

LHC AND THE ORIGIN OF NEUTRINO MASS

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experiment



tiny neutrino mass: $1/30 \text{ eV} \leq m_\nu \leq 1 \text{ eV}$



$m_\nu \ll m_e$ - interesting puzzle?

window to new physics

Neutrino mass in the standard model

Fermi-like effective: Weinberg $d = 5$ operator

$$\mathcal{L} = Y_{eff} \frac{LHHL}{\Lambda}$$

L - lepton doublet H - Higgs doublet

$$m_\nu = Y_{eff} \frac{v^2}{\Lambda} \quad \text{neutrino mass} - \text{Majorana}$$

perturbative cut-off: $\Lambda \simeq 10^{14} GeV$

- clue why $m_\nu \ll m_e$?

- case for new physics

Violation of Lepton Number : $\Delta L = 2$

- neutrino-less double beta decay $\nu 0\beta\beta$ a text-book fact

Racah '37; Furry '38

- same sign charged lepton pairs in colliders

Keung, G.S., '83

at which scale: $\Lambda=?$

Fermi \rightarrow renormalizable W picture

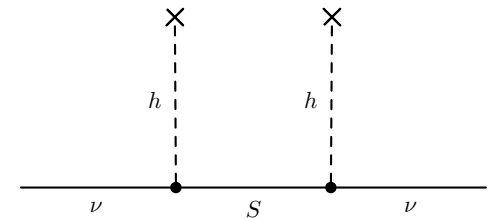
Origin of neutrino mass: seesaw

Single set of new particles

- fermion singlet $S_F = N$ ($Y = 0$): Type I

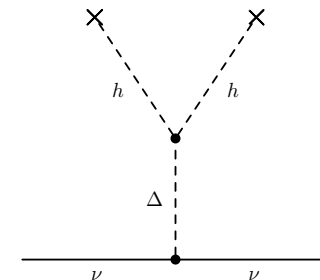
Minkowski, '77; Mohapatra, G.S., '79

Gell-Mann et al, '79; Glashow, '79; 'Yanagida, 79



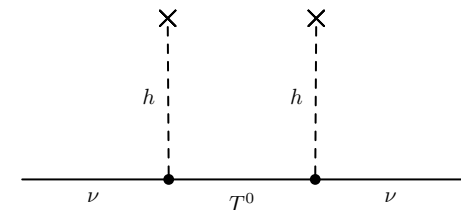
- boson weak triplet Δ ($Y = 2$): Type II

Lazarides et al, '80; Mohapatra, G.S., '80



- fermion weak triplet T_F ($Y = 0$): Type III

Foot et al, '86



standard lore (especially I and II)

More useful than Weinberg $d=5$ operator?

weak interactions: perturbative cut-off 300 GeV

effective Fermi V-A theory

Marshak, Sudarshan; Feynman, Gell-Mann '58

until we had a theory behind \rightarrow Standard Model

The theory behind seesaw?

L-R symmetry

Pati, Salam; Mohapatra, G.S. '74

$$SU(2)_L \times U(1)_{B-L} \times SU(2)_R$$

- $W_L \Rightarrow W_R$: LR at high energy
- $\nu_L \Rightarrow \nu_R$: massive neutrinos
- seesaw: connects m_ν to scale of LR restoration

$$m_N \propto M_{W_R} \quad \Rightarrow \quad m_\nu \propto M_{W_L}^2 / M_{W_R}$$

Minkowski, Mohapatra, G.S.

V-A limit:

M_{W_R} infinite \Rightarrow neutrinos massless

idea of P restoration as old as of P violation

Lee, Yang '56

mirror fermions

$$f_{L(R)} \leftrightarrow F_{R(L)}$$

3 more generations: V + A?

Gell-Mann, Minkowski '75

Gell-Mann, Ramond, Slansky '79

Wilczek, Zee '82

G.S., Wilczek, Zee '84

Bereziani, Mohapatra '90's

L - R theory: neutrino mass and seesaw

(both type I and II)

L number violation

LHC?

Model content: **bidoublet** $\phi \sim (h_{light}, H_{heavy})$, **triplets** Δ_L, Δ_R ,

$$\langle \Delta_L \rangle = \begin{pmatrix} \\ v_L \end{pmatrix}, \quad \langle \Delta_R \rangle = \begin{pmatrix} \\ v_R \end{pmatrix}, \quad \langle \phi \rangle = \begin{pmatrix} v \\ v' \end{pmatrix}$$

spontaneously with $v_L \ll v' < v \ll v_R$.

Mohapatra GS '75 '81

- Quark, Dirac lepton masses from $\bar{\psi}_L (Y_\phi \phi + \tilde{Y}_\phi \tilde{\phi}) \psi_R$

$$M_u = |v| Y_\phi + |v'| e^{i\alpha} \tilde{Y}_\phi$$

$$M_d = |v'| Y_\phi + |v| e^{i\alpha} \tilde{Y}_\phi$$

- Majorana neutrino masses, in addition to Dirac

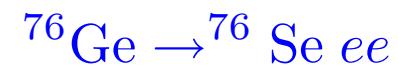
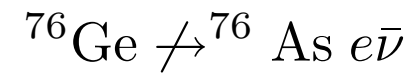
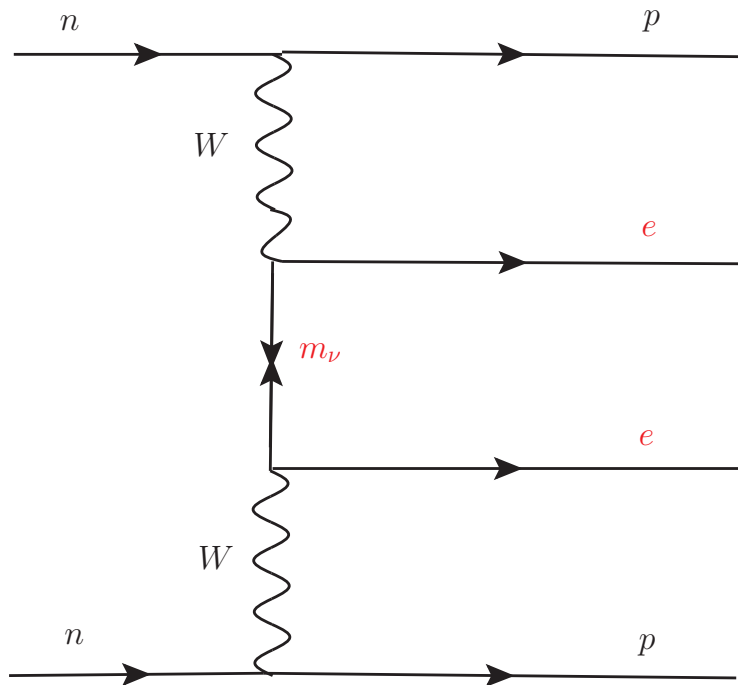
$$m_{LL} = Y_\Delta \langle \Delta_L \rangle \ll m_{RR} = Y_\Delta \langle \Delta_R \rangle$$

- Spectrum: $W_R, \nu_R, \Delta_{L,R}$ in the TeV region?

