

The CLIC Physics Potential for long-lived and exotic signatures

A summary of results from the Yellow Report “CLIC physics potential for new physics”

JAN. 23 2019

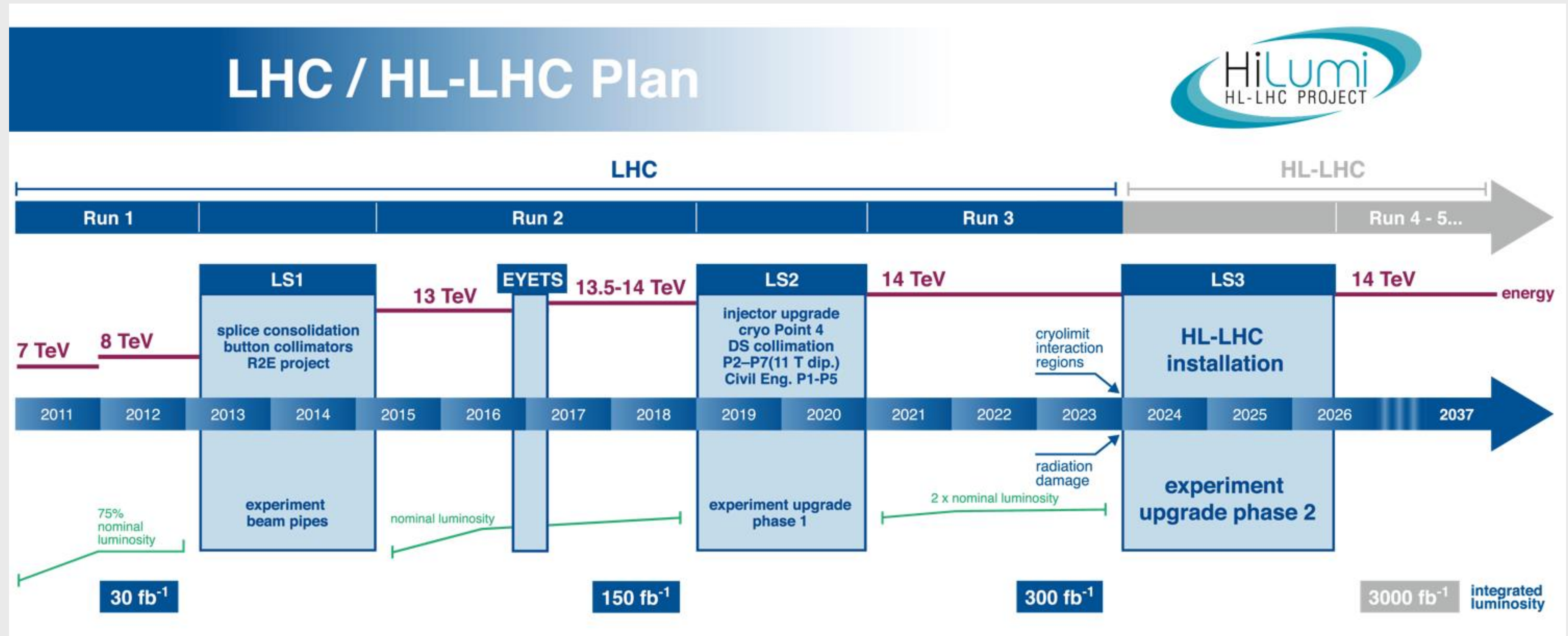
ROBERTO FRANCESCHINI (ROMA 3 UNIVERSITY)



CLIC Week 2019 - <https://indico.cern.ch/event/753671/>

Still looking for new physics

SUBTLE IS THE LORD



LHC: New Physics is either heavy or shows up in subtle ways

↑ ICHEP14

↑ ICHEP16

↑ ICHEP18

LHC: New Physics is either even heavier or shows up in even subtler ways

No guarantee for discovery

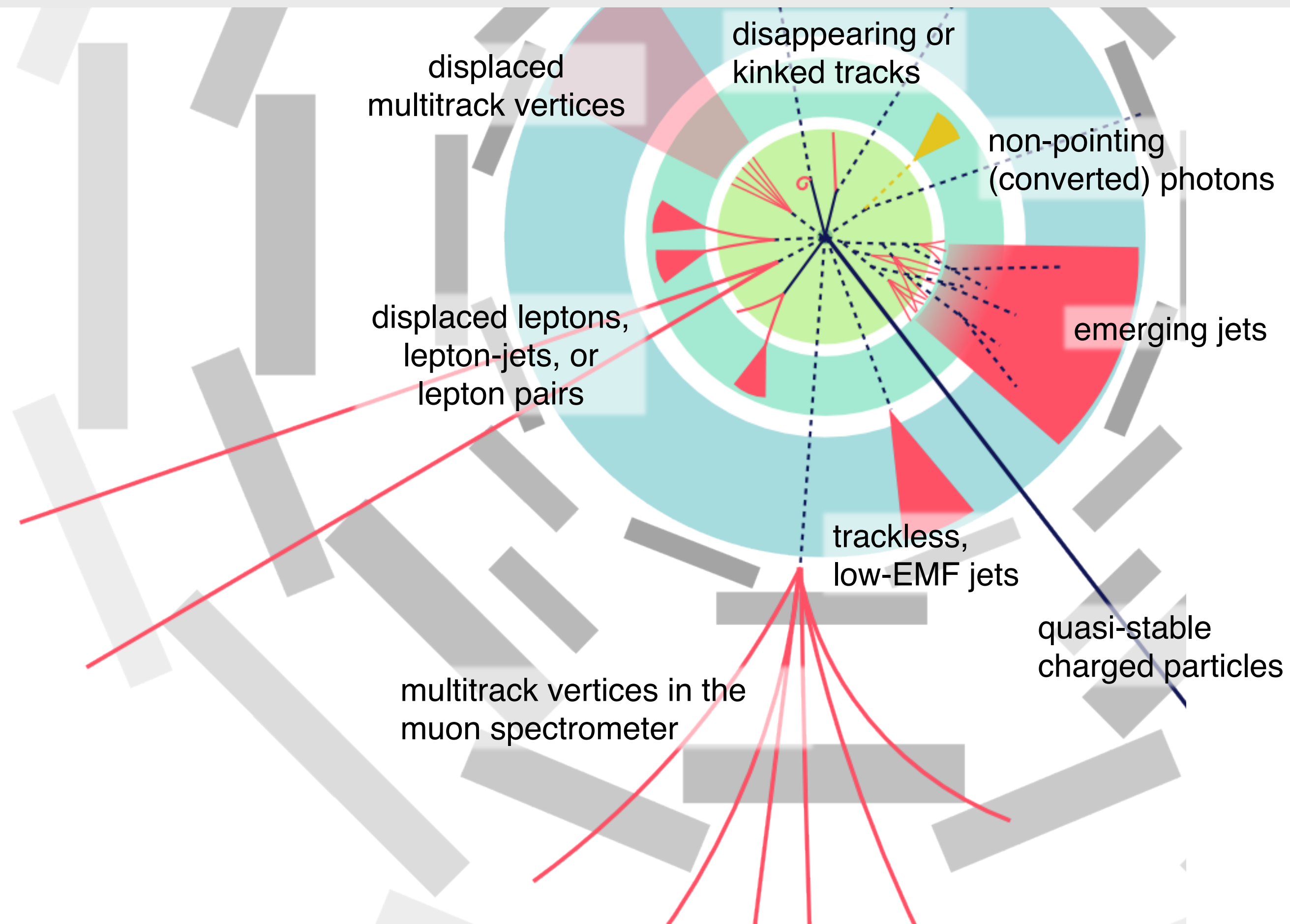
- neither at the LHC
- nor at any future collider

A demonstrated capability to cope with a large and diverse set of experimental signatures is a key item in the wishlist of features for any future collider project

Sensitivity to a broad spectrum of signatures is a key asset, that *will survive changes in the theory knowledge/prejudice*

Challenges for detectors

SUBTLE IS THE LORD



Heather Russell, McGill University

24 April 2017

Current “general purpose” experiments were really designed for pretty standard signals, e.g. $H \rightarrow \gamma \gamma$ or jets+mET

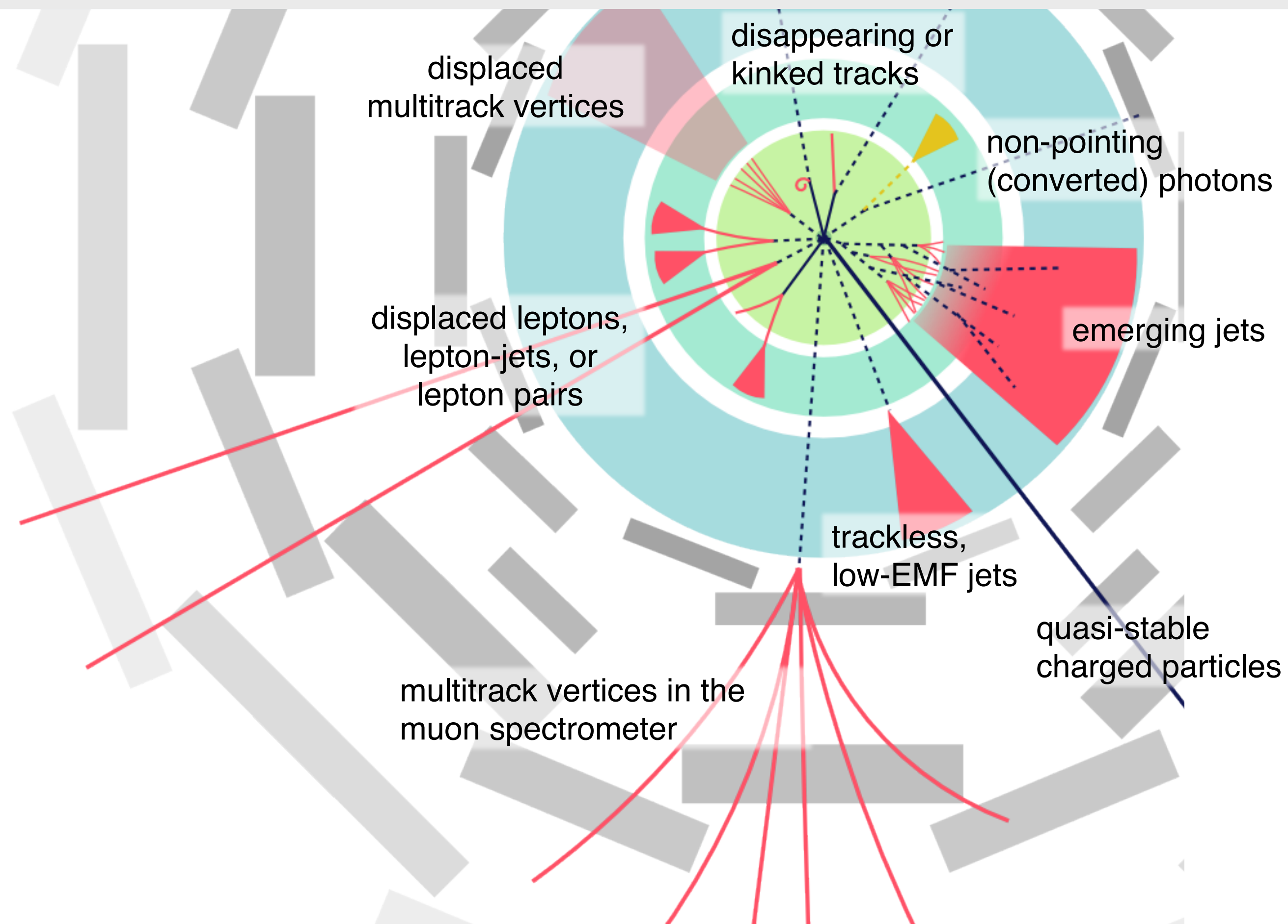
- Now is the time to think about what can be done in future detectors.
- CLIC (and ILC) can exploit full detector simulation

Please note these signals are not necessarily a way to “jump off the ship”. Most of them, if not all(!), can arise in the MSSM if one relaxes some (motivated?) UV assumptions.

Long Lived Particles Forum

NEXT MEETING AT CERN 27 TO 29 MAY 2019

latest meeting in Amsterdam



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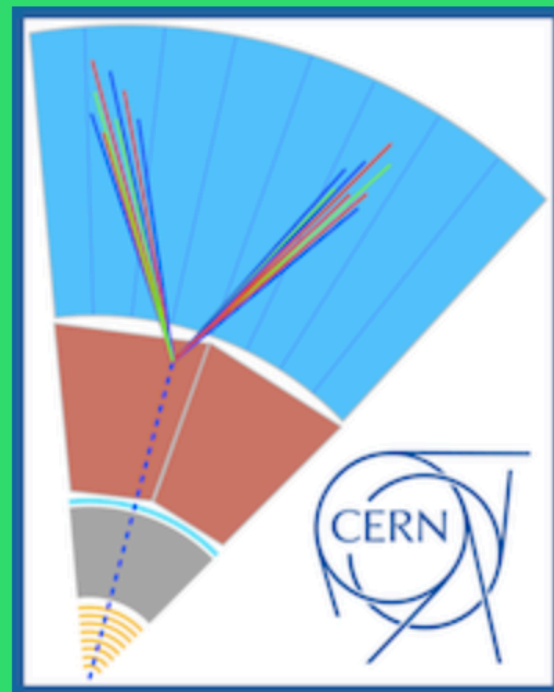
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Searching for long-lived particles at the LHC: Fifth workshop of the LHC LLP Community

27-29 May 2019
CERN
Europe/Zurich timezone

Search... 

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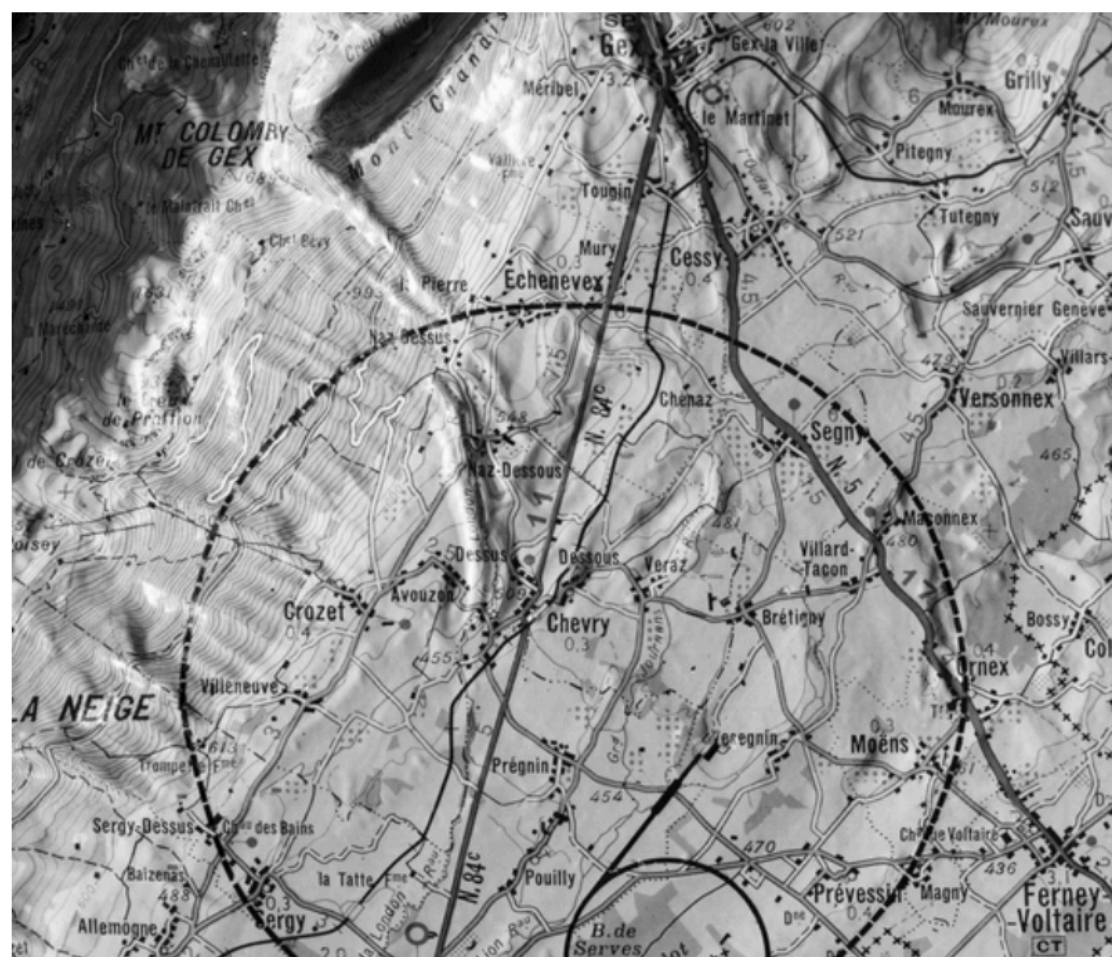
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Overview

Timetable

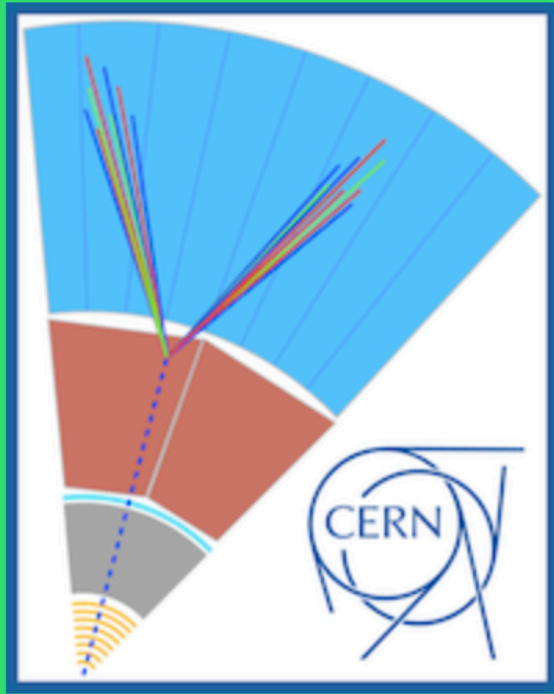
Participant List



Long Lived Particles Forum

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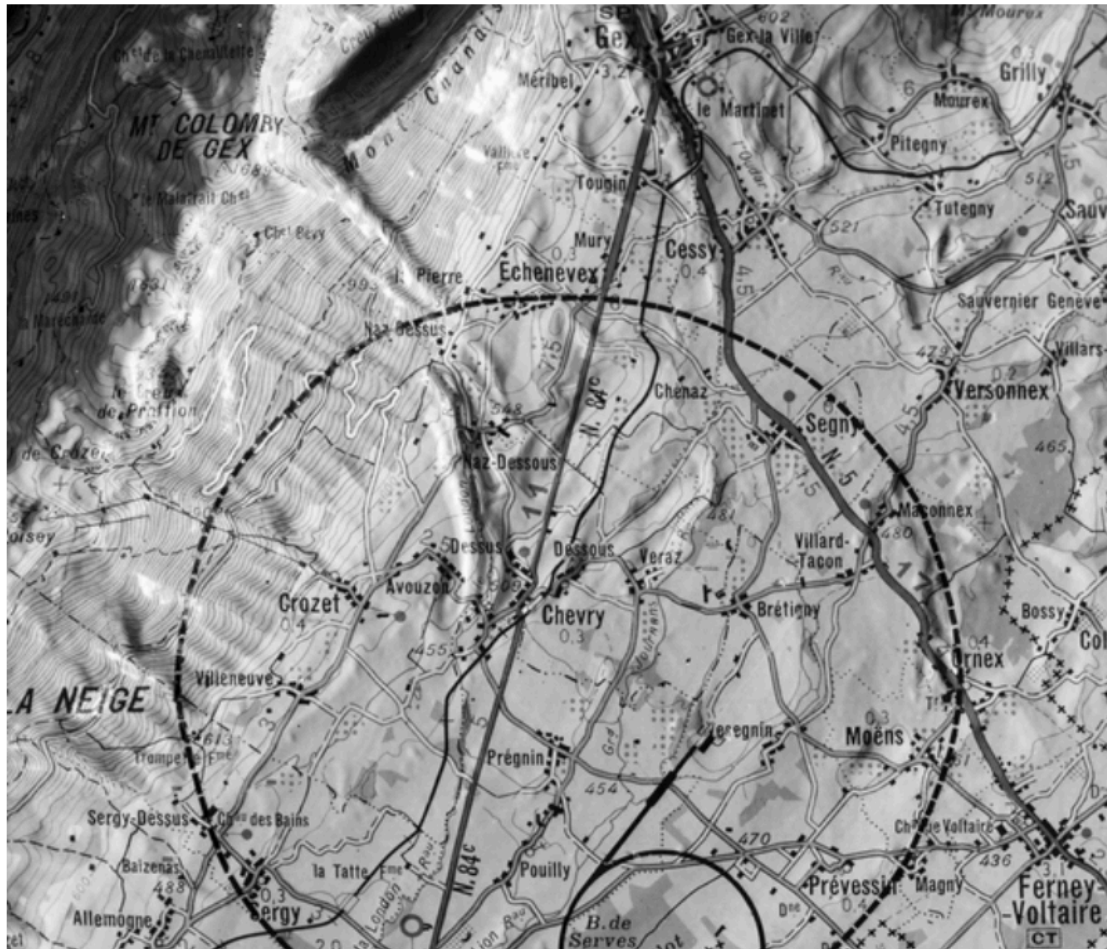
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LLP at Future Colliders “Breakout session”

27-29 May 2019
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Search...

- Overview
- Timetable
- Participant List



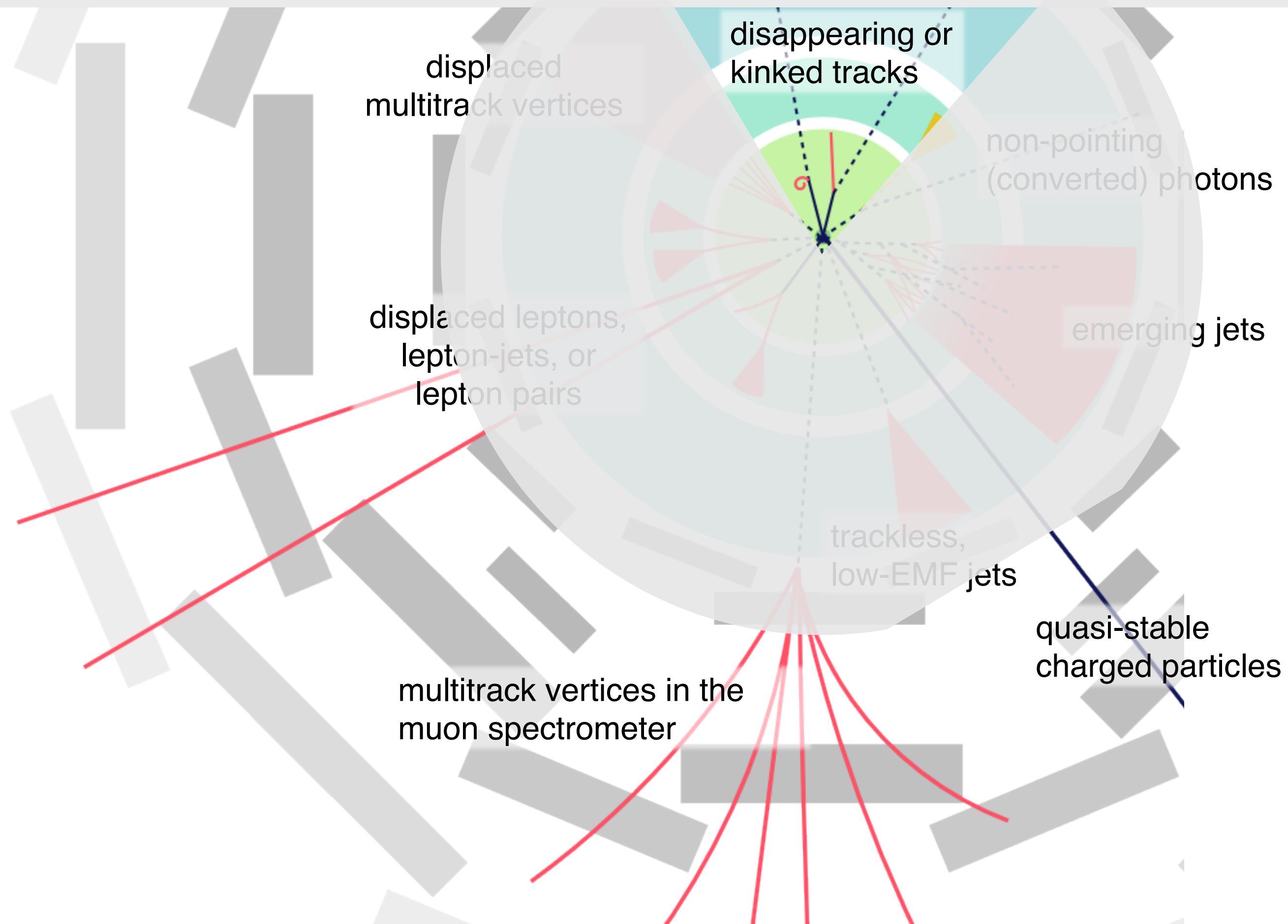
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Challenges for detectors

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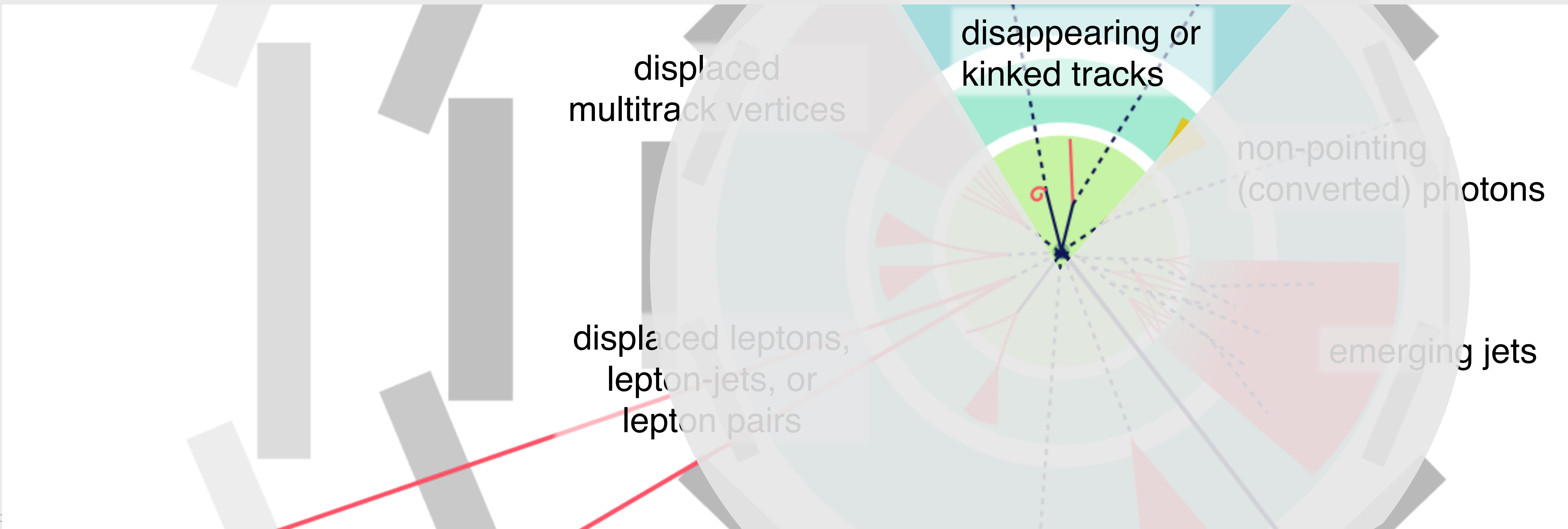


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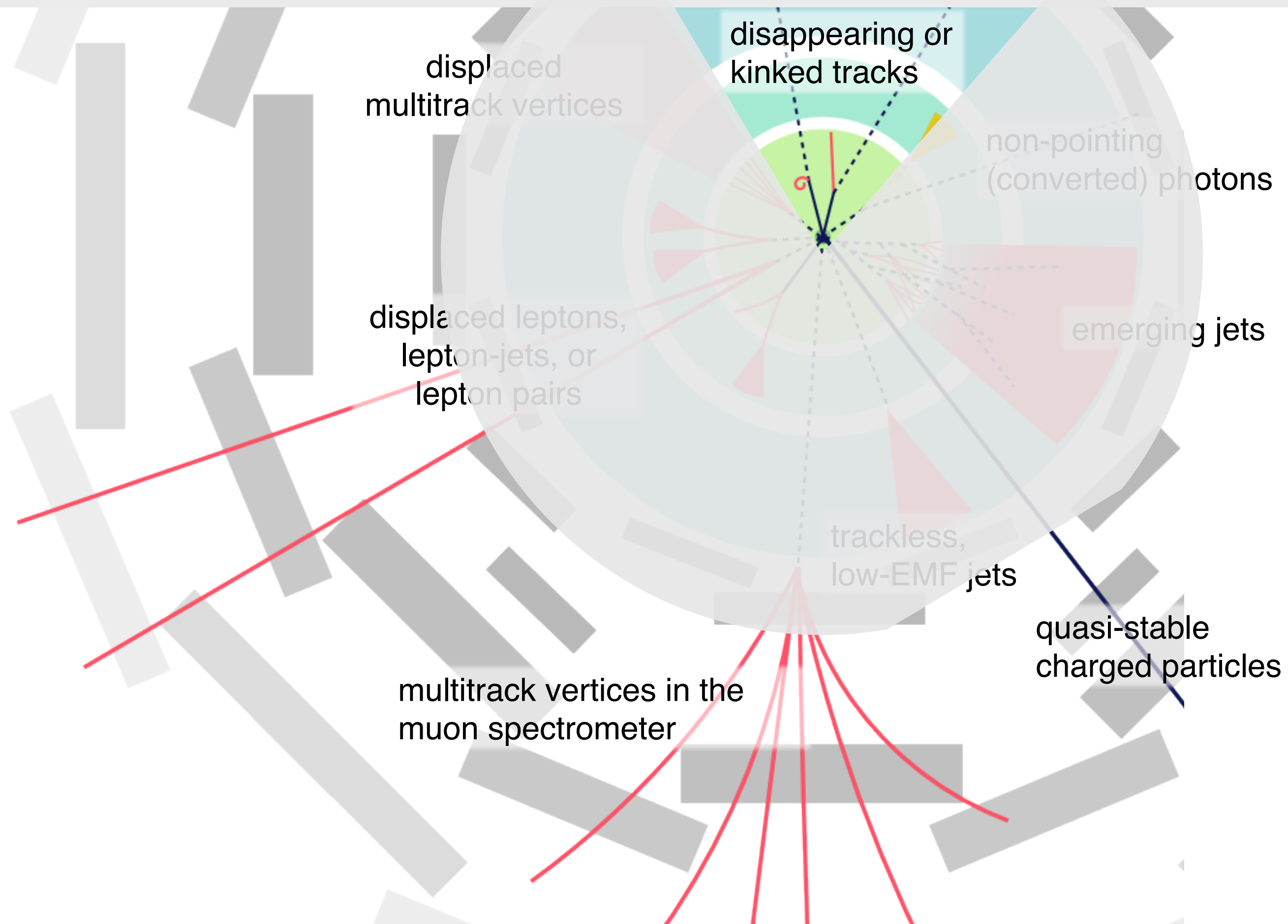
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SUBTLE IS THE LORD



Challenges for detectors

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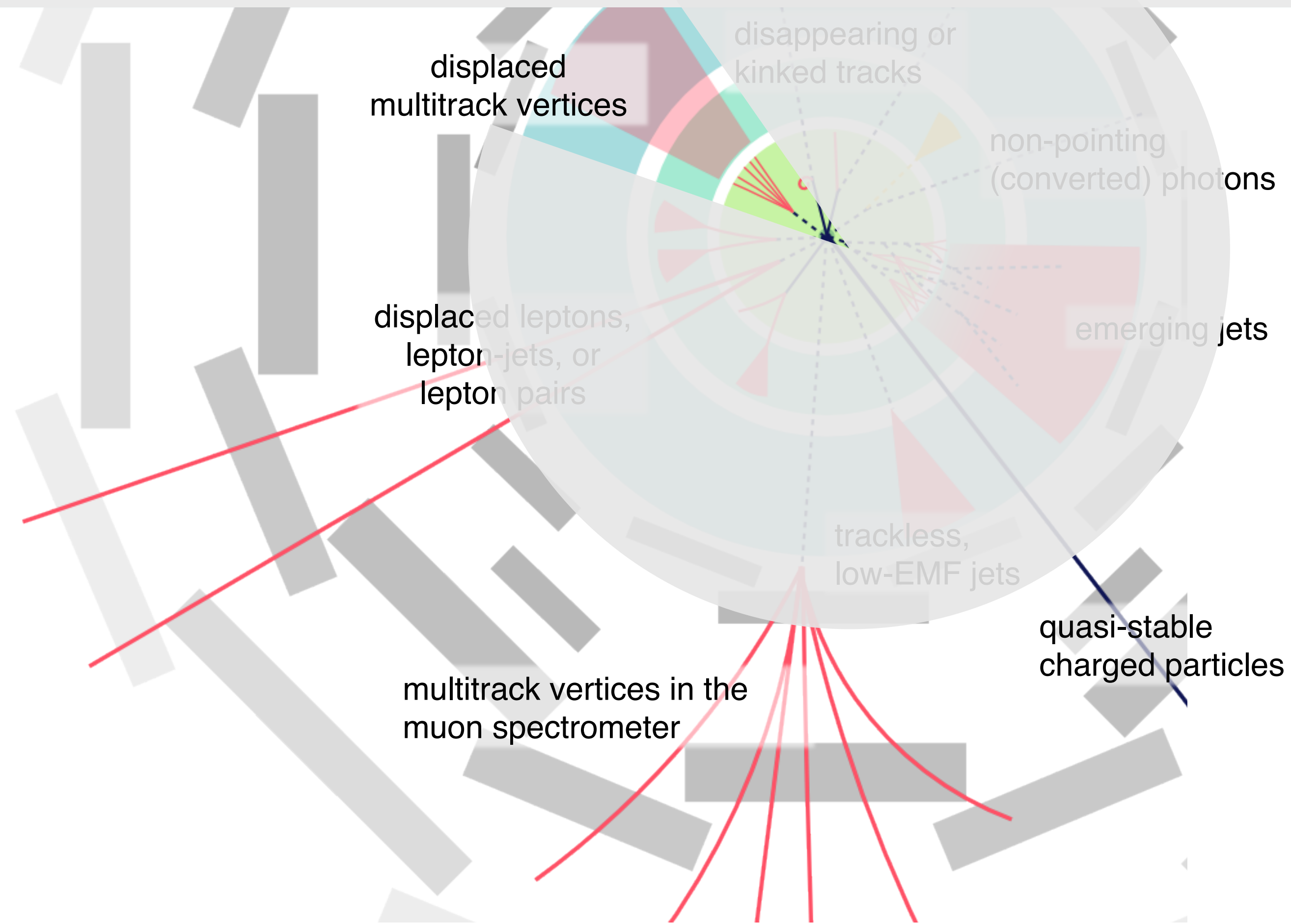


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Challenges for detectors

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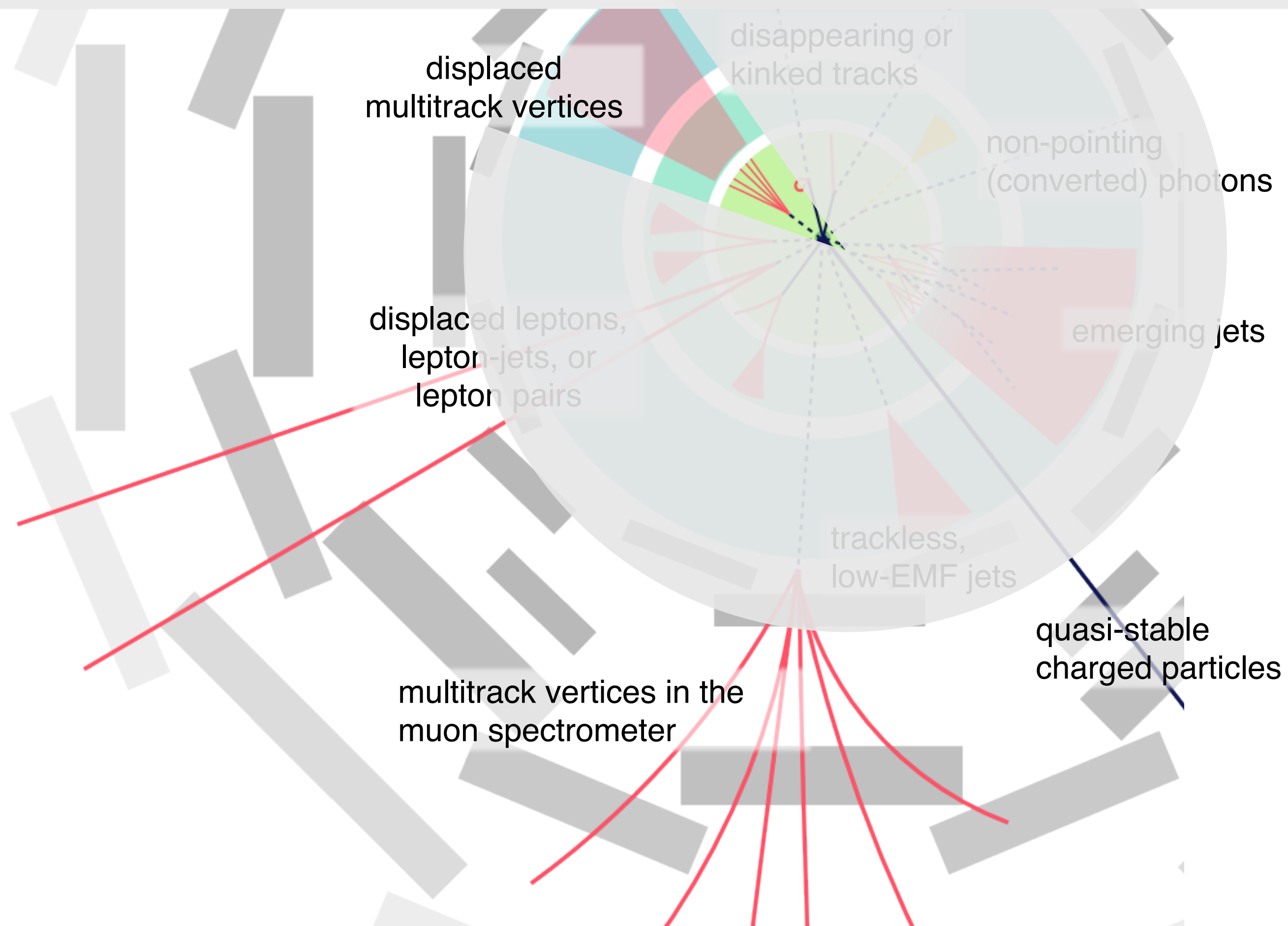


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Challenges for detectors

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Challenges for detectors

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displaced
multitrack vertices

disappearing or
kinked tracks

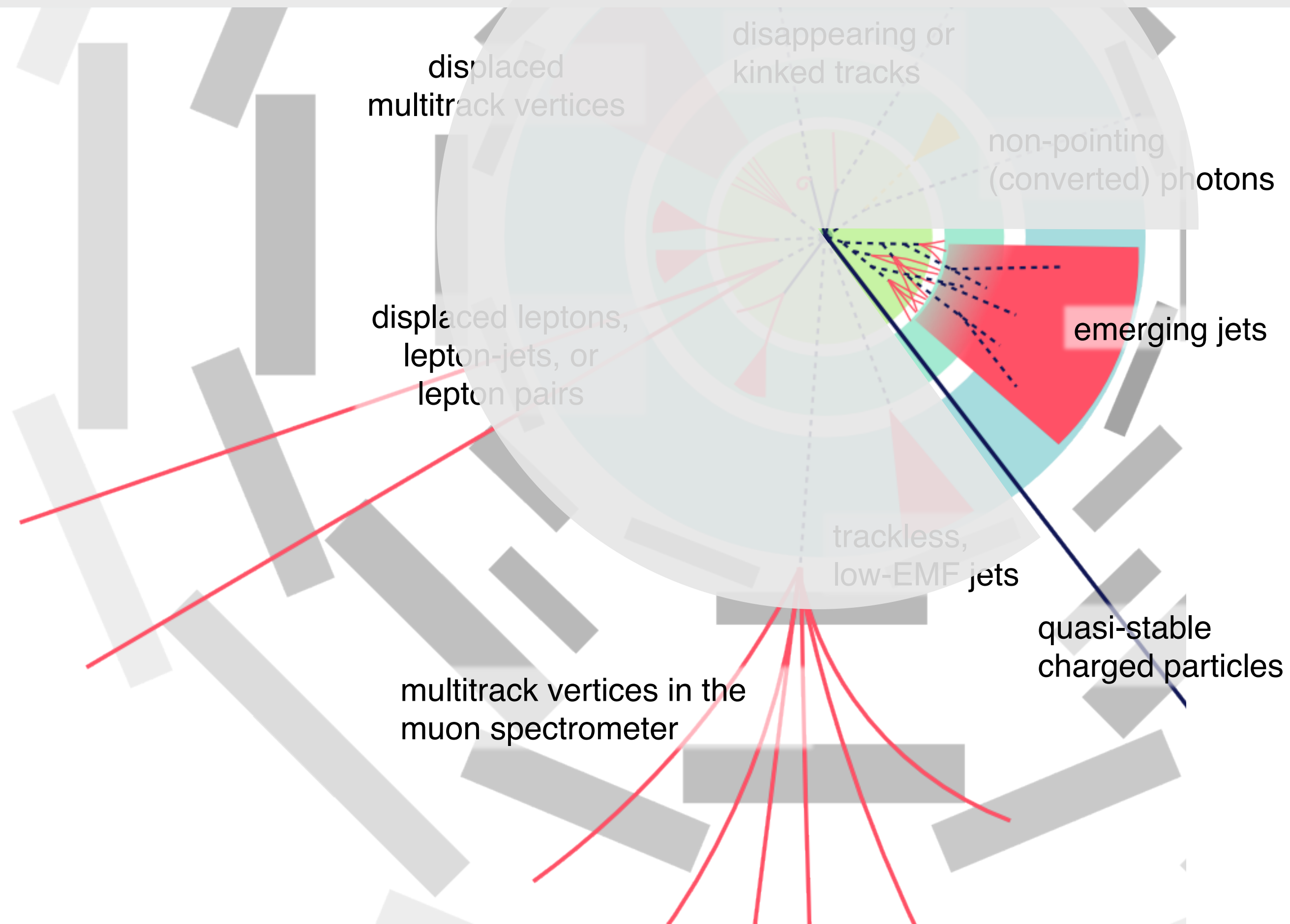
non-pointing
(converted) photons

displaced leptons,
lepton-jets, or
lepton pairs

emerging jets

Challenges for detectors

SUBTLE IS THE LORD

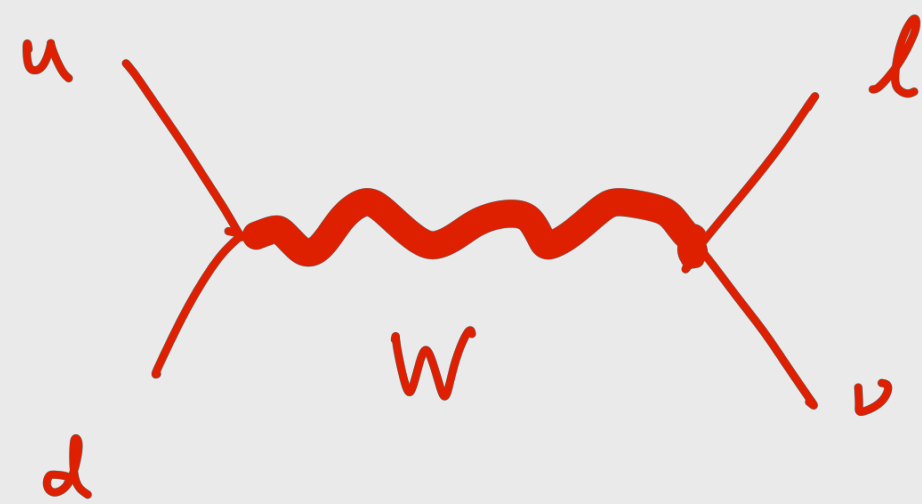


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... plenty of other options. A white paper is coming from the Long Lived Forum to summarize.

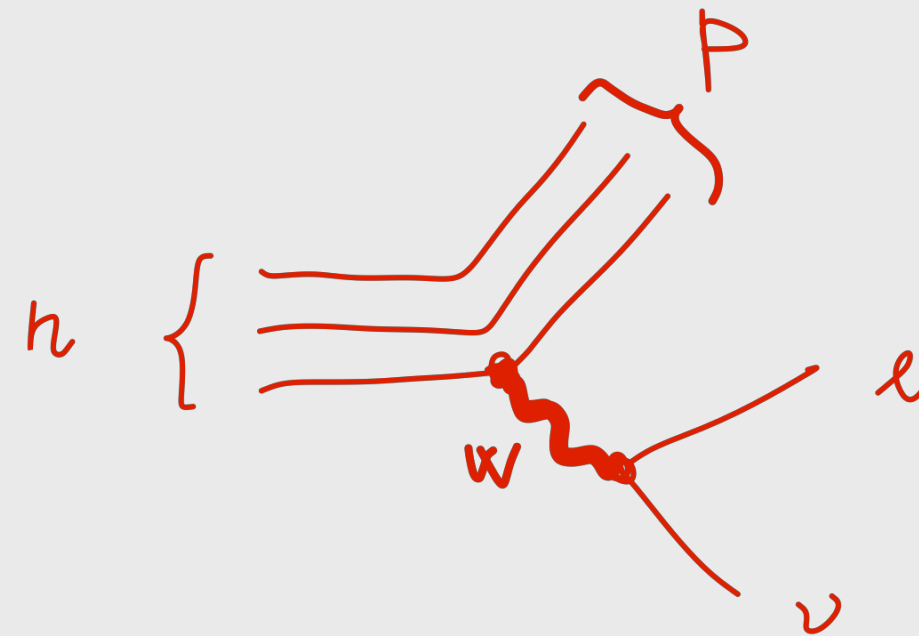
Three generic reasons why this can happen



$$\pi \rightarrow \ell \nu$$

heavy mediator

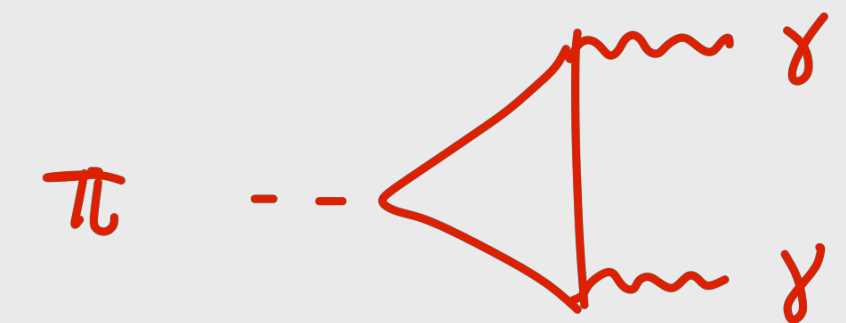
from a hierarchical larger mass scale



$$n \rightarrow p e \bar{\nu}$$

small phase-space

from symmetry



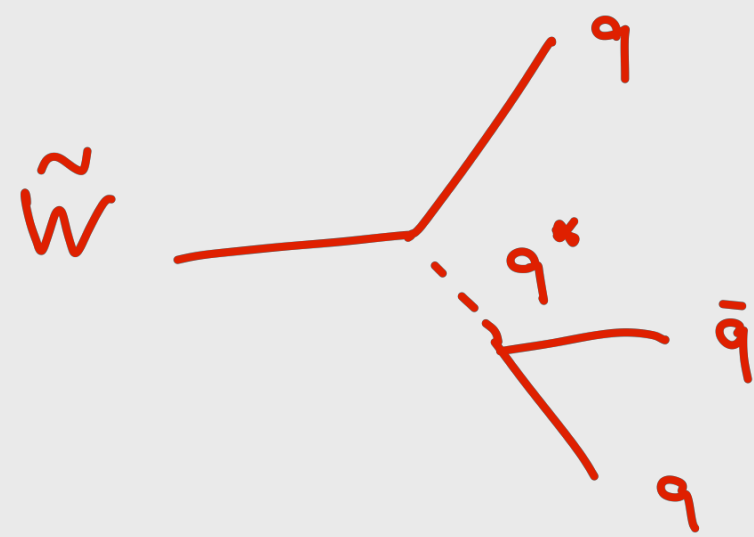
$$\pi \rightarrow \gamma \gamma$$

small coupling

from symmetry breaking

- it happens in QCD (!)
- the necessary ingredients are just the same as for the formulation of any model of particle physics (mass scales and their hierarchies, symmetries and their breaking, ...)

