High precision for vector-boson pair production at the LHC

Marius Wiesemann

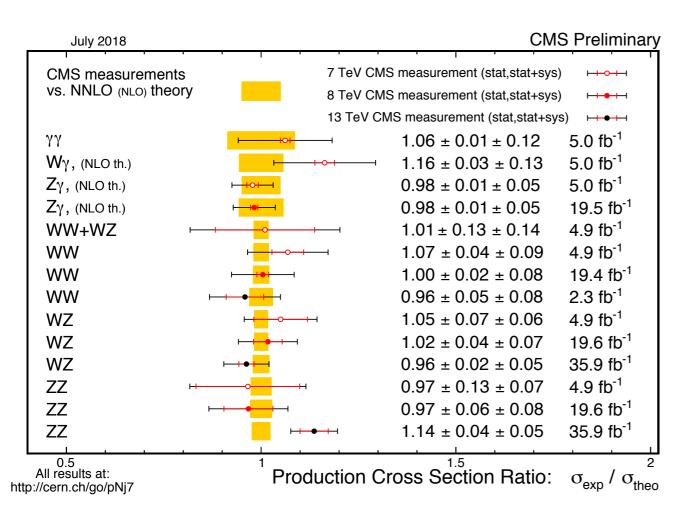


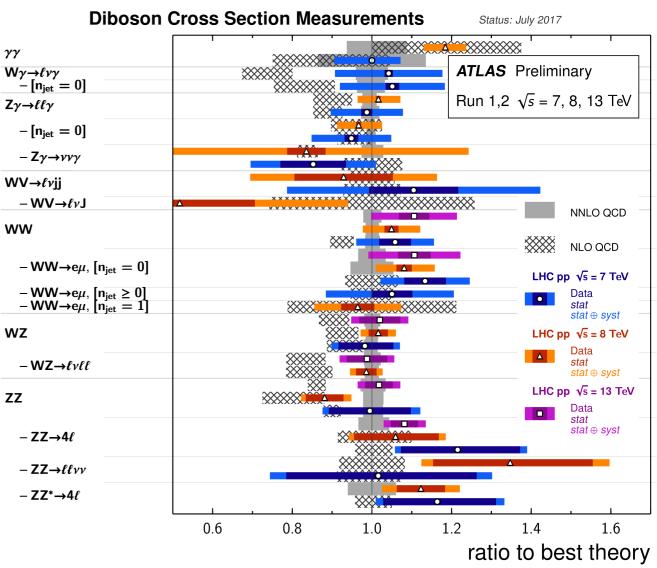
Standard Model at the LHC 2019

Zurich (Switzerland), April 23th-26th, 2019

Measured diboson cross sections







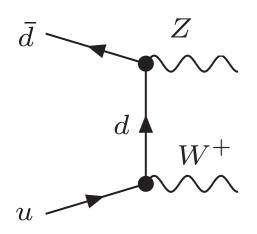
→ talks by Evgenii Baldin and Pietro Vischia...

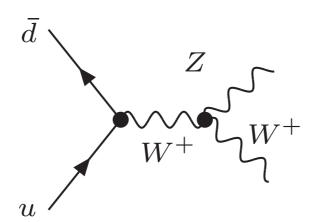


example: WZ production



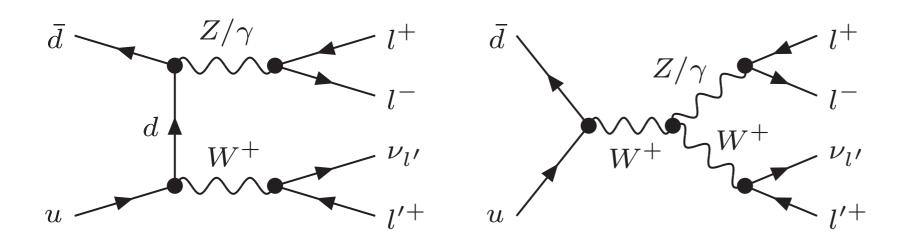
example: WZ production (on-shell)







example: WZ production (off-shell)

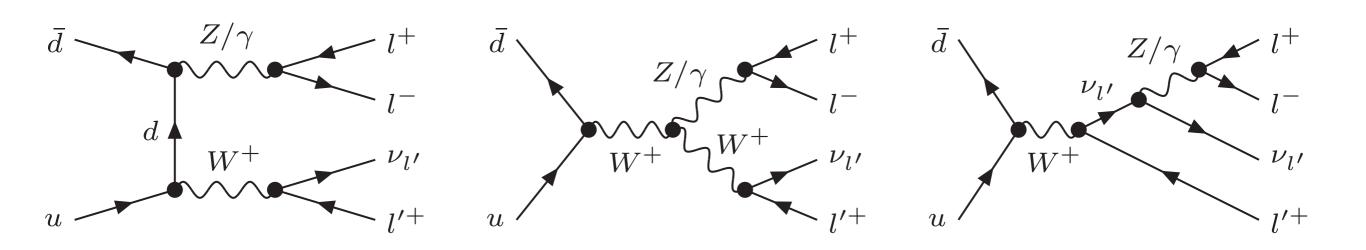


EW decays of heavy bosons (W, Z, γ^*) \checkmark (only isolated photons in the final state)





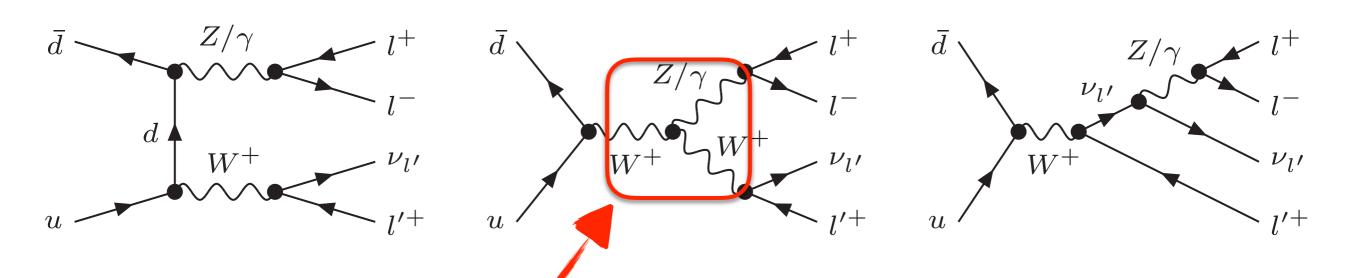
example: WZ production (off-shell)



- [®] EW decays of heavy bosons (W, Z, γ^*) \checkmark (only isolated photons in the final state)
- $^{ ilde{p}}$ all topologies to same leptonic final state (with spin correlations & off-shell effects) \checkmark



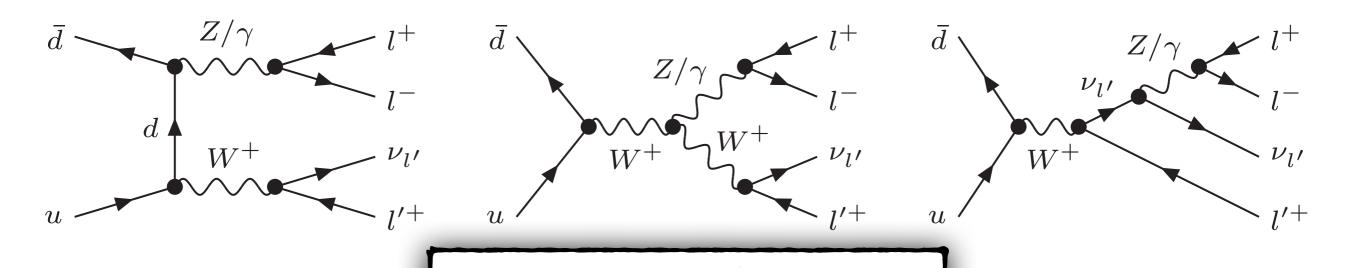
example: WZ production (off-shell)



- EW decays of heavy bosons (W, Z, γ^*) (only isolated photons in the final state)
- $^{ ext{@}}$ all topologies to same leptonic final state (with spin correlations & off-shell effects) \checkmark
 - \rightarrow access to triple gauge couplings (TGCs) \rightarrow high relevance for BSM physics



example: WZ production (off-shell)

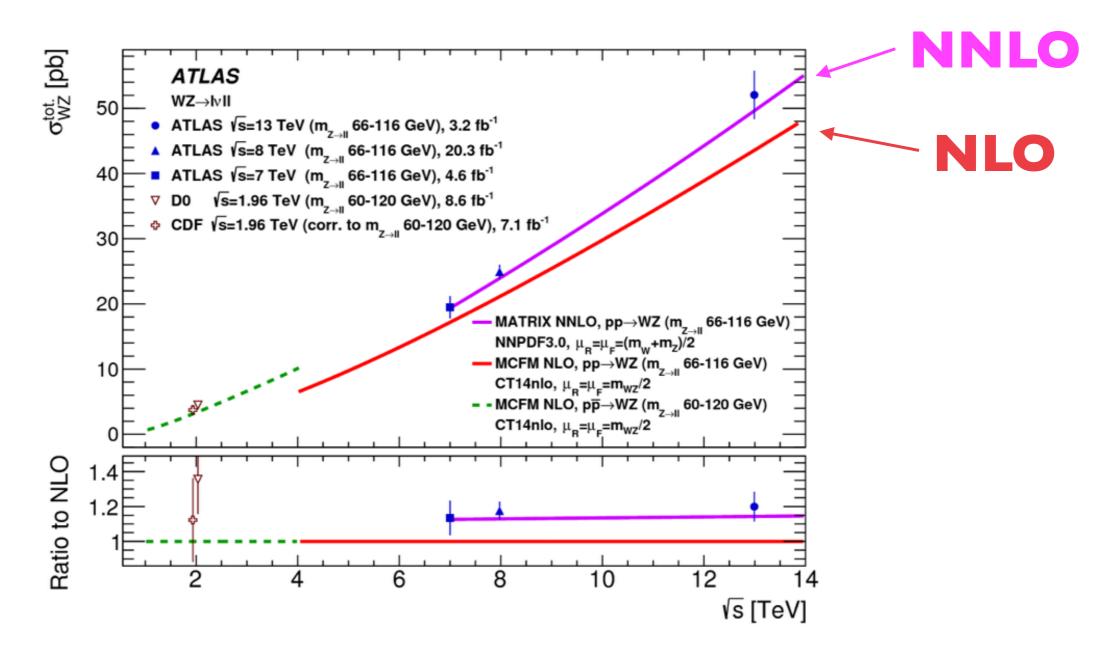


- EW decays of heavy bose
- all topologies to same le
 - → access to triple gauge couplings mgn relevance for BSM physics
- loop-induced gg channel enters NNLO for charge-neutral processes \checkmark (eg, for ZZ)

photons in the final state)

tions & off-shell effects) 🗸

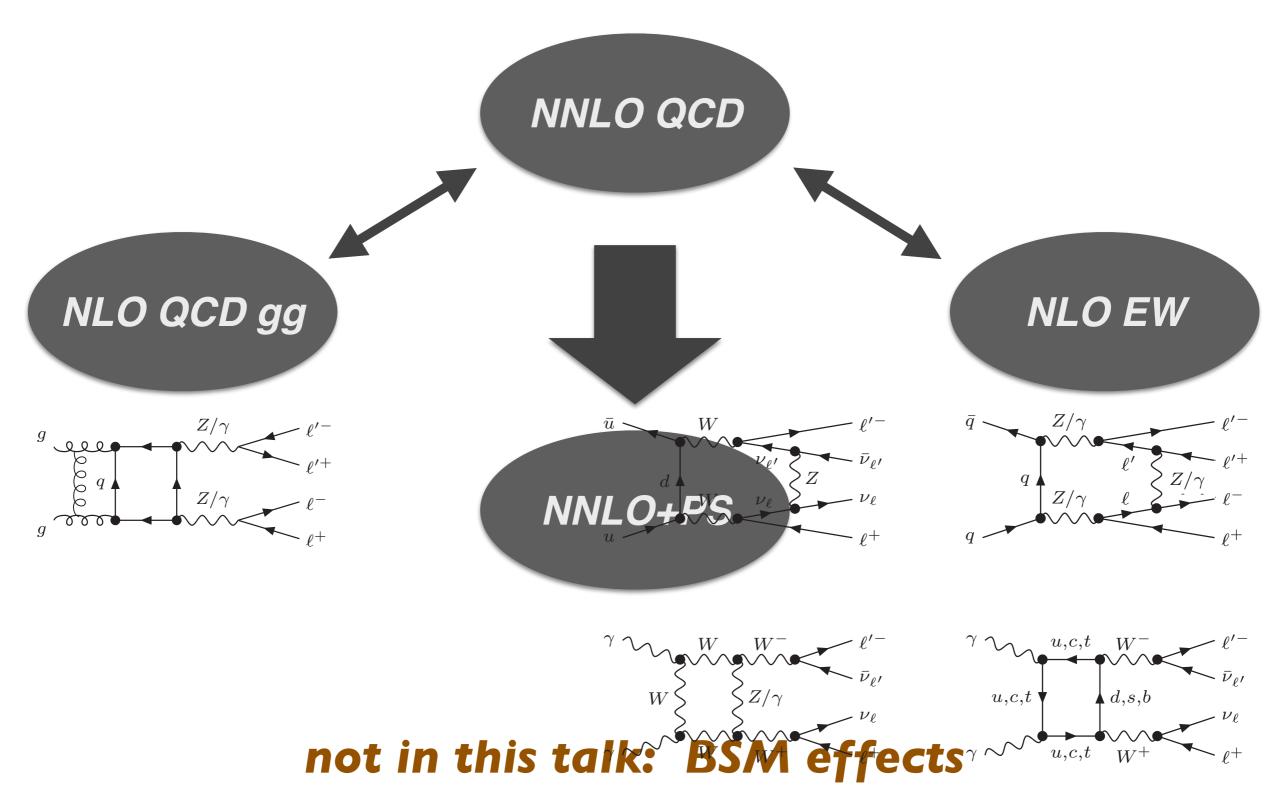
Importance of QCD corrections (example WZ)



NNLO crucial for accurate description of data

SM predictions: what is there?



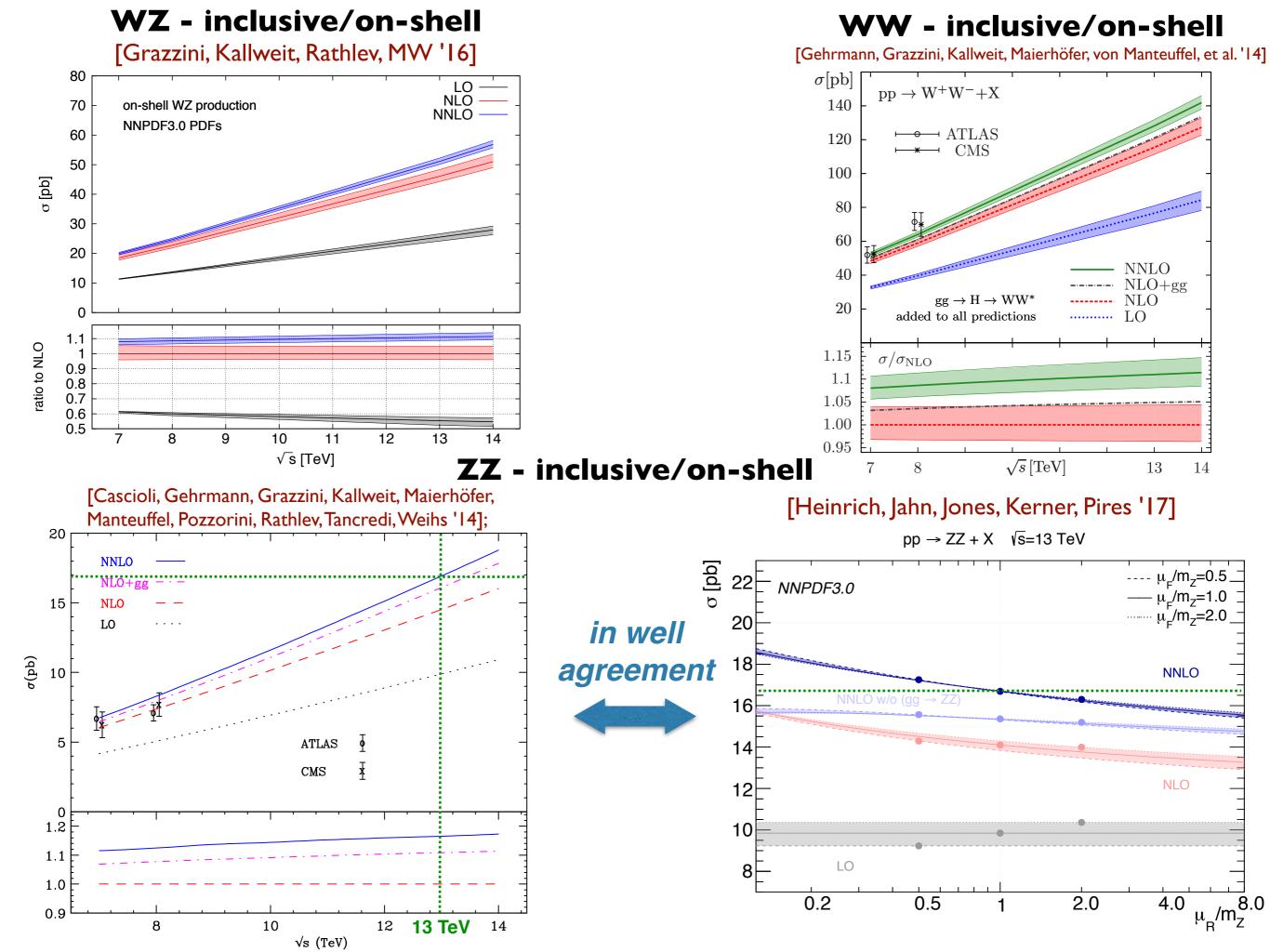


Higher-order QCD corrections



All VV processes known through NNLO QCD:

- → inclusive/on-shell Z,W & differential/off-shell Z,W (leptonic)
- inclusive and differential [Catani, Cieri, de Florian, Ferrera, Grazzini '12], [Campbell, Ellis, Li, Williams '16], [Grazzini, Kallweit, MW '17]
- **Zγ** inclusive/on-shell and differential/off-shell [Grazzini, Kallweit, Rathlev, Torre '13], [Grazzini, Kallweit, Rathlev '15]; see also: [Campbell et al. '17]
- Wγ inclusive/on-shell and differential/off-shell
 [Grazzini, Kallweit, Rathlev, Torre '13], [Grazzini, Kallweit, Rathlev '15]
- **TZ** inclusive/on-shell [Cascioli, Gehrmann, Grazzini, Kallweit, Maierhöfer, von Manteuffel, Pozzorini, Rathlev, Tancredi, Weihs '14]; see also: [Heinrich et al. '17]
 - differential/off-shell [Grazzini, Kallweit, Rathlev '15], [Kallweit, MW '18]
- WW inclusive/on-shell [Gehrmann, Grazzini, Kallweit, Maierhöfer, von Manteuffel, et al. '14]
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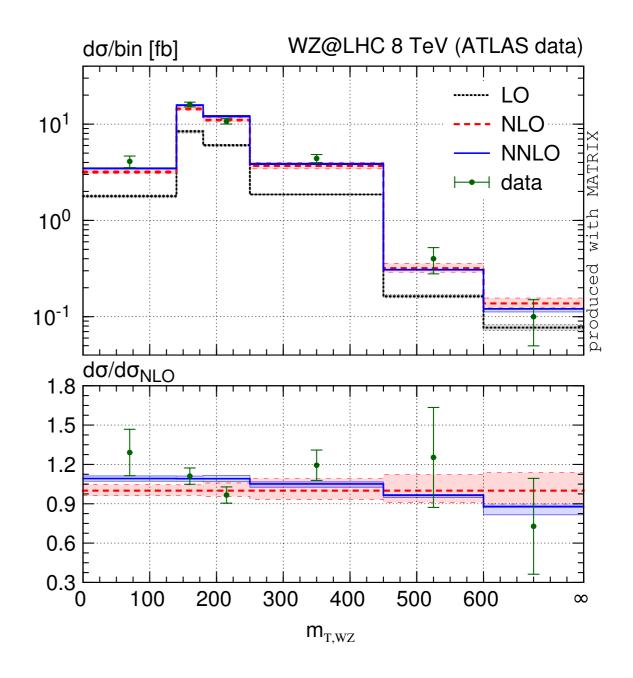
WW - differential/off-shell