- H should be very heavy (tree-level FC)

Senjanović, GS '80, ..., Zhang et al '07

Neutrino-less double beta decay



Racah, '37; Furry, '38

Majorana '37

$$\mathcal{A}_{LL} \propto \frac{m_\nu}{p^2} \leq 10^{-8} \text{GeV}^{-1}$$

$$(p \simeq 100 \text{MeV})$$

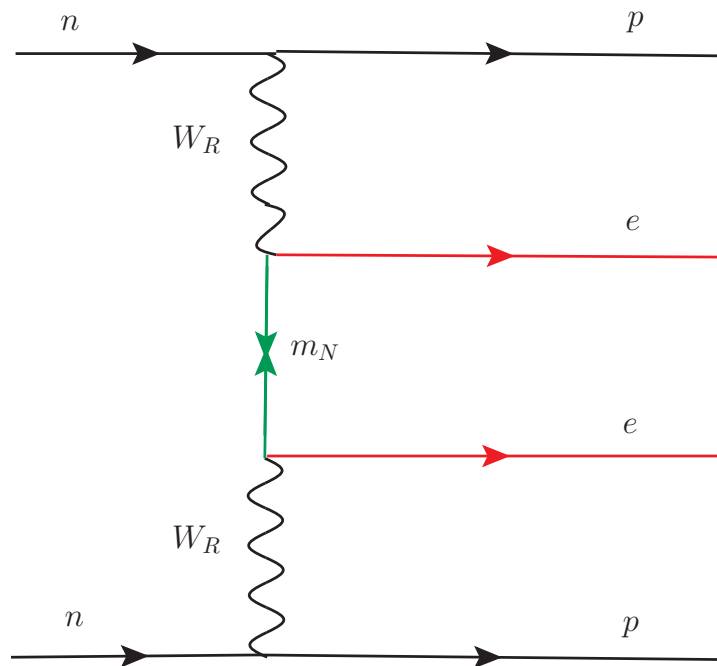
In general m_ν not directly connected to $\nu 0\beta\beta$ decay:

depends on the completion of the SM

Example:

LR symmetry with low W_R, ν_R masses: a nonzero $\nu 0\beta\beta$ decay

even with $y_D, m_\nu \rightarrow 0$



$$\mathcal{A}_{RR} \propto \left(\frac{M_L}{M_R} \right)^4 \frac{1}{M_N} \simeq 10^{-8} \text{GeV}^{-1}$$

$$M_R \simeq 2.5 \text{ TeV} \ \& \ M_N \simeq 100 \text{ GeV}$$

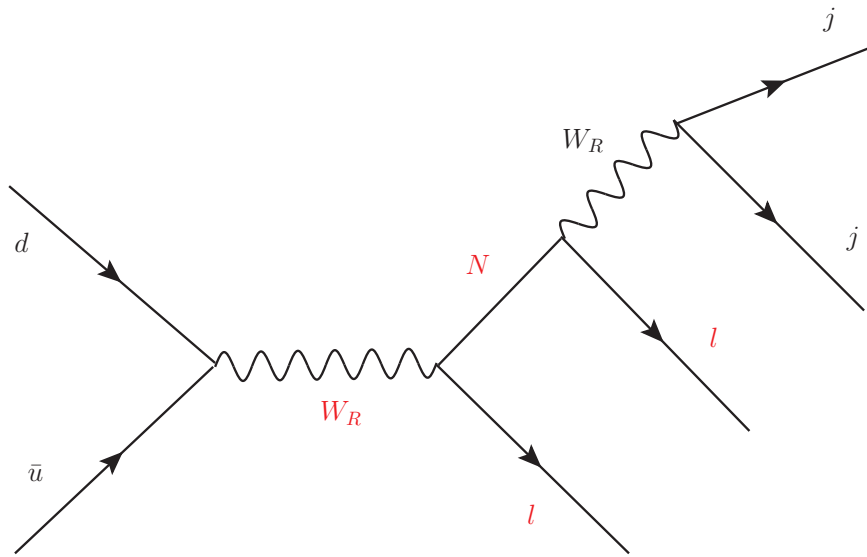
Mohapatra, G.S. '80

Colliders

To trace see-saw : measure $\Delta L = 2$ in colliders

Keung, G.S., '83

produce right-handed neutrinos
through W_R

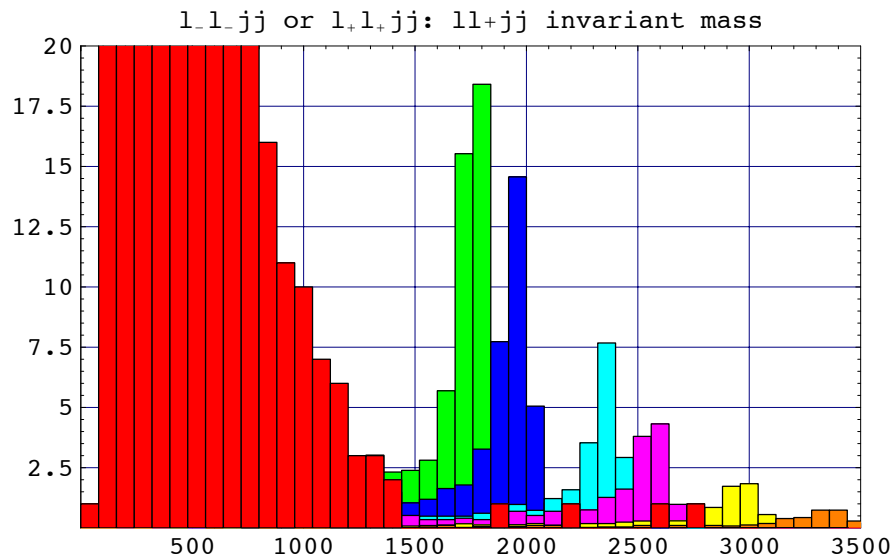


- direct test of parity restoration
- direct test of lepton number violation
- determination of W_R and N masses

