

Charmonia production with a density operator model

Hard Probes 2020

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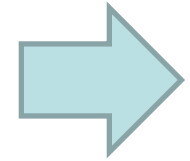
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with

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Motivation & Background



The two faces of quarkonia production

Sequential dissociation

Main goal: describe with good precision the dissociation / melting of some quarkonia wannabe $Q\bar{Q}$ pair : recent progresses due to Open Quantum System approach + improved knowledge of the dissociation rate



Q-Qbar **recombination**

How do we understand and describe the recombination of ≈ 100 c and 100 \bar{c} all together at LHC ?

Usual approaches:

- Detailed balance (chemistry)
- Statistical hadronization

More recently: J.P. Blaizot & M. A. Escobedo JHEP06 (2018) 034

Keywords:

- Interaction of heavy quarks with the bulk particles
- Expansion of the medium
- **Off-diagonal contributions**

Motivation & Background

Generic idea : describe charmonia (Ψ) production using density matrix

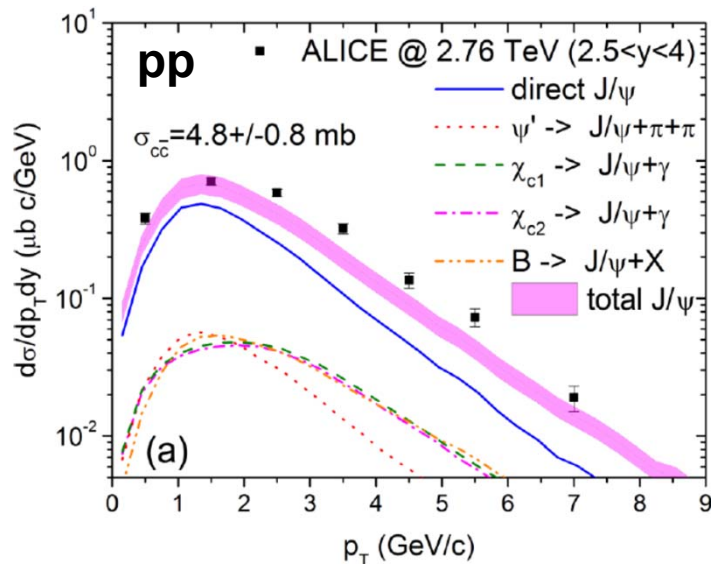
$$P^\Psi(t) = \text{Tr}[\rho_{QQ}^\Psi \rho_N(t)]$$

$$\rho_{QQ}^\Psi = \sum_i |\Psi_{QQ}^i\rangle \langle \Psi_{QQ}^i|$$

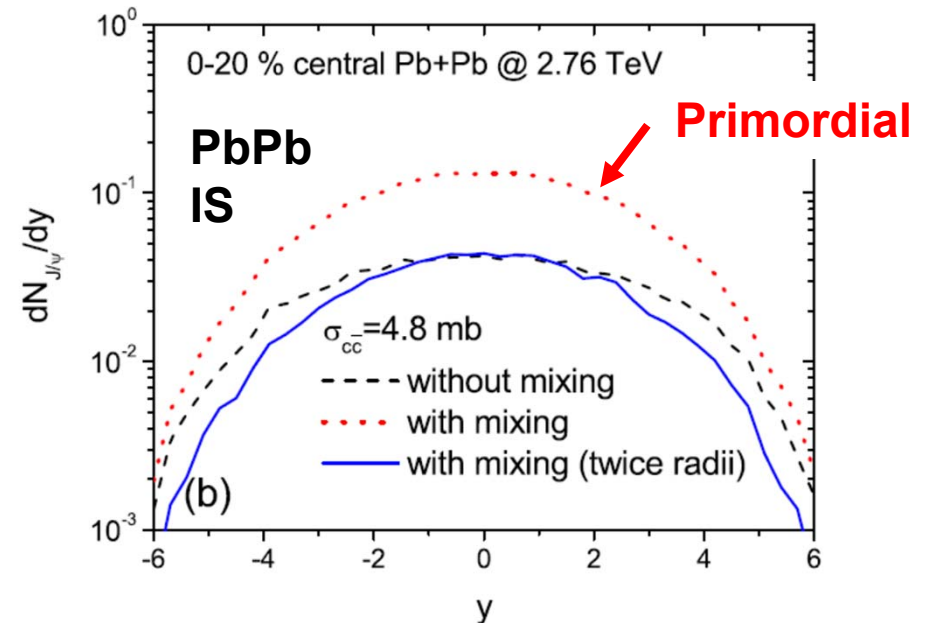
N-body density matrix (bulk partons + many c and many cbar)

Single quarkonia density operator

“Just” looking at the **initial stage** brings interesting features: Taesoo .S, J.Aichelin and E.Bratkovskaya , PRC 96. 014907 (2017)



Good reproduction of $pp \rightarrow J/\psi + x$!!!



considerable enhancement of primordial J/ψ (in the initial state): **large off-diagonal contributions**

Motivation & Background

$$P^\Psi(t) = \text{Tr}[\rho_{QQ}^\Psi \rho_N(t)]$$

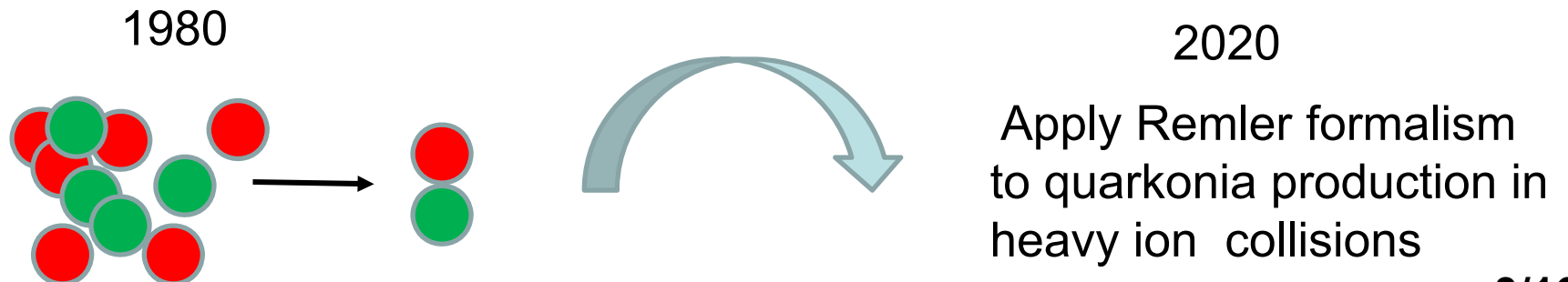
$\rho_{QQ}^\Psi = \sum_i |\Psi_{QQ}^i\rangle \langle \Psi_{QQ}^i|$

$\frac{\partial \rho_N(t)}{\partial t} = -i[H_N, \rho_N(t)]$

Dealing with the dynamics ?

