



*Characterization of silicon  $n^+ - p - p^+$  detectors  
with  $Al_2O_3$  passivating layers  
grown by Atomic Layer Deposition method*

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# *Al<sub>2</sub>O<sub>3</sub> layers processed by ALD*

Atomic Layer Deposition (ALD) Al<sub>2</sub>O<sub>3</sub> layers – suggested for surface passivation in devices on p-Si with low and high Si resistivity

1. B. Hoex, et al., Appl. Phys. Lett., 91, 112107 (2007)
2. F. Werner, et al., J. Appl. Phys. 109, 113701 (2011)
3. M. Christophersen and B. F. Philips, 2011 IEEE Nucl. Sci. Symp. Conf. Rec., p. 113
4. J. Härkönen, et al., Nucl. Instrum. Meth. A 828 (2016) 46
5. J. Härkönen, et al., Nucl. Instrum. Meth. A 831 (2016) 2

## ALD method of Al<sub>2</sub>O<sub>3</sub> processing:

Low temperatures (200-300°C)

After standard planar technology process

Thin films (tens of nm)

Negative charge,  $Q_f = -(10^{11}-10^{13}) \text{ cm}^{-2}$  (referred data)

– allows to isolate n<sup>+</sup> segments and simplify design and processing of n-on-p detectors

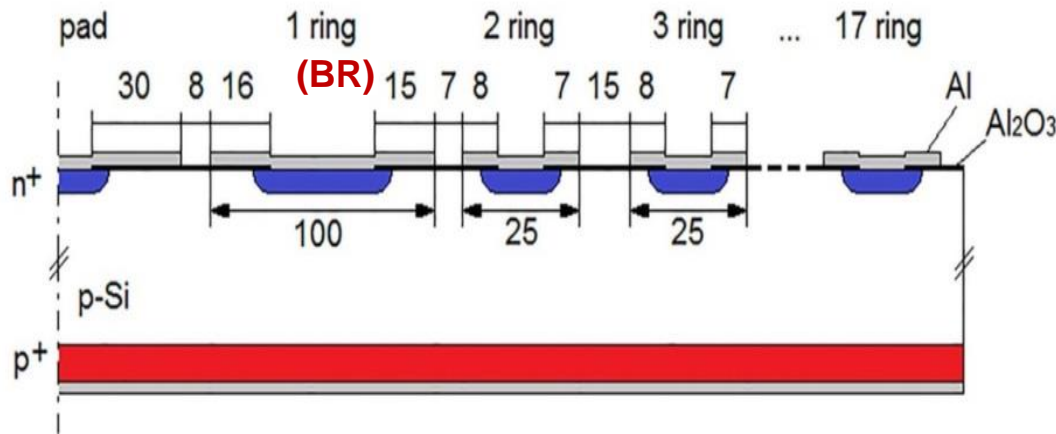
# Goals of the study

**Approach: to combine experimental data and TCAD simulations for characterizing n-on-p Si detectors with alumina isolating layers** 

- Measure the characteristics of n<sup>+</sup>-p-p<sup>+</sup> silicon detectors with Al<sub>2</sub>O<sub>3</sub> isolation films designed for high voltage operation
- Simulate the potentials and electric field in the n-on-p detector with multiple n<sup>+</sup> rings
- Extract the fixed charge density  $Q_f$  in Al<sub>2</sub>O<sub>3</sub> in the detector under study
- Estimate the impact of the charge density  $Q_f$  in Al<sub>2</sub>O<sub>3</sub> on the electric field
- Define the parameters of the detector design allowing to maintain high voltage operation

The results are published in: V. Eremin, et al., 2018 JINST 13 P11009

# Si n<sup>+</sup>-p-p<sup>+</sup> detectors with Voltage Termination Structure (VTS)



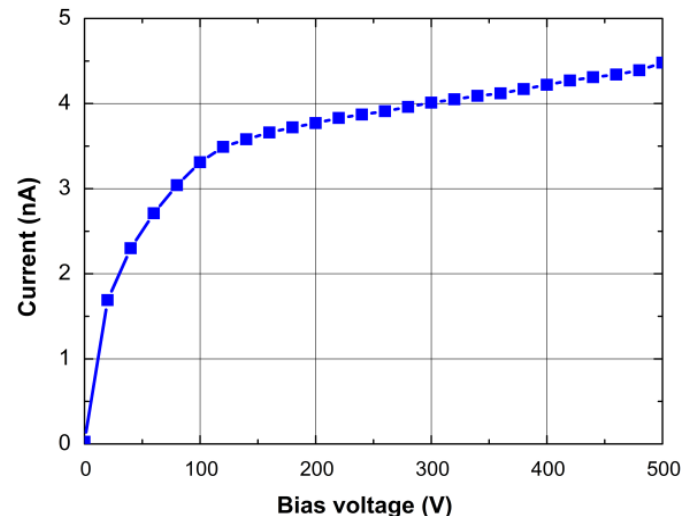
MCz p-type silicon,  $\rho = 7 \text{ k}\Omega\text{cm}$ ,  
 $(N_{\text{eff}} = 1.9 \times 10^{12} \text{ cm}^{-3})$ ,  $d = 300 \text{ }\mu\text{m}$   $V_{\text{fd}} = 125 \text{ V}$   
 active area  $5 \times 5 \text{ mm}^2$

