

Dynamical energy loss: exploring the QGP with high pt theory and data

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Motivation

- **Energy loss of high-pt particles traversing QCD medium is an excellent probe of QGP properties.**
- **Theoretical predictions can be compared with a wide range of data, coming from different experiments, collision systems, collision energies, centralities, observables...**
- **Can be used together with low-pt theory and experiments to study the properties of created QCD medium, i.e. for precision QGP tomography.**

The dynamical energy loss formalism

- **Finite size medium of dynamical (moving) partons**
 - **Based on finite T field theory and generalized HTL approach**
M. D., PRC74 (2006), PRC 80 (2009), M. D. and U. Heinz, PRL 101 (2008).



- **Same theoretical framework for both radiative and collisional energy loss**
- **Finite magnetic mass effects** (M. D. and M. Djordjevic, PLB 709:229 (2012))
 - **Running coupling** (M. D. and M. Djordjevic, PLB 734, 286 (2014)).

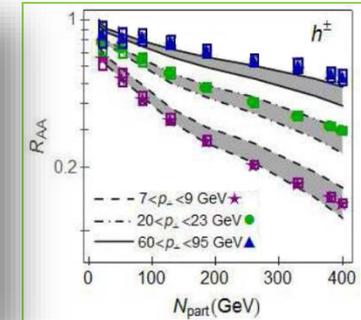
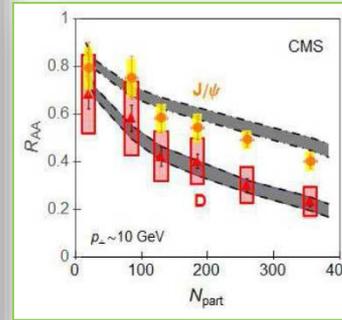
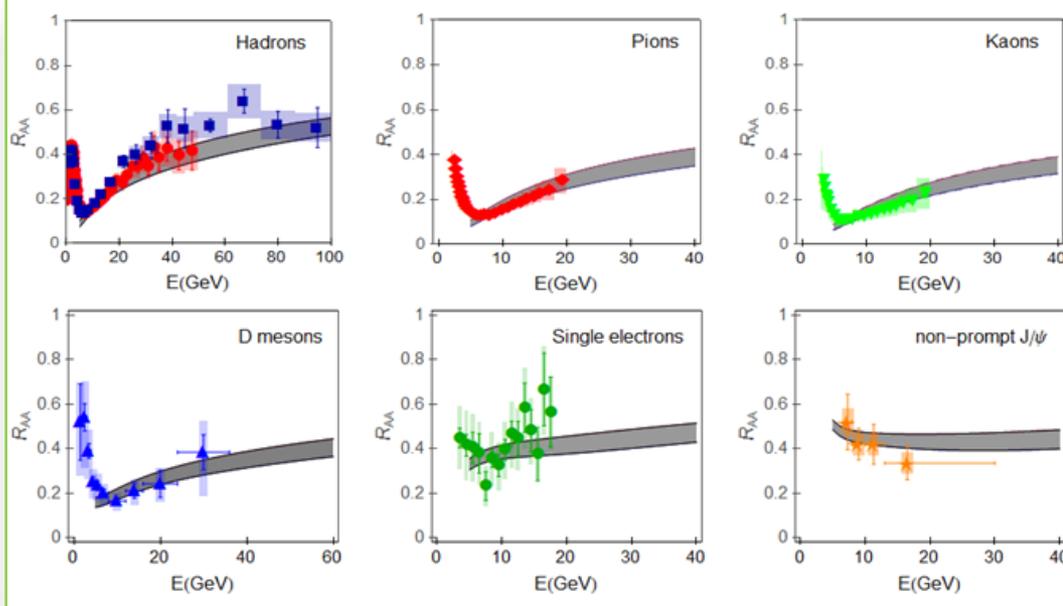
Most recently: Relaxed soft-gluon approximation, see poster 87 by B. Ilic (Blagojevic)
(B. Blagojevic, M.D. and M. Djordjevic, PRC 99, 024901 (2019))



Integrated in a numerical procedure including parton production, fragmentation functions, path-length and multi-gluon fluctuations



- **No fitting parameters**
- **Treats both light and heavy flavor partons**



- Explains R_{AA} for different probes, collision energies, and centralities.
- Resolved the longstanding “heavy flavour puzzles at RHIC and LHC”.
- Good agreement with subsequent measurements.
- Clear predictions for future experiments.
- Agreement obtained by the same model and parameter set, no fitting parameters introduced.
- All steps in the suppression scheme are important, and have to be kept in all future framework developments.



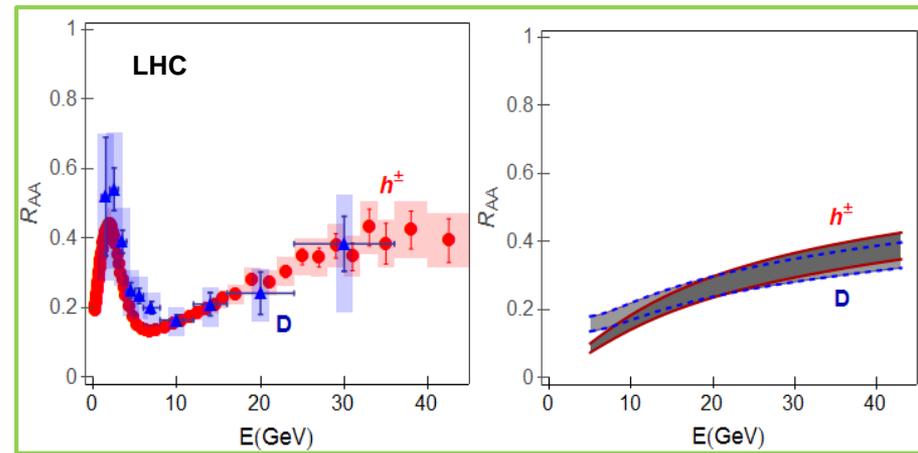
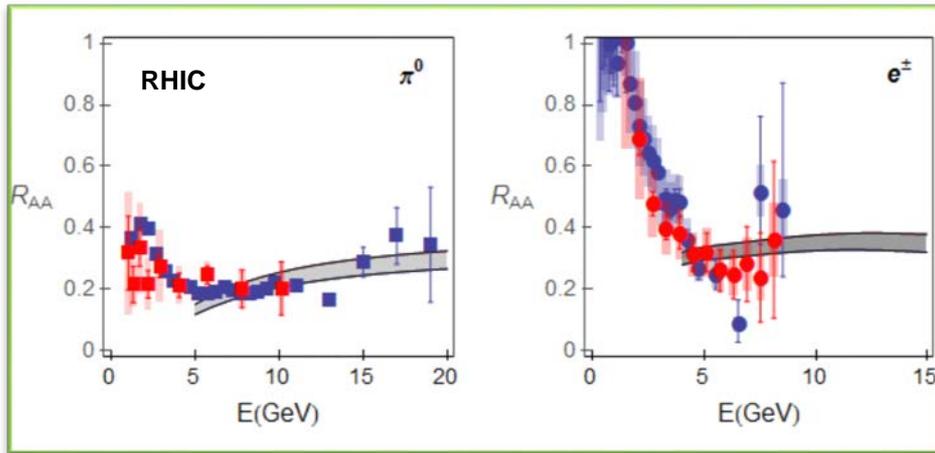
A realistic description
parton-medium interactions.



However, until recently, the model
did *not* include QGP evolution.



Predictions only for the
observables weakly sensitive
to QGP evolution (i.e. R_{AA}).