- The idea of the formalism goes back to Remler's work
- General scheme connecting composite-particle cross section and rates with time-dependent density operators
- Applied by Remler et al to the deuteron production in (low energy) AA collisions The formalism is able to deal with many particles (nucleons \rightarrow deuterium)

E.A. Remler, ANNALS OF PHYSICS 136, 293-316 (1981)



Remler formalism at work

The effective rate for quarkonia state creation (dissociation) in the medium is

$$\Gamma^{\Psi}(t) = \frac{\partial P^{\Psi}(t)}{\partial t} = \text{Tr} \left[\rho_{Q\bar{Q}}^{\Psi} \frac{\partial \rho_N(t)}{\partial t} \right]$$

Working in the phase space through Wigner distribution

$$W^{\Psi} = \int d^3 y e^{i p y} \left\langle r - \frac{y}{2} \right| \Psi^i \rangle \left\langle \Psi^i \left| r + \frac{y}{2} \right. \right\rangle$$

Quarkonia: Double Gaussian approximation

$$W_{Q\bar{Q}}^{\Psi}(r_{rel}, p_{rel}) = C e^{r_{rel}^2 \sigma^2} e^{-\frac{p_{rel}^2}{\sigma^2}}$$

Parameter: The Gaussian width σ

$$\left[\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right] \Psi_{Q\bar{Q}}(r) = E_{Q\bar{Q}} \Psi_{Q\bar{Q}} \rightarrow \langle r^2 \rangle \rightarrow W^{\Psi}$$

W_N : Semi-classical approach

$$W_N = \prod_i \hbar^3 \delta(x_i - x_{i0}(t)) \delta(p_i - p_{i0}(t))$$

... but no explicit description of W_N required (as it appears in the trace)

and (less trivial) : generalisation at finite velocity; fully relativistic

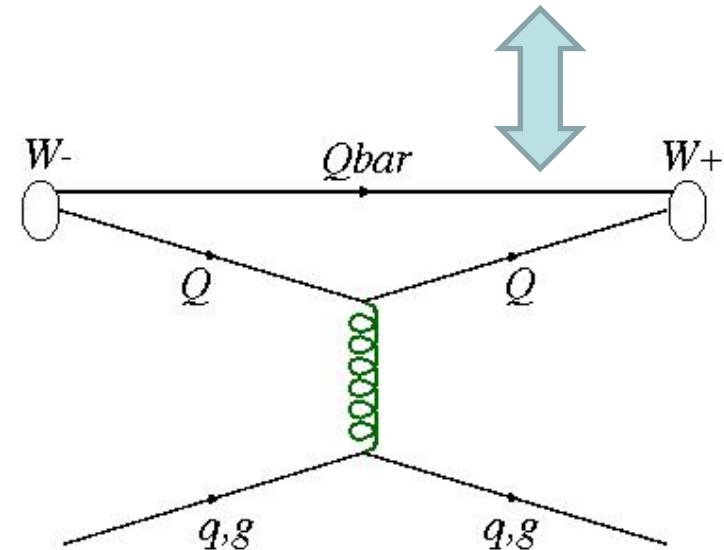
Remler formalism at work

Combining the expression of the Wigner's functions and substituting in the **effective rate equation** :

$$\Gamma(t) = \sum_{i=1,2} \sum_{j \geq 3} \delta(t - t_{ij}(v)) \int \frac{d^3 p_i d^3 x_i}{h^3} W_{QQ}^{\Psi}(p_1, x_1; p_2, x_2) [W_N(t+\epsilon) - W_N(t-\epsilon)]$$

- The quarkonia production in our model is a three body process, the HQ (anti-quark) interact only by collision !!!
- The “details” of H_{int} between HQ and bulk partons are incorporated into the evolution of W_N after each collision / time step (nice feature for the MC simulation)
- $W_N(t+\epsilon)$ and $W_N(t-\epsilon)$ are NOT the equivalent of gain and loss terms in usual rate equation
- Dissociation and recombination treated in the same scheme

Then:
$$P^{\Psi}(t) = P^{\Psi}(t_{\text{init}}^{\Psi}) + \int_{t_{\text{init}}}^t \Gamma(t) dt$$



