Current theory status for jet production

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Jet production at the LHC

- look at production of jets of hadrons with large transverse energy in proton-proton collisions
- for sufficiently high transverse momentum p_T > 20 GeV high rates and clean and simple cross section definition
- anti-kt jet definition infrared and collinear safe used in experimental measurements and theory predictions



Jet production: uncertainties at the LHC

Jet measurements have become very precise:

- dominant systematic uncertainty JES ~ 1-2% corresponds to < 10% uncertainty on single jet inclusive cross section
- similar for ATLAS and CMS: 5% systematic on a wide range, and sub-% statistical → allow jet precision physics at the LHC

Ideal testing ground for QCD:

- test perturbative QCD description of jet data at the LHC
- constrain PDFs (sensitive to gluon at LO)





Gluon PDF fractional uncertainty with LHC jet data included CMS (arXiv:1609.5331)

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- determine α_s → huge range in jet p_T directly tests energy dependence/running of α_s from a single experiment





 $\alpha_{\rm s}$ running in the TeV range from LHC jet data

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- a window to new physics e.g. → bumps in dijet mass distributions → precise BSM searches in dijet angular distributions



[[]ATLAS, arXiv:1711.026926]



[ATLAS, arXiv:1703.09127]

Jet observables

- Single jet inclusive cross section: $pp \rightarrow jet + X$
 - inclusive sum of individual jet contributions in the event that pass the jet fiducial cuts
 - → each jet in the event contributes separately, leading to multiple entries of a single event in distributions
 - differential in transverse momentum p_T and rapidity y

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- Di-jet inclusive cross section: $pp \rightarrow 2 jet + X$
 - consider only the two leading jets (in p_T) in the event
 - \rightarrow single entry per event in distributions
 - multi-differential measurements possible (M_{jj} , $p_{T,avg}$, y^* , $y_{boost,...}$)

Jet observables: theory status

- Much progress in fixed-order calculations/resummation and parton shower predictions
 - NLO QCD [Ellis, Kunszt, Soper '92][Giele, Glover, Kosower '94] [Nagy 02]
 - NLO QCD + PS [Alioli, Hamilton, Nason, Oleari, Re '11] [Hoeche, Schonherr '12] [Herwig7 '15]
 - NLO QCD + Resummation (threshold+jet radius)

[Dasgupta, Dreyer, Salam, Soyez '14] [Liu, Moch, Ringer '17]

- NLO EW [Dittmaier, Huss, Speckner '13] [Campbell, Wackeroth, Zhou '16]
 NLO QCD+EW [Frederix, Frixione, Hirschi, Pagani, Shao, Zaro '17]
- NNLO QCD [Gehrmann-De Ridder, Gehrmann, Glover, JP '13] [Currie, Glover, JP '16] [Currie, Gehrmann-De Ridder, Gehrmann, Glover, Huss, JP '17]

NLO QCD + PS (POWHEG)

Nason [arXiv:hep-ph/0409146] JHEP 0411 (2004) 040 Alioli, Hamilton, Nason, Oleari, Re [arXiv: 1012.3380] JHEP 1104 (2011) 081

hardest emission generated first according to:

$$d\sigma = \bar{B}\left(\Phi_{B}\right) d\Phi_{B} \left[\Delta_{R}\left(p_{T}^{\min}\right) + \frac{R\left(\Phi_{R}\right)}{B\left(\Phi_{B}\right)}\Delta_{R}\left(k_{T}\left(\Phi_{R}\right)\right) d\Phi_{\mathrm{rad}}\right]$$

 with the born cross section replaced with the NLO differential cross section at fixed Born kinematics integrated over single emission → preserves NLO accuracy for inclusive quantities

$$\bar{B}(\Phi_B) = B(\Phi_B) + \left[V(\Phi_B) + \int d\Phi_{\rm rad} R(\Phi_R)\right]$$

• probability of not having an emission harder than $p_T \rightarrow POWHEG$ Sudakov

$$\Delta_{R}(p_{T}) = \exp\left[-\int d\Phi_{\text{rad}} \frac{R(\Phi_{R})}{B(\Phi_{B})} \theta\left(k_{T}(\Phi_{R}) - p_{T}\right)\right]$$

 interface with shower generator to develop the rest of the shower vetoing emissions harder than the first one

