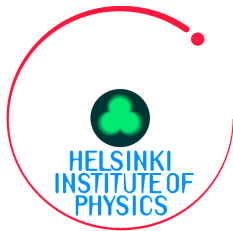


# Cosmological constraints on light flavons

Venus Keus

University of Helsinki & Helsinki Institute of Physics



In collaboration with  
M. Heikinheimo & K. Huitu & N. Koivunen  
[arXiv:1812.10963](https://arxiv.org/abs/1812.10963) [hep-ph] to appear in JHEP

HPNP 2019

- 1 The Froggatt-Nielsen mechanism
- 2 Light flavons as Dark Matter candidates
- 3 Cosmological constraints on the abundance of light flavons
- 4 Summary

# The SM and beyond

## SM predictions confirmed:

- the W & Z (1983)
- the top quark (1995)
- the tau neutrino (2000)
- “a” Higgs boson (2012)

## What is missing:

- Fermion mass hierarchy explanation
- Dark Matter candidate

⇒ Non-minimal Higgs frameworks



# The Froggatt-Nielsen mechanism

Yukawa interactions as effective operators:

$$\mathcal{L}_Y \supset c_{ij} \left( \frac{\Phi}{\Lambda} \right)^{n_{ij}} \bar{f}_{L,i} f_{R,j} H \quad \text{where} \quad H = \begin{pmatrix} 0 \\ \frac{v_h + h}{\sqrt{2}} \end{pmatrix}, \quad \Phi = \left( \frac{v_\phi + \phi}{\sqrt{2}} \right)$$

U(1) charges:  $q_{\bar{L},i} + q_{R,j} + q_h + q_\phi n_{ij} = 0$

$$\begin{cases} q_\phi = -1 \\ q_h = 0 \end{cases} \Rightarrow Y_{ij} = c_{ij} \left( \frac{v_\phi}{\sqrt{2}\Lambda} \right)^{(q_{\bar{L},i} + q_{R,j})} \equiv c_{ij} \epsilon^{n_{ij}}$$

Leading to flavour violating interactions

$$\mathcal{L}_Y \supset \frac{v_h}{\sqrt{2}} Y_{ij} \left( 1 + \frac{h}{v_h} + n_{ij} \frac{\phi}{v_\phi} \right) \bar{f}_{L,i} f_{R,j}$$

$$\tilde{\kappa}_{ij} = \frac{1}{v_\phi} \frac{v_h}{\sqrt{2}} (Y_{ij} \cdot n_{ij})$$

# A leptophilic flavon

0.511 MeV	105.7 MeV	1.777 GeV
$-1$	$-1$	$-1$
$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$
<b>e</b>	<b>μ</b>	<b>τ</b>
electron	muon	tau

$$Y \sim \begin{pmatrix} \epsilon^6 & \epsilon^5 & \epsilon^4 \\ \epsilon^5 & \epsilon^4 & \epsilon^3 \\ \epsilon^4 & \epsilon^3 & \epsilon^2 \end{pmatrix} \quad \text{and} \quad \tilde{\kappa} \sim \frac{1}{v_\phi} \begin{pmatrix} m_e & m_\mu \epsilon & m_\tau \epsilon^2 \\ m_\mu \epsilon & m_\mu & m_\tau \epsilon \\ m_\tau \epsilon^2 & m_\tau \epsilon & m_\tau \end{pmatrix}$$

$\bar{e}_L$	$e_R$	$\bar{\mu}_L$	$\mu_R$	$\bar{\tau}_L$	$\tau_R$	$H$	$\Phi$
3	3	2	2	1	1	0	-1

# The softly broken scalar potential

$$V = -\frac{\mu_h^2}{2}(H^\dagger H) + \frac{\lambda_h}{2}(H^\dagger H)^2 - \frac{\mu_\phi^2}{2}(\Phi^\dagger \Phi) + \frac{\lambda_\phi}{2}(\Phi^\dagger \Phi)^2 + \lambda_{h\phi}(H^\dagger H)(\Phi^\dagger \Phi) - \frac{\mu_\phi'^2}{4}(\Phi^2 + \Phi^{\dagger 2}).$$

$$H = \begin{pmatrix} 0 \\ \frac{v_h + h}{\sqrt{2}} \end{pmatrix}, \quad \Phi = \begin{pmatrix} v_\phi + \phi \\ \sqrt{2} \end{pmatrix}, \quad \text{with } \phi = \sigma + i\eta$$