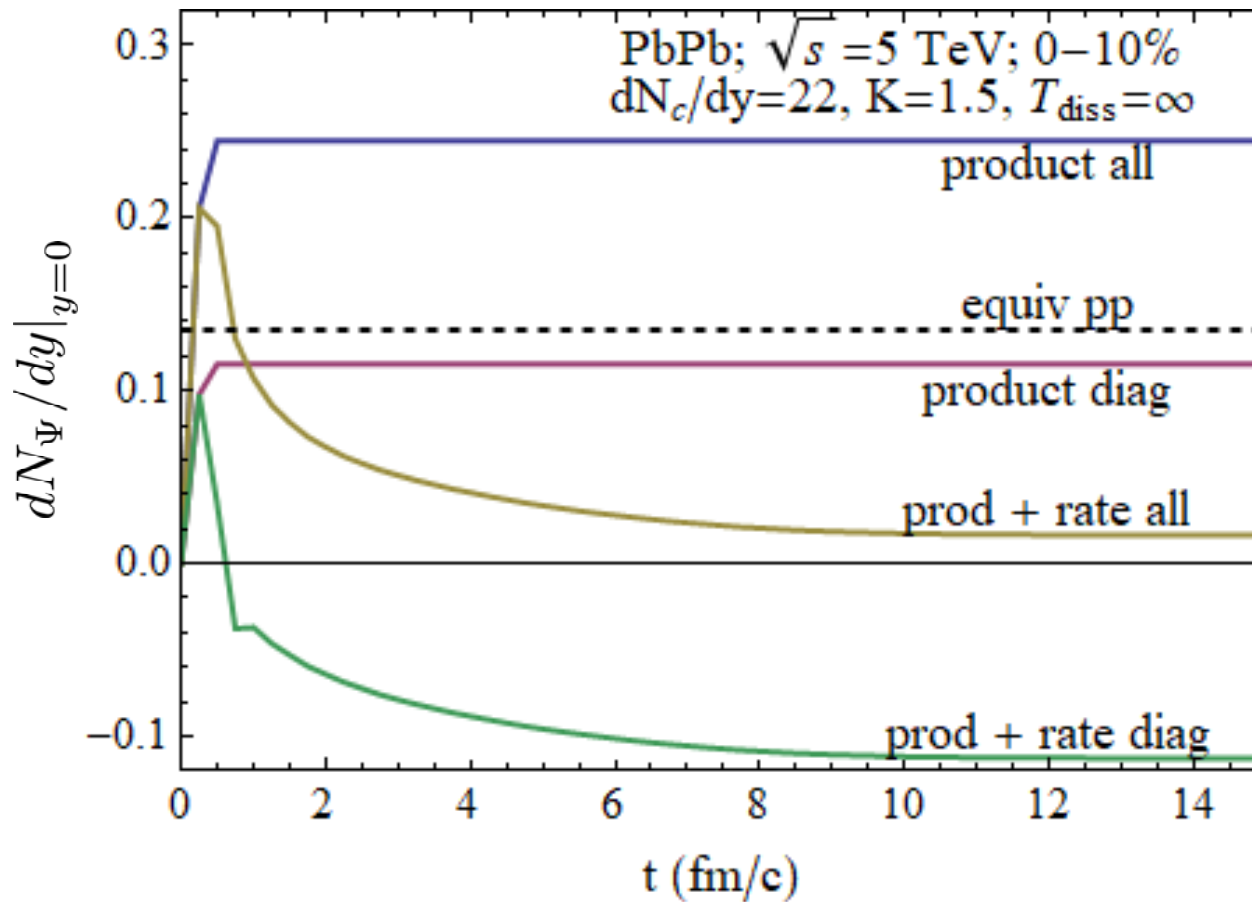
Interaction of HQ with the QGP are carried out by EPOSHQ (good results for D and B mesons production)

NB: Also possible to generate similar relations for differential rates

Preliminary results for J/ψ

Word of caution: Exploratory phase => not meant to have an exact comparison with exp. data

$$P^\Psi(t) = \boxed{P^\Psi(t_{init}^\Psi)} + \int_{t_{init}^\Psi}^t \Gamma(t) dt$$



Cumulated « production » (if no rate equation), indeed overshoots pp due to off-diagonal contributions

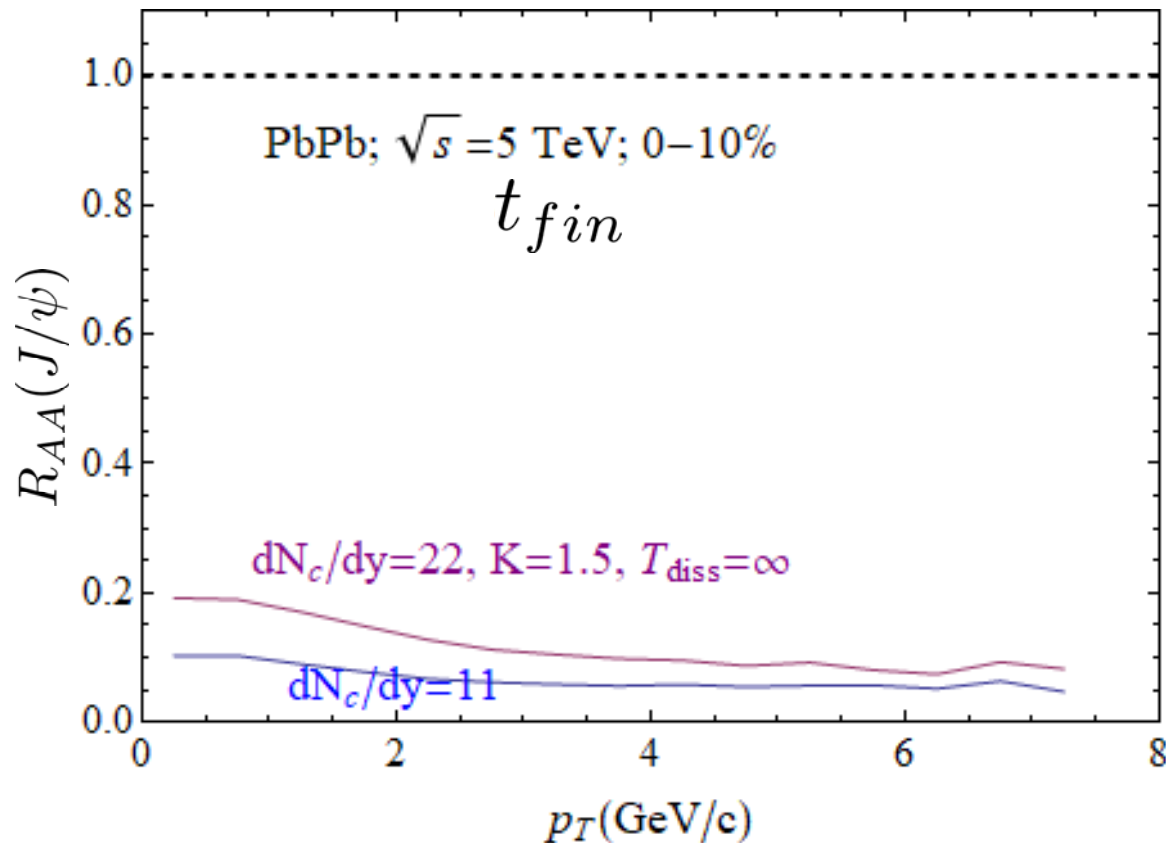
The denominator in the R_{AA}

The full production (i.e. the numerator in the R_{AA})

First answer to puzzle found in Song et al: the primordial production is killed rather fast by the « loss » rate.

