

# Signatures of quark-hadron phase transitions in neutron-star mergers

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# The two-body problem in GR

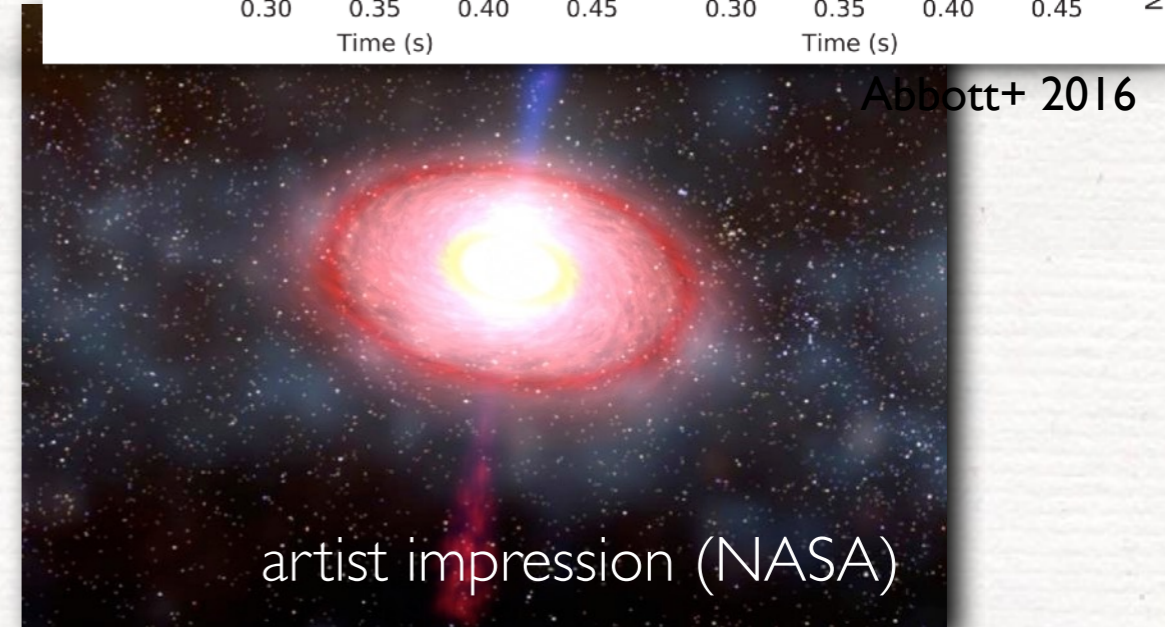
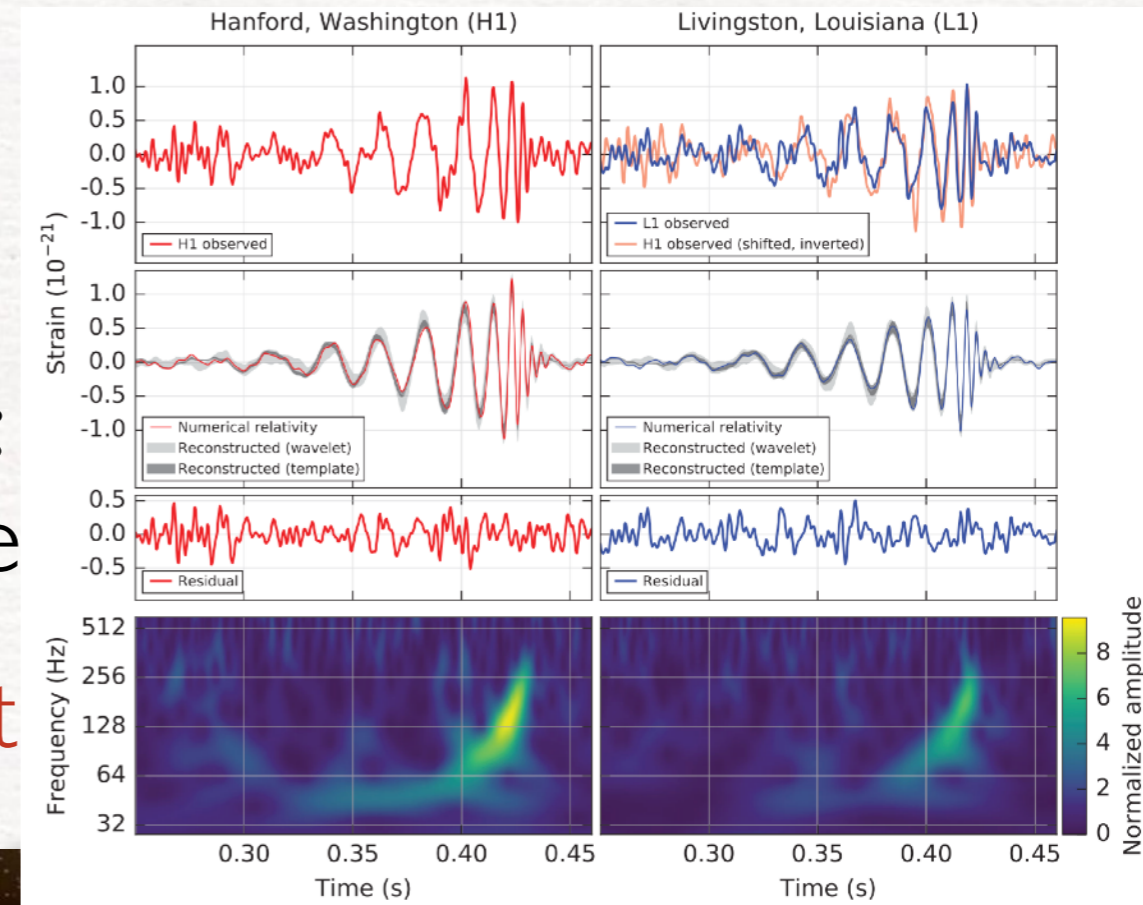
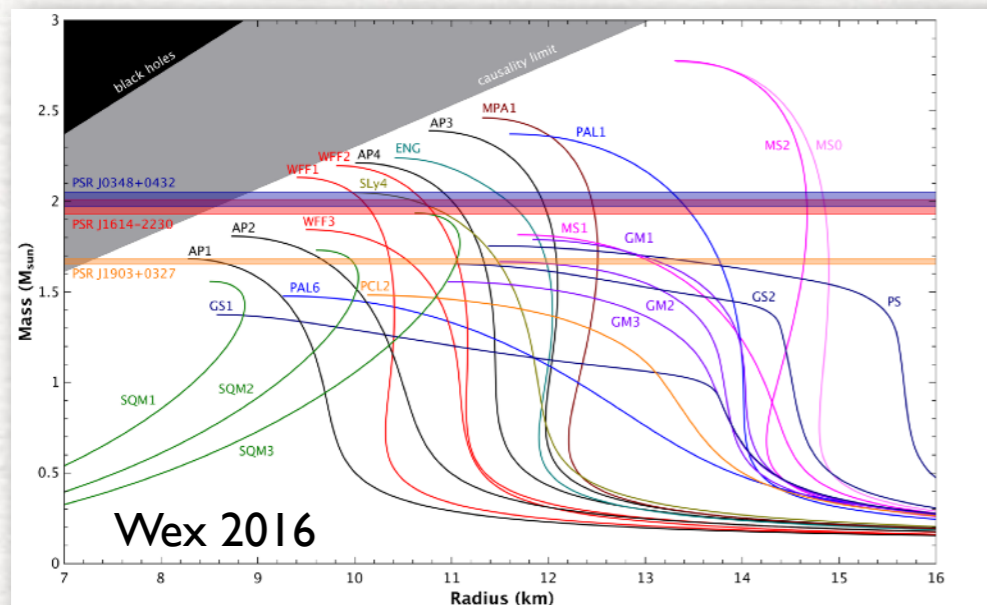
- For BHs we know what to **expect**:

$$\text{BH} + \text{BH} \longrightarrow \text{BH} + \text{GWs}$$

- For NSs the question is more **subtle**: hyper-massive neutron star (HMNS), ie

$$\text{NS} + \text{NS} \longrightarrow \text{HMNS} + \dots ? \longrightarrow \text{BH} + \text{t}$$

- **HMNS** phase can provide clear information on **EOS**



- **BH+torus** system may tell us on the central engine of **GRBs**

# The two-body problem in GR

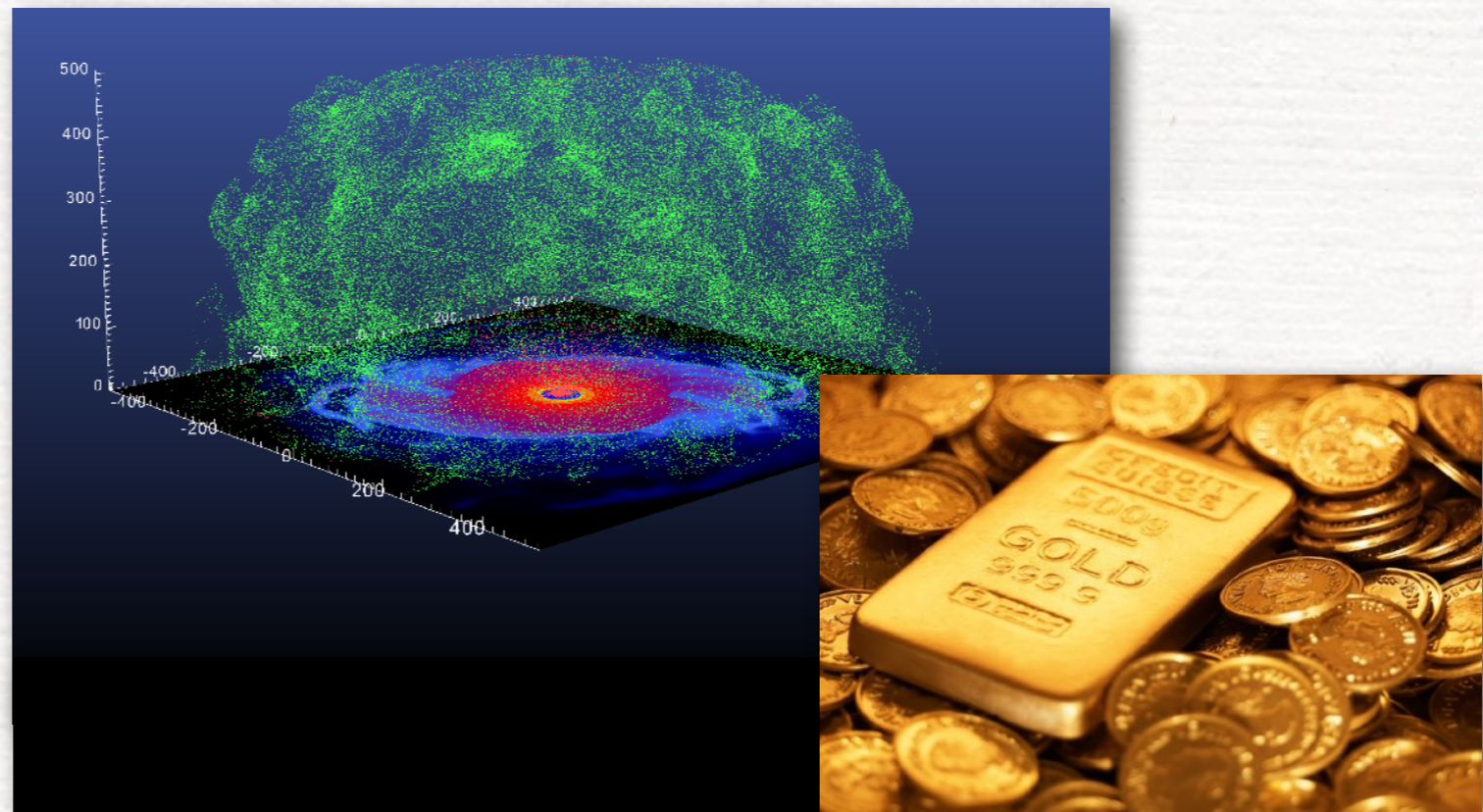
- For BHs we know what to **expect**:



- For NSs the question is more **subtle**: the merger leads to an hyper-massive neutron star (HMNS), ie a metastable equilibrium:



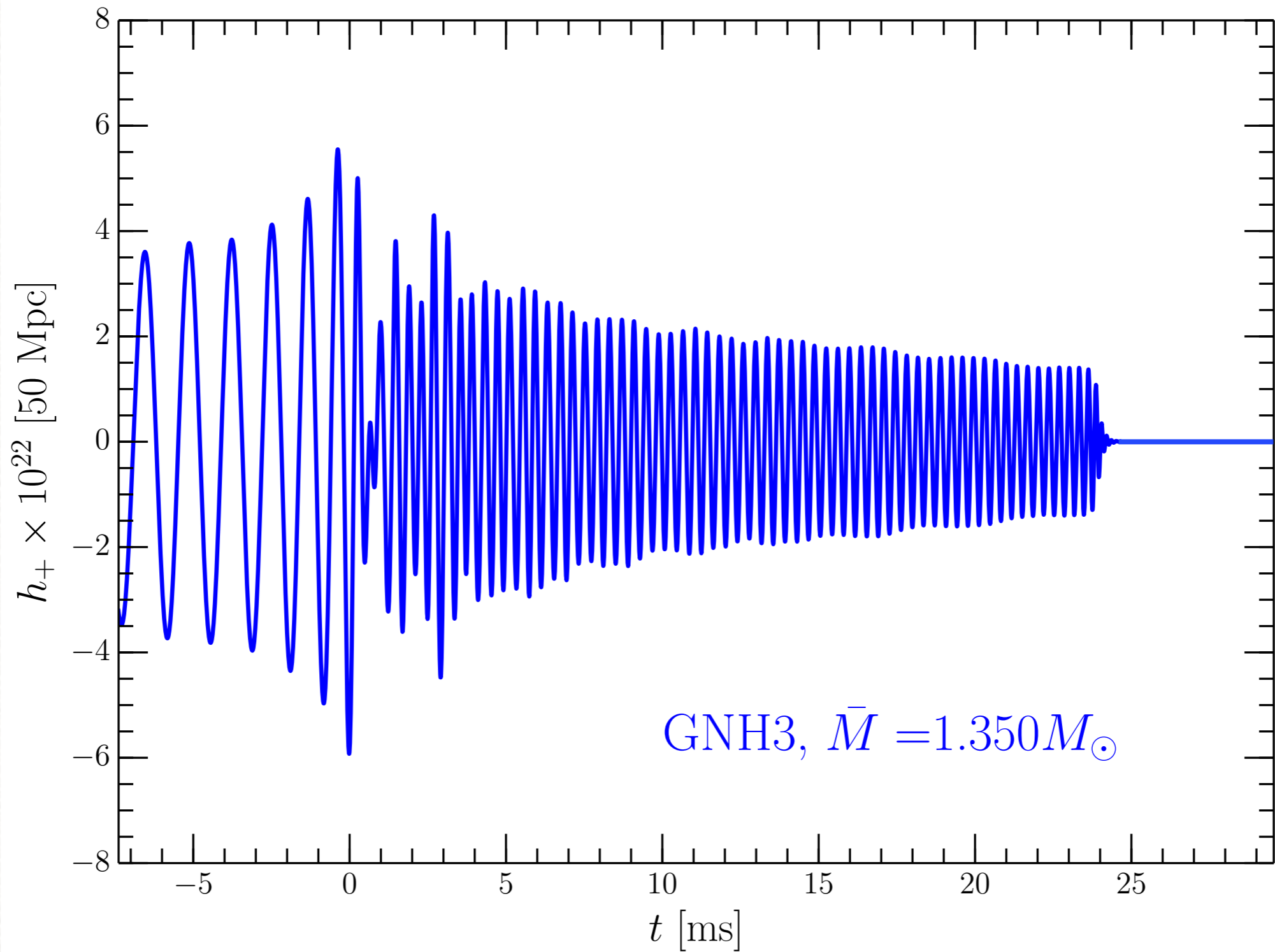
- **ejected matter** undergoes nucleosynthesis of heavy elements



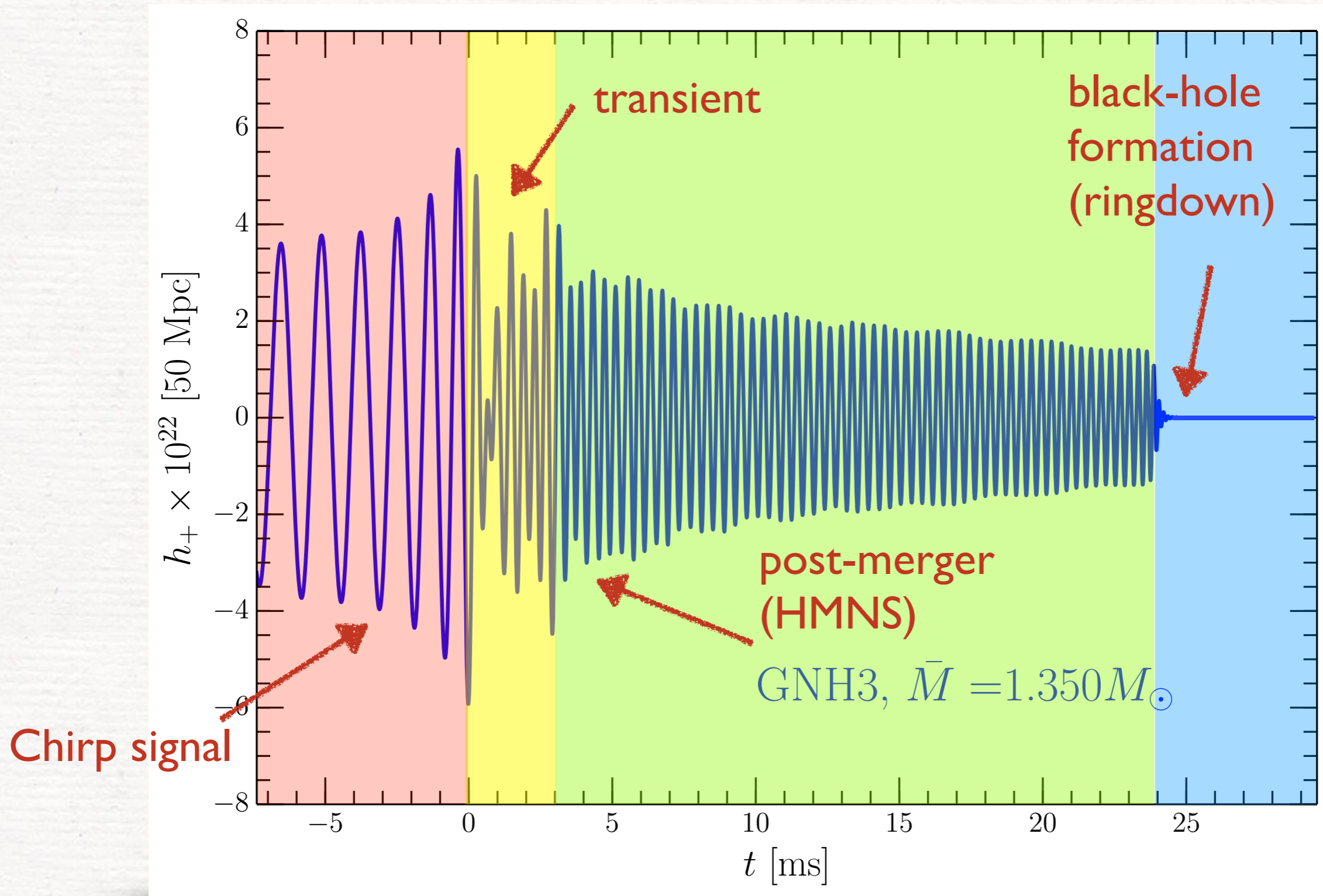
A prototypical simulation with possibly  
the best code looks like this...