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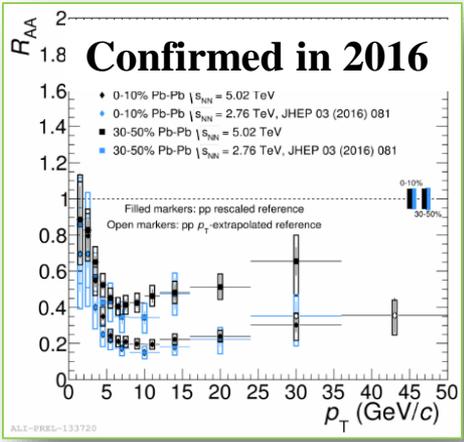
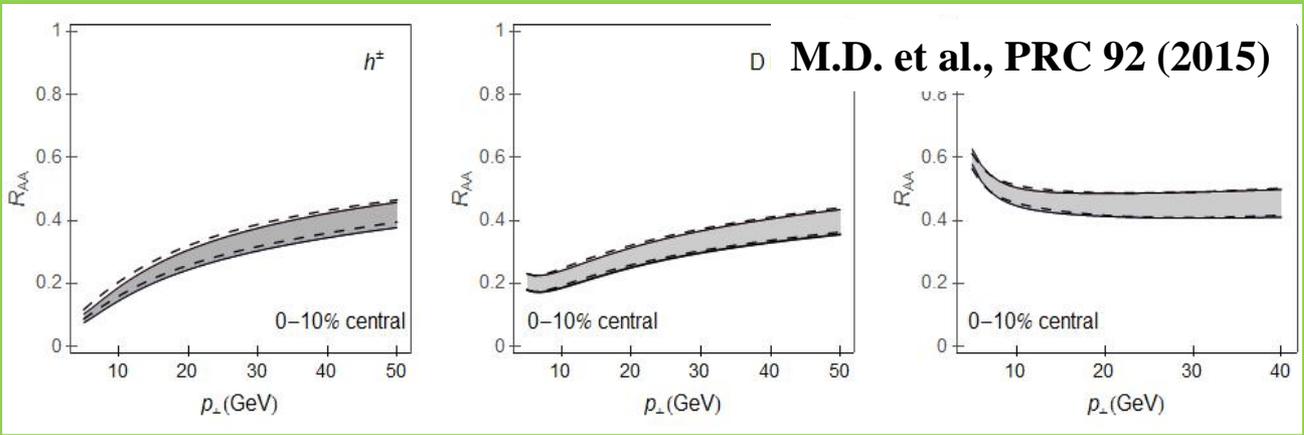
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Goals

- **Allow systematic comparison of experimental data and theoretical predictions, obtained by the same formalism and the same parameter set. In particular:**
 - For different observables (e.g. R_{AA} and v_2)
 - Different collision systems (e.g. Pb+Pb, Xe+Xe, etc.)
 - Different probes (light and heavy)
 - Different collision energies
 - Different centralities
- **Introduce medium evolution in the model, for now through Bjorken expansion.**
- **Better understand energy loss mechanisms.**
- **Infer the bulk QGP medium properties from high pt data.**

DREENA framework

DREENA (**D**ynamical **R**adiative and **E**lastic **E**nergy loss **A**pproach) is a computational framework in which dynamical energy loss is implemented.

Version **C** – **C**onstant temperature medium

D. Zigic, I. Salom, J. Auvinen, M. Djordjevic and M.D., JPG (2019), in press

Version **B** – **1+1D Bjorken expansion**

D. Zigic, I. Salom, J. Auvinen, M. Djordjevic and M.D., PLB 791 (2019) 236

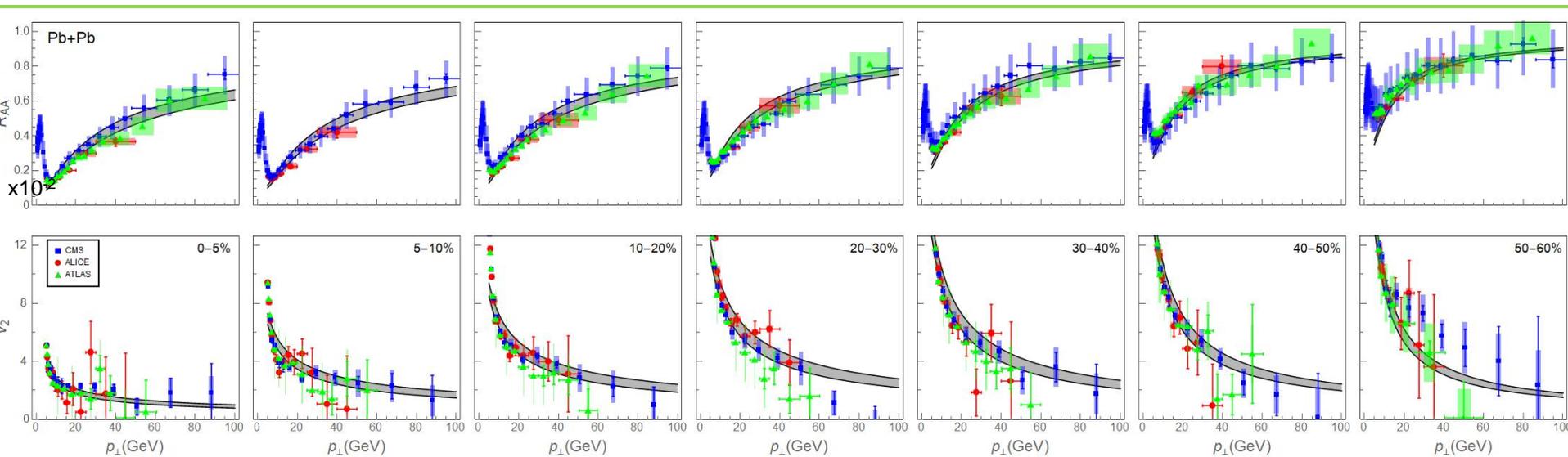
Outline:

- I. Comparison with experimental data
- II. Differentiate between different energy loss mechanisms
- III. Infer the shape of the QGP droplet from high pt data

