

A Possible Connection Between Dark Matter and Neutrino Masses and Its Implications

Michael A. Schmidt

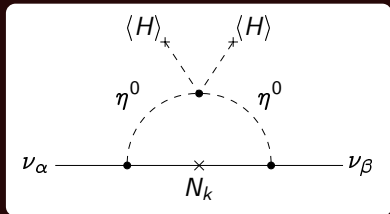
Institute for Particle Physics Phenomenology
Durham

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based on:
Y. Farzan, S. Pascoli, MS [1005.5323] and [1006.xxxx]

Radiative Mass Generation

Ma Model_[Ma (2006)]



$$(m_\nu)_{\alpha\beta} = \sum_k \frac{Y_{\alpha k} Y_{\beta k} M_k}{16\pi^2} \left[\frac{m_R^2}{m_R^2 - M_k^2} \ln \frac{m_R^2}{M_k^2} - \frac{m_I^2}{m_I^2 - M_k^2} \ln \frac{m_I^2}{M_k^2} \right]$$

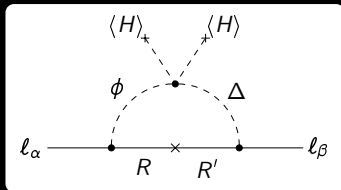
Conditions for Majorana Neutrino Mass Term

- Particle in loop have to couple to $SU(2)_L$ doublet
- Massive scalar and fermion in loop
- Mass splitting between scalar and pseudoscalar in loop
- $\Delta L = 2$ lepton number violation
- Discrete symmetry to avoid FCNCs and Dirac mass term
 $\Rightarrow \eta^0$ stable particle \Rightarrow DM candidate
- But no symmetry explanation for smallness of couplings

Global U(1)

Particle Content

	SU(2)	U(1)	U(1) _X	\mathbb{Z}_2
$\ell_L^{(i)}$	2	-1/2	0	+
R_R	2	-1/2	1	-
R'_R	2	1/2	-1	-
Δ	3	1	1	-
ϕ	1	0	-1	-



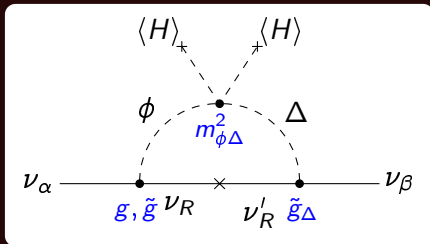
- Symmetry explanation for smallness of couplings $U(1) \rightarrow \mathbb{Z}_2$
- \Rightarrow here explicit, later spontaneous
- Symmetry protects smallness from large quantum corrections

Fermion Sector

$$- m_{RR} (R'^C)^\dagger \cdot R - g_\alpha \phi^\dagger R^\dagger \ell_{L\alpha} - \tilde{g}_\alpha \phi R^\dagger \ell_{L\alpha} - (\tilde{g}_\Delta)_\alpha R'^\dagger \cdot \Delta \cdot \ell_{L\alpha} + \text{h.c.}$$

Neutrino Masses

One Loop Diagram Generating Neutrino Masses



- neutral scalar mass eigenstates δ_i
- with scalar masses M_i
- α_1 mixing between $\delta_{1,3}$
- α_2 mixing between $\delta_{2,4}$

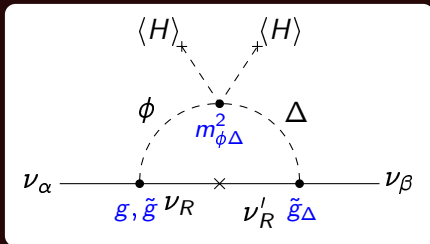
$$(m_\nu)_{\alpha\beta} = [g_\alpha(\tilde{g}_\Delta)_\beta + g_\beta(\tilde{g}_\Delta)_\alpha]\tilde{\eta} + [\tilde{g}_\alpha(\tilde{g}_\Delta)_\beta + \tilde{g}_\beta(\tilde{g}_\Delta)_\alpha]\eta$$

$$\tilde{\eta} = \frac{m_{RR}}{64\pi^2} \left(\frac{M_3^2}{m_{RR}^2 - M_3^2} \ln \frac{m_{RR}^2}{M_3^2} - \frac{M_1^2}{m_{RR}^2 - M_1^2} \ln \frac{m_{RR}^2}{M_1^2} \right) \sin 2\alpha_1 + [(1, 3) \rightarrow (2, 4)]$$

$$\eta = \frac{m_{RR}}{64\pi^2} \left(\frac{M_3^2}{m_{RR}^2 - M_3^2} \ln \frac{m_{RR}^2}{M_3^2} - \frac{M_1^2}{m_{RR}^2 - M_1^2} \ln \frac{m_{RR}^2}{M_1^2} \right) \sin 2\alpha_1 - [(1, 3) \rightarrow (2, 4)]$$