Search at LHC

$W_R \rightarrow 4$ (2) TeV and ν_R in 100 - 1000 GeV
 14 (7) TeV and integrated luminosity of 30 fb^{-1}

Ferrari et al, '99; Gninenko et al, '07

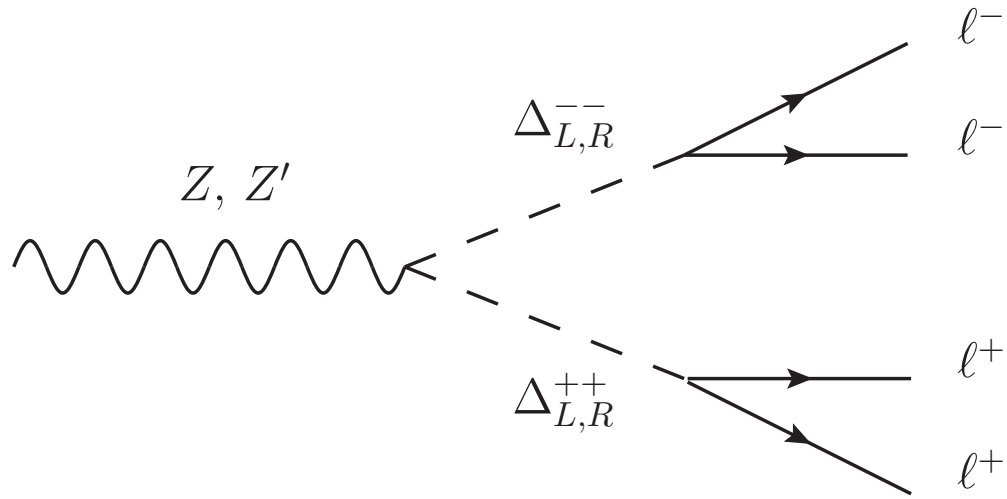


of events for $L = 8\text{fb}^{-1}$

$M_R(\text{TeV})$

1.8; 2, 0; 2.4; 2.6; 3.0; 3.4

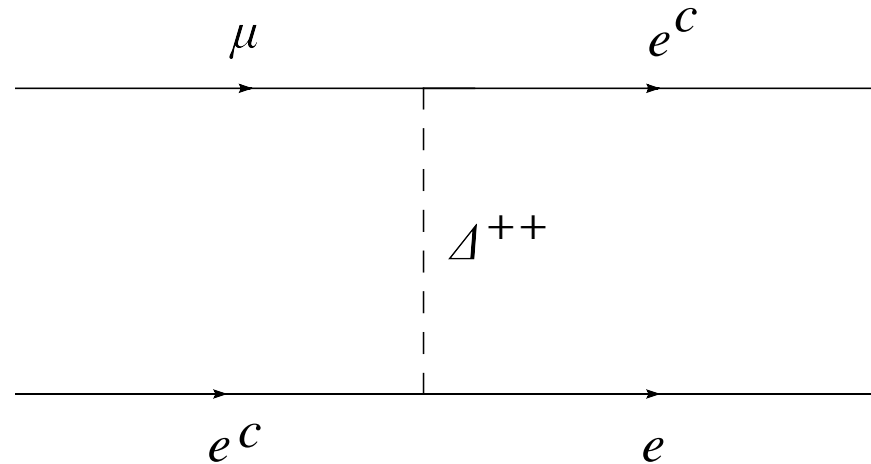
Nesti, talk at CERN Collider meeting '09 and parallel talk



$$\propto (Y_{\Delta})_{ij} (Y_{\Delta}^*)_{kl}$$

$$Y_{\Delta} = \frac{g_R}{M_{W_R}} V_R^T M_N V_R$$

LFV and $0\nu\beta\beta$



$$B(\mu \rightarrow 3e) = \frac{|Y_{e\mu} Y_{ee}^*|^2}{4G_F^2} \left(\frac{1}{M_{\Delta_L}^4} + \frac{1}{M_{\Delta_R}^4} \right)$$

$$Y_{\Delta} = \frac{g_R}{M_{W_R}} V_R^T M_N V_R$$

No predictions - Requires LHC

Ghosh, Nemevšek, Nesti, Tello, Zhang, GS - in progress

Type II

$$V_R = K_e V_L^*$$

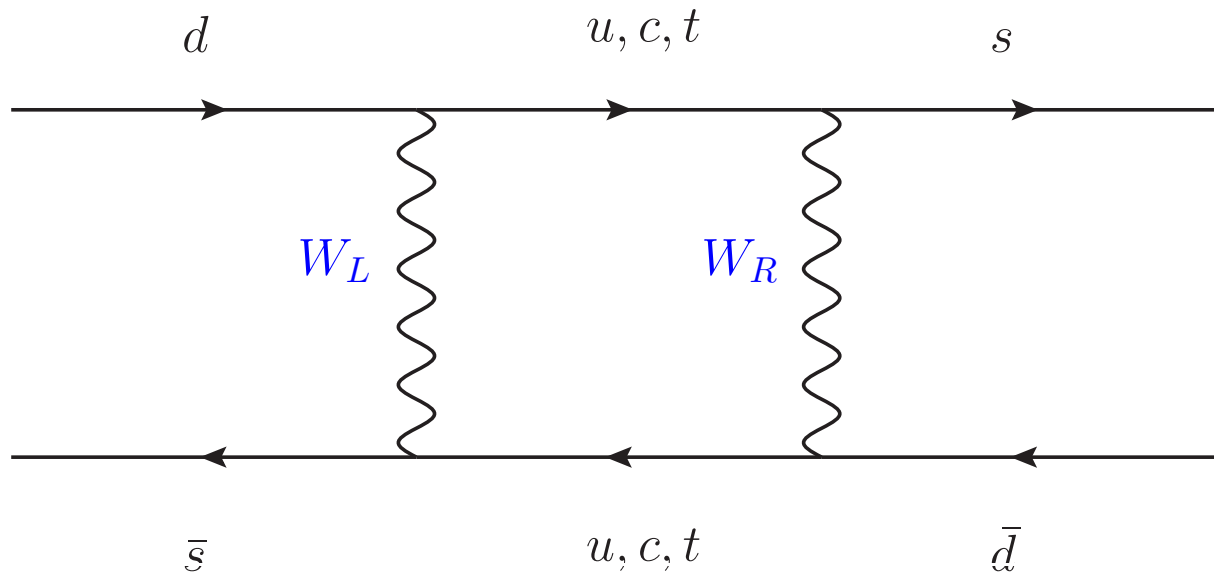


predictions

Nemevšek, Nesti, Tello, GS to appear

Limits on M_R

- direct limit: $M_R \geq 800 \text{ GeV}$ (from dijets@DO)
- strong sensibility from $K_L - K_S$ mass difference



