

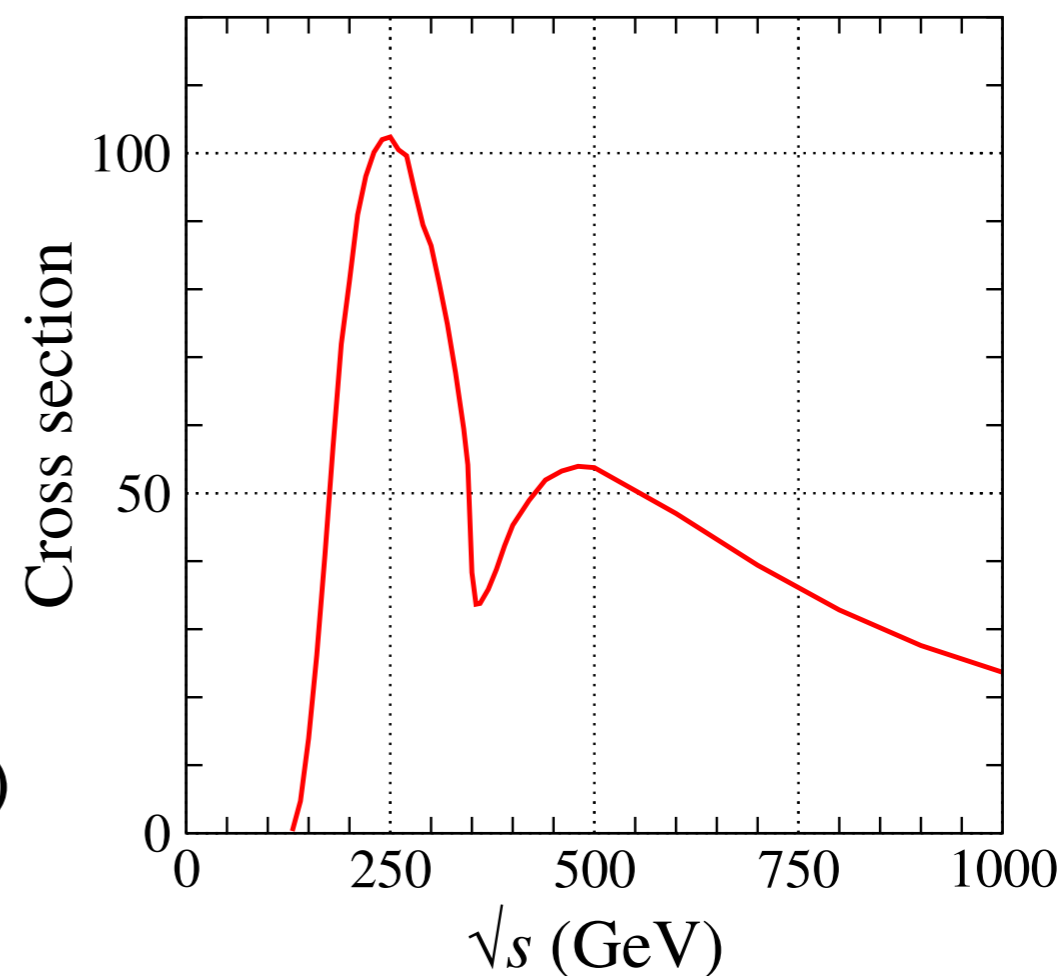
# Higgs production in association with a photon in extended Higgs models at ILC250



Kentarou Mawatari

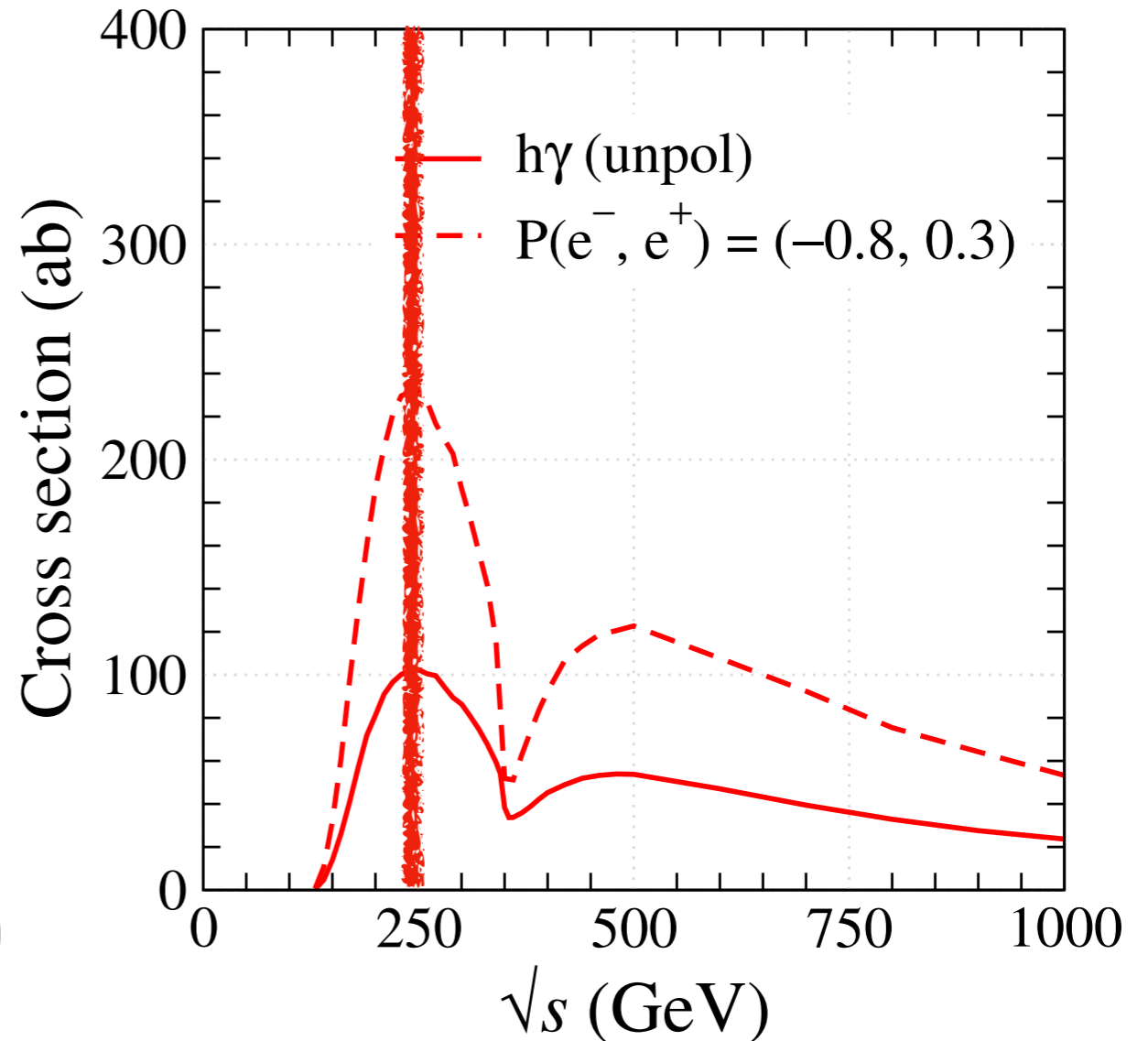
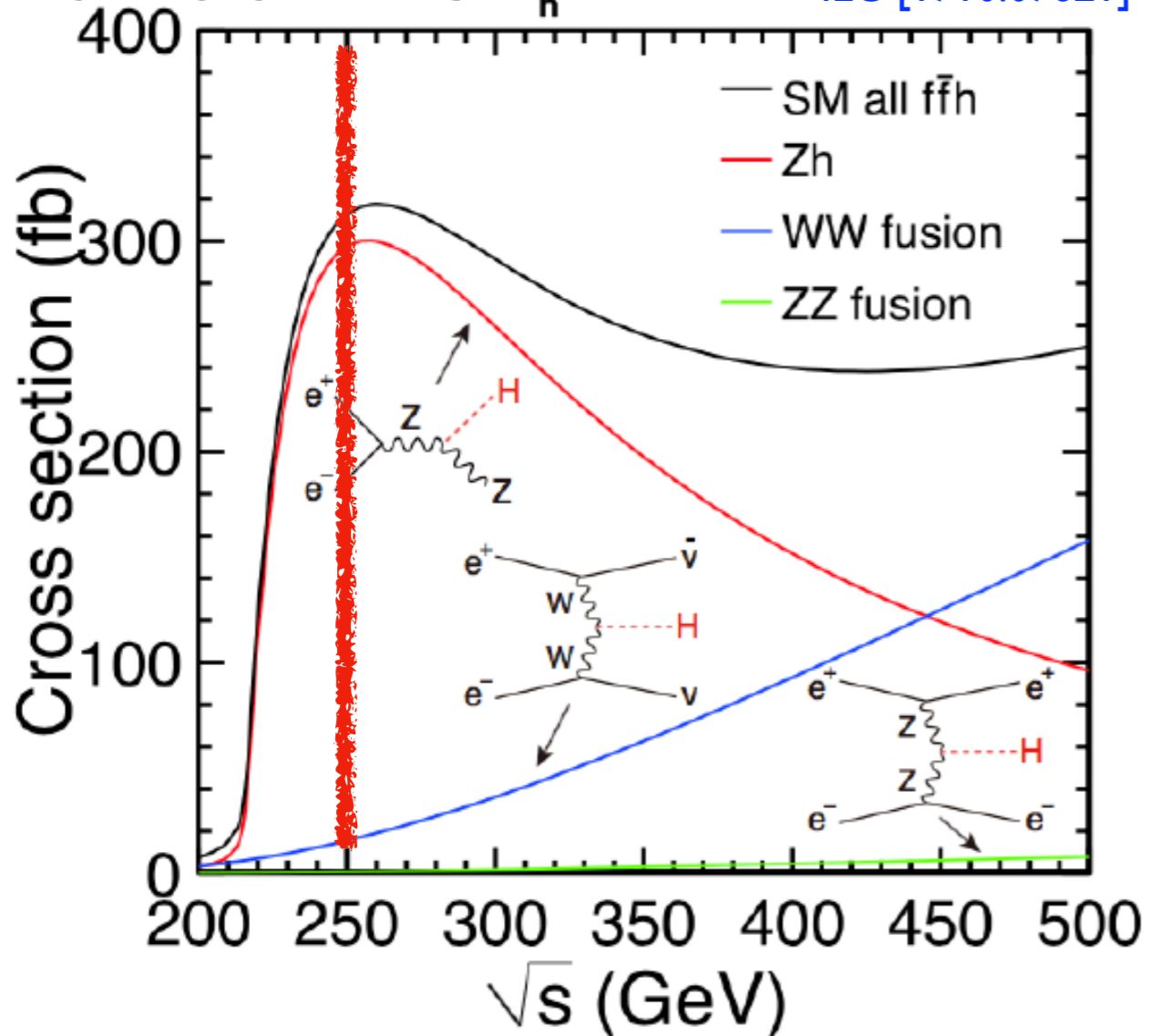
arXiv:1808.10268, PRD

with Prof. Shinya Kanemura (Osaka)  
and **Dr.** Kodai Sakurai (Toyama/Osaka → Karlsruhe)



