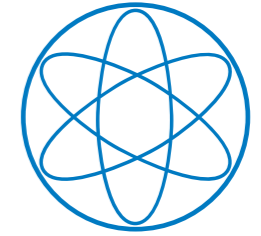


Nature vs. BSM

Andreas Weiler
TU München

20th annual RDMS CMS Collaboration Conference
Tashkent
13/9/2018





“Is nature natural?”

Andreas Weiler
TU München

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Shouldn't we believe that the Standard Model is exactly correct?

The ability to fit data is not everything. The SM fails to **explain** many of the most important features of particle physics.

The SM is incapable of explaining the **phase transition** to an ordered state that breaks $SU(2)_L \times U(1)_Y$.

1. The most general renormalizable potential for the Higgs field is

$$V = \mu^2 |\Phi|^2 + \lambda |\Phi|^4$$

2. For some reason, $\mu^2 < 0$.

The value of μ^2 receives large quantum corrections with both signs of order of the fundamental scale $\Lambda \gg \mu$.

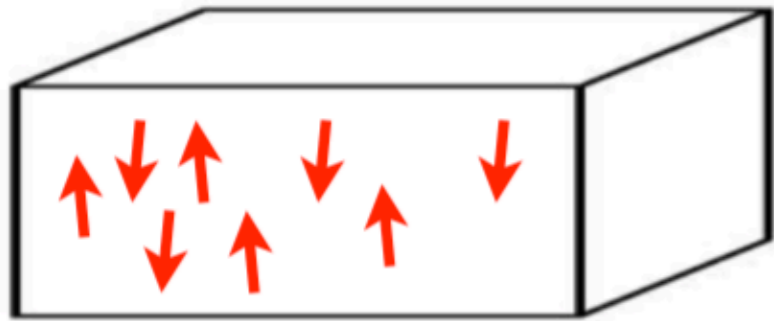
Sophisticated people call this the “gauge hierarchy problem”. We have **no clue as where μ^2 comes from.**

We have experimentally performed the tuning

$$\mu \ll \Lambda$$

many times...

For generic T , a ferromagnet is not at critical point T_c



$$\Lambda^{-1} \longleftrightarrow a$$

atomic spacing

$$\mu^{-1} \longleftrightarrow \xi$$

coherence length

$$\Lambda \gg m \longleftrightarrow \xi \gg a \quad \text{critical point}$$

For generic T , a ferromagnet is not at critical point T_c



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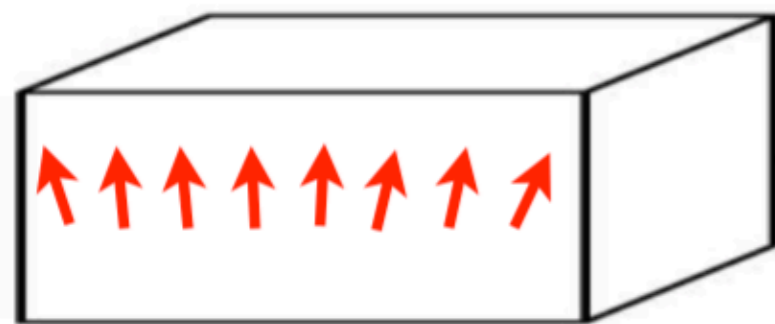
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$T \rightarrow T_c$ requires to finetune the temperature:



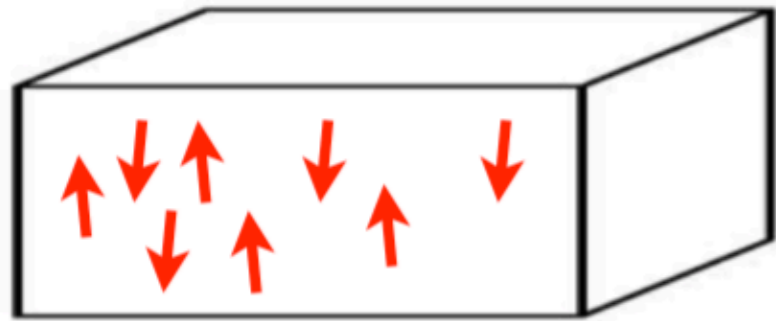
ξ

$$T \approx T_c$$



experimenter

For generic T , a ferromagnet is not at critical point T_c



$$\Lambda^{-1} \longleftrightarrow a$$

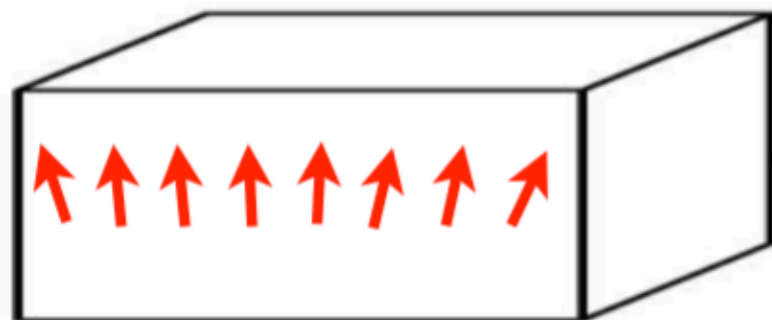
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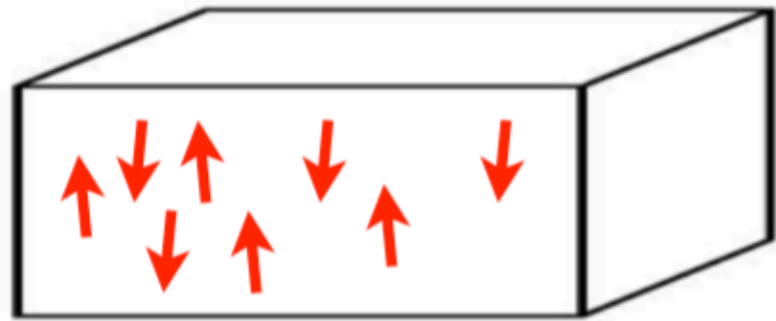
$T \approx T_c$



ξ



For generic T , a ferromagnet is not at critical point T_c



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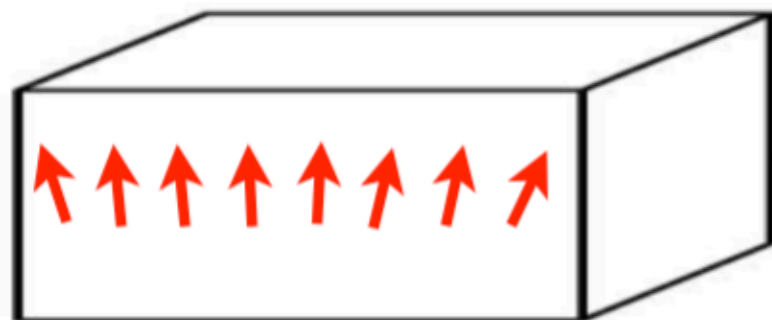
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$$\Lambda \gg m \longleftrightarrow \xi \gg a \quad \text{critical point}$$

$T \rightarrow T_c$ requires to finetune the temperature:



$$T \approx T_c$$



ξ



There is a strong analogy here to the theory of superconductivity.

$$F = F_0 + \alpha |\psi|^2 + \frac{\beta}{2} |\psi|^4 + \dots$$

(Landau-Ginzburg
free energy)

In **1950**, Landau and Ginzburg proposed a phenomenological theory that explained many properties: the thermodynamics of the phase transition, the critical magnetic field, the presence of Type I and Type II, etc.

However, Landau Ginzburg gave no fundamental understanding of superconductivity. That was found only in 1957, by Bardeen, Cooper, and Schrieffer.

For the electroweak phase transition, we are still in the Landau-Ginzburg era.

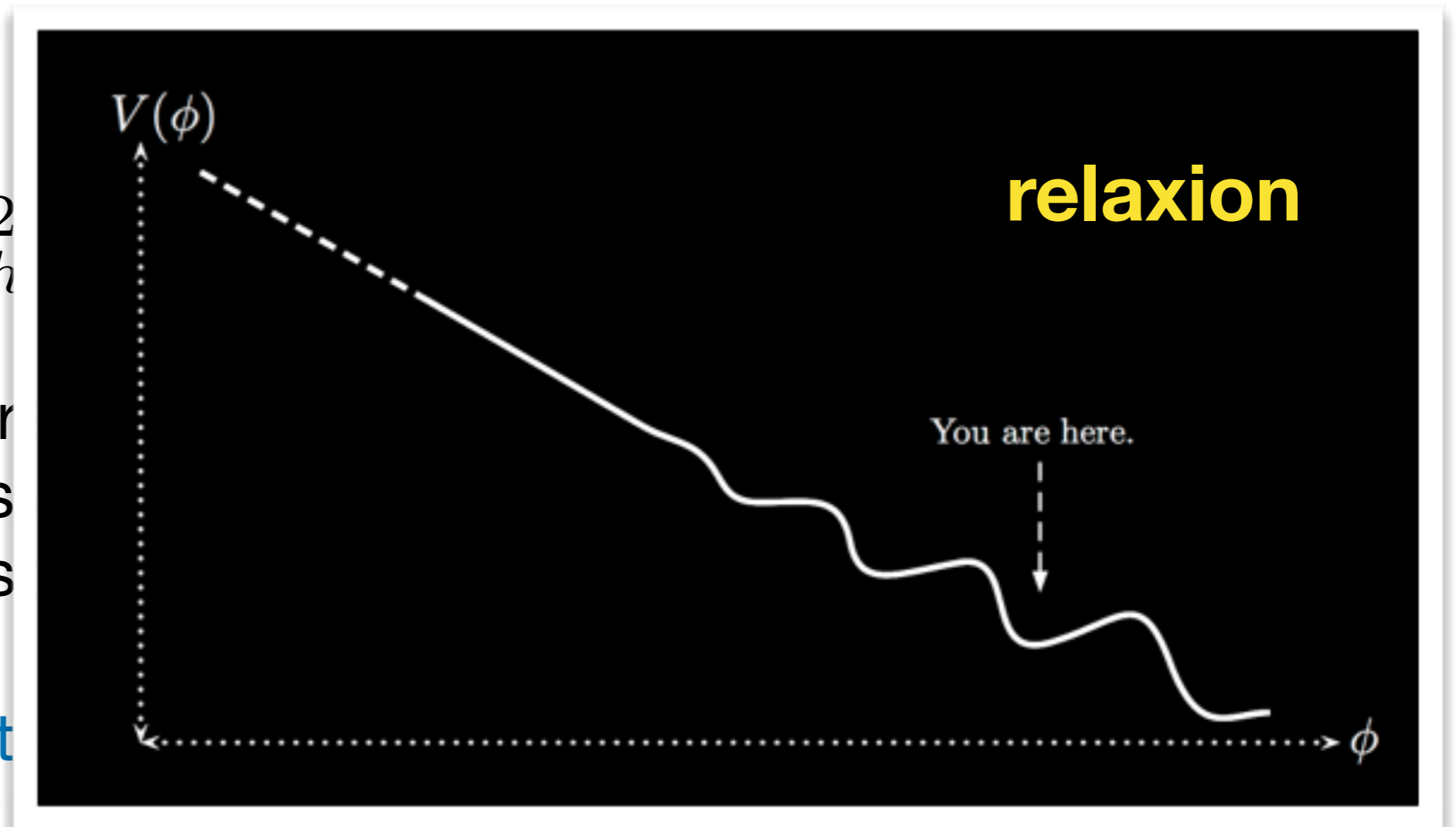
What are the options?

Solutions

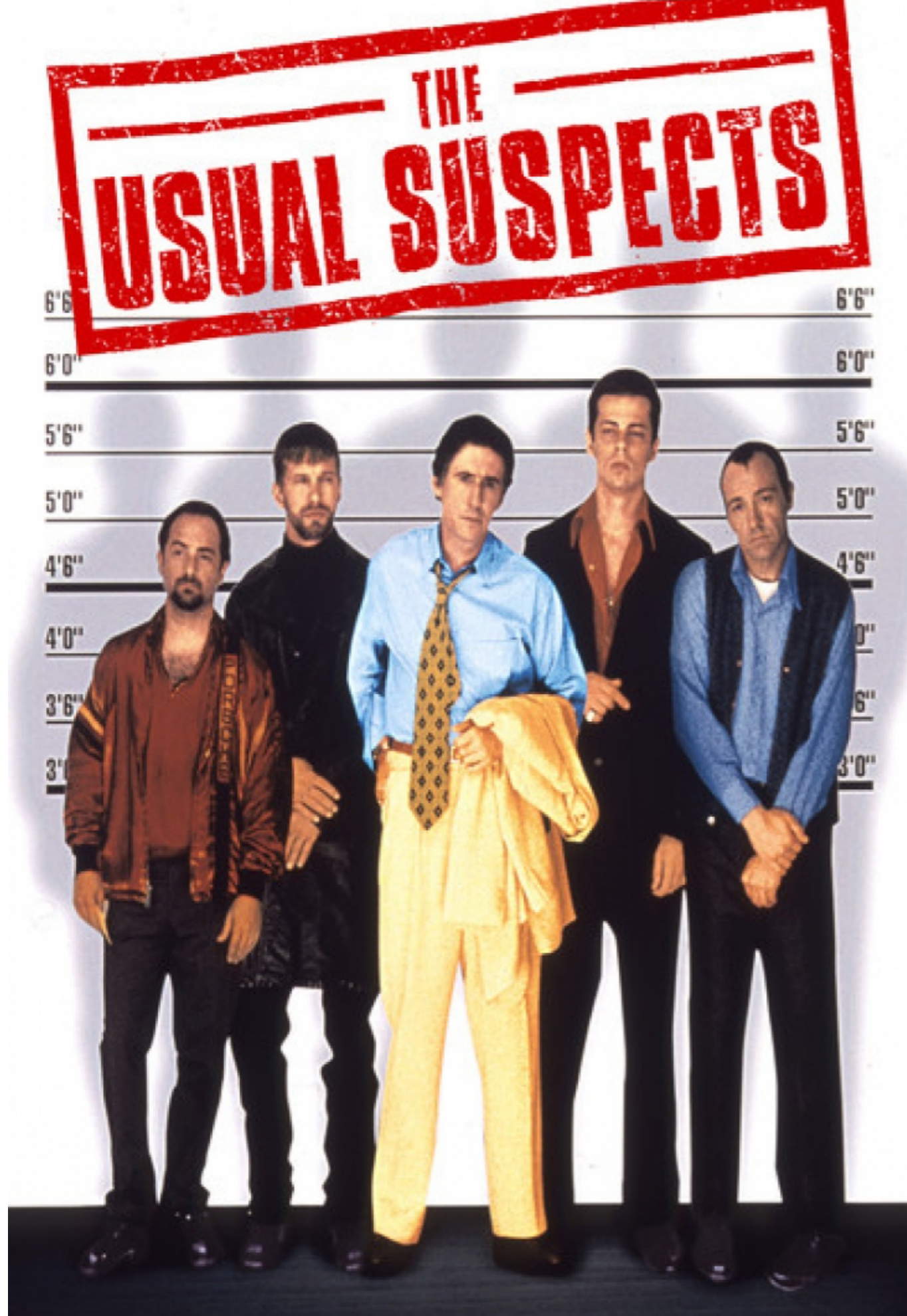
- Protect $m_h^2 \phi^* \phi$ by some symmetry
 - **Susy**: chiral symmetry. Super-multiplet (higgs,higgsino). Exact susy guarantees equal mass of fermion and higgs boson (also at quantum level).
 - **Composite pGB Higgs**: shift symmetry $\phi \rightarrow \phi + \alpha$
- Make Higgs mass dynamical and select $(\Lambda^2 - \varphi) \phi^* \phi$ vacuum with cosmological dynamics (**Relaxion**)

Solutions

- Protect m_h^2
- **Susy**: chiral
Exact susy
higgs boson
- **Compositeness**



- Make Higgs mass dynamical and select $(\Lambda^2 - \varphi) \phi^* \phi$ vacuum with cosmological dynamics (**Relaxion**)



xMSSM

scalar top partners

DM ✓

composite pGB Higgs

fermionic top partners

DM ?

Minimally natural susy

sparticle
masses

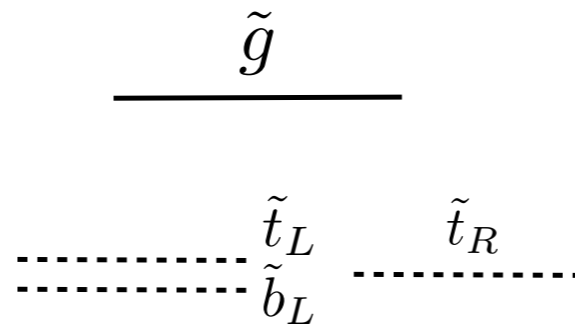
$$\delta m_H^2$$

\ll

vs current limits
(13 TeV, 36 fb⁻¹)

