



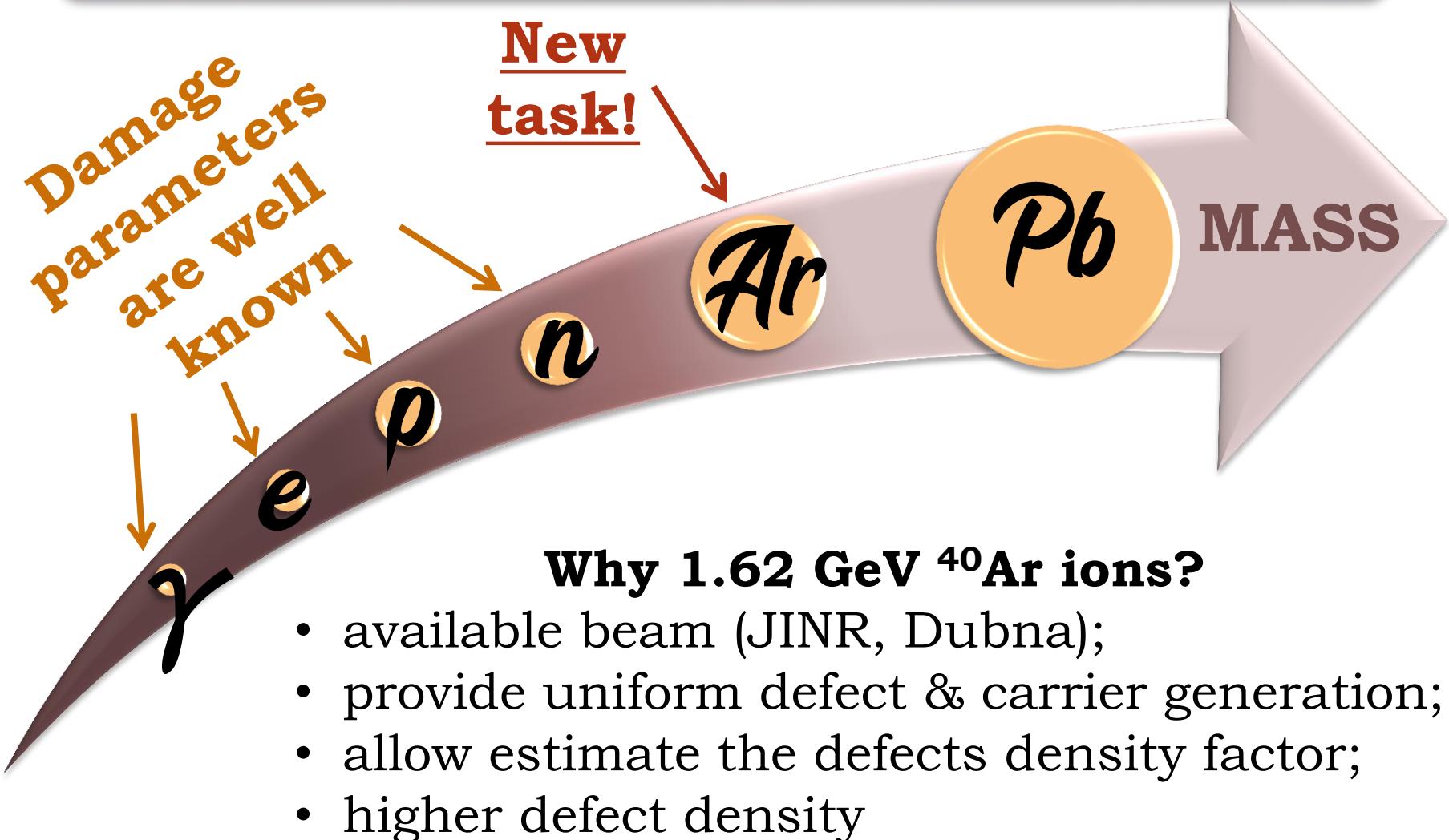
Comparative analysis of proton and ion damages in Si detectors supplemented with SRIM simulations

D. Mitina, V. Eremin, E. Verbitskaya

*Ioffe Institute
St. Petersburg, Russia*

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Motivation



Goal

- To find macroscopic parameters (from I-V, I-T, TCT)
- To find microscopic parameters (from DLTS)
- To compare obtained damage parameters with parameters for proton irradiation
- To add SRIM as a tool for microscopic damage study

Results partially published in:

V. Eremin, D. Mitina, et al., A comparative study of silicon detector degradation under irradiation by heavy ions and relativistic protons, 2018 *JINST* **13** P01019

Detectors

Structure:

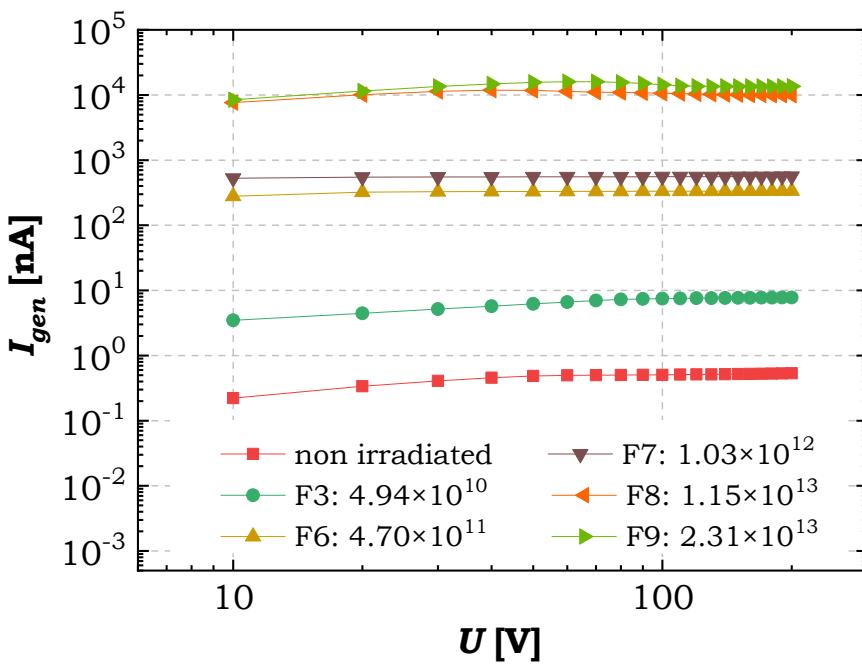
- p -on- n pad FZ detectors; p - n junction area: 25 mm²
- thickness: 300 μm
- resistivity: 10 kΩcm and 50 Ωcm

Irradiation:

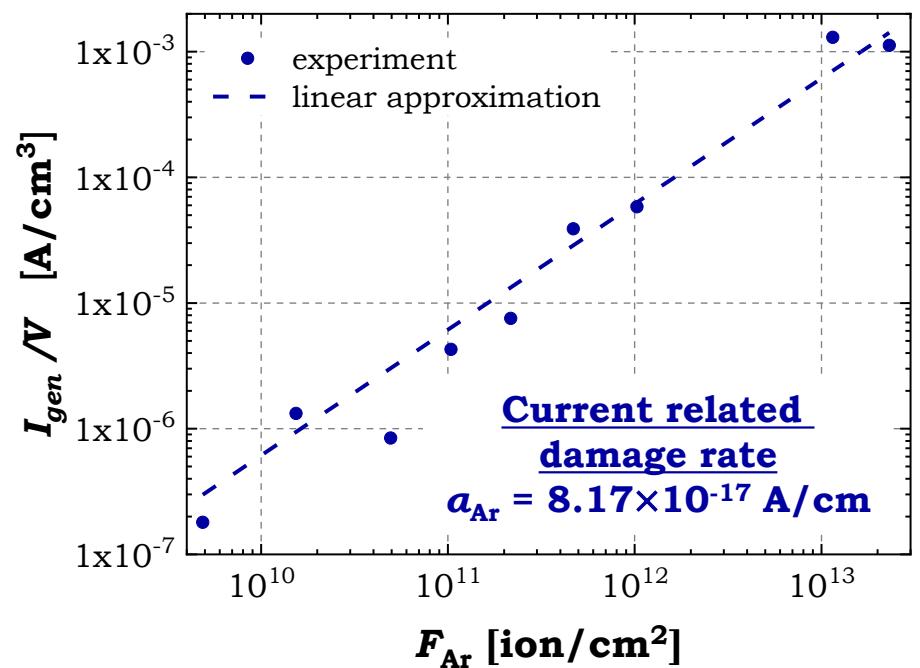
- JINR, Dubna, Russia
- 1.62 GeV ^{40}Ar ions (40.5 MeV/u)
- Beam intensity: 10⁸ to 10¹⁰ ion/cm² per second
- Fluence: 5×10^{10} – 2.3×10^{13} ion/cm²

