

Generating Fermion Masses in Emergent EWSB

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What is the origin of mass?

I. Higgs mechanism

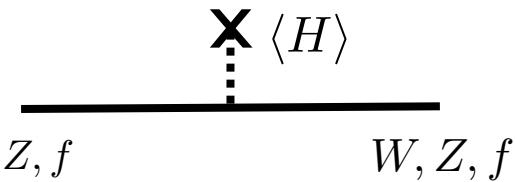
[Guralnik, Hagen, Kibble '64; Englert, Brout '64; Higgs '64]

Higgs boson: $\langle H \rangle$ vacuum expectation value

- Elementary fermion and W, Z boson masses

W, Z-boson: $m_{W,Z} \propto g\langle H \rangle$

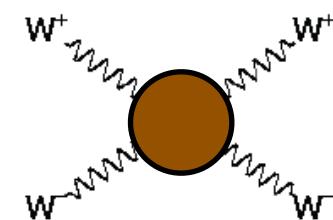
Fermion: $m_f \propto \lambda\langle H \rangle$



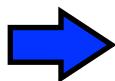
- WW scattering

Bad high-energy behaviour

$$\mathcal{A}(E) \stackrel{E \rightarrow \infty}{\sim} \frac{g^2 E^2}{32\pi m_W^2}$$



But, spin-0 Higgs boson



$$\mathcal{A}(E) \stackrel{E \rightarrow \infty}{\sim} \text{constant}$$

2. Strong dynamics

e.g. QCD: strong coupling at $\Lambda_{QCD} = e^{-\frac{8\pi^2}{g^2 b_i}} M_P$

- Hadron mass spectrum

proton: $m_P \propto \Lambda_{QCD}$

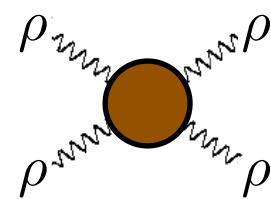
vector-mesons:

e.g. SU(2) isospin-triplet $\rho^{0,\pm}$ $m_\rho \propto \Lambda_{QCD}$

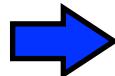
isospin SU(2)  global symmetry

- No unitarity violation

$\rho\rho$ scattering



chiral Lagrangian



QCD Lagrangian

Higgs boson not yet seen...maybe soon...

...or are there surprises in store?

Question:

Can one generate mass in the Standard Model without the Higgs mechanism?

Early work: Bjorken, 1977; Hung, Sakurai 1978

Abbot, Farhi 1981; Fritzsch, Schildknecht, Kogerler, 1982

Model building limited due to nonperturbative nature

New approach:

Use AdS/CFT correspondence!

[Maldacena, 97; Gubser, Klebanov, Polyakov, 98; Witten 98]

Effective 4D chiral
Lagrangian of massive
 W, Z , fermions



5D Lagrangian in
warped dimension!

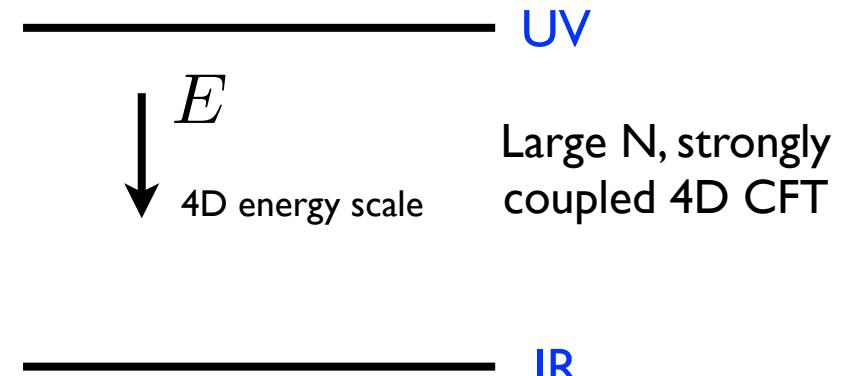
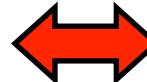
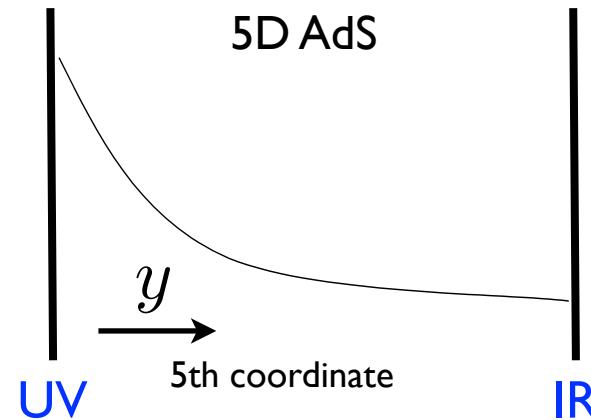
[Randall, Sundrum 99]

