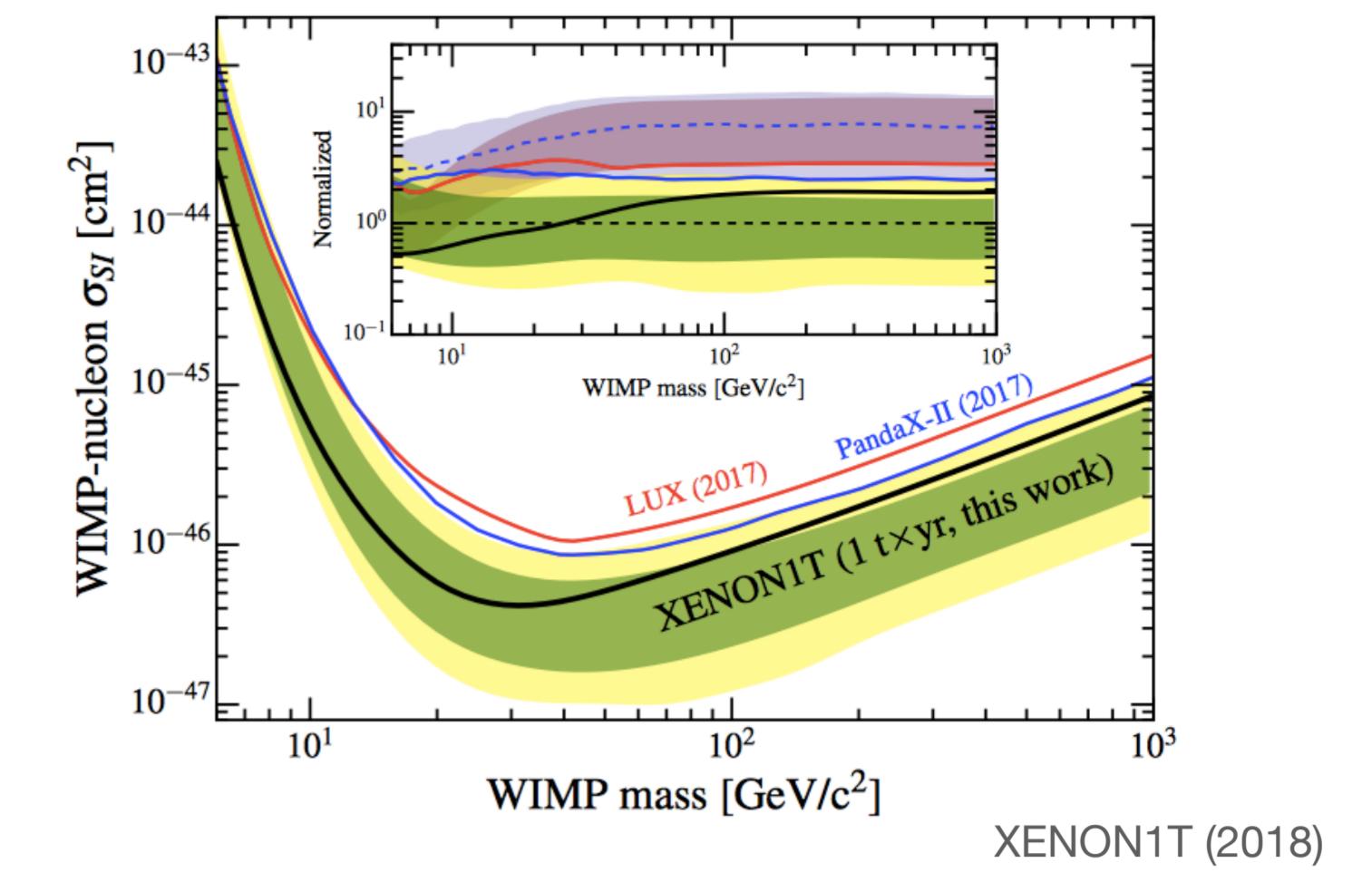


# Suppression of dark matter-nucleon scattering in singlet Higgs extension of the SM and its signature at collider experiments

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## §0 Strong constraint on the DM models

Dark matter (DM) plays an important role in the galaxy formation and evolution / have not been found directly yet. One of the most popular DM candidates is a Weakly Interacting Massive Particle (WIMP). The DM direct-detection experiments give a severe upper bound on the DM-nucleon scattering cross section.