# Anatomy of GW signal

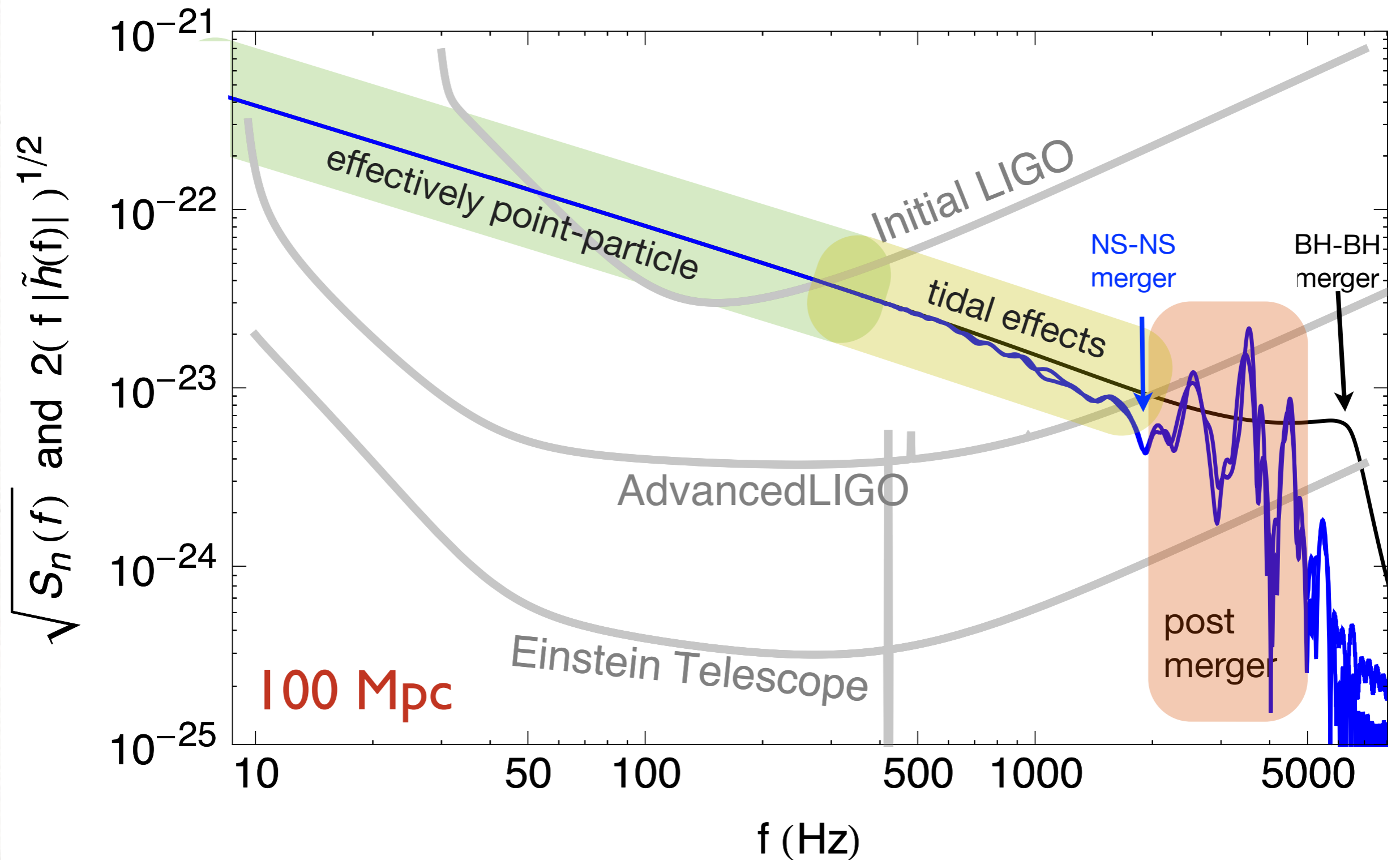


# Anatomy of GW signal



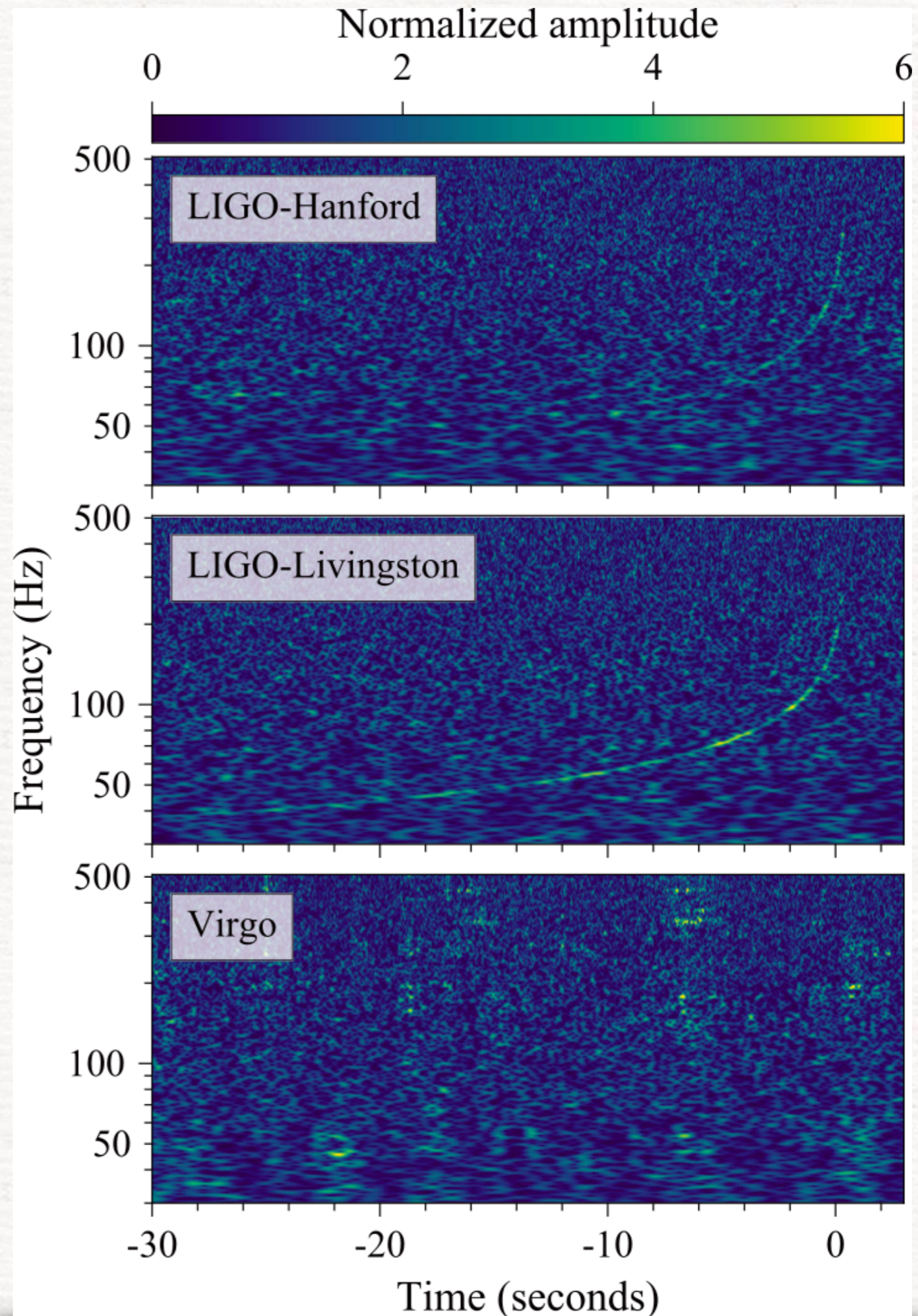
**Postmerger signal:** peculiar of binary NSs

# In frequency space



# GW170817: the first binary neutron-star system

- \* **Unfortunately** only the **inspiral** signal was detected.
- \* **Fortunately** this was **sufficient** to set a number of constraints on max. mass, tidal deformability, radii, etc.
- \* **Gold mine** is in the **post-merger** for there is a lot of information there, also about **phase transitions!**

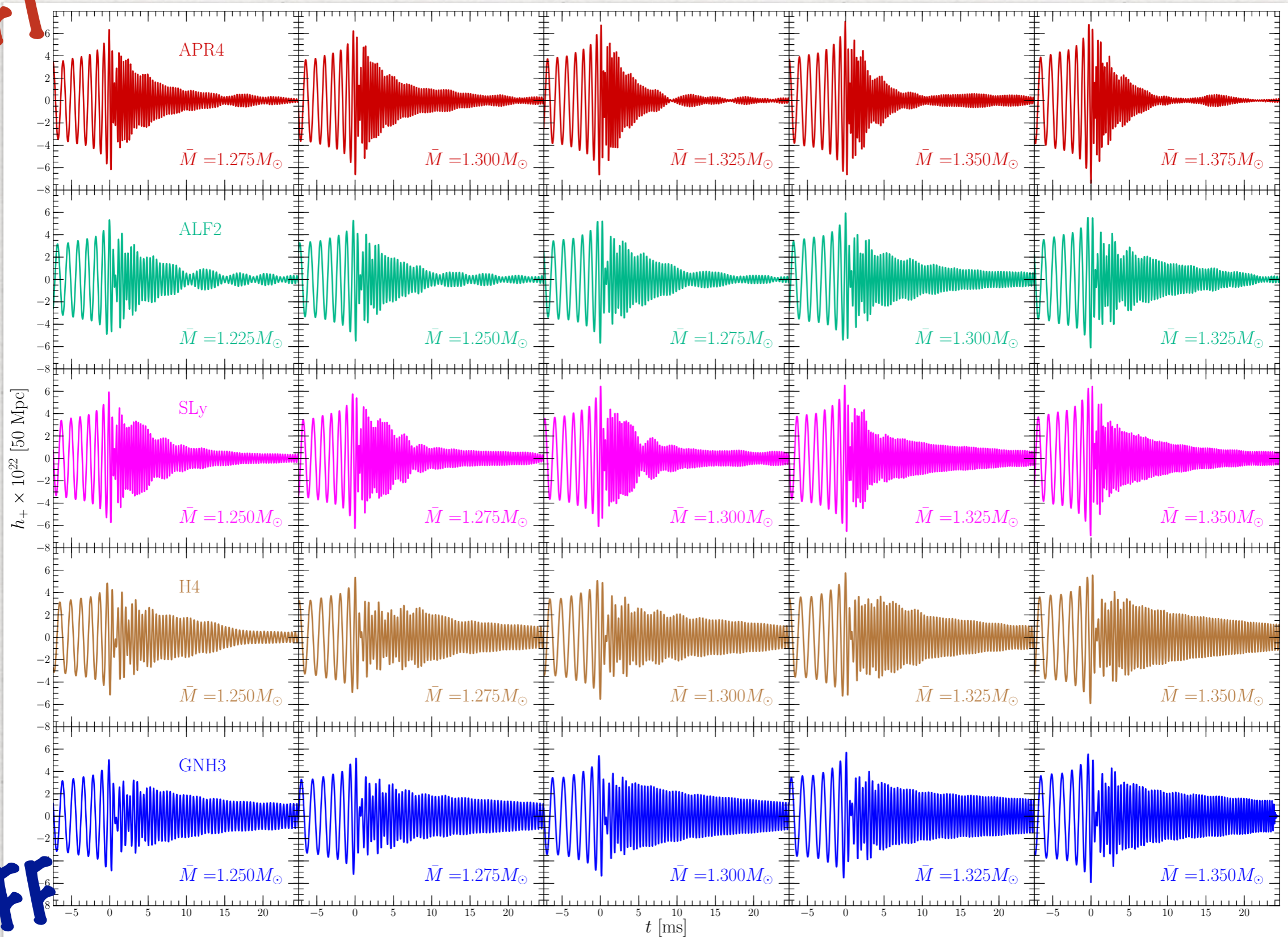




# What we can do nowadays

Takami, LR, Baiotti (2014, 2015), LR+ (2016)