huge enhancement - around 1000 - over LL

Beall, Bander, Soni '81
Mohapatra, G.S., Tran '83
parallel talk by Nemevšek

depends on **LR symmetry**

- **LR = P (parity)** - canonical

hermitian Yukawas: $V_L = V_R$ - up to phases

best limit claimed: combined EDM and ϵ'

$$M_{W_R} \gtrsim 4 \text{ TeV}$$

Zhang, An, Ji, Mohapatra '07

even stronger limit claimed

$$M_{W_R} \gtrsim 10 \text{ TeV}$$

Xu, An, Ji '09

⇒ LR **NOT**@LHC?

recently revisited:

$$M_{W_R} \gtrsim 3 \text{ TeV}$$

Maiezza, Nemevšek, Nesti, G.S. '10

- **LR = C (charge conjugation)**

symmetric Yukawas: $V_L = V_R^*$ - up to phases

best limit from $K_L - K_S$:

$$M_{W_R} \geq 2.5 \text{ TeV} \quad M_{Z_R} \geq 4 \text{ TeV}$$

Maiezza, Nemevšek, Nesti, G.S. '10

parallel talks by Nesti, Nemevšek



LHC reach

additional motivation:

the tension regarding CP violation in the B_s and B_d
systems

disagreement with SM at the level of few sigma, but going
up and down

UTfit: Bona et al '08

could be relieved in LR theory - needs more work

it would correlate the W_R and new Higgs masses

parallel talks by Nemevšek, Nesti

Simple predictive GUT with measurable seesaw?

What is the minimal realistic GUT?

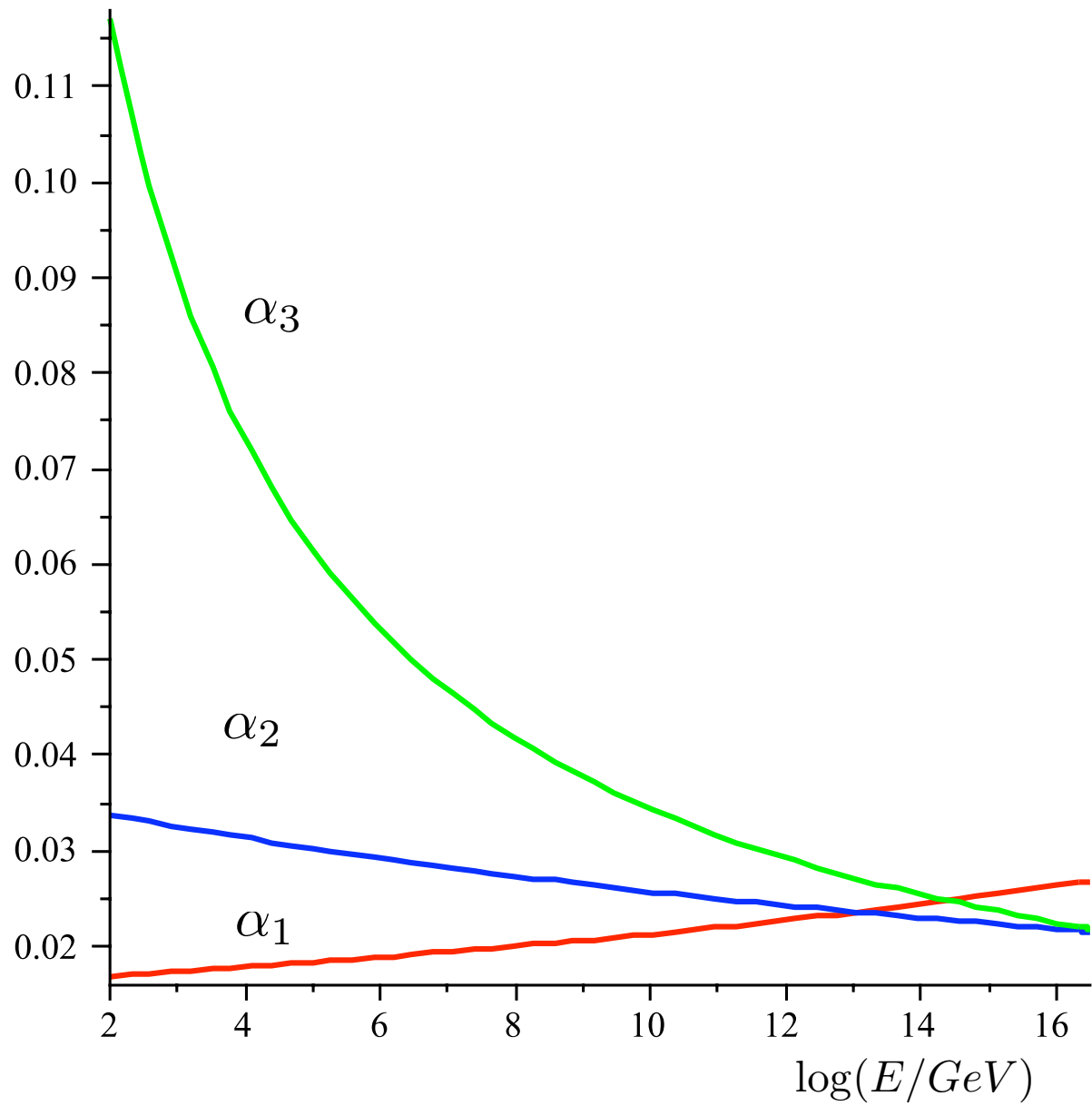
Minimal SU(5)

