

Strong coupling, discrete symmetry and flavour

Planck 2010

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June 3rd 2010



based on S.Abel, JB - arXiv:1005.1668 [hep-ph]

Some observations

Several open questions in SM flavour structure:

- Origin of large hierarchies in masses and quark mixing.
- Origin of tribimaximal mixing in neutrino sector.
- Grand unification?

Attempt to address these problems using two observations:

Observation

Neutrino mixing matrix suggests a **discrete flavour symmetry**.

Observation

Large hierarchies arise naturally in theories with **strong coupling**.

Mixing patterns from discrete symmetry

Observation

Neutrino mixing matrix suggests a **discrete flavour symmetry**.

- Highly symmetrical – **tribimaximal mixing**.
- Many successful models¹ apply a discrete flavour symmetry.
- Examples of groups: A_4 , S_4 , $PSL_2(7)$, $\Delta(27)$...

But

- Need extended Higgs sector with particular **vacuum alignment** to break flavour symmetry.
- Hard to apply to quark sector
⇒ **difficulties with GUTs**.

¹G.Altarelli, F.Feruglio - arXiv:1002.0211 [hep-ph]

Hierarchies from dynamical scales

Observation

Large hierarchies arise naturally in theories with **strong coupling**.

- Dynamical scales generated by renormalisation group flow:

$$\Lambda \sim E e^{-8b\pi^2/g^2(E)}$$

- Λ is dynamical **strong coupling** scale where $g \rightarrow \infty$.
- Exponential dependence \implies naturally large hierarchies.
- Readily applicable to SM mass hierarchies².

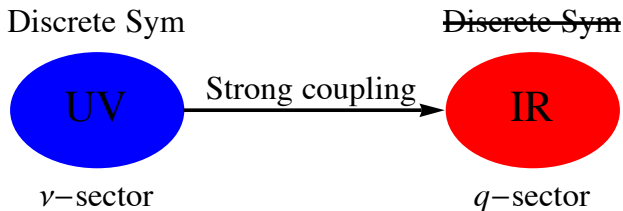
²A.E.Nelson, M.J.Strassler - hep-ph/9607362, hep-ph/0006251

Addressing the flavour problem

Postulate

The flavour problem can be solved using two principles:

- 1 **Discrete flavour symmetry** for mixing.
- 2 **Strong coupling** for **all hierarchies** and flavour symmetry breaking.



(C.F. A.E.Nelson, M.J.Strassler - hep-ph/0006251 with discrete flavour symmetry in UV rather than anarchy)

Comments

- **Multiple** strongly coupled sectors for hierarchies – one per generation.
- Should be a generic feature of strong coupling.

Note:

- Strongly coupled \equiv asymptotically free.
- Will use $\mathcal{N} = 1$ SUSY but ideas should apply to non-SUSY theories.

Yukawa couplings

How does it work?

- Suppose SU(5) $\mathbf{10}_j$ is actually a **bound state** of the strongly coupled gauge group G_j

$$\Lambda_j \mathbf{10}_j = Y_j Y_j$$

- **Elementary** 'Yukawa' couplings are e.g.

$$W_{UV} \supset \frac{\xi_{ij}}{M_X} \bar{\mathbf{5}}_i (Y_j Y_j) \bar{\mathbf{5}}_H$$

G_j s-confines³, match to low energy couplings

$$\frac{\xi_{ij}}{M_X} \bar{\mathbf{5}}_i (Y_j Y_j) \bar{\mathbf{5}}_H = \lambda_{ij} \bar{\mathbf{5}}_i \mathbf{10}_j \bar{\mathbf{5}}_H \quad \Longrightarrow \quad \lambda_{ij} = \frac{\Lambda_j}{M_X} \xi_{ij}$$

³K.A.Intriligator, N.Seiberg - hep-th/9509066

Yukawa couplings

Ten centred model with Z_3 permutation symmetry yields

$$\lambda_{ij} = \frac{\Lambda_j}{M_X} \xi_{ij} = \frac{\Lambda_3}{M_X} \begin{pmatrix} a\varepsilon\eta & b\varepsilon & c \\ c\varepsilon\eta & a\varepsilon & b \\ b\varepsilon\eta & c\varepsilon & a \end{pmatrix} \quad \text{where } \varepsilon, \eta \ll 1$$

- a, b, c are complex, order unity parameters (from Z_3).
- Λ_j are dynamical scales, $\varepsilon = \Lambda_2/\Lambda_3$ and $\eta = \Lambda_1/\Lambda_2$.
- $\Lambda_1 \ll \Lambda_2 \ll \Lambda_3$ natural.
- Dynamical scales produce **hierarchies in quark Yukawas** and **break flavour symmetry**.