SOFT

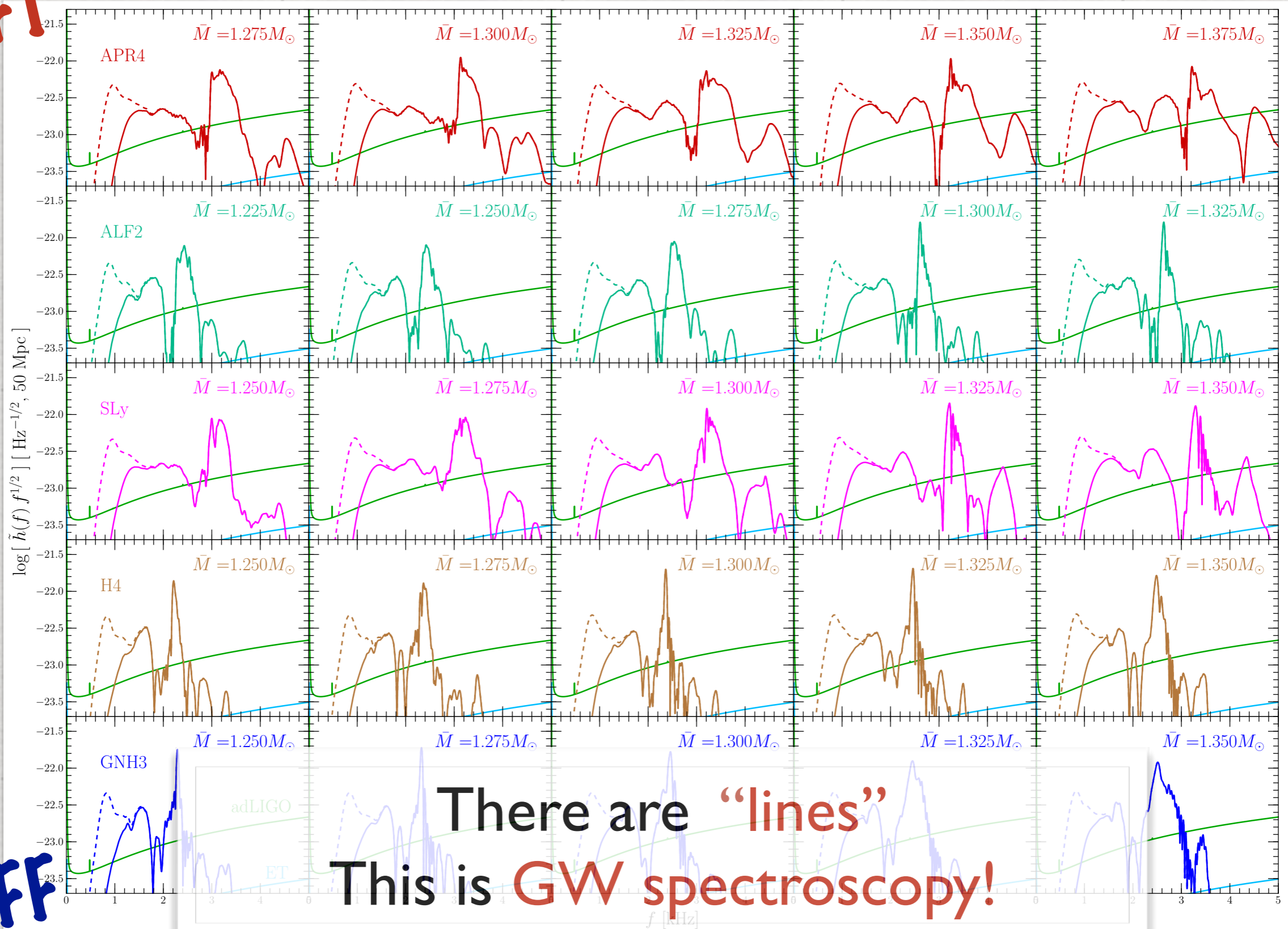


STIFF

# Extracting information from the EOS

Takami, LR, Baiotti (2014, 2015), LR+ (2016)

SOFT

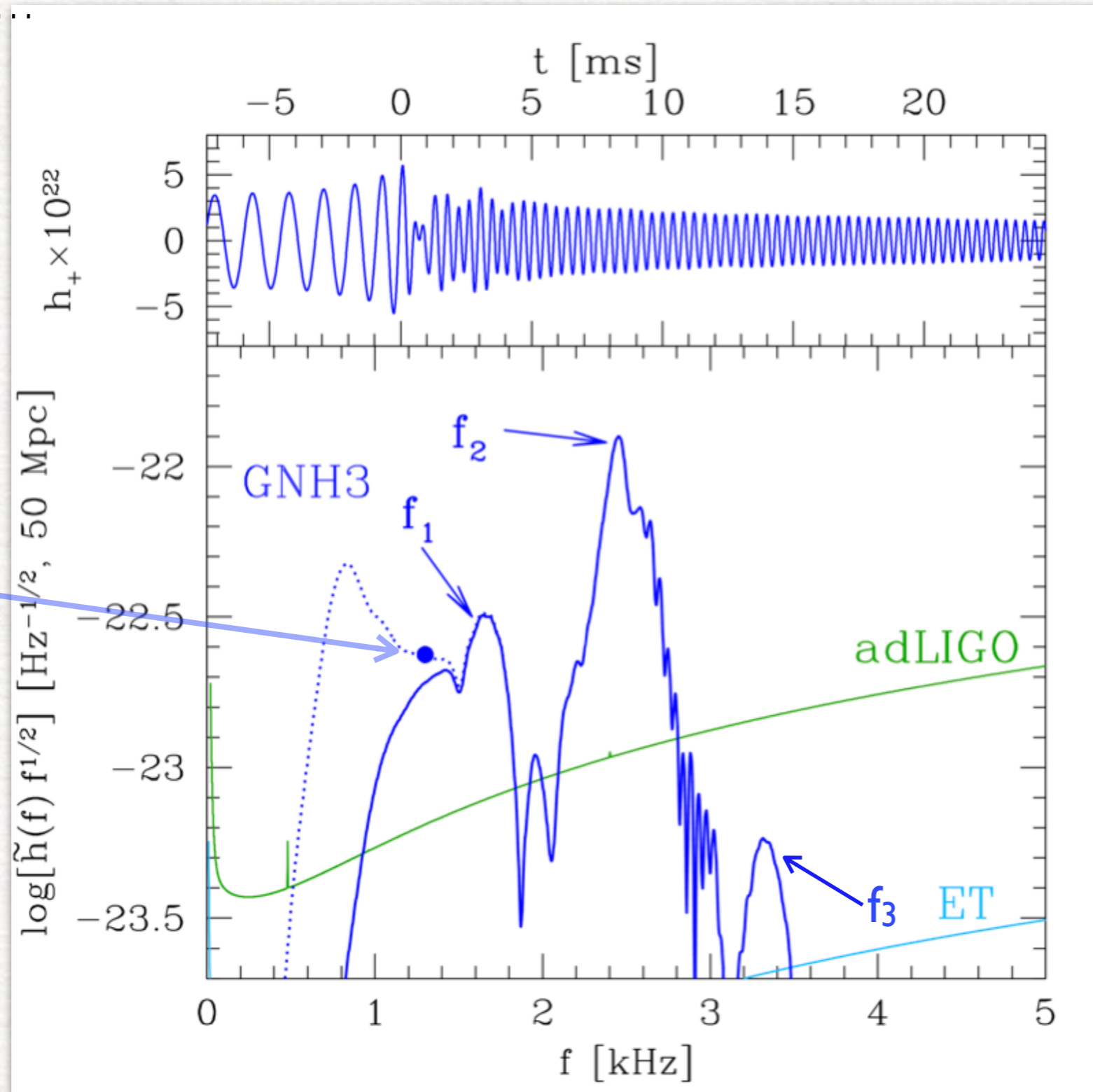


STIFF

# A spectroscopic approach to the EOS

Oechslin+2007, Baiotti+2008, Bauswein+ 2011, 2012, Stergioulas+ 2011, Hotokezaka+ 2013, Takami 2014, 2015, Bernuzzi 2014, 2015, Bauswein+ 2015, Clark+ 2016, LR+2016, de Pietri+ 2016, Feo+ 2017, Bose+ 2017 ...