# Higgs productions at ILC250

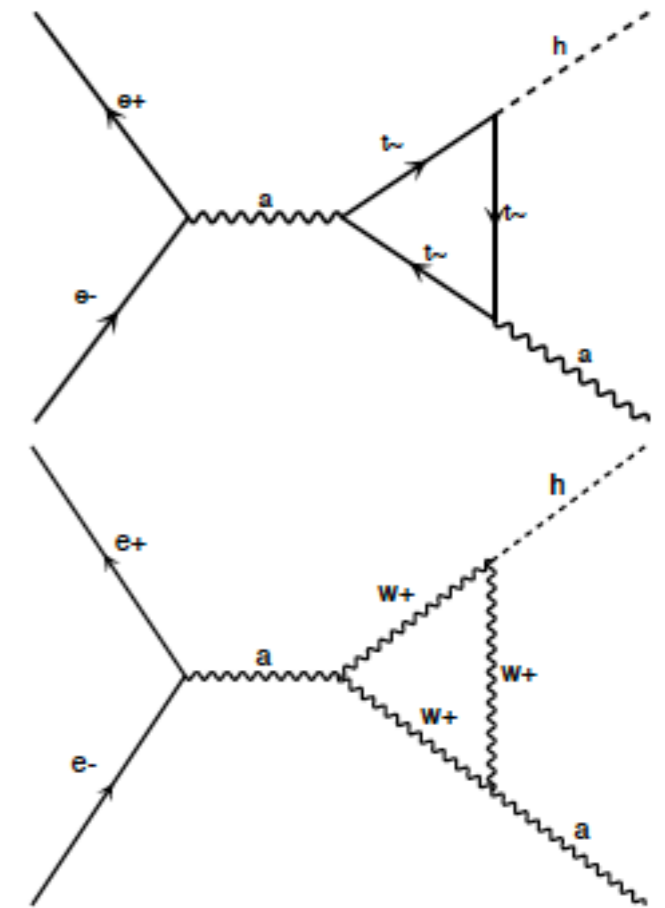
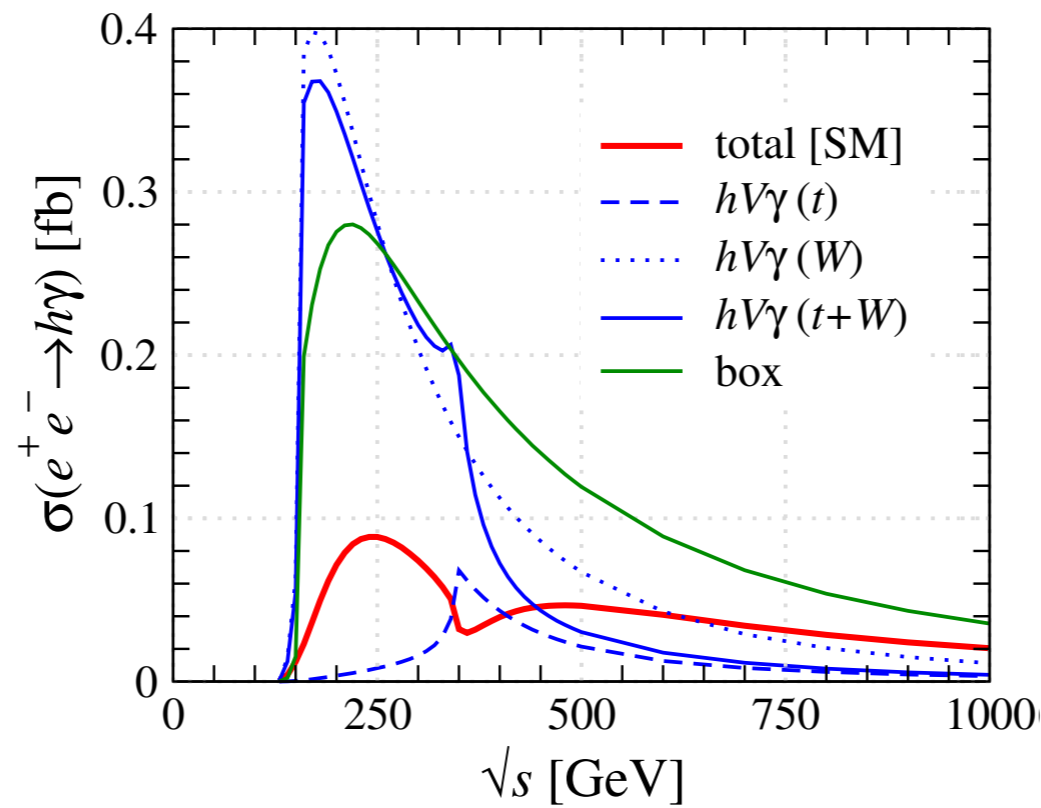
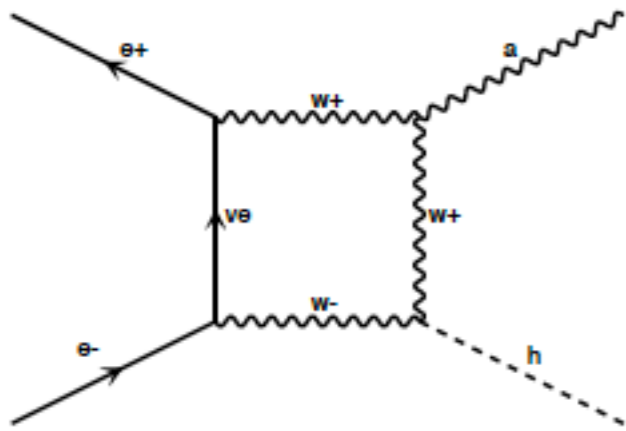
$P(e^-, e^+) = (-0.8, 0.3)$ ,  $M_h = 125 \text{ GeV}$  ILC [1710.07621]



This is not the only Higgs production which has a peak at 250 GeV at  $e^+e^-$  colliders.

- $\sigma(h\gamma) \sim \sigma(hZ) \times 10^{-3}$  because of loop-induced.
- Beam polarization enhances.
- A monochromatic photon.  $E_\gamma = \frac{\sqrt{s}}{2} \left(1 - \frac{m_h^2}{s}\right) \sim 93.8 \text{ GeV} @ \sqrt{s} = 250 \text{ GeV}$
- **Sensitive to New Physics!**

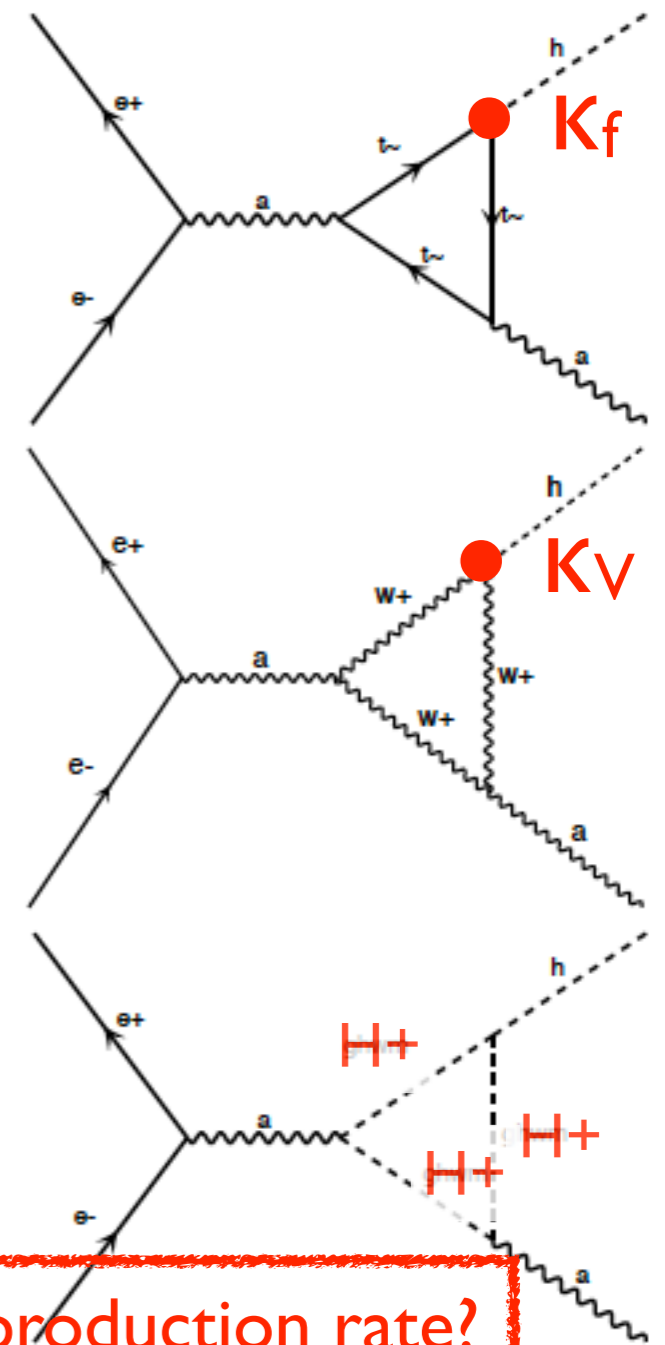
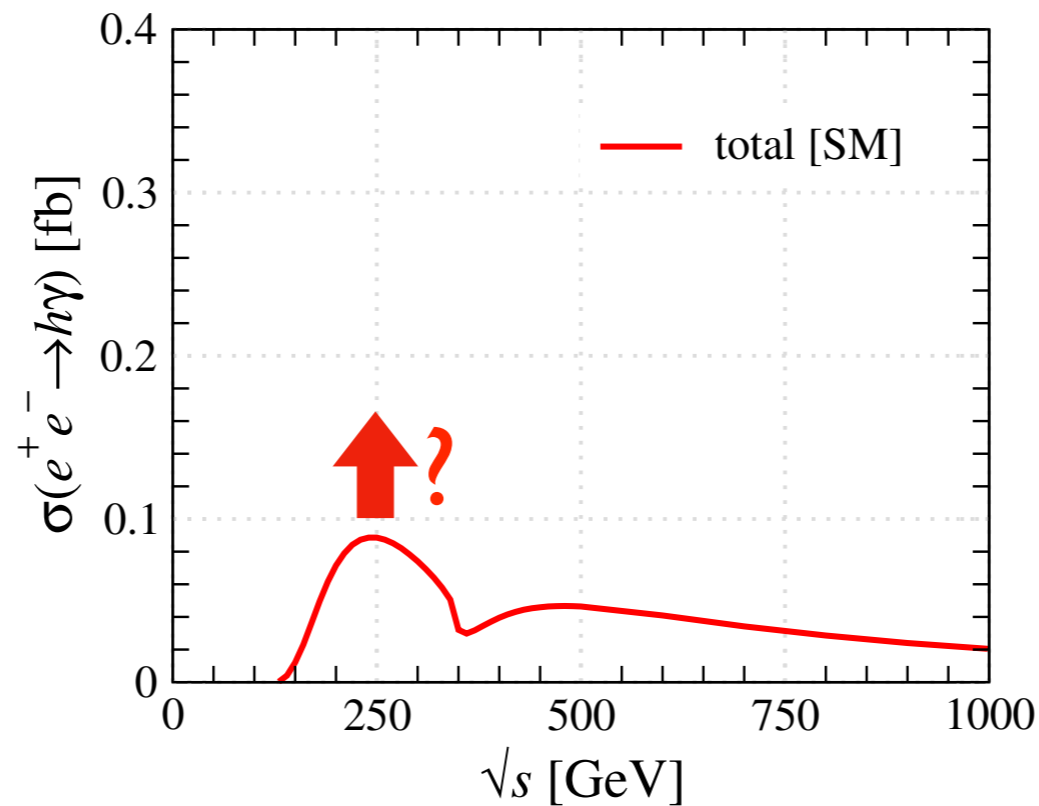
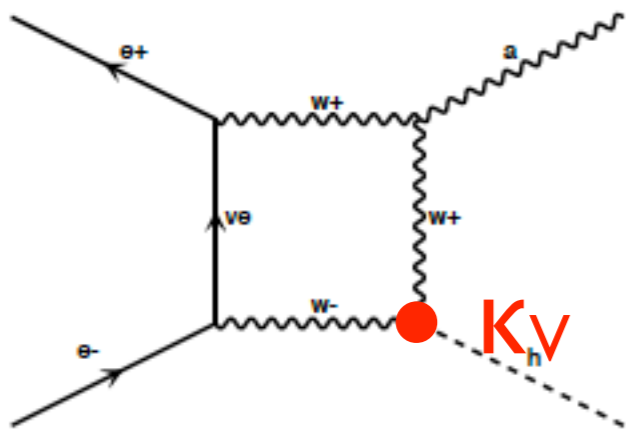
# $e^+e^- \rightarrow h\gamma$ in the SM



