

Fitting the strong coupling constant for soft-drop thrust

Vincent Theeuwes

Georg-August-University Göttingen

In Collaboration with:

Jeremy Baron and Simone Marzani [[arXiv:1803.04719](https://arxiv.org/abs/1803.04719)]

Simone Marzani, Daniel Reichelt, Steffen Schumann, Gregory Soyez [[arXiv:1906.10504](https://arxiv.org/abs/1906.10504)]

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Federal Ministry
of Education
and Research

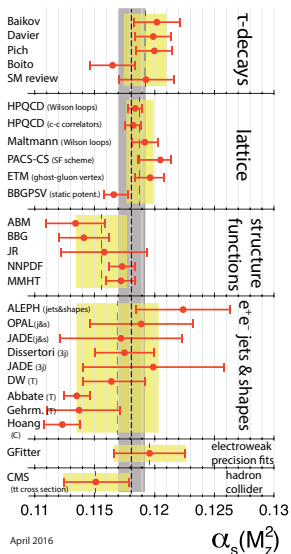


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The importance of α_s

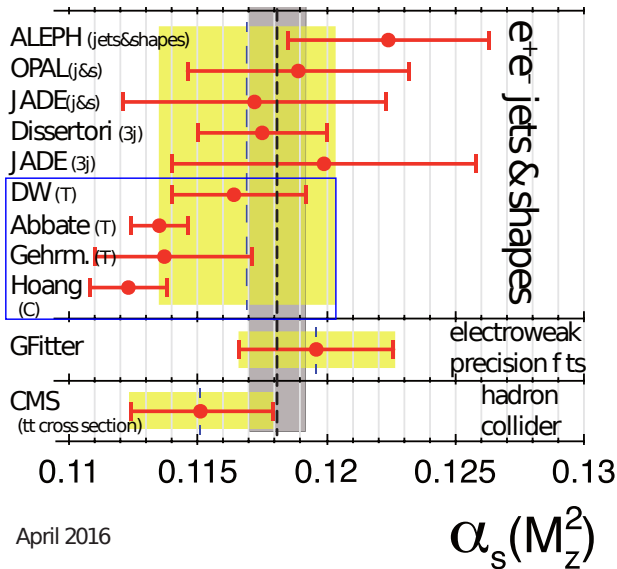
- No new physics at LHC \rightarrow requires **higher theoretical precision**
- Jet physics of importance to the LHC
- Higher order perturbative corrections shown to be important scale with higher powers of α_s
- Higgs boson production scales as α_s^2

An accurate measurement of α_s is necessary for LHC precision measurements

α_s Measurement

[Particle Data Group; '16]

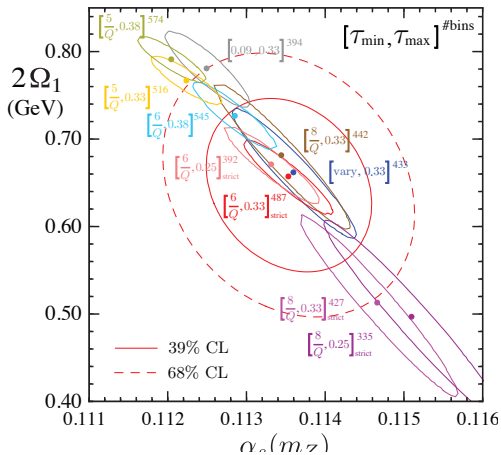
α_s Measurement



April 2016

[Particle Data Group; '16]

NP contributions



[Abbate, Fickinger, Hoang, Mateu, Stewart; '10]

mMDT & Soft drop

Technique we will deal with is soft drop: [Larkoski, Marzani, Soyez, Thaler; '14]

$$\frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}} > z_c \left(\frac{\Delta R_{12}}{R} \right)^\beta$$

or at e^+e^- colliders:

$$\frac{\min[E_i, E_j]}{E_i + E_j} > z_{\text{cut}} (1 - \cos \theta_{ij})^{\beta/2}$$

Makes use of Cambridge/Aachen clustering.