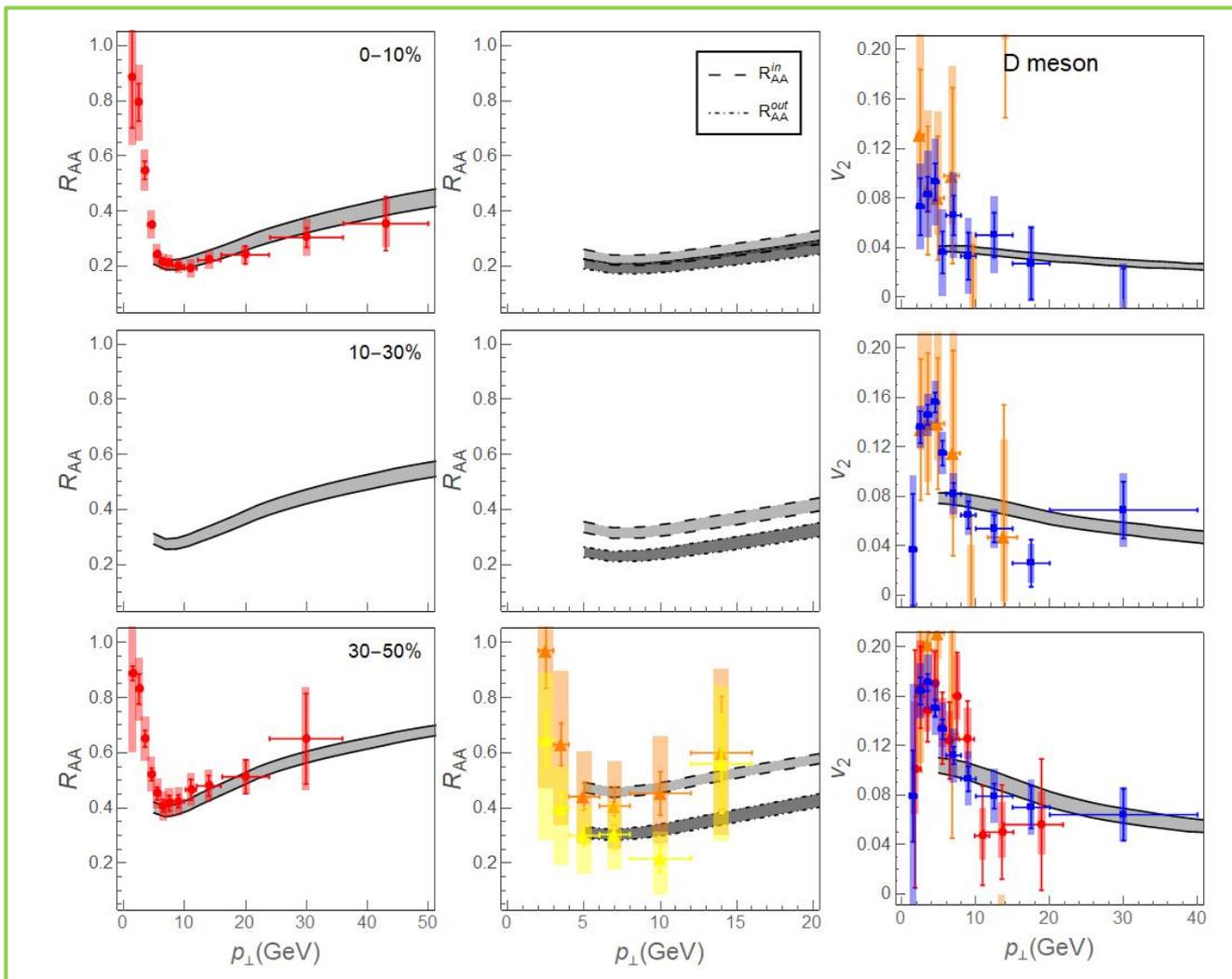
I. Comparison with the data

5.02 TeV Pb+Pb collisions

D. Zigic, I. Salom, J. Auvinen, M. Djordjevic and M.D., PLB 791 (2019) 236



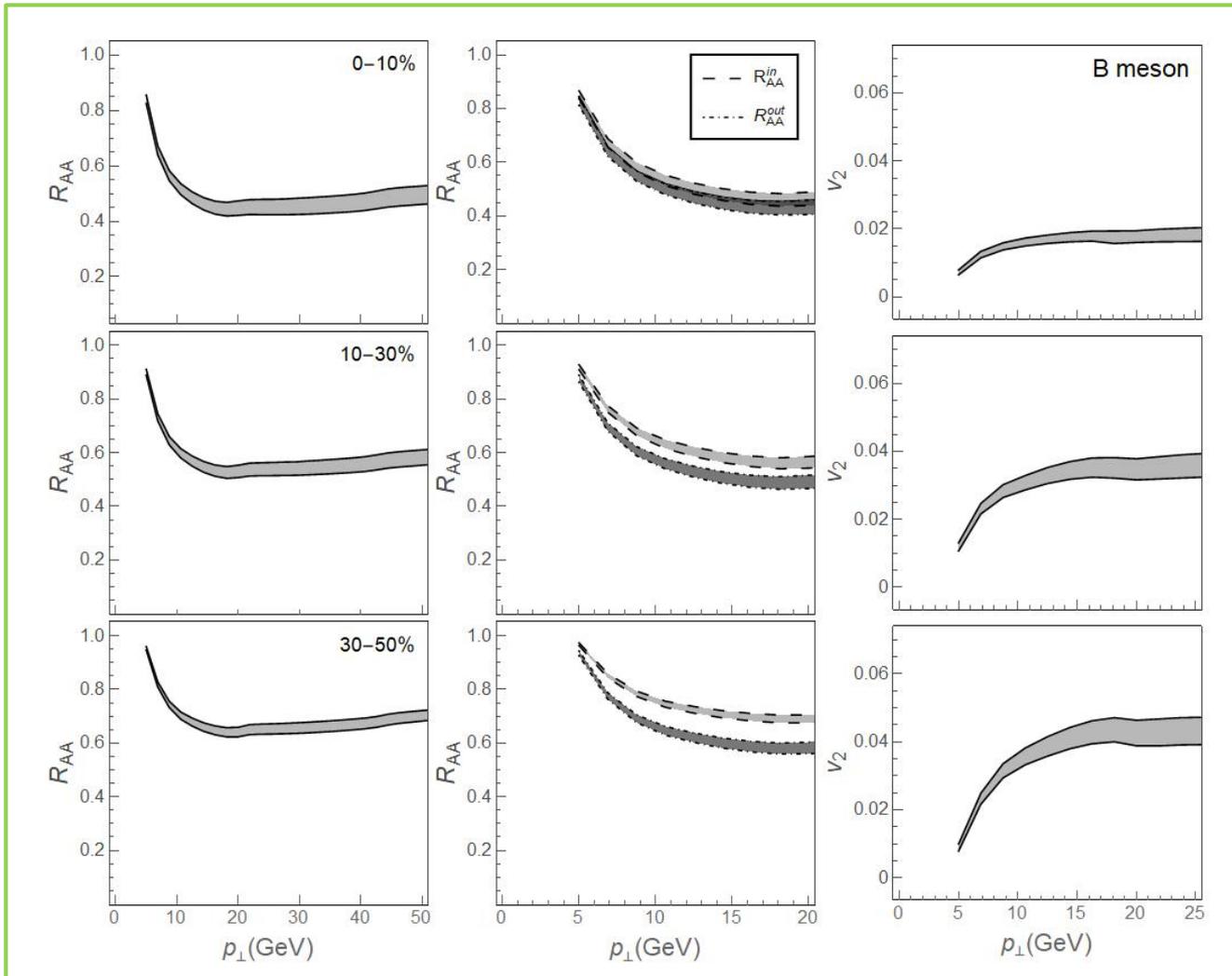
Very good joint agreement with R_{AA} and v_2 data for charged hadrons!



ALICE and CMS data



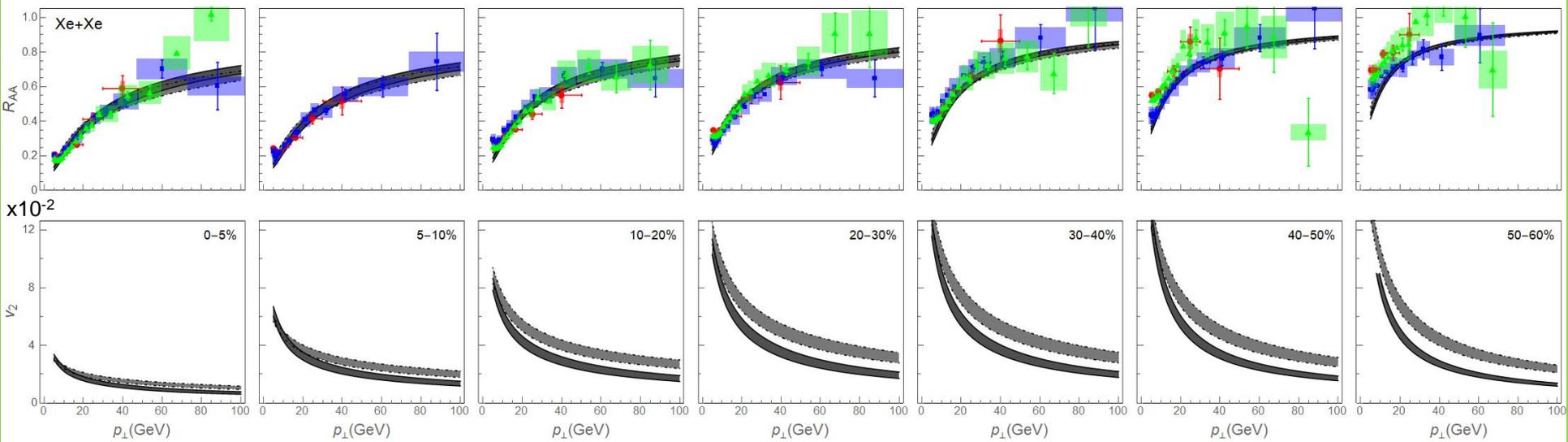
Good joint agreement for D mesons as well!



