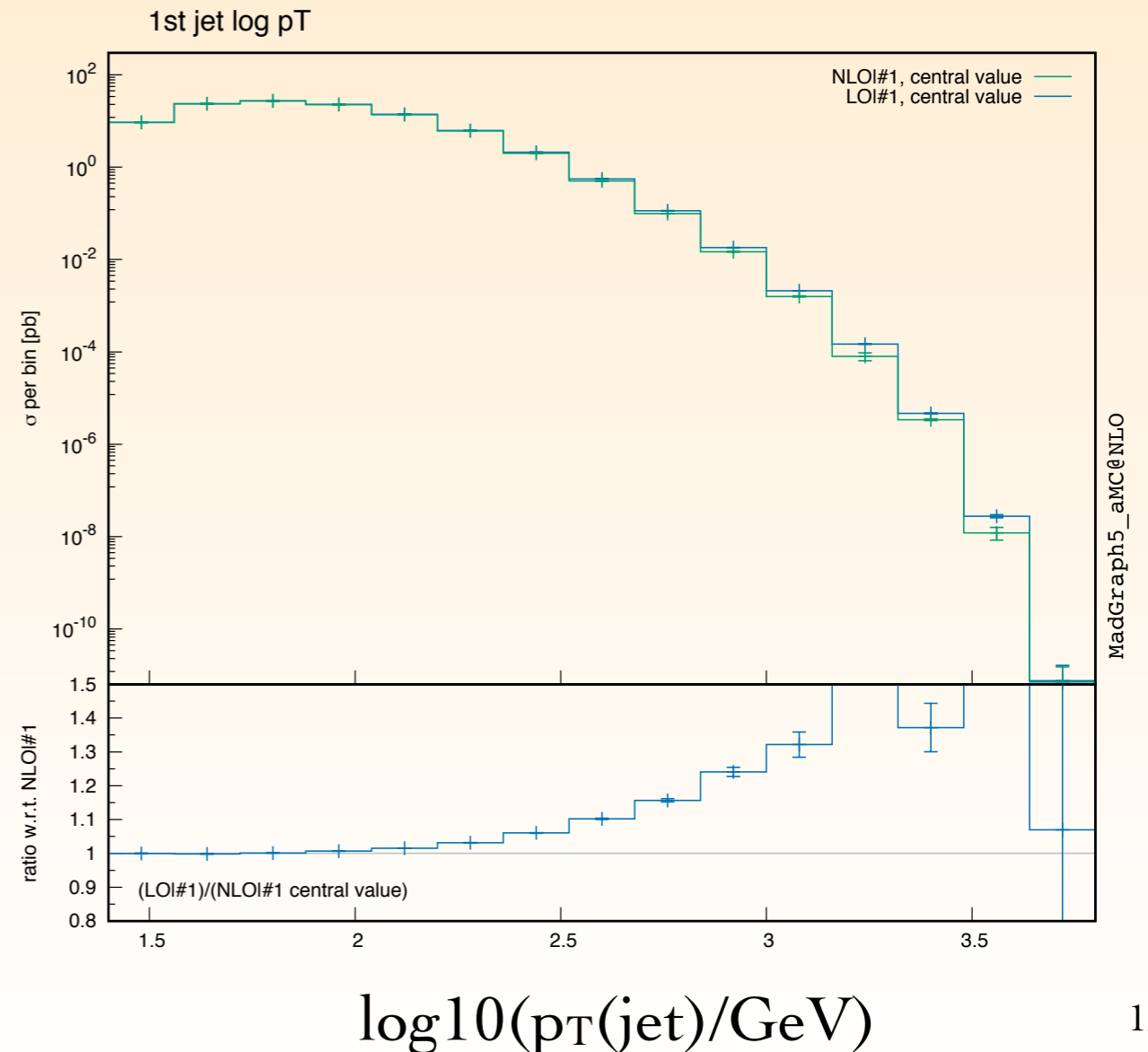