Georgi-Glashow model: $24_H + 5_H$
 $3(10_F + \bar{5}_F)$

ugly: asymmetric matter, fine-tuning
- but predictive

RULED OUT

- 1. gauge couplings do not unify**
 - 2 and 3 meet at 10^{16} GeV (as in susy),
 - but 1 meets 2 too early at $\approx 10^{13}$ GeV
- 2. neutrinos massless** (as in the SM)



Add just one **extra fermionic** 24_F

Bajc, G.S., '06 ; Bajc, Nemevsek, G.S., '07

maintains the ugliness of the minimal model

asymmetric matter,
even more fine tuning,

\Rightarrow but also its predictivity

$SU(3)_C \times SU(2)_L \times U(1)_Y$ decomposition

$$24_F = (1, 1)_0 + (1, 3)_0 + (8, 1)_0 + (3, 2)_{5/6} + (\bar{3}, 2)_{-5/6}$$

singlet S triplet T

$$\mathcal{L}_{Y\nu} = L_i (y_T^i T + y_S^i S) H + h.c.$$

Mixed Type I and Type III seesaw:

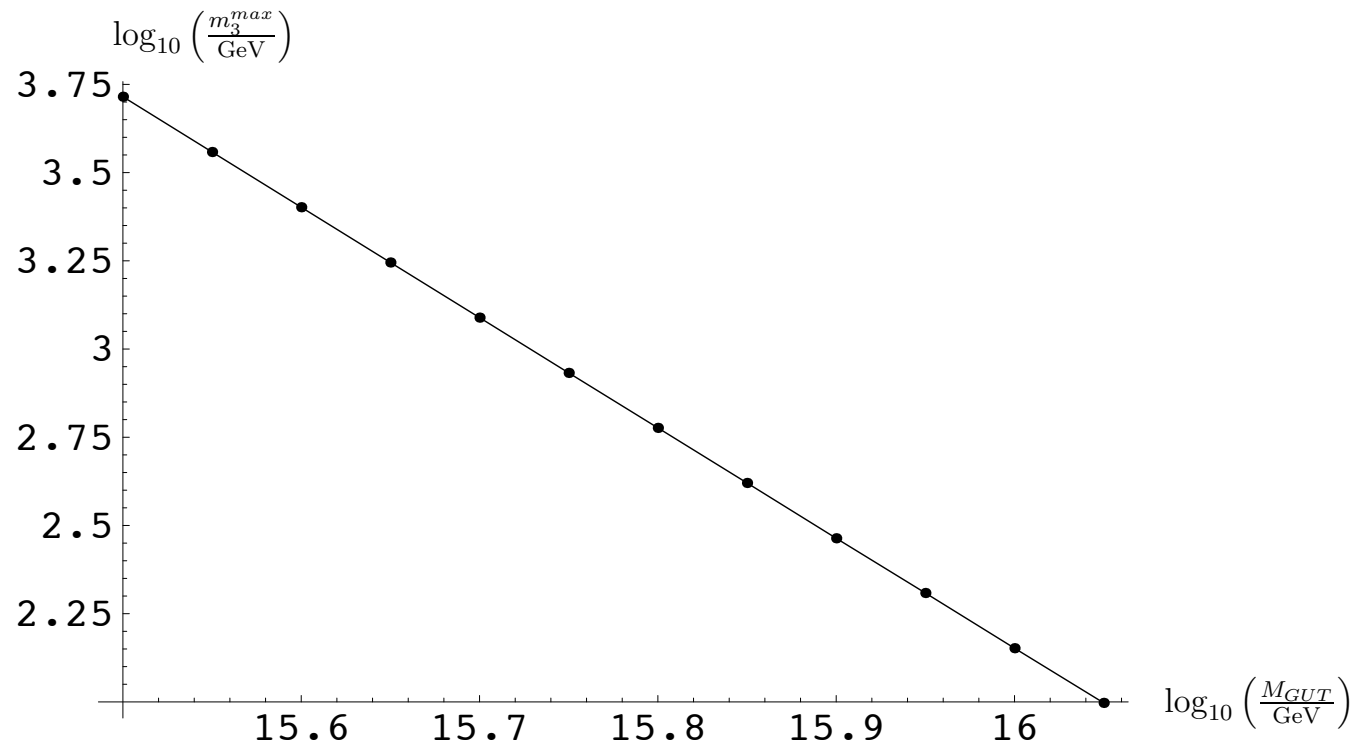
$$M_\nu^{ij} = v^2 \left(\frac{y_T^i y_T^j}{m_T} + \frac{y_S^i y_S^j}{m_S} \right)$$

\Rightarrow one massless neutrino \rightarrow hierarchical spectrum

$$M_{GUT} \gtrsim 10^{15.5} \text{ GeV (p decay)}$$

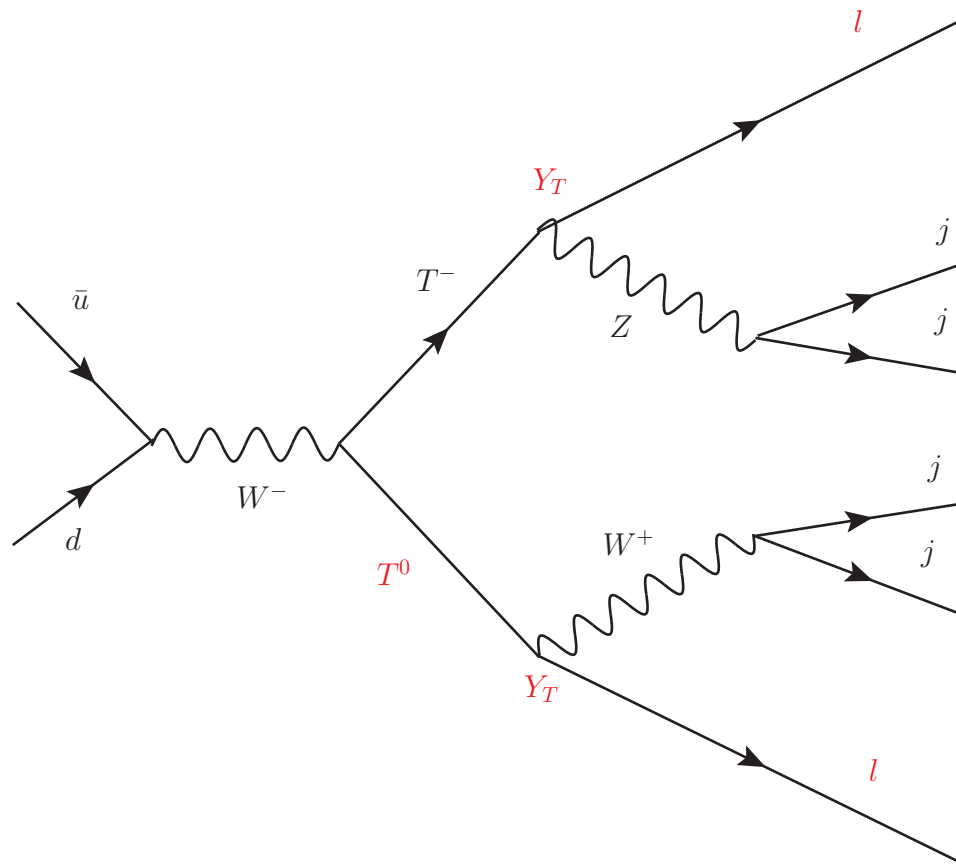
$$\Rightarrow m_3 \lesssim 1 \text{ TeV}$$