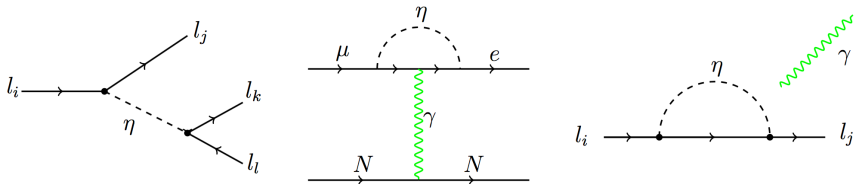
The physical states  $H_1, H_2$

$$\begin{pmatrix} H_1 \\ H_2 \end{pmatrix} \equiv \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} h \\ \sigma \end{pmatrix}$$

with  $m_{H_2} \approx m_{H_1} = 125$  GeV and a light  $\eta$

CLFV processes mediated by light  $\eta$ 

$$m_{H_1} = 125\text{GeV}, \quad m_{H_2} = 500\text{GeV}, \quad m_\eta < 2m_e$$

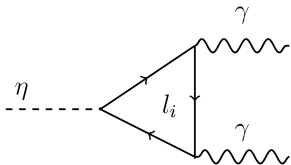


$$BR(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13} \Rightarrow v_\phi > 14.4 \text{ TeV}$$

The MEG collaboration, *Hyperfine Interact.* 239 (2018) no.1, 58

Lifetime of  $\eta$  ( $m_\eta < 2m_e$ )

$$\tau_\eta = \frac{1}{\Gamma_{\eta \rightarrow \gamma\gamma}} = \frac{32\pi m_\eta}{|\mathcal{M}_{\eta \rightarrow \gamma\gamma}|^2} > 10^{17} \text{ s} \quad \text{age of the universe}$$



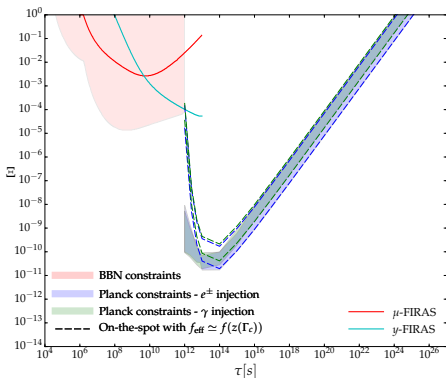
$$|\mathcal{M}_{\eta \rightarrow \gamma\gamma}|^2 = \frac{e^4}{32\pi^4} m_\eta^4 \left( \frac{\tilde{\kappa}_{ee}}{m_e} + \frac{\tilde{\kappa}_{\mu\mu}}{m_\mu} + \frac{\tilde{\kappa}_{\tau\tau}}{m_\tau} \right)^2$$

- for  $m_\eta = 10 \text{ KeV} \Rightarrow v_\phi \sim 10^6 \text{ GeV}$
- for  $m_\eta = 1 \text{ MeV} \Rightarrow v_\phi \sim 10^{13} \text{ GeV}$



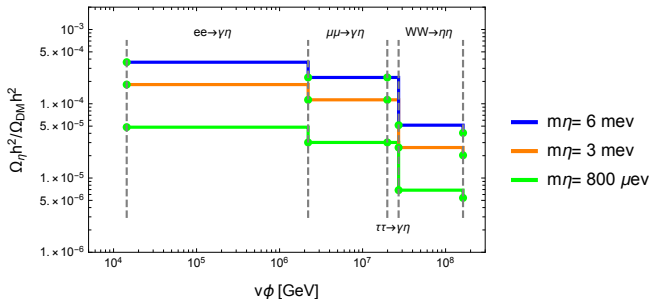
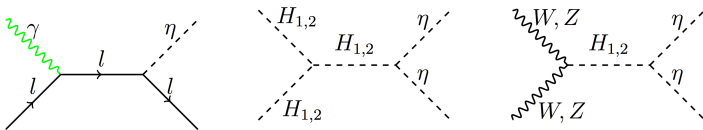
# CMB constraints on long-lived decaying $\eta$

$$\text{for } \tau_\eta \sim 10^{17} \text{ s} \Rightarrow \Omega_\eta h^2 \lesssim 10^{-8} \Omega_{DM} h^2$$



$$\text{For } \Omega_\eta h^2 = \Omega_{DM} h^2 \Rightarrow \begin{cases} m_\eta = 10 \text{ KeV} \rightarrow v_\phi \sim 10^{10} \text{ GeV} \\ m_\eta = 1 \text{ MeV} \rightarrow v_\phi \sim 10^{17} \text{ GeV} \end{cases}$$