## §1 Complex singlet extension of the SM (CxSM) w/ softly broken global U(1)

$$V = \frac{m^2}{2}|H|^2 + \frac{\lambda}{4}|H|^4 + \frac{\delta_2}{2}|H|^2|S|^2 + \frac{b_2}{2}|S|^2 + \frac{d_2}{4}|S|^4 + \left(a_1 S + \frac{b_1}{4}S^2 + \text{c.c.}\right)$$

$$H = \begin{pmatrix} 0 \\ v+h \\ \sqrt{2} \end{pmatrix} \quad S = (v_S + s + i\chi)/\sqrt{2}$$

assumption : CP conservation

ATLAS, EPJC75 (2016) 476  
CMS, PRD92, 012004 (2015)

CP-even scalars;  $h, s \rightarrow h_1, h_2$  (mass eigenstates)

$$\begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} h \\ s \end{pmatrix}$$

$$M^2 = \begin{pmatrix} \frac{\lambda v^2}{2} & \frac{\delta_2 v v_S}{2} \\ \frac{\delta_2 v v_S}{2} & \Lambda^2 \end{pmatrix}$$

$$\tan 2\alpha = 2 \frac{\frac{\delta_2}{2} v v_S}{\frac{\lambda}{2} v^2 - \Lambda^2}$$

$$m_{h_1}^2 = \frac{1}{2} \left( \frac{\lambda}{2} v^2 + \Lambda^2 \mp \sqrt{\left( \frac{\lambda}{2} v^2 - \Lambda^2 \right)^2 + 4 \left( \frac{\delta_2}{2} v v_S \right)^2} \right)$$

$$m_{h_1} = m_{h_{SM}} = 125 \text{ GeV}$$

$$\Lambda^2 \equiv \frac{d_2}{2} v_S^2 - \sqrt{2} \frac{a_1}{2 v_S}$$

CP-odd scalar;  $\chi$  (Goldstone boson) = DM

$$m_\chi^2 = -b_1 - \sqrt{2} \frac{a_1}{v_S} \quad (\text{massive})$$

## §2 Direct detection

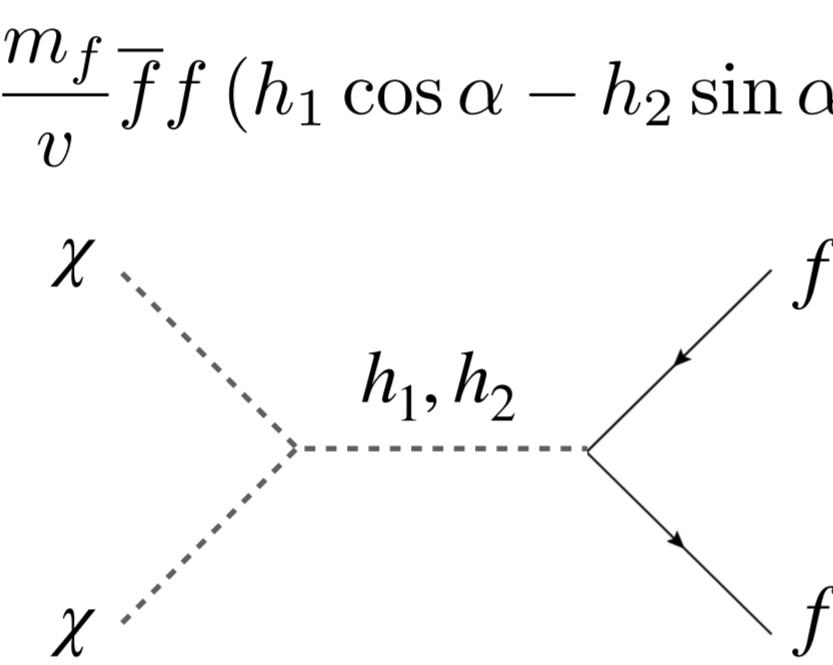
scalar trilinear interactions

$$\mathcal{L}_S = g_{h_1 \chi \chi} h_1 \chi^2 + g_{h_2 \chi \chi} h_2 \chi^2$$

$$g_{h_1 \chi \chi} \equiv -\frac{m_{h_1}^2 + \frac{\sqrt{2}a_1}{v_S}}{2v_S} \sin \alpha$$

$$g_{h_2 \chi \chi} \equiv -\frac{m_{h_2}^2 + \frac{\sqrt{2}a_1}{v_S}}{2v_S} \cos \alpha$$

yukawa interactions



$$\frac{\sqrt{2}a_1}{v_S} \left( \frac{1}{m_{h_1}^2} - \frac{1}{m_{h_2}^2} \right) \simeq 0 \iff a_1 \rightarrow 0 \text{ or } m_{h_1} \simeq m_{h_2}$$