Bulk Generation Current

I-V characterization for high-resistivity detectors irradiated by ^{40}Ar ions



Fluence dependence of generation current



Temperature Dependences of I_{gen}

According to the
Shockley-Read-Hall
theory:

Approximation was
performed assuming
that

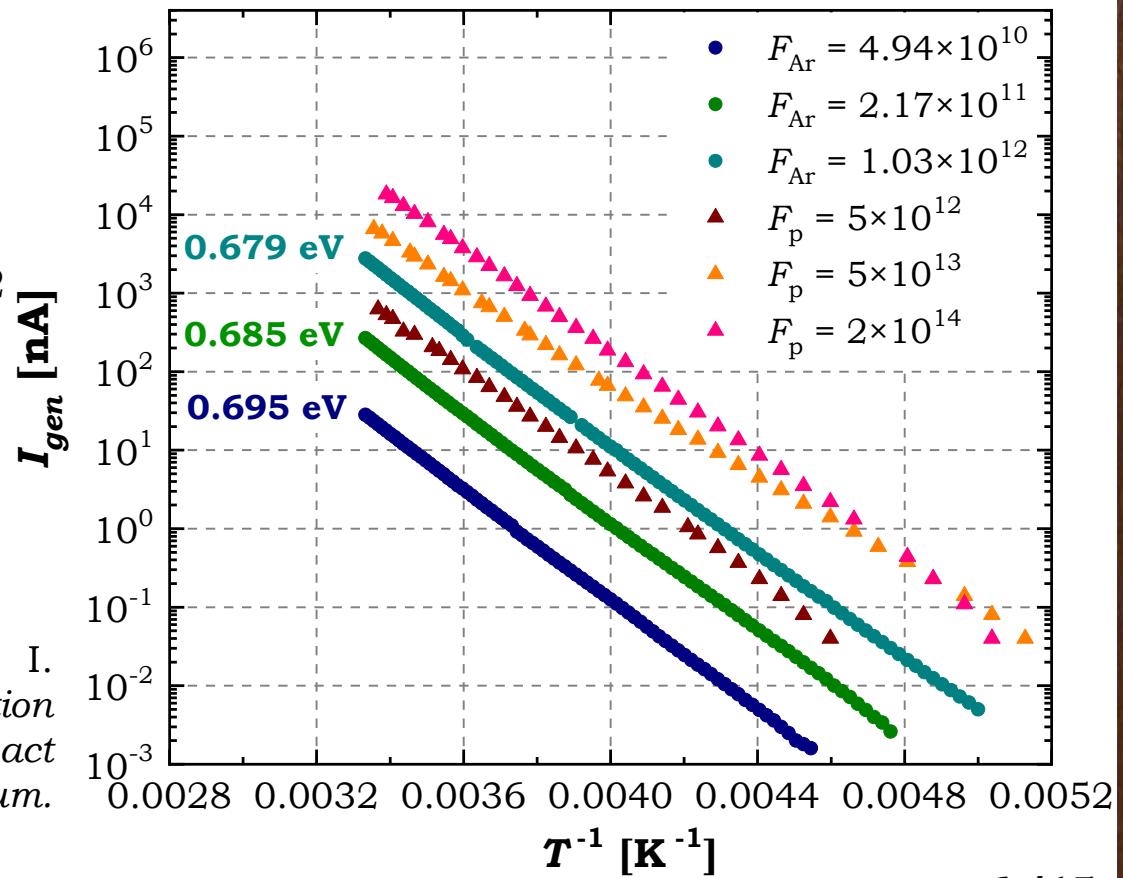
$$\sigma_e = \sigma_p = 8 \times 10^{-14} \text{ cm}^2$$

Activation energies:

protons	^{40}Ar ions
0.65 eV	0.68 eV

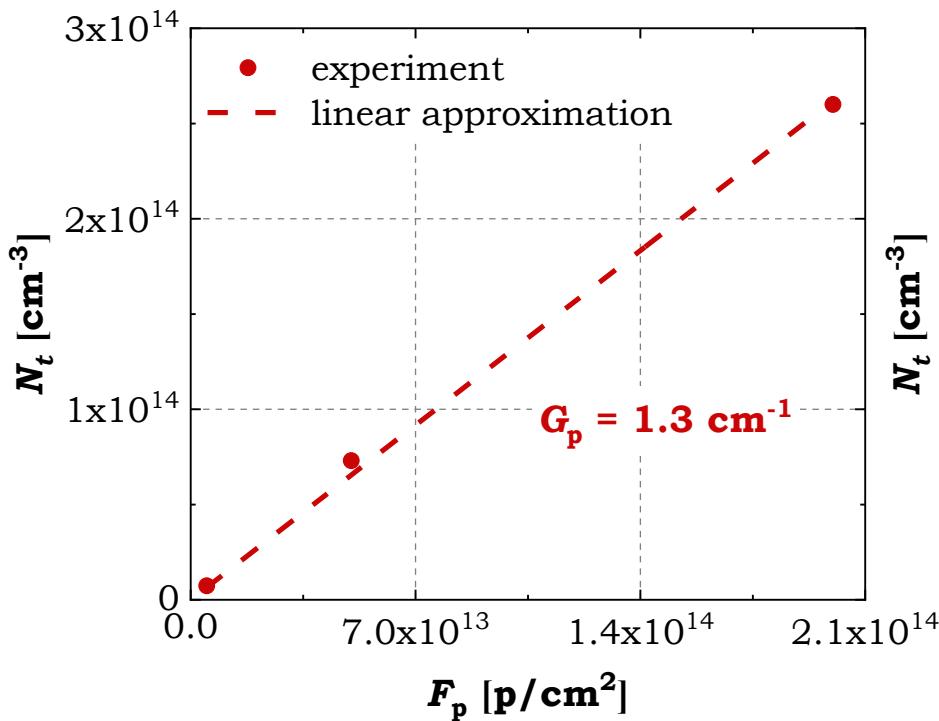
E. Verbitskaya, V. Eremin, I. Ilyashenko and Z. Li, *Carrier generation in irradiated Si detectors and its impact on the electric field profile, Nucl. Instrum. Meth. A* **754** (2014) 63

$$I_{gen} \approx \frac{\sigma_e \sigma_h N_t}{\sigma_e \exp\left(-\frac{E_g - E_t}{k_B T}\right) + \sigma_h \exp\left(-\frac{E_t}{k_B T}\right)}$$

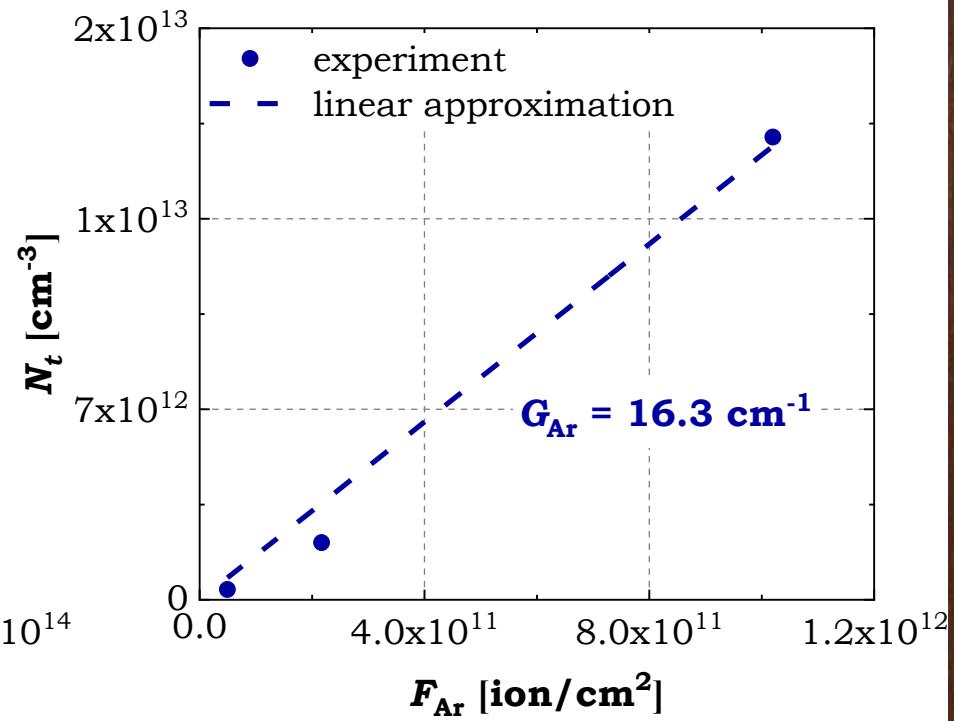


Introduction Rates

protons



⁴⁰Ar ions



E. Verbitskaya, V. Eremin, I. Illyashenko and Z. Li, *Carrier generation in irradiated Si detectors and its impact on the electric field profile*, Nucl. Instrum. Meth. **A 754** (2014) 63