≈ 1000 GeV

2loop



2020 GeV

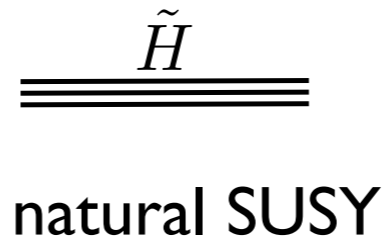
≈ 500 GeV

1loop

1050 GeV

≈ 250 GeV

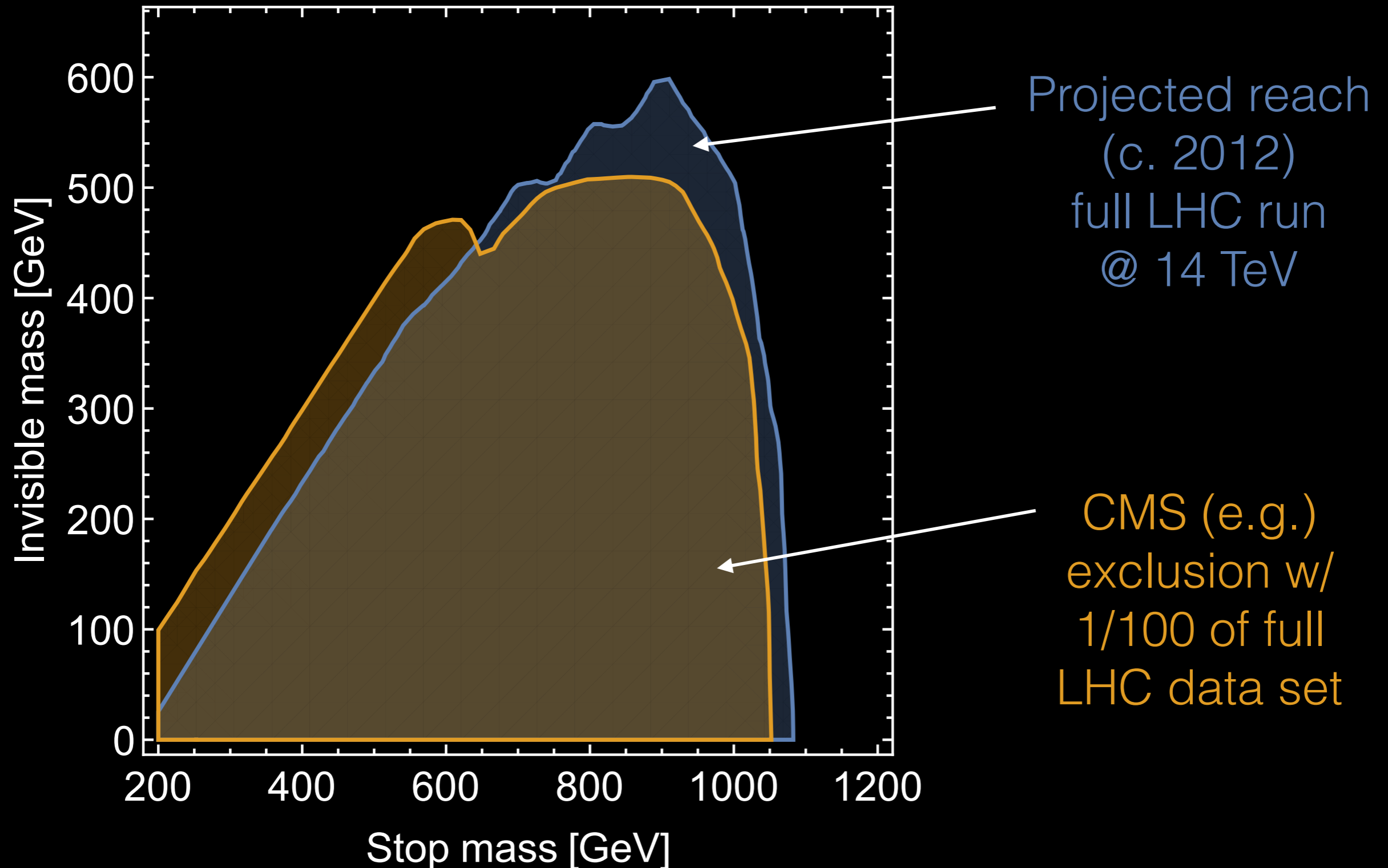
tree



145 GeV

**O(1) departure from
natural expectation!**

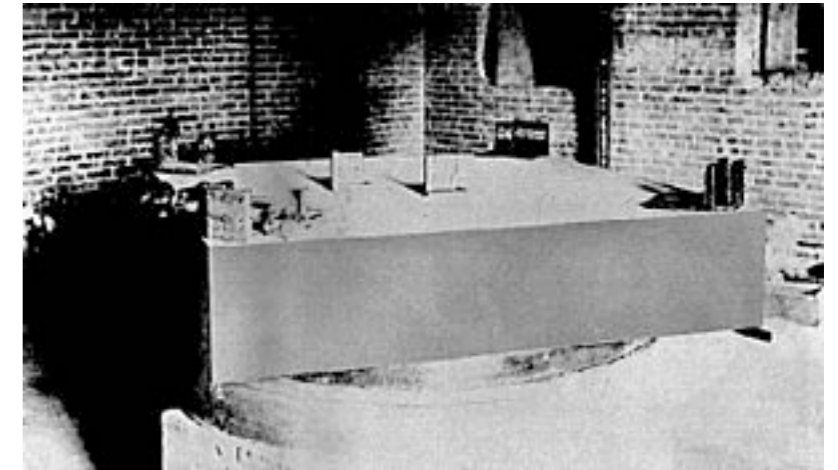
Not for lack of trying





?

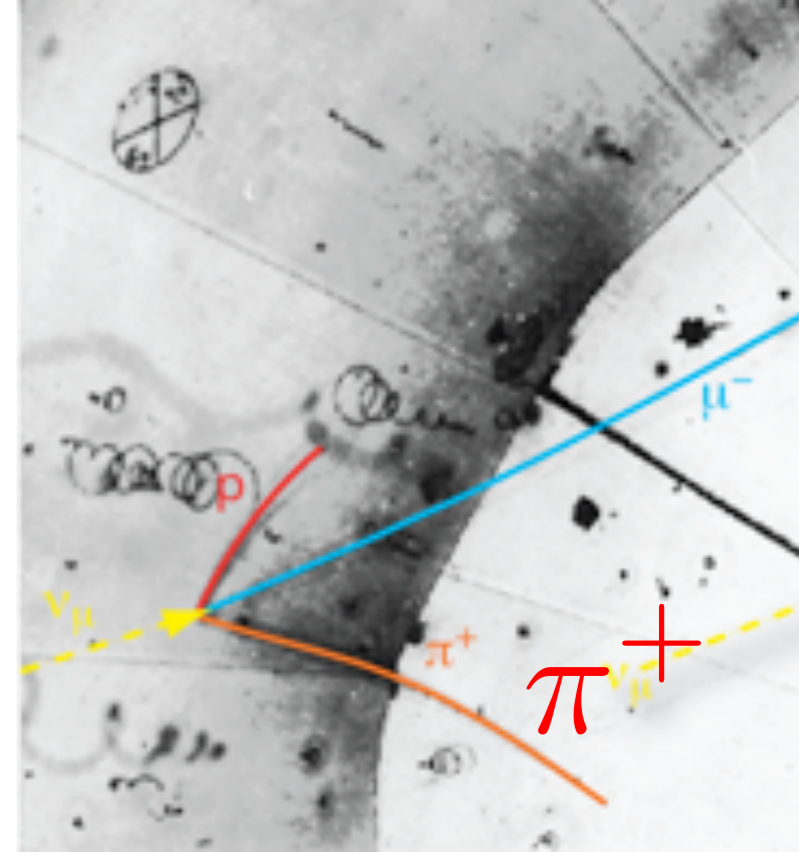
||



Recent development: DM vs. composite Higgs

Composite Higgs & composite DM

Balkin, Ruhdorfer, Salvioni, **AW**



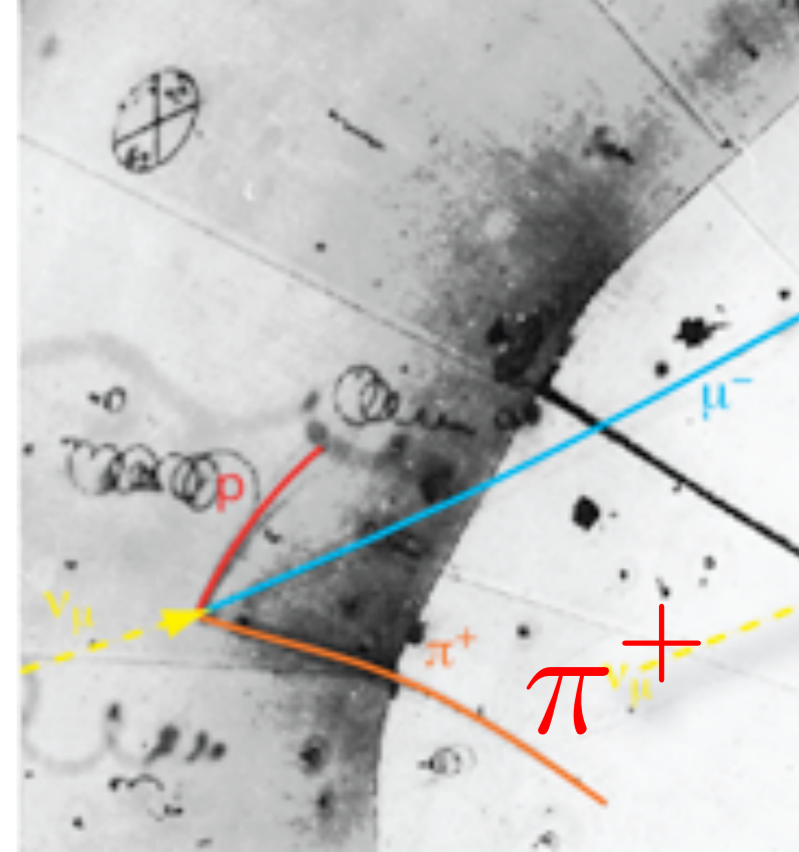
$$U = e^{i\pi/f} \langle \phi \rangle \quad \mathcal{L}_{\text{eff}} = \frac{f^2}{4} \text{tr}(\partial_\mu U^\dagger \partial^\mu U) + \dots$$

- Higgs and DM are pseudo Goldstone bosons
- Same symmetry that protects Higgs mass also stabilises DM

(analog to R parity in SUSY)

Composite Higgs & composite DM

Balkin, Ruhdorfer, Salvioni, **AW**



$$U = e^{i\pi/f} \langle \phi \rangle \quad \mathcal{L}_{\text{eff}} = \frac{f^2}{4} \text{tr}(\partial_\mu U^\dagger \partial^\mu U) + \dots$$

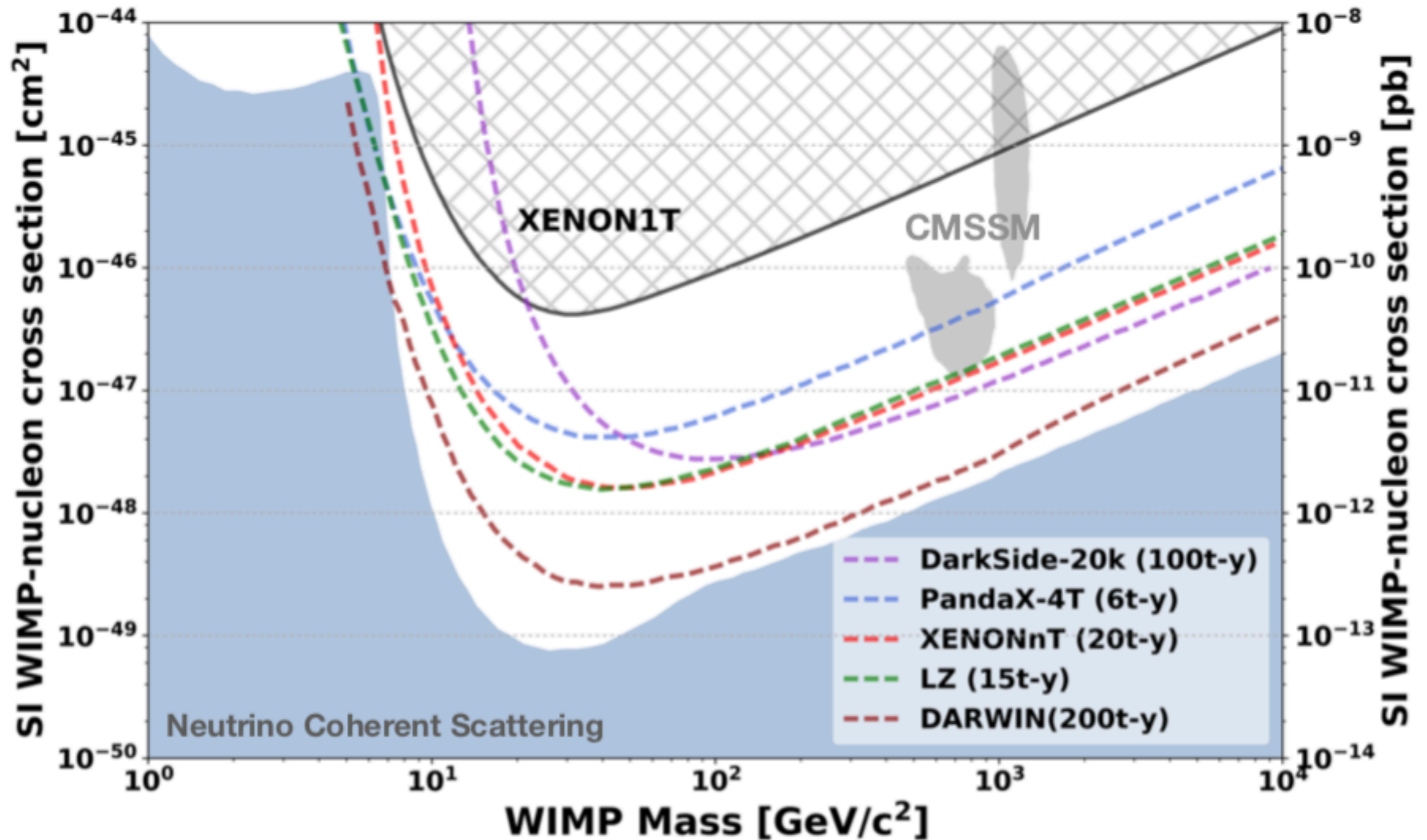
- Higgs and DM are pseudo Goldstone bosons
- Same symmetry that protects Higgs mass also stabilises DM

$$\pi \longrightarrow -\pi$$

$$(U \longrightarrow U^\dagger)$$

(analog to R parity in SUSY)

WIMPs are getting squeezed



Goldstone DM

$$\chi \rightarrow \chi + f \quad \text{GB shift symmetry}$$

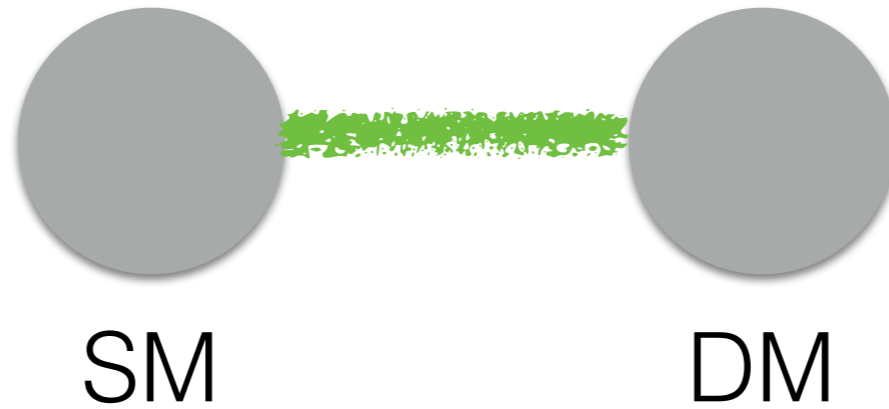
Only sizeable coupling between DM and SM:

$$\frac{1}{f^2} \partial_\mu (h^2) \partial^\mu (\chi^* \chi) + \dots$$

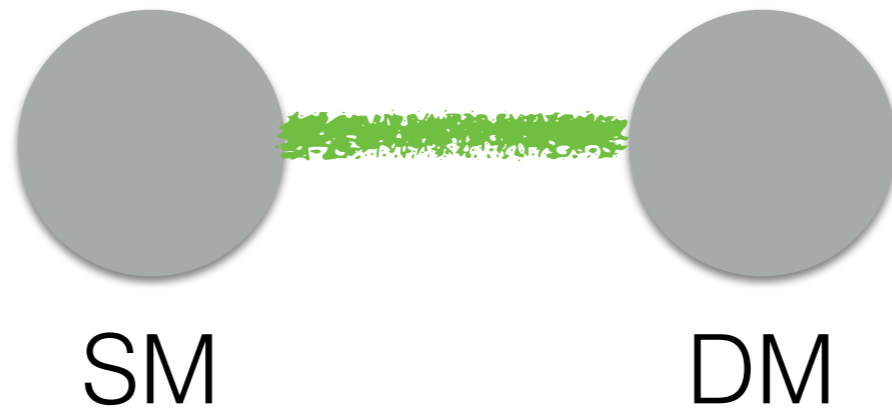
Negligible direct detection signal:

$$\frac{E_{\text{recoil}}^2}{f^2} \sim \frac{(100\text{MeV})^2}{(1\text{TeV})^2} \sim 10^{-8}$$

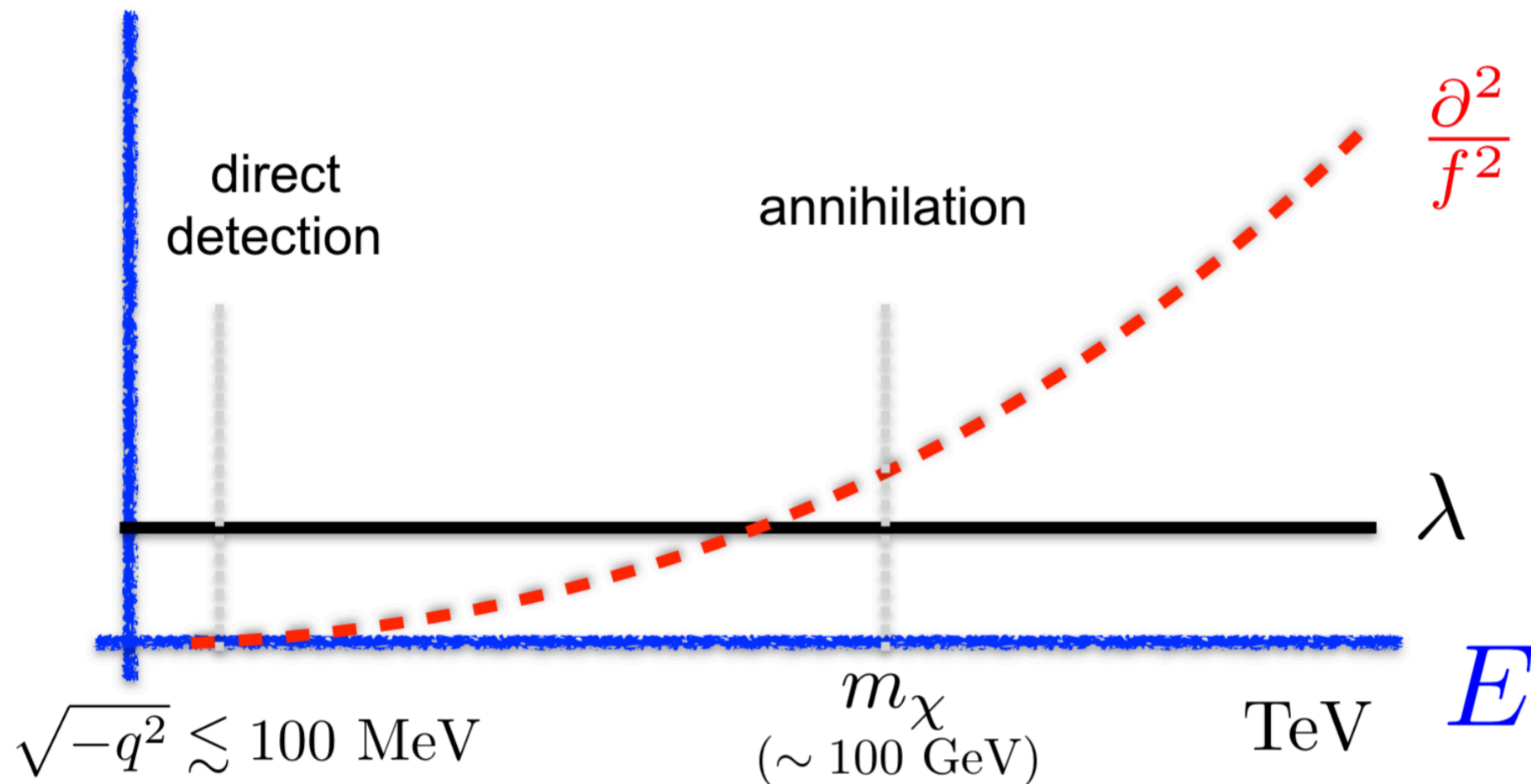
$$\mathcal{L}_{SM}(h) + \mathcal{L}_{DM}(\chi) + \lambda h^2 \chi^* \chi + \frac{1}{f^2} \partial_\mu (h^2) \partial^\mu (\chi^* \chi) + \dots$$