Many generic BSM reasons why this can happen



$$\tilde{\lambda}_i^0 \rightarrow q\bar{q}^* \rightarrow q\bar{q}q$$

heavy mediator

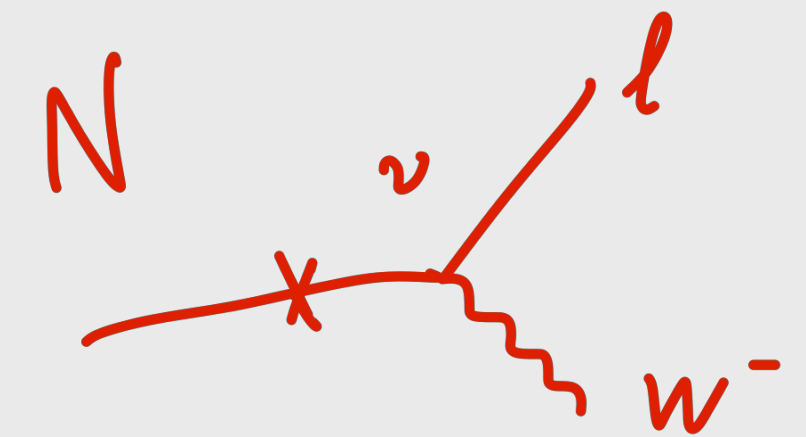
from a hierarchical larger mass scale



$$\tilde{h}^+ \rightarrow \tilde{h}^0 W^{+*}$$

small phase-space

from symmetry



$$N^0 \rightarrow \ell^+ W^-$$

small coupling

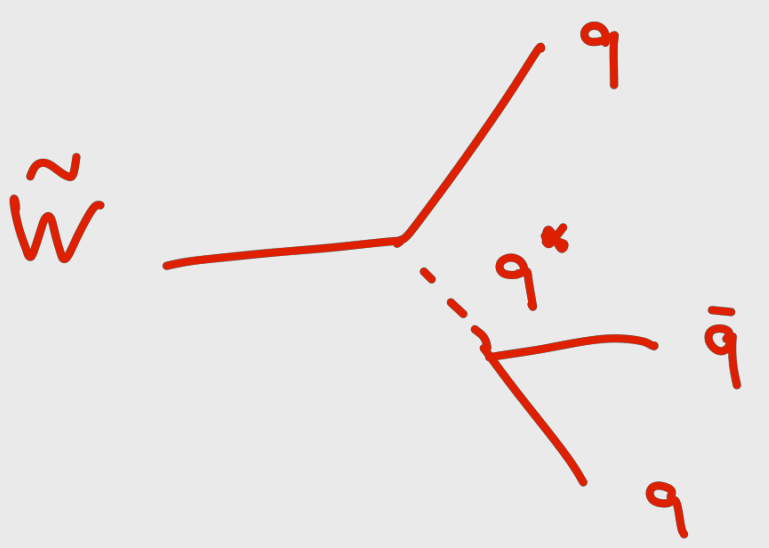
from symmetry breaking

- it happens in QCD (!)
- the necessary ingredients are just the same as for the formulation of any model of particle physics (mass scales and their hierarchies, symmetries and their breaking, ...)

Many generic BSM reasons why this can happen

long life-time may be dictated by other physics requirements as well

late-time decay in the Early Universe to fulfill out-of-equilibrium condition for baryon number generation



$$\tilde{\lambda}_i^0 \rightarrow q\tilde{q}^* \rightarrow q\bar{q}q$$

heavy mediator

from a hierarchical larger mass scale

Dark Matter candidate



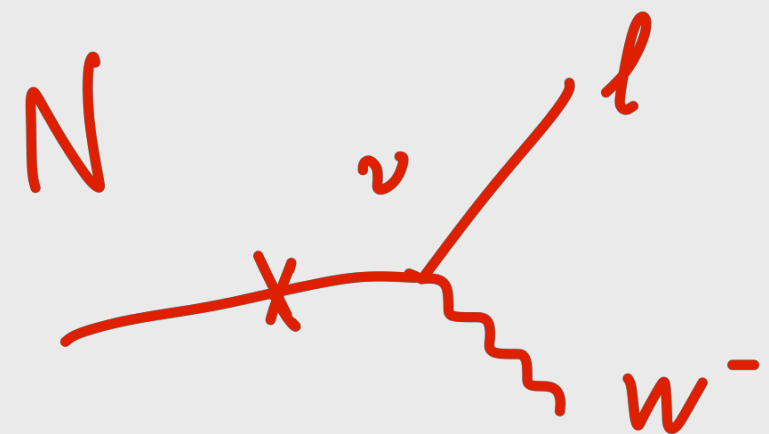
$$\tilde{h}^+ \rightarrow \tilde{h}^0 W^{+*}$$

small phase-space

from symmetry

relic abundance $\Omega \sim M_{DM}/g$
so $g \ll 1$ for light particles

$$\gamma_D \rightarrow \ell^+\ell^-$$



$$N^0 \rightarrow \ell^+ W^-$$

small coupling

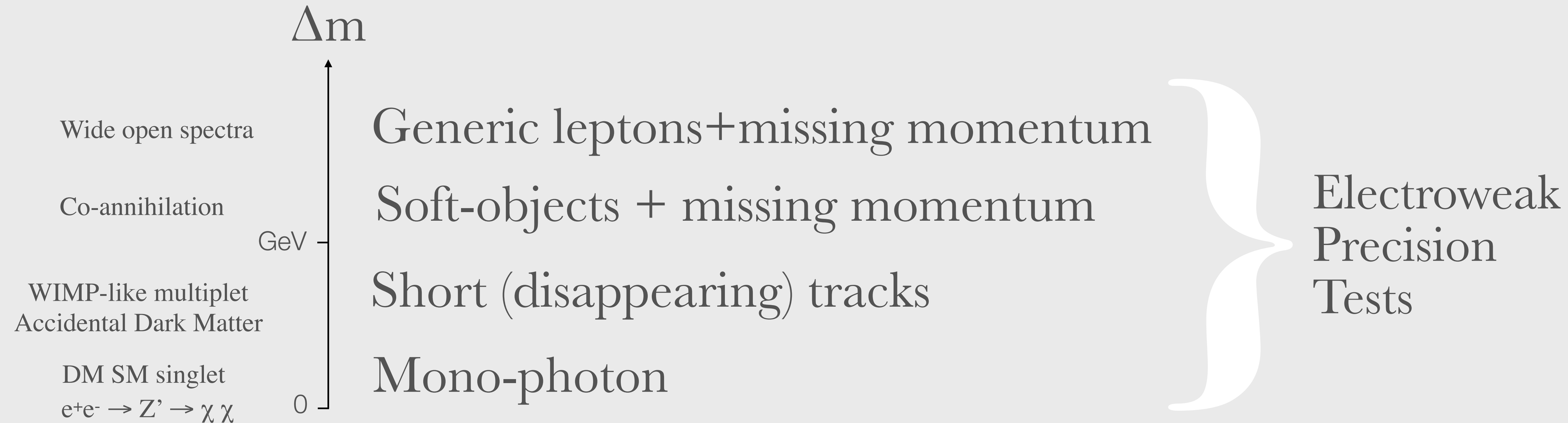
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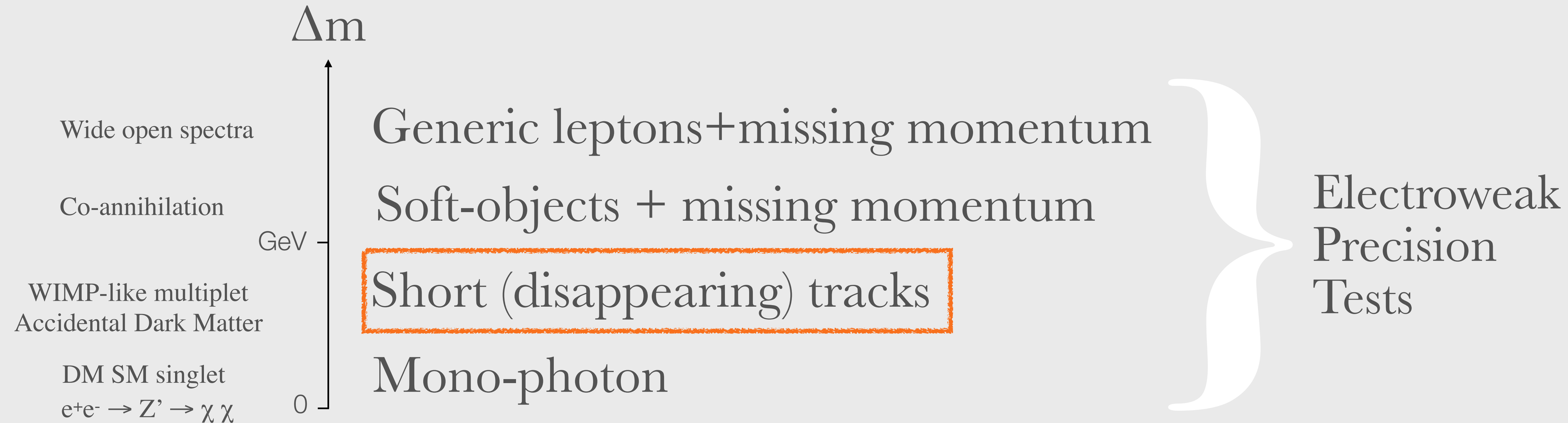
Direct searches

Dark Matter

Electroweak Dark Matter: LSP (+NLSP)



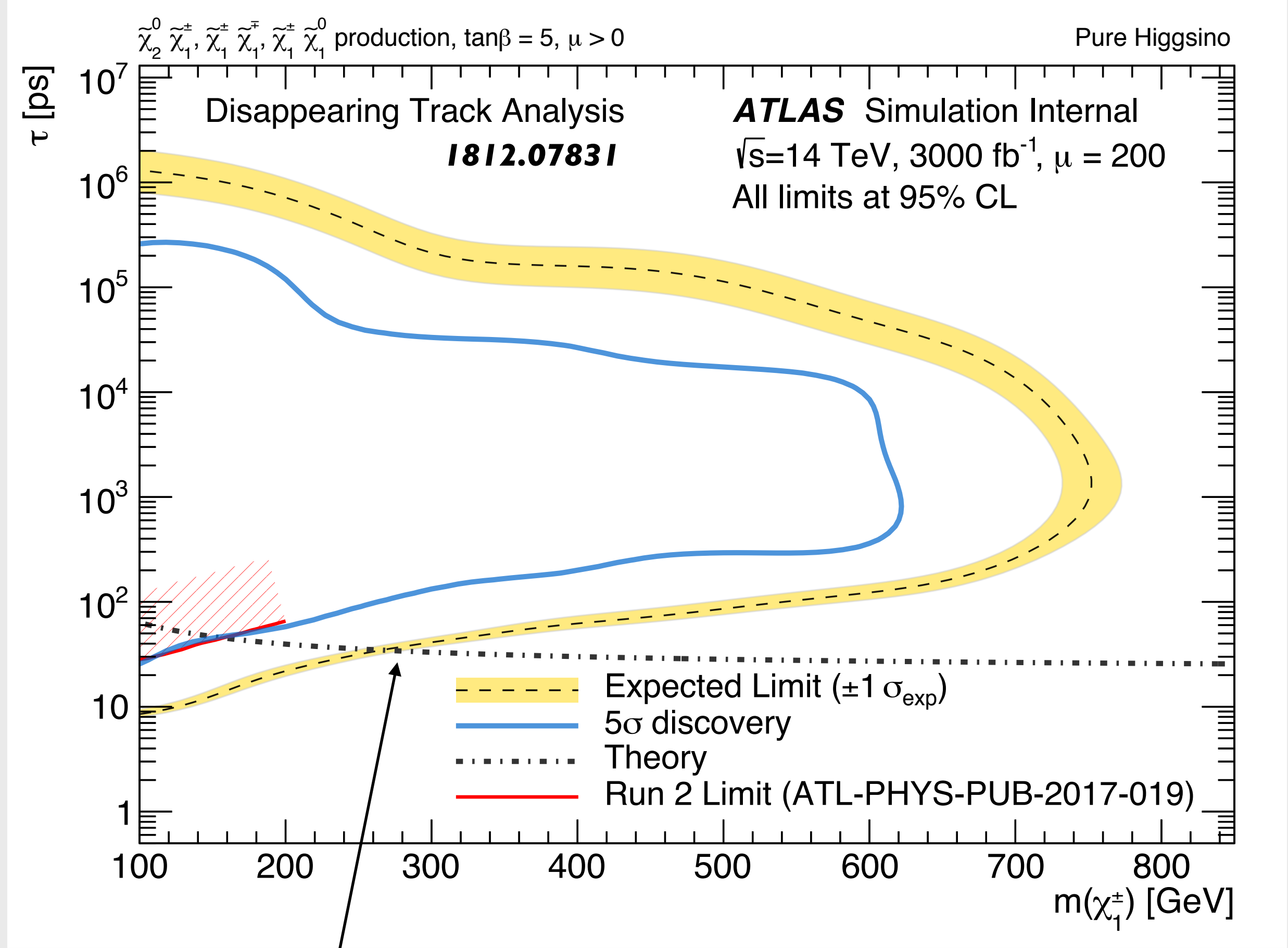
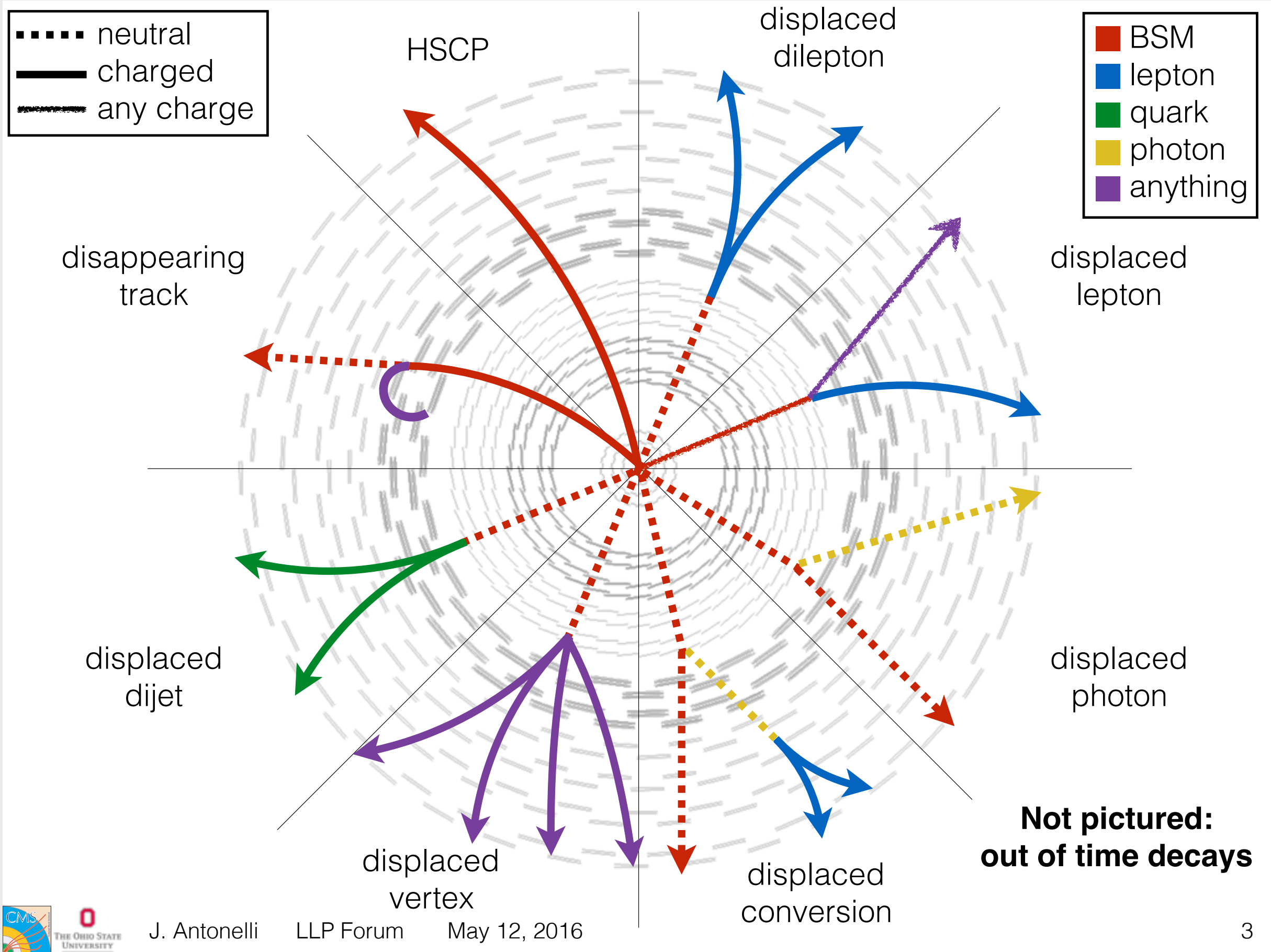
Electroweak Dark Matter: LSP (+NLSP)



Short (disappearing) tracks

HIGGSINO DM

O(CM) DISAPPEARING TRACKS

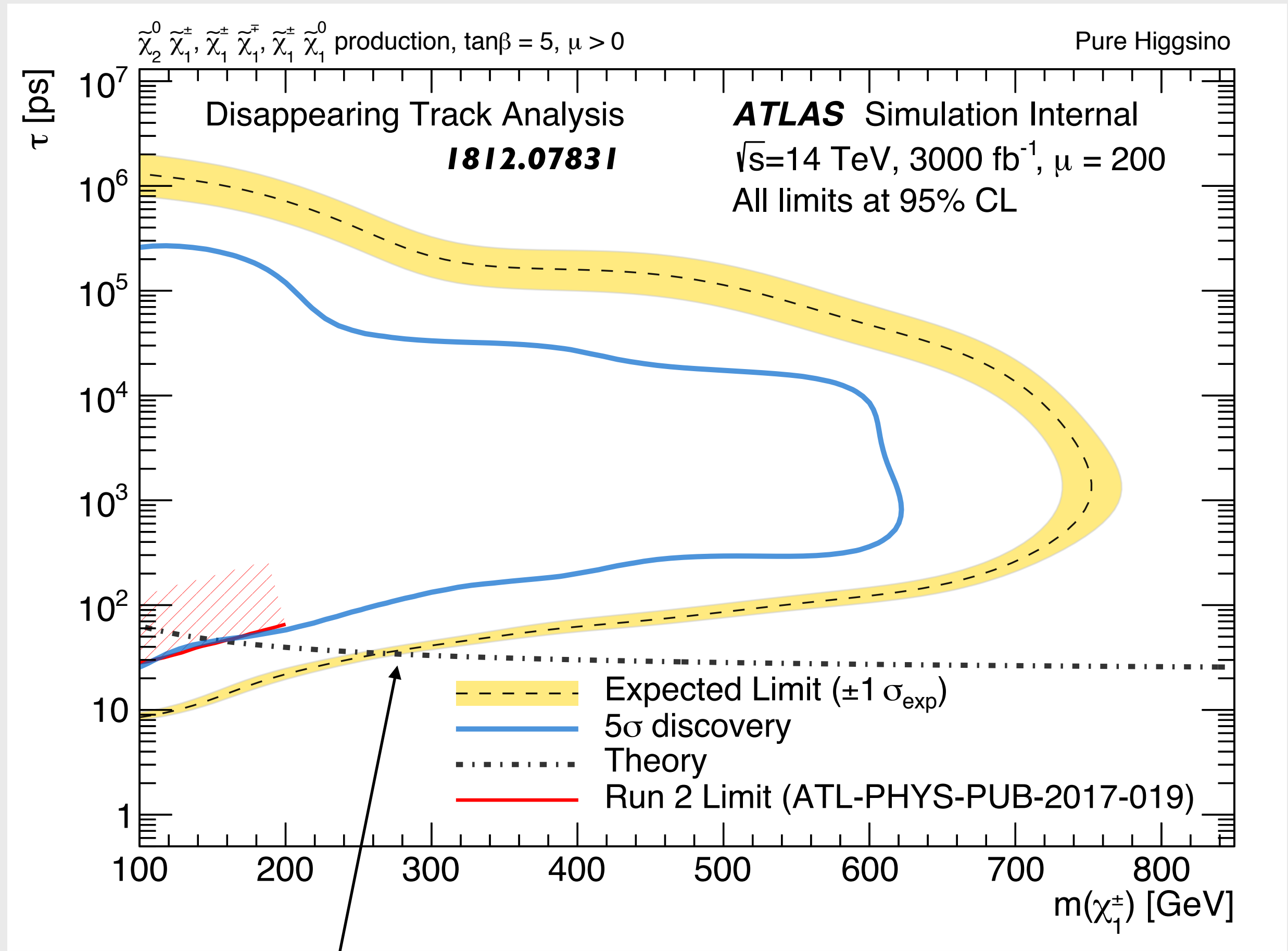
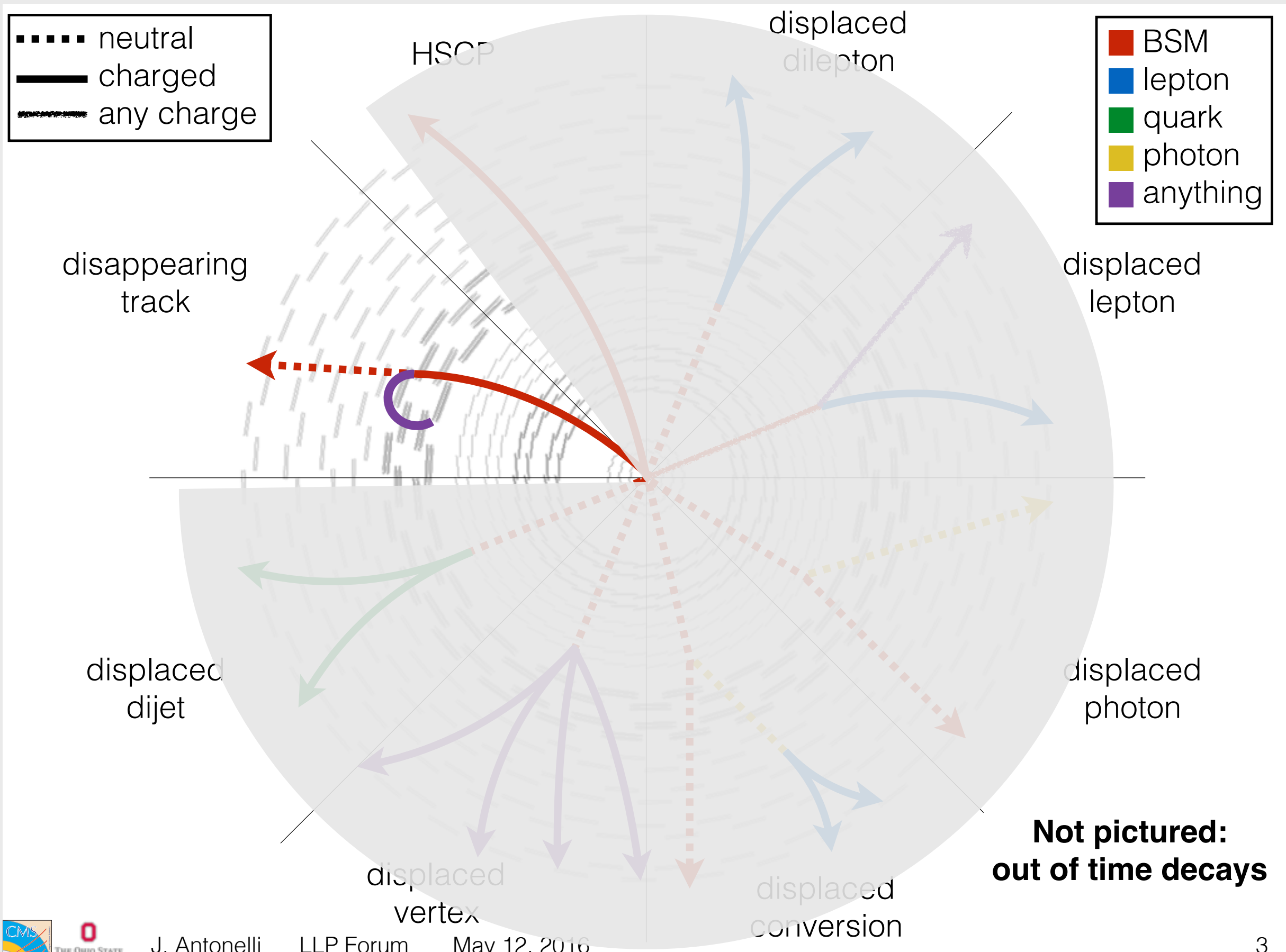


HL-LHC up to 300 GeV

Short (disappearing) tracks

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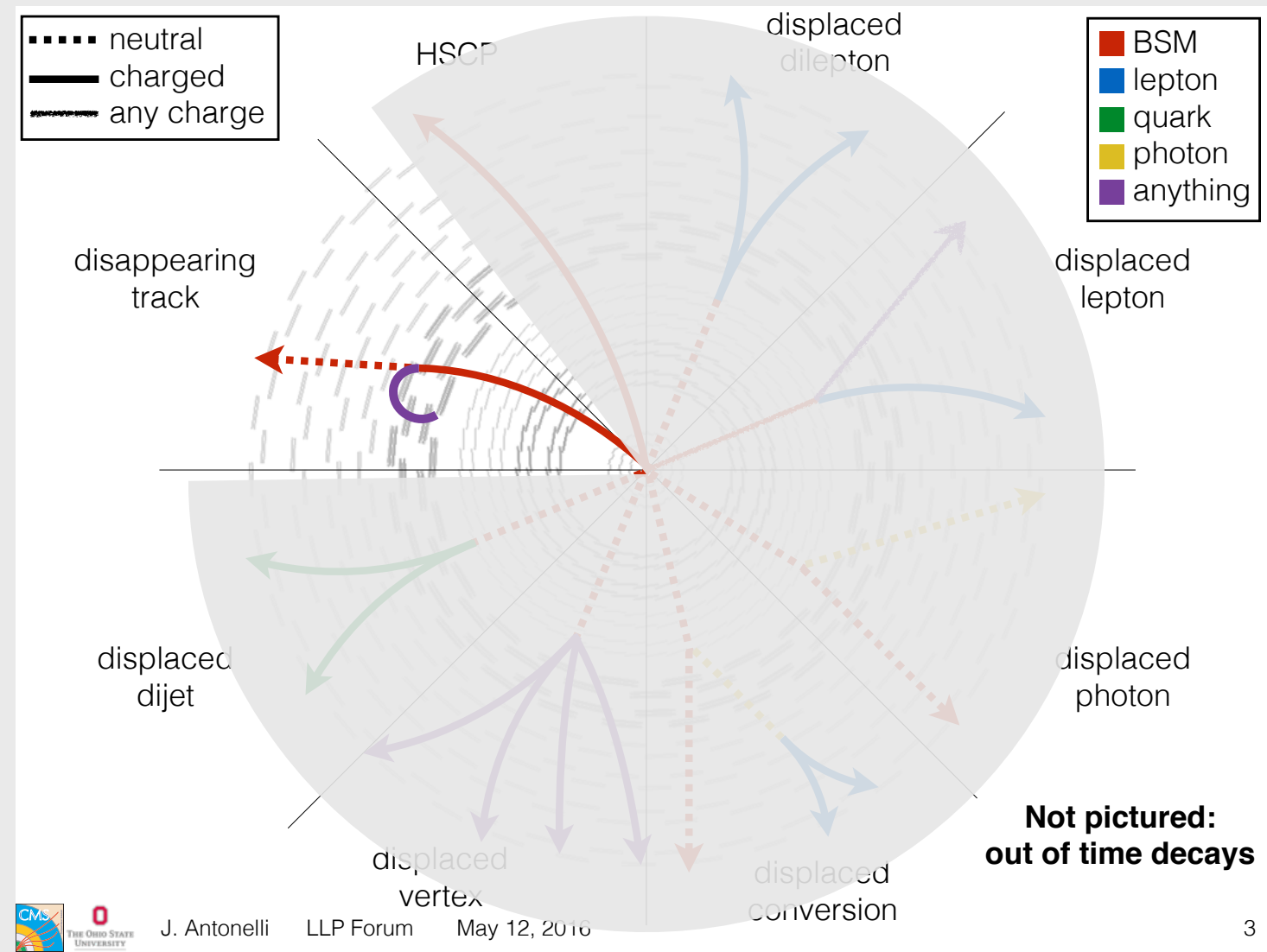
O(CM) DISAPPEARING TRACKS



Short (disappearing) tracks

HIGGSINO DM

O(CM) DISAPPEARING TRACKS



Clean experimental environment:

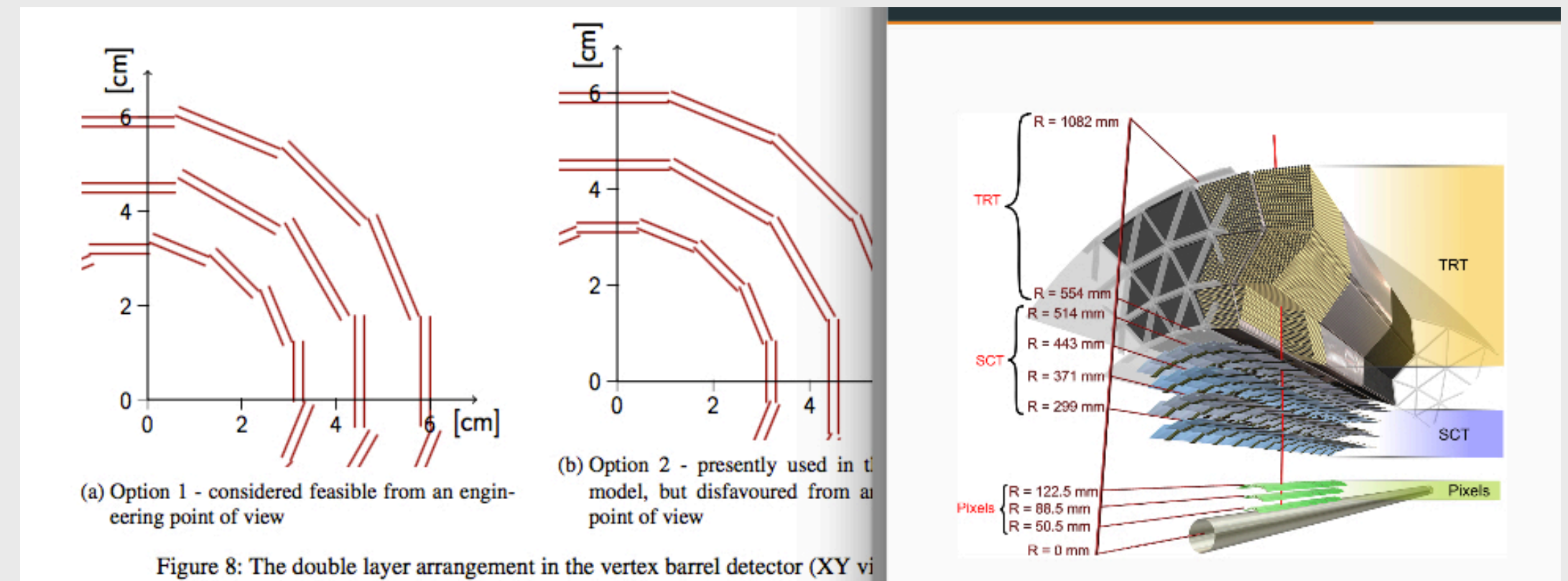
- No trigger
- No QCD background
- Tracker is closer to the beam

Challenges:

- boost cannot make \tilde{h} tracks longer

CLIC Efficient at $d \geq 4$ cm

LHC Efficient at $d \geq 10$ cm



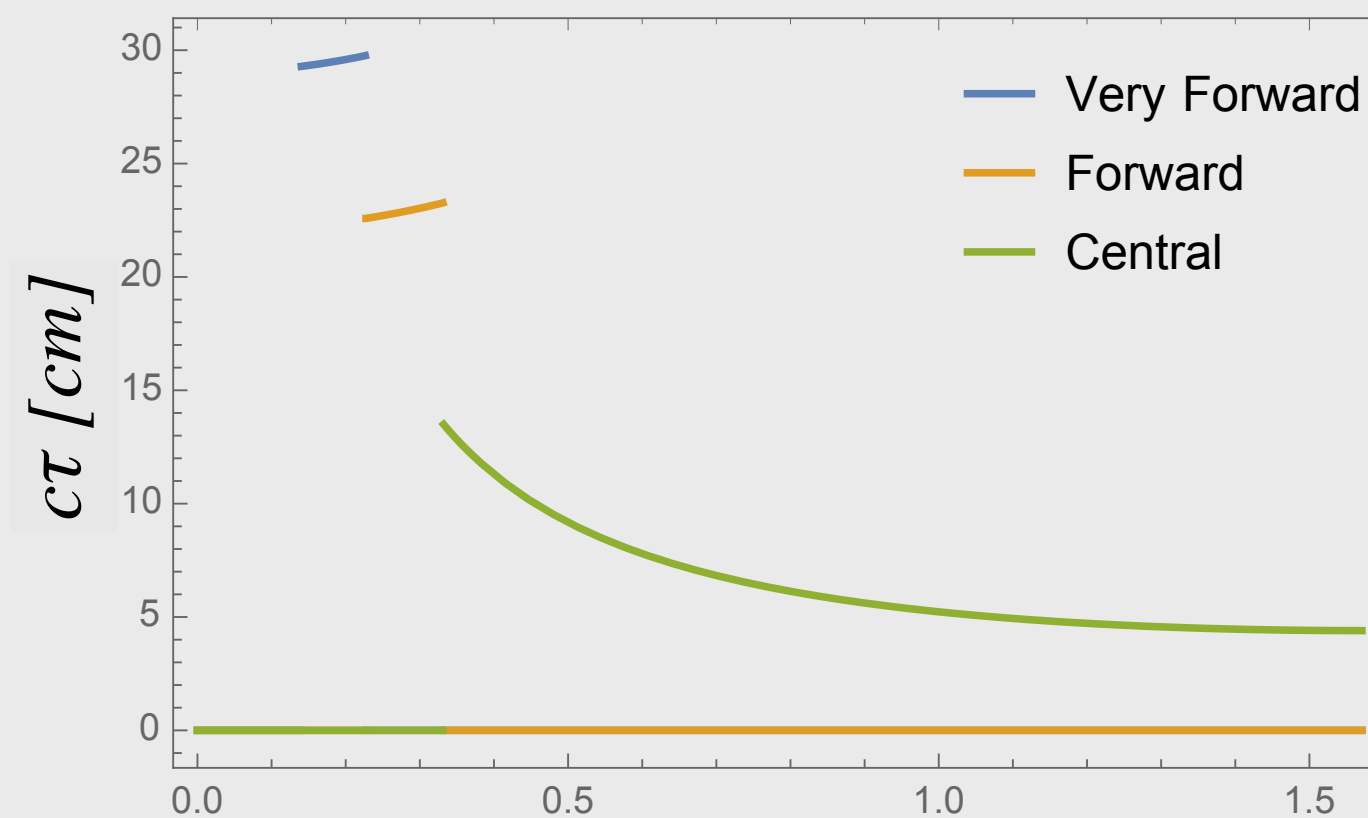
HL-LHC can put bounds on Higgsino up to ~ 300 GeV

Exponential rate gain when $c\tau < 10$ cm

SIMPLIFIED RECONSTRUCTION

Assume track is seen when

- $c\tau > 4.4 \text{ cm}/\sin\theta$ ($19^\circ < \theta < 90^\circ$)
- $c\tau > 22 \text{ cm}/\cos\theta$ ($13^\circ < \theta < 19^\circ$)
- $c\tau > 29 \text{ cm}/\cos\theta$ ($0^\circ < \theta < 13^\circ$)



CLICdp-Note-2017-001

Short (disappearing) tracks

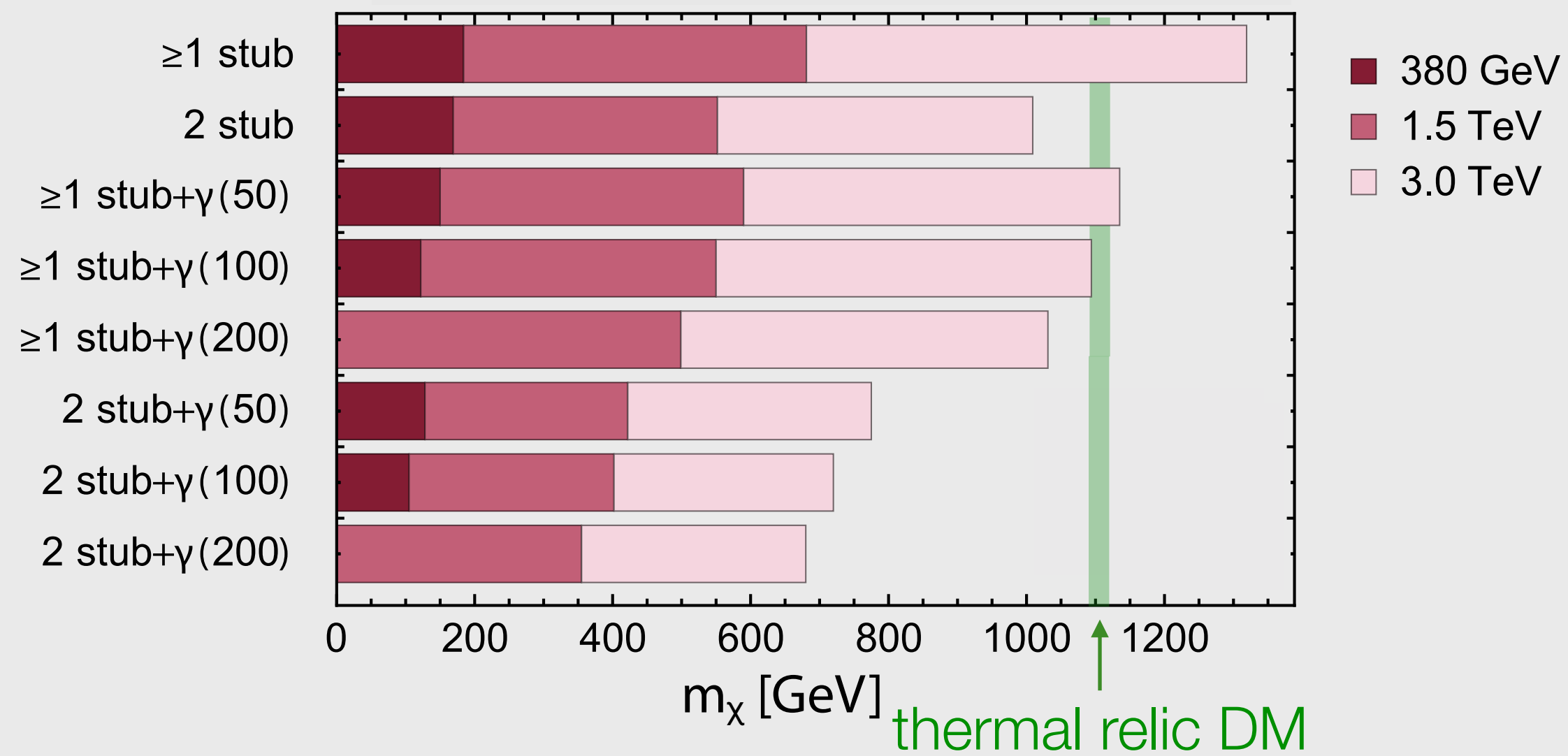
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O(CM) DISAPPEARING TRACKS

ASSUME PURE HIGGSINO LIFETIME

$ct = 1.2 \text{ cm @ } 200 \text{ GeV} \rightarrow 0.7 \text{ cm @ } 1 \text{ TeV}$

95% C.L. (Assuming Zero Background)



Short (disappearing) tracks

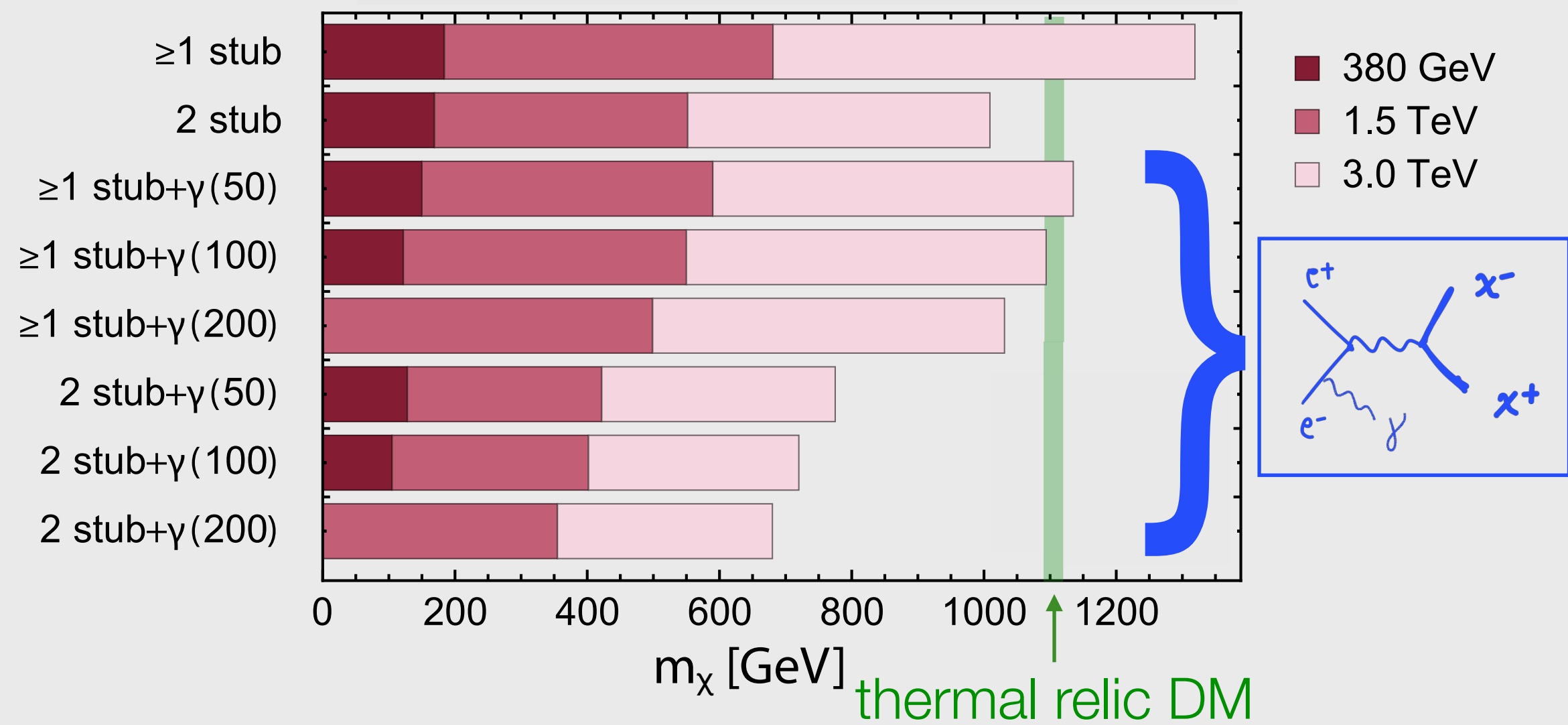
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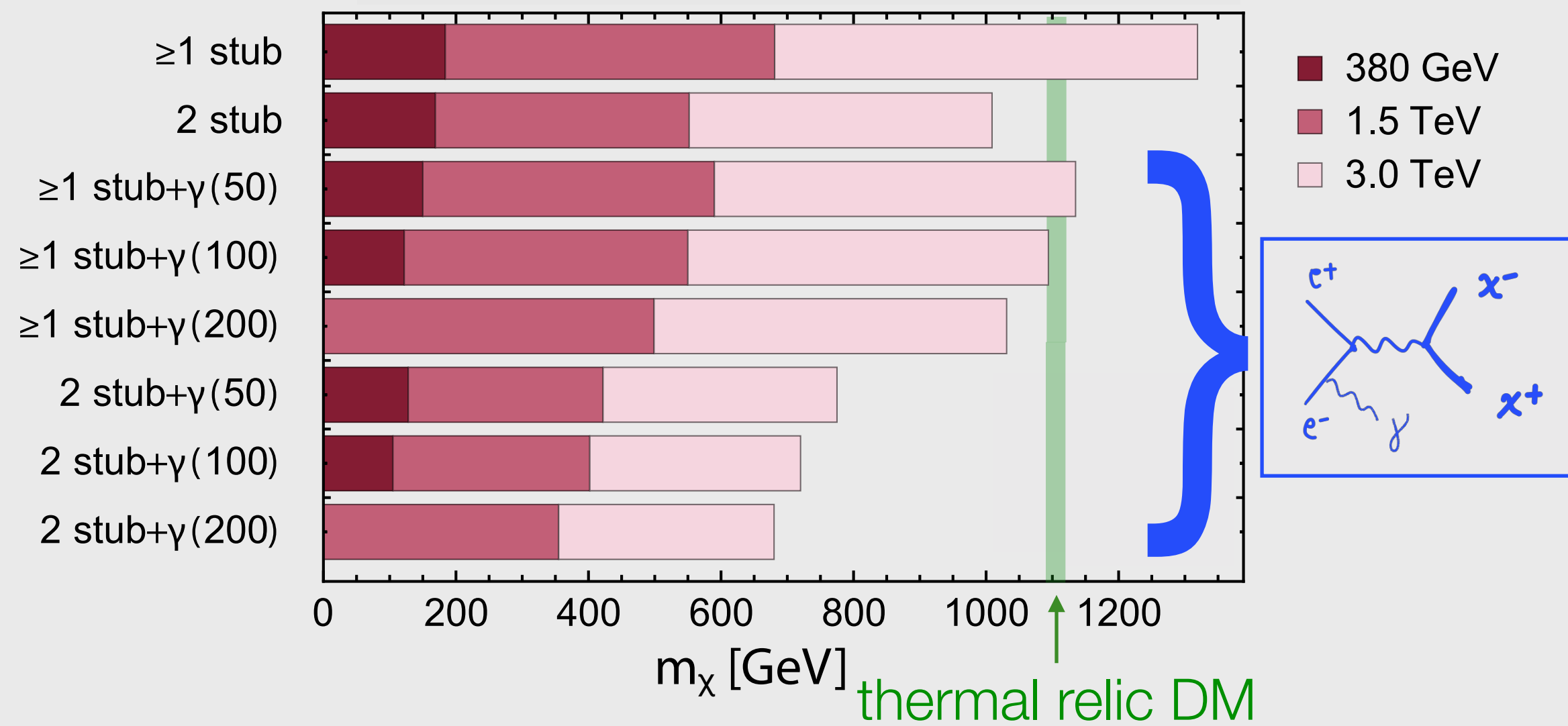
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background evaluation needed