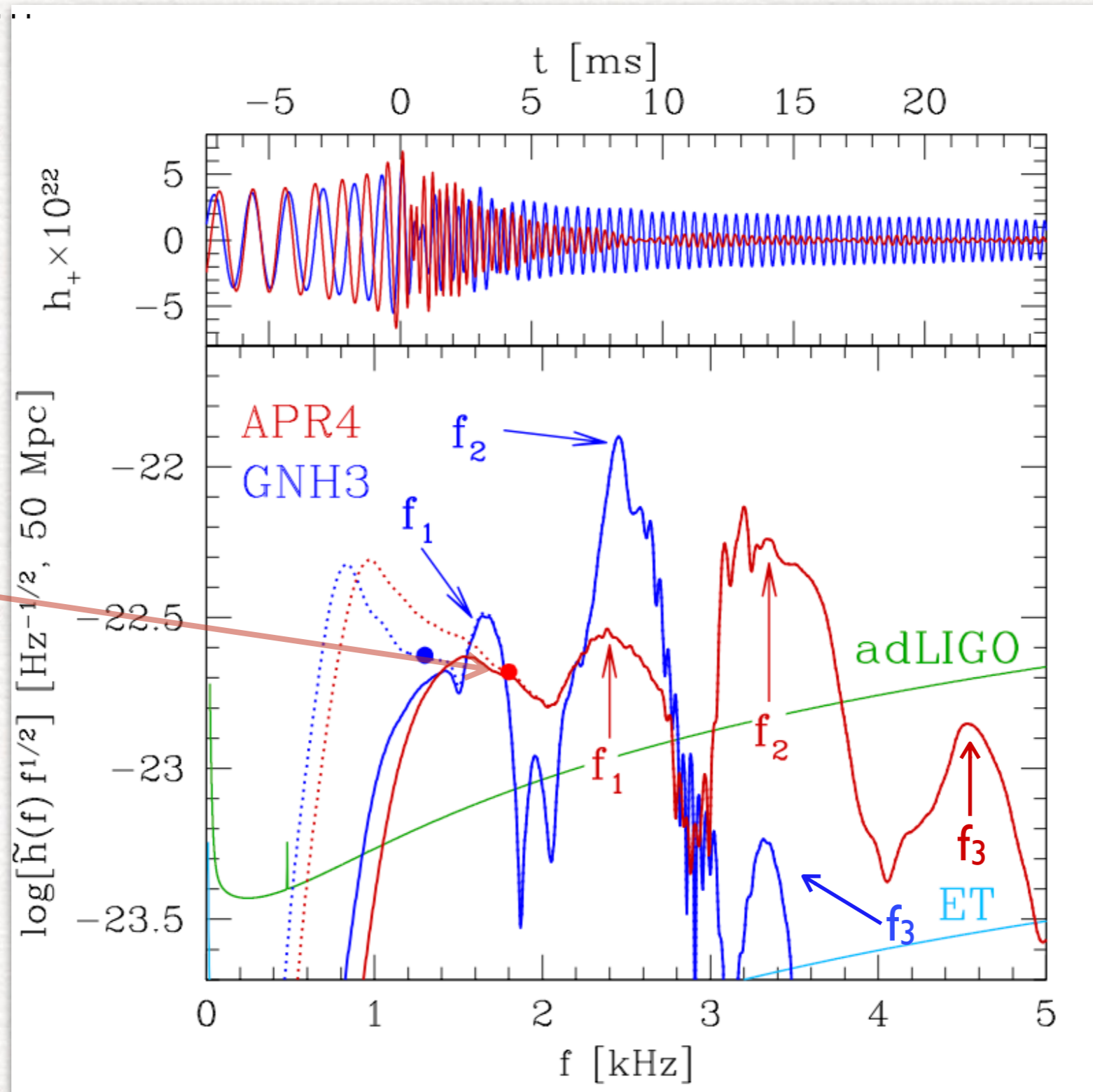
merger  
frequency



# A spectroscopic approach to the EOS

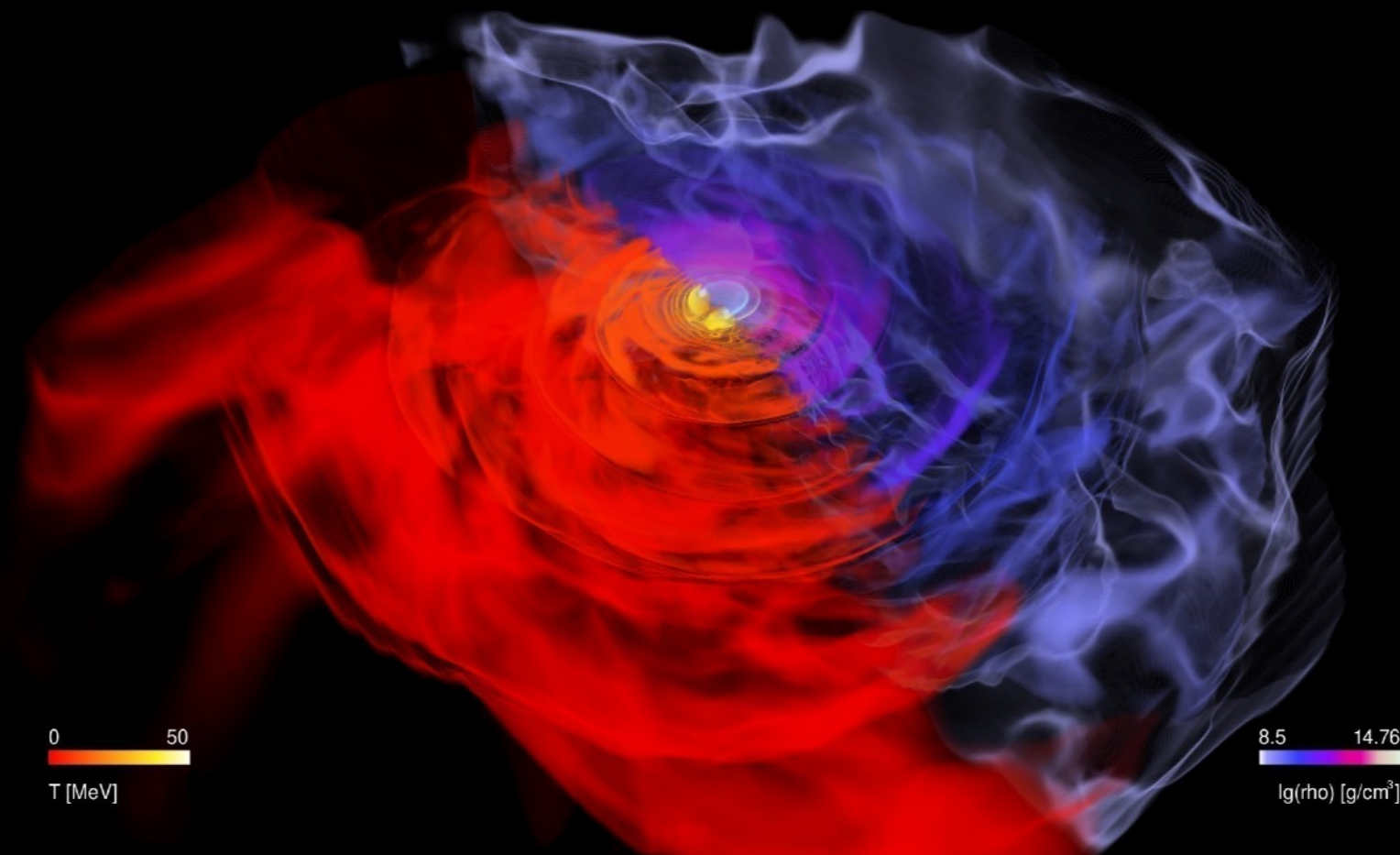
Oechslin+2007, Baiotti+2008, Bauswein+ 2011, 2012, Stergioulas+ 2011, Hotokezaka+ 2013, Takami 2014, 2015, Bernuzzi 2014, 2015, Bauswein+ 2015, Clark+ 2016, LR+2016, de Pietri+ 2016, Feo+ 2017, Bose+ 2017 ...

merger  
frequency

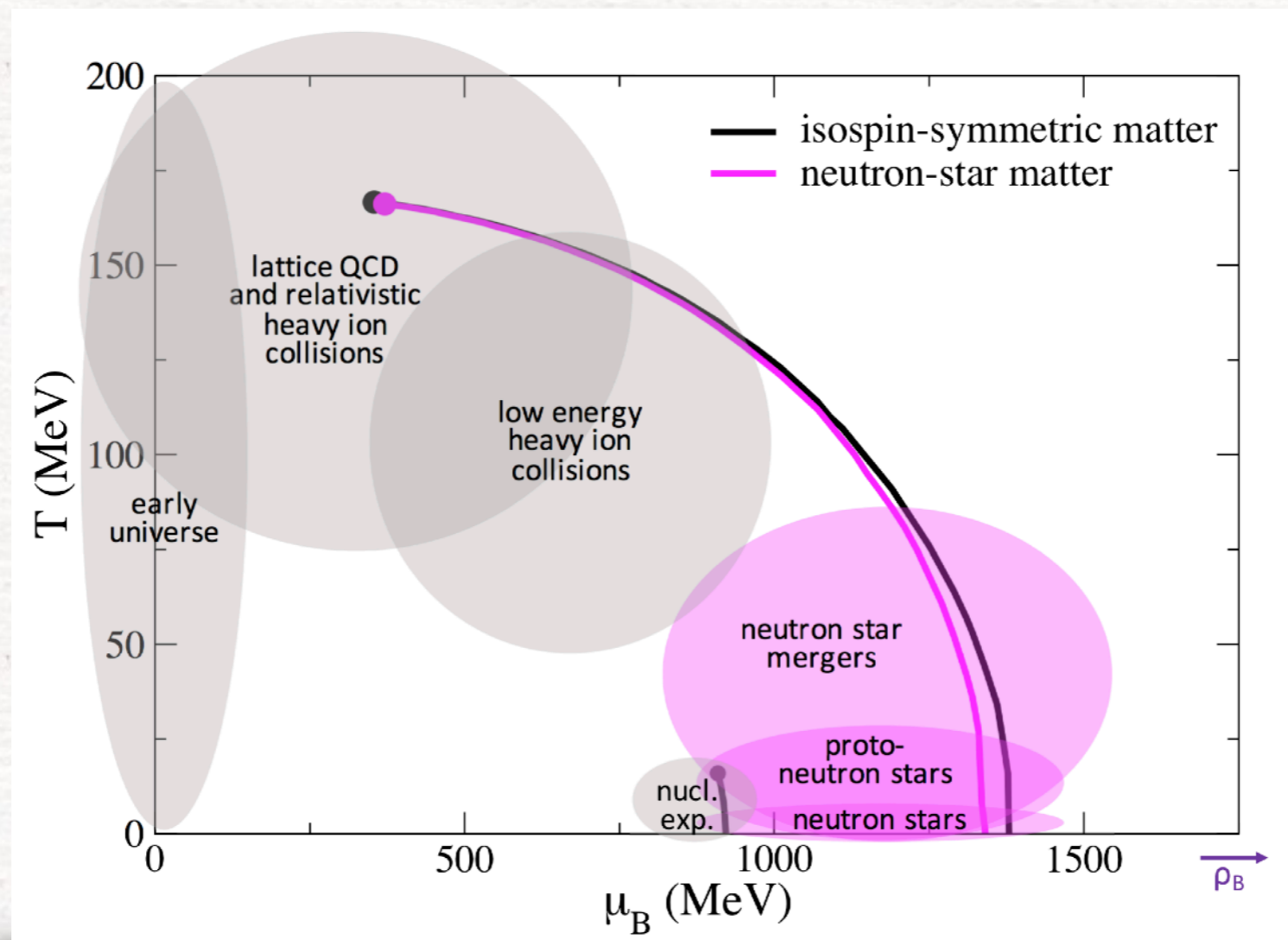


# Phase transitions and their signatures

Most, Papenfort, Dexheimer, Hanauske, Schramm, Stoecker, LR (2019)  
see also Bauswein, Bastian, Blaschke, Chatziioannou, Clark, Fischer, Oertel (2019)



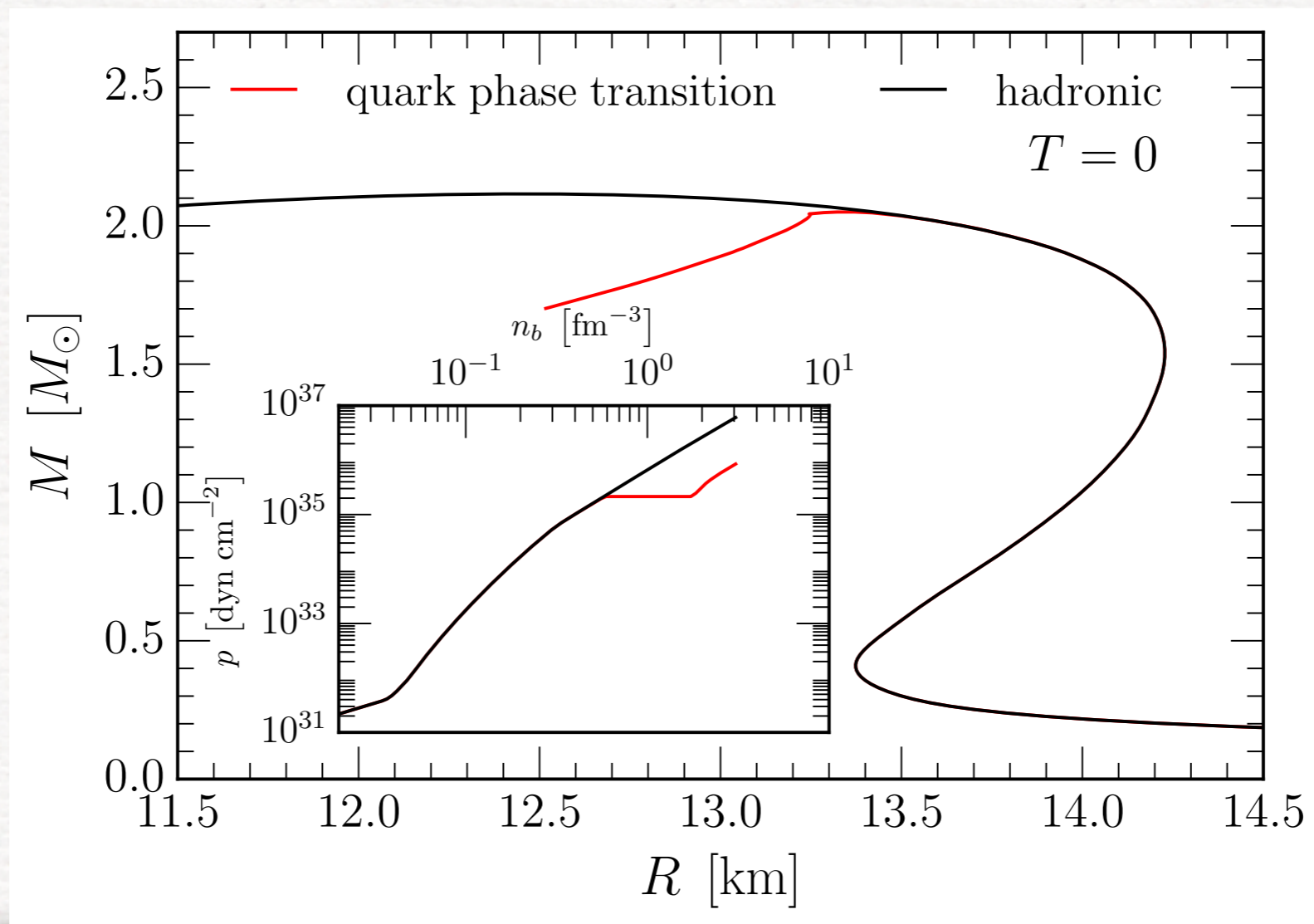
- \* **Strangeness** is expected in neutron stars both in the inspiral (hyperons) and possibly after merger (strange quarks?)
- \* **Isolated** neutron stars probe a small fraction of phase diagram