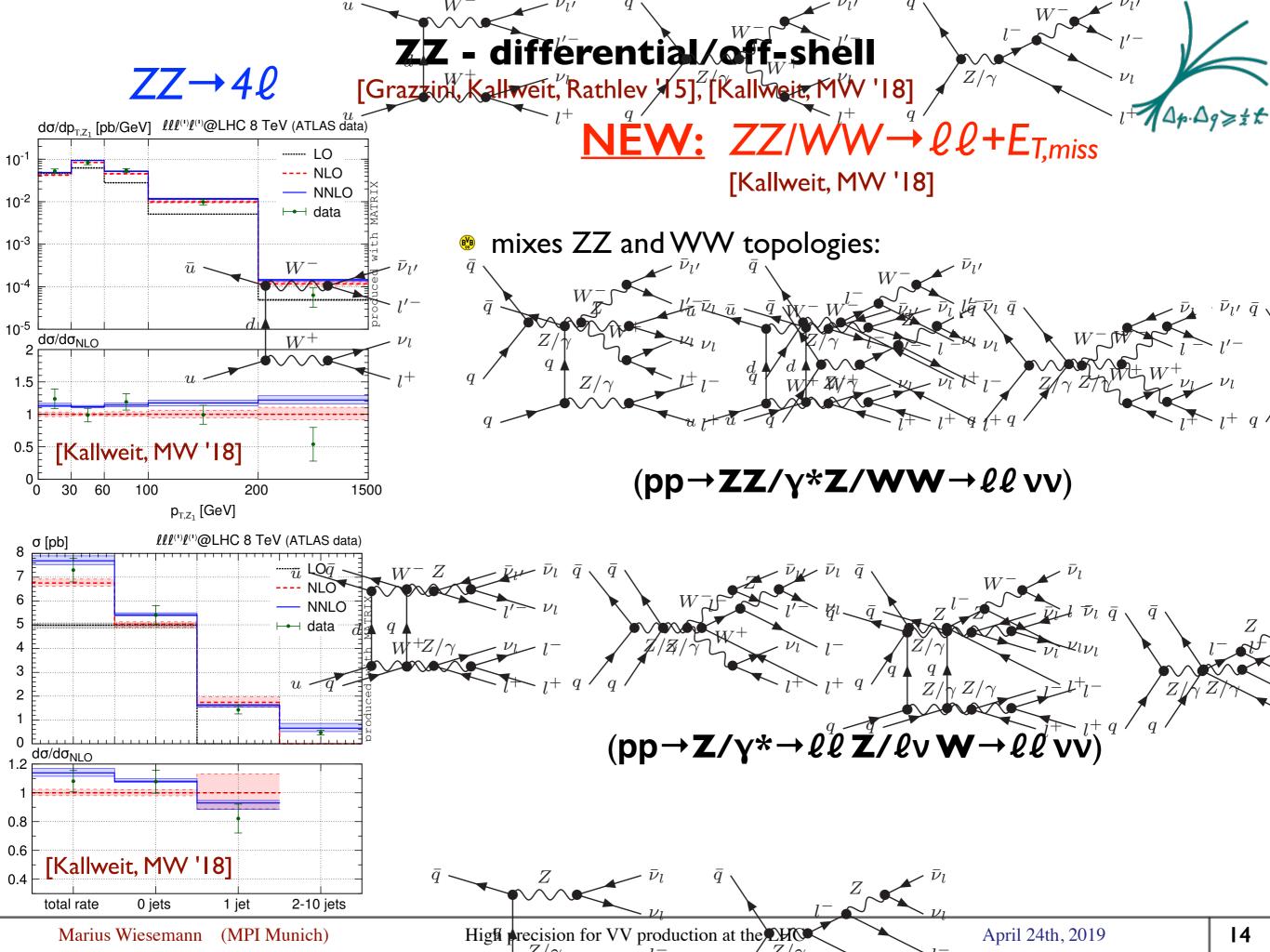
[Grazzini, Kallweit, Pozzorini, Rathlev, MW '15]

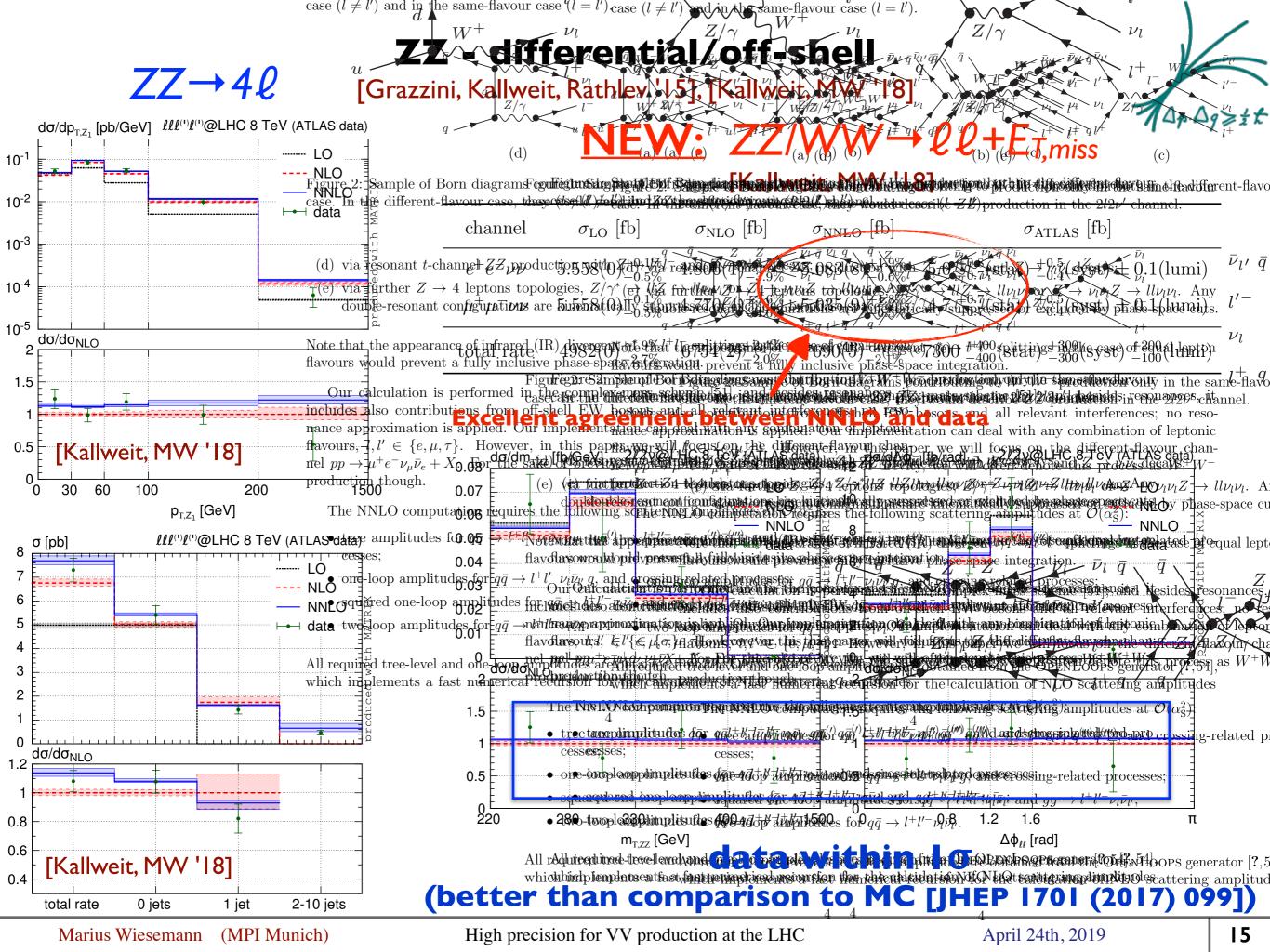
$\mu^+e^-\nu_{\mu}\bar{\nu}_{e}$ (inclusive)@LHC 8 TeV dσ/dm_{WW} [fb/GeV] 10¹ 10⁰ 10^{-1} 10⁻² LO 10⁻³ **NLO NNLO** 10⁻⁴ $d\sigma/d\sigma_{NLO}$ NLO'+gg 1.2 1.1 500 200 250 300 350 450 400 150 m_{WW} [GeV]

WZ - differential/off-shell

[Grazzini, Kallweit, Rathlev, MW '17]







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Higher-order QCD corrections



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all publicly available within MATRIX

process	status	comment
pp→ Z /γ*(→ℓℓ/νν)	√	validated analytically + FEWZ
$pp \rightarrow W(\rightarrow \ell \nu)$ single proces	√	validated with FEWZ, NNLOjet
pp→H	√	validated analytically (by SusHi)
рр→үү	√	validated with 2γNNLO
pp→ Z γ→ℓℓγ photo	n 🗸	[Grazzini, Kallweit, Rathlev '15]
pp→ Z γ→ννγ proces	sses 🗸	[Grazzini, Kallweit, Rathlev '15]
pp→Wγ→ℓνγ	✓	[Grazzini, Kallweit, Rathlev '15]
pp→ ZZ	√	[Cascioli et al. 'I4]
pp→ ZZ →ℓℓℓℓ		[Grazzini, Kallweit, Rathlev '15], [Kallweit, MW '18]
pp→ ZZ →ℓℓℓ'ℓ'		[Grazzini, Kallweit, Rathlev '15], [Kallweit, MW '18]
pp→ ZZ →ℓℓν'ν'		[Kallweit, MW '18]
pp→ZZ/WW→ℓℓvv mas	ssive	[Kallweit, MW '18]
nn-\\\\\\	cesses 🗸	[Gehrmann et al. '14]
pp→ WW →ℓν ℓ'ν'	1	[Grazzini, Kallweit, Pozzorini, Rathlev, MW ' 6
pp→ WZ	1	[Grazzini, Kallweit, Rathlev, MW '16]
pp→ WZ →ℓνℓℓ	1	[Grazzini, Kallweit, Rathlev, MW '17]
pp→ WZ →ℓ'v'ℓℓ	1	[Grazzini, Kallweit, Rathlev, MW '17]
рр→НН	(/)	not in public release

The MATRIX framework

[Grazzini, Kallweit, MW '17]

https://matrix.hepforge.org/

Amplitudes

OPENLOOPS (COLLIER, CUTTOols, ...)

Munich

MUlti-chaNnel Integrator at Swiss (CH) precision

 q_{T} subtraction $\iff q_{\mathrm{T}}$ resummation

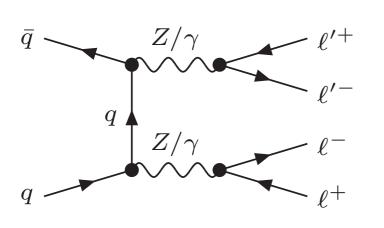
MATRIX

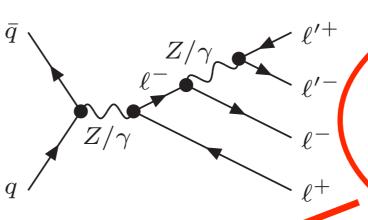
Munich Automates qT Subtraction and Resummation to Integrate X-sections.

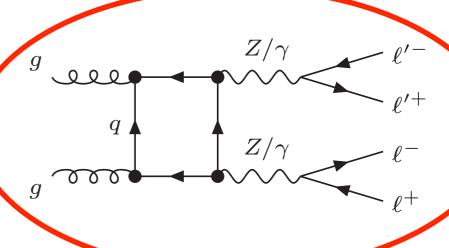
Recent developments for VV production

$gg \rightarrow ZZ \rightarrow 4\ell$ at NLO [Grazzini, Kallweit, MW, Yook '18]





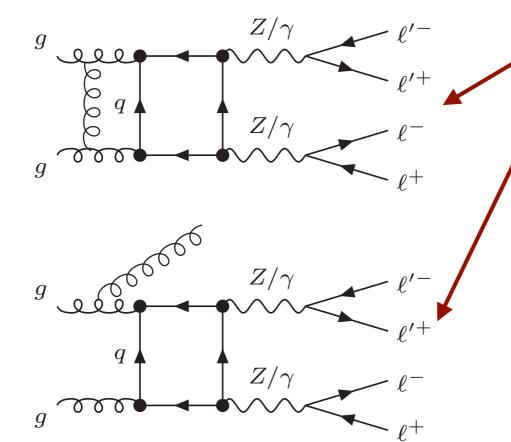




gg NLO:

virtuals:

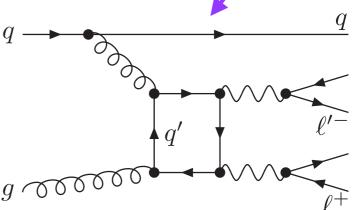
reals:



see also:

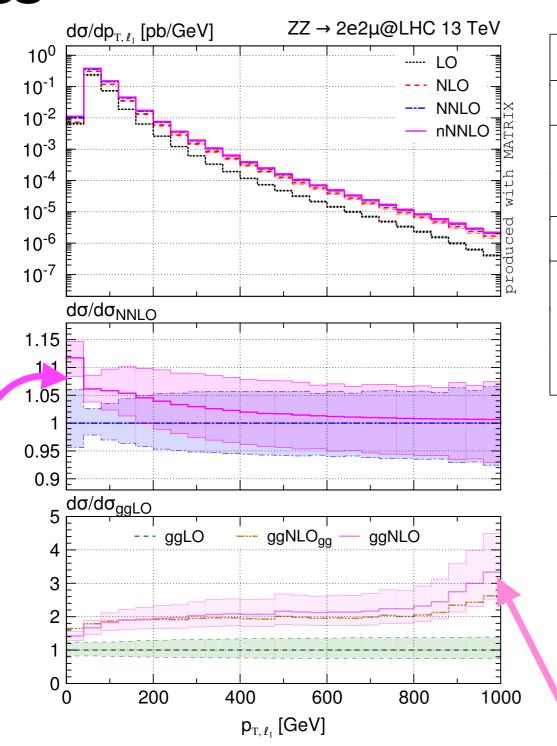
[Caola, Melnikov, Röntsch, Tancredi '15]

NEW: qg contributions



$+ZZ \rightarrow 4\ell$ at NLO [Grazzini, Kallweit, MW, Yook '18]





				•
\sqrt{s}	8 TeV	$13\mathrm{TeV}$	8 TeV	13 TeV
	σ [fb]		$\sigma/\sigma_{ m ggLO}-1$	
ggLO	$0.79354(8)_{-20.9\%}^{+28.2\%}$	$2.0054(2)_{-17.9\%}^{+23.5\%}$	0%	0%
ggNLO	$1.3901(9)_{-13.6\%}^{+15.4\%}$	$3.423(3)_{-12.0\%}^{+13.9\%}$	+75.2%	+70.7%
	σ [fb]		$\sigma/\sigma_{ m NLO}-1$	
NLO	$11.2958(4)_{-2.0\%}^{+2.5\%}$	$19.8454(7)_{-2.1\%}^{+2.5\%}$	0%	0%
NNLO	$12.87(3)_{-2.1\%}^{+2.8\%}$	$23.55(2)_{-2.6\%}^{+3.0\%}$	+13.9%	+18.7%
nNNLO	$13.47(3)_{-2.2\%}^{+2.6\%}$	$24.97(2)_{-2.7\%}^{+2.9\%}$	+19.2%	+25.8%

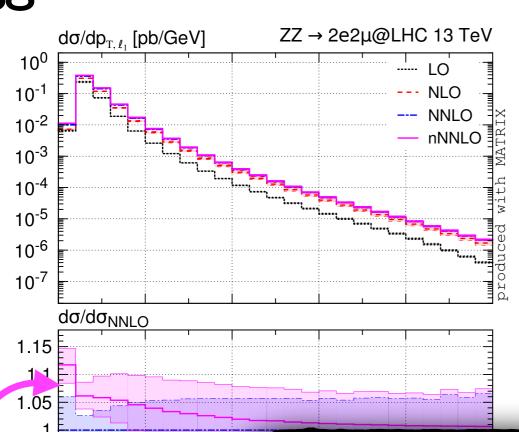
+5-6% effect due to NLO correction to gg compared to NNLO

NLO gg correction large+not flat; moves nNNLO outside uncertainty band of NNLO

huge NLO gg K-factor (~2 & more); impact of newly computed fermionic channels clearly visible

$gg \rightarrow ZZ \rightarrow 4\ell$ at NLO [Grazzini, Kallweit, MW, Yook '18]





\sqrt{s}	8 TeV	$13\mathrm{TeV}$	$8\mathrm{TeV}$	$13\mathrm{TeV}$
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QCD corrections beyond nNNLO?

- 1. Full N3LO (qq + mixing with gg)?
 - → covered by qq scale variations
- 2. beyond NLO gg: "ggNNLO"?
 - → NLO K-factor ~1.7
 - \rightarrow if NNLO K-factor \sim O(1.2-1.3), then nnNNLO/nNNLO~O(1.02-1.04)

& more); d fermionic

orrection to gg

NLO gg correction moves nNNLO of band of NNLO

0.95

0.9

 $d\sigma/d\sigma_{qqLO}$

ggLO

200

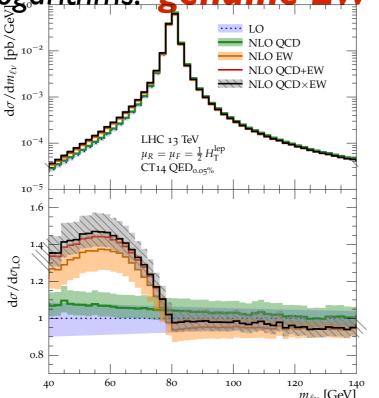
 p_{T,ℓ_1} [GeV]

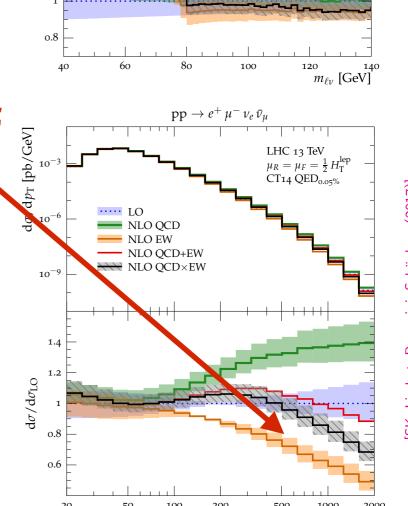
Combination: NNLO QCD and NLO EW

[Grazzini, Kallweit, Lindert, Pozzorini, MW]

- why adding NLO EW?
 - → needed for percent-level precision (inclusive)
 - → distortion of line shapes (EW resonances) due to photon FSR from leptons: QED effect
 - → sizable negative correction in tails

due to large Sudakov logarithms: genuine EW effect





 $pp \rightarrow e^+ \mu^- \nu_e \bar{\nu}_{\mu}$

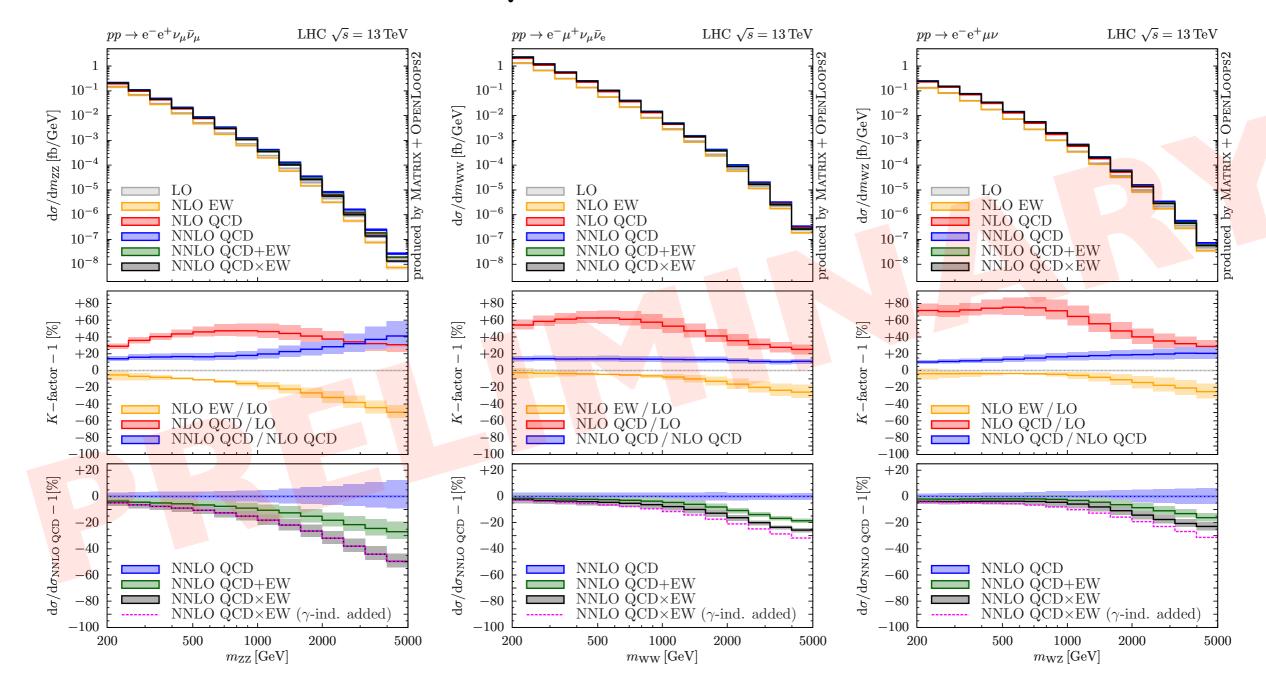
CT14 QED_{0.05%}



[Grazzini, Kallweit, Lindert, Pozzorini, MW]

- all vector-boson pair production processes
- **©** combination? additive: $d\sigma_{\rm QCD+EW}^{\rm (N)NLO} = d\sigma^{\rm LO}(1 + \delta_{\rm QCD} + \delta_{\rm EW})$