CLIC 3 TeV yields 10 events per ab^{-1} for 1.1 TeV Higgsino thermal DM candidate

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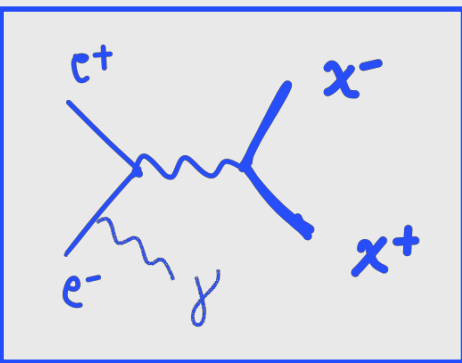
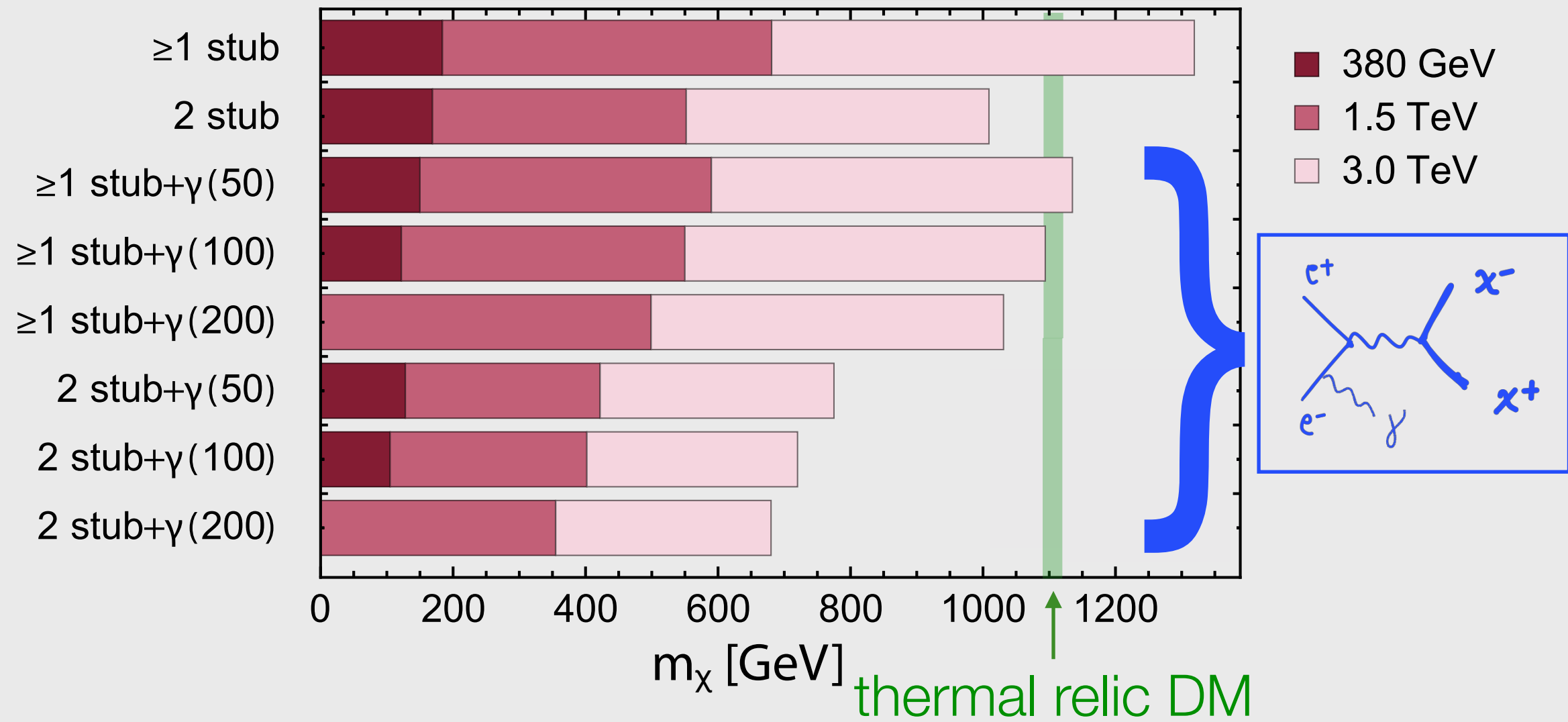
HIGGSINO DM

O(CM) DISAPPEARING TRACKS

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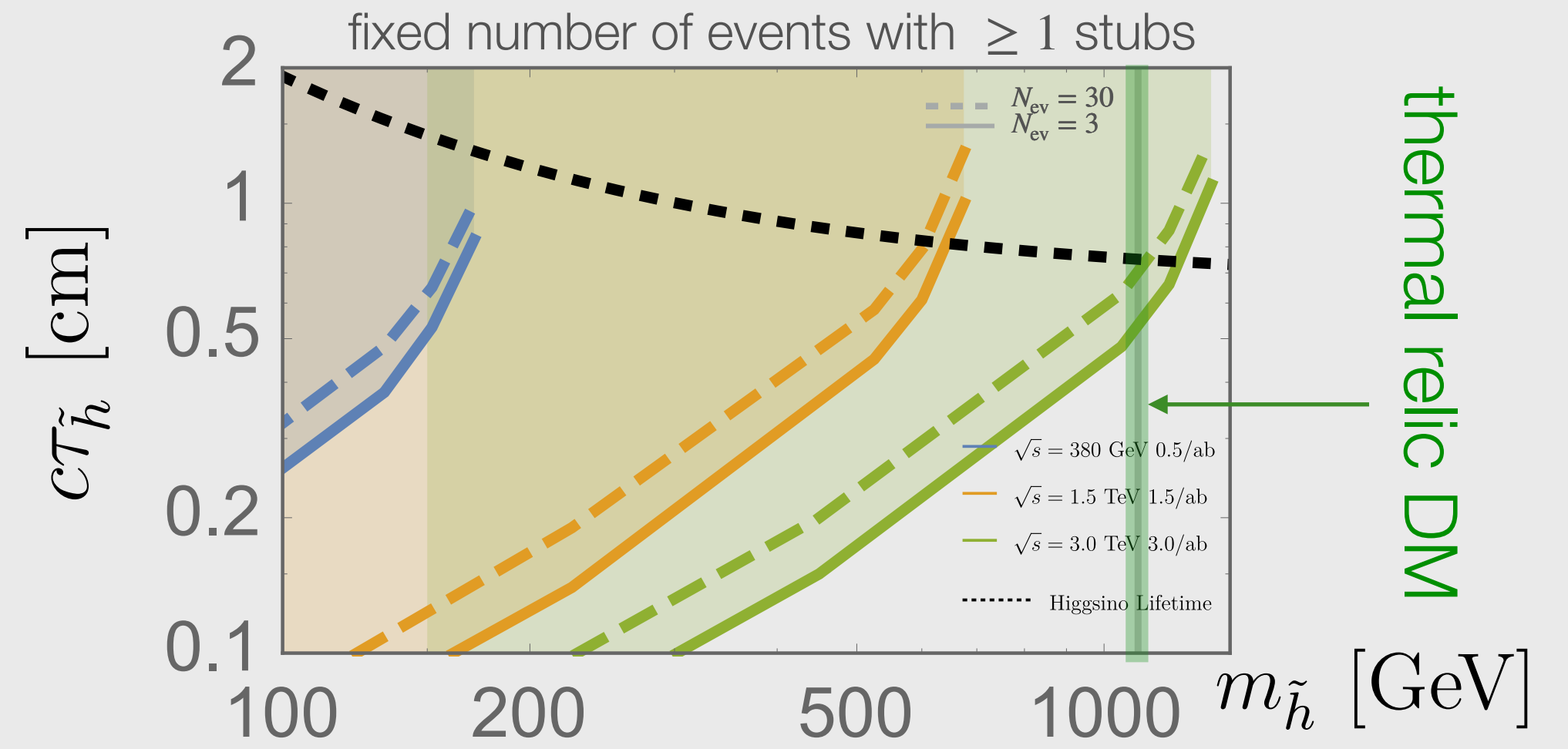
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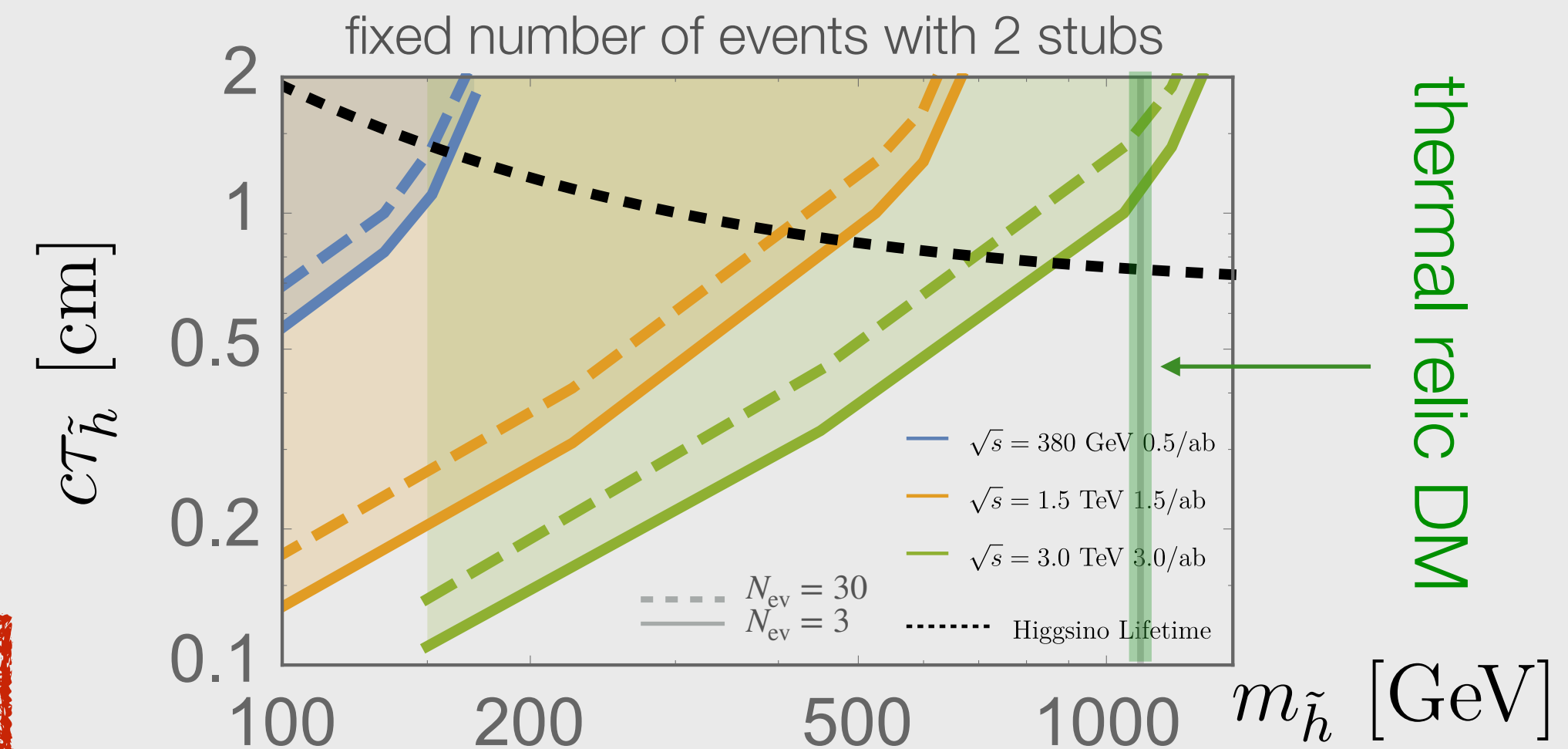
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LIFETIME AS PARAMETER (NON-PURE HIGGSINO)



thermal relic DM



thermal relic DM

Hidden Sector

via simplified models

Hidden Valley Displaced Vertex

CLICDP-NOTE-2018-001

$e^+e^- \rightarrow h \nu\nu$

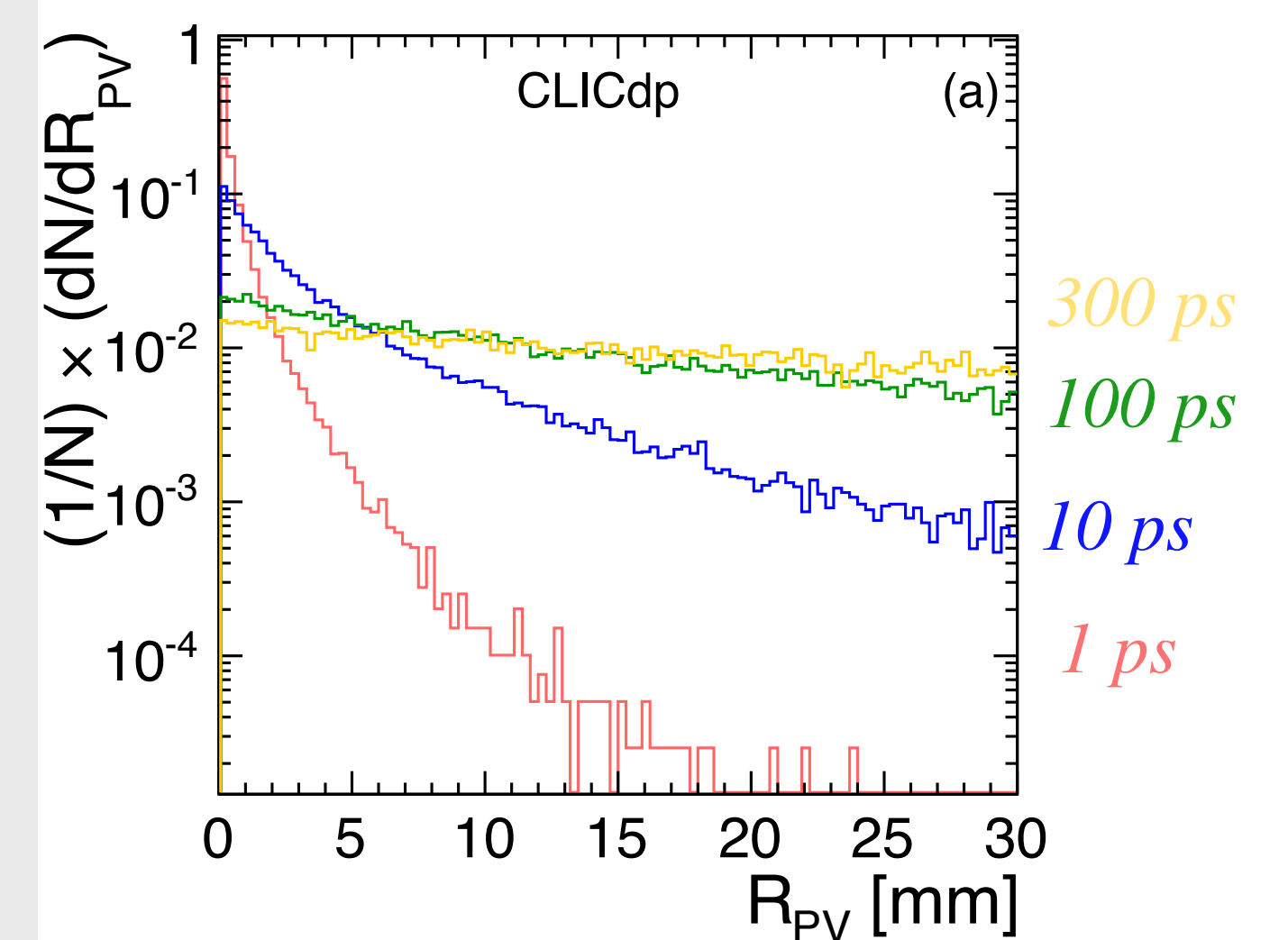
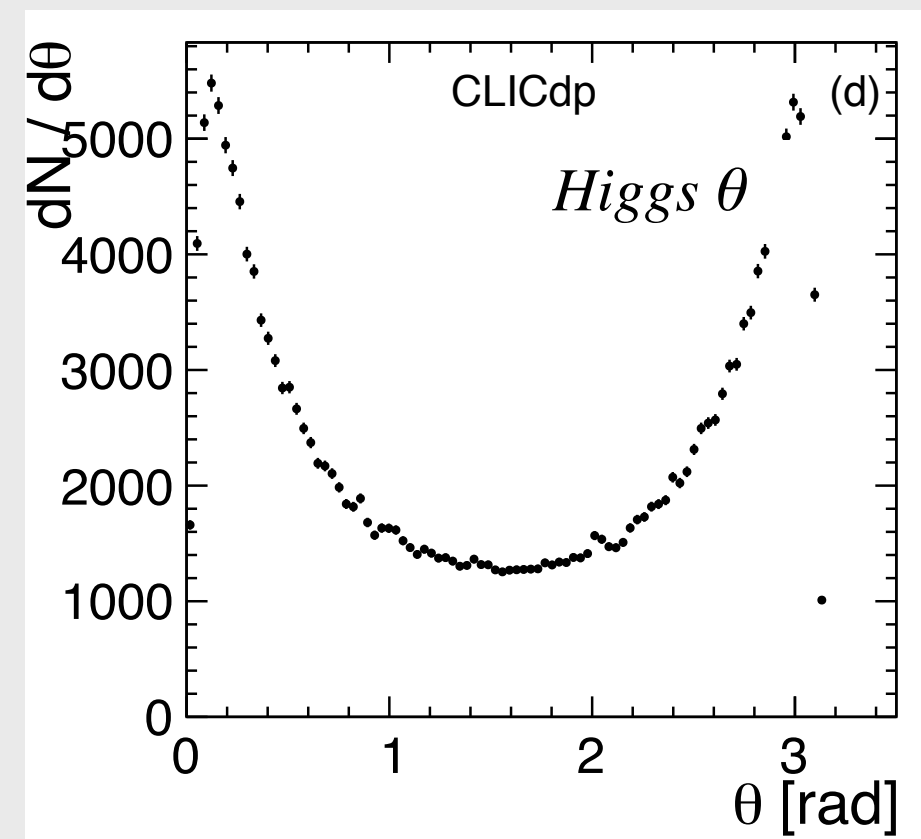
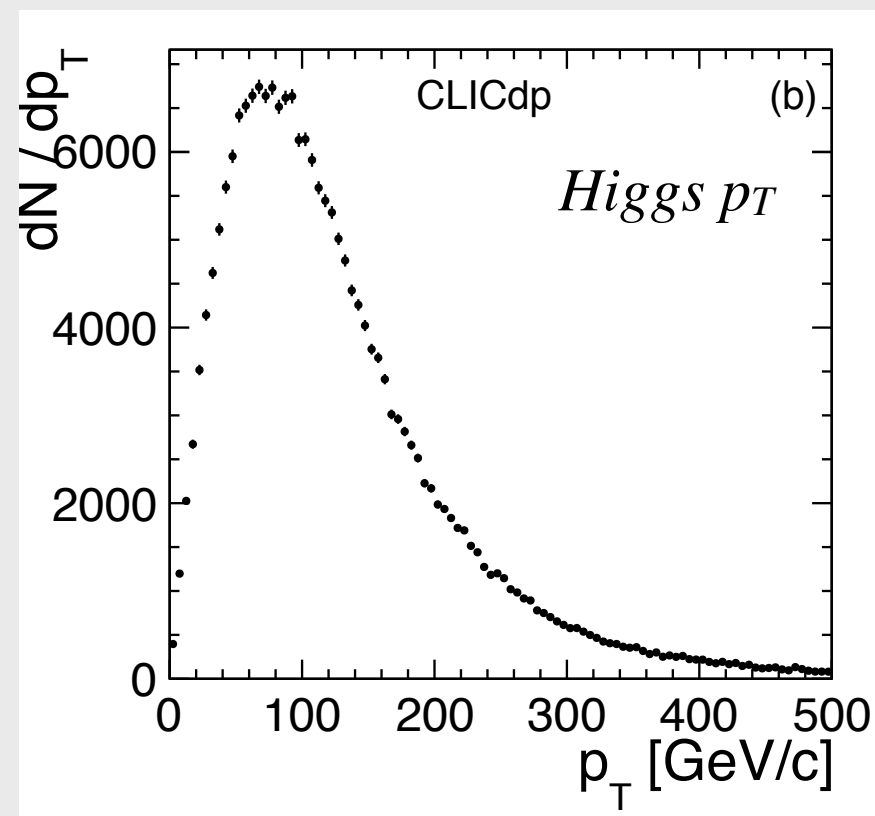
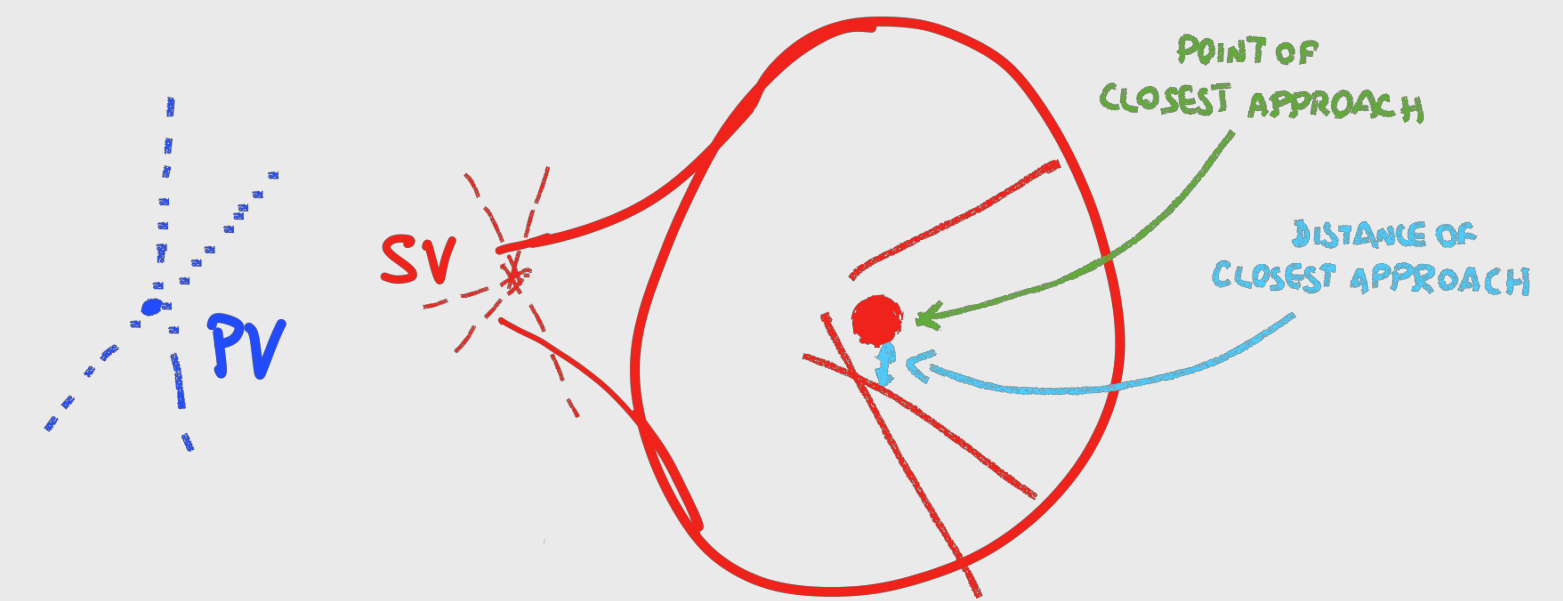
$h \rightarrow \pi_\nu \pi_\nu$

$c\beta\gamma\tau_0$

$\pi_\nu \rightarrow bb$

Point Of Closest Approach +
Distance Of Closest Approach

Process	π_ν^0 lifetime [ps]	π_ν^0 mass [GeV/c ²]	cross section [pb]
$h^0 \rightarrow \pi_\nu^0 \pi_\nu^0$	1,10,100,300	25,35,50	$0.42 \cdot BR$
$e^+e^- \rightarrow q\bar{q}$	-	-	2.95
$e^+e^- \rightarrow q\bar{q}\nu\bar{\nu}$	-	-	0.55
$e^+e^- \rightarrow q\bar{q}q\bar{q}$	-	-	1.32
$e^+e^- \rightarrow q\bar{q}q\bar{q}\nu\bar{\nu}$	-	-	0.07



Hidden Valley Displaced Vertex

CLICDP-NOTE-2018-001

$e^+e^- \rightarrow h \nu\nu$

$h \rightarrow \pi\nu \pi\nu$

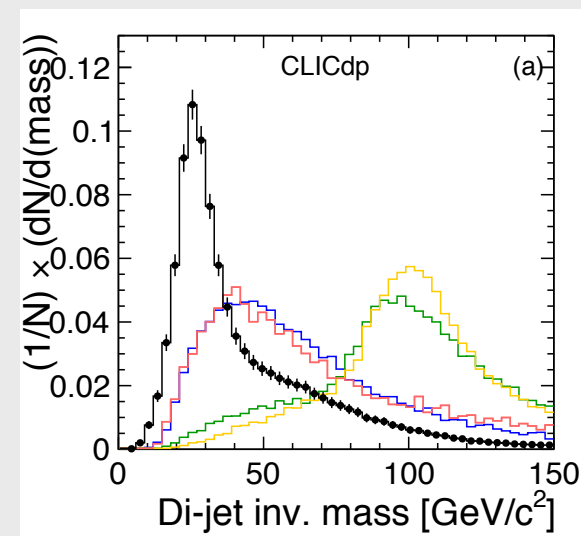
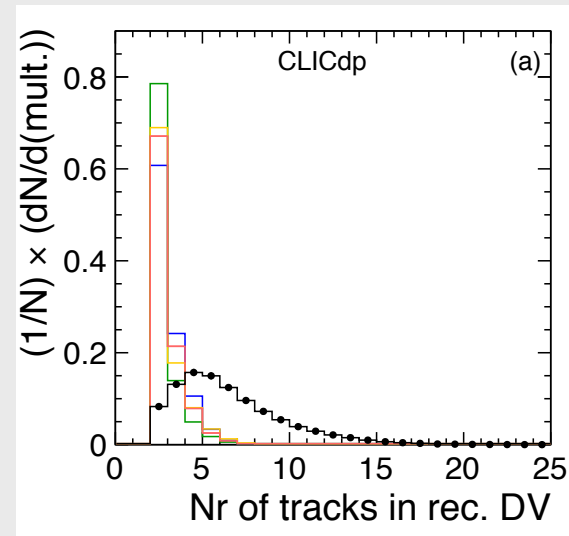
$c\beta\gamma\tau_0$

$\pi\nu \rightarrow bb$

N=4 exclusive k_T jets

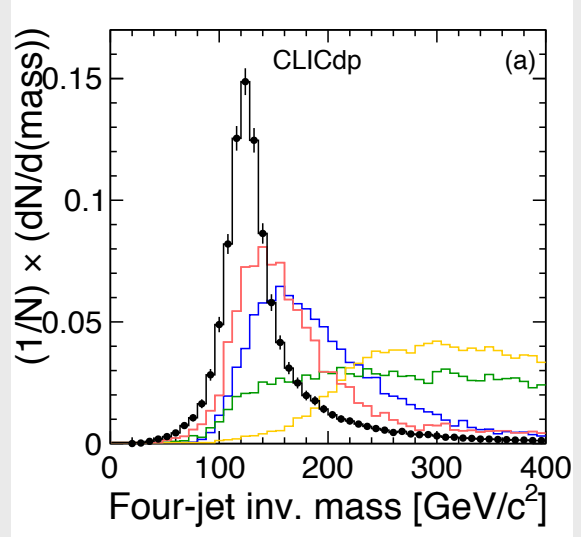
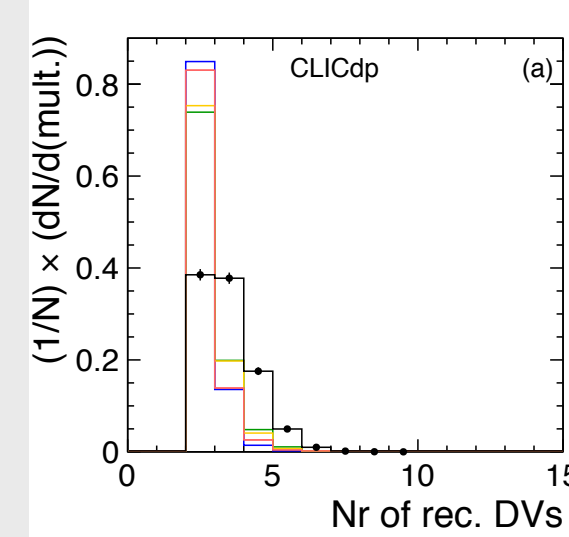
qq $qq\nu\nu$ $qqqq$ $qqqq\nu\nu$

of tracks



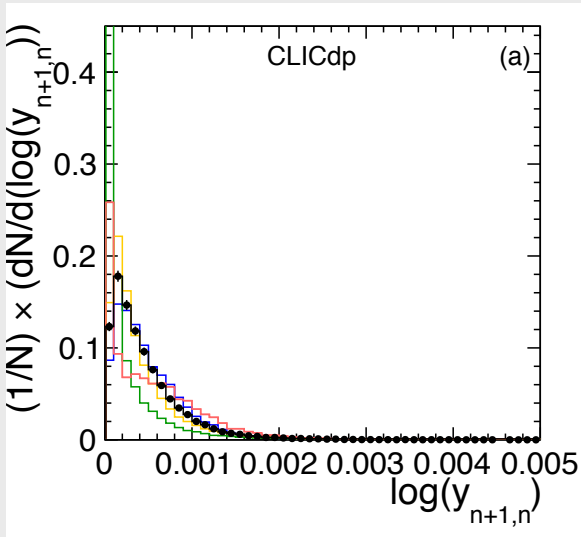
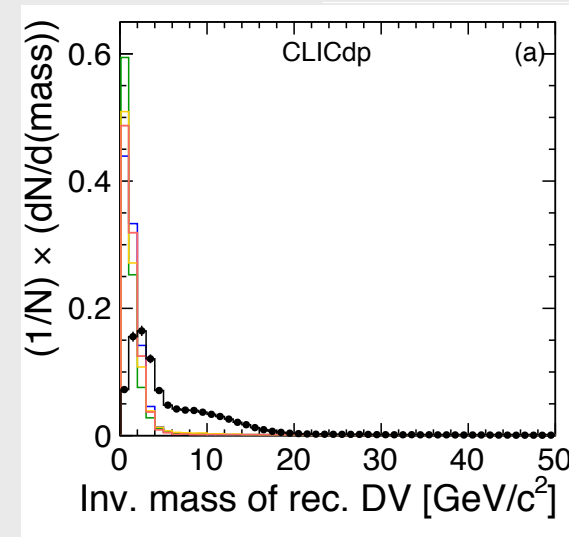
Mass of jj

of DV



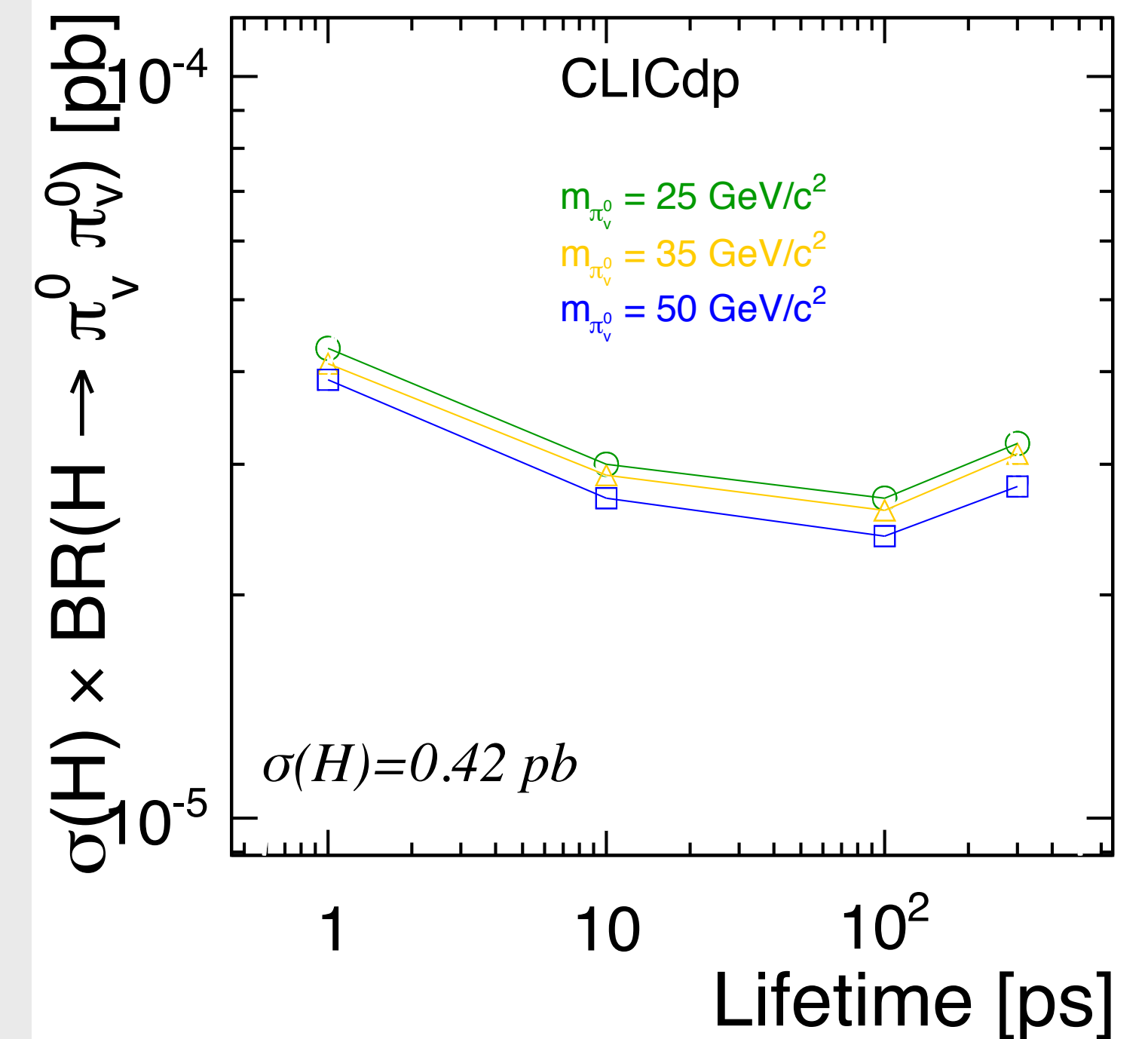
Mass of 4j

Mass of DV



Jets y_{34} and y_{23}

Boosted Decision Tree: $\epsilon_S \geq 0.1$



re-use of simplified
models

Heavy Higgs Displaced Decay

BASED ON CLICDP-NOTE-2018-001

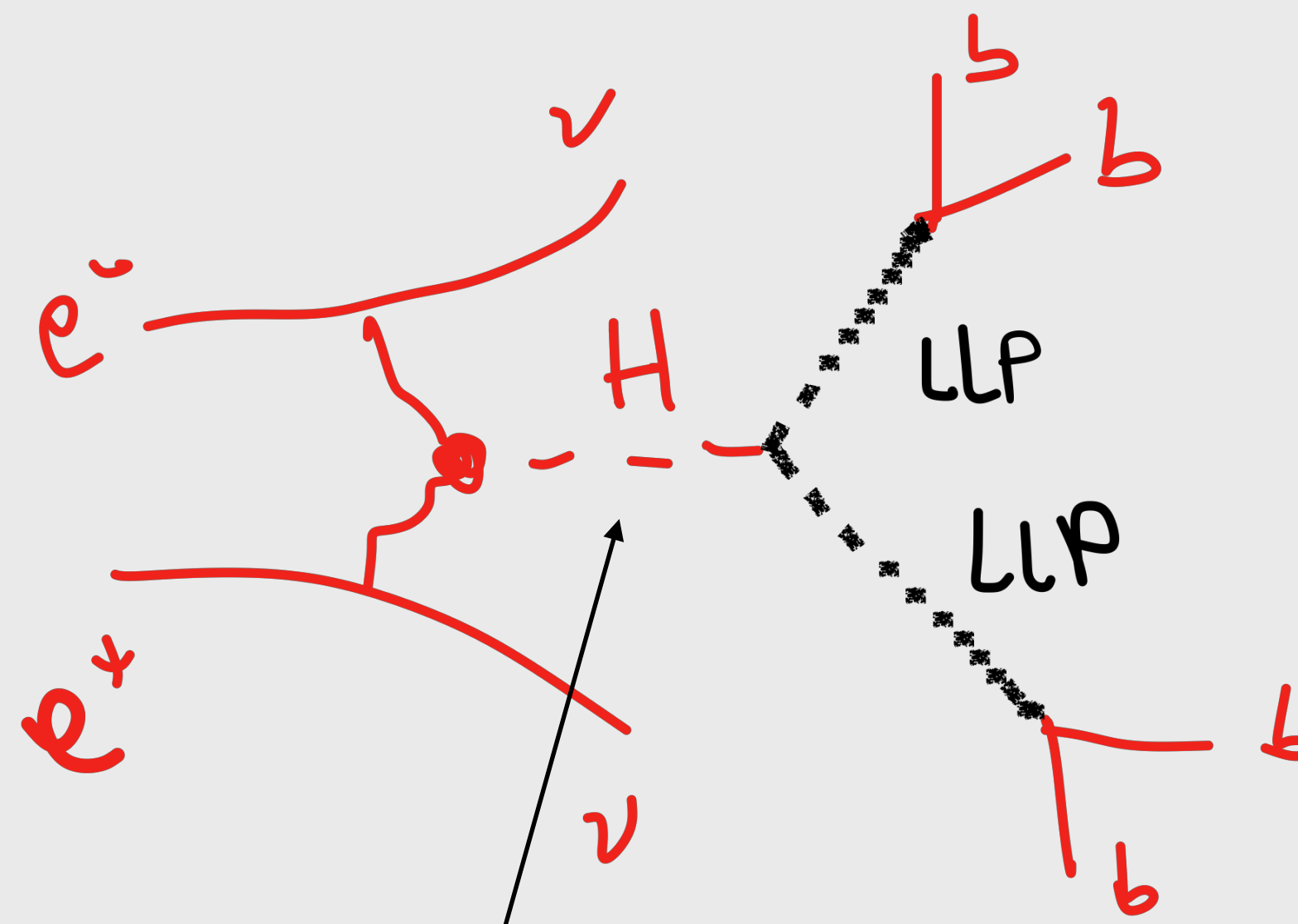
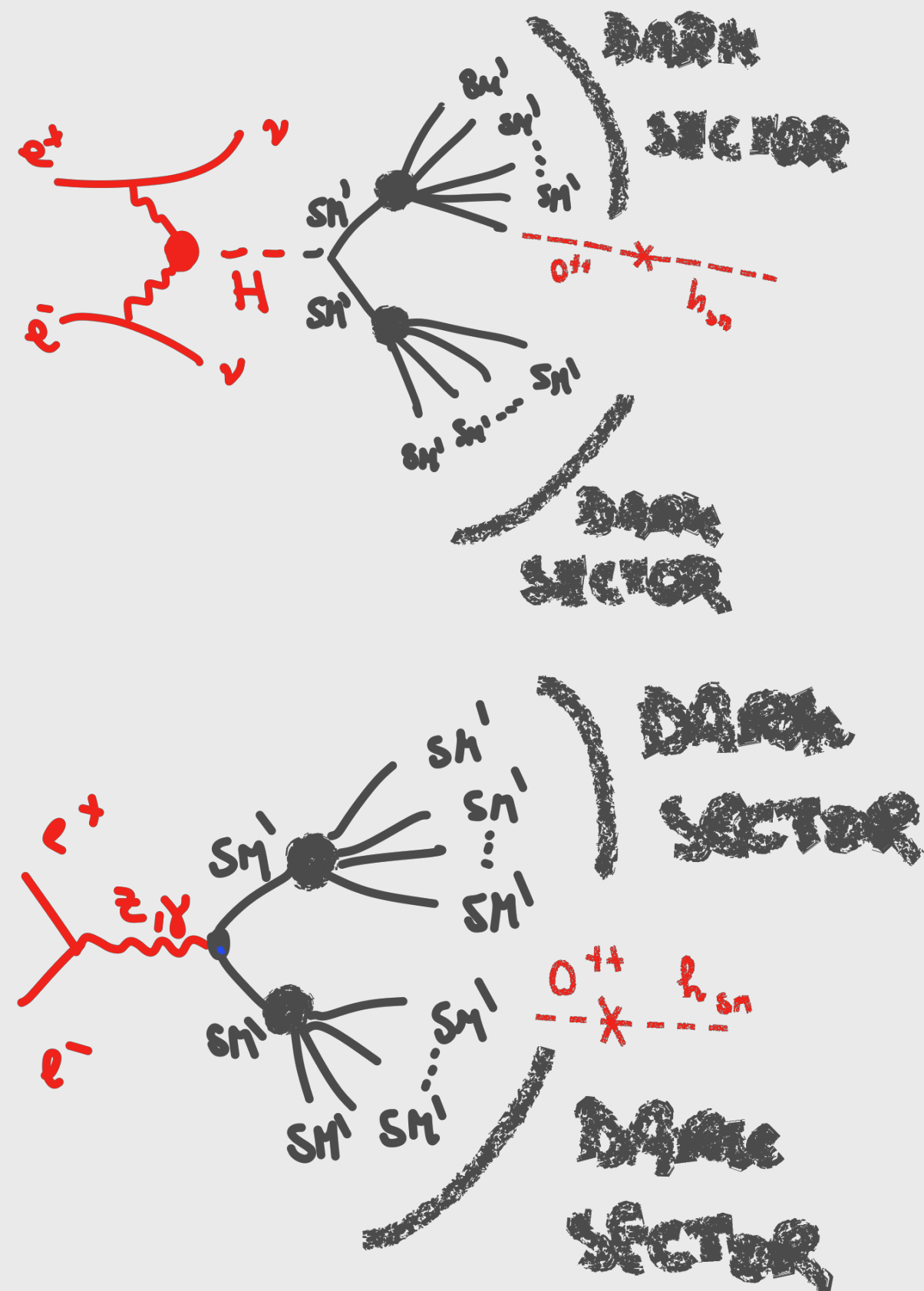
$$e^+e^- \rightarrow H \nu\nu$$

$$H \rightarrow LLP LLP$$

$c\beta\gamma\tau_0$

$$LLP \rightarrow bb$$

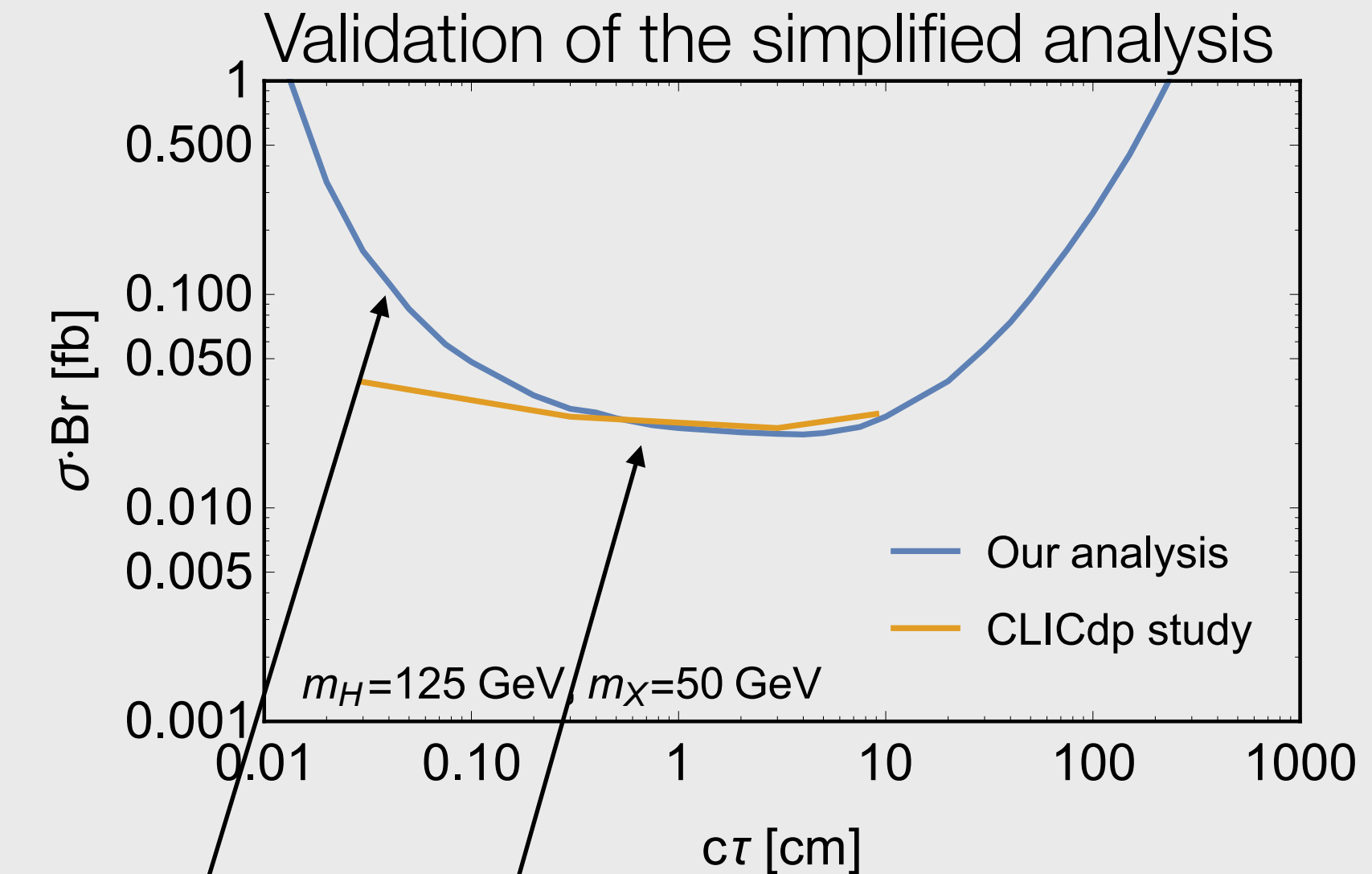
“Neutral Naturalness” scenarios: Folded SUSY, fraternal Twin Higgs, ...



