



**Universität
Zürich**^{UZH}

Establishing connections between B-physics & High-pT

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SM@LHC Workshop 2019, 25th of April 2019

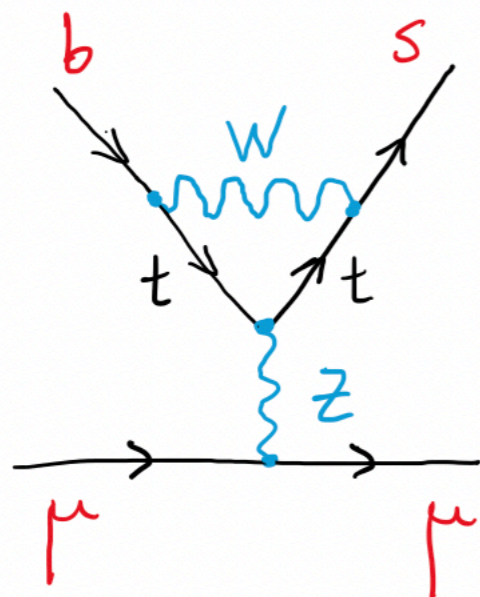
The B-physics anomalies

See talk by Thibaud Humair

Hints of **L**epton **F**lavour **U**niversality **V**iolation in semileptonic B decays

$$b \rightarrow s \ell^+ \ell^-$$

μ/e universality

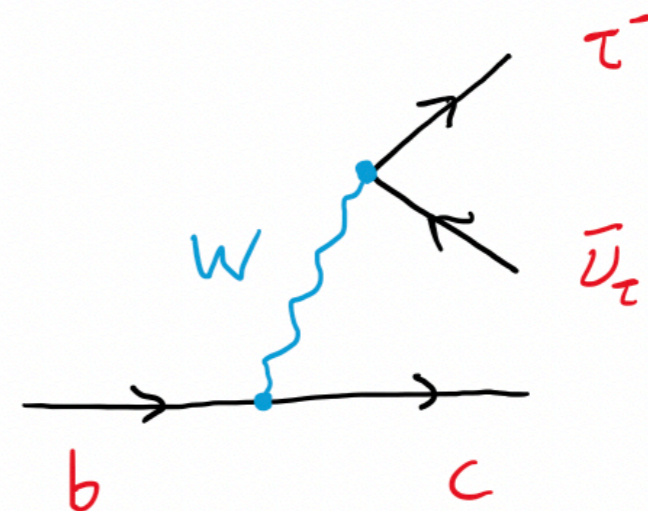


$> 4\sigma$

★ New R_K measurement by LHCb and R_{K^*} by Belle

$$b \rightarrow c \tau \nu$$

$\tau/\mu, e$ universality



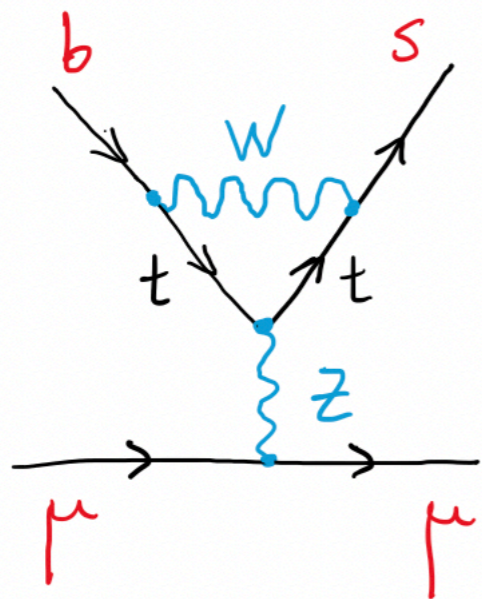
$\sim 3\sigma$

★ New $R_{D^{(*)}}$ measurement by Belle

Towards a combined explanation of the anomalies

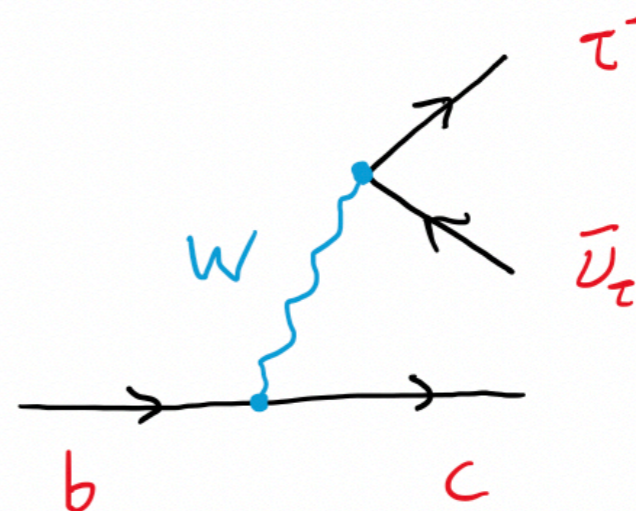
Taken together, these are a very significant set of deviations from the SM

→ It is worth looking for a **combined explanation** in terms of NP!



$$3_Q \rightarrow 2_Q 2_L 2_L$$

~20% of a SM **loop** effect



$$3_Q \rightarrow 2_Q 3_L 3_L$$

~15% of a SM **tree-level** effect

The only source of **lepton flavor universality violation** in the SM (Yukawas) follows a similar trend: $y_e \ll y_\mu \ll y_\tau \dots$. Are the anomalies connected to them?

What are the anomalies telling us?

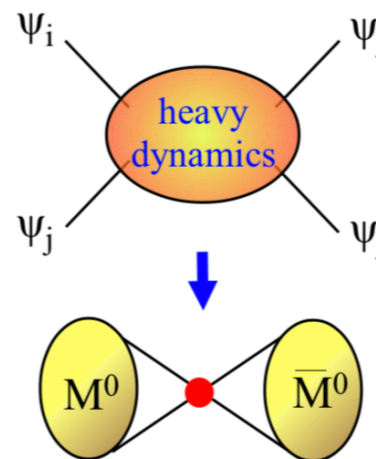
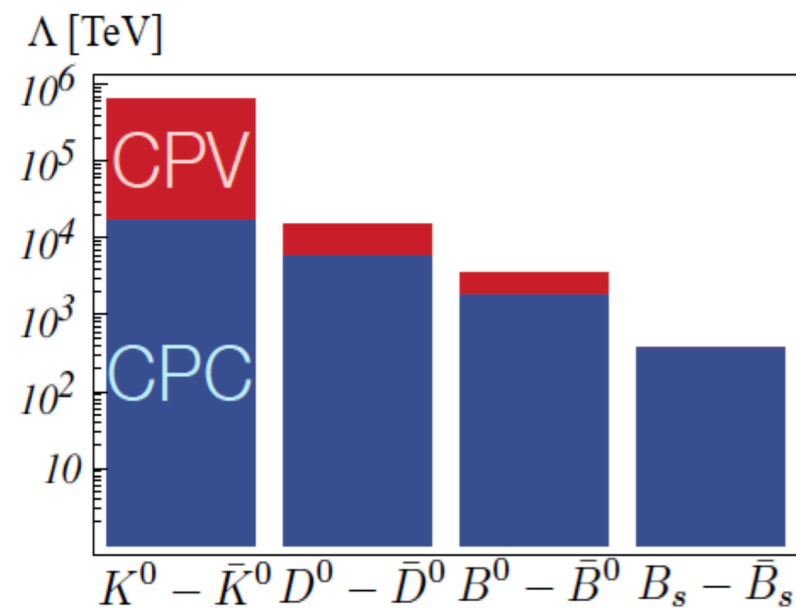
A combined explanation calls for NP: (*)

★ Coupled dominantly to the **3rd generation**

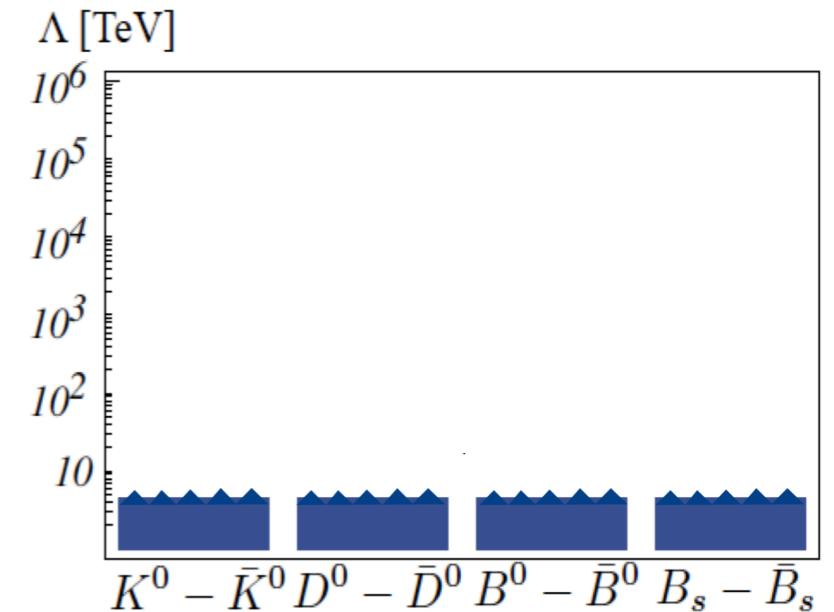
★ $\Lambda_{\text{NP}} \sim \mathcal{O}(1 \text{ TeV})$

(*) N.B.: conclusions driven (mostly) by $R(D^{(*)})$

Anarchical couplings



Hierarchical couplings



Severe constraints on generic new (BSM) flavor breaking sources
(mis)interpreted as indication of a high flavor scale

A NP hint to the SM flavor puzzle?

The SM Yukawa sector is characterized by **13** parameters

[**3** lepton masses + **6** quark masses + **3+1** CKM parameters]

... whose values do **not** look at all accidental

$$M_{u,d,e} \sim \begin{array}{|c|c|c|} \hline \text{light} & \text{light} & \text{light} \\ \hline \text{light} & \text{medium} & \text{medium} \\ \hline \text{light} & \text{medium} & \text{dark} \\ \hline \end{array}$$

