

# New perspectives on baryogenesis

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# Many thanks

Many thanks to our conference organizers and local hosts!

Many thanks to my baryogenesis co-conveners: Daniele Teresi & Sacha Davidson

Many thanks to our baryogenesis parallel speakers:

Monday:

<b>Neutrino Option Leptogenesis</b> <i>Grey Room 2, IFT</i>	<i>Jessica Turner</i> 16:40 - 17:00
<b>Type-I Seesaw as the Common Origin of Neutrino Mass, Baryon Asymmetry, and the Electroweak Scale</b> <i>Grey Room 2, IFT</i>	<i>Kai Ruven Schmitz</i> 17:20 - 17:40

Tuesday:

<b>Baryogenesis from Primordial Helical Hypermagnetic Fields</b> <i>Red Room, IFT</i>	<i>Kohei Kamada</i> 17:10 - 17:30
<b>The baryon asymmetry from a composite Higgs</b> <i>Red Room, IFT</i>	<i>Oleksii Matsedonskyi</i> 17:30 - 17:50
<b>High Scale Electroweak Baryogenesis</b> <i>Red Room, IFT</i>	<i>Dr Iason Baldes</i> 17:50 - 18:10
<b>Parameter space of baryogenesis in the vMSM</b> <i>Grey Room 2, IFT</i>	<i>Inar Timiryasov</i> 14:15 - 14:35

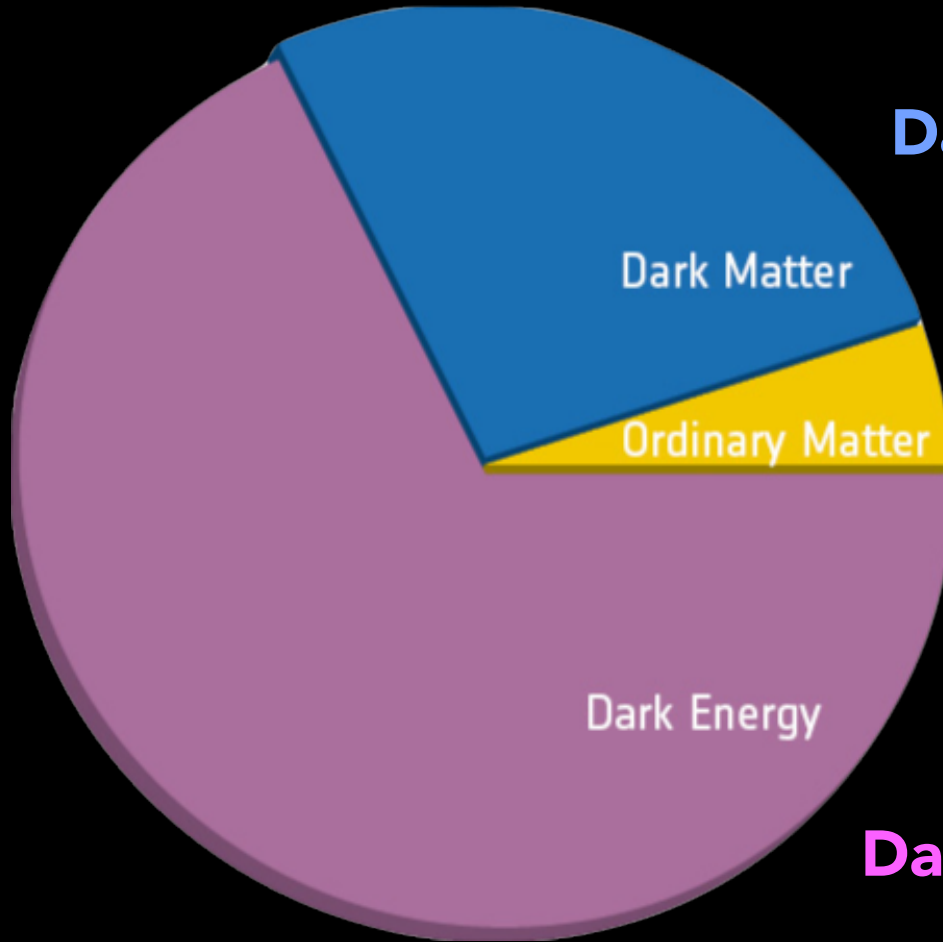
Wednesday:

<b>Low-scale leptogenesis with three heavy neutrinos</b> <i>Grey Room 2, IFT</i>	<i>Juraj Klaric</i> 14:35 - 14:55
<b>Leptogenesis in the Scotogenic model</b> <i>Grey Room 2, IFT</i>	<i>Vedran Brdar</i> 14:55 - 15:15

Thursday:

<b>Baryogenesis and Dark Matter from B Mesons</b> <i>Grey Room 2, IFT</i>	<i>Miguel Escudero</i> 15:40 - 16:00
<b>All-in-one relaxation</b> <i>Grey Room 2, IFT</i>	<i>Jakub Scholtz</i> 16:00 - 16:20
<b>Leptogenesis without Loops</b> <i>Grey Room 2, IFT</i>	<i>Amab Dasgupta</i> 16:20 - 16:40

# Why baryogenesis?



**Dark Matter: what is it?**

**Ordinary Matter:  
where does it come from?**

**( interesting: uniquely connected  
with elementary particle physics )**

**Dark Energy: what is it?**

baryogenesis = a dynamical process that occurred in the early universe and led to the cosmological excess of matter over antimatter

# Baryogenesis needs B-number violation

[Planck (2018)]

There is more matter than antimatter

cosmology says there are  
more protons than anti-protons

$$\Omega_b h^2 \simeq 0.02230 \pm 0.00014$$

$$n_B/s \simeq (0.861 \pm 0.005) \times 10^{-10}$$

We use baryon number to quantify the matter excess

charge assignments:

$$B(\text{quark}) = +1/3$$

$$B(\text{anti-q}) = -1/3$$

$$B(\text{proton}) = B(\text{neutron}) = +1$$

$$B(\text{anti-p}) = B(\text{anti-n}) = -1$$

$$B(\text{mesons}) = B(\text{leptons}) = 0$$

Baryogenesis needs B-number violation

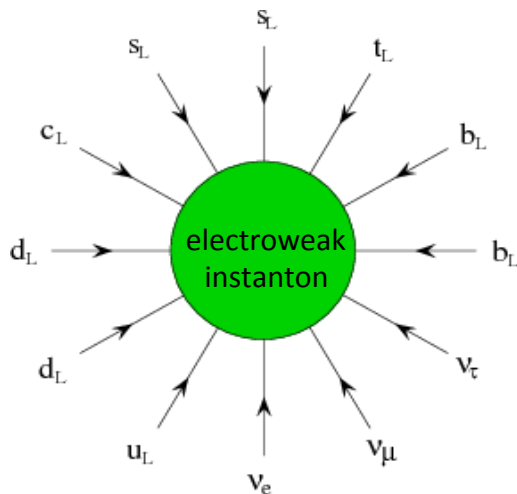
$$B(\text{universe}) = 0 \quad \xrightarrow{\text{baryogenesis}} \quad B(\text{universe}) \neq 0$$



The chiral EW interactions lead to anomalies in B and L

$$\partial_\mu j_B^\mu = \partial_\mu j_L^\mu = 3 \times \left( \frac{g^2}{32\pi^2} W_{\mu\nu}^a \widetilde{W}^{a\mu\nu} - \frac{g'^2}{32\pi^2} Y_{\mu\nu} \widetilde{Y}^{\mu\nu} \right)$$

Large quantum fluctuations of  $W_\mu^a$  (aka electroweak instanton) induce reactions among the quarks and leptons that violate baryon number



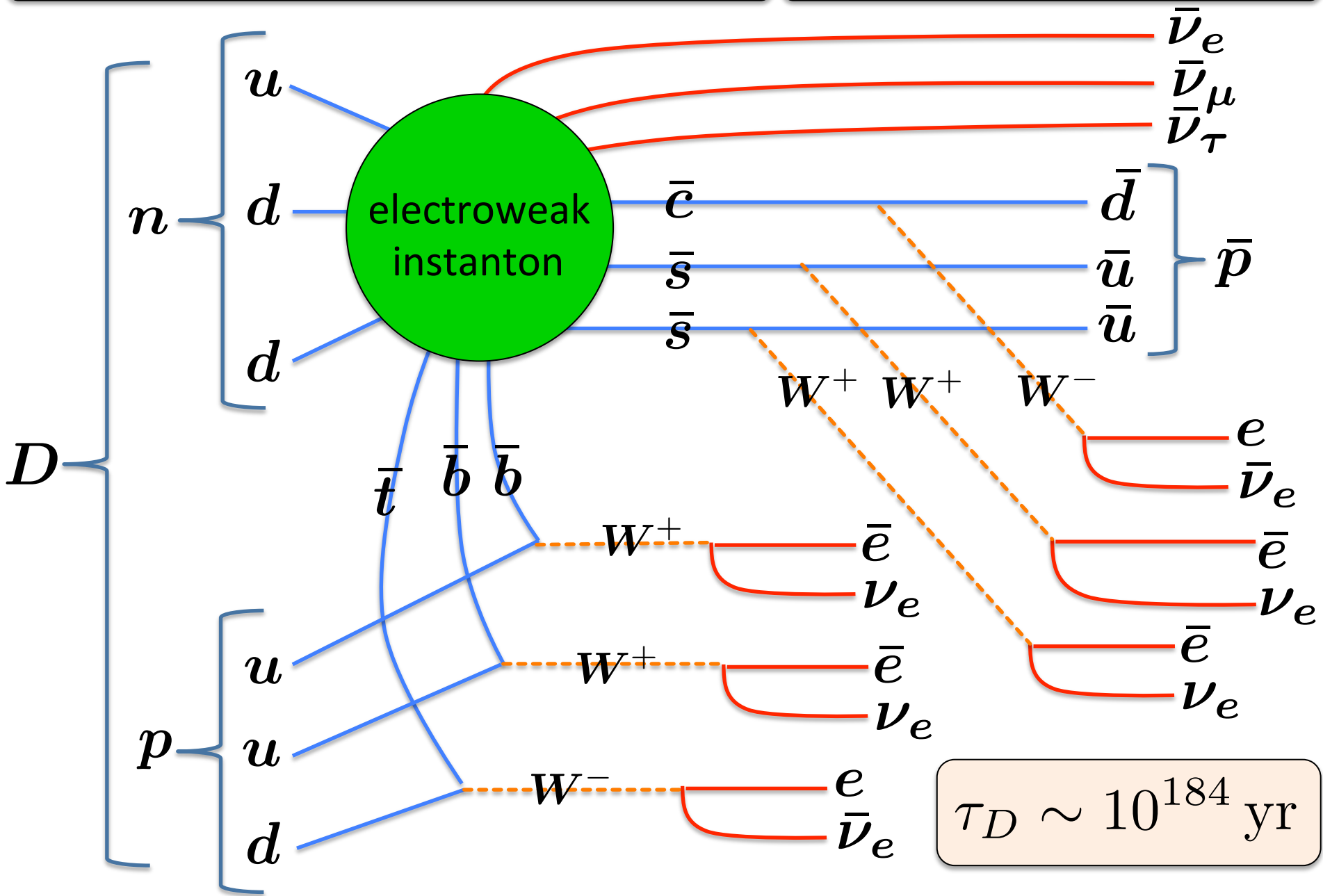
$$\langle W \widetilde{W} \rangle = \text{quantum}$$

