



Production and flow of light-flavor hadrons and nuclei in small and large systems with ALICE

- Particle multiplicity, spectra and yields
- Radial flow and azimuthal anisotropy
- Nuclear modification factors

Jacek Otwinowski (IFJ PAN, Krakow)
On behalf of the ALICE Collaboration

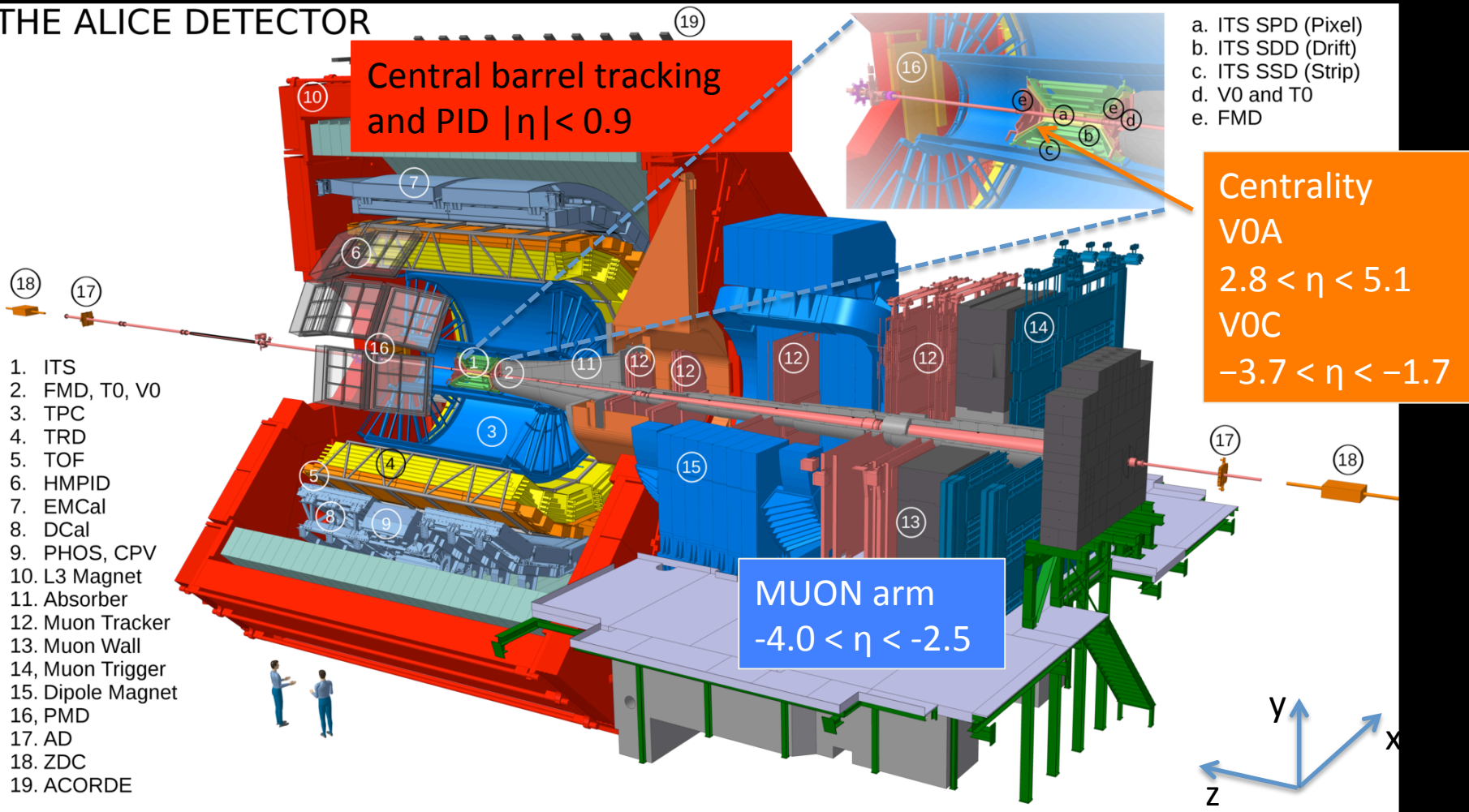


A Large Ion Collider Experiment



- Excellent particle identification capabilities over a wide p_T range 0.1-20 GeV/c
- Good momentum resolution $\sim 1-5\%$ for $p_T = 0.1-50$ GeV/c

THE ALICE DETECTOR



ALICE at work since 2009



System	Year	$\sqrt{s_{NN}}$ (TeV)	L_{int}
Pb-Pb	2010-2011	2.76	$\sim 75 \mu\text{b}^{-1}$
	2015	5.02	$\sim 250 \mu\text{b}^{-1}$
	2018	5.02	$\sim 0.9 \text{nb}^{-1}$
Xe-Xe	2017	5.44	$\sim 0.3 \mu\text{b}^{-1}$
p-Pb	2013	5.02	$\sim 15 \text{nb}^{-1}$
	2016	5.02, 8.16	$\sim 3 \text{nb}^{-1}, \sim 25 \text{nb}^{-1}$
pp	2009-2013	0.9, 2.76, 7, 8	$\sim 200 \mu\text{b}^{-1}, \sim 100 \mu\text{b}^{-1}, \sim 1.5 \text{pb}^{-1}, \sim 2.5 \text{pb}^{-1}$
	2015-2018	5.02, 13	$\sim 1.3 \text{pb}^{-1}, \sim 59 \text{pb}^{-1}$

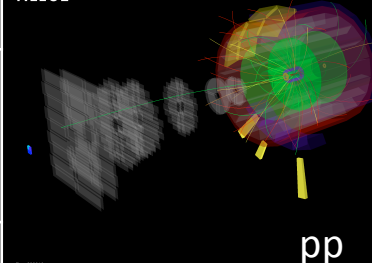
- Energy and system dependence studies of particle production are possible
- Large statistics of pp, p-Pb and Pb-Pb collisions at the same $\sqrt{s_{NN}}$

→ Precise comparison studies

Quest for the Quark-Gluon Plasma

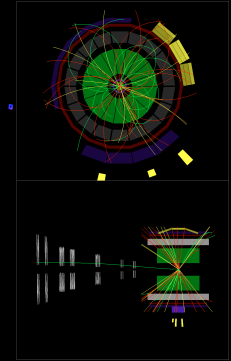


ALICE

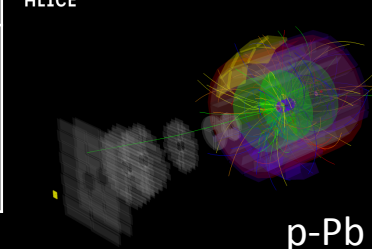


pp

Run 2017
Triggering: 2017-11-11 21:38:21 (UTC)
Collision system: p-p
Energy: 5.02 TeV

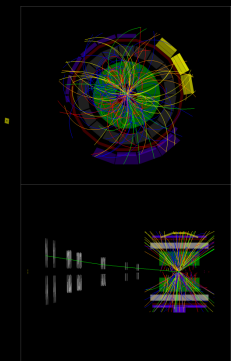


ALICE

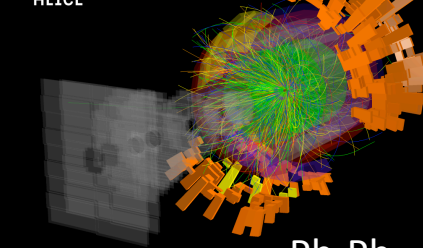


p-Pb

Run 2015
Triggering: 2015-11-14 14:12:07 (UTC)
Collision system: p-Pb
Energy: 5.02 TeV

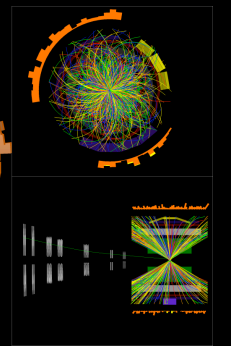


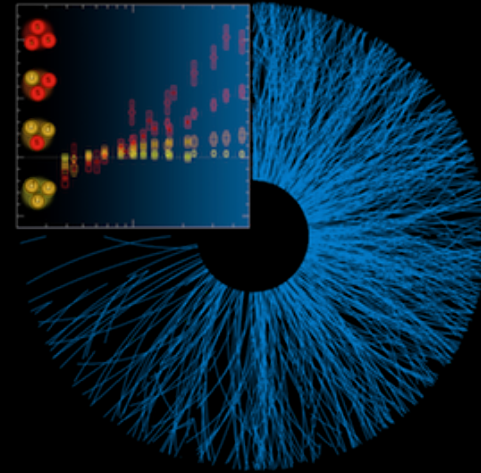
ALICE



Pb-Pb

Run 2015
Triggering: 2015-11-04 20:38:06 (UTC)
Collision system: Pb-Pb
Energy: 5.02 TeV



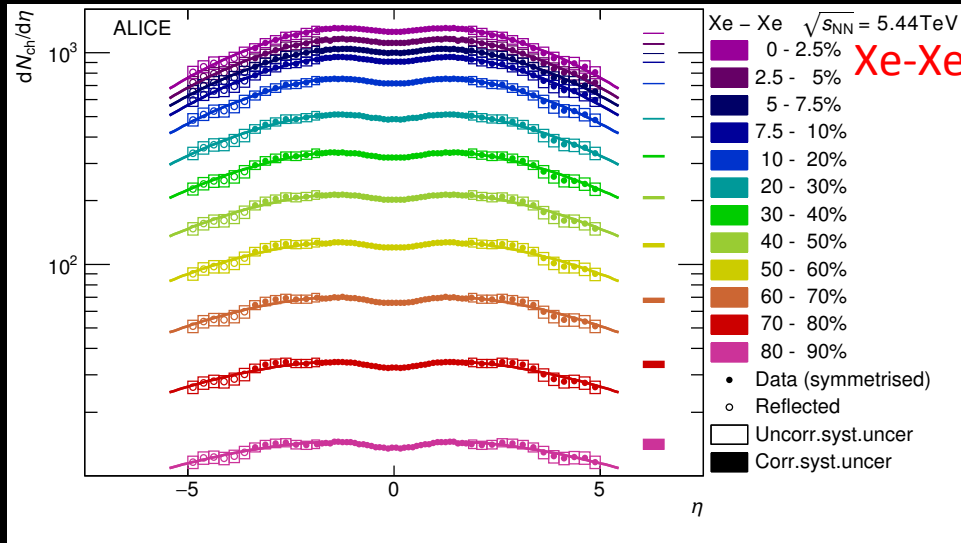


PARTICLE MULTIPLICITY, SPECTRA AND YIELDS

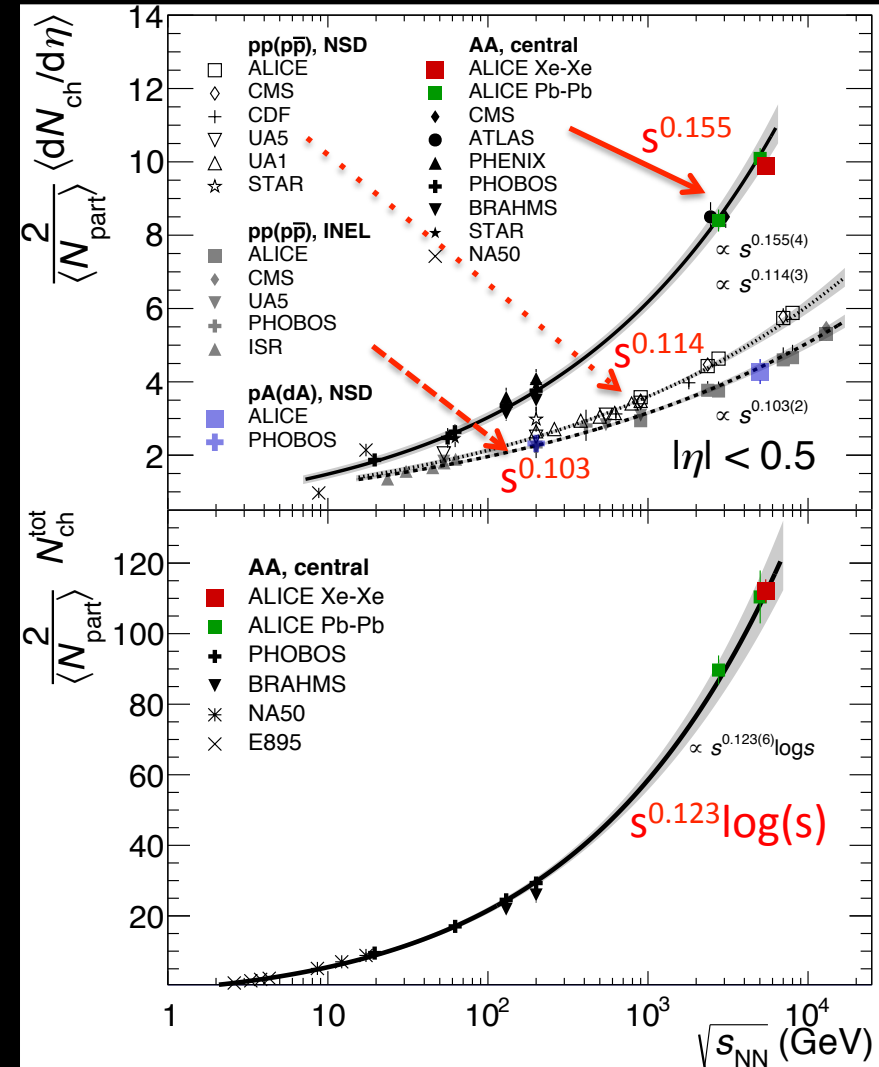
Charged-particle multiplicity in pp, p-Pb, Pb-Pb and Xe-Xe



Phys. Lett. B 790 (2019) 35



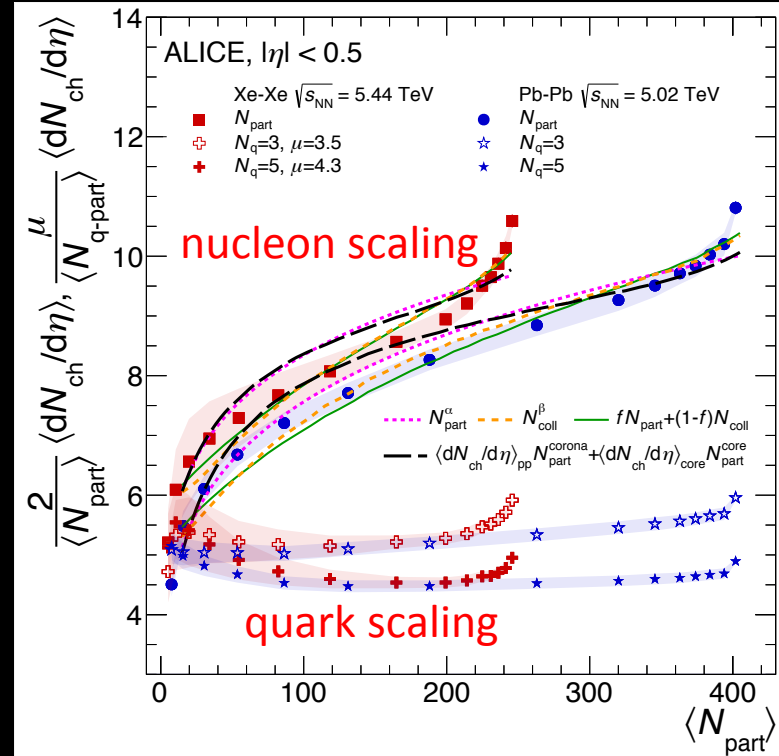
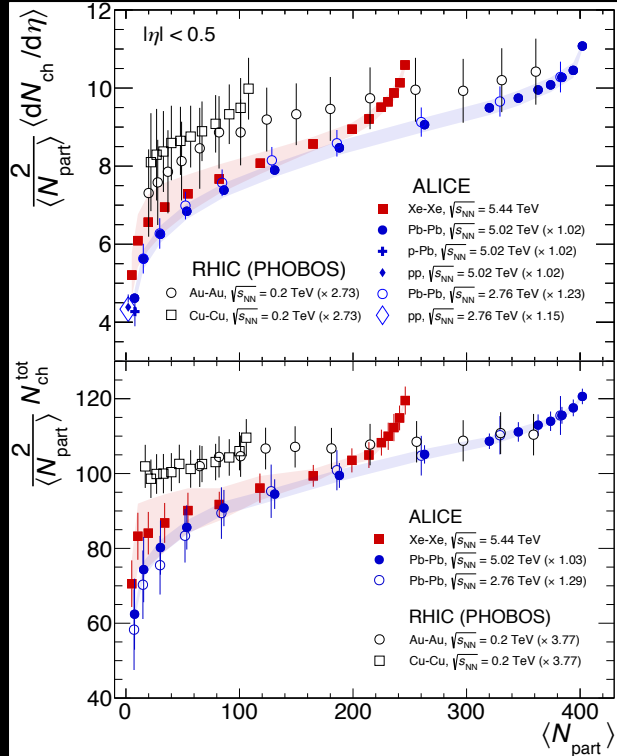
- Measurement in broad rapidity range and several centrality intervals
- \sqrt{s} dependence in A-A collisions differs from pp and p-Pb (**no universal scaling**)
- Total charged-particle multiplicity from extrapolation up to $\eta = \pm y_{\text{beam}}$ in unmeasured region



Charged-particle multiplicity vs N_{part} in pp, p-Pb, Pb-Pb and Xe-Xe

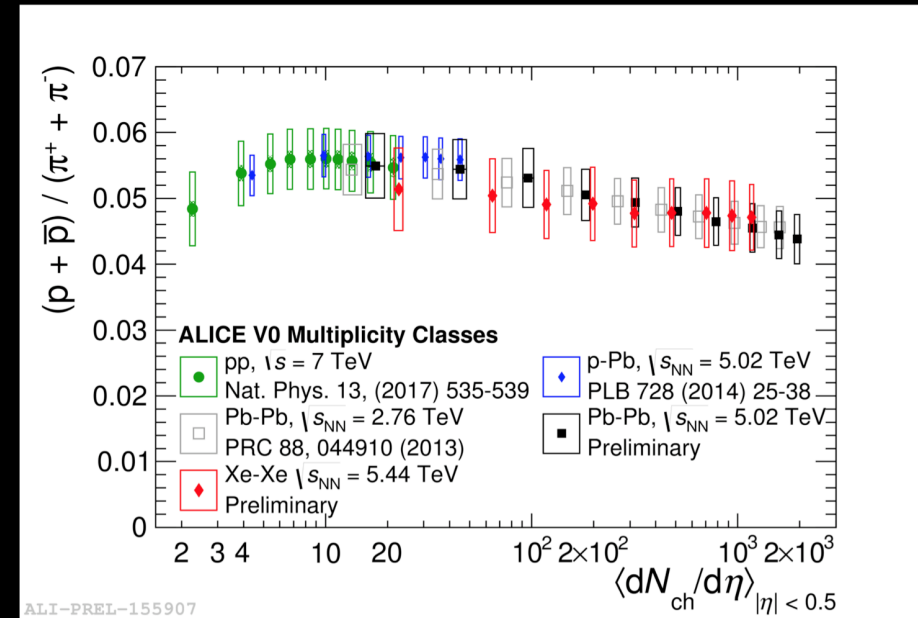
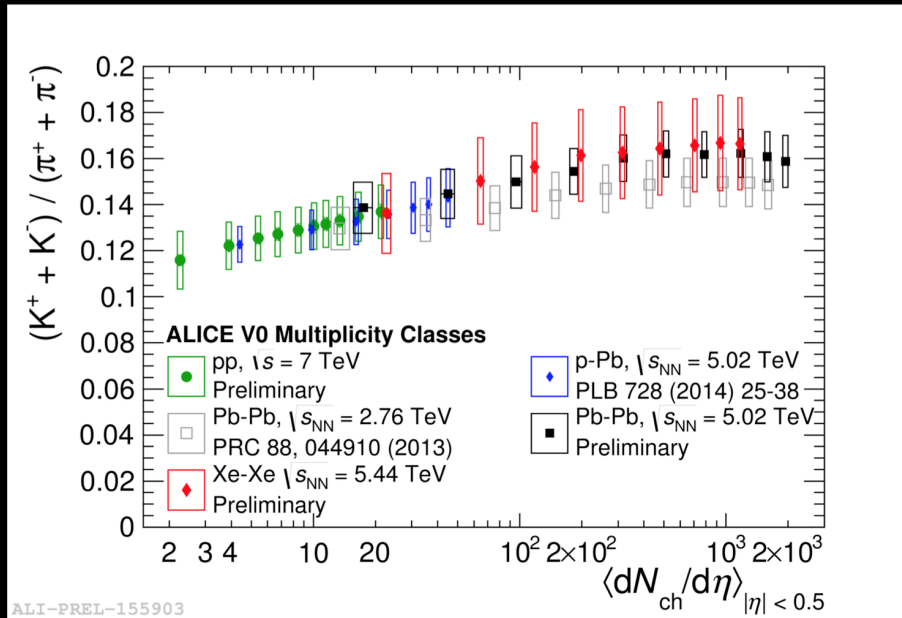


