

Testing effective Yukawa couplings in Higgs searches at Tevatron & LHC

Emidio Gabrielli

CERN PH-TH

in collaboration with B. Mele

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Yukawa couplings

- Hierarchy of fermion masses is still a puzzle
- in SM, problem just shifted to the Yukawa sector
- difficult to explain all fermion spectrum and CKM mixing by means of few parameters
- many SM extensions proposed → no one can be considered as conclusive
- maybe ChSB and EWSB have not the same origin → different symmetry-breaking mechanisms
- examples: compositeness, extradimensions, technicolor, unknown mechanism of ChSB ?

- However, there are indications that SM Higgs mechanism is behind the EWSB
- EW precision tests favour a light SM Higgs boson
- perturbative unitarity in $WW \rightarrow WW$ satisfied

what if the SM Higgs boson is only responsible of EWSB \rightarrow W and Z masses ?

- fermion masses put in by hand, no tree-level Yukawa couplings \rightarrow explicit ChSB
- SM becomes non-renormalizable, but still it can be considered as an effective field theory

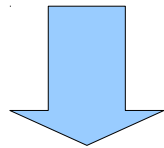
suppose Λ is the scale where Yukawas are vanishing (i.e. scale of fermion mass generation)

■ Yukawa couplings not protected against radiative corrections due ChSB → radiatively generated

■ we need to re-sum large logs $g_i^{2n} \log^n (\Lambda/m_H)$

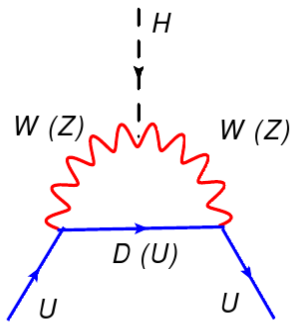
■ efficient tool → Renormalization Group Equations

■ SM RGE are not valid (tree-level Yukawa couplings)

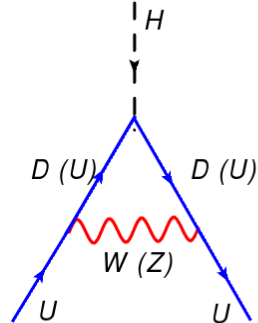


RGE must be derived by keeping Yukawa couplings and fermion masses separate

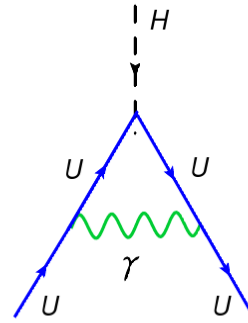
Relevant Feynman diagrams for the Up Yukawa's beta-functions. Unitary gauge



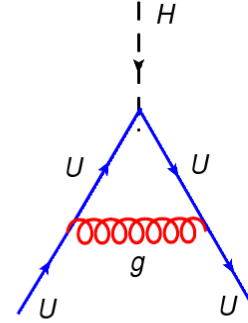
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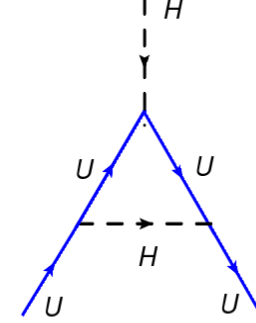
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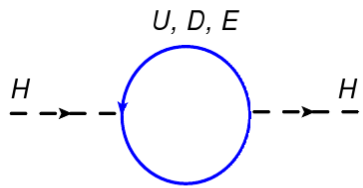
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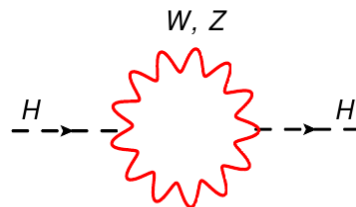
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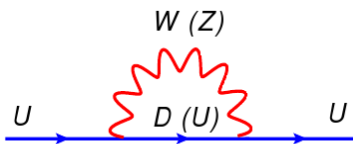
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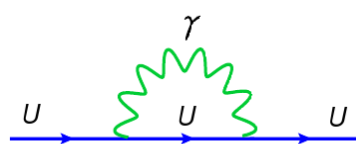
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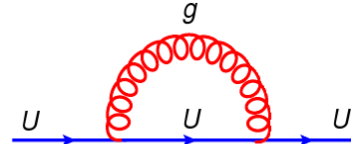
g)



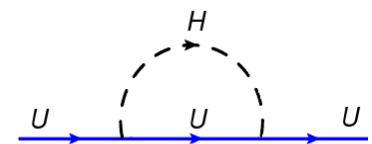
h)



i)



l)



m)