D. Mitina et al., 33 RD50 Workshop, CERN, Geneva, Nov 26-28, 2018

TCT Data

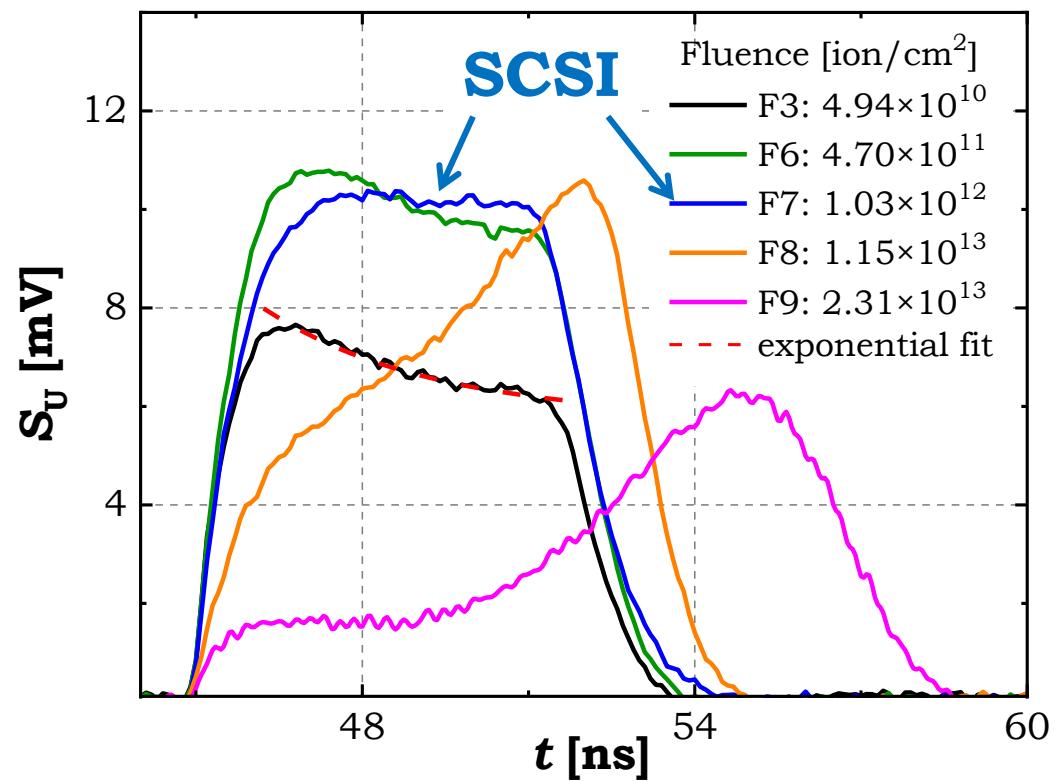
Current response,
generated by electrons:

$$i(t, T) = \frac{Q_0}{d} \mu(T) E_0 \exp\left(-\frac{t}{\tau_{dr}}\right)$$

$$\tau_{dr} = \frac{\epsilon \epsilon_0}{\epsilon \mu(T) N_{eff}}$$

Last equation allows one
to calculate N_{eff} at a fixed
mobility via
approximation of the
experimental current
pulse top by an
exponential function.

Pulse shapes of detectors
irradiated by ^{40}Ar ions



Effective Concentration

Hamburg model:

$$N_{eff}(F) = N_0 e^{-cF} - \beta F$$

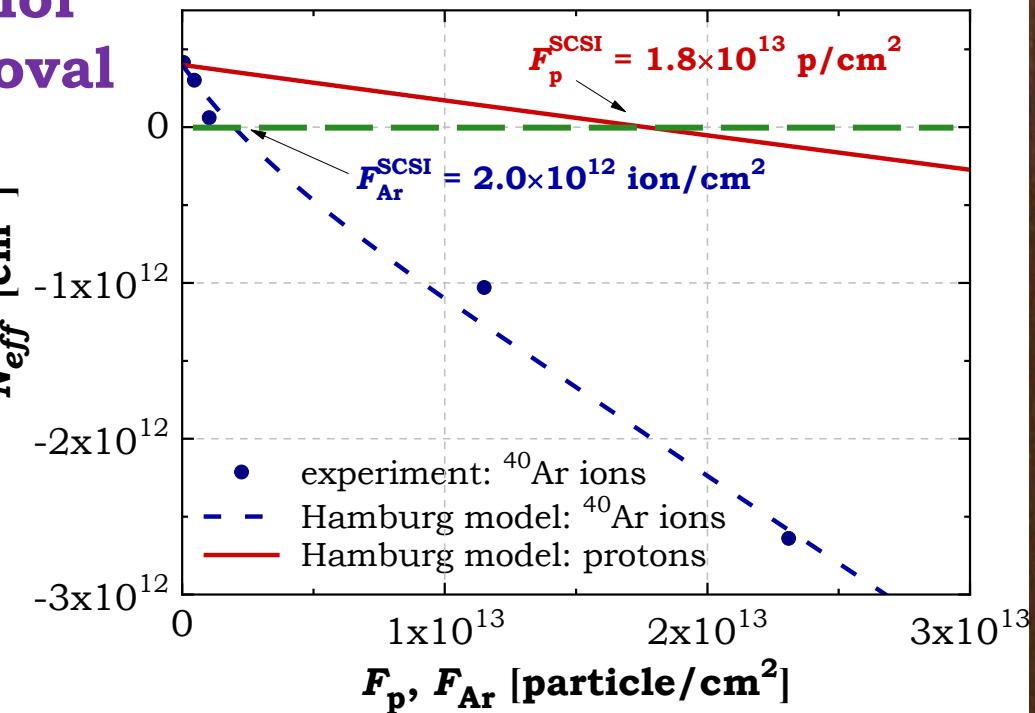
acceptor introduction donor removal

Approximation results:

	^{40}Ar	p
c [cm ⁻²]	3×10^{-13}	1×10^{-14}
β [cm ⁻¹]	0.11	0.019

G. Kramberger, *Signal development in irradiated silicon detectors*, Ph.D. Thesis, University of Ljubljana, Ljubljana Slovenia (2001), CERN-THESIS-2001-038