Phys. Lett. B 790 (2019) 35



- Charged-particle multiplicity density and total multiplicity as a function of centrality
 - Deviation from N_{part} scaling
 - Steeper rise in most central Xe-Xe and Pb-Pb collisions due to upward fluctuations
- Comparison to scaling ($\sim N_{\text{part}}$, $\sim N_{\text{coll}}$, core+corona, $\sim N_{\text{qpart}}$) shows that collision geometry plays an important role in particle production

K/ π and p/π in pp, p-Pb, Pb-Pb and Xe-Xe



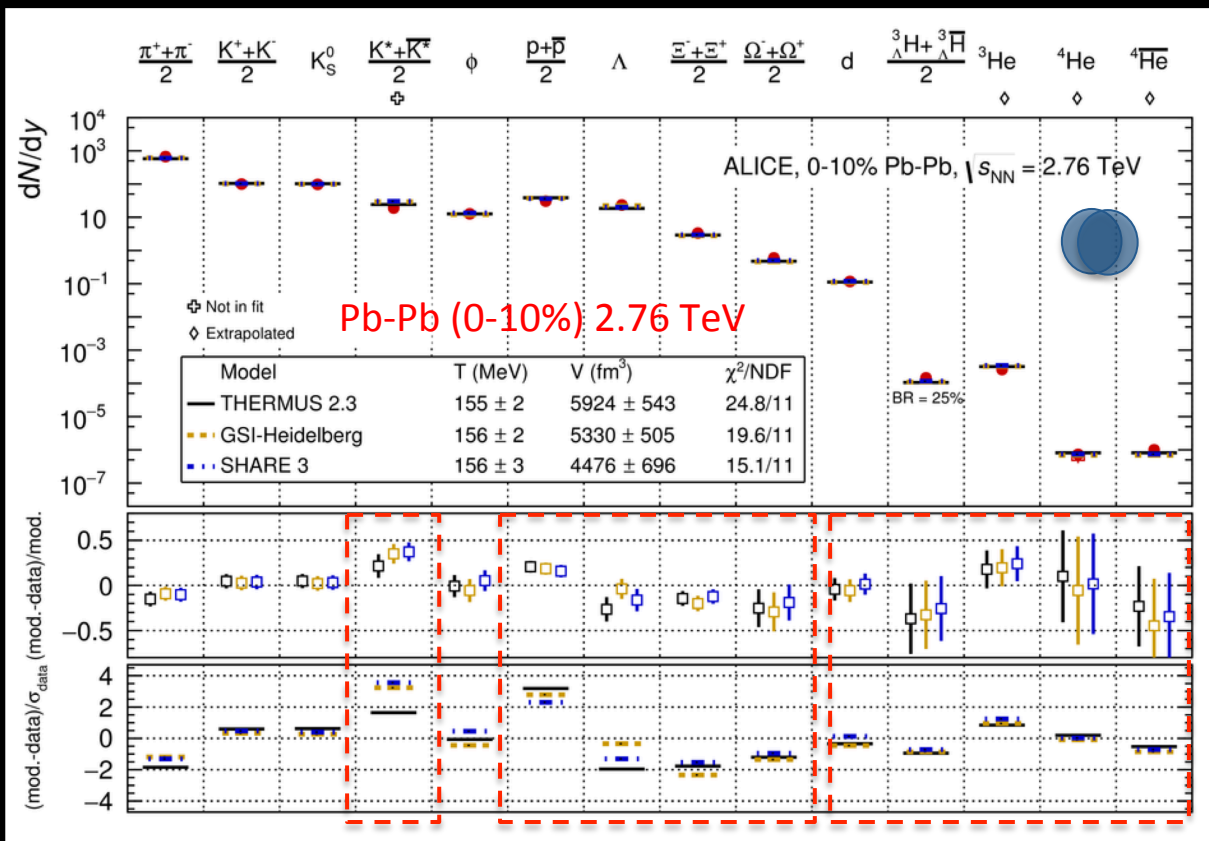
- No significant energy dependence is observed
- K/ π and p/π are consistent for all collision systems at similar multiplicity

→ Particle production is driven by the characteristics of final state

M. Sefcik 11/06

Particle yields vs thermal models in Pb-Pb at $\sqrt{s_{NN}} = 2.76$ TeV

Nucl. Phys. A 971 (2018) 1



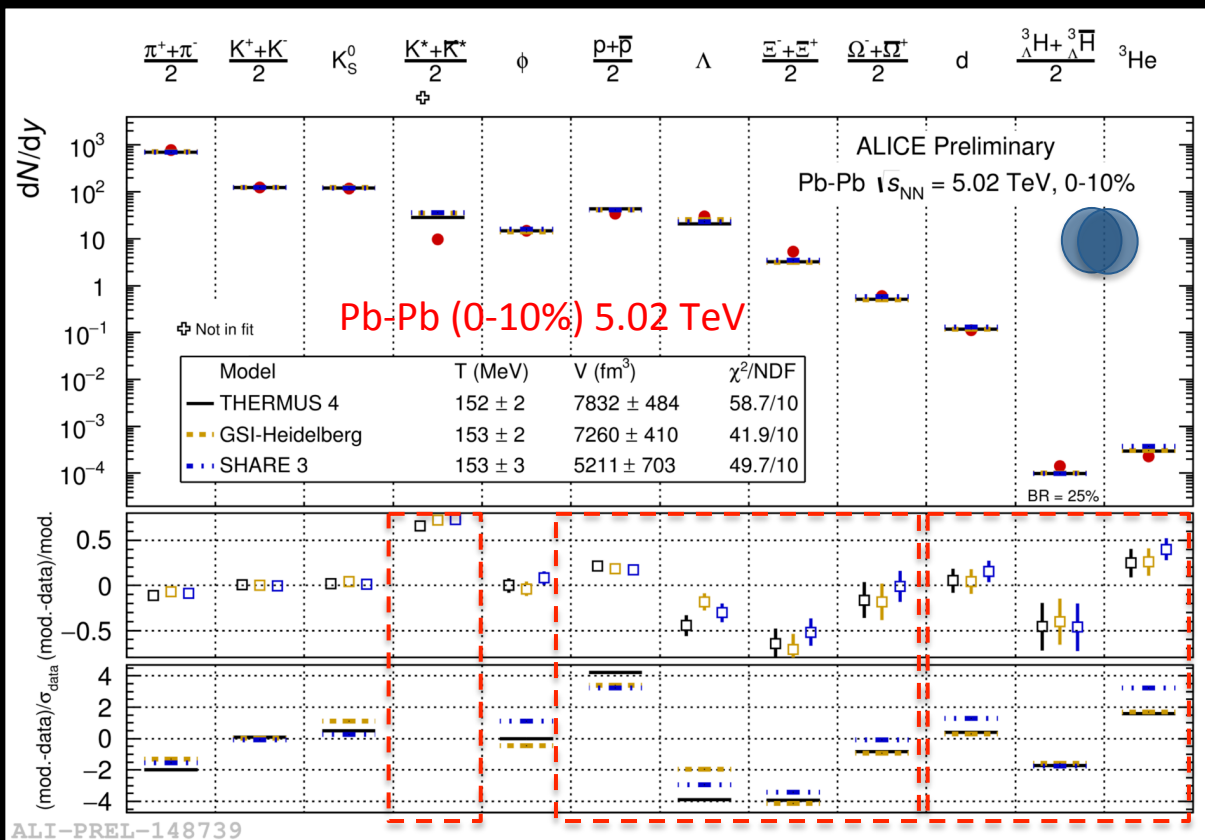
- Production of (most) hadrons well described at freezeout temperature $T_{ch} \sim 156$ MeV
- K^{*0} resonance (not included in the fit): production overestimated by thermal models
- Tension for protons and multi-strange baryons
- Light nuclei production described by thermal models (binding energy $\ll T_{ch}$)?

THERMUS: Wheaton et al., Comput. Phys. Commun, 180 (2009) 84

GSI-Heidelberg: Andronic et al., Phys. Lett. B 673 142

SHARE: Petran et al., Comp. Phys. Commun. 195 (2014) 2056

Particle yields vs thermal models in Pb-Pb at $\sqrt{s}_{NN} = 5.02$ TeV



- Production of (most) hadrons well described at freeze-out temperature $T_{ch} \sim 153$ MeV (caused by protons)
- K^{*0} resonance (not included in the fit): production overestimated by thermal models
- Tension for protons and multi-strange baryons
- Light nuclei production described by thermal models (binding energy $\ll T_{ch}$)?

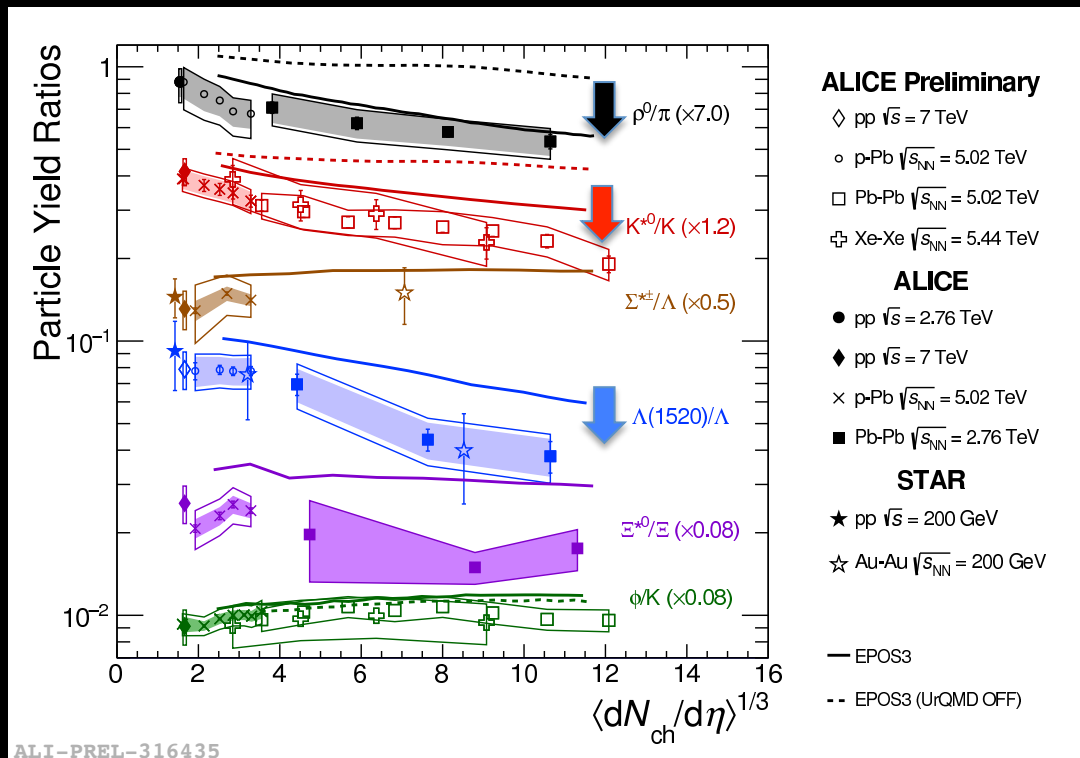
Fit χ^2 worst than at 2.76 TeV (more precise data)

THERMUS: Wheaton et al., Comput. Phys. Commun, 180 (2009) 84
 GSI-Heidelberg: Andronic et al., Phys. Lett. B 673 142
 SHARE: Petran et al., Comp. Phys. Commun. 195 (2014) 2056

Relative resonance production in pp, p-Pb, Pb-Pb and Xe-Xe collisions



Resonance	ρ^0	K^{*0}	$\Sigma^{*\pm}$	$\Lambda(1520)$	Ξ^{*0}	ϕ
Lifetime (fm/c)	1.3	4.16	5.5	12.6	21.7	46.2



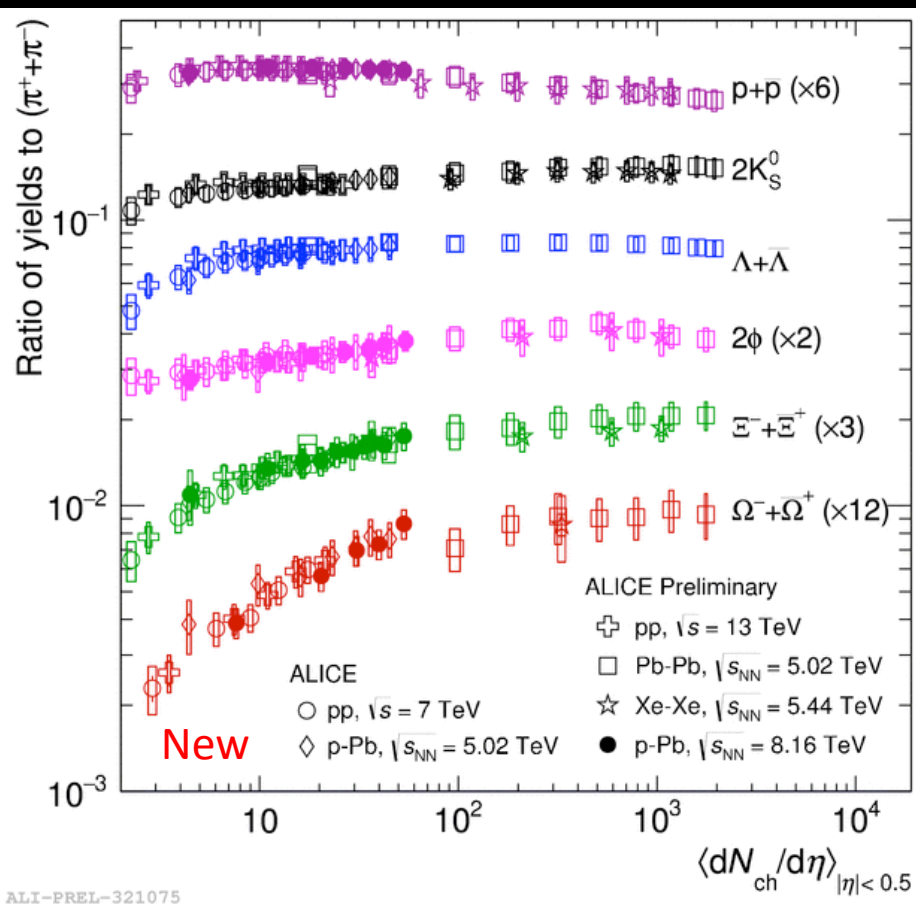
- Relative suppression of ρ^0 , K^{*0} , $\Lambda(1520)$ with increasing multiplicity
- Σ^*/Λ , Ξ^*/Ξ , ϕ/K is independent of multiplicity
- Similar trend seen in all collision systems
- EPOS3 + UrQMD describes the trend of data

S. Tripathy 13/06

ρ^0 : Phys. Rev. C99, (2019) 064901
 K^{*0} : Phys. Rev. C95, (2017) 064606
 $\Lambda(1520)$: Phys. Rev. C99, (2019) 024905
 STAR, Phys.Rev.C78 (2008) 044906
 STAR, Phys. Rev. Lett. 97 (2006) 132301
 ϕ : Phys. Rev. C91 (2015) 024609
 Σ^* , Ξ^* : Eur. Phys. J. C 77 (2017) 389

→ Dominance of rescattering over (re)combination in hadronic phase

Relative strangeness production in pp, p-Pb, Pb-Pb and Xe-Xe



Historically a signature of the QGP
Rafelski & Mueller, Phys. Rev. Lett. 48 (1982) 1066

- Smooth evolution from pp to Pb-Pb
- Enhancement increases with strangeness content
- No significant energy and system dependence is observed at similar multiplicity

→ Strangeness production is driven by the characteristics of final state

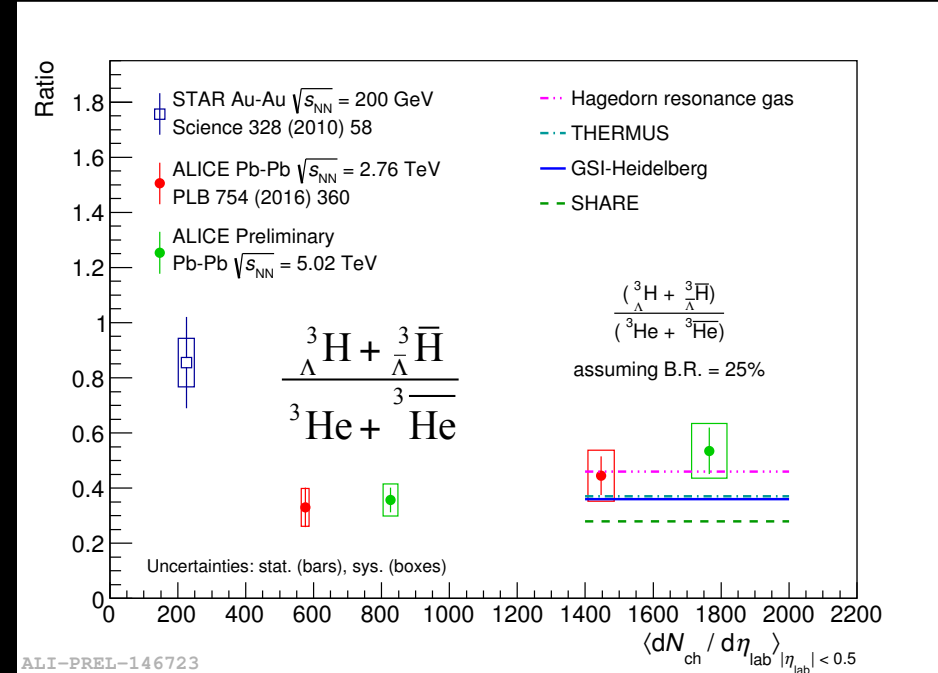
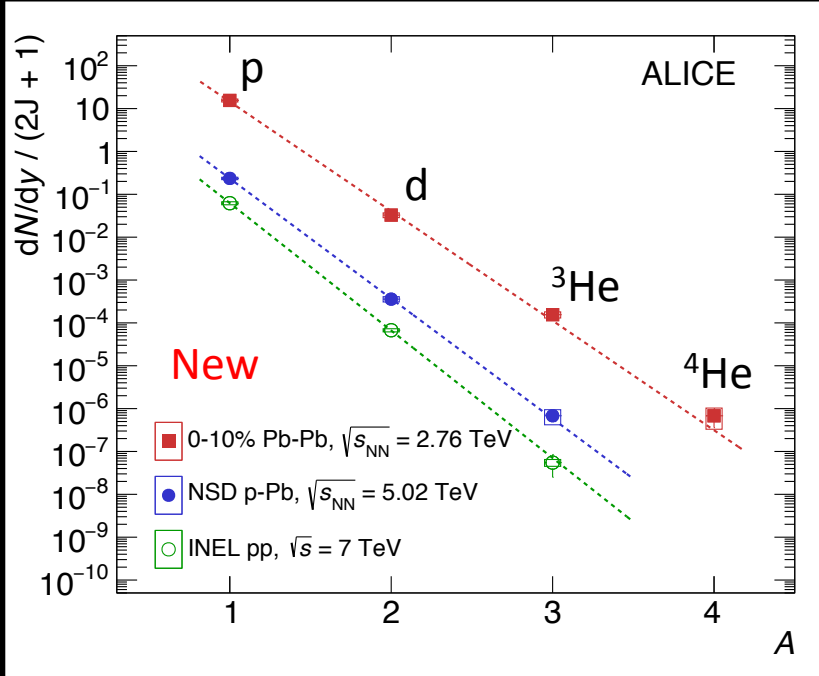
pp 7 TeV: Nature Physics 13 (2017) 535
p-Pb 5.02 TeV: Phys. Lett. B728 (2014) 25, Phys. Lett. B758 (2016) 389

E. Willsher 11/06

Nuclei production in pp, p-Pb and Pb-Pb



p-Pb: arXiv:1906.03136



- Exponential decrease in nuclei rate in agreement with thermal model predictions
- Nuclei production rate decrease by factor of ~ 300 (Pb-Pb), ~ 600 (p-Pb) and ~ 1000 (pp) for each additional nucleon
- ${}^3_{\Lambda}\text{H}$ production consistent with thermal model (binding energy ~ 0.13 MeV $\ll T_{ch}$)
 \rightarrow Production mechanisms: thermal vs coalescence?