NLO QCD + PS (POWHEG)

ATLAS [arXiv: 1706.03192] JHEP 1709 (2017) 020

- Comparison to cross section data R = 0.4; 0.6 @ 8 TeV
- POWHEG prediction lower than fixed-order NLO QCDxNPxEW
- Ratio to data less sensitive to the jet radius in POWHEG



NLO QCD + Resummation (threshold+jet radius)

Liu, Moch, Ringer **[arXiv: 1708.04641] Phys. Rev. Lett. 119, 212001 (2017)** Liu, Moch, Ringer **[arXiv: 1801.07284] Phys.Rev. D97 (2018) no.5, 056026** Dasgupta, Dreyer, Salam, Soyez **[arXiv: 1602.01110] JHEP 1606 (2016) 057**

Double differential single jet inclusive cross section:

$$\frac{p_T^2 d^2 \sigma}{dp_T^2 dy} = \sum_{i_1 i_2} \int_0^{V(1-W)} dz \int_{\frac{VW}{1-z}}^{1-\frac{1-V}{1-z}} dv x_1^2 f_{i_1}(x_1) x_2^2 f_{i_2}(x_2) \frac{d^2 \hat{\sigma}_{i_1 i_2}}{dv dz}(v, z, p_T, R)$$

• $z = s_4/s$ invariant mass recoiling against the jet

- in the small R-limit and threshold limit $z \to 0$ cross section factorizes within SCET

$$\frac{\mathrm{d}^2 \hat{\sigma}_{i_1 i_2}}{\mathrm{d} v \, \mathrm{d} z} = s \int \mathrm{d} s_X \, \mathrm{d} s_c \, \mathrm{d} s_G \, \delta(zs - s_X - s_G - s_c) \mathrm{Tr} \left[\mathbf{H}_{i_1 i_2}(v, p_T, \mu_h, \mu) \, \mathbf{S}_G(s_G, \mu_{sG}, \mu) \right] \\ \times J_X(s_X, \mu_X, \mu) \sum_m \mathrm{Tr} \left[J_m(p_T R, \mu_J, \mu) \otimes_\Omega S_{c,m}(s_c R, \mu_{sc}, \mu) \right],$$

- perturbative contributions know at least to NLO
- joint resummation at NLL accuracy $\alpha_s^n \left(\ln^k(z)/z \right)_+$; $\alpha_s^n \ln^k(R)$
- additive matching to fixed-order NLO QCD result

$$\sigma_{\rm NLO+NLL} = \sigma_{\rm NLO} - \sigma_{\rm NLO_{sing}} + \sigma_{\rm NLL}$$

• resummation scales μ_i evolved through RGE equations to common hard scale $\mu=\mu_h=p_T^{max}$

NLO QCD + Resummation (threshold+jet radius)

Liu, Moch, Ringer **[arXiv: 1801.07284] Phys.Rev. D97 (2018) no.5, 056026** Liu, Moch, Ringer **[arXiv: 1808.04574]**



• threshold logs lead to an enhancement of the cross section for large p_T $~\mu=\mu_h=p_T^{max}$

resummation of small R-logs lead to a decrease in the cross section in the entire range p_T

NNLO QCD

• Perturbative QCD expansion of the inclusive jet cross section at hadron colliders

$$\mathrm{d}\sigma = \sum_{i,j} \int \left[\mathrm{d}\hat{\sigma}^{LO}_{ij} + \left(rac{lpha_s}{2\pi}
ight) \mathrm{d}\hat{\sigma}^{NLO}_{ij} + \left(rac{lpha_s}{2\pi}
ight)^2 \mathrm{d}\hat{\sigma}^{NNLO}_{ij} + \mathcal{O}(lpha_s^3)
ight] f_i(x_1) f_j(x_2) dx_1 dx_2$$

• NNLO gluonic contributions



- tree-level 2→4 matrix elements
- one-loop 2→3 matrix elements
- two-loop 2→2 matrix elements
- NNLO DGLAP evolution
- NNLO PDF's

[Berends, Giele '87] [Mangano, Parke, Xu '87] [Britto, Cachazo, Feng '06]

[Bern, Dixon, Kosower '93] [Kunszt, Signer, Trocsanyi '94]