$$\Delta B = \Delta L = \pm 3$$

The Standard Model predicts  
**catastrophic nuclear decay!**

$$D \rightarrow \bar{p} + 4e^+ + 2e^- + 4\nu_e + 3\bar{\nu}_e + \bar{\nu}_\mu + \bar{\nu}_\tau$$

$$\Gamma \sim G_F^{12} m_D^{25} V_{td}^2 V_{ub}^4 V_{cd}^2 V_{us}^4 e^{-16\pi^2/g^2}$$





early time <--> high temperature

10 ps

13.7 Gyr

$10^{184}$  yr

increasing  
time

↑  
thermal B-violation  
 $\langle W\widetilde{W} \rangle = \text{thermal}$

↑  
today



↑  
quantum B-violation  
 $\langle W\widetilde{W} \rangle = \text{quantum}$

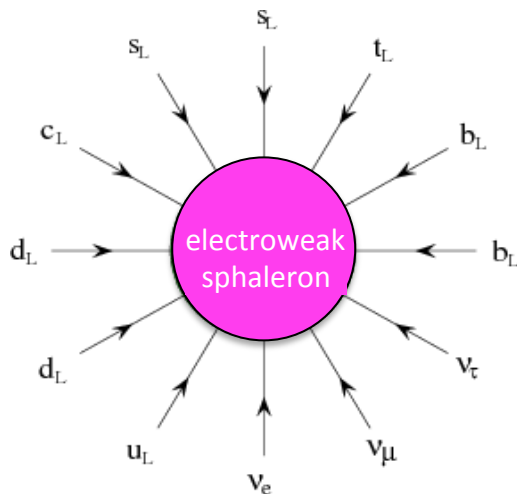
# Standard Model baryon-number violation @ high temperature

[Klinkhamer & Manton (1984); Kuzmin, Rubakov, & Shaposhnikov (1985)]

The chiral EW interactions lead to anomalies in B and L

$$\partial_\mu j_B^\mu = \partial_\mu j_L^\mu = 3 \times \left( \frac{g^2}{32\pi^2} W_{\mu\nu}^a \widetilde{W}^{a\mu\nu} - \frac{g'^2}{32\pi^2} Y_{\mu\nu} \widetilde{Y}^{\mu\nu} \right)$$

At high temperature large thermal fluctuations of  $W_\mu^a$  (aka electroweak sphaleron) mediate unsuppressed baryon-number violation



at  $T > 100$  GeV

$\langle W \widetilde{W} \rangle = \text{thermal}$

$\Delta B = \Delta L = \pm 3$

$\Gamma_B \sim \alpha_W^5 T$

SM B-number violation plays an important role in many models of baryogenesis

# outline

(1) testing electroweak baryogenesis  
on the collider & cosmological frontiers

(2) a new perspective on  
baryogenesis from (lepto-)bubbles

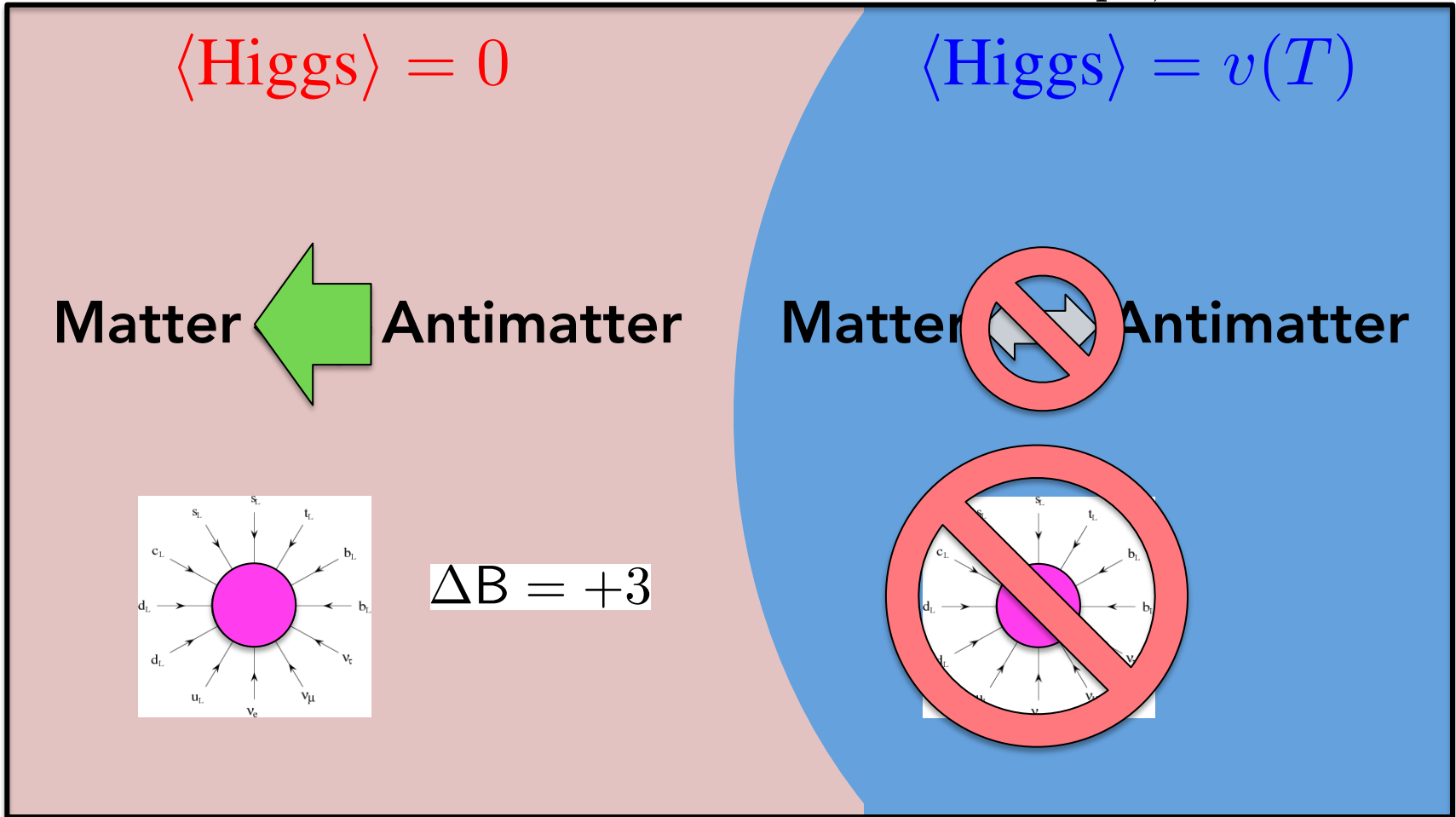
(3) baryon asymmetry  
from magnetic helicity

# Electroweak baryogenesis

[Kuzmin, Rubakov, & Shaposhnikov (1985); Cohen, Kaplan, & Nelson (1990)]

If the SM admits B-number violation does it admit baryogenesis?

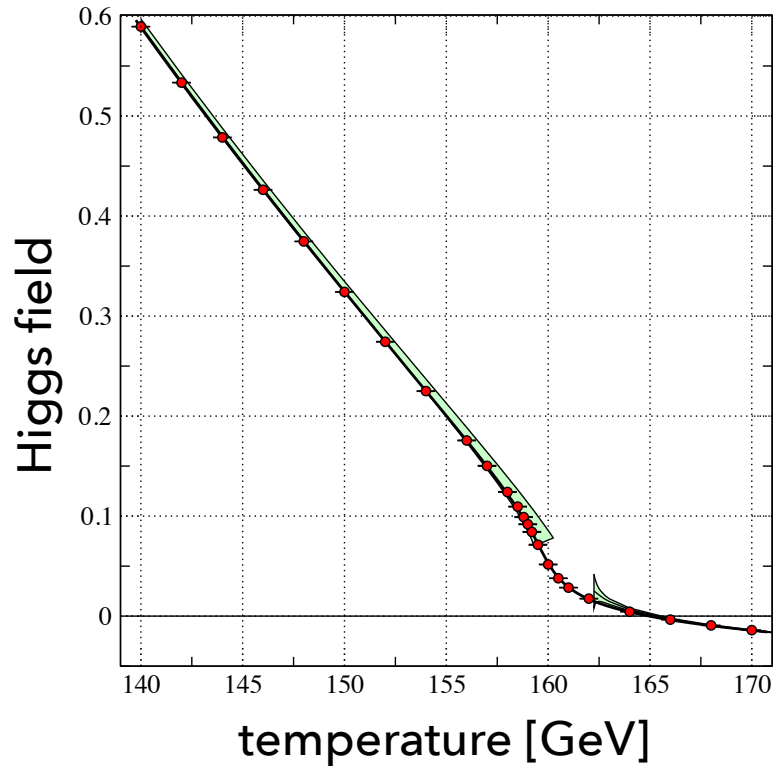
$t \approx 10 \text{ ps}$ ,  $T \approx 100 \text{ GeV}$



# To bubble or not to bubble? ... that's the trouble!

[D'Onofrio & Rummakainen (2015)]

the SM predicts



continuous crossover  
(no bubbles)

if nature allows



then new physics is  
required!