$m_H = 125, 200, 400, 600, 800, 1000$ GeV

CLIC 3 TeV simplified analysis:

- Interaction point Significance > 16
- “Loose” 90% b-tag efficiency
- $\Delta R_{bb} > 0.5$ for isolation
- 0.5 efficiency for $N_{\text{track}} > 5$



In general conservative good agreement

Heavy Higgs Displaced Decay

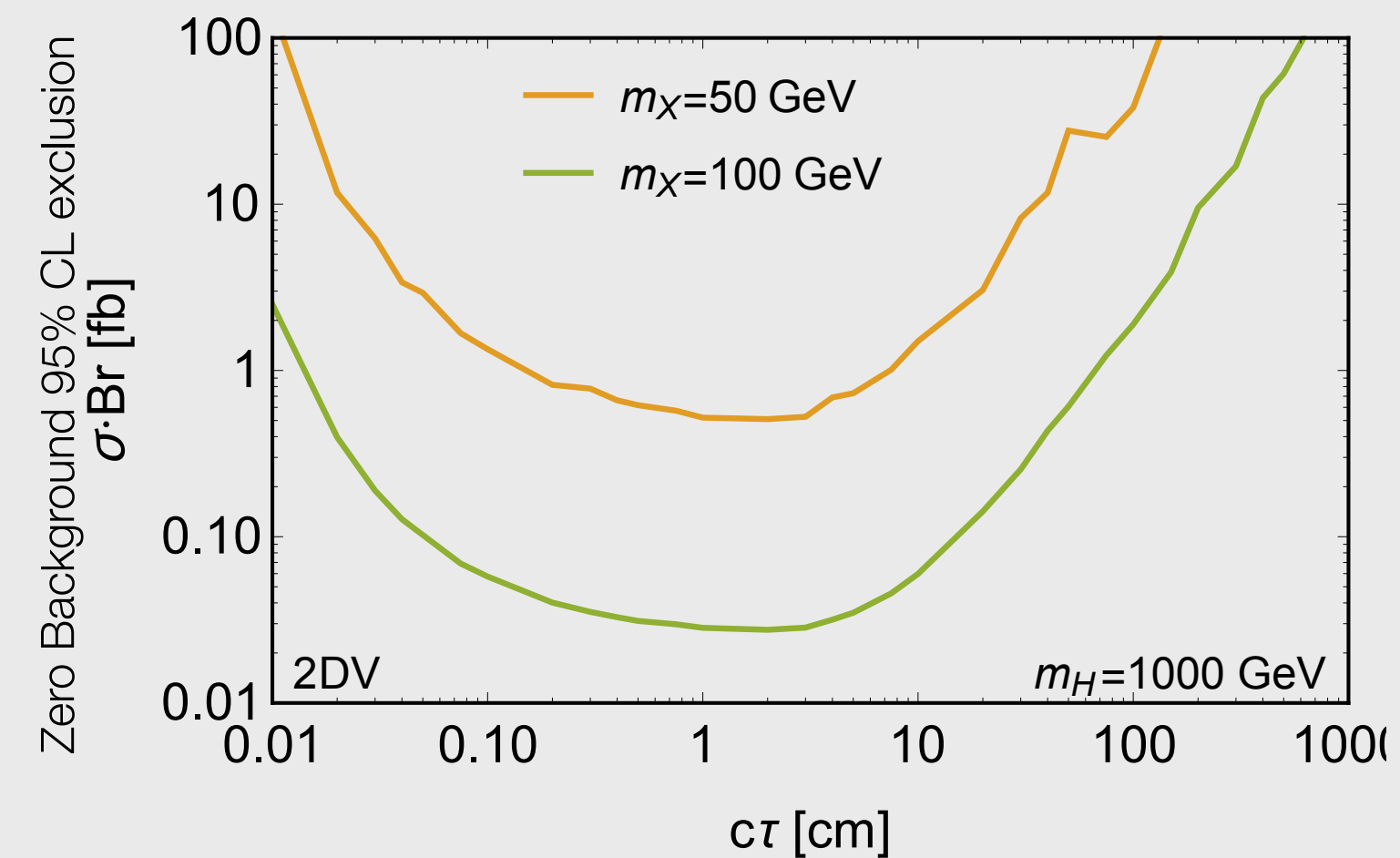
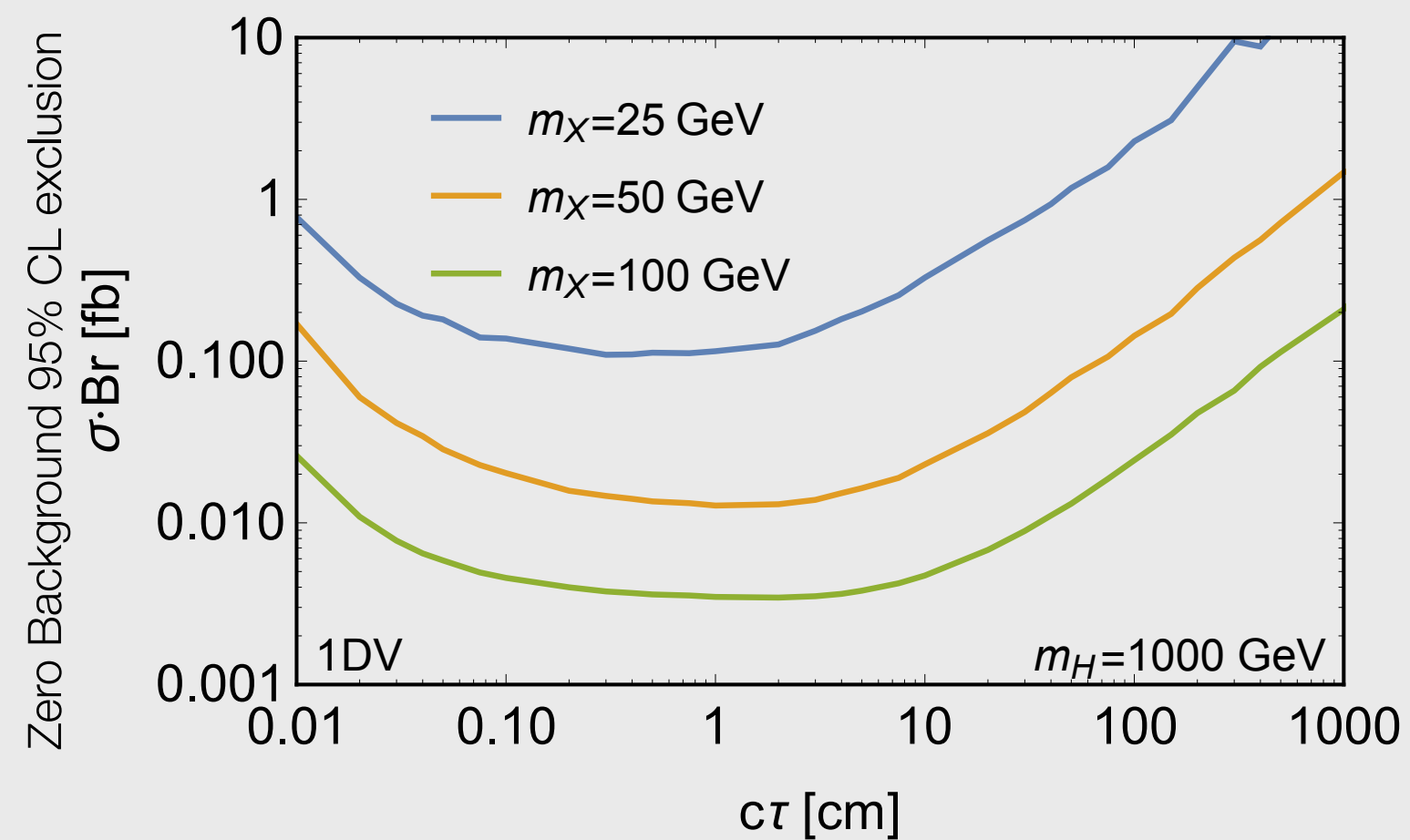
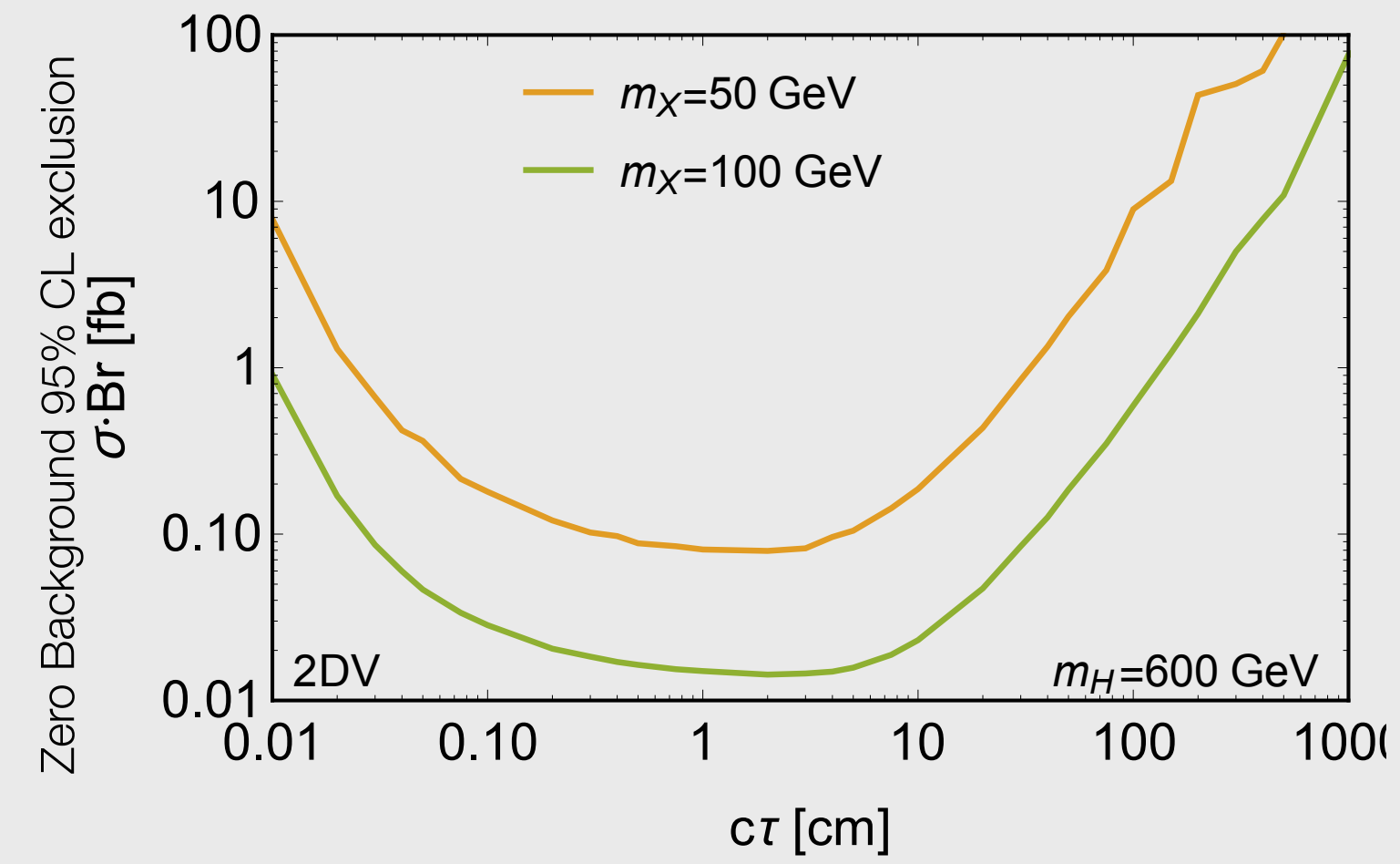
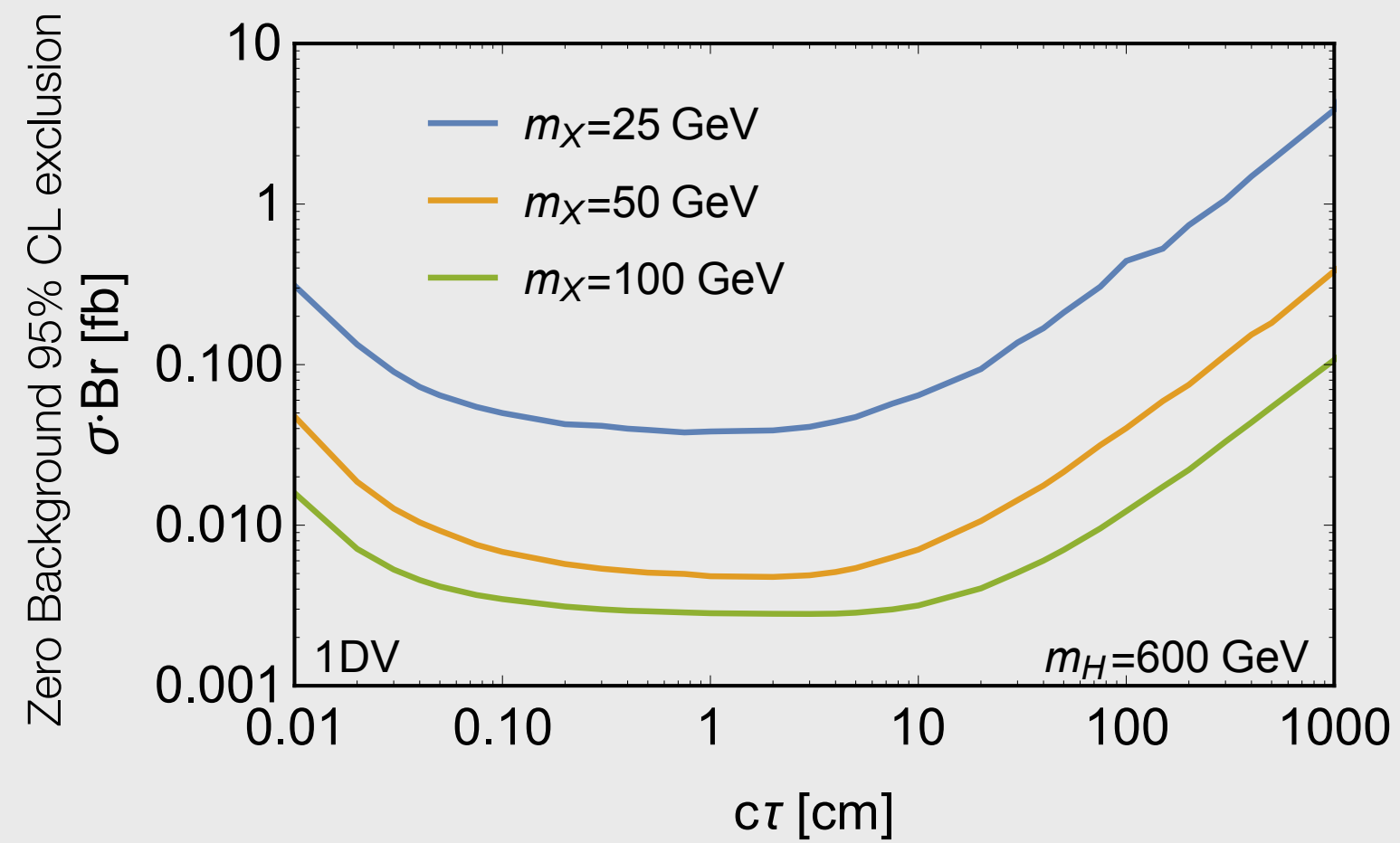
BASED ON CLICDP-NOTE-2018-001

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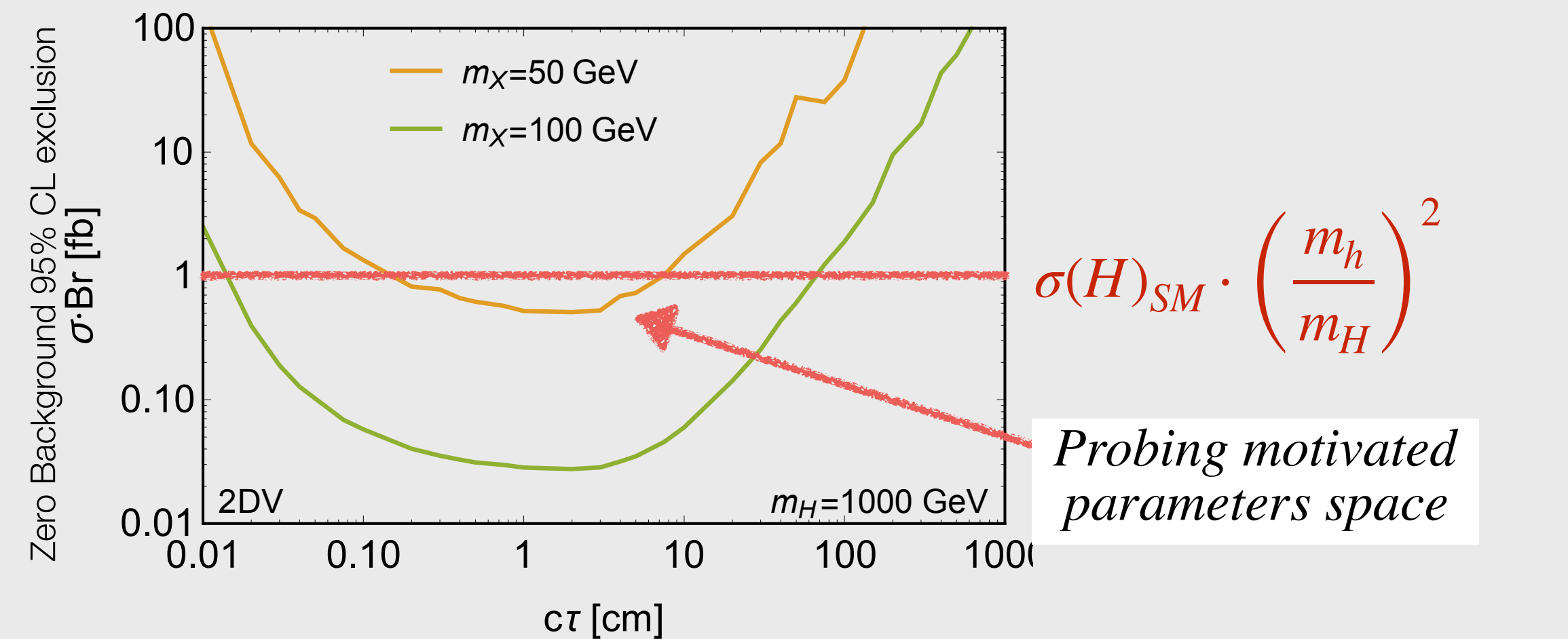
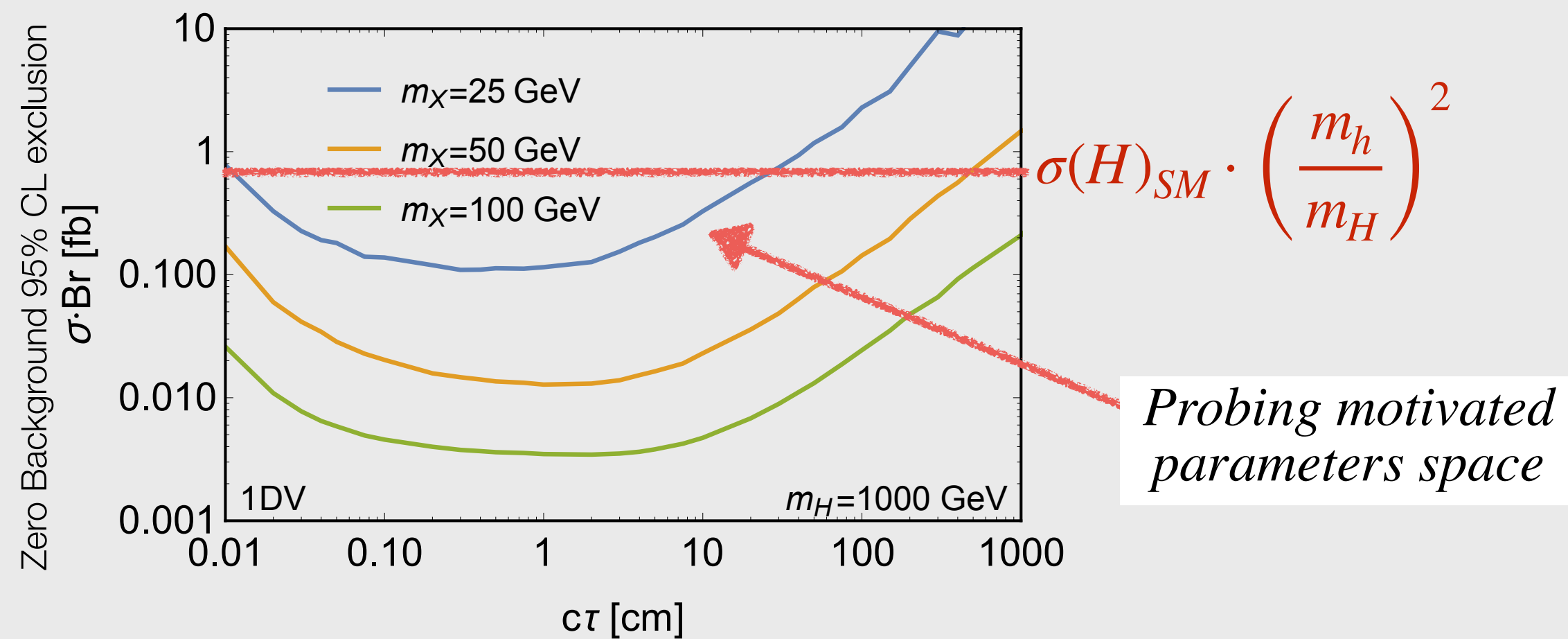
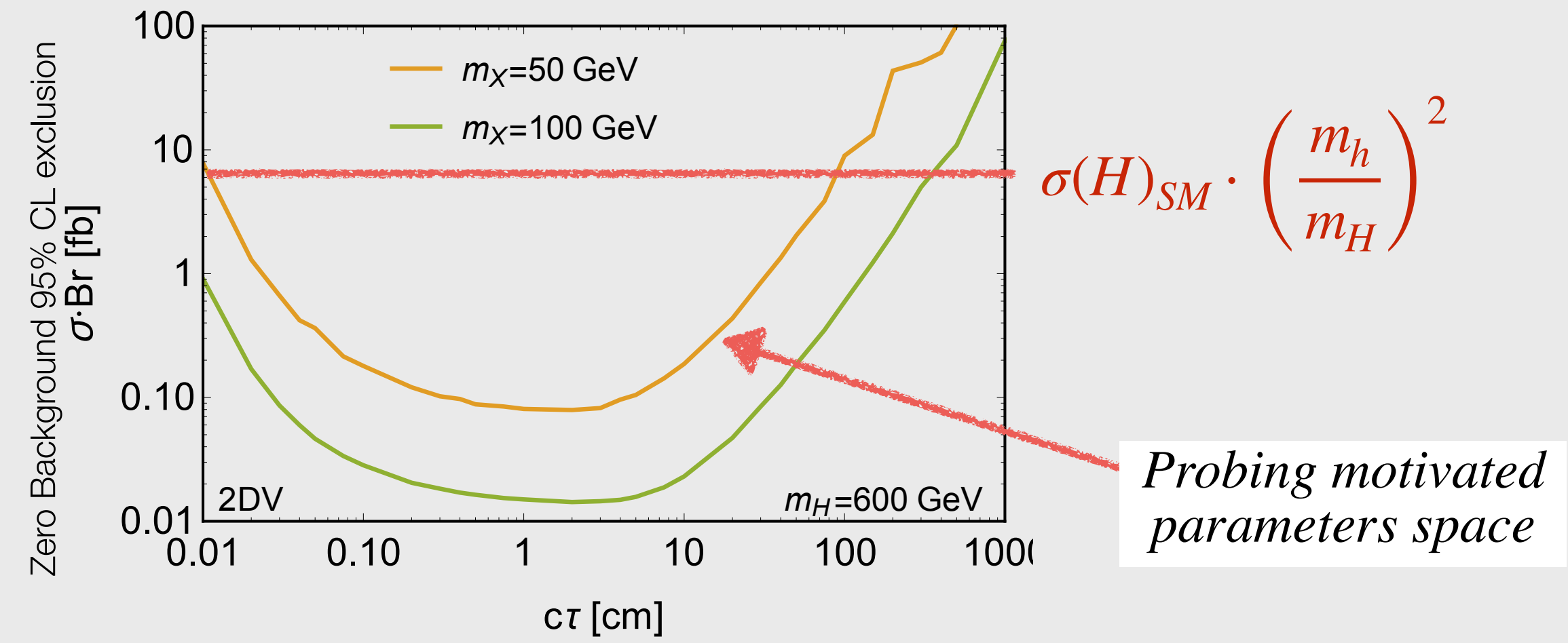
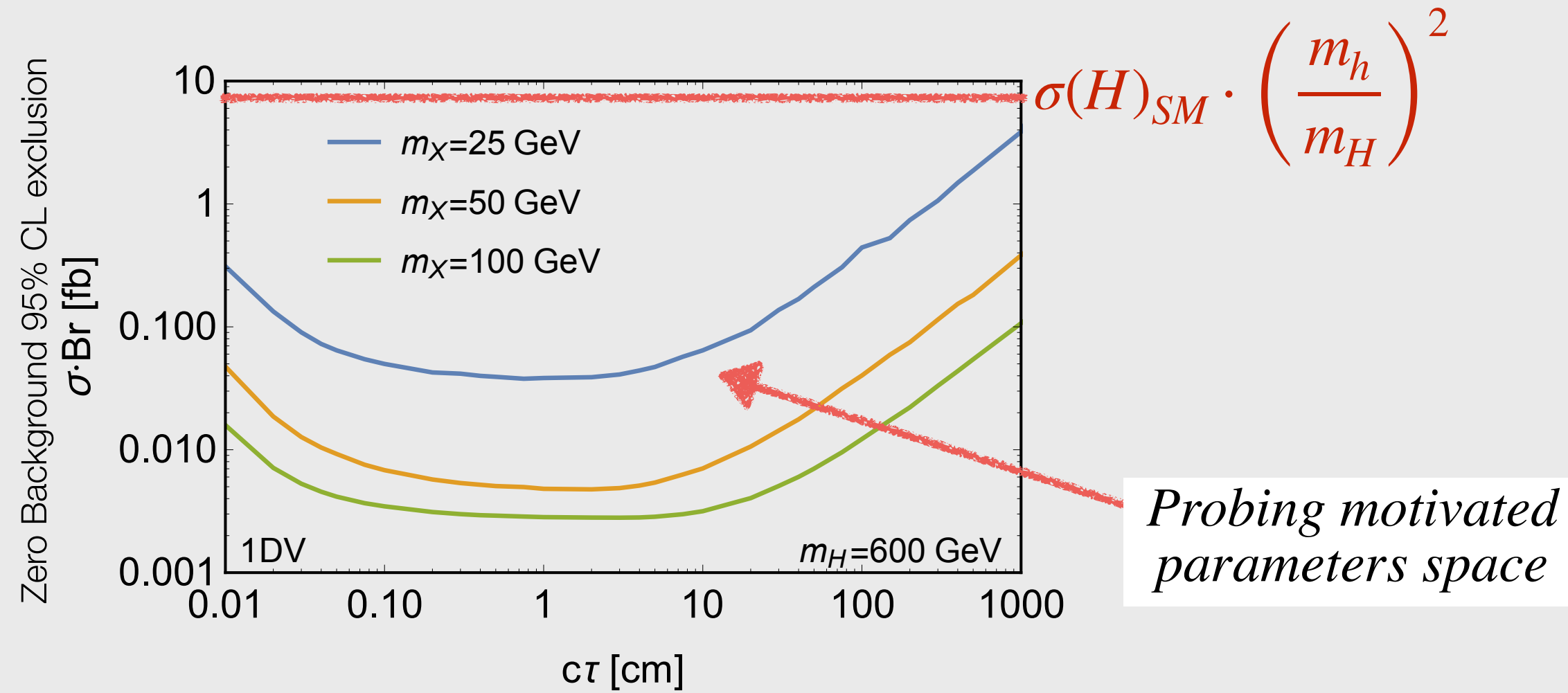
BASED ON CLICDP-NOTE-2018-001

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$H \rightarrow LLP LLP$

$c\beta\gamma\tau_0$

$LLP \rightarrow bb$



Baryogenesis

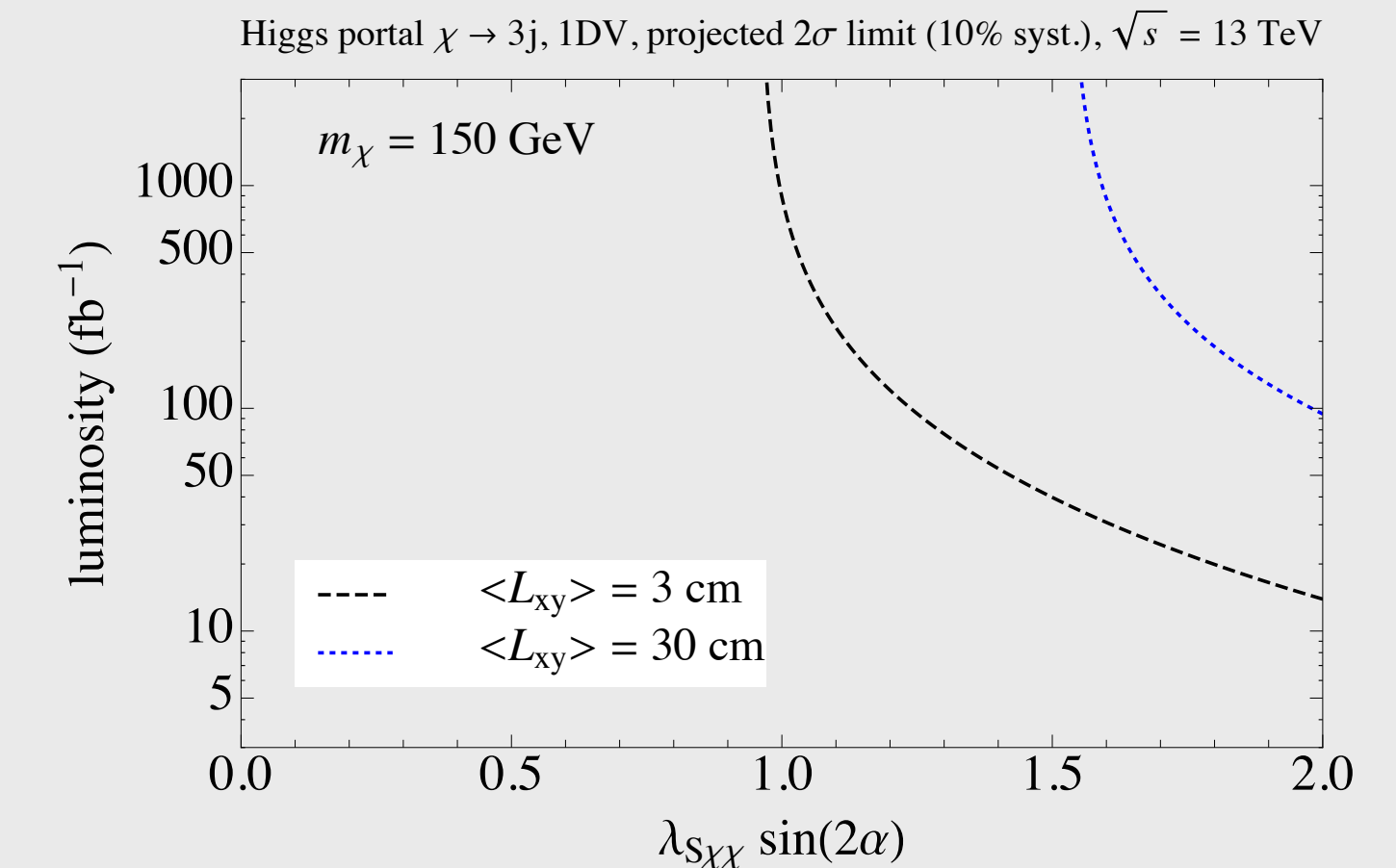
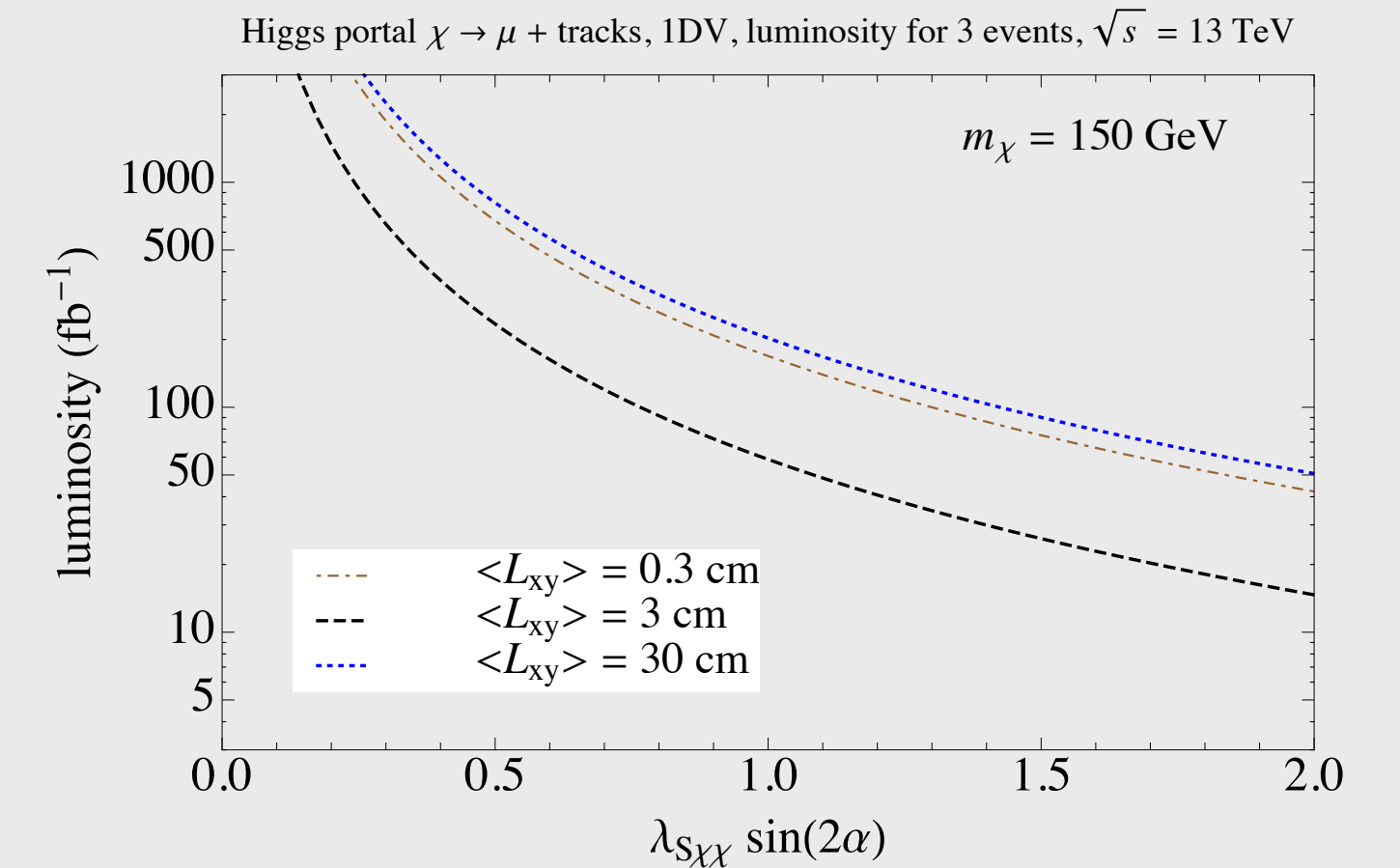
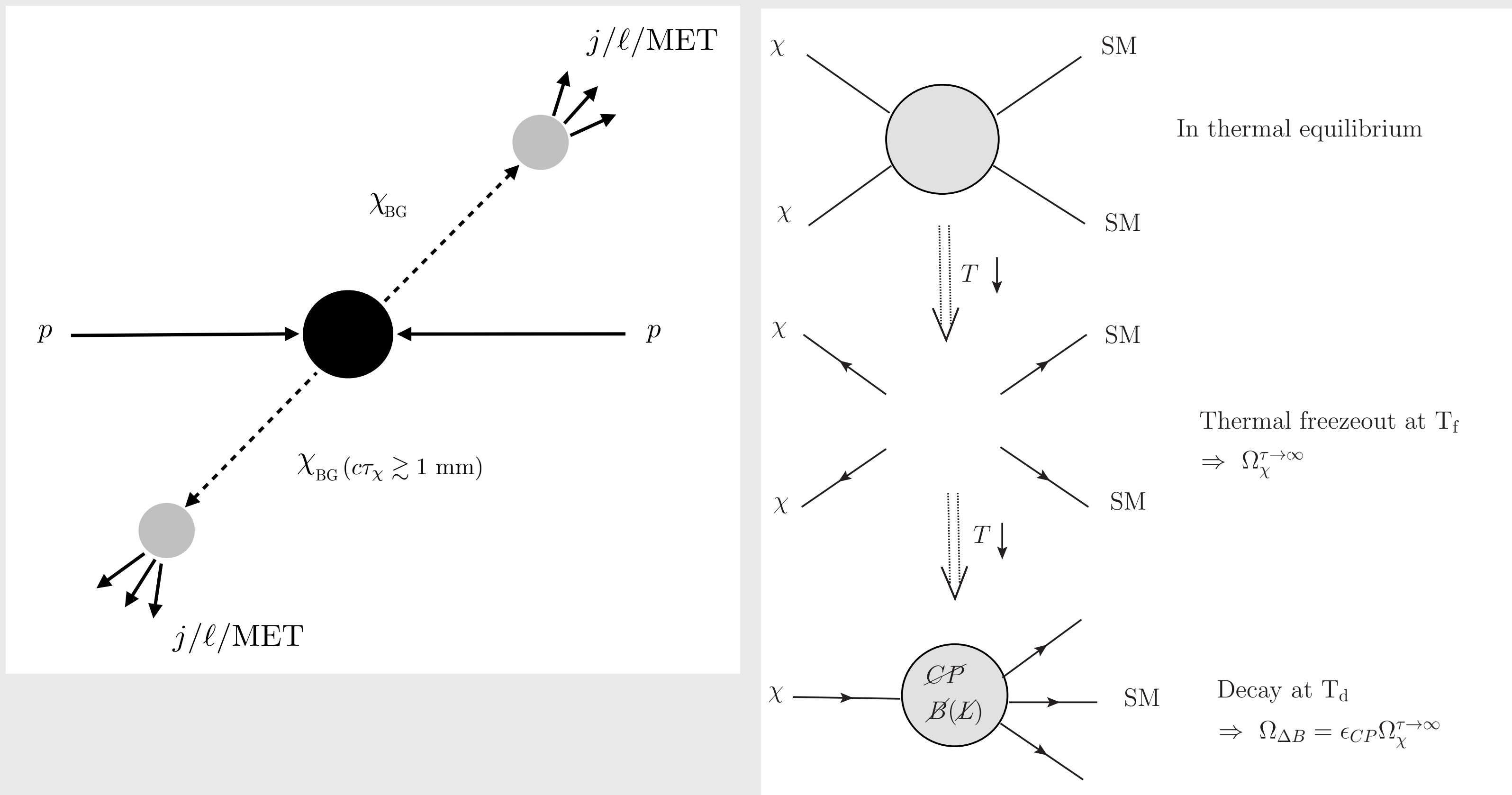
WIMP Baryogenesis

SMALL COUPLINGS

N_B FROM OUT-OF-EQUILIBRIUM DECAY OF WOULD-BE WIMP

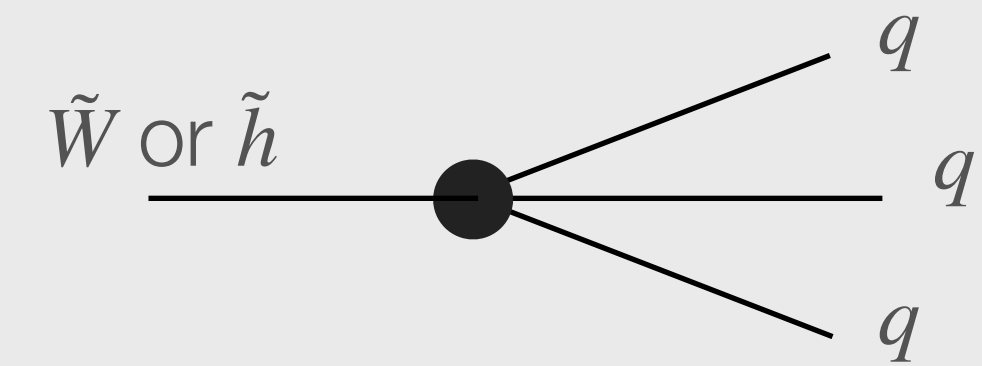
Would-be WIMP can be SM-singlet or SM-charged

$$\mathcal{L}_{\text{prod}}^{\text{singlet}} \supset \frac{c_H}{\Lambda_H} \chi^2 |H|^2 + \frac{c_q}{\Lambda_q^2} (\bar{\chi} \Gamma \chi) (\bar{q} \Gamma' q) + \frac{c_g}{\Lambda_g^3} \chi^2 (G_{\mu\nu})^2 + \dots,$$



Possible $e^+e^- \rightarrow \nu\nu + 2 \text{ DV}$ signal

WIMP Baryogenesis



BASED ON CLICDP-NOTE-2018-001

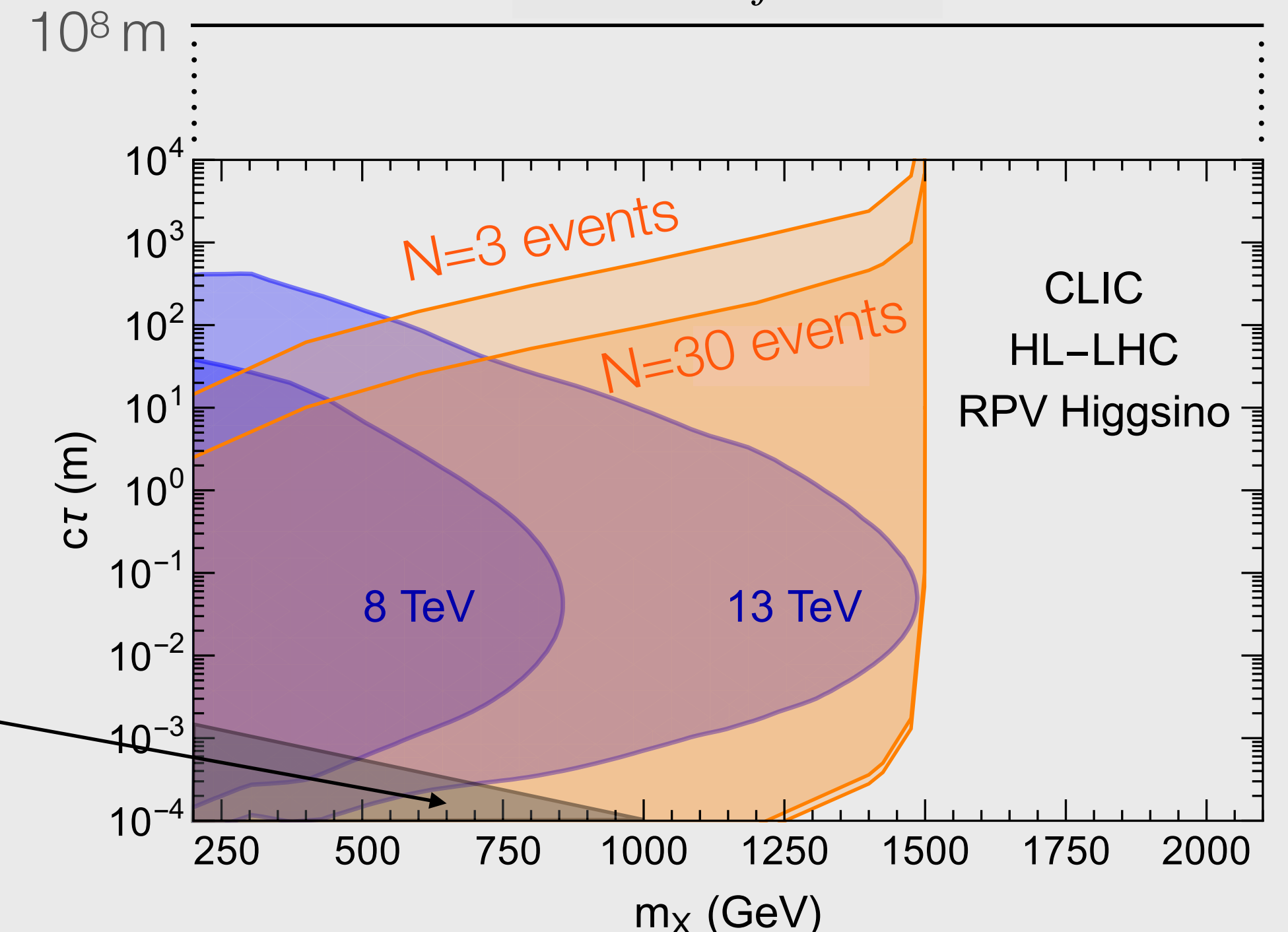
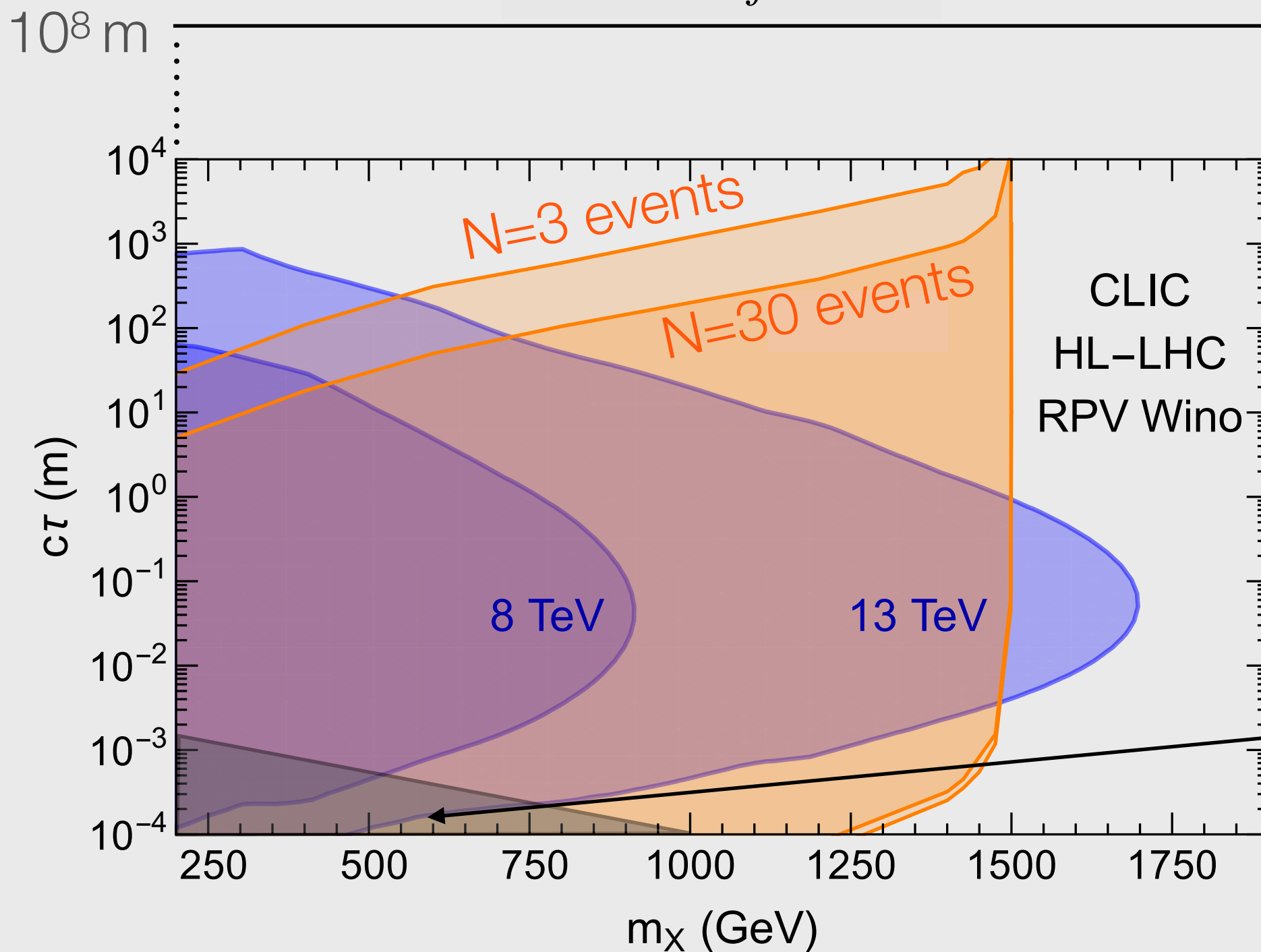
DISPLACED DECAYS TO HADRONS

$$e^+e^- \rightarrow \chi\chi \quad \nu\nu \rightarrow \nu\nu + 2 (DV \rightarrow jjj)_\chi$$

assume 100% efficient vertex finder in $3 \cdot 10^{-3} m < c\tau < 0.1 m$ (CLICdp-Note-2018-001)