Challenges:

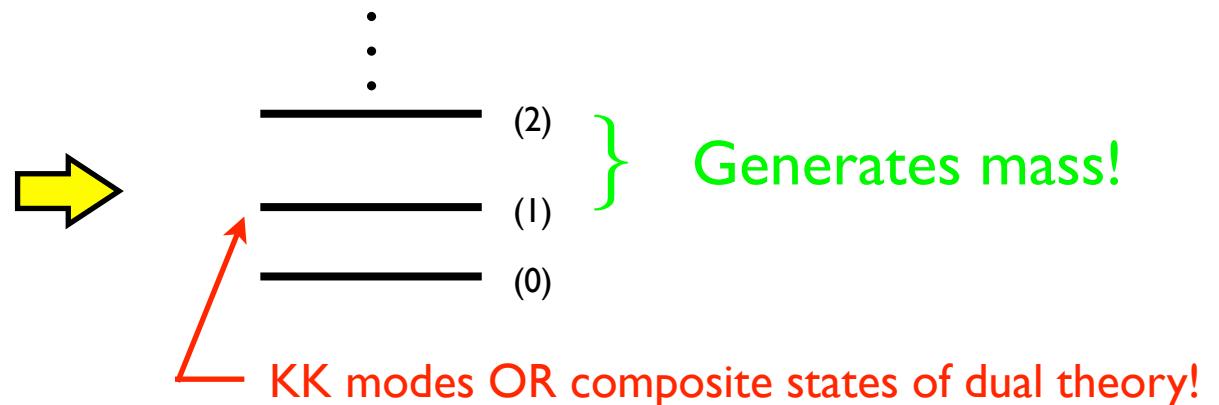
- Generate fermion and W, Z boson masses
- Universality of gauge couplings
- Consistent with EW precision tests
- Natural, no fine tuning.....

AdS/CFT dictionary

[Arkani-Hamed, Randall, Poratti 00; Rattazzi, Zaffaroni 00; Perez-Victoria 01]