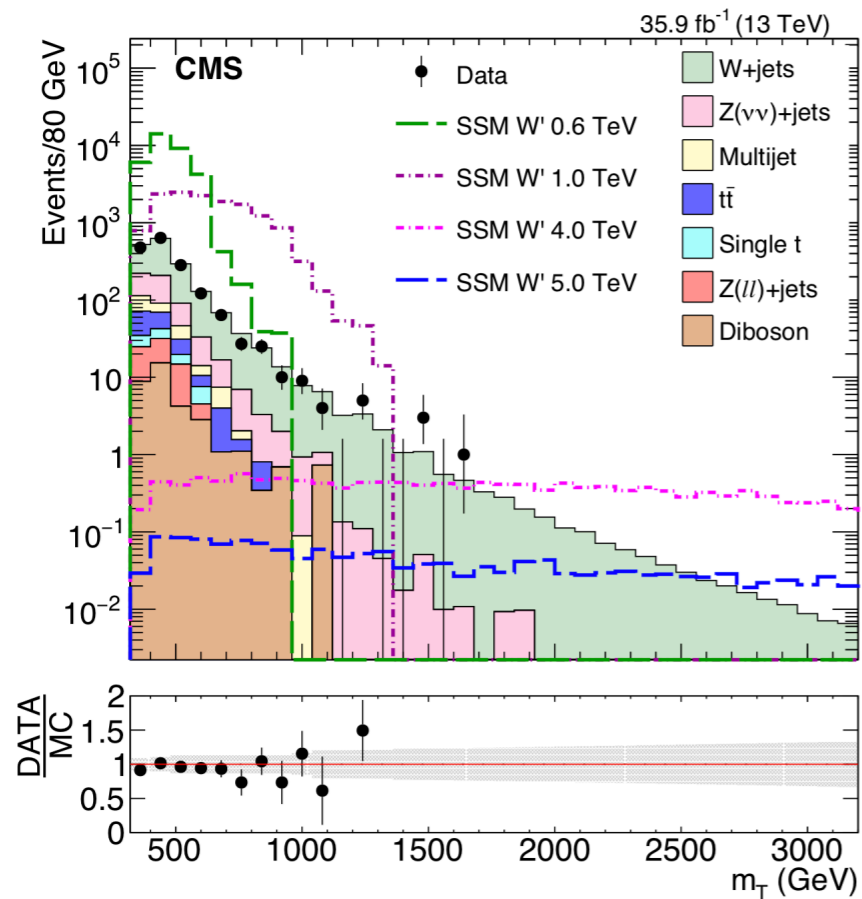
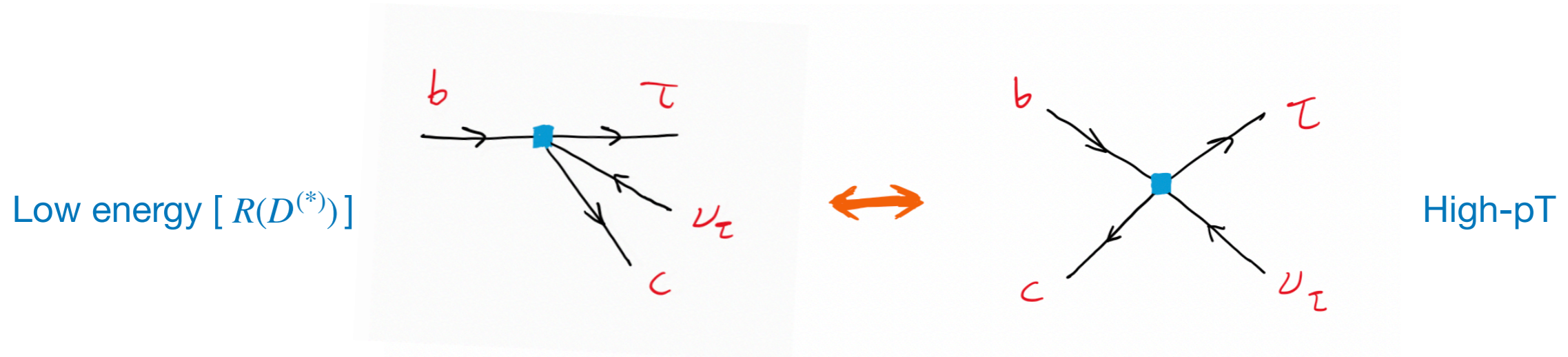
$$V_{\text{CKM}} \sim \begin{array}{|c|c|c|} \hline \text{dark} & \text{medium} & \text{light} \\ \hline \text{medium} & \text{dark} & \text{light} \\ \hline \text{light} & \text{light} & \text{dark} \\ \hline \end{array}$$

- ✓ The flavor anomalies seem to suggest a similar trend: large **NP effects in 3rd generation**, gradually smaller effects in the light generations
- ✓ Recent theoretical progress connecting the anomalies to the SM flavor hierarchies

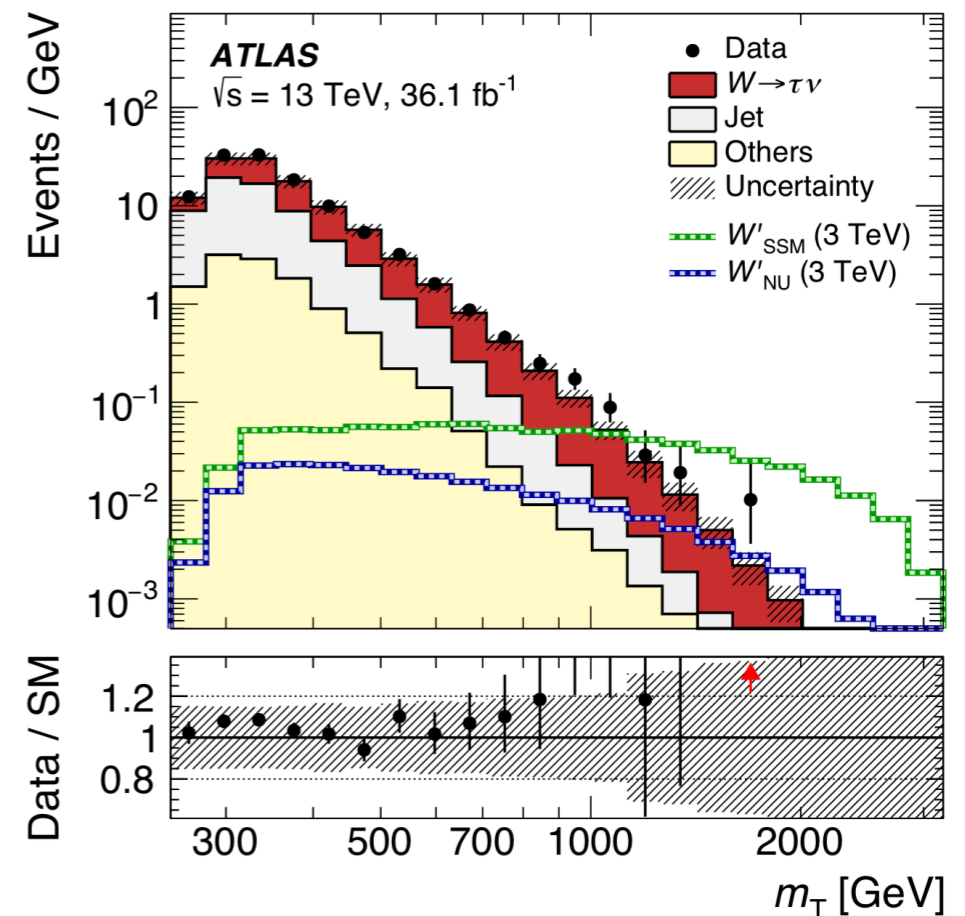
[Bordone, Cornella, JFM, Isidori 1712.01368; Greljo, Stefanek 1802.04274; Allanach, Davighi 1809.01158; Barbieri, Ziegler 1904.04121]

Complementarity with high-pT data

The same NP (specially in $R(D^{(*)})$) is within the reach of high-pT searches at LHC!



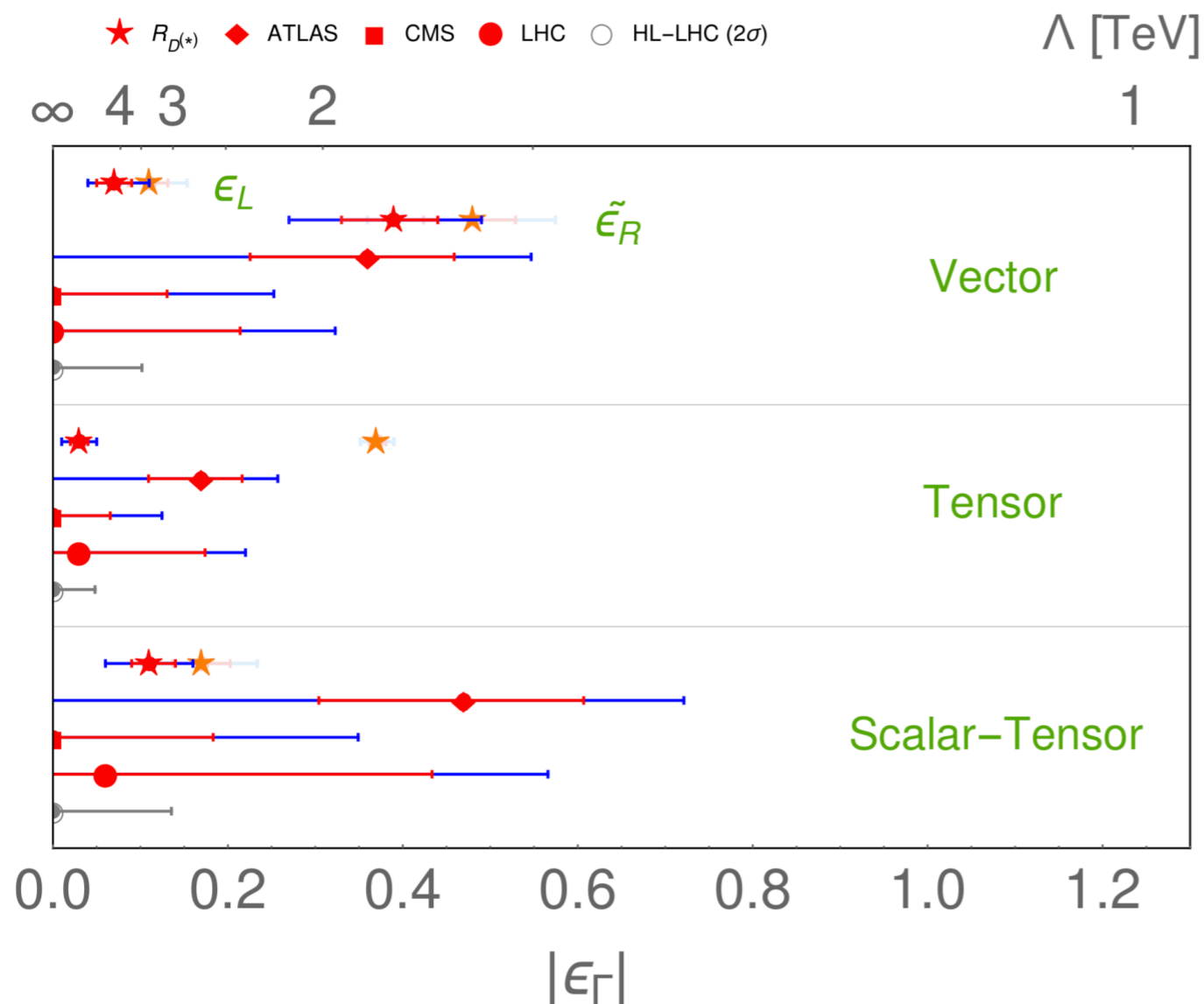
1807.11421



1801.06992

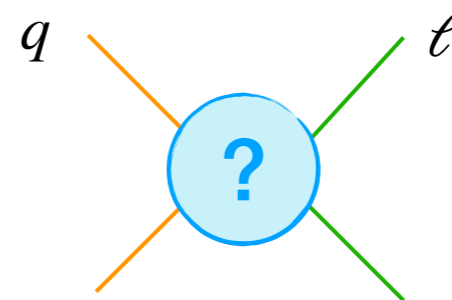
High-pT limits on EFT operators

$$\mathcal{L}_{\text{EFT}} \supset -\frac{2V_{cb}}{v^2} \left[(1 + \epsilon_L) (\bar{\tau} \gamma_\mu P_L \nu_\tau) (\bar{c} \gamma^\mu P_L b) + \epsilon_R (\bar{\tau} \gamma_\mu P_L \nu_\tau) (\bar{c} \gamma^\mu P_R b) + \epsilon_{S_L} (\bar{\tau} P_L \nu_\tau) (\bar{c} P_L b) \right. \\ \left. + \epsilon_{S_R} (\bar{\tau} P_L \nu_\tau) (\bar{c} P_R b) + \epsilon_T (\bar{\tau} \sigma_{\mu\nu} P_L \nu_\tau) (\bar{c} \sigma^{\mu\nu} P_L b) \right] + h.c.$$



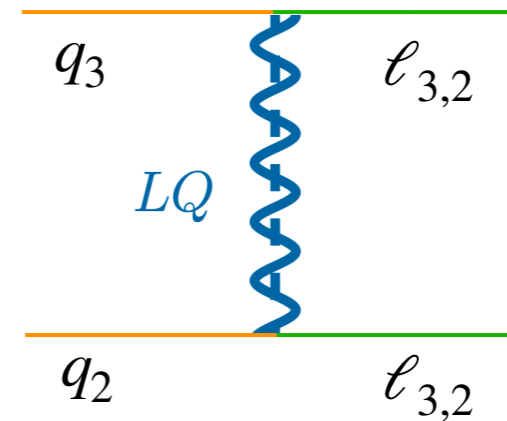
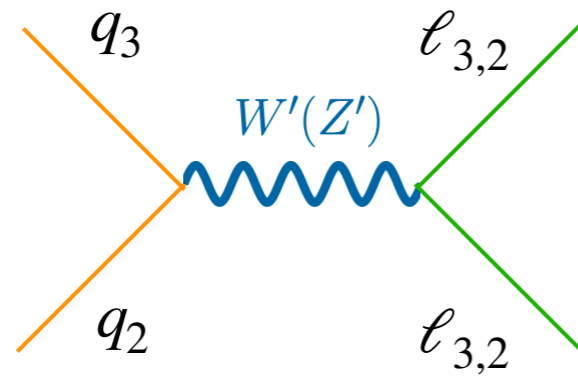
Caveats

- ★ EFT stops being valid when $E \approx \Lambda$
- ★ UV models typically involve additional effects



...Which mediator?