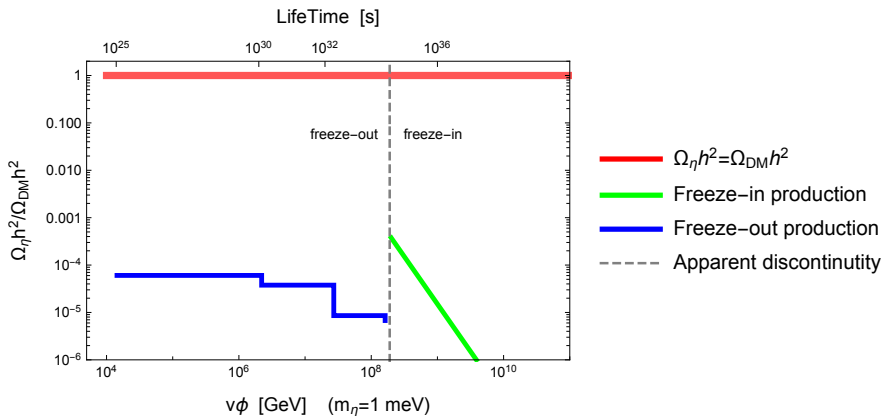
# Production of $\eta$ in the early universe

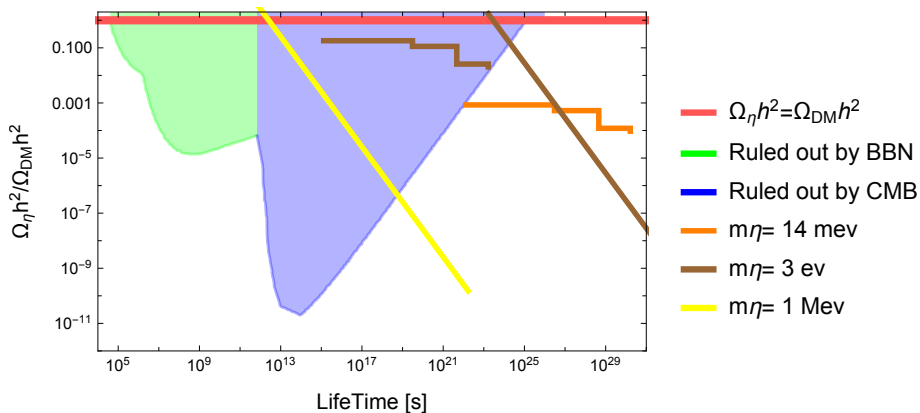


for  $v_\phi \sim \mathcal{O}(10^4 - 10^8) \text{ GeV} \Rightarrow \Gamma_{\alpha X \rightarrow \eta Y}(T_{fo}) = H(T_{fo})$  where  $T_{fo} \sim m_\alpha$

# Production of $\eta$ in the early universe

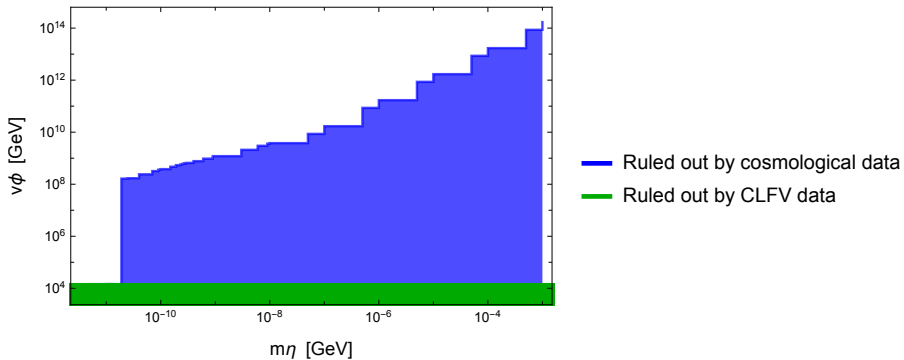
for  $v_\phi \gtrsim 10^8 \text{ GeV} \Rightarrow \eta$  does not thermalise



Cosmological constraints on the abundance of  $\eta$ 

- $m_\eta \leq 14 \text{ meV}$ : no cosmological constraints
- $14 \text{ meV} < m_\eta < 3 \text{ eV}$ : hot relic, moderately constrained
- $3 \text{ eV} < m_\eta < 2m_e$ : severely constrained

# Cosmological constraints on the abundance of $\eta$



# Summary

- Froggatt-Nielsen mechanism predicts a complex flavon responsible for producing the SM fermion Yukawa couplings.
- The real part of flavon mixes with the higgs and leads to CLFV processes.
- The abundance of the imaginary part of flavon is constrained by BBN, CMB and small scale structure observations.