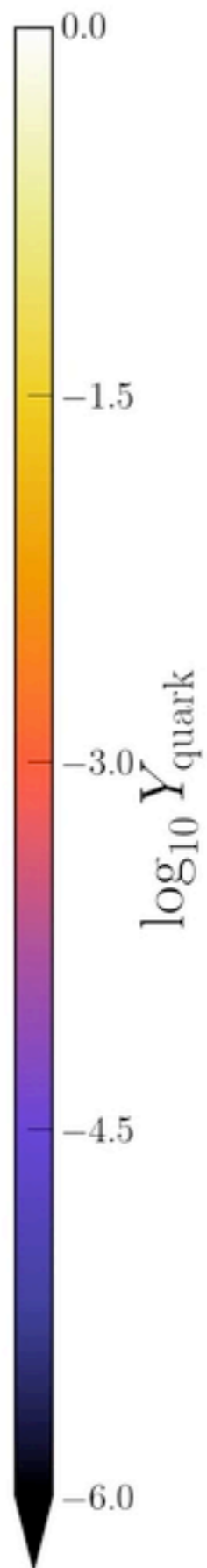
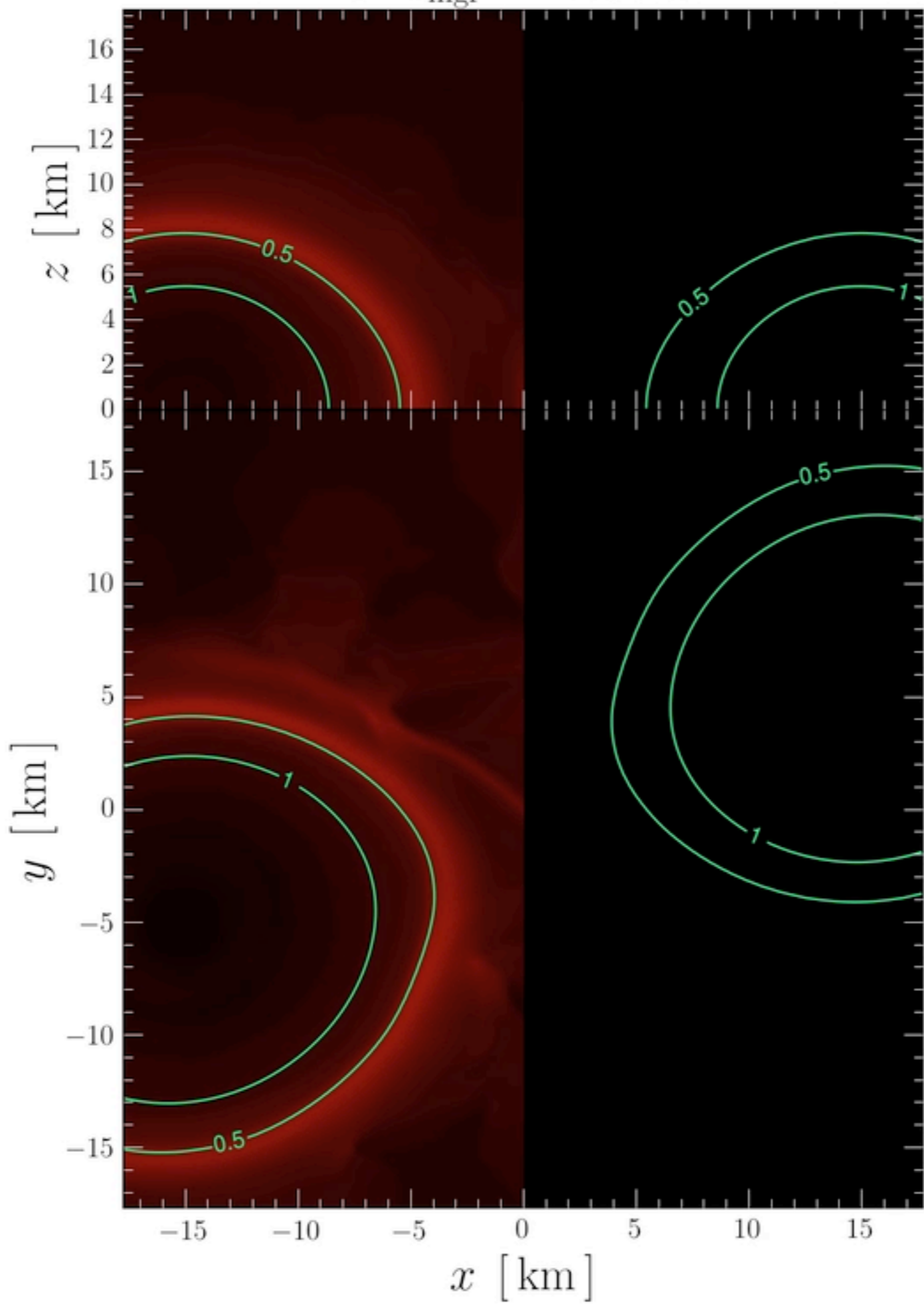
- \* Neutron-star **binary** mergers reach temperatures up to **80 MeV** and probe regions complementary to experiments

# Modelling the EOS

- EOS based on Chiral Mean Field (CMF) and nonlinear SU(3) sigma model
- Includes hyperons and quarks that can be turned on/off
- Uses Polyakov loop to implement a strong first order phase transition
- Includes a cross-over transition at high temperatures

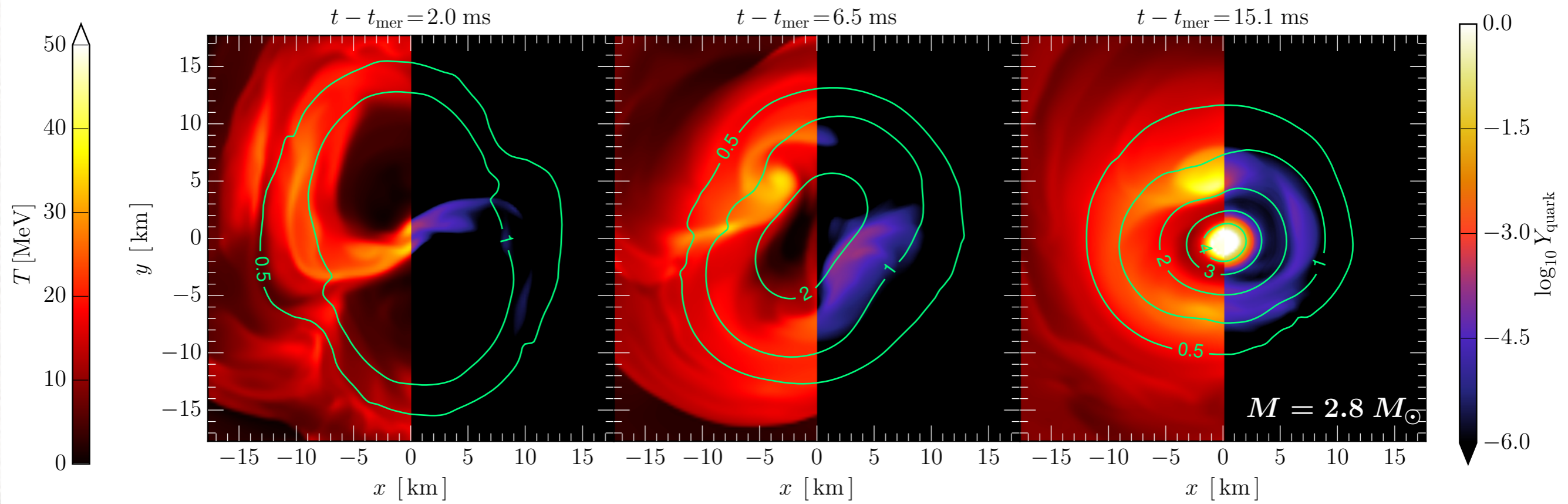


# Temperature

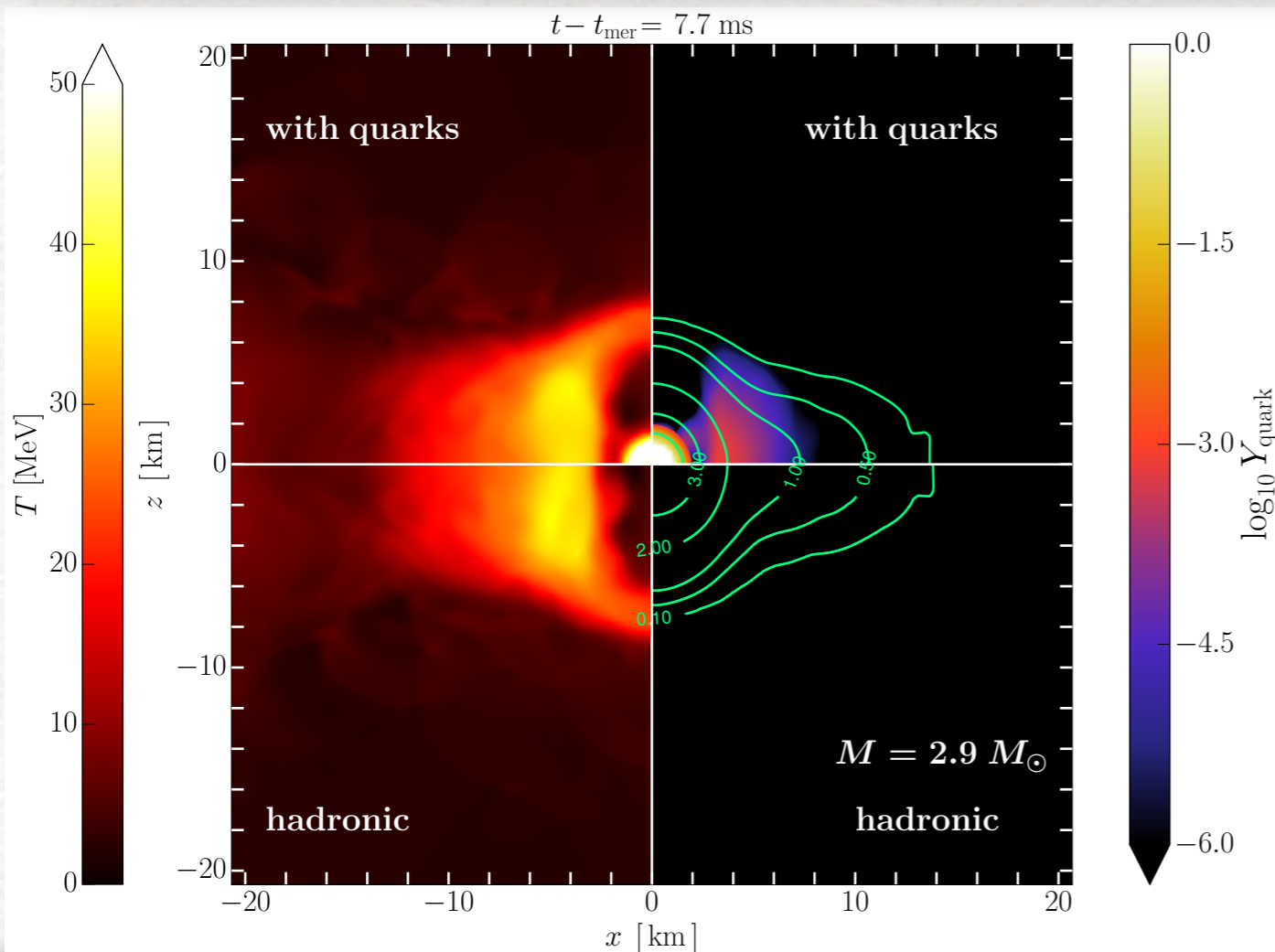
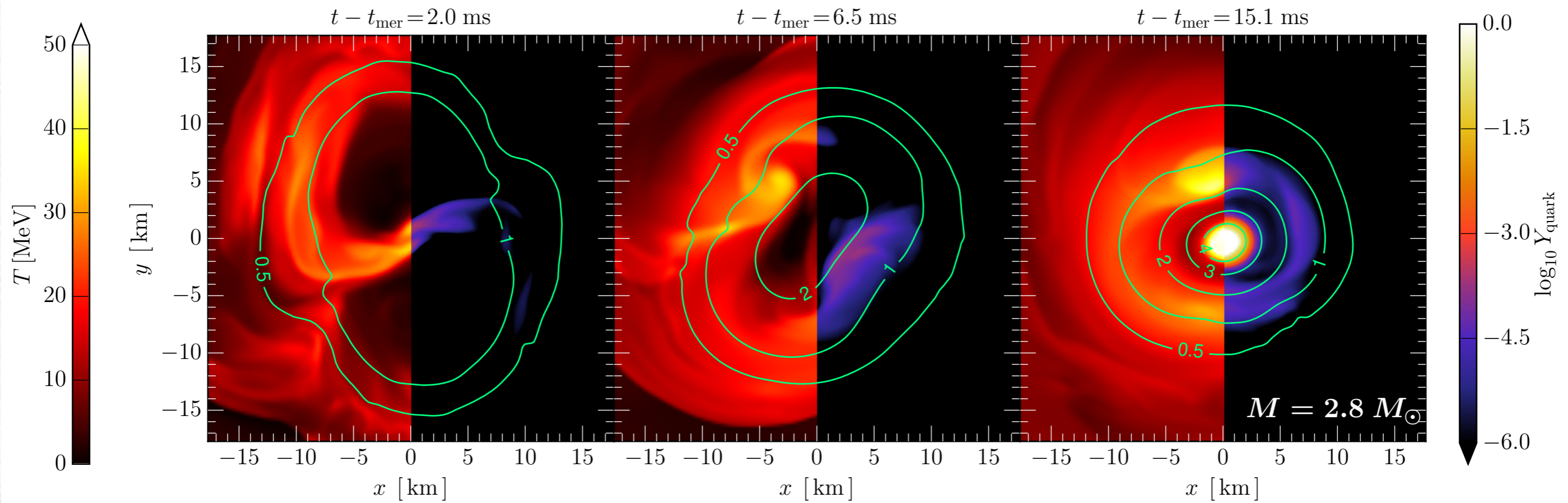