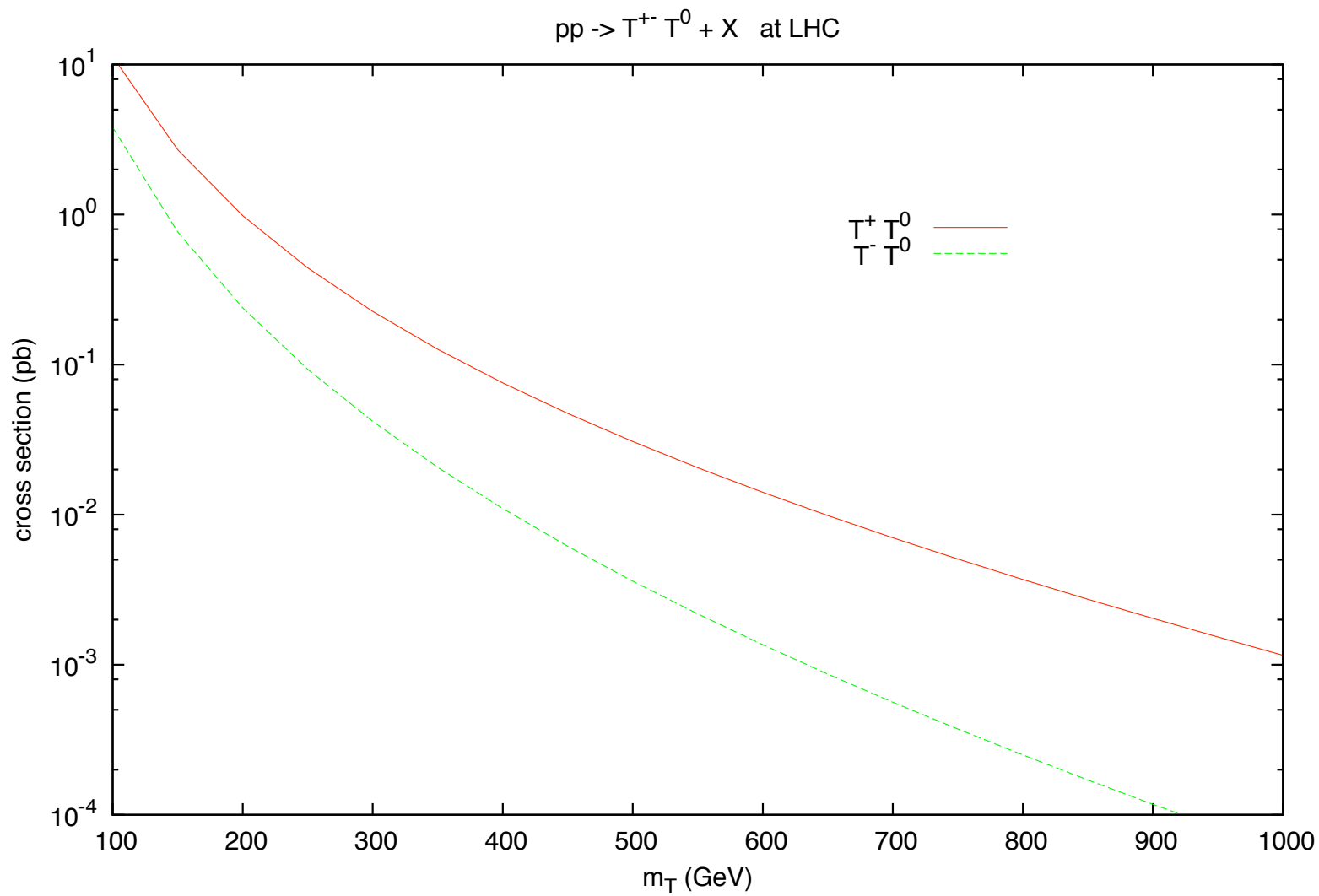
$m_3^{max} - M_{GUT}$ at two loops



Seesaw at LHC

(T^+, T^0, T^-) : weak triplet





Probing neutrino parameters

$$\Gamma_T \approx m_T |y_T|^2 \qquad \tau_T \leq \left(\frac{200 \text{ GeV}}{M_T} \right)^2 0.5 \text{ mm}$$

The best channel is like-sign dileptons + jets

$$BR(T^\pm T^0 \rightarrow l_i^\pm l_j^\pm + 4 \text{ jets}) \approx \frac{1}{20} \times \frac{|y_T^i|^2 |y_T^j|^2}{(\sum_k |y_T^k|^2)^2}$$