See presentation of Jennifer Ott

- ALD at 200-300°C after ion implantation and annealing
- mean  $\text{Al}_2\text{O}_3$  thickness about 50 nm
- $Q_f$  in the range of  $-(10^{11}-10^{12}) \text{ cm}^{-2}$

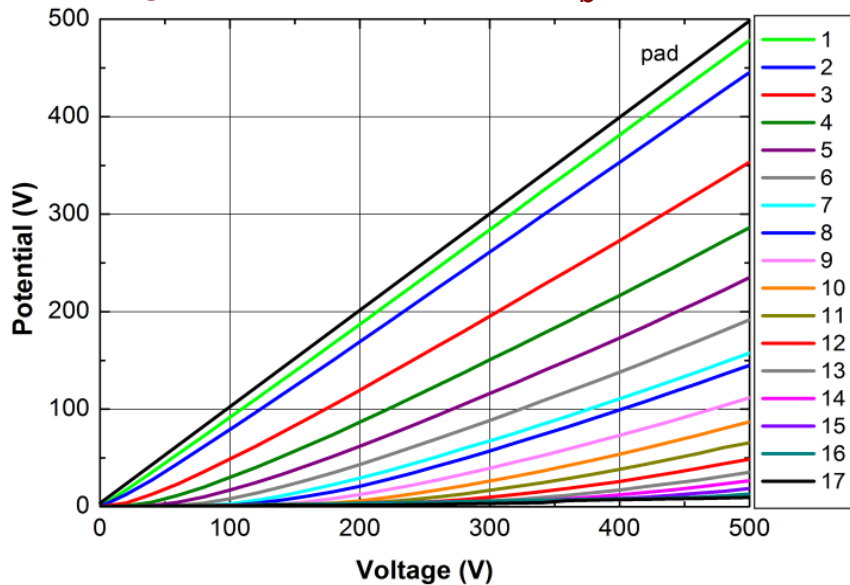
n-on-p detector with **bias ring (BR) + 16 floating guard rings - Voltage Terminating Structure (VTS)** distributing the bias /potential over the rings

## I-V characteristic



# Potential distribution in Si n-on-p detectors with VTS

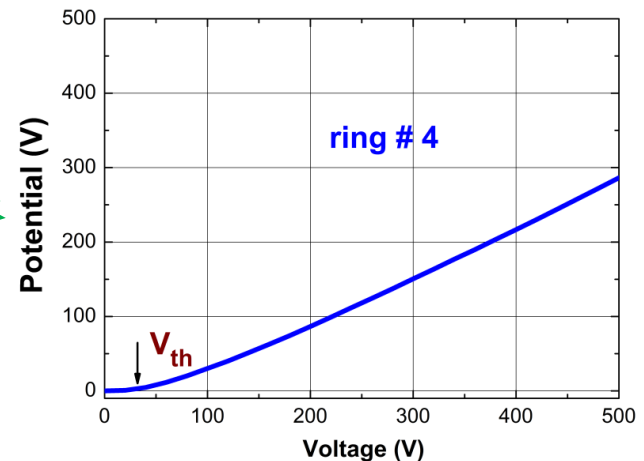
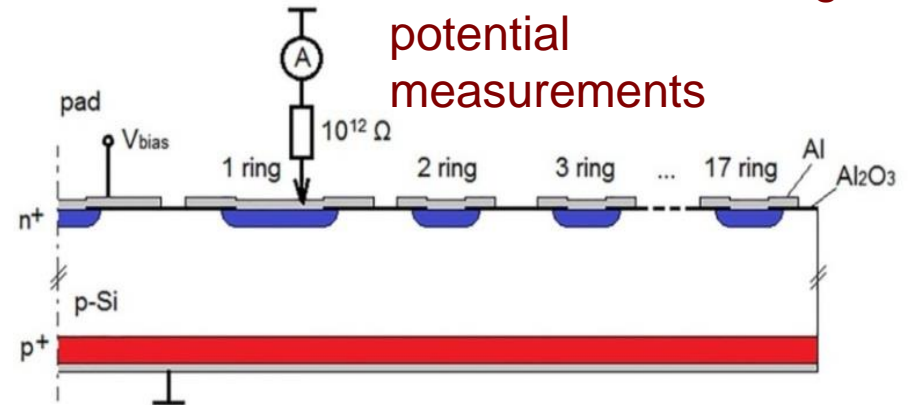
Distribution of potential  $\phi$  over the rings vs. applied bias  $V_b$



$\phi$  vs.  $V_b$  curves show two regions:

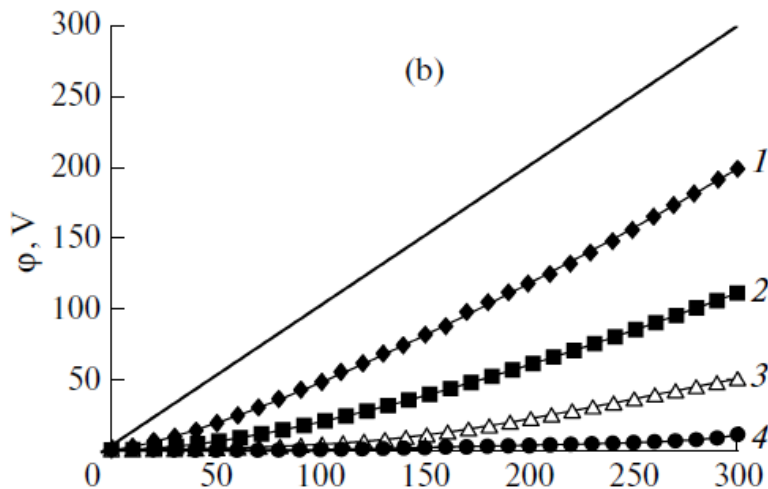
- the flat region with a zero potential,
- region of the gradual potential rise starting at threshold voltage  $V_{th}$ .

Schematic of the ring potential measurements



# Punch-through mechanism of potential distribution over the rings

$\varphi(V_b)$  curves in p-on-n detectors with  $\text{SiO}_2$  isolation (positive charge in  $\text{SiO}_2$ )



V. Eremin et al., *Semiconductors* (2011) 45, No. 4, 536 )

Similar shapes of  $\varphi(V_b)$  curves in p-on-n detectors with  $\text{SiO}_2$  isolation (positive charge in  $\text{SiO}_2$ ) and in n-on-p detectors with  $\text{Al}_2\text{O}_3$  isolation (negative charge in alumina) evidences **the same punch-through mechanism of potential transfer:**

coupling of neighboring rings occurs via carrier injection from inner ring to the outer ring that switches one of the junctions in forward direction

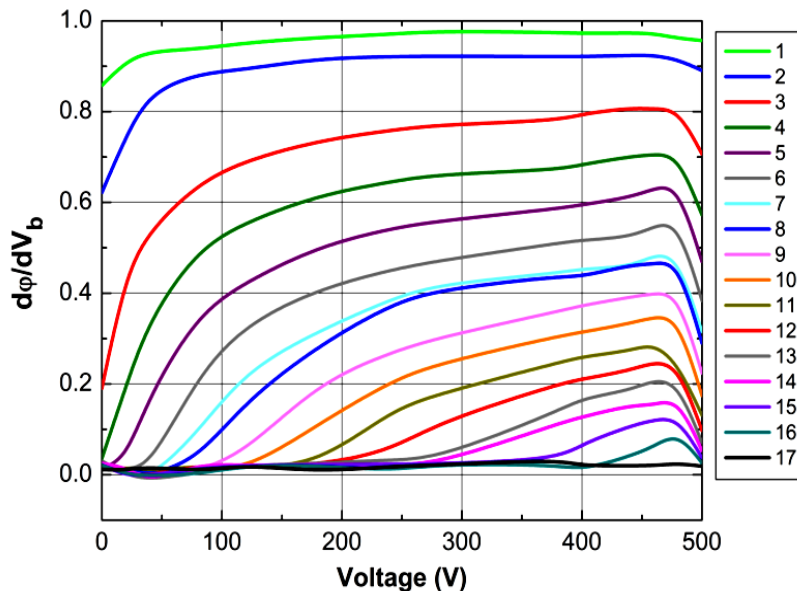
## Difference

p-on-n with  $\text{SiO}_2$  – hole injection from  $p^+$

n-on-p with  $\text{Al}_2\text{O}_3$  – electron injection from  $n^+$

# Rate of potential growth

$V_{th}$  is visualized in the dependence of derivative  $d\phi/dV_b$  vs.  $V_b \rightarrow$  shows the rate of potential growth



$V_{th}$  is about zero for the rings 1-4;

At  $V_b > V_{th}$  the rate of potential growth is not constant:

first increases and then is low or close to zero  $\rightarrow$  manifests reaching the steady-state potential value;

At  $V_b > 450$  V the rate of the  $\phi$  rise drops.

# Simulations of potentials and 2D electric field

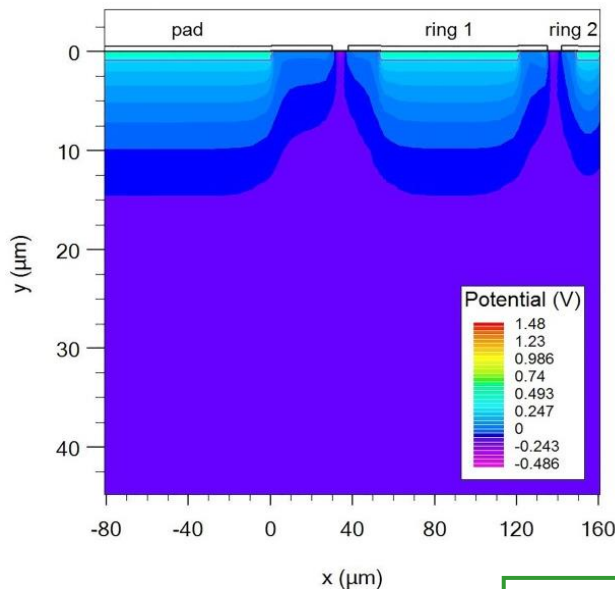
2D Athena and Atlas (Silvaco) TCAD tool

- ◆ basic equations solved via iteration procedure
- ◆ Selberherr model of impact ionization
- ◆ the n<sup>+</sup>-p junction depth of 1 μm
- ◆ Al<sub>2</sub>O<sub>3</sub> thickness of 52 nm
- ◆ Q<sub>f</sub> as a parameter  $-(1-10) \times 10^{11} \text{ cm}^{-2}$