We predict non-zero v_2 for high pt B mesons.

5.44 TeV Xe+Xe predictions

D. Zigic, I. Salom, J. Auvinen, M. Djordjevic and M.D., PLB 791 (2019) 236



Good agreement with available charged hadrons R_{AA} data!

II. How to differentiate between different energy loss models?

Temperature and path-length dependences (energy loss $\sim T^a L^b$) straightforwardly differentiate different energy loss models.

- Many energy loss models have linear ($b=1$), quadratic ($b=2$) or cubic ($b=3$) path-length dependence.
- Temperature dependence is commonly assumed to be cubic ($a=3$).
- For the dynamical energy loss, both path-length and temperature dependences are between linear and quadratic.
- Therefore, these dependences provide excellent signatures differentiating between different energy loss models.

How to extract path-length and temperature dependences from high pt data?

A non-trivial problem: Have to propose both appropriate observables and appropriate systems to allow the extractions (e.g. standard R_{AA} and v_2 observables are not suitable for this purpose).

“How to test path-length dependence in energy loss mechanisms: analysis leading to a new observable” – poster 189 by Dusan Zigic.
(M.D., D. Zigic, M. Djordjevic and J. Auvinen, PRC Rapid Communications, in press, 2019)

“From R_{AA} to energy loss temperature proportionality factor” – poster 142 by Stefan Stojku.

III. How to infer the shape of the QGP droplet from the data?

Initial spatial anisotropy is one of the main properties of QGP.

A major limiting factor for precision QGP tomography.

Still not possible to directly infer the initial anisotropy from experimental measurements.

Several theoretical studies (MC-Glauber, EKRT, IP-Glasma, MC-KLN) infer the initial anisotropy; lead to notably different predictions, affecting predictions of both low and high pt observables.



Alternative approaches for inferring anisotropy are necessary!

Optimally, these approaches should be complementary to existing predictions.

Based on a method that is fundamentally different to models of early stages of QCD matter.

A novel approach to extract the initial state anisotropy

- **Inference from already available high pt R_{AA} and v_2 measurements** (also to be measured with much higher precision in the future).
- **Use experimental data** (rather than on calculations of early stages of QCD matter).
- **Exploit information from interactions of rare high-pt partons with QCD medium.**
- **Advances the applicability of high pt data.**
- **Up to now, these data mainly used to study the jet-medium interactions, rather than inferring bulk QGP parameters, such as spatial asymmetry.**

What is appropriate observable?

M.D., S. Stojku, M. Djordjevic and P. Huovinen, arXiv:1903.06829

The initial state anisotropy is quantified in terms of eccentricity parameter ϵ_2 :

$$\epsilon_2 = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle} = \frac{\int dx dy (y^2 - x^2) \rho(x, y)}{\int dx dy (y^2 + x^2) \rho(x, y)}$$

where $\rho(x,y)$ is the initial density distribution of the QGP droplet.

High pt v_2 is sensitive to both the anisotropy of the system and its size.

R_{AA} is sensitive only to the size of the system.



Can we extract eccentricity from high pt v_2 and R_{AA} data?

Anisotropy observable

Use a scaling arguments for high pt (M.D., *et al.*, arXiv:1805.04030; M. D. and M. Djordjevic, PRC 92, 024918 (2015))

$$\Delta E/E \sim \langle T \rangle^a \langle L \rangle^b$$

where within our model $a \approx 1.2$, $b \approx 1.4$, consistent with the data.

$$\begin{aligned} R_{AA} &\approx 1 - \xi \langle T \rangle^a \langle L \rangle^b \\ 1 - R_{AA} &\approx \xi \langle T \rangle^a \langle L \rangle^b \end{aligned}$$

$$\begin{aligned} v_2 &\approx \frac{1}{2} \frac{R_{AA}^{in} - R_{AA}^{out}}{R_{AA}^{in} + R_{AA}^{out}} \\ &\approx \xi \langle T \rangle^a \langle L \rangle^b \left(\frac{b}{2} \frac{\Delta L}{\langle L \rangle} - \frac{a}{2} \frac{\Delta T}{\langle T \rangle} \right) \end{aligned}$$



$$\frac{v_2}{1 - R_{AA}} \approx \left(\frac{b}{2} \frac{\Delta L}{\langle L \rangle} - \frac{a}{2} \frac{\Delta T}{\langle T \rangle} \right)$$

This ratio carries information on the asymmetry of the system, but through both spatial and temperature variables.

Anisotropy parameter ζ

$$\frac{v_2}{1 - R_{AA}} \approx \left(\frac{b \Delta L}{2 \langle L \rangle} - \frac{a \Delta T}{2 \langle T \rangle} \right)$$



$$\frac{v_2}{1 - R_{AA}} \approx \frac{1}{2} \left(b - \frac{a}{c} \right) \frac{\Delta L}{\langle L \rangle} \approx 0.57 \zeta$$

$$\zeta = \frac{\Delta L}{\langle L \rangle} = \frac{\langle L_{out} - L_{in} \rangle}{\langle L_{out} + L_{in} \rangle}$$

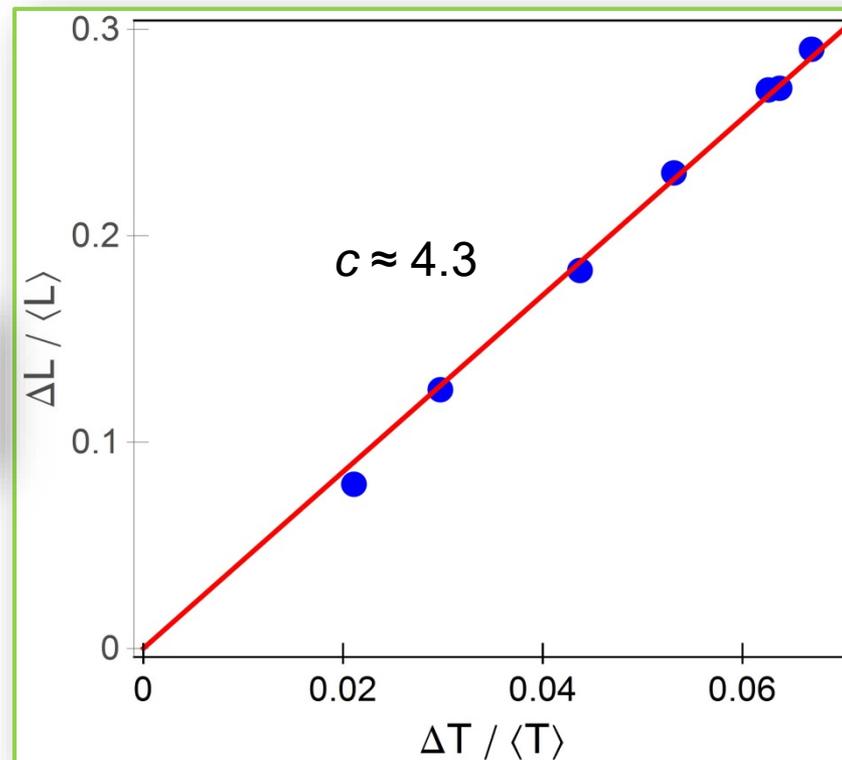


At high pt v_2 over $1 - R_{AA}$ ratio is dictated *solely* by the geometry of the initial fireball.