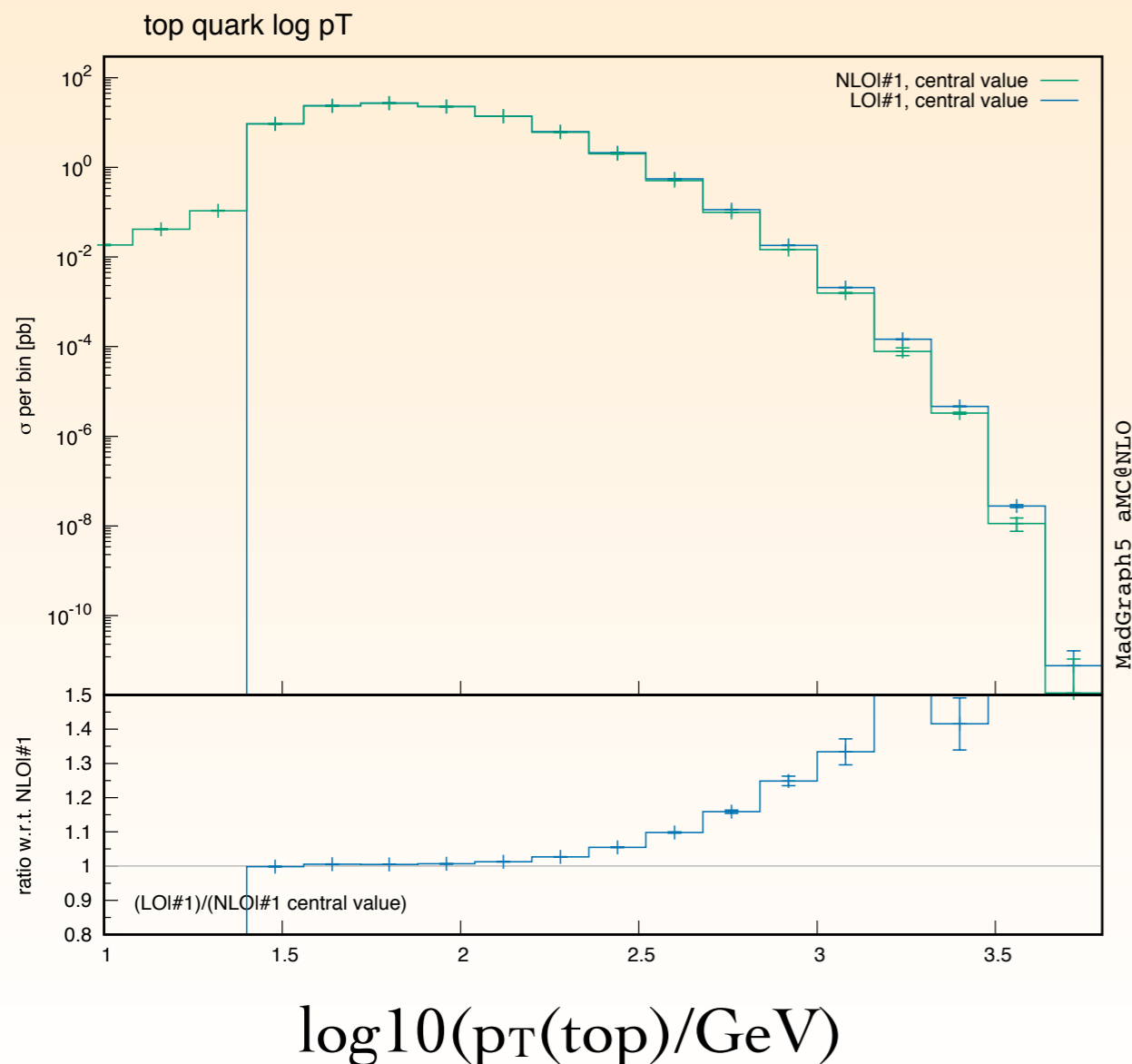