# Many examples of new physics for 1<sup>st</sup> order EWPT

Model	References
SM + Scalar Singlet	Espinosa & Quiros, 1993; Benson, 1993; Choi & Volkas, 1993; McDonald, 1994; Vergara, 1996; Branco, Delepine, Emmanuel-Costa, & Gonzalez, 1998; Ham, Jeong, & Oh, 2004; Ahriche, 2007; Espinosa & Quiros, 2007; Profumo, Ramsey-Musolf, & Shaughnessy, 2007; Noble & Perelstein, 2007; Espinosa, Konstandin, No, & Quiros, 2008; Ashoorioon & Konstandin, 2009; Das, Fox, Kumar, & Weiner, 2009; Espinosa, Konstandin, & Riva, 2011; Chung & AL, 2011; Wainwright, Profumo, & Ramsey-Musolf, 2012; Barger, Chung, AL, & Wang, 2012; Huang, Shu, Zhang, 2012; Katz & Perelstein, 2014; Jiang, Bian, Huang, Shu, 2015; Huang & Li 2015; Cline, Kainulainen, Tucker-Smith, 2017; Kurup & Perelstein, 2017; Chen, Kozaczuk, & Lewis, 2017
SM + Scalar Doublet	Davies, Froggatt, Jenkins, & Moorhouse, 1994; Huber, 2006; Fromme, Huber, & Seniuch, 2006; Cline, Kainulainen, & Trott, 2011; Kozhushko & Skalozub, 2011;
SM + Scalar Triplet	Patel, Ramsey-Musolf, 2012; Patel, Ramsey-Musolf, Wise, 2013; Huang, Gu, Yin, Yu, Zhang 2016
SM + Chiral Fermions	Carena, Megevand, Quiros, Wagner, 2005
MSSM	Carena, Quiros, & Wagner, 1996; Delepine, Gerard, Gonzales Felipe, & Weyers, 1996; Cline & Kainulainen, 1996; Laine & Rummukainen, 1998; Cohen, Morrissey, & Pierce,; Carena, Nardini, Quiros, & Wagner, 2012;
NMSSM / nMSSM / $\mu\nu$ SMS	Pietroni, 1993; Davies, Froggatt, & Moorhouse, 1995; Huber & Schmidt, 2001; Ham, Oh, Kim, Yoo, & Son, 2004; Menon, Morrissey, & Wagner, 2004; Funakubo, Tao, & Toyoda, 2005; Huber, Konstandin, Prokopec, & Schmidt, 2006; Chung, AL, 2010, Huang, Kang, Shu, Wu, Yang, 2014; Bian, Guo, Shu (2017)
EFT-like Approach ( $H^6$ operator)	Grojean, Servant, Wells, 2005; Huang, Gu, Yin, Yu, Zhang 2015; Huang, Joglekar, Li, Wagner, 2015; Huang, Wan, Wang, Cai, Zhang 2016; Huang, Gu, Yin, Yu, Zhang 2016; Cao, F.P. Huang, Xie, Zhang (2017)



# An example of new physics

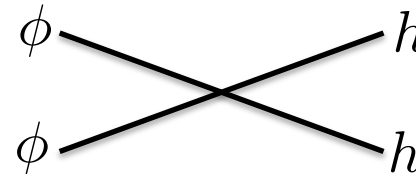
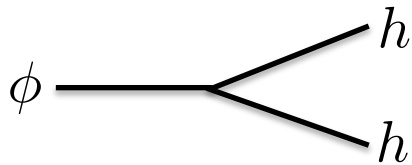
[Espinosa & Quiros (1993); Choi & Volkas (1993); McDonald (1994); Profumo, Ramsey-Musolf, & Shaughnessy (2007)]

[Chung, AL, & Wang (2012)]

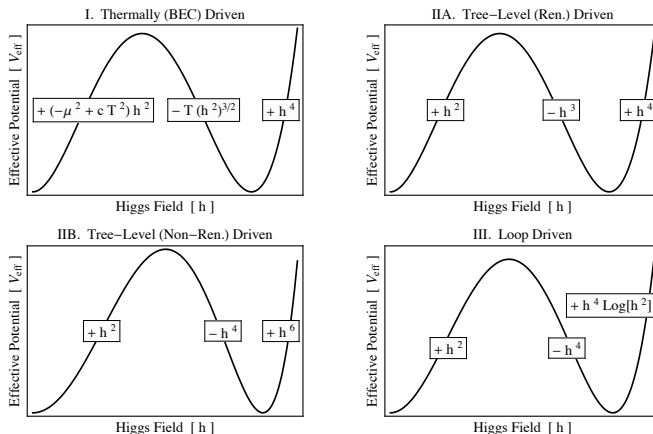
## Consider the theory

SM + a spin-0 colorless uncharged particle [i.e., real scalar singlet]

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{2}(\partial\phi)^2 - a_2\phi^2 - a_3\phi^3 - a_4\phi^4 - \underbrace{b_1\phi H^\dagger H - b_2\phi^2 H^\dagger H}$$



Its interactions with the Higgs can lead to a 1<sup>st</sup> order EW PT



However such an inert particle is especially challenging to detect!

# Indirect signatures: Higgs precision

[Craig, Englert, & McCullough (2013); McCullough (2014); Curtin, Meade, & Yu (2014)]

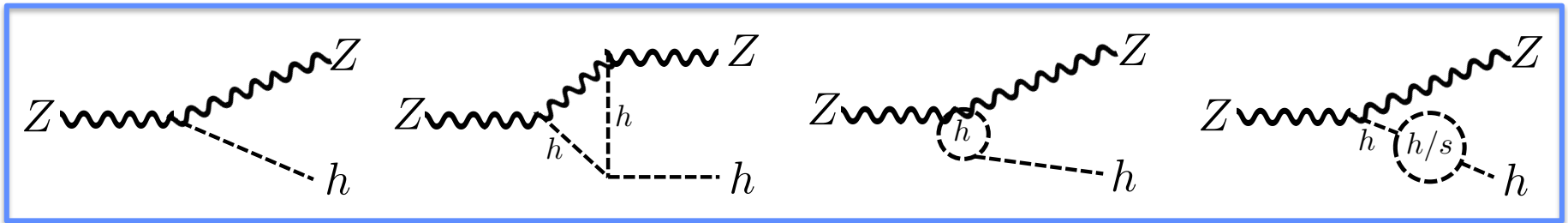
The new physics affects how strongly the Higgs couples to other SM particles such as the Z-boson

$$\frac{g_{hZZ}}{g_{hZZ,SM}} \approx \cos \theta + 0.006 \left( \frac{\lambda_3}{\lambda_{3,SM}} - 1 \right) - 2 \frac{|a_{hs} + \lambda_{hs} v_s|^2}{16\pi^2} I_B(M_h^2; M_h^2, M_s^2) - \frac{|\lambda_{hs}|^2 v^2}{16\pi^2} I_B(M_h^2; M_s^2, M_s^2).$$

(leading effect is from mixing)

(triangle probes self-coupling)

(Higgs wfcn. renorm.)

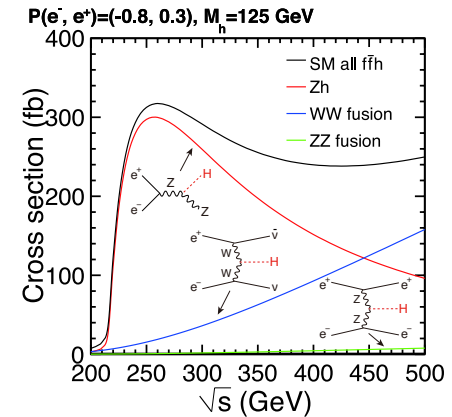
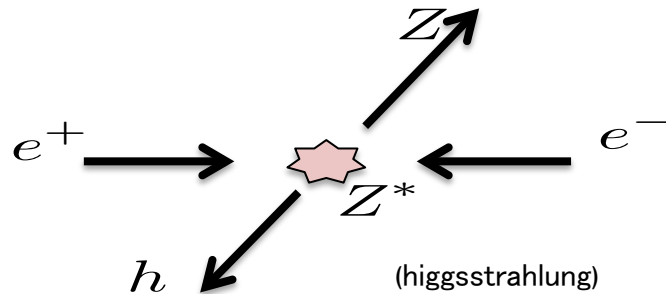


Current LHC measurements @ 27% precision [1606.02266]

Projected HL-LHC sensitivity is ~7% precision [1307.7135]

# The future collider program

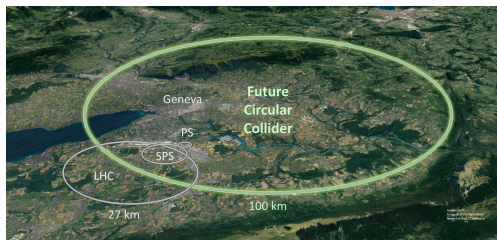
Higgs factory:  $e^+e^- @ E_{\text{com}} = 250 \text{ GeV}$



a primary goal is to measure the  $hZZ$  coupling with sub-% precision

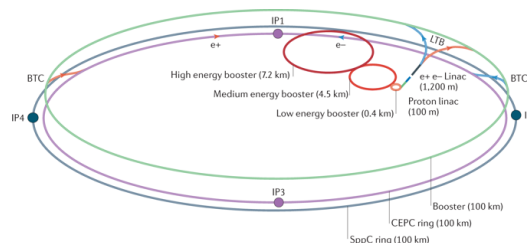
## Future Circular Collider (FCC)

Europe / CERN



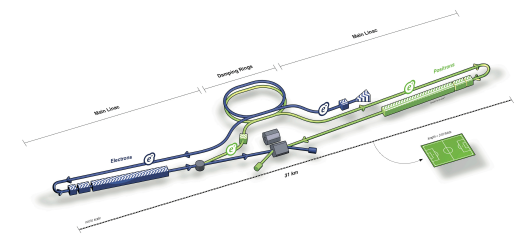
## Circular Electron-Positron Collider (CEPC)