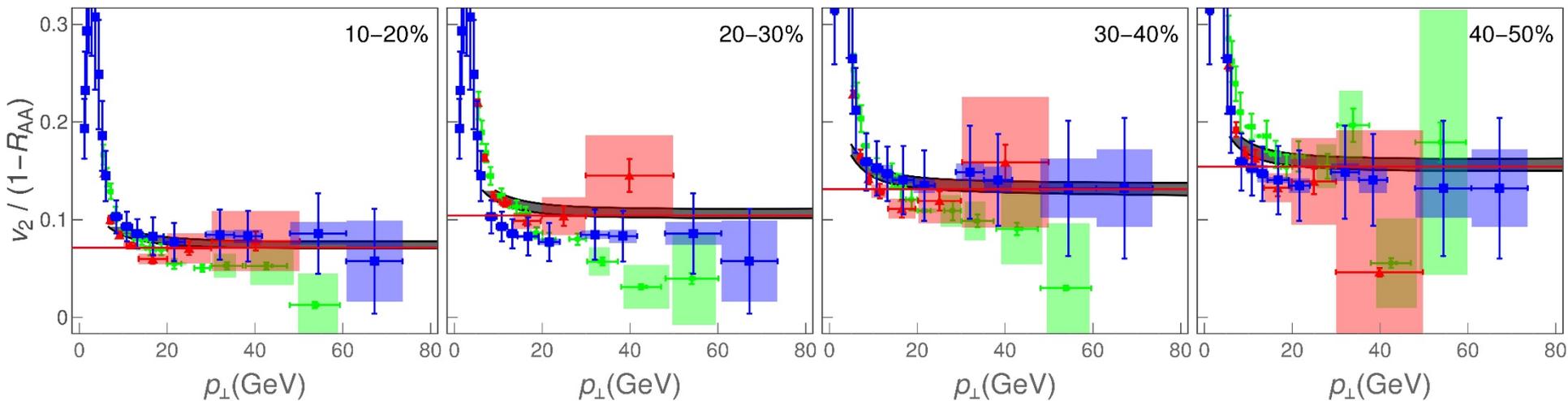


Anisotropy parameter ζ can be *directly* extracted from the high-pt experimental data.

Temperature and spatial assymetry:



Predictions vs. data



- **Solid red line – analytically derived asymptote.**
- **For each centrality and from $p_{\perp} \sim 20$ GeV, $v_2/(1-R_{AA})$ does not depend on p_{\perp} , but is determined by the geometry of the system.**
- **The experimental data for **ALICE**, **CMS** and **ATLAS**, show the same tendency, though the error bars for the data are still large.**
- **In the LHC Run 3, the error bars should reduce by two orders of magnitude.**



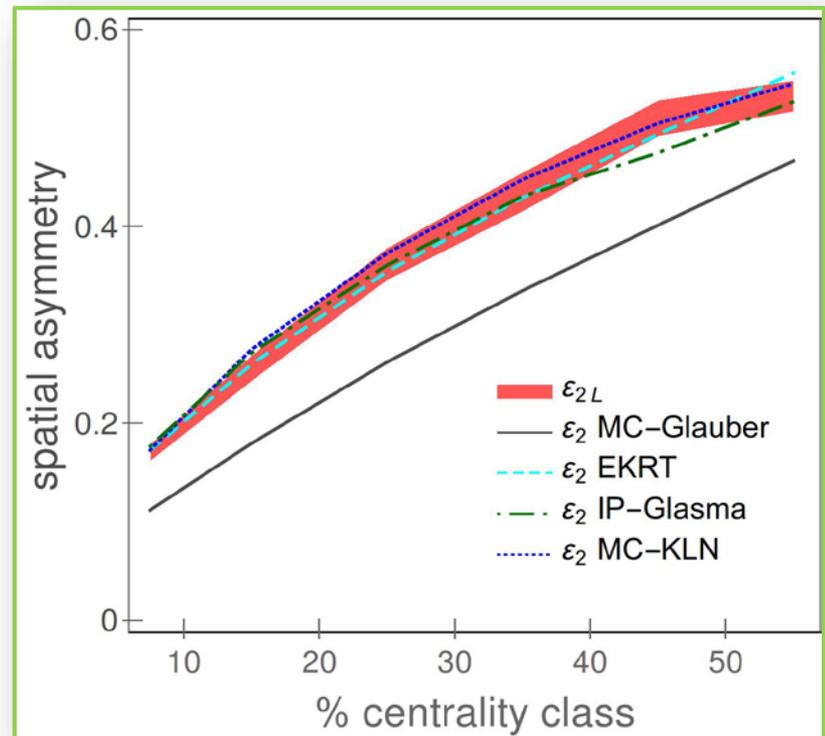
$v_2/(1-R_{AA})$ indeed carries the information about the system's anisotropy, which can be simply (from the straight line high-pt limit) and robustly (in the same way for each centrality) inferred from experimental data.

Eccentricity

Note that the anisotropy parameter ζ is not the commonly used anisotropy parameter ϵ_2 . To facilitate comparison with ϵ_2 values in the literature, we define:

$$\epsilon_{2L} = \frac{\langle L_{out} \rangle^2 - \langle L_{in} \rangle^2}{\langle L_{out} \rangle^2 + \langle L_{in} \rangle^2} = \frac{2\zeta}{1 + \zeta^2}$$

and compare with results in the literature.



ϵ_{2L} is in an excellent agreement with ϵ_2 from which we started from.



$v_2/(1-R_{AA})$ – reliable/robust procedure to recover initial state anisotropy.

The width of our ϵ_{2L} band is smaller than the difference in the ϵ_2 values obtained by using different models (e.g. MC-Glauber vs. MC-KLN).



Resolving power to distinguish between different initial state models, although it may not be possible to separate the finer details of more sophisticated models.

Summary

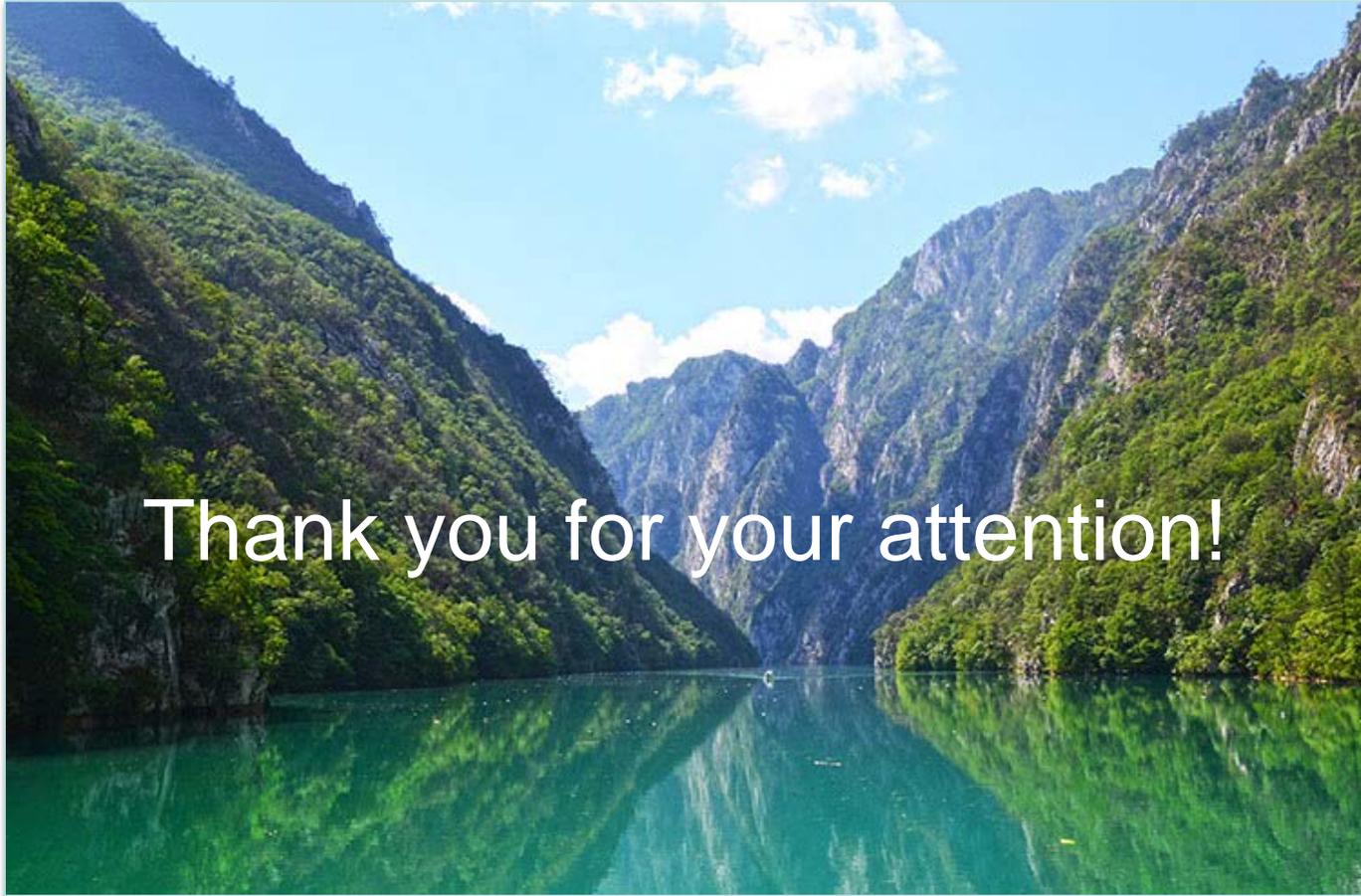
DREENA-C and DREENA-B frameworks are introduced, where DREENA is a computational implementation of the dynamical energy loss formalism.