- ◆ Single-top production (with on-shell top quark) is a purely EW process. Hence, no difficulties in defining NLO QCD & EW
- ◆ However, t- and s-channel differentiation needs to be revisited
 - At NLO_{EW} , Initial state photon results in diagrams that contain both an t-channel and an s-channel W-boson (but one can probably still use parton flavours for differentiation)
- ◆ In the next results, no attempt in updating the differentiation will be made. We will only consider the sum
 - If necessary, one could always subtract the s-channel contribution at LO to obtain an NLO t-channel prediction
- ◆ NLO EW corrections for single-top production first studied by M. Beccaria et al. (2006), Mirabella (2008) and Bardin et al. (2011).



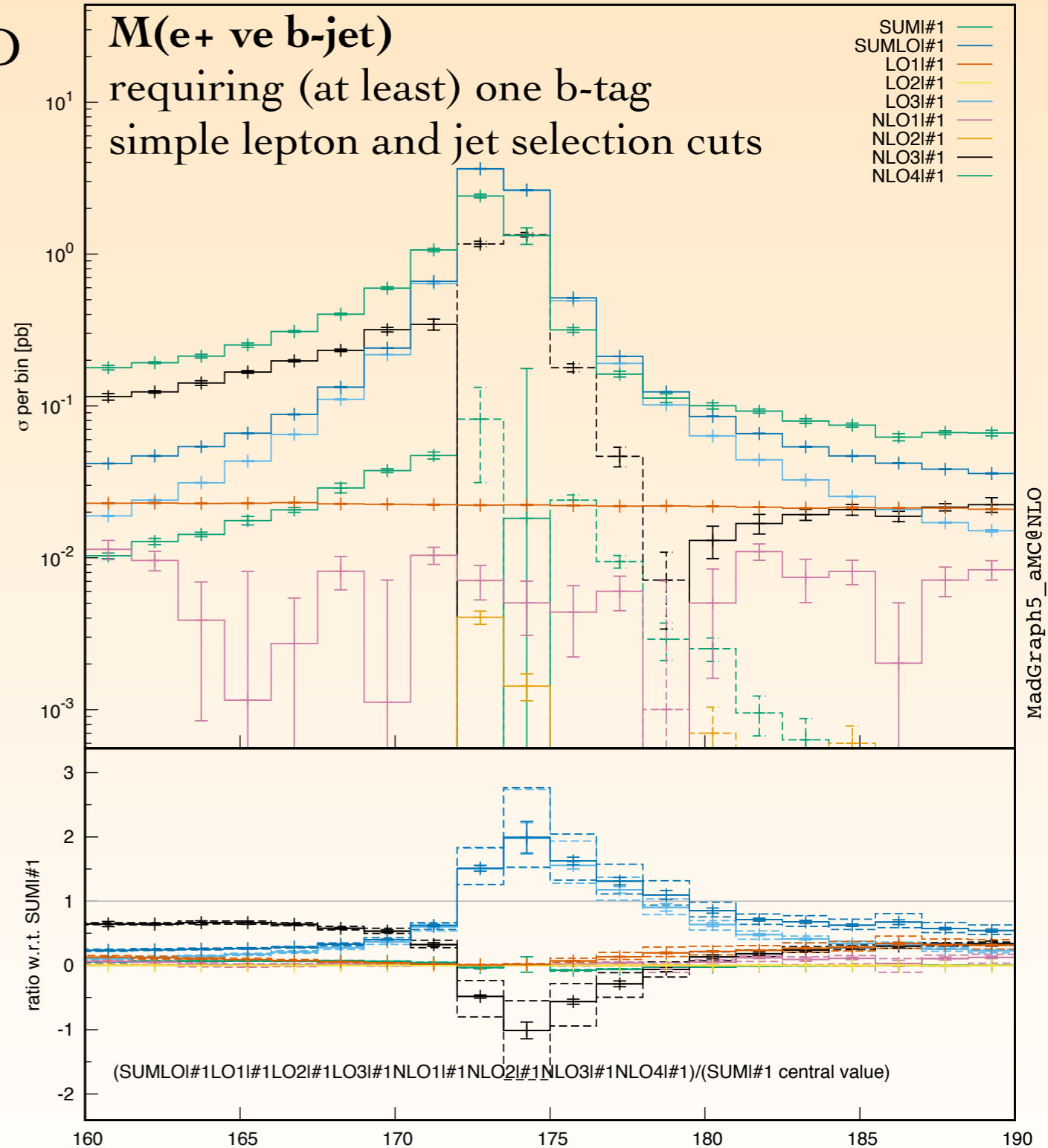
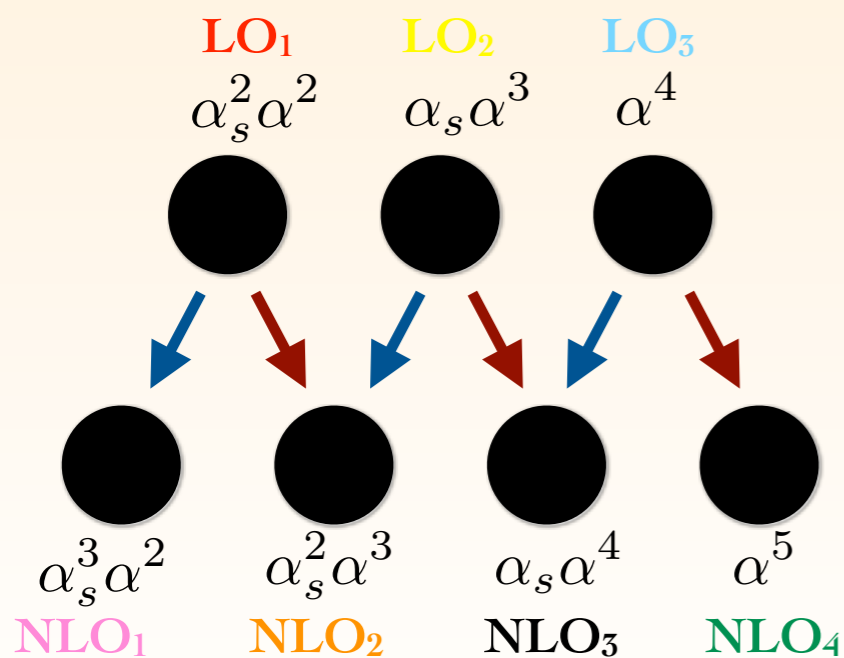
INCLUSIVE RATES