$c\tau < 10^8 m$ for BBN

$c\tau < 10^8 m$ for BBN



Cosmology requires

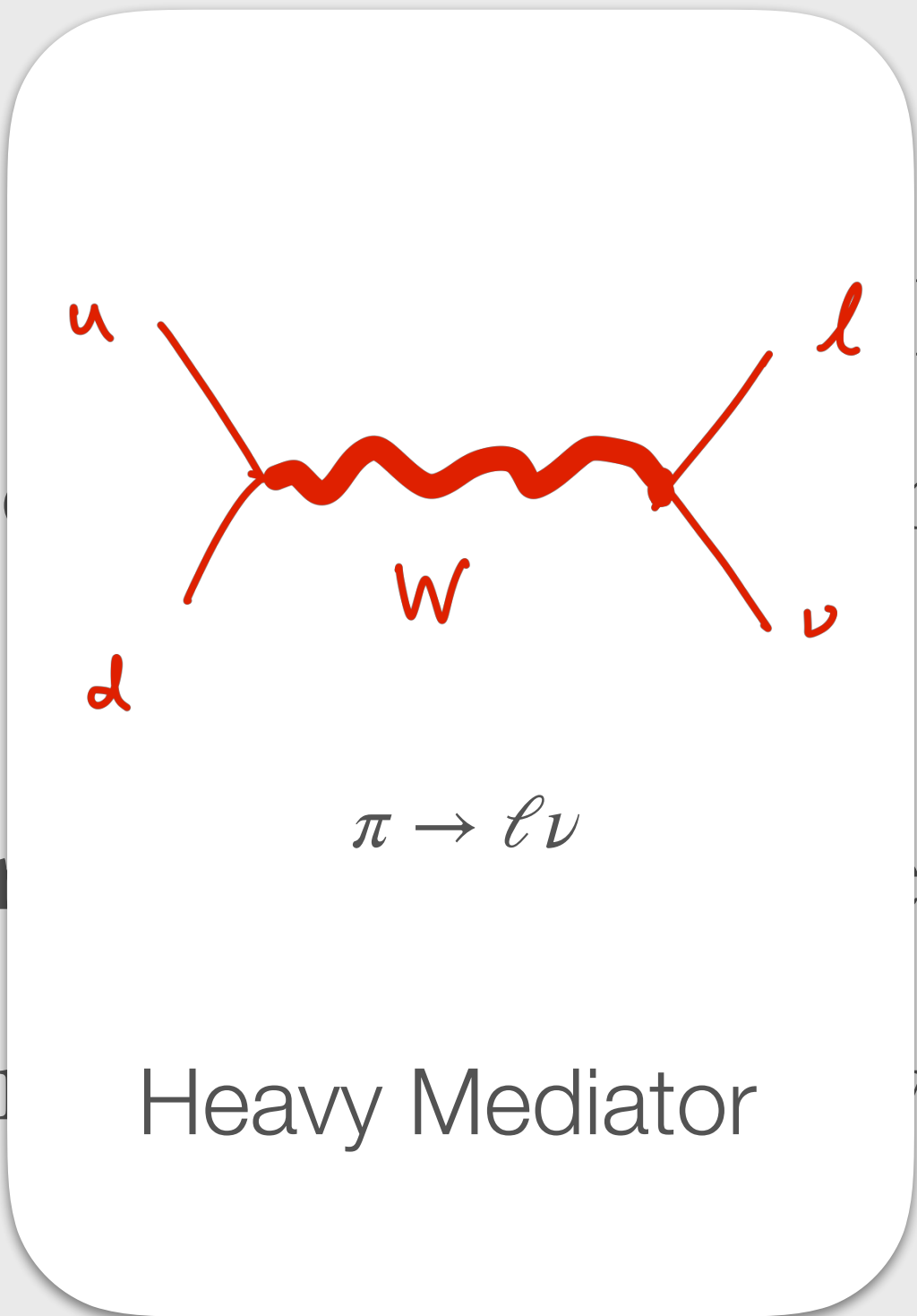
- before nucleosynthesis
 $c\tau_\chi < 10^8 m$
- out-of-equilibrium
 $\tau_\chi > 1/H(T \sim m_\chi)$

- Standard “targets” such as vanilla SUSY, compositeness of Higgs and other states, sub-TeV WIMPs are all being probed and are under a fair amount of pressure
- Motivations for new physics to be out there are stronger than ever

- New physics can be accessible at future colliders, but no “no-loose theorem” as for LHC
 - ★ next generation of new physics exploration has to be necessarily “broad-band”
 - ★ macro/meso-scopic life-times are possible and motivated from several standpoints

- New physics can be accessible at future colliders, but no “no-loose theorem” as for LHC

- ★ next physics exploration
- has to “road-band”
- ★ mac -times are possible
- and Heavy Mediator eral standpoints



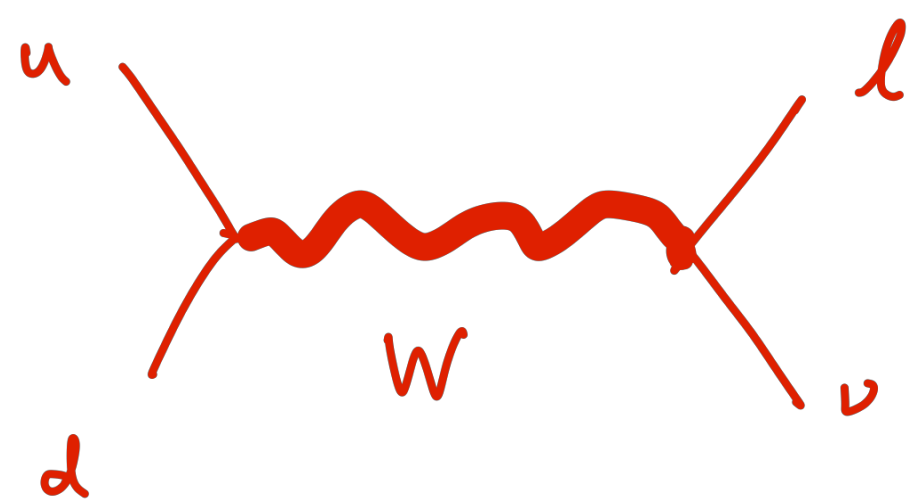
- New physics can be accessible at future colliders, but no “no-loose theorem” as for LHC

★ next

has t

★ macr

and



$$\pi \rightarrow \ell \nu$$

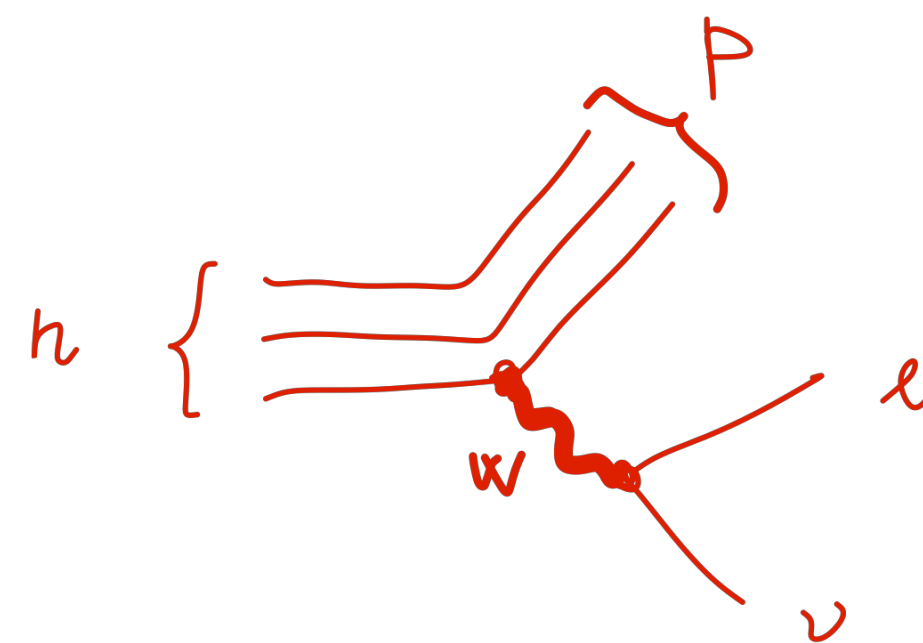
Heavy Mediator

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$$n \rightarrow p e \bar{\nu}$$

Small Phase-Space

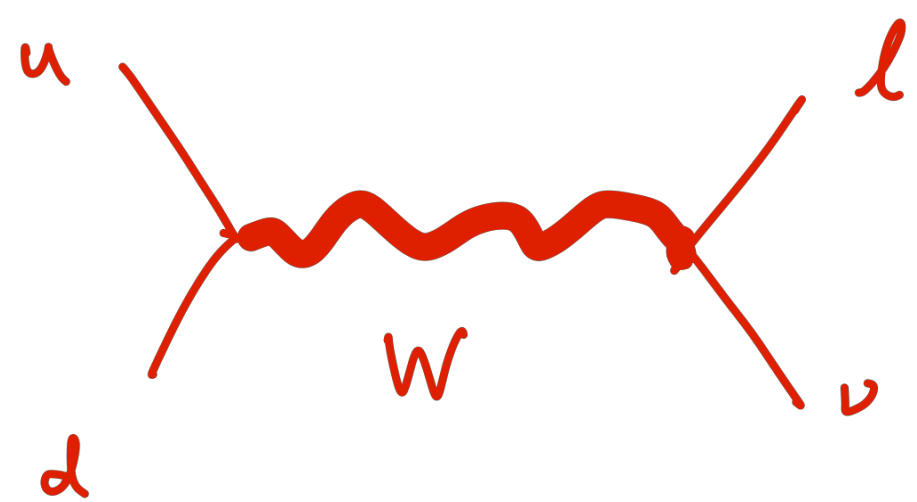
- New physics can be accessible at future colliders, but no “no-loose theorem” as for LHC

★ next

has t

★ macr

and



$$\pi \rightarrow \ell \nu$$

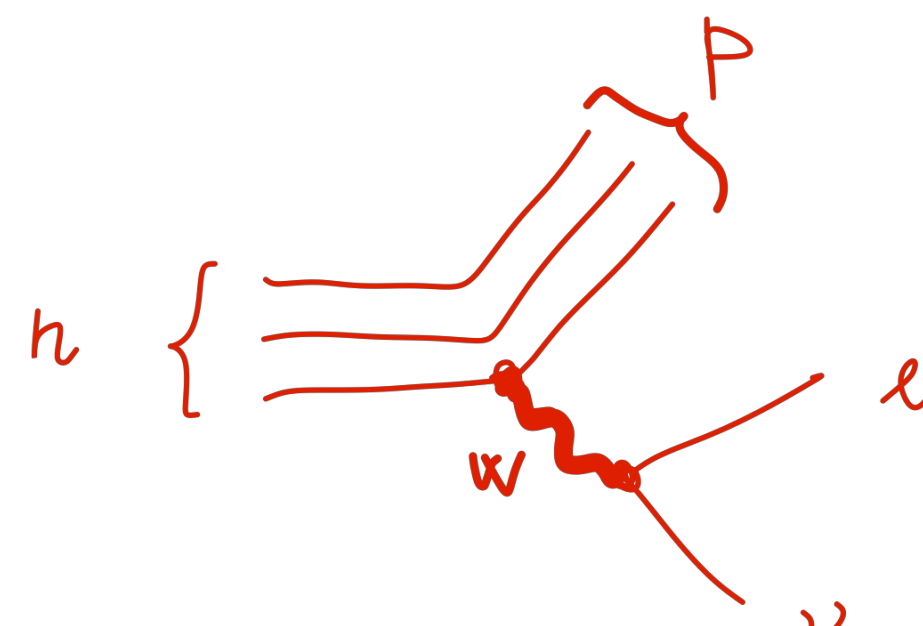
Heavy Mediator

p

ro

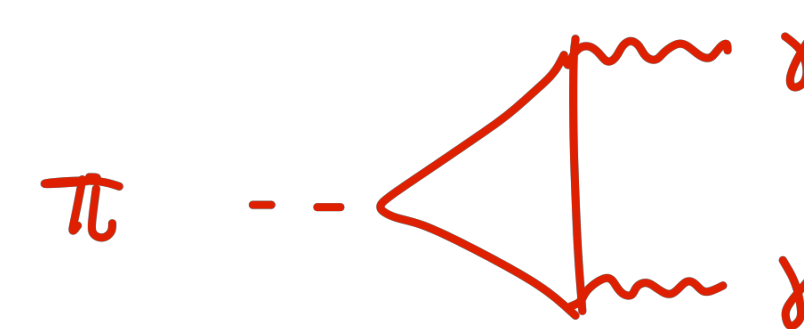
-1

re



$$n \rightarrow p e \bar{\nu}$$

Small Phase-Space



$$\pi \rightarrow \gamma \gamma$$

Small Coupling

CONCLUSIONS

WORK DONE

- long-lived decays of a 125 GeV Higgs (CLICdp)
- long-lived decays of a Heavy Higgs (pheno-level Yellow Report)
- long-lived decays of a “baryo-genitor” (pheno-level Yellow Report)
- short-tracks from a Higgsino DM candidate (pheno-level Yellow Report)

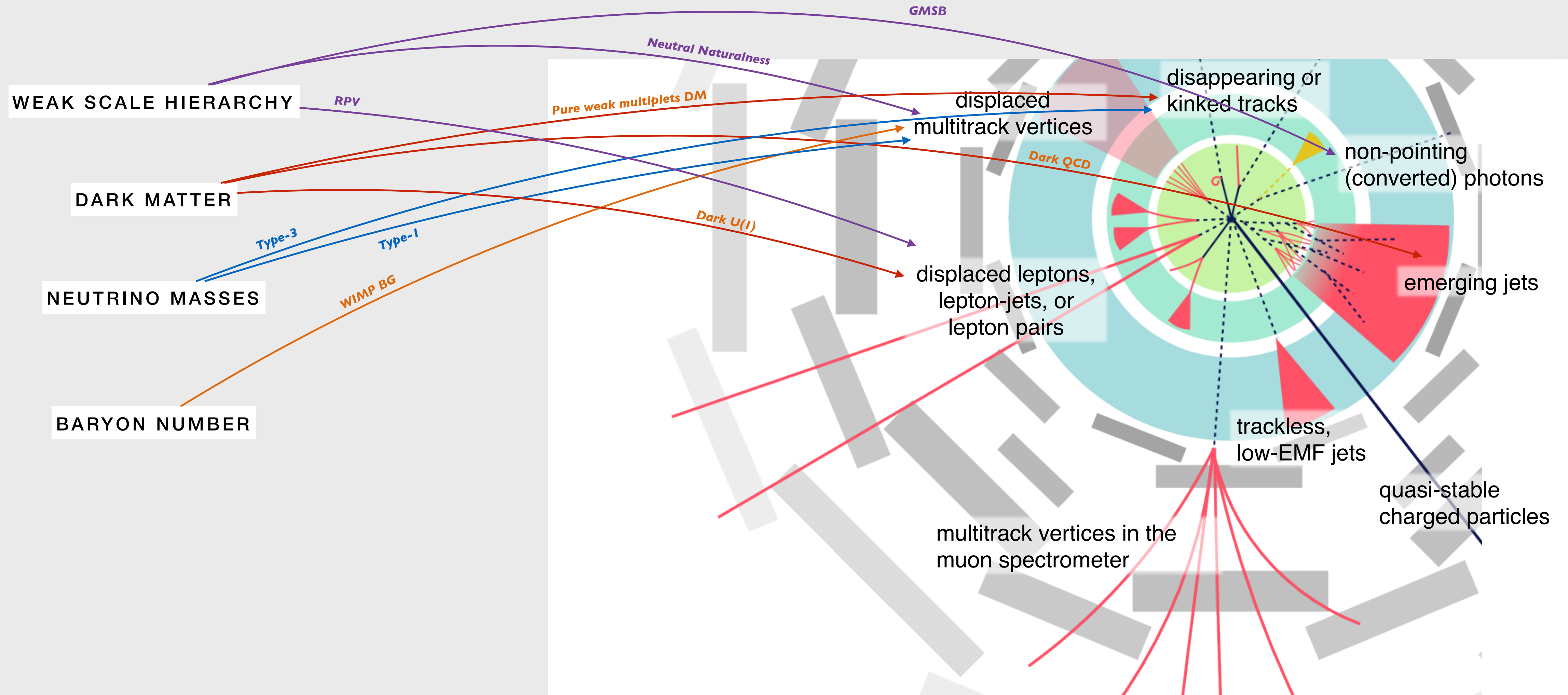
CONCLUSIONS

ACTION ITEMS

- short-tracks from a Higgsino DM candidate (CLICdp inputs, especially on background)
- long-lived decays of a Heavy Higgs (CLICdp validation of recast)
- long-lived decays of a “baryo-genitor” (CLICdp input and effort for the 3j DV)
- long-lived decays of a Heavy Neutrino (open chapter)
- long-lived decays of a dark photons and dark Z' (open chapter)

CONCLUSIONS

MORE SIGNATURES ON THE TABLE*



Heather Russell, McGill University

24 April 2017

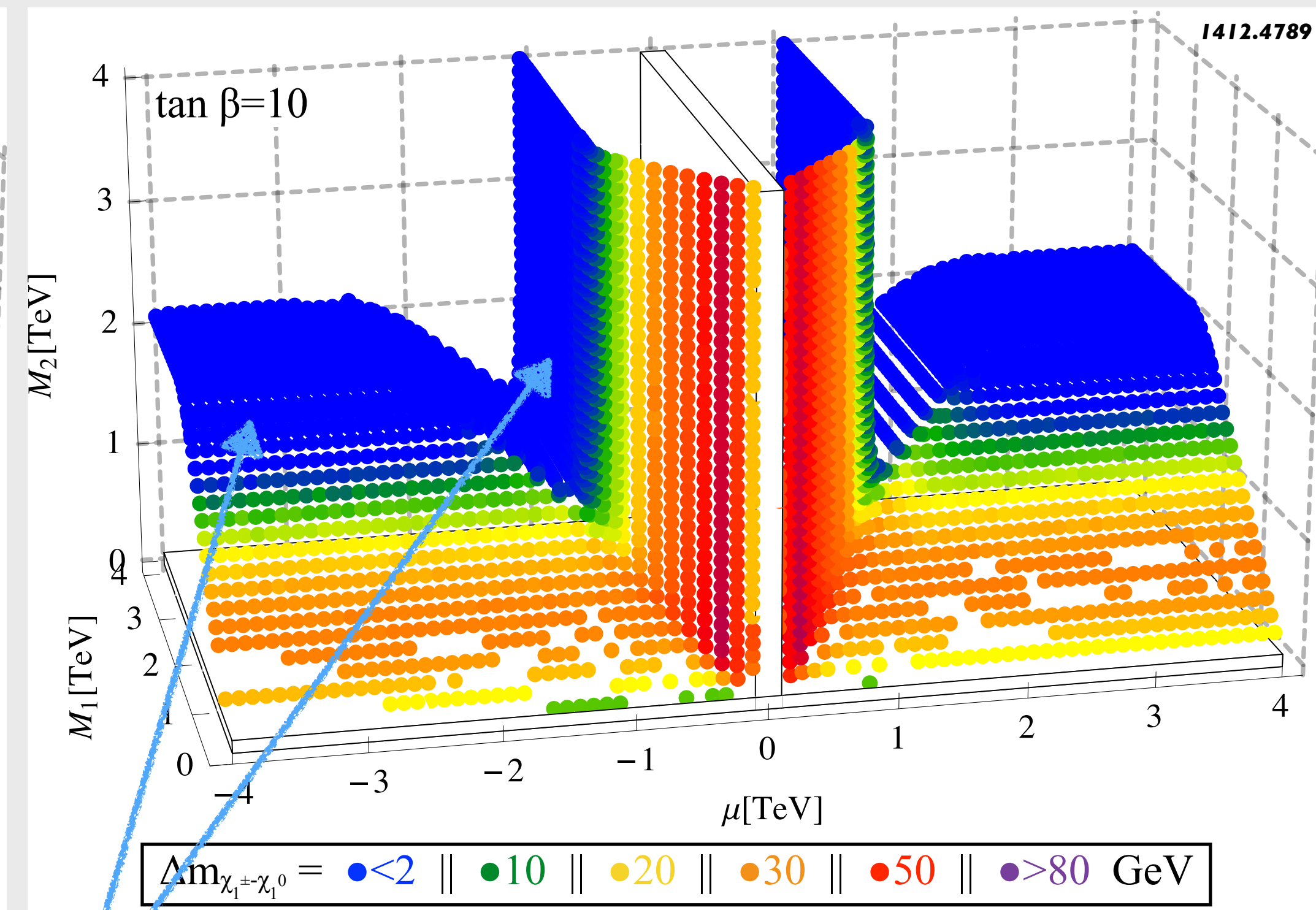
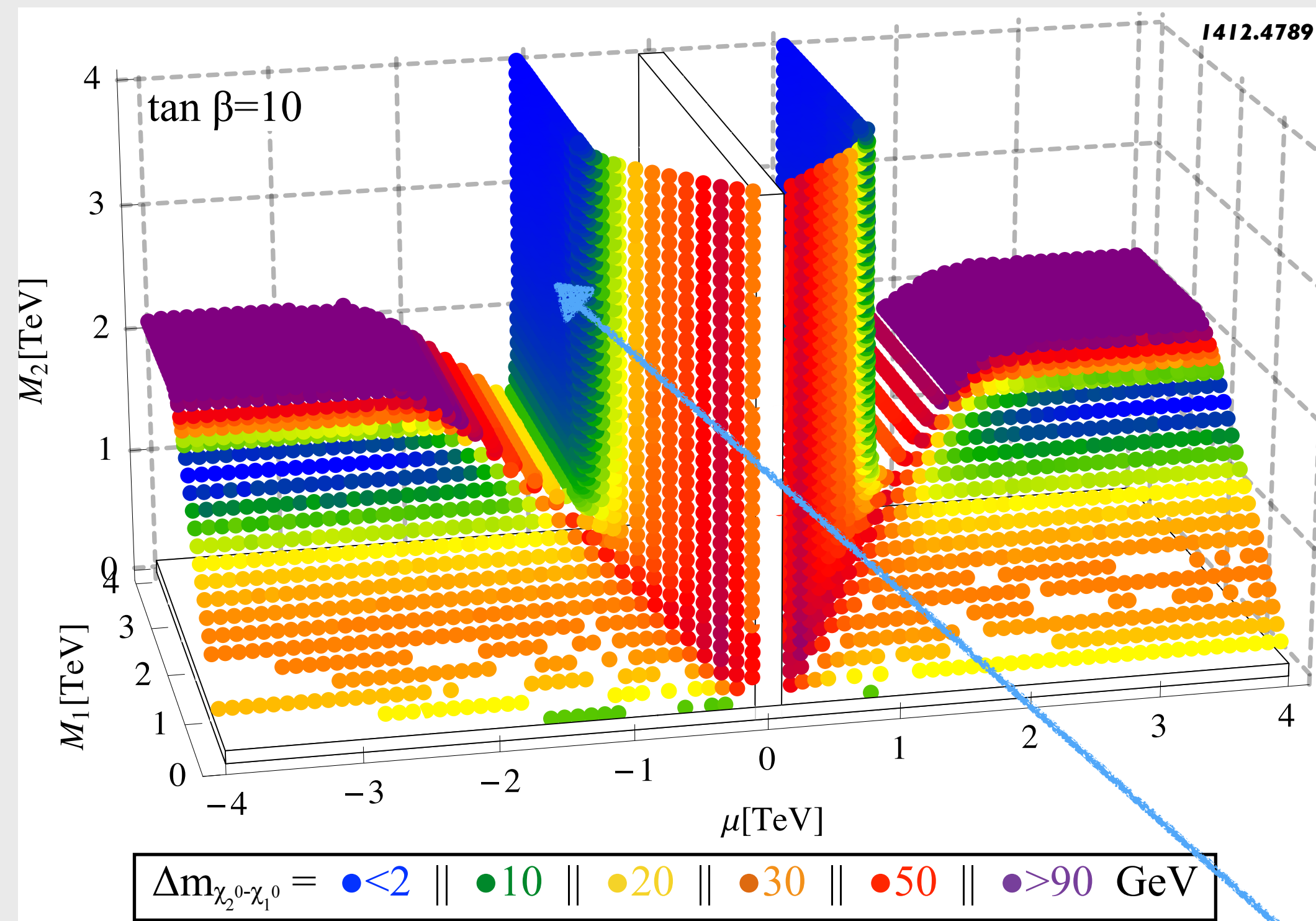
Thank you!

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WIMP: EW-ino as Dark Matter

MSSM

WITH R-PARITY

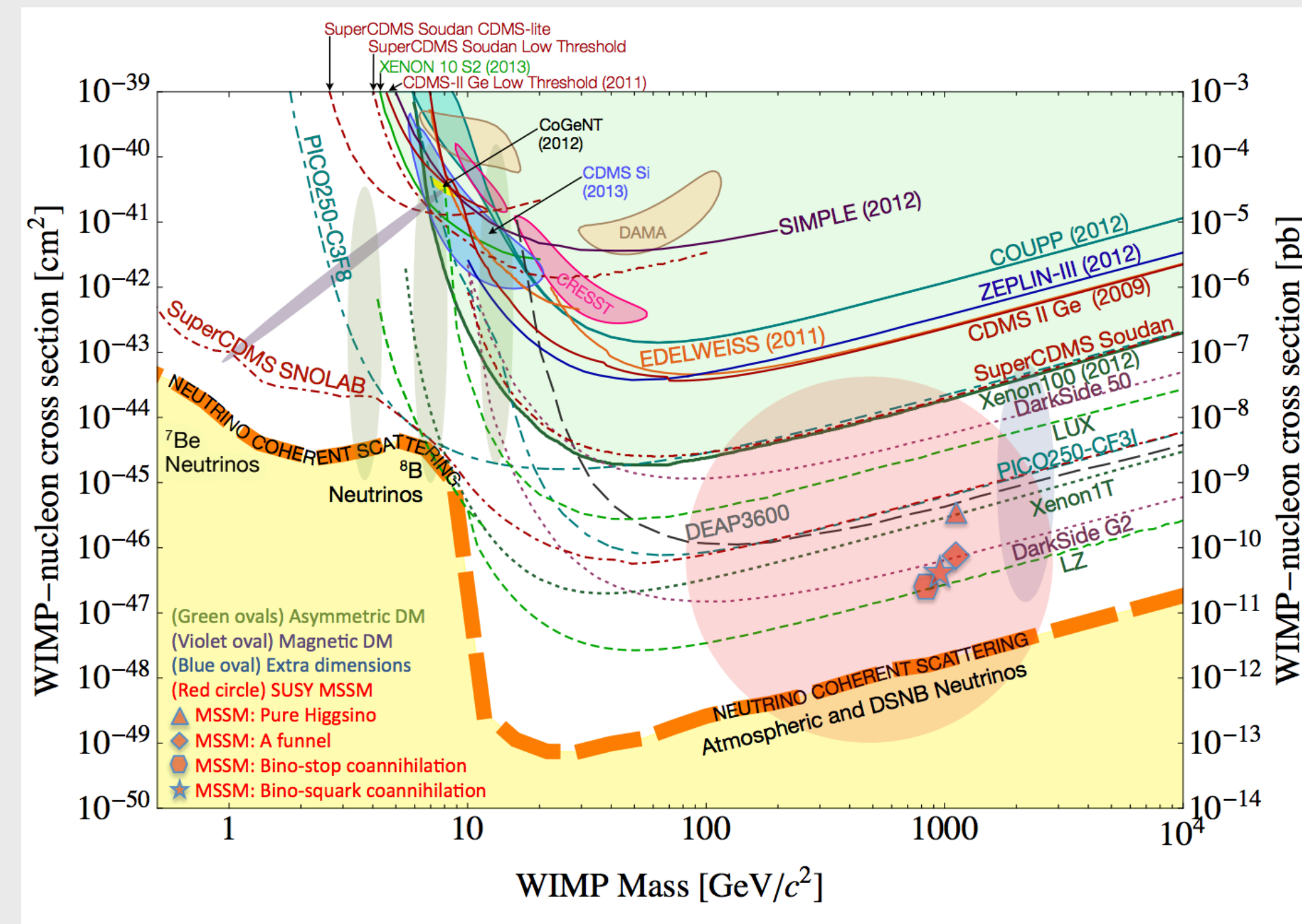
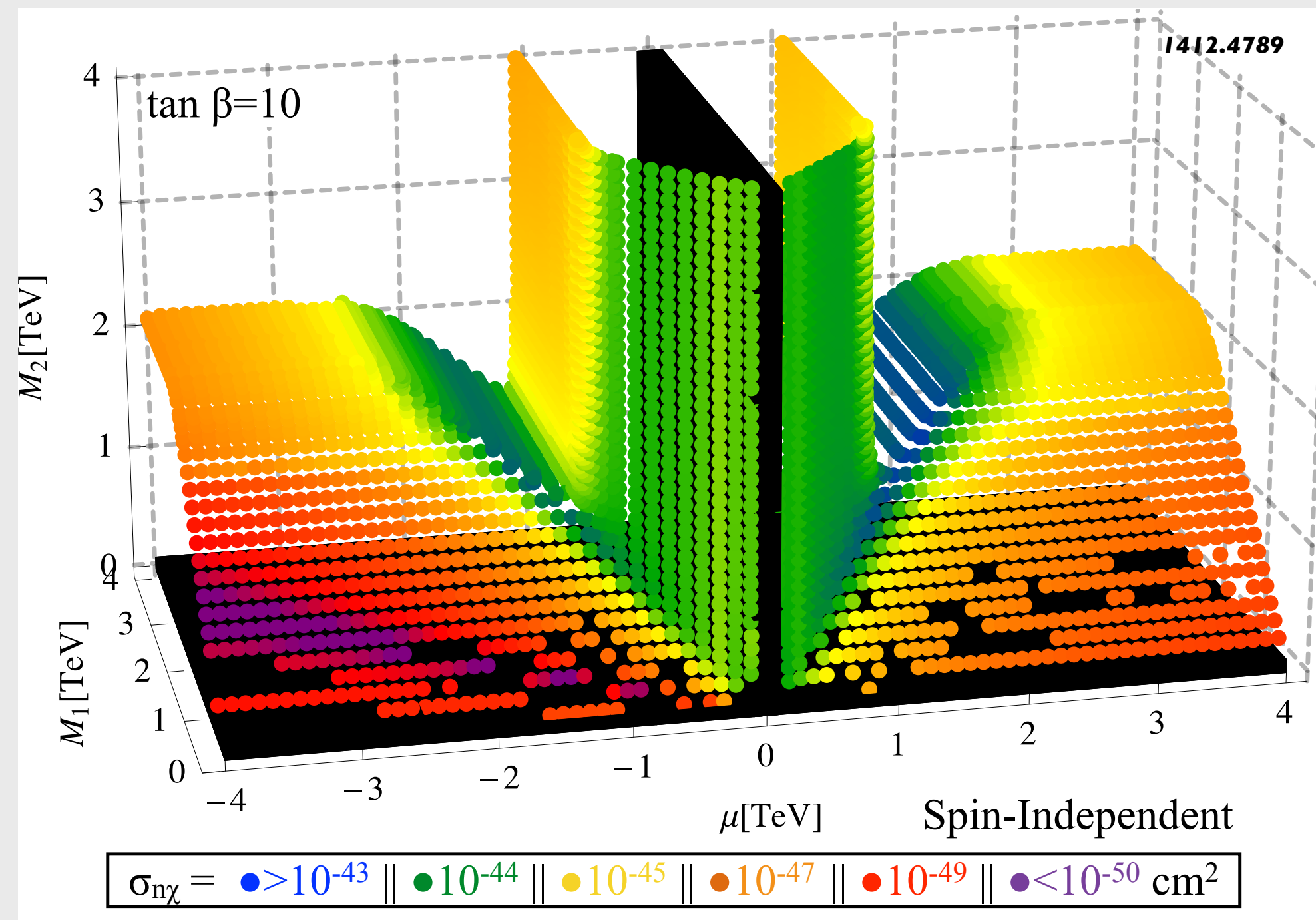


large parts of parameters space have almost degenerate multiplets

Direct Searches

WIMP

UNDER PRESSURE

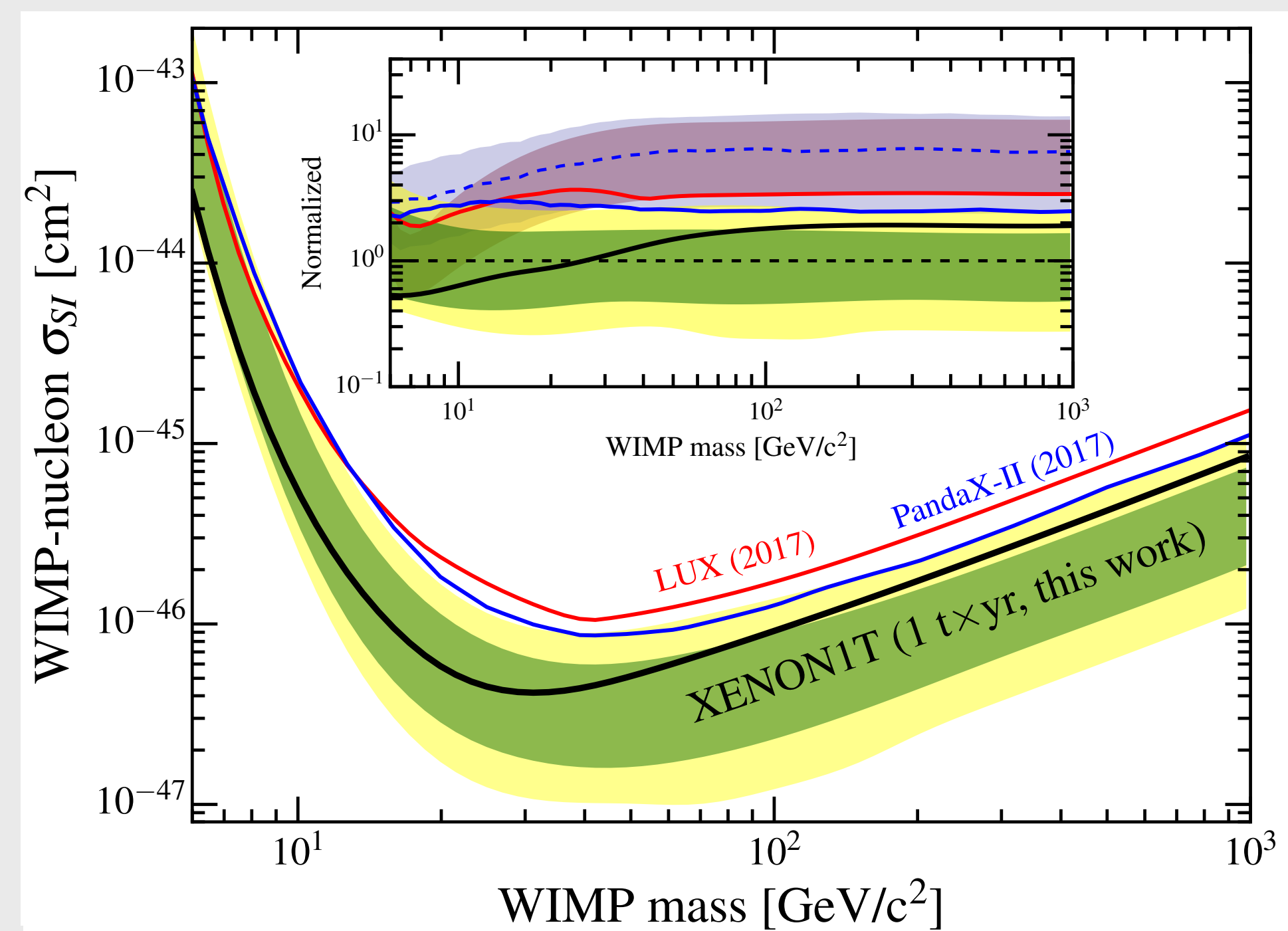
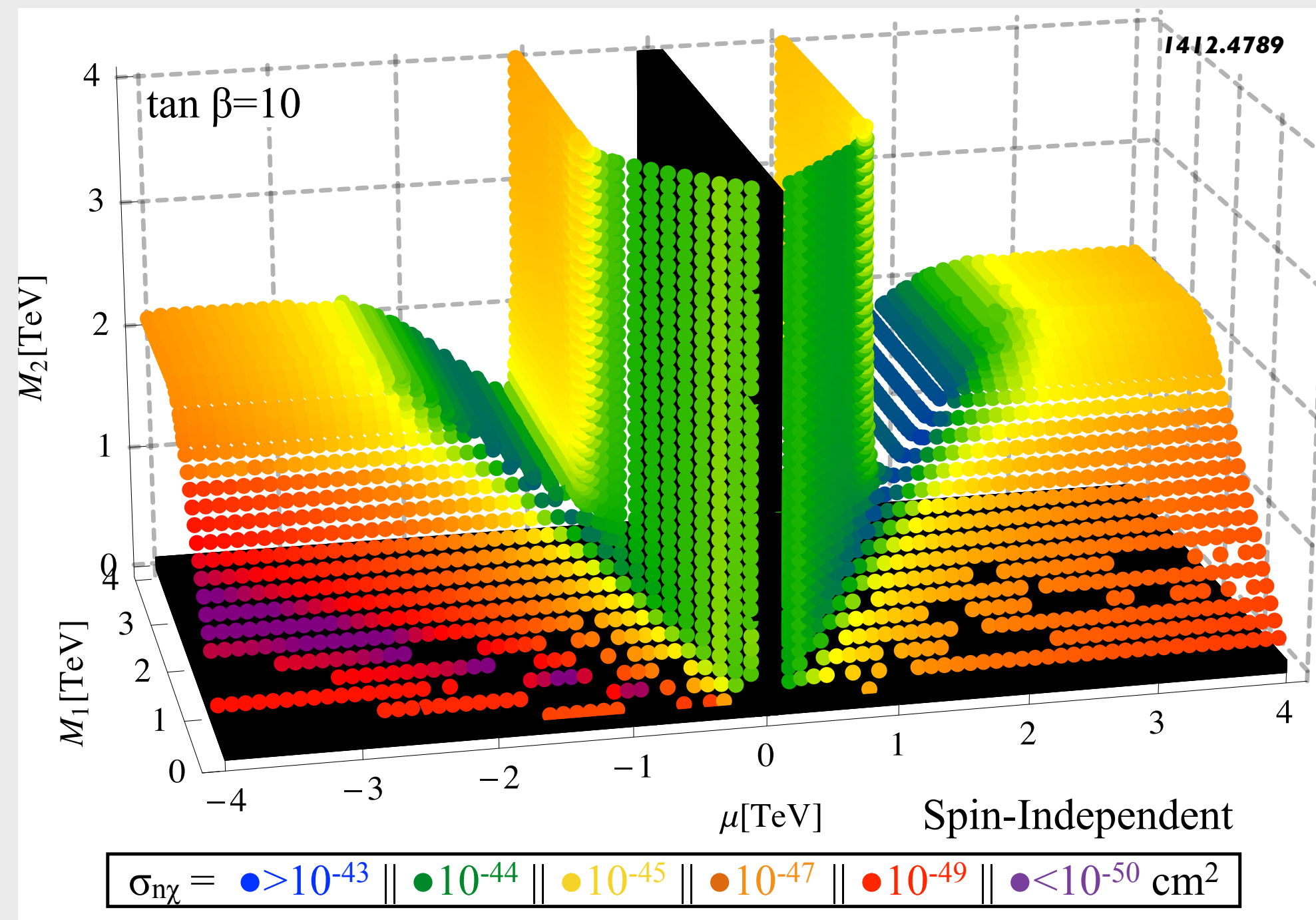


WIMP is under a fair amount of pressure, still an interesting candidate to test the capabilities of a future collider

Direct Searches

WIMP

UNDER PRESSURE



WIMP is under a fair amount of pressure, still an interesting candidate to test the capabilities of a future collider

Short (disappearing) tracks

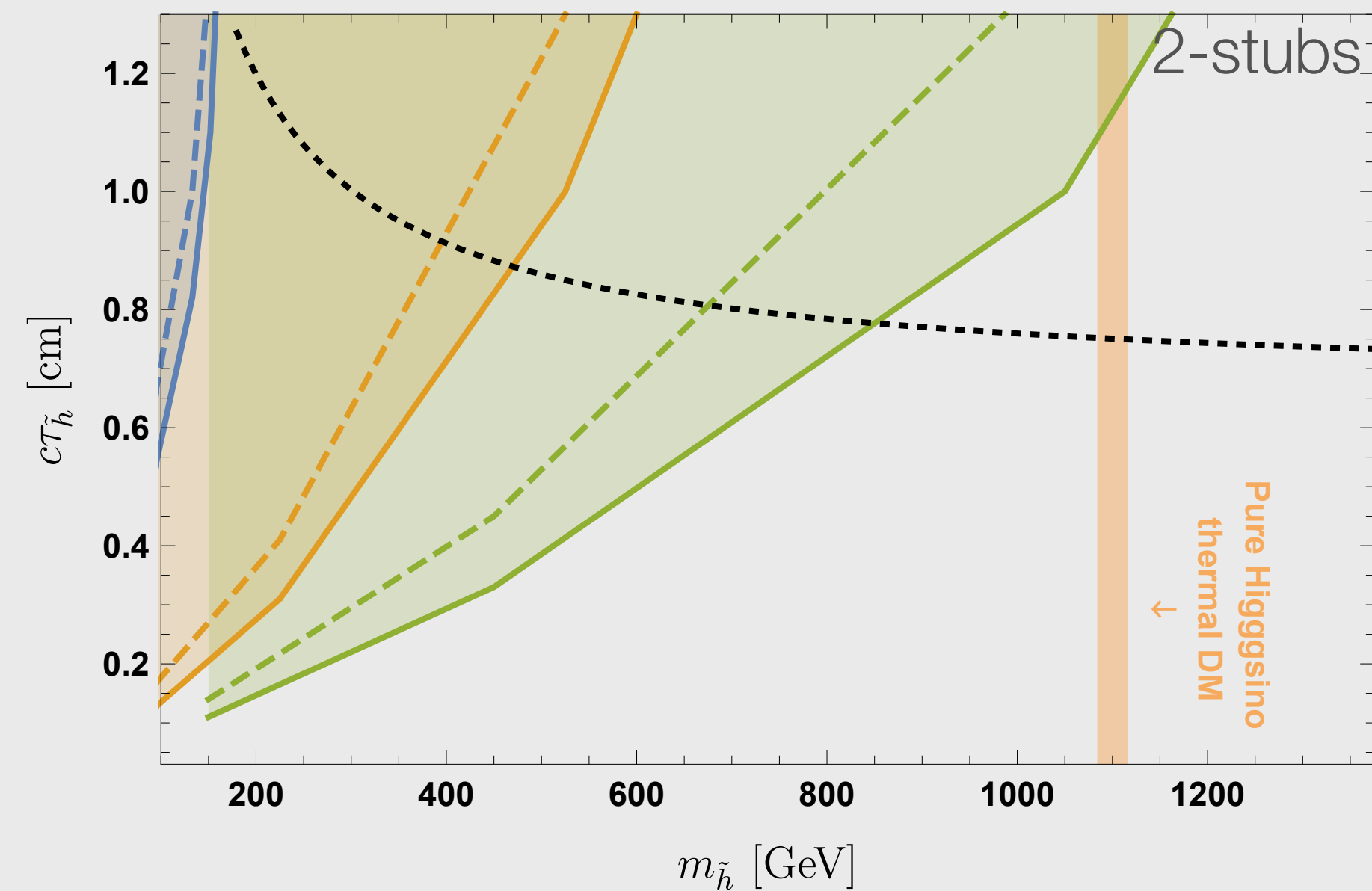
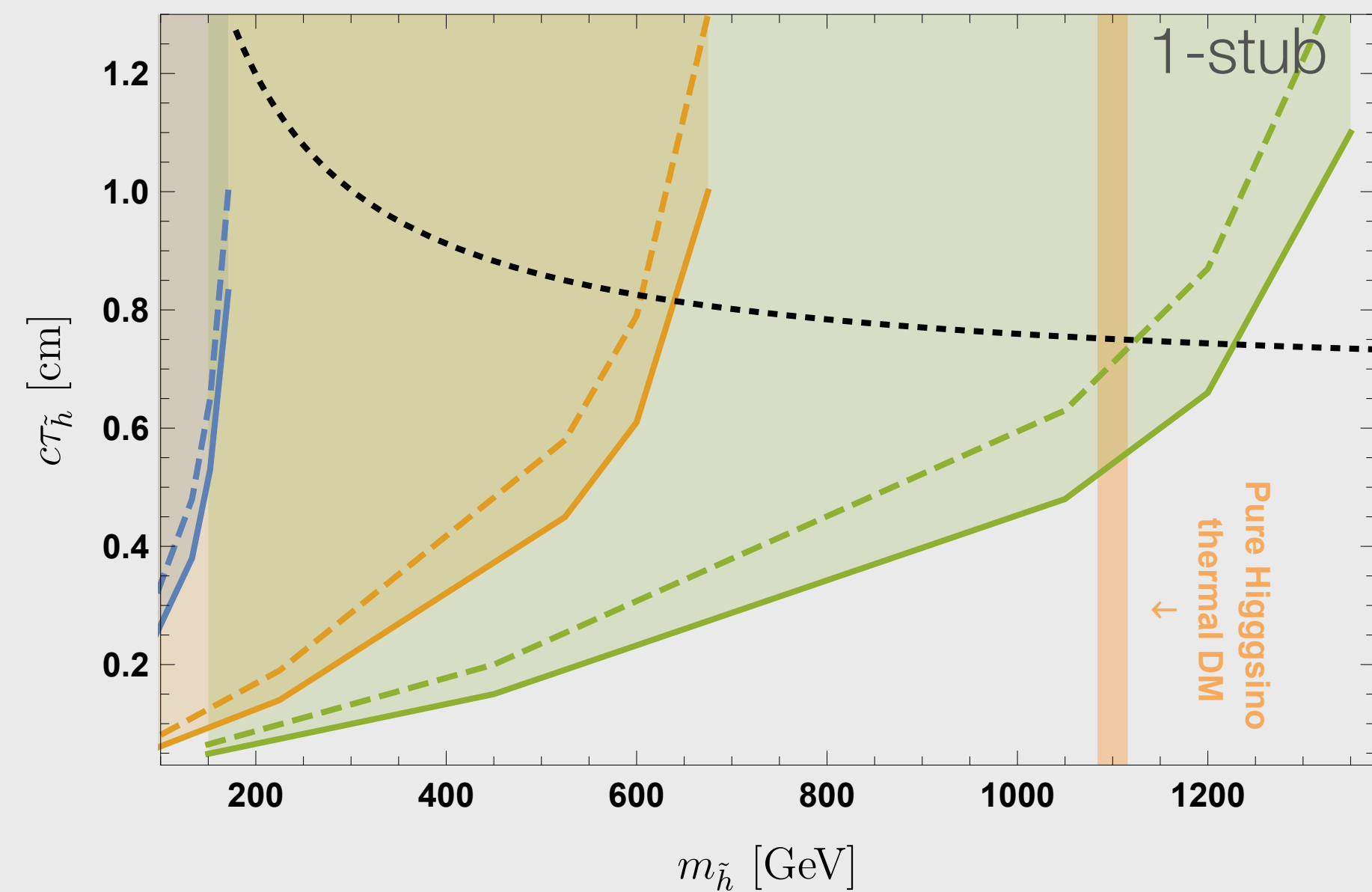
HIGGSINO DM