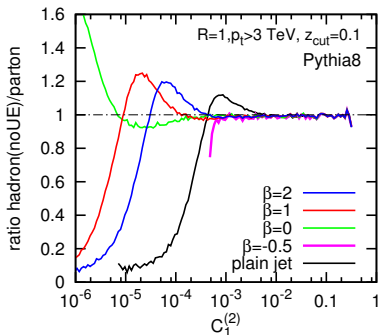
[Dokshitzer, Leder, Moretti, Webber; '97][Wobisch, Wengler; '99]

Reduces to modified Mass Drop Tagger (mMDT) for $\beta = 0$

[Dasgupta, Fregoso, Marzani, Salam; '13]

Reduction in NP corrections

[Larkoski, Marzani, Soyez, Thaler; '14]



Reduces non-perturbative effects, such as hadronization

Further work

- Computation of soft drop using SCET at NNLL accuracy approximated for $e_2^{(2)} \ll z_{cut}$ [Frye, Larkoski, Schwartz, Yan; '16]
- Calculation in dQCD including finite z_{cut} effects [Marzani, Schunk, Soyez; '17]
- Including jet radius resummation in SCET [Kang, Lee, Liu, Ringer; '18]
- Good agreement to experiments [CMS;'17] [ATLAS;'17]
- Application to top quark mass measurements [Hoang, Mantry, Pathak, Stewart;'17] [Andreassen, Schwartz;'17]
- And α_s measurements at LHC [Les Houches;'18] and e^+e^- [Baron, Marzani, VT; '18][Marzani, Reichelt, Schumann, Soyez, VT; 19]
- Calculation of non-perturbative effects [Hoang, Mantry, Pathak, Stewart, '19]

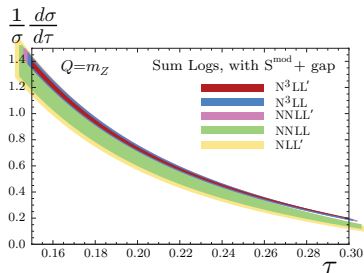
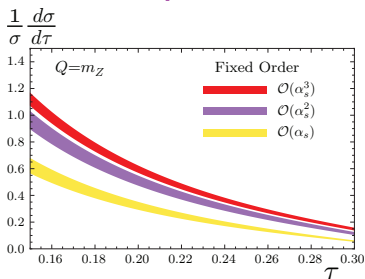
Talk by Aditya Pathak

Thrust

$$\tau = 1 - T = \min_{\vec{n}} \left(1 - \frac{\sum_i |\vec{n} \cdot \vec{p}_i|}{\sum_i |\vec{p}_i|} \right)$$

Minimize for thrust axis \vec{n}

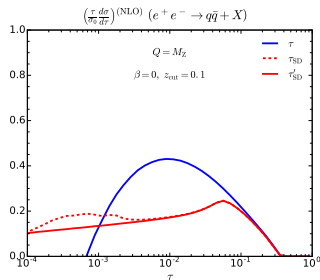
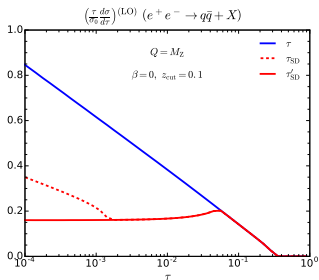
[Abbate, Fickinger, Hoang, Mateu, Stewart; '10]



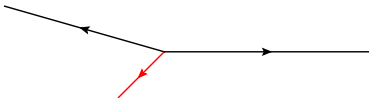
SD Distribution

Hemisphere jets at an e^+e^- collider \rightarrow Different soft drop condition:

$$\frac{\min[E_i, E_j]}{E_i + E_j} > z_{\text{cut}} (1 - \cos \theta_{ij})^{\beta/2}$$