Which mediator?



Only few possibilities for a combined explanation

★ **Minimal W'/Z' models** in tension with high- p_T data ($pp \rightarrow \tau\tau$ tails)

[Faroughy et al. 1609.07138]

W' + light ν_R in better shape but still in tension with $pp \rightarrow \tau\nu$ tails

[Greljo et al. 1811.07920]

Talk by Y. Takahashi

★ **Leptoquarks** (scalars or vectors) are the best candidates so far

✓ no 4-lepton (LFV, LFUV) and 4-quark processes ($\Delta F = 2$) at tree level

The main suspects

Faroughi @ CKM18

	Model	$R_{K(*)}$	$R_{D(*)}$	$R_{K(*)}$ & $R_{D(*)}$
Scalars	$S_1 = (\mathbf{3}, \mathbf{1})_{-1/3}$	✗	✓	✗
	$R_2 = (\mathbf{3}, \mathbf{2})_{7/6}$	✗	✓	✗
	$\tilde{R}_2 = (\mathbf{3}, \mathbf{2})_{1/6}$	✗	✗	✗
	$S_3 = (\mathbf{3}, \mathbf{3})_{-1/3}$	✓	✗	✗
Vector	$U_1 = (\mathbf{3}, \mathbf{1})_{2/3}$	✓	✓	✓
	$U_3 = (\mathbf{3}, \mathbf{3})_{2/3}$	✓	✗	✗

Angelescu, Becirevic, DAF, Sumensari [1808.08179]

Three viable options in the market:

★ $U_1 + UV$ completion

[di Luzio, Greljo, Nardecchia 1708.08450;
Calibbi, Crivellin, Li 1709.00692;
Bordone, Cornella, JF, Isidori 1712.01368;
Barbieri, Tesi, 1712.06844...]

★ $S_1 + S_3$

[Crivellin, Muller, Ota 1703.09226;
Buttazzo et al. 1706.07808;
Marzocca 1803.10972]

★ $S_3 + R_2$

[Bečirević et al., 1806.05689]

The vector leptoquark (U_1) brings some interesting theoretical features into the game

- ✓ Low-scale bottom-tau unification. Possible link to Pati-Salam unification
- ✓ Connections to the SM flavor puzzle

The U_1 leptoquark solution

$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \left[\beta_{i\alpha}^L (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) - \beta_{i\alpha}^R (\bar{d}_R^i \gamma_\mu e_R^\alpha) \right] + \text{h.c.}$$

Flavor structure (*)

$$\beta^L \approx \begin{pmatrix} 0 & 0 & \beta_{d\tau}^L \\ 0 & \beta_{s\mu}^L & \beta_{s\tau}^L \\ 0 & \beta_{b\mu}^L & \beta_{b\tau}^L \end{pmatrix}$$

$$\beta^R \approx \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \beta_{b\tau}^R \end{pmatrix}$$

$$\beta_{b\tau}^L = 1$$

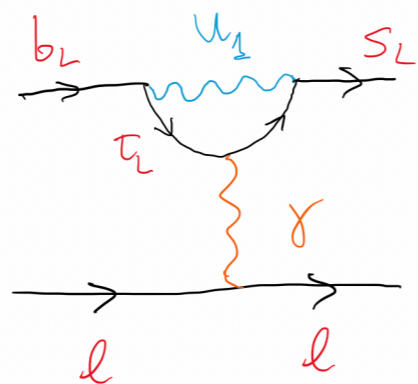
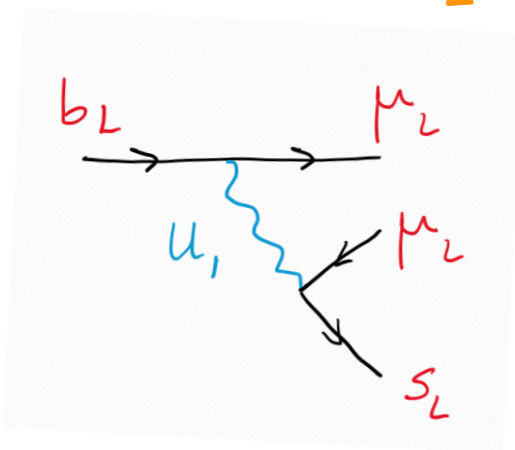
$$\beta_{b\tau}^R \sim \mathcal{O}(1)$$

$$\beta_{s\tau}^L, \beta_{b\mu}^L \sim \mathcal{O}(0.1)$$

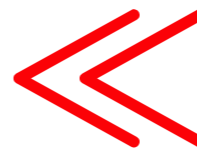
$$\beta_{s\mu}^L, \beta_{d\tau}^L \sim \mathcal{O}(0.01)$$

ϵ_{VL}

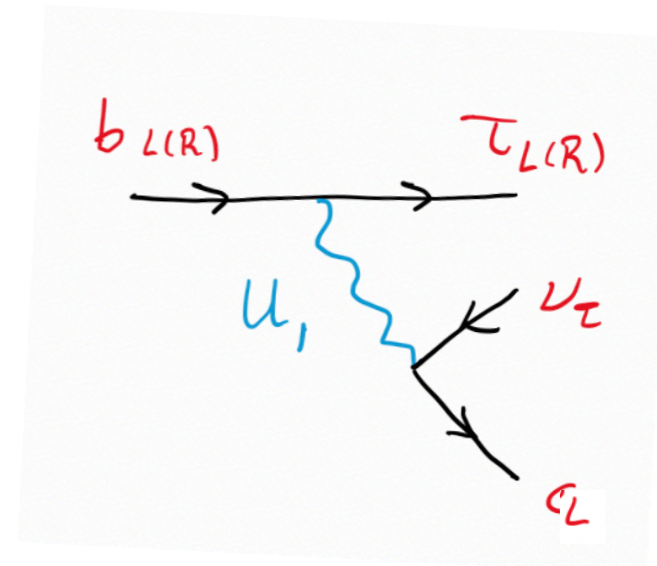
ϵ_{SR}



$b \rightarrow sl\ell$



$R(D^{(*)})$



$$\Delta C_9^\mu = -\Delta C_{10}^\mu$$

$$\Delta C_9^{e,\mu,\tau}$$

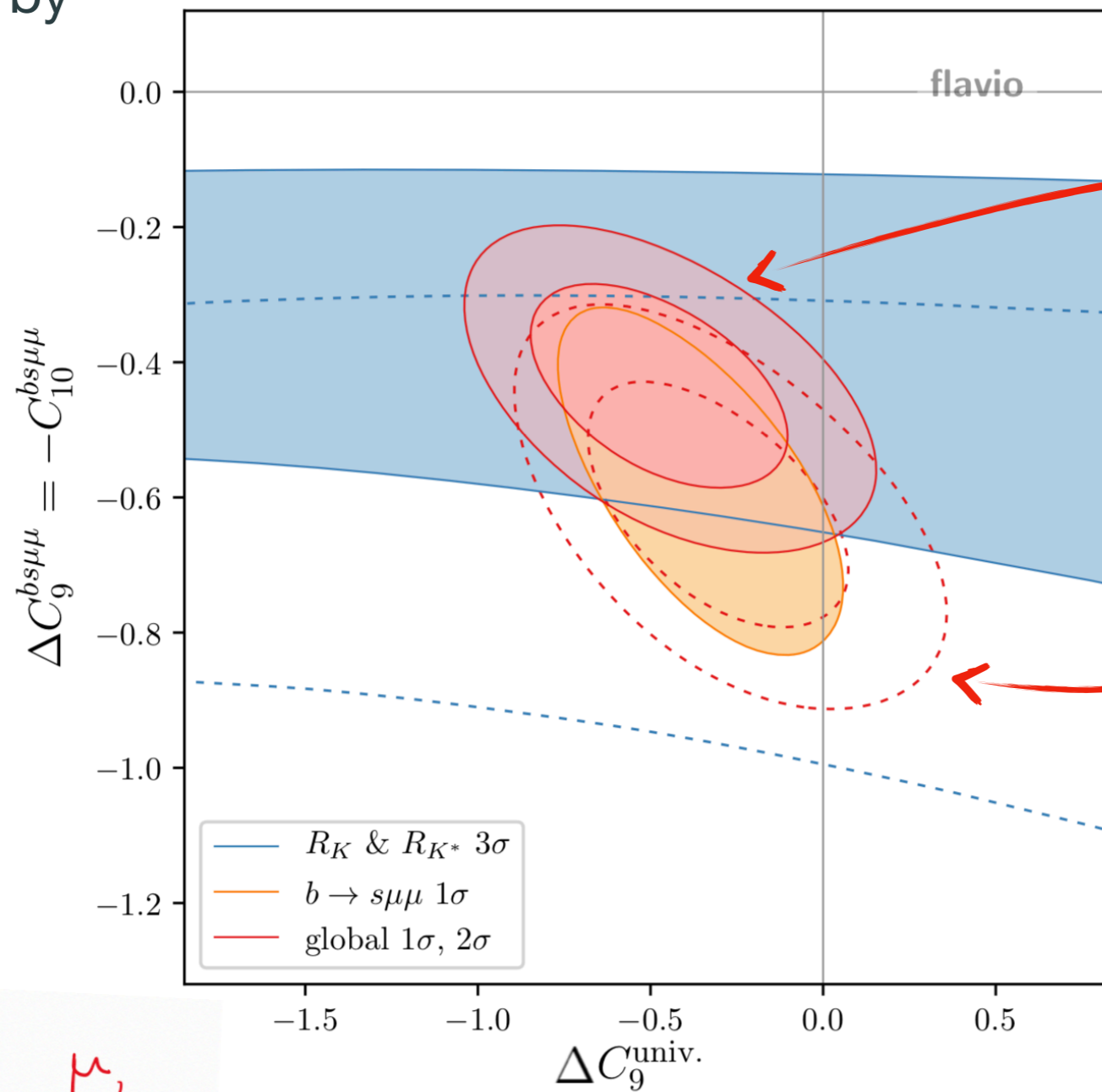
[Crivellin et al., 1807.02068]

(*) N.B.: Deviations from this structure highly constrained by low-energy flavor data