RGE for Yukawa couplings of U , D quarks

SM RGE recovered for $\mathbf{Y}_f^{\text{SM}} \rightarrow \mathbf{Y}_f$

$$\frac{d\mathbf{Y}_U}{dt} = \frac{1}{16\pi^2} \left\{ 3\xi_H^2 (\mathbf{Y}_U - \mathbf{Y}_U^{\text{SM}}) - 3\mathbf{Y}_U^{\text{SM}}\mathbf{Y}_D^{\text{SM}} (\mathbf{Y}_D - \mathbf{Y}_D^{\text{SM}}) + \frac{3}{2}\mathbf{Y}_U (\mathbf{Y}_U\mathbf{Y}_U - \mathbf{Y}_D^{\text{SM}}\mathbf{Y}_D^{\text{SM}}) - \mathbf{Y}_U \left(\frac{17}{20}g_1^2 + \frac{9}{4}g_2^2 + 8g_3^2 - \text{Tr}(\mathbf{Y}) \right) \right\},$$

$$\frac{d\mathbf{Y}_D}{dt} = \frac{1}{16\pi^2} \left\{ 3\xi_H^2 (\mathbf{Y}_D - \mathbf{Y}_D^{\text{SM}}) - 3\mathbf{Y}_D^{\text{SM}}\mathbf{Y}_U^{\text{SM}} (\mathbf{Y}_U - \mathbf{Y}_U^{\text{SM}}) + \frac{3}{2}\mathbf{Y}_D (\mathbf{Y}_D\mathbf{Y}_D - \mathbf{Y}_U^{\text{SM}}\mathbf{Y}_U^{\text{SM}}) - \mathbf{Y}_D \left(\frac{1}{4}g_1^2 + \frac{9}{4}g_2^2 + 8g_3^2 - \text{Tr}(\mathbf{Y}) \right) \right\},$$

$$\mathbf{Y} \equiv N_c \mathbf{Y}_U \mathbf{Y}_U + N_c \mathbf{Y}_D \mathbf{Y}_D + \mathbf{Y}_E \mathbf{Y}_E$$

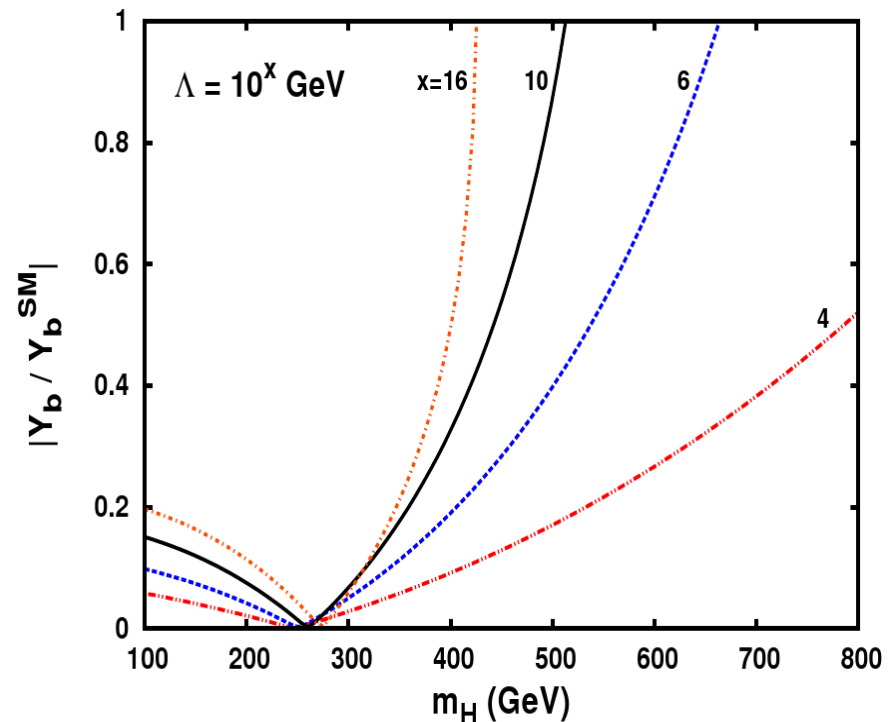
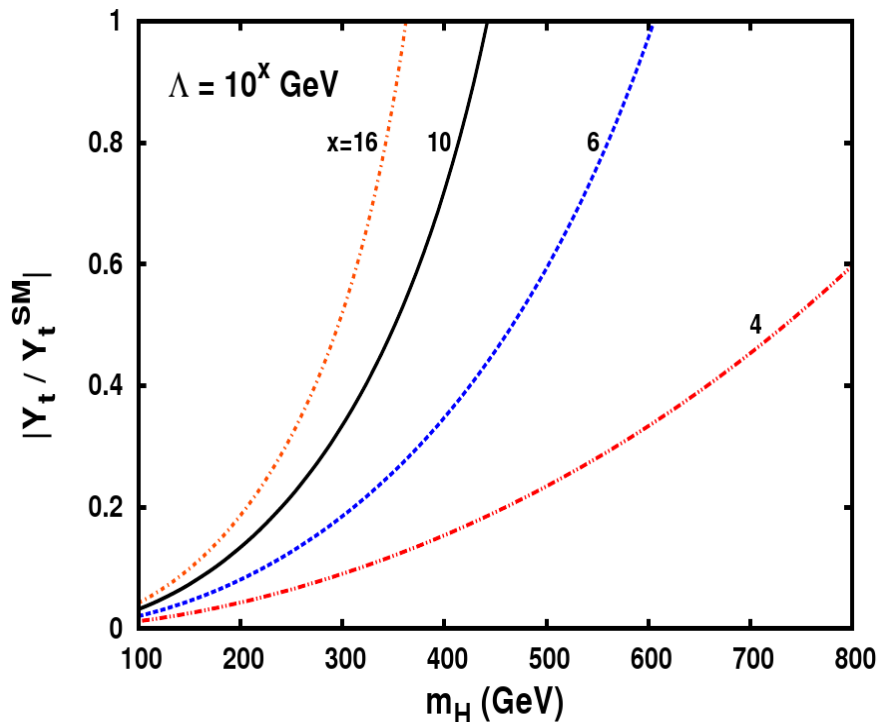
W(L) polarizations