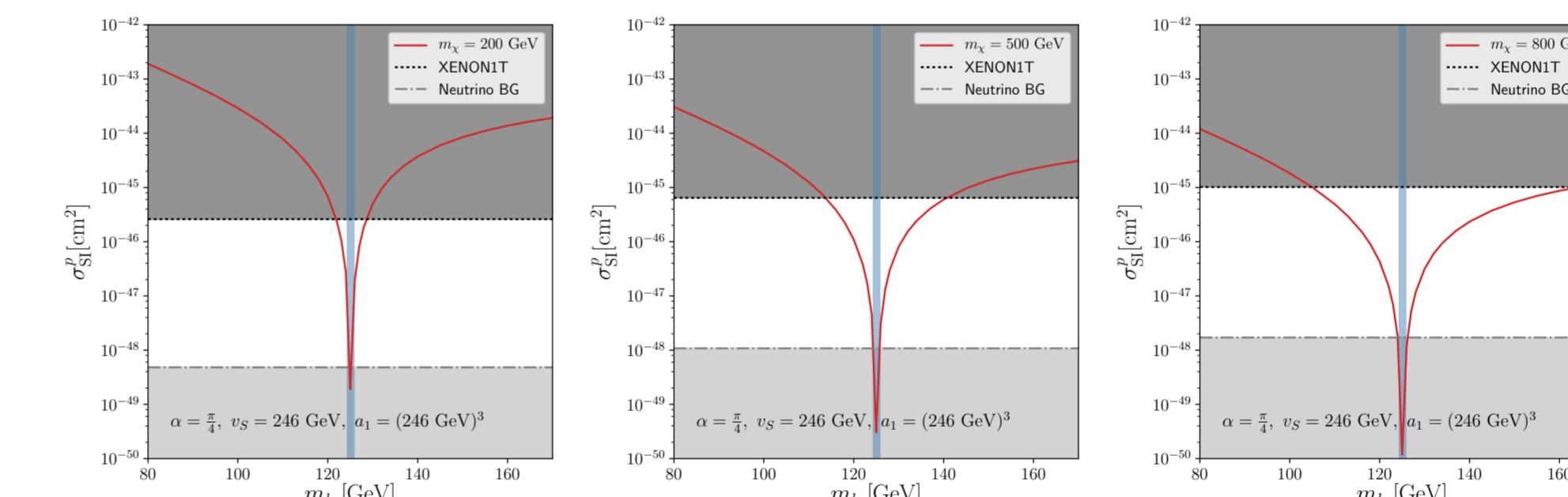
$$\begin{cases} a_1 \rightarrow 0 & \text{Z2 symmetry : } S \rightarrow -S \text{ (domain wall problem)} \\ m_{h_1} \simeq m_{h_2} & 2 \text{ degenerate 125 GeV Higgs bosons} \end{cases}$$

assumption : degenerate within the uncertainty of the measurement

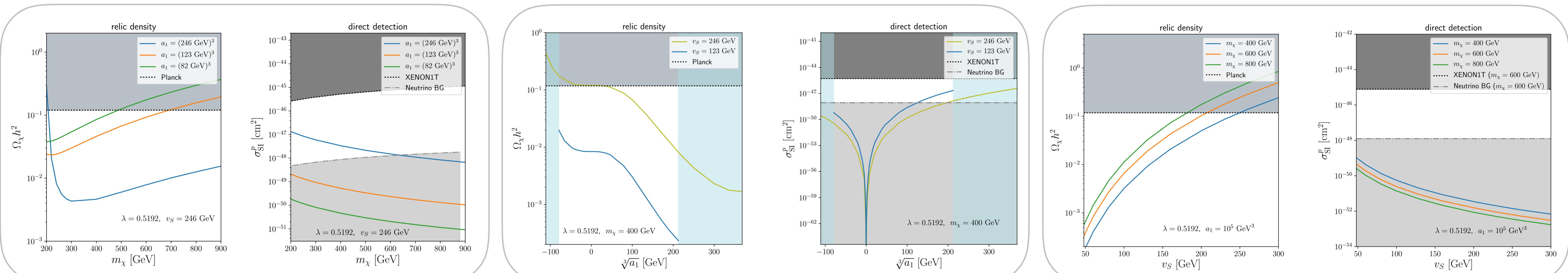
$$m_H = 125.09 \pm 0.24 \text{ GeV}$$

ATLAS (Run1, 2) + CMS(Run1)

$$\rightarrow \Delta m = |m_{h_1} - m_{h_2}| \leq 0.72 \text{ GeV} \quad (\text{blue bands in the graphs below})$$