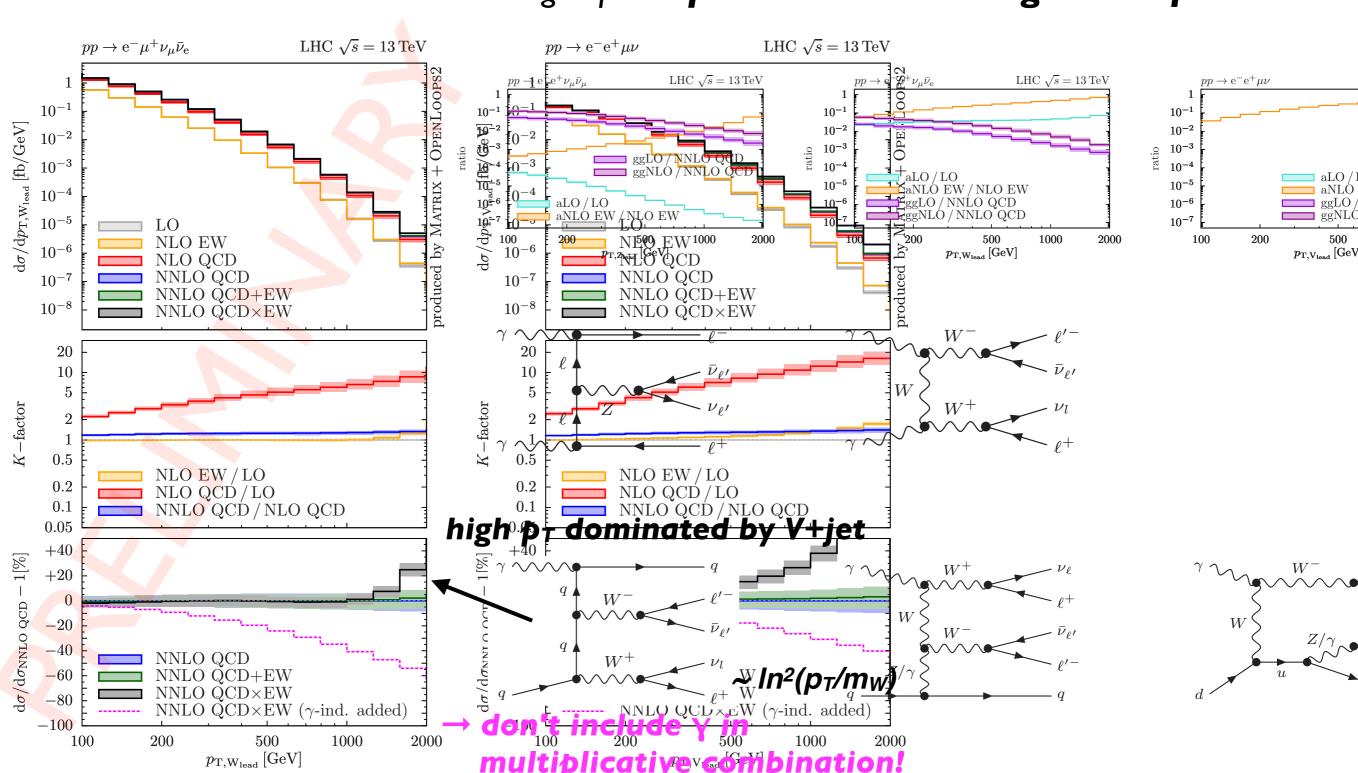
multiplicative: $d\sigma_{\mathrm{QCD} \times \mathrm{EW}}^{\mathrm{(N)NLO}} = d\sigma^{\mathrm{LO}} (1 + \delta_{\mathrm{QCD}}) (1 + \delta_{\mathrm{EW}})$





[Grazzini, Kallweit, Lindert, Pozzorini, MW]

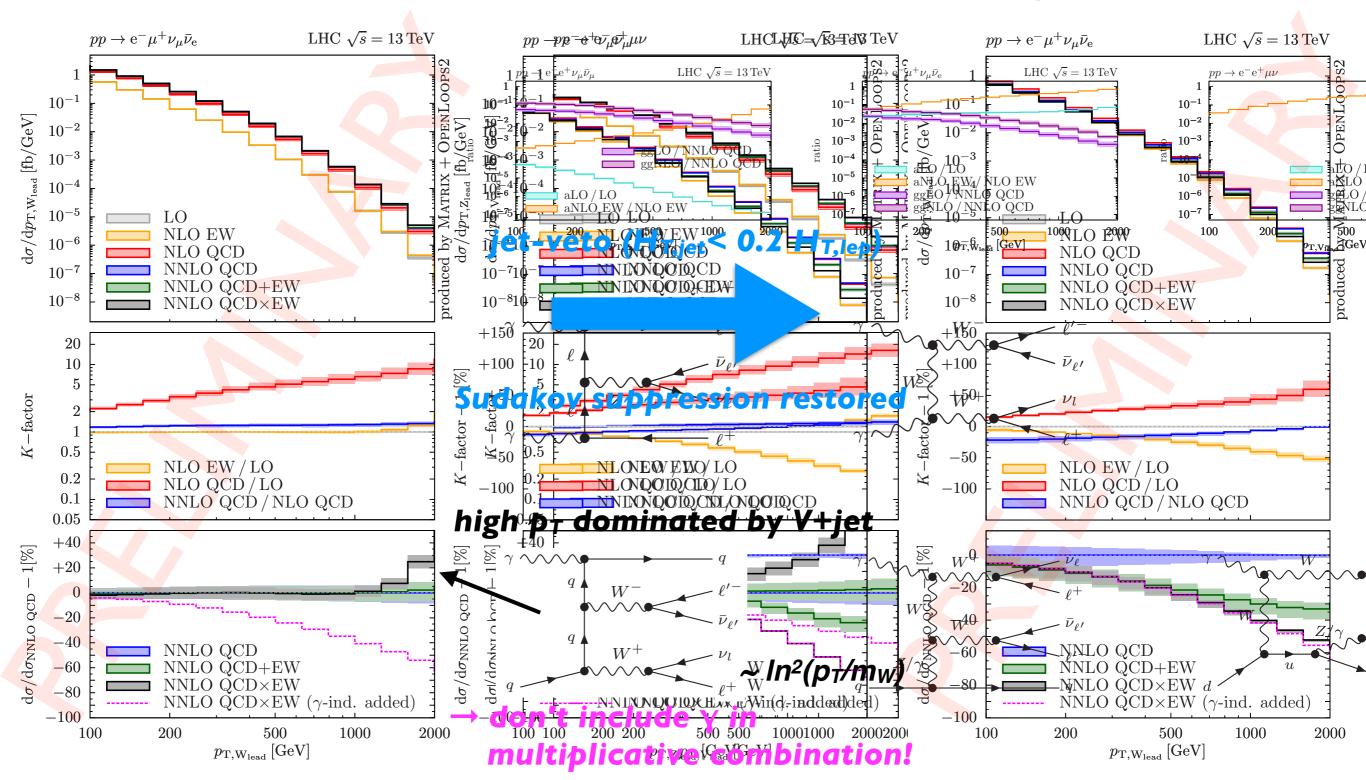
let's look in detail on one interesting aspect: photon-induced + giant K-factor





[Grazzini, Kallweit, Lindert, Pozzorini, MW]

let's look in detail on one interesting aspect: photon-induced + giant K-factor





[Grazzini, Kallweit, Lindert, Pozzorini, MW]

Many unanswered questions for final combination

$$\sigma \sim \sigma_{q\bar{q}}^{\rm LO} + \sigma_{\gamma\gamma}^{\rm LO} + \Delta_{q\bar{q}}^{\rm NLO_{QCD}} + \Delta_{q\bar{q}}^{\rm NNLO_{QCD}} + \Delta_{q\bar{q}/\gamma\gamma}^{\rm NLO_{EW}} + \Delta_{\rm gg}^{\rm LO} + \Delta_{\rm gg}^{\rm NLO_{QCD}}$$

$$\sigma \sim \left(\sigma_{q\bar{q}}^{\mathrm{LO}} + \sigma_{\gamma\gamma}^{\mathrm{LO}}\right) \left(1 + \frac{\Delta_{q\bar{q}}^{\mathrm{NLO_{QCD}}} + \Delta_{q\bar{q}}^{\mathrm{NNLO_{QCD}}} + \Delta_{\mathrm{gg}}^{\mathrm{LO}} + \Delta_{\mathrm{gg}}^{\mathrm{NLO_{QCD}}}}{\sigma_{q\bar{q}}^{\mathrm{LO}} + \sigma_{\gamma\gamma}^{\mathrm{LO}}}\right) \left(1 + \frac{\Delta_{q\bar{q}/\gamma\gamma}^{\mathrm{NLO_{QCD}}}}{\sigma_{q\bar{q}}^{\mathrm{LO}} + \sigma_{\gamma\gamma}^{\mathrm{LO}}}\right)$$

- 1. best approach to combine QCD and EW?
- 2. how to treat photon initial states?
- 3. how to treat QCD-EW mixing in q\(\gamma\) channels?
- 4. how to treat (NLO) gg contributions?

 ${
m l}\sigma/{
m d}p_{
m T,W_{
m lead}} \left[{
m fb/GeV}
ight]$

NNLOPS for WW

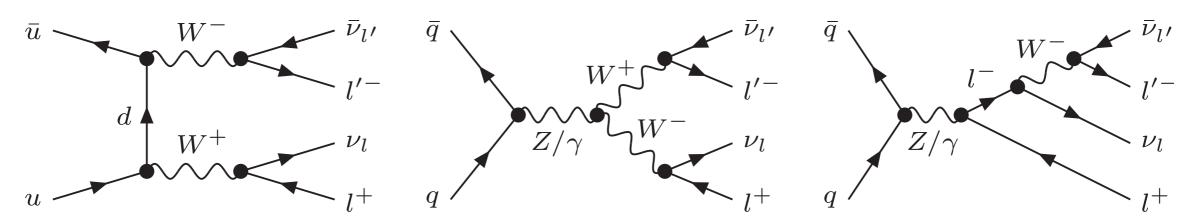


[Re, MW, Zanderighi '18]

using MiNLO+reweighting approach to reach NNLOPS:

$$W(\Phi_B) = \frac{\left(\frac{d\sigma}{d\Phi_B}\right)_{\text{NNLO}}}{\left(\frac{d\sigma}{d\Phi_B}\right)_{\text{XJ-MiNLO'}}} = \frac{c_0 + c_1\alpha_{\text{S}} + c_2\alpha_{\text{S}}^2}{c_0 + c_1\alpha_{\text{S}} + d_2\alpha_{\text{S}}^2} \simeq 1 + \frac{c_2 - d_2}{c_0}\alpha_{\text{S}}^2 + \mathcal{O}(\alpha_{\text{S}}^3)$$

- NNLO for WW production from: $pp \to \ell^- \bar{\nu}_\ell \; \ell'^+ \nu_{\ell'} + X$
 - → first fully differential NNLO computation for [Grazzini, Kallweit, Rathlev, MW '16]
 - → publicly available within MATRIX [Grazzini, Kallweit, MW '17]
- matched to parton showering within MiNLO [Hamilton, Monni, Re, Zanderighi '12]



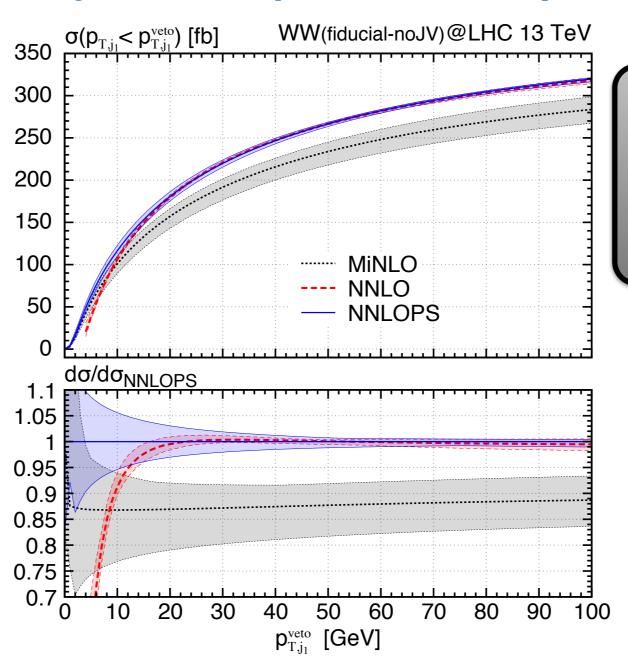
First NNLOPS accurate computation for 2→4 process

NNLOPS for WW



[Re, MW, Zanderighi '18]

jet veto (IR sensitive)



Many other distributions shown in Re, MW, Zanderighi; arXiv:1805.09857

→ NNLO provides adequate description of jet veto down to ~15 GeV, below NNLO is unphysical

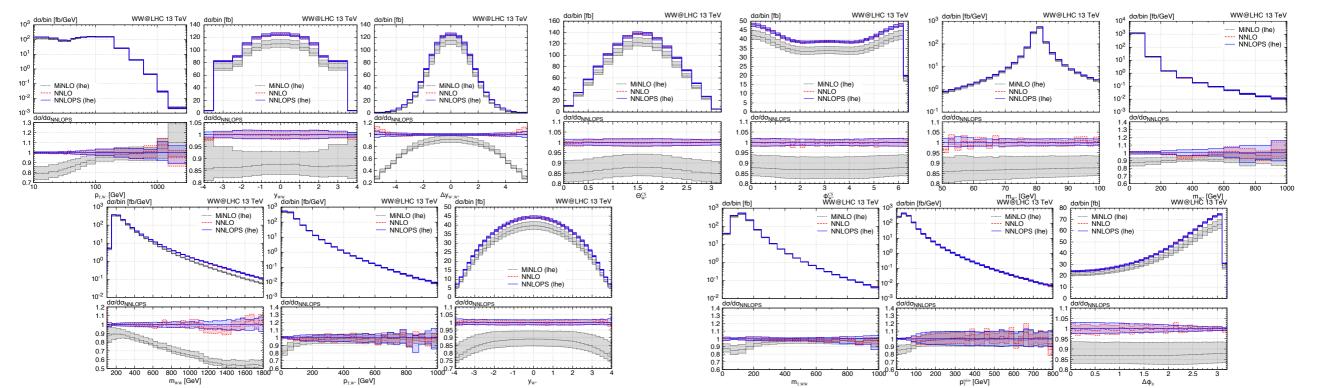
NNLOPS without reweighting



- despite physical sound results, numerical reweighting induces limitations:
 - ⇒ approximation: mw flat & CS angles $\frac{\mathrm{d}\sigma}{\mathrm{d}\Phi_{\mathrm{B}}} = \frac{\mathrm{d}\sigma}{\mathrm{d}p_{T,W} \mathrm{d}y_{WW}\mathrm{d}\Delta y_{W^{+}W^{-}}\mathrm{d}\cos\theta_{W^{+}}^{\mathrm{CS}}\mathrm{d}\phi_{W^{-}}^{\mathrm{CS}}\mathrm{d}\phi_{W^{-}}^{\mathrm{CS}}\mathrm{d}m_{W^{+}}\mathrm{d}m_{W^{-}}}$ to convert the 9D Born phase space into the computation of 81 3D moments
 - ⇒ discrete binning limits p_{T,W^-} : $[0.,17.5,25.,30.,35.,40.,47.5,57.5,72.5,100.,200.,350.,600.,1000.,1500.,\infty]$; applicability in less y_{WW} : $[-\infty,-3.5,-2.5,-2.0,-1.5,-1.0,-0.5,0.0,0.5,1.0,1.5,2.0,2.5,3.5,\infty]$; populated phase- $\Delta y_{W^+W^-}$: $[-\infty,-5.2,-4.8,-4.4,-4.0,-3.6,-3.2,-2.8,-2.4,-2.0,-1.6,-1.2,-0.8,-0.4,0.0,0.4,0.8,1.2,1.6,2.0,2.4,2.8,3.2,3.6,4.0,4.4,4.8,5.2,\infty]$.
 - → reweighting still numerically intensive
 - → thorough validation required



Issue in experimental NNLOPS event production already for DY

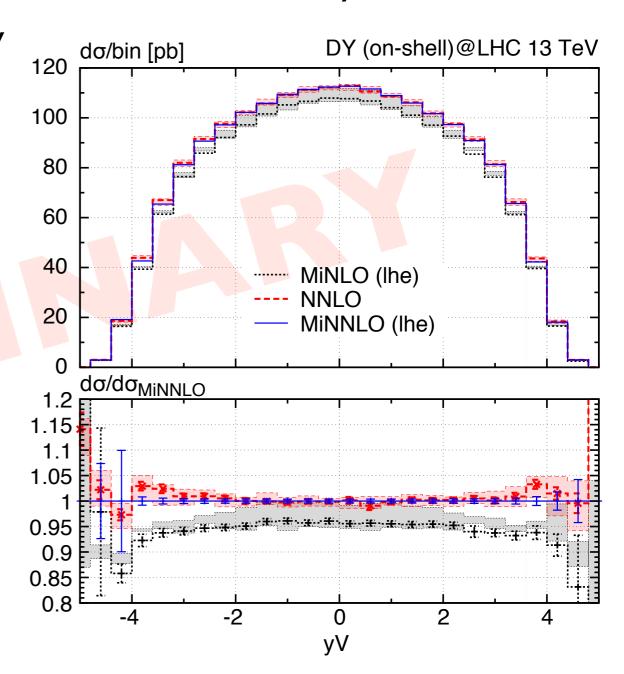


NNLOPS without reweighting



- despite physical sound results, numerical reweighting induces limitations
- new idea to avoid reweighting: [Monni, Nason, Re, MW, Zanderighi]
 - → compare MiNLO with resummation formula at sufficient accuracy
 - → identify missing terms for NNLO accuracy
 - → add them explicitly to MiNLO

first results for DY look promising:



Conclusions



Theory predictions under excellent control:

- NNLO QCD for VV done! → publicly available within MATRIX
- $\ell\ell$ +ET,miss signature studied at NNLO, mixes ZZ and WW resonances
- combination with NLO QCD for loop-induced gg and with NLO EW ongoing
- first NNLOPS computation for a 2→4 process (WW)

Several open issues/questions remain:

- corrections beyond nNNLO (NLO gg) relevant? especially "NNLO gg"?
- how to combine NNLO, NLO gg and NLO EW? photon contributions?
- how to avoid numerically heavy reweighting for NNLOPS? (most important for complicated processes such as dibosons)

FREE YOUR MIND

[[wiesemann:~/different-branch-munich/MATRIX] ./matrix



<<MATRIX-MAKE>> This is the MATRIX process compilation.

[|=====>> list

process_id	П	process	П	description
pph21 ppz01 ppw01 ppwx01 ppeex02 ppnenex02 ppenex02 ppenex02 ppena02 ppaa02 ppeexa03		p p> H p p> Z p p> W^- p p> W^+ p p> e^- e^+ p p> e^- v_e^+ p p> e^- v_e^- p p> gamma gamma p p> e^- e^+ gamma	>> >> >> >> >> >> >> >> >> >>	Z production with decay

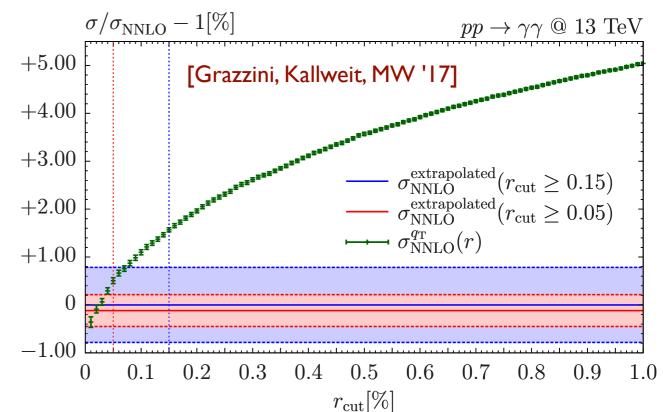
Thank You !