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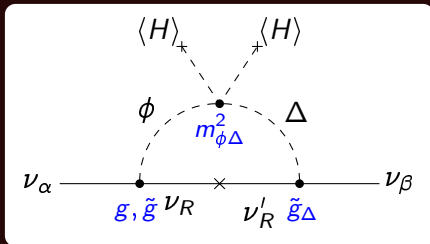
$$(m_\nu)_{\alpha\beta} = [g_\alpha(\tilde{g}_\Delta)_\beta + g_\beta(\tilde{g}_\Delta)_\alpha]\tilde{\eta} + [\tilde{g}_\alpha(\tilde{g}_\Delta)_\beta + \tilde{g}_\beta(\tilde{g}_\Delta)_\alpha]\eta$$

$$\tilde{\eta} \simeq \frac{m_{RR}}{16\pi^2} \left(\frac{\tilde{m}_\phi^2 m_{\phi\Delta}^2}{m_{RR}^2 m_\Delta^2} \left(\frac{m_{RR}^2}{m_{RR}^2 - m_\Delta^2} \ln \frac{m_{RR}^2}{m_\Delta^2} + 1 - \ln \frac{m_{RR}^2}{M_1^2} \right) - \frac{\tilde{m}_\phi^2}{m_{RR}^2 - m_\Delta^2} \ln \frac{m_{RR}^2}{m_\Delta^2} \right)$$

$$\eta \simeq - \frac{m_{RR}}{16\pi^2} \frac{m_{\phi\Delta}^2}{m_{RR}^2 - m_\Delta^2} \ln \frac{m_{RR}^2}{m_\Delta^2}$$

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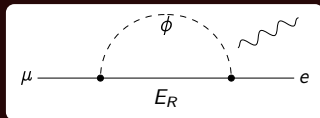
$$g\tilde{g}_\Delta \simeq 4.0 \times 10^{-6} \frac{m_\nu}{0.05 \text{ eV}} \frac{70 \text{ GeV}}{M_1} \frac{50 \text{ MeV}}{\delta} \frac{m_{RR}}{300 \text{ GeV}} \frac{0.1}{|\sin \alpha_1|} \left(\frac{m_{RR}^2}{m_{RR}^2 - m_\Delta^2} \dots \right)^{-1}$$

$$g\tilde{g}_\Delta \simeq 4.5 \times 10^{-6} \frac{m_\nu}{0.05 \text{ eV}} \frac{300 \text{ GeV}}{m_{RR}} \frac{1 \text{ GeV}^2}{\tilde{m}_{\phi\Delta}^2} \left(\frac{m_\Delta}{500 \text{ GeV}} \right)^2 \frac{m_{RR}^2 - m_\Delta^2}{m_\Delta^2} \left(\log \frac{m_{RR}^2}{m_\Delta^2} \right)^{-1}$$

$$\tilde{g}\tilde{g}_\Delta \simeq 1.8 \times 10^{-10} \frac{m_\nu}{0.05 \text{ eV}} \frac{300 \text{ GeV}}{m_{RR}} \frac{0.1}{\sin \alpha_1} \frac{m_{RR}^2 - m_\Delta^2}{m_\Delta^2} \left(\log \frac{m_{RR}^2}{m_\Delta^2} \right)^{-1}$$

Lepton Flavour Violation

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$$\text{Br}(\mu \rightarrow e \gamma) = 2.5 \cdot 10^{-9} \left(\frac{300 \text{ GeV}}{m_{RR}} \right)^4 \left| \frac{g_\mu^* g_e}{0.1 \cdot 0.1} \right|^2$$

$$\text{Br}(\tau \rightarrow \alpha \gamma) = 4.5 \cdot 10^{-10} \left(\frac{300 \text{ GeV}}{m_{RR}} \right)^4 \left| \frac{g_\tau^* g_\alpha}{0.1 \cdot 0.1} \right|^2$$

Experimental Limits

(90% CL)_[PDG 2009]

$$\text{Br}(\mu \rightarrow e \gamma) < 1.2 \cdot 10^{-11}$$

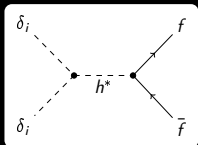
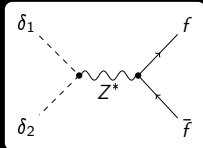
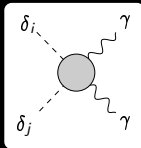
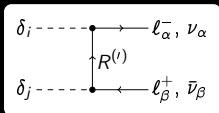
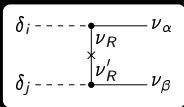
$$\text{Br}(\tau \rightarrow e \gamma) < 1.1 \cdot 10^{-7}$$

$$\text{Br}(\tau \rightarrow \mu \gamma) < 4.5 \cdot 10^{-8}$$

Solutions

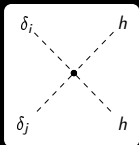
- $m_{RR}/g > 6 \text{ TeV}$
- $g_e \ll g_\mu$ or $g_\mu \ll g_e$
(allowed by flavour structure)