Unfortunate destructive interference among the different contributions...

$$\kappa_X = \frac{g_{hXX}^{\text{NP}}}{g_{hXX}^{\text{SM}}}$$

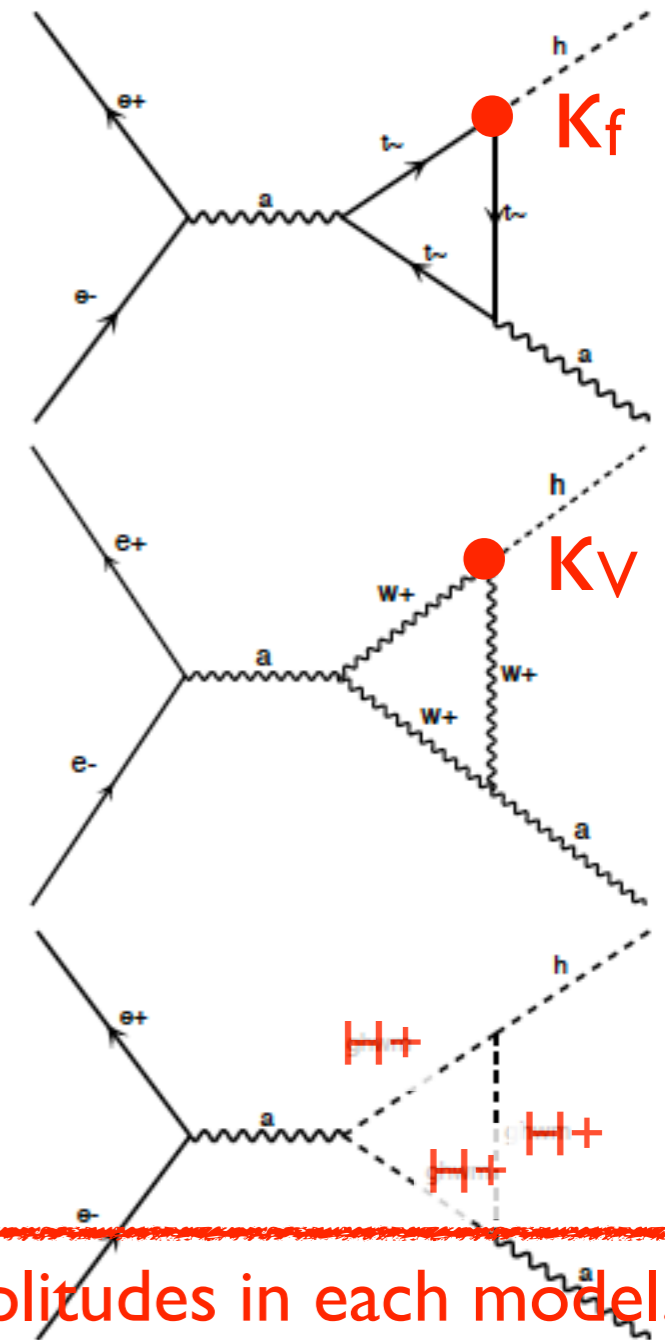
# $e^+e^- \rightarrow h\gamma$ in extended Higgs models



How much can new physics enhance (or reduce) the production rate?

# 3 benchmark extended Higgs models

	$K_f (=t)$	$K_V$	$H^+$	$H^{++}$
IDM (Inert doublet model)	1	1	○	×
ITM (Y=1) (Inert triplet model)	1	1	○	○
THDM (Two Higgs doublet model)	$S_{\beta-\alpha}$ $-C_{\beta-\alpha}/t_\beta$	$S_{\beta-\alpha}$	○	×



We employ the H-COUP program to compute the loop amplitudes in each model.

Kanemura, Kikuchi, Sakurai, Yagyu [1710.04603, CPC]

# H-COUP

\* See Kei Yagyu's talk and Kodai Sakurai's poster

$$\hat{\Gamma}_{hVV}^{\mu\nu}(p_1^2, p_2^2, q^2) = g^{\mu\nu} \hat{\Gamma}_{hVV}^1 + \frac{p_1^\nu p_2^\mu}{m_V^2} \hat{\Gamma}_{hVV}^2 + i\epsilon^{\mu\nu\rho\sigma} \frac{p_{1\rho} p_{2\sigma}}{m_V^2} \hat{\Gamma}_{hVV}^3$$

H-COUP is a calculation tool composed of a set of Fortran codes to compute the renormalized Higgs boson couplings with radiative corrections in various non-minimal Higgs models, such as the Higgs singlet model, four types of two Higgs doublet models and the inert doublet model. The involved on-shell renormalization scheme is adopted, where the gauge dependence is eliminated.

Authors: Shinya Kanemura, Mariko Kikuchi, Kodai Sakurai and Kei Yagyu **+K. Mawatari**

The manual for H-COUP version 1.0 can be taken on [arXiv:1710.04603 \[hep-ph\]](https://arxiv.org/abs/1710.04603).

Loop effects on the Higgs decay widths  
in extended Higgs models [1803.01456, PLB]