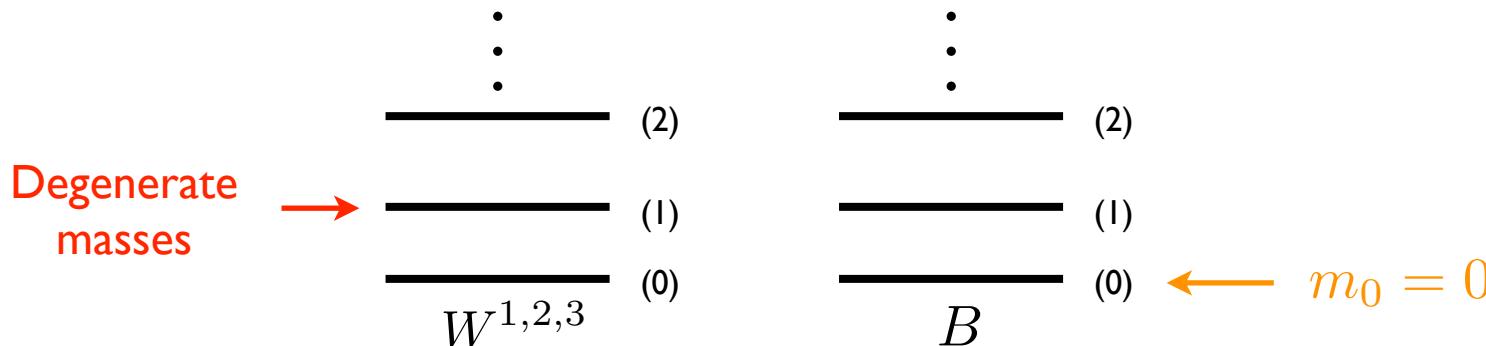


IR brane breaks
conformal symmetry



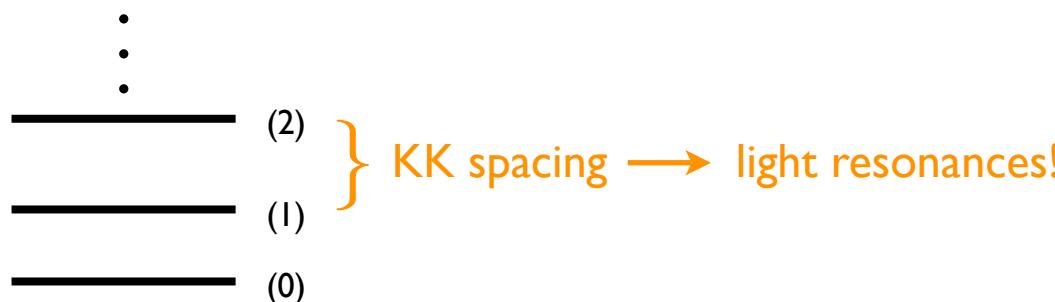
But,

- Massless mode



→ Break EW symmetry at Planck scale
to decouple massless modes!

- Heavy KK modes

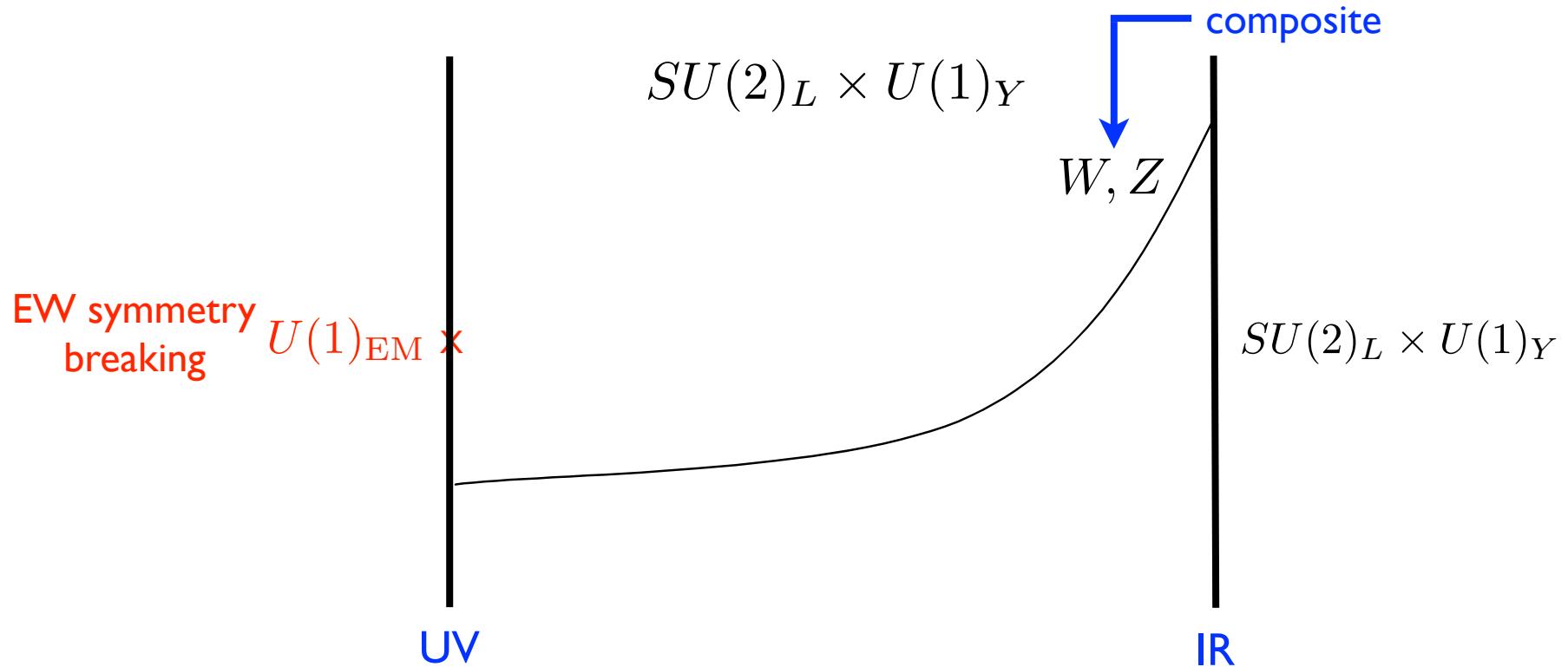


→ Separate lightest KK mode from rest of tower
with brane kinetic terms!