O(CM) DISAPPEARING TRACKS

*Charged-Neutral mass splitting
can be different if Higgsino Mixed with other states (e.g. Wino)*

TAKE LIFETIME AS FREE PARAMETER

ISOLINES FOR NUMBER OF EVENTS ASSUMED FOR DISCOVERY



- $\sqrt{s} = 380 \text{ GeV } 0.5/\text{ab}$
- $\sqrt{s} = 1.5 \text{ TeV } 1.5/\text{ab}$
- $\sqrt{s} = 3.0 \text{ TeV } 3.0/\text{ab}$
- - - Higgsino Lifetime

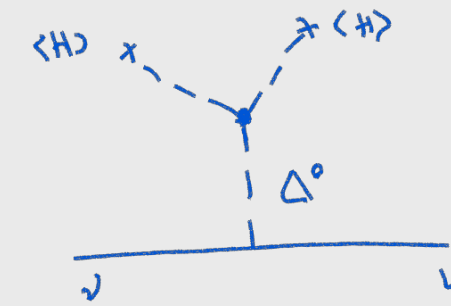
Neutrino mass mechanisms

LEPTON

NUMBER BREAKING

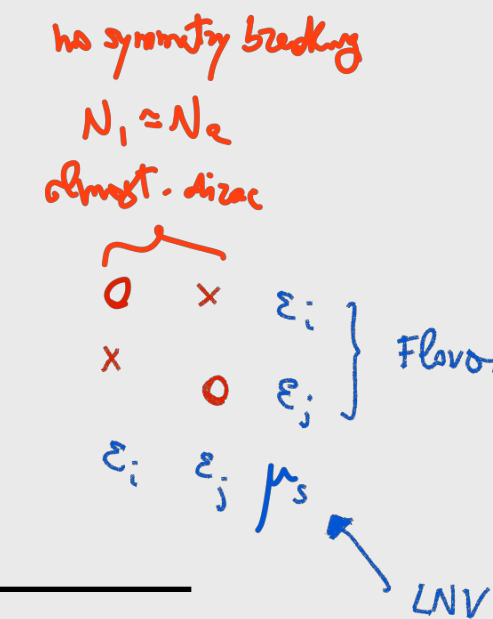
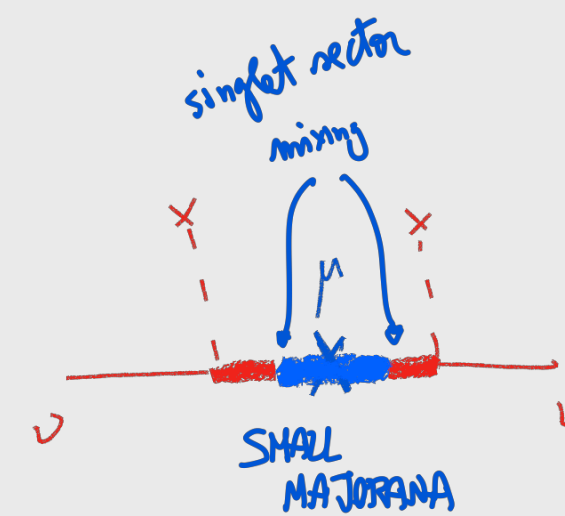
L – violation

(1,1,0) (at least 2)



(1,3,1) (1 is enough)

(1,1,0) (at least 2+1)



L – not accidental

new physics before 2012

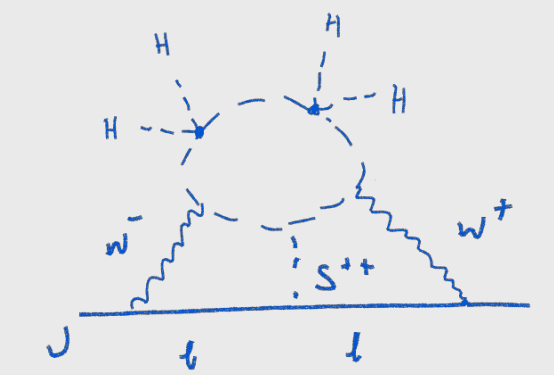
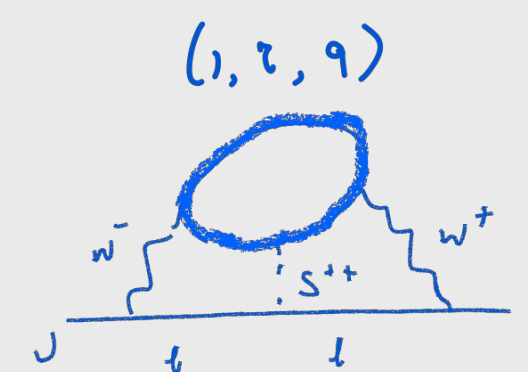
$$d = 5 \quad (1,2,1/2) \quad \frac{(LH)^2}{\Lambda}$$

UV



$$d = 7 \quad (1,1,2) \quad \frac{(DH\sigma_2 H)^2 S^{--}}{\Lambda^3}$$

UV



L – gauged, SSB

$$SU(3) \otimes SU(2)_L \otimes SU(2)_L \otimes U(1)_{B-L}$$

$$(1,2,1,1), (1,1,2,1), (1,2,2,1), (1,1,1,2),$$

Neutrinos, See-saw

Mediator of Neutrino mass mechanism

DOUBLY CHARGED

Generically $S^{++} \rightarrow \ell^+ \ell^+$ (or W^+W^+)

$$v_T = \frac{\kappa v^2}{\sqrt{2}M_T^2}$$

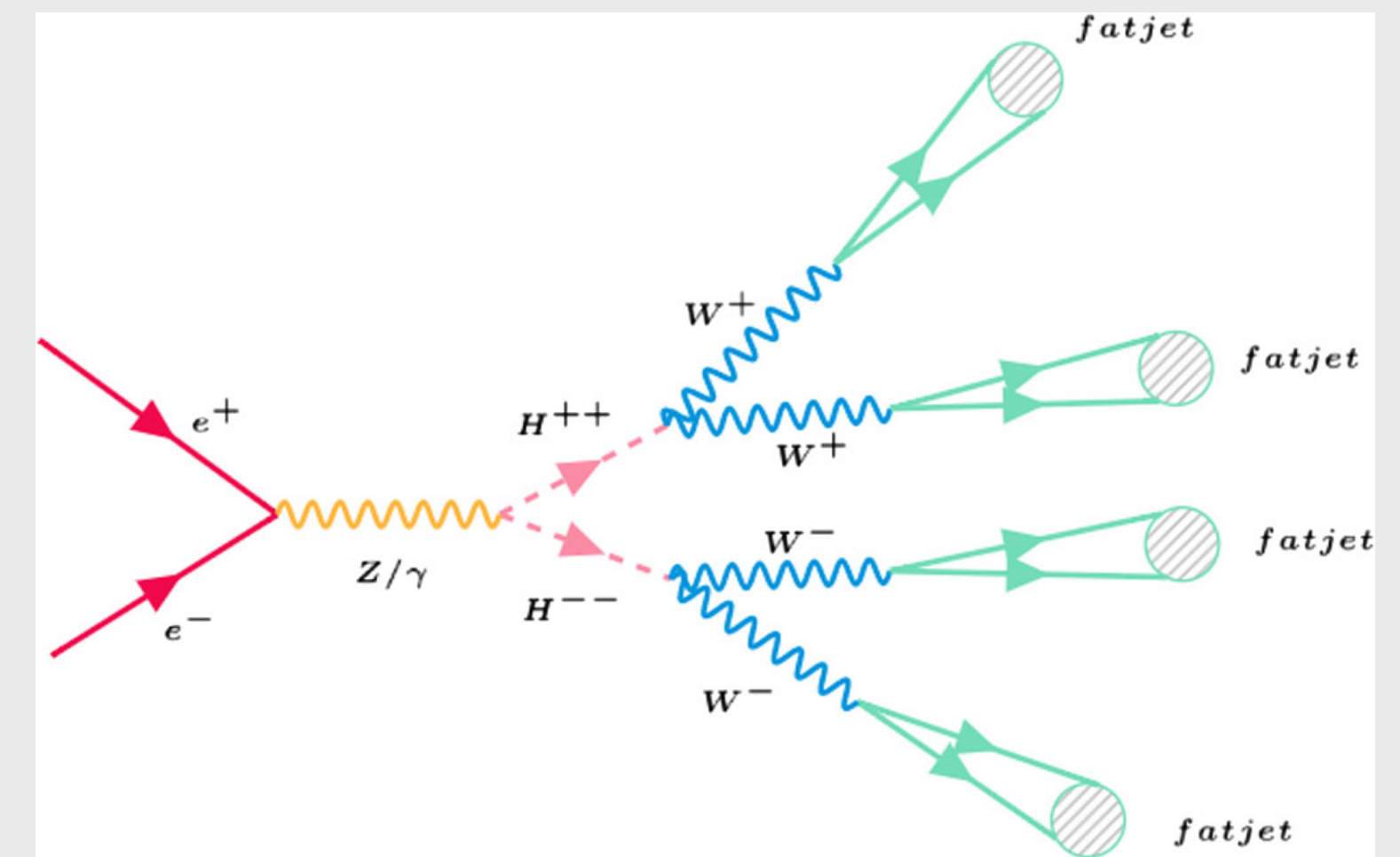
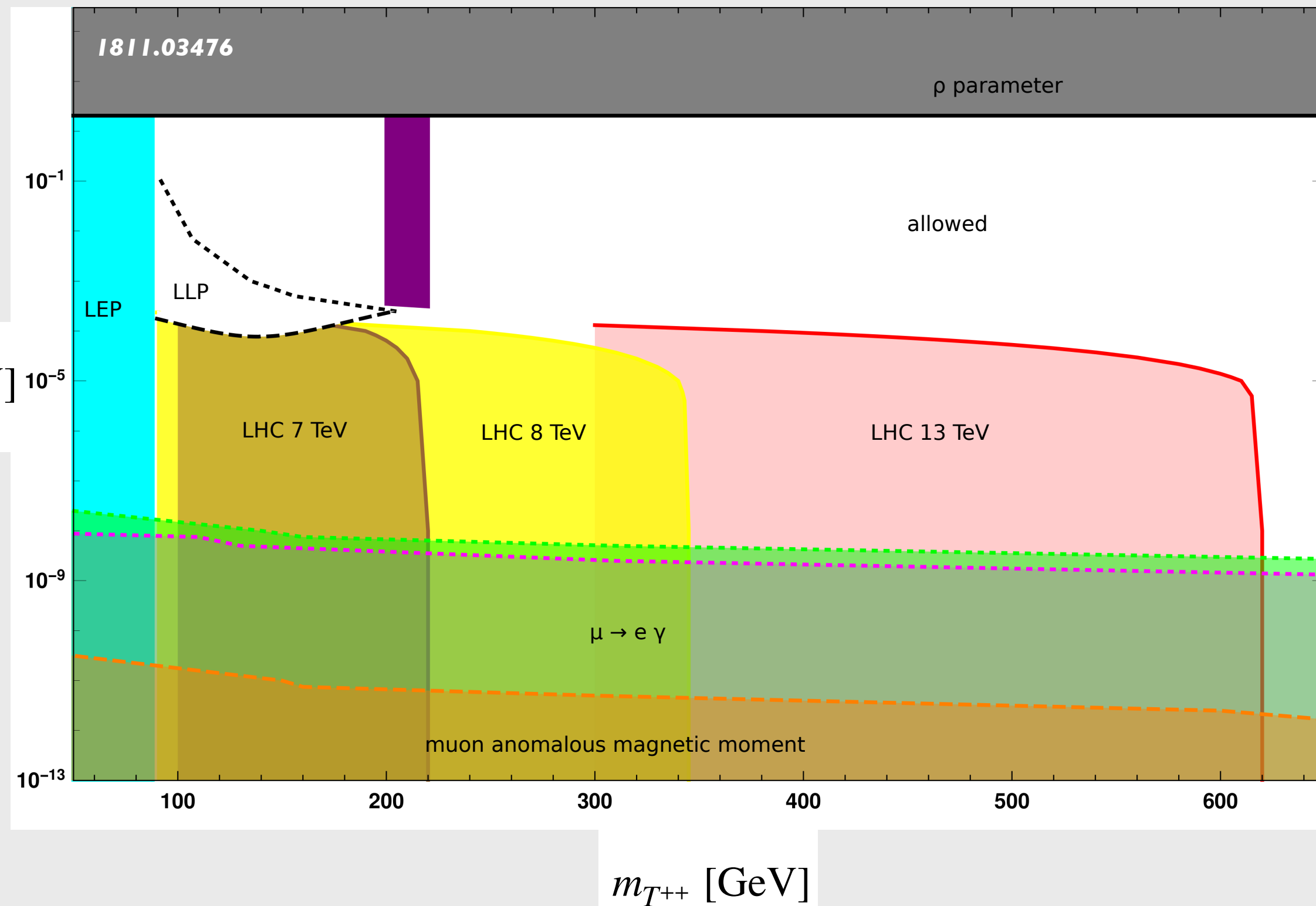
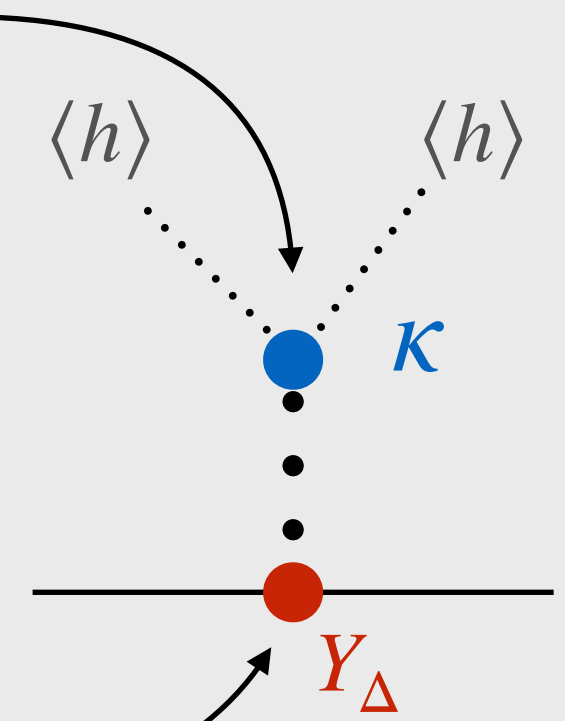
$v_T < 100 \text{ KeV} \quad H^{++} \rightarrow \ell^+ \ell^+$

$v_T > 100 \text{ KeV} \quad H^{++} \rightarrow W^+W^+$

$$m_\nu \simeq \sqrt{2}v_T Y_\Delta$$

$$V(\Phi, \Delta) = -(\kappa \Phi^\top i\sigma^2 \Delta^\dagger \Phi + h.c.)$$

$$\mathcal{L}_{Y_\Delta} = Y_\Delta \bar{\ell}^c i\sigma^2 \Delta \ell + H.c..$$



1.5 TeV

Exclude Type-2 seesaw below 1.5 TeV **for any Triplet VEV**

Mediator of Neutrino mass mechanism

DOUBLY CHARGED

Generically $S^{++} \rightarrow \ell^+ \ell^+$ (or W^+W^+)

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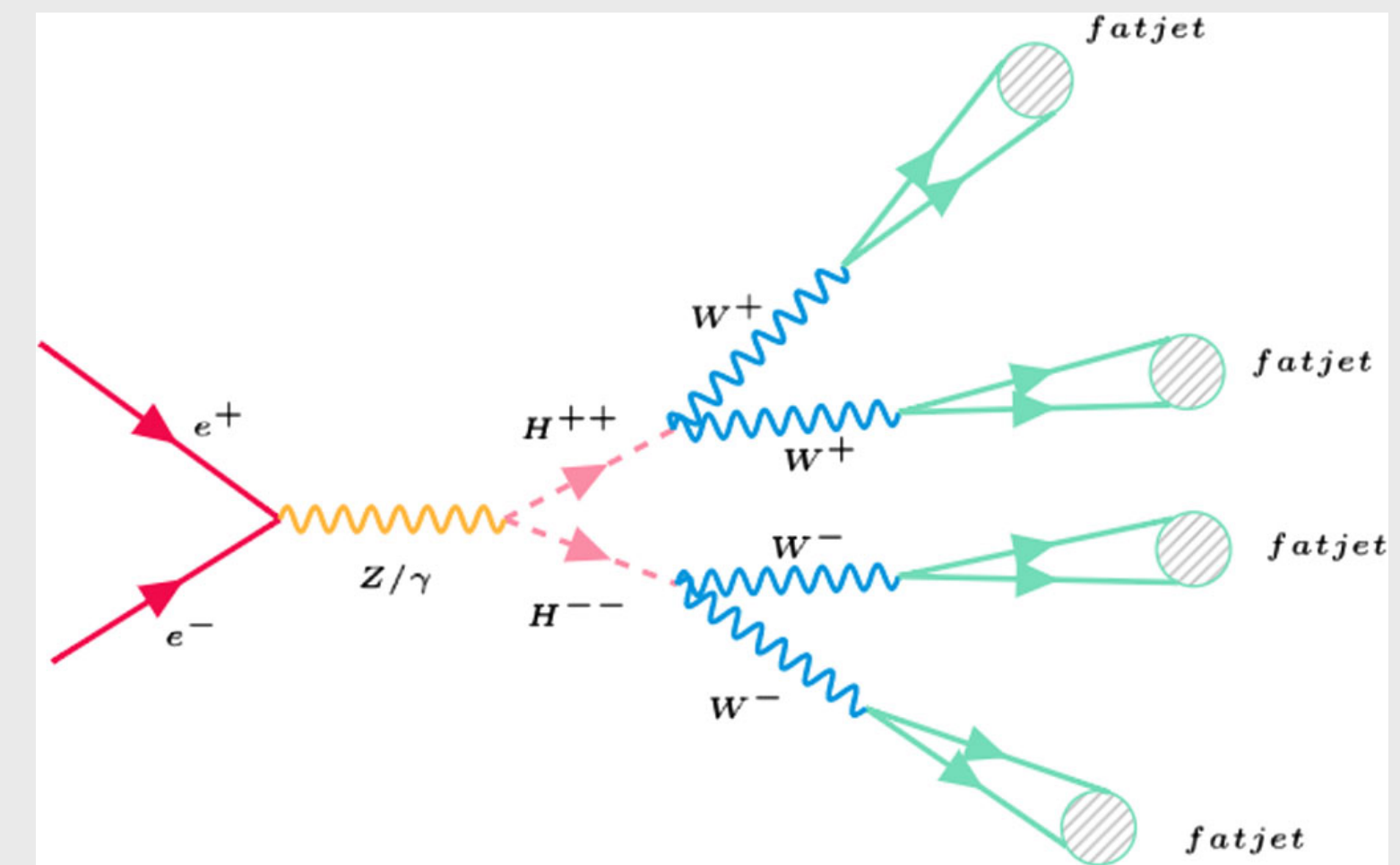
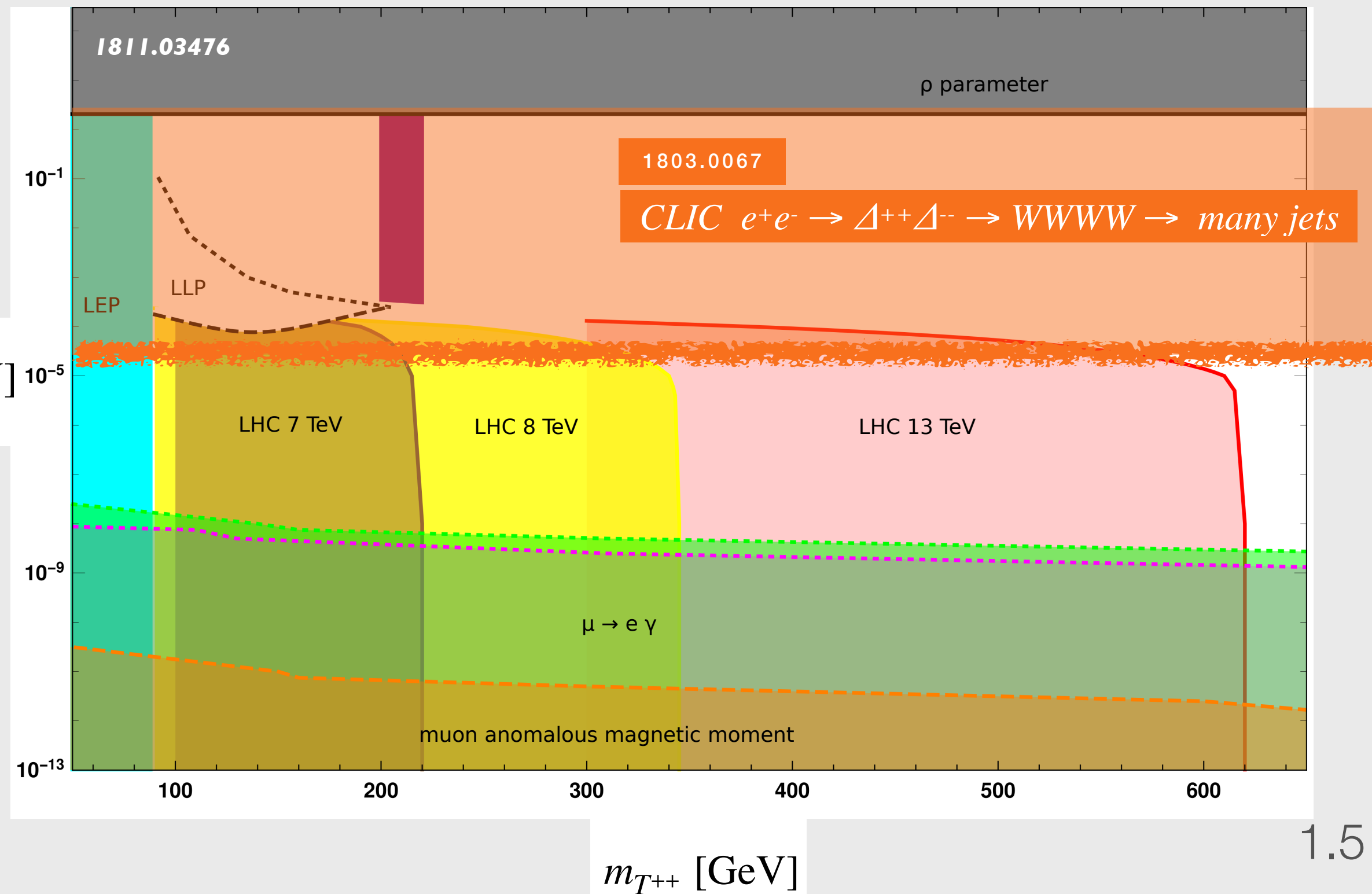
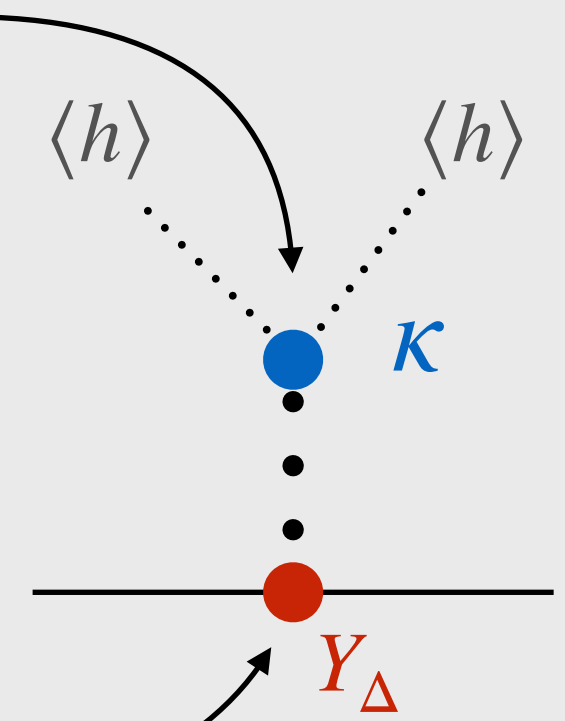
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Exclude Type-2 seesaw below 1.5 TeV **for any Triplet VEV**

Mediator of Neutrino mass mechanism

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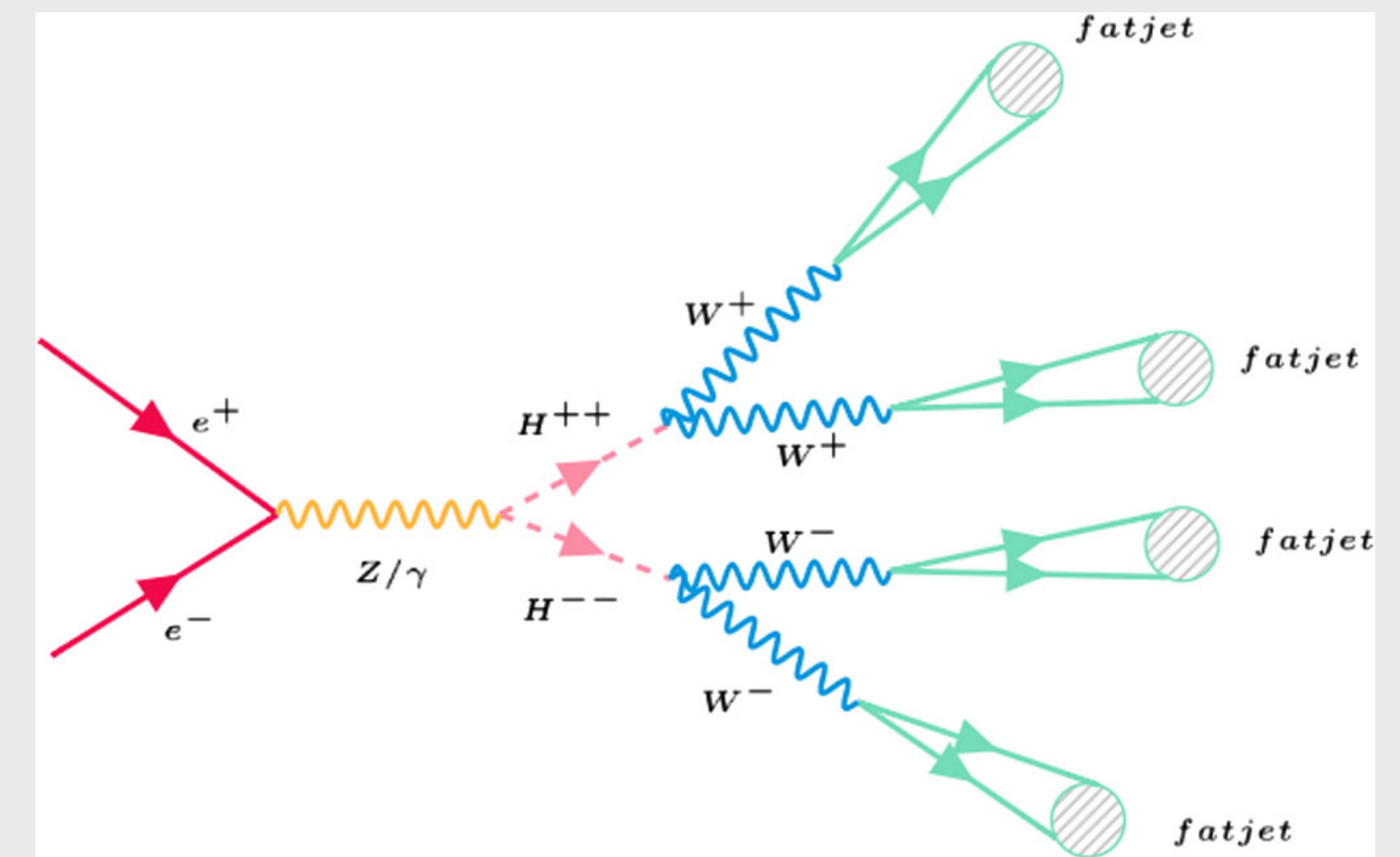
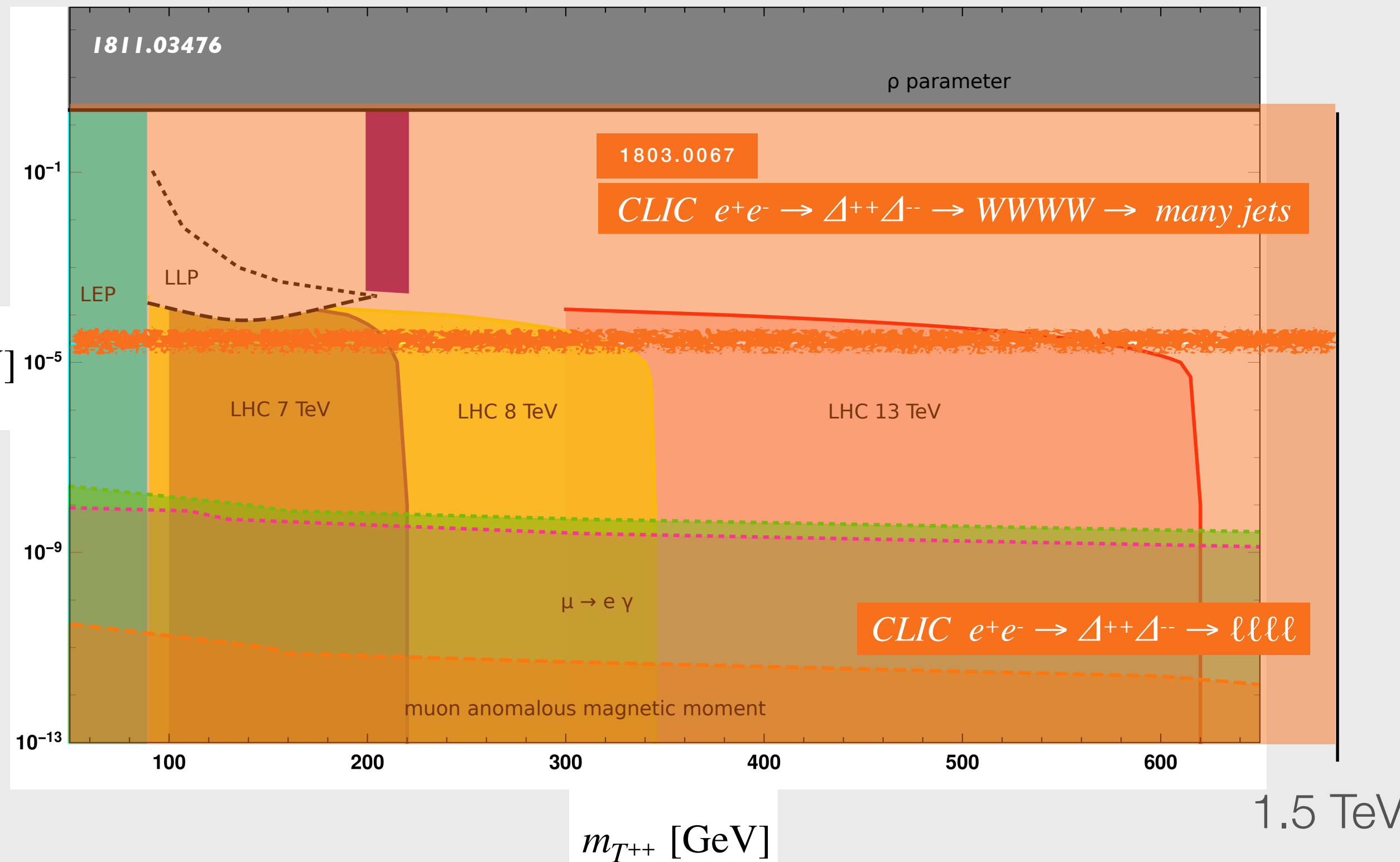
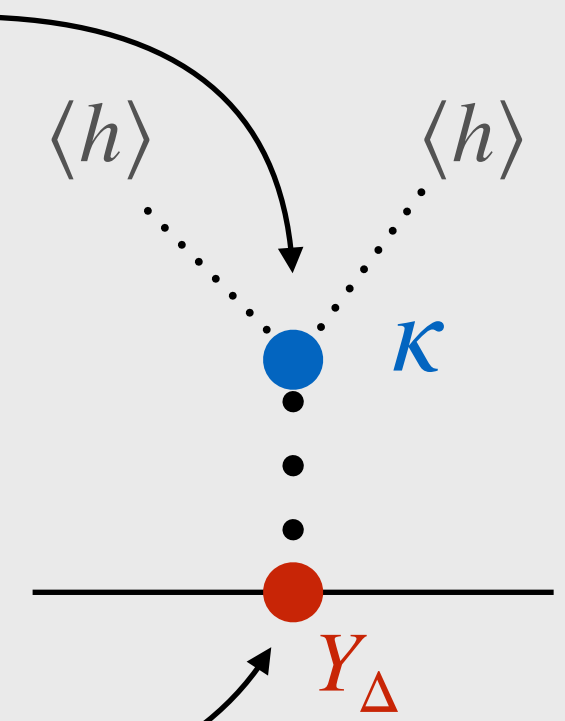
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$v_T > 100 \text{ KeV} \quad H^{++} \rightarrow W^+W^+$

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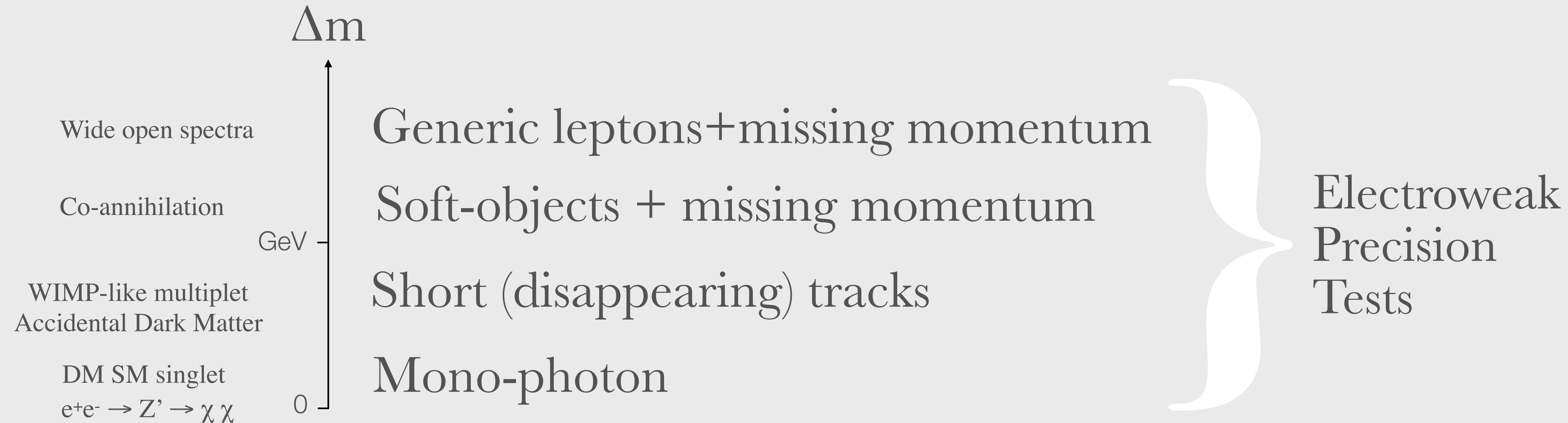
$$V(\Phi, \Delta) = -(\kappa \Phi^\top i\sigma^2 \Delta^\dagger \Phi + h.c.)$$

$$\mathcal{L}_{Y_\Delta} = Y_\Delta \bar{\ell}^c i\sigma^2 \Delta \ell + H.c..$$

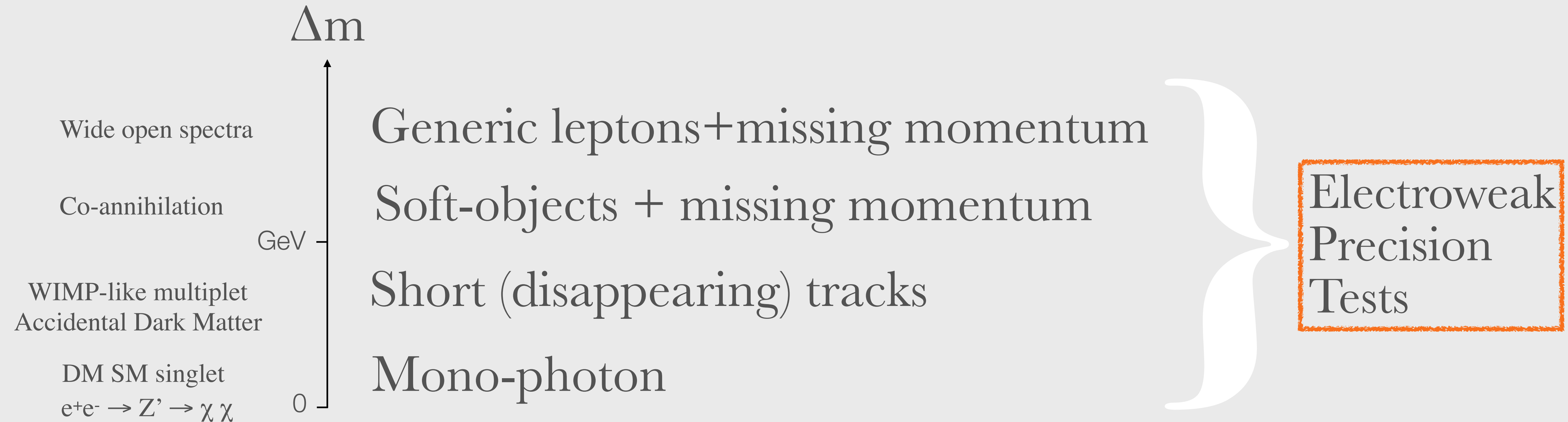


Exclude Type-2 seesaw below 1.5 TeV **for any Triplet VEV**

Electroweak Dark Matter: LSP (+NLSP)



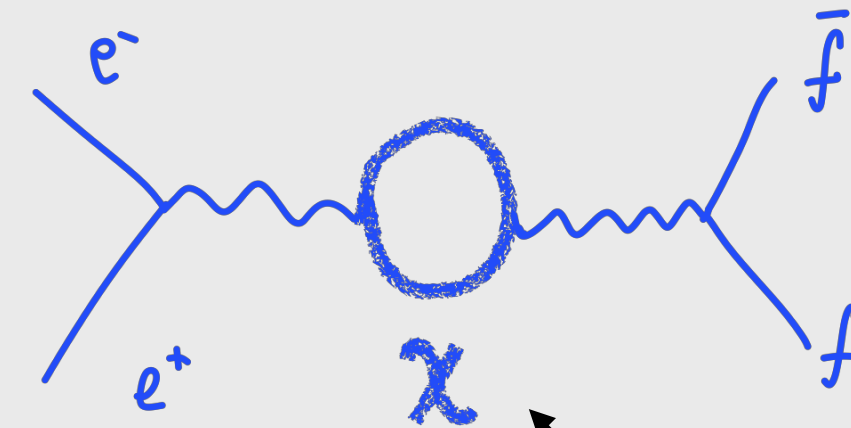
Electroweak Dark Matter: LSP (+NLSP)



$$e^+ e^- \rightarrow f \bar{f}$$

PRECISION

ANGULAR DISTRIBUTION



χ is heavy/light new physics

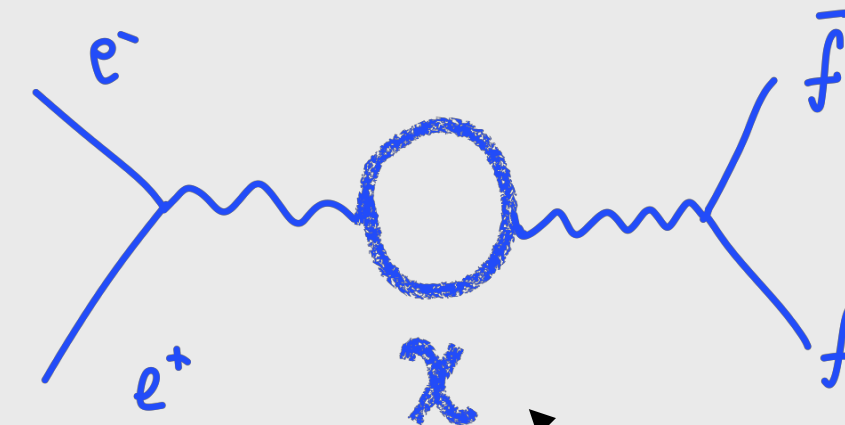
$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{g^2 C_{WW}^{\text{eff}}}{8} W_{\mu\nu}^a \Pi(-D^2/m_\chi^2) W^{a\mu\nu} + \frac{g'^2 C_{BB}^{\text{eff}}}{8} B_{\mu\nu} \Pi(-\partial^2/m_\chi^2) B^{\mu\nu}$$

1504.03402

$$e^+ e^- \rightarrow f \bar{f}$$

PRECISION

ANGULAR DISTRIBUTION



χ is heavy/light new physics

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{g^2 C_{WW}^{\text{eff}}}{8} W_{\mu\nu}^a \Pi(-D^2/m_\chi^2) W^{a\mu\nu} + \frac{g'^2 C_{BB}^{\text{eff}}}{8} B_{\mu\nu} \Pi(-\partial^2/m_\chi^2) B^{\mu\nu}$$

1504.03402

HEAVY NEW PHYSICS (EFT LIMIT)

$$\Pi\left(\frac{s}{m^2}\right) \sim \frac{1}{480\pi^2} \cdot \frac{s}{m^2}$$

$$C_{WW}^{\text{eff}} = \kappa(n^3 - n)/6, \quad C_{BB}^{\text{eff}} = \kappa 2nY^2$$

$$\kappa = \frac{1}{2}, 1, 4, 8 \text{ for RS, CS, MF, DF}$$

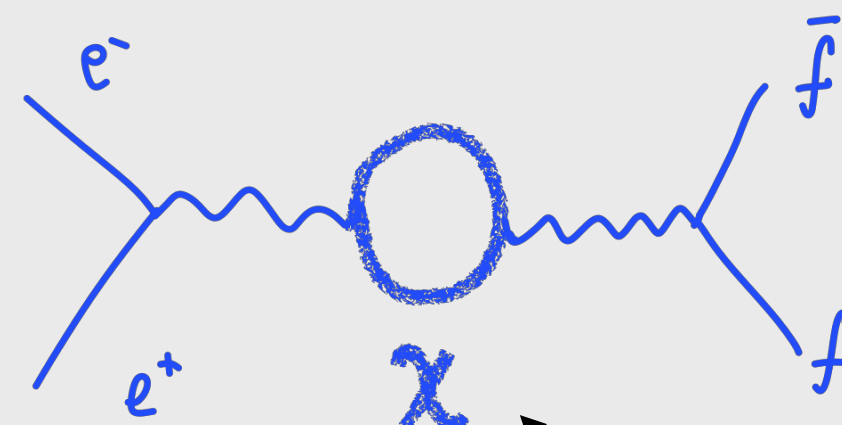
$$W = \frac{g^2 C_{WW}^{\text{eff}}}{960\pi^2} \frac{m_W^2}{m_\chi^2} \quad Y = \frac{g'^2 C_{BB}^{\text{eff}}}{960\pi^2} \frac{m_W^2}{m_\chi^2}$$

EFFECTS GROW WITH ENERGY

$e^+e^- \rightarrow f\bar{f}$

PRECISION

ANGULAR DISTRIBUTION



χ is heavy/light new physics

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{g^2 C_{WW}^{\text{eff}}}{8} W_{\mu\nu}^a \Pi(-D^2/m_\chi^2) W^{a\mu\nu} + \frac{g'^2 C_{BB}^{\text{eff}}}{8} B_{\mu\nu} \Pi(-\partial^2/m_\chi^2) B^{\mu\nu}$$