## Downloads

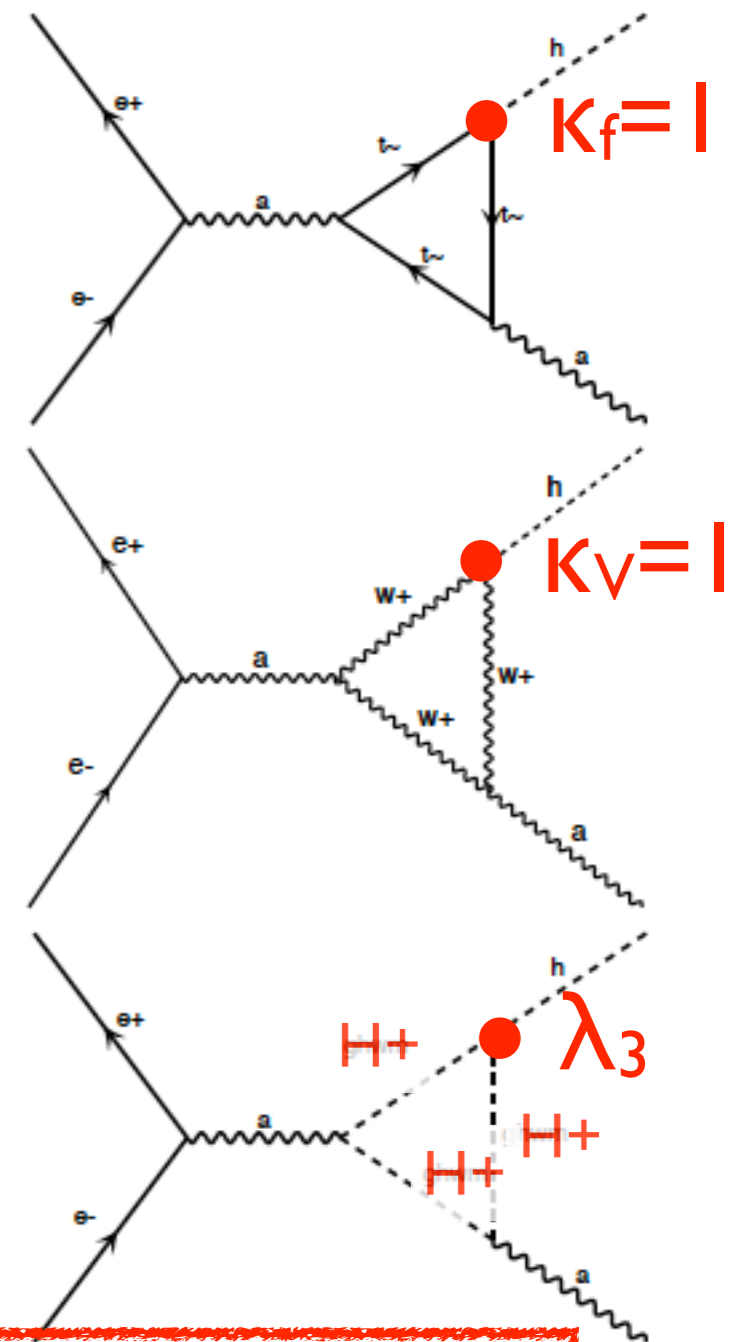
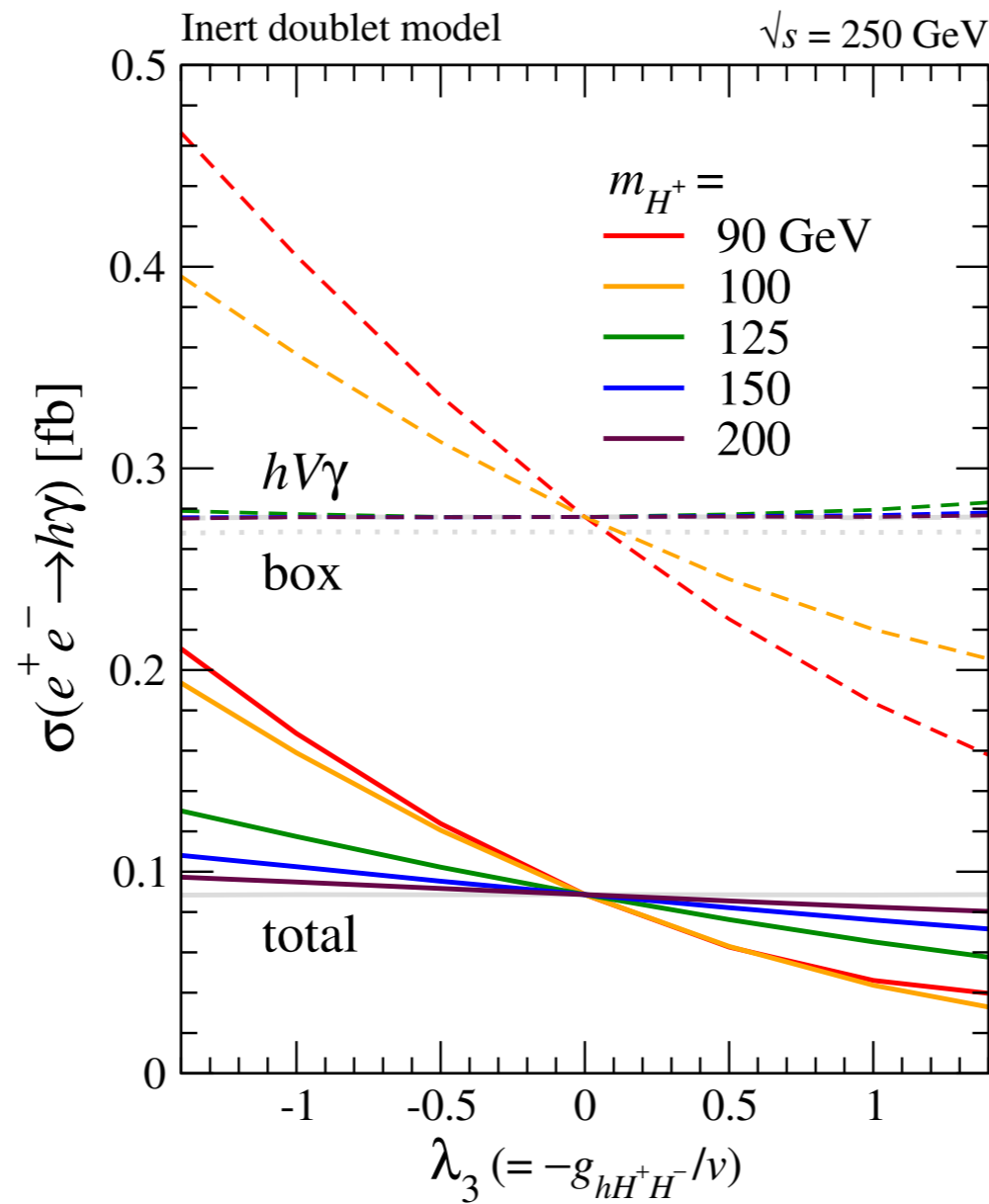
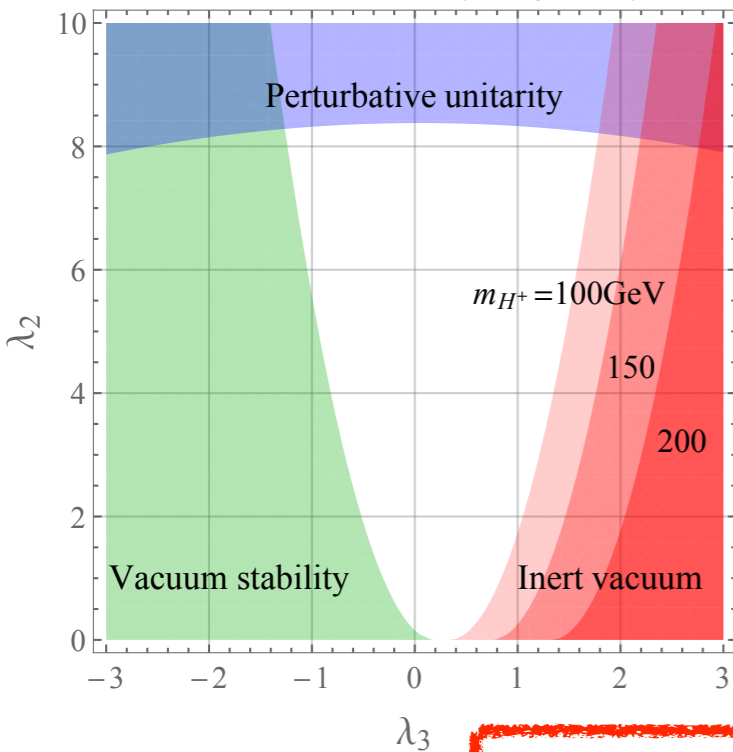
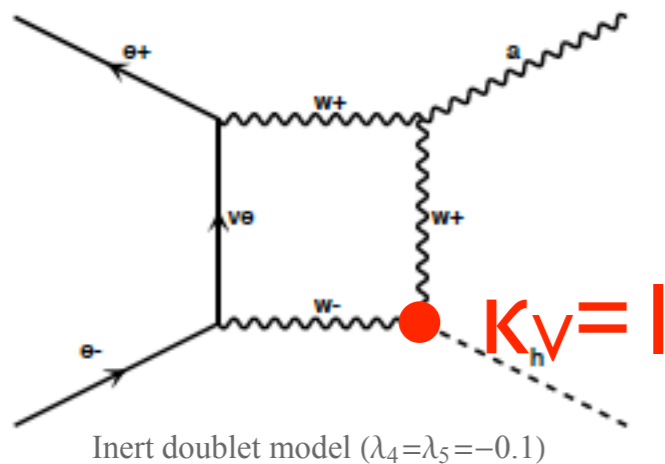
- H-COUP version 1.0 : [\[HCOUP-1.0.zip\]](#) [The manual is [here](#)]

\* An additional Higgs doublet with an exact  $Z_2$  symmetry

$$\begin{aligned}
 V = & \mu_1^2 |\Phi_1|^2 + \mu_2^2 |\Phi_2|^2 \\
 & + \frac{1}{2} \lambda_1 |\Phi_1|^4 + \frac{1}{2} \lambda_2 |\Phi_2|^4 \\
 & + \lambda_3 |\Phi_1|^2 |\Phi_2|^2 + \lambda_4 |\Phi_1^\dagger \Phi_2|^2 \\
 & + \frac{1}{2} \lambda_5 \{ (\Phi_1^\dagger \Phi_2)^2 + \text{h.c.} \}
 \end{aligned}$$

# $e^+e^- \rightarrow h\gamma$ in the IDM

Arhrib, Benbrik, Yuan  
[1401.6698, EPJC]

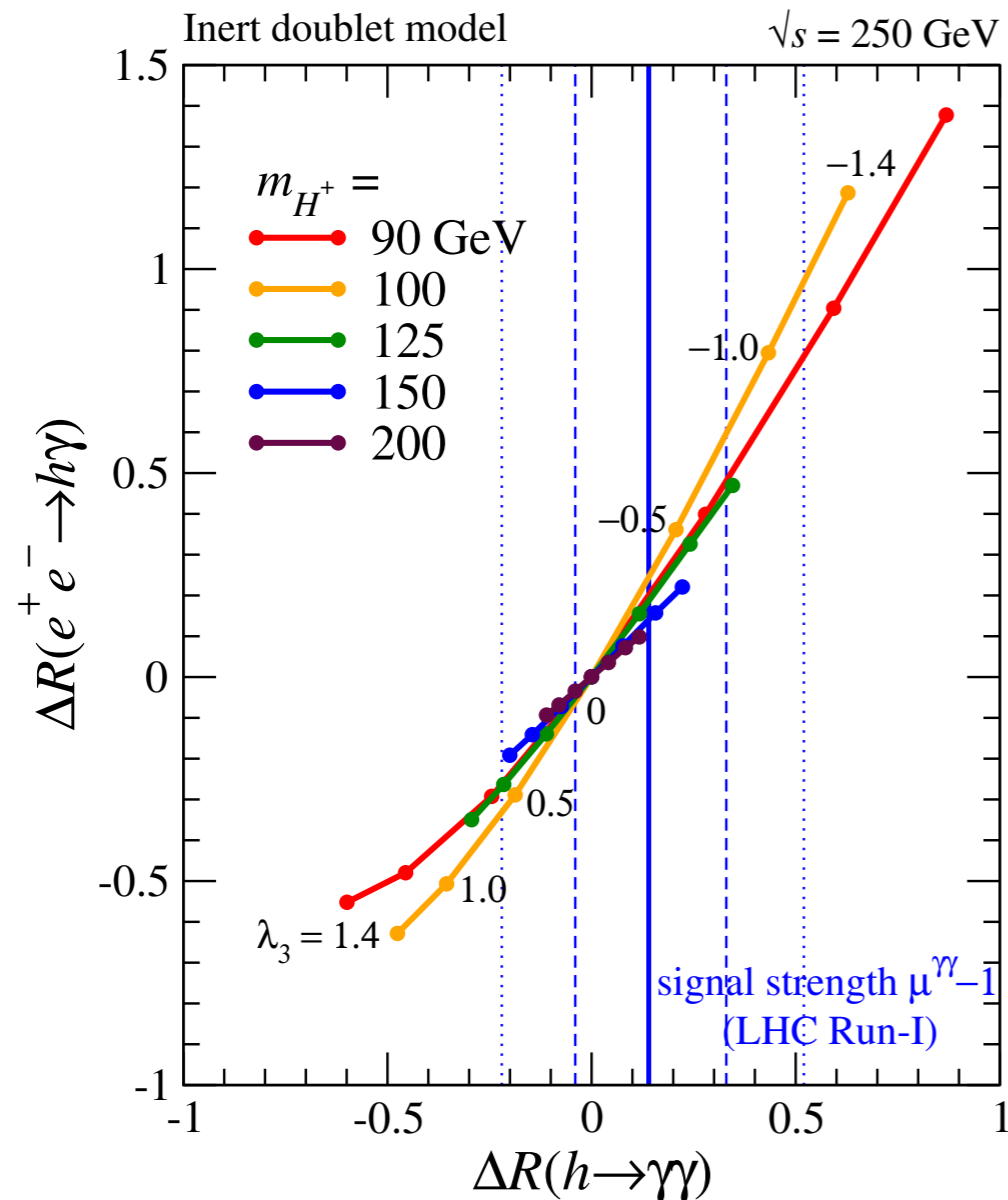


Lighter charged scalars with a negative  $\lambda_3$  can enhance the production rate.