pp: Phys. Rev. C97 (2018) 024615

Pb-Pb: Nucl. Phys. A 971 (2018) 1

11-06-2019

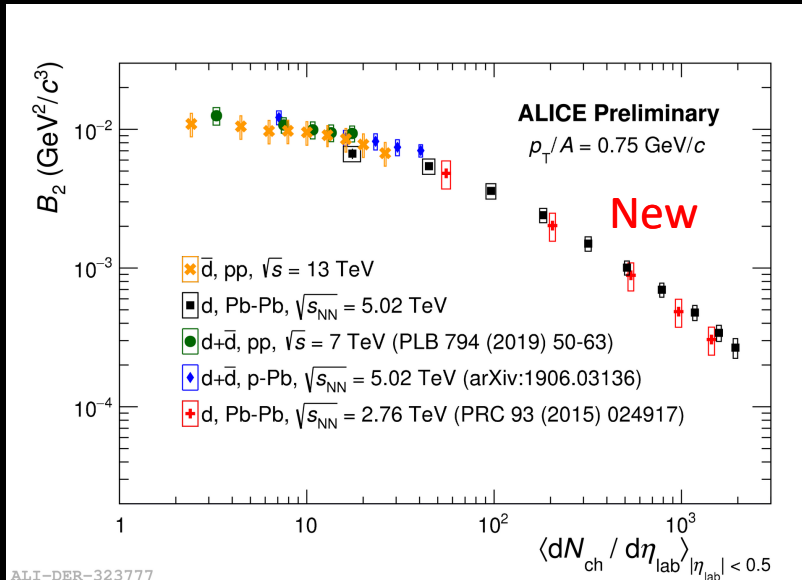
SQM2019 - Jacek Otwinowski

L. Barioglio 11/06

S. Bufalino 13/06 (${}^3_{\Lambda}\text{H}$ life time)

12

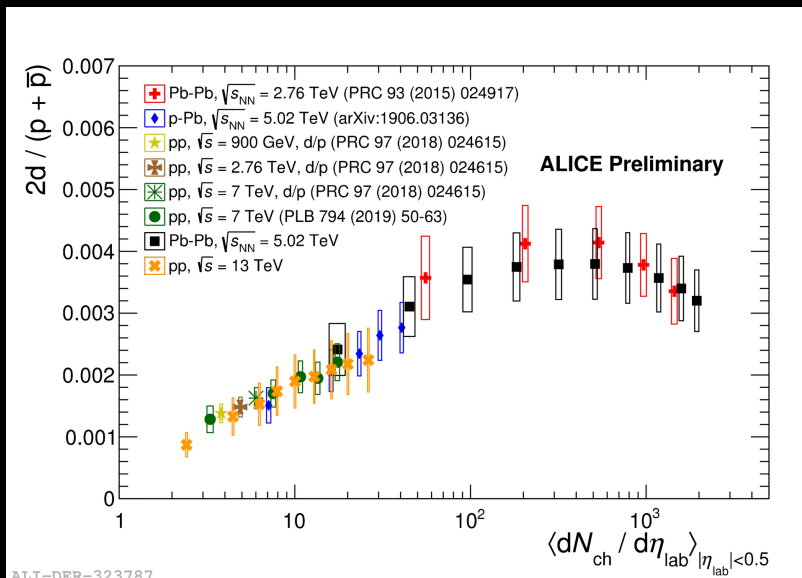
Formation of light nuclei: (anti) deuterons



- Coalescence of baryons close in phase space (A - mass number)

$$E_A \frac{d^3 N_A}{dp_A^3} = B_A \left(E_p \frac{d^3 N_p}{dp_p^3} \right)^A$$

- B_2 shows dependence on multiplicity (no dependence on p_T)

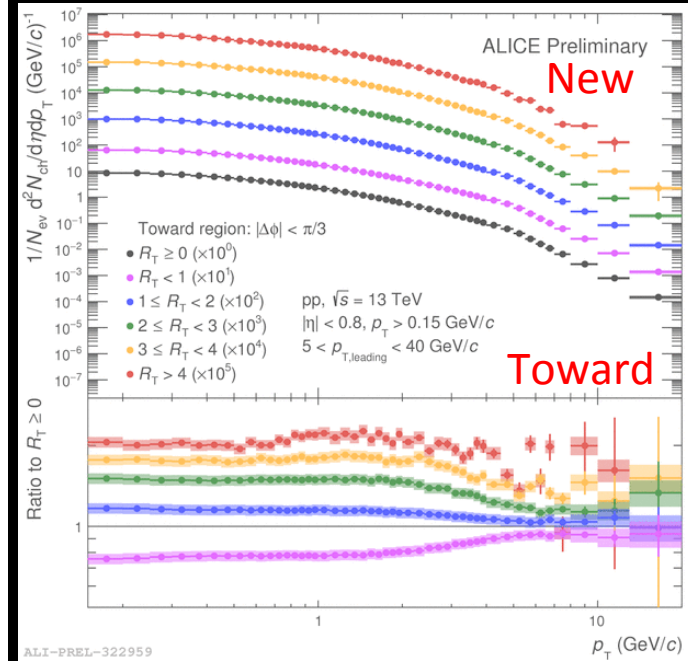
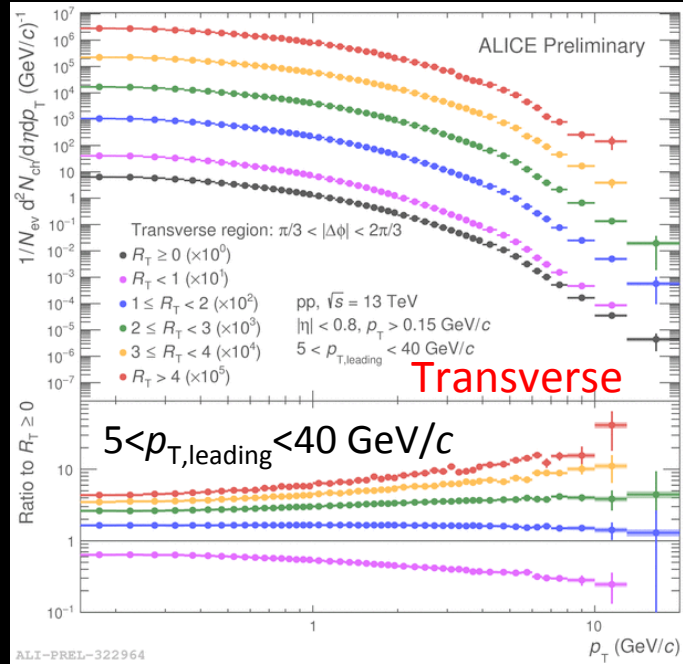
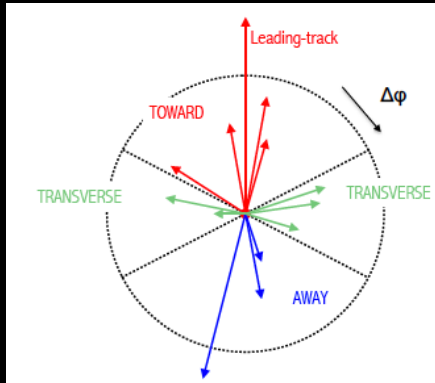


- d/p vs multiplicity
 - Increase from pp to peripheral Pb-Pb consistent with coalescence model
 - No centrality dependence in high multiplicity Pb-Pb (yields consistent with thermal model)

→ Production mechanisms: thermal vs coalescence?

SPECTRA

Transverse/toward p_T spectra in pp



Transverse activity classifier:

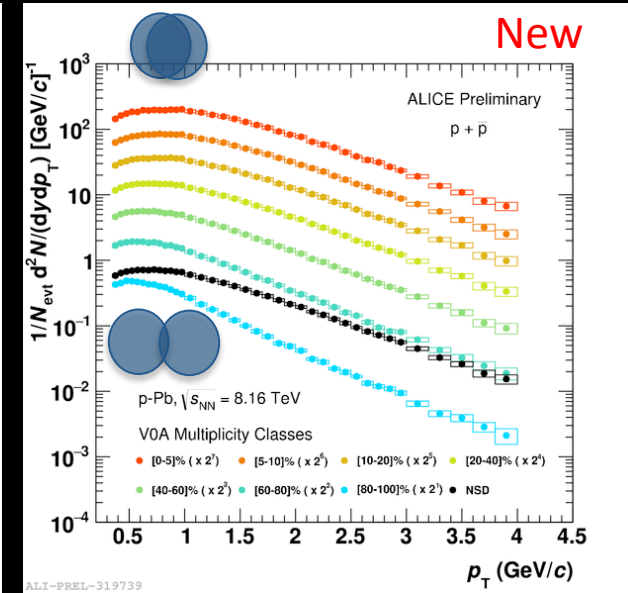
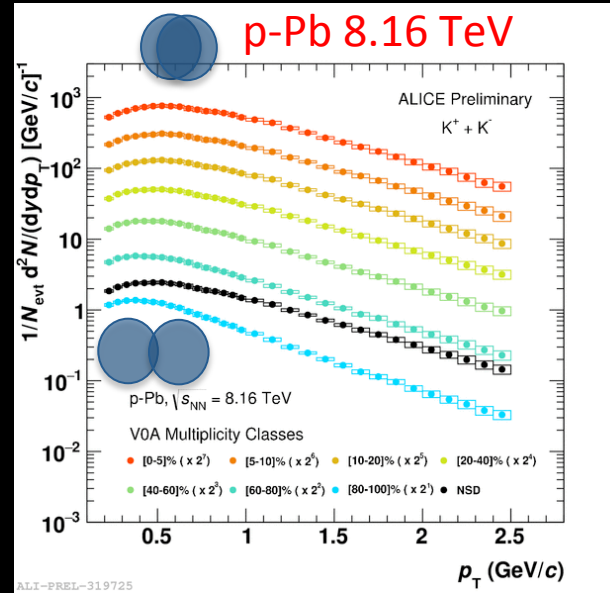
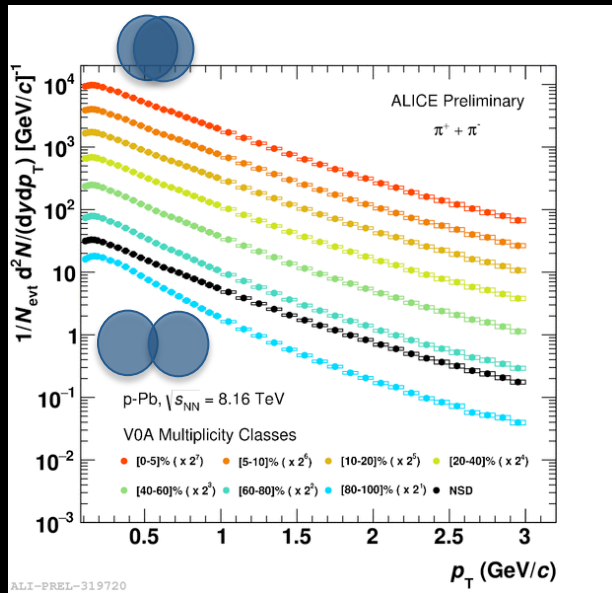
- $R_T = N_{\text{inc.}} / \langle N_{\text{inc.}} \rangle$

Martin et al.
 Eur. Phys. J.C 76 (2016) 299

- p_T spectra in multiplicity intervals are biased due to correlations between low- and high- p_T particles \rightarrow harder spectra in higher multiplicity events [ALICE, arXiv:1905.07208]
- R_T selection (**jet free multiplicity estimator**) can be used to reduce such biases and help to understand these correlations
- Opposite trend at high p_T for the transverse and toward spectra
- Convergence to the inclusive jets for the toward spectra at high $p_T \rightarrow$ **separation between soft (UE) and hard (jet) part of the spectrum**

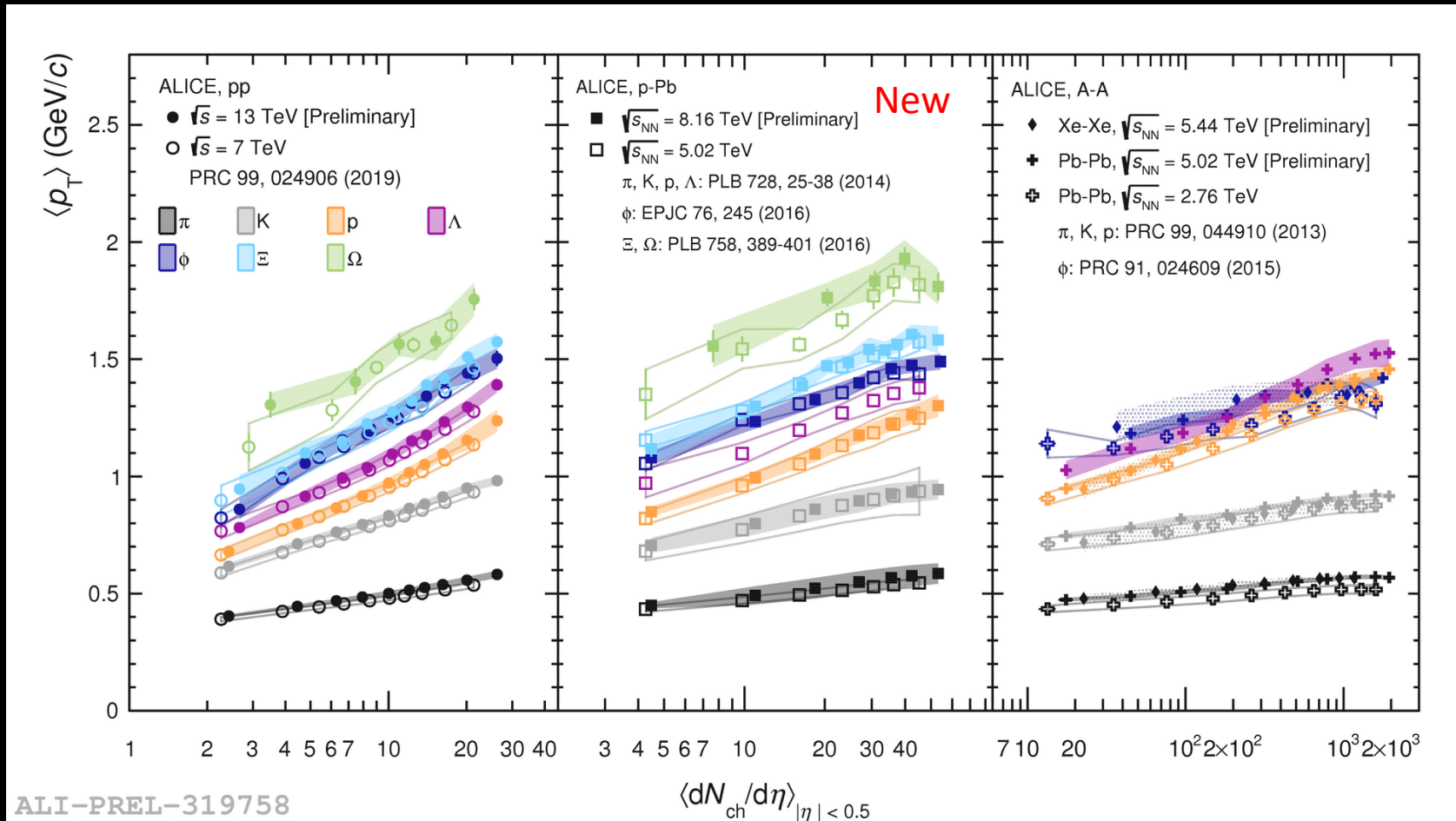
V. Zacco 11/06

p_T spectra of charged π , K and p in p-Pb



- Measured and identified with different analysis techniques: ITS, TPC, TOF, HMPID and topological identification of decaying charged kaons
- Mass dependent hardening of the spectra with increasing centrality
→ Collective radial expansion

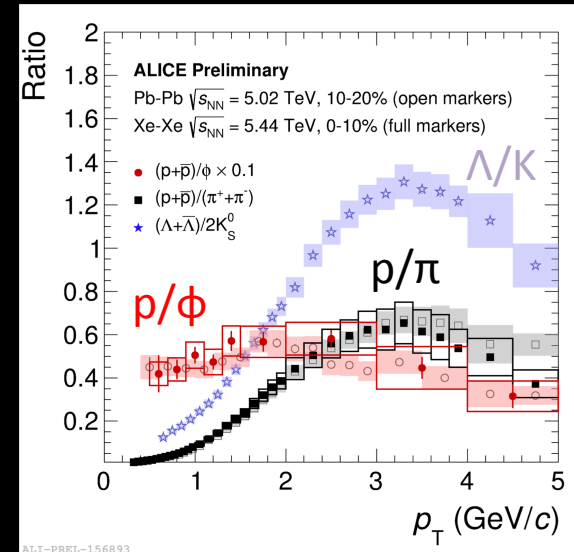
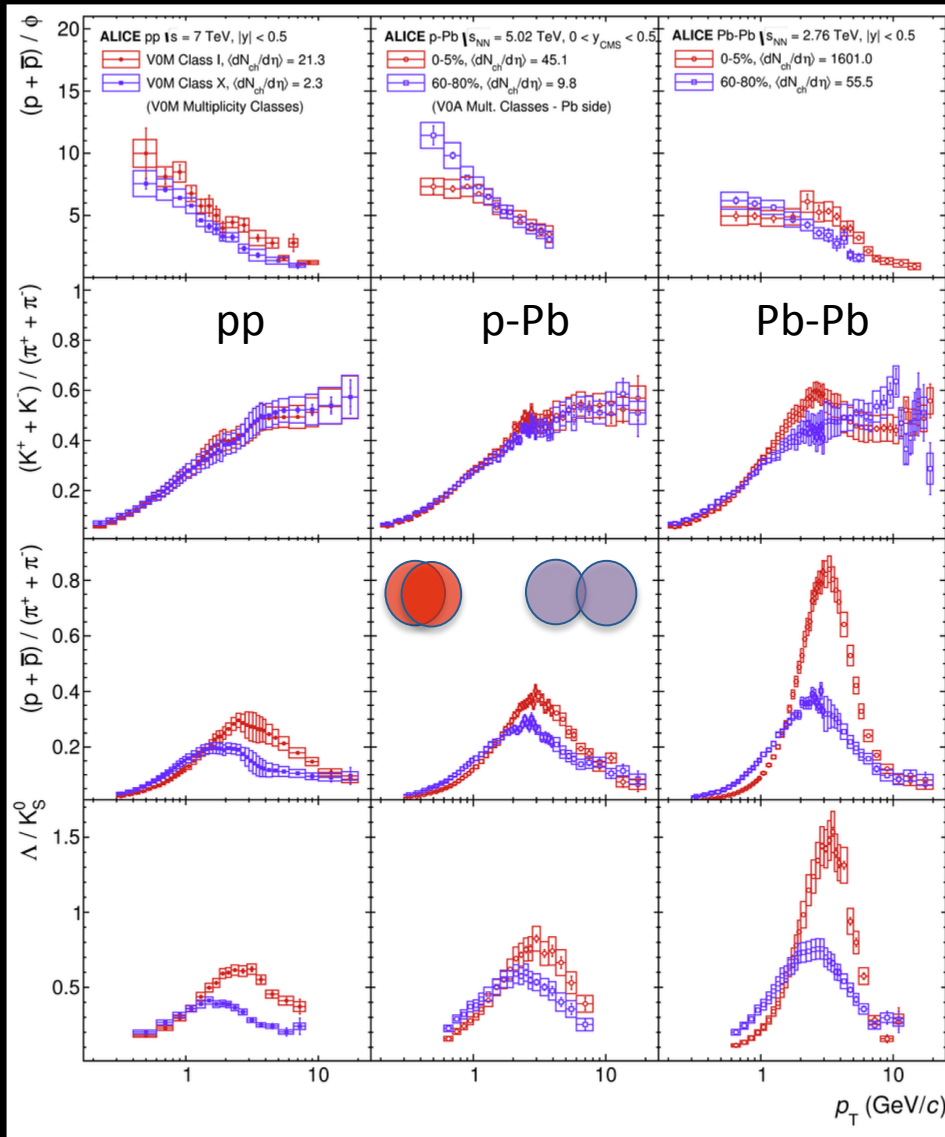
$\langle p_T \rangle$ vs centrality of hadrons in pp, p-Pb, Pb-Pb and Xe-Xe



- $\langle p_T \rangle$ increases with increasing centrality and mass
- Larger increase in smaller systems
- Collective radial expansion

Particle ratios in pp, p-Pb, Pb-Pb and Xe-Xe

Phys. Rev. C99, 024906 (2019)



- Similar evolution pattern of particle ratios from pp to Pb-Pb
- Similar values for Pb-Pb and Xe-Xe at similar multiplicity
- p/ϕ consistent with radial flow but also with (re)combination [Phys.Rev. C92 (2015)054904]
- Particle production is driven by the characteristics of final state
- Spectra at intermediate p_T determined by flow or recombination?