$$\xi_H \equiv \frac{g_2 m_H}{2M_W}, \quad \mathbf{Y}_f^{\text{SM}} \equiv \frac{g_2}{\sqrt{2}M_W} \text{diag}[\mathbf{m}_f]$$

ChSB terms

Theoretical Framework

- Yukawa couplings set to zero at scale Λ and connected to m_H scale by solving the RGE
 - only SM degrees of freedom below Λ assumed
 - $|Y(m_H)| < 1$ \rightarrow satisfied for light Higgs masses
- top-Yukawa (mH)** **b-Yukawa (mH)**



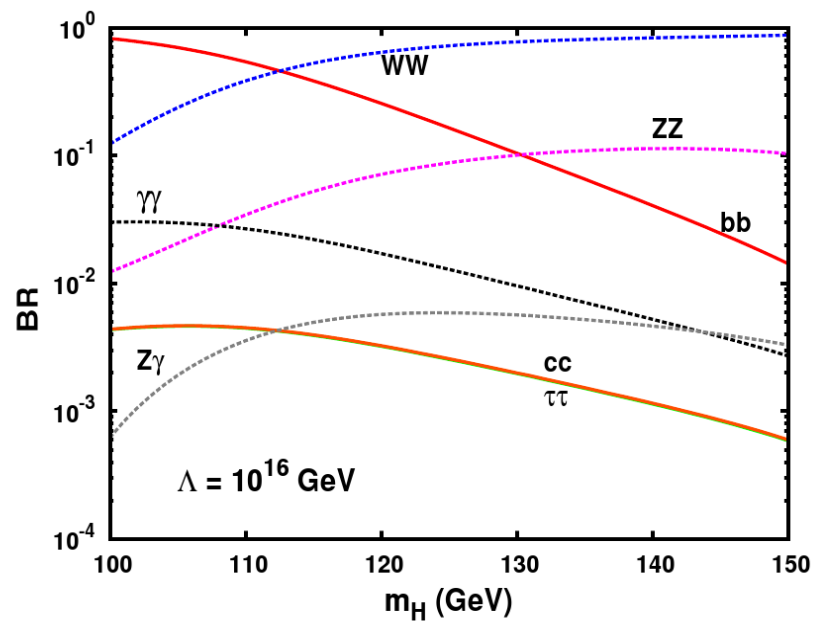
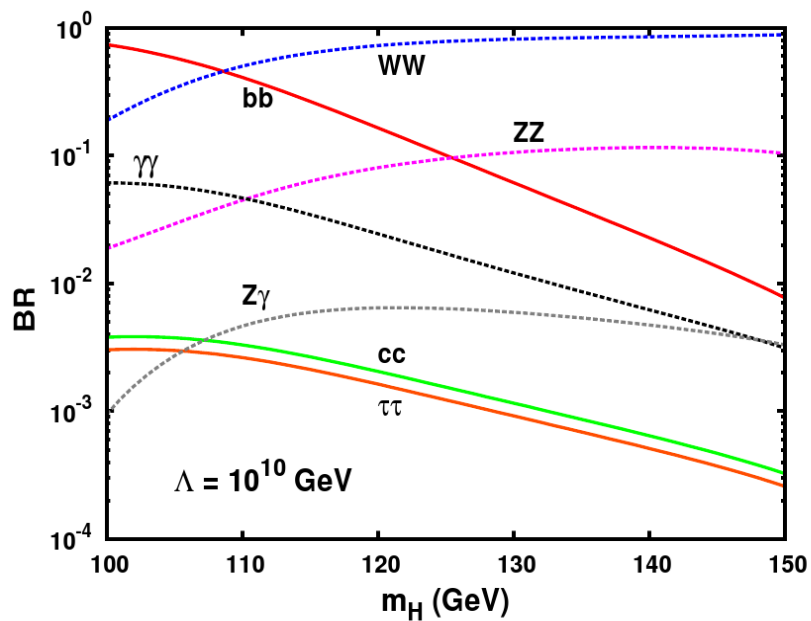
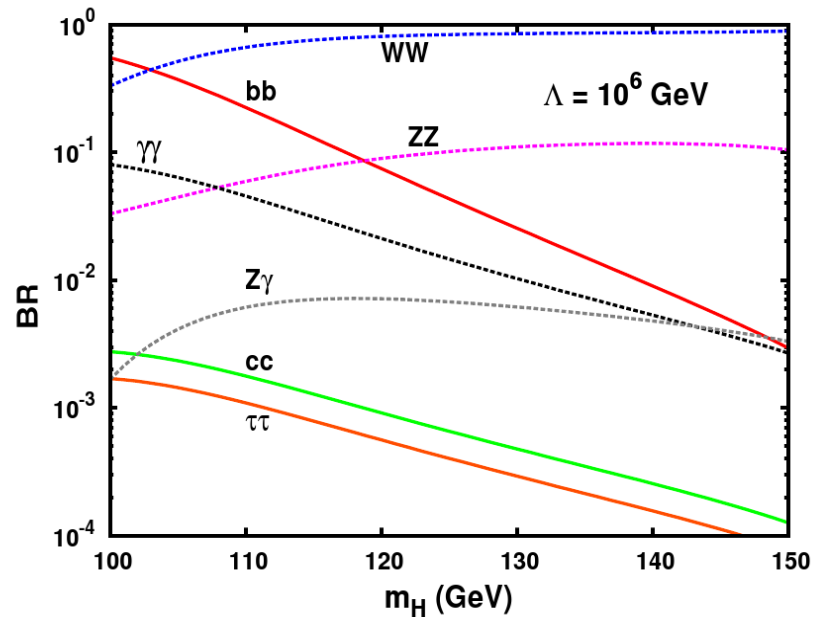
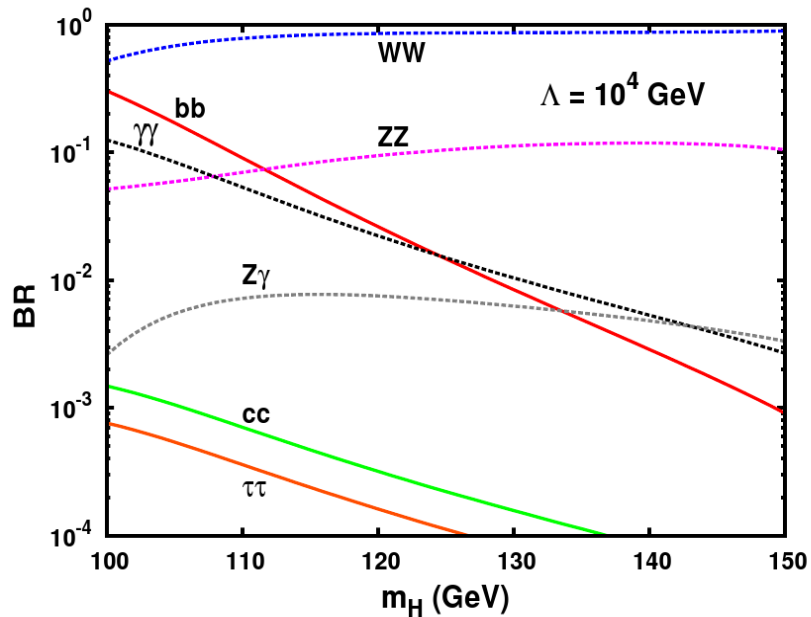
Higgs decay modes dramatically affected