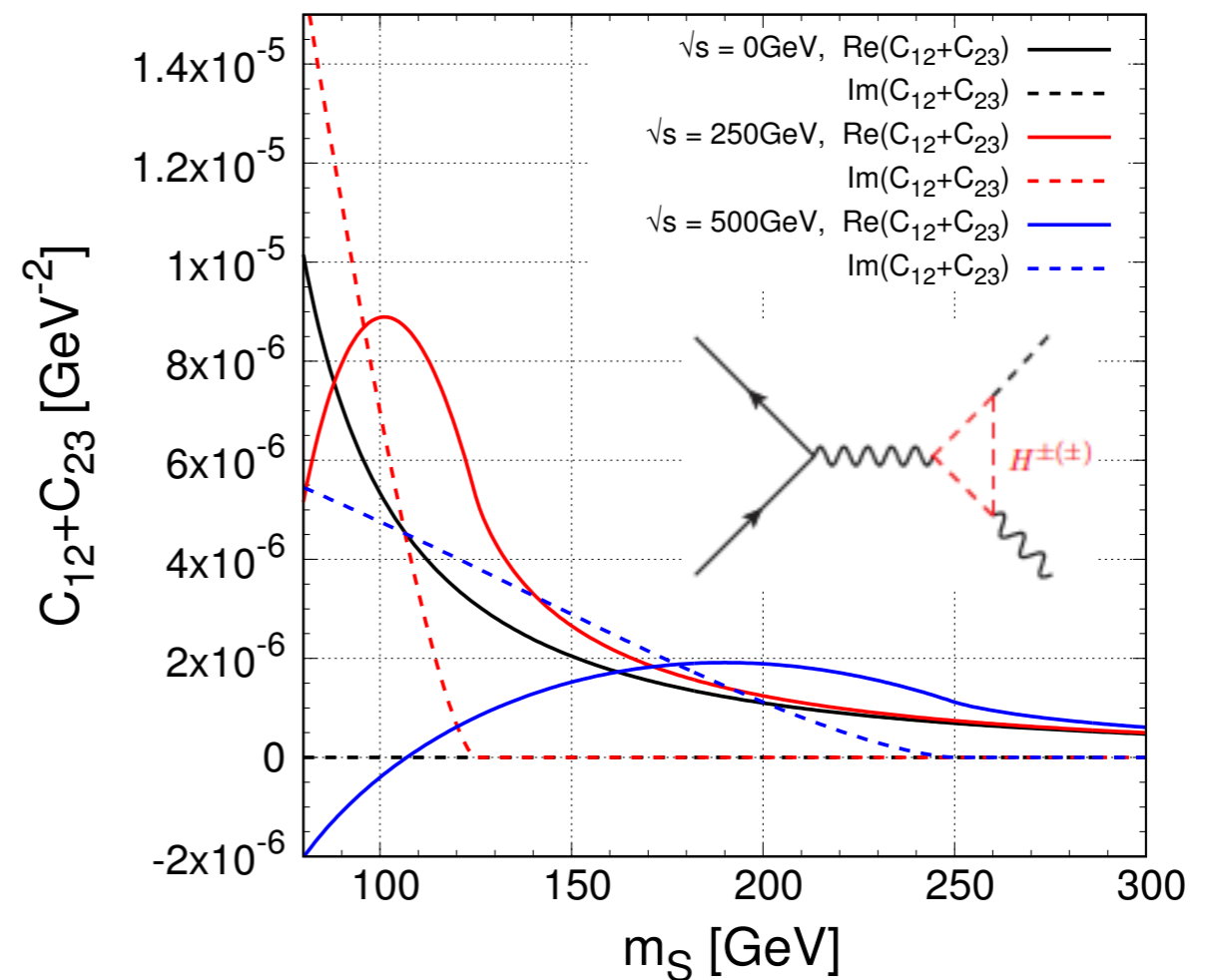


$$\Delta R(e^+e^- \rightarrow h\gamma) = \frac{\sigma_{\text{NP}}(e^+e^- \rightarrow h\gamma)}{\sigma_{\text{SM}}(e^+e^- \rightarrow h\gamma)} - 1, \quad \Delta R(h \rightarrow \gamma\gamma) = \frac{\Gamma_{\text{NP}}(h \rightarrow \gamma\gamma)}{\Gamma_{\text{SM}}(h \rightarrow \gamma\gamma)} - 1$$

# $(e^+e^- \rightarrow h\gamma)$ vs. $(h \rightarrow \gamma\gamma)$ in the IDM



$$\Gamma_{h\gamma\gamma}^{2,1\text{PI}}(p_1^2, p_2^2, p_h^2)_S = -\frac{8e^2}{16\pi^2} \sum_S g_{hSS} Q_S^2 [C_{12}(S, S, S) + C_{23}(S, S, S)]$$

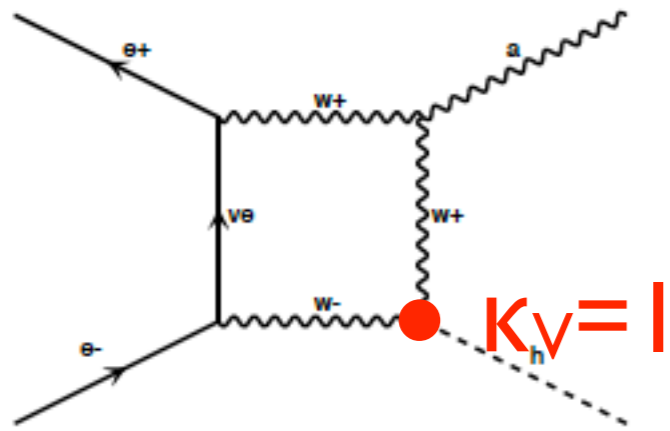


- Strong positive correlation.
- Stronger constraints for light  $H^+$  by the Higgs measurement.

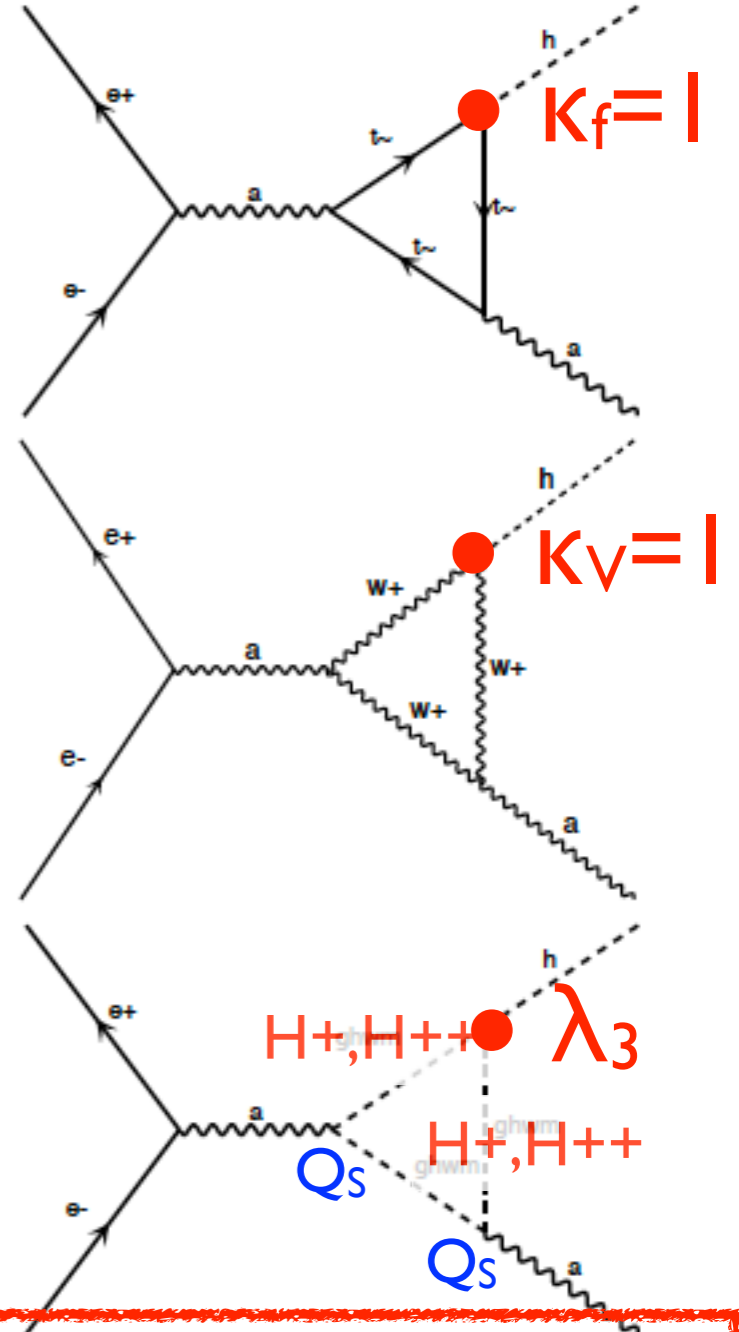
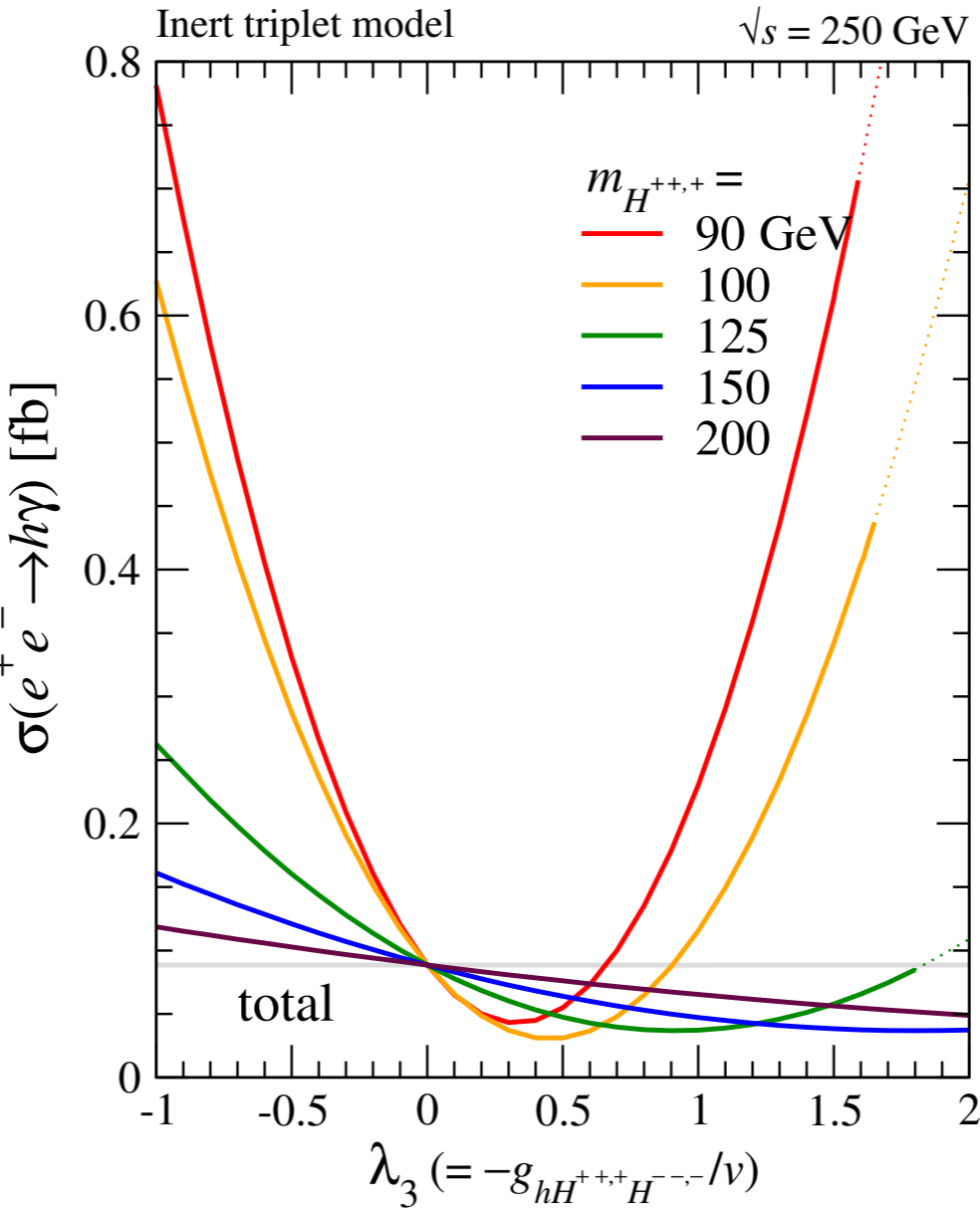
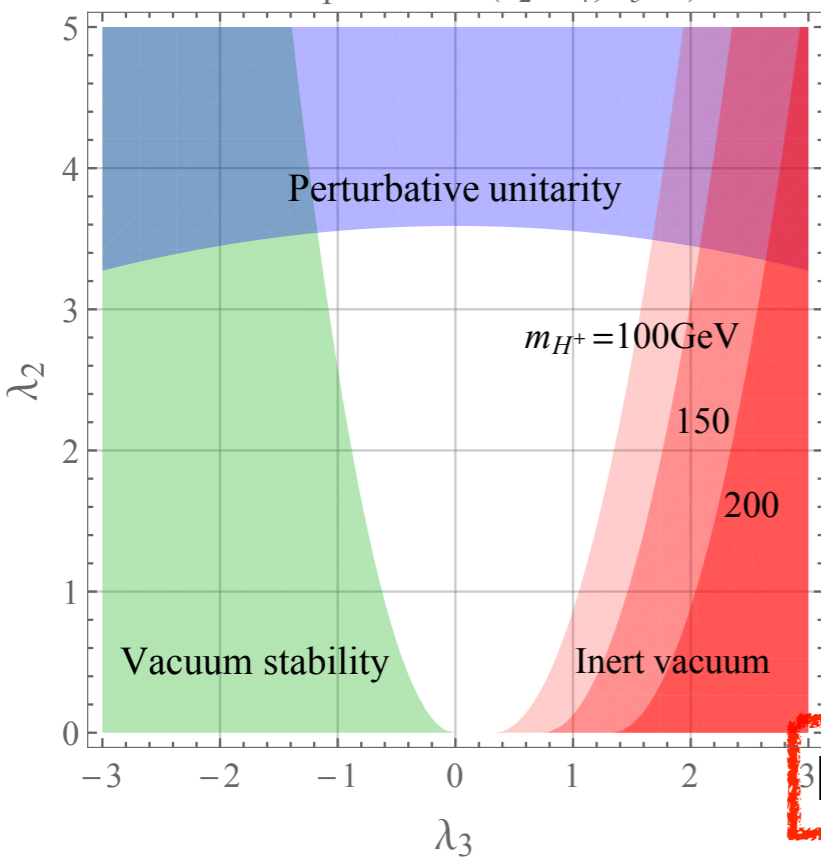
\* An additional Higgs triplet with an exact  $Z_2$  symmetry