Preliminary results for J/ψ

Word of caution: Exploratory phase => not meant to have an exact comparison with exp. data

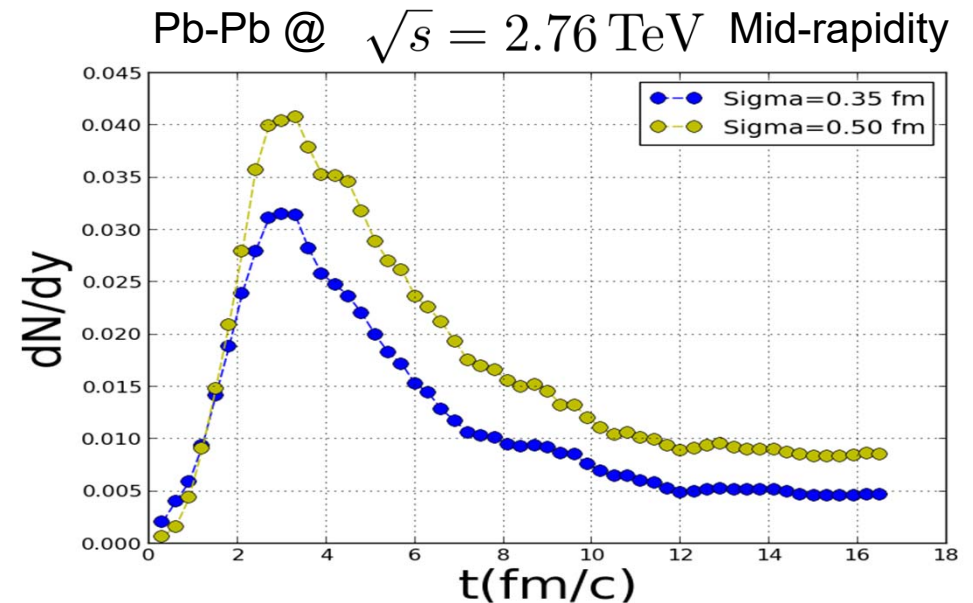
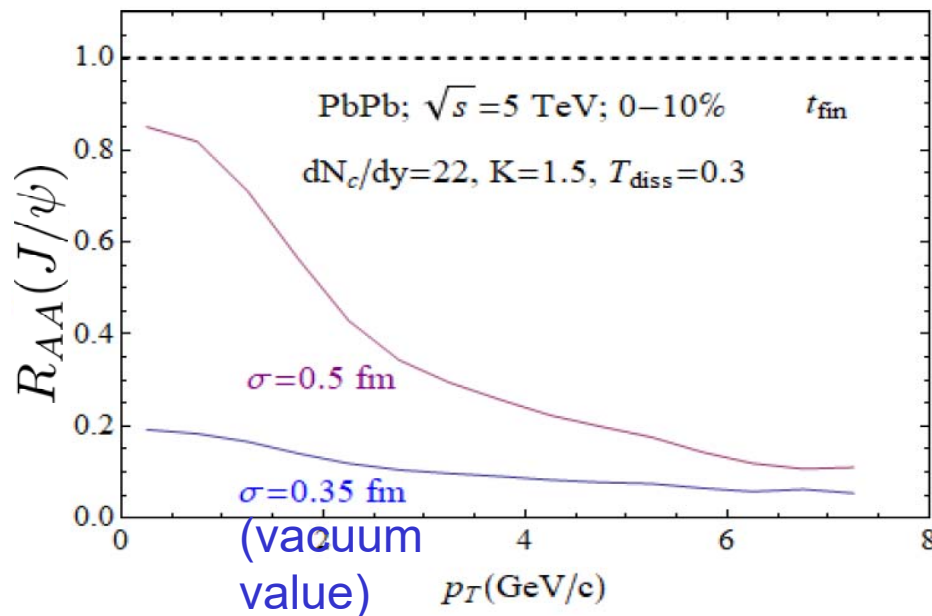


Effect of charm abundance in phase space (x2):

- Correct trends for charm recombination
- Absolute value too small

Preliminary results for J/ψ

Exploring the effect of « in medium » Wigner distribution through the variation of the σ parameter (spatial width):