Back Up

yy - inclusive and differential

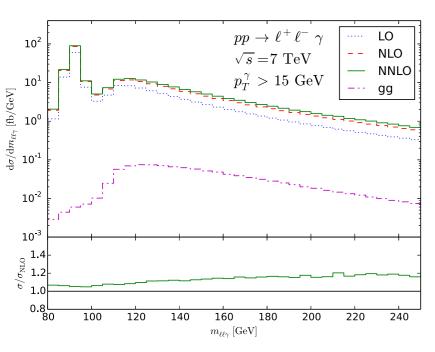
[Catani, Cieri, de Florian, Ferrera, Grazzini '12], [Campbell, Ellis, Li, Williams '16], [Grazzini, Kallweit, MW '17]

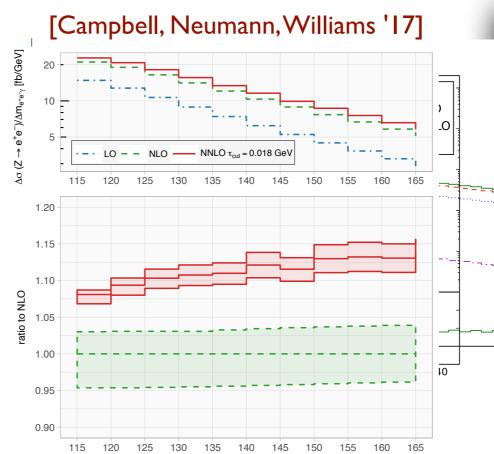
- known only with slicing techniques
- photon processes quite delicate dependence on slicing parameter due to photon isolation
- well under control in state-of-the-art tools like MATRIX (see plot on the right)
- systematic uncertainties still larger than for other diboson processes, but few permille possible
- agreement among computation within respective uncertainties



Zy - inclusive/on-shell and differential/off-shell

[Grazzini, Kallweit, Rathlev 'I 5]



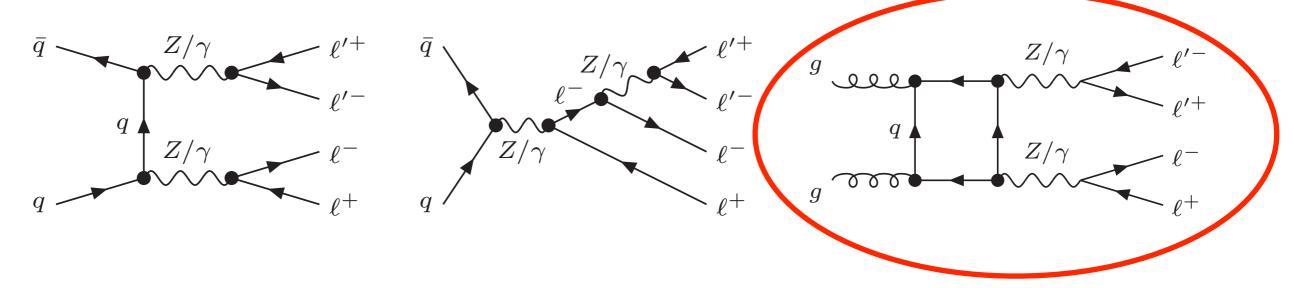


 $m_{e^+e^-\gamma}$ [GeV]

		[Grazzini, Kallweit, MW '17]		
(<pre>process \${process_id})</pre>	$\sigma_{ m NNLO}^{ m extrapolated}$	$K_{ m NLO}$	$K_{ m NNLO}$
	$pp o \gamma \gamma$ (ppaa02)	$40.28(30)^{+8.7\%}_{-7.0\%} \mathrm{pb}$	+361%	+56.4%
	$pp \rightarrow e^-e^+\gamma$ (ppeexa03)	$2316(5)^{+1.1\%}_{-1.2\%}$ fb	+44.3%	+9.29%
	$pp o u_e \bar{ u}_e \gamma$ (ppnenexa03)	$113.5(6)^{+2.9\%}_{-2.4\%}$ fb	+55.2%	+15.0%
	$pp o e^- \bar{\nu}_e \gamma$ (ppenexa03)	$2256(15)^{+3.7\%}_{-3.5\%}$ fb	+155%	+22.0%
	$pp \to e^+ \nu_e \gamma$ (ppexnea03)	$2671(35)^{+3.8\%}_{-3.6\%}$ fb	+154%	+22.1%

$gg \rightarrow ZZ \rightarrow 4\ell$ at NLO QCD





- gg→ZZ contribution part of NNLO corrections to ZZ production

 BUT: significant impact ~50% of the (large) NNLO correction, due to gluon PDFs

 Output

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 Description

 BUT: significant impact
- NLO corrections to gg→ZZ are formally part of N3LO cross section
 BUT: still important due to large K-factor
- computed without fermionic channels [Caola, Melnikov, Röntsch, Tancredi '15]
- so far NNLO qq and NLO gg (uncertainty) combined as independent processes HOWEVER: their diagrams mix/interfere already at NNLO



$gg \rightarrow ZZ \rightarrow 4\ell$ at NLO [Grazzini, Kallweit, MW, Yook '18]



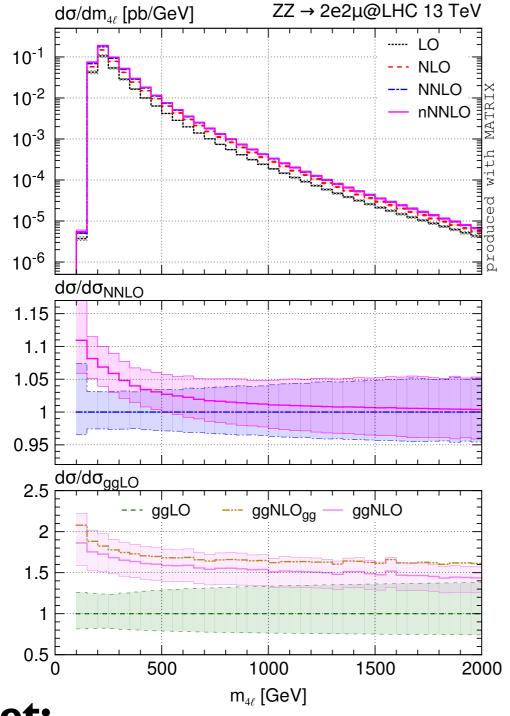
included in the MATRIX framework

\sqrt{s}	8 TeV	$13\mathrm{TeV}$	8 TeV	$13\mathrm{TeV}$
	σ [fb]		$\sigma/\sigma_{ m NI}$	$L_{0} - 1$
LO	$8.1881(8)_{-3.2\%}^{+2.4\%}$	$13.933(7)_{-6.4\%}^{+5.5\%}$	-27.5%	-29.8%
NLO	$11.2958(4)_{-2.0\%}^{+2.5\%}$	$19.8454(7)_{-2.1\%}^{+2.5\%}$	0%	0%
$q\bar{q}$ NNLO	$12.08(3)_{-1.1\%}^{+1.1\%}$	$21.54(2)_{-1.2\%}^{+1.1\%}$	+6.9%	+8.6%
	σ [fb]		$\sigma/\sigma_{ m ggLO}-1$	
ggLO	$0.79354(8)_{-20.9\%}^{+28.2\%}$	$2.0054(2)_{-17.9\%}^{+23.5\%}$	0%	0%
$ggNLO_{gg}$	$1.4810(9)_{-13.2\%}^{+16.0\%}$	$3.627(3)_{-12.8\%}^{+15.2\%}$	+86.6%	+80.9%
ggNLO	$1.3901(9)^{+15.4\%}_{-13.6\%}$	$3.423(3)_{-12.0\%}^{+13.9\%}$	+75.2%	+70.7%
	σ [fb]		$\sigma/\sigma_{ m NI}$	$L_{0} - 1$
NNLO	$12.87(3)_{-2.1\%}^{+2.8\%}$	$23.55(2)_{-2.6\%}^{+3.0\%}$	+13.9%	+18.7%
nNNLO	$13.47(3)_{-2.2\%}^{+2.6\%}$	$24.97(2)_{-2.7\%}^{+2.9\%}$	+19.2%	+25.8%

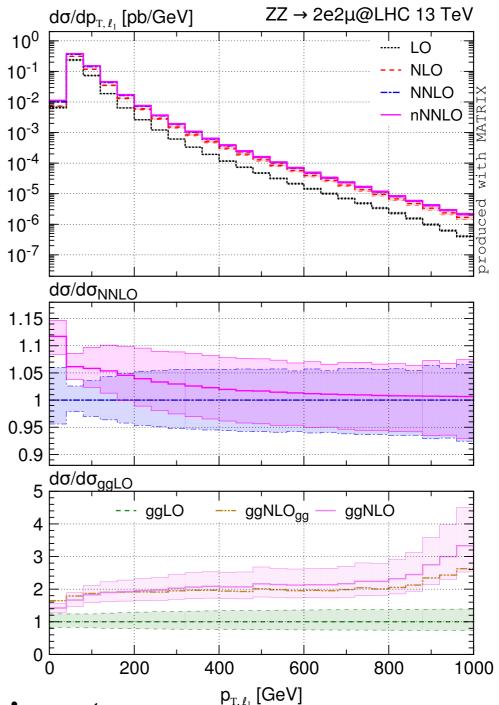
+5-6% effect due to NLO correction to gg compared to NNLO

$gg \rightarrow ZZ \rightarrow 4\ell$ at NLO [Grazzini, Kallweit, MW, Yook '18]





I. inset: moves nNNLO outside uncertainty band of NNLO



2. inset: **NLO** gg correction large+not flat; huge NLO gg K-factor (~2 & more); impact of newly computed fermionic channels clearly visible

$gg \rightarrow ZZ \rightarrow 4\ell$ at NLO [Grazzini, Kallweit, MW, Yook '18]



Setup (8 TeV ATLAS ZZ measurement):

definition of the fiducial volume for $pp \to e^+e^-\mu^+\mu^- + X$

$$p_{T,e/\mu} > 7 \,\text{GeV}$$
, one electron with $|\eta_e| < 4.9$, the others $|\eta_e| < 2.5$, $|\eta_{\mu}| < 2.7$
 $\Delta R_{ee/\mu\mu} > 0.2$, $\Delta R_{e\mu} > 0.2$, $66 \,\text{GeV} \le m_{e^+e^-/\mu^+\mu^-} \le 116 \,\text{GeV}$,

Top quark, and Higgs (interference) contribution:

- full top-quark dependence everywhere, but in the 2-loop amplitudes (true for both NNLO qq and NLO gg)
- 2-loop NNLO qq: closed (massless) fermion loops small → top neglected
- 2-loop NLO gg: top dependence approximated by rescaling massless 2-loop amplitude with the full Born-level amplitude
- Higgs (interference) contribution fully included following this approximation
 - → applicable to Higgs studies; validity of approximation around and beyond top threshold is to be tested; in ZZ signal region, Higgs contribution about -5%

Combination: NNLO QCD and NLO EW



[Grazzini, Kallweit, Lindert, Pozzorini, MW]

Setup of calculation (simplified selection, only for illustration purposes):

- EW input and renormalization in G_{μ} scheme.
- Uncertainties from 7-point scale variation around $\mu_R = \mu_F = \frac{1}{2} (E_{T,V_1} + E_{T,V_2})$.
- PDF set NNPDF31_nnlo_as_0118_luxqed(_nf_4) at each order.
- Only very basic set of phase-space selections:
 - $60 \, \mathrm{GeV} < m_{\ell\ell} < 120 \, \mathrm{GeV}$ for each SFOS lepton pair,
 - $p_{\mathrm{T},\ell} > 25 \,\mathrm{GeV}$ for each lepton, $p_{\mathrm{T,miss}} > 25 \,\mathrm{GeV}$,
 - lepton-photon recombination for $dR(\gamma, \ell) < 0.1$.
- Two setups: inclusive and with jet-veto $(H_{\rm T,jet} < 0.2 H_{\rm T,lep})$.

NNLOPS for VV production



- Use knowledge of fully differential NNLO cross section
- Use MiNLO to consistently merge VV+0, I jet at NLO [Hamilton, Nason, Zanderighi '12]
 - → no merging scale
 - \rightarrow NLO accurate in the entire VV+0, I phase space
- Use reweighting to promote MiNLO samples to NNLOPS



yields fully exclusive (hadron-level) events with NNLO accuracy

(the same way as POWHEG/MC@NLO are NLO+PS accurate)

NNLOPS procedure



MiNLO is general for colour singlets (X=H,V, VH,VV, ...)

[Hamilton, Nason, Zanderighi '12]

(only B2 NNLL coefficient becomes non-trivial starting from $2\rightarrow 2$)

	X	X+jet	X+2jets	X+nj (n>2)
XJ (NLO)		NLO	LO	
XJ-MiNLO	NLO	NLO	LO	PS
X@NNLO	NNLO	NLO	LO	

⊗ XJ-MiNLO already almost right accuracy, only $\mathcal{O}(\alpha_s^2)$ terms missing to obtain NNLO accuracy at Born level

NNLOPS procedure



MiNLO is general for colour singlets (X=H,V, VH,VV, ...)

[Hamilton, Nason, Zanderighi '12]

(only B2 NNLL coefficient becomes non-trivial starting from 2→2)

	X	X+jet	X+2jets	X+nj (n>2)
XJ (NLO)		NLO	LO	
XJ-MiNLO	NLO	NLO	LO	PS
X@NNLO	NNLO	NLO	LO	
X@NNLOPS	NNLO	NLO	LO	PS

**solution: reweighting MiNLO events in Born phase space to correct for $\mathcal{O}(\alpha_s^2)$ terms by the following ratio:

$$W(\Phi_B) = \frac{\left(\frac{d\sigma}{d\Phi_B}\right)_{\text{NNLO}}}{\left(\frac{d\sigma}{d\Phi_B}\right)_{\text{XJ-MiNLO'}}} = \frac{c_0 + c_1\alpha_{\text{S}} + c_2\alpha_{\text{S}}^2}{c_0 + c_1\alpha_{\text{S}} + d_2\alpha_{\text{S}}^2} \simeq 1 + \frac{c_2 - d_2}{c_0}\alpha_{\text{S}}^2 + \mathcal{O}(\alpha_{\text{S}}^3)$$

NNLOPS procedure



already used for processes with simpler Born kinematics:

- Higgs (only rapitidy) → ID reweighting [Hamilton, Nason, Re, Zanderighi 'I3], [Hamilton, Nason, Zanderighi 'I5],
- Drell-Yan (dilepton system) → 3D reweighting [Karlberg, Hamilton, Zanderighi '14]
- → Higgsstrahlung (Hℓℓ/Hℓν system) → 6D reweighting [Astill, Bizoń, Re, Zanderigh '16 '18]



[Re, MW, Zanderighi '18]

in principle, requires 9D reweighting (numerically impossible)

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Phi_{\mathrm{B}}} = \frac{\mathrm{d}^{9}\sigma}{\mathrm{d}p_{T,W^{-}}\mathrm{d}y_{WW}\mathrm{d}\Delta y_{W^{+}W^{-}}\mathrm{d}\cos\theta_{W^{+}}^{\mathrm{CS}}\mathrm{d}\phi_{W^{+}}^{\mathrm{CS}}\mathrm{d}\cos\theta_{W^{-}}^{\mathrm{CS}}\mathrm{d}m_{W^{+}}\mathrm{d}m_{W^{-}}}$$

drop invariant masses which are flat (in particular around the peak)

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Phi_{\mathrm{B}}} = \frac{\mathrm{d}^{9}\sigma}{\mathrm{d}p_{T,W^{-}}\mathrm{d}y_{WW}\mathrm{d}\Delta y_{W^{+}W^{-}}\mathrm{d}\cos\theta_{W^{+}}^{\mathrm{CS}}\mathrm{d}\phi_{W^{+}}^{\mathrm{CS}}\mathrm{d}\cos\theta_{W^{-}}^{\mathrm{CS}}\mathrm{d}\phi_{W^{-}}^{\mathrm{CS}}\mathrm{d}m_{W^{+}}\mathrm{d}m_{W^{-}}^{\mathrm{CS}}$$

Collins-Soper angles [Collins, Soper '77] are used to describe both W decays in terms of spherical harmonics and parametrize angular dependence by

$$\frac{d\sigma}{d\Phi_{\rm B}} = \frac{9}{256\pi^2} \sum_{i=0}^{8} \sum_{j=0}^{8} AB_{ij} f_i(\theta_{W^-}^{\rm CS}, \phi_{W^-}^{\rm CS}) f_j(\theta_{W^+}^{\rm CS}, \phi_{W^+}^{\rm CS})$$

$$AB_{ij}(p_{T,W^{-}}, y_{WW}, \Delta y_{W^{+}W^{-}}) = \int \frac{d\sigma}{d\Phi_{B}} g_{i}(\theta_{W^{-}}^{\text{CS}}, \phi_{W^{-}}^{\text{CS}}) g_{j}(\theta_{W^{+}}^{\text{CS}}, \phi_{W^{+}}^{\text{CS}}) d\cos\theta_{W^{-}}^{\text{CS}} d\phi_{W^{-}}^{\text{CS}} d\cos\theta_{W^{+}}^{\text{CS}} d\phi_{W^{+}}^{\text{CS}}$$

$$f_0(\theta,\phi) = (1 - 3\cos^2\theta)/2$$
, $f_1(\theta,\phi) = \sin 2\theta \cos \phi$, $f_2(\theta,\phi) = (\sin^2\theta \cos 2\phi)/2$,

$$f_3(\theta,\phi) = \sin\theta\cos\phi$$
, $f_4(\theta,\phi) = \cos\theta$, $f_5(\theta,\phi) = \sin\theta\sin\phi$,

$$f_6(\theta,\phi) = \sin 2\theta \sin \phi$$
, $f_7(\theta,\phi) = \sin^2 \theta \sin 2\phi$, $f_8(\theta,\phi) = 1 + \cos^2 \theta$.



compute 81 3D distributions (numerically feasible)

April 24th, 2019



[Re, MW, Zanderighi '18]

Setup:

The remaining three variables and their binning chosen to be

 $p_{T,W^-}: [0., 17.5, 25., 30., 35., 40., 47.5, 57.5, 72.5, 100., 200., 350., 600., 1000., 1500., \infty];$

 $y_{WW}: [-\infty, -3.5, -2.5, -2.0, -1.5, -1.0, -0.5, 0.0, 0.5, 1.0, 1.5, 2.0, 2.5, 3.5, \infty];$

 $\Delta y_{W^+W^-} \ : \quad [\, -\infty, -5.2, -4.8, -4.4, -4.0, -3.6, -3.2, -2.8, -2.4, -2.0, -1.6, -1.2, \\$

 $\hspace*{35pt} -0.8, -0.4, 0.0, 0.4, 0.8, 1.2, 1.6, 2.0, 2.4, 2.8, 3.2, 3.6, 4.0, 4.4, 4.8, 5.2, \infty] \, .$

Cuts inspired by ATLAS 13 TeV study (1702.04519):

lepton cuts	$p_{T,\ell} > 25 \mathrm{GeV}, \eta_\ell < 2.4, m_{\ell^-\ell^+} > 10 \mathrm{GeV}$		
lepton dressing	add photon FSR to lepton momenta with $\Delta R_{\ell\gamma} < 0.1$ (our results do not include photon FSR, see text)		
neutrino cuts	$p_T^{ m miss} > 20{ m GeV}, p_T^{ m miss,rel} > 15{ m GeV}$		
	anti- k_T jets with $R = 0.4$;		
jet cuts	$N_{ m jet}=0 ext{ for } p_{T,j}>25 ext{ GeV}, \eta_j <2.4 ext{ and } \Delta R_{ej}<0.3$		
	$N_{ m jet}=0 ext{ for } p_{T,j}>30 ext{ GeV}, \eta_j <4.5 ext{ and } \Delta R_{ej}<0.3$		

NNLO uses the central scale

$$\mu_R = \mu_F = \mu_0 \equiv rac{1}{2} \, \left(\sqrt{m_{e^-ar{
u}_e}^2 + p_{T,e^-ar{
u}_e}^2} + \sqrt{m_{\mu^+
u_\mu}^2 + p_{T,\mu^+
u_\mu}^2}
ight)$$

All uncertainty bands are the envelop of 7-scales. In the NNLOPS scales in MiNLO and NNLO are varied in a correlated way

gg-channel not included in our study, as it can it is know at one-loop and can be added incoherently



[Re, MW, Zanderighi '18]

Validation:

- 1. Total inclusive NNLO cross section reproduced by NNLOPS sample
- 2. NNLO distributions for observables used for reweighting reproduced $\sqrt{}$
- 3. NNLO distributions for CS angles reproduced $\sqrt{}$
- 4. NNLO distributions for invariant masses of W's reproduced $\sqrt{}$
- 5. NNLO distributions for other Born-level observables reproduced 🎻

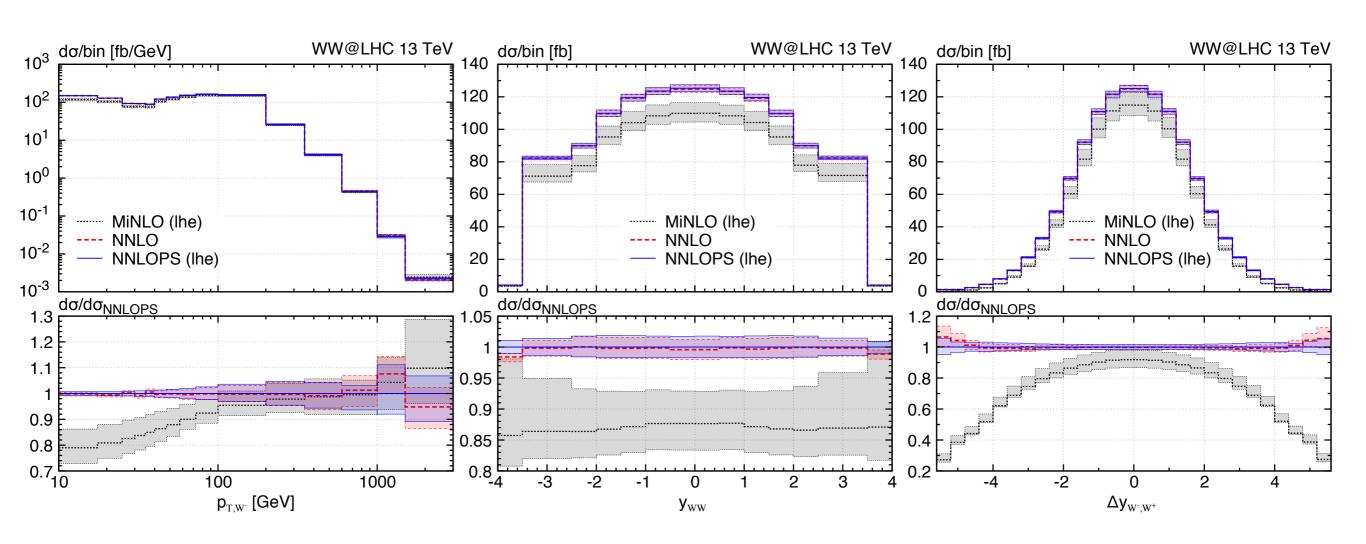


[Re, MW, Zanderighi '18]

Validation at LHE level:

2. NNLO distributions for observables used for reweighting reproduced





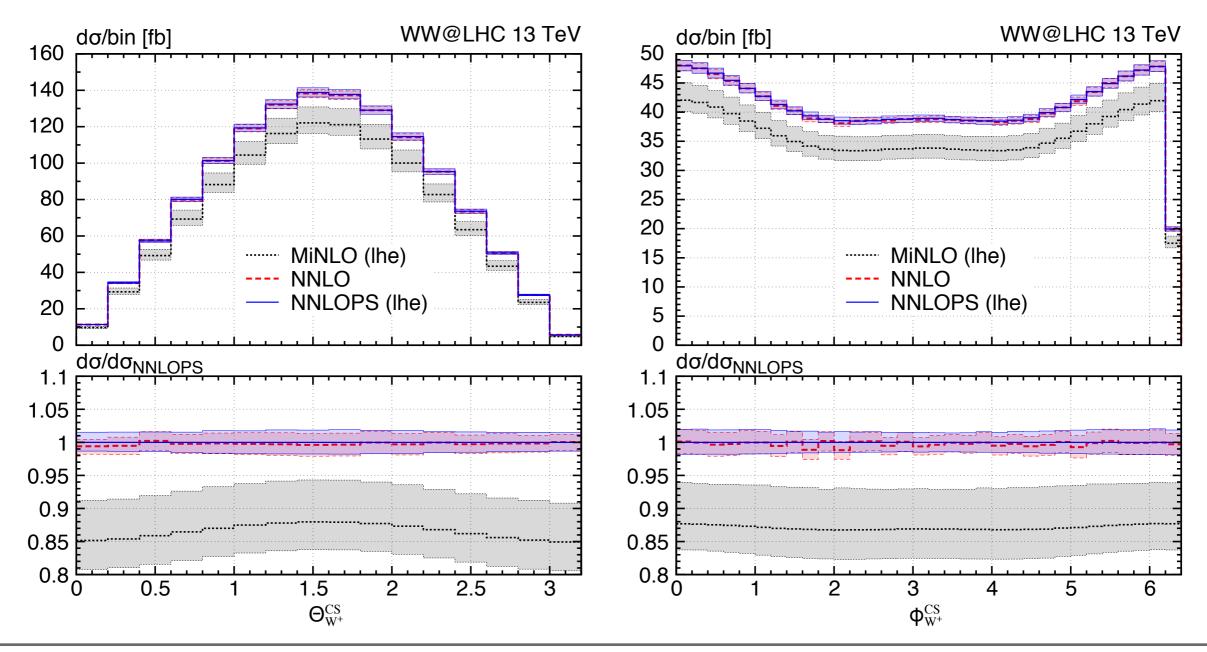


[Re, MW, Zanderighi '18]

Validation at LHE level:

3. NNLO distributions for CS angles reproduced





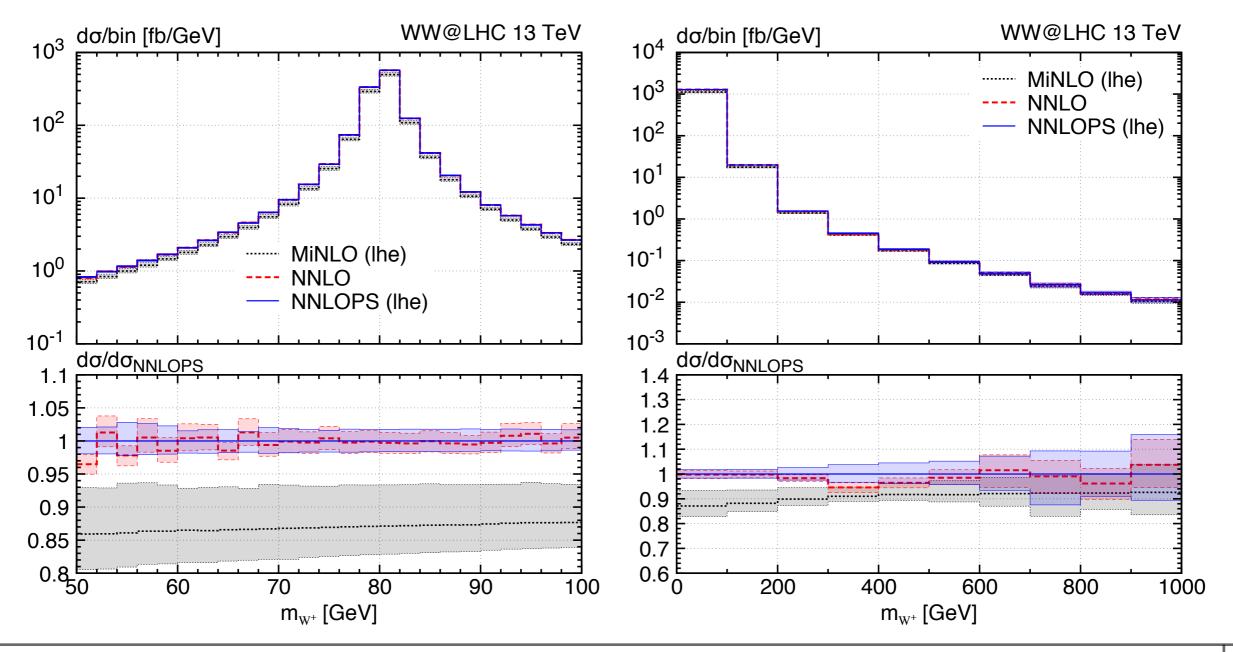


[Re, MW, Zanderighi '18]

Validation at LHE level:

4. NNLO distributions for invariant masses of W's reproduced



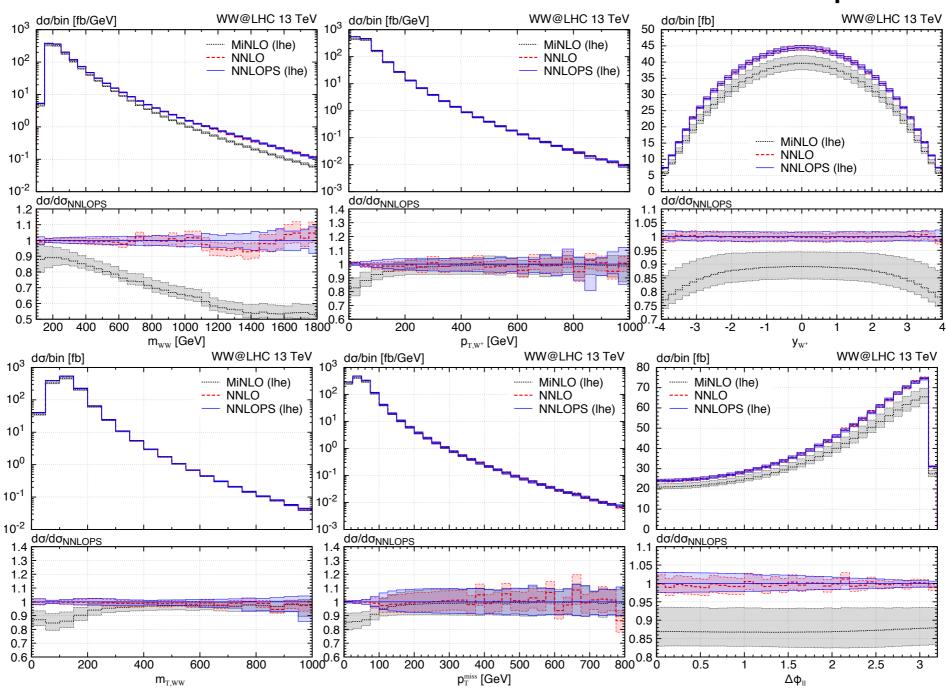


[Re, MW, Zanderighi '18]

Validation at LHE level:

5. NNLO distributions for other Born-level observables reproduced







[Re, MW, Zanderighi '18]

Phenomenological results:

Integrated cross sections

$q\bar{q}$ (no loop ² gg)	$\sigma_{\rm incl}(pp \to W^+W^-)$ [pb]	$\sigma_{\rm fid}(pp \to e^{\mp} \nu_e \mu^{\pm} \nu_{\mu})$ [fb]	$A = \sigma_{fid}/\sigma_{incl}$ [%]
LO	$70.66(1)_{-6.2\%}^{+5.1\%}$	$440.5(0)^{+6.0\%}_{-7.1\%}$	0.623
NLO	$ \begin{array}{c} 70.66(1)_{-6.2\%}^{+5.1\%} \\ 99.96(3)_{-2.8\%}^{+3.5\%} \\ 110.0(1)_{-1.6\%}^{+1.6\%} \end{array} $	$ \begin{array}{c} 440.5(0) - 7.1\% \\ -6.5\% \\ 411.8(1) + 2.7\% \\ +0.3\% \\ 413.1(2) + 1.0\% \\ -0.7\% \end{array} $ +0.3%	0.412
NNLO	$110.0(1)_{-1.6\%}^{+1.6\%}$	$413.1(2)_{-0.7\%}^{+1.0\%}$	0.376
MINLO	$96.05(1)_{-4.9\%}^{+7.1\%}$ $110.2(2)_{-1.6\%}^{+1.7\%}$	$359.6(1)^{+5.4\%}_{-8.3\%}$	0.374
NNLOPS	$110.2(2)_{-1.6\%}^{+1.7\%}$	$413.0(2)_{-2.3\%}^{+2.2\%}$	0.375
ATLAS-gg [9]	$124.7 \pm 5 (\mathrm{stat}) \pm 13 (\mathrm{syst}) \pm 3 (\mathrm{lumi})$	$473 \pm 20 (\mathrm{stat}) \pm 50 (\mathrm{syst}) \pm 11 (\mathrm{lumi})$	0.379
CMS-gg [10]	$108.5 \pm 5.8 (\mathrm{stat}) {}^{\pm 5.7 (\mathrm{exp. syst})}_{\pm 6.4 (\mathrm{theo. syst})} \pm 3.6 (\mathrm{lumi})$		

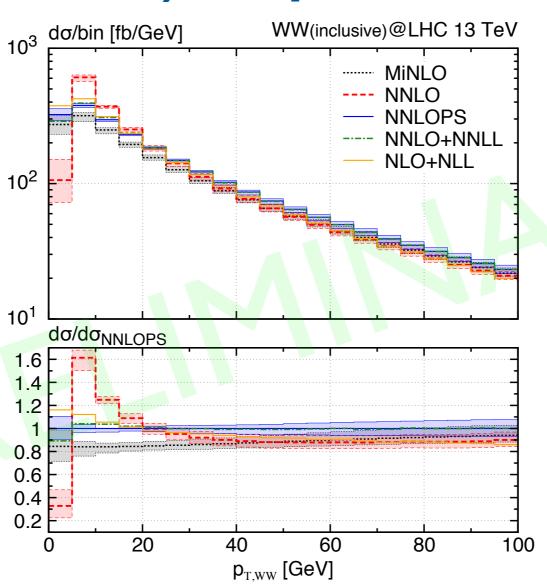
- → inclusive: large QCD corrections; fiducial: strongly reduced by the jet veto
- → MiNLO, NLO quite different in fiducial ← poor jet-veto modelling at NLO (acceptance too high)
- → NNLO, NNLOPS in excellent agreement (by construction for inclusive, but also in fiducial)
- → MiNLO, NNLO, NNLOPS yield very similar acceptances, in agreement with data



[Re, MW, Zanderighi '18]

Phenomenological results:

рт,ww (IR sensitive) compared to NNLO+NNLL



not completely fair comparison yet:

- on-shell WW for analytic resummation
- slightly different setups
- → full comparison will be done

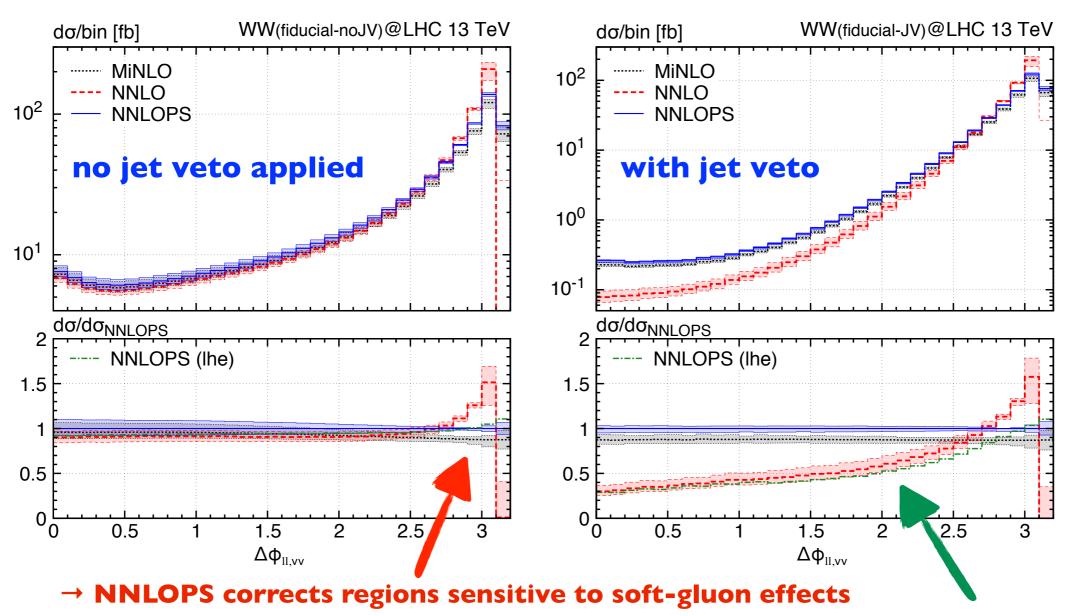
ightarrow Resummation (analytic or shower) crucial at low p $_{
m T}$; NNLOPS in decent agreement with NNLL



[Re, MW, Zanderighi '18]

Phenomenological results:

 $\Delta \Phi_{\ell\ell,vv}$ (IR sensitive)

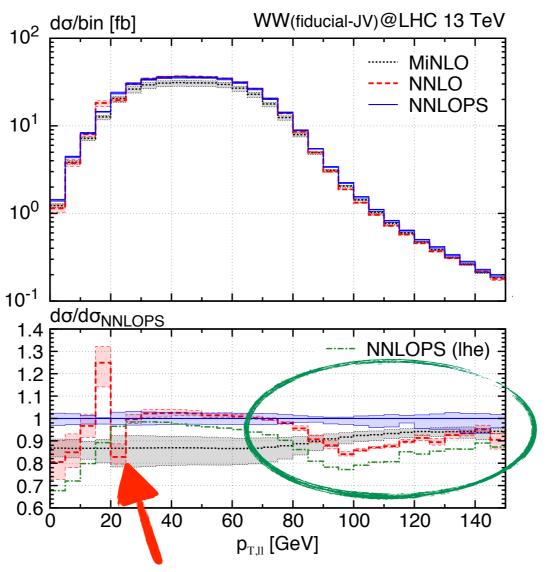


→ jet veto can turn observables sensitive soft-gluon emissions everywhere



[Re, MW, Zanderighi '18]

Phenomenological results: p_T of dilepton system



Many other distributions shown in Re, MW, Zanderighi; arXiv:1805.09857

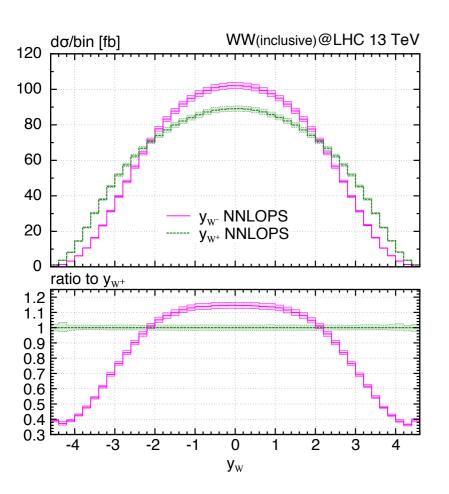
- → NNLOPS cures perturbative instabilities (sudakov shoulder; due to fiducial missing p_T cut)
- → NNLOPS can induce additional shape effects due to recoil/migration

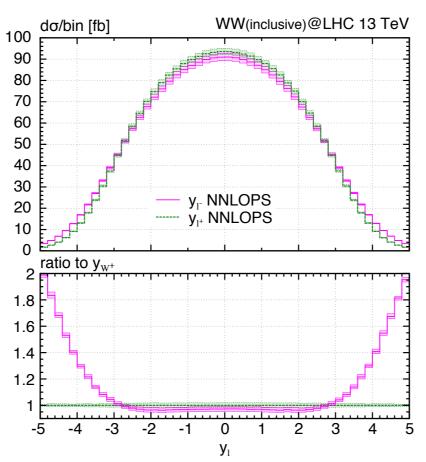
Marius Wiesemann (MPI Munich)

 $\Delta_{p}.\Delta_{q} \geqslant \pm t$

[Re, MW, Zanderighi '18]

Phenomenological results: Charge asymmetry





- W momentum cannot be reconstructed → use leptons
- lepton asymmetry smaller; almost vanishes in fiducial
- can be recovered by widening rapidity range of leptons or by considering boosted regime
- sensitive to W polarizations→ powerful probe of new physics

$$A_C^W = \frac{\sigma(|y_{W^+}| > |y_{W^-}|) - \sigma(|y_{W^+}| < |y_{W^-}|)}{\sigma(|y_{W^+}| > |y_{W^-}|) + \sigma(|y_{W^+}| < |y_{W^-}|)},$$

$$A_C^{\ell} = \frac{\sigma(|y_{\ell^+}| > |y_{\ell^-}|) - \sigma(|y_{\ell^+}| < |y_{\ell^-}|)}{\sigma(|y_{\ell^+}| > |y_{\ell^-}|) + \sigma(|y_{\ell^+}| < |y_{\ell^-}|)}.$$

NNLOPS	inclusive phase space	fiducial phase space
A_C^W	$0.1263(1)_{-1.8\%}^{+2.1\%}$	$0.0726(3)_{-2.6\%}^{+2.0\%}$
A_C^ℓ	$-[0.0270(1)_{-6.4\%}^{+5.0\%}]$	$-[0.0009(4)^{+72\%}_{-87\%}]$

q_T subtraction



$$d\sigma_{\mathrm{NNLO}} = \left[d\sigma_{\mathrm{NLO}}^{F+1\mathrm{jet}} - \Sigma_{\mathrm{NNLO}} \otimes d\sigma_{\mathrm{LO}} \right] + \mathcal{H}_{\mathrm{NNLO}} \otimes d\sigma_{\mathrm{LO}}$$

[Catani, Grazzini '07]

subtraction terms known from resummation:

$$d\sigma^{F+1\text{jet}} \xrightarrow{p_T \ll Q} \left[d\sigma^{(\text{res})} \right]_{\text{f.o.}} \equiv \Sigma(p_T/Q) \otimes d\sigma_{\text{LO}}$$

NNLO accuracy consequence of unitarity:

$$\int dp_T^2 \frac{d\sigma^{(\text{res})}}{dp_T^2 dy dM d\Omega} = \mathcal{H} \otimes d\sigma_{\text{LO}} \qquad \left(\ln(Q^2 b^2 / b_0^2) \to \ln(Q^2 b^2 / b_0^2 + 1)\right)$$

Resummation formula:

$$rac{d\sigma^{({
m res})}}{dp_T^2\,dy\,dM\,d\Omega} \sim \int db\,rac{b}{2}\,J_0(b\,p_T)\,S(b,A,B)\,{\cal H}_{N_1,N_2}\,f_{N_1}\,f_{N_2}$$

[Collins, Soper, Sterman '85], [Bozzi, Catani, de Florian, Grazzini '06]

q_T subtraction



$$d\sigma_{\mathrm{NNLO}} = \left[d\sigma_{\mathrm{NLO}}^{F+1\mathrm{jet}} - \Sigma_{\mathrm{NNLO}} \otimes d\sigma_{\mathrm{LO}} \right] + \mathcal{H}_{\mathrm{NNLO}} \otimes d\sigma_{\mathrm{LO}}$$

[Catani, Grazzini '07]

practical implementation:

- subtraction not local
- both terms in squared brackets separately divergent
- introduce lower cut-off r_{cut} on dimensionless quantity $r = p_{T,WW}/m_{WW}$
- $^{\textcircled{6}}$ use very small r_{cut} value and integrate both terms separately down to $r \geq r_{cut}$
- $^{\otimes}$ assumption: for $r \leq r_{cut}$ terms cancel (true up to power-suppressed terms)
- numerics forbids arbitrarily small r_{cut} values: use fit towards $r_{cut} \rightarrow 0$ limit
- MATRIX uses extrapolation $r_{cut} \rightarrow 0$ to obtain the final prediction

q_T subtraction



$$d\sigma_{\rm NNLO} = \begin{bmatrix} d\sigma_{\rm NLO}^{F+1\rm jet} - \Sigma_{\rm NNLO} \otimes d\sigma_{\rm LO} \\ \text{(r > r_{\rm cut)}} \end{bmatrix} + \mathcal{H}_{\rm NNLO} \otimes d\sigma_{\rm LO}$$

[Catani, Grazzini '07]

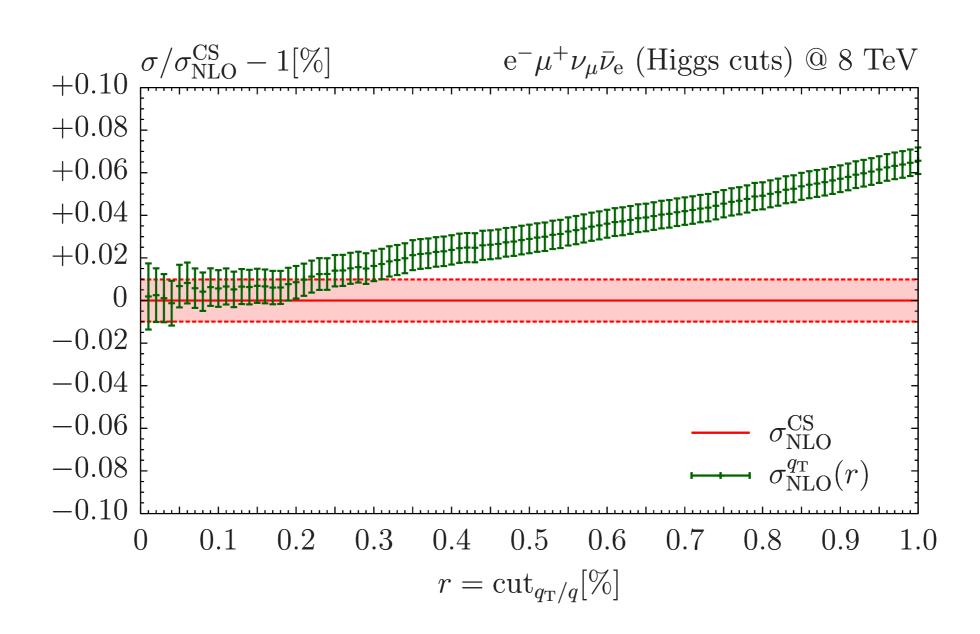
practical implementation:

- subtraction not local
- both terms in squared brackets separately divergent
- introduce lower cut-off r_{cut} on dimensionless quantity $r = p_{T,WW}/m_{WW}$
- $^{\odot}$ use very small r_{cut} value and integrate both terms separately down to $r \ge r_{cut}$
- \bullet assumption: for $r \le r_{cut}$ terms cancel (true up to power-suppressed terms)
- numerics forbids arbitrarily small r_{cut} values: use fit towards $r_{cut} \rightarrow 0$ limit
- MATRIX uses extrapolation $r_{cut} \rightarrow 0$ to obtain the final prediction

r_{cut} dependence at NLO



(example from WW)

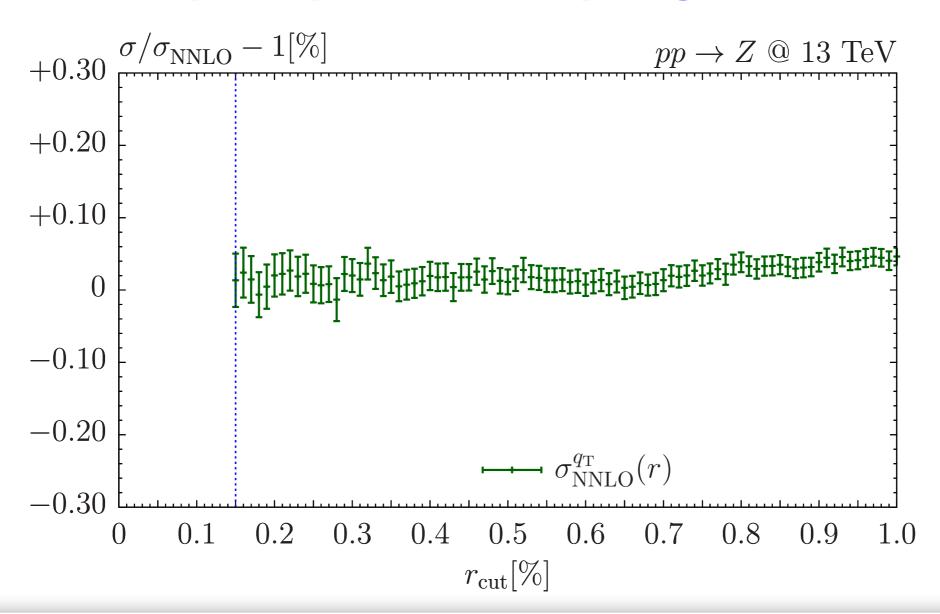


$$d\sigma_{\mathrm{NNLO}}(r_{\mathrm{cut}}) = \left[d\sigma_{\mathrm{NLO}}^{F+\mathrm{jet}}(r > r_{\mathrm{cut}}) - \Sigma_{\mathrm{NNLO}} \otimes d\sigma_{\mathrm{LO}}(r > r_{\mathrm{cut}}) \right] + \mathcal{H}_{\mathrm{NNLO}} \otimes d\sigma_{\mathrm{LO}}$$



[Grazzini, Kallweit, MW '17]

automatically computed in every single MATRIX NNLO run

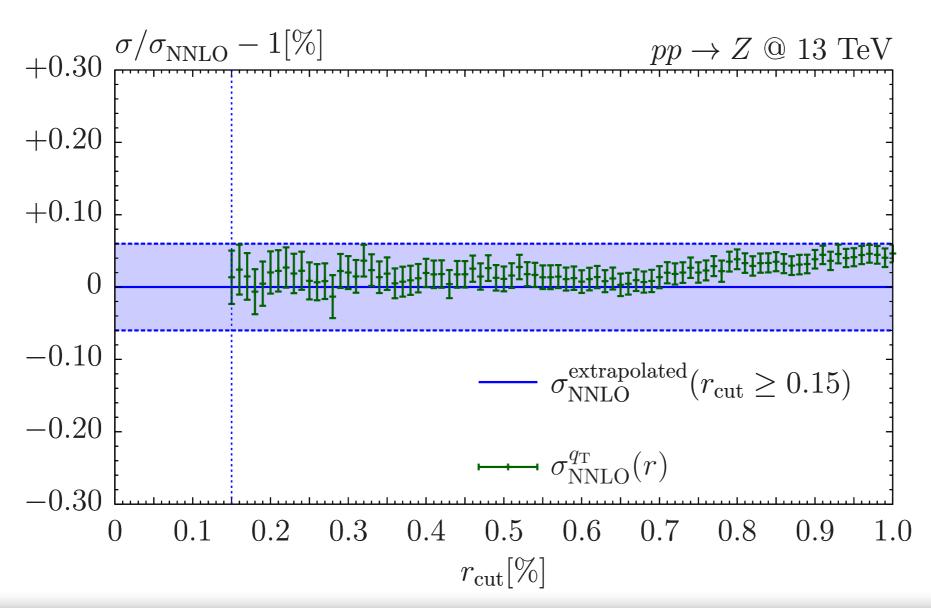


$$d\sigma_{\mathrm{NNLO}}(r_{\mathrm{cut}}) = \left[d\sigma_{\mathrm{NLO}}^{F+\mathrm{jet}}(r > r_{\mathrm{cut}}) - \Sigma_{\mathrm{NNLO}} \otimes d\sigma_{\mathrm{LO}}(r > r_{\mathrm{cut}}) \right] + \mathcal{H}_{\mathrm{NNLO}} \otimes d\sigma_{\mathrm{LO}}$$



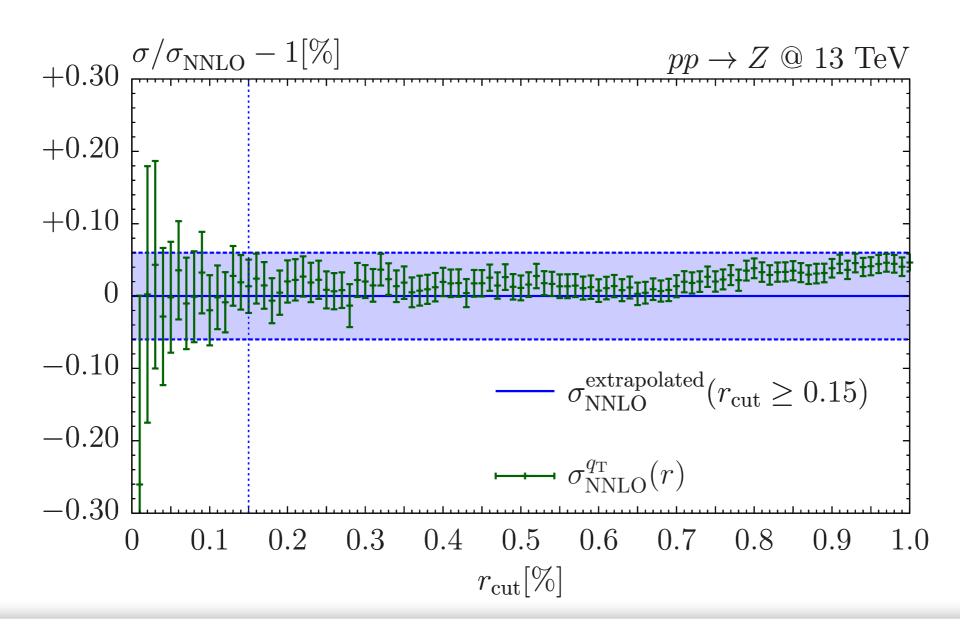
[Grazzini, Kallweit, MW '17]

simple quadratic fit (A * r_{cut}^2 + B * r_{cut} + C) to extrapolate to r_{cut} =0



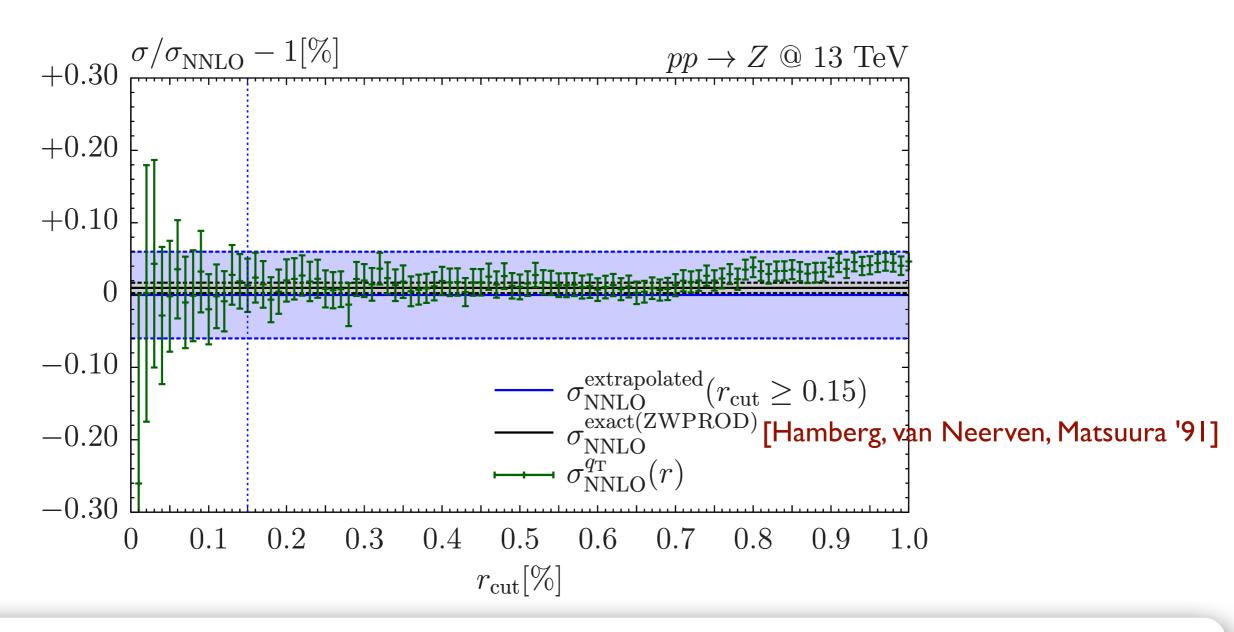
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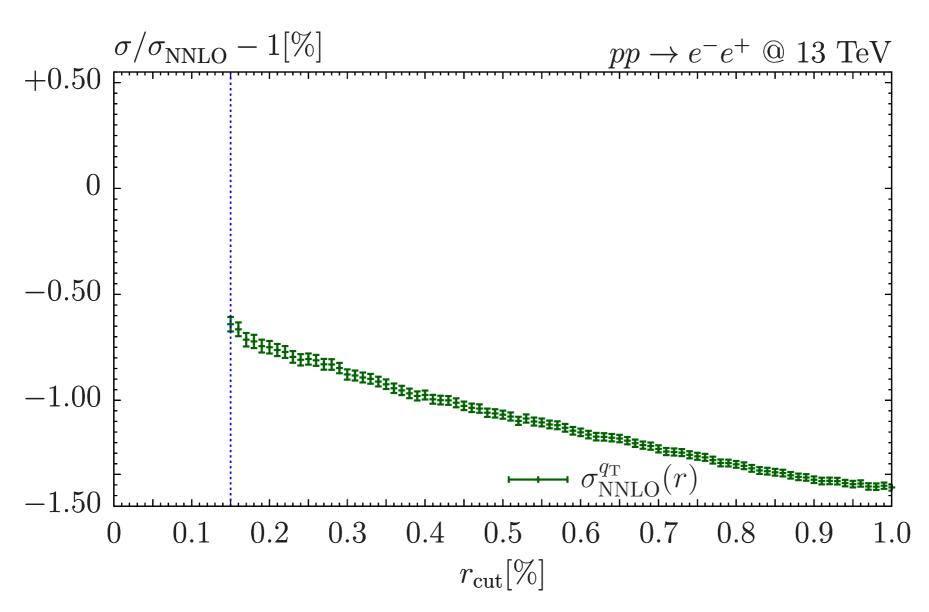


$$d\sigma_{\mathrm{NNLO}}(r_{\mathrm{cut}}) = \left[d\sigma_{\mathrm{NLO}}^{F+\mathrm{jet}}(r > r_{\mathrm{cut}}) - \Sigma_{\mathrm{NNLO}} \otimes d\sigma_{\mathrm{LO}}(r > r_{\mathrm{cut}}) \right] + \mathcal{H}_{\mathrm{NNLO}} \otimes d\sigma_{\mathrm{LO}}$$

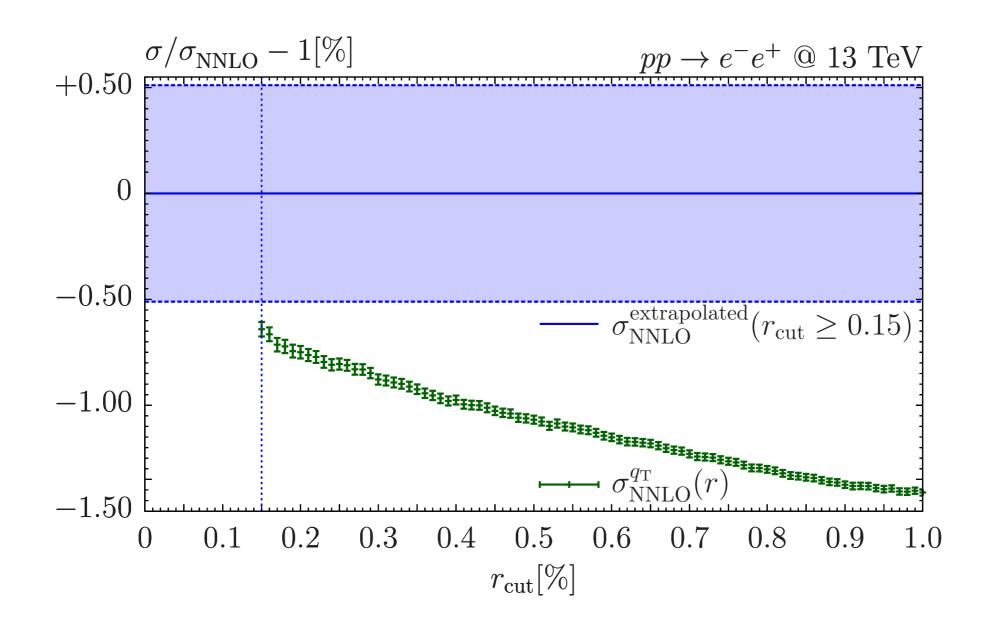


[Grazzini, Kallweit, MW '17]

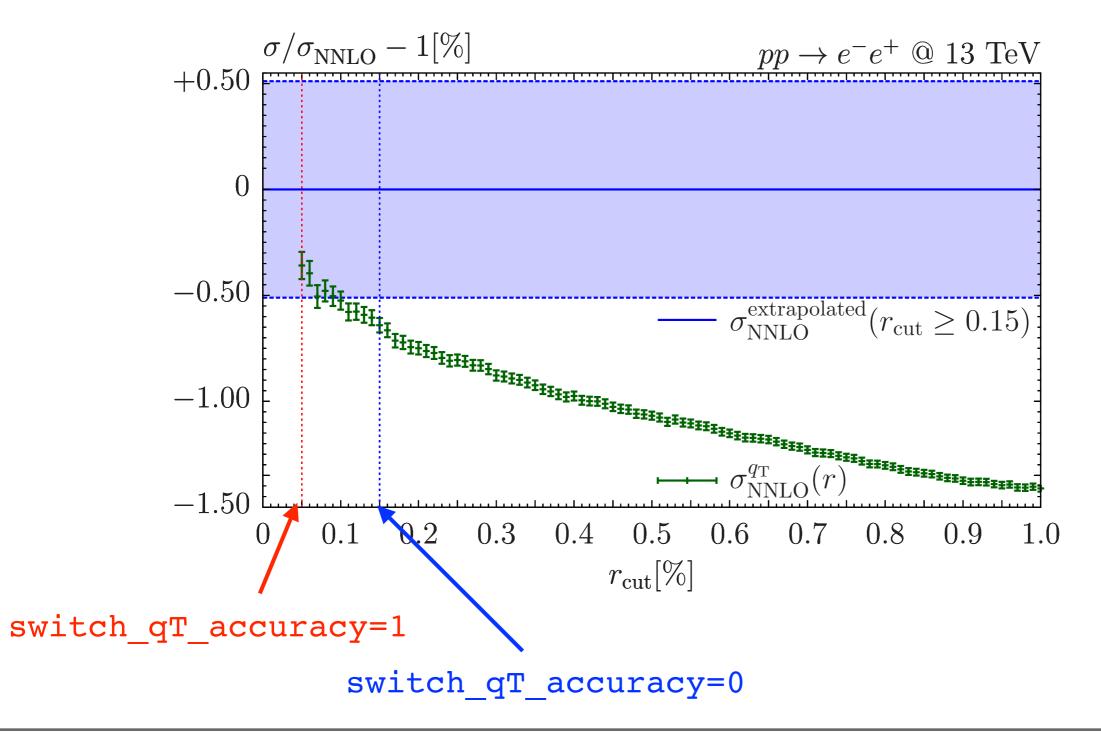
dileptons with certain cuts (and photon final states) are special