DREENA provides a very good agreement with experimental data, both in terms of predictions and postdictions.

It leads to a novel insight into energy loss mechanisms, allowing to distinguish between different energy loss models.

We argue that high-pt probes are also powerful tomography tools, as they are sensitive to global QGP properties, which we here show in the case of spatial anisotropy of QCD matter.

Important as high-pt theory and data are commonly used to explore high-pt parton interactions with QGP, while QGP bulk properties are traditionally explored through low-pt data and corresponding models.



Thank you for your attention!

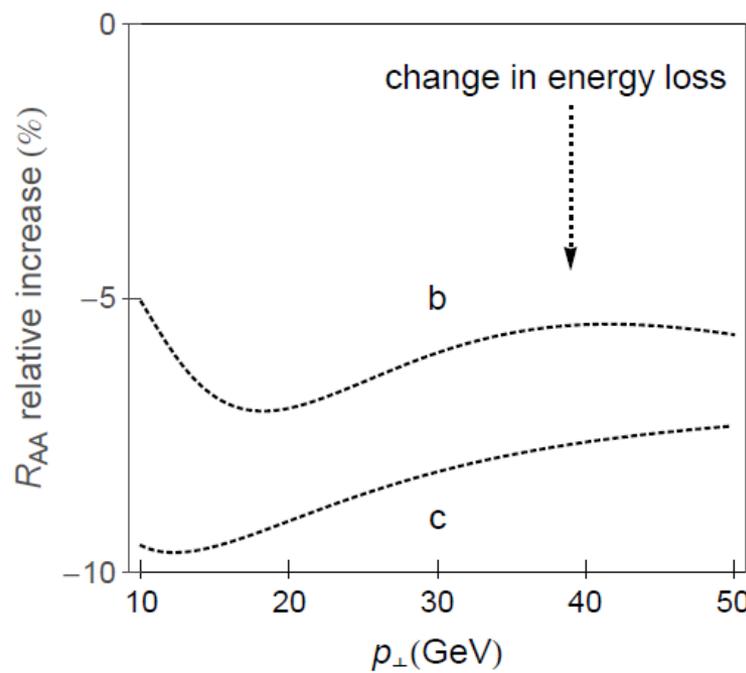
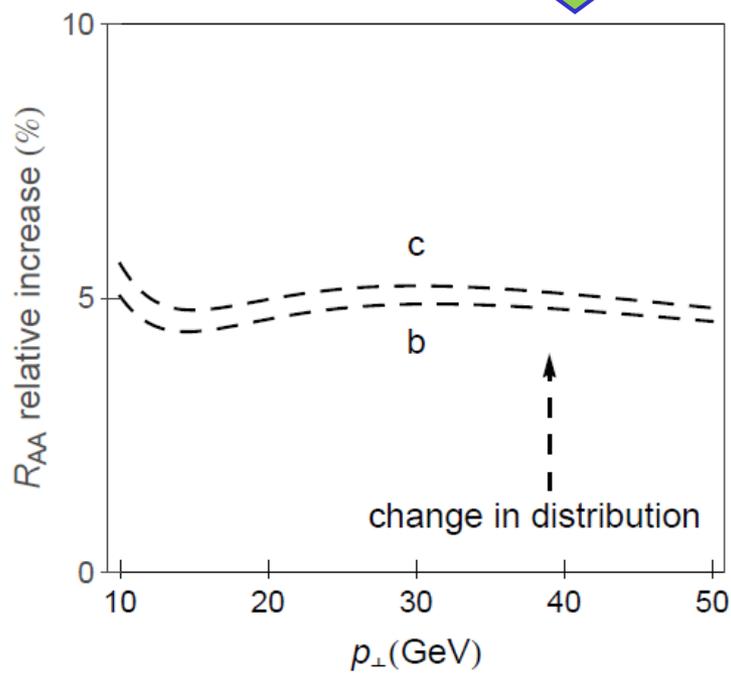
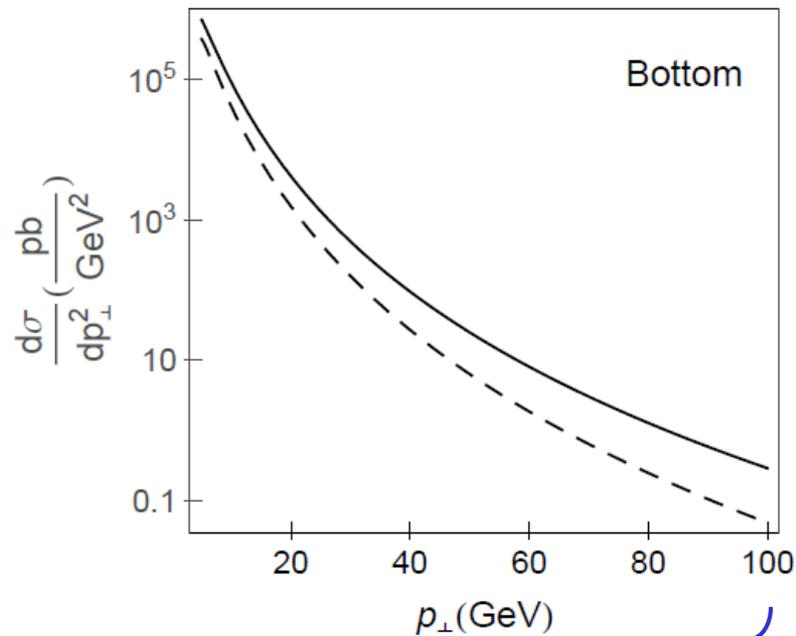
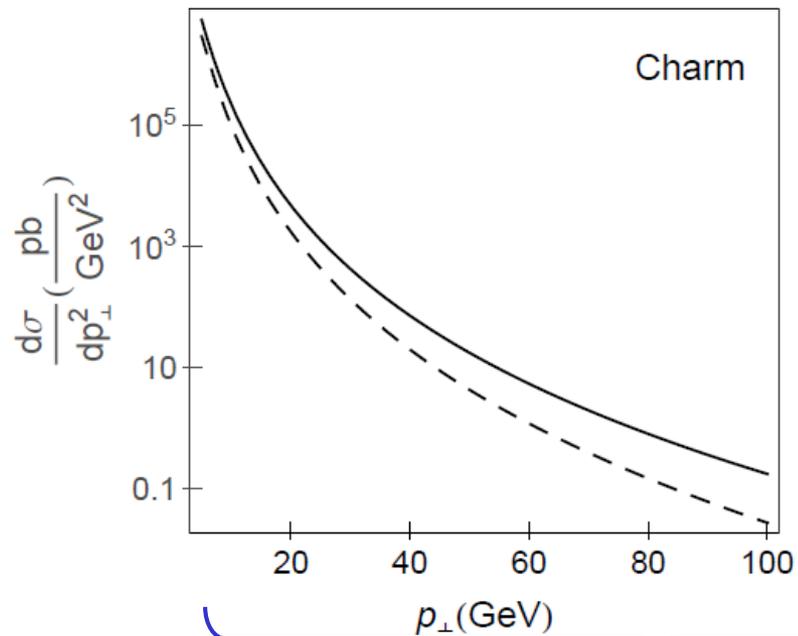
Canyon of river DREENA in Serbia