## §3 Parameter analysis



: excluded by Planck observation PLANCK 2018, arXiv 1807.06209

: excluded by XENON1T experiment

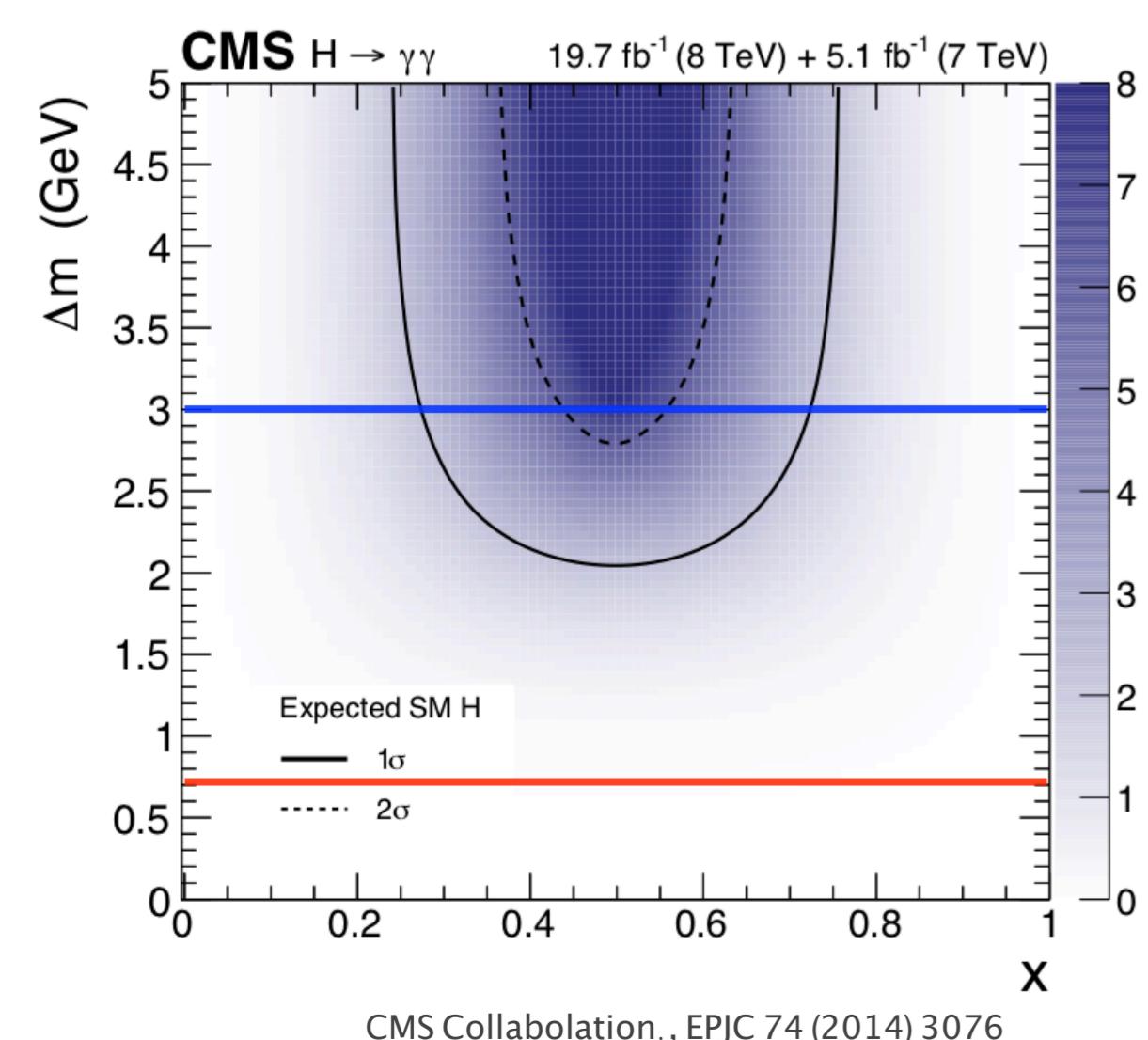
: Neutrino BG J.Billard et al., PRD 89, 023524 (2014)

Other constraints : perturbativity / stability of the scalar potential

$$\lambda < \frac{16\pi}{3}, \quad d_2 < \frac{16\pi}{3}, \quad \lambda \left( d_2 + \frac{2\sqrt{2}a_1}{v_S^3} \right) > d_2^2$$