Universal contribution to ΔC_9

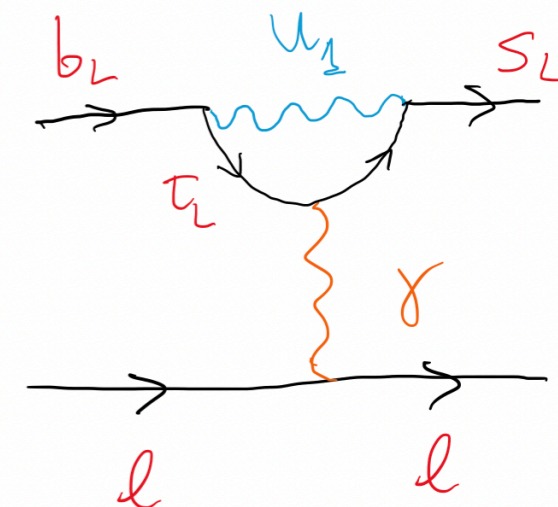
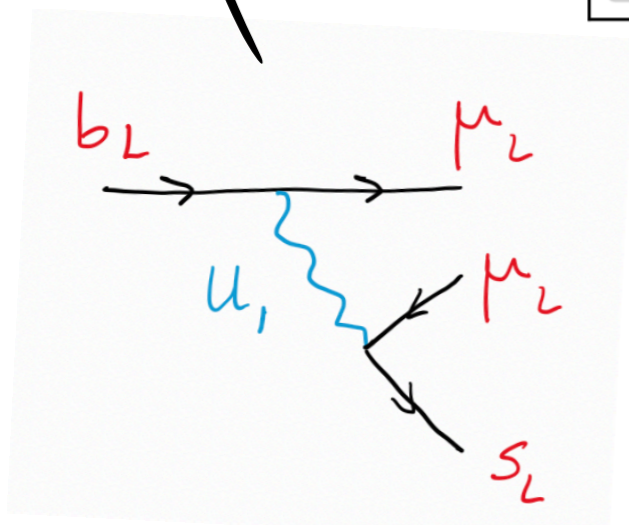
... as discussed by David Straub

[Aebischer et al., 1903.10434]



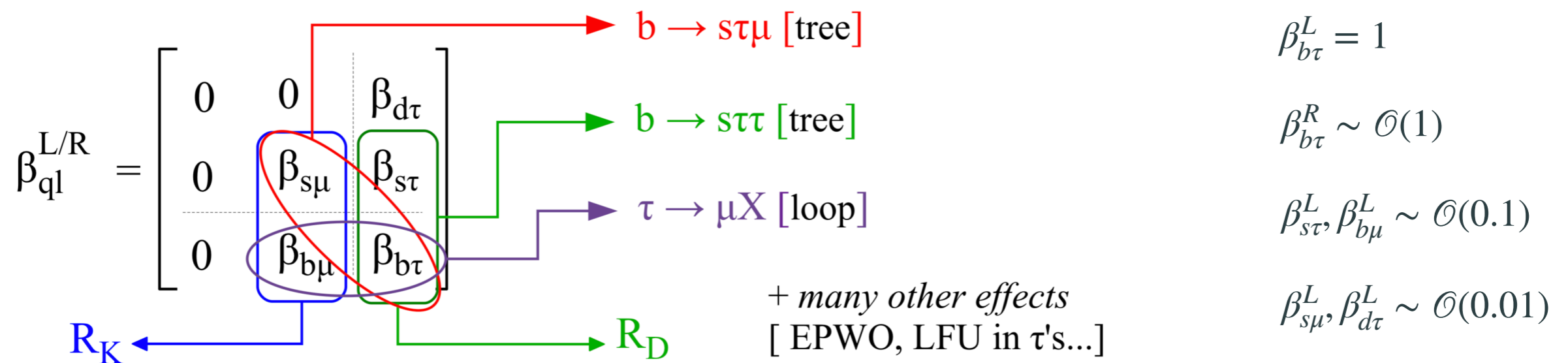
Now

Pre-Moriond



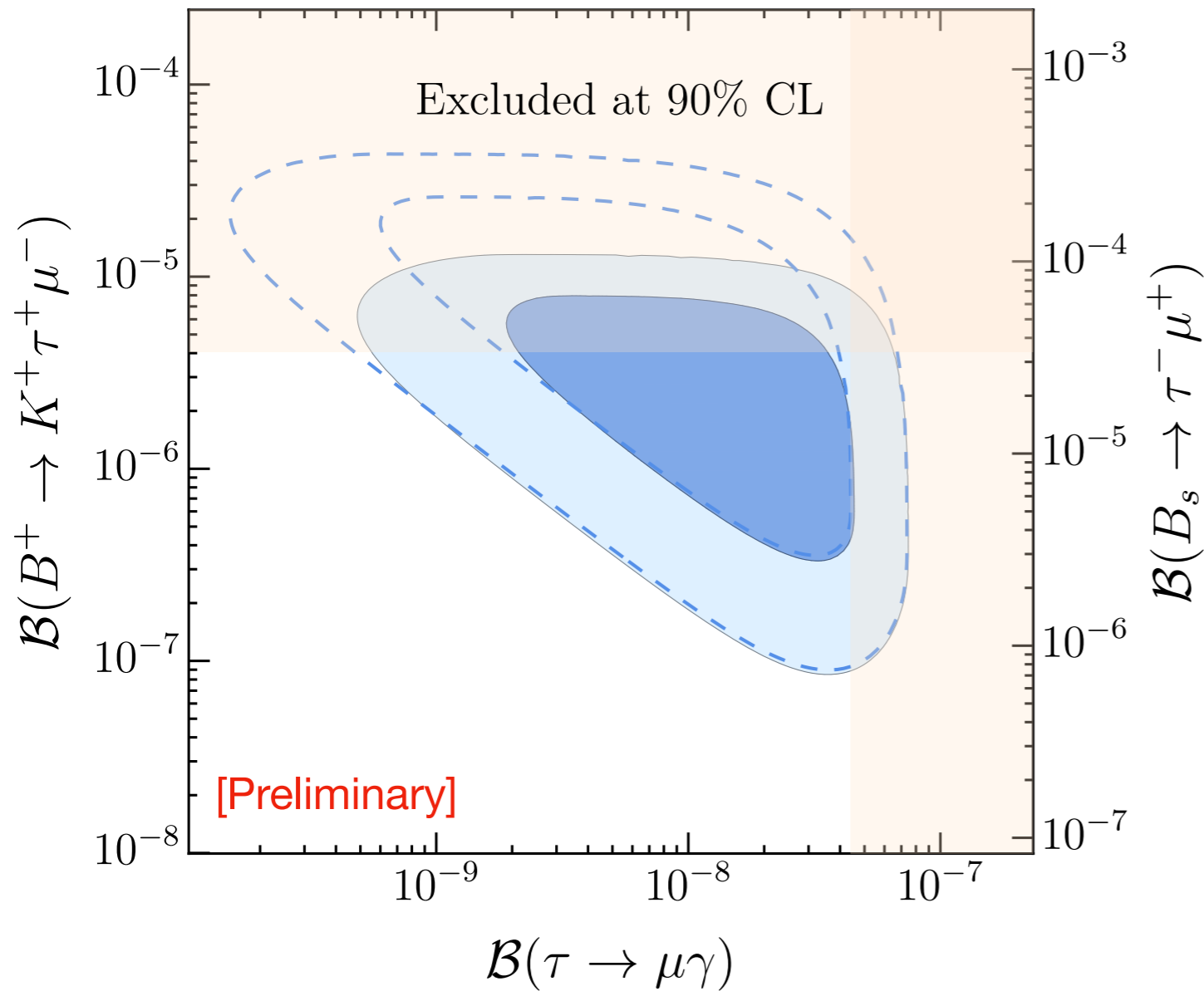
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A non-zero $\beta_{b\tau}^R$ leads to a series of interesting effects both at low-energy and at high-pT

LFV in $\tau \rightarrow \mu$ transitions ($\beta_{b\tau}^R = 1$)



Large $\tau\mu$ LFV is expected

→ strong enhancement of $B_s \rightarrow \tau\mu$, $B \rightarrow K\tau\mu$, $\tau \rightarrow \mu\gamma$

$$\mathcal{B}(\tau \rightarrow \mu\gamma) \sim 10^{-9}$$

$$\mathcal{B}(B \rightarrow K\tau\mu) \sim 10^{-6}$$

$$\mathcal{B}(B_s \rightarrow \tau\mu) \sim 10^{-5}$$

★ First measurement of $\mathcal{B}(B_s \rightarrow \tau\mu)$ by LHCb presented here!

See talk by Giampiero Mancinelli

Low-energy fit results

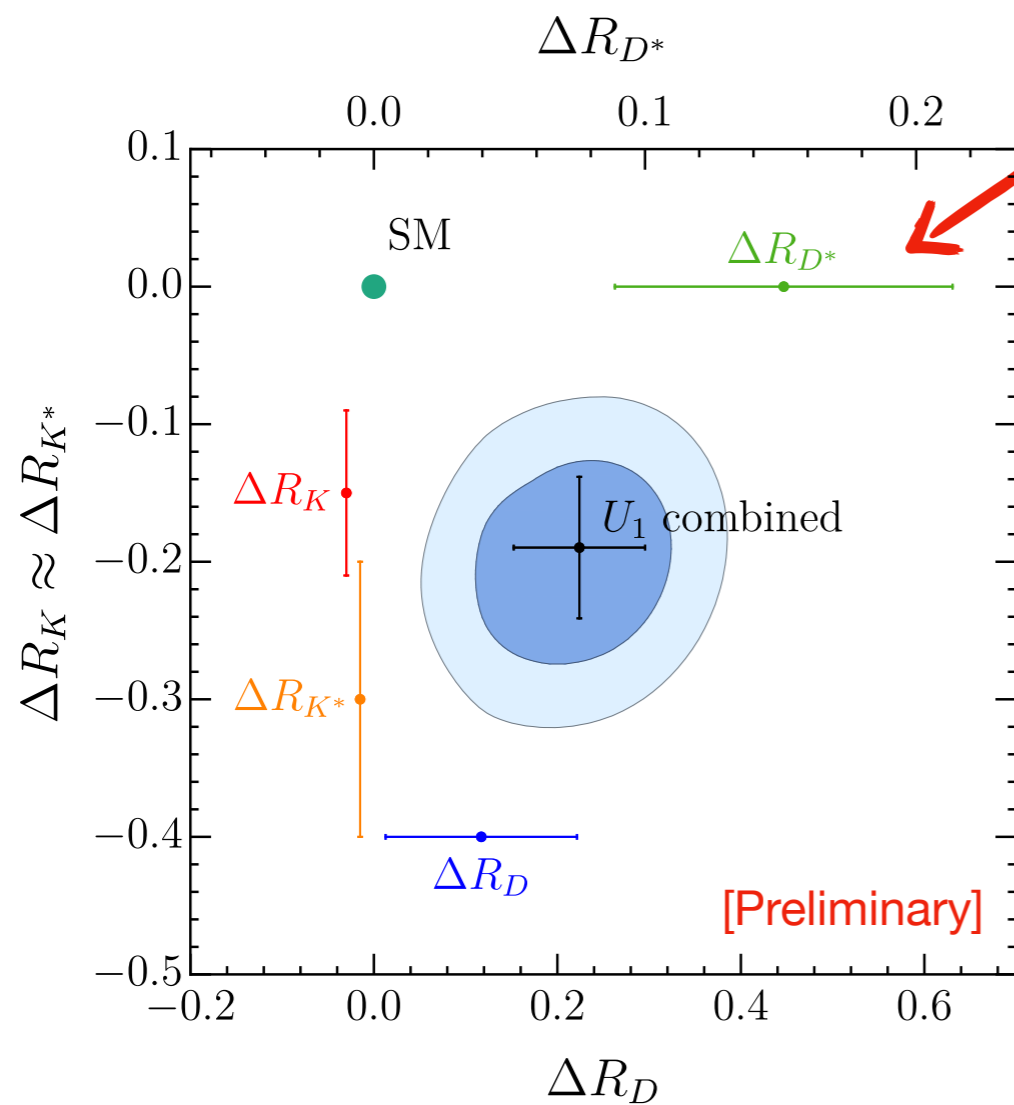
Both options give a good the low-energy fit (in particular to the anomalies)