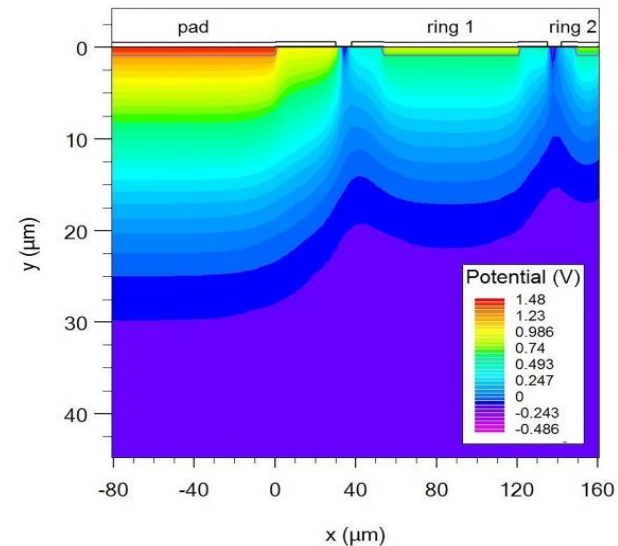
Initial stage of potential ( $\phi$ ) transfer

$$Q_f = -4 \times 10^{11} \text{ cm}^{-2}$$

$V_b = 0 \text{ V}$ , but  $V_{\text{cont}} \neq 0$

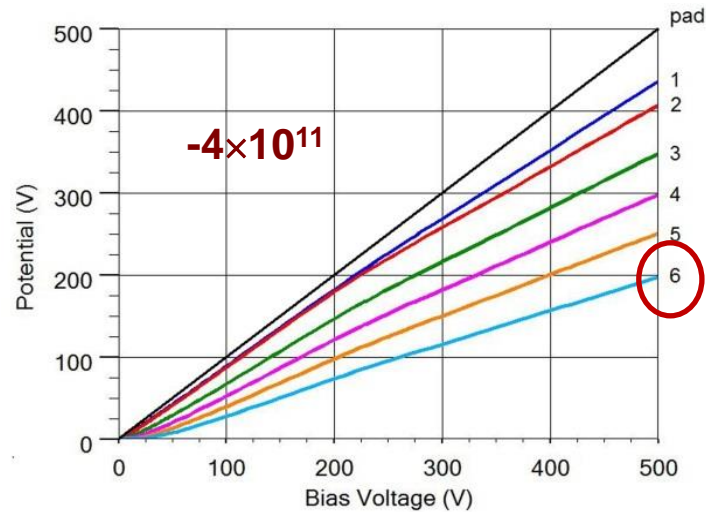
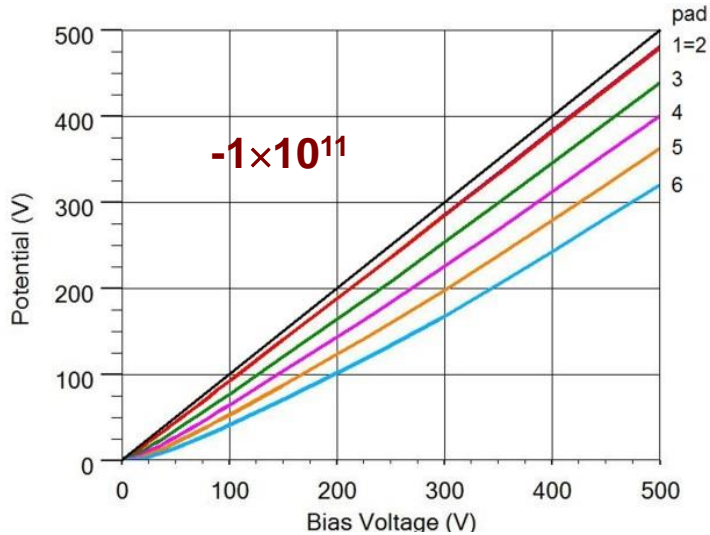