[Carena, Ponton, Tait, Wagner 2002; Davoudiasl, Hewett, Rizzo 2002]

5D Model

[Cui,TG,Wells, arXiv:0907.0906]



5D action:

$$S = \int d^4x dz \sqrt{-g} \left[-\frac{1}{4}(F_{MN}^{La})^2 - \frac{1}{4}(F_{MN}^Y)^2 - \frac{1}{2}(kz)\delta(z - z_{UV}) \frac{\zeta_Q}{g_{Y5}^2 + g_{L5}^2} (g_{Y5} F_{\mu\nu}^{L3} + g_{L5} F_{\mu\nu}^Y)^2 - \frac{1}{2}(kz)\delta(z - z_{IR}) (\zeta_L (F_{\mu\nu}^{La})^2 + \zeta_Y (F_{\mu\nu}^Y)^2) \right]$$

$\zeta_Q, \zeta_L, \zeta_Y$ = boundary kinetic term coefficients

Mass spectrum:

Boundary conditions:

$$z = z_{UV} : \begin{cases} \partial_z(g_{Y5}A_\mu^{L3} + g_{L5}B_\mu) + \zeta_Q \square(g_{Y5}A_\mu^{L3} + g_{L5}B_\mu) = 0, \\ g_{L5}A_\mu^{L3} - g_{Y5}B_\mu = 0, \\ A_\mu^{L1,2} = 0, \end{cases}$$

$$z = z_{IR} : \begin{cases} \partial_z A_\mu^{La} - \zeta_L k z_{IR} \square A_\mu^{La} = 0, \\ \partial_z B_\mu - \zeta_Y k z_{IR} \square B_\mu = 0, \end{cases}$$

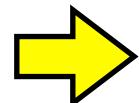
Obtain:

$$m_\gamma = 0$$

$$m_W \simeq \sqrt{\frac{2}{\zeta_L k}} m_{IR}$$

$$m_Z \simeq \sqrt{\frac{2}{\zeta_L k} + \frac{2}{\zeta_Q k(1 + g_{L5}^2/g_{Y5}^2)}} m_{IR}$$