Blast-Wave fit to hadron p_T spectra

$$E \frac{d^3 N}{d p^3} \propto \int_0^R m_T I_0 \left(\frac{p_T \sinh(\rho)}{T_{Kin}} \right) K_1 \left(\frac{m_T \cosh(\rho)}{T_{Kin}} \right) r dr$$

$$m_T = \sqrt{m^2 + p_T^2} \quad \rho = \tanh^{-1}(\beta_T) \quad \beta_T = \beta_s \left(\frac{r}{R} \right)^n$$

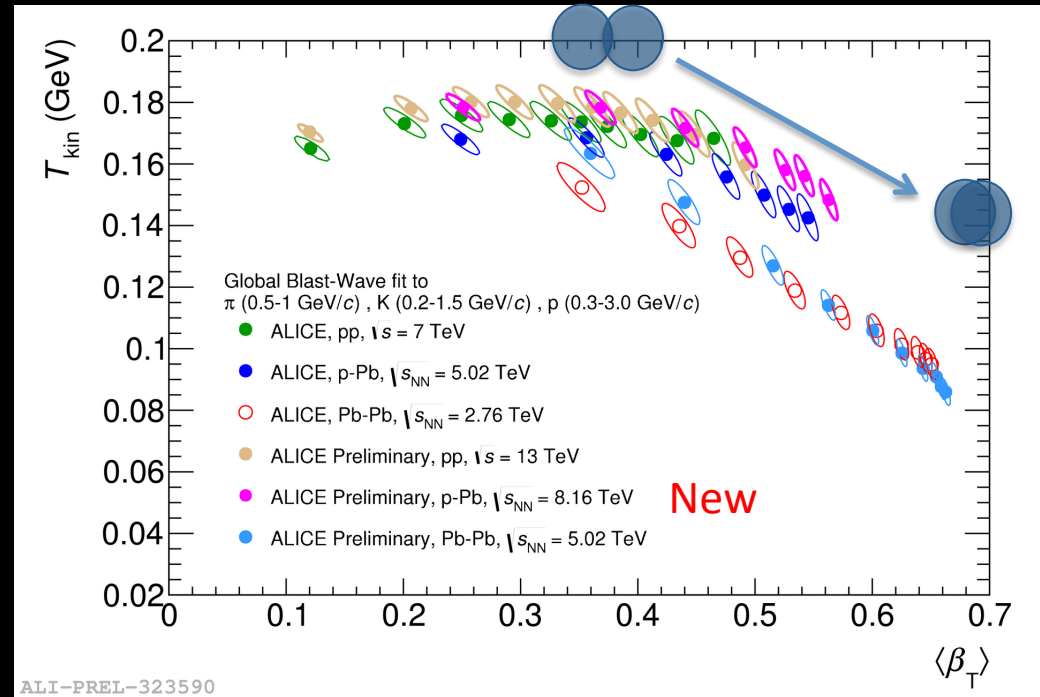
Schnedermann, Sollfrank and Heinz Phys. Rev. C 48, 2462

Simplified hydrodynamic model with 3 parameters:

- β_T - radial expansion velocity
- T_{kin} - kinetic freeze-out temperature
- n - velocity profile

- Similar trend observed in pp, p-Pb and Pb-Pb collisions
- Larger β_T in small systems at similar multiplicity

Simultaneous fit to the π , K, ρ spectra

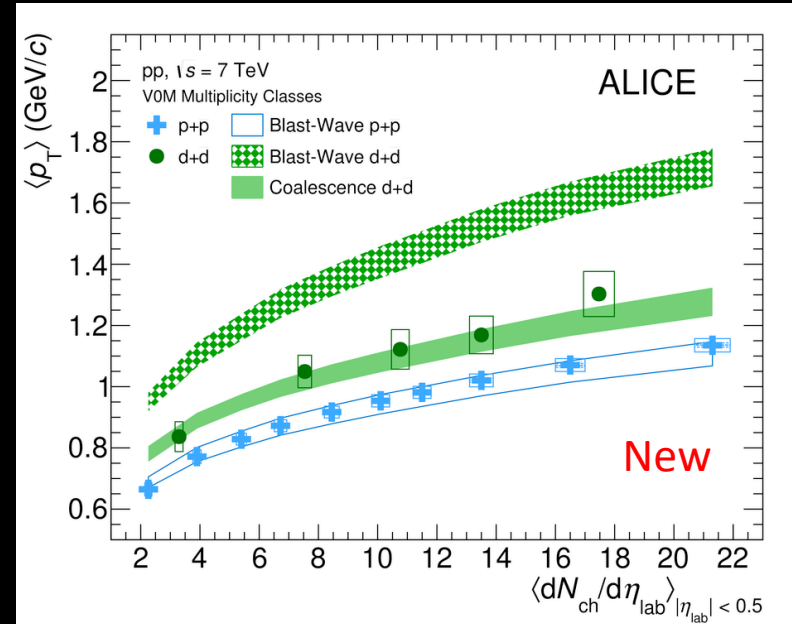
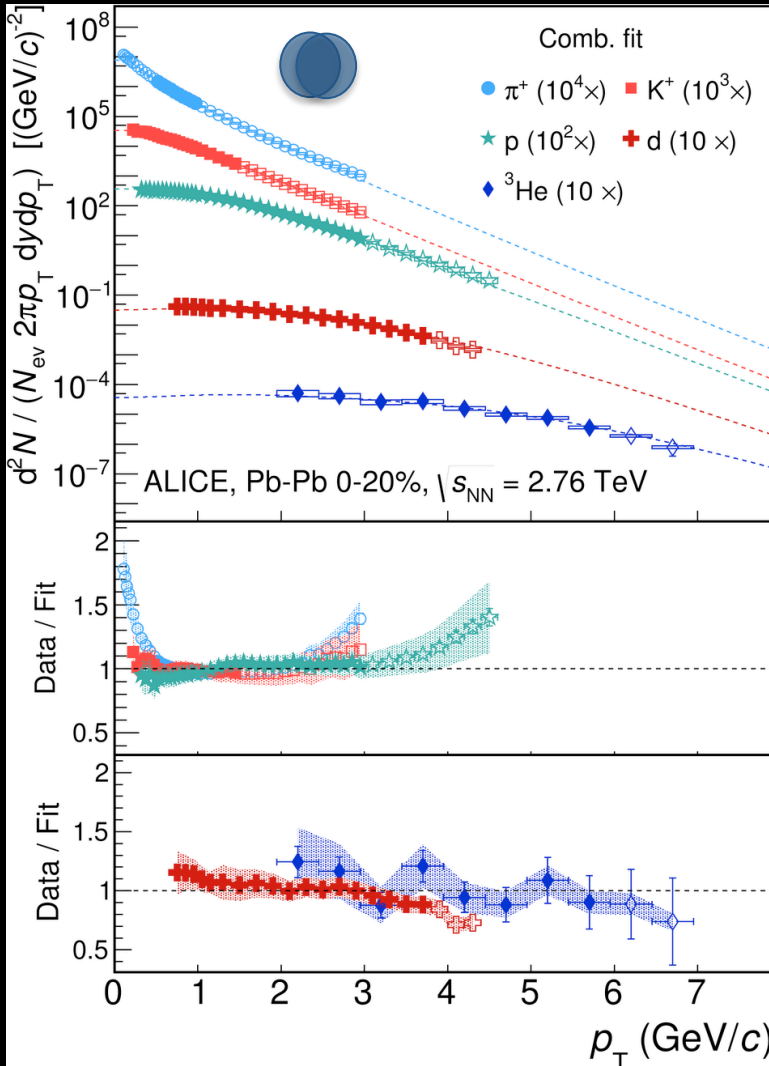


pp 7 TeV: Eur. Phys. J. C75 (2015) 226
 p-Pb 5.02 TeV: Phys. Lett. B 760 (2016) 720
 Pb-Pb 2.76 TeV: Phys. Rev. C88 (2013) 044910

Blast-Wave fit including nuclei

Phys. Rev. C 93 (2015) 024917

Phys.Lett. B794 (2019) 50



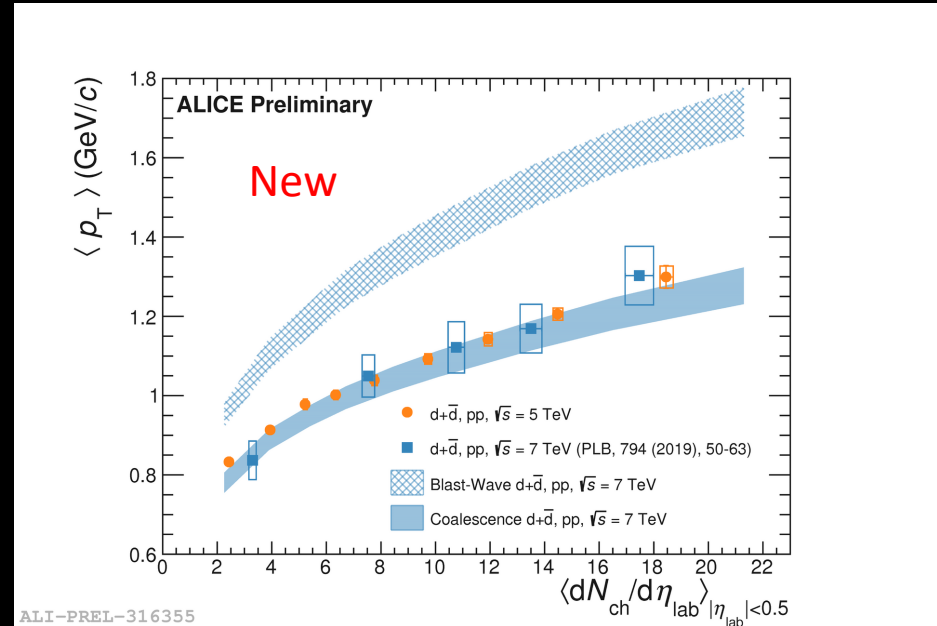
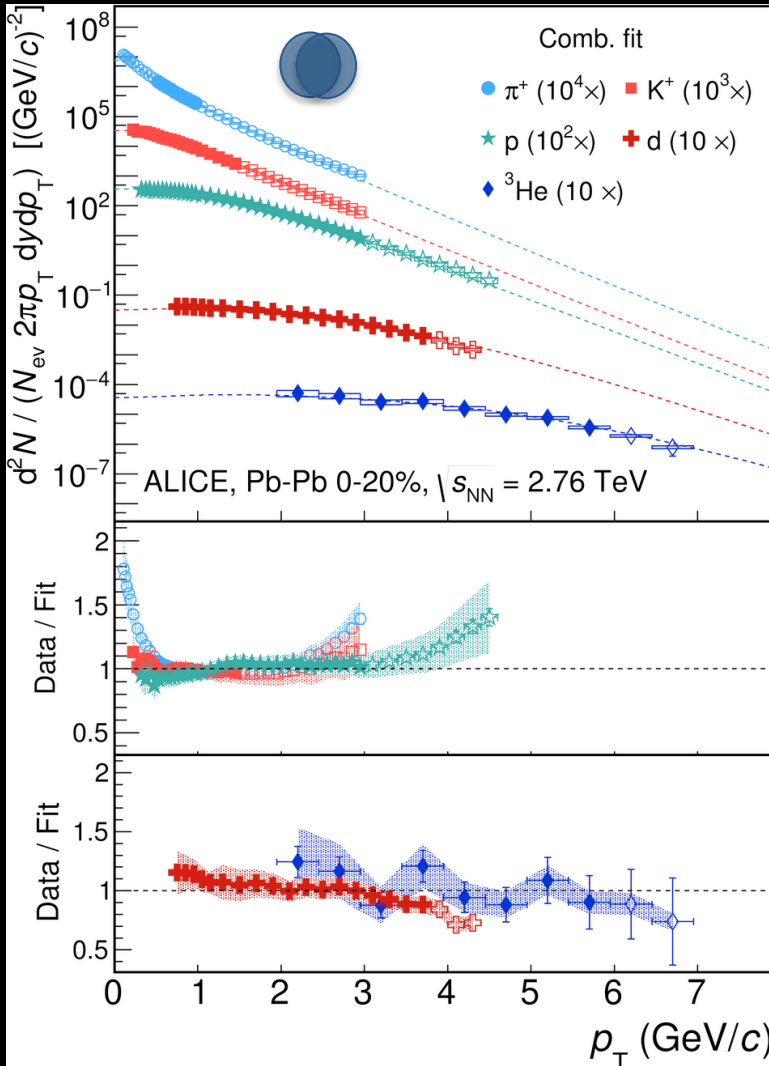
- Pb-Pb: nuclei participate in collective expansion with $\langle \beta_T \rangle \sim 0.63$ and $T_{kin} \sim 113$ MeV (consistent with values for π , K, p)
- pp: Blast-Wave model does not reproduce $\langle p_T \rangle$ of (anti)deuterons in contrast to Pb-Pb
 → coalescence scenario favored for deuteron production in pp

Blast-Wave fit including nuclei

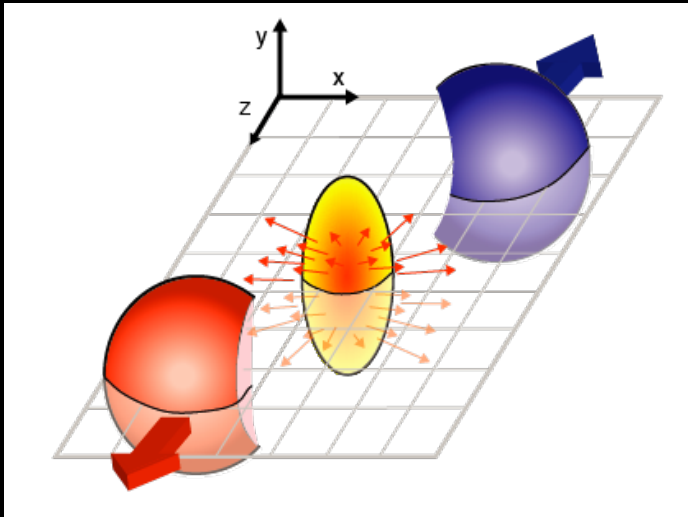


Phys. Rev. C 93 (2015) 024917

A. Balbino - Poster 11/06



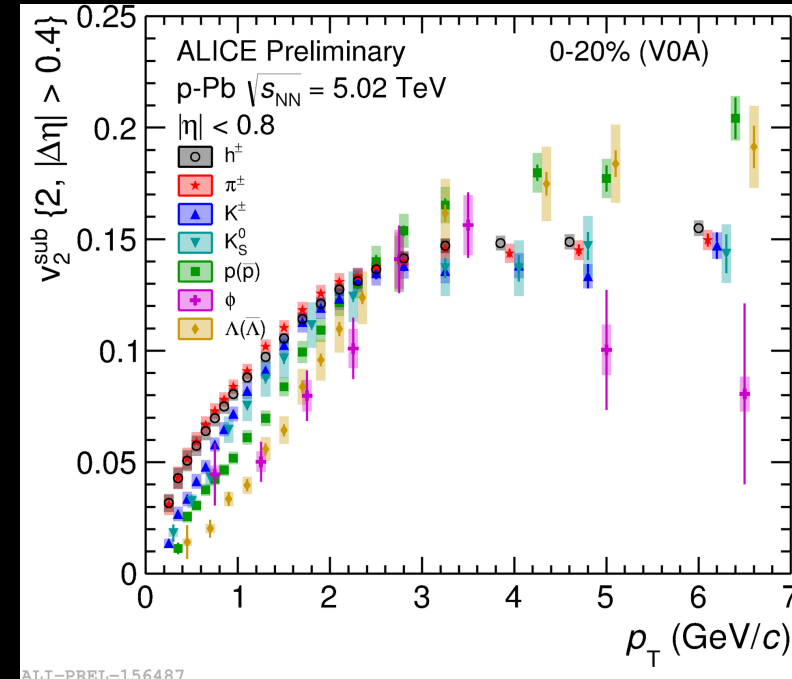
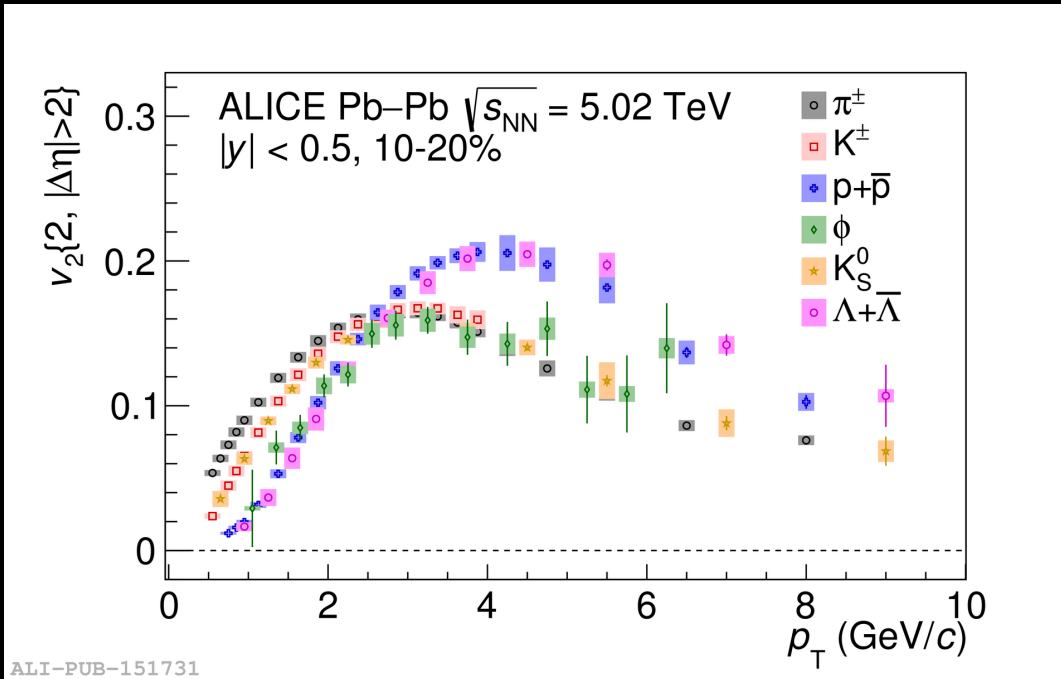
- Pb-Pb: nuclei participate in collective expansion with $\langle \beta_T \rangle \sim 0.63$ and $T_{\text{kin}} \sim 113$ MeV (consistent with values for π , K , p)
 - pp: Blast-Wave model does not reproduce $\langle p_T \rangle$ of (anti)deuterons in contrast to Pb-Pb
- coalescence scenario favored for deuteron production in pp



ANISOTROPIC FLOW

Identified particle v_2 in Pb-Pb and p-Pb

JHEP09 (2018) 006



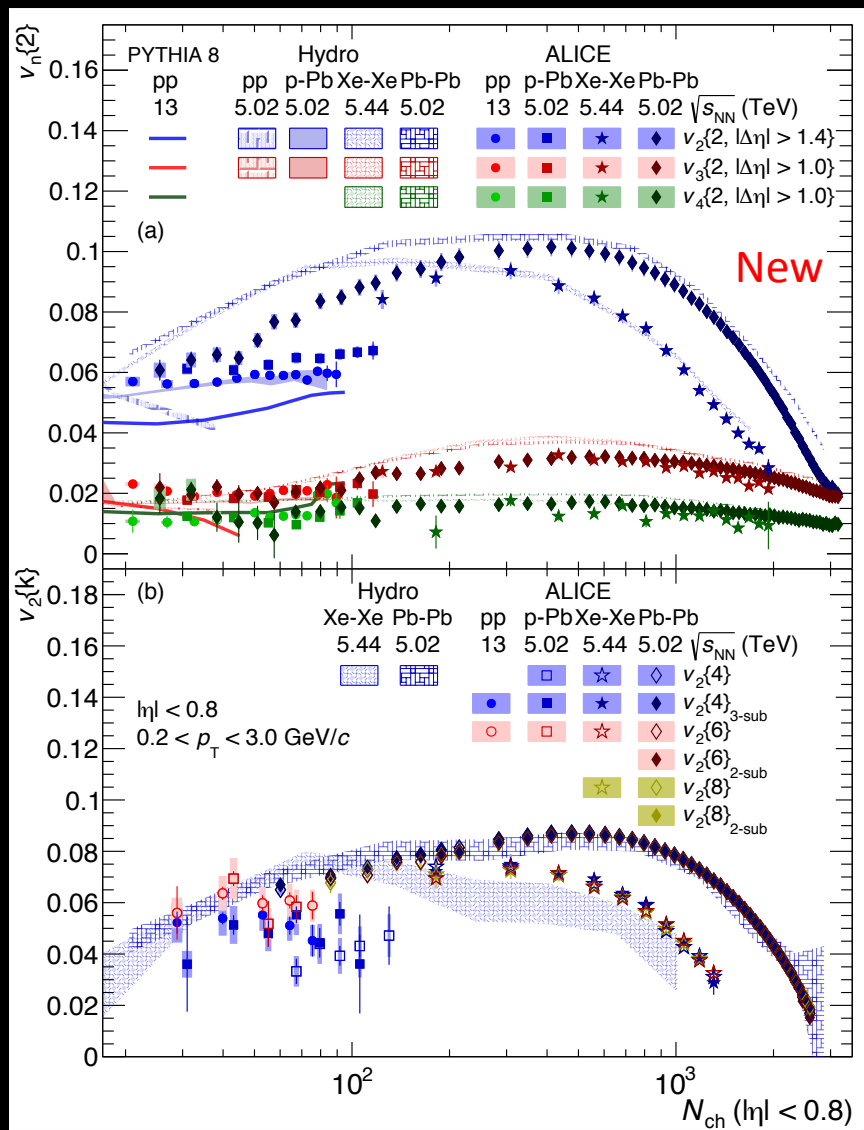
- $p_T < 2$ GeV/c: hadron mass ordering seen in Pb-Pb and p-Pb collisions
- $p_T \sim 2.5$ GeV/c: crossing between v_2 of baryons and mesons
- $p_T > 2.5$ GeV/c: baryons $v_2 >$ mesons v_2 (flow driven by quark content)
- ϕ meson follows mass ordering at low p_T and quark content at intermediate p_T