$V_b = 1 \text{ V}$



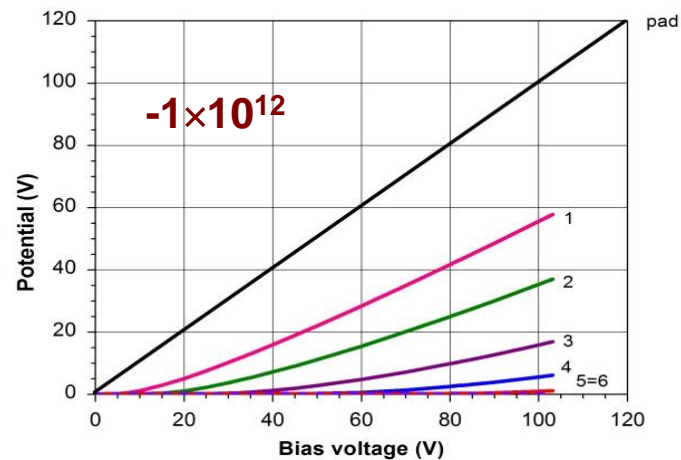
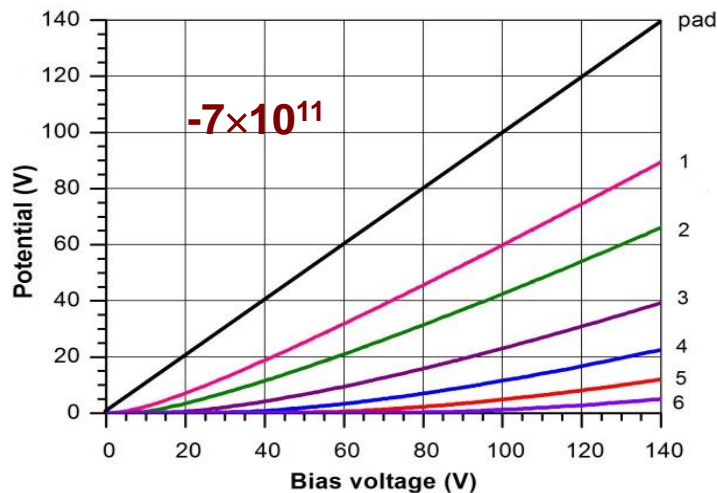


# Simulated potentials at various $Q_f$



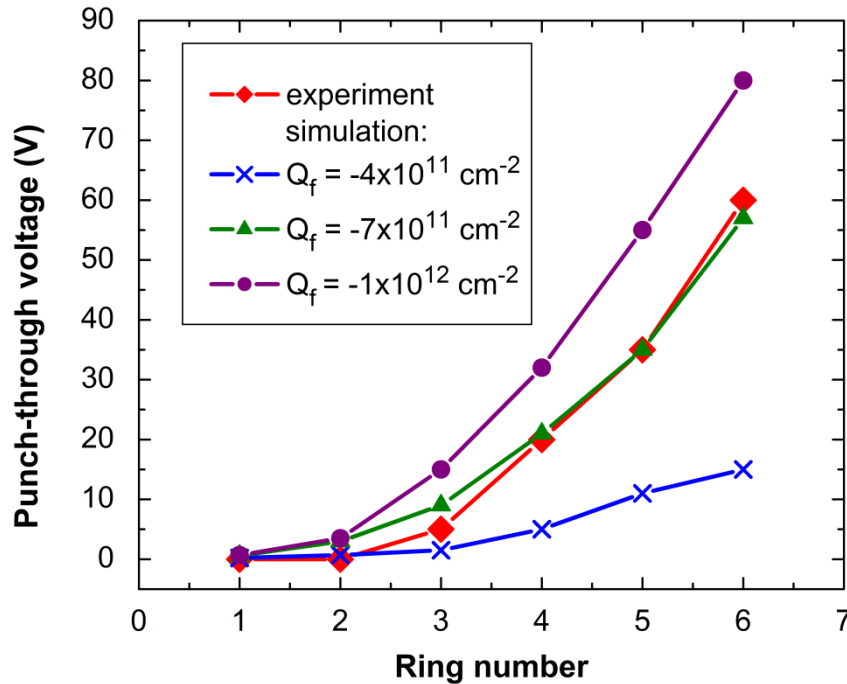
N.B. Scales in fragments differ

The same  $\phi$  as in the experiment



Potential of any ring goes down with  $Q_f$  rise

# Extraction of $Q_f$ from comparison of experimental and simulated data on $V_{th}$

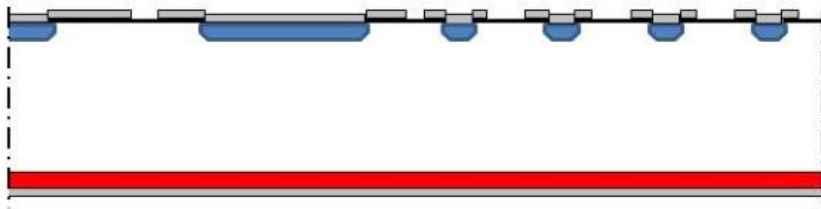


- $Q_f = -1 \times 10^{11} \text{ cm}^{-2}$  - no agreement
- $Q_f = -4 \times 10^{11} \text{ cm}^{-2}$  -  $V_{th}$  agree for the 1<sup>st</sup> and 2<sup>nd</sup> rings
- $Q_f = -7 \times 10^{11} \text{ cm}^{-2}$  - agreement for rings 4-6 better than 10%
- $Q_f = -1 \times 10^{12} \text{ cm}^{-2}$  - all simulated  $V_{th}$  values exceed the experimental data