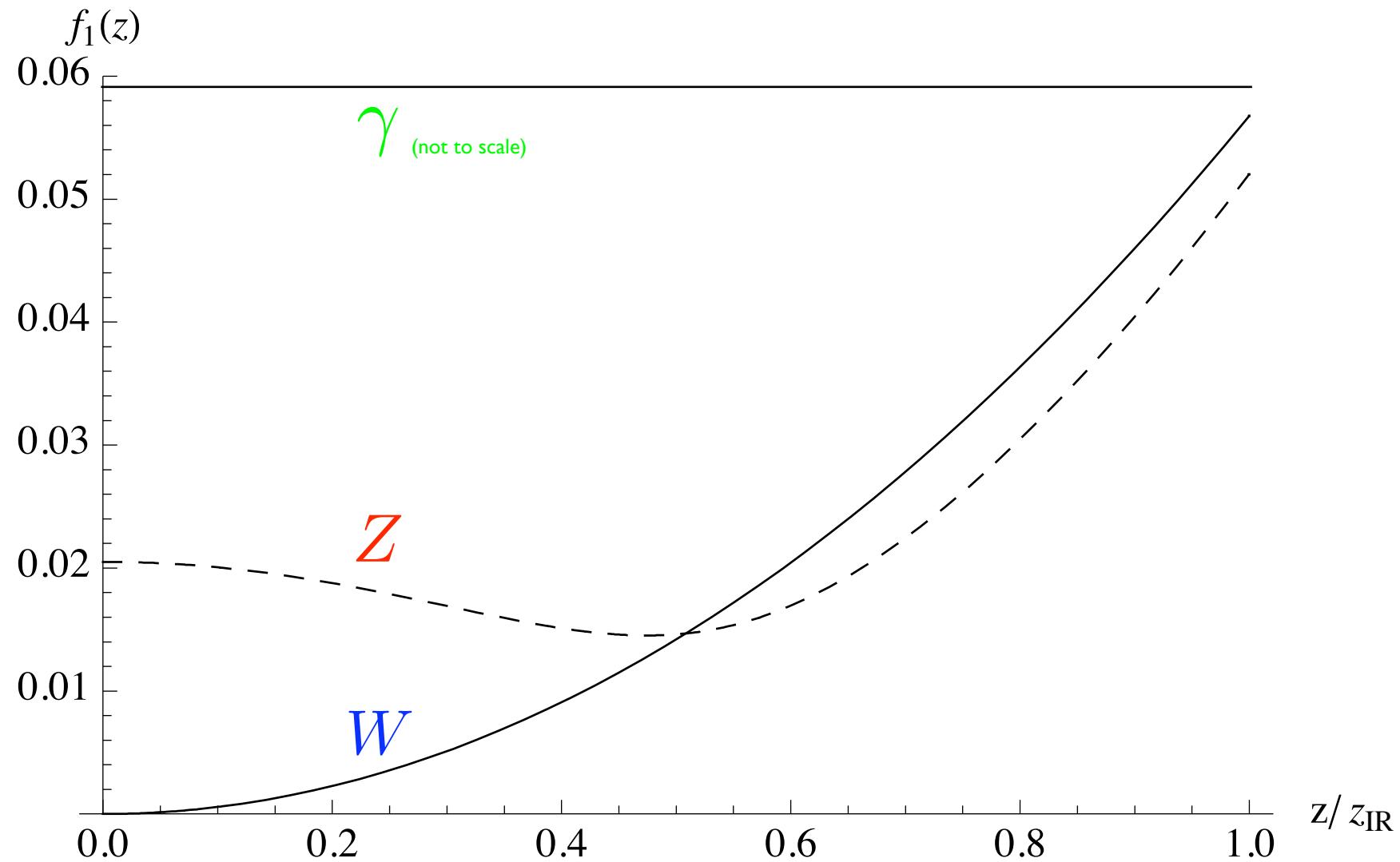
For: $m_{IR} = \text{TeV}$ $\zeta_Q k \simeq 500, \zeta_L k \simeq 310, \zeta_Y k \simeq 0.1$



$$m_W \simeq 80.4 \text{ GeV}, \quad m_Z \simeq 91.2 \text{ GeV}$$

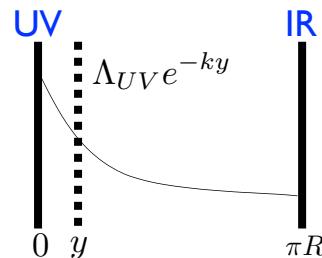
$$(m_{KK} \gtrsim 2 \text{ TeV})$$

W, Z , photon profiles

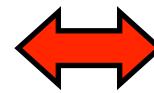


NB: Boundary kinetic terms introduce discontinuity at endpoints.

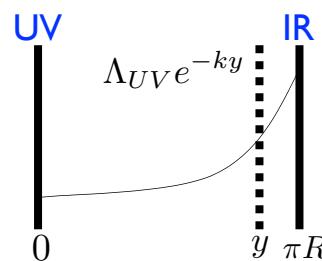
Dual 4D interpretation



UV localized field



elementary “source” state

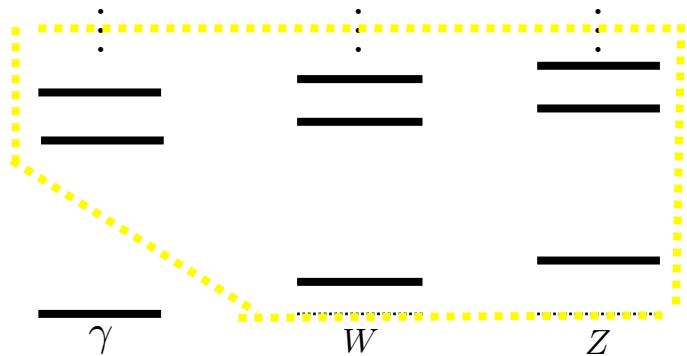


IR localized field

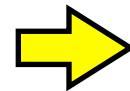


CFT bound state

∴ Composite W, Z but elementary photon!