Effective concentration dependences on fluence



Deep Level Transient Spectroscopy

DLTS spectra for 50 Ωcm irradiated detectors

Neutrality + Emission

$$U_e = 20 \text{ V}$$

$$U_f = 0$$

$$t_e = 50 \text{ ms}$$

$$t_p = 5 \text{ ms}$$

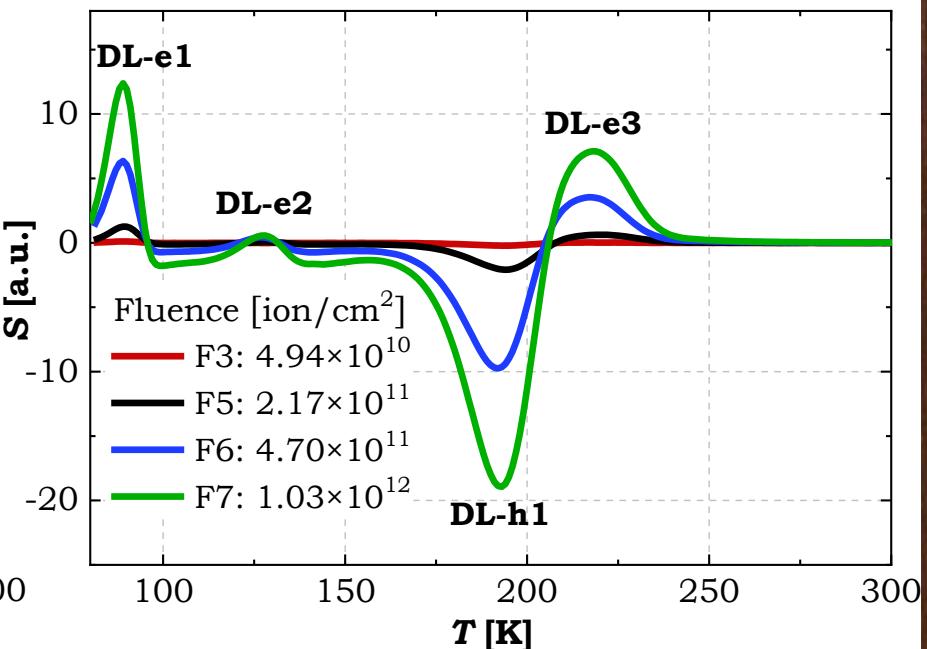
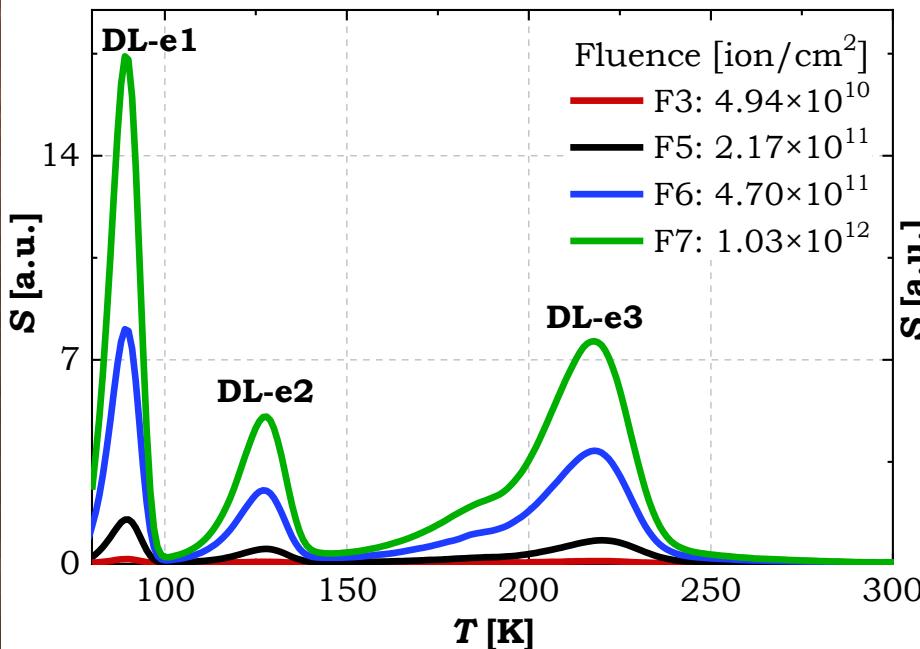
Injection + Emission