# e<sup>+</sup>e<sup>-</sup> → hγ in the ITM

$$\begin{aligned}
 V = & \mu_1^2 |\Phi|^2 + \mu_2^2 \text{Tr}[\Delta^\dagger \Delta] \\
 & + \frac{1}{2} \lambda_1 |\Phi|^4 + \frac{1}{2} \lambda_2 (\text{Tr}[\Delta^\dagger \Delta])^2 \\
 & + \lambda_3 |\Phi|^2 \text{Tr}[\Delta^\dagger \Delta] + \frac{1}{2} \lambda_4 \text{Tr}[(\Delta^\dagger \Delta)^2] \\
 & + \lambda_5 \Phi^\dagger \Delta \Delta^\dagger \Phi
 \end{aligned}$$



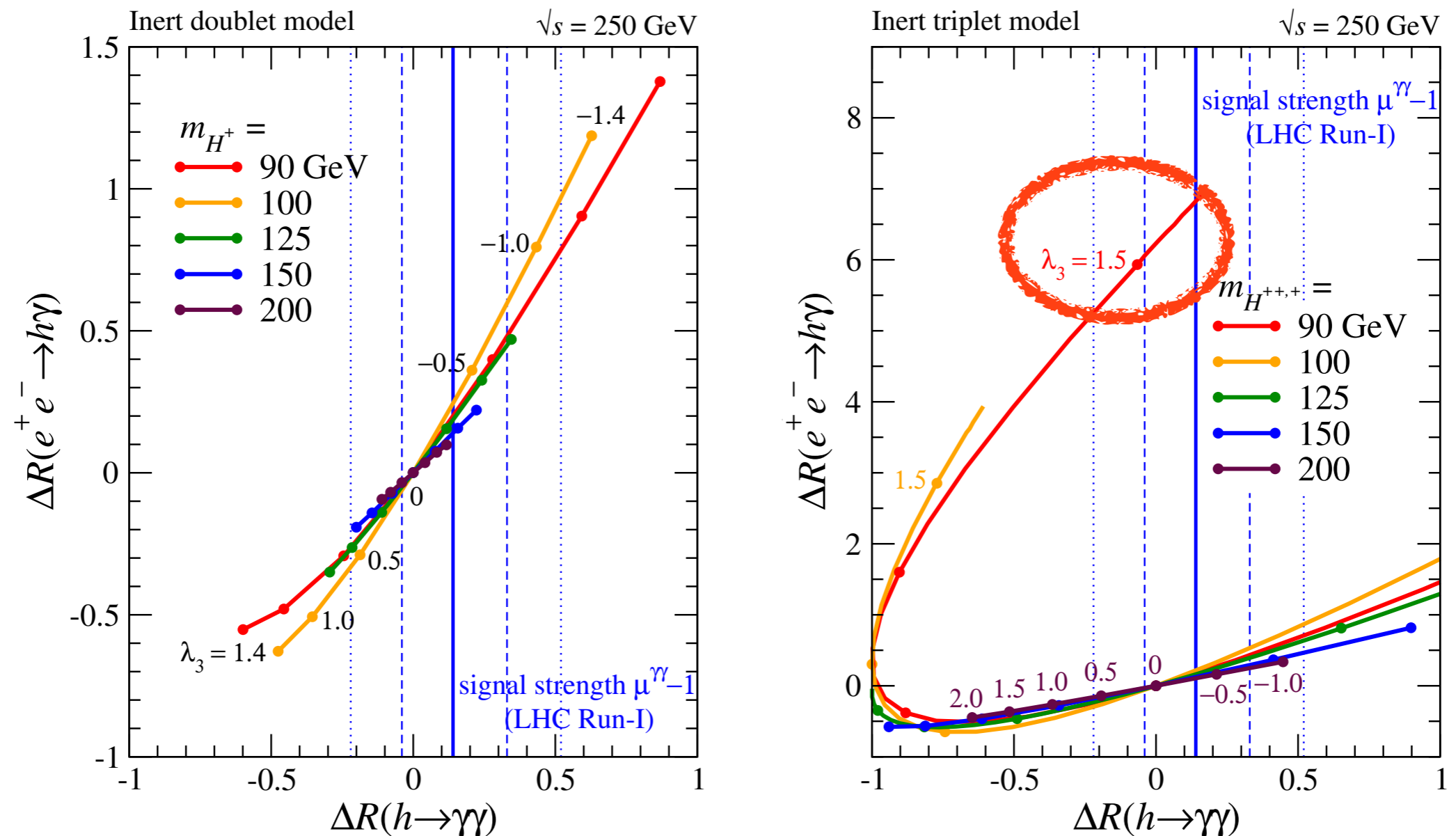
Inert triplet model ( $\lambda_2 = \lambda_4, \lambda_5 = 0$ )



If  $m_{H^{++,+}} \sim 100 \text{ GeV}$ , a positive  $\lambda_3$  can also enhance the production rate.

$$\Delta R(e^+e^- \rightarrow h\gamma) = \frac{\sigma_{\text{NP}}(e^+e^- \rightarrow h\gamma)}{\sigma_{\text{SM}}(e^+e^- \rightarrow h\gamma)} - 1, \quad \Delta R(h \rightarrow \gamma\gamma) = \frac{\Gamma_{\text{NP}}(h \rightarrow \gamma\gamma)}{\Gamma_{\text{SM}}(h \rightarrow \gamma\gamma)} - 1$$

# $(e^+e^- \rightarrow h\gamma)$ vs. $(h \rightarrow \gamma\gamma)$ in the IDM/ITM

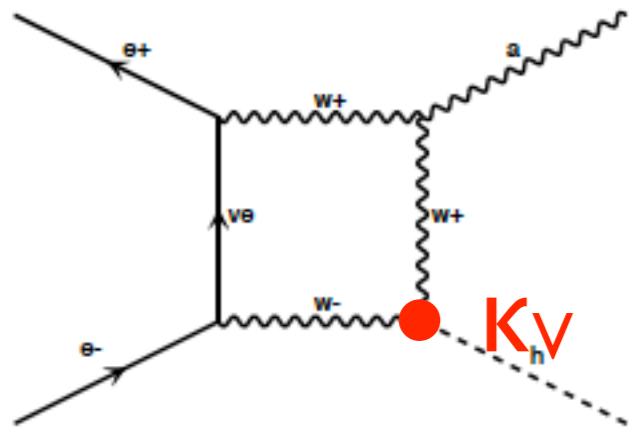


In the ITM, we can find a particular parameter region where the  $h\gamma$  production significantly increases, but the  $h \rightarrow \gamma\gamma$  decay still remains as in the SM.