China



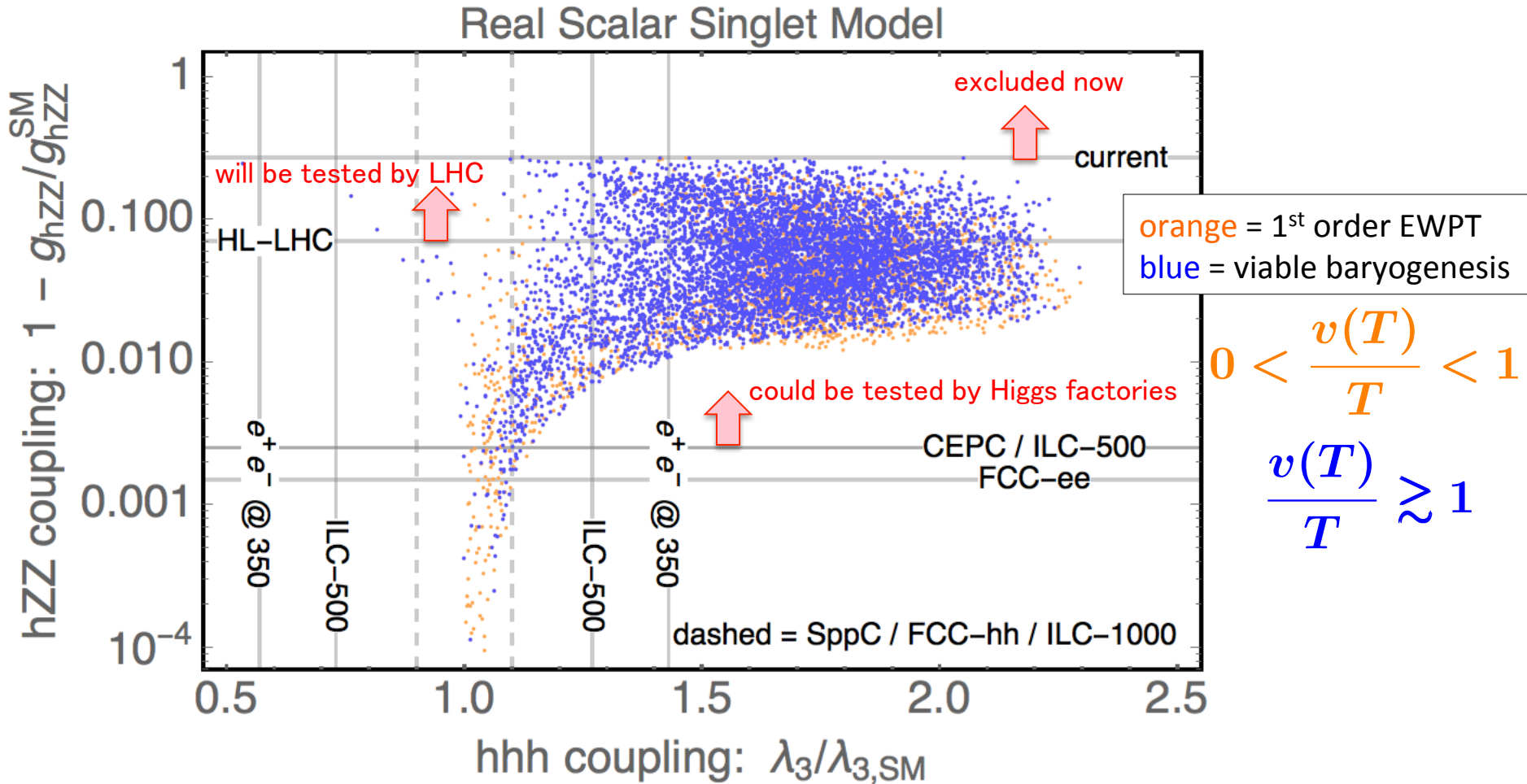
## International Linear Collider (ILC)

Japan



# Prospects for testing EW phase transition

[Huang, AL, & Wang (2016)]

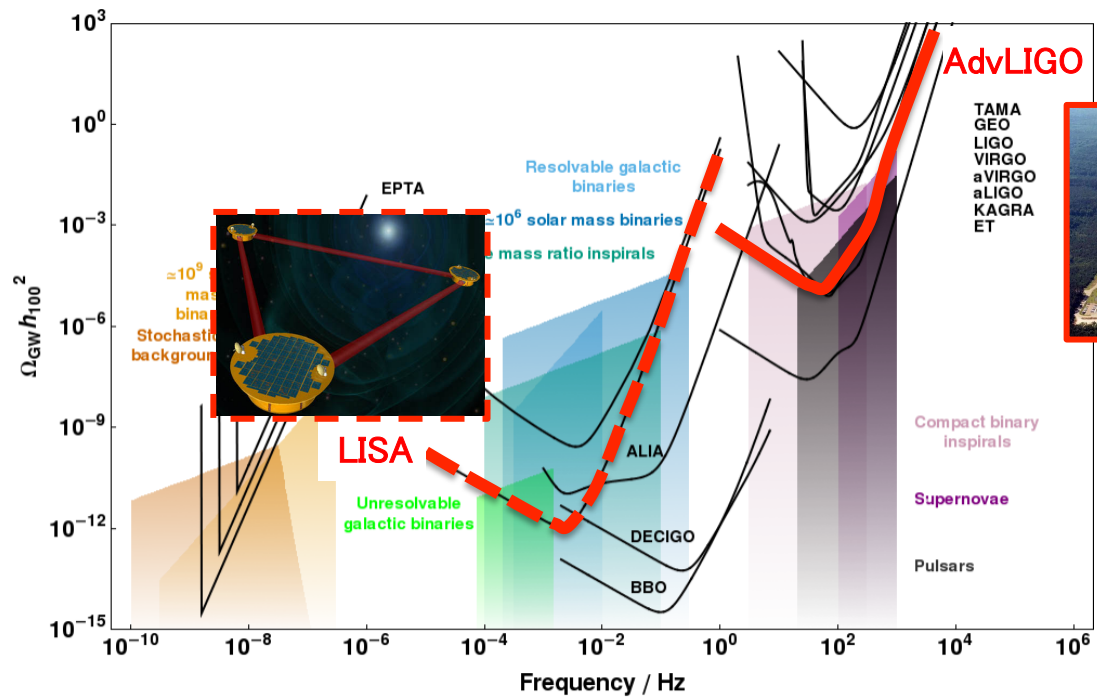
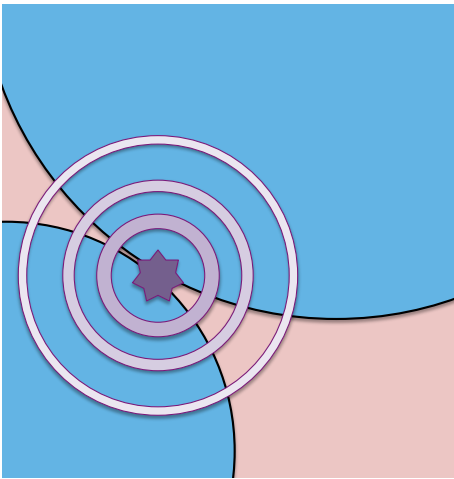


# Complementarity with gravitational waves

[Caprini et. al. (2015)]

## Bubble collisions & fluid motion create a gravitational wave noise

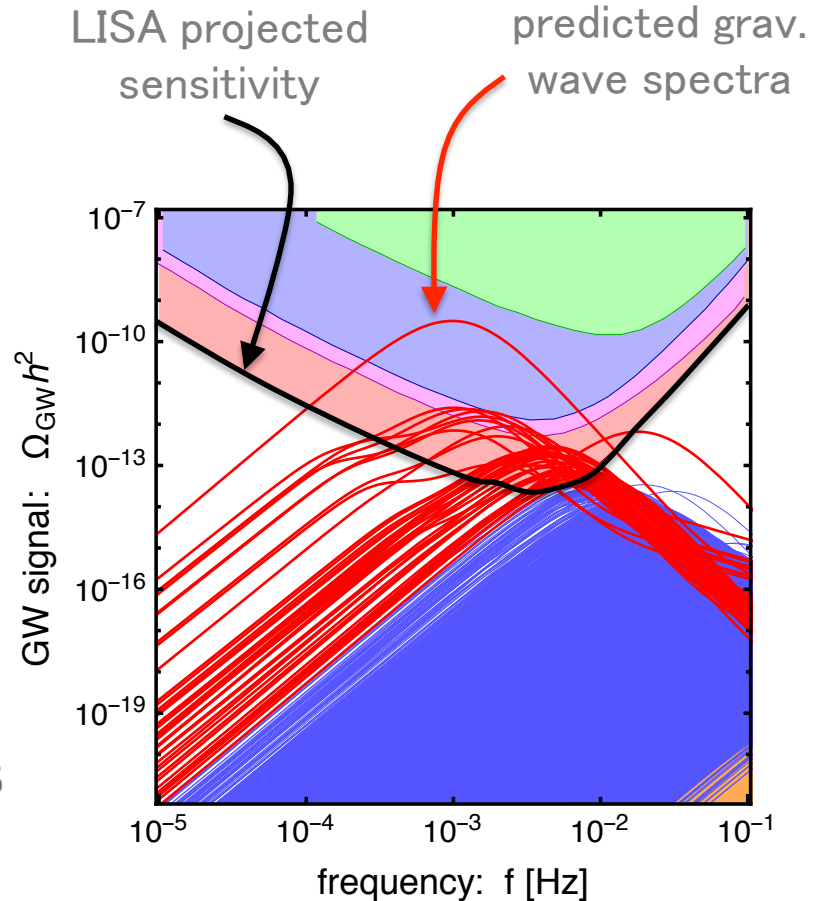
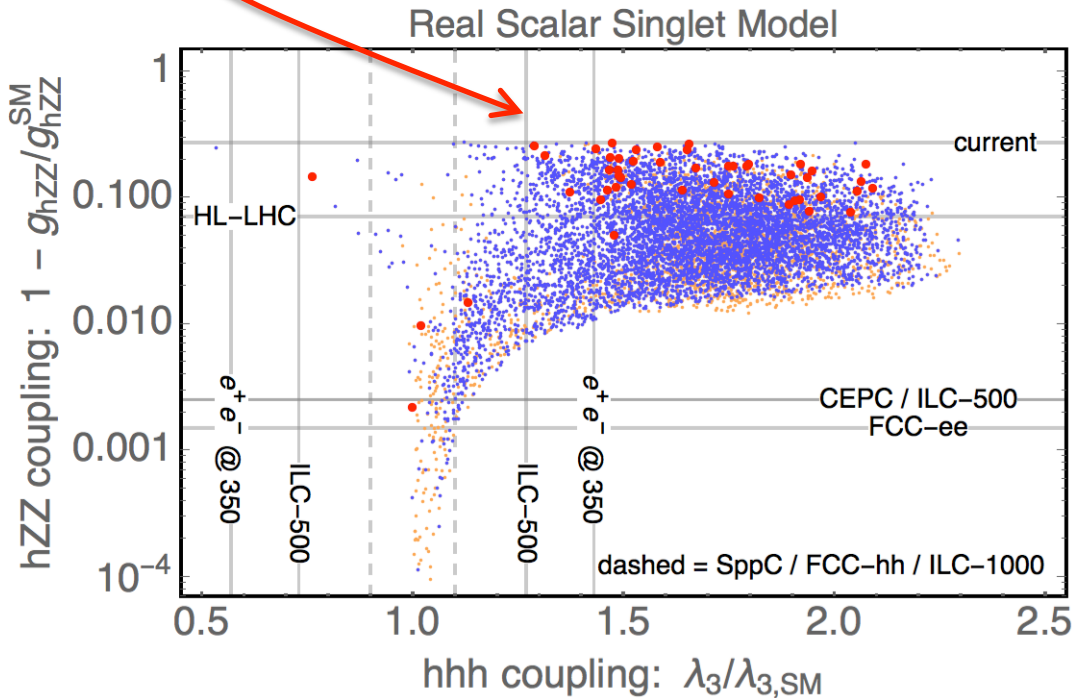
GW spectrum is set by the bubble size at the time of collision & typically falls right into the sensitivity band of future space-based GW interferometer telescopes



# Complementarity with gravitational waves

[Huang, AL, & Wang (2016)]

orange = 1<sup>st</sup> order EWPT  
blue = also viable baryogenesis  
red = also detectable GWs



# outline

## (I) testing electroweak baryogenesis on the collider & cosmological frontiers

- The new physics required for EWBG has thus far evaded our grasp at the LHC.
- With precision measurements of Higgs couplings ( $hZZ$ ) at the HL-LHC upgrade and a future Higgs factory, we can test many of even the most elusive models.
- A deviation from SM predictions would give *indirect* evidence for the kind of new physics that leads to a first order EW phase transition, but a detection of the associated gravitational wave signal is the best *direct* evidence for a cosmological first order phase transition.