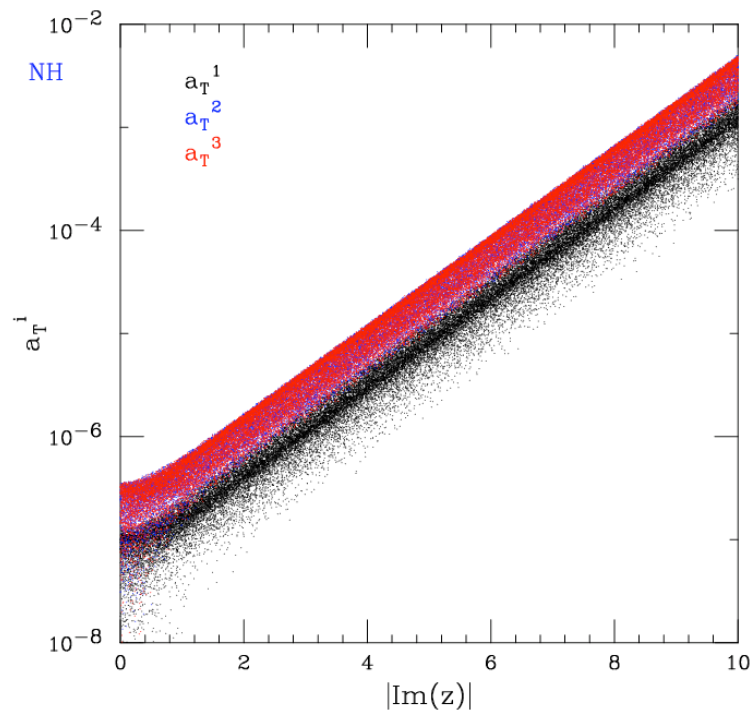
Same couplings y_T^i contribute to

ν mass matrix and T decays

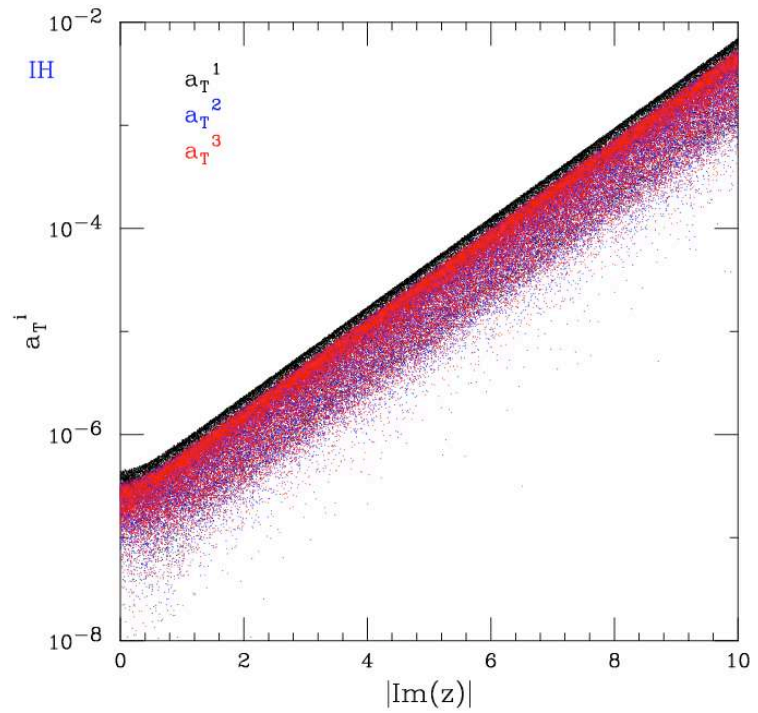
$$vy_T^{i*} = \begin{cases} i\sqrt{M_T} (U_{i2}\sqrt{m_2^\nu} \cos z + U_{i3}\sqrt{m_3^\nu} \sin z), & \text{NH} \quad (m_1^\nu = 0), \\ i\sqrt{M_T} (U_{i1}\sqrt{m_1^\nu} \cos z + U_{i2}\sqrt{m_2^\nu} \sin z), & \text{IH} \quad (m_3^\nu = 0). \end{cases}$$