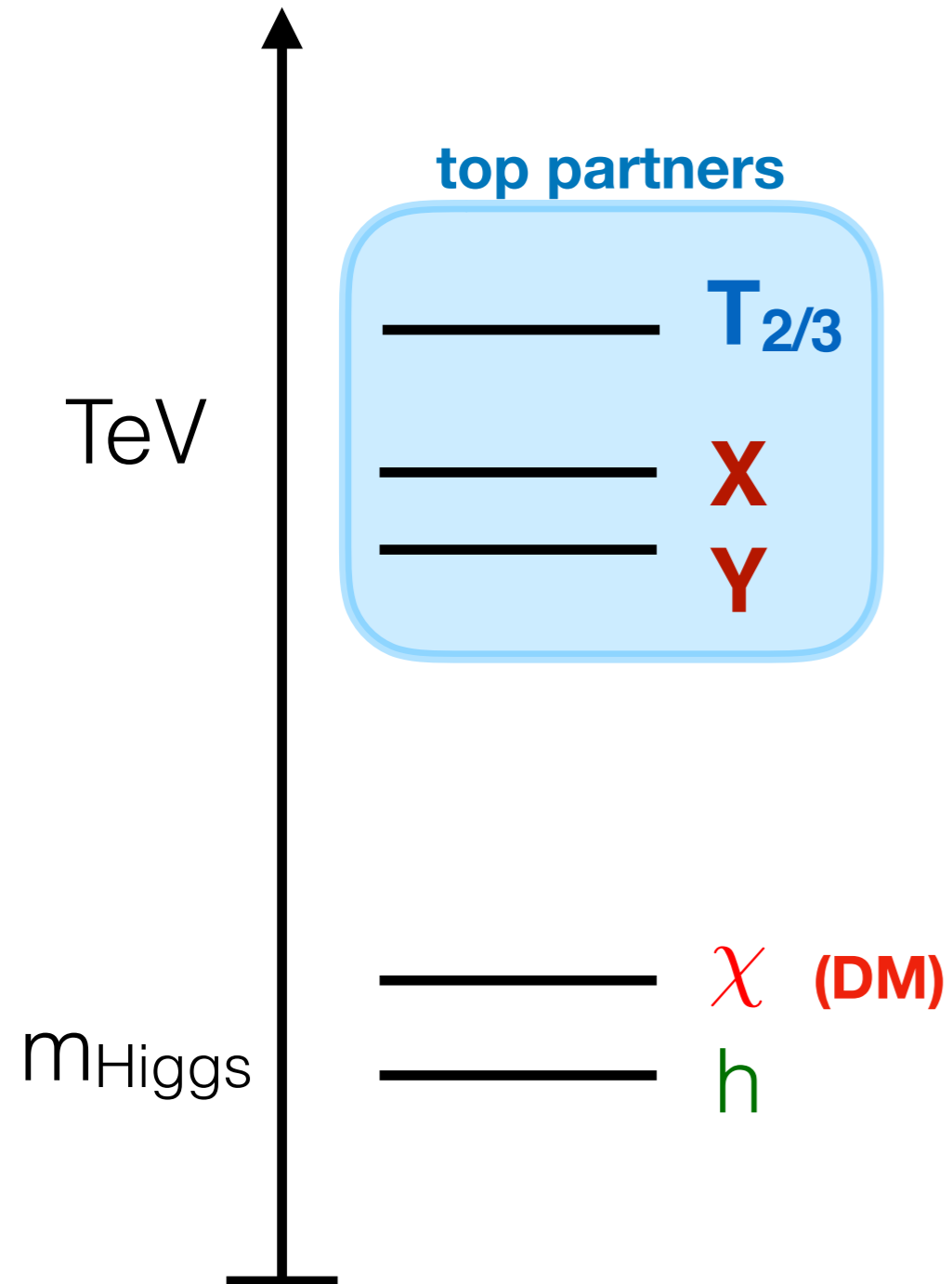
$$\mathcal{L}_{SM}(h) + \mathcal{L}_{DM}(\chi) + \lambda h^2 \chi^* \chi + \frac{1}{f^2} \partial_\mu (h^2) \partial^\mu (\chi^* \chi) + \dots$$



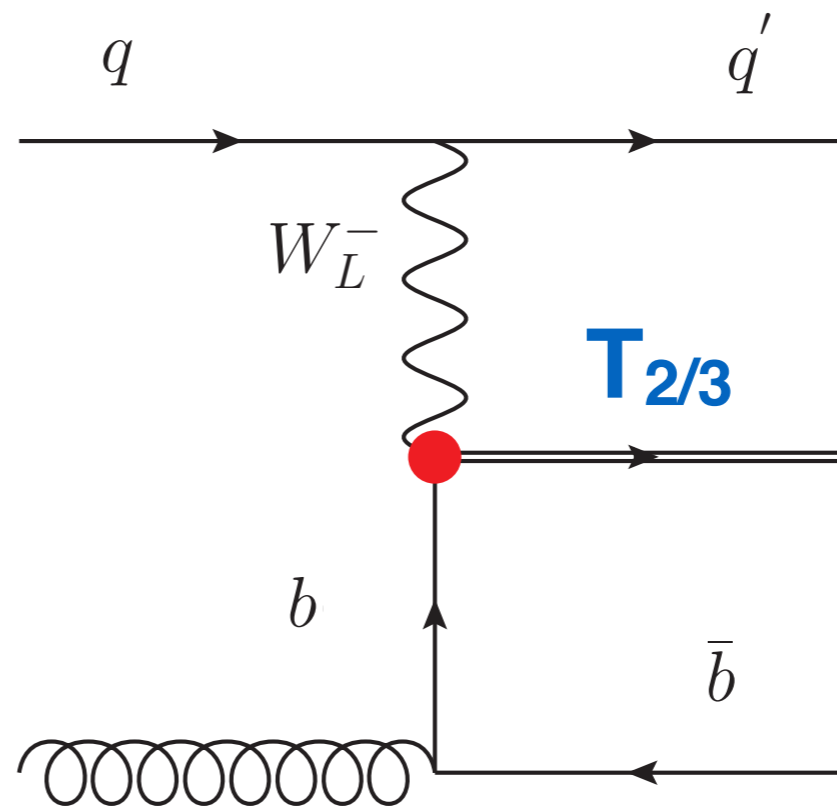
$$g_{DM-SM}^2(E)$$



Fermionic top partners+ MET

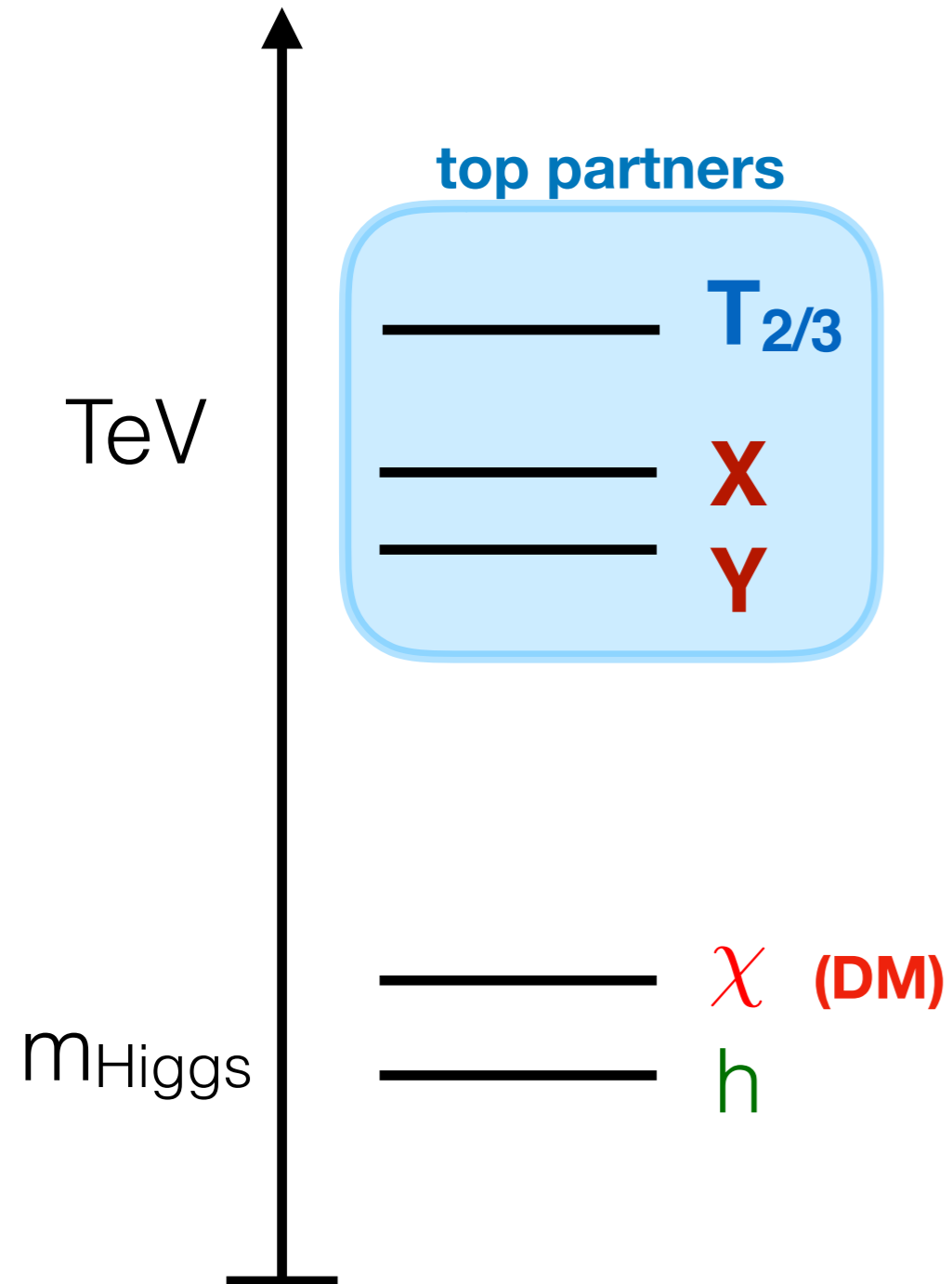


new signatures (single top +MET)

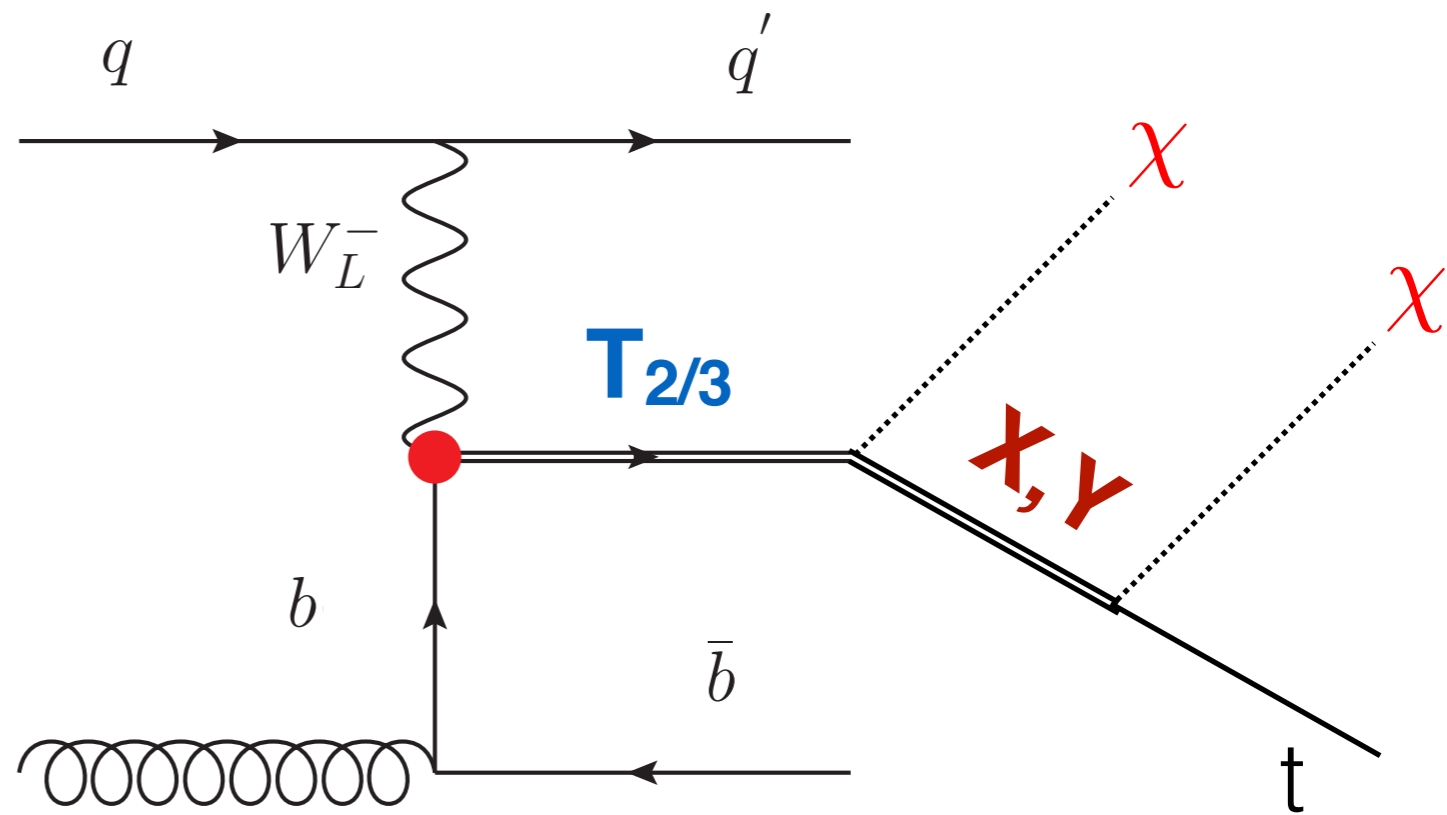


$$\text{BR}(S \rightarrow \chi^* \mathcal{Y}) = \text{BR}(S \rightarrow \chi \mathcal{Z}) \simeq \frac{1}{6}.$$

Fermionic top partners+ MET



new signatures (single top +MET)

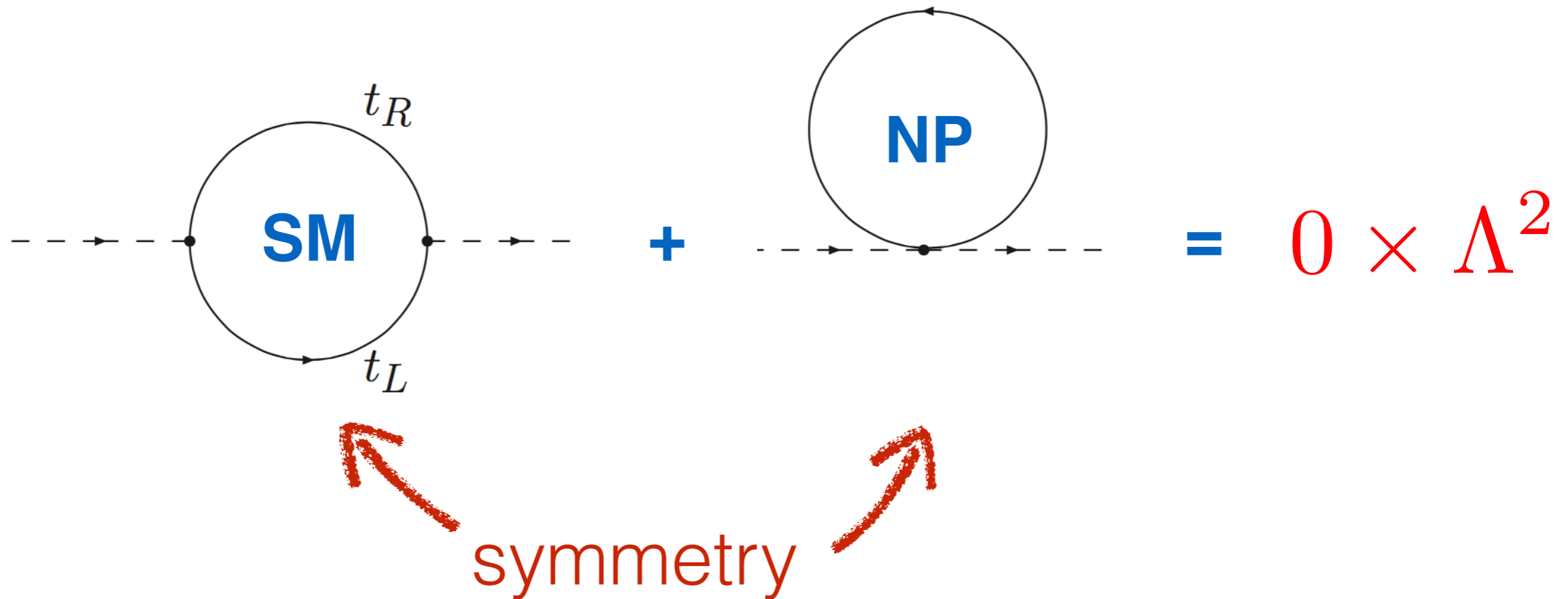


$$\text{BR}(S \rightarrow \chi^* \mathcal{Y}) = \text{BR}(S \rightarrow \chi \mathcal{Z}) \simeq \frac{1}{6}.$$