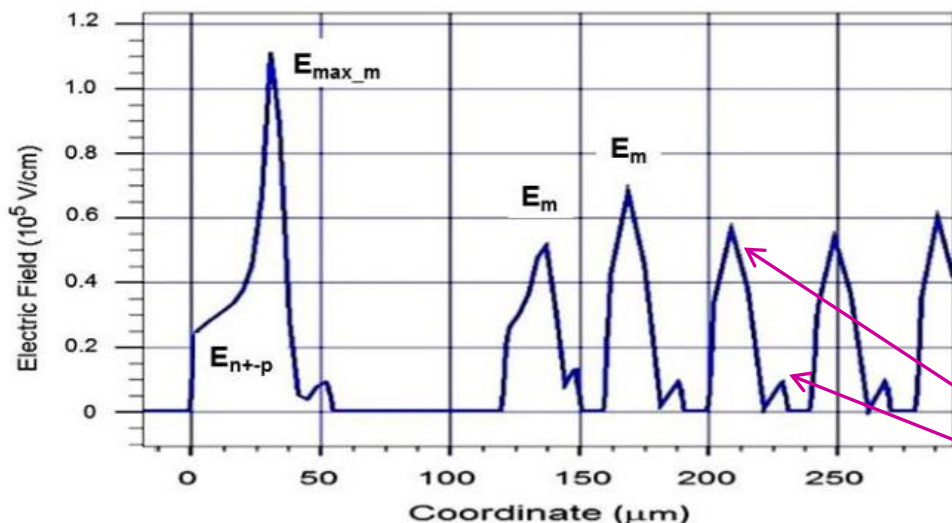
**In the detector under study (used design and technological process)**

**$Q_f$  is within the range  $-(4-7) \times 10^{11} \text{ cm}^{-2}$**

# Simulation of electric field distributions



Top: detector schematic topology  
Bottom: electric field at the distance of  $0.1 \mu\text{m}$  beneath  $\text{Al}_2\text{O}_3/\text{Si}$  interface at 500 V and  $Q_f$  of  $-4 \times 10^{11} \text{ cm}^{-2}$ ; taken from 2D electric field



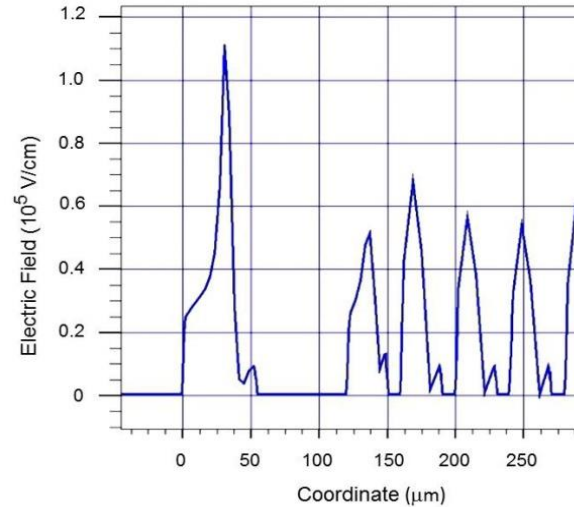
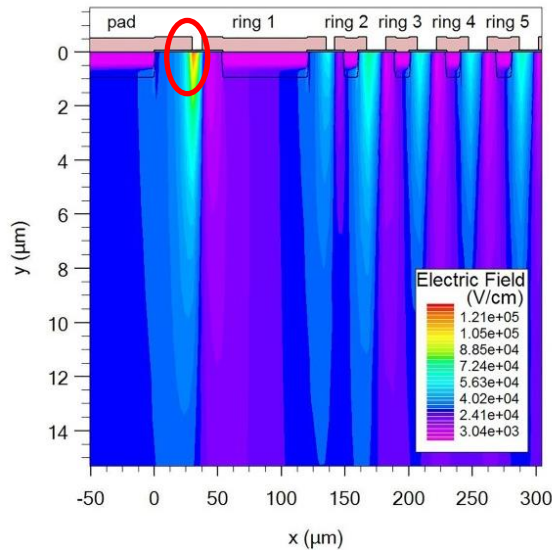
## Specific points in $E(x)$

$E_{n+p}$  - at pad junction edge ( $x = 0$ )  
 $E_m$  - E maxima at outward edges of the  $n^+$  segment field plates  
 $E_{\text{max}_m}$  - absolute E maximum, at the edge of the pad field plate

E around passivated gap - a structure with two separate peaks and a minimum *in between*