Mass and mixing hierarchies

$\mathbf{10} \ni Q, u^c, e^c$ and $\bar{\mathbf{5}} \ni L, d^c$ hence

- $\lambda_{ij}^e \sim \Lambda_j$ for charged lepton mass hierarchy.
- $\lambda_{ij}^u \sim \Lambda_i \Lambda_j$ (from $\mathbf{10}_i \mathbf{10}_j \mathbf{5}_H$ term) for **squared** up quark mass hierarchy.

Result – mass hierarchies

$$\frac{m_e}{m_\mu} \sim \frac{m_d}{m_s} \sim \sqrt{\frac{m_u}{m_c}} \sim \eta \quad \text{and} \quad \frac{m_\mu}{m_\tau} \sim \frac{m_s}{m_b} \sim \sqrt{\frac{m_c}{m_t}} \sim \varepsilon$$

Diagonalise to find CKM matrix:

Result – hierarchies in quark mixing

$$V_{\text{CKM}} \sim \begin{pmatrix} 1 & \eta & \varepsilon\eta \\ \eta & 1 & \varepsilon \\ \varepsilon\eta & \varepsilon & 1 \end{pmatrix}$$

Mass and mixing hierarchies

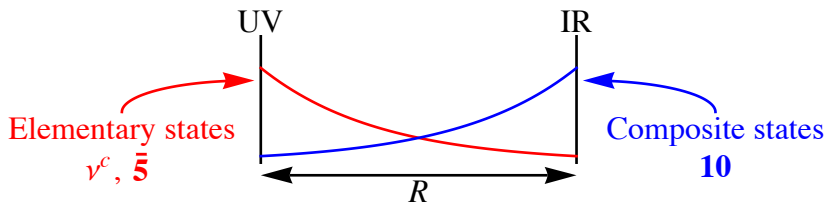
Different story in neutrino⁴ sector:

- Assume **elementary** neutrinos \implies no dynamical scales in Yukawa couplings.
- Neutrinos **do not** see effects of strong coupling.
- Mixing determined **only** by elementary coupling constants.
- Flavour symmetry preserved \implies **tribimaximal mixing**.
- Absence of mass hierarchies in neutrino sector.

⁴C.F. N.Haba - arXiv:hep-ph/9807552

Geometric interpretation

Use AdS/CFT to gain **equivalent** geometrical description. Picture is dual to **5D gravity theory**:

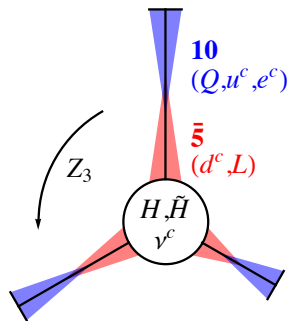


strongly coupled sector \leftrightarrow warped extra dimension
 dynamical scale \leftrightarrow warp factor $M_X e^{-kR}$

Geometric interpretation

Three strongly coupled sectors \implies multithroat background⁵.

- Central UV brane with rotational Z_3 flavour symmetry.
- Strongly coupled sectors in 3 throats.
- Throat lengths generate hierarchies, break flavour symmetry.
- Quarks (in throats) feel hierarchies and flavour symmetry breaking.
- Neutrinos (on UV brane) do not.
- **All mixing** occurs on UV brane.
- **GUT compatible**.

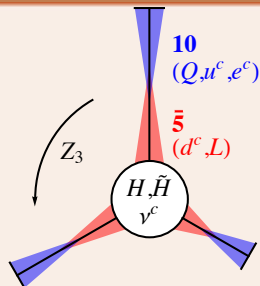


⁵G.Cacciapaglia, C.Csaki, C.Grojean, J.Terning - hep-ph/0604218

Summary

The flavour problem can be addressed using two principles:

- Discrete flavour symmetry for mixing.
- Strong coupling for hierarchies and flavour symmetry breaking.
- Quarks see strong coupling effects, neutrinos do not.
- Approach is GUT compatible.



Simplest examples⁶ use Z_3 permutation symmetry with strong coupling described by:

- s -confinement
- AdS/CFT correspondence

with same results.

⁶S.Abel, JB - arXiv:1005.1668 [hep-ph]