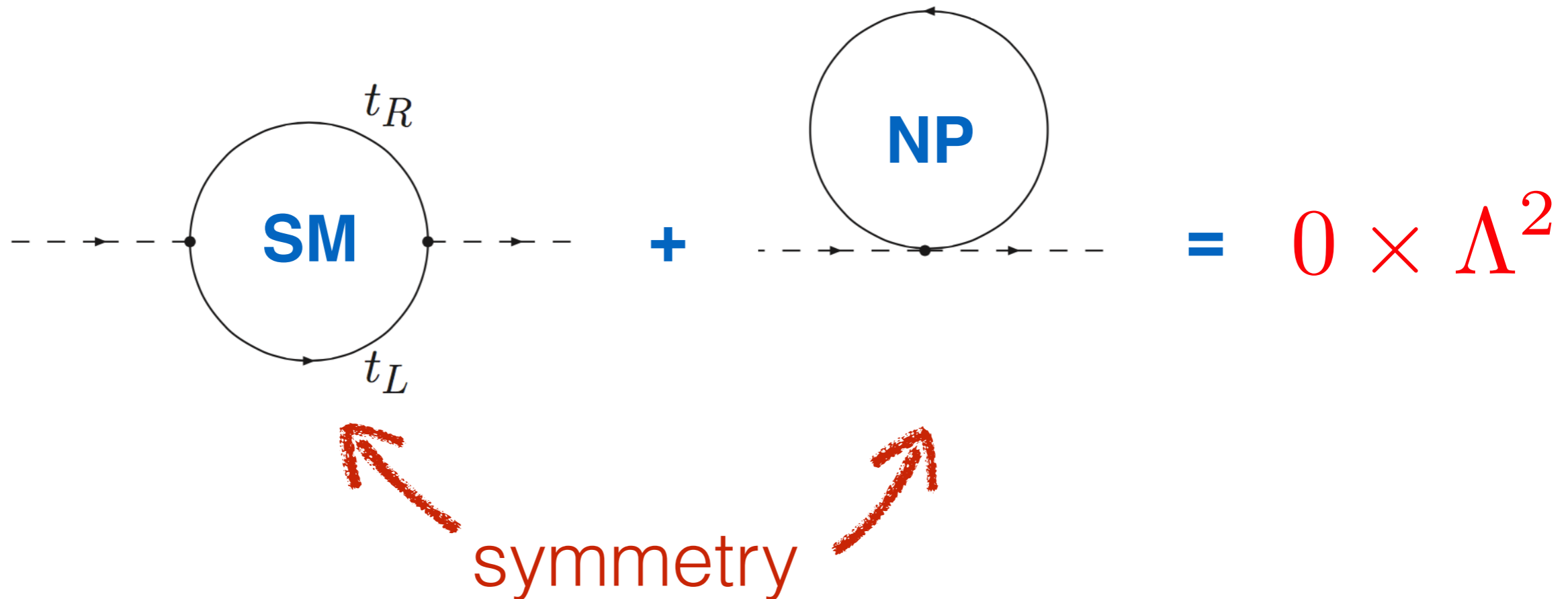
Unusual suspects



The unusual suspects



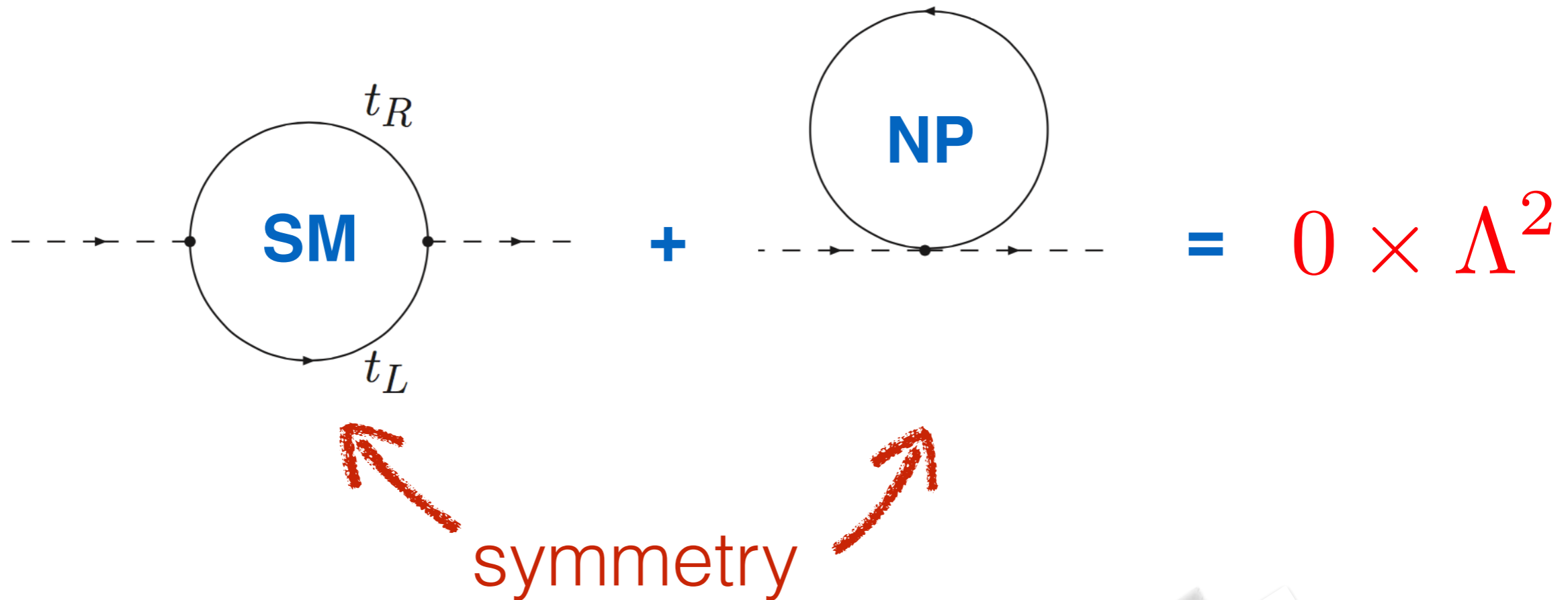
The unusual suspects



NP is related to the top by a symmetry, natural new particle mass around TeV

Symmetry commutes with color: NP will be produced copiously at the LHC!

The unusual suspects



NP is related to the top by a symmetry, natural new particle mass around TeV

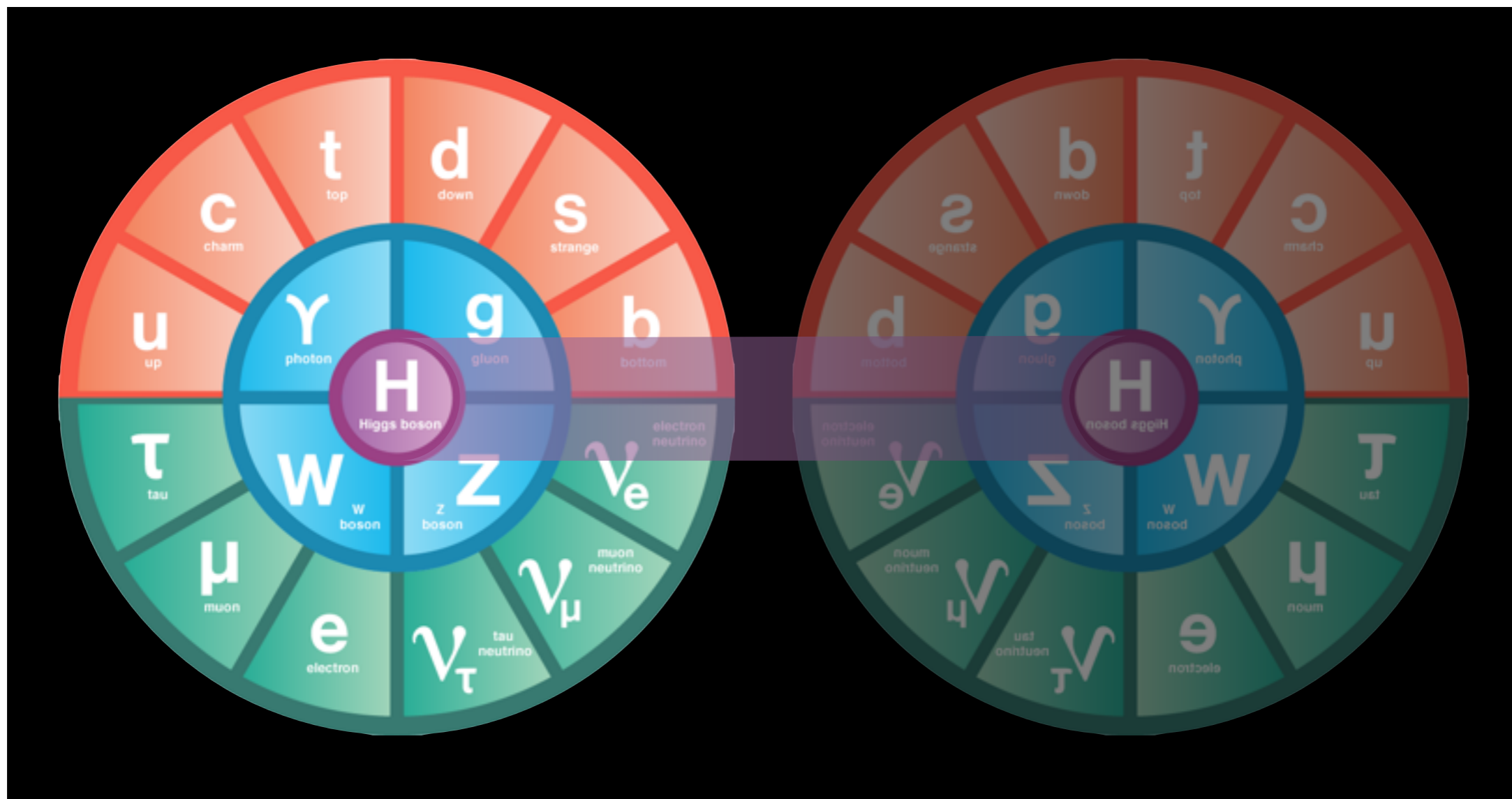
Symmetry breaks with color: NP produced copiously at the LHC!

Not true!

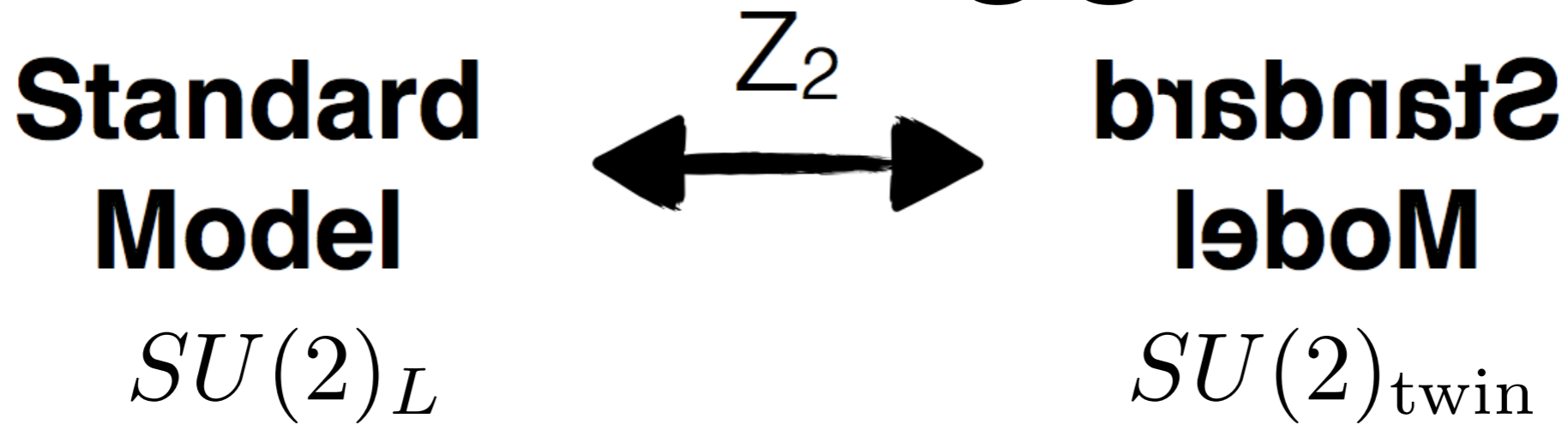
Twin Higgs & co

Chacko, Goh, Harnik '05

New ingredient: discrete symmetry



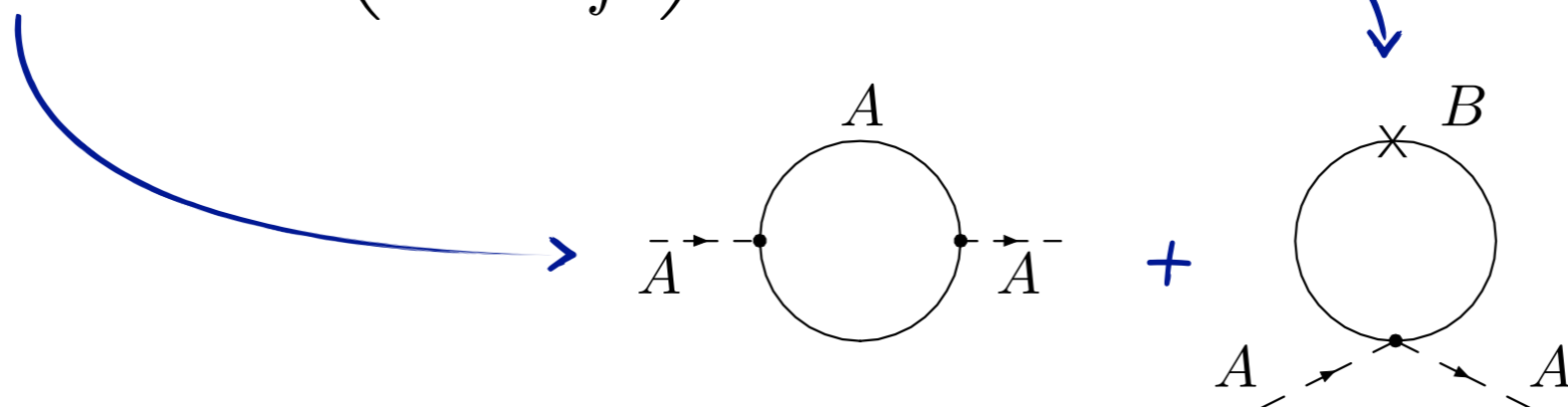
Twin Higgs



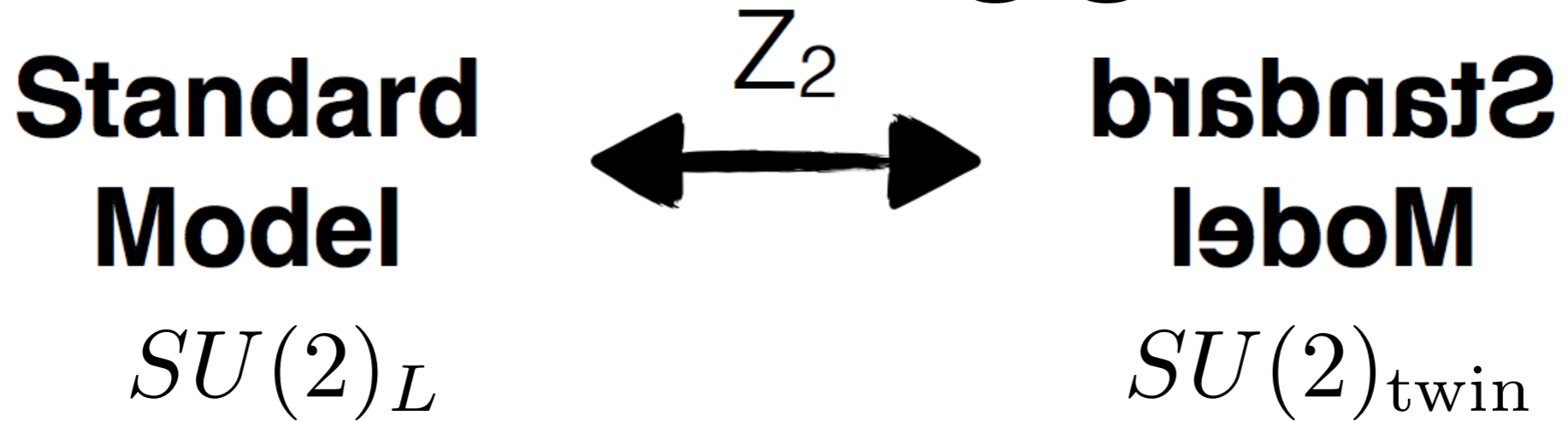
Quadratic divergences from SM top quark loops cancelled by loops of “Twin” top quarks.

$$\mathcal{L} \supset y_t H_A \bar{t}_A t_A + y_t H_B \bar{t}_B t_B$$

$$= y_t h \bar{t}_A t_A + y_t \left(f - \frac{|h|^2}{2f} \right) \bar{t}_B t_B + \dots$$



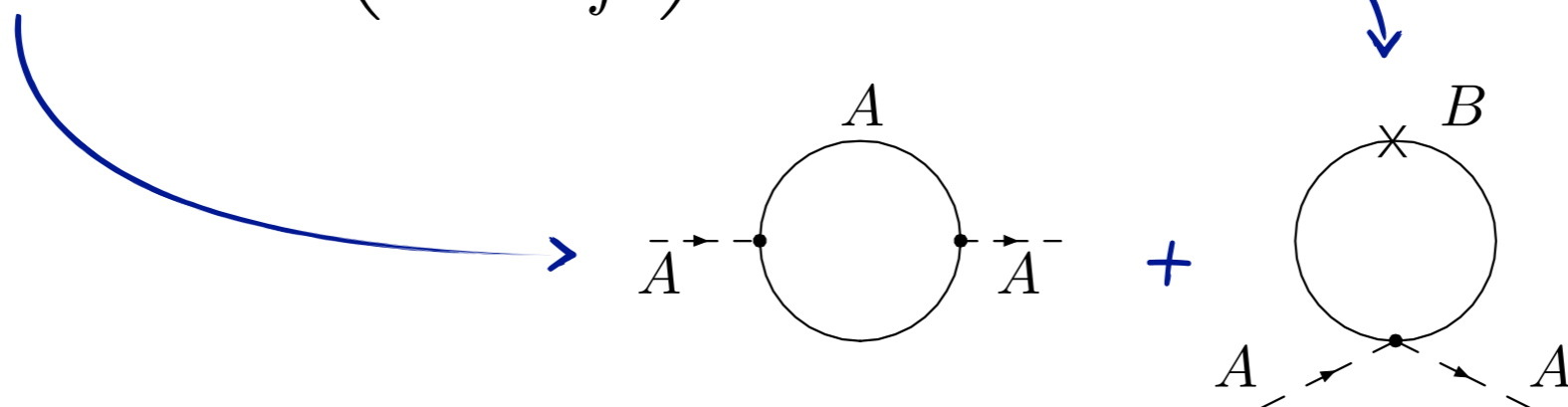
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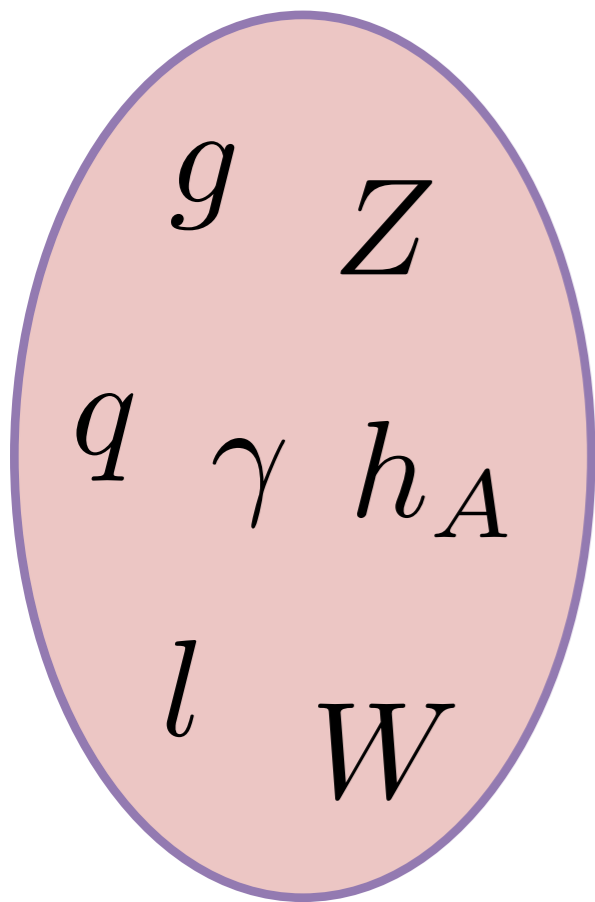
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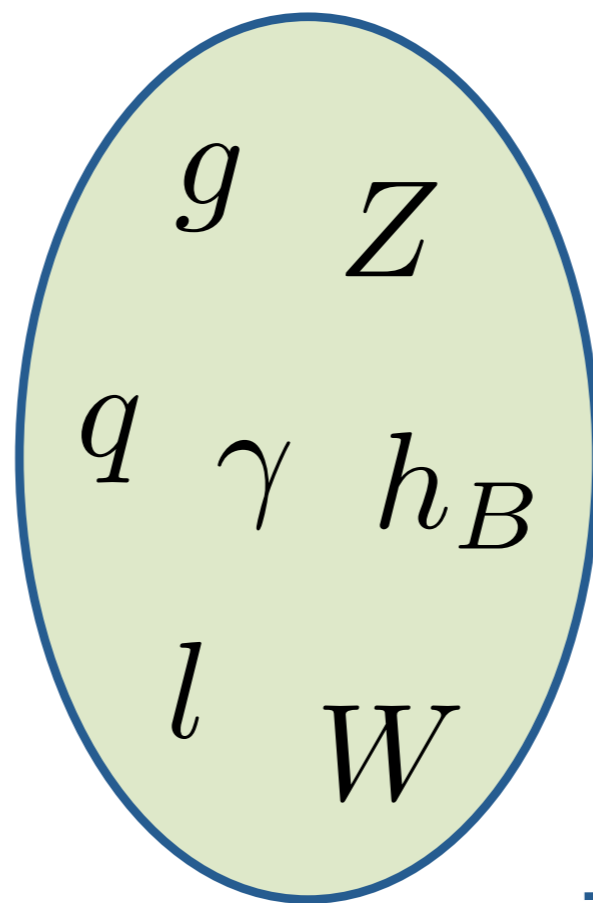