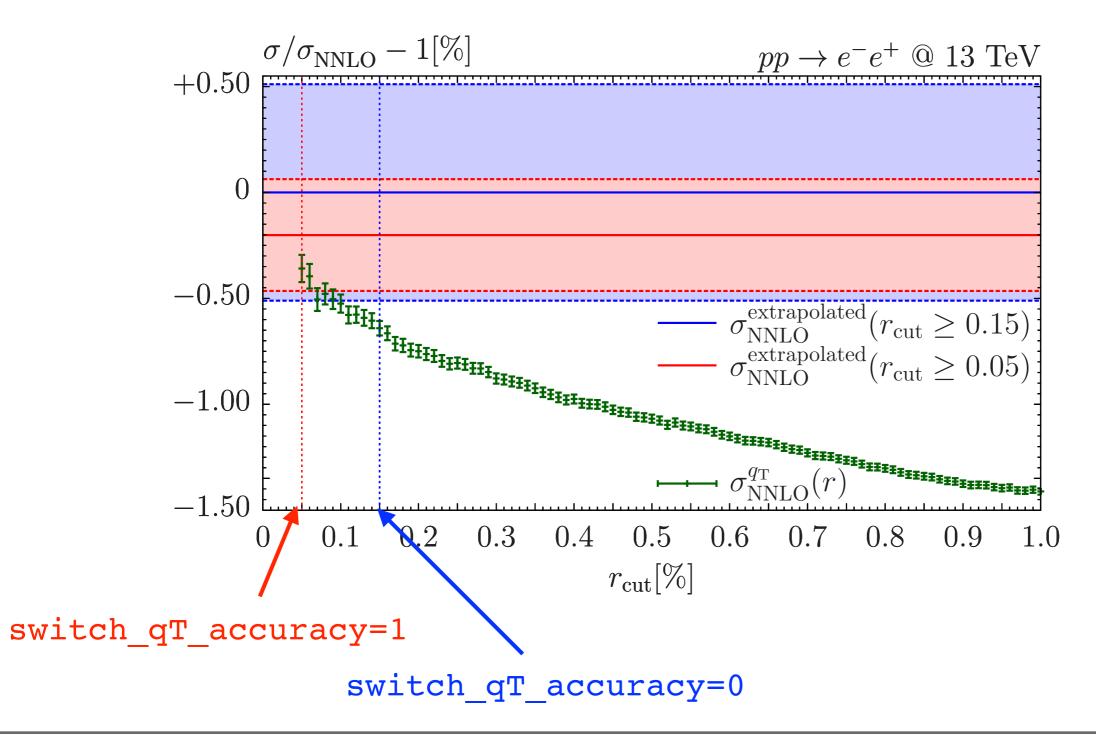




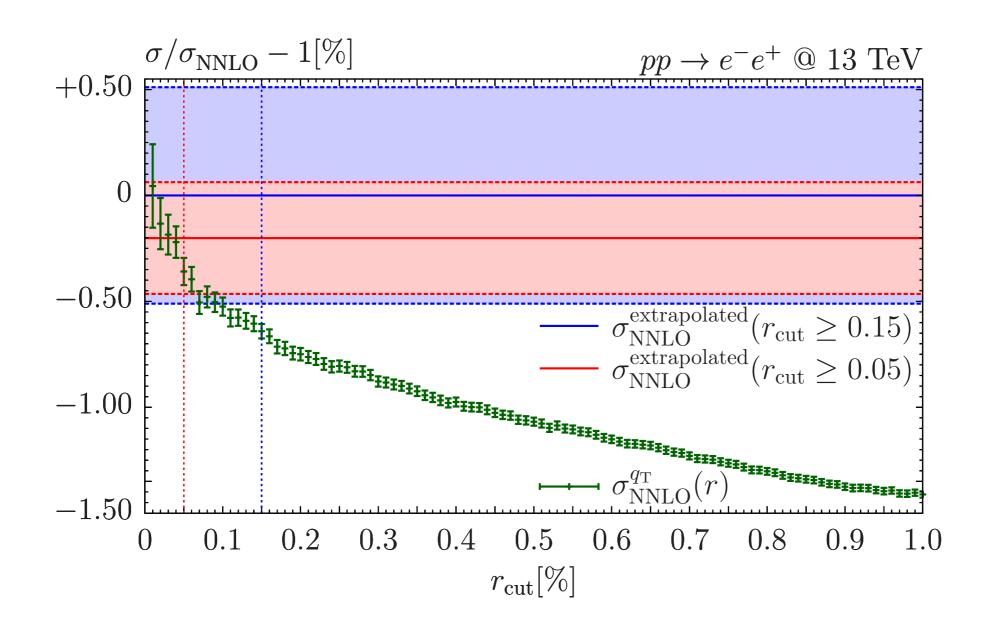




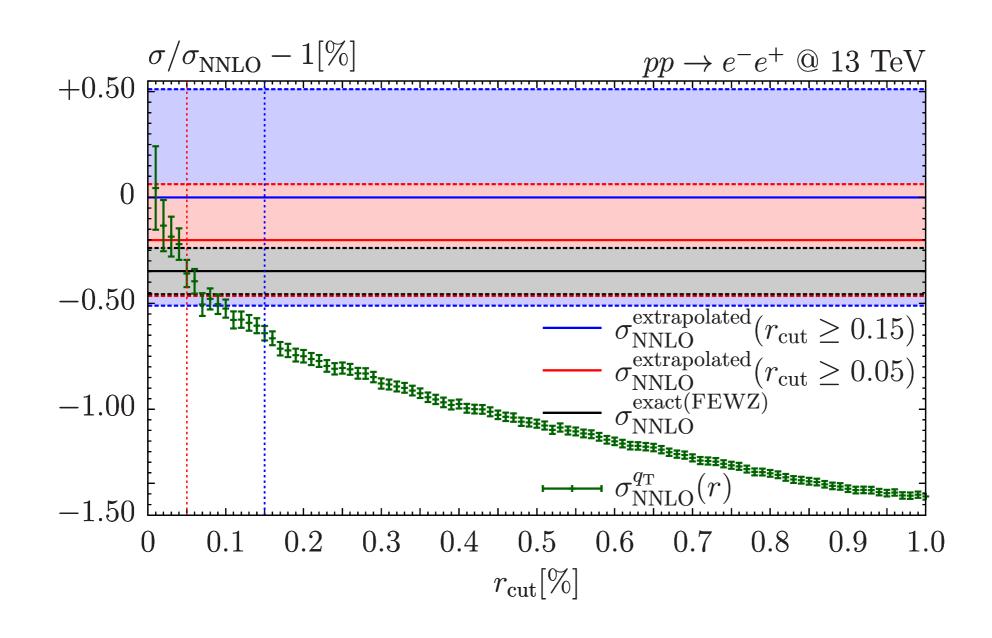








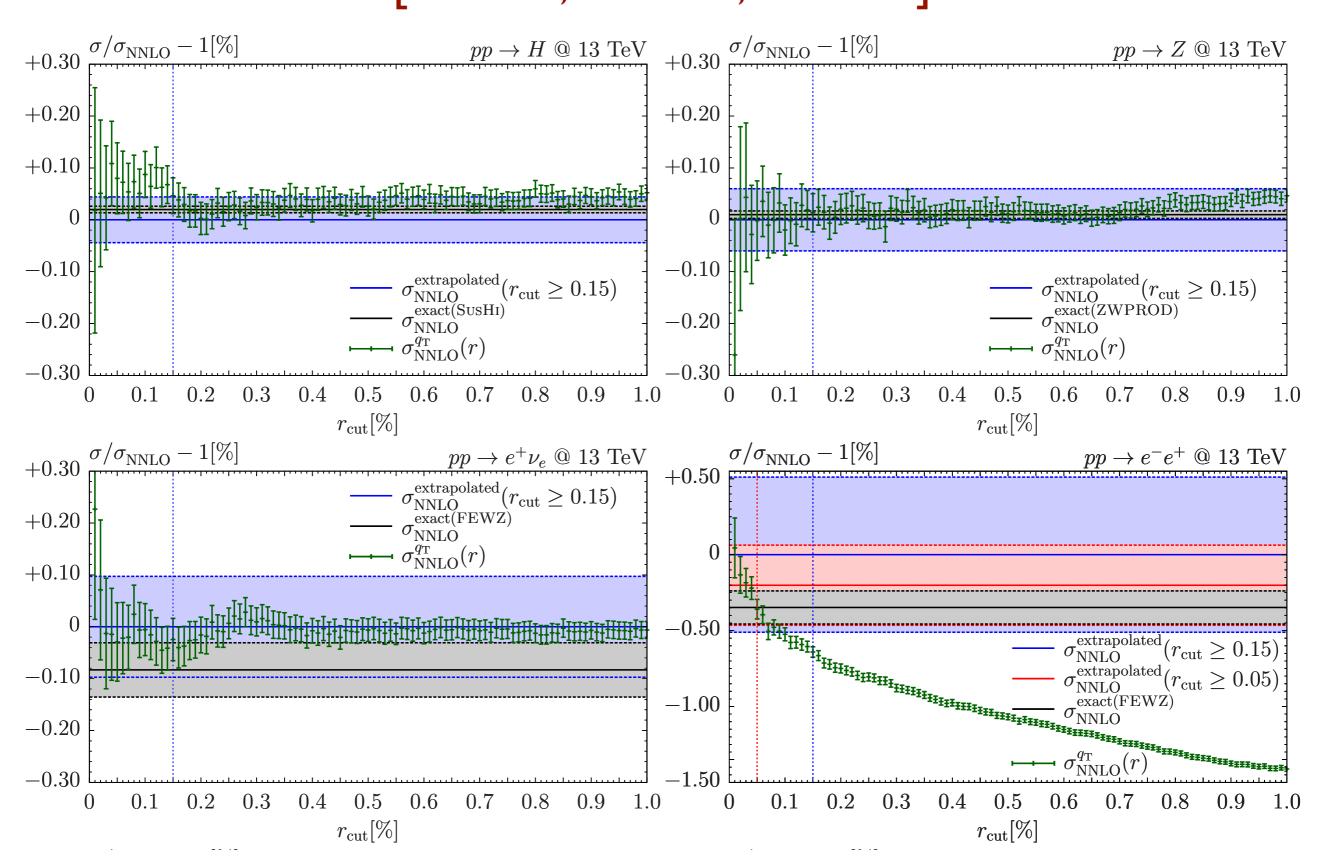


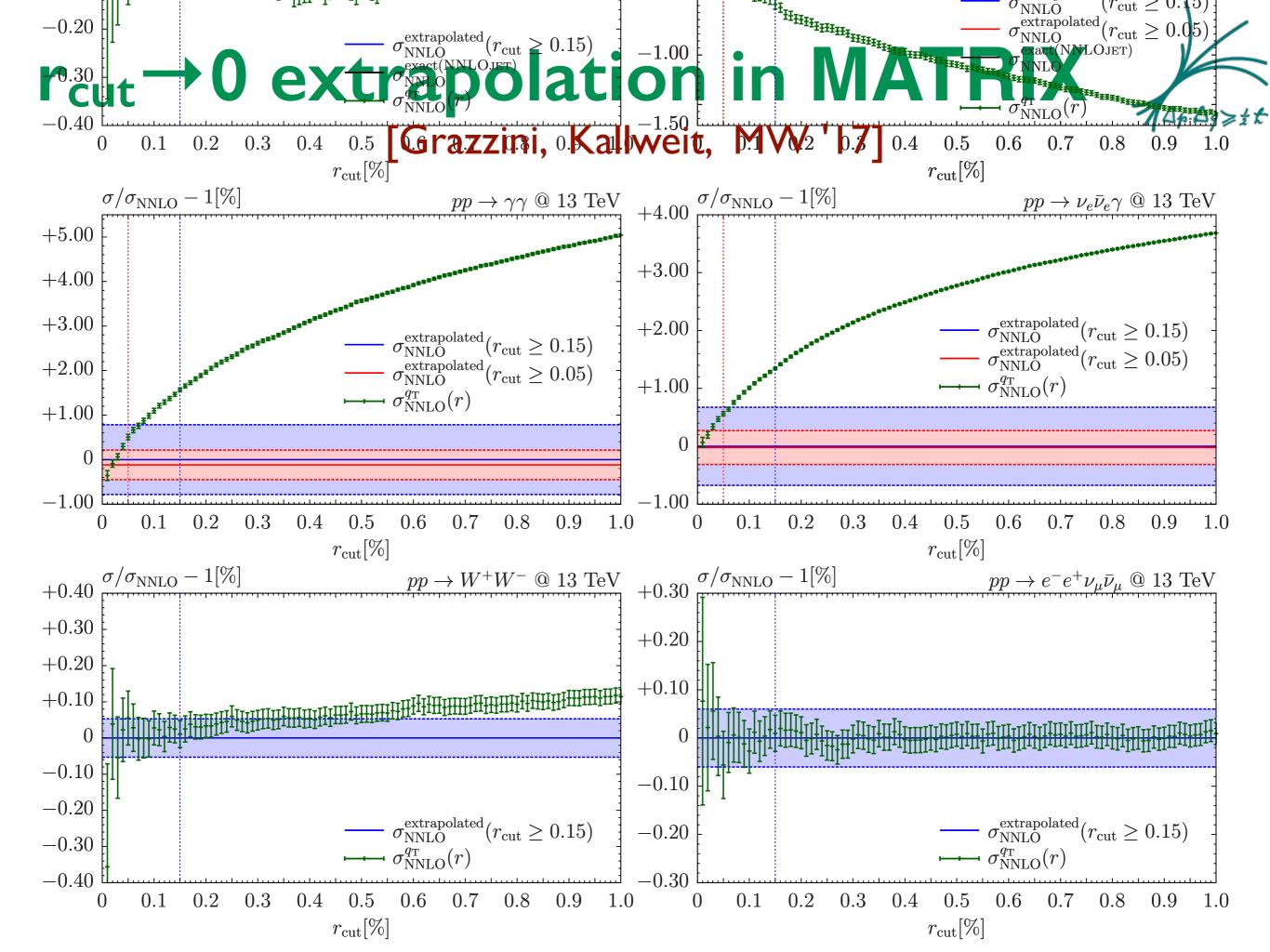


r_{cut}→0 extrapolation in MATRIX



[Grazzini, Kallweit, MW '17]





MATRIX features on one slide



 \circ Colourless 2→ I and 2→2 reactions (decays, off-shell effects, spin correlations; previous slide)

physics features:

- NNLO accuracy based on q_T subtraction
- loop-induced gg component part of NNLO cross section (effectively LO accurate)
- © CKM for W-boson production
- essential fiducial cuts, dynamical scales and distributions already pre-defined for each process
- final-state particles directly accessible (for distributions, cuts, scales)
- scale uncertainty estimated automatically estimated (7- or 9-point) with every run
- NEW: automatic extrapolation of q_T-subtraction cut-off to zero (with extrapolation uncertainty)

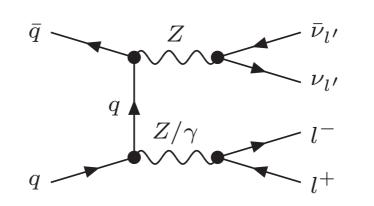
technical features:

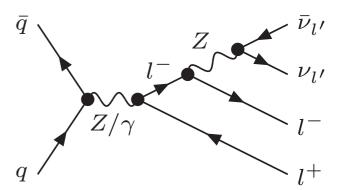
- © Core: C++ code; steered by Python interface (compilation/running/job submission/result collection)
- only requirements: LHAPDF 5 or 6 pre-installed & Python 2.7 with numpy
- Otherwise fully automatic! (download/compilation of external packages; inputs via interface etc.)
- local and cluster support: LSF (Ixplus), HT-Condor (Ixplus), condor, SLURM, Torque/PBS, SGE
 - → missing your favourite cluster? Let us know!
- option to reduce workload (output) on slow file systems
- all relevant references in CITATION.bib (provided with every run)
- comprehensive manual shipped with the code



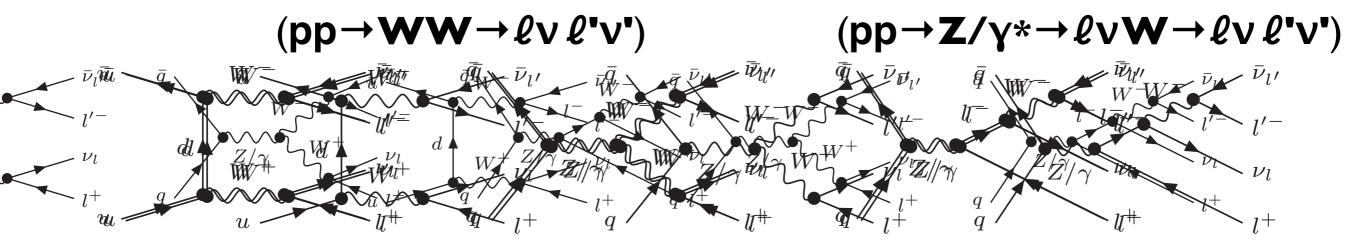
DF (ℓℓν'ν'): double-resonant ZZ

(pp
$$\rightarrow$$
ZZ/ $\gamma*Z\rightarrow\ell\ell\nu'\nu'$)
and single-resonant DY-like
(pp \rightarrow Z/ $\gamma*\rightarrow\ell\ell$ Z $\rightarrow\ell\ell\nu'\nu'$)





• DF $(\ell \nu \ell' \nu')$: double-resonant WW and single-resonant DY-like



• SF ($\ell\ell\nu\nu$): mixes all ZZ, WW and DY-like contributions

$$(pp \rightarrow ZZ/\gamma * Z/WW \rightarrow \ell\ell vv) \qquad (pp \rightarrow Z/\gamma * \rightarrow \ell\ell Z/\ell vW \rightarrow \ell\ell vv)$$



Comparison against ATLAS 8 TeV data [JHEP 1701 (2017) 099]

fiducial cuts:

definition of the fiducial volume for $pp \to \ell^+\ell^-\nu\bar{\nu} + X$, $\ell \in \{e, \mu\}$ and $\nu \in \{\nu_e, \nu_\mu, \nu_\tau\}$

$$p_{T,\ell} > 25 \,\text{GeV}, \quad |\eta_{\ell}| < 2.5, \quad \Delta R_{\ell\ell} > 0.3, \quad 76 \,\text{GeV} \le m_{\ell^+\ell^-} \le 106 \,\text{GeV},$$

Axial-
$$p_T^{\text{miss}} > 90 \,\text{GeV}, \quad p_T\text{-balance} < 0.4$$

$$Axial-p_T^{miss}>90\,\text{GeV},\quad p_T\text{-balance}<0.4,$$

$$N_{\text{jets}}=0,\quad \text{anti-}k_T \text{ jets with } R=0.4, \, p_{T,j}>25\,\text{GeV}, \, |\eta_j|<4.5 \text{ and } \Delta R_{ej}>0.3$$

jet veto

prefers Z bosons in backto-back like configurations



Comparison against ATLAS 8 TeV data [JHEP 1701 (2017) 099]

our setup:

- $\mu_R = \mu_F = \frac{m_{ZZ}}{2}$ (uncertainties: 7-point scale variation)
- NNPDF 3.0 (set consistent at each perturbative order) with $\alpha_s(m_Z) = 0.118$
- 5-flavour scheme



Comparison against ATLAS 8 TeV data [JHEP 1701 (2017) 099]

our setup:

- $\mu_R = \mu_F = \frac{m_{ZZ}}{2}$ (uncertainties: 7-point scale variation)
- NNPDF 3.0 (set consistent at each perturbative order) with $\alpha_s(m_Z) = 0.118$
- 5-flavour scheme
 - \longrightarrow top-quark contamination included in SF $(\ell\ell\nu_{\ell}\nu_{\ell})$ channel

BUT: suppressed by jet veto

AND: $W^+W^ (\ell\nu_\ell\ell'\nu_{\ell'})$ and top contributions subtracted as background

- \longrightarrow compute SF $\ell\ell\nu_{\ell}\nu_{\ell}$ channel and subtact DF $\ell\nu_{\ell}\ell'\nu_{\ell'}$ (also in 5FS)
- → removes subtracted backgrounds, but keeps all interference contributions!

$$\sigma(\ell\ell \nu_{e/\mu/\tau}\nu_{e/\mu/\tau}) = \sigma(\ell\ell \nu_{\ell}\nu_{\ell}) - \sigma(\ell\nu_{\ell}\ell'\nu_{\ell'}) + 2 \cdot \sigma(\ell\ell \nu_{\ell'}\nu_{\ell'})$$



Comparison against ATLAS 8 TeV data

fiducial cross section (POWHEG+gg+EW): [JHEP 1701 (2017) 099]

	Measurement	Prediction
$\sigma_{ZZ\to e^-e^+\nu\bar{\nu}}^{\mathrm{fid}}$	= $5.0^{+0.8}_{-0.7}$ (stat) $^{+0.5}_{-0.4}$ (syst) ± 0.1 (lumi) fb	$3.7 \pm 0.3 \text{ fb}$
$\sigma^{\mathrm{fid}}_{ZZ \to \mu^- \mu^+ \nu \bar{\nu}}$	= $4.7 \pm 0.7 \text{ (stat)}$ $^{+0.5}_{-0.4} \text{ (syst)} \pm 0.1 \text{ (lumi)}$ fb	$3.5 \pm 0.3 \text{ fb}$
$\sigma_{pp\to ZZ}^{\text{total}}$	= $7.3 \pm 0.4 \text{ (stat)} \pm 0.3 \text{ (syst)}$ $^{+0.2}_{-0.1} \text{ (lumi) pb}$	6.6 ^{+0.7} _{-0.6} pb



Comparison against ATLAS 8 TeV data

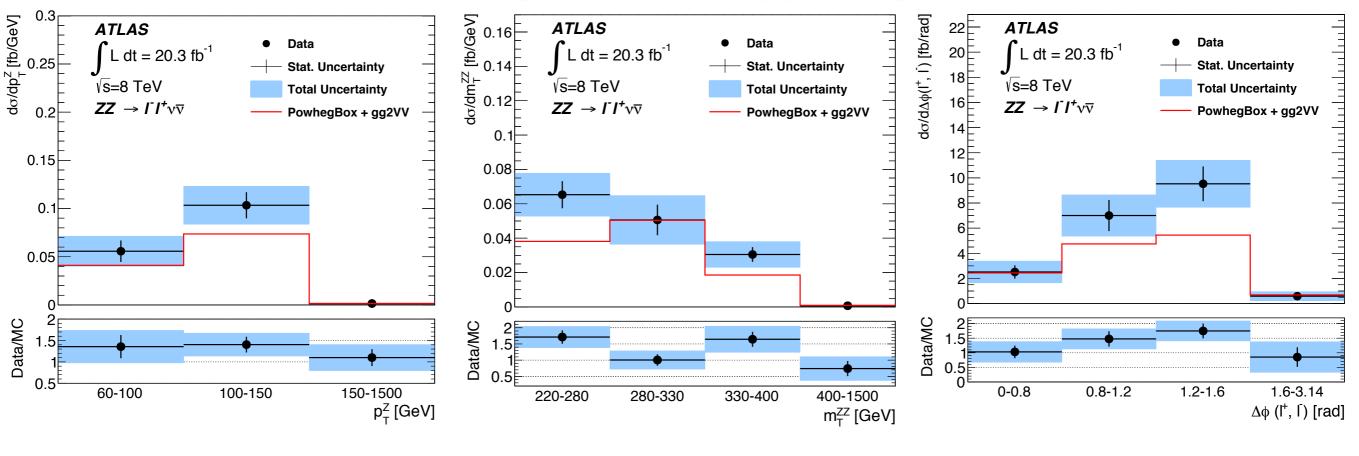
fiducial cross section (POWHEG+gg+EW): [JHEP 1701 (2017) 099]

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fiducial cross section (fixed order): [Kallweit, MW '18]

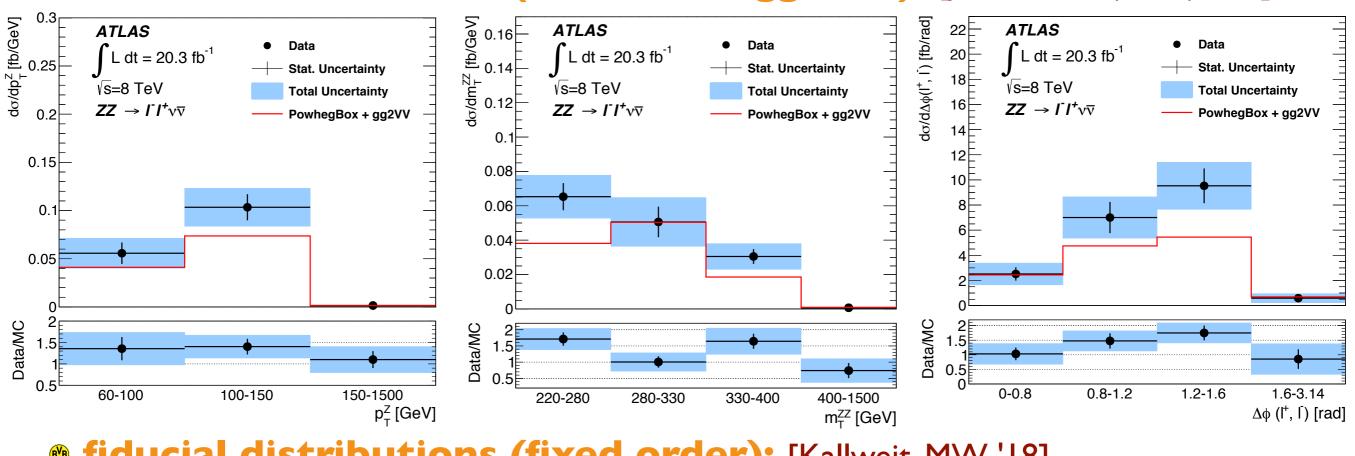
channel	$\sigma_{ m LO} \ [{ m fb}]$	$\sigma_{ m NLO}$ [fb]	$\sigma_{ m NNLO}$ [fb]	$\sigma_{ m ATLAS} \ [{ m fb}]$
$e^+e^- u u$	$5.558(0)_{-0.5\%}^{+0.1\%}$	$4.806(1)_{-3.9\%}^{+3.5\%}$	$5.083(8)_{-0.6\%}^{+1.9\%}$	$5.0^{+0.8}_{-0.7}(\mathrm{stat})^{+0.5}_{-0.4}(\mathrm{syst}) \pm 0.1(\mathrm{lumi})$
$\mu^+\mu^- u u$	$5.558(0)_{-0.5\%}^{+0.1\%}$	$4.770(4)_{-4.0\%}^{+3.6\%}$	$5.035(9)_{-0.5\%}^{+1.8\%}$	$4.7^{+0.7}_{-0.7}(\mathrm{stat})^{+0.5}_{-0.4}(\mathrm{syst}) \pm 0.1(\mathrm{lumi})$
total rate	$4982(0)^{+1.9\%}_{-2.7\%}$	$6754(2)^{+2.4\%}_{-2.0\%}$	$7690(5)^{+2.7\%}_{-2.1\%}$	$7300^{+400}_{-400}(\text{stat})^{+300}_{-300}(\text{syst})^{+200}_{-100}(\text{lumi})$

fiducial distributions (POWHEG+gg+EW): [JHEP 1701 (2017) 099]

