

Complex Singlet Extension of the SM (CxSM)

Ref: V. Barger et al. Phys. Rev. D79(2009)015018

SM + S : Gauge Singlet
Complex Scalar Field

$$V = \frac{m^2}{2}|H|^2 + \frac{\lambda}{4}|H|^4 + \frac{\delta_2}{2}|H|^2|S|^2 + \frac{b_2}{2}|S|^2 + \frac{d_2}{4}|S|^4 + \left(a_1 S + \frac{b_1}{4} S^2 + \text{c.c.} \right)$$

soft breaking of the global U(1)

H, S take the vacuum expectation values (vev)

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v+h \end{pmatrix}, \quad S = \frac{1}{\sqrt{2}}(v_s + s + i\xi)$$

mass eigenstates (h1, h2) described in the mixing angle α

$$\begin{pmatrix} h \\ s \end{pmatrix} = \begin{pmatrix} \cos\alpha & -\sin\alpha \\ \sin\alpha & \cos\alpha \end{pmatrix} \begin{pmatrix} h_1 \\ h_2 \end{pmatrix}$$

diagonalization of M^2

$$O^T M^2 O = \text{diag.} (m_{h_1}^2, m_{h_2}^2), \quad O \equiv \begin{pmatrix} \cos\alpha & -\sin\alpha \\ \sin\alpha & \cos\alpha \end{pmatrix}$$

$$m_{h_{1,2}}^2 = \frac{1}{2} \left(\frac{\lambda}{2} v^2 + \bar{\Lambda}_S^2 \mp \sqrt{\left(\frac{\lambda}{2} v^2 - \bar{\Lambda}_S^2 \right)^2 + 4 \left(\frac{\delta_2}{2} v v_S \right)^2} \right)$$

$$\cos 2\alpha = \frac{\frac{\lambda}{2} v^2 - \bar{\Lambda}_S^2}{m_{h_1}^2 - m_{h_2}^2}$$