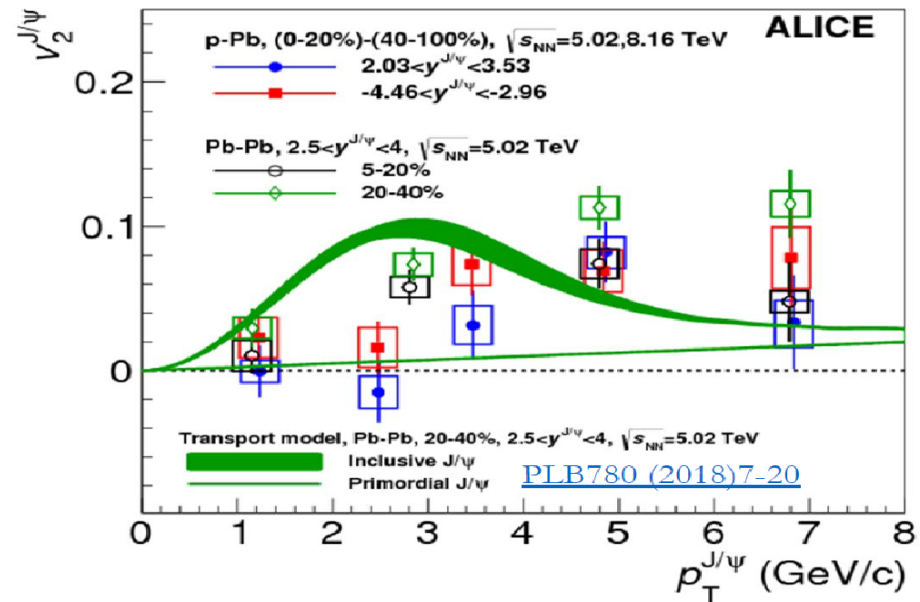
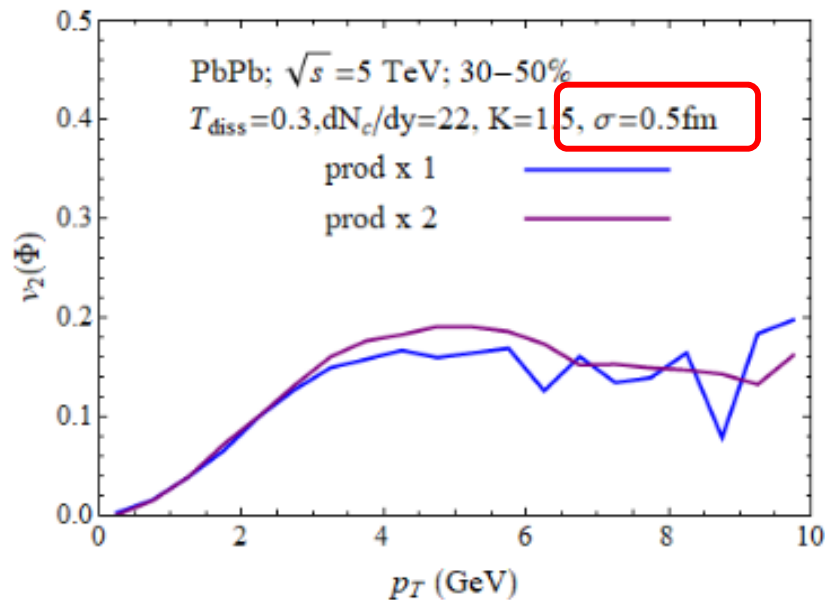


Missing ingredient for quantitative agreement:

- temperature dependent $\sigma(T)$ in the Wigner function ?
- Interactions between Q & \bar{Q} (real part of the potential, not implemented in EPOSHQ)
- ... ?

Preliminary results for J/ψ

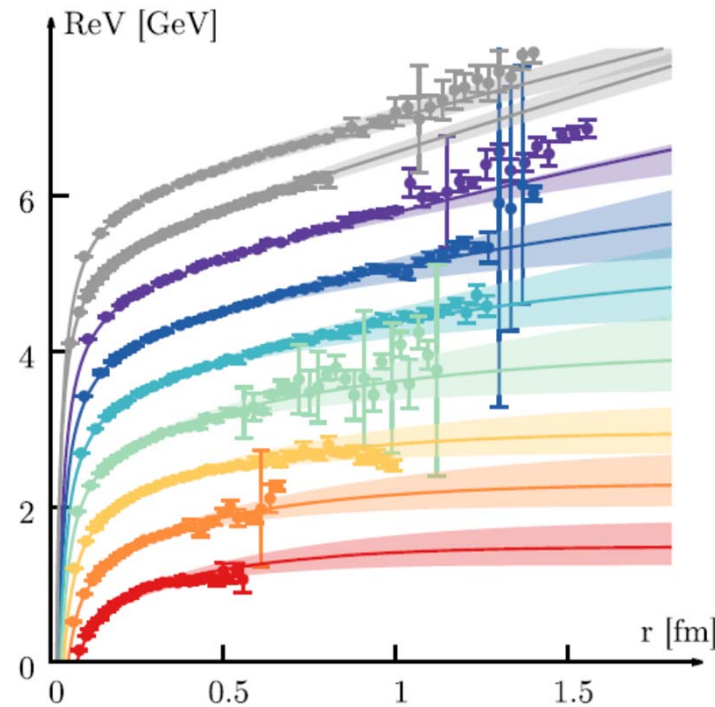
Despite quantitative underprediction in the R_{AA} ... taking the most appropriate parameter:



Interesting feature: Within our approach, v_2 extends at rather large p_T

The Q-Qbar interaction

- Not implemented up to now in EPOSHQ
- More and more reliable calculations are becoming available for the real part of the potential, thanks to lattice calculations:

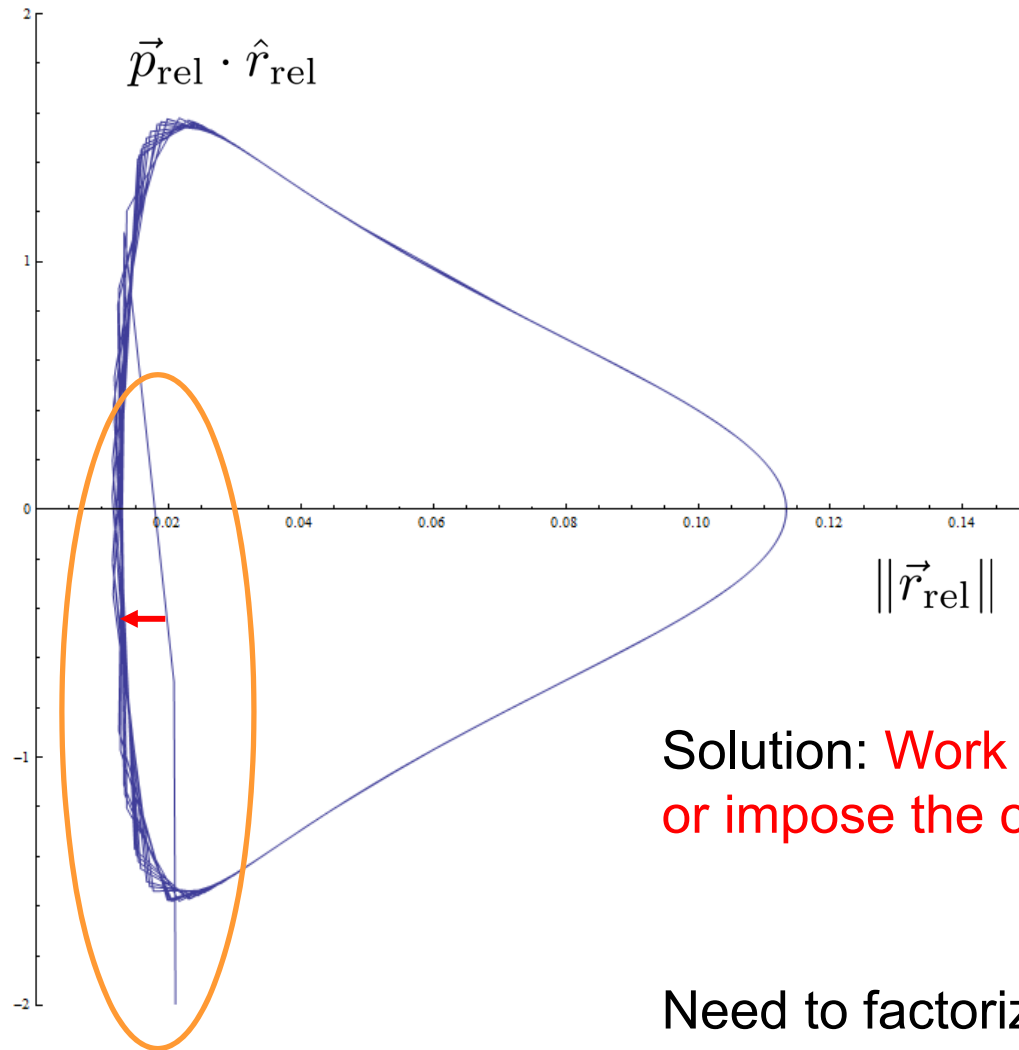


D. Lafferty and A. Rothkopf,
PHYS. REV. D 101, 056010 (2020)

- Go for it !

The Q-Qbar interaction: strategy 1

- “Minor problem” #1: Classical equations of motion are **unstable** (in the CM):



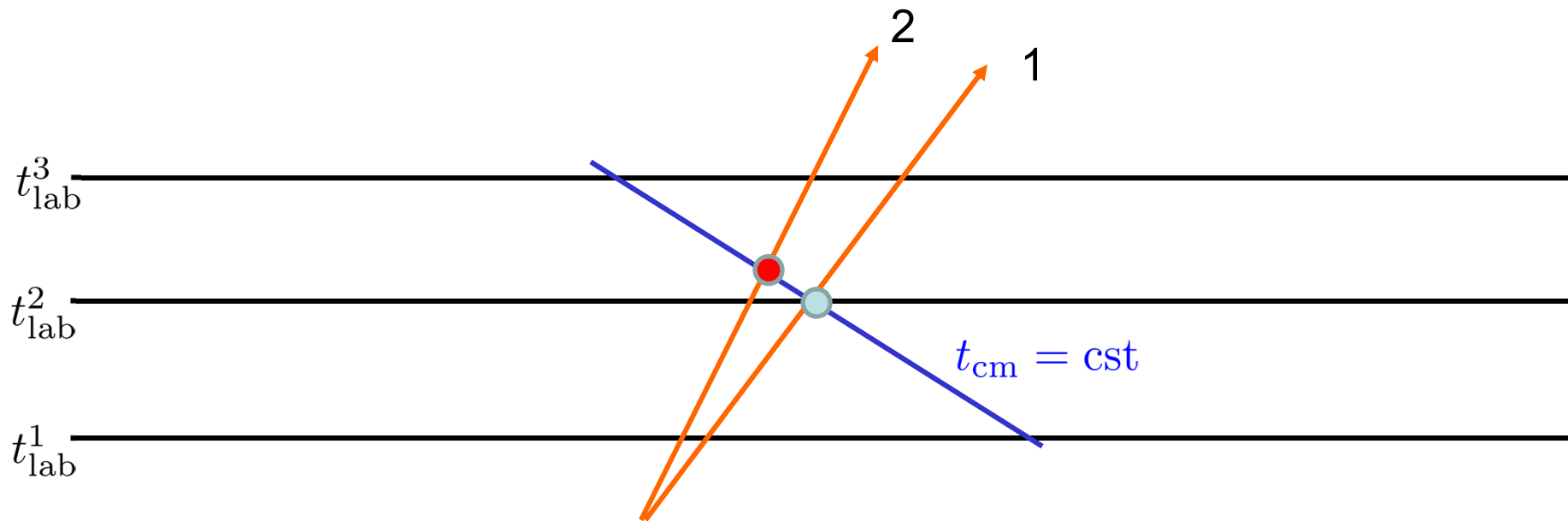
Solution: **Work in Hamilton – Jacobi coordinates or impose the conserved quantities (L and Etot)**