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(1) testing electroweak baryogenesis  
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(2) a new perspective on  
baryogenesis from (lepto-)bubbles

(3) baryon asymmetry  
from magnetic helicity



# Alternate applications of our EWBG tool kit

While studying EW baryogenesis we've developed a big tool kit

We've learned how to calculate ...

- ... when the phase transition will be 1st order
- ... phase transition temperature & order parameter
- ... phase transition duration & latent heat
- ... stochastic gravitational wave spectrum
- ... CP-violating scattering at the bubble wall (source term)
- ... charge diffusion in front of the wall
- ... charge transport & washout

Can we use our tool kit to study BG at other phase transitions?

Symmetry breaking associated with

- ... grand unification
- ... lepton-number
- ... flavor
- ... dark sector
- ... chiral symmetry
- ... composite Higgs

see works by: Iason Baldes, Jim Cline, Eleanor Hall, Kimmo Kainulainen, Thomas Konstandin, Robert McGehee, Hitoshi Murayama, Kai Schmitz, Geraldine Servant, David Tucker-Smith, ...

# High-scale lepton-number violation

Tiny neutrino mass scale from huge L-violation mass scale

type-I seesaw: 
$$m_\nu \sim (0.1 \text{ eV}) \left( \frac{m_{\text{L-violation}}}{10^{12} \text{ GeV}} \right)^{-1} \left( \frac{\lambda_N}{0.1} \right)^2$$

Consider the theory [singlet Majoron model]

SM + complex scalar  $\phi$  + Weyl fermion  $N$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} - \frac{1}{2} \kappa \phi N N - \lambda_N L H N + \text{h.c.}$$

Lepton-number is broken spontaneously by  $\langle \phi \rangle$ :  $m_N = \kappa v_\phi / 2$ .

If the corresponding phase transition is 1<sup>st</sup> order, we can do leptogenesis in the style of electroweak baryogenesis!

[Cohen, Kaplan, & Nelson (1990)]

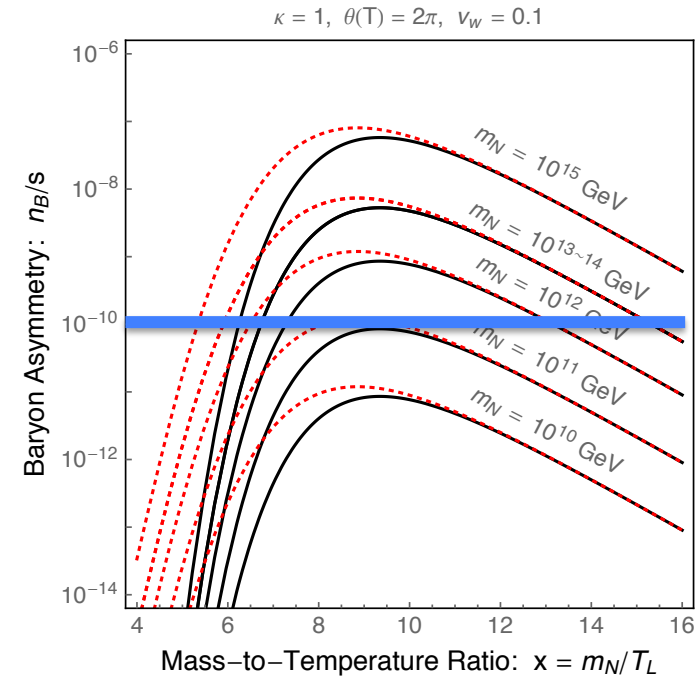
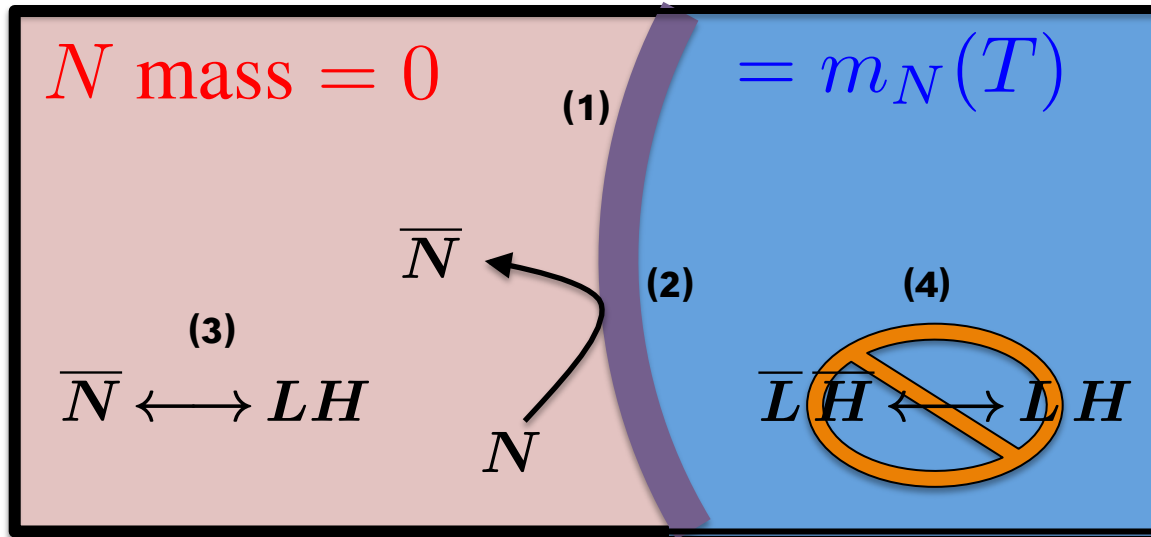
# Leptogenesis from a L-number breaking phase transition

[AL, Tesi, Wang (2016)]

see also: [Pascoli, Turner, & Zhou (2016)]

The phase transition induces a mass for  $N$

CP-violating scattering of  $N$ 's at the bubble wall leads to an L-asymmetry



- (1) First order phase transition
- (2) CP-violating source
- (3) Charge transport & diffusion
- (4) Washout avoidance
- (5) Electroweak sphaleron converts  $L \rightarrow B$

Washout avoidance

requires  $m_N \gg T$

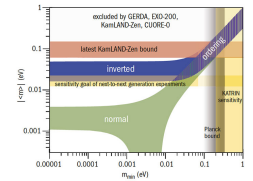
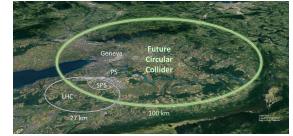
→ large couplings

→ or strong supercooling

# Observational prospects are challenging

## Particle physics

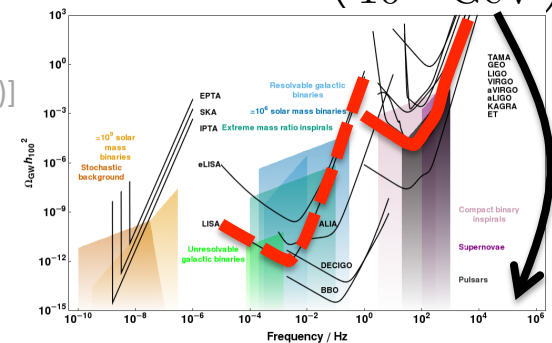
- Predicts: no collider signature of  $m \sim 10^{12}$  GeV particles!
- Predicts: light neutrinos are Majorana &  $0\nu\beta\beta$  may occur.  
(but this is not really probing the physics of baryogenesis)
- Predicts: light pseudo-NGB *majoron* ... constrained by astrophysics



## Cosmological relics

- Predicts: stochastic gravitational waves background
- Predicts: cosmic string network & GW signal  
see also: [Dror, Hiramatsu, Kohri, Murayama, & White (2019)]
- Majorons might survive in the Universe today as candidates for dark matter or dark radiation.

$$f \gtrsim (10^5 \text{ Hz}) \left( \frac{T_{\text{PT}}}{10^{11} \text{ GeV}} \right)$$



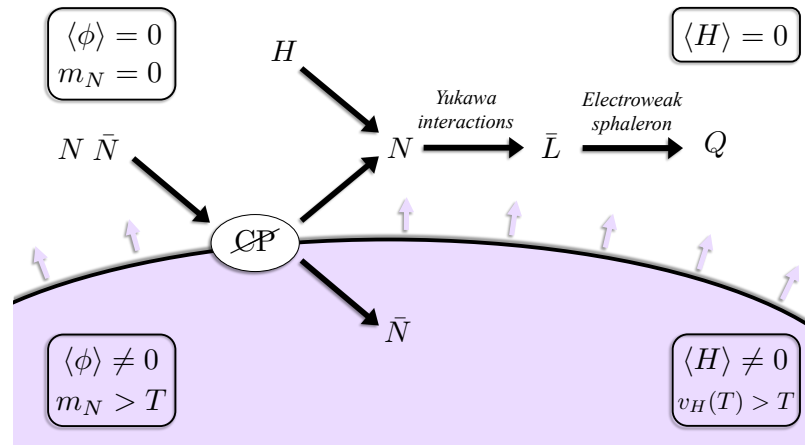
The tiny  $m_\nu$  might instead result from low-scale L-number violation

inverse seesaw: 
$$m_\nu \sim (0.1 \text{ eV}) \left( \frac{\mu_{\text{L-violation}}}{0.01 \text{ keV}} \right) \left( \frac{\kappa\phi}{100 \text{ GeV}} \right)^2 \left( \frac{M_{\text{pseudo-Dirac}}}{1 \text{ TeV}} \right)^{-2}$$

Consider the theory

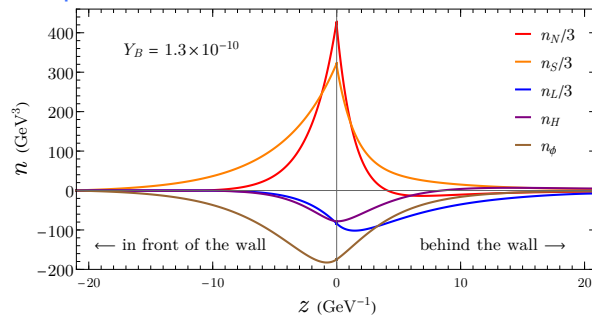
SM + complex scalar  $\phi$  + Weyl fermion  $N$  + **Weyl fermion  $S$**

$$\mathcal{L} = \mathcal{L}_{\text{SM}} - \kappa\phi SN - \lambda_N LHN - \frac{1}{2}\mu SS + \text{h.c.}$$

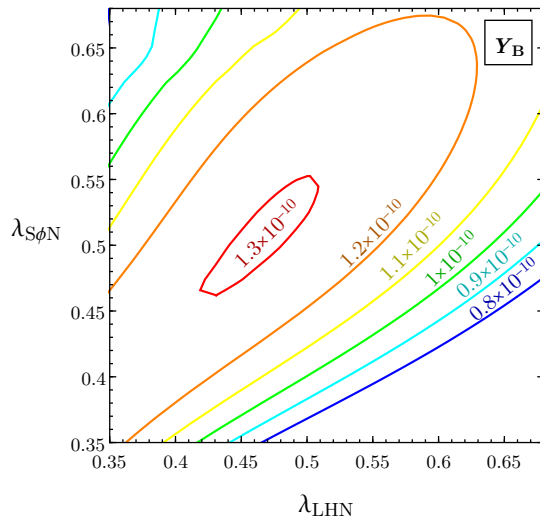


## Baryogenesis

density profiles ...