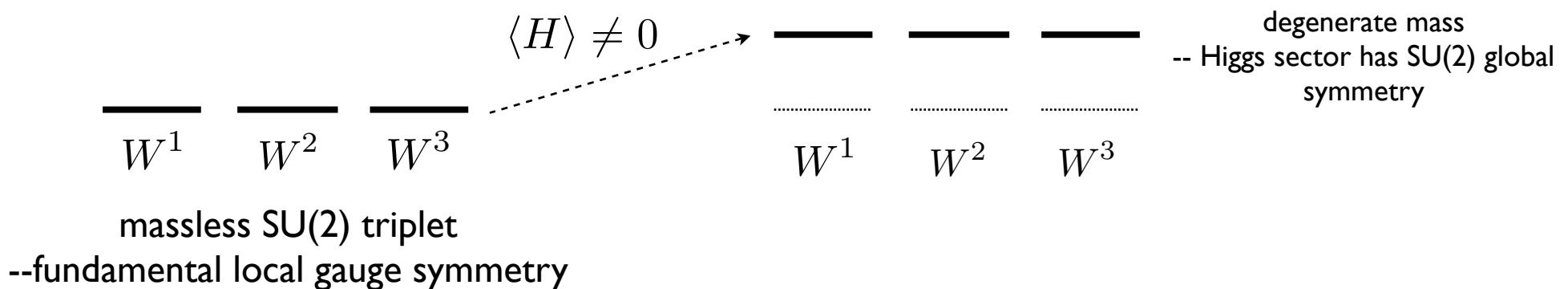
EWSB emerges at IR scale



“Emergent” EWSB

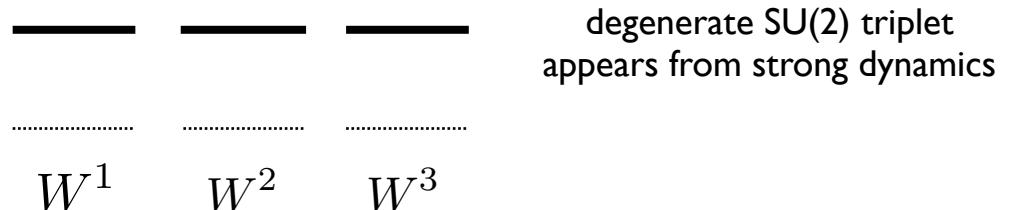
Custodial Symmetry

- Higgs Mechanism in SM



- Emergent EWSB

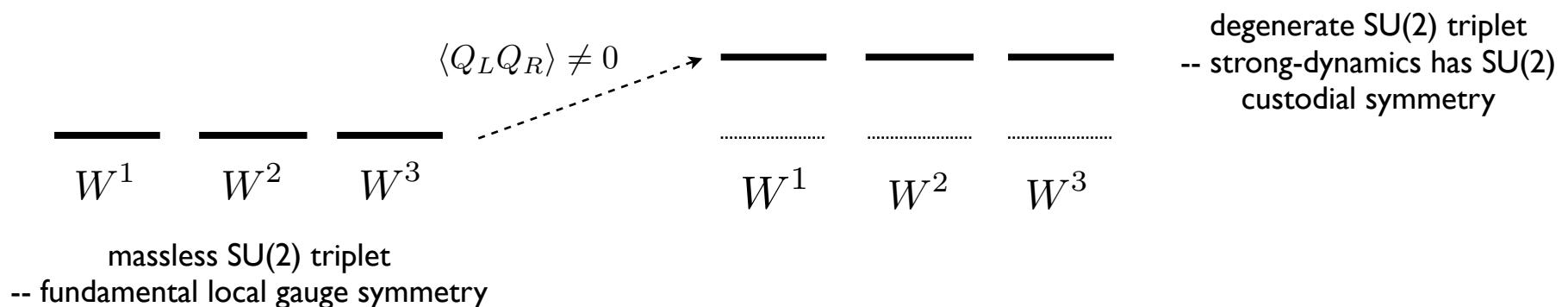
SU(2) symmetry is
NOT fundamental



Emergent EWSB is Higgsless... but no Higgs mechanism!