# Quark fraction



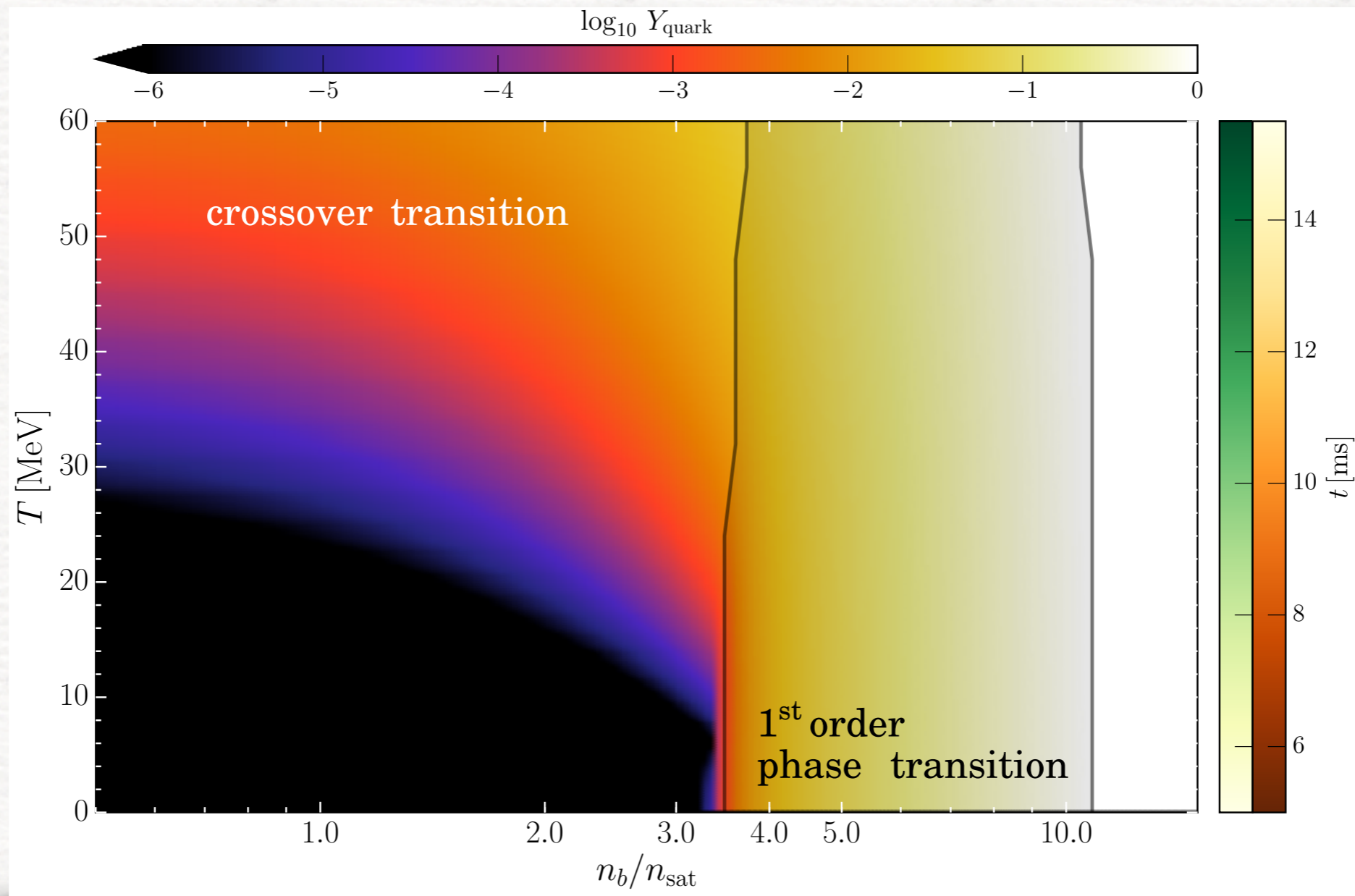


- EOS based on Chiral Mean Field (CMF) model, based on a nonlinear SU(3) sigma model.
- Quarks appear at sufficiently large temperatures and densities.



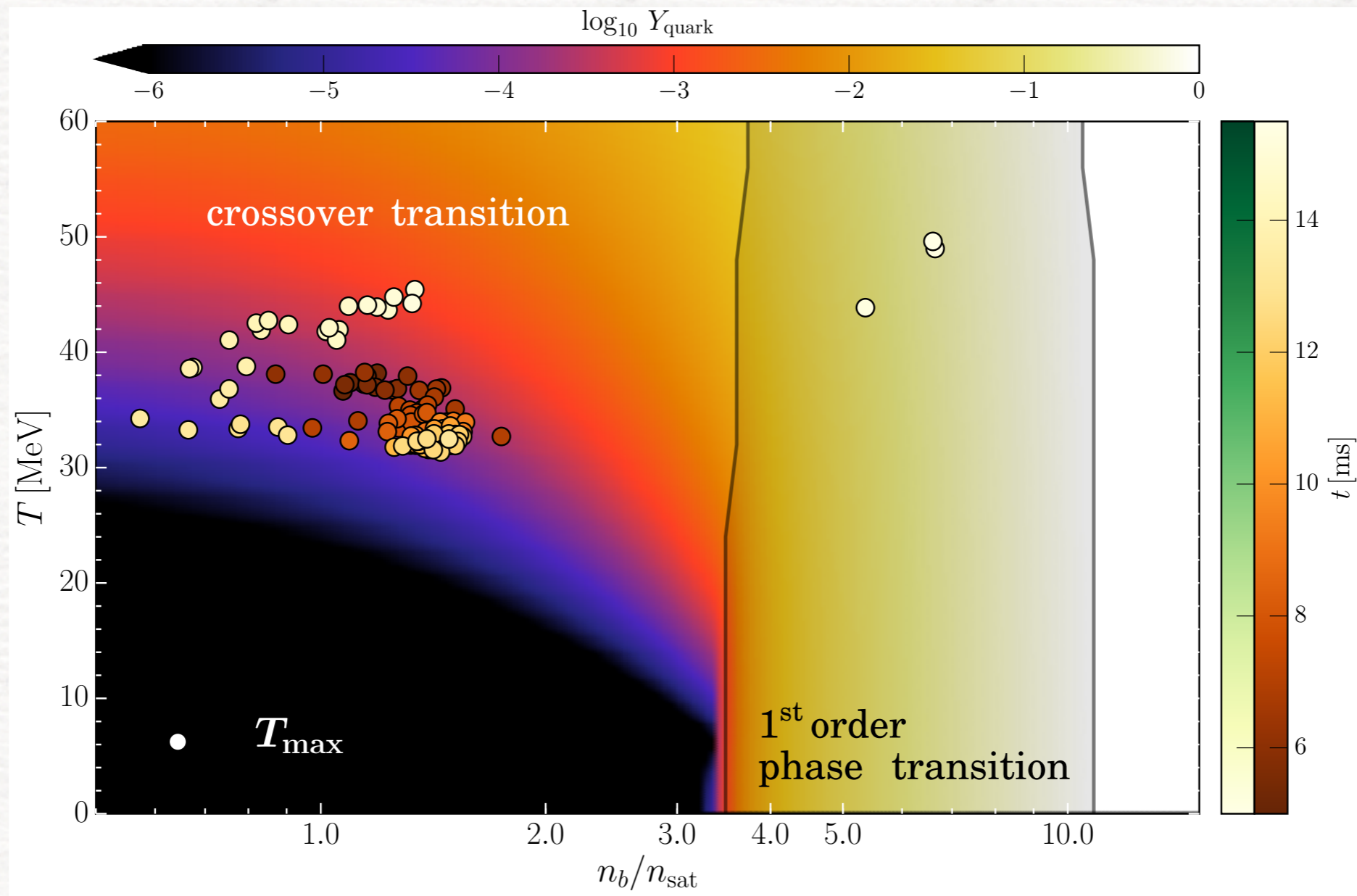
- EOS based on Chiral Mean Field (CMF) model, based on a nonlinear SU(3) sigma model.
- Quarks appear at sufficiently large temperatures and densities.
- For EOS without quarks, the dynamics is very similar, but no PT.