Dark Matter Annihilation



$$\Rightarrow \langle \sigma_{f\bar{f}}^H v \rangle \simeq N_c \frac{|\lambda_L|^2}{\pi} \frac{m_f^2}{(4M_1^2 - m_h^2)^2} \frac{(M_1^2 - m_f^2)^{3/2}}{M_1^3}$$

$$\sim 1.5 \cdot 10^{-26} \frac{\text{cm}^3}{\text{sec}} \text{ for } \lambda_L \simeq 0.044; m_h = 120 \text{ GeV}$$

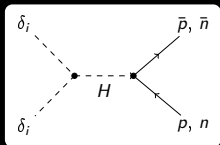


$$\Rightarrow \langle \sigma_{hh}^H v \rangle \simeq \frac{|\lambda_L|^2 (M_1^2 - m_h^2)^{1/2}}{16\pi M_1^3}$$

- Experimentally required value $\langle \sigma(\delta_1 \delta_1 \rightarrow \dots) v \rangle \sim 3 \cdot 10^{-26} \text{ cm}^3/\text{sec}$
- $\sigma_{12} \ll \sigma_{11}, \sigma_{22} \Rightarrow \delta_1$ and δ_2 produced and later $\delta_2 \rightarrow \delta_1 \nu \bar{\nu}$ with

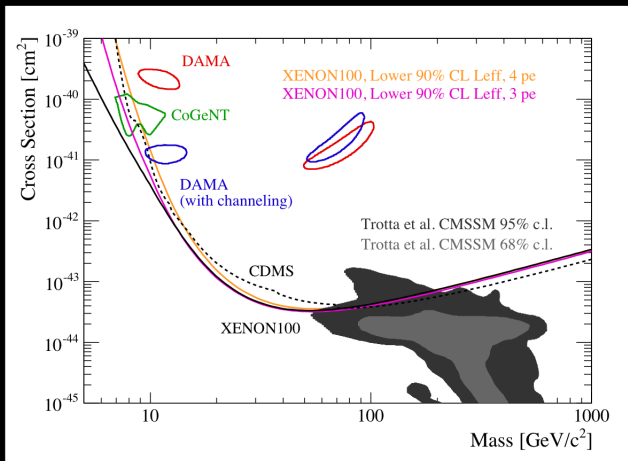
$$\Gamma(\delta_2 \xrightarrow{Z} \delta_1 \nu \bar{\nu}) \approx 14 \left(\frac{\delta}{50 \text{ MeV}} \right)^5 \left(\frac{\sin \alpha_1}{0.1} \right)^4 \text{ sec}^{-1}$$

Direct DM Detection



$$\sigma_n = \frac{|\lambda_L|^2}{\pi} \frac{\mu_{\delta_1 n}^2 m_p^2}{M_1^2 m_h^4} f^2$$

$$\approx 1.7 \cdot 10^{-44} \text{ cm}^2 \left(\frac{\lambda_L}{0.044} \right)^2 \left(\frac{70 \text{ GeV}}{M_1} \right)^2 \left(\frac{120 \text{ GeV}}{m_h} \right)^4 \left(\frac{f}{0.3} \right)^2$$



Constraints

Electroweak Precision Tests: Higgs Triplet

$$\hat{S} = \frac{g_{\text{SU}(2)}^2}{24\pi^2} \xi \quad \hat{T} = \frac{25g_{\text{SU}(2)}^2}{576\pi^2} \frac{m_{\Delta}^2}{m_W^2} \xi^2$$

with $\xi := 1 - m_{\Delta^{++}}^2/m_{\Delta}^2$ and $m_{\Delta^{++}}^2 = m_{\Delta}^2 + 2m_{\Delta^+}^2$

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Invisible Z-Decay Width

- DM particle δ_1 couples to Z-boson via mixing of Δ^0
- If $M_1 + M_2 < m_Z$, the corresponding Z-decay width is

$$\Gamma(Z \rightarrow \delta_1 \delta_2) = \frac{G_F \sin^2 \alpha_1 \sin^2 \alpha_2}{6\sqrt{2}\pi} m_Z^3$$

- Bound on mixing angle in scalar sector: $\sin \alpha_1 \sin \alpha_2 < 0.07$,
i.e. $m_{\phi\Delta}^2 \ll m_{\Delta}^2$ (protected by $U(1)_{\phi} \times U(1)_{\Delta}$)