LH + RH ($\beta_{b\tau}^L = \beta_{b\tau}^R$)

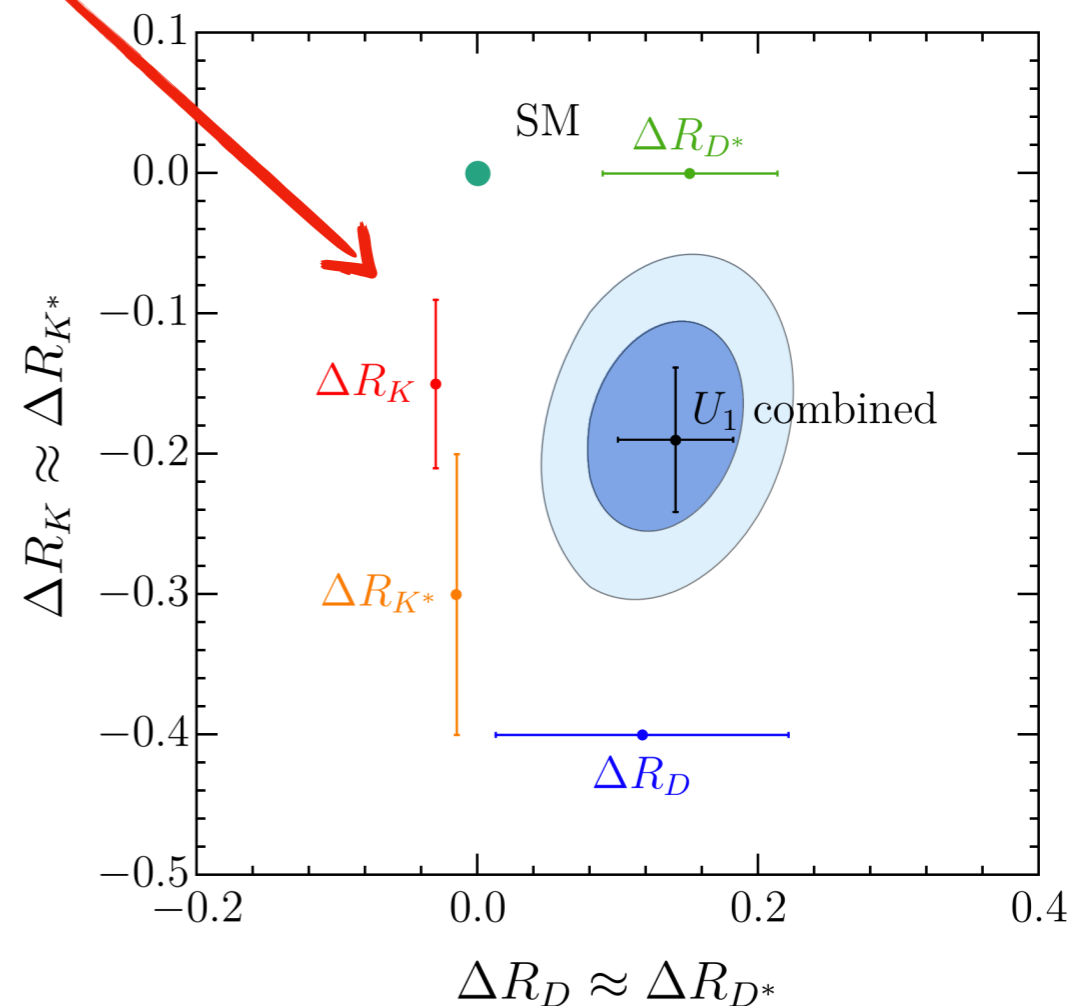
Experiment

LH only ($\beta_{b\tau}^R = 0$)

@ 1 σ



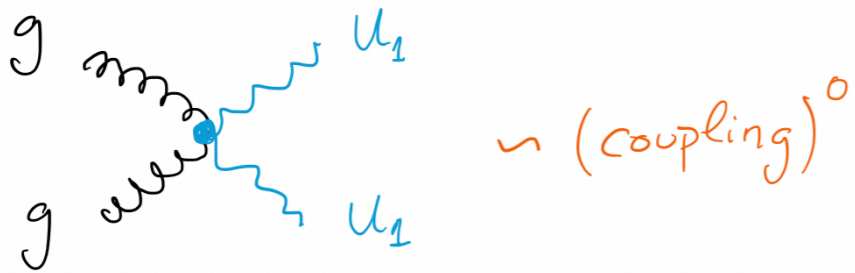
NP scale naturally higher
[thanks to the additional NP contribution from $\beta_{b\tau}^R$]



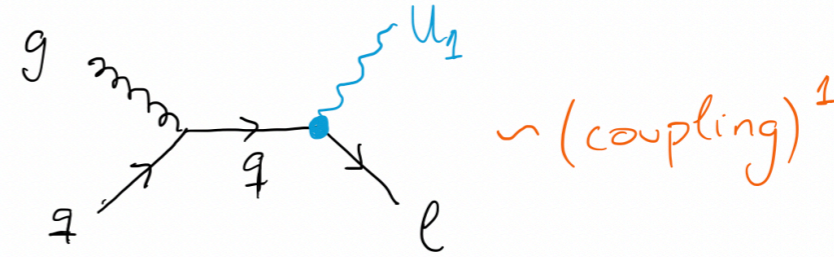
Slightly better fit
[driven by the new (preliminary) Belle measurement]

High-pT leptoquark limits

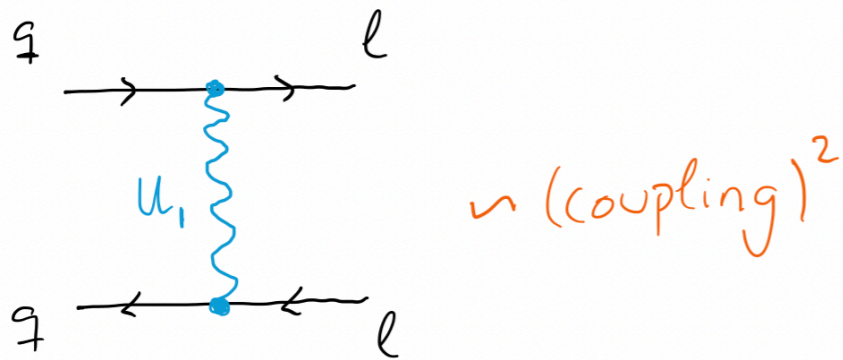
Pair production (PP)



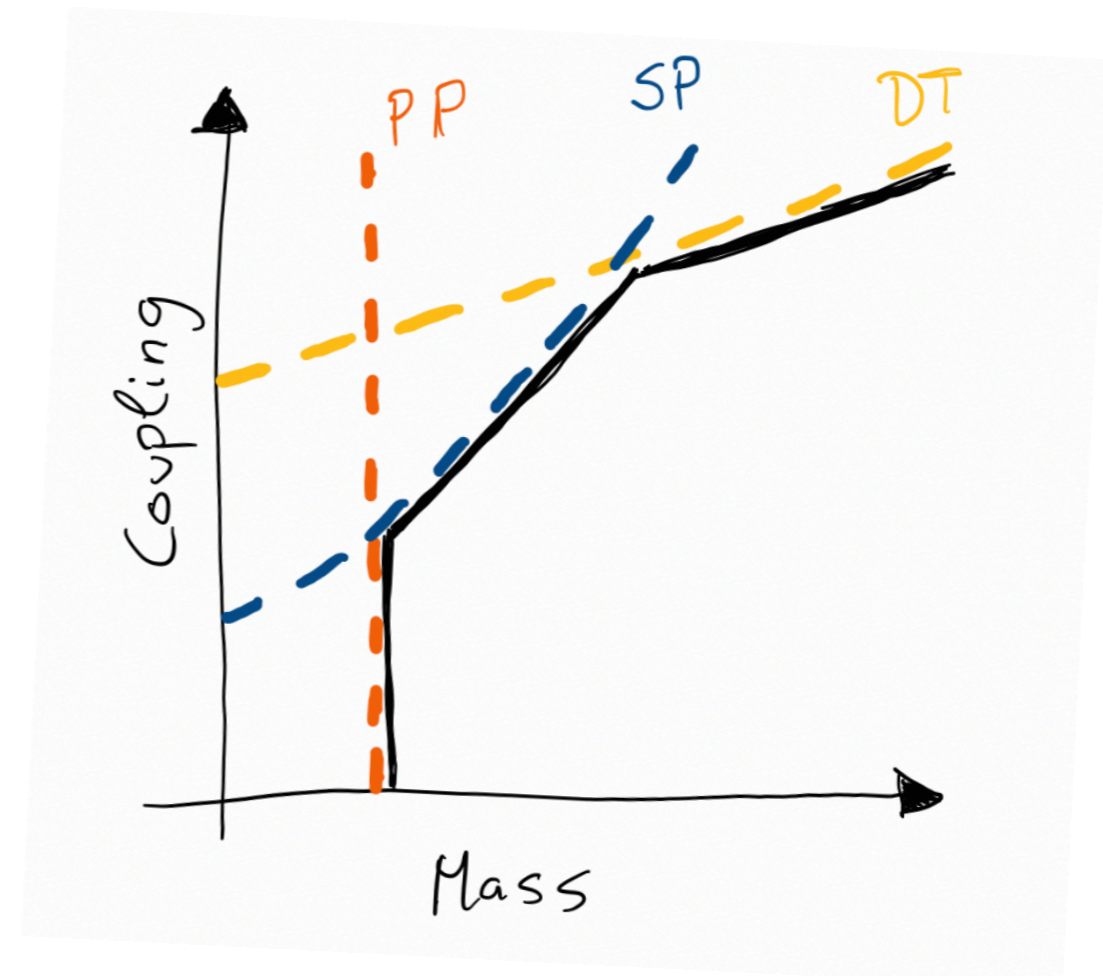
Single production (SP)



Dilepton tails (DT)