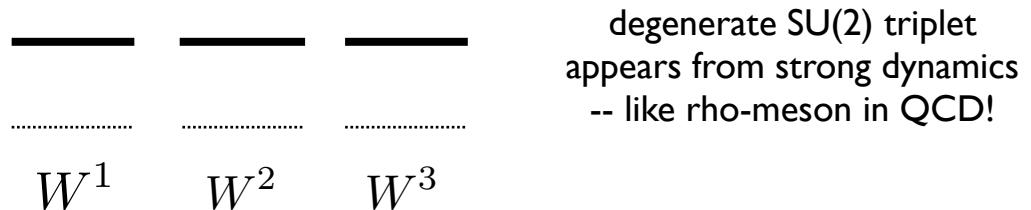
- Usual ``Higgsless'' (technicolor-like)

[Csaki, Grojean, Pilo, Terning 04]



- Emergent EWSB

No Higgs mechanism



Electroweak constraints

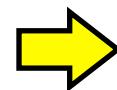
Assume fermions predominantly on IR brane

- T parameter Custodial symmetry in limit $\zeta_Y \rightarrow 0, \zeta_Q \rightarrow \infty$

i.e. same boundary condition for $A^{L1,2,3}$

- S parameter $S \simeq \frac{8\pi}{g^2 + g'^2} \cos 2\theta_w \sin^2 \theta_w (1 + \beta^2) (m_Z z_{IR})^2$

$$\begin{aligned}\zeta_L k &\simeq 1000, \zeta_Q k \simeq 1700, \zeta_Y k \simeq 0.2 \\ m_{IR} &\simeq 1.8 \text{ TeV}\end{aligned}$$



$$S \simeq 0.1, \quad T \simeq 0.02$$

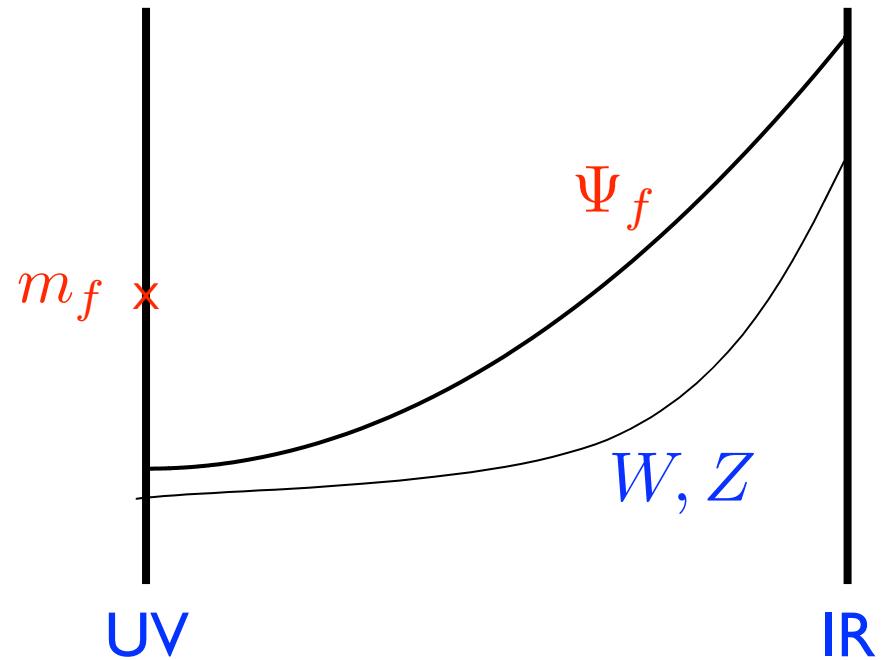
What about fermion masses?

[Cui, TG, Stokes, to appear]

Assume universal bulk fermion profile

- Add UV boundary fermion masses

Fermion mass hierarchy via
Froggatt-Nielsen mechanism



$$S_\psi = i \int d^5x \sqrt{-g} \left[\frac{1}{2} (\bar{\Psi}_i^{(L)} \Gamma^M D_M \Psi_i^{(L)} - D_M \bar{\Psi}_i^{(L)} \Gamma^M \Psi_i^{(L)}) + \underbrace{m_L^{(i)} \bar{\Psi}_i^{(L)} \Psi_i^{(L)}}_{= c k} + (L \leftrightarrow R) \right]$$

$$S_m^{(UV)} = i \int d^5x \sqrt{-g} \lambda_5^{(i)} \left[\bar{\Psi}_i^{(L)} \Psi_i^{(R)} + \bar{\Psi}_i^{(R)} \Psi_i^{(L)} \right] (kz) \delta(z - z_{UV})$$

$$S_{KE}^{(IR)} = i \int d^5x \sqrt{-g} \left[\frac{1}{2} \eta_{iL} (\bar{\Psi}_i^{(L)} \Gamma^\mu D_\mu \Psi_i^{(L)} - D_\mu \bar{\Psi}_i^{(L)} \Gamma^\mu \Psi_i^{(L)}) + (L \leftrightarrow R) \right] (kz) \delta(z - z_{IR})$$