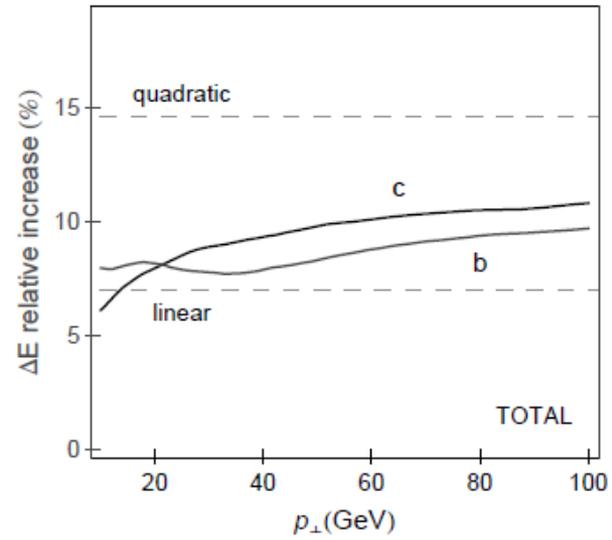
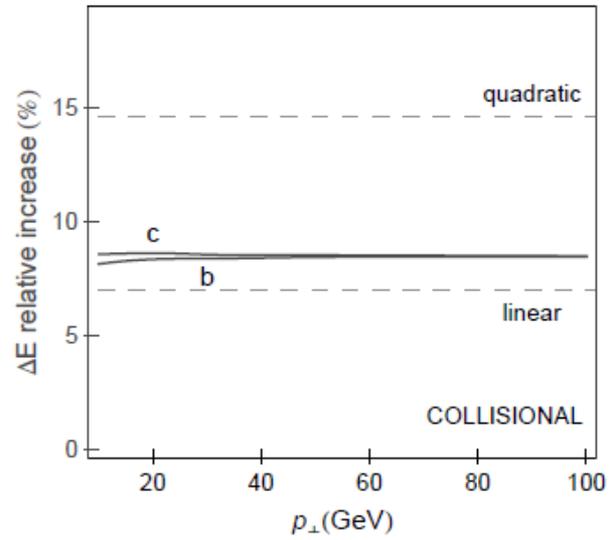
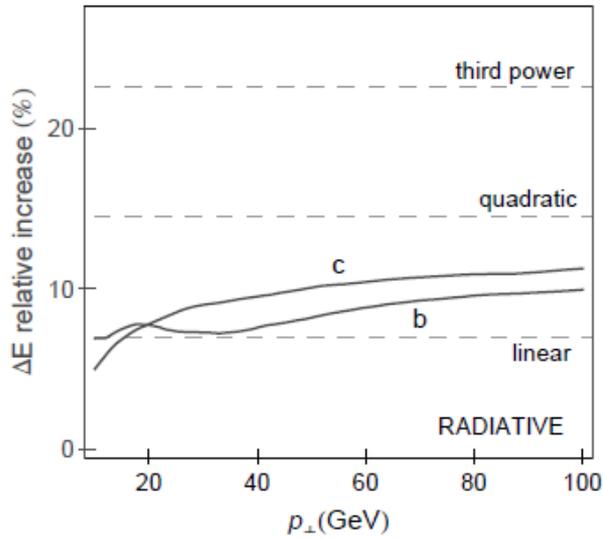
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МИНИСТАРСТВО ПРОСВЕТЕ,
НАУКЕ И ТЕХНОЛОШКОГ РАЗВОЈА



Temperature dependence of the energy loss



What is an appropriate system?

Measurements on 5.02 TeV Pb+Pb already available, 5.44 TeV smaller systems (e.g. Xe+Xe) are also becoming available.

The main property differentiating the two systems is its size ($A_{\text{PbPb}}=208$, $A_{\text{XeXe}}=129$).

All other properties basically remain the same:

- i. Initial momentum distribution
- ii. Average temperature for each centrality region
- iii. Path length distributions (up to rescaling factor $A^{1/3}$)

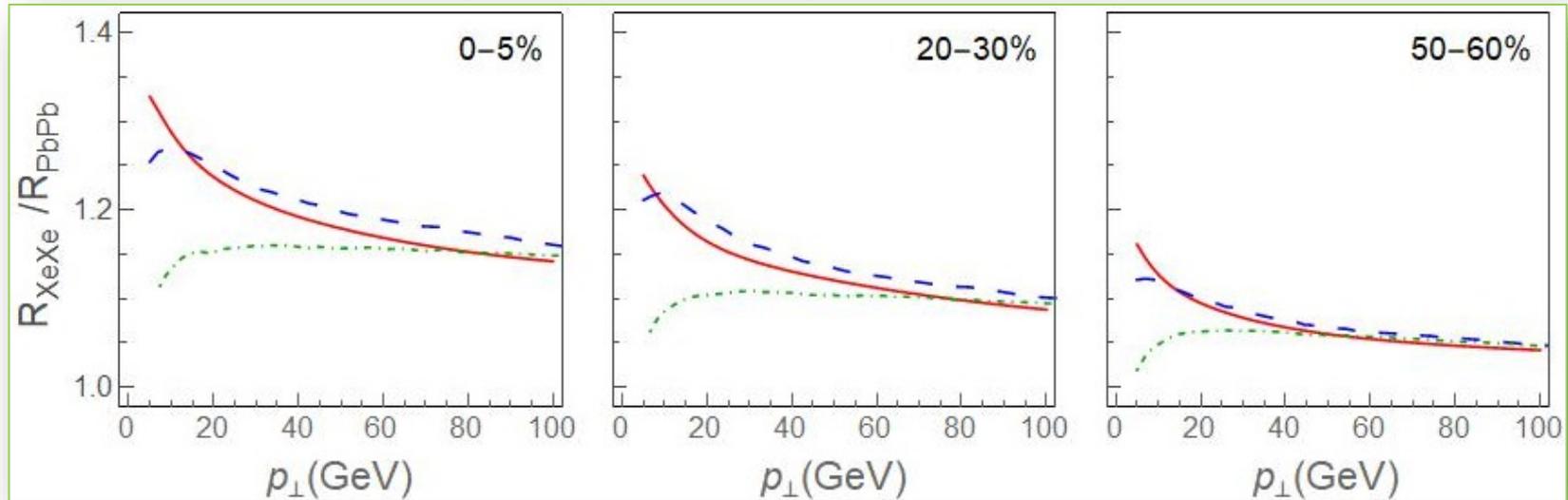
Consequently, comparison of suppressions in Pb+Pb and Xe+Xe is an excellent way to study the path length dependence.

What is appropriate observable?

M.D., D. Zigic, M. Djordjevic and J. Auvinen, PRC Rapid Communications, in press, 2019.

The ratio of the two R_{AA} s seems a natural choice, and has been proposed before.