$$= 0 \times \Lambda^2$$

Standard
Model



“Twin”
Standard
Model

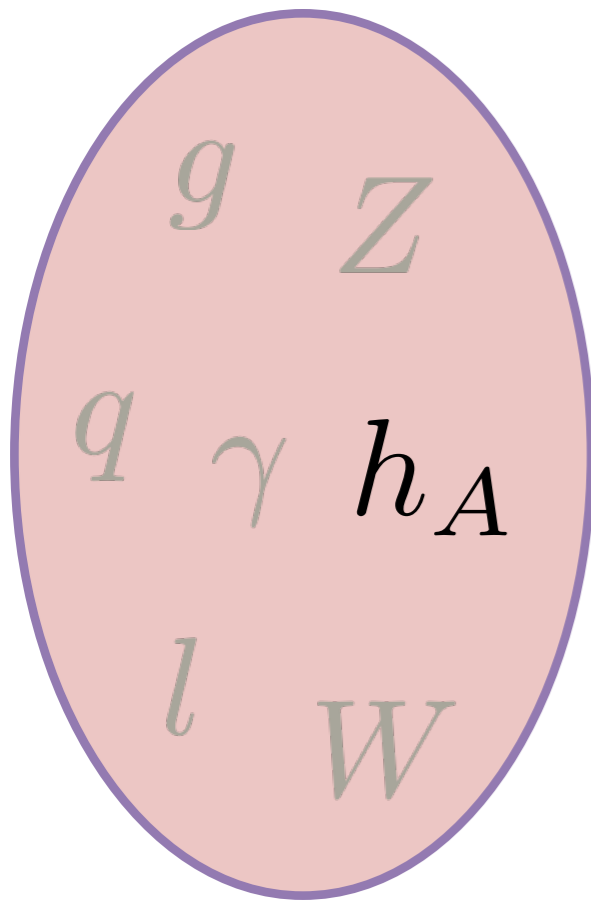


These fields
completely
neutral:
“Neutral
Naturalness”

Predictions for Twin sector most robust for the Twins
of the SM fields that couple most strongly to Higgs.

Standard Model

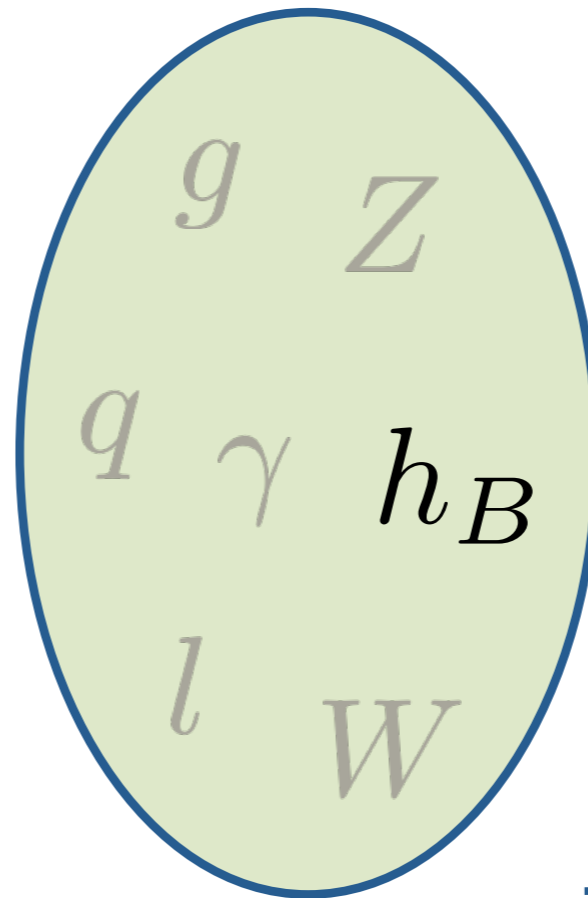
“Twin” Standard Model



$$\sim m^2 h_A h_B$$



Only communication through small “Higgs Portal” mixing



These fields completely neutral: “Neutral Naturalness”

LHC production

top partner spin

		<i>scalar</i>	<i>fermion</i>
<i>strong direct production</i> {	<i>QCD</i>	SUSY	Composite Higgs/ RS
<i>DY direct production</i> {	<i>EW</i>	folded SUSY	Quirky Little Higgs
<i>Higgs portal direct production</i> {	<i>singlet</i>	hyperbolic Higgs/ accidental Susy*	Twin Higgs

Mirror Glueballs

Higgs portal observables

Higgs coupling shifts

~ tuning

* Cohen... '18, Cheng, Li, Salvioni, ... '18

Why not?

Higgs portal maintains equilibrium down to $T \sim \text{GeV}$

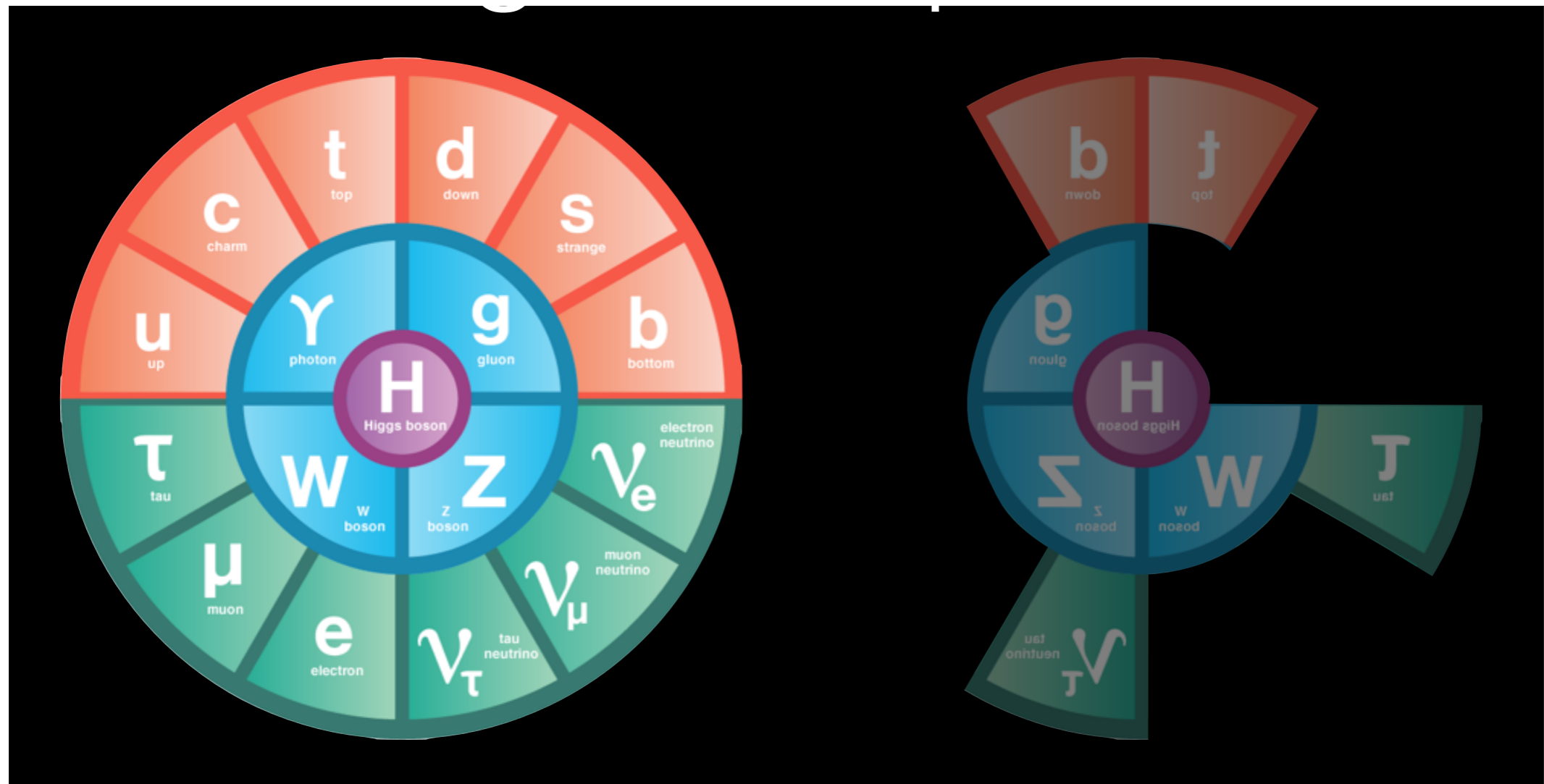
$$\Delta N_{\text{eff}} \gg 1$$

$$\Delta N_{\text{eff}} \lesssim 0.6$$

This is excluded by CMB measurements...

- 1) change cosmology
- 2) change spectrum

only 3rd generation twinned

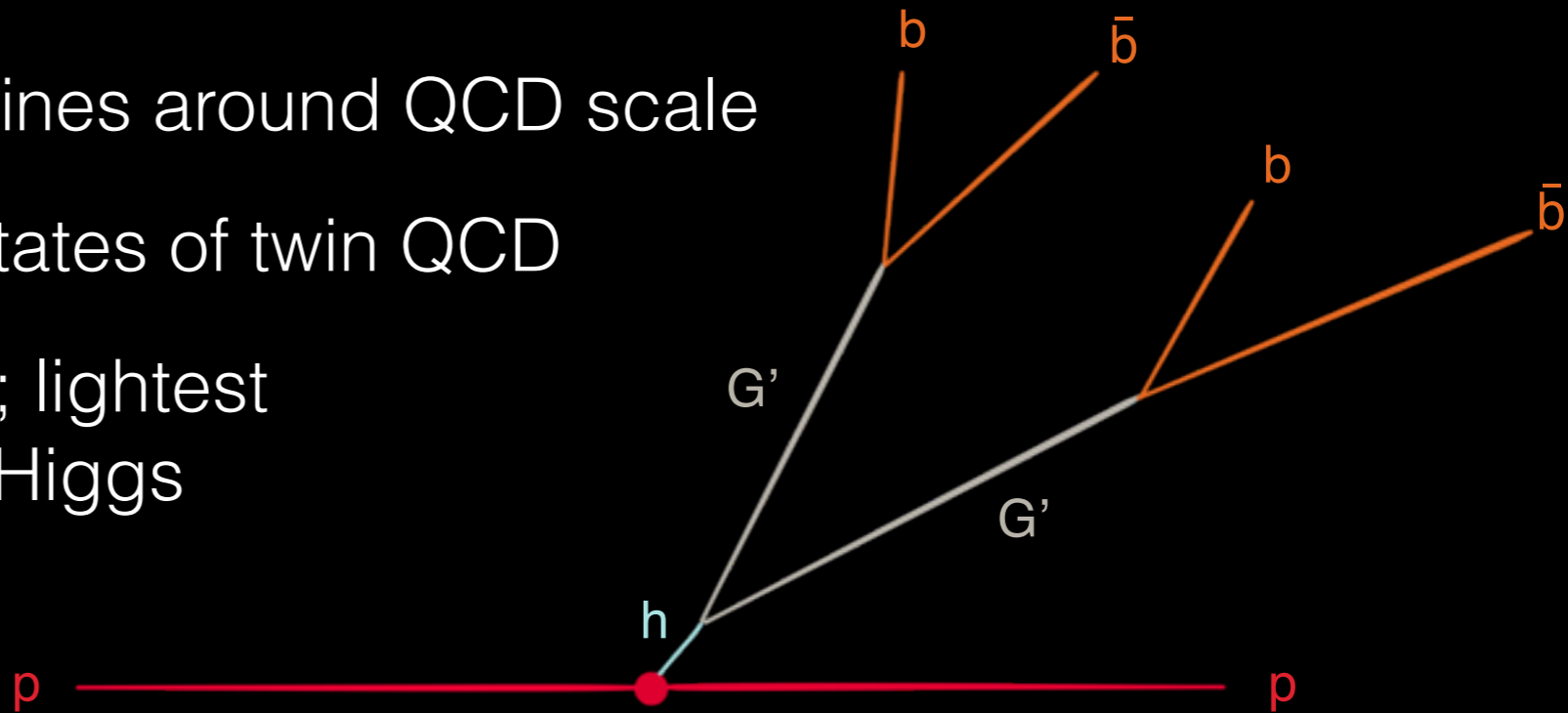


“fraternal twin Higgs”

Exotic Higgs Decays

- Must have twin QCD, confines around QCD scale
- Higgs couples to bound states of twin QCD
- Glueballs most interesting; lightest have same quantum # as Higgs

$$\mathcal{L} \supset \frac{v}{f} \frac{h}{f} G'_{\mu\nu} G'^{\mu\nu}$$



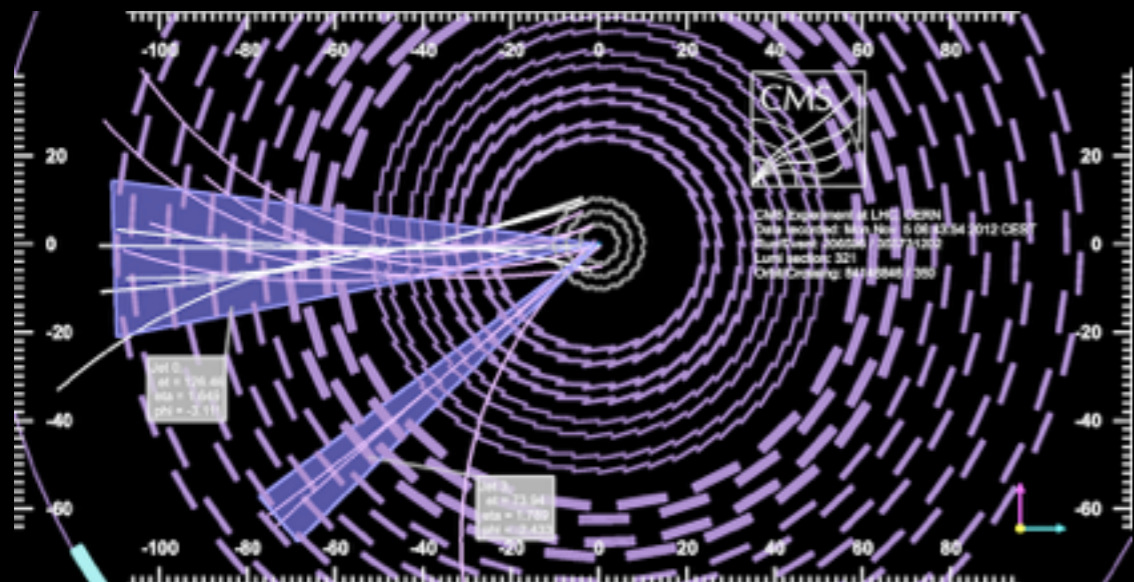
Produce in rare Higgs decays ($\text{BR} \sim 10^{-3} - 10^{-4}$)

$$gg \rightarrow h \rightarrow 0^{++} + 0^{++} + \dots$$

Decay back to SM via Higgs

$$0^{++} \rightarrow h^* \rightarrow f\bar{f}$$

Long-lived, length scale \sim LHC detectors
Hidden Valley signature [Strassler, Zurek '06]