# BACKUP SLIDES

# Yukawa interactions

$$\mathcal{L}_{\text{eff}} \supset \frac{v_h}{\sqrt{2}} Y_{ij} \left( 1 + \frac{h}{v_h} + n_{ij} \frac{\phi}{v_\phi} \right) \bar{l}'_{L,i} l'_{R,j}$$

Diagonalising the Yukawa matrix

$$Y_{\text{diag}} = U_L Y_{ij} U_R^\dagger$$

In the lepton mass basis

$$\mathcal{L}_{\text{eff}} \supset \bar{l}_L M_{\text{diag}} l_R + \frac{h}{\sqrt{2}} \bar{l}_L Y_{\text{diag}} l_R + \frac{v_h}{v_\phi} \frac{\phi}{\sqrt{2}} \bar{l}_{L,i} \kappa_{ij} l_{R,j} + h.c.$$

CLFV processes with coupling

$$\tilde{\kappa}_{ij} = \frac{v_h \kappa_{ij}}{\sqrt{2} v_\phi} = \left[ y_j \sum_{k=1}^3 q_{\bar{L},k} (U_L)_{ik} (U_L)_{jk}^* + y_i \sum_{k=1}^3 q_{R,k} (U_R)_{ik} (U_R)_{jk}^* \right]$$



# The scalar potential

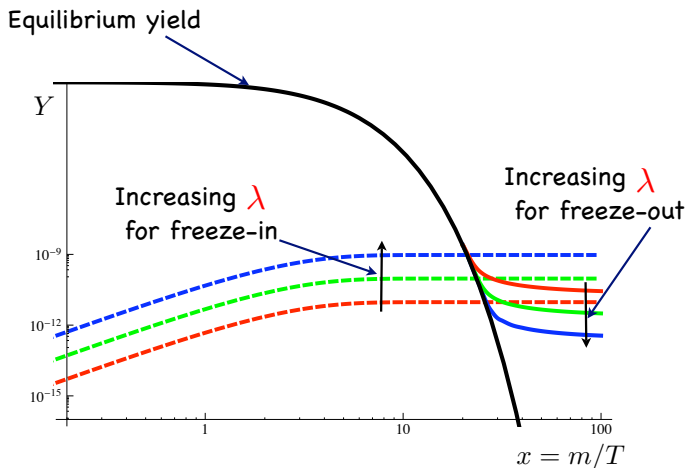
$$V = -\frac{\mu_h^2}{2}(H^\dagger H) + \frac{\lambda_h}{2}(H^\dagger H)^2 - \frac{\mu_\phi^2}{2}(\Phi^\dagger \Phi) + \frac{\lambda_\phi}{2}(\Phi^\dagger \Phi)^2 + \lambda_{h\phi}(H^\dagger H)(\Phi^\dagger \Phi) - \frac{\mu_\phi'^2}{4}(\Phi^2 + \Phi^{\dagger 2}).$$

$$H = \begin{pmatrix} 0 \\ \frac{v_h + h}{\sqrt{2}} \end{pmatrix}, \quad \Phi = \begin{pmatrix} v_\phi + \phi \\ \sqrt{2} \end{pmatrix}, \quad \text{with } \phi = \sigma + i\eta$$

The physical states  $H_1$ ,  $H_2$  (mixture of  $h$  and  $\sigma$ ) and  $\eta$

$$\begin{aligned} \mathcal{L} \supset & \left[ \cos \theta \frac{Y_{ij}^{\text{diag}}}{\sqrt{2}} + \sin \theta \tilde{\kappa}_{ij} \right] \bar{l}_i P_R l_j H_1 \\ & + \left[ -\sin \theta \frac{Y_{ij}^{\text{diag}}}{\sqrt{2}} + \cos \theta \tilde{\kappa}_{ij} \right] \bar{l}_i P_R l_j H_2 \\ & + i\tilde{\kappa}_{ij} \bar{l}_i P_R l_j \eta + \text{h.c.} \end{aligned}$$

# Freeze-in vs. Freeze-out



(borrowed from S. West's talk in ULB in 2010)