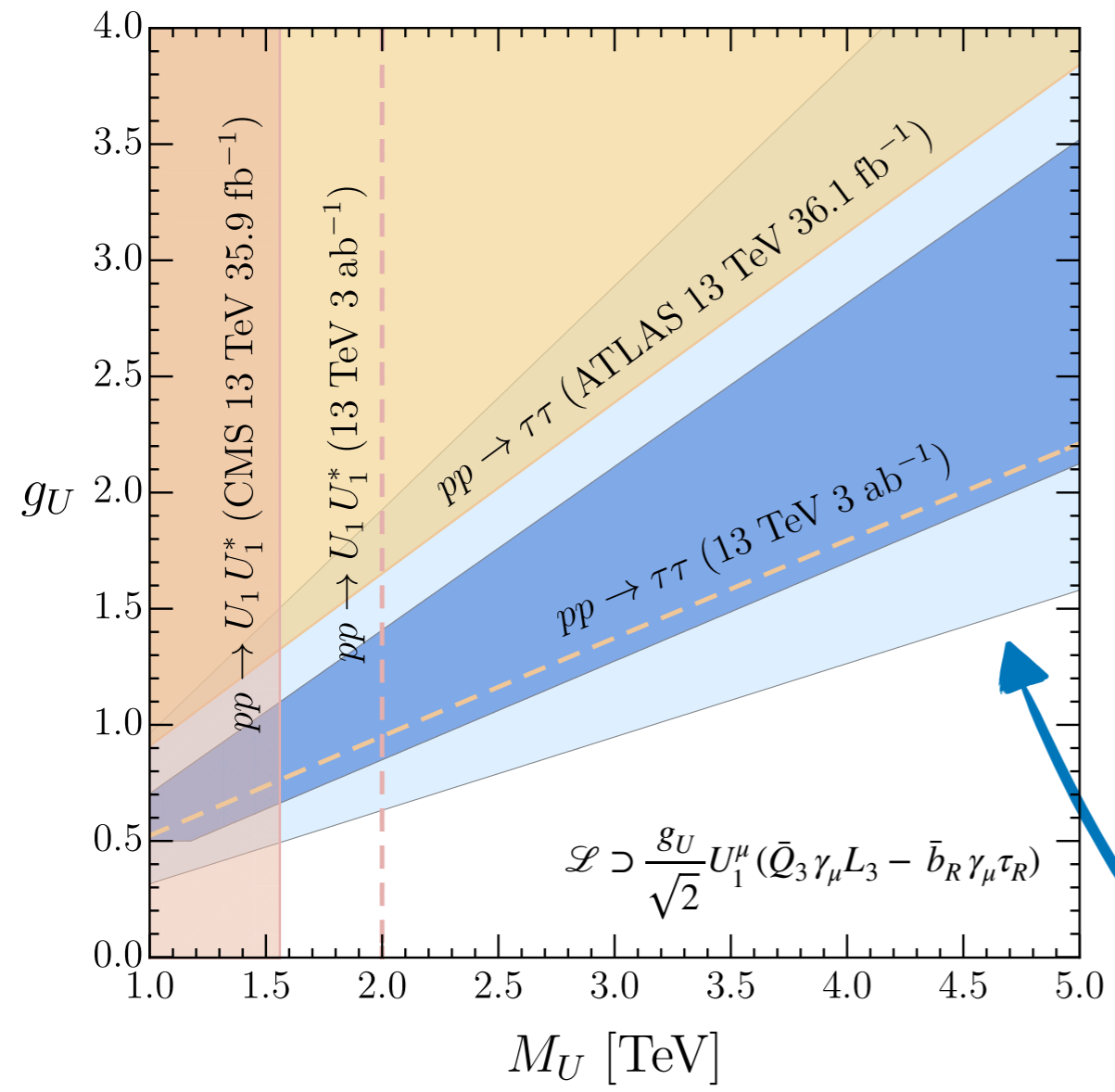
This talk!



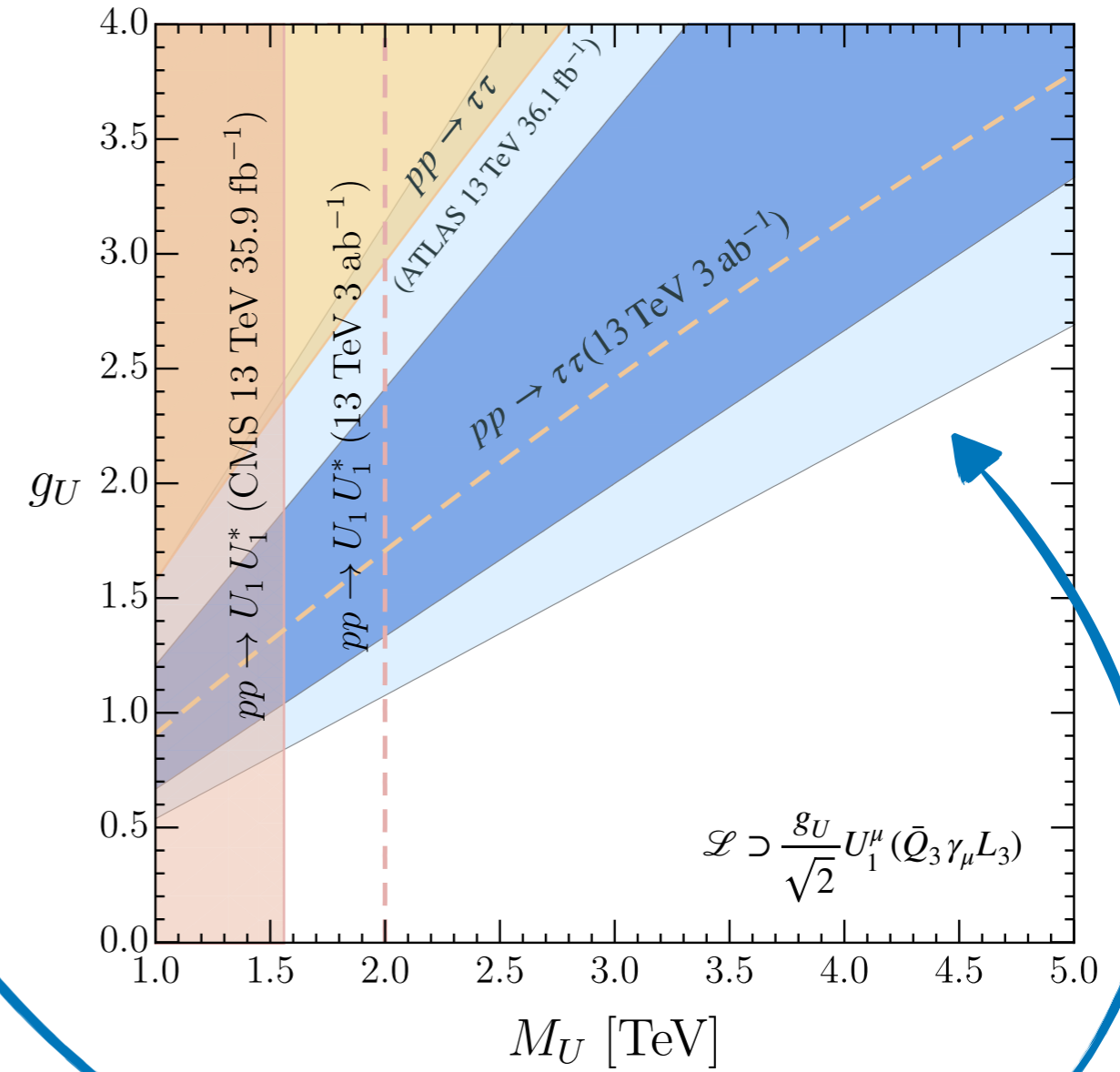
High-pT + low energy

[High-pT bounds from Baker, JFM, Isidori, König, 1901.10480]

LH + RH ($\beta_{b\tau}^L = \beta_{b\tau}^R$)



LH only ($\beta_{b\tau}^R = 0$)



[Cornella, JFM, Isidori, 1903.11517]

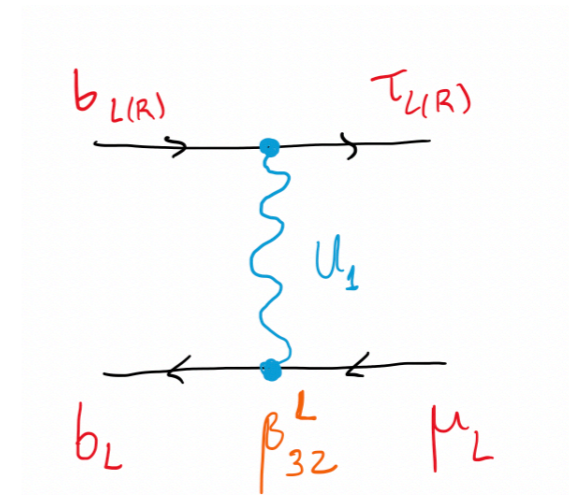
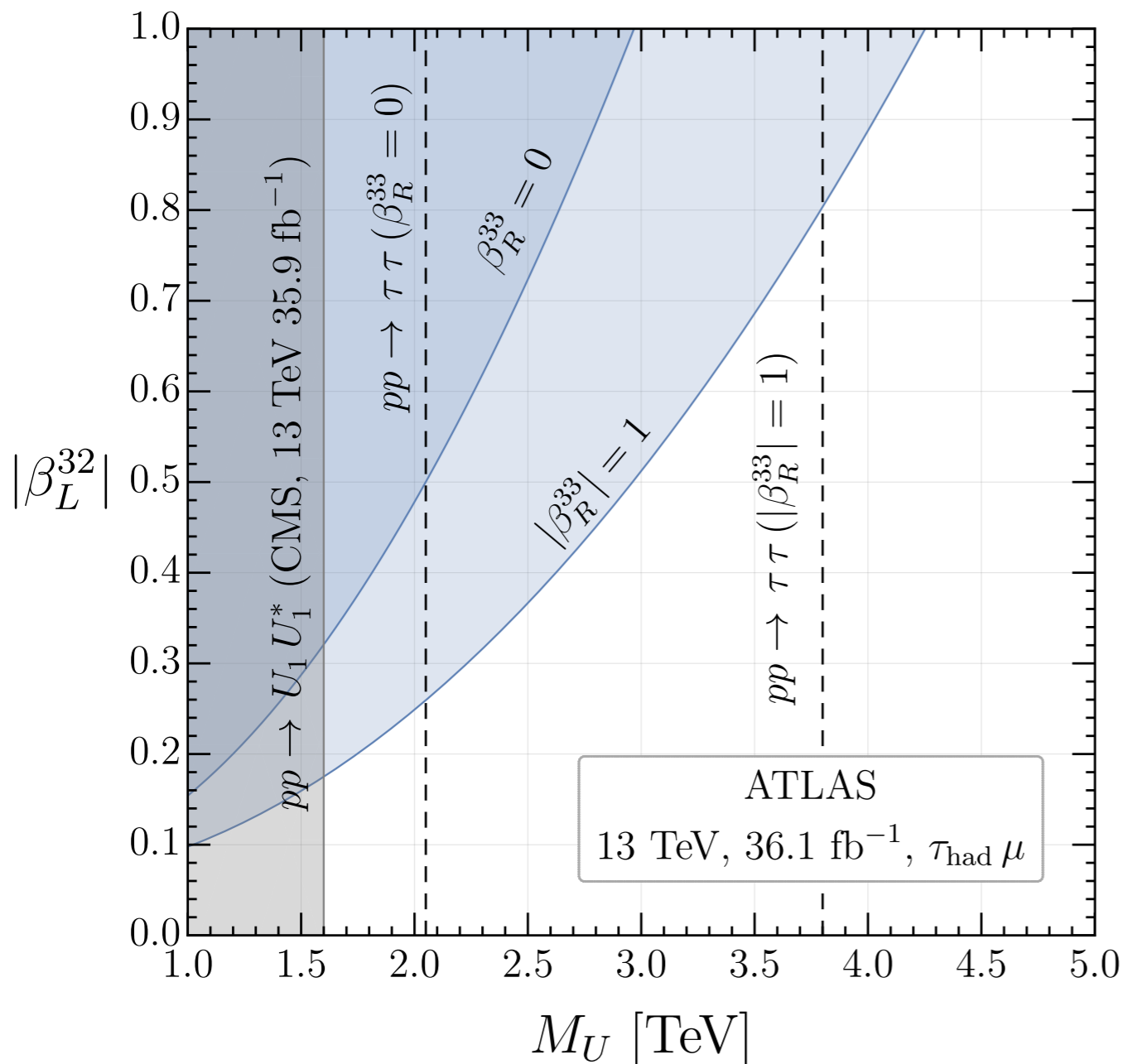
(*) N.B.: Assuming $\beta_{s\tau} < 0.25$