Need to factorize the N-body problem as an $\{$ of 2-body problems

The Q-Qbar interaction: strategy 1

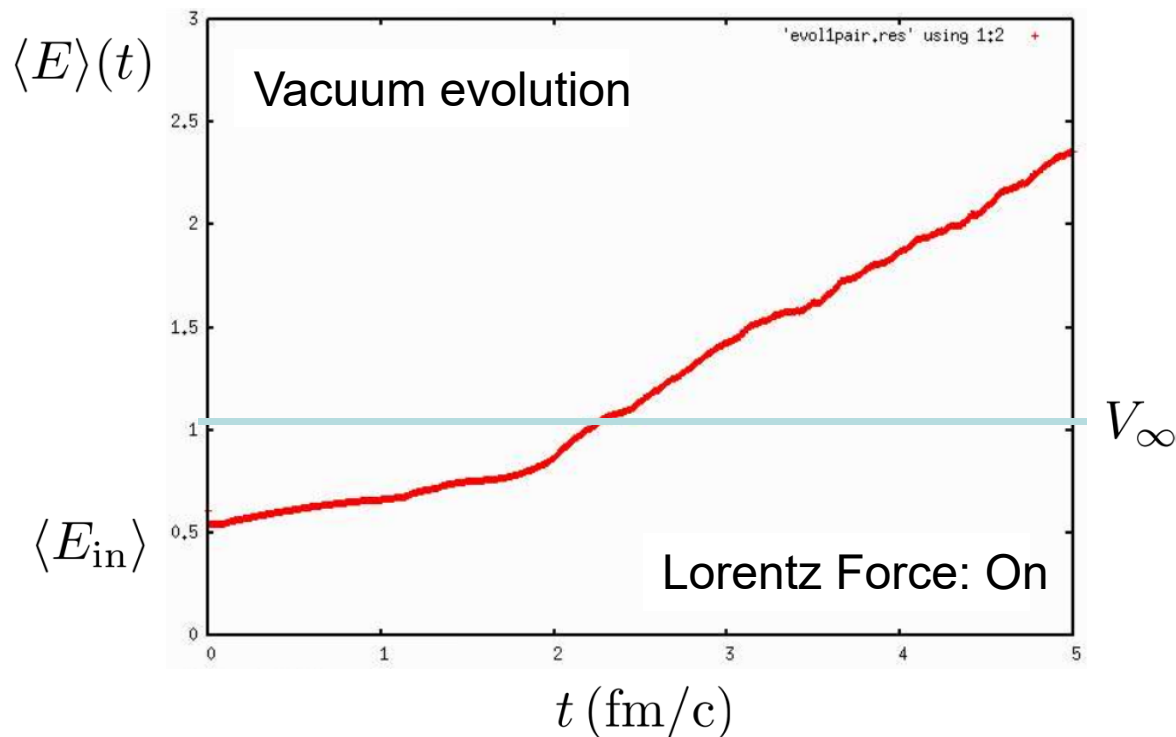
- “Major problem” : slicing the global time evolution (usual strategy in MC) is not compatible with passing to c.m. frame for **each individual pair**...



Generic need to store / describe the trajectory of particle 2 at a time $t_{\text{lab}} > t_{\text{lab}}^2$ if one propagates particle 1 up to t_{lab}^2 by resorting to evolution in the c.m.

The Q-Qbar interaction: strategy 2

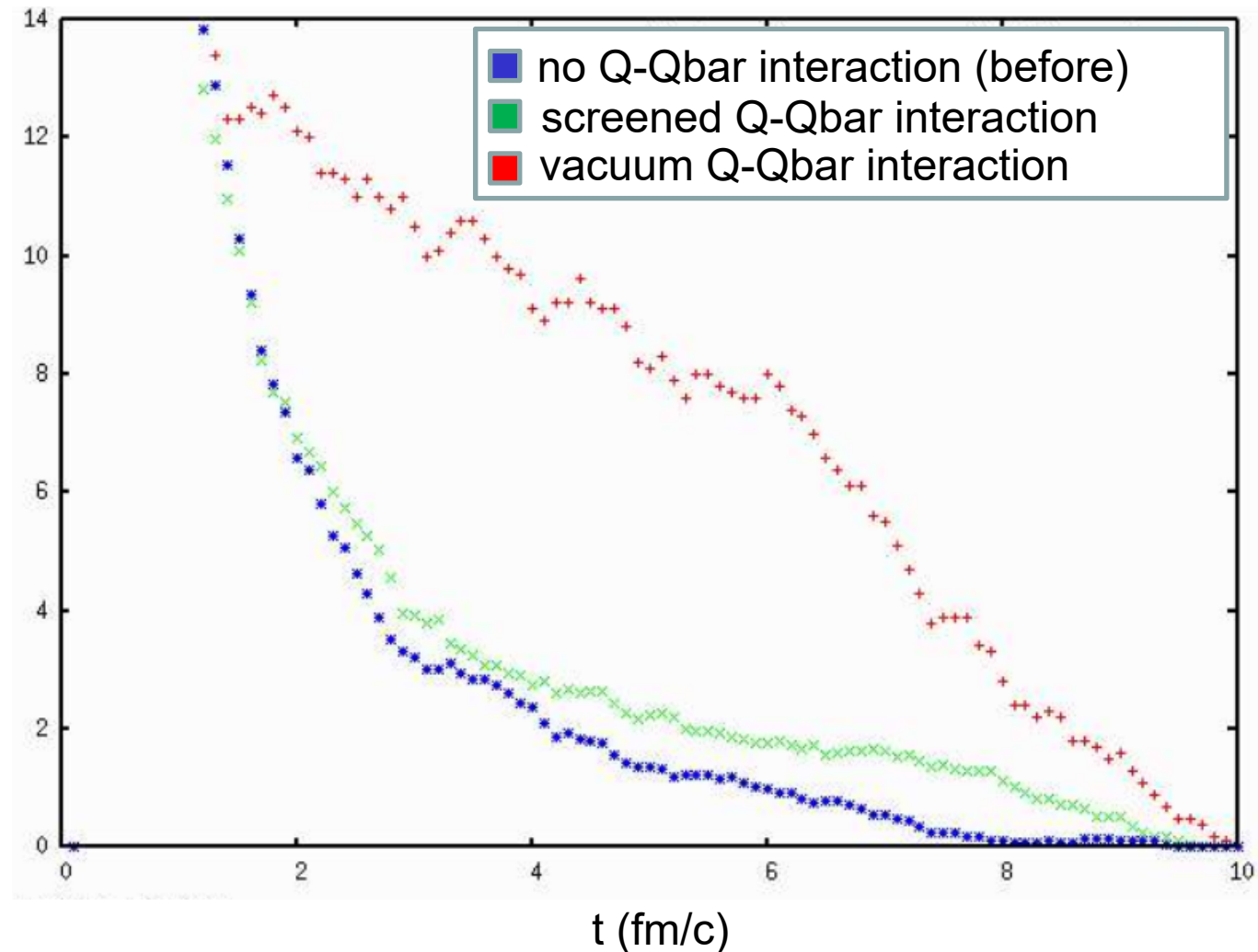
- Describe dynamics through retarded interactions... calibrated to map to the static potential in the infinite mass static case... obviously several prescriptions available, need discussion with IQCD experts !
- Cures all problems encountered with strategy 1 ☺ ... but (to my knowledge): No invariant associated with the retarded force (even without considering Abraham Lorentz force) => need specific methods to tame the instabilities