- main Higgs decays into $\gamma\gamma$, WW , ZZ , γZ
- $BR(H \rightarrow bb)$ can be comparable to enhanced $\gamma\gamma$
- differs from naïve fermiophobic scenarios where Yukawa couplings are set to zero at EW scale

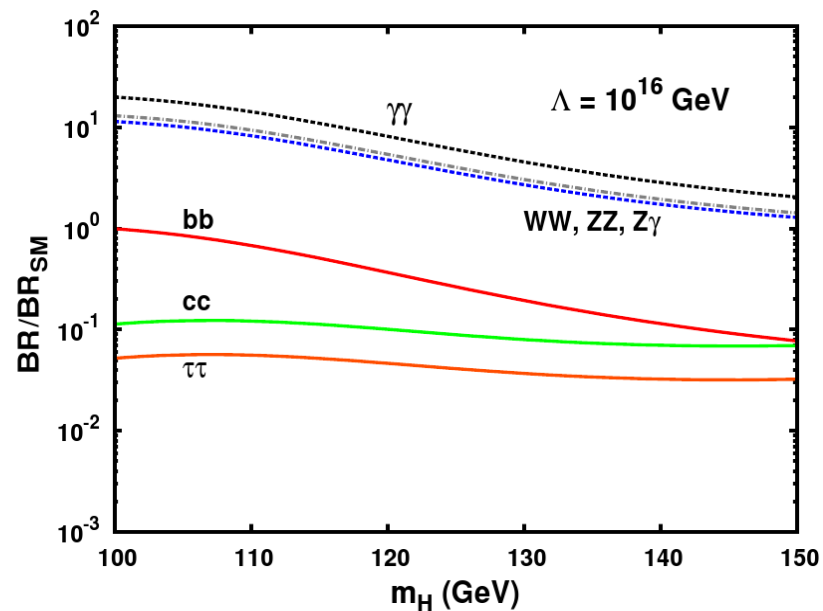
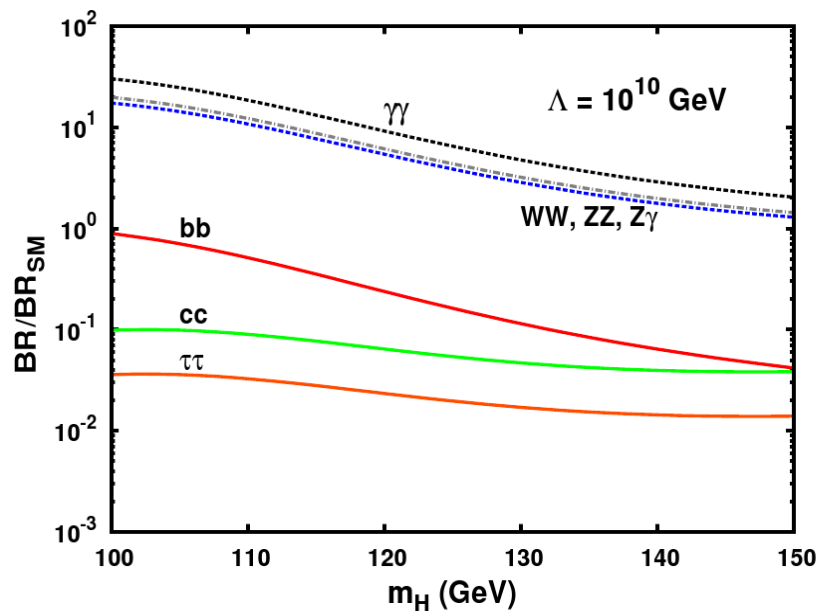
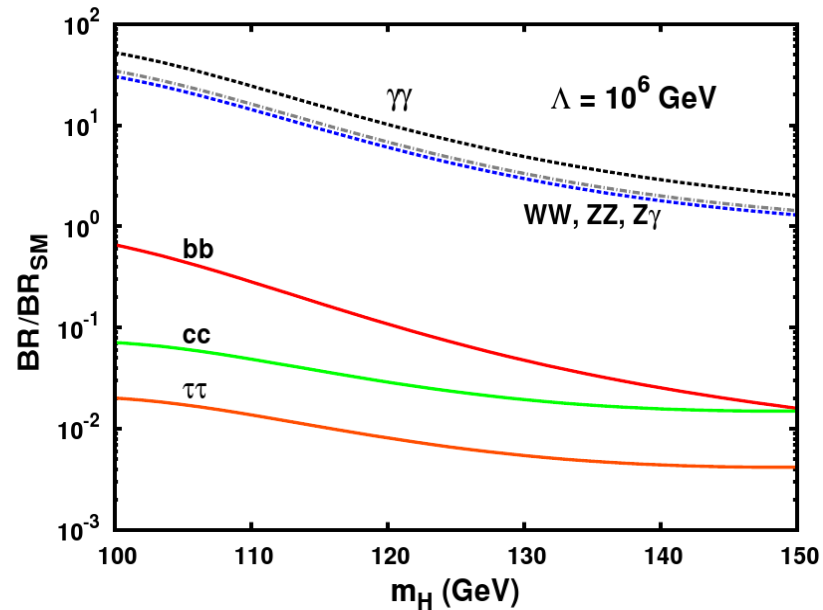
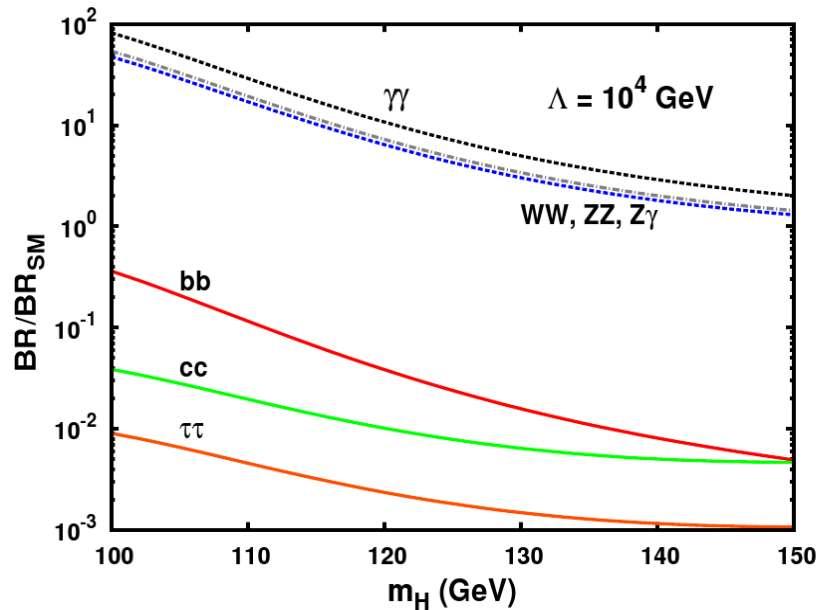
Higgs production mechanisms change

- top-Yukawa coupling radiatively induced (small)
- gluon-gluon fusion suppressed
- Vector boson fusion (VBF) becomes the dominant production mechanism