Collider Physics

Higgs Search

- Higgs might dominantly decay invisibly if $2 M_1 < m_h$

$$H \rightarrow \delta_1 \delta_1, \quad H \rightarrow \delta_2 \delta_2 \rightarrow (\delta_1 \nu \bar{\nu})(\delta_1 \nu \bar{\nu})$$

- Displaced vertex if mass splitting large $\delta_2 \rightarrow \delta_1 \mu^+ \mu^-$

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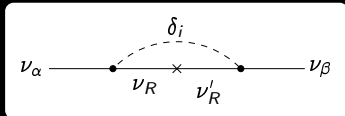
New Particles

- New particles accessible at LHC
- ... decay into the SM particles and DM \Rightarrow missing energy
- Mass relation of triplet $2m_{\Delta^+}^2 = m_{\Delta^{++}}^2 + m_{\Delta^0}^2$
- Expect small mass splitting $m_{\Delta^{++}}^2 - m_{\Delta^+}^2$
- Determination of g_α : $\text{Br}(E_R^- \rightarrow \ell_\alpha^- \delta_{1,2}) \propto |g_\alpha|^2$
- Determination of \tilde{g}_Δ : decay modes of Δ^+ and Δ^{++} , especially $\Gamma(\Delta^{++} \rightarrow \ell_\alpha^+ \ell_\beta^+ \delta_{1,2}) \propto |(\tilde{g}_\Delta)_\alpha g_\beta + (\tilde{g}_\Delta)_\beta g_\alpha|^2$

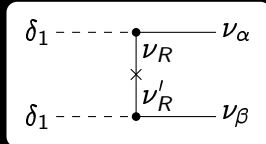
Alternative Scenario [Boehm, Farzan, Hambye, Palomares-Ruiz, Pascoli (2006)]

- Not $U(1)_X$ symmetry but lepton number L or $B - L \Rightarrow \tilde{g}_\Delta$ not small

Neutrino Mass generation



DM annihilation

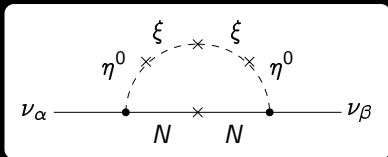


- Direct relation between neutrino mass and DM annihilation \Rightarrow light dark matter $M_1 \sim \mathcal{O}(\text{MeV})$
- MeV scale dark matter might explain the SPI/INTEGRAL signal

Gauge Model

Particle Content

	SU(2)	U(1)	U(1) _X	\mathbb{Z}_2
$\ell_L^{(i)}$	2	-1	0	+
N_R	1	0	1	-
N_L	1	0	1	-
η	2	-1	-1	-
ξ	1	0	-1	-
ϕ	1	0	-2	+



- Neutrino masses from pseudo Dirac neutrinos N

$$m_{\alpha\beta} \simeq \frac{\kappa \lambda_{H\eta\xi\phi}^2}{16\pi^2} \frac{\langle\phi\rangle^4}{m_\eta^4} \frac{\langle H\rangle^2}{m_N^2} \sum_{ij} (Y_N)_{i\alpha} (Y_N)_{j\beta} f_{ij}(m_N, Y_{LL}, Y_{RR}, m_\xi)$$

- LFV similar: $(m_N, m_{\eta^-})/Y_N \gtrsim 6 \text{ TeV}$ (unless special flavour structure)

Conclusions

- Neutrino mass can be generated radiatively at the TeV scale and linked to dark matter
- Discrete symmetry seems necessary for radiative neutrino mass generation
- U(1) gauge symmetry might be origin of discrete \mathbb{Z}_2 [See also talk by Batell]
- It is testable and the interplay of different experiments imposes strong constraints
- Upcoming experiments may exclude or confirm these models
- Strongest constraints from lepton flavour violating rare decays
- Higgs searches might be more challenging

Thank you very much for your attention.

More Phenomenology

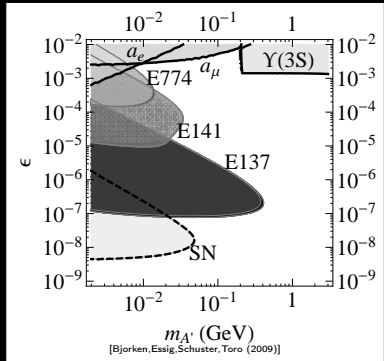
Z'

Gauge kinetic Lagrangian

$$\mathcal{L}_{U(1)} = \frac{1}{4}F_Y^2 + \frac{1}{4}F_X^2 + \frac{\epsilon}{2}F_X F_Y$$

η in loop induces

$$\epsilon = \frac{g_X g_Y}{3\pi^2} \ln \frac{\Lambda^2}{\mu_\eta^2} \lesssim 1.3 g_X g_Y$$



Collider

- SM Higgs mixes with $U(1)_X$ -Higgs \Rightarrow Higgs mass bounds weakened
- Invisible Higgs decay like in all DM models in which Higgs exchange dominates