1504.03402

LIGHT NEW PHYSICS

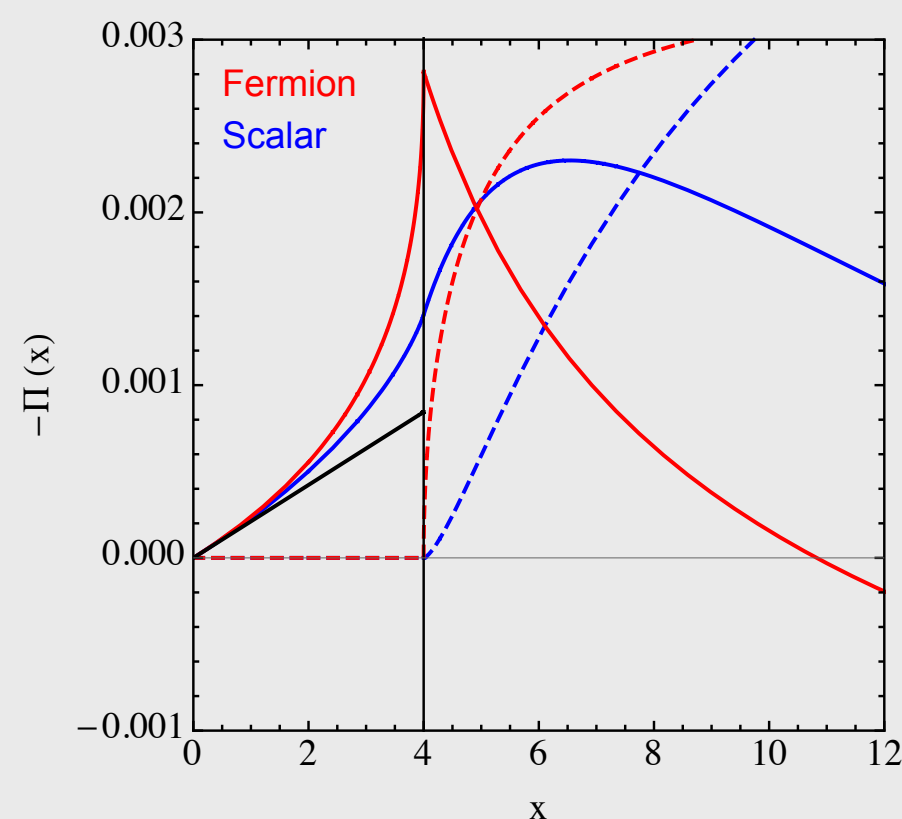
HEAVY NEW PHYSICS (EFT LIMIT)

$$\Pi(x) = \begin{cases} -\frac{8(x-3)+3x\left(\frac{x-4}{x}\right)^{3/2} \log\left(\frac{1}{2}\left(\left(\sqrt{\frac{x-4}{x}}-1\right)x+2\right)\right)}{144\pi^2 x} & \text{(scalars)} \\ -\frac{12+5x+3\sqrt{\frac{x-4}{x}}(x+2) \log\left(\frac{1}{2}\left(\left(\sqrt{\frac{x-4}{x}}-1\right)x+2\right)\right)}{288\pi^2 x} & \text{(fermions)} \end{cases}$$

$$\Pi\left(\frac{s}{m^2}\right) \sim \frac{1}{480\pi^2} \cdot \frac{s}{m^2}$$

$$C_{WW}^{\text{eff}} = \kappa(n^3 - n)/6, \quad C_{BB}^{\text{eff}} = \kappa 2nY^2$$

$$\kappa = \frac{1}{2}, 1, 4, 8 \text{ for RS, CS, MF, DF}$$



$$W = \frac{g^2 C_{WW}^{\text{eff}}}{960\pi^2} \frac{m_W^2}{m_\chi^2} \quad Y = \frac{g'^2 C_{BB}^{\text{eff}}}{960\pi^2} \frac{m_W^2}{m_\chi^2}$$

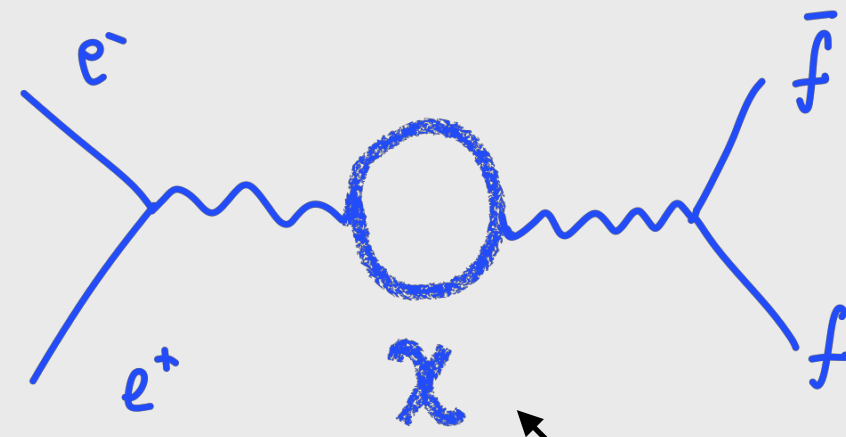
PRECISION PHYSICS

EFFECTS GROW WITH ENERGY

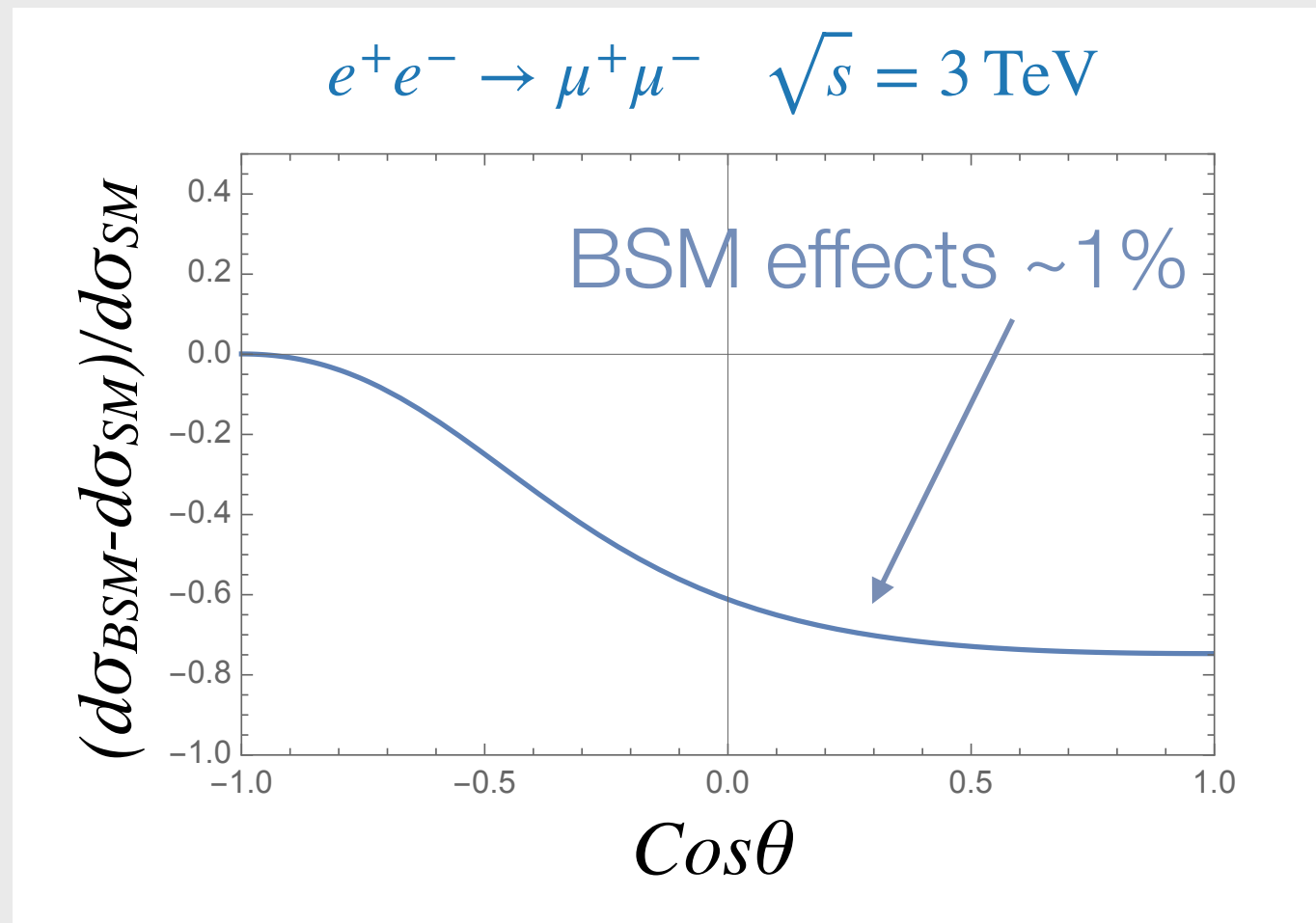
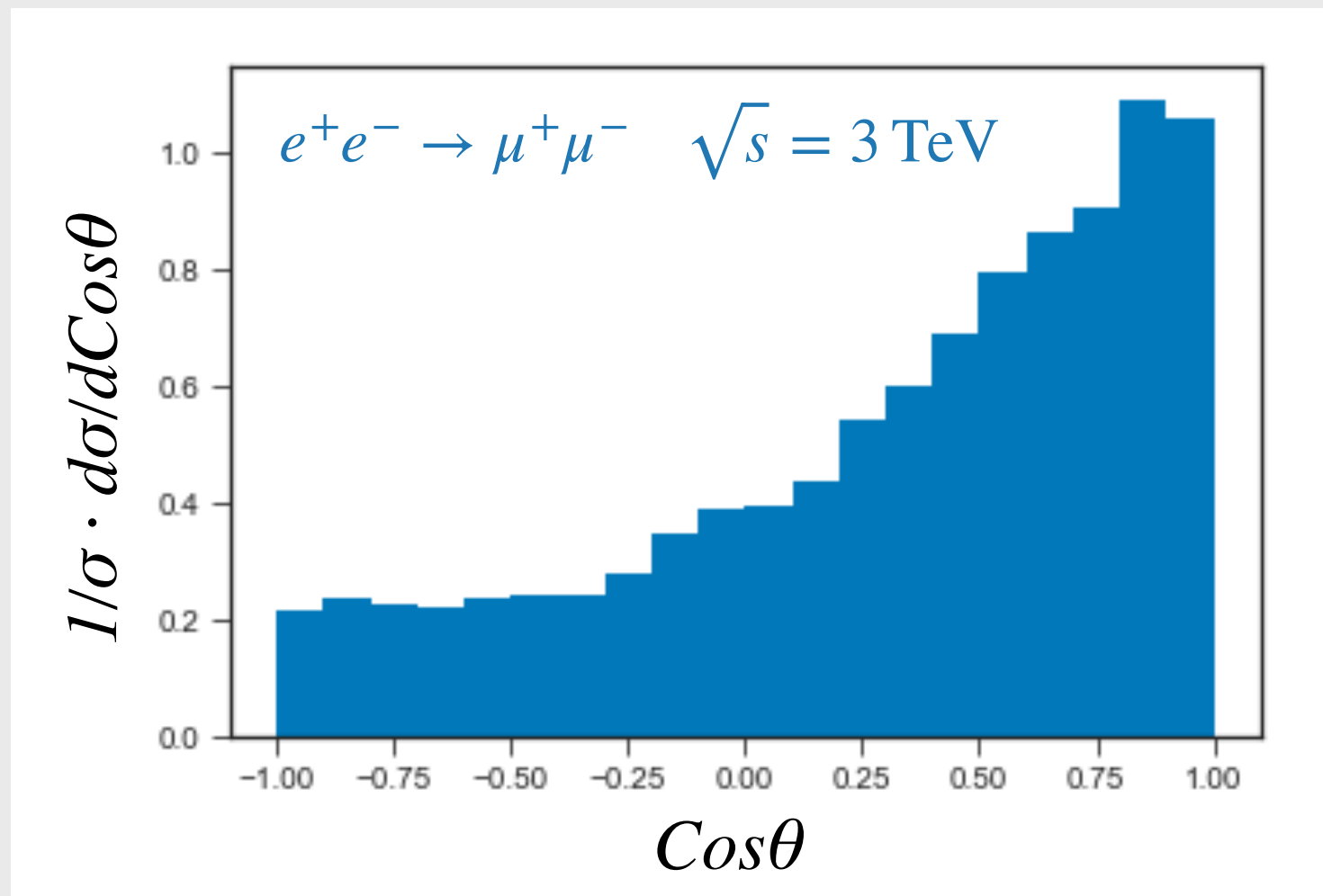
$$e^+ e^- \rightarrow f \bar{f}$$

PRECISION

ANGULAR DISTRIBUTION



χ is heavy/light new physics

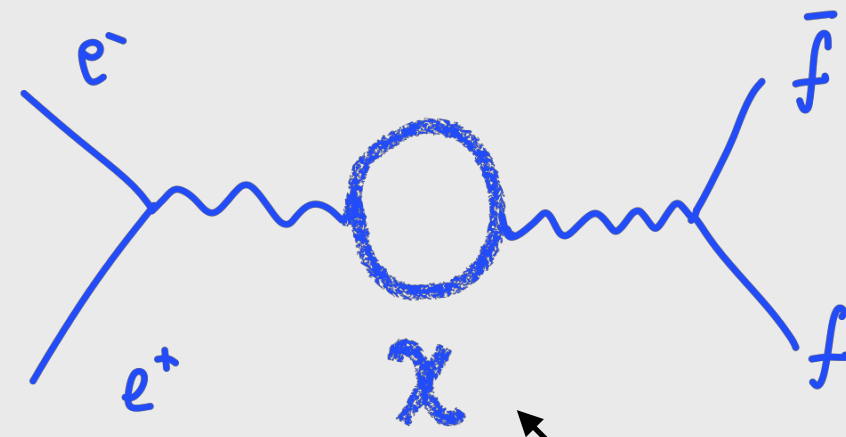


beams polarization is beneficial to increase NP effects

$$e^+ e^- \rightarrow f \bar{f}$$

PRECISION

ANGULAR DISTRIBUTION



χ is heavy/light new physics

$|\cos\theta| < 0.95$ for all final states

χ^2 over 10 bins

$$\chi^2 = \sum_{i=1}^{10} \frac{(N_i^{\text{SM+BSM}} - N_i^{\text{SM}})^2}{N_i^{\text{SM}} + (\epsilon_i N_i^{\text{SM}})^2}$$

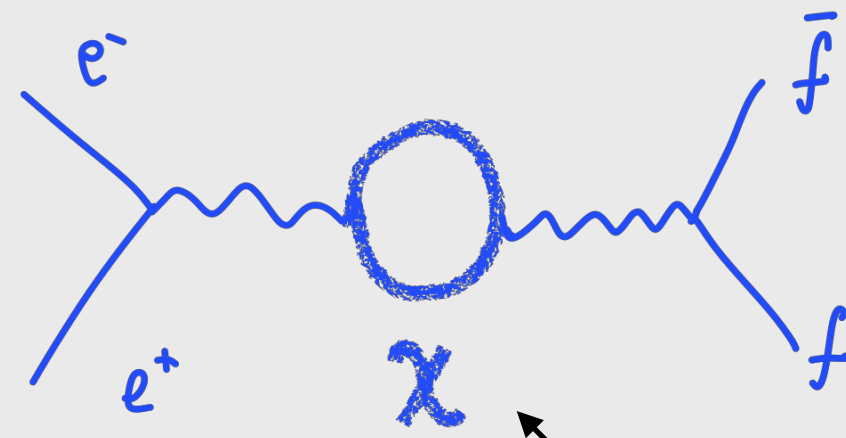
Systematic Unc.

beams polarization is beneficial to increase NP effects

$$e^+ e^- \rightarrow f \bar{f}$$

PRECISION

ANGULAR DISTRIBUTION



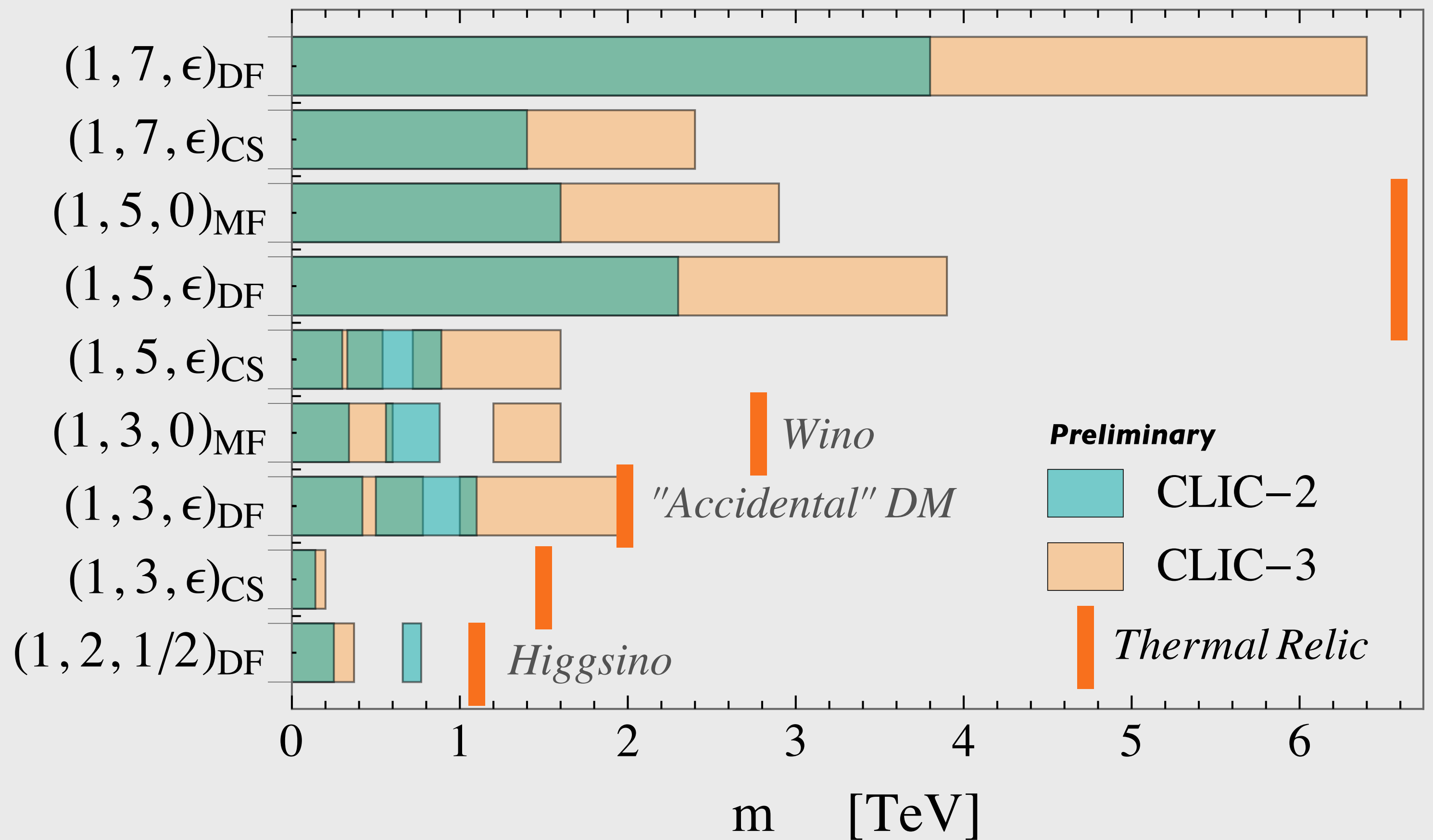
χ is heavy/light new physics

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Systematic Unc.



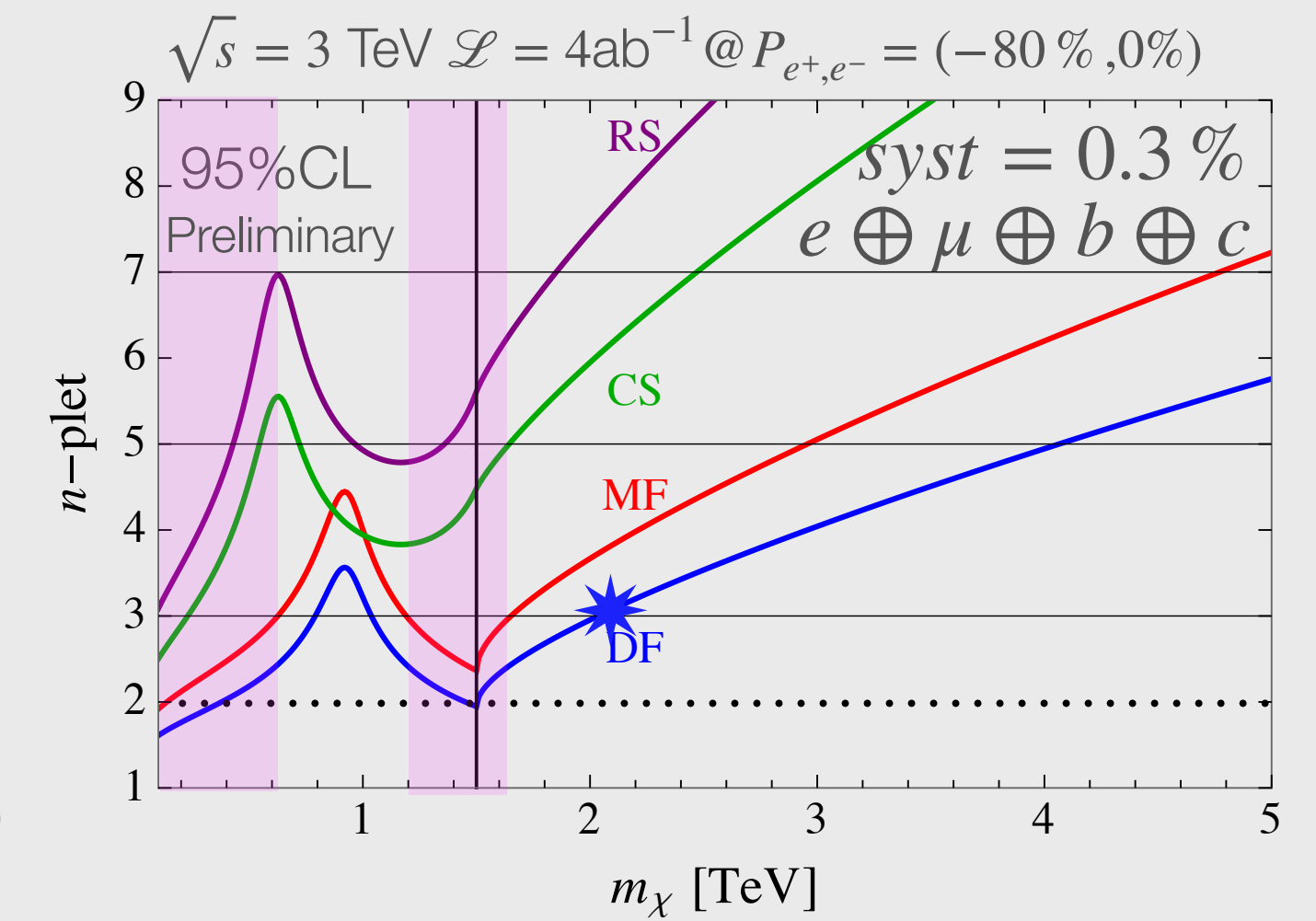
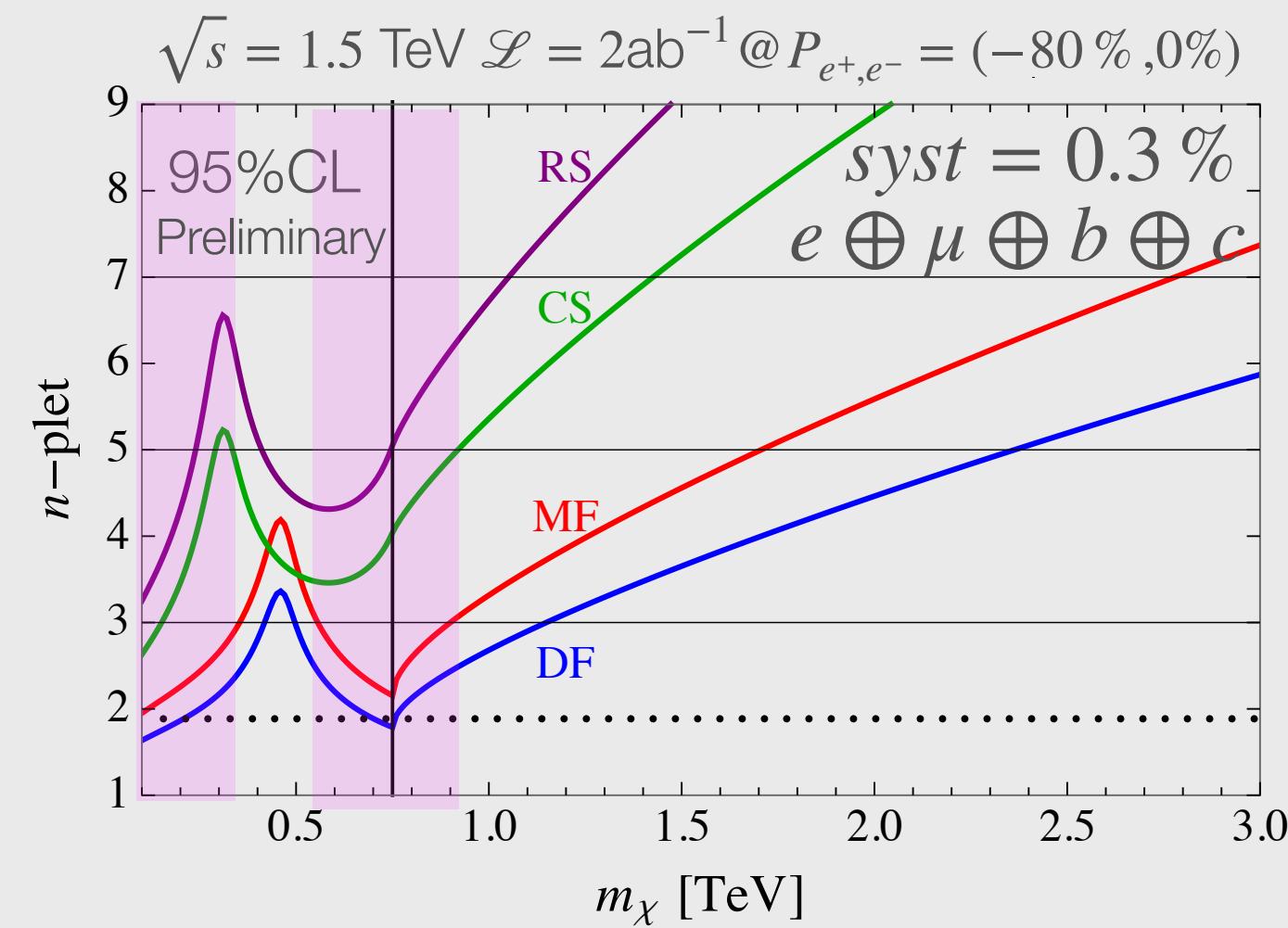
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PRECISION

ANGULAR DISTRIBUTION

χ	$m_\chi^{(DM)}$ [TeV]	$m_\chi^{(CLIC-3)}$ [TeV]
$(1, 2, 1/2)_{DF}$	1.1	1.5
$(1, 3, \epsilon)_{CS}$	1.55	-
$(1, 3, \epsilon)_{DF}$	2.0	2.1 *
$(1, 3, 0)_{MF}$	2.8	1.7
$(1, 5, \epsilon)_{CS}$	6.6	1.7
$(1, 5, \epsilon)_{DF}$	6.6	4.1
$(1, 5, 0)_{MF}$	11	3.0
$(1, 7, \epsilon)_{CS}$	16	2.5
$(1, 7, \epsilon)_{DF}$	16	6.8



Higgsino of split-SUSY (heavy sfermions)

Wino of split-SUSY (heavy sfermions)

Accidental Dark Matter 3-plet Dirac Fermion

$$e^+ e^- \rightarrow f \bar{f}$$

PRECISION

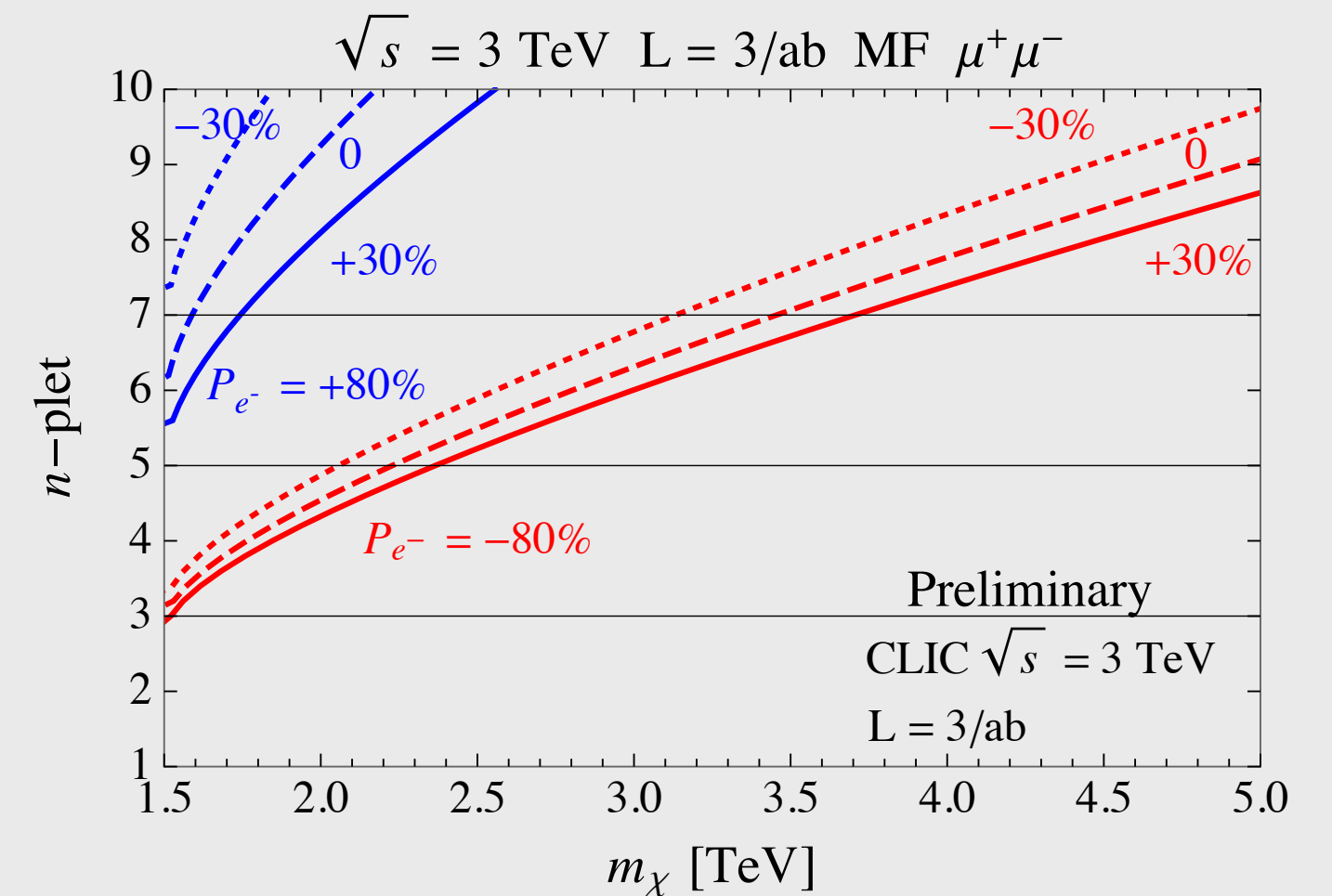
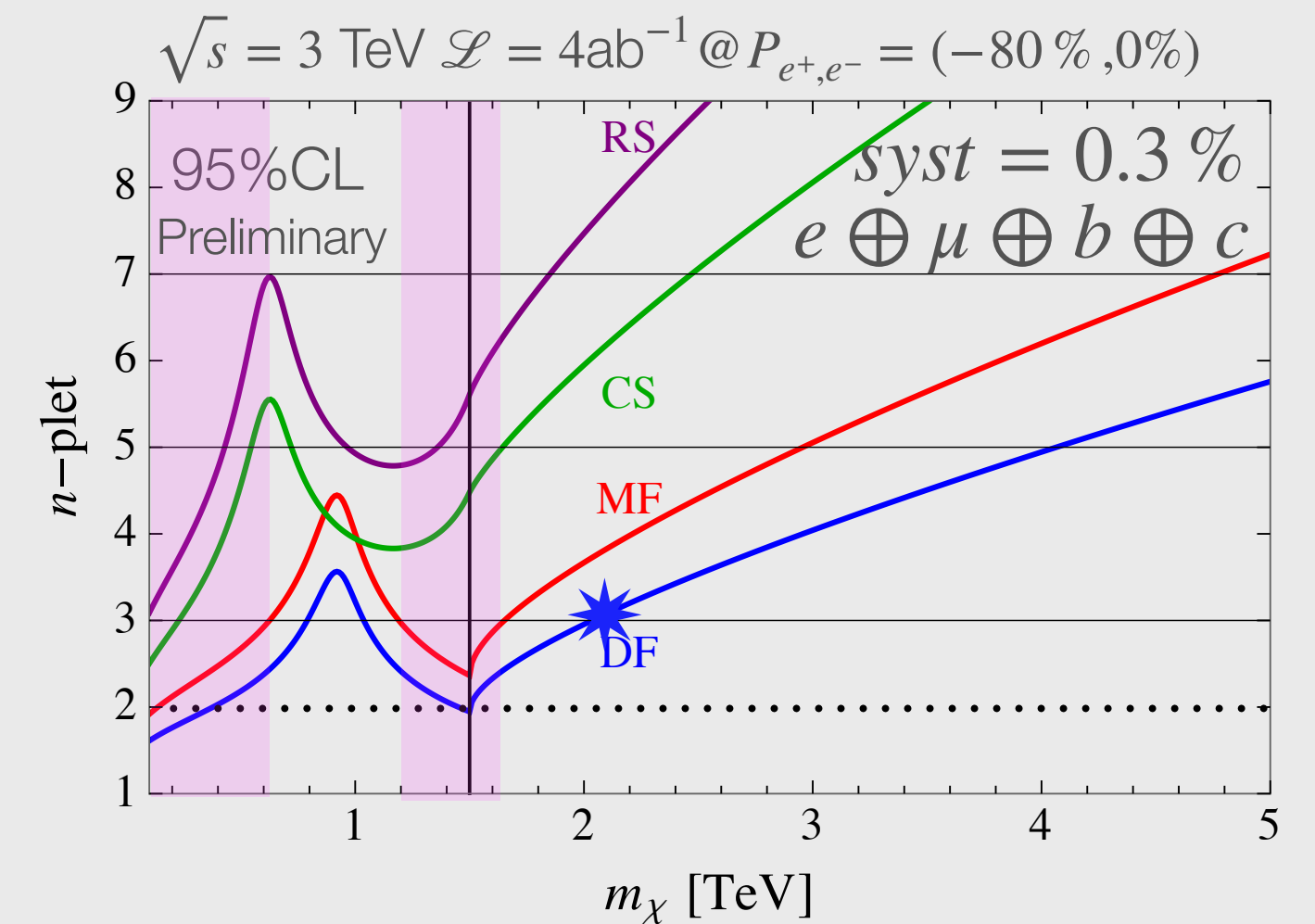
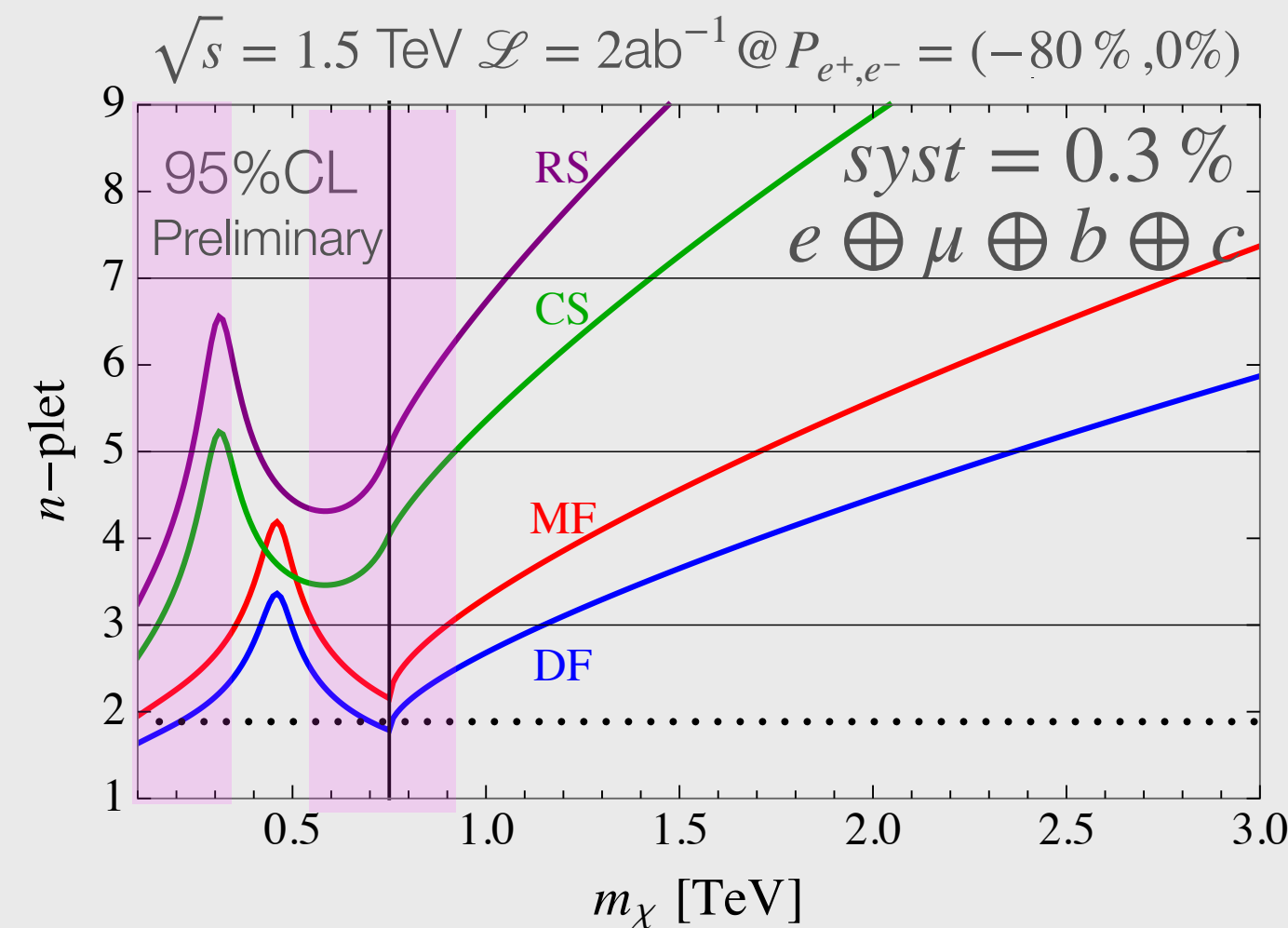
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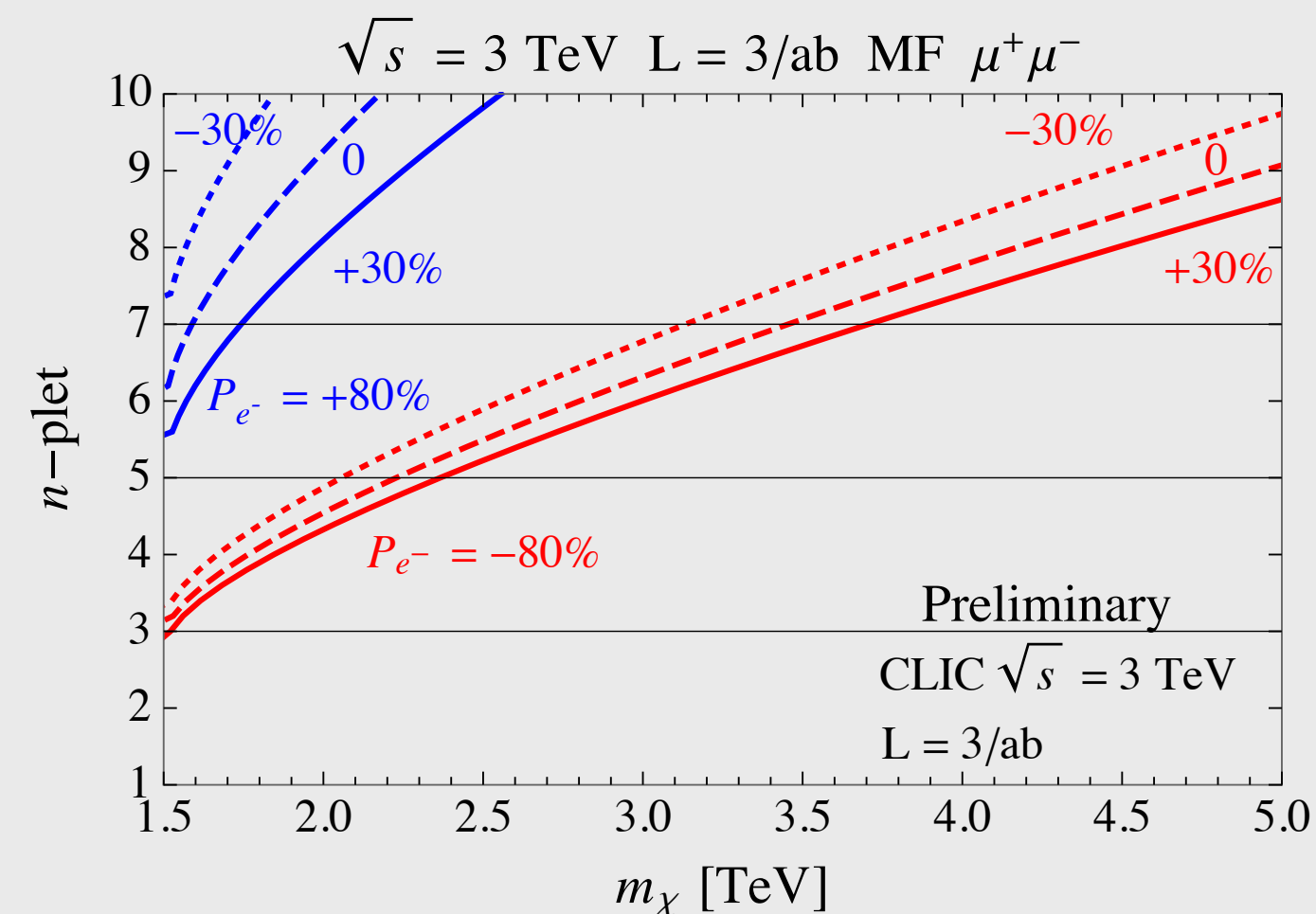
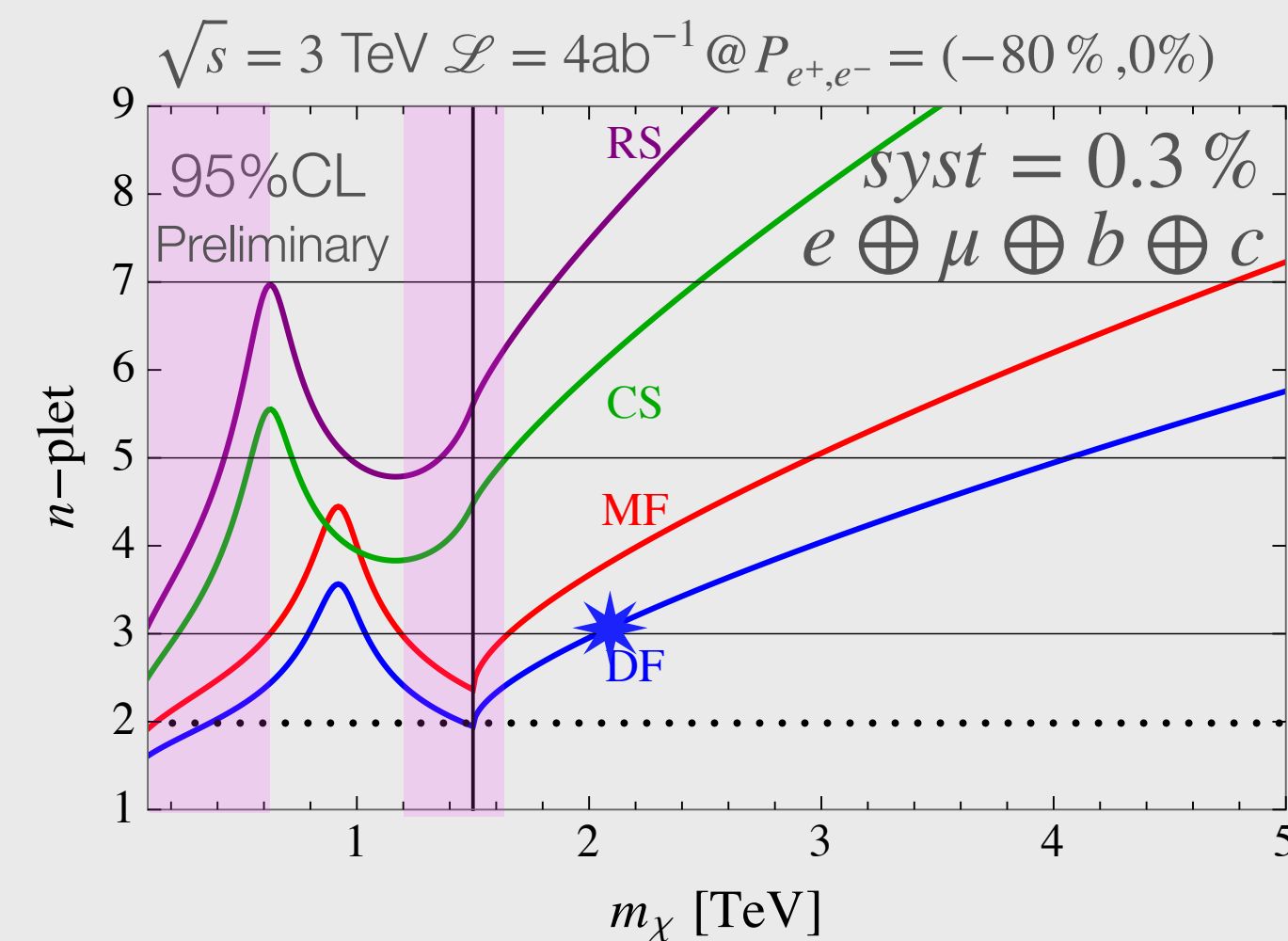
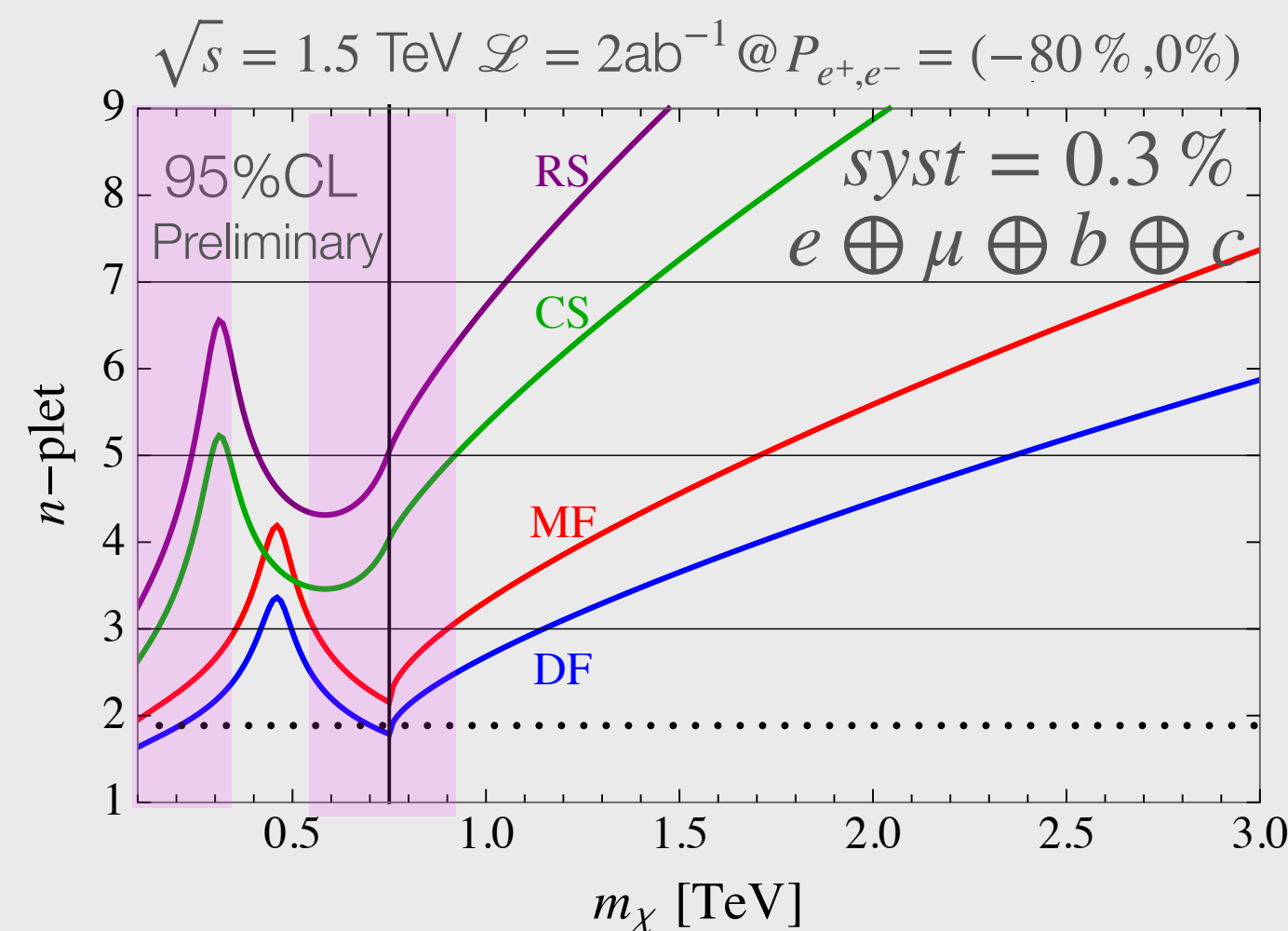
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PRECISION