Higgs branching ratios



SM normalized branching ratios



Higgs production at hadron colliders

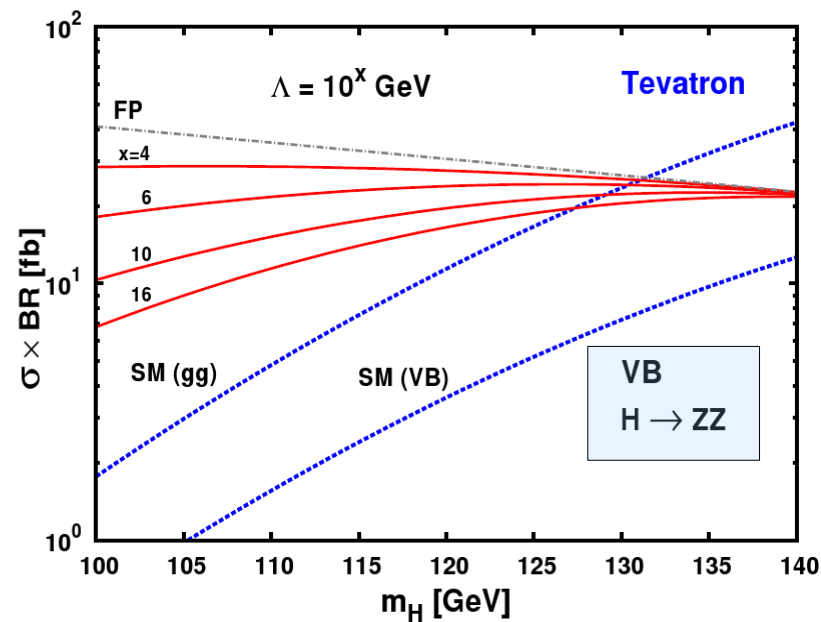
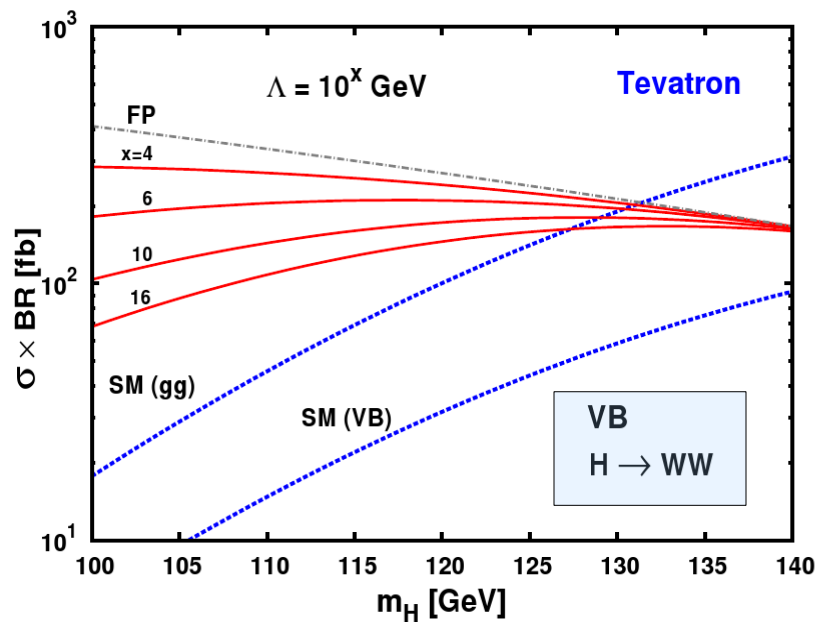
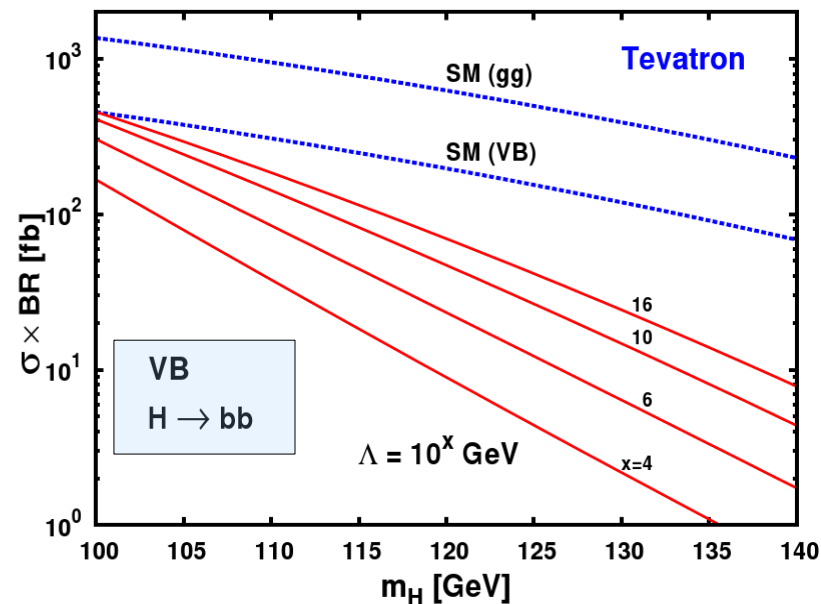
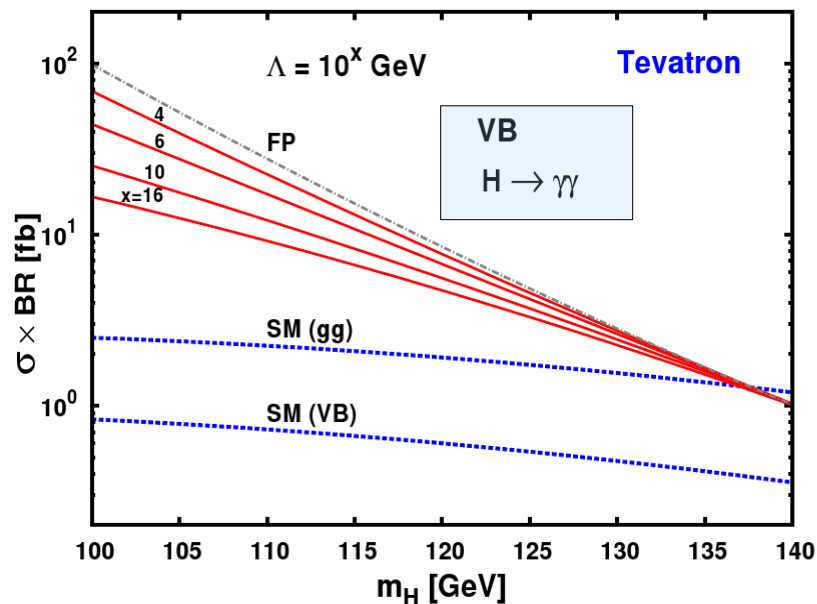
Tevatron $\sqrt{S} = 1.96 \text{ TeV}$