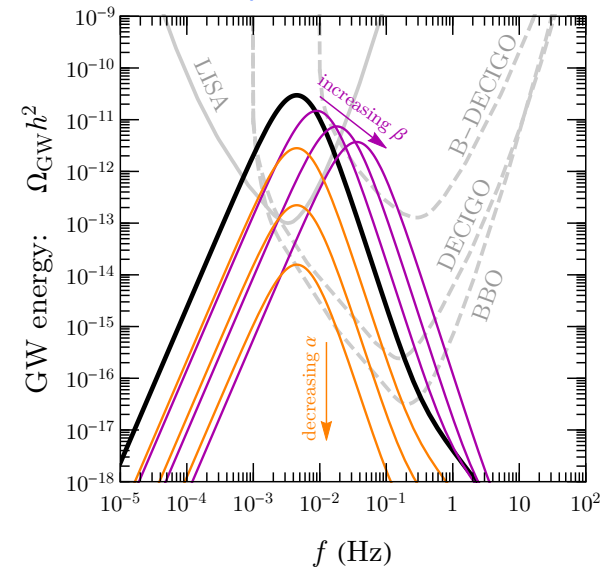


predicted baryon asymmetry ...



## Gravitational Waves

predicted GW spectrum ...



# outline

(1) testing electroweak baryogenesis  
on the collider & cosmological frontiers

(2) a new perspective on  
baryogenesis from (lepto-)bubbles

- The essence\* of electroweak baryogenesis transcends the SM Higgs.  
(\*CP-violating scattering & charge transport at bubble walls)
- An interesting direction to explore: EWBG-like dynamics at other symmetry-breaking phase transitions. E.g., lepton-number, flavor, dark sector, ...

# outline

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from magnetic helicity



# Let's go back to SM baryon-number violation

The chiral EW interactions lead to anomalies in B and L

$$\partial_\mu j_B^\mu = \partial_\mu j_L^\mu = 3 \times \left( \underbrace{\frac{g^2}{32\pi^2} W_{\mu\nu}^a \widetilde{W}^{a\mu\nu}}_{\text{(EW instanton \& sphaleron)}} - \frac{g'^2}{32\pi^2} Y_{\mu\nu} \widetilde{Y}^{\mu\nu} \right)$$

Can we use the  $U(1)_Y$  term to do baryogenesis also?

It vanishes in vacuum but becomes nonzero when a  $B_Y$  field is present.

# Changing $U(1)_Y$ helicity sources B-number

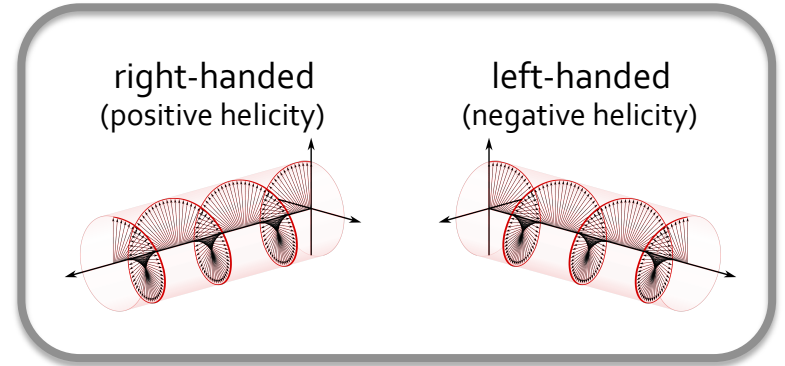
[Joyce & Shaposhnikov (1997); Giovannini & Shaposhnikov (1997); Fujita & Kamada (2016)]

A helical hypermagnetic field provides support for the  $U(1)_Y$  term

$$\langle Y_{\mu\nu} \tilde{Y}^{\mu\nu} \rangle = -4 \mathbf{E}_Y \cdot \mathbf{B}_Y = 2 \left[ \partial_t (\mathbf{A} \cdot \mathbf{B}) + \nabla \cdot (\phi \mathbf{B} + \mathbf{E} \times \mathbf{A}) \right]$$

Changing helicity induces B-number

$$\Delta Q_B = -3 \frac{g'^2}{16\pi^2} \Delta \mathcal{H}_Y$$



The helicity decays in a plasma because of ohmic losses (finite  $\sigma_Y$ )

$$\mathbf{E}_Y = \mathbf{j}_Y / \sigma_Y \approx \nabla \times \mathbf{B}_Y / \sigma_Y$$

$$\langle Y_{\mu\nu} \tilde{Y}^{\mu\nu} \rangle = -4 \langle \mathbf{B}_Y \cdot \nabla \times \mathbf{B}_Y \rangle / \sigma_Y$$

Decaying magnetic helicity provides an avenue for baryogenesis!

# Baryogenesis from decaying magnetic helicity

[Kamada & AL (2016a,b)]

## Recipe for baryogenesis

- (1) Make a helical magnetic field (e.g., axion inflation)
- (2) Universe reheats above the EW scale
- (3) B-field evolution given by MHD (inverse cascade)
- (4) Decaying helicity sources B-number (EW anomaly)
- (5) Solve transport equations to calculate baryon asymmetry

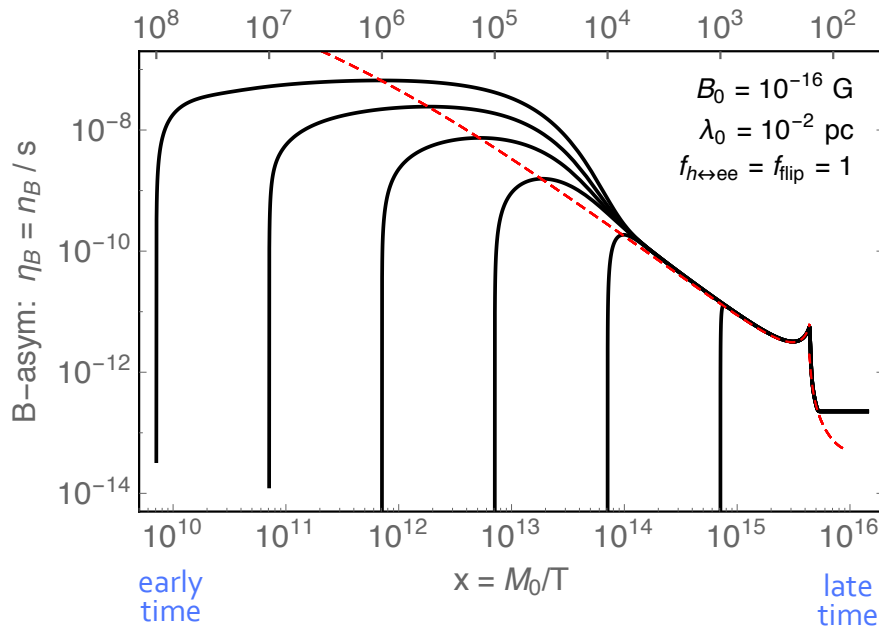
basic idea:

$$\dot{n}_B = -\Gamma_{\text{sphaleron}} n_B + S_{\text{helicity}}$$

➔  $n_B \sim S_{\text{helicity}} / \Gamma_{\text{sphaleron}}$

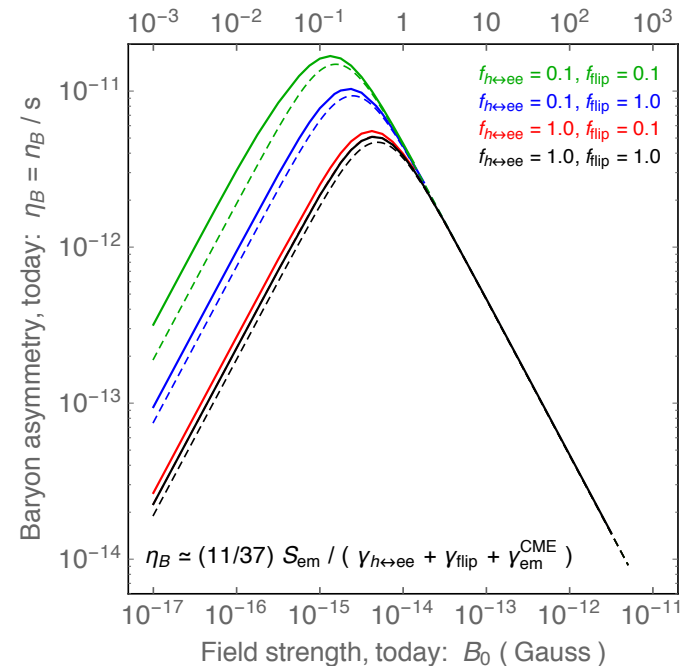
## time evolution

Temperature:  $T$  (GeV)



## relic baryon asymmetry

Coherence length, today:  $\lambda_0$  (pc)