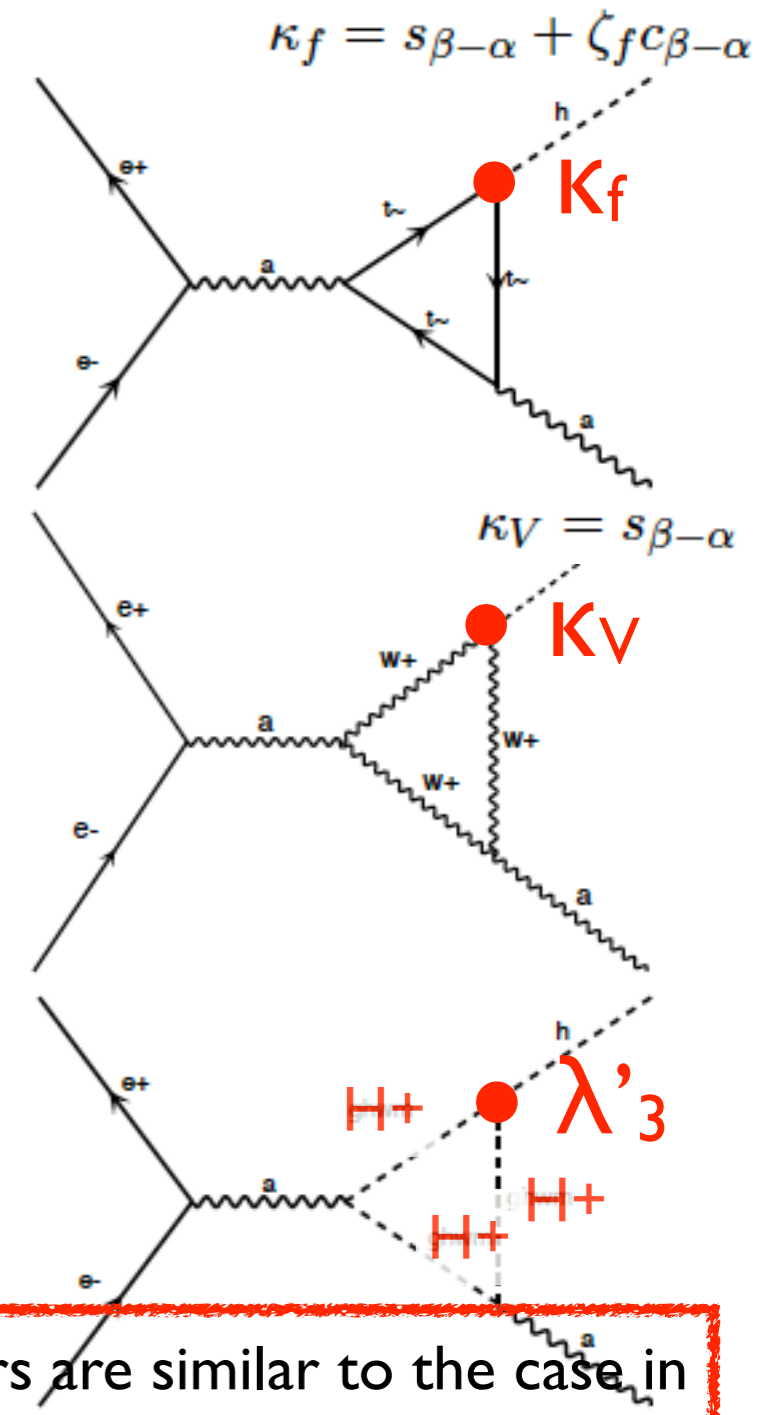
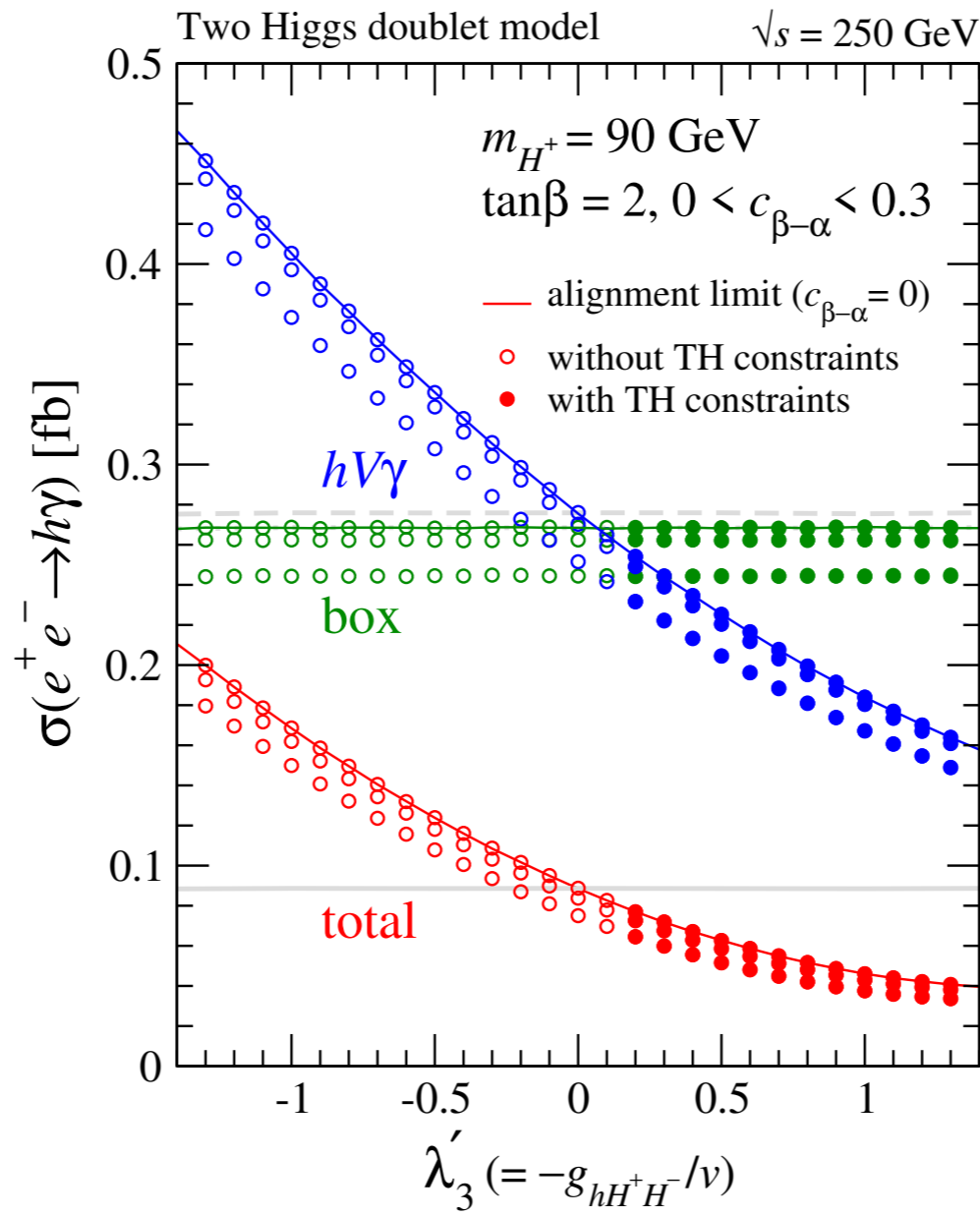
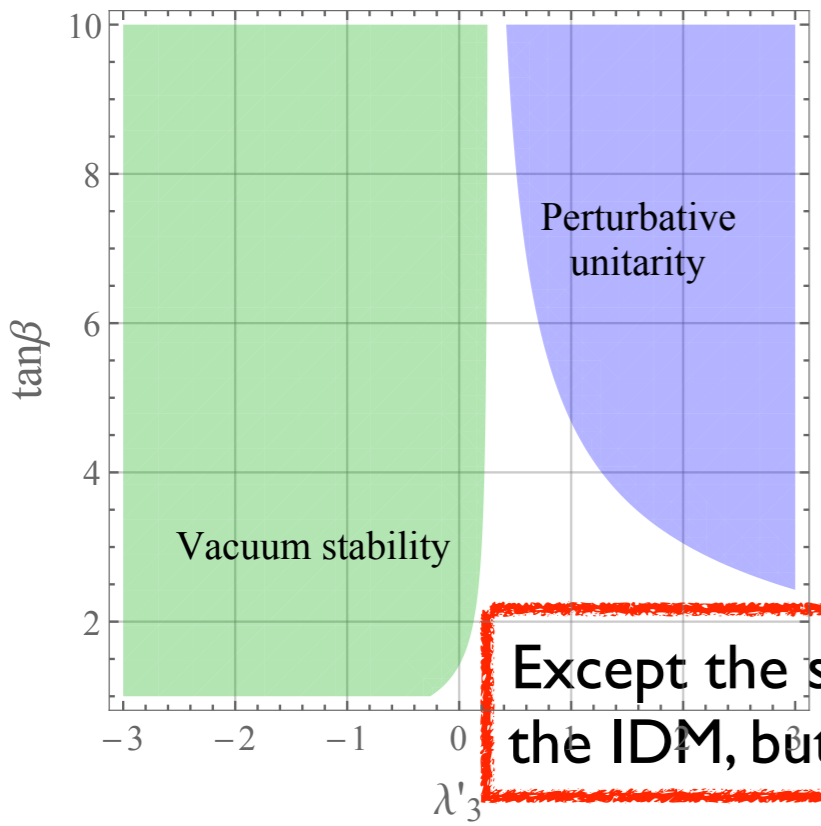
\* An additional Higgs doublet with a softly broken  $Z_2$  symmetry

$$\begin{aligned}
 V = & m_1^2 |\Phi_1|^2 + m_2^2 |\Phi_2|^2 \\
 & - m_3^2 (\Phi_1^\dagger \Phi_2 + \text{h.c.}) \\
 & + \frac{1}{2} \lambda_1 |\Phi_1|^4 + \frac{1}{2} \lambda_2 |\Phi_2|^4 \\
 & + \lambda_3 |\Phi_1|^2 |\Phi_2|^2 + \lambda_4 |\Phi_1^\dagger \Phi_2|^2 \\
 & + \frac{1}{2} \lambda_5 \{ (\Phi_1^\dagger \Phi_2)^2 + \text{h.c.} \}
 \end{aligned}$$

# $e^+ e^- \rightarrow h \gamma$ in the THDM



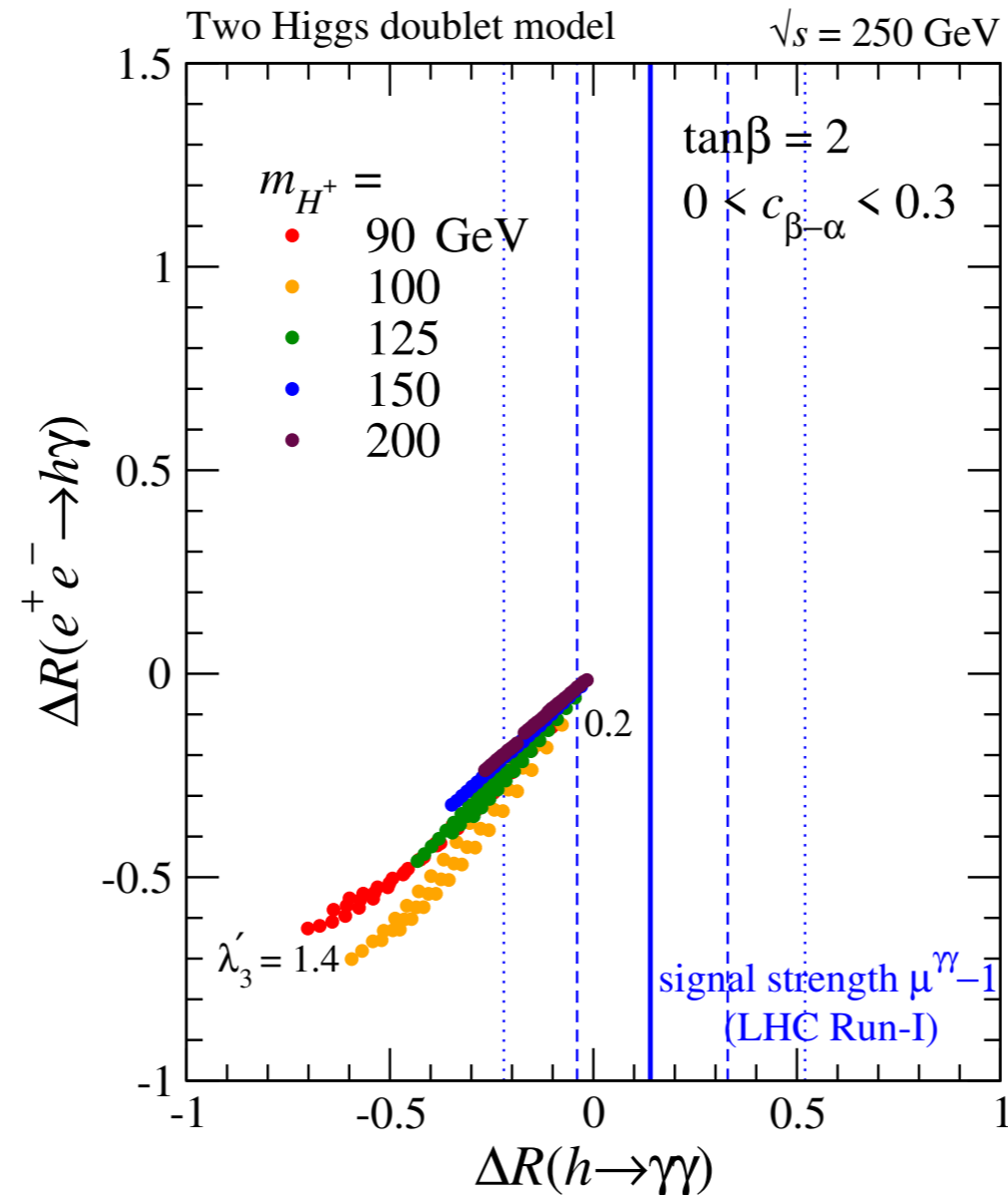
Two Higgs doublet model ( $m_H = m_A = m_{H^+}$ ,  $c_{\beta-\alpha} = 0$ )