→ Quantitative differences between systems

Charged-particle v_2 , v_3 , v_4 in pp, p-Pb, Xe-Xe and Pb-Pb

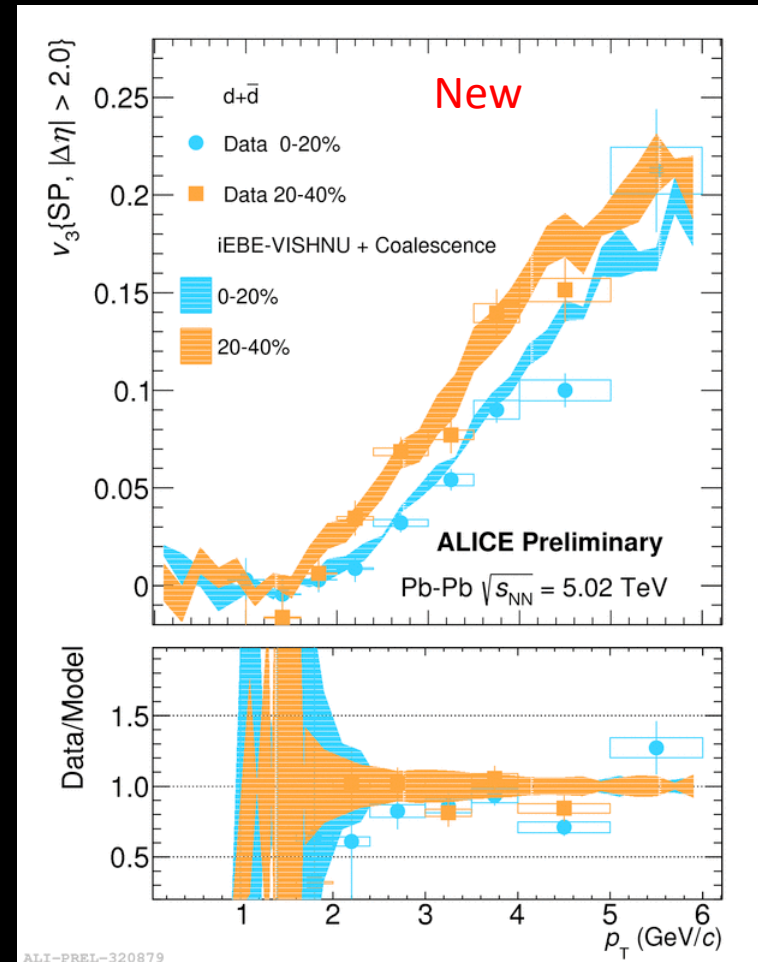
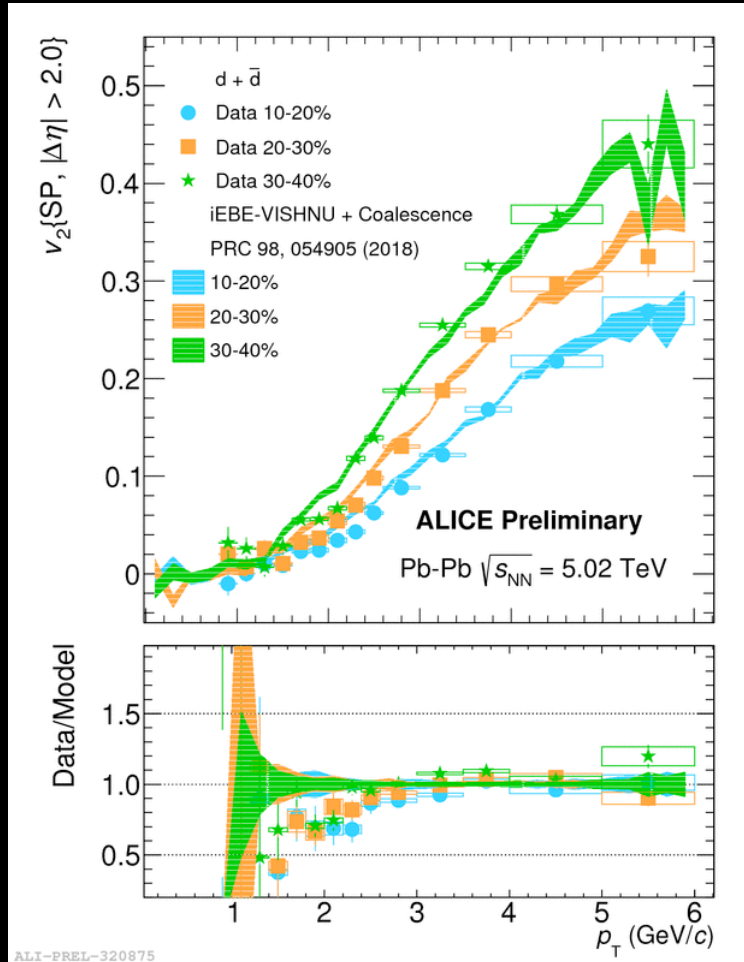


arXiv:1903.01790



- Two and multi-particle azimuthal correlations
- Similar ordering of $v_2 > v_3 > v_4$ seen in small and large systems
- Weak v_2 multiplicity dependence seen in pp and p-Pb
- Novel subevent method supports long-range multi-particle azimuthal correlations in high multiplicity pp and p-Pb
- PYTHIA8 and IP-Glasma+MUSIC+UrQMD cannot describe correlations in pp and p-Pb

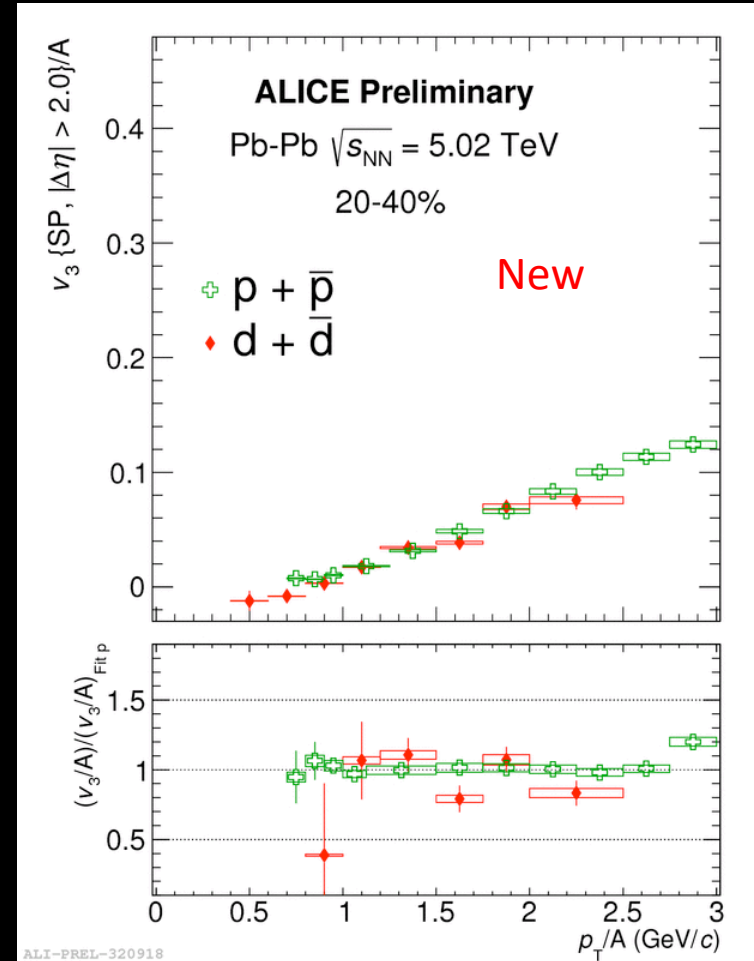
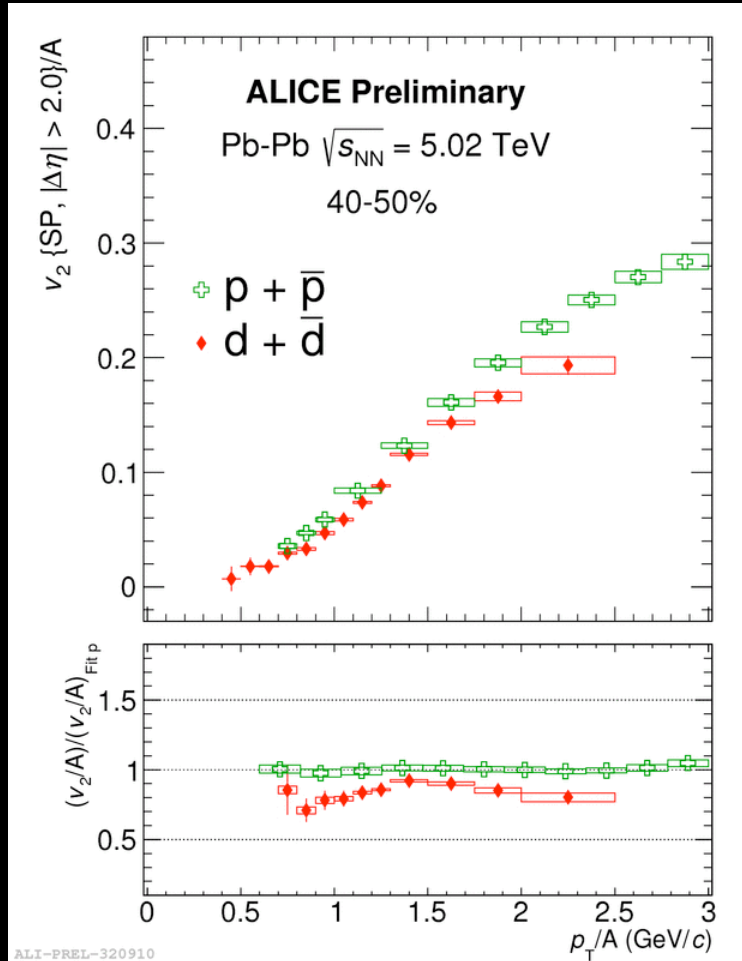
(Anti)deuteron v_2 and v_3 in Pb-Pb



- Coalescence model with phase-space distributions of protons and neutrons from iEBE-VISHNU in agreement with data

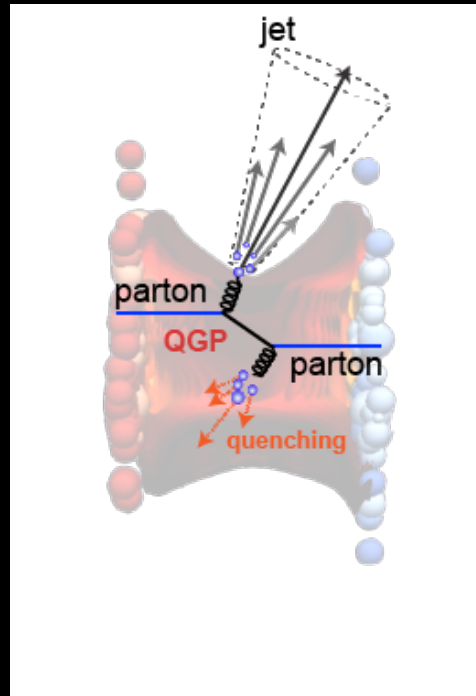
A. Caliva 11/06

(Anti)deuteron v_2 and v_3 in Pb-Pb



- Coalescence model with phase-space distributions of protons and neutrons from iEBE-VISHNU in agreement with data
- (Anti)deuteron v_2 and v_3 comparable to protons

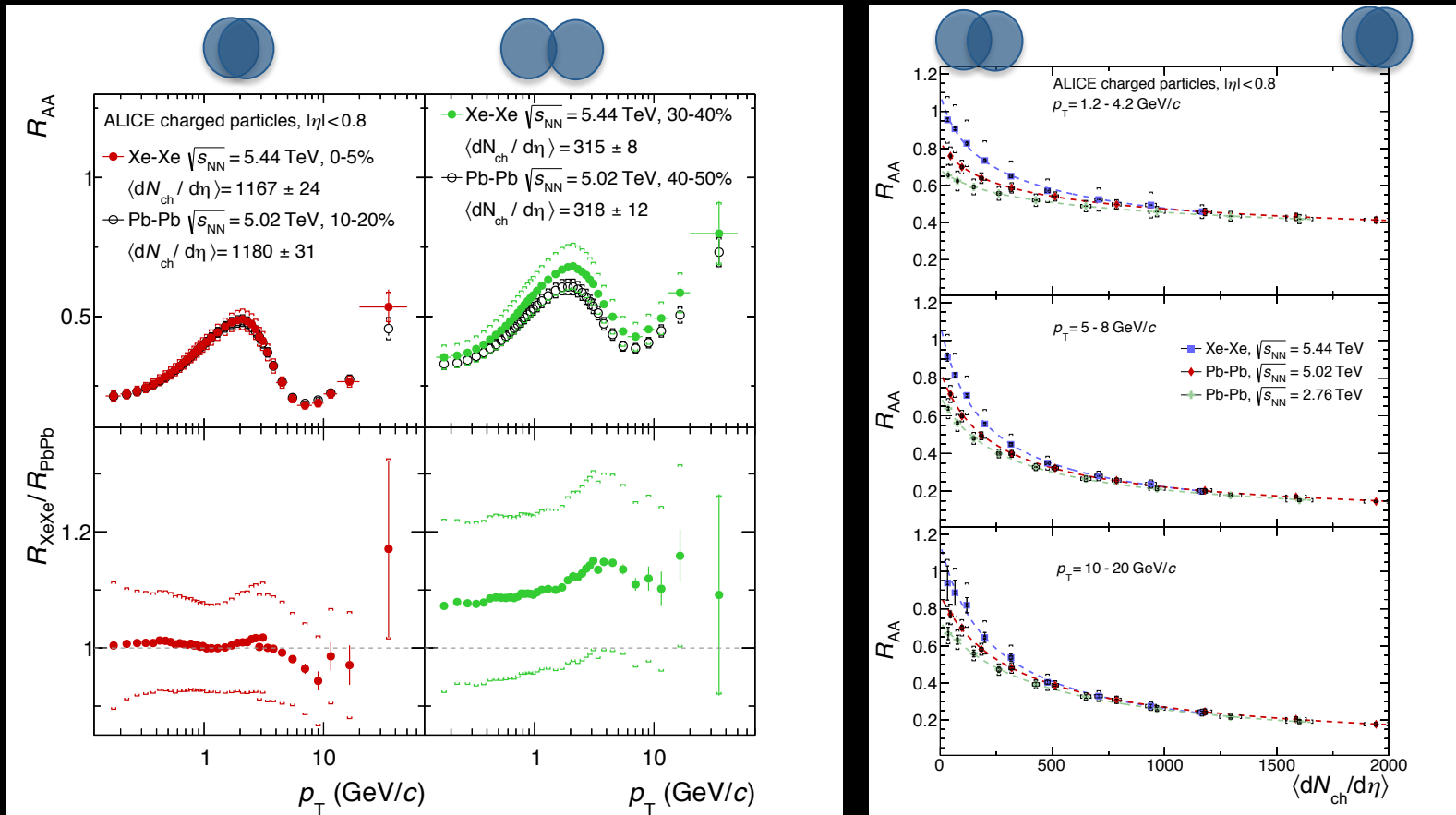
A. Caliva 11/06



NUCLEAR MODIFICATION FACTORS AND JET QUENCHING

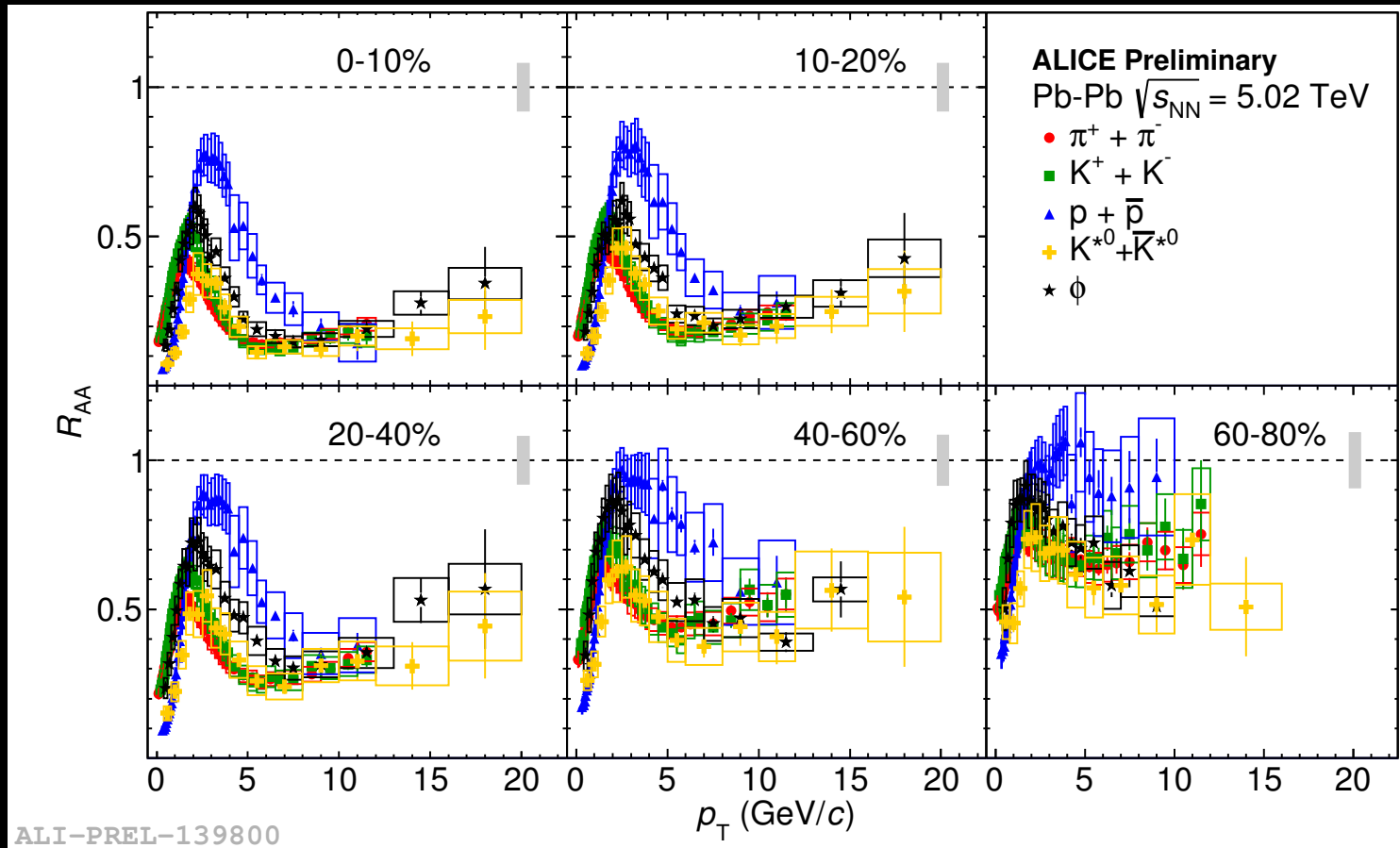
Charged-particle R_{AA} in Xe-Xe and Pb-Pb

Phys. Lett. B 788 (2019) 166



- Similar R_{AA} in central Xe-Xe and Pb-Pb collisions at similar multiplicity
- Different R_{AA} in more peripheral collisions
- Result of interplay between geometry and path length dependence of parton energy loss

R_{AA} of identified hadrons in Pb-Pb at $\sqrt{s_{NN}} = 5.02$ TeV



- $p_T > 10$ GeV/c: the same strong suppression of light-flavor hadrons
 \rightarrow high- p_T fragmentation function not affected by the medium
- $p_T < 10$ GeV/c: differences between light-flavor hadrons (difficult to interpret due to other effects e.g. radial flow, (re)combination,...)