- ◆ For inclusive rates, the contributions from NLO EW corrections are small (less than a percent)
- ◆ This does no longer hold for the (extreme) tails of distributions, where the corrections can reach tens of percents



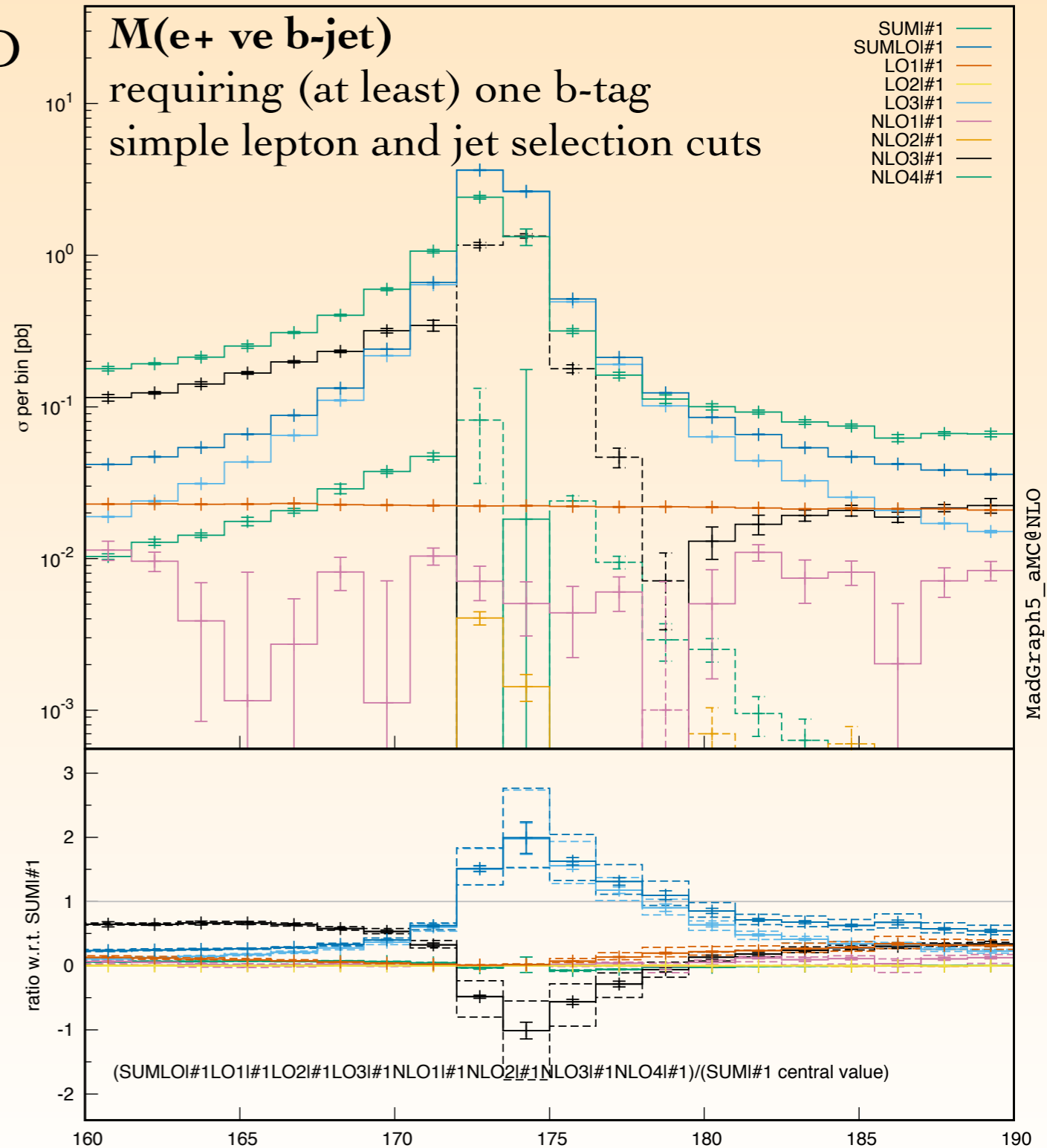
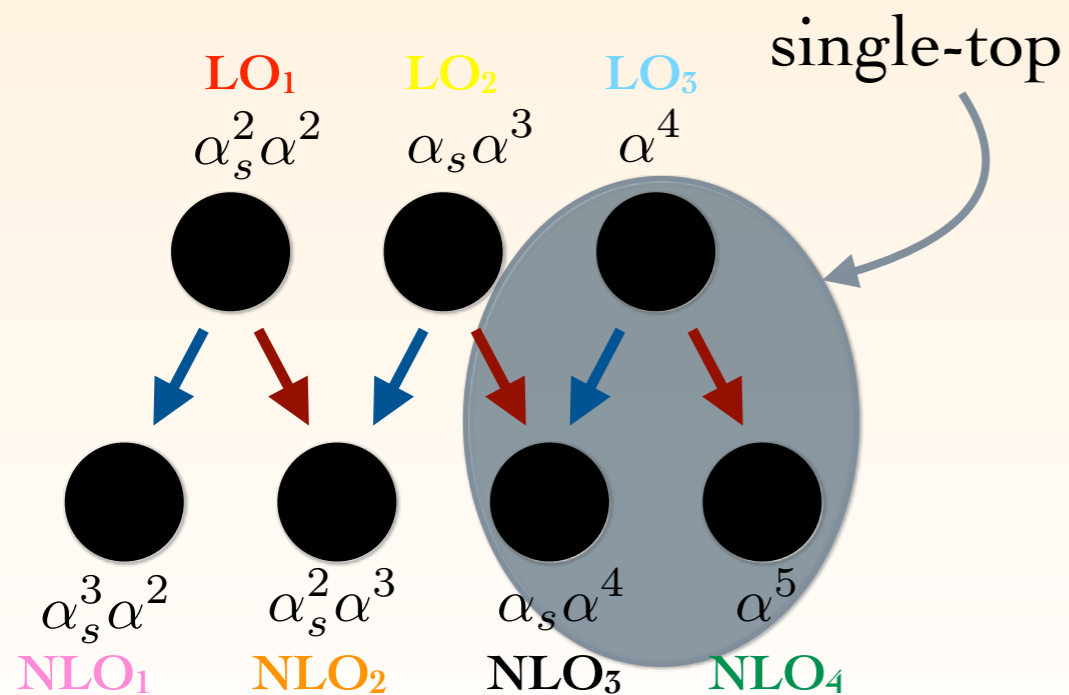
OFF-SHELL EFFECTS

- ◆ Generate the process at complete-NLO
 $pp \rightarrow e^+ \nu_e jj$
- ◆ This includes single-top production, but also background processes, with possible interferences
- ◆ Straightforward to generate, but difficult to interpret, assess uncertainties and to make use of