# Comparing with the phase diagram



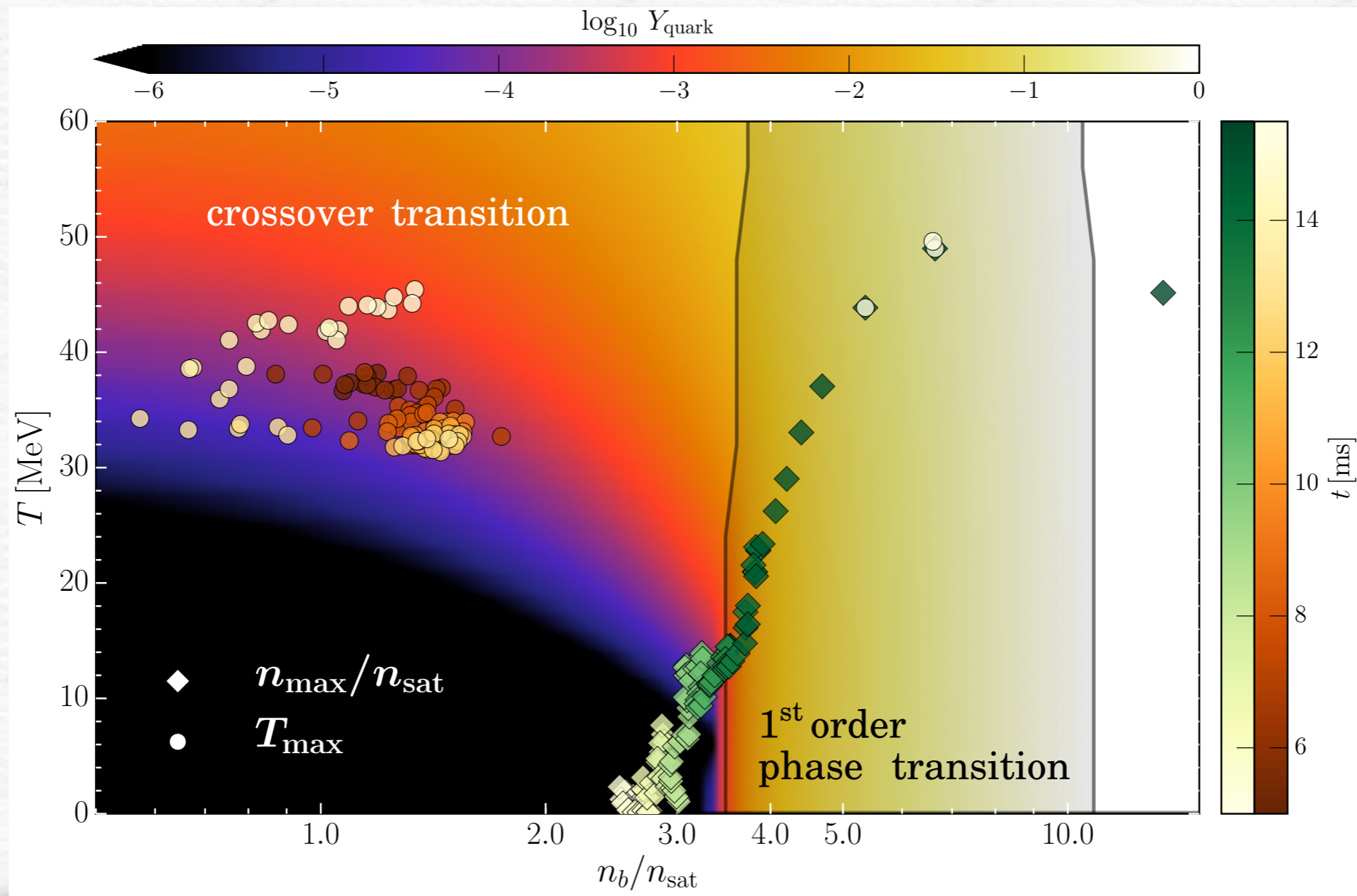
- Phase diagram with quark fraction

# Comparing with the phase diagram



- Phase diagram with quark fraction
- Circles show the position in the diagram of the maximum temperature as a function of time

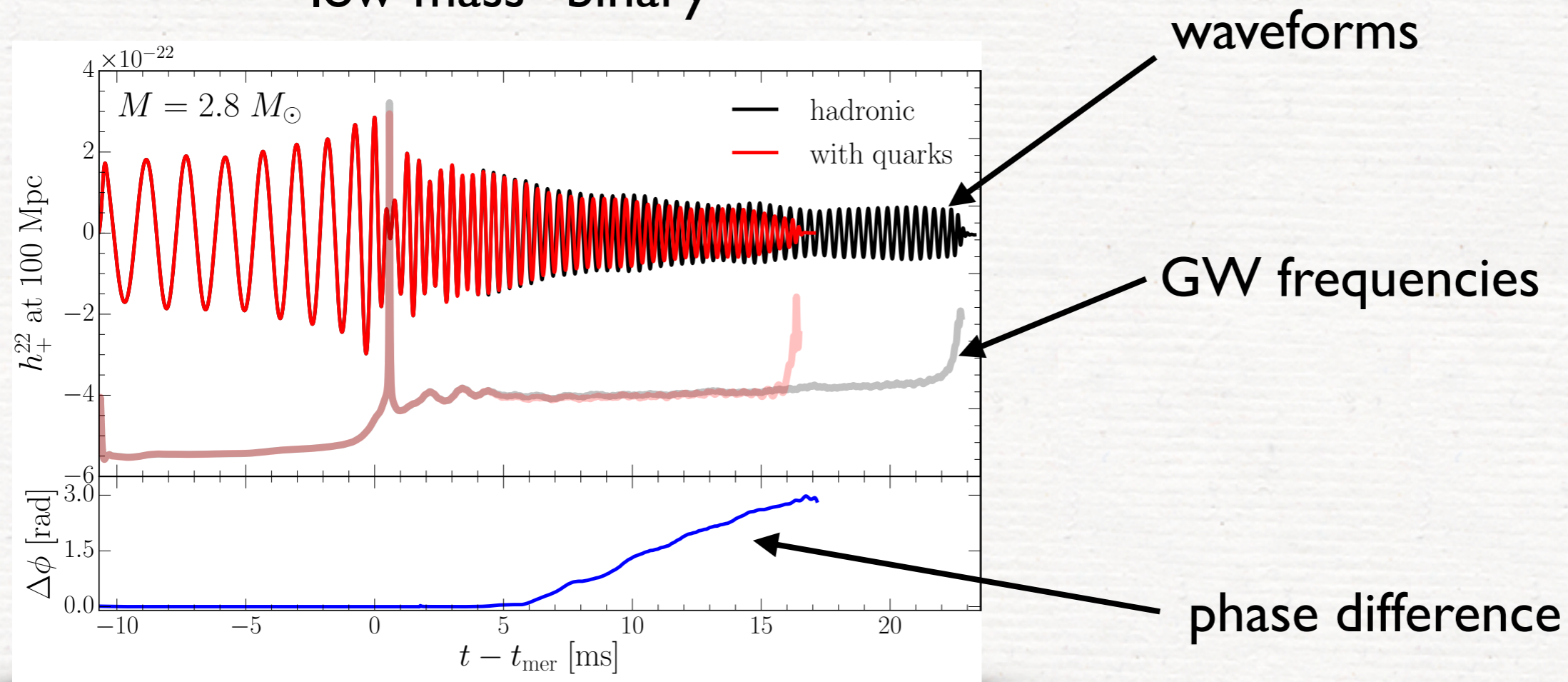
# Comparing with the phase diagram



- Reported are the evolution of the max. temperature and density.
- Quarks appear already early on, but only in small fractions.
- Once sufficient density is reached, a full phase transition takes place.

# Gravitational-wave emission

“low-mass” binary

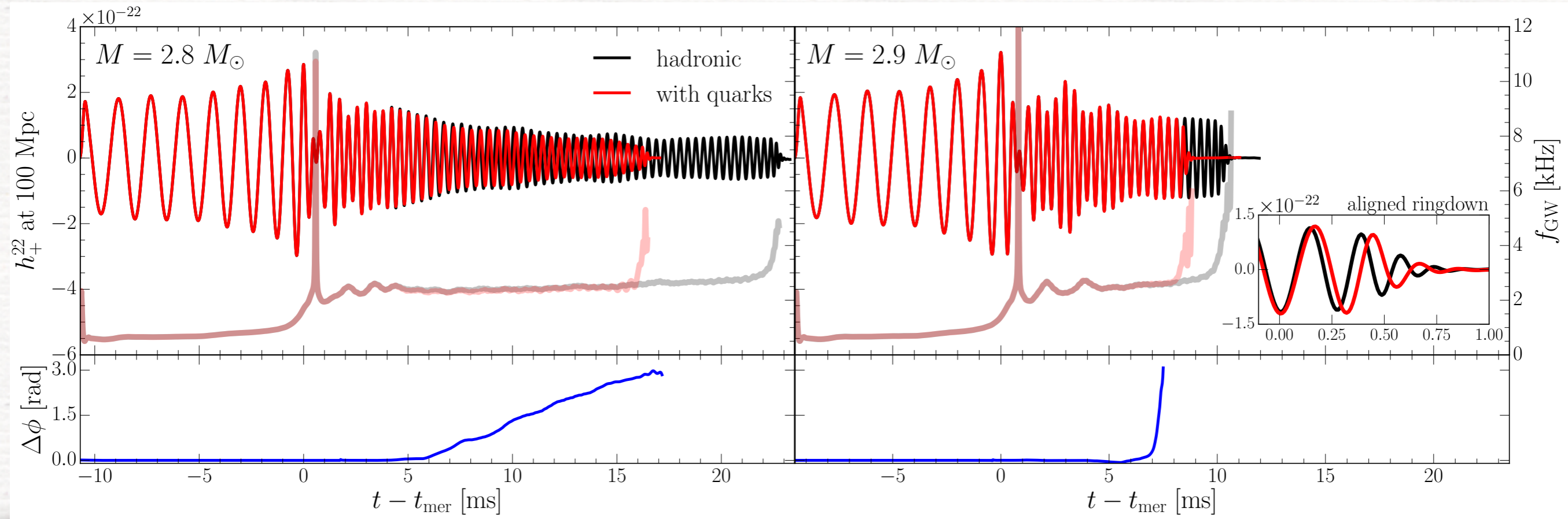


- In **low-mass binary**, after  $\sim 5$  ms, quark fraction is large enough to change quadrupole moment and yield differences in the waveforms.
- Note the phase difference is **zero** in the inspiral.
- Sudden softening of the phase transition leads to collapse and **large difference** in phase evolution.

# Gravitational-wave emission

“low-mass” binary

“high-mass” binary



- In **low-mass binary**, after  $\sim 5$  ms, quark fraction is large enough to change quadrupole moment and yield differences in the waveforms.
- In **high-mass binary**, phase transition takes place rapidly after  $\sim 5$  ms. Waveforms are similar but **ringdown** is **different** (free fall for PT). Observing mismatch between **inspiral** (fully hadronic) and **post-merger** (phase transition): clear **signature** of a **PT**.

# Conclusions

\* Spectra of post-merger shows peaks, some **"quasi-universal"**.

\* When used together with tens of observations, they will set tight constraints on EOS: radius known with  $\sim 1$  km precision.

\* **GW170817** has already provided new limits on

$$2.01_{-0.04}^{+0.04} \leq M_{\text{TOV}}/M_{\odot} \lesssim 2.16_{-0.15}^{+0.17} \quad \text{maximum mass}$$

$$12.00 < R_{1.4}/\text{km} < 13.45 \quad \tilde{\Lambda}_{1.4} > 375 \quad \text{radius, tidal deformability}$$

\* **GWs**: new probe of high-density matter and of **QCD** to be used in conjunction with **laboratory experiments**.

\* A phase transition after a BNS merger leaves a GW signature and opens a gate to access quark matter beyond accelerators