[Anastasiou, Glover, Oleari, Tejeda-Yeomans '01] [Bern, De Freitas, Dixon '02]

[Moch, Vermaseren, Vogt '04]

[ABMP, CT, NNPDF, MMHT]

NNLO antenna subtraction

$$\begin{aligned} \mathrm{d}\hat{\sigma}_{NNLO} &= \int_{\mathrm{d}\Phi_4} \left(\mathrm{d}\hat{\sigma}_{NNLO}^{RR} - \mathrm{d}\hat{\sigma}_{NNLO}^{S} \right) \\ &+ \int_{\mathrm{d}\Phi_3} \left(\mathrm{d}\hat{\sigma}_{NNLO}^{RV} - \mathrm{d}\hat{\sigma}_{NNLO}^{T} \right) \\ &+ \int_{\mathrm{d}\Phi_2} \left(\mathrm{d}\hat{\sigma}_{NNLO}^{VV} - \mathrm{d}\hat{\sigma}_{NNLO}^{U} \right) \end{aligned}$$

$$d\hat{\sigma}^{S}_{NNLO} \quad d\hat{\sigma}^{T}_{NNLO}$$

• mimic RR,RV in unresolved limits

$${
m d}\hat{\sigma}_{NNLO}^T$$
 ${
m d}\hat{\sigma}_{NNLO}^U$

 analytically cancel the poles in RV and VV matrix elements

For inclusive jet and dijet production: $pp \rightarrow jet + X$; $pp \rightarrow 2jet + X$

- NNLO corrections known [Currie, Gehrmann-De Ridder, Gehrmann, Glover, Huss, JP '17]
- all channels $\{gg, qg, q\bar{q}, qq, q\bar{q}', q\bar{q}'\}$ at leading colour $\alpha_s^2 N^2, \alpha_s^2 N N_F, \alpha_s^2 N_F^2$
- gg-channel with full colour
- sub-leading colour contributions: (suppressed by 1/N²)
 - below two percent at NLO (all channels)

NNLOJET

NNLO fully differential parton-level generator

• Based on antenna subtraction for the analytic cancellation of IR singularities at NNLO

Infrastructure

- Process management
- Phase space, histogram routines
- Validation and testing
- ApplFast interface in progress

Processes implemented at NNLO

- Z+(0,1) jet, W+(0,1) jet
- H+(0,1) jet
- DIS-2jet
- VBF H+2jet
- Inclusive jet production

*X.Chen,J.Cruz-Martinez, J.Currie, R.Gauld, T.Gehrmann, A.Gehrmann-De Ridder,E.W.N.Glover, M.Höfer, A.Huss, T.Morgan, I.Majer, J.Niehues, D.Walker, JP **[arXiv: 1801.06415]** and references therein

Jets at the LHC: Fixed-order and Parton Shower descriptions



A thorough study of jet cross sections at the LHC [Bellm, Buckley, et.al., arXiv:1903.12563]

• PS MC's matched or merged to NLO accuracy and fixed-order LO,NLO,NNLO for:

 \rightarrow Higgs+jet; Z+jet; inclusive jet production at the LHC

 Observe good agreement between different PS predictions for inclusive jet production reflecting the underlying fixed-order NLO results

Jets at the LHC: Fixed-order and Parton Shower descriptions



A thorough study of jet cross sections at the LHC [Bellm, Buckley, et.al., arXiv:1903.12563]

- Complete analysis of **Hadronization** and **MPI** corrections as a function of the jet size for measurements in:
 - \rightarrow Higgs+jet; Z+jet; inclusive jet production at the LHC
- Compatible results for the combined hadronization and MPI corrections between Sherpa and HERWIG, small for R
 values around R=0.4

Jets at the LHC: Fixed-order and Parton Shower descriptions



Inclusive jet cross section *R*-dependence fits to the functional form $f(R) = a + b \log(R) + cR^2$

- larger slope at NNLO and NLO+PS compared to NLO
- Iower panels show scale uncertainty *R*-dependence and alternative uncertainty estimates
 - ansatz 1 uncertainty estimate combines scale uncertainties in $\sigma(R) = \sigma(R_0) \frac{\sigma(R)}{\sigma(R_0)}$ in quadrature
 - ansatz 2 uncertainty estimate combines scale uncertainties of *a*,*b*,*c* in quadrature