Example: massless bulk fermions (c=0)

$$|f_{L-}^{(n)}(z)| = |f_{R+}^{(n)}(z)| = N_n^{(0)} (kz)^2 [\cos(\hat{m}_n - m_n z) - (\eta k) \hat{m}_n \sin(\hat{m}_n - m_n z)]$$

where $N_n^{(0)} \simeq \frac{1}{\sqrt{z_{\text{IR}}}} \sqrt{\frac{1}{1 + (\eta k)/2 + (\eta k)^2 \hat{m}_n^2}}$ $\hat{m} = \frac{m}{m_{IR}}$

Obtain for ($\eta k = 10$)

$$m_e \leq m \leq m_t \quad \text{with} \quad 3.1 \times 10^{-6} \leq \lambda_5 \leq 1.15$$

W,Z boson couplings

Determined by wavefunction overlap:

e.g. $g_W \simeq \frac{g_{5L}}{\sqrt{\zeta_L}} \frac{2}{\eta k} \left[\frac{1}{3} + \xi_{IR} k (1 - 2(\eta k) \hat{m}_i \hat{m}_j) \right]$



$$\hat{m} = \frac{m}{m_{IR}} \ll 1$$

$$\xi_{IR} = \sqrt{\frac{\zeta_L}{2}} \Lambda \eta$$

→ { Light fermions: nonuniversality at the per-mille level!
3rd generation: nonuniversality at 15%-25% level

$$\frac{g_{W-}(\text{tb})}{g_{W-}(\text{ud})} = 0.854$$

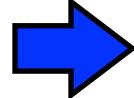
$$\frac{g_{Z-}(\text{top})}{g_{Z-}^{(SM)}(\text{top})} = 0.746$$

$$\frac{g_{Z+}(\text{top})}{g_{Z+}^{(SM)}(\text{top})} = 0.745$$

Wtb: 20% level @Tevatron arXiv:0903.0850 Single top production

Ztt: 40% level @LHC with 300 fb⁻¹

Gauge coupling universality due to light fermion masses!

Anomalous top couplings  electroweak corrections

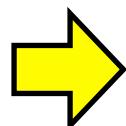
[Larios, Perez, Yuan '99]

Obtain:

$$\epsilon_1^{SM} + \delta\epsilon_1 \simeq 19 \times 10^{-3}; \quad \epsilon_b^{SM} + \delta\epsilon_b \simeq -13 \times 10^{-3}$$

This compares with:

$$4.4 \times 10^{-3} \leq \epsilon_1^{exp} \leq 6.4 \times 10^{-3}, \quad 68\% \text{ C.L.}$$
$$-6.2 \times 10^{-3} \leq \epsilon_b^{exp} \leq -3.1 \times 10^{-3} \quad 68\% \text{ C.L.}$$

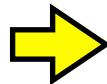


Requires special treatment of top quark

e.g. separate brane for top [Cacciapaglia, Csaki, Grojean, Reece, Terning '05]

WW scattering

Composite W,Z boson



Momentum dependent
form factor

$$F_{WWZ}(q^2) = \frac{1}{N_Z(q^2)N_W^2} \left\{ \left[\int_{z_{UV}}^{z_{IR}} \frac{dz}{kz} f^{L3}(q^2, z) (f_W(z))^2 \right] + \zeta_L f^{L3}(q^2, z_{IR}) (f_W(z_{IR}))^2 \right\}$$

→ Possible deviation in W, Z-boson vertices at LHC
(in progress)

Interestingly, in large N theory there
are no partons inside hadrons!

[Polchinski-Strassler 02]

i.e. composite W, Z bosons are unlike vector-mesons in QCD!

Summary

- Generate W, Z boson and fermion masses from strong dynamics, not Higgs mechanism
- Electroweak symmetry breaking “emerges” at IR scale
- Composite W, Z bosons lead to deviations in couplings at the LHC
- Requires further model-building to successfully incorporate top quark....