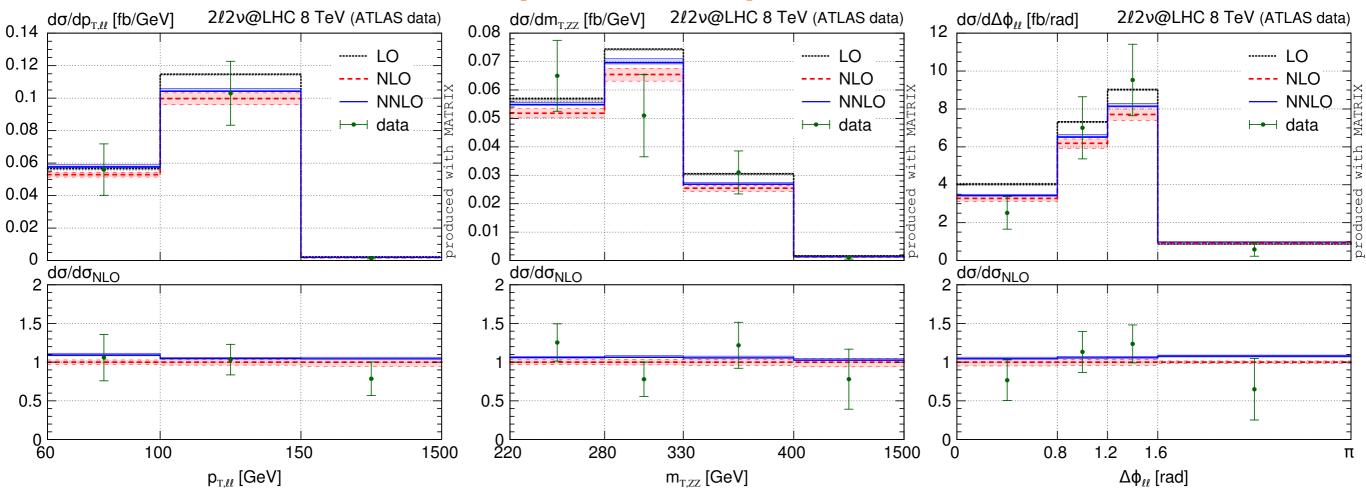


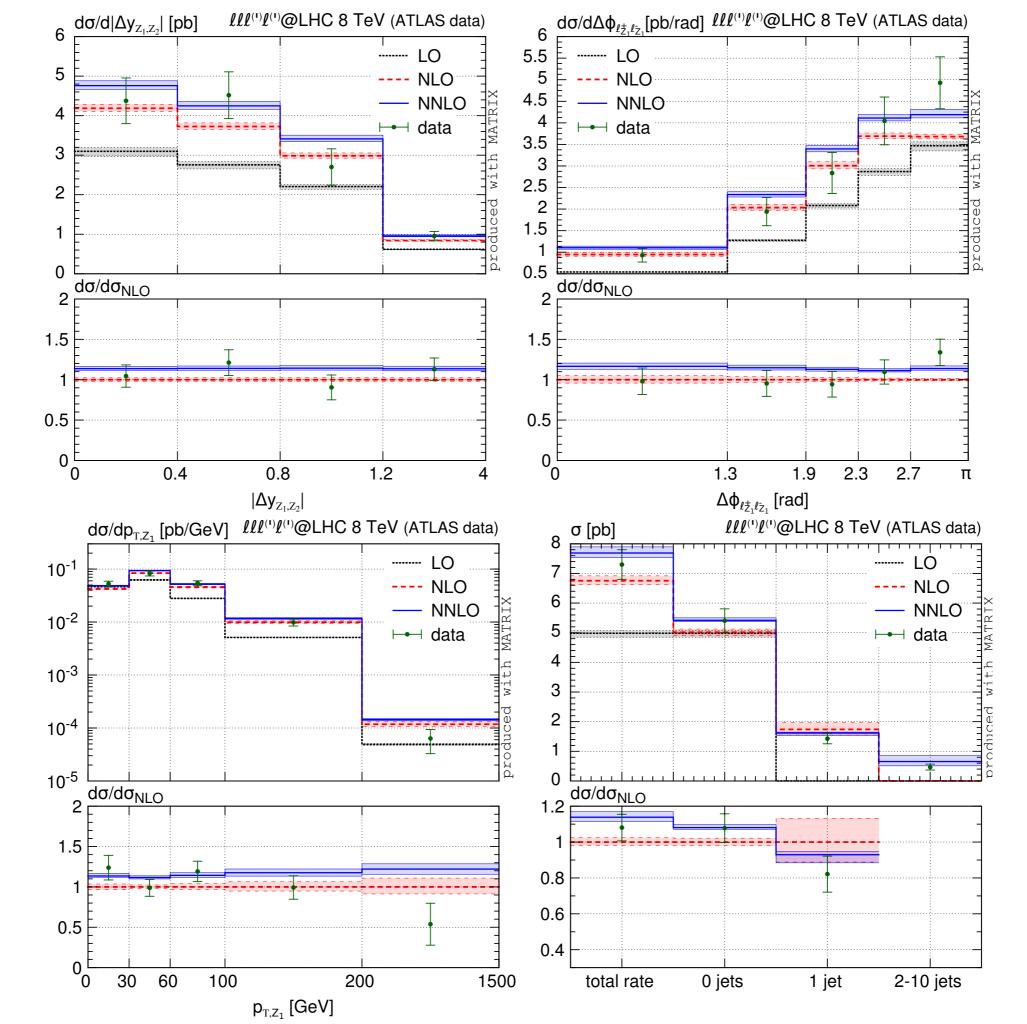
$$m_{T,ZZ} = \sqrt{\left(\sqrt{p_{T,\ell\ell}^2 + m_Z^2} + \sqrt{(p_T^{\text{miss}})^2 + m_Z^2}\right)^2 - (\mathbf{p_T}_{\ell\ell} + \mathbf{p_T^{\text{miss}}})^2}$$

fiducial distributions (POWHEG+gg+EW): [JHEP 1701 (2017) 099]



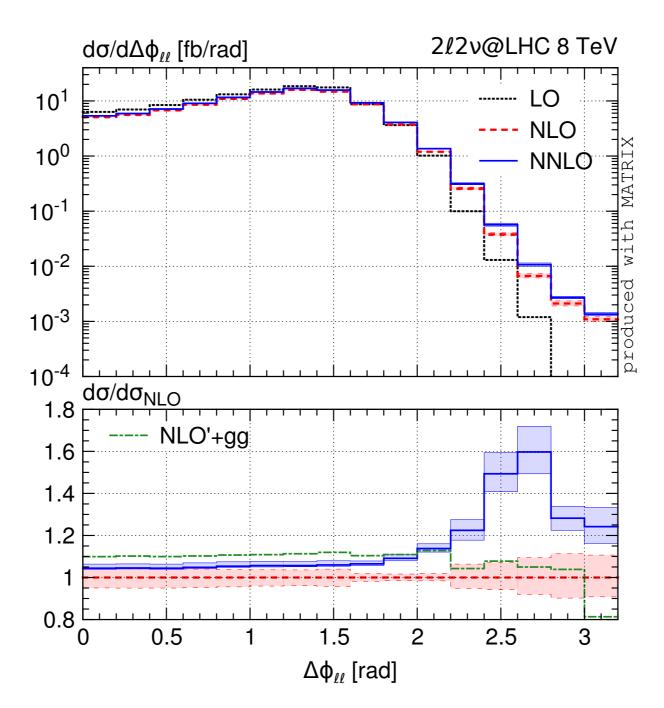
fiducial distributions (fixed order): [Kallweit, MW '18]







Impact of NNLO on fiducial distributions (same cuts as before)





Marius Wiesemann

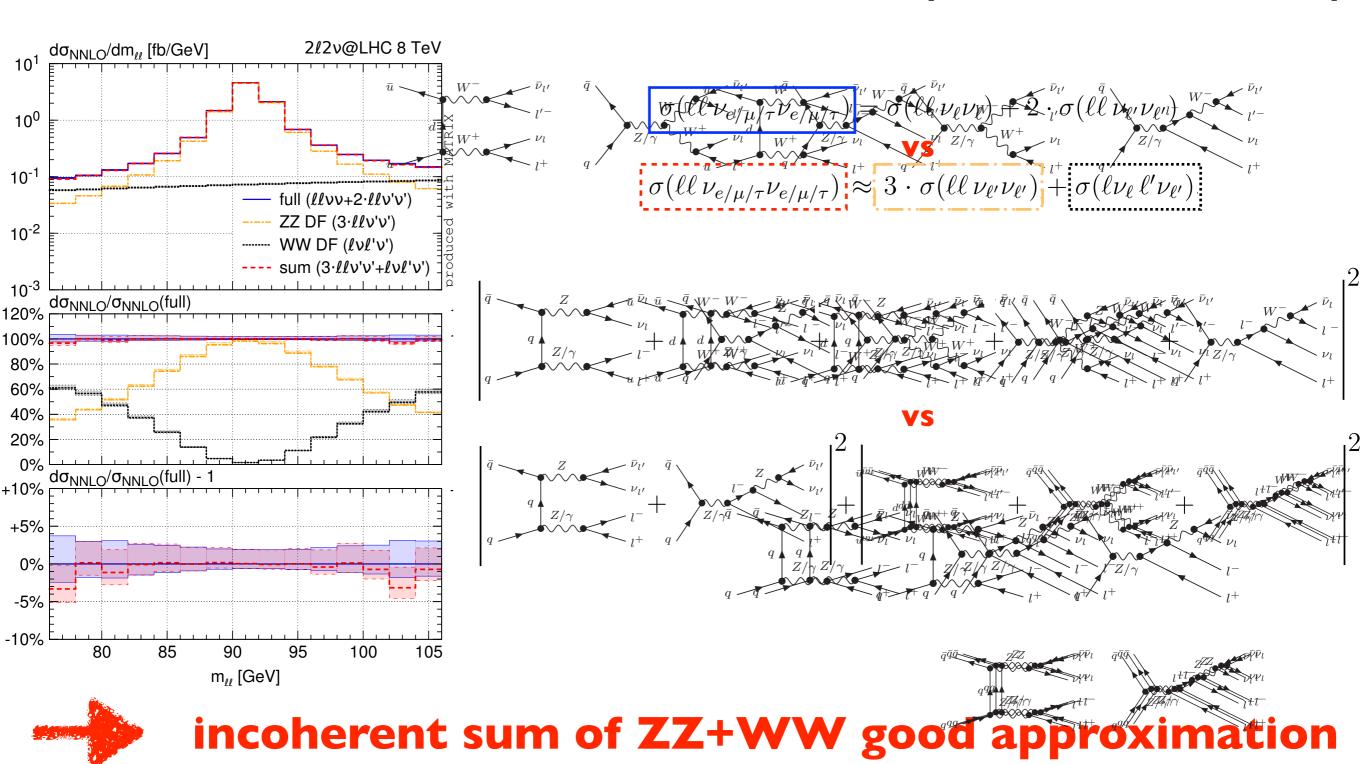
despite jet veto, NNLO corrections quite large

April 24th, 2019

(MPI Munich)

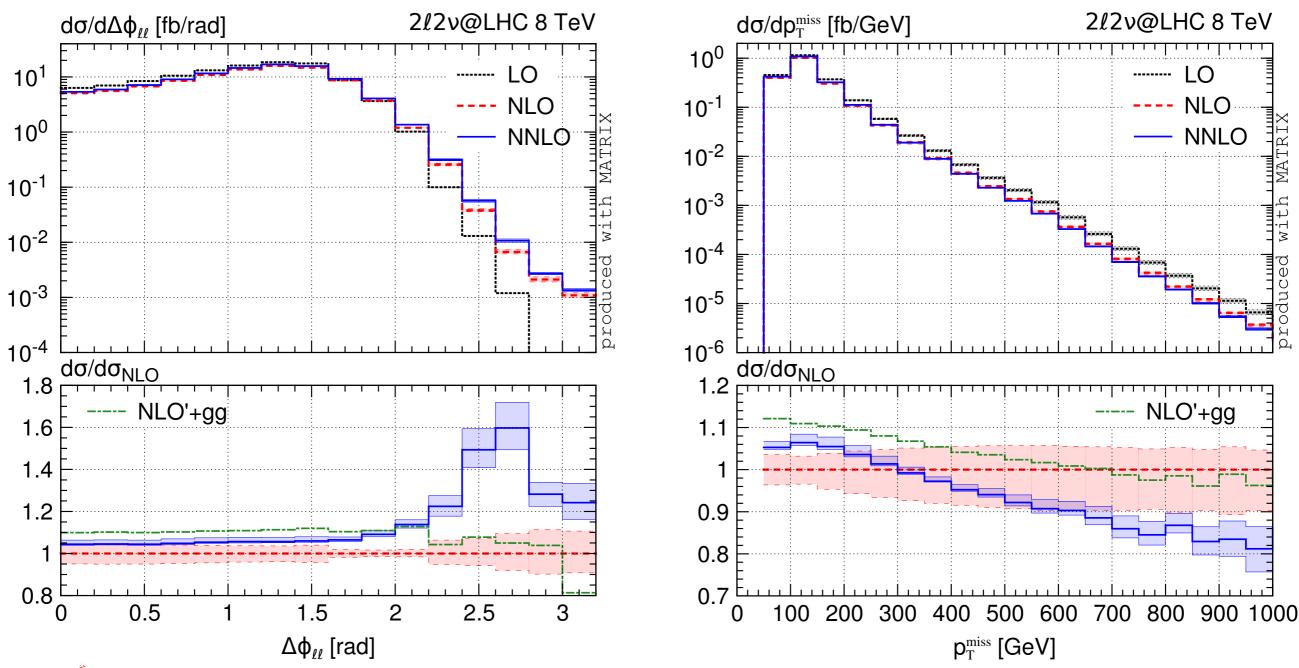


Relative size of ZZ and WW contributions (same cuts as before)





Impact of NNLO on fiducial distributions (same cuts as before)



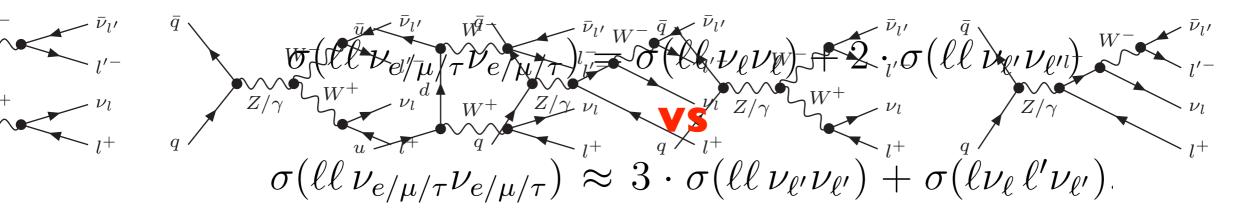


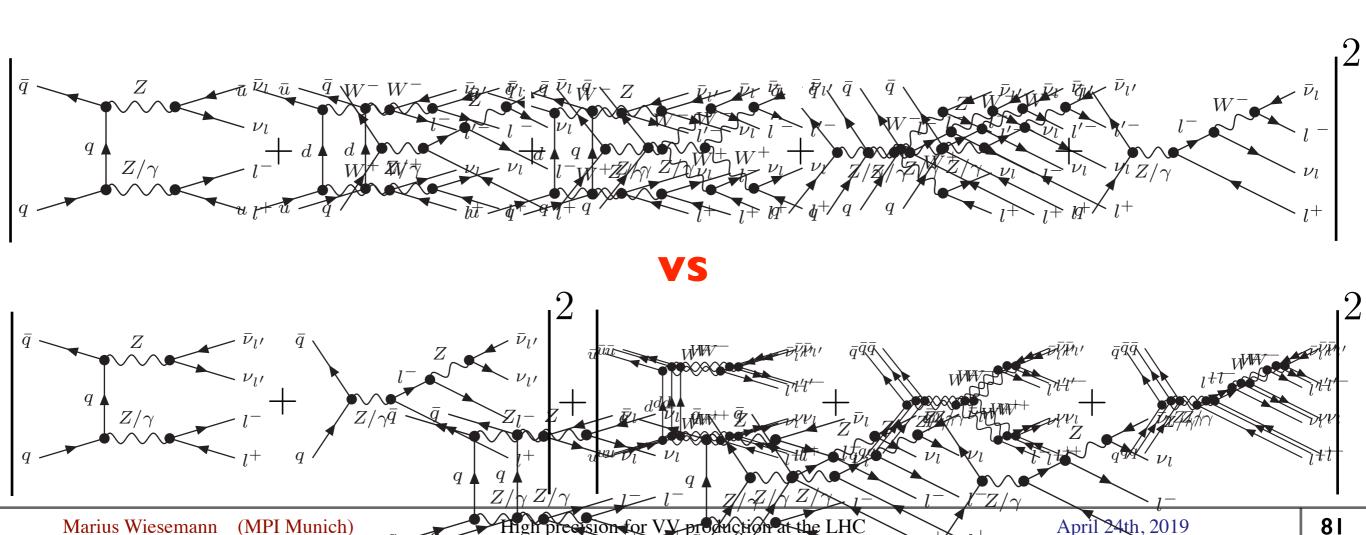
despite jet veto, NNLO corrections quite large

\$\ell \ell \ell + E_{T,miss} at NNLO [Kallweit, MW '18]\$



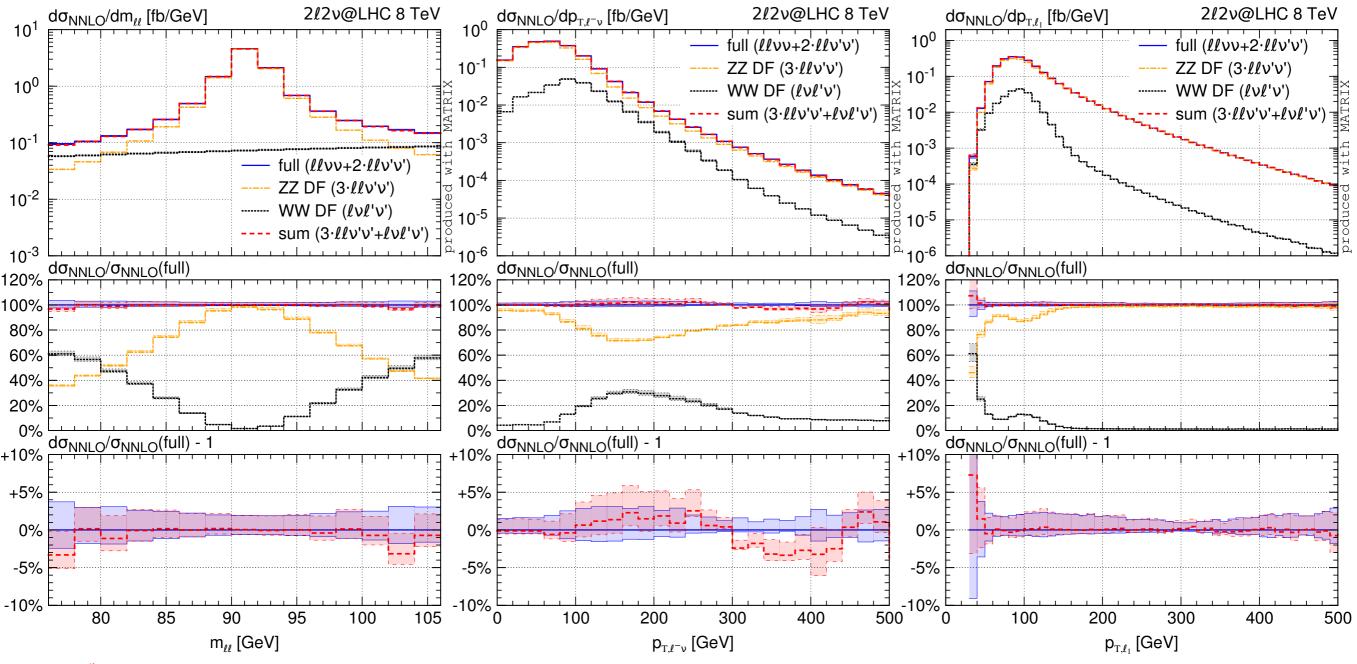
Relative size of ZZ and WW contributions (same cuts as before)







Relative size of ZZ and WW contributions (same cuts as before)





incoherent sum of ZZ+WW good approximation