However, in this way the path length dependence cannot be naturally extracted (also a strong centrality dependence):



What is the reason for this? – use scaling arguments:

$$\Delta E/E \sim T^a L^b \longrightarrow \frac{R_{XeXe}}{R_{PbPb}} \approx 1 - \xi T^a L_{Pb}^b \left(1 - \left(\frac{A_{Xe}}{A_{Pb}} \right)^{b/3} \right)$$

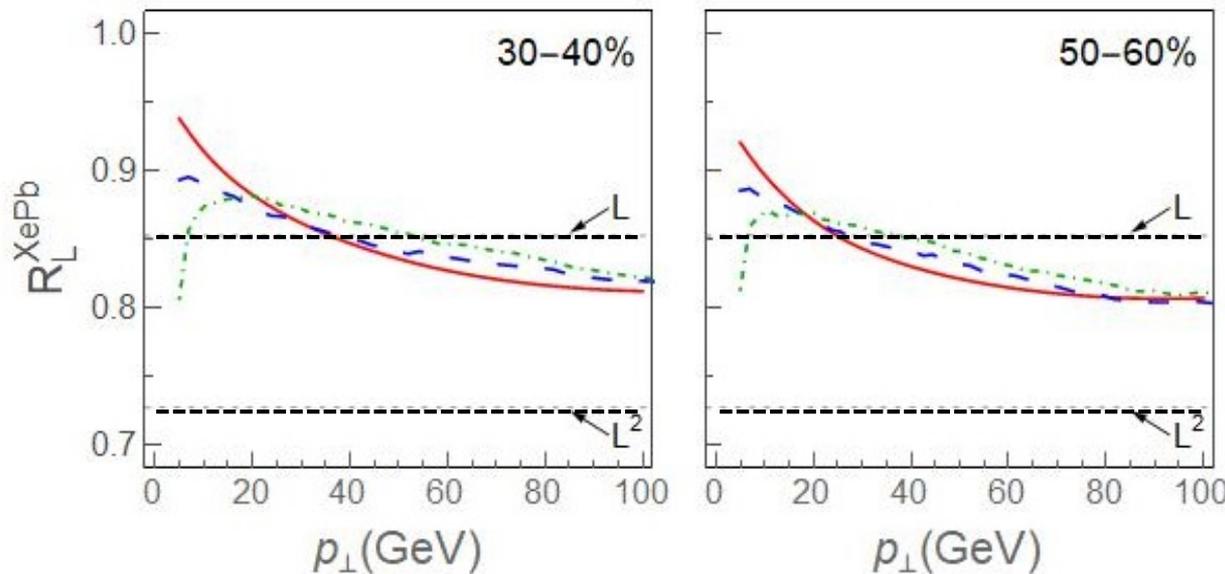
We see that the ratio includes a complicated relationship.

What we propose?

Use $1-R_{AA}$ ratio instead:

$$R_L^{XePb} \equiv \frac{1 - R_{XeXe}}{1 - R_{PbPb}} \approx \frac{\xi T^a L_{Xe}^b}{\xi T^a L_{Pb}^b} \approx \left(\frac{A_{Xe}}{A_{Pb}} \right)^{b/3}$$

We see a simple dependence on only size of the medium ($A^{1/3}$ ratio) and the path length dependence (exponent b).



The path length dependence can be extracted in a simple way, and there is only a weak centrality dependence.

M.D., D. Zigic, M. Djordjevic and J. Auvinen, PRC Rapid Communications, in press, 2019.

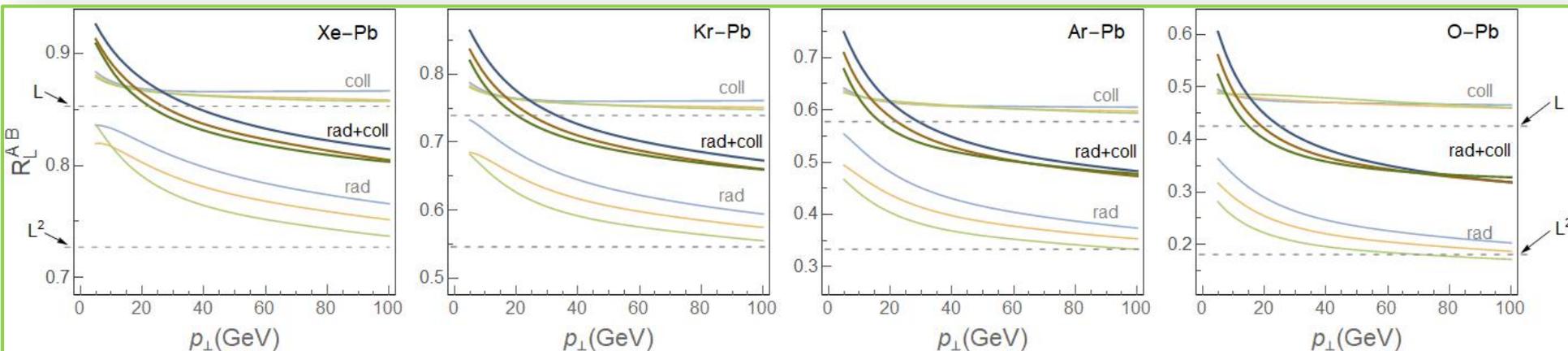
$1-R_{AA}$ ratio therefore seems as a **natural observable**, which we call **path-length sensitive suppression ratio**.

What about other smaller systems ?

- BSS at the LHC

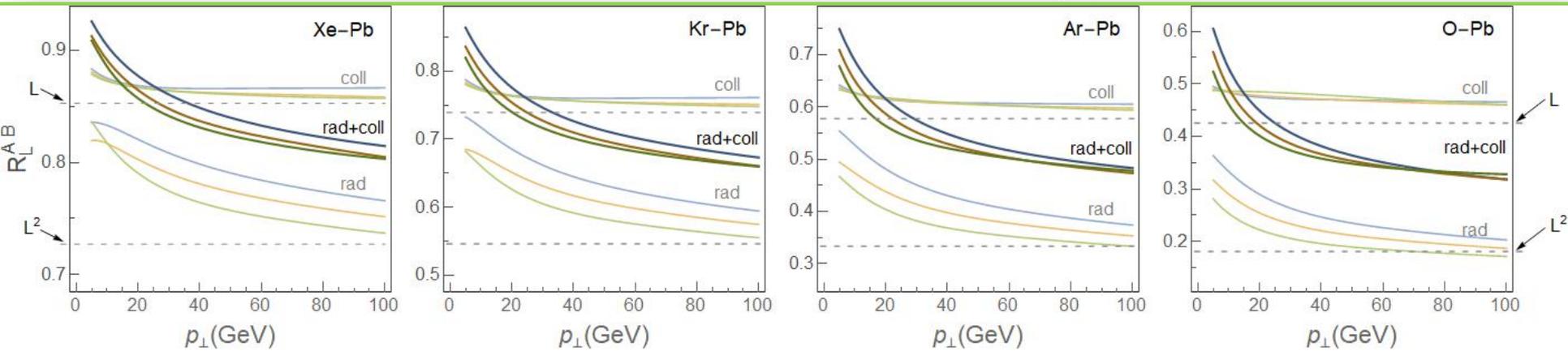
Precision measurements of smaller systems, i.e. Kr+Kr, Ar+Ar and O+O, are expected to become available in the future Beam Size Scan (BSS) at the LHC.

Can these systems be also used to extract path length dependence of the energy loss?



M.D., D. Zigic, M. Djordjevic and J. Auvinen, PRC Rapid Communications, in press, 2019.

What about other smaller systems?



- i.* R_L^{AB} is almost independent on centrality for the 30-60% centrality region.
- ii.* For all four systems, R_L^{AB} shows the same behavior, i.e. it is very robust with respect to extracting path-length dependence.
- iii.* Reliably recovers collisional and radiative energy loss path-length dependences.
- iv.* From experimental perspective, smaller systems might be more convenient for applying this observable, as the region of applicability (between L and L^2) increases with decreasing system size.

Consequently, we propose that R_L^{AB} is simple, robust and reliable observable for extracting the path-length of the energy loss.