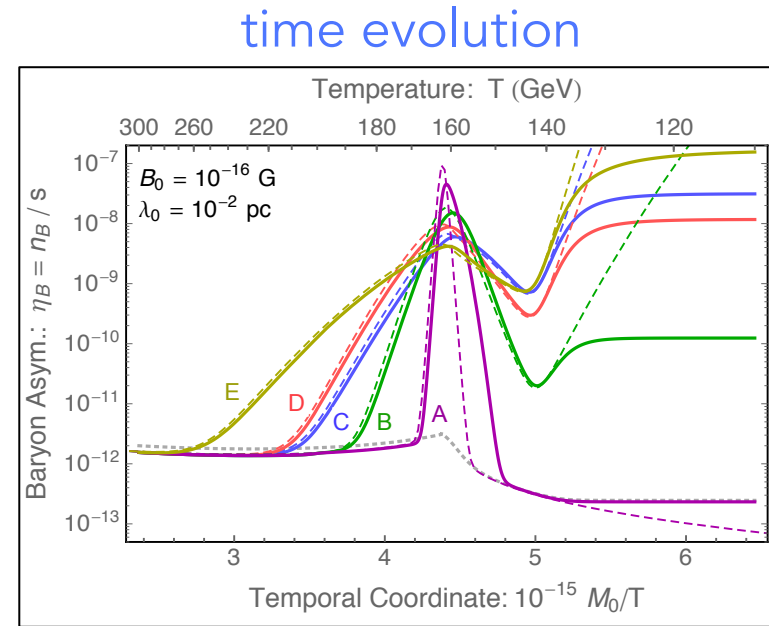
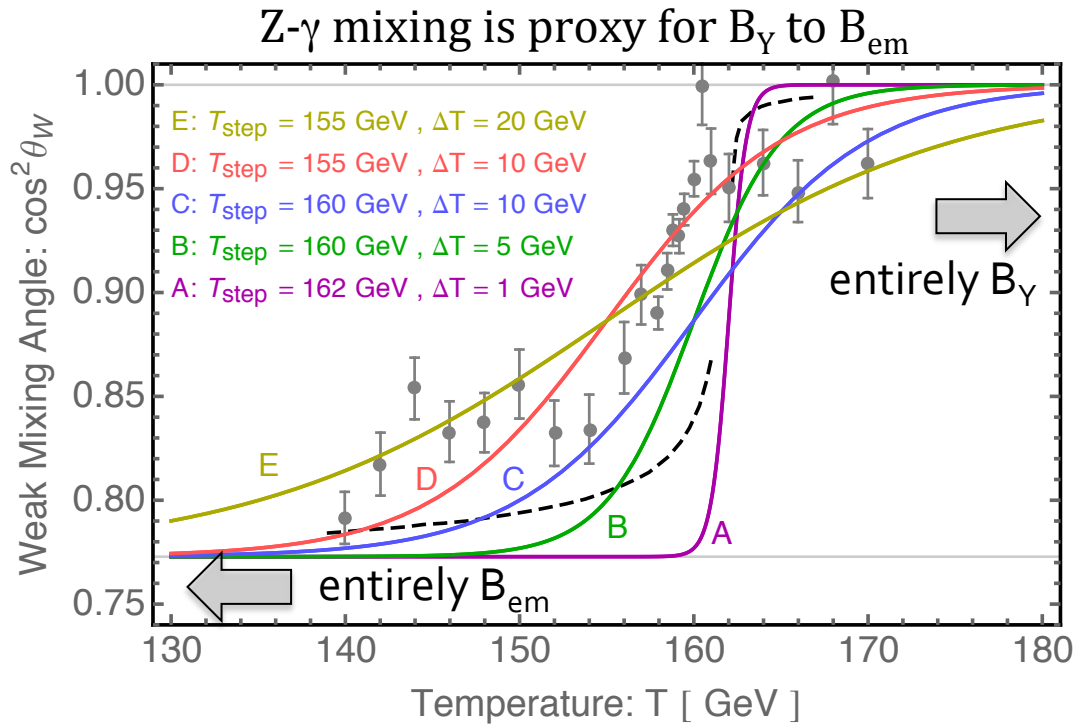


for additional details, see Kohei Kamada's talk on Tuesday afternoon  
 see also: Jimenez, Kamada, Schmitz, & Xu (2017)

# Modeling the field evolution during EW phase transition

[Kamada & AL (2016a,b)]

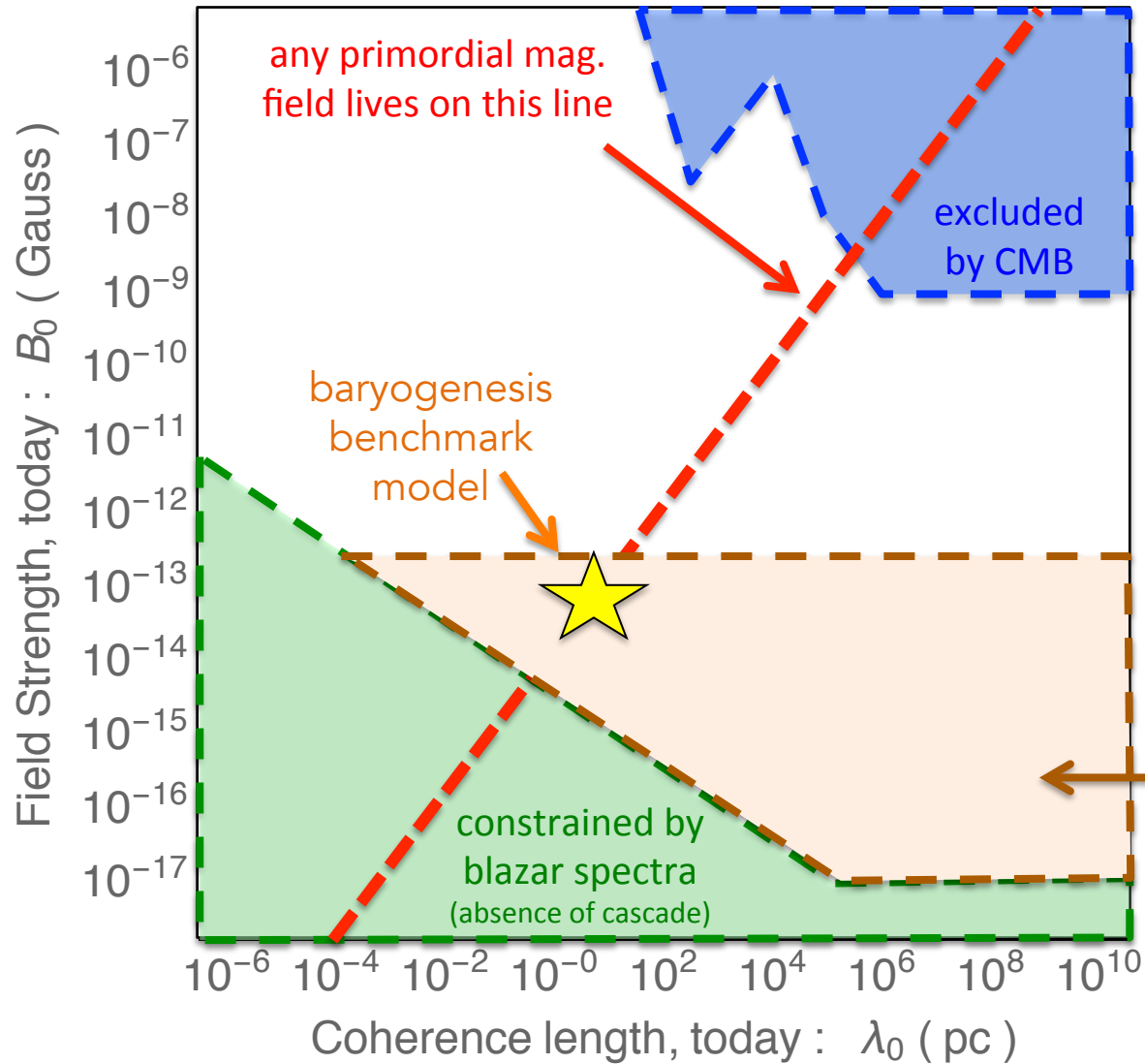
Predicted asymmetry depends on how you model  $U(1)_Y \rightarrow U(1)_{EM}$



for additional details, see Kohei Kamada's talk on Tuesday afternoon

# Testing BG-from-PMF

adapted from: [Durrer & Neronov (2013)]



The relic of the primordial magnetic field is an intergalactic magnetic field today.

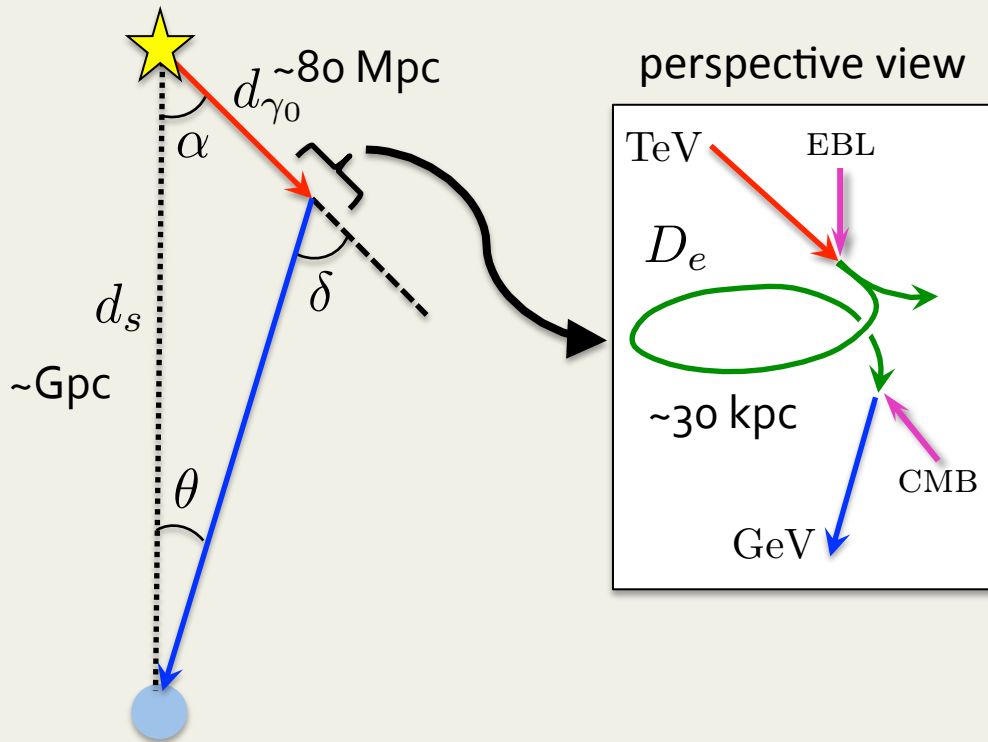
probed by measurements of TeV blazars

# Detecting IGMF with TeV blazars

[Aharonian, Coppi, & Voelk (1994); Neronov & Semikoz (2006); AL & Vachaspati (2016)]

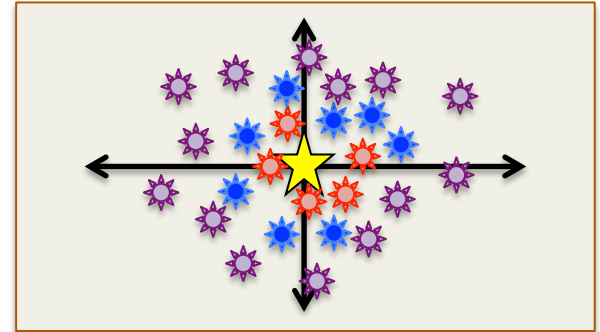
Gamma rays from TeV blazars develop an electromagnetic cascade by scattering on starlight (EBL) and cosmic microwave background (CMB) photons.

The presence of an IGMF deflects the cascade.