ANGULAR DISTRIBUTION

Polarization is very advantageous

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Accidental Dark Matter 3-plet Dirac Fermion

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PRECISION

ANGULAR DISTRIBUTION

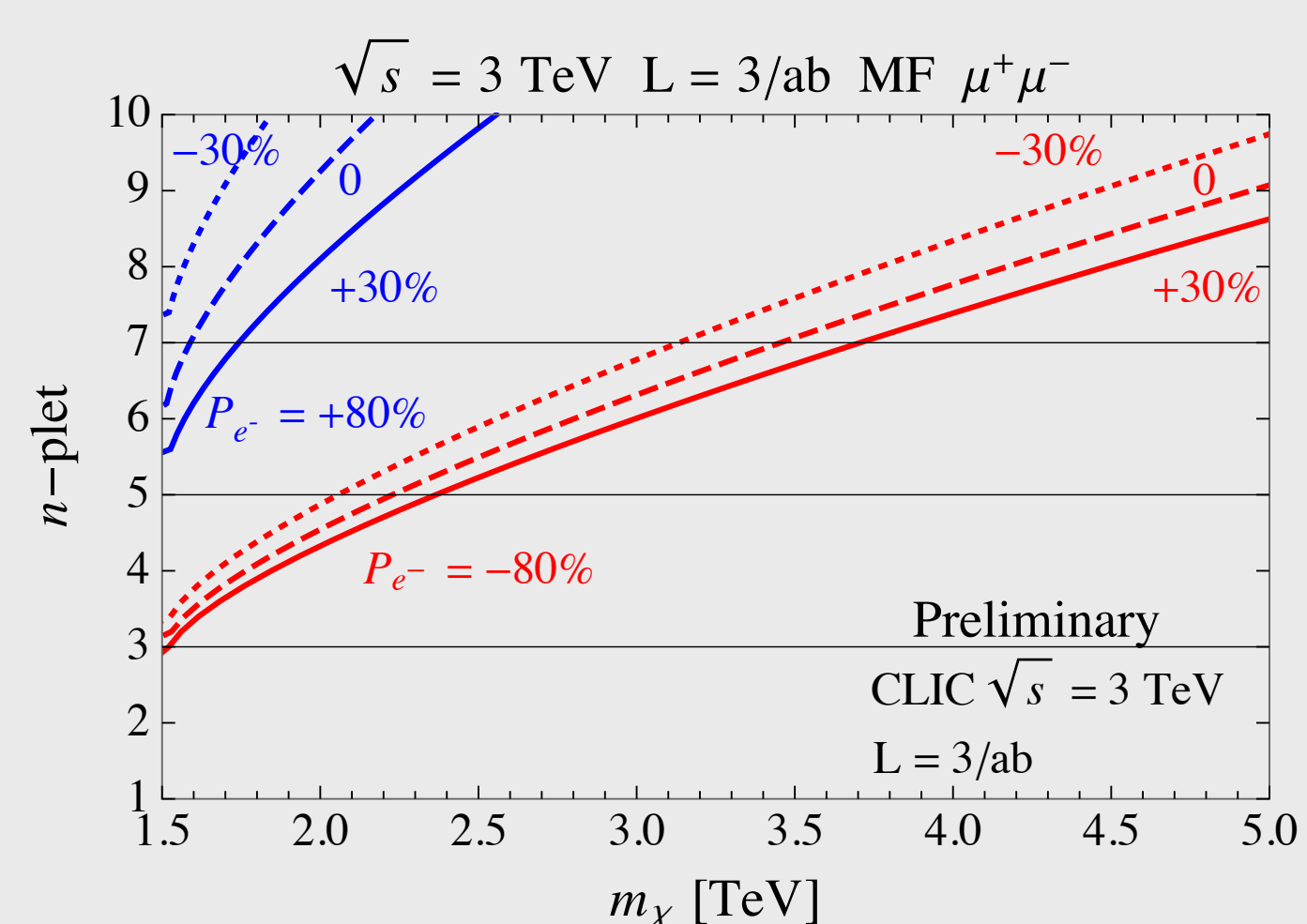
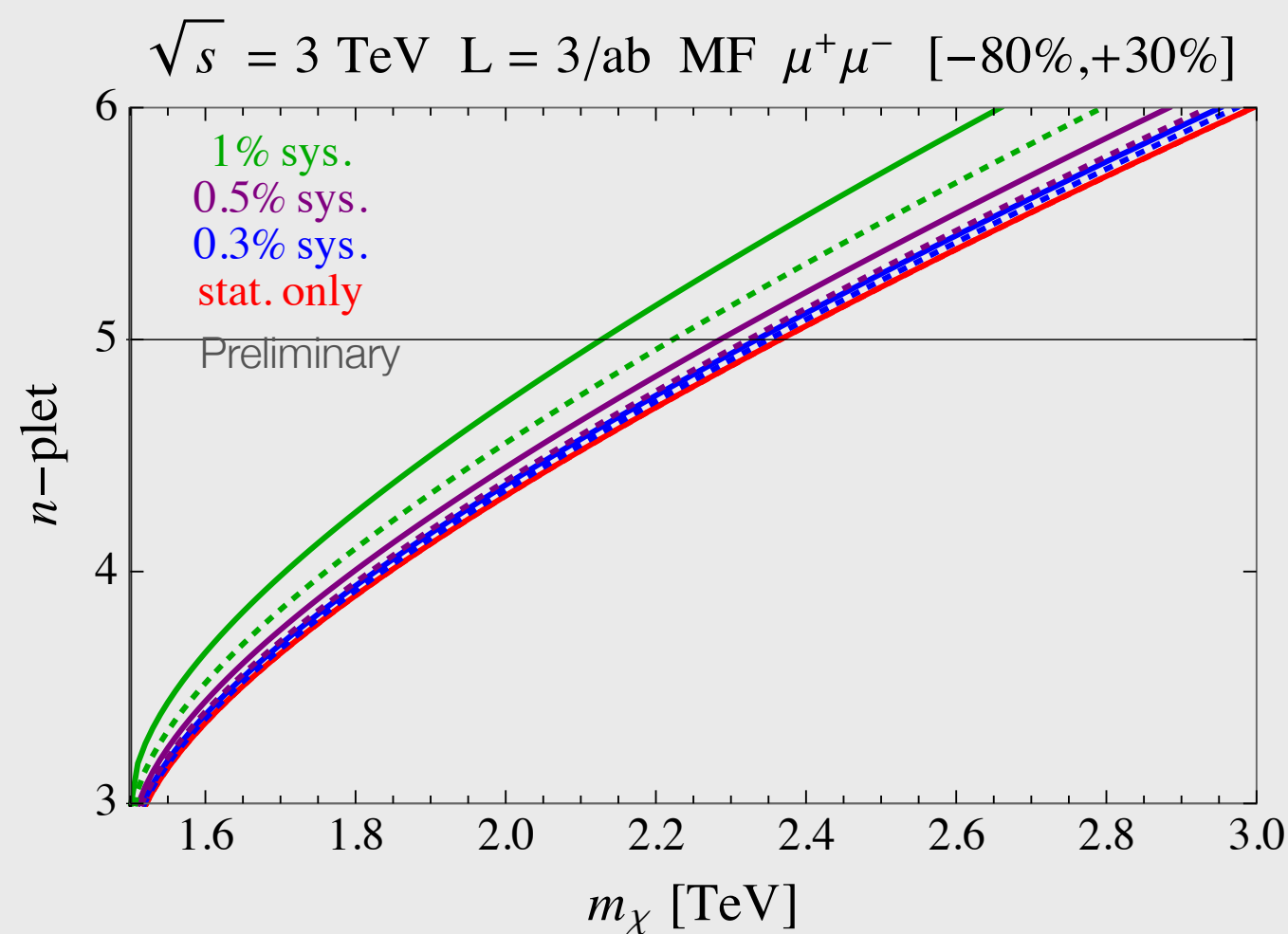
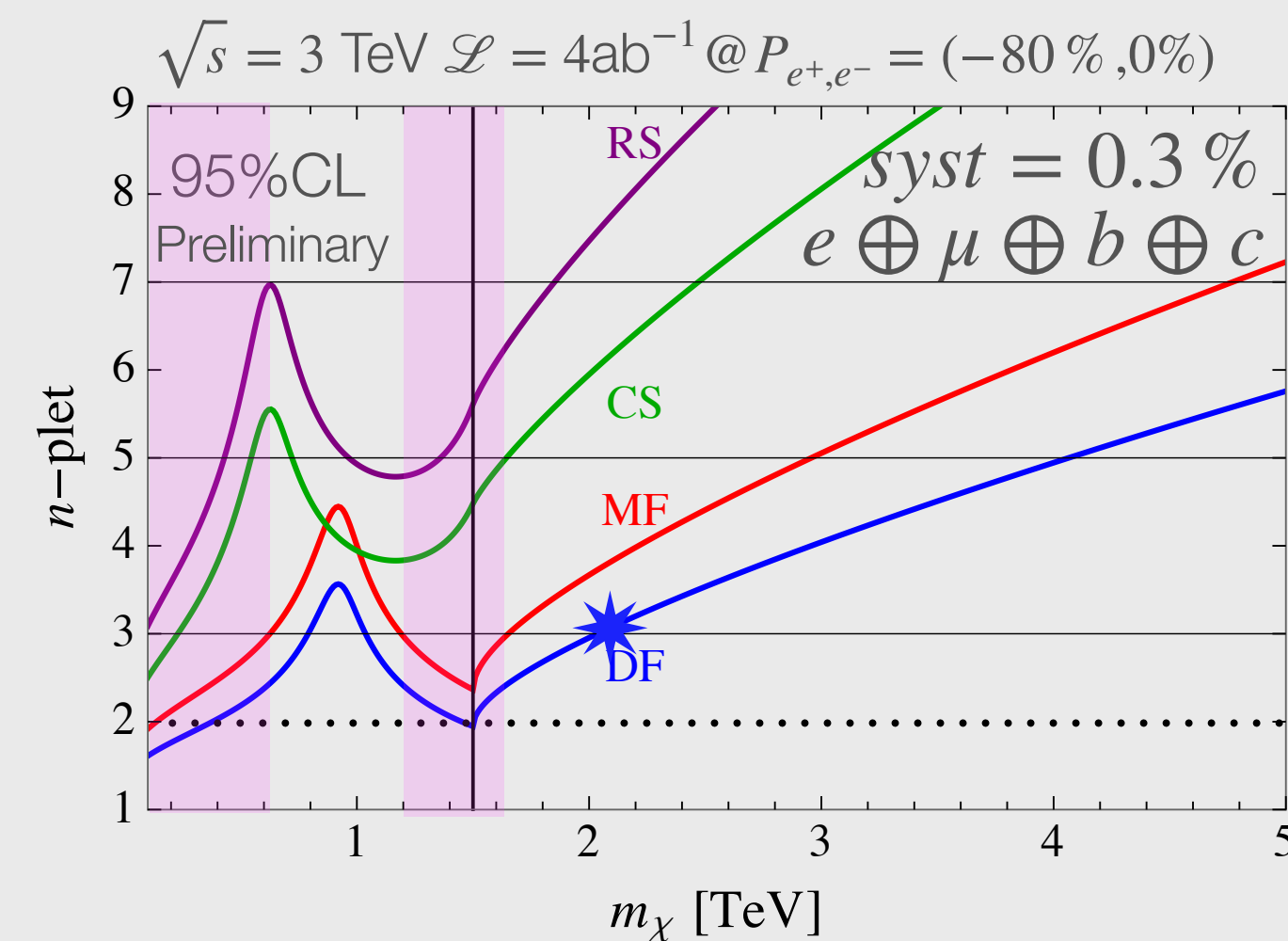
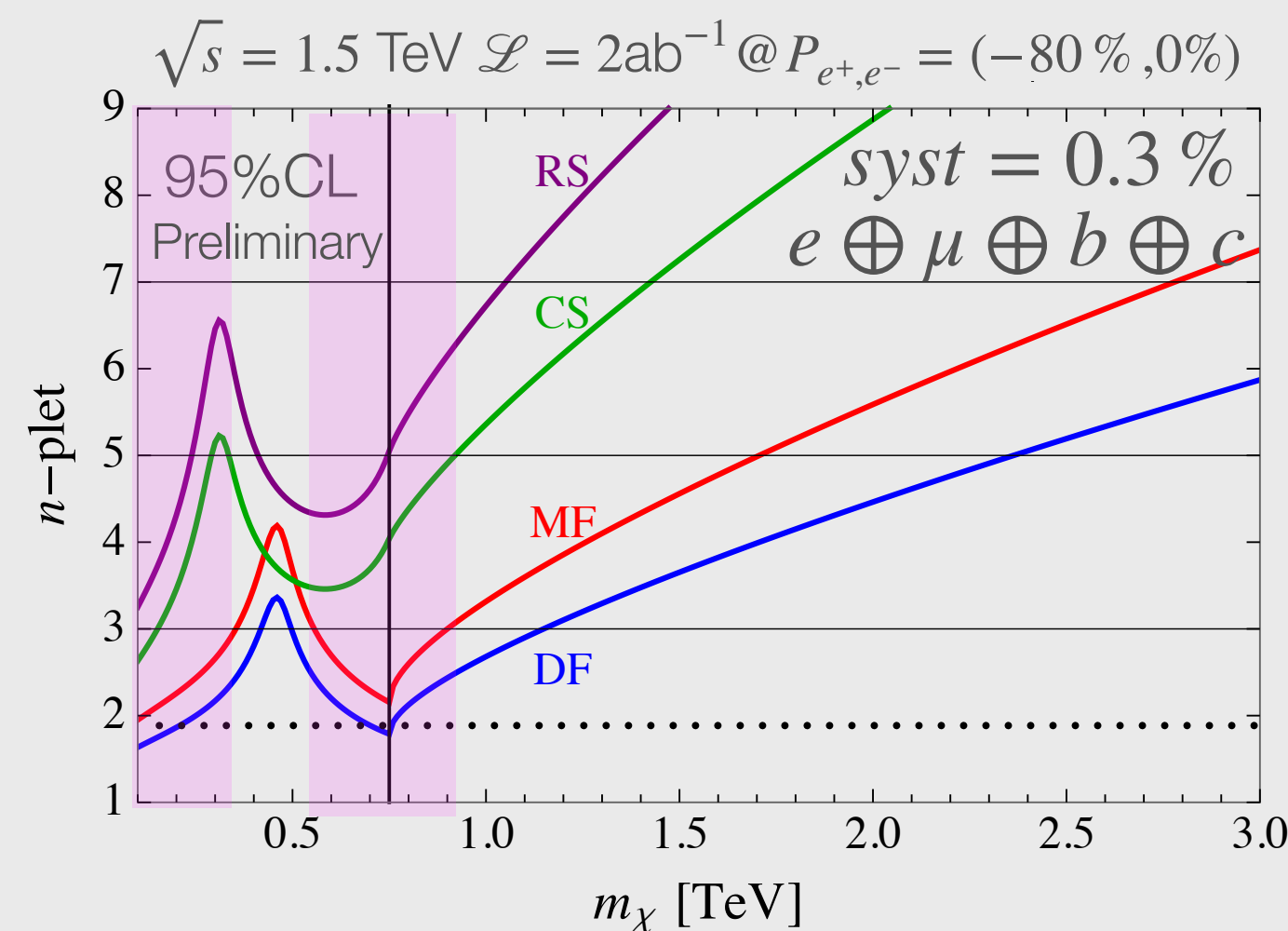
Polarization is very advantageous

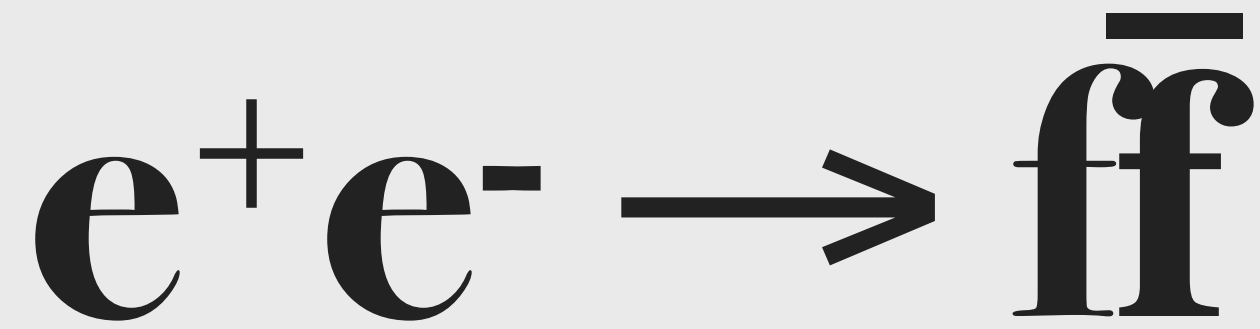
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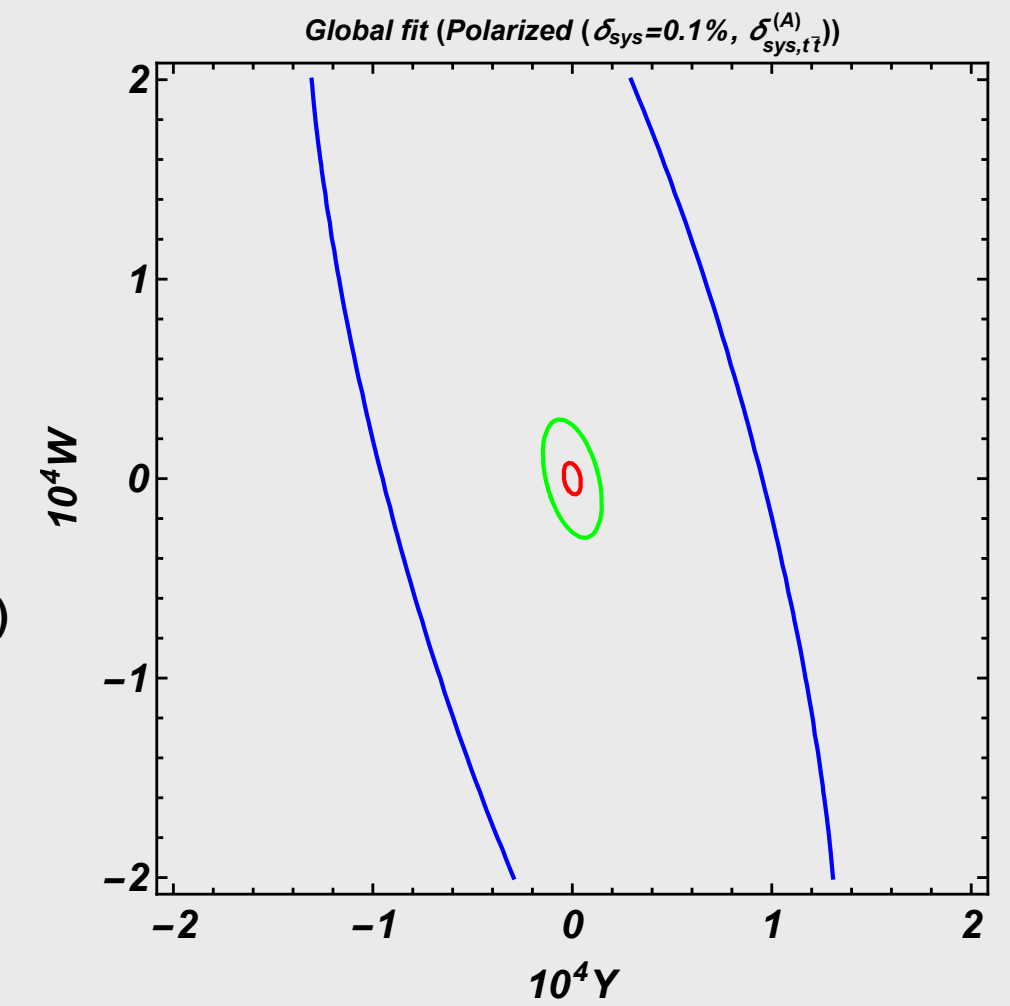
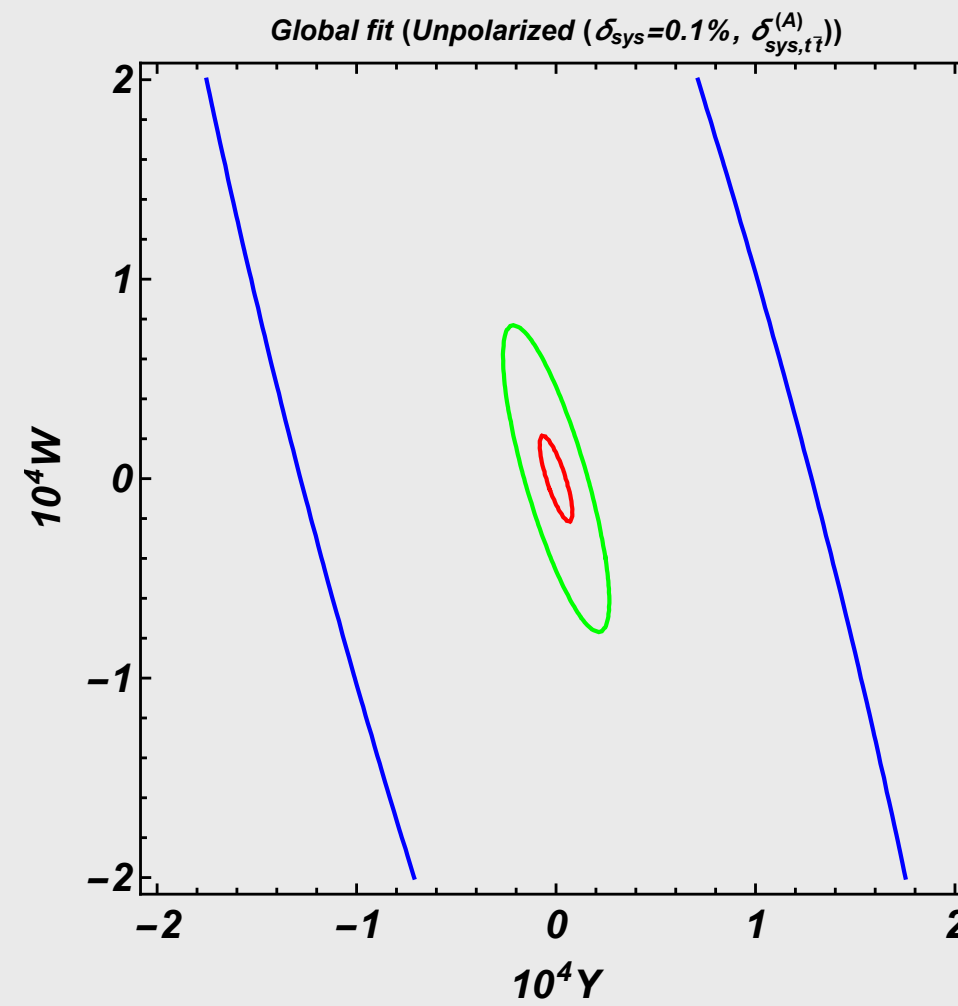
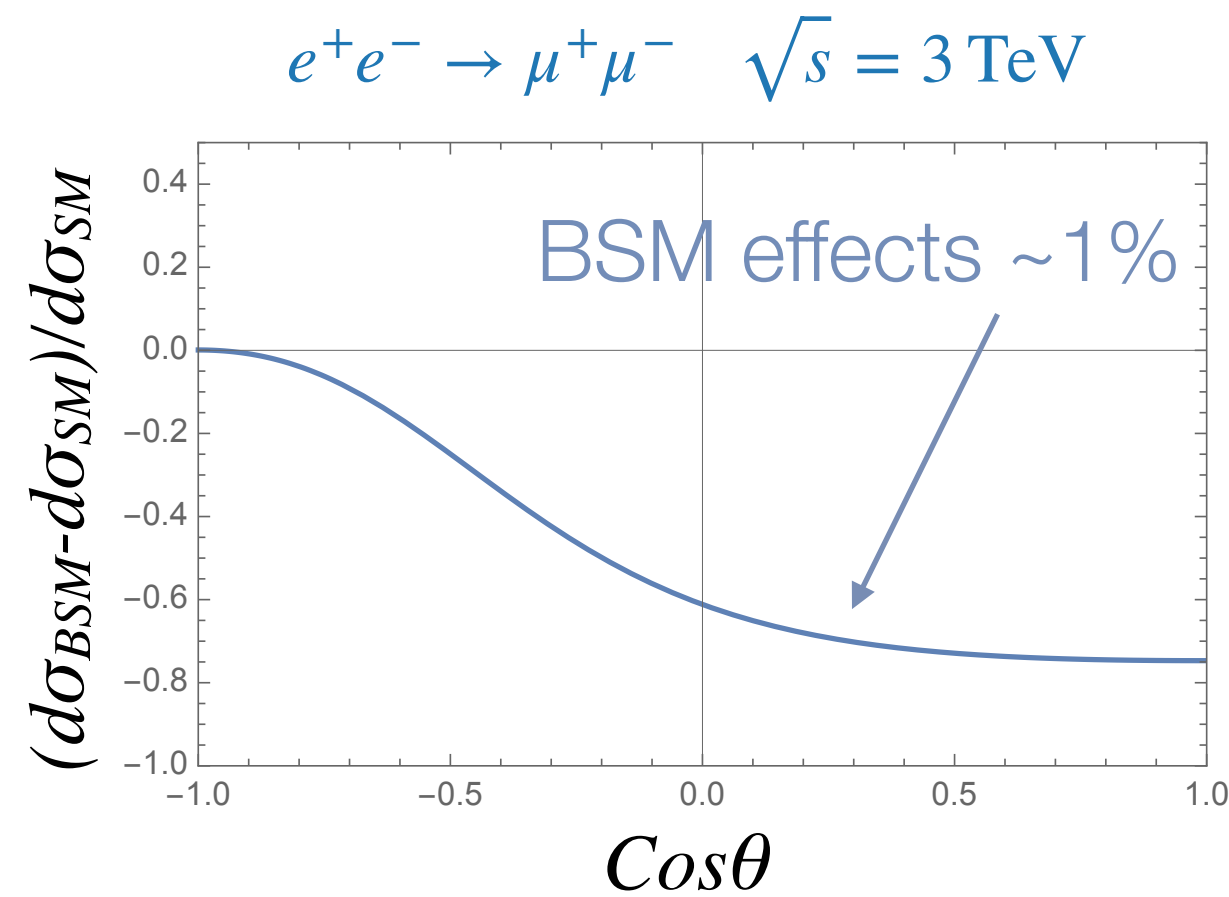
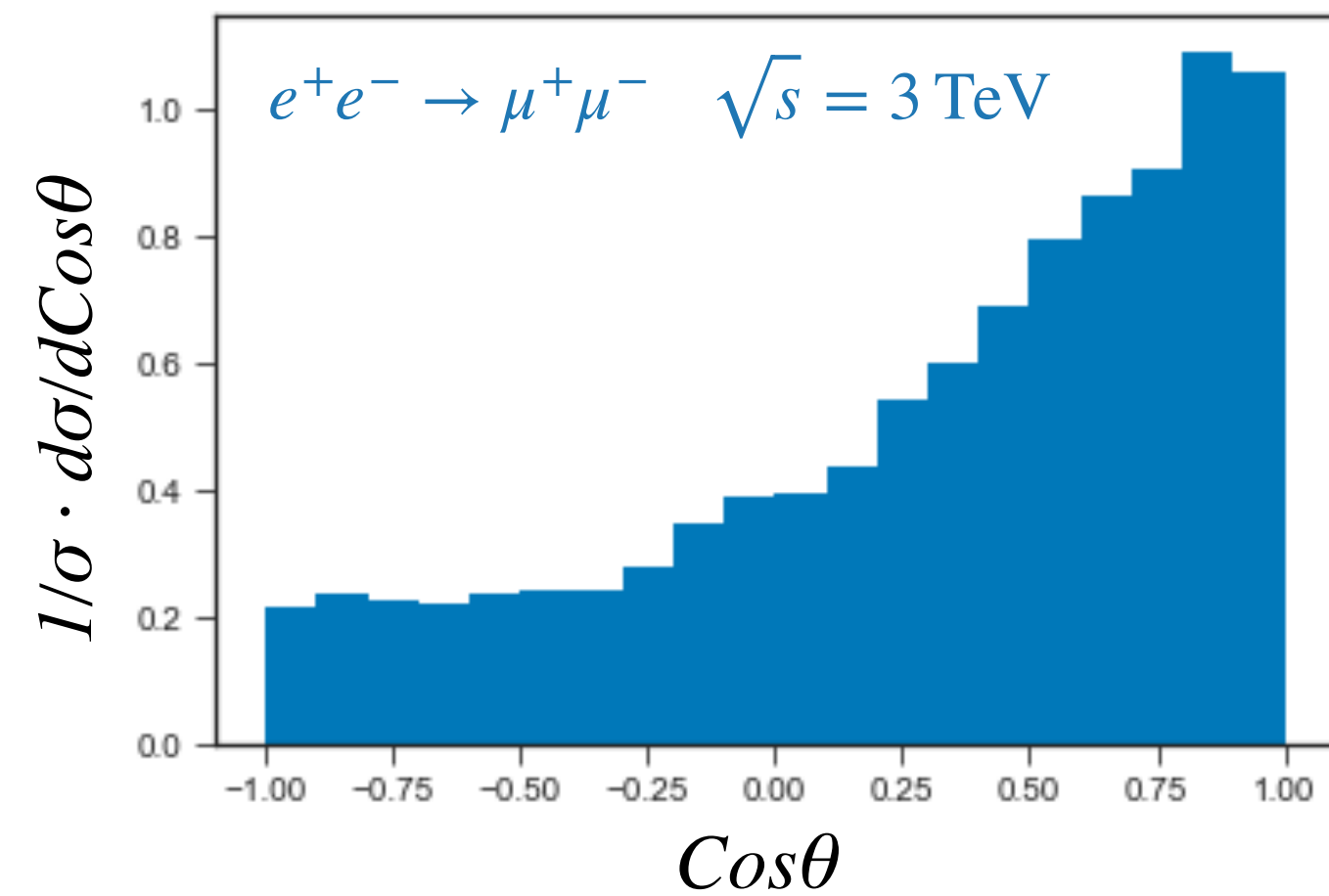
Accidental Dark Matter 3-plet Dirac Fermion



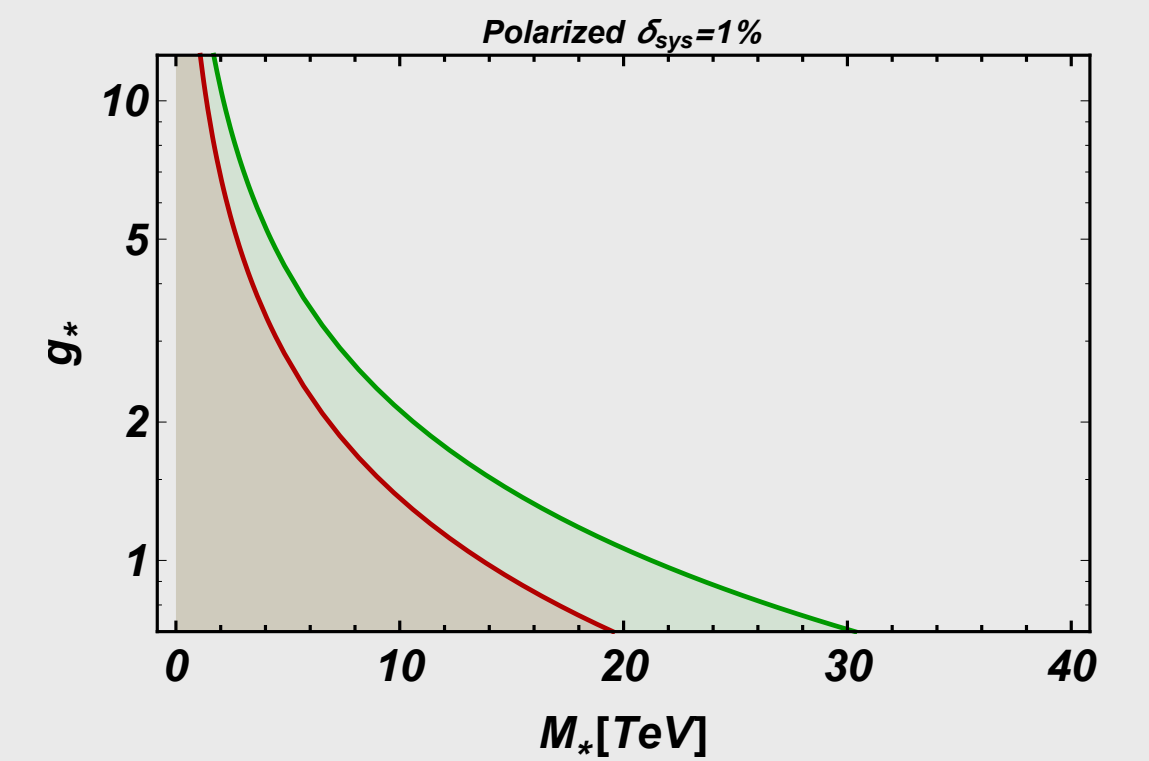
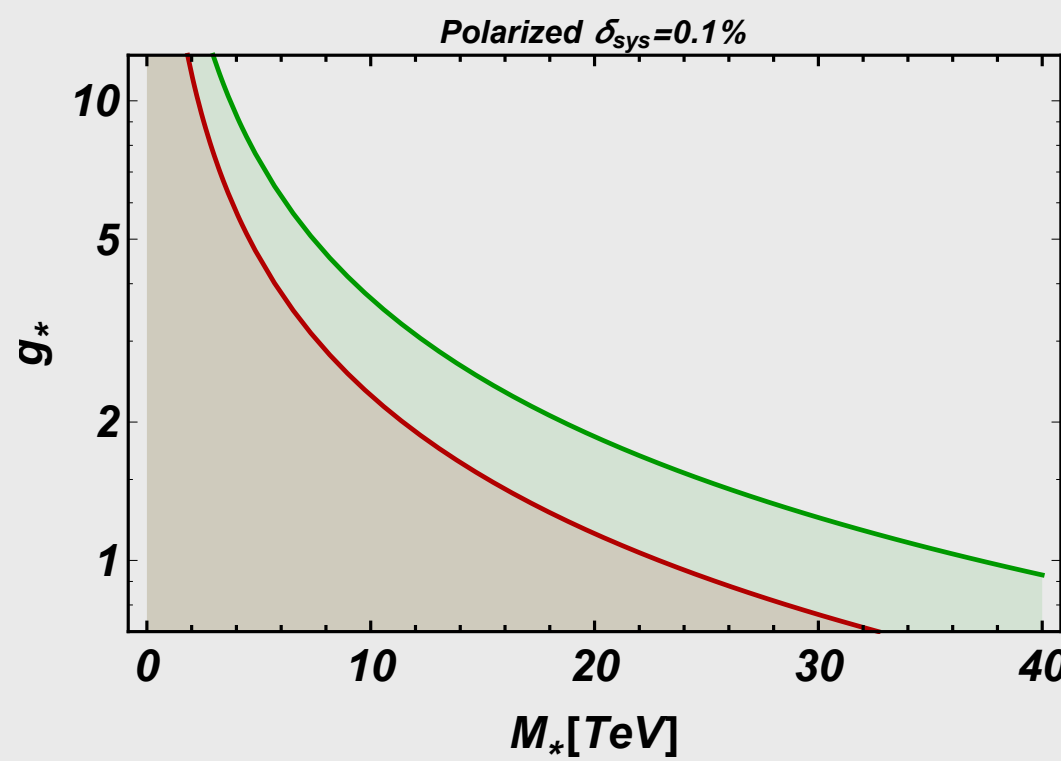


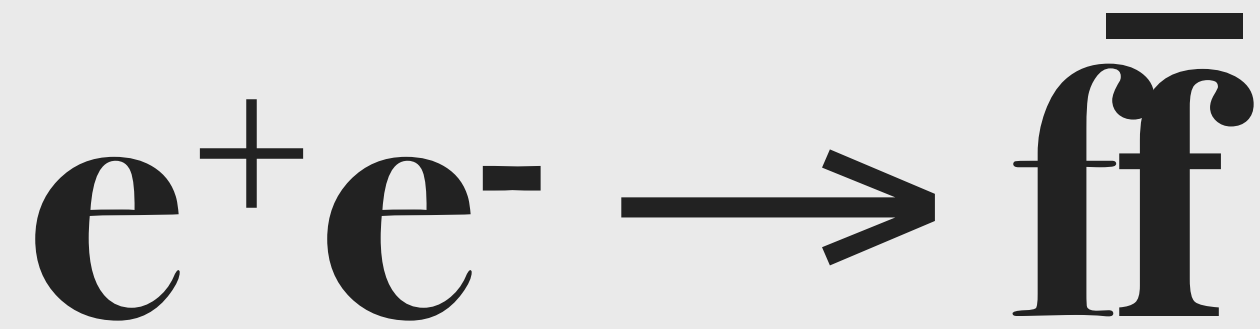
PRECISION

ANGULAR DISTRIBUTION



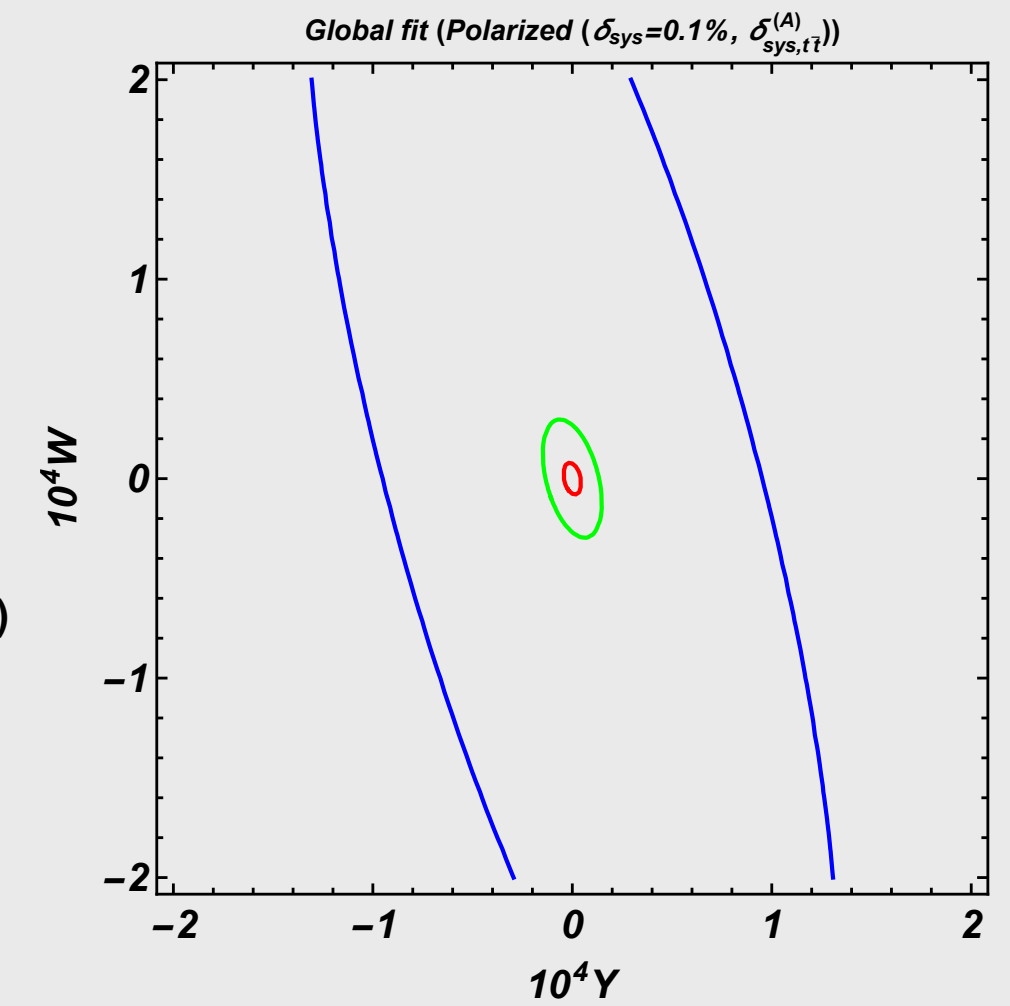
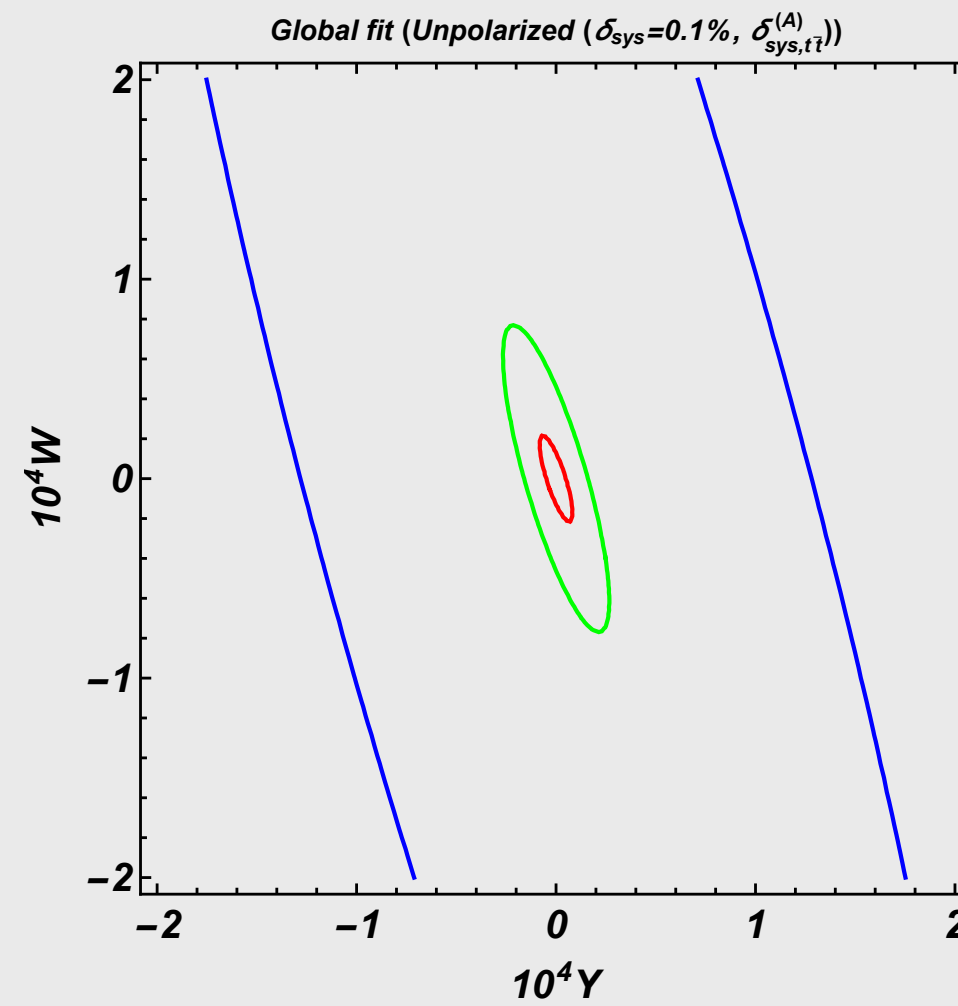
$$W = 2 \frac{g^2}{g_\star^2} \frac{M_W^2}{M_\star^2}, \quad Y = 2 \frac{g'^2}{g_\star^2} \frac{M_W^2}{M_\star^2}.$$





PRECISION

ANGULAR DISTRIBUTION



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χ^2 over 10 bins

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Systematic Unc.