1 and 2 σ regions preferred by the low-energy fit

Recast of the $pp \rightarrow \tau\mu$ data

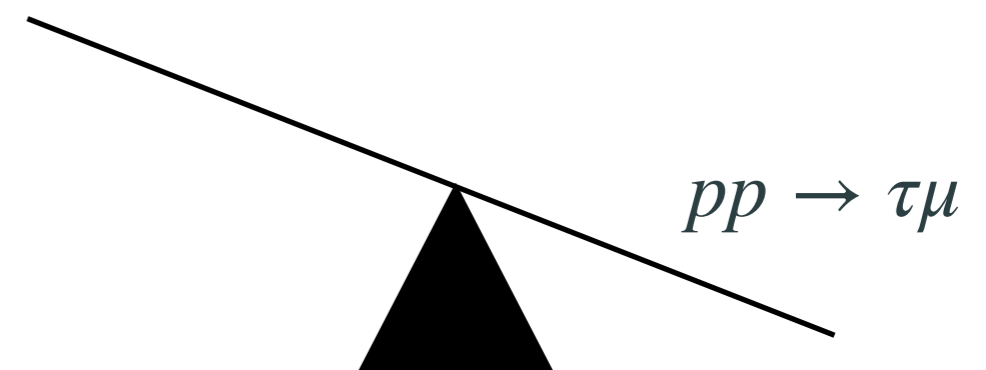
[Baker, JFM, Isidori, König, 1901.10480]

$$\mathcal{L} \supset \frac{3}{\sqrt{2}} U_\mu (\beta_{32}^L \bar{Q}_3 \gamma^\mu L_2 + \bar{Q}_3 \gamma^\mu L_3 - \beta_R \bar{b}_R \gamma^\mu \tau_R)$$



Stronger limits from $pp \rightarrow \tau\tau$
 ($\beta_{32}^L \sim 0.1$ preferred by $R(D^{(*)})$ fit)

$\Upsilon \rightarrow \tau\mu$ [1806.07403]



Beyond the (over-)simplified U_1 picture

The U_1 leptoquark is a massive vector...

→ It is not well-defined without specifying a UV origin

There are two ways to UV-complete the vector leptoquark

- U_1 as a gauge boson from an **extended symmetry**, like the W and Z in the SM

Minimal working setup: $SU(4) \times SU(3) \times SU(2)_L \times U(1) \rightarrow SU(3)_c \times SU(2)_L \times U(1)_Y$

[di Luzio, Greljo, Nardecchia 1708.08450; Calibbi, Crivellin, Li 1709.00692; Bordone, Cornella, JFM, Isidori 1712.01368; Greljo, Stefanek 1802.04274...]

- U_1 as a composite resonance from a **strongly coupled sector**, similar to QCD resonances

[Barbieri, Tesi, 1712.06844, Barbieri, Murphy, Senia, 1611.04930, Barbieri, Isidori, Pattori, Senia, 1512.01560,...]

Leptoquark vector companions

... but no matter the approach (or the model details), the U_1 leptoquark always comes along with vector companions



$$U^\alpha \sim (3, 1)_{2/3}$$

$$G'^a \sim (8, 1)_0 \text{ (heavy "gluon")}$$

$$Z' \sim (1, 1)_0$$

$$M_U \sim M_{Z'} \sim M_{G'} \sim \mathcal{O}(\text{TeV})$$

As a gauge boson

$$SU(4) \sim \begin{pmatrix} G'^a & U^\alpha \\ (U^\alpha)^* & Z' \end{pmatrix}$$

As a quark-lepton bound state

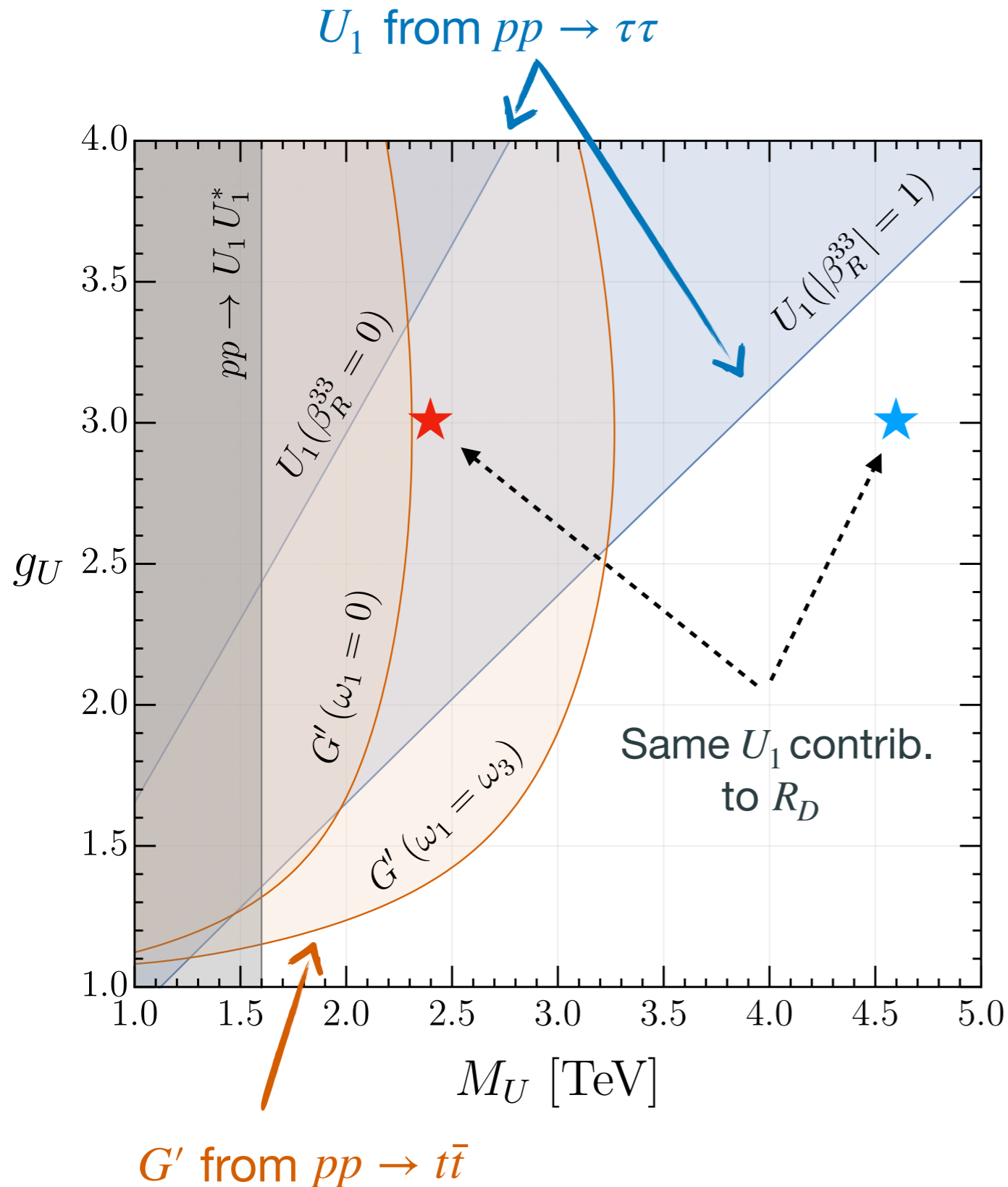
$$U \sim \langle \bar{Q}L \rangle$$

$$G' + Z' \sim \langle \bar{Q}Q \rangle$$

$$Z' \sim \langle \bar{L}L \rangle$$

Other states, e.g. **vector-like fermions** (important later!), are commonly needed

High-pT interplay with the new vectors



In particular models the U_1 , G' and Z' masses are related

$$M_{G'} = M_U \frac{g_U}{\sqrt{g_U^2 - g_c^2}} \sqrt{\frac{2\omega_3^2}{\omega_1^2 + \omega_3^2}}$$

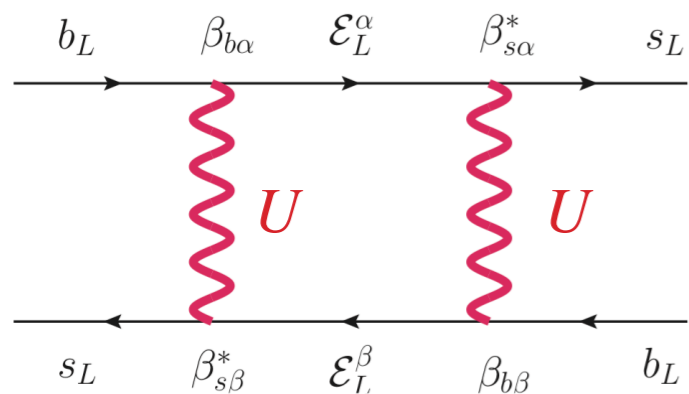
ω_i : scalar vevs

G' searches are very important for the LH leptoquark ($\beta_R = 0$)... but not so much for $\beta_R = 1$

Z' searches typically less relevant

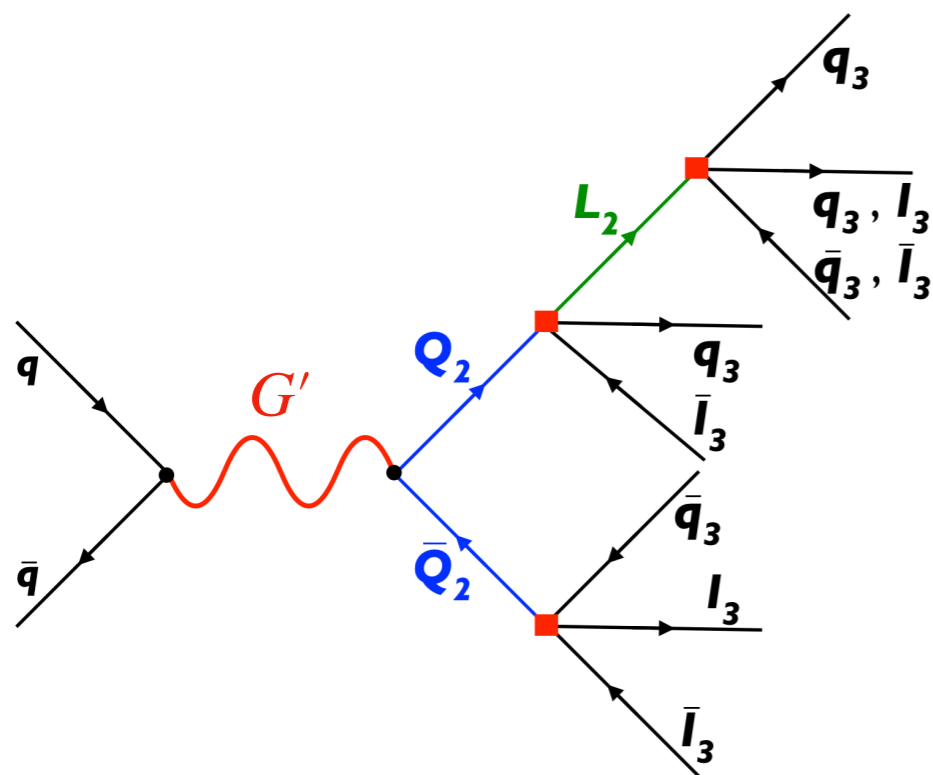
Exotic multi-jet plus multi-lepton signatures

Rather generically, vector-like fermions are needed to make U_1 loops finite. Similar to the SM case with the W and the prediction of the charm quark



$$C_{B_s\text{-mixing}}^{LL} \sim \Delta R_{D^{(*)}}^2 M_\chi^2$$

Vector-like fermions are expected to be among the lightest states in the theory!