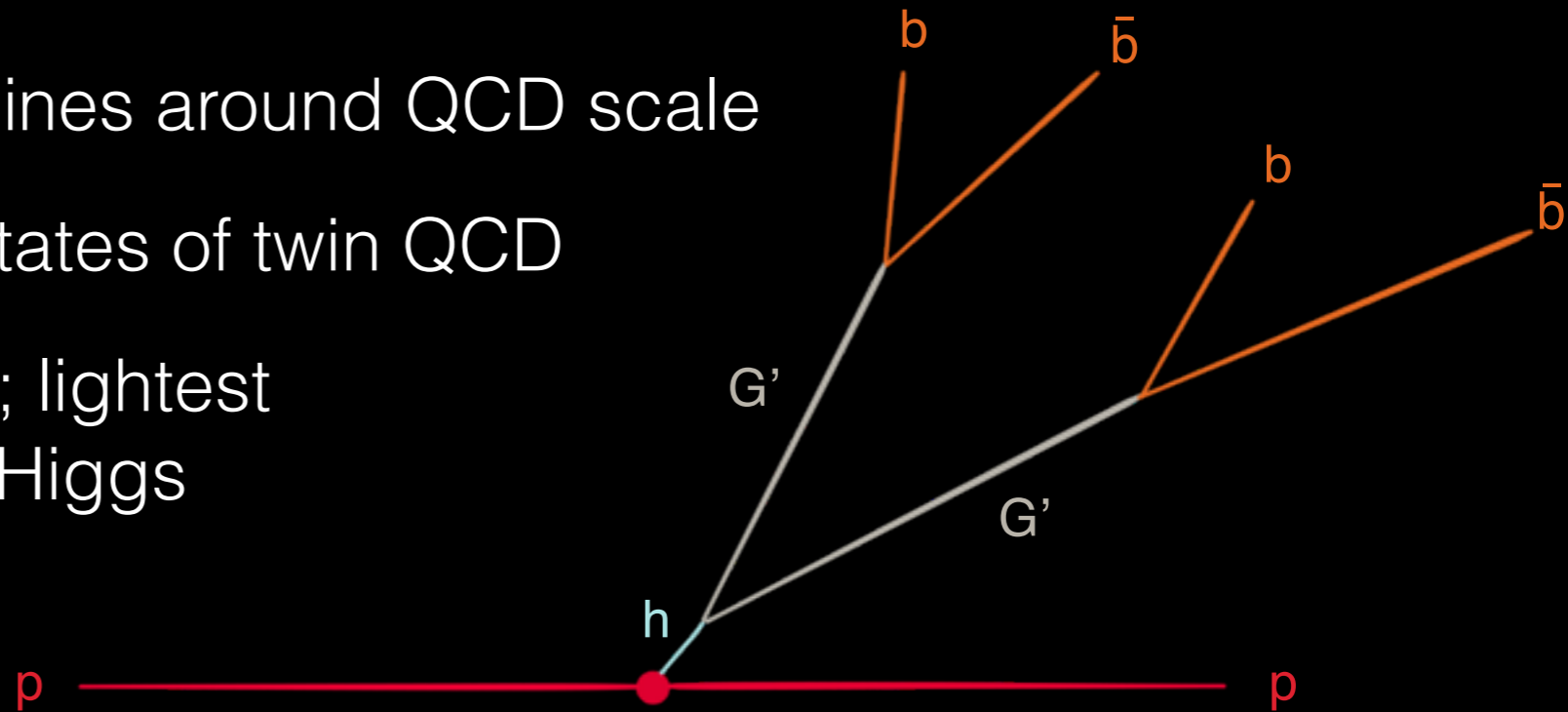


important signature

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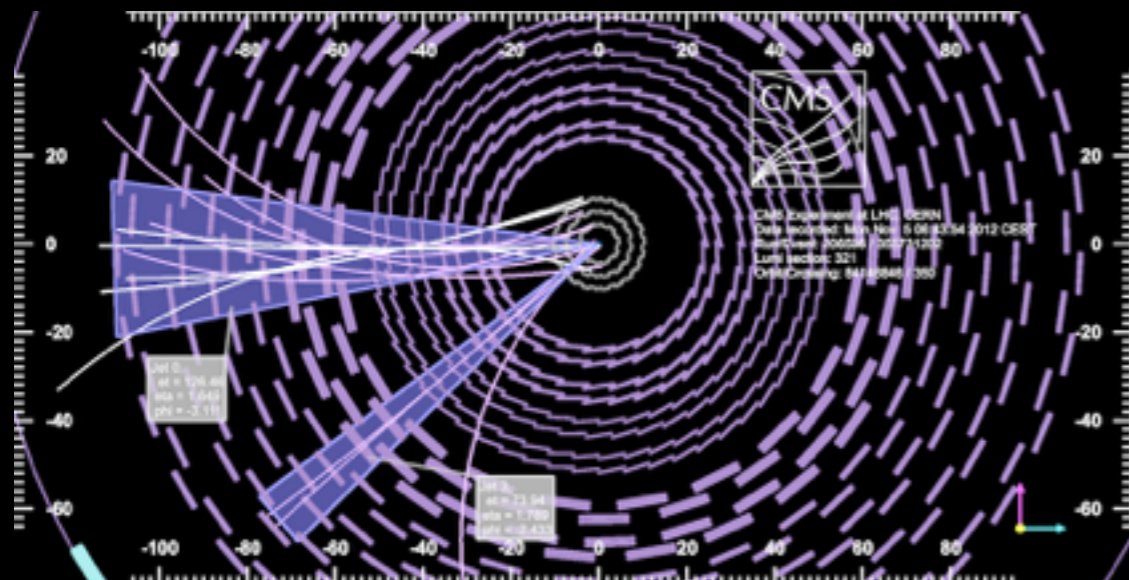
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important signature



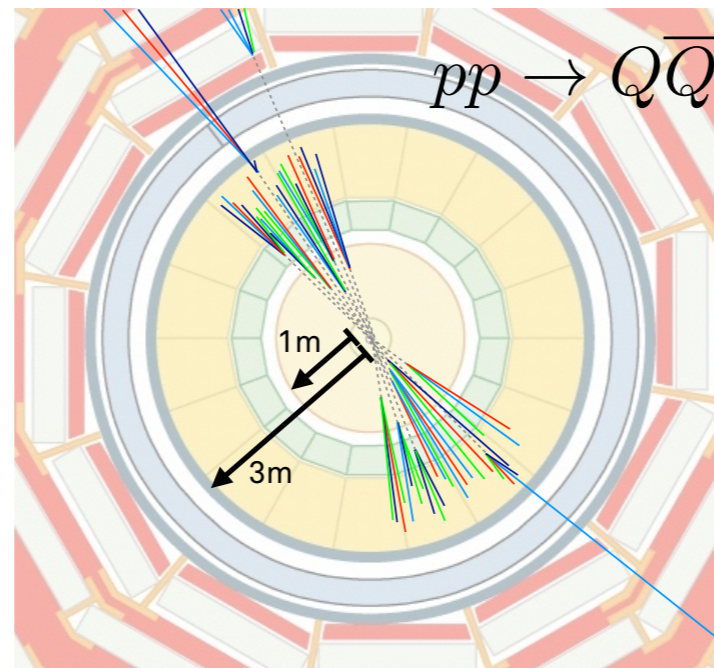
LLPs

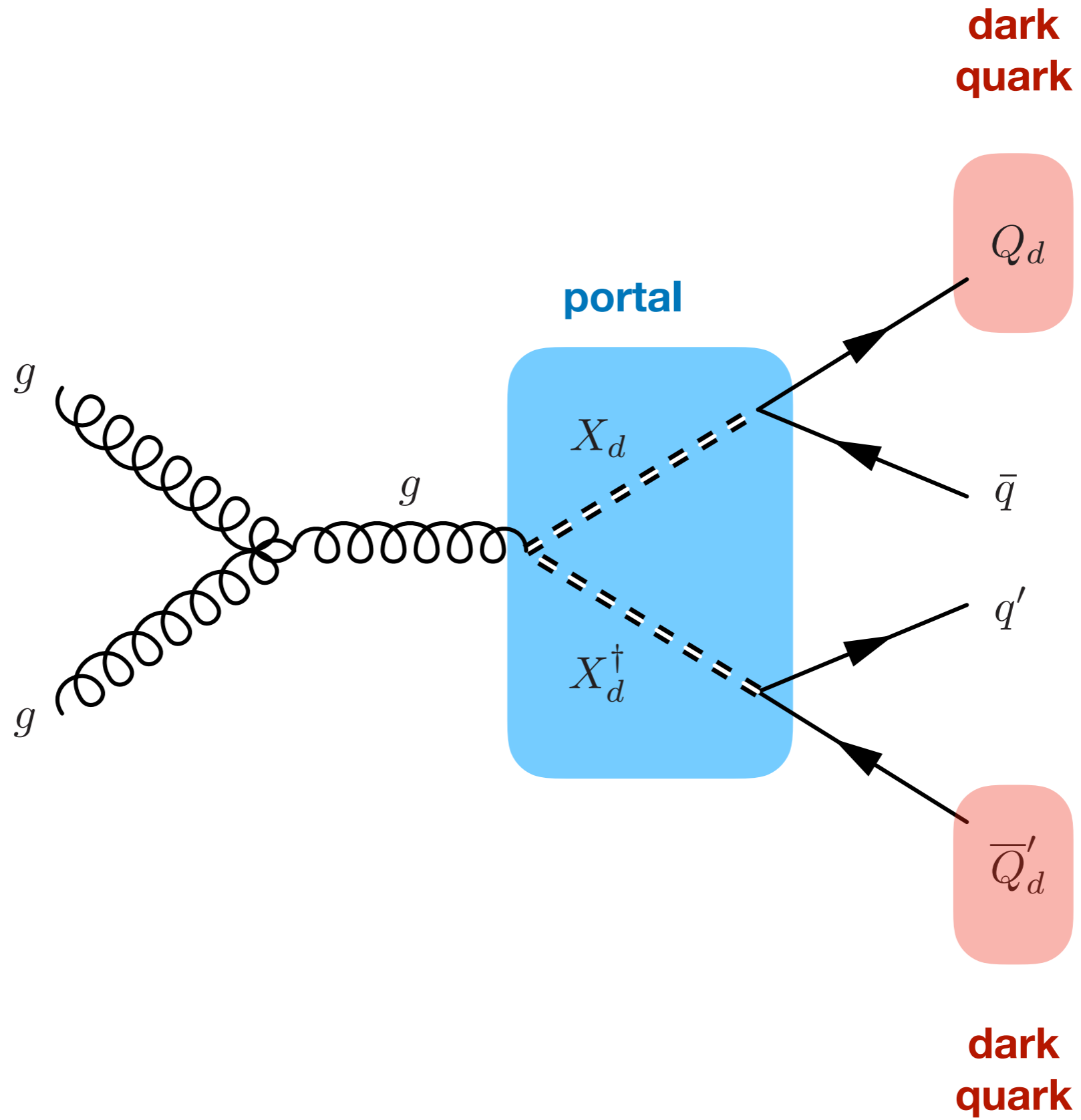
- Long Lived, length scale: LHC detectors
- New confining gauge group
- These ingredients are present in many new BSM extensions, e.g. “emerging jets”
- **Emerging jets**: motivated by asymmetric DM, co genesis of baryon asymmetry and DM

Emerging Jets

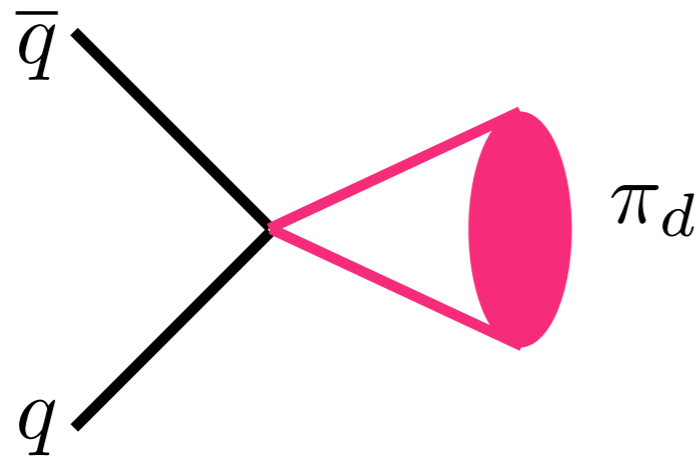
Schwaller, Stolarski, **AW** '15

NEW: First search by **CMS PAS EXO-18-001**





Dark Pion Lifetime



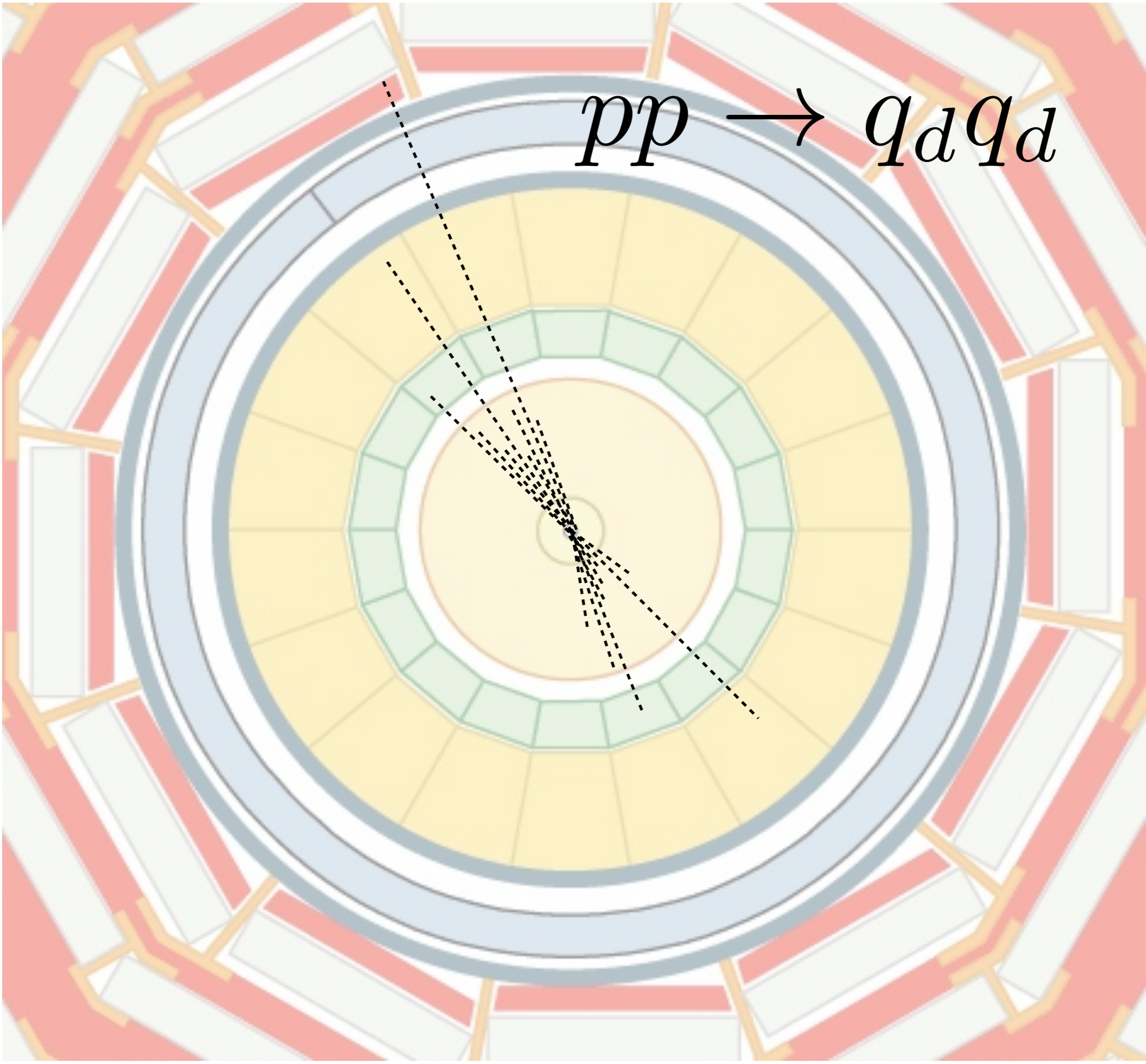
$$\frac{1}{M_X^2} \bar{Q} \gamma_\mu Q \bar{d}_R \gamma^\mu d_R$$

Use chiral Lagrangian to estimate

$$\Gamma(\pi_d \rightarrow \bar{d}d) \approx \frac{f_{\pi_d}^2 m_d^2}{32\pi M_{X_d}^4} m_{\pi_d}$$

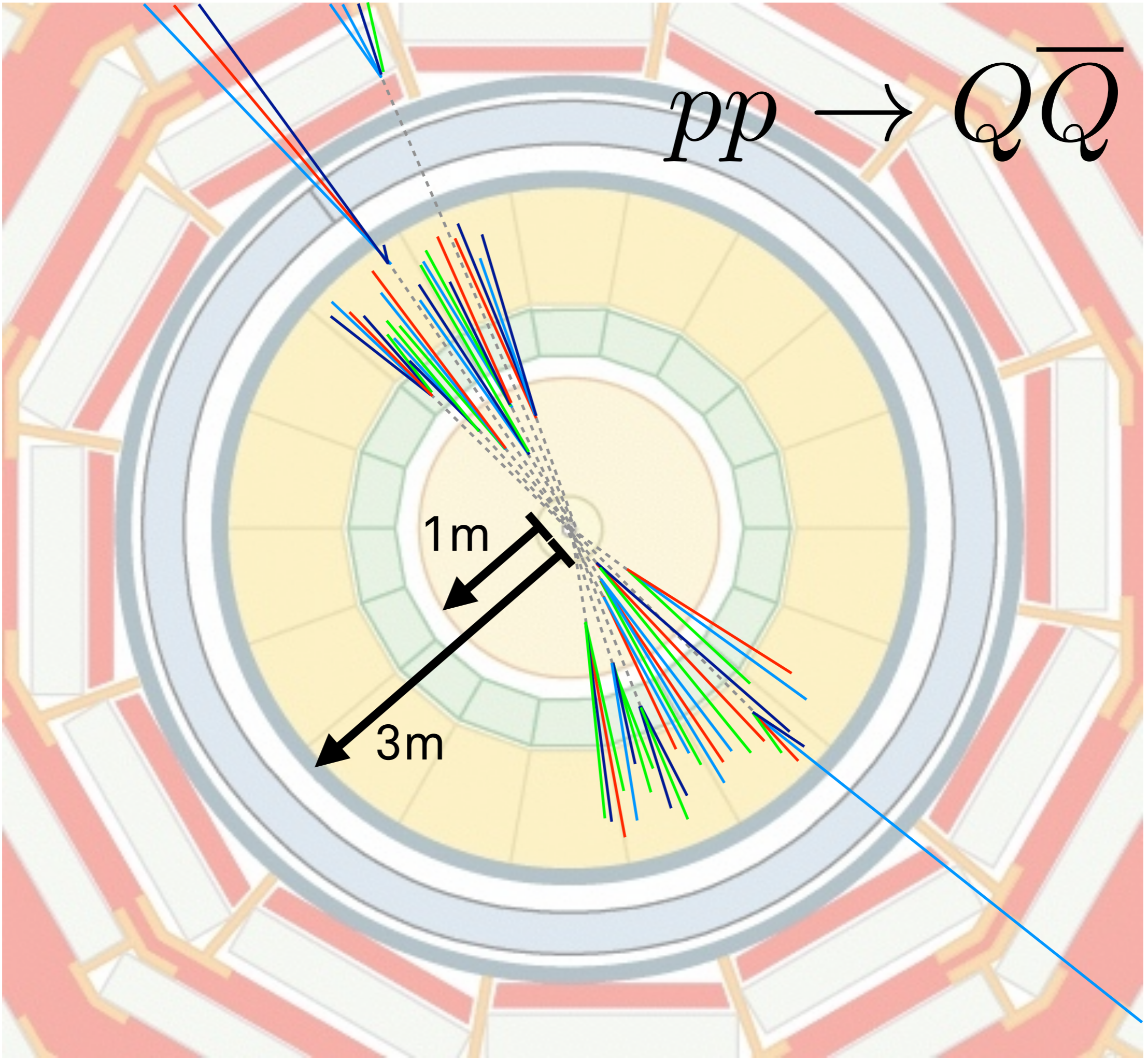
$$c\tau \approx 5 \text{ cm} \times \left(\frac{1 \text{ GeV}}{f_{\pi_d}} \right)^2 \left(\frac{100 \text{ MeV}}{m_d} \right)^2 \left(\frac{1 \text{ GeV}}{m_{\pi_d}} \right) \left(\frac{M_{X_d}}{1 \text{ TeV}} \right)^4$$