# Simulated electric field at various $Q_f$

$-4 \times 10^{11}$   
 $\text{cm}^{-2}$

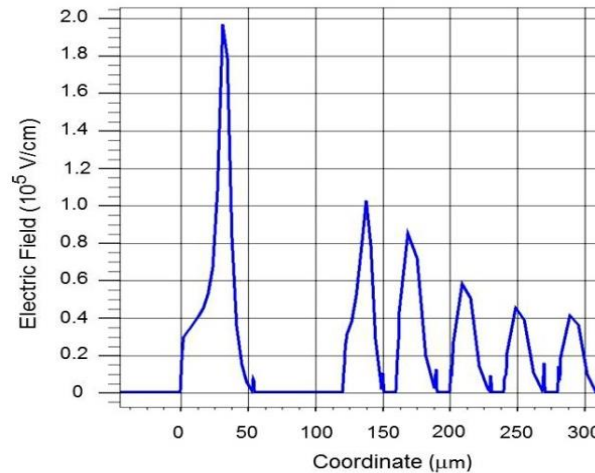
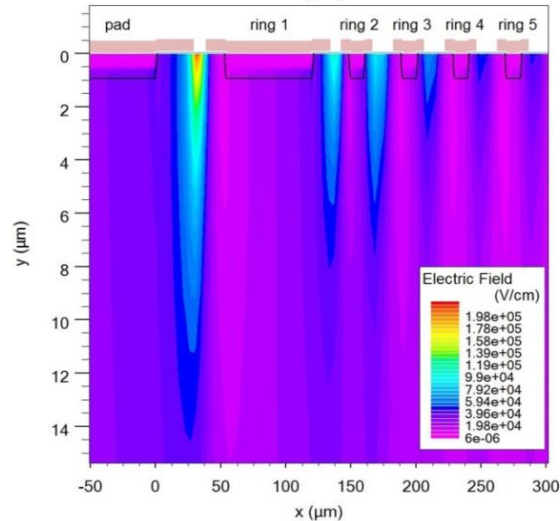


In oval:  
regions of  $E_{\text{max}_m}$   
and  $E_{\text{min}}$  in between

$$E_{\text{max}_m} = 1.1 \times 10^5 \text{ V/cm}$$



$-7 \times 10^{11}$   
 $\text{cm}^{-2}$



$\phi$  and  $E$  redistribu-  
tion, extension of  
high field region

$$E_{\text{max}_m} = 2 \times 10^5 \text{ V/cm}$$

Critical for carrier  
avalanche  
multiplication

# *Impact of design parameters on the electric field in the gaps: goal*

## **Goal: reduction of E at the edge of the pad field plate**

Two parameters were changed in simulations for n<sup>+</sup>-p-p<sup>+</sup> detectors with Al<sub>2</sub>O<sub>3</sub> isolation:

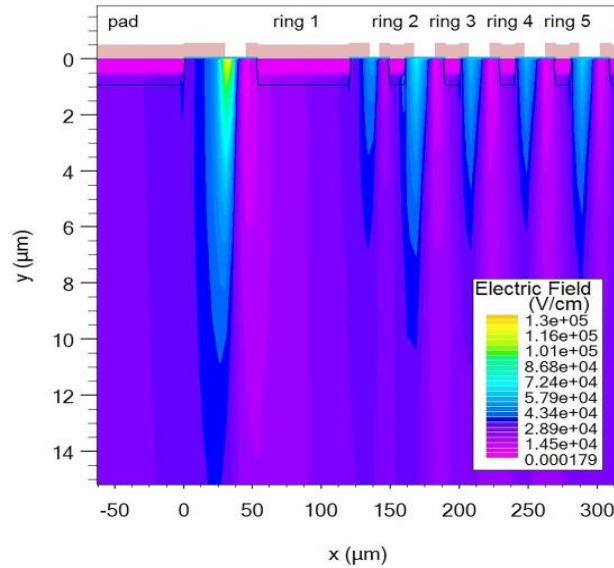
- Width of the gap between the edges of the pad field plate and inner metal plate of the bias ring ( $w_{\text{pad-BR}}$ )
- Silicon resistivity  $\rho$

### **Known for p-on-n detectors:**

- $w_{\text{pad-BR}}$  growth increases E at the pad field plate edge; our task is estimating the scale of this increase;
- use of n-type 20 kΩcm silicon (FZ Si) allowed to realize detectors with a thickness of **1 mm**, operating **up to 1000 V** with a current density  $\leq 10$  nA/cm<sup>2</sup> (see V. Eremin et al., *Semiconductors* (2011) **45**, No. 4, 536)

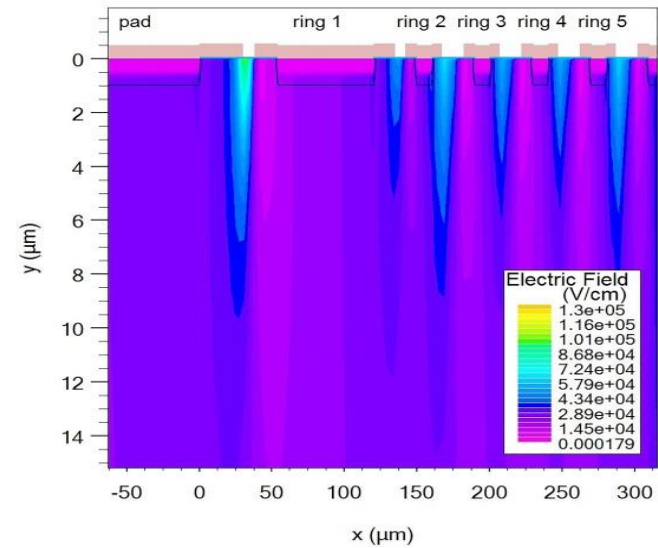
# Impact of design parameters on the electric field in the gaps: results