S : interaction between H

→ h_1, h_2 : interaction with SM just through H

$$h = h_1 \cos\alpha - h_2 \sin\alpha$$

$$s = h_1 \sin\alpha + h_2 \cos\alpha$$

$$\begin{pmatrix} h_1 : 10 \text{ GeV} \\ h_2 : 125 \text{ GeV} \end{pmatrix}$$

$$h_2 \rightarrow h_{SM} \text{ @ } \alpha = -\frac{\pi}{2}$$

theoretical restriction on parameters

① For the potential V to have a global minimum

$$\lambda > 0, \quad d_2 > 0, \quad \lambda d_2 > \delta_2^2$$

$$\lambda \left(d_2 + \frac{2\sqrt{2}a_1}{v_S^3} \right) > \delta_2^2$$

② perturbativity limit

$$\lambda, d_2 \leq \frac{16\pi}{3}$$

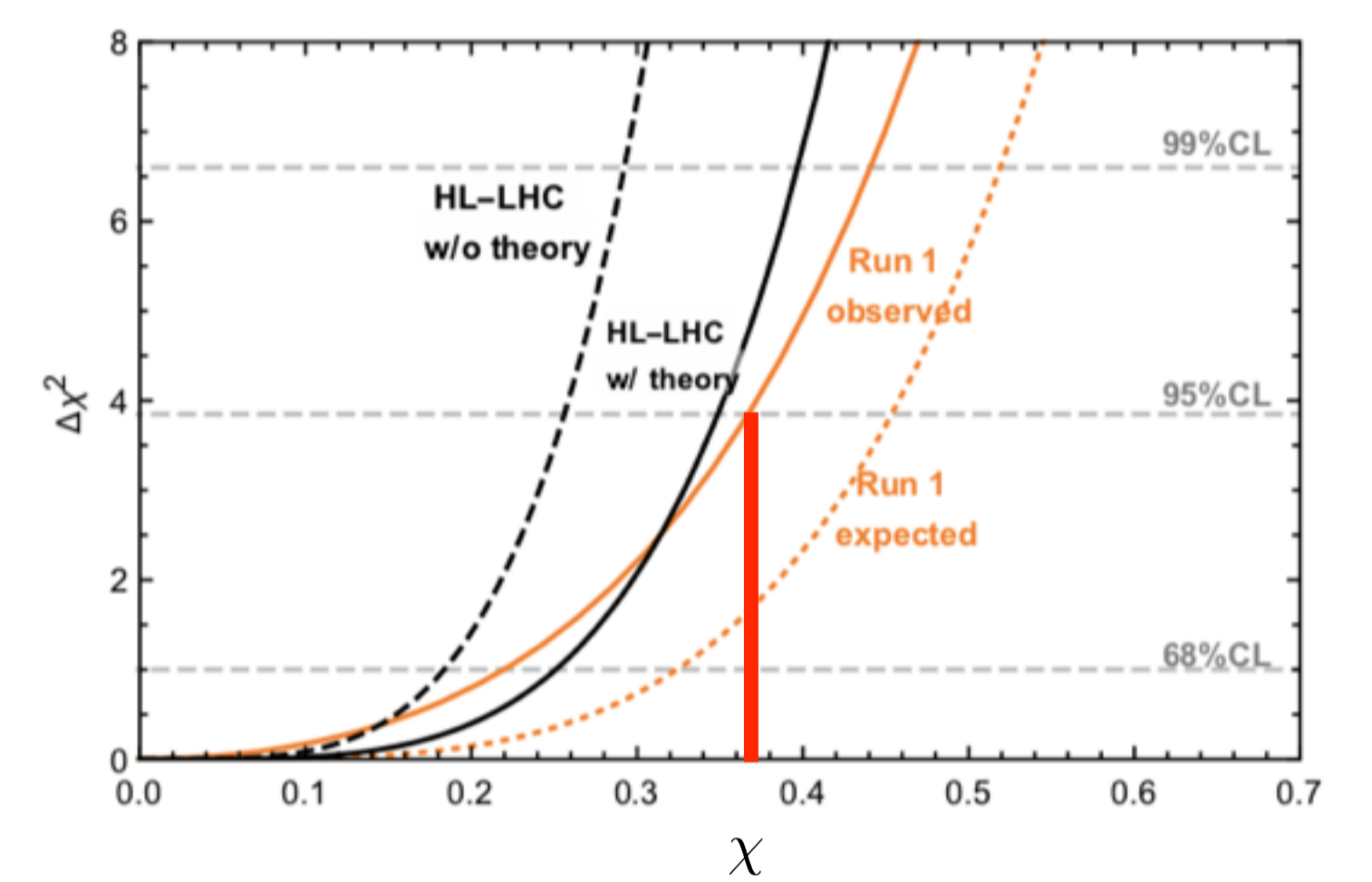
allowed region of mixing angle α

signal strength

$$\mu \equiv \frac{[\sigma(\text{XX} \rightarrow h)\text{Br}(h \rightarrow \text{YY})]_{\text{exp}}}{[\sigma(\text{XX} \rightarrow h)\text{Br}(h \rightarrow \text{YY})]_{\text{SM}}} = 1 - \chi^2$$