Stability achieved for a couple of cycles

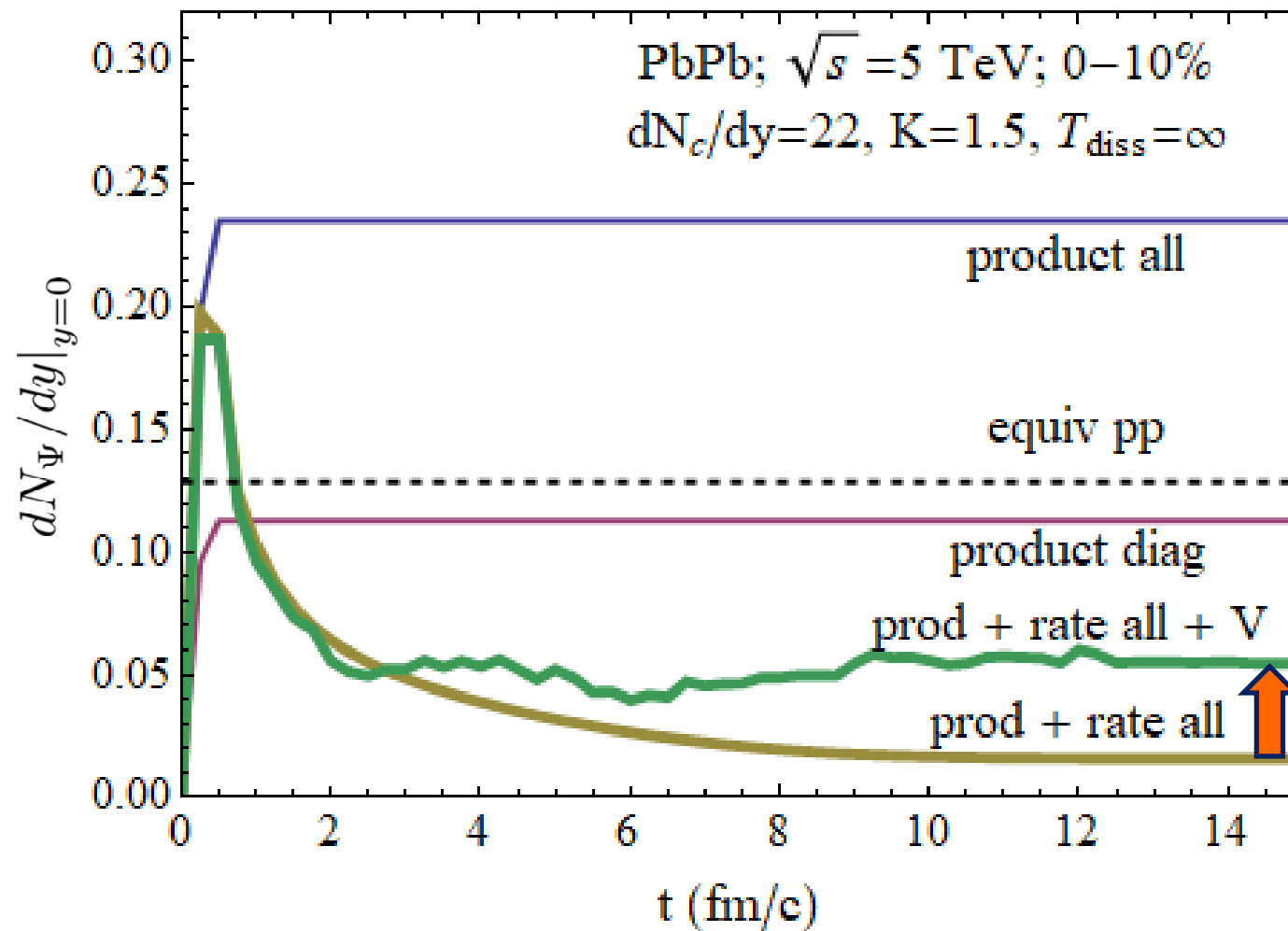
Consequences on the charmonia in AA collisions

Instantaneous # of
Q-Qbar at (invariant)
distance $< 1\text{fm}$



Although screened, the Q-Qbar interaction has important consequences on the probability to find a Q-Qbar at close distance in the final stage of the evolution

Consequences on the charmonia in AA collisions



Although screened, the **Q-Qbar** interaction has important consequences on the probability to find a J/ψ in the final stage of the evolution

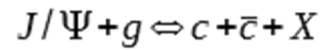
Conclusions

- We have presented a model based on the probability density operator that allows to obtain the evolution of the effective rate of Quarkonia production in QGP.
- This scheme treats both the dissociation and the recombination consistently
- We have implemented a first version of the Q-Qbar interaction in the EPOSHQ scheme; effect on charmonia production seems to be large, even with a screened potential
- Some further developments are still nevertheless mandatory.

Further remark : first Semi-Classical simulation of multiple c-cbar pairs in QGP: C Young et E Shuryak Phys.Rev.C81:034905 (2010)

Back up

The usual detailed balance relation



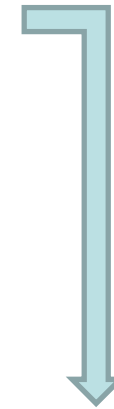
Rates Equation



$$\frac{dN_{\Psi}}{d\tau} = \Gamma_{recomb} N_c N_{\bar{c}} [V_{FB}(\tau)]^{-1} - \Gamma_{diss} N_{\Psi}$$



$$\frac{dN_{\Psi}}{d\tau} = -N_{\Psi} L \tau + G(\Psi) = \frac{-1}{\tau_{\Psi}} [N_{\Psi}(\tau) - N_{\Psi}^{eq}(\tau)]$$



Remler

$$\Gamma(t) = \sum_{i=1,2} \sum_{j \geq 3} \delta(t - t_{ij}(v)) \int \frac{d^3 p_i d^3 x_i}{h^3} W_{QQ}^{\Psi}(p_1, x_1; p_2, x_2) [W_N(t+\epsilon) - W_N(t-\epsilon)]$$

No one to one correspondance of the gain
& loss term !