$$U_e = 20 \text{ V}$$

$$U_f = -10 \text{ V}$$

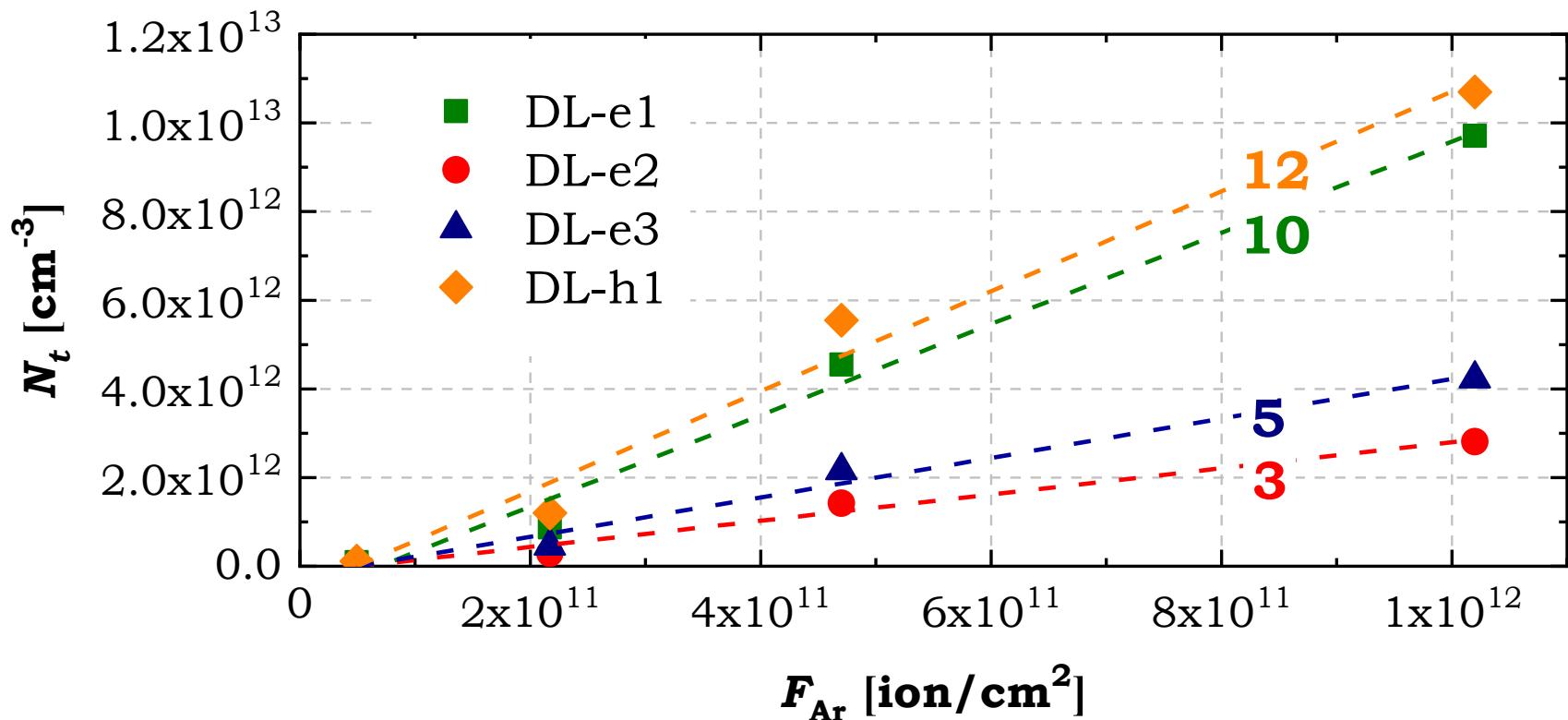
$$t_e = 50 \text{ ms}$$

$$t_p = 5 \text{ ms}$$



Trap Parameters

	DL-e1	DL-e2	DL-e3	DL-h1
E_t [eV]	0.168	0.240	0.415	0.366
σ_t [cm ²]	4.28×10^{-15}	3.31×10^{-15}	1.74×10^{-15}	1.18×10^{-15}



Trap Structure

A comparison of the introduction rates and activation energies for protons and ^{40}Ar ions

	protons		^{40}Ar ions	
	$g \text{ [cm}^{-1}]$	$\langle E_t \rangle \text{ [eV]}$	$g \text{ [cm}^{-1}]$	$\langle E_t \rangle \text{ [eV]}$
DL-e1: V-O	0.73	0.176	10	0.168
DL-e2: VV--	0.37	0.241	3	0.240
DL-e3: VV-	0.37	0.410	5	0.415
DL-h1: C_i-O_i	1.30	0.320	12	0.366

I. Pintilie, *Investigation of radiation defects*, in proceedings of the *3rd MC-PAD Network Training Event on Radiation Hardness and Silicon processing and the Project Midterm Review*, Ljubljana, Slovenia, 26–30 September 2010.

Comparison

	protons	^{40}Ar ions	ratio
$\alpha [\text{A}/\text{cm}]$	0.83×10^{-17}	8.17×10^{-17}	9.81
$G [\text{cm}^{-1}]$	1.3	16.3	12.5
F scsi [particle/cm ²]	1.8×10^{13}	2.0×10^{12}	9
$g [\text{cm}^{-1}]$	$V-O$	0.73	10
	VV^{--}	0.37	3
	VV^-	0.37	5
	C_i-O_i	1.30	12
			9.23



SRIM simulation

	50 MeV protons	1.62 GeV ^{40}Ar ions
collisions	8	132
displacements	8	1725
vacancies	1	1594
collision events for 3 incident particles		
<p>COLLISION EVENTS Target Displacements</p> <p>Number / (Angstrom - Ion)</p> <p>Silicon</p> <p>- Target Depth -</p> <p>0 A 300 um</p> <p>Y-axis scale: 0 to 45×10^{-6} in increments of 5×10^{-6}.</p> <p>The plot shows three distinct peaks at approximately 0, 150, and 280 micrometers depth.</p>		
<p>COLLISION EVENTS Target Displacements</p> <p>Number / (Angstrom - Ion)</p> <p>Silicon</p> <p>- Target Depth -</p> <p>0 A 300 um</p> <p>Y-axis scale: 0 to .016 in increments of .002.</p> <p>The plot shows a very broad distribution with a primary peak at ~10 micrometers and smaller peaks extending to 300 micrometers.</p>		

Summary

- The activation energy of the bulk generation current in ^{40}Ar irradiated detectors is only 5% higher than the corresponding energy in the case of proton irradiation;
- The increase of ion fluence led to space charge sign inversion at $F \sim 1 \times 10^{12} \text{ ion/cm}^2$;
- DLTS-spectra demonstrated a set of deep levels which are typical of proton irradiation;
- The difference in effects on silicon detectors of hadrons and heavy ions can be scaled by the Scaling Coefficient parameter;
- SRIM simulation demonstrated that the divacancy production rate is significantly lower for ^{40}Ar irradiation than for protons.



*Thank you
for attention!*