Except the small mixing effects, the qualitative behaviors are similar to the case in the IDM, but the enhanced parameter region is excluded by the TH constraints...

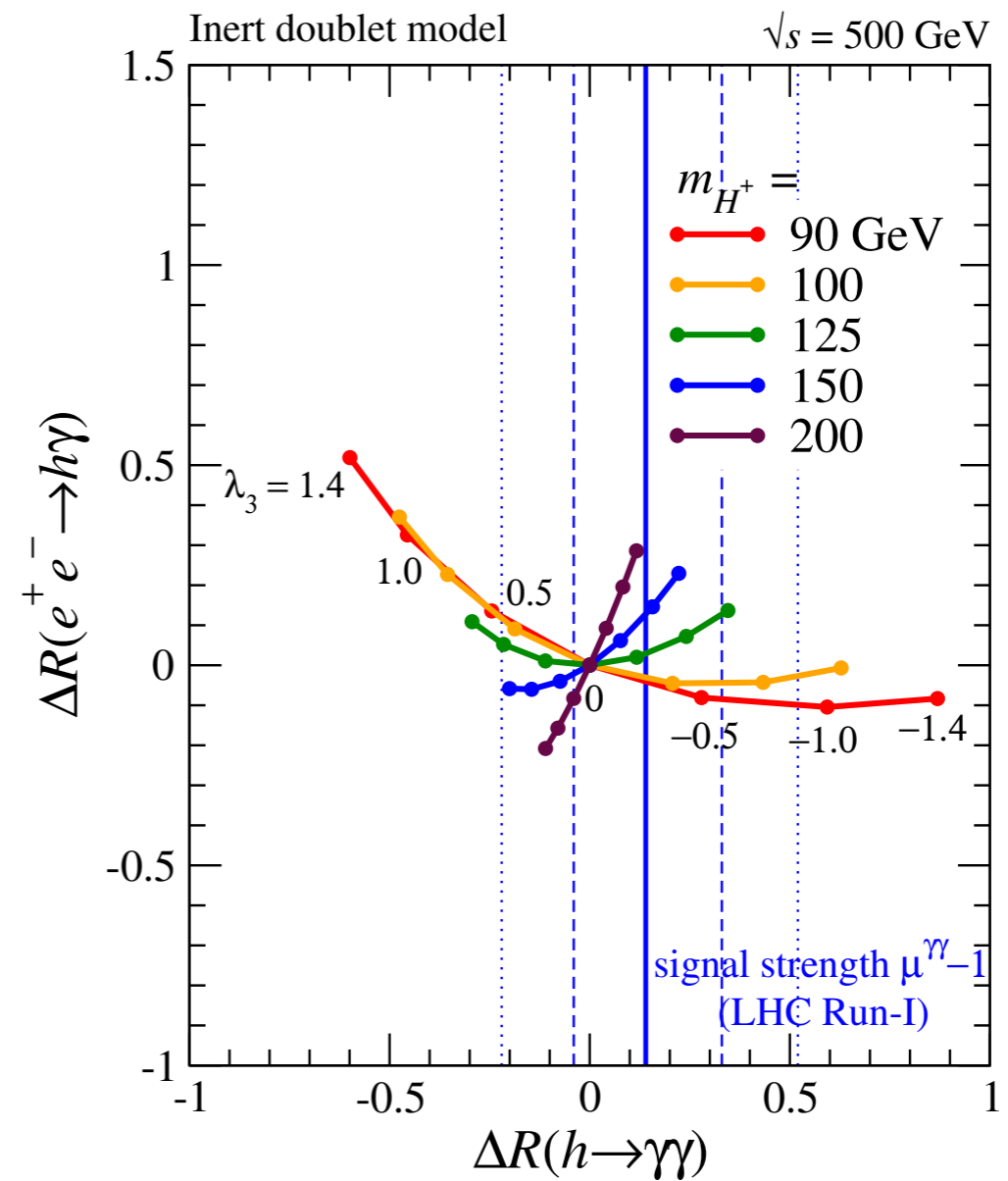
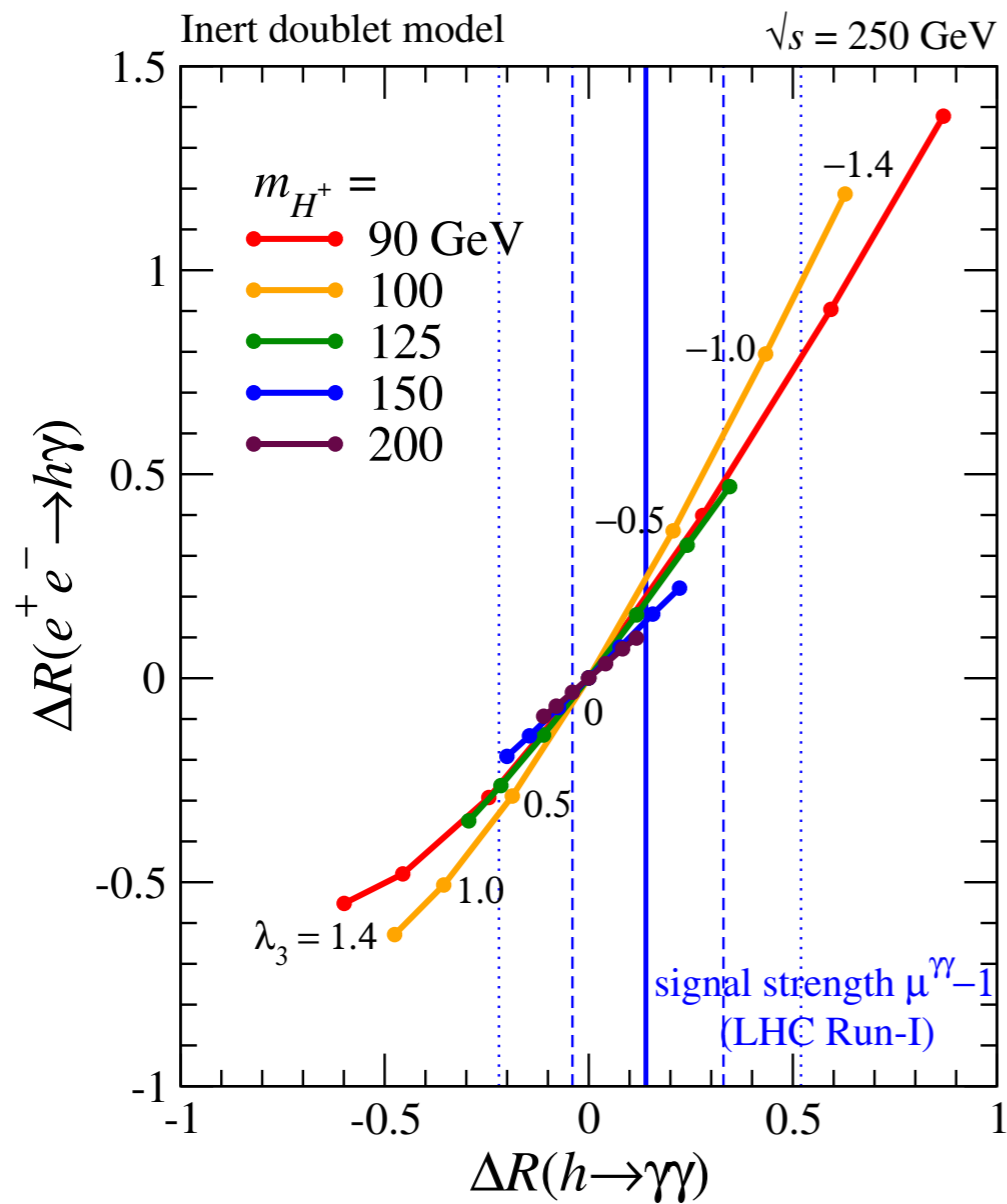
$$\Delta R(e^+e^- \rightarrow h\gamma) = \frac{\sigma_{\text{NP}}(e^+e^- \rightarrow h\gamma)}{\sigma_{\text{SM}}(e^+e^- \rightarrow h\gamma)} - 1, \quad \Delta R(h \rightarrow \gamma\gamma) = \frac{\Gamma_{\text{NP}}(h \rightarrow \gamma\gamma)}{\Gamma_{\text{SM}}(h \rightarrow \gamma\gamma)} - 1$$

# $(e^+e^- \rightarrow h\gamma)$ vs. $(h \rightarrow \gamma\gamma)$ in the THDM



Possible deviations from the SM prediction are minor in the viable parameter space.

# 250 GeV vs. 500 GeV (IDM)



The correlations are different.  $\Rightarrow$  A possibility to access more information on the Higgs sector!

# Summary

- **$h\gamma$  production at lepton colliders:**

- ▶ The cross section is **peaked at  $E=250\text{GeV}$** .
- ▶ **Beam polarization** can enhance the cross section.
- ▶ The signal is clean and very sensitive to **New Physics**.

- ▶ Background study is needed. [Aoki, Fujii, Jung, Lee, Tian, Yokoya](#)  
“Study of the  $h\gamma Z$  coupling” [1902.06029]

- With help from the **H-COUP** program, we studied the process in the benchmark extended Higgs models, such as IDM/ITM/THDM, systematically.

- ▶ Light charged scalars ( $m_{H^+} \sim 100\text{GeV}$ ) can enhance the event rates by a factor of 2 at most under the theoretical and experimental constraints.
- ▶ In the ITM, thanks to doubly charged Higgs bosons, we can also find a particular parameter region where **the  $h\gamma$  production significantly increases**, but the  $h \rightarrow \gamma\gamma$  decay still remains as in the SM.