## 1. Increase of $w_{\text{pad-BR}}$

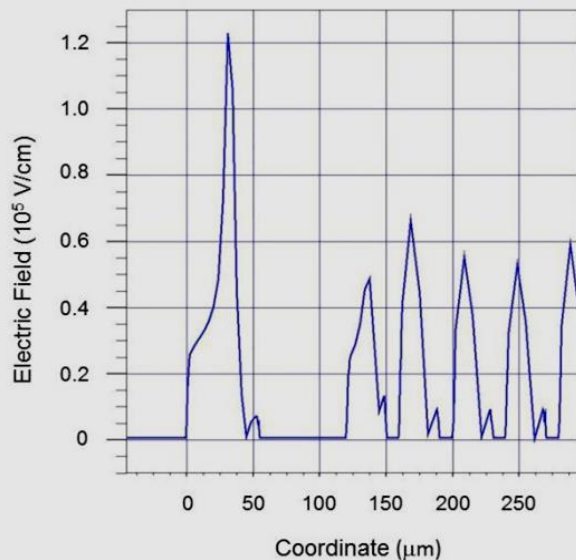


$-4 \times 10^{11}$   
 $\text{cm}^{-2}$ ;  
 500 V

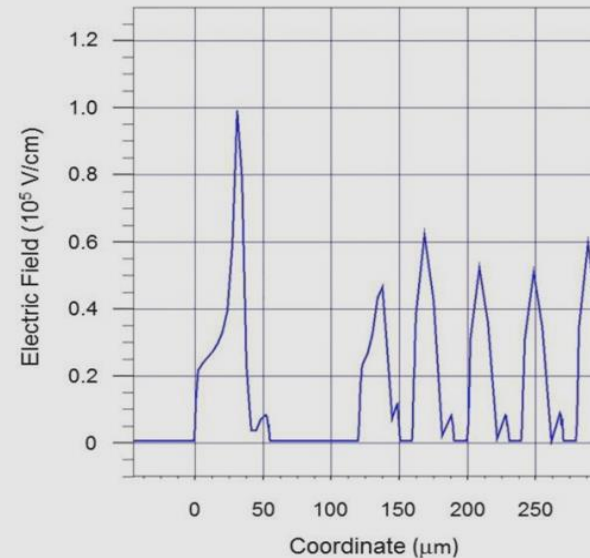
## 2. Increase of silicon resistivity to 20 k



$E_{\text{max}_m} \uparrow$   
 10%



$E_{\text{max}_m} \downarrow$   
 in 2 times





# Parameters of the detector design and characteristics derived from simulations

$$\alpha_e = A_e \exp\left(-\frac{B_e}{E}\right)$$

$$j(x) = \int eGdx + \int_0^d \alpha_e(x)j_e(x)dx$$

Si grade	$\rho$ (k $\Omega$ cm)	width of the pad field-plate ( $\mu$ m)	width of pad-BR gap ( $\mu$ m)	$Q_f$ (cm <sup>-2</sup> )	$E_{\max\_m}$ at pad field-plate edge (V/cm)	$\alpha_e$ (cm <sup>-1</sup> )
<b>experiment</b>						
MCZ	7	30	8	-(4-7) $\times 10^{11}$		
<b>simulations</b>						
MCZ	7	30	8	$-1 \times 10^{11}$	$5 \times 10^4$	$3.3 \times 10^{-5}$
MCZ	7	30	8	$-4 \times 10^{11}$	$1.1 \times 10^5$	14.7
MCZ	7	30	8	$-7 \times 10^{11}$	$2 \times 10^5$	$2 \times 10^3$
MCZ	7	30	15*	$-4 \times 10^{11}$	$1.2 \times 10^5$	36.4
FZ	20	30	8	$-4 \times 10^{11}$	$1 \times 10^5$	4.93

\* the width of the BR field plate is reduced to 9  $\mu$ m.

## With respect to avalanche multiplication:

- $Q_f$  of  $-7 \times 10^{11}$  cm<sup>-2</sup> is critical
- Application of 20 k $\Omega$ cm p-Si allows reducing the electric field and  $\alpha_e$   
Experiment: see [5] in slide 2, n-on-p detectors operating at V up to 900 V.

# Summary

1. The measurements showed **efficient VTS operation**: gradual potential attenuation to less than 10% of the applied bias on the last n<sup>+</sup> ring.
2. **Punch-through model** built for the p<sup>+</sup>-n-n<sup>+</sup> detectors with SiO<sub>2</sub> passivation is valid for n<sup>+</sup>-p-p<sup>+</sup> detectors having Al<sub>2</sub>O<sub>3</sub> isolation layer with the opposite oxide charge.
3. Simulations allowed extracting the **Q<sub>f</sub> charge of  $-(4-7)\times 10^{11} \text{ cm}^{-2}$**  in Al<sub>2</sub>O<sub>3</sub> in the detector with the used design and technology ( $-(10^{11}-10^{13}) \text{ cm}^{-2}$  is referred to in literature).
4. **Q<sub>f</sub> is shown to be crucial** for VTS operation and detector performance. For the studied detector, **an upper limit of Q<sub>f</sub> is  $-7\times 10^{11} \text{ cm}^{-2}$** .
5. A promising approach to reducing the electric field is to **use p-type Si with increased resistivity** of about 20 kΩcm.



*Thank you for attention!*