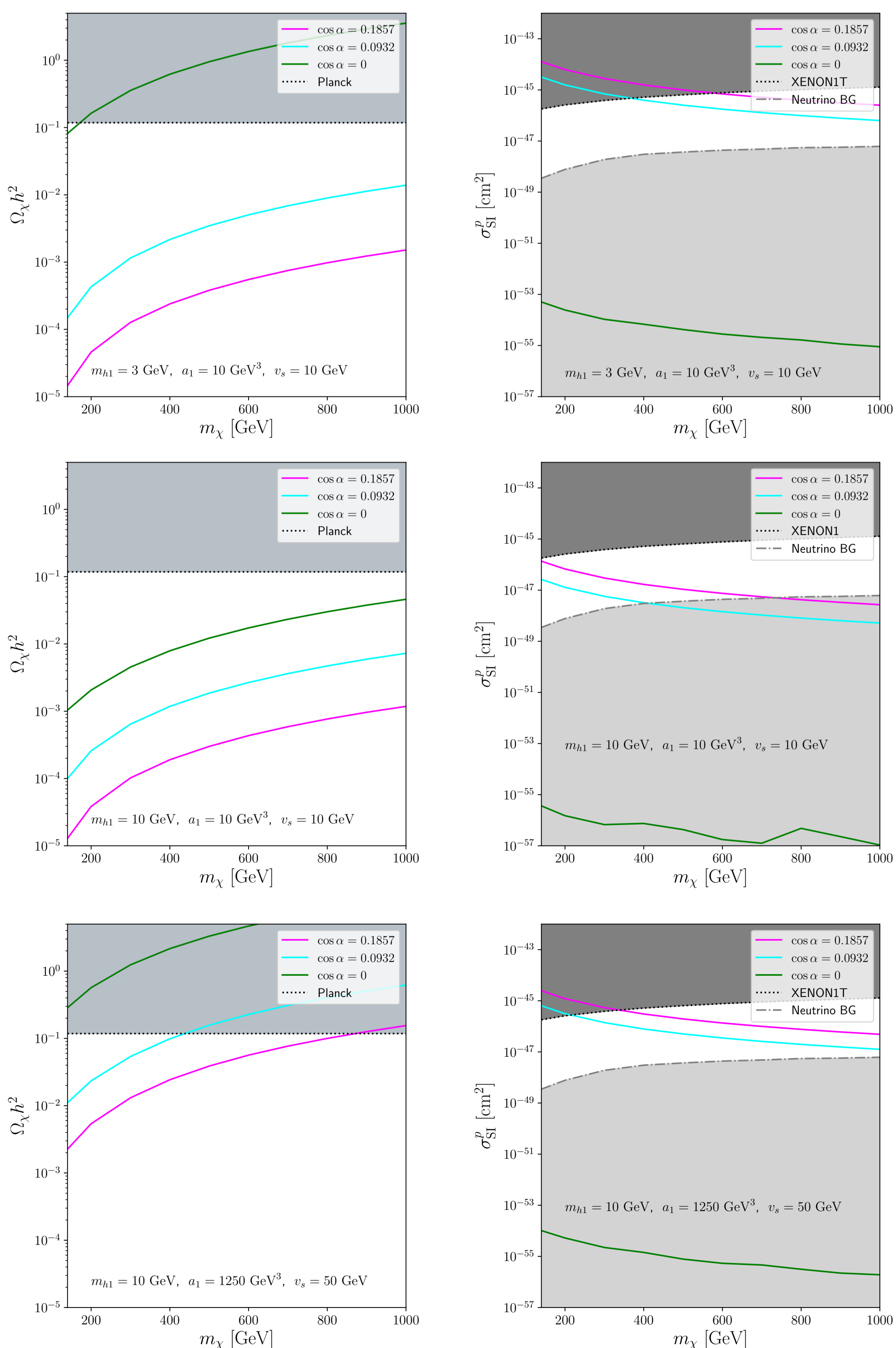
$$\hat{\mu} \equiv \frac{[\sigma(\text{XX} \rightarrow h)\text{Br}(h \rightarrow \text{YY})]_{\text{CxSM}}}{[\sigma(\text{XX} \rightarrow h)\text{Br}(h \rightarrow \text{YY})]_{\text{SM}}}$$

$$= \sin^2 \alpha \cdot \frac{1}{1 + \frac{1}{\sin^2 \alpha} \frac{\Gamma(h_2 \rightarrow h_1 h_1)}{\Gamma_{\text{SM}}^{\text{tot}}}} \rightarrow \frac{1 - \chi^2}{\sin^2 \alpha} = \frac{1 - \chi^2}{1 - \cos^2 \alpha}$$