Summary

- The collision geometry plays an important role in particle production
- Particle production is driven by characteristics of final state
- Several mechanisms of hadron and nuclei production: rescattering, (re)combination, coalescence etc.
- Strangeness enhancement seen in small and large systems (droplet of QGP?)
- NEW results: charged-particle p_T spectra vs. R_T , π, K, p in p-Pb at 8.16 TeV, nuclei production in pp and p-Pb
- Collective effects are observed in small and large systems
- Quantitative differences between flow coefficients in small and large systems (different underlying mechanisms?)
- NEW results: charged-particle v_2, v_3 and v_4 in pp, p-Pb, Xe-Xe and Pb-Pb, nuclei v_2 and v_3 in Pb-Pb
- Strong suppression of particle production in central Pb-Pb and Xe-Xe collisions (interplay between geometry and path length dependence)
- Suppression at high p_T independent of particle species

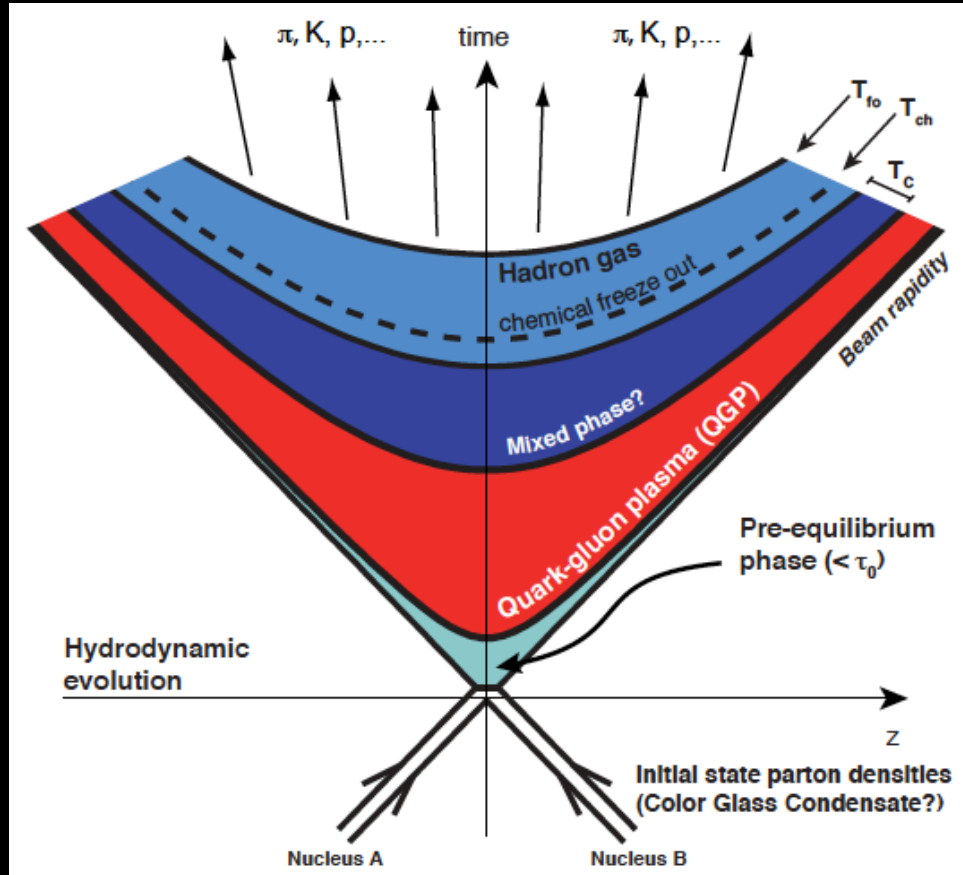
Backup



Stages of the relativistic heavy-ion collision



Bjorken 1983



Chemical and thermal freeze-out

Hadronization ($T \sim T_c$)

Quark-Gluon Plasma
(thermalized matter?)

Pre-equilibrium, fast
thermalization ~ 1 fm/c, glasma?

Hard collisions

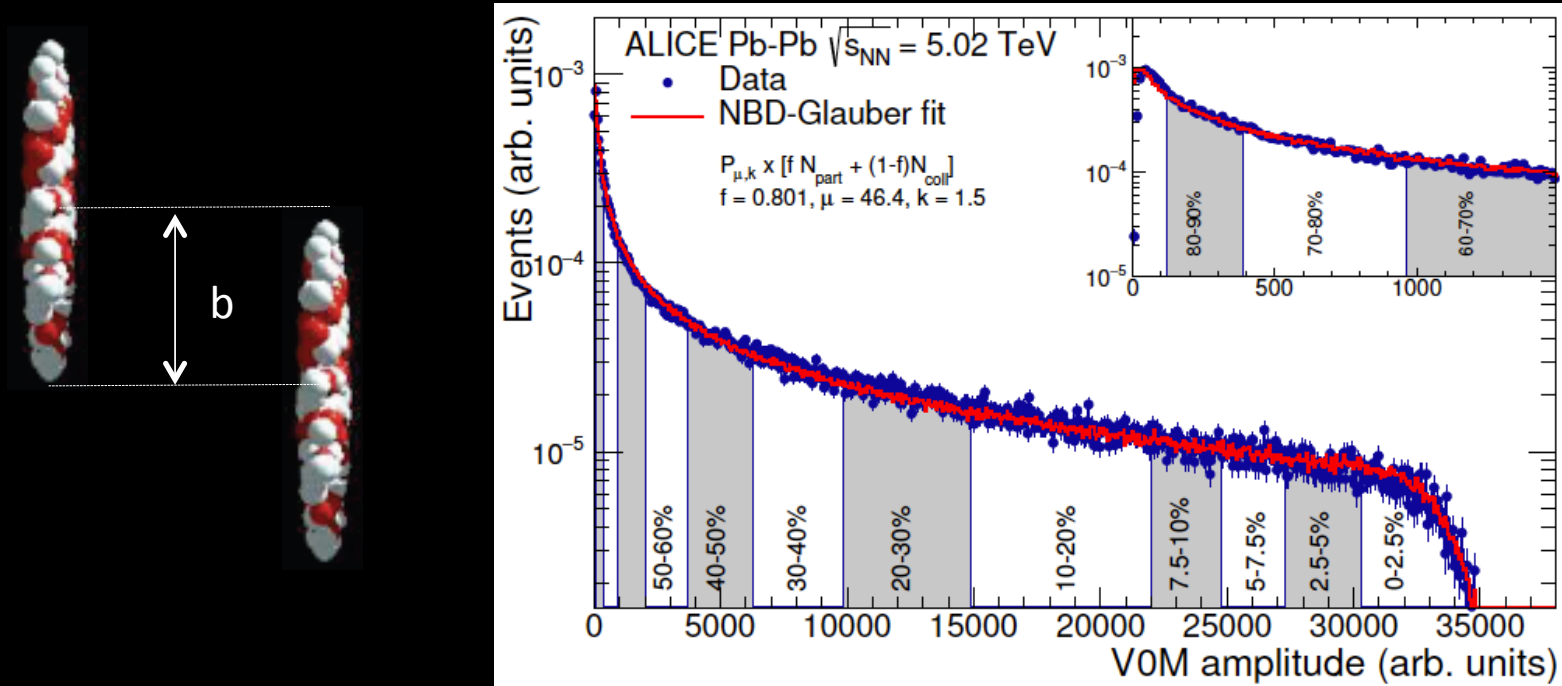
Lorentz-contracted ions
(dense gluonic matter, Color
Glass Condensate?)



- All stages in models of nuclear collisions
- QGP expansion modeled by relativistic hydrodynamics

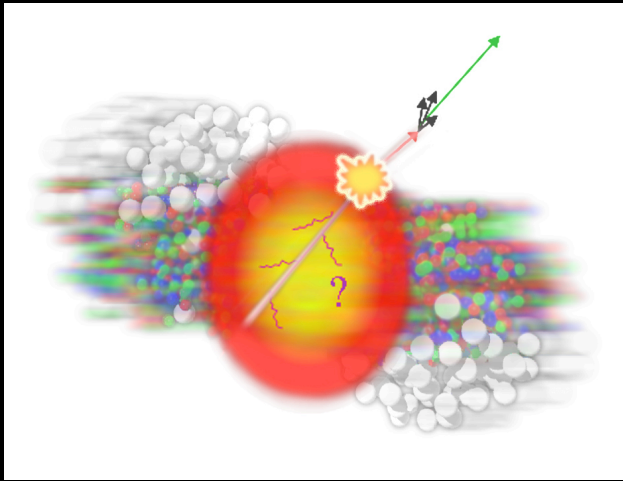
Event Centrality Selection

ALICE-PUBLIC-2018-011



- Correlate particle multiplicity with collision geometry i.e. impact parameter, volume and shape (A. Białas et al. APPB 8 (1977) 389)
- N_{coll} , N_{part} and $T_{AA} = N_{coll} / \sigma_{INEL}^{NN}$ values determined by fitting NBD-Glauber coupled to two parameter model

Parton energy loss and jet quenching



Radiative and collisional parton energy loss:

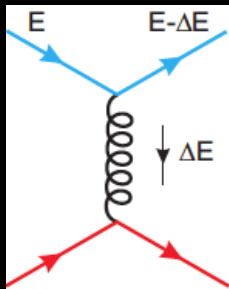
$$\Delta E = \Delta E_{\text{coll}} + \Delta E_{\text{rad}}, \quad \Delta E (E, m, C_R; \rho_g, \alpha_s, T, L)$$

D. d'Enterria, arXiv:0902.2011

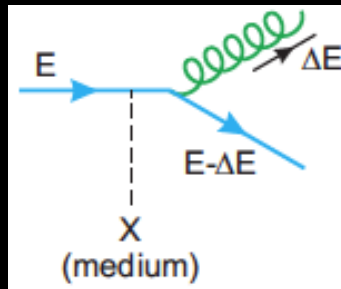
Radiative energy loss dominates at high- p_T :

- Color charge dependence C_R : $C_{R,g} > C_{R,q,Q}$
 $\rightarrow \Delta E_g > \Delta E_{q,Q}$
- Mass dependence “dead cone”: gluon radiation suppression at $\Theta < m_Q/E$
 $\rightarrow \Delta E_q > \Delta E_Q$

L. Dokshitzer & D.E. Kharzeev, PLB 519 (2001) 199



collisional

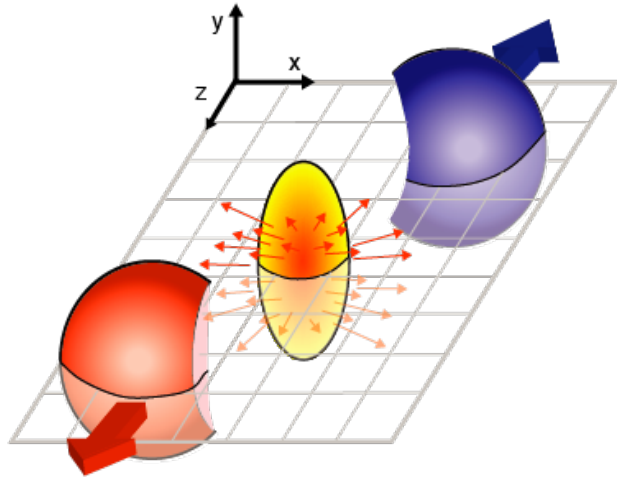


radiative

- In static medium: $\Delta E_{\text{coll}} \sim L, \Delta E_{\text{rad}} \sim L^2$
- Characterize medium transport properties via parton energy loss

$$\hat{q} \equiv \frac{m_D^2}{\lambda} = m_D^2 \rho \sigma$$

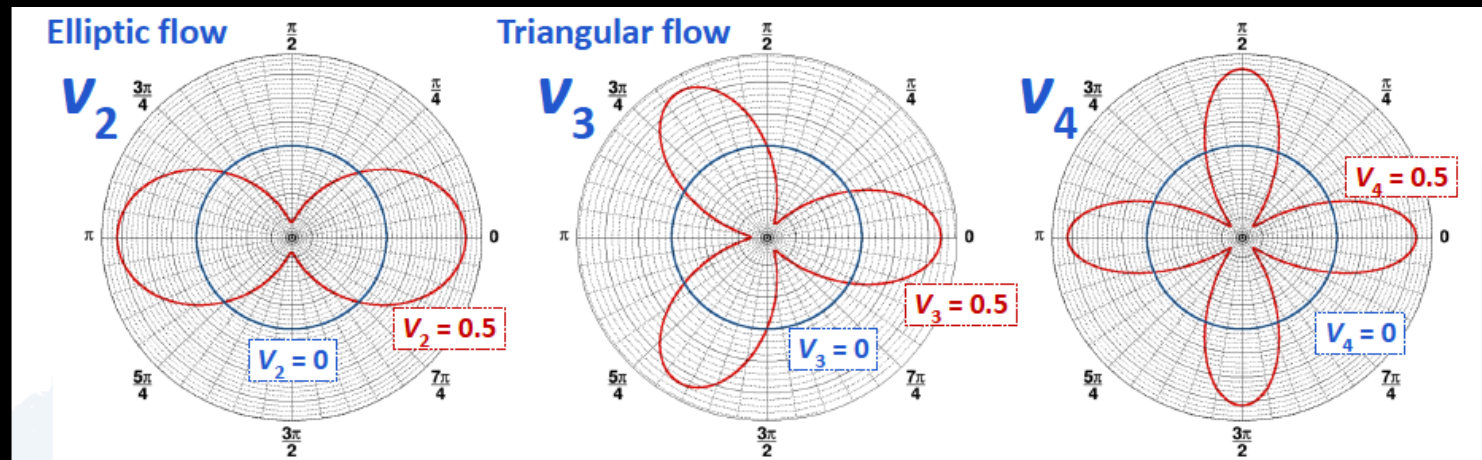
Anisotropic flow



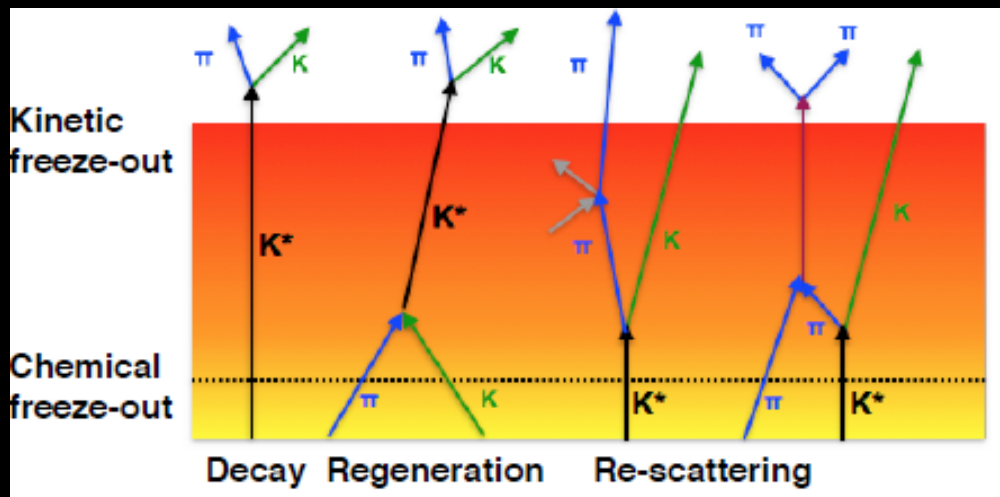
- Strongly interacting system:
spatial anisotropy \rightarrow momentum anisotropy
- Quantified in terms of Fourier coefficients v_n

$$E \frac{d^3 N}{d^3 p} = \frac{1}{2\pi p_T dp_T dy} \left(1 + 2 \sum_{n=1}^{\infty} v_n \cos[(\varphi - \Psi_n)] \right)$$

$$v_n(p_T, y) = \langle \cos[n(\varphi - \Psi_n)] \rangle$$



Resonance production



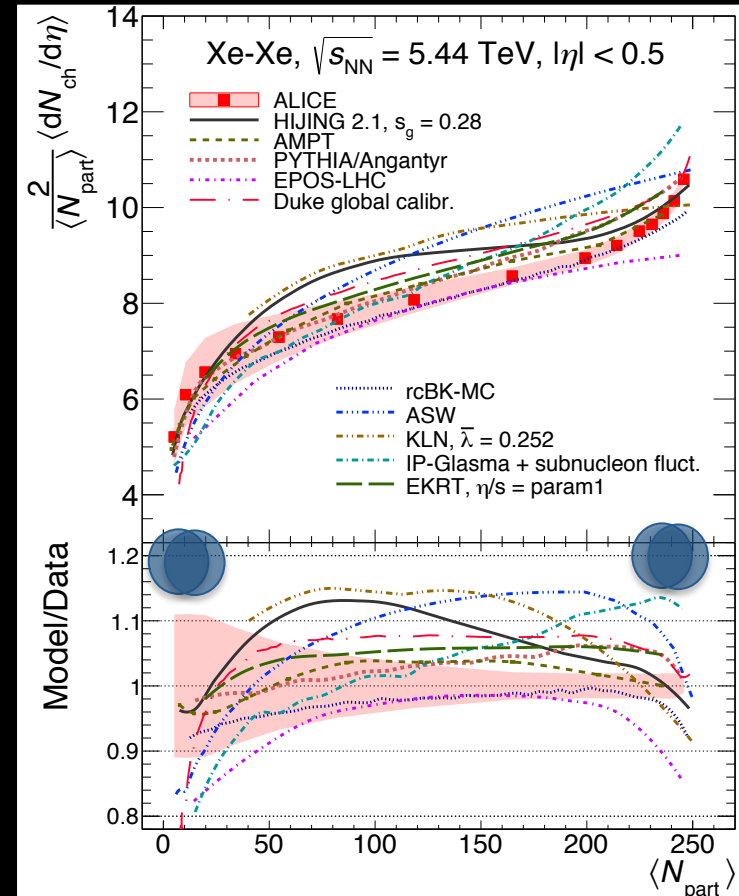
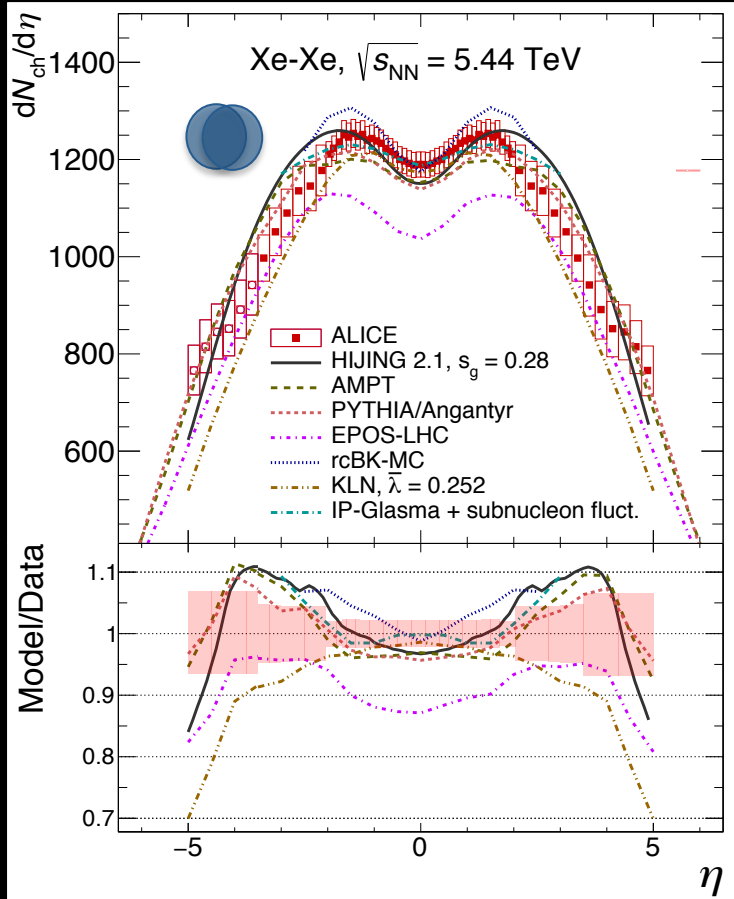
Rescattering and (re)combination in the hadronic phase influence the measured resonance yields