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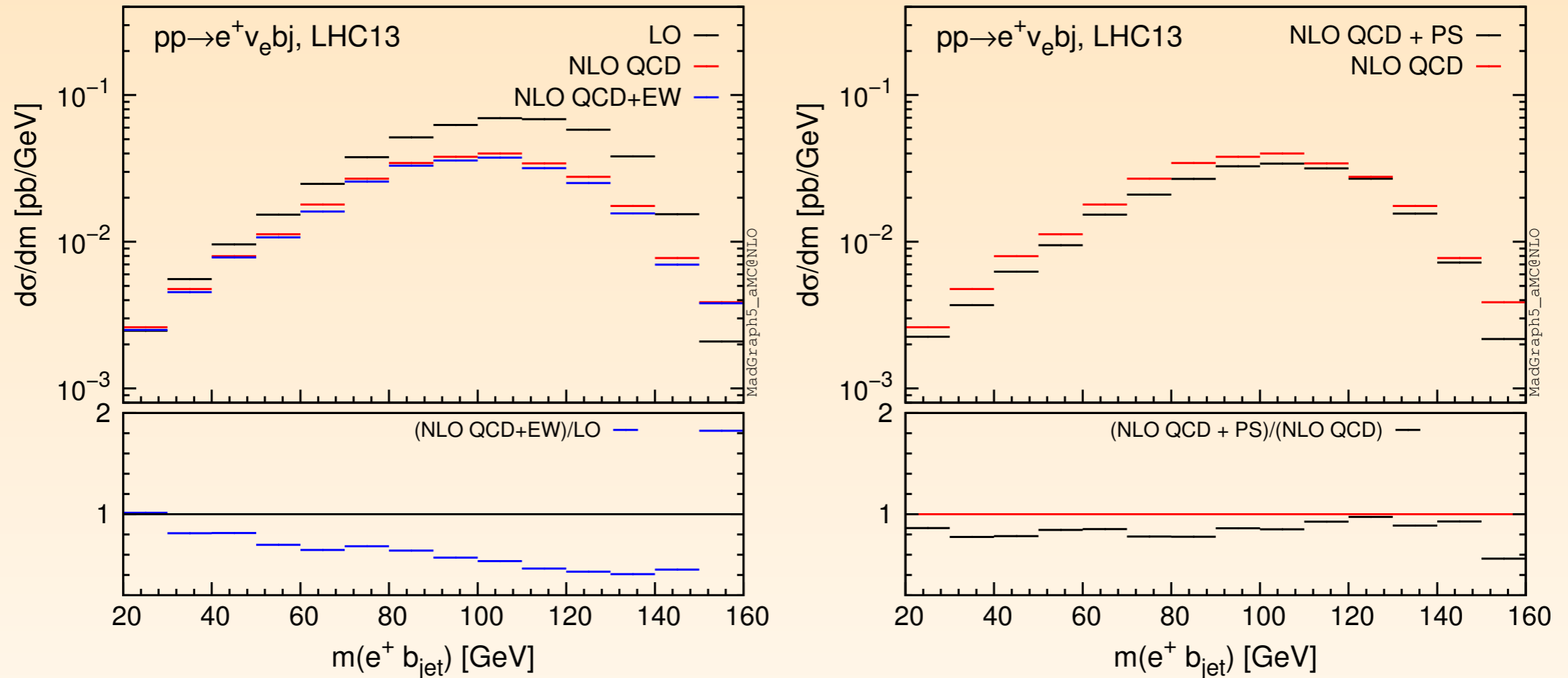


JET VETO ENHANCES CORRECTIONS

- ◆ Let's ignore possible interferences and focus again on single-top as signal: $LO_3 + NLO_3$ & NLO_4
- ◆ To enhance single-top signal, typically (b-)jet-veto is applied
 - require exactly one lepton, one b-jet, and one additional non-b-tagged jet
- ◆ In particular, the jet vetos enhance the effects from NLO corrections enormously
- ◆ Including higher orders does not solve the problem. Also at NNLO QCD ([Berger et al.]) the corrections remain large
- ◆ Resummation through parton shower improves the situation considerably, however not available for the EW corrections