## §4 Collider search for degenerate Higgs bosons

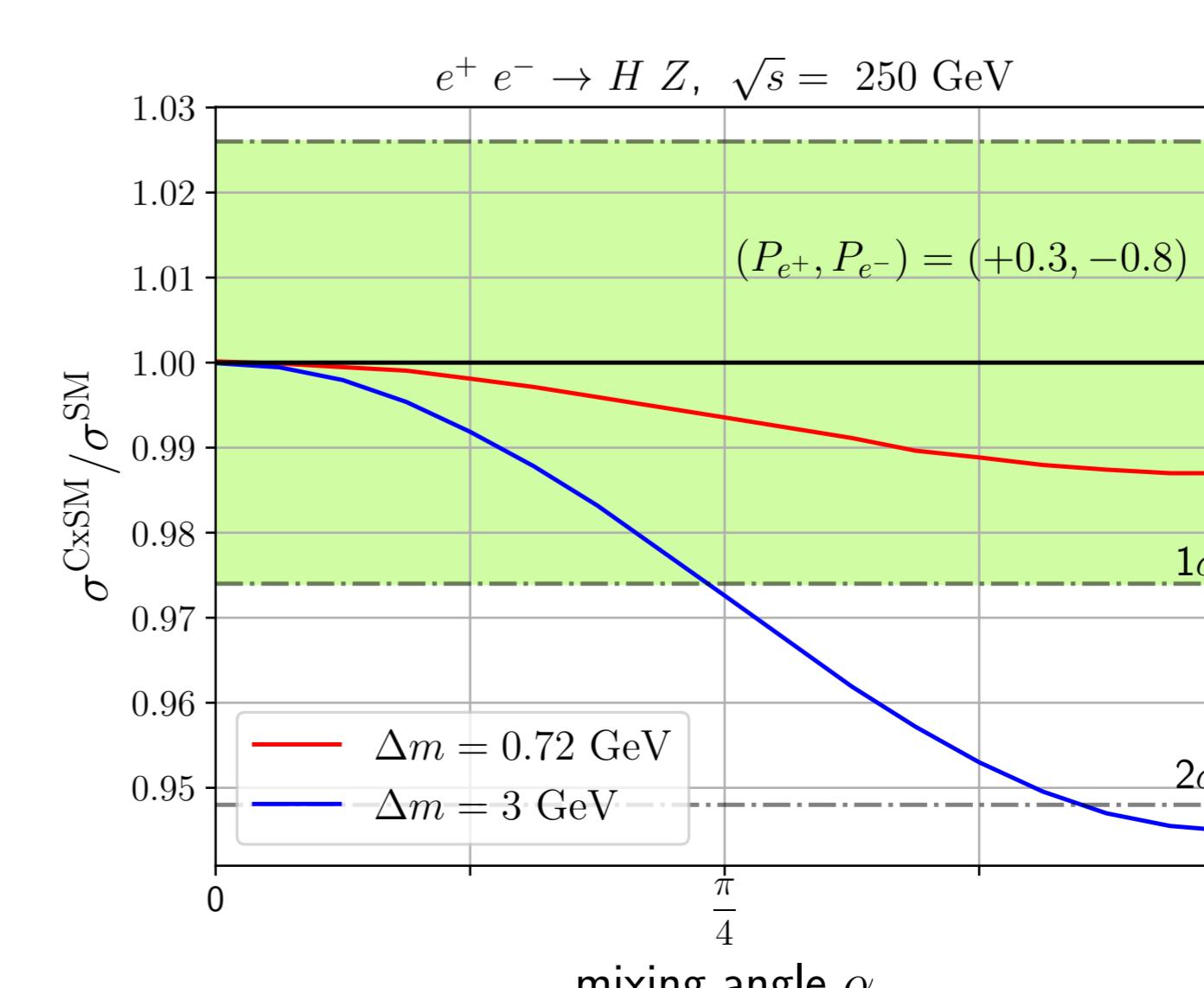
compare cross sections of  $e^+ e^- \rightarrow H Z (\sqrt{s} = 250 \text{ GeV})$  in the SM / CxSM



It is difficult to distinguish two degenerate ( $\Delta m \lesssim 3 \text{ GeV}$ )

Higgs produced at the Large Hadron Collider (LHC).

→ How about at the International Linear Collider (ILC)?



( $h_1 \rightarrow h_{SM}$  @ $\alpha \rightarrow 0$ )

precision of cross section :

$$\frac{\Delta \sigma}{\sigma} = 2.6\%$$

D.M. Asner et al., arXiv 1310.0763

→ It is hard to distinguish two models at the ILC.