Final yields at kinetic freeze-out depend on:

- Initial yield after chemical freezeout
- Lifetime of hadronic phase
- Resonance lifetime
- Scattering cross-section of decay products

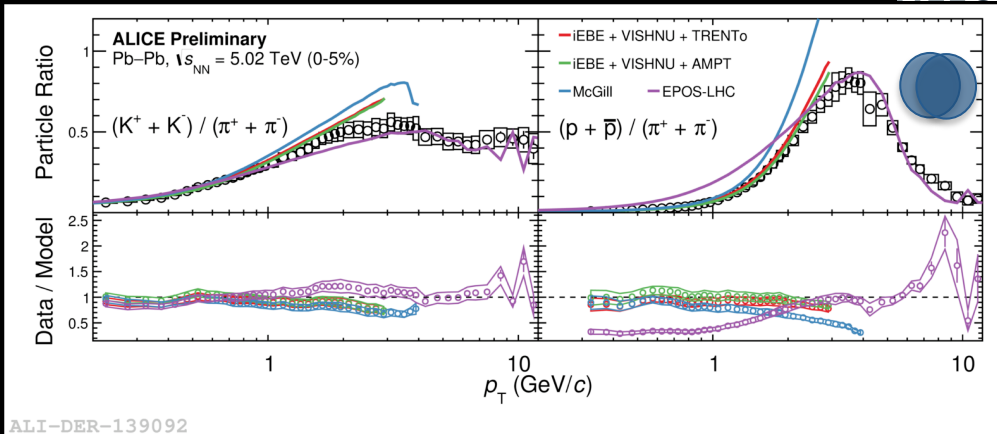
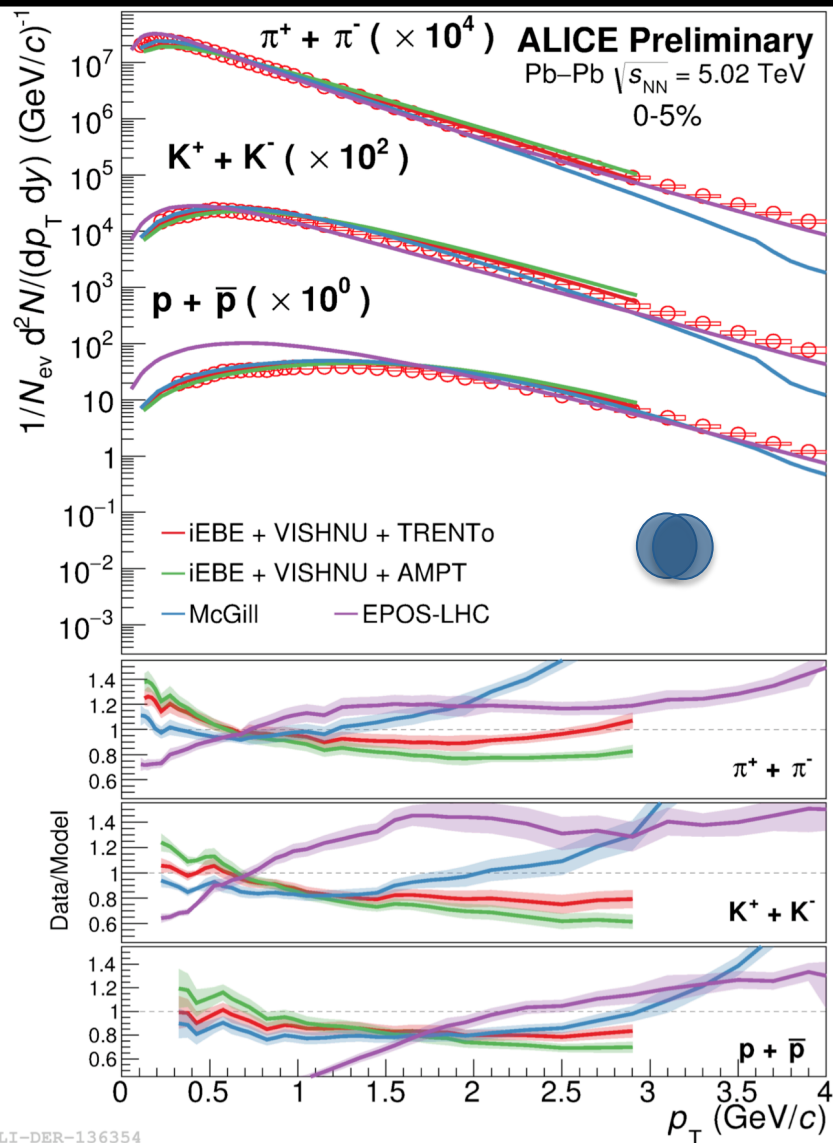
Charged-particle multiplicity vs models in Xe-Xe

Phys. Lett. B 790 (2019) 35



- Models do not describe charged-particle production in the whole rapidity range
- N_{part} dependence is best described by rcBK-MC: CGC saturation model based on Balitsky-Kovchegov gluon evolution equation

Comparison to hydro models – Pb-Pb (0-5%)

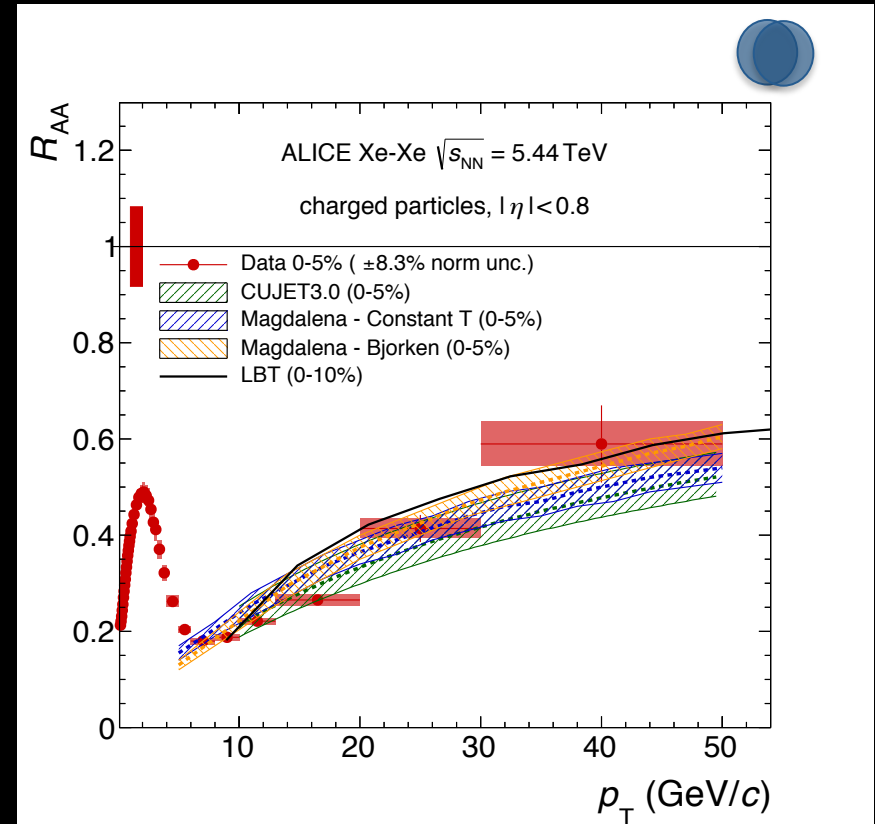
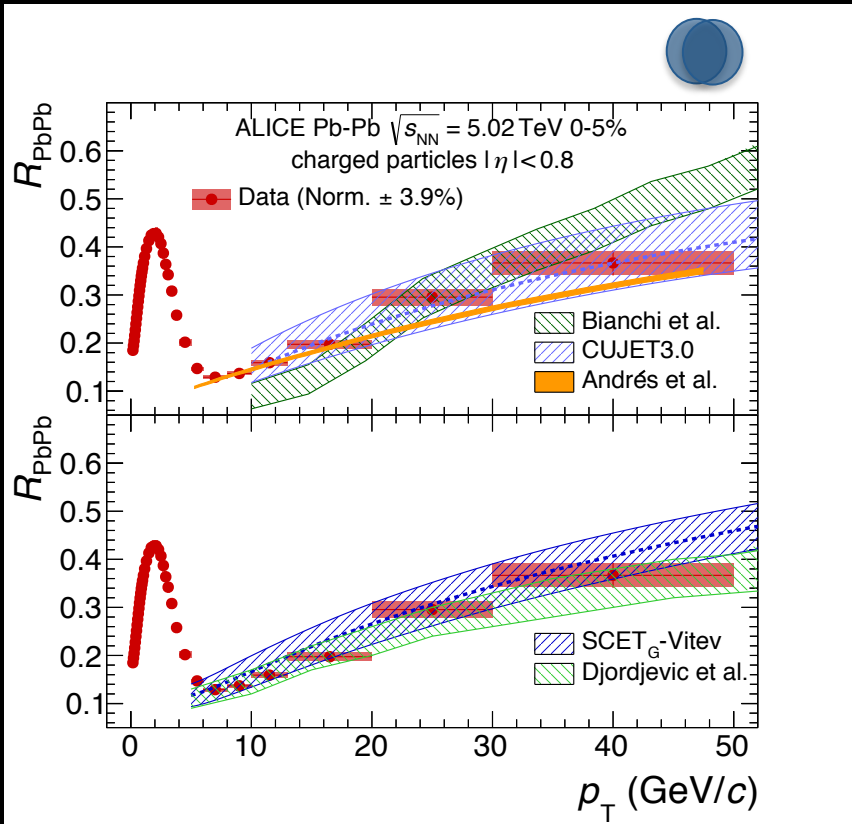


- iEBE-VISHNU (hydro+UrQMD) + TRENTo/AMPT initial conditions
 - McGill: MUSIC (hydro) + IP glasma initial conditions
 - EPOS LHC: hydro param. + hadronization [3]
 - Models description within 20-40% for Pb-Pb central collisions
 - Presented models give much worse description in peripheral collisions
- radial flow (hydro) - Navier-Stokes equations
 iEBE-VISHNU: Eur. Phys. J C77 (2017) 645
 McGill: Phys. Rev. C95 (2017) 064913
 EPOS LHC: Phys. Rev. C92 (2015) 034906

Charged-particle R_{AA} vs models in Pb-Pb and Xe-Xe

JHEP 1811 (2018) 013

Phys. Lett. B 788 (2019) 166

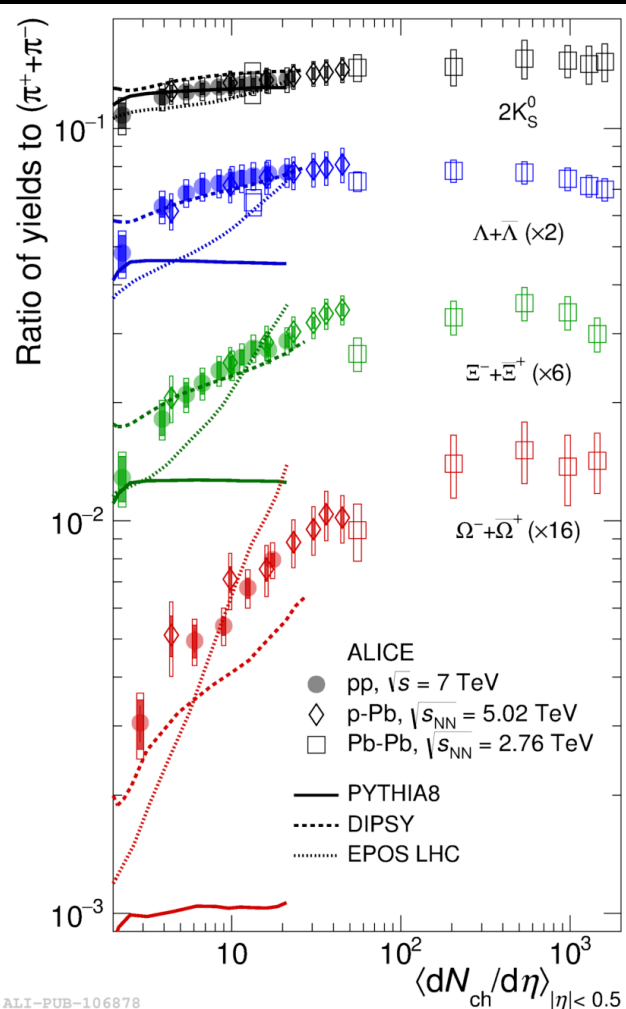


Models:

- All models include radiative energy loss
- CUJET3.0 and Magdalena Djordjevic models also include elastic energy loss
- Calculations are performed in dynamically expanding medium except that of Vitev et al.

Relative strangeness production in pp, p-Pb and Pb-Pb

ALICE, Nature Physics 13 (2017) 535



Historically a signature of the QGP
Rafelski & Mueller, PRL 48 (1982) 1066

- Smooth evolution from pp to Pb-Pb
- Enhancement observed in A-A which increases with strangeness content
- Enhancement also seen in the smaller systems pp and p-Pb
- No significant energy and system dependence is observed at similar multiplicity
- DIPSY model (rope hadronization) describes data best
 - But DIPSY overestimates p/π (not shown)

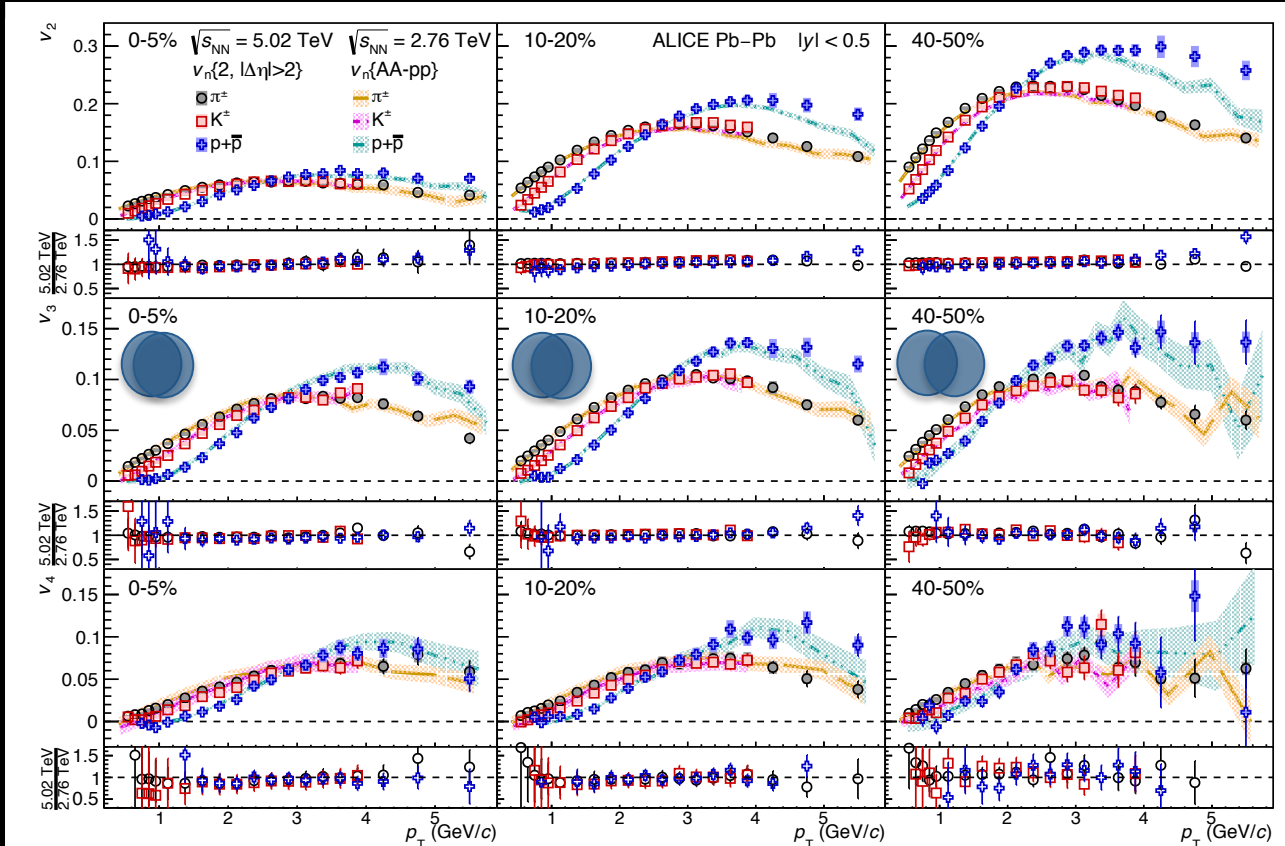
PYTHIA 8: Comput. Phys. Commun. 178 (2008) 852867

EPOS LHC: PRC92 (2015) 034906

DIPSY: PRD92 (2015) 094010

Identified particle v_2 , v_3 , v_4 in Pb-Pb

JHEP09 (2018) 006



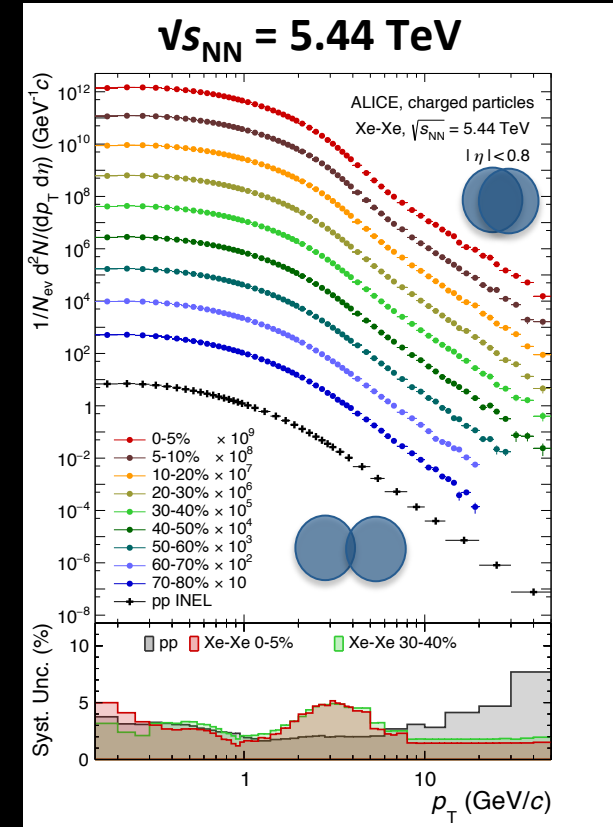
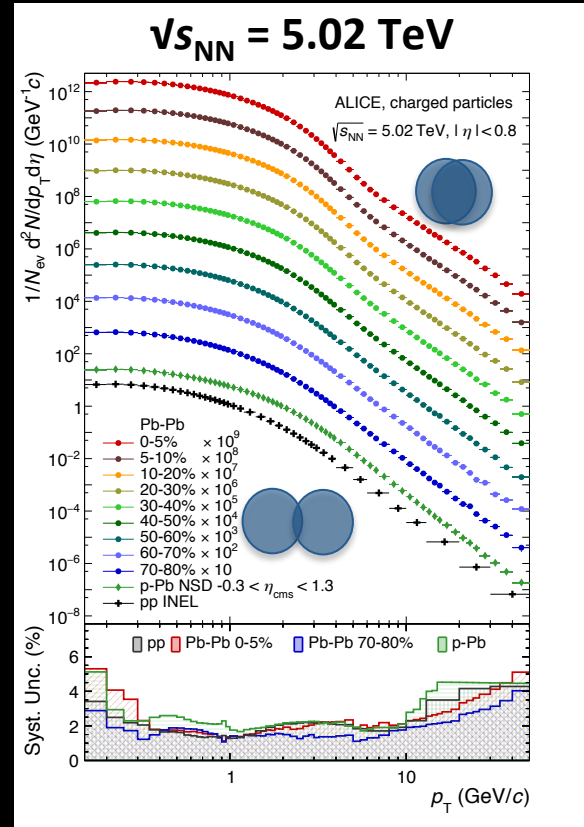
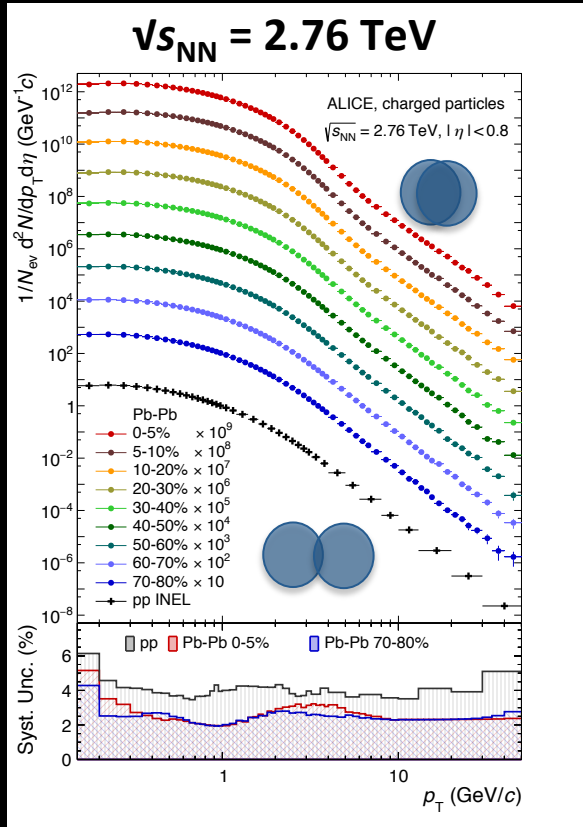
- Weak energy dependence of flow coefficients
- Hadron mass ordering at low p_T
- Strong centrality dependence of v_2
- v_3 is driven by initial state fluctuations