Single jet inclusive production: scale choices μ_{R} , μ_{F}

- $p_T \rightarrow$ transverse momentum of the individual jets
- $p_{T1} \rightarrow$ transverse momentum of the leading jet
- $H_T \rightarrow$ scalar sum of the transverse momenta of the reconstructed jets
- \check{H}_T -> scalar sum of the transverse momenta of all partons
- µ_R, µ_F are arbitrary and unphysical parameters and are absent from the true result → a priori each scale above is an equally valid scale choice

However, a suitable scale choice would

- minimize ratios of Q^2/μ^2 ,i.e, faster perturbative convergence and smaller scale uncertainties
- avoid scales that introduce pathological behaviours in the prediction, i.e, σ < 0
- avoid scales that are discontinuous on the phase space of the observable, i.e, no kinks in k-factors

→ recently derived NNLO predictions for inclusive jet production allow for the first time a robust study on scale setting, making use of the knowledge of three orders in the perturbative expansion of the observable

 $\mu \sim p_T$ $\mu \sim p_{T1}$ $\mu \sim H_T$ $\mu \sim \hat{H}_T$

Individual jet contributions and jet fractions

• Single jet inclusive observable receives contributions from all jets in the event, at $O(\alpha_s^4)$

$$\frac{d\sigma}{dp_T}(\mu = p_T) = \frac{d\sigma}{dp_{T1}}(\mu = p_{T1}) + \frac{d\sigma}{dp_{T2}}(\mu = p_{T2}) + \frac{d\sigma}{dp_{T3}}(\mu = p_{T3}) + \frac{d\sigma}{dp_{T4}}(\mu = p_{T4})$$

<u>NLO:</u>

- leading jet dominates
- third jet negligible
- second jet sizeable at high p_T negligible at low p_T

<u>NNLO:</u>

- leading and second jet fractions similar over the whole $\ensuremath{\mathsf{p}_{\mathsf{T}}}$ range
- significant increase in second jet p_T contribution to the inclusive jet sample at NNLO with respect to NLO



Second jet transverse momentum distribution



Corrections to second jet distribution integrated over rapidity R=0.4

- NLO: large and negative with huge uncertainty → potentially large logs sensitive on IR effects; NNLO: large and positive
- Stabilization of the predictions at NNLO (in line with the LO) → functional form of the scale matters

Second jet transverse momentum distribution



Decomposition of events contributing to a single bin in p_{T2}

- bin content above constrained to add up to one by construction
- instability in the second jet p_T distribution from events when additional radiation is not recombined into the outgoing jet
- → Large cancellations between positive real emission and large negative virtual correction provide a guide to understand the behaviour of the scale choice
- worse for $\mu=p_{T1}$ that changes event-by-event in the distribution; $\mu=p_T$ remains constant
- detailed single jet inclusive scale study in [J.Currie, T.Gehrmann, A.Gehrmann-De Ridder, E.W.N.Glover, A.Huss, JP '18]
 [arXiv: 1807.03692] JHEP 1810 (2018) 155

 $\mu = p_{T,1}$ \Rightarrow strongly disfavoured in NNLO predictions for single jet inclusive p_T distributions

ATLAS pp→jet+X, √s=13 TeV

ATLAS [arXiv: 1711.02692] JHEP 05 (2018) 195



inclusive p_T distributions

ATLAS pp→jet+X, √s=13 TeV

ATLAS [arXiv: 1711.02692] JHEP 05 (2018) 195



ATLAS [arXiv: 1711.02692] JHEP 05 (2018) 195

- Conclusion in agreement with recent measurements from ATLAS $\sqrt{s}=13$ TeV
- NNLO theory from NNLOJET
- Theoretical understanding of the scale choice and precision phenomenology with jet observables at NNLO is just starting

Understanding single jet inclusive production at the LHC - first steps of the NNLO era

<u>Results have appeared in the context of individual theory publications looking at published</u> <u>LHC data</u>



Understanding single jet inclusive production at the LHC - dissemination of NNLO results

Today:

- tables of *k*-factors from NNLO runs, fixed binning, fixed PDF set and scale choice
- determination of (PDF, α_s) require NNLO predictions to be computed multiple times (varying PDF sets, scales, etc.)
- would like to store the perturbative coefficients of the (N)NLO QCD calculations of final state observables in lookup tables or grids once and for all with very high statistical precision

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Outlook/steps for the coming months:

- single jet inclusive grids generation using APPLFAST-NNLOJET interface in progress [D.Britzger, C.Gwenlan, K. Rabbertz, M.Sutton] cross sections at NNLO available in both formats
 - FASTNLO
 - APPLGRID
 - → First application: H1 determination of α_s from jet production in DIS (H1: arXiv:1709.07251)
 - grids hosted at CERN → ploughshare project: ploughshare.web.cern.ch
 - many grids already available \rightarrow source for NNLO grids from APPLfast

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Ultimate goal:

 have a consistent description of jet data at NNLO for all jet datasets (ATLAS,CMS,LHCb,ALICE) at low and high-p_T in the central and forward rapidity regions for multiple jet R cone sizes

Dijet triple differential measurement - CMS

- Obtain the maximal sensitivity from the dijet cross section to the parton densities from multidifferential distributions
- Explore the shape of the dijet cross section in triple differential form to constrain PDFs at small to moderate x-values: interplay between the parton luminosities and the scattering matrix elements



- 2 jet observables built from the two-leading jets in the event, with the caveat that the assignment of the hardest jet in the event is not infra-red safe
- Follow the setup of the 8 TeV CMS measurement that overcomes the problem

$$p_{T,\text{avg}} = (p_{T,1} + p_{T,2})/2$$

$$y^* = |y_1 - y_2|/2$$
$$y_b = |y_1 + y_2|/2$$



- density plot in the partonic fractions (x₁,x₂) plane of the triple differential cross section as a function of y* and y_b for the fiducial cuts of the CMS measurement p_{T,avg} >133 GeV
- cuts on y* and y_b directly map to surfaces on the x-Q⁻ plane where the PDFs are determined; CMS setup

State of the art

Previously studied:

- The NLO inclusive two jet triply differential cross section *W.Giele, E.W.N.Glover, D.A.Kosower* [hep-ph/9403347] Phys.Rev.Lett. 73 (1994) 2019
 - Scale uncertainties and missing higher order corrections limit the achievable precision
- Measurement from CDF at 1.8 TeV based on an integrated luminosity of 86pb⁻¹
 CDF collaboration [hep-ex/0012013 Phys.Rev.D64 (2001) 012001

<u>Goal:</u>

- Compute for the first time the triple differential observable at NNLO
- Detailed comparison with CMS 8 TeV data 19.7fb⁻¹ dataset at the percent level
- Include an assessment of NP and EW effects

CMS √s=8 TeV anti-k_T R=0.7



• NP correction 10% in the lowest $p_{T,avg}$ bins, negligible above 1 TeV

- center of the envelop of HERWIG, PYTHIA and POWHEG with and without hadronization and MPI
- EW correction relevant in the low y* and low y_b rapidity slice reaching 8% at high p_{T,avg}
 - EW effects visible in the LHC measurements that reach percent-level precision



- size of the corrections varies significantly as a function of p_{T,avg} and the applied cuts on y* and y_b
- NNLO correction changes both the shape and normalisation of the NLO result
- QCD scale choice $\mu_R=\mu_F=m_{jj}$

$CMS \sqrt{s}=8 \text{ TeV} 0 < y_b < 1$



$CMS \sqrt{s}=8 \text{ TeV } y_b > 1$



Summary/Outlook

Summary:

- Significant theoretical progress in the description of inclusive jet and dijet production at the LHC
- Theoretical developments driven by the increase in precision of the experimental measurements
- New theoretical calculations available that can be used to understand the impact of the effects of higher order corrections in the description of jet data at hadron colliders

Outlook

- Perform further quantitative comparisons between data and theory (different energies, covering wide jet p_T and rapidity and jet cone sizes) → APPLFAST interface in progress
- Use new data to understand effects in tuning of hadronization and underlying event parameters and respective uncertainties
- Extend existing phenomenological predictions to triple-differential measurements and angular observables and jet shapes
- Study sensitivity of jet-based observables to α_s and PDF extractions and assess ultimate reach in precision in a combined fit