Casas, Ibarra '01; Ibarra, Ross '03

Scanning over whole parameter space



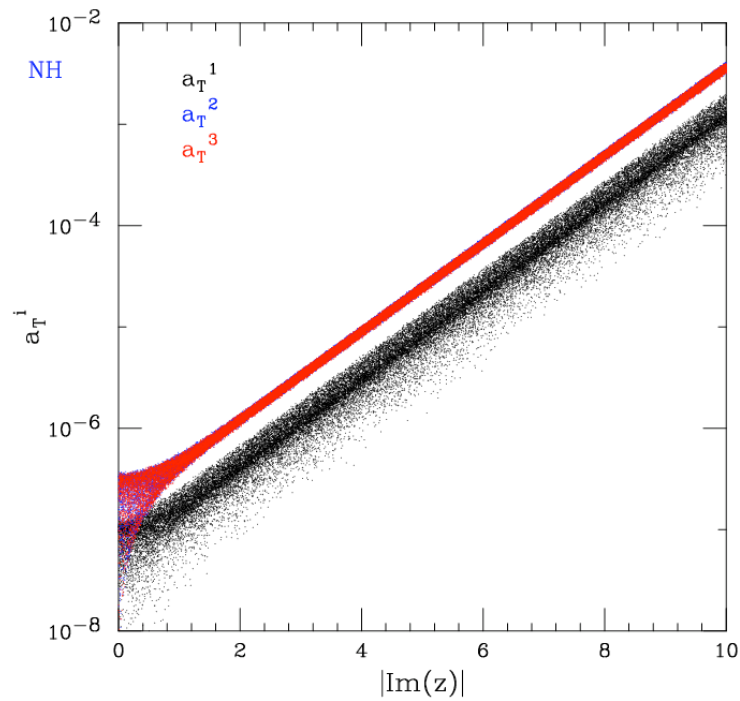
normal hierarchy



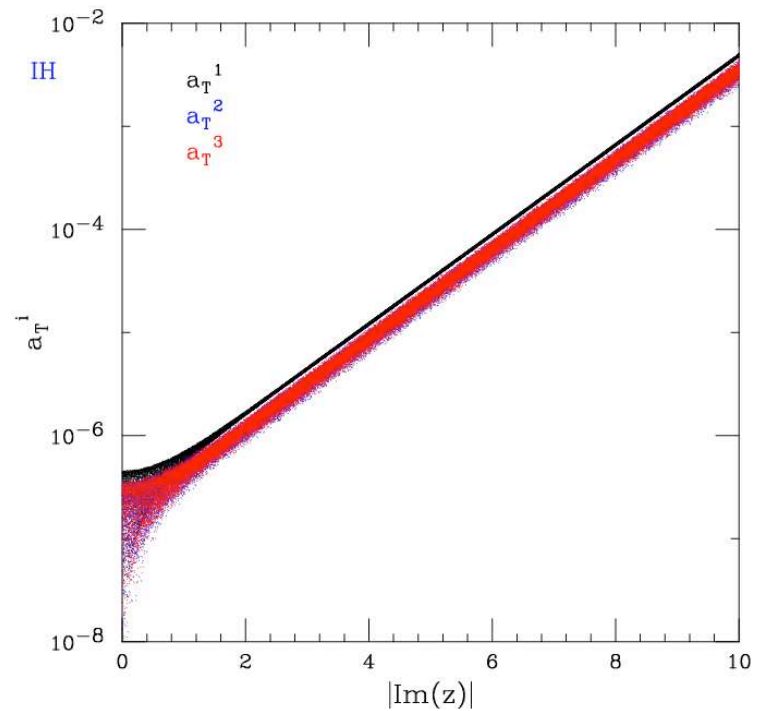
inverse hierarchy

$$a_T^i \equiv |y_T^i| \sqrt{\frac{v}{2M_T}}$$

Assuming Majorana phase $\Phi = 0$



normal hierarchy



inverse hierarchy

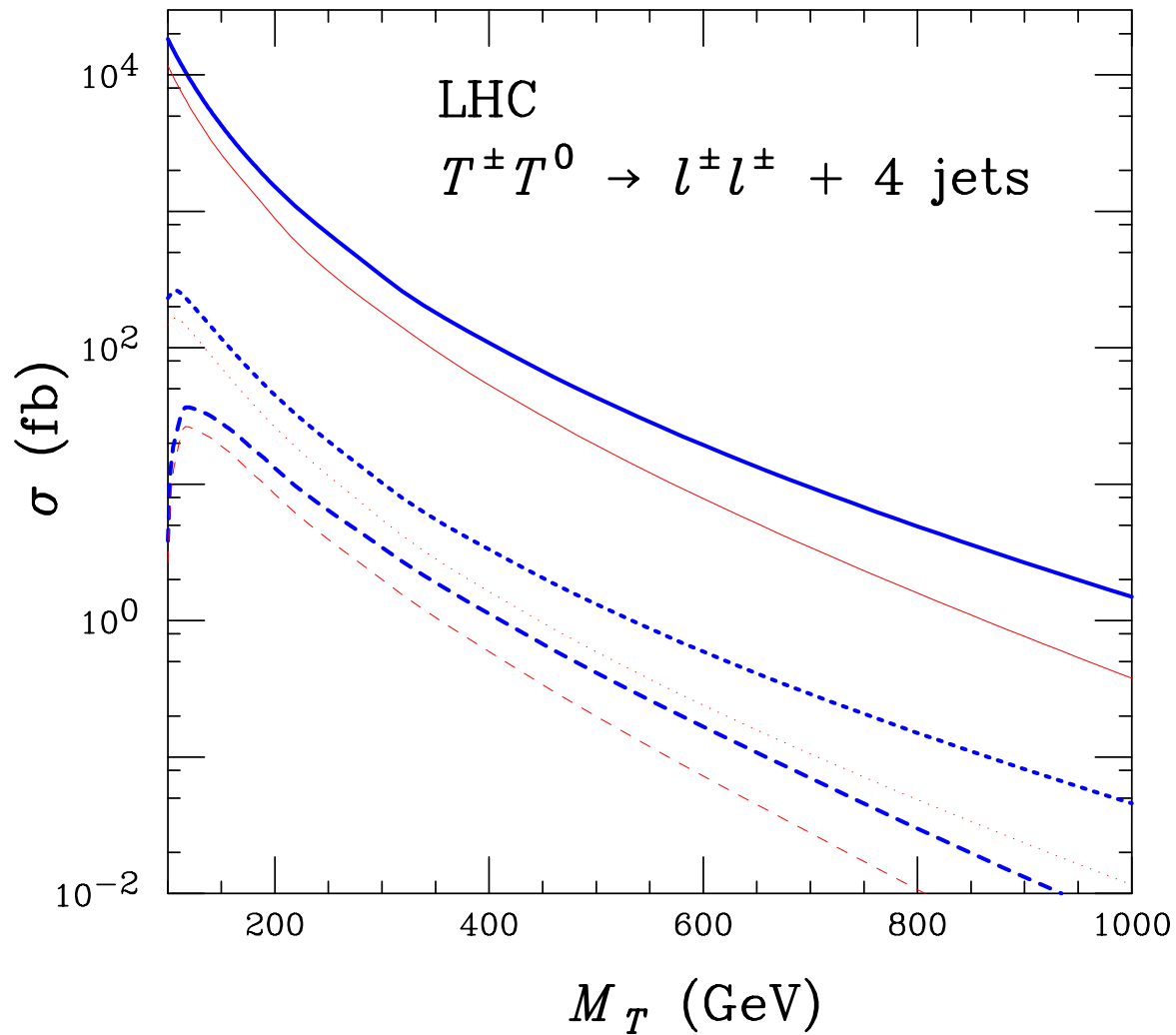
$$a_T^i \equiv |y_T^i| \sqrt{\frac{v}{2M_T}}$$

Incremental increase of cuts on the signal ($M_T = 200$ GeV):

Cuts ↓	$\sigma_{\text{sig.}}$ (fb)
$p_T(\ell) > 15(\text{GeV})$	35
$p_T(jets) > 20$ (GeV)	20
$ \eta(\ell) < 2.5$	16
$ \eta(jets) < 3$	14
$\Delta R_{\ell\ell} > 0.3$	13.8
$\Delta R_{\ell j} > 0.5$	12
$\Delta R_{jj} > 0.5$	10

Francheschini, Hambye, Strumia '08; del Aguila, Aguilar-Saveedra '08

Arhrib, Bajc, Ghosh, Han, Huang, Puljak, G.S. '09 '10



14 TeV 10 TeV

Dotted:
 branching fraction.

Dashed:
 selection cuts.

Conclusions

- experimental probe of Majorana neutrino mass origin
lepton number violation at LHC – a high energy analogue of neutrino-less double beta decay
- **L-R theory**: possible discovery of $W_R \rightarrow 4 \text{TeV}$
 $\nu_R \rightarrow \text{TeV}$ through LR restoration and L violation
- explicit example of a predictive GUT theory
SU(5) with a fermionic adjoint
- **weak fermionic triplet predicted in the TeV range**
LHC@14 TeV \rightarrow 450 (700) GeV with $L = 10(100) \text{ft}^{-1}$
- possible to get information on unmeasured neutrino parameters

R measures separations

$$R = [(\Delta\phi)^2 + (\Delta\eta)^2]^{1/2}$$

where $\Delta\phi$ and $\Delta\eta$ are the azimuthal angular separation and (pseudo) rapidity difference between two particles