Charged-particle p_T spectra in pp, p-Pb, Pb-Pb and Xe-Xe



JHEP 1811 (2018) 013

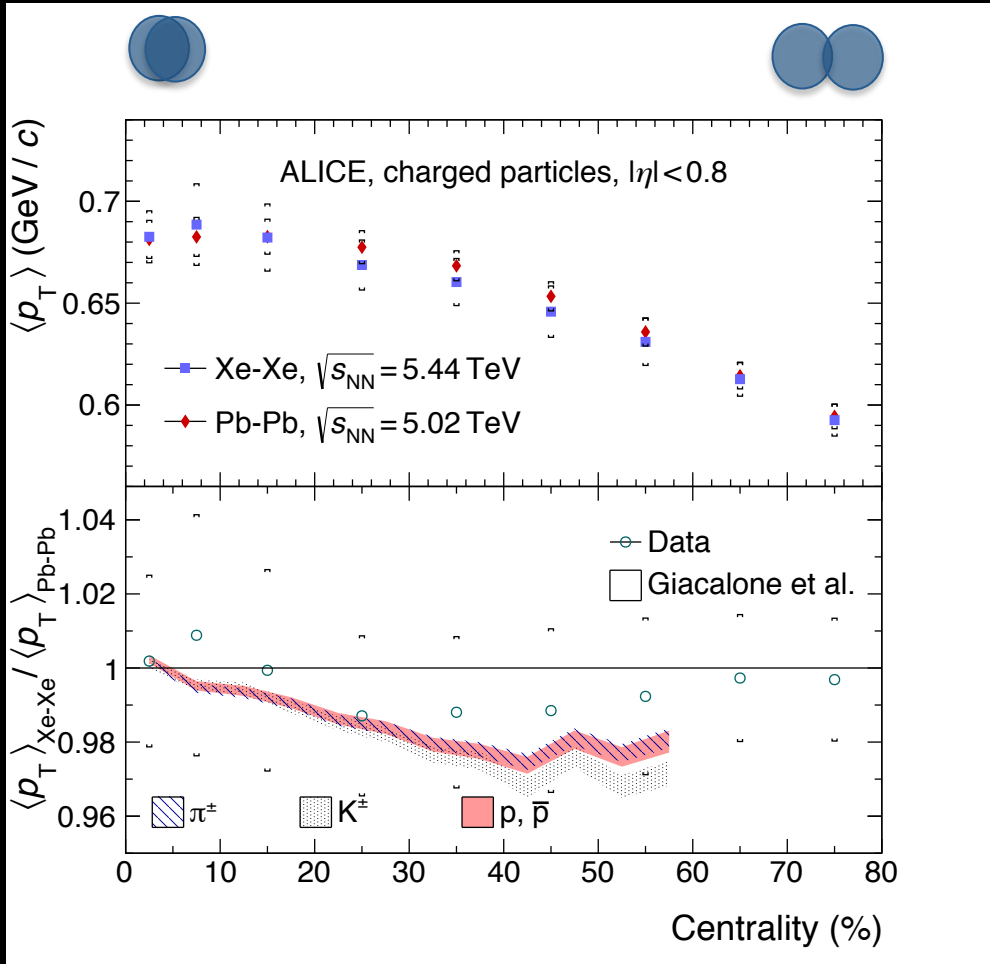
Phys. Lett. B 788 (2019) 166



- p_T spectra in Pb-Pb and Xe-Xe measured in nine centrality intervals
- p_T reference spectra measured in pp and p-Pb collisions
- All spectra obtained using updated corrections (MC tuned on data)

Mean p_T spectra in Xe-Xe and Pb-Pb

Phys. Lett. B 788 (2019) 166



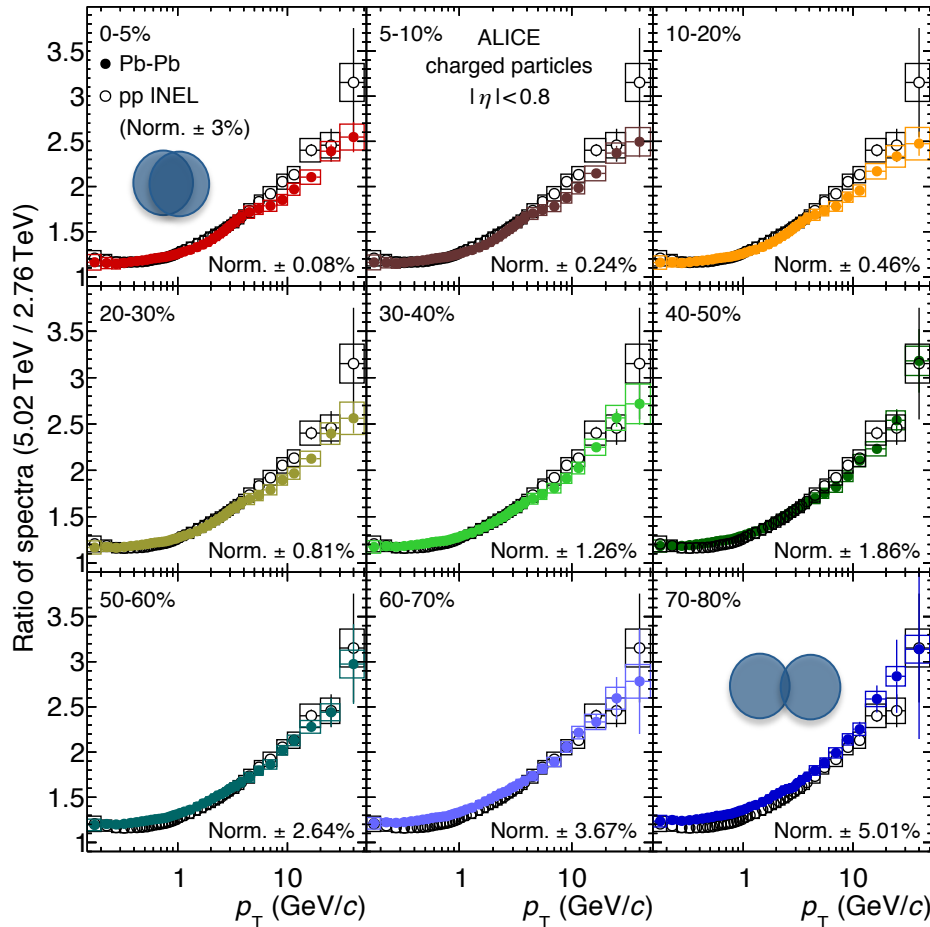
Testing system size (A) dependence

- Similar $\langle p_T \rangle$ as function of centrality in Xe-Xe and Pb-Pb collisions
- $\langle p_T \rangle$ increases with centrality due to radial flow
- Predictions by Giacalone et al. [Phys. Rev. C 97, 034904 (2018)] describes trend in the data
 - Event-by-event simulations: T_R ENTo initial condition + viscous hydro

→ Strong constraints on the hydrodynamic evolution of the system

Ratios of spectra $\nu_{s_{NN}} = 5.02 / 2.76$ TeV

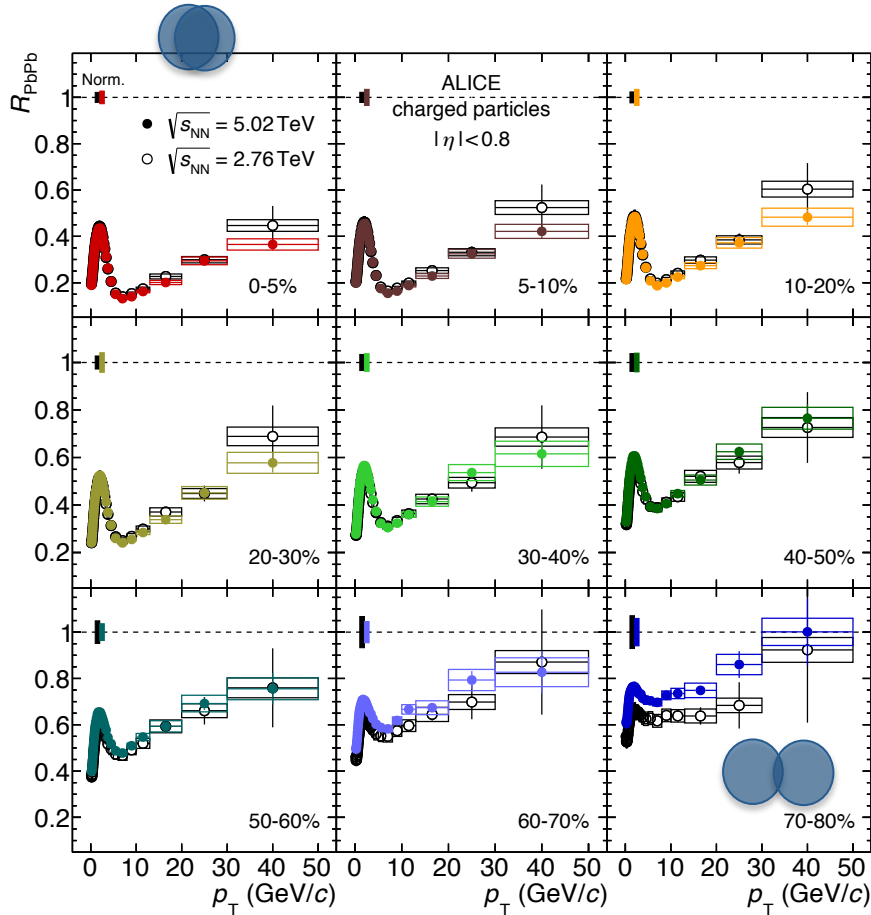
JHEP 1811 (2018) 013



- Spectra get significantly harder with collision energy
- Similar increase with energy in pp and peripheral Pb-Pb collisions
- Gradual reduction of the ratio towards central Pb-Pb collisions

Charged-particle R_{AA} at $\sqrt{s_{NN}} = 2.76$ and 5.02 TeV

JHEP 1811 (2018) 013



$$R_{AA} = \frac{1}{\langle T_{AA} \rangle} \frac{dN_{AA} / dp_T}{d\sigma_{pp} / dp_T} \equiv \frac{[medium]}{[vacuum]}$$

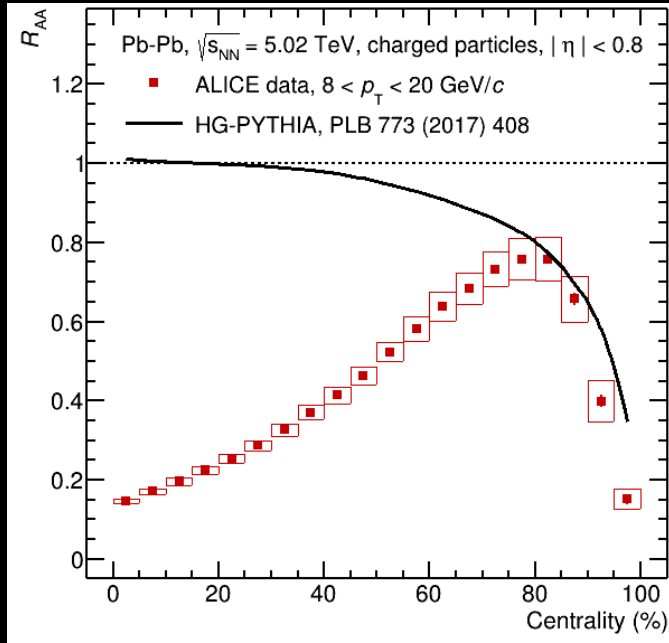
Nuclear overlap function $\langle T_{AA} \rangle$
from Glauber MC

- Different suppression pattern depending on Pb-Pb collision centrality
- Maximum suppression by a factor ~ 7 ($6 < p_T < 7$ GeV/c) in 0-5% collisions
- No significant evolution with collision energy

→ Indication of larger parton energy loss at $\sqrt{s_{NN}} = 5.02$ TeV

Suppression in peripheral Pb-Pb collisions?

ALICE, arXiv:1805.05212



- R_{AA} average over $8 < p_T < 20$ GeV/c
- R_{AA} never reach unity
- HG-Pythia contains no nuclear effects
 - no need for jet quenching
 - centrality selection is biased by fluctuations in particle production towards smaller #MPIs

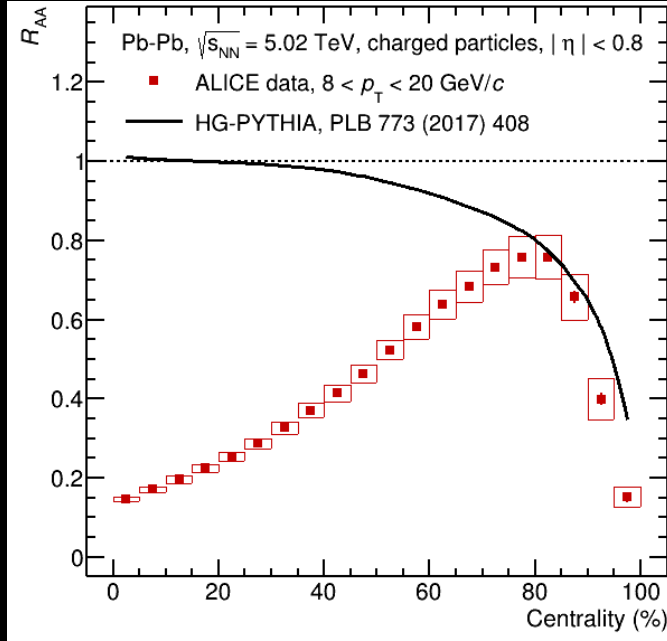
HG-Pythia model:

- incoherent superposition of Pythia pp collisions with #MPIs from HIJING-Glauber

A. Morsh & C. Loizides, PLB 773 (2017) 408

Suppression in peripheral Pb-Pb collisions?

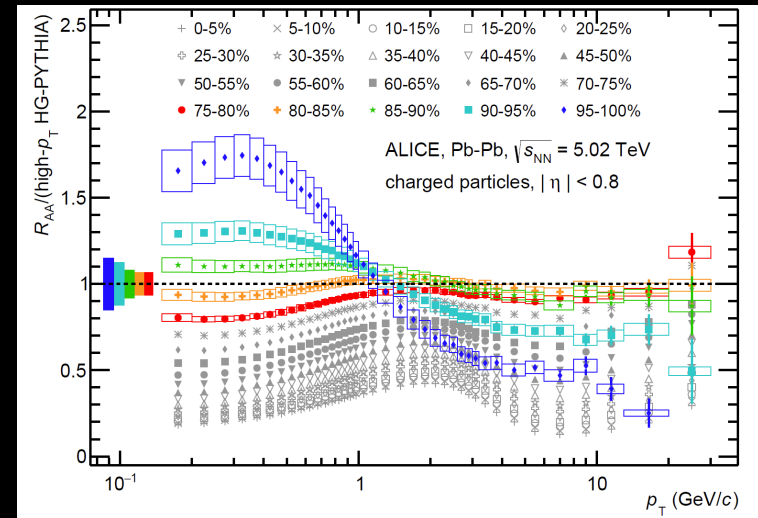
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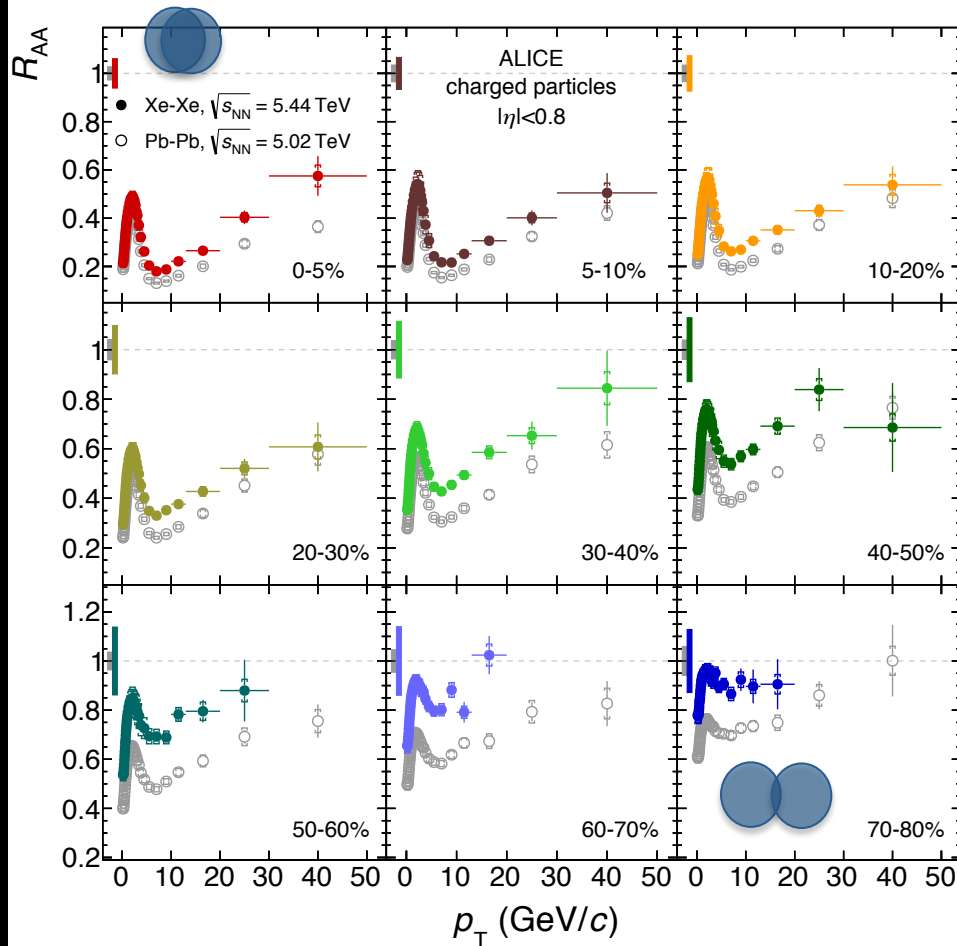


R_{AA} scaled with high- p_T bias from HG-Pythia

→ Indication that $R_{AA} \sim 1$ for 75-90% central collisions

Charged-particle R_{AA} in Pb-Pb and Xe-Xe

Phys. Lett. B 788 (2019) 166



- Similar suppression pattern in Xe-Xe and Pb-Pb
 - Larger suppression in Pb-Pb than in Xe-Xe collisions at high p_T at the same centrality
 - Normalization uncertainty (T_{AA} and pp norm.) are much larger for Xe-Xe
 - less precisely known nuclear-charge-density distribution of deformed ^{129}Xe nucleus
- Result of interplay between geometry and path length dependence of parton energy loss

Charged-particle R_{pPb} and R_{PbPb} at $\sqrt{s_{NN}} = 5.02$ TeV

JHEP 1811 (2018) 013

$$R_{AA} = \frac{1}{\langle T_{AA} \rangle} \frac{dN_{AA} / dp_T}{d\sigma_{pp} / dp_T} \equiv \frac{[medium]}{[vacuum]}$$

For $p_T > 7$ GeV/c:

- Strong suppression in central Pb-Pb collisions
- Small suppression in peripheral Pb-Pb collisions
- No modification in p-Pb collisions (no centrality selection)

→ Suppression in central Pb-Pb collisions is due to final state effects!

Confirmed also by jet measurements
ALICE, Phys. Lett. B749 (2015) 68