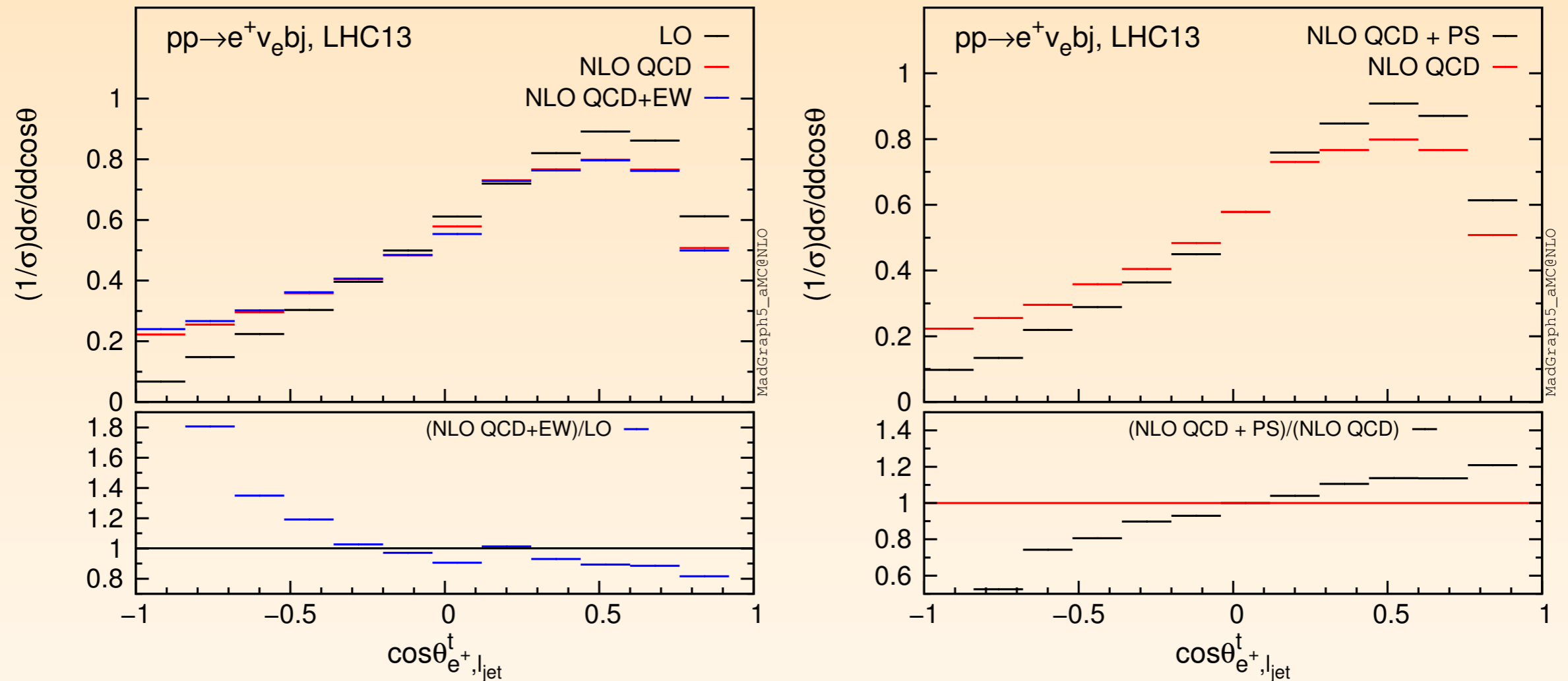
$pp \rightarrow e^+ v_e b j$, PDFs=LUXQED17 (82200)	
$\mu_f^0 = \mu_r^0 = H_T/2$	
Order	σ [fb]
LO QCD	4.616(4) $^{+0.415(+9.0\%)}_{-0.532(-11.5\%)}$
NLO QCD	2.75(3) $^{+0.22(+8.2\%)}_{-0.24(-8.8\%)}$
NLO QCD+EW	2.57(3) $^{+0.22(+8.5\%)}_{-0.25(-9.6\%)}$
LO QCD + PS	3.038(6) $^{+0.280(+9.2\%)}_{-0.357(-11.7\%)}$
NLO QCD + PS	2.36(2) $^{+0.12(+5.0\%)}_{-0.10(-4.0\%)}$