- Inclusive cross sections used **VB=ZH+WH+VBF**
- gg fusion contribution neglected (suppressed by more than **SM/100**)
- computed at NLO, central values presented

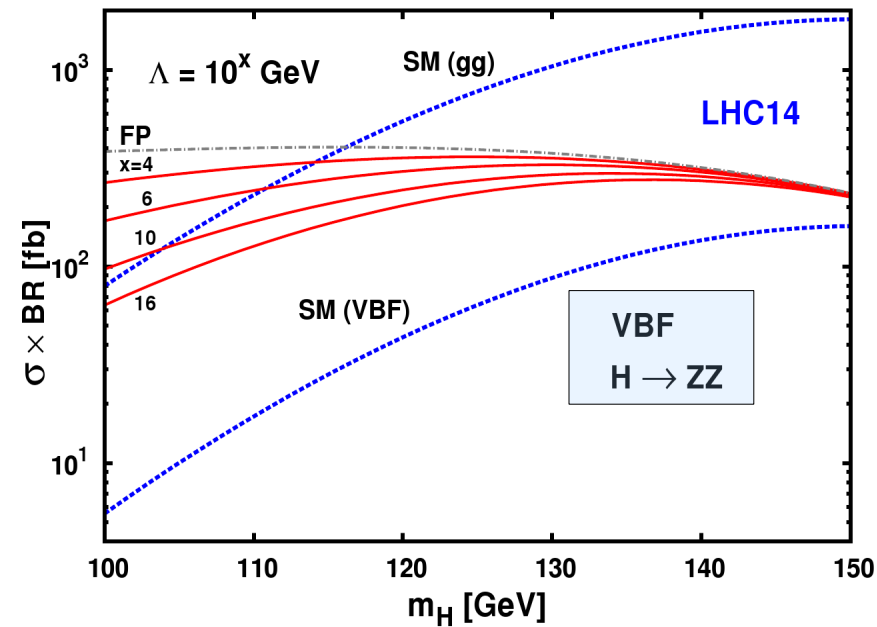
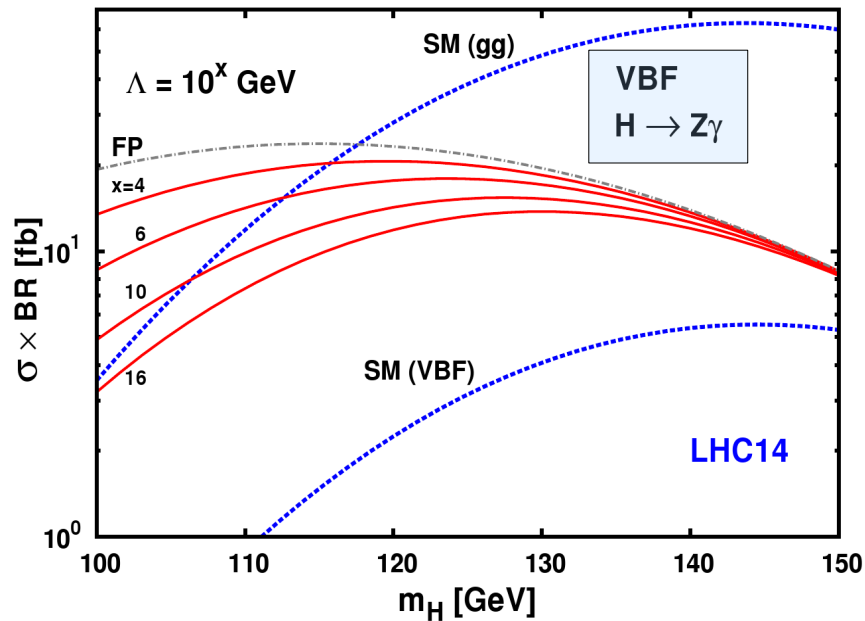
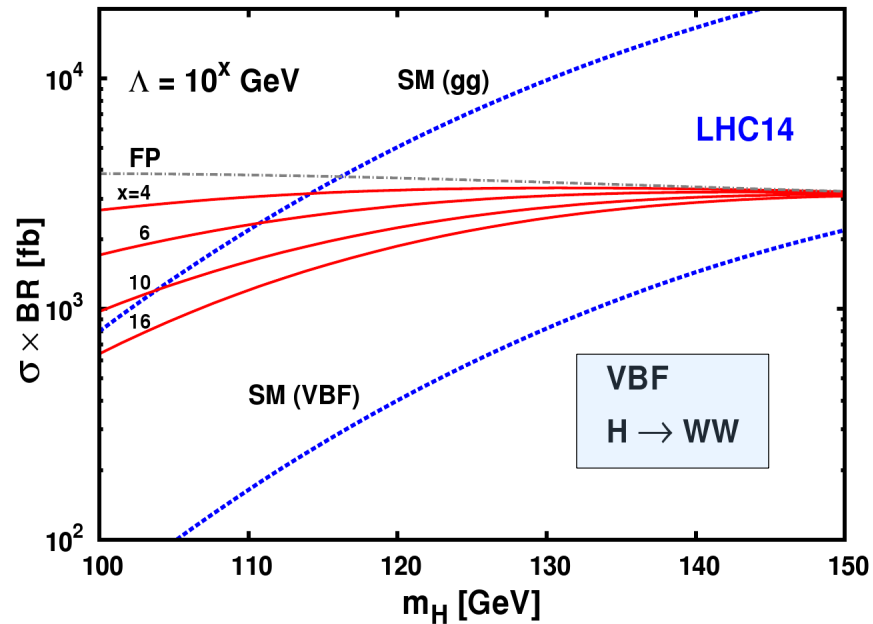
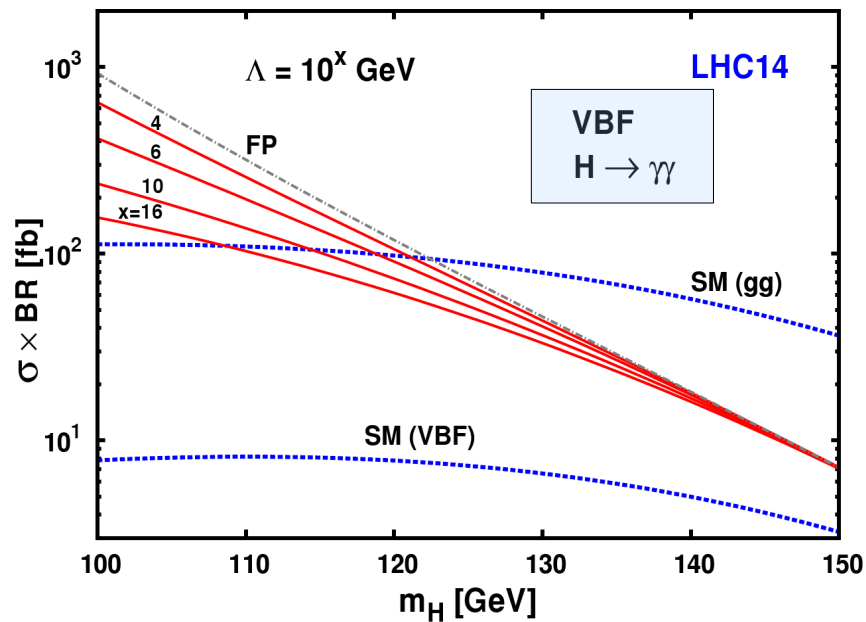
LHC $\sqrt{S} = 14 \text{ TeV}$

- cross section used \rightarrow **VBF** at NLO
- gg fusion neglected (a few percent of VBF)
- central values presented

Tevatron cross sections X BR (fb)



LHC14 cross sections X BR (fb)



Conclusions

- **Scenario: Yukawa vanishing at a scale Λ**
- **radiatively generated at EW scale**
- **enhanced $H \rightarrow \gamma\gamma, WW, ZZ, Z\gamma$**
- **non-trivial depletion of bb decay**
- **Tevatron: could probe $m_H < 110$ GeV**
- **LHC 14 TeV**
 - **excellent probe of this scenario**
 - **VBF signature larger than SM**
 - **better S/B ratio compared to SM**
 - **better theoretical accuracy (VBF)**
- **LC $e+e-$ best probe: allow to directly test radiatively induced Yukawa couplings**