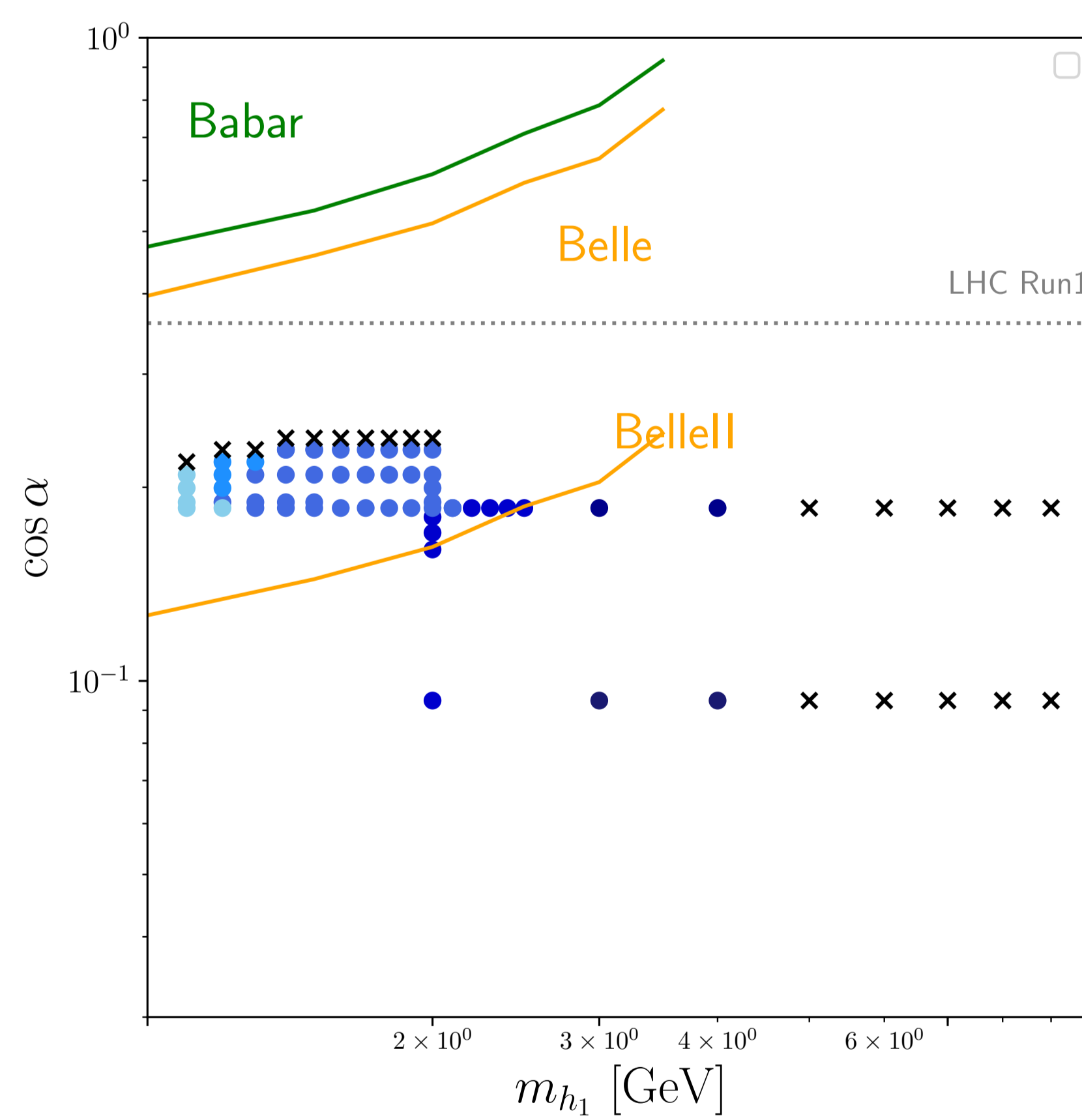
Ref: Marcela Carena et al. (arXiv : 1801.00794)

① allowed region of the relic density @Planck and the direct detection @XENON1T



Ref : XENON1T(2018)
Ref : PLANCK 2018 (arXiv : 1807.06209)