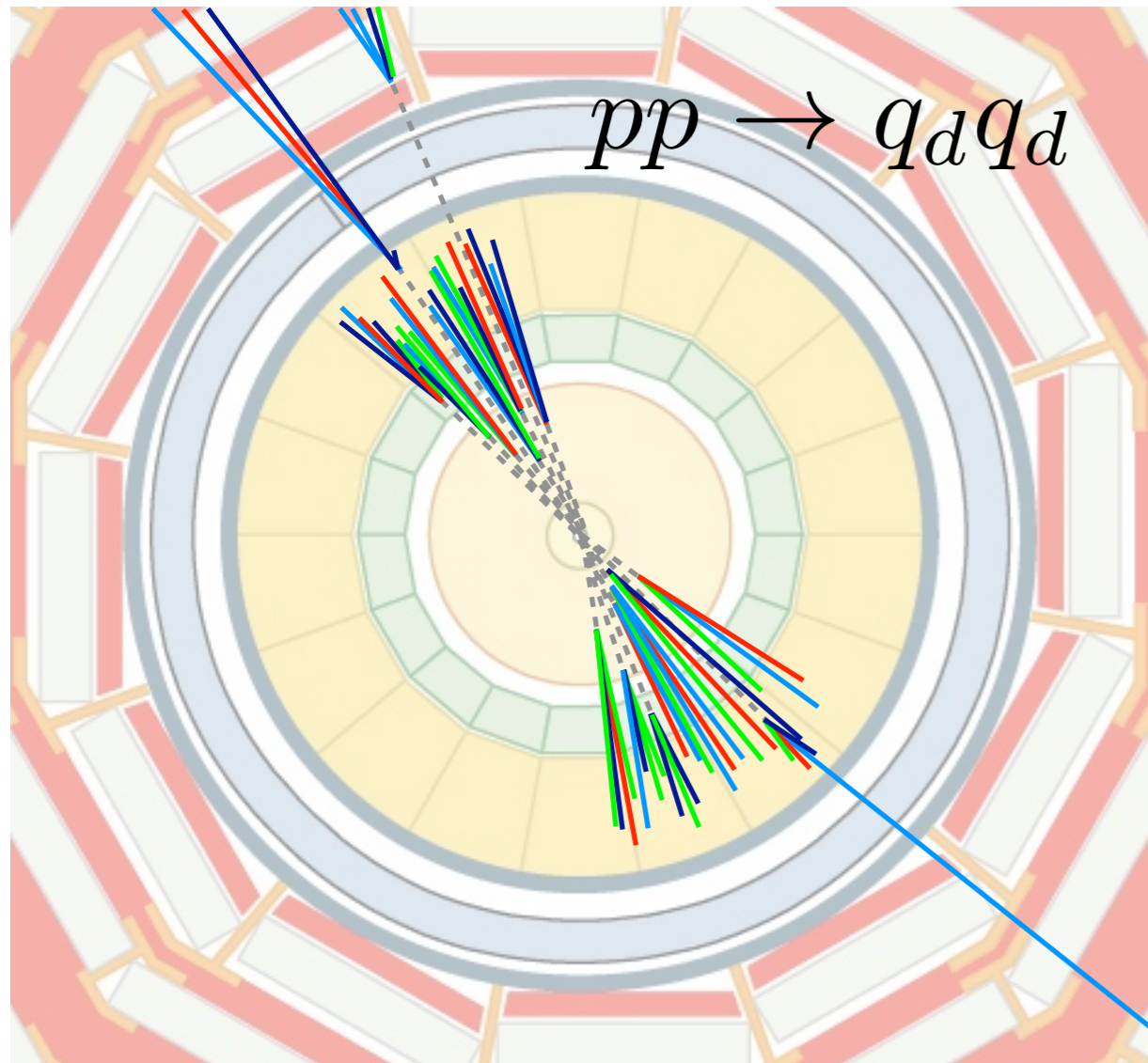
$pp \rightarrow qdqd$



$pp \rightarrow Q\bar{Q}$

1m
3m

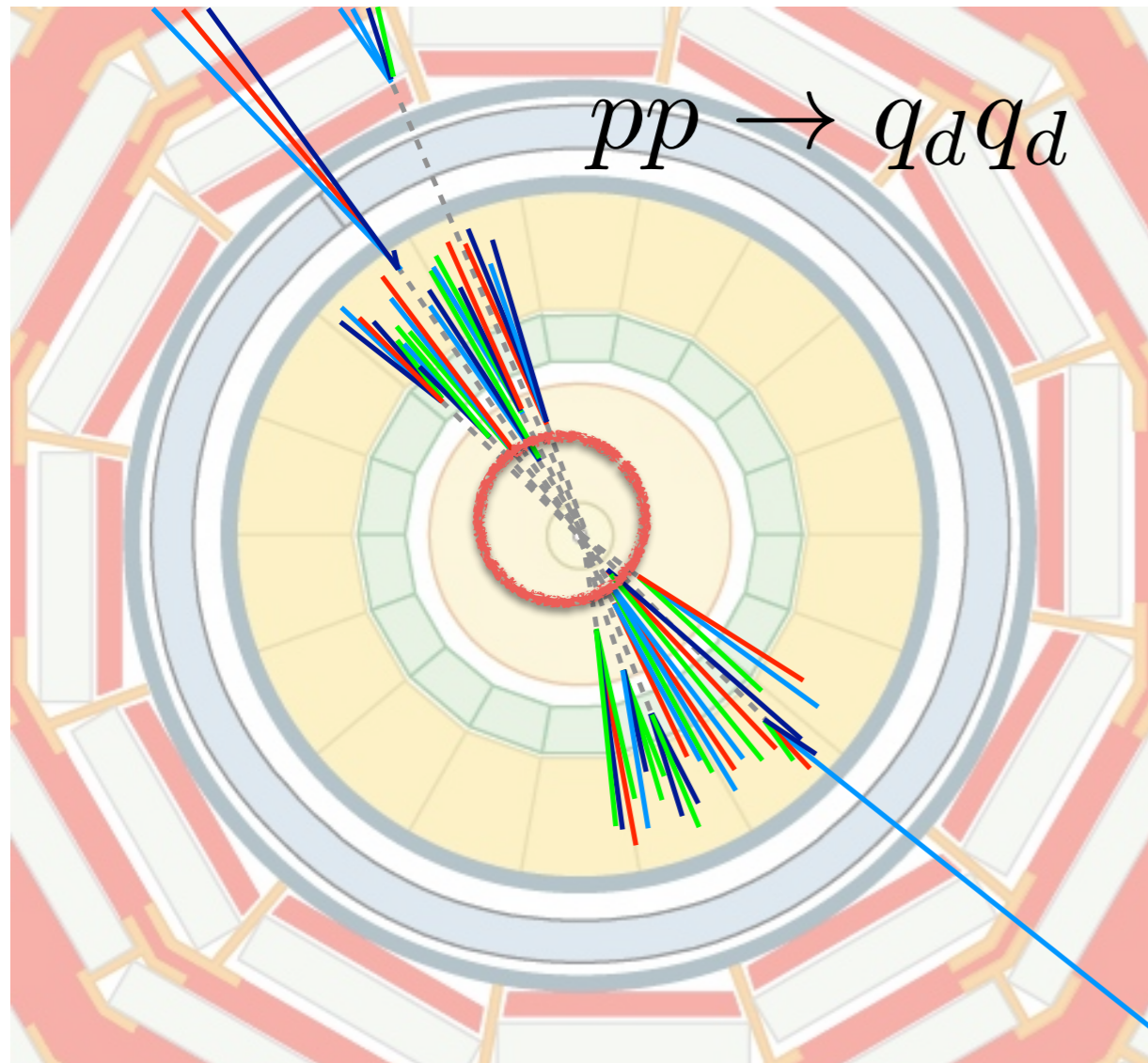




Decay lifetime of \sim cm

Exponential decay profile: Several displaced vertices inside a jet “cone” (or calo-jet)

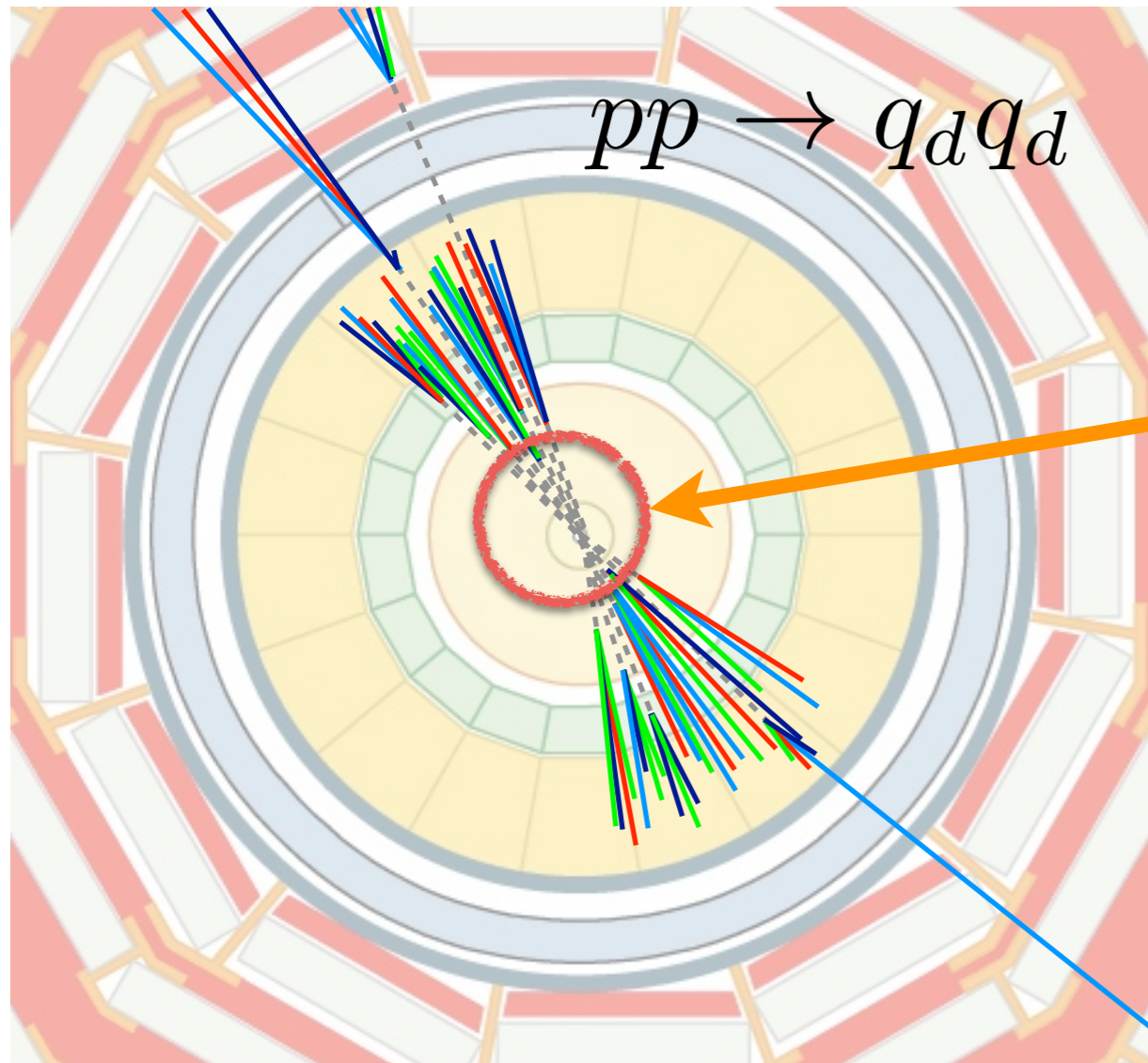
No/few tracks originating from interaction point



Look for Hcal-jets with no/few tracks below distance to interaction point (inside **circle**)

New **'track-less'** signature

Universal for a large class of displaced physics



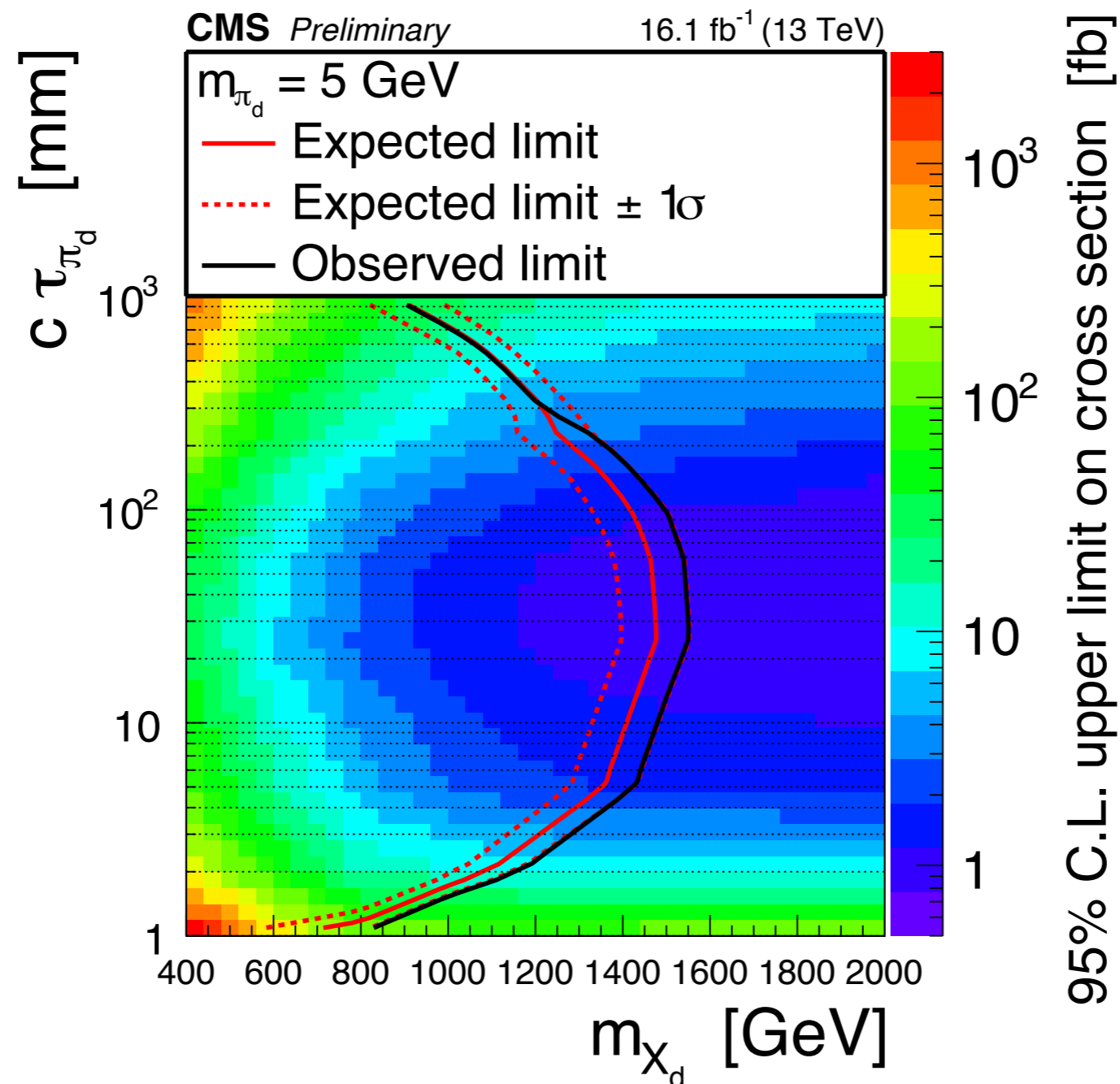
Look for Hcal-jets with no/few tracks below distance to interaction point (inside **circle**)

New **‘track-less’** signature

Universal for a large class of displaced physics

Emerging jets search

“Mediator particles with masses between 400 and 1250 GeV are excluded for dark hadron decay lengths between 5 and 225 mm.”



[CMS PAS EXO-18-001]

Amazing work by UMD CMS team (Belloni, Eno, Jeng, ...)

Relaxion paradigm

P.W. Graham, D.E. Kaplan, S.Rajendran '15
(earlier work by Abbott 85, G.Dvali, A.Vilenkin 04, G.Dvali 06)

A **technically natural** solution to the hierarchy problem

Uses dynamics, not symmetries

At the drafting stage, but currently extensive development

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$$m^2 |H|^2$$

Higgs mass

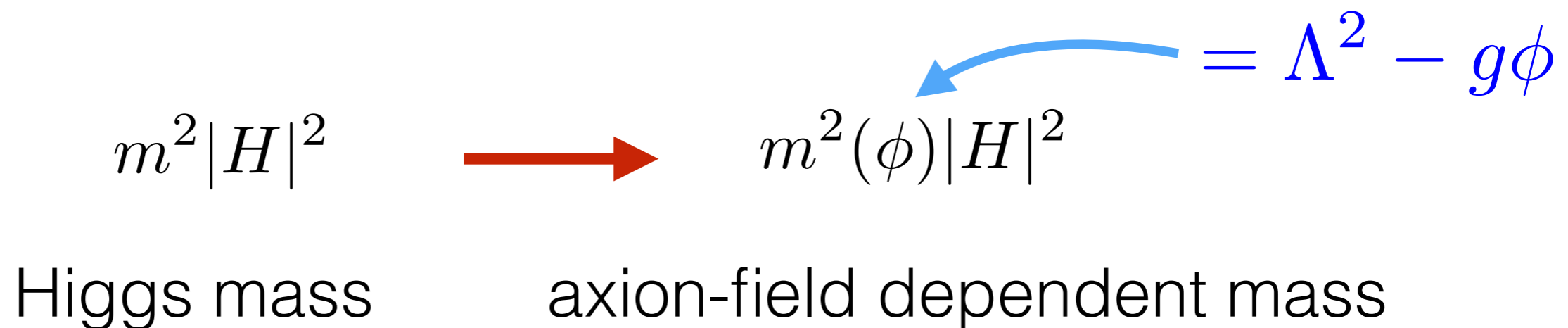
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Uses dynamics, not symmetries

At the drafting stage, but currently extensive development

$$m^2|H|^2 \quad \longrightarrow \quad m^2(\phi)|H|^2 \quad = \Lambda^2 - g\phi$$

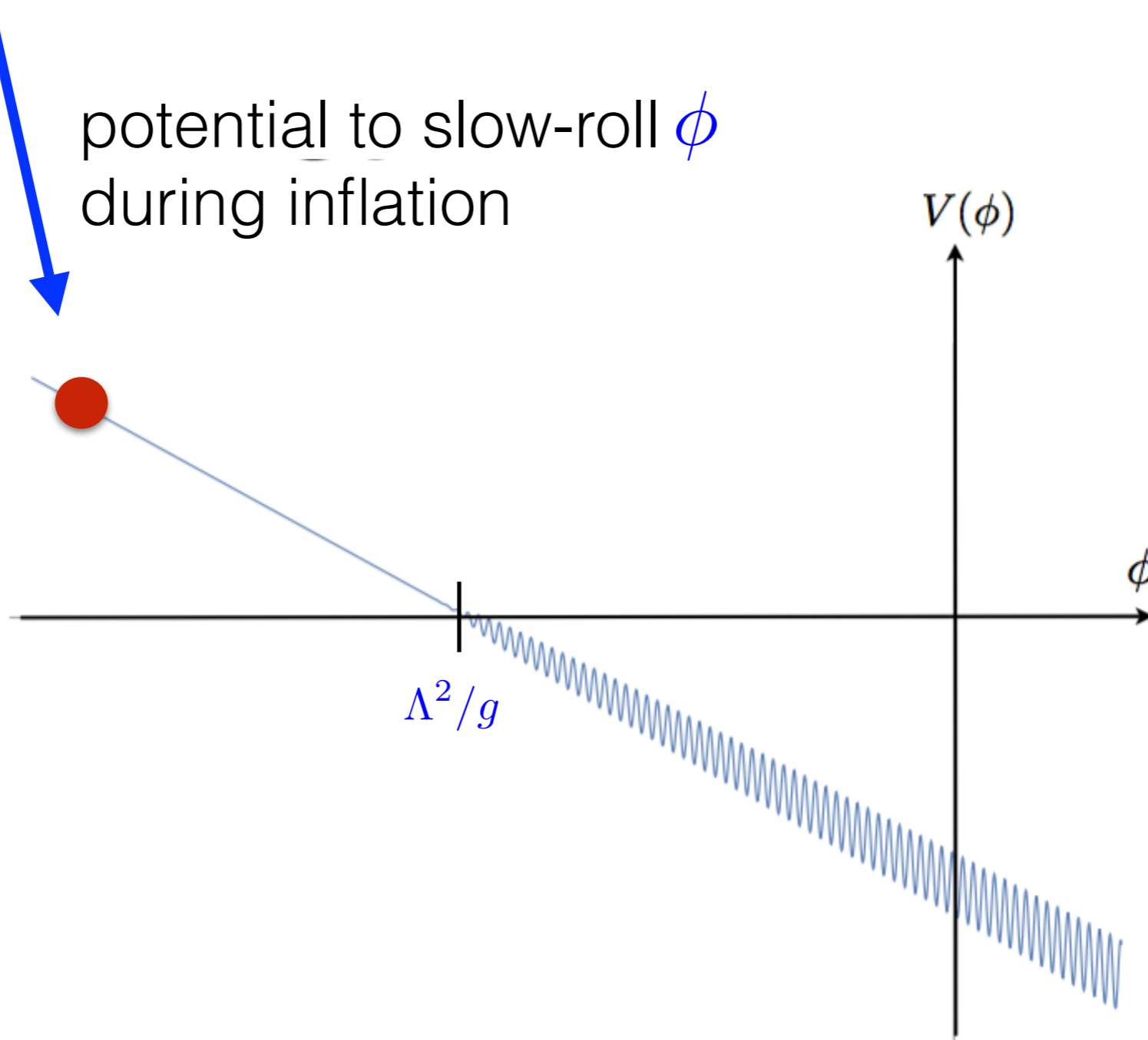
Higgs mass

axion-field dependent mass

Clever dynamics stabilizes ϕ at values: $m^2(\phi) \ll \Lambda^2$

$$V(g\phi) + (\Lambda^2 - g\phi)|H|^2 + \epsilon\Lambda_{QCD}^3 h \cos \frac{\phi}{f}$$