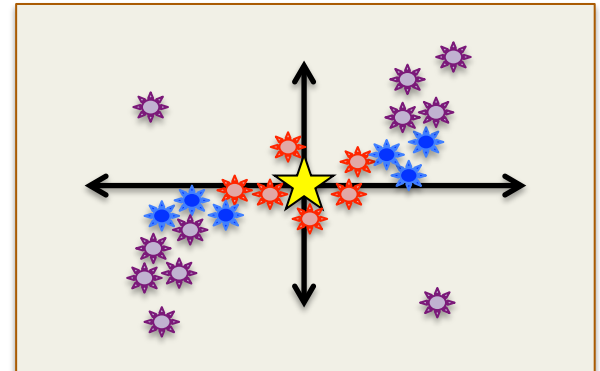


$$\Theta_{\text{ext}} \simeq (0.68^\circ) \left( \frac{B_0}{10^{-15} \text{ G}} \right) \left( \frac{E_\gamma}{10 \text{ GeV}} \right)^{-3/2} \left( \frac{d_s}{1 \text{ Gpc}} \right)^{-1}$$

the blazar acquires a halo

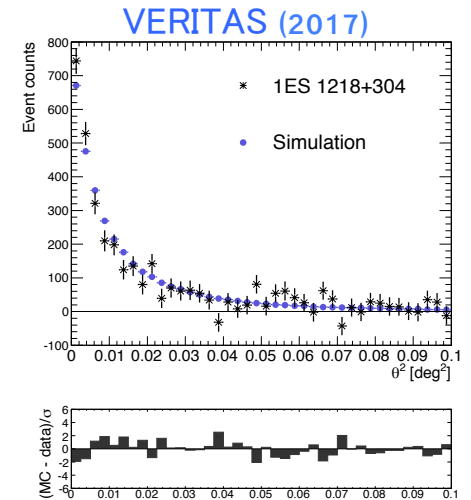
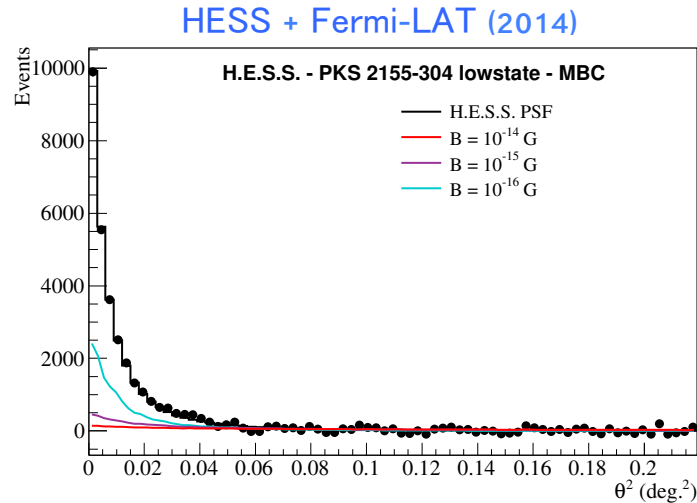
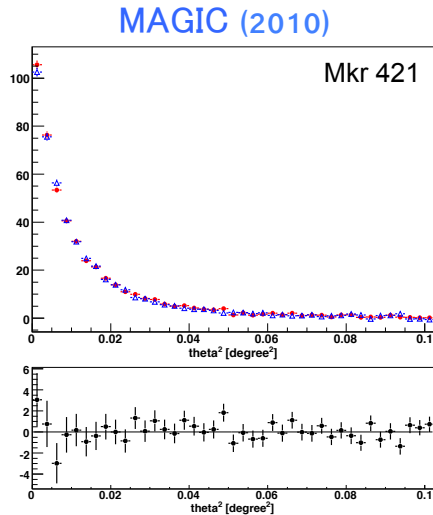


the blazar halo becomes "twisted" by a *helical* IGMF



# Prospects for finding an intergalactic magnetic field

## Ongoing experimental efforts ...



A halo is not observed → some of the IGMF parameter space is excluded:

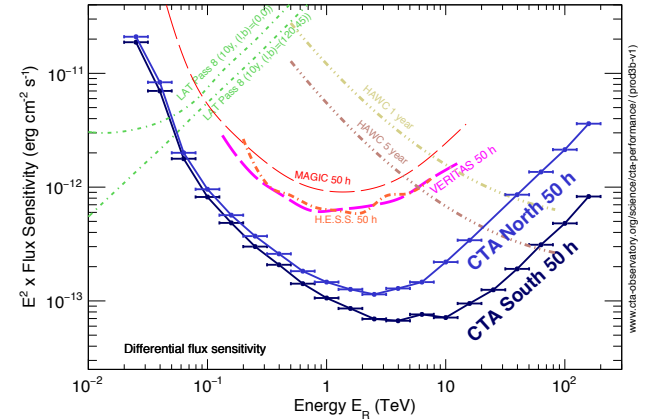
$$B_0 \sim (0.3 - 70) \times 10^{-15} \text{ G} \quad \text{for} \quad \lambda_0 = 1 \text{ Mpc}$$

# Prospects for finding an intergalactic magnetic field

[Meyer, Conrad, & Dickinson (2016); CTA Science Book (2017)]

On the horizon...

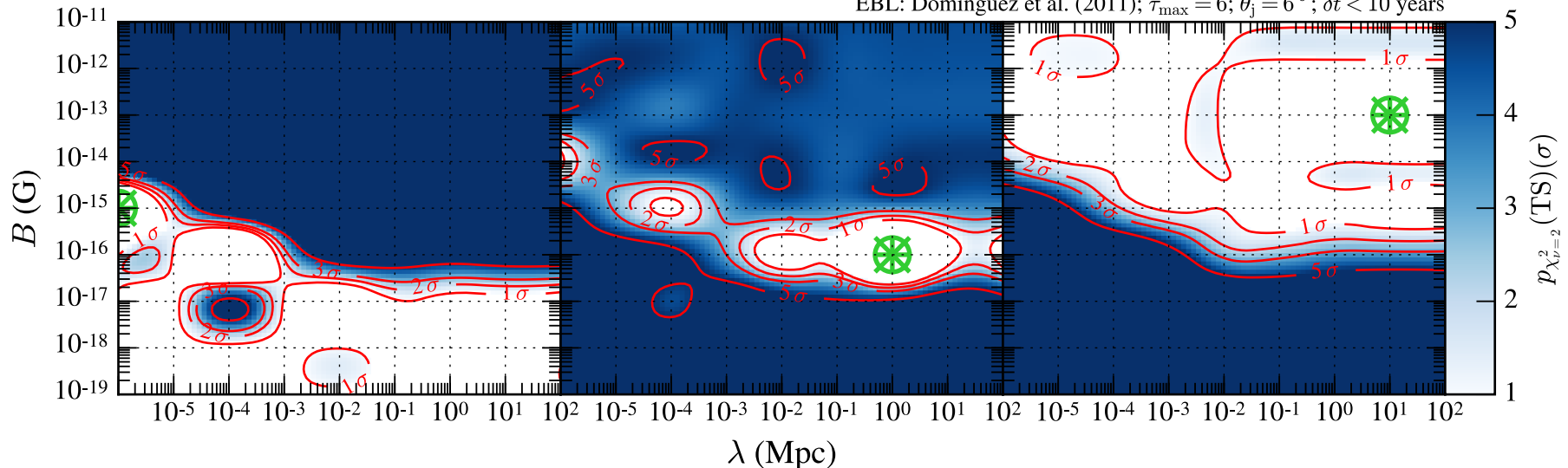
The Cherenkov Telescope Array (CTA) will dramatically improve the flux sensitivity & angular resolution (PSF).



Implications for IGMF constraints:

Benchmark models (3, green). Corresponding exclusions (blue).

EBL: Dominguez et al. (2011);  $\tau_{\text{max}} = 6$ ;  $\theta_j = 6^\circ$ ;  $\delta t < 10$  years





# outline

(1) testing electroweak baryogenesis  
on the collider & cosmological frontiers

(2) a new perspective on  
baryogenesis from (lepto-)bubbles

(3) baryon asymmetry  
from magnetic helicity

- In addition to the EW instanton & sphaleron, the SM predicts B-number violation through changing hypermagnetic helicity as well.
- Lots of interesting directions: How did the helical  $U(1)_Y$  field arise in the first place? How do these fields evolve through the EW crossover? What's the best way to detect their presence in the universe today?

# My Perspective on Baryogenesis

If the EW phase transition is first order, as required by electroweak baryogenesis, then LHC upgrades & future colliders are well-suited to find the new physics.

The essence of EW baryogenesis lives on by clever application of the EWBG tool kit to other cosmological phase transitions.

The SM predicts B-number violation from changing helicity of hypermagnetic fields, which provides a new avenue for BG that can be tested with high-energy astroparticle observables.

# New Perspectives on Baryogenesis

<b>Neutrino Option Leptogenesis</b>	<i>Jessica Turner</i>
<i>Grey Room 2, IFT</i>	16:40 - 17:00
<b>Type-I Seesaw as the Common Origin of Neutrino Mass, Baryon Asymmetry, and the Electroweak Scale</b>	<i>Kai Ruven Schmitz</i>
<b>Preheating confronts leptogenesis</b>	<i>Djuna Croon</i>
<i>Grey Room 2, IFT</i>	17:20 - 17:40
<b>Baryogenesis from Primordial Helical Hypermagnetic Fields</b>	<i>Kohei Kamada</i>
<i>Red Room, IFT</i>	17:10 - 17:30
<b>The baryon asymmetry from a composite Higgs</b>	<i>Oleksii Matsedonskyi</i>
<i>Red Room, IFT</i>	17:30 - 17:50
<b>High Scale Electroweak Baryogenesis</b>	<i>Dr Iason Baldes</i>
<i>Red Room, IFT</i>	17:50 - 18:10
<b>Parameter space of baryogenesis in the vMSM</b>	<i>Inar Timiryasov</i>
<i>Grey Room 2, IFT</i>	14:15 - 14:35
<b>Low-scale leptogenesis with three heavy neutrinos</b>	<i>Juraj Klaric</i>
<i>Grey Room 2, IFT</i>	14:35 - 14:55
<b>Leptogenesis in the Scotogenic model</b>	<i>Vedran Brdar</i>
<i>Grey Room 2, IFT</i>	14:55 - 15:15
<b>Baryogenesis and Dark Matter from B Mesons</b>	<i>Miguel Escudero</i>
<i>Grey Room 2, IFT</i>	15:40 - 16:00
<b>All-in-one relaxation</b>	<i>Jakub Scholtz</i>
<i>Grey Room 2, IFT</i>	16:00 - 16:20
<b>Leptogenesis without Loops</b>	<i>Amab Dasgupta</i>
<i>Grey Room 2, IFT</i>	16:20 - 16:40