Alternative definition

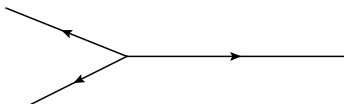


- Separation into two jets at the hand of thrust axis pre-softdrop
- After softdrop each hemisphere will have its own axis
- Each thrust axis is the jet axis

$$\tau_{SD} = \frac{\sum_{i \in \mathcal{E}_{SD}} |\vec{p}_i|}{\sum_{i \in \mathcal{E}} |\vec{p}_i|} \left[1 - \frac{\sum_{i \in \mathcal{H}_{SD}^L} |\vec{n}_L \cdot \vec{p}_i| + \sum_{i \in \mathcal{H}_{SD}^R} |\vec{n}_R \cdot \vec{p}_i|}{\sum_{i \in \mathcal{E}_{SD}} |\vec{p}_i|} \right]$$

Factor for collinear safety

Collinear safety



Failure of cancellation between virtual and real:
Virtual corrections at

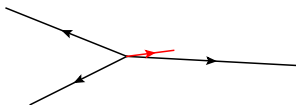
$$\tau = 1 - \frac{2E_R}{Q} = \frac{Q - 2E_R}{Q}$$

While a collinear emission that is groomed results in:

$$\tau_{SD}^{unsc.} = 1 - \frac{E_R}{Q - zE_R} - \frac{E_R(1-z)}{Q - zE_R} = \frac{Q - 2E_R}{Q - zE_R}$$

Rescaled soft-drop thrust fixes this. The issue is unique to a single particle in one hemisphere

Collinear safety



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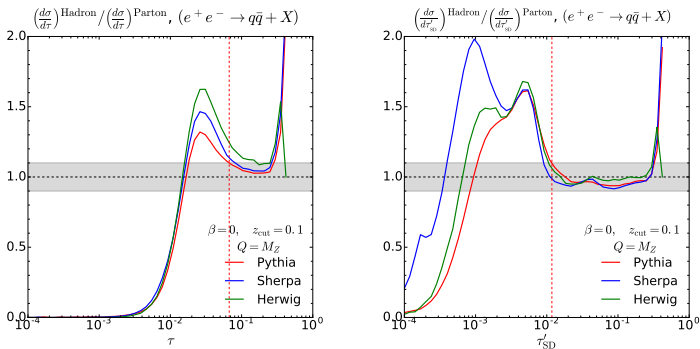
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Reduction of hadronization with soft drop

[Baron, Marzani, VT; '18]



Non-perturbative corrections above 10% from $\tau \simeq 0.07$ to $\tau \simeq 0.001$
 Reduction in hadronization corrections.

Hadronic thrust UE \rightarrow poster by Jeremy Baron

Factorization

Factorization for $\tau \ll z_{\text{cut}} \ll 1$ [Frye, Larkoski, Schwartz, Yan; '16]:

$$\frac{d\sigma}{d\tau} = H(Q) S_G(z_{\text{cut}}, \beta) [S_C(\tau, z_{\text{cut}}, \beta) \otimes J(\tau)]^2$$

Computed in Laplace space and inverted leading to:

$$\Sigma(\tau) = \left[1 + \left(\frac{\alpha_s}{\pi} \right) C^{(1)} + \dots \right] \exp \left[\frac{1}{\alpha_s} g_1(-\lambda_\tau, \lambda_{z_{\text{cut}}}) + g_2(-\lambda_\tau, \lambda_{z_{\text{cut}}}) + \dots \right]$$

for $\lambda_x = \alpha_s b_0 \log x$ and confirmed using dQCD.