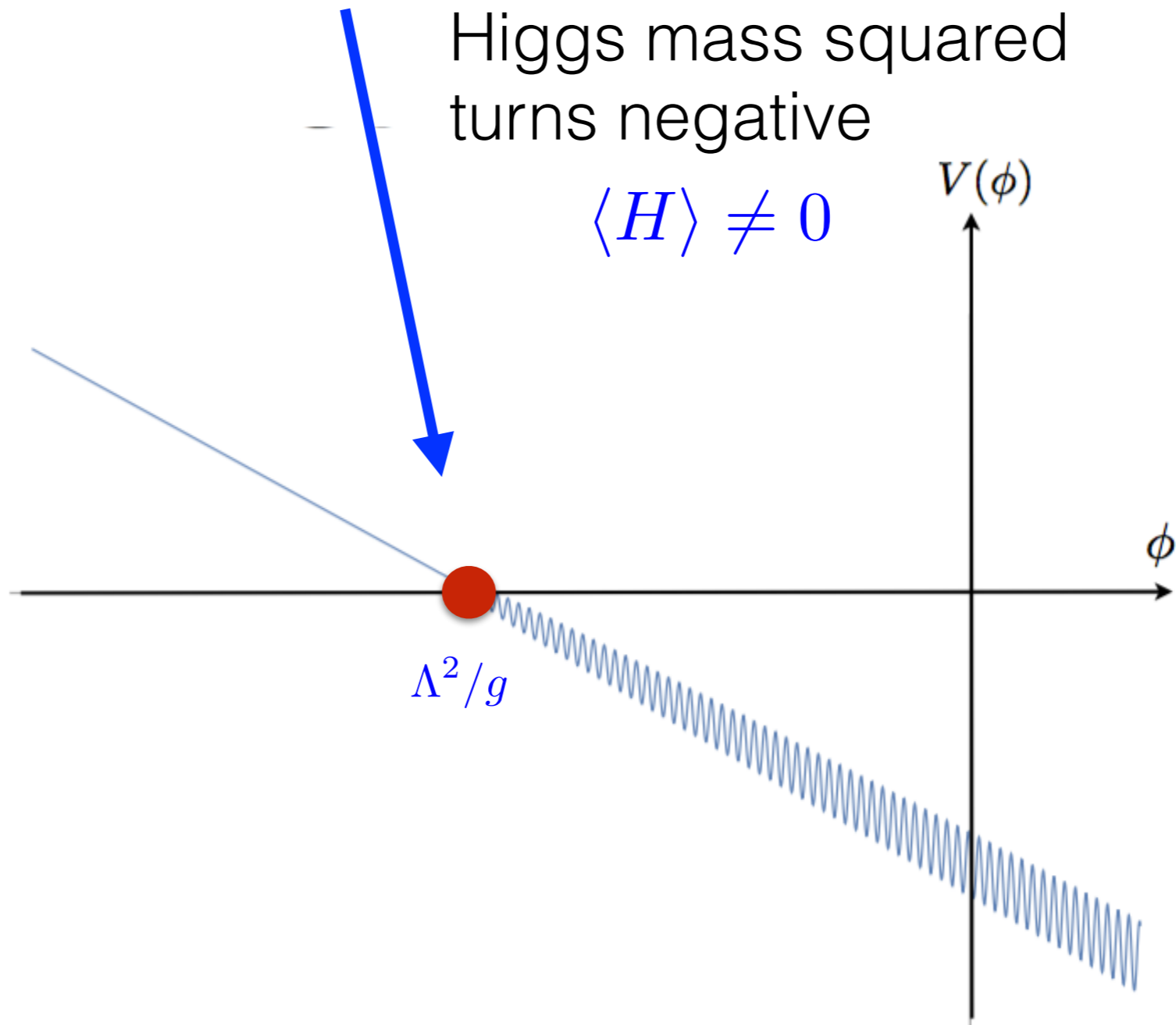
potential to slow-roll ϕ
during inflation



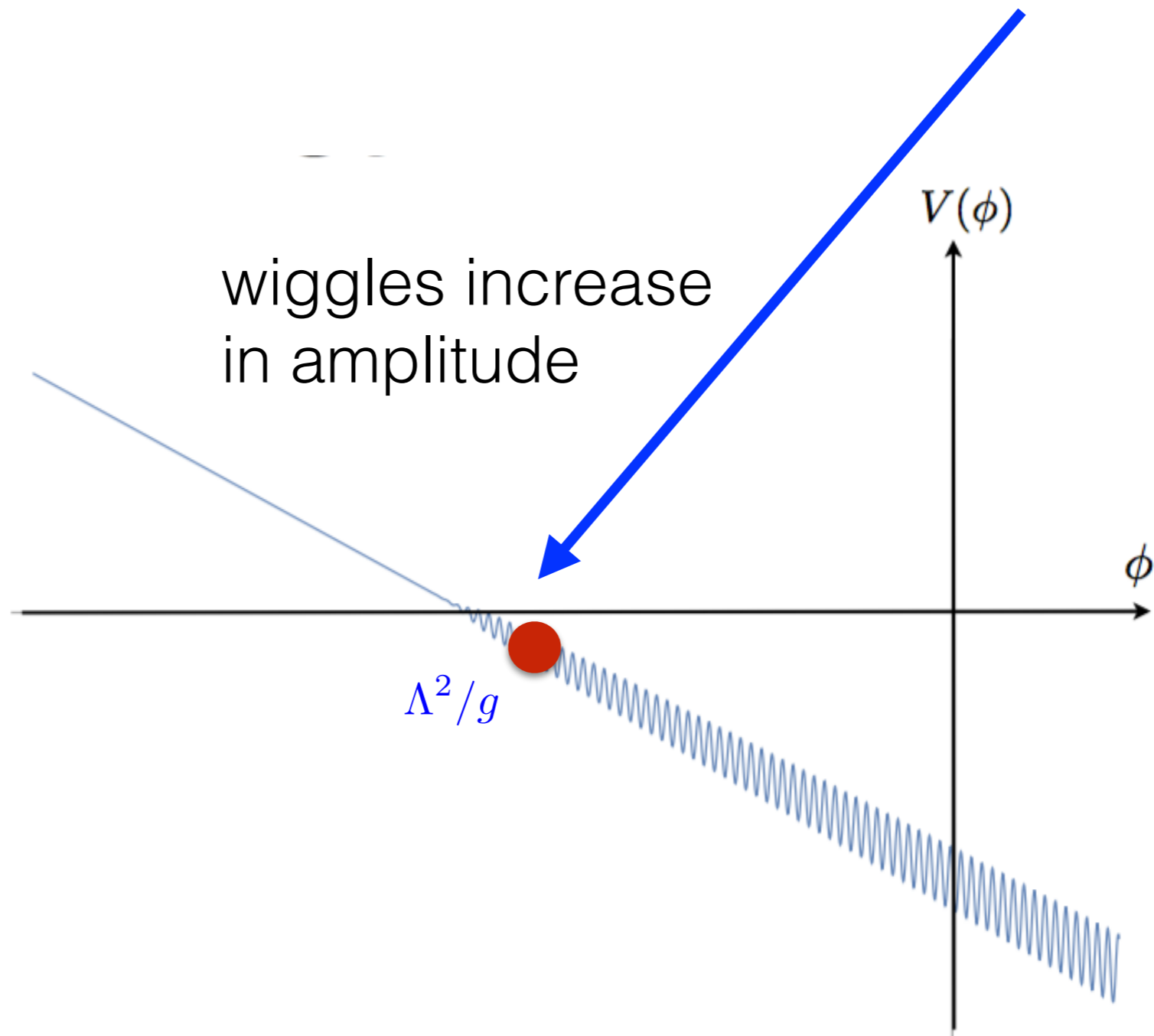
$$V(g\phi) + (\Lambda^2 - g\phi)|H|^2 + \epsilon\Lambda_{QCD}^3 h \cos \frac{\phi}{f}$$

Higgs mass squared
turns negative

$$\langle H \rangle \neq 0$$

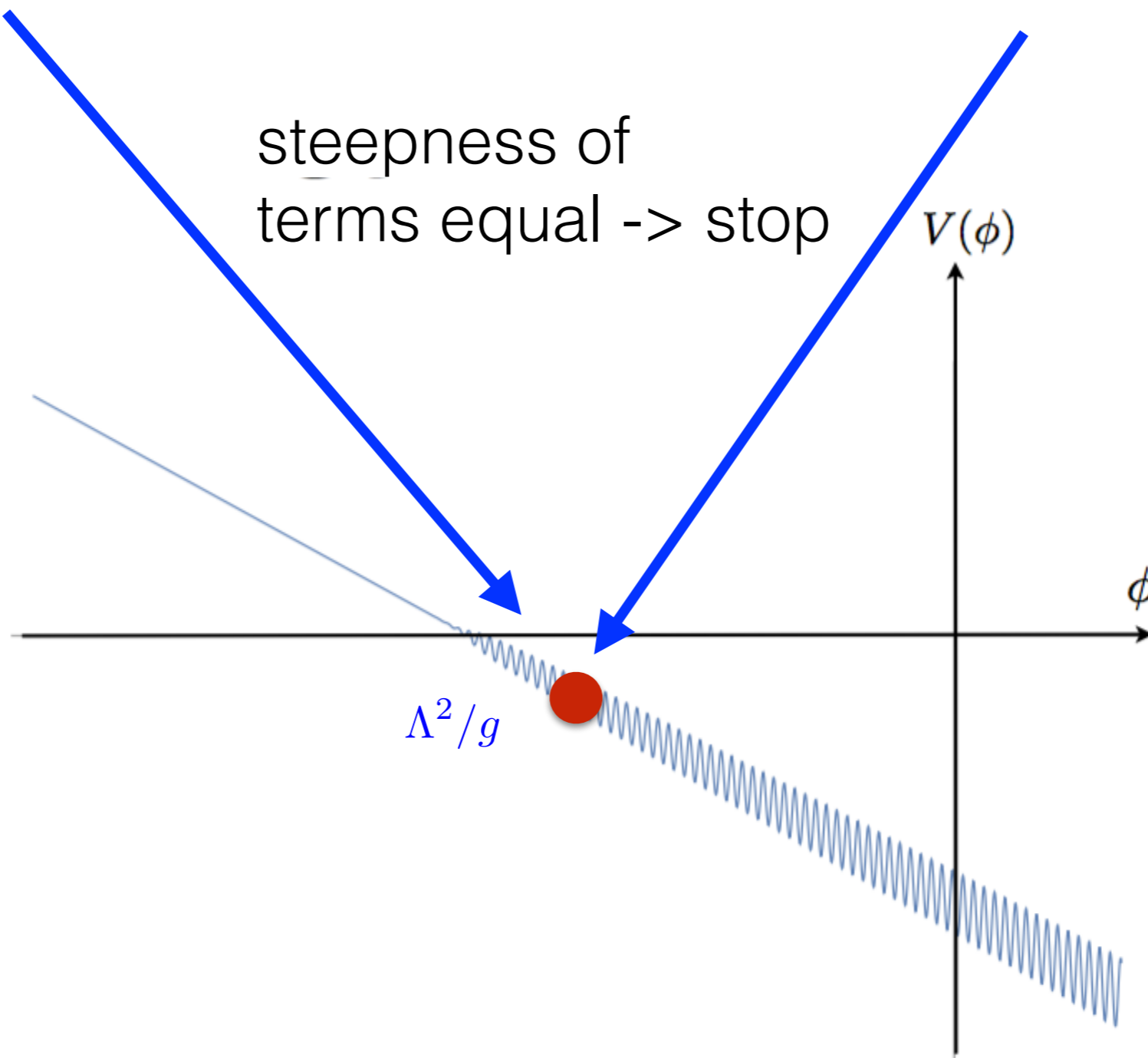


$$V(g\phi) + (\Lambda^2 - g\phi)|H|^2 + \epsilon\Lambda_{QCD}^3 h \cos \frac{\phi}{f}$$



$$V(g\phi) + (\Lambda^2 - g\phi)|H|^2 + \epsilon\Lambda_{QCD}^3 h \cos \frac{\phi}{f}$$

steepness of
terms equal \rightarrow stop



$$V(g\phi) + (\Lambda^2 - g\phi)|H|^2 + \epsilon\Lambda^3_{QCD}h \cos \frac{\phi}{f}$$

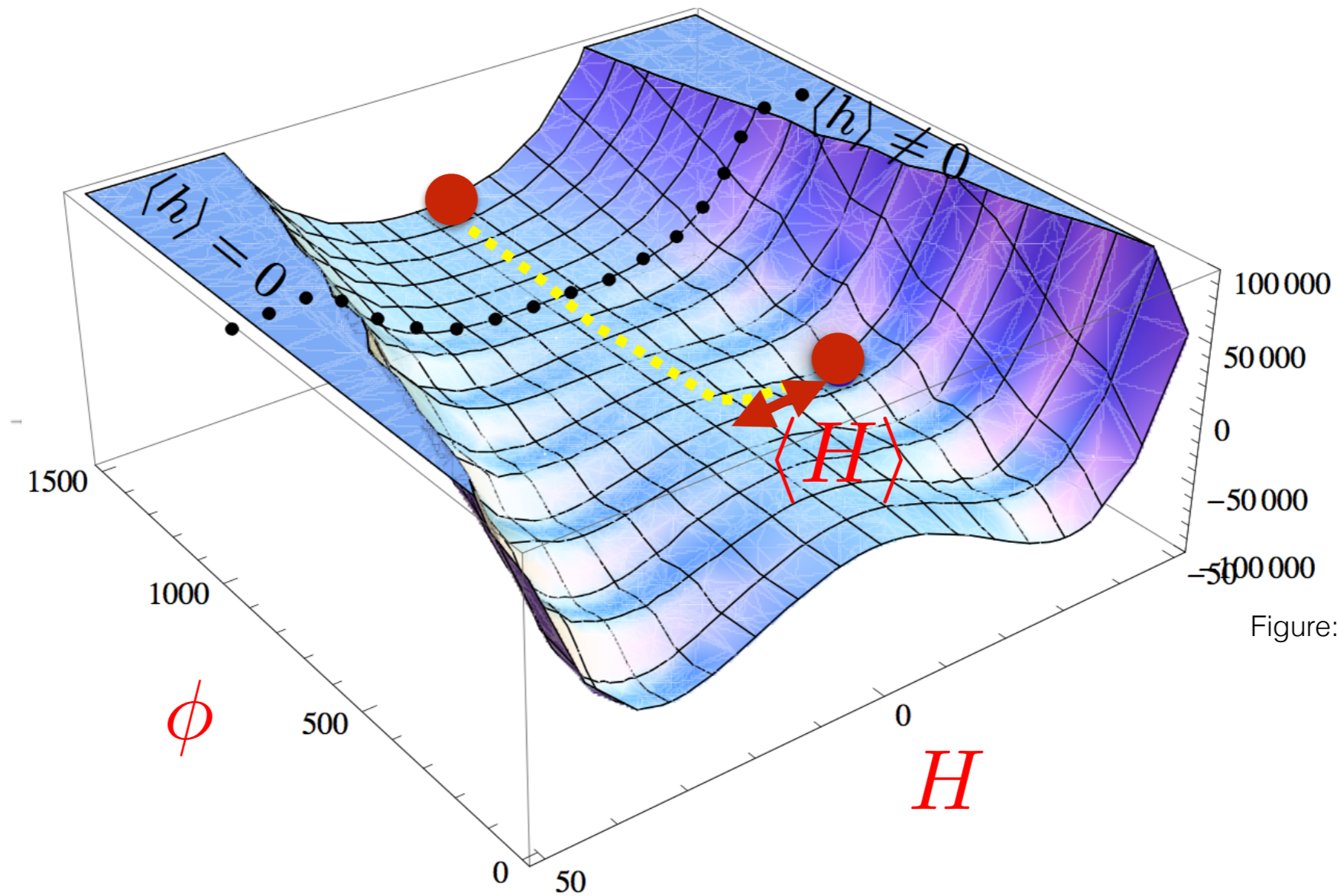


Figure: C. Grojean

- * QCD axion doesn't work: $\theta_{QCD} \sim 1$ due to tilt
- * Add new QCD' group => new weak-scale signals!
- * Add additional scanning field => **no** collider signals!

Espinosa, Grojean, Panico, Pomarol, Pujolas, Servant '15

More work: Hardy '15; Gupta et al '15; Batell, Giudice, McCullough '15; Choi, Im '15; Kaplan, Rattazzi '15; Di Chiara et al. '15; Ibanez et al. '15; Hook, Marques-Tavares '16; Nelson, Prescod-Weinstein '17; ...]

Some points of concern:

$$g \sim 10^{-27} \text{ GeV}$$

UV completion !?

$$N > H^2 / g^2 \sim 10^{45}$$

inflation !?

$$\Delta\Phi \simeq 10^{41} \text{ GeV}$$

large field excursions

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YEP, THAT'S HIM ALL RIGHT...
THAT'S THE GUY I MUGGED



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WHAMOND



Conclusions

- No signs of new physics have appeared so far.
- The Higgs fine-tuning puzzle is as puzzling as ever. Do we simply live in a (mildly?) fine-tuned universe? Or is there a subtle solution?
- Themes of recent years: search for electroweak or neutral new particles at colliders to exhaust possibilities; intriguing possibilities for connections of the weak scale with cosmology.
- Amazing landscape of experiments: LHC, dark matter, EDMs, flavor physics. New physics discovery could come at any time!