DIFFERENTIAL DISTRIBUTIONS 1



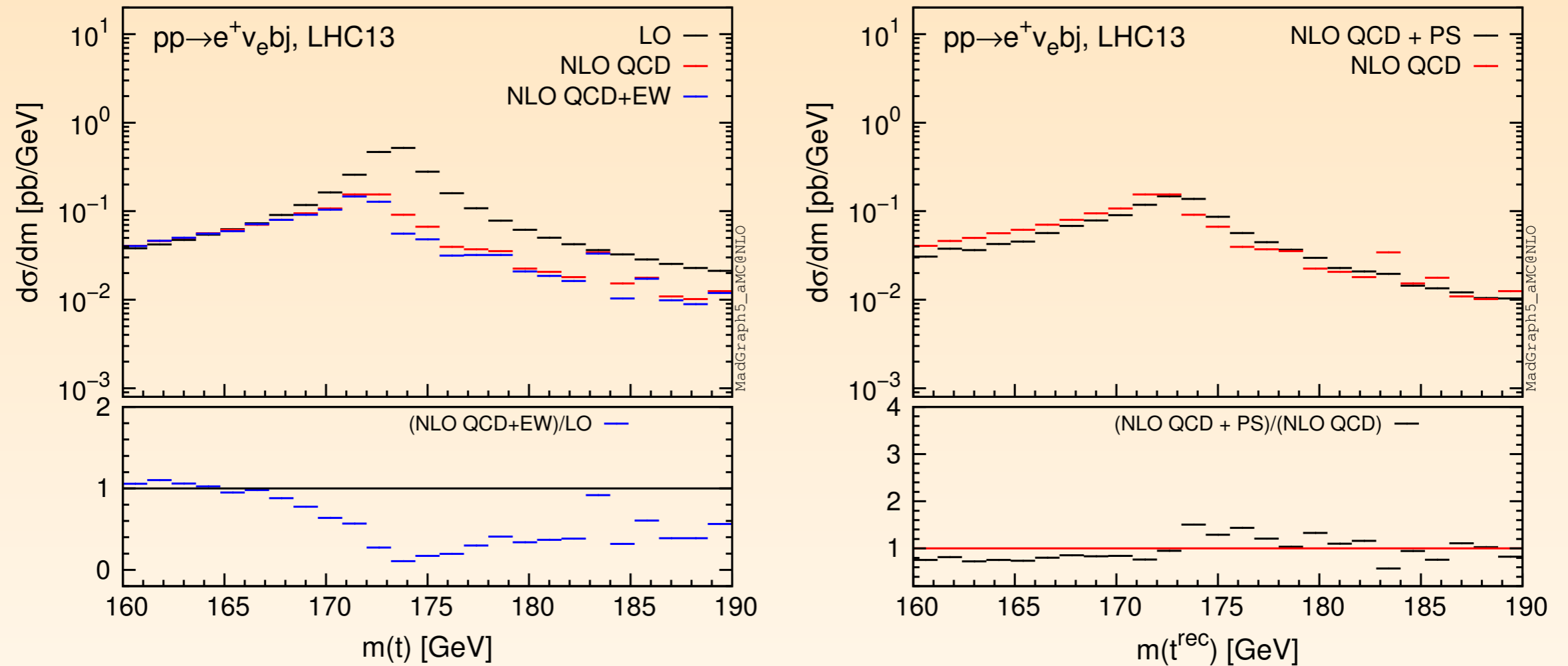
- ◆ Lepton + b-jet invariant mass
- ◆ *left*: Fixed order comparison; *right*: NLO vs NLO+PS (with QCD corrections)
- ◆ EW corrections small compared to other effects

DIFFERENTIAL DISTRIBUTIONS 2



- ◆ Angle between lepton, in the top rest-frame, and light jet: very sensitive to spin correlations
- ◆ Effects from parton shower again larger than from EW corrections

DIFFERENTIAL DISTRIBUTIONS 3



- ◆ Reconstructed to quark mass from lepton, b-jet and missing energy, using W -boson mass constraint
- ◆ EW corrections are of similar size as compared to effects from parton shower

CONCLUSIONS

- ◆ **NLO EW** corrections are a part of a family of NLO corrections due to the mixed coupling expansion of the perturbative series (**complete-NLO**)
- ◆ Automation of **complete-NLO** for all* relevant SM processes (e.g. in MadGraph5_aMC@NLO v3_beta)
- ◆ Not covered: beyond **NLO_{QCD}** the distinction between jets, photons and leptons becomes non-trivial without fragmentation functions (work in progress)
- ◆ Work-in-progress: consistent matching to parton showers when including **NLO_{EW}** corrections
- ◆ EW corrections to single-top production are small, but enhanced in tails of distributions. Also applying a jet-veto enhances the effects from higher-order corrections enormously, but here the EW corrections remain smaller than other effects