Multiple emission correction

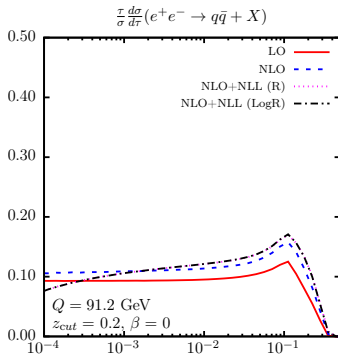
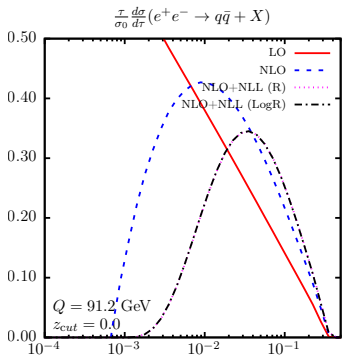
Additional transition point effect from multiple emissions.

$$\Sigma_{\text{res}}|_{\tau > \text{tr}} = C \frac{e^{R(\tau)}}{\Gamma(1 - R'(\tau))} \exp \left[R'(\tau) \frac{\text{tr}}{\tau} {}_3F_2 \left(1, 1, 1 + R'(\tau); 2, 2; \frac{\text{tr}}{\tau} \right) \right] \\ \times \exp \left[\log \left(\frac{\tau}{\text{tr}} \right) \{ R''(\tau) - R''(\tau = \text{tr}, z_{\text{cut}}) \} \right]$$

Originates from extrapolation of derivative in the inversion across the transition point.

Resummation results

[Marzani, Reichelt, Schumann, Soyez, VT; '19]



- Significant corrections from both fixed order and resummation
- Agreement between matching schemes
- In order to take into account $\tau \sim z_{cut}$ shift transition point to $\text{tr} = z_{cut} 2^{\beta/2}$, which holds at NLL

Fitting setup

- Compared to ALEPH data for thrust [*ALEPH;'04*]
- Fit to MEPS@NLO 2-5j Sherpa result for soft drop [*Schumann, Krauss;'07*] [*Gehrmann, Hoche, Krauss, Schonherr, Siegert;'12*]
- Assume ALEPH uncertainty with $\sqrt{\sigma}$ scaling for statistical uncertainty
- Same systematic uncertainty with minimal correlation

- Minimization of chi-squared
- Using range $\tau \in [0.06, 0.25]$

Setup is similar to a combination of methods from

[*Abbate, Fickinger, Hoang, Mateu, Stewart;'10*] and [*Gehrmann, Luisoni, Monni;'12*]

Non-perturbative corrections Analytical

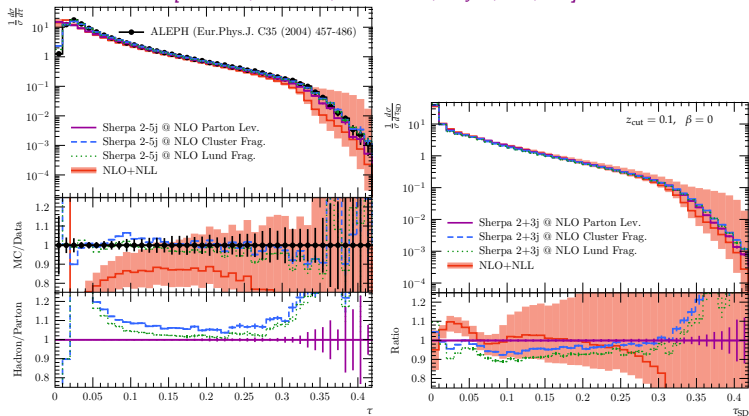
For the non-perturbative model use a shift in τ and z_{cut} from

[Dasgupta, Fregoso, Marzani, Salam;'13] [Marzani, Schunk, Soyez;'17]

- Shift in τ from non-perturbative emission within a cone defined by thrust
- Shift in z_{cut} from non-perturbative emission that reduces energy leading to it being groomed
- Both are computed in a $2 \rightarrow 2 + \text{NP}$ with small τ limit configuration
- More detailed calculation can be performed as a better approximation