② comparison with Model-independent constraints

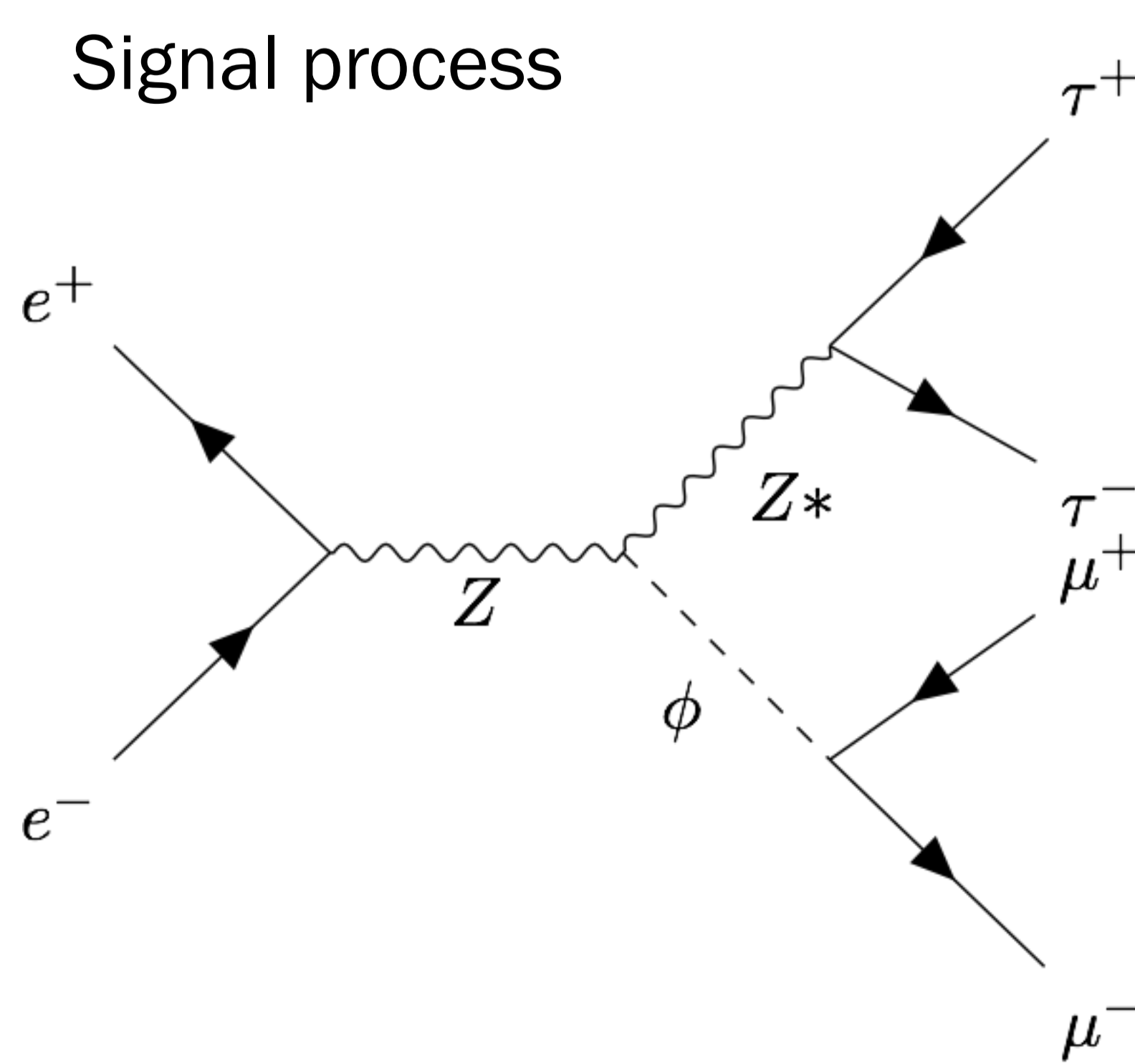


fix : $a_1 = 10 \text{ GeV}^3, v_s = 10 \text{ GeV}$

DMmass : 400GeV - 3TeV

Ref : The Belle II Physics Book (arXiv : 1808.10567)

③ calculation (event generation by MadGraph5)



Background process

$$e^+ e^- \rightarrow \mu^+ \mu^- \tau^+ \tau^-$$

Background cross-section :

$$5.833 \times 10^{-3} \pm 2.4 \times 10^{-5} \text{ pb}$$

set electron beams $\sqrt{s} = 10 \text{ GeV}$

set default cuts $p_t^j > 1 \text{ GeV}$
 $p_t^l > 1 \text{ GeV}$