Non-perturbative corrections Monte Carlo

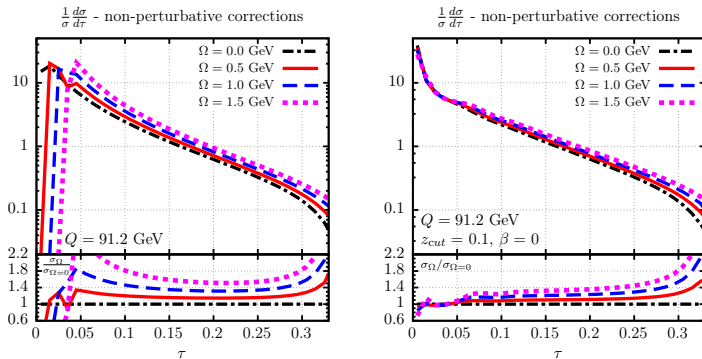
[Marzani, Reichelt, Schumann, Soyez, VT; '19]



- Based on MEPS@NLO Sherpa 2+3j to have comparable accuracy
- Significant differences between hadronization models
- Better agreement analytical model to PL

Comparison hadronization models

[Marzani, Reichelt, Schumann, Soyez, VT; '19]



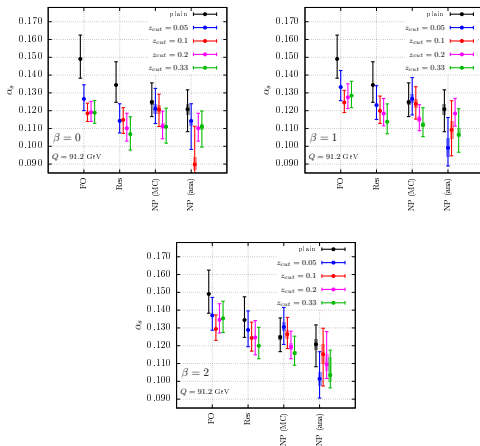
- Plain thrust replicates shape
- Increase instead of decrease above transition point

Uncertainties

- Experimental uncertainty at the hand of $\Delta\chi^2 = 1$ for fixed hadronization
- Theoretical uncertainty based all extreme variations:
 - Variation of μ_R and x_L with $1/2 < \mu_R x_L / Q < 2$
 - Variation of end-point p variable 1 or 2
 - Switch between matching scheme: R or LogR matching
- Hadronization uncertainty:
 - Analytical model: $\Delta\chi^2 = 1$ for varying Ω
 - Monte Carlo model: Switch between Clustering [*Winter, Krauss, Soff, '03*] and Lund String [*Sjöstrand, '82*]

Variation of z_{cut} and β

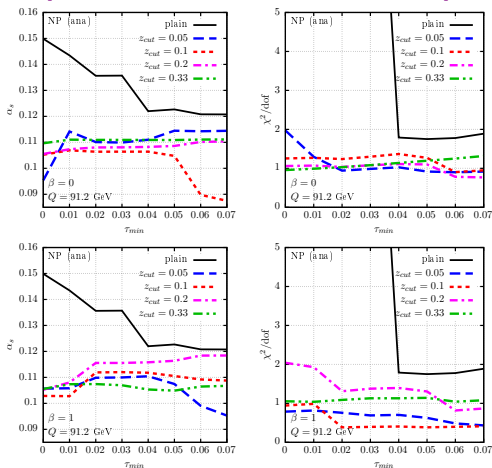
[Marzani, Reichelt, Schumann, Soyez, VT; '19]



Significant reduction in shift from hadronization corrections

Variation of minimum of fit range

[Marzani, Reichelt, Schumann, Soyez, VT; '19]

Stabilization of the fit over a longer range of τ_{\min}

Summary

Conclusions

- Better agreement between analytical results and PL MC
- Soft drop helps reduce non-perturbative effects for α_s with thrust
- Fits are more stable with respect to an increase in fitting range
- Conclusions do not depend on hadronization model

Future work

- Study other observables and definitions of soft drop
- More precise calculation of analytical hadronization model
Possibly through means of *[Hoang, Mantry, Pathak, Stewart, '19]*
- Higher accuracy fixed order and resummation
- Actual experimental measurement
- Extension to LHC

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Thank you for your attention