Exotic multi-jet plus multi-lepton signatures within the reach of the LHC are predicted [rich signal with b-tags and τ -tags]

Similar existing SUSY searches by ATLAS (1706.03731) but a dedicated analysis is needed

Conclusions

Current data is inconclusive and the overall picture may change but...

... it is possible to find solutions to the flavor anomalies while remaining consistent with all the other data

Interesting **connections to the SM Yukawa structure** (hinting to a possible solution of the **SM flavor puzzle**)

Great **complementarity** between low-energy and high-pT

Going beyond simplified models is important

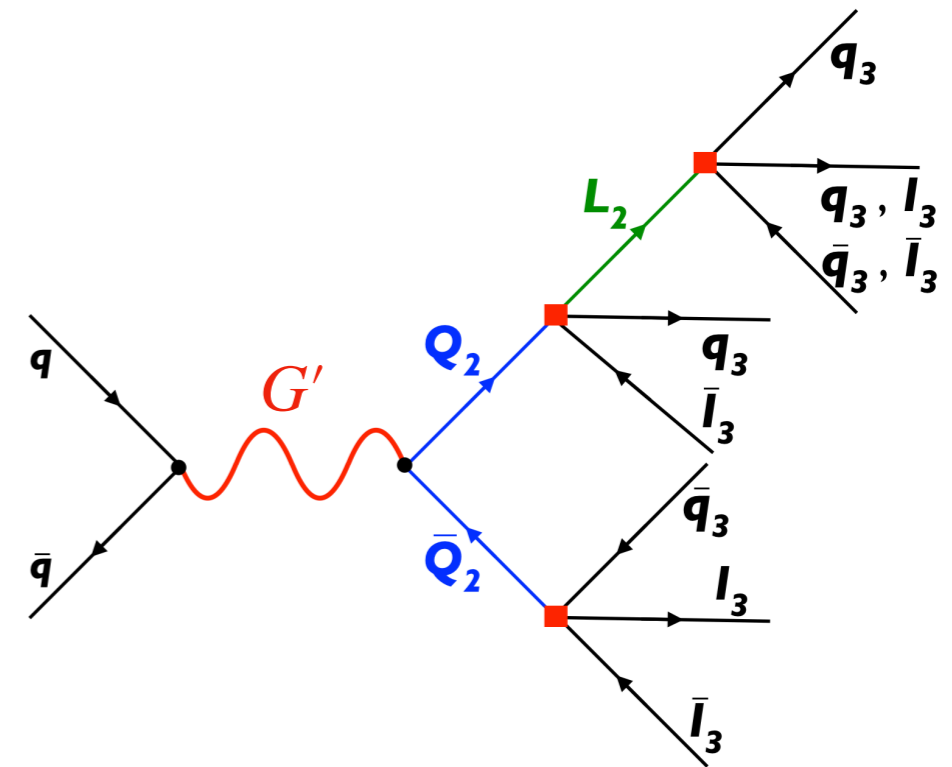
→ unexpected experimental signatures (G' , Z' , VL fermions,...)

If the anomalies are really pointing to NP, **new experimental indications** (both in high-pT and at low energies) should show up soon in several observables

...and summary

High-pT signatures

	LQ, Z'	G'	LQ LQ*
$pp \rightarrow$	$\tau\tau, \tau\nu$	$t\bar{t}, b\bar{b}$	$b\tau b\tau, \nu\nu, \dots$
	$\tau\mu$	jj	$b\tau b\mu, \nu b\mu, \dots$
	$\mu\mu$		$b\mu b\mu, \dots$



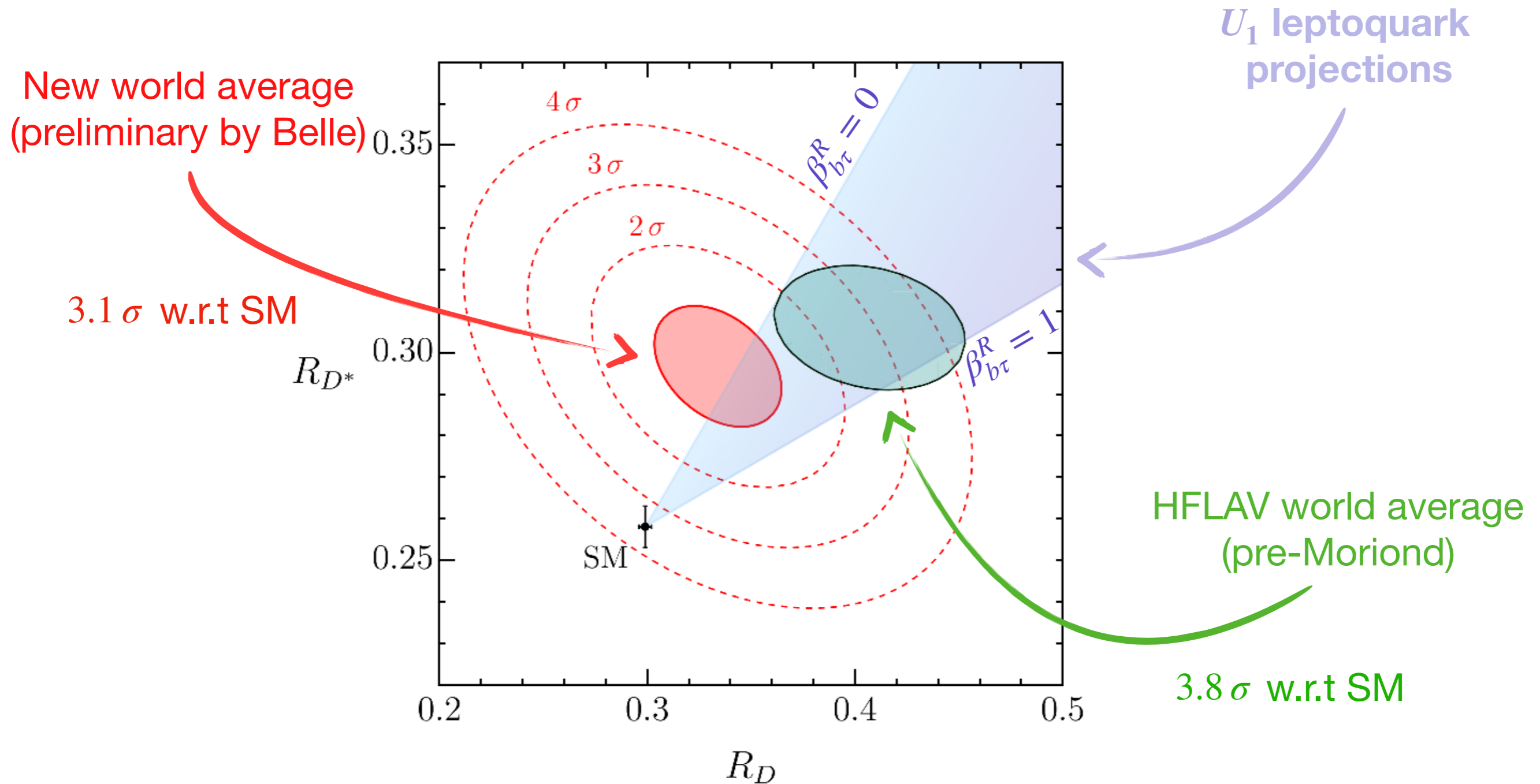
Generic characteristics of the new exotic states:

- Masses **around the TeV scale** [2-6 TeV]
- Flavor non-universal, with much **stronger couplings to 3rd generation**
- **Large widths** for Z' and G'

UFO model publicly available at <https://feynrules.irmp.ucl.ac.be/wiki/LeptoQuark>

Backup slides

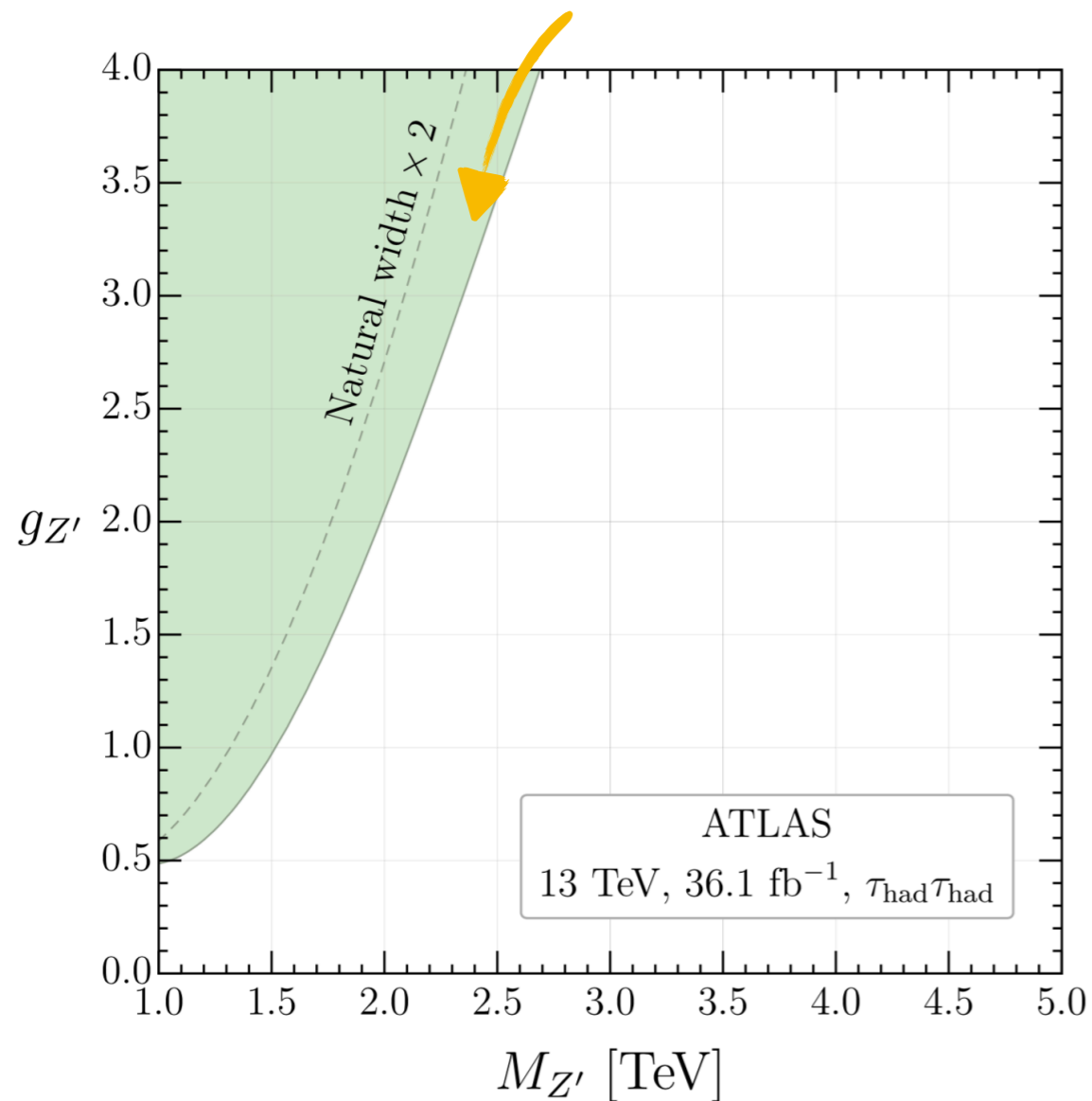
Which value of $\beta_{b\tau}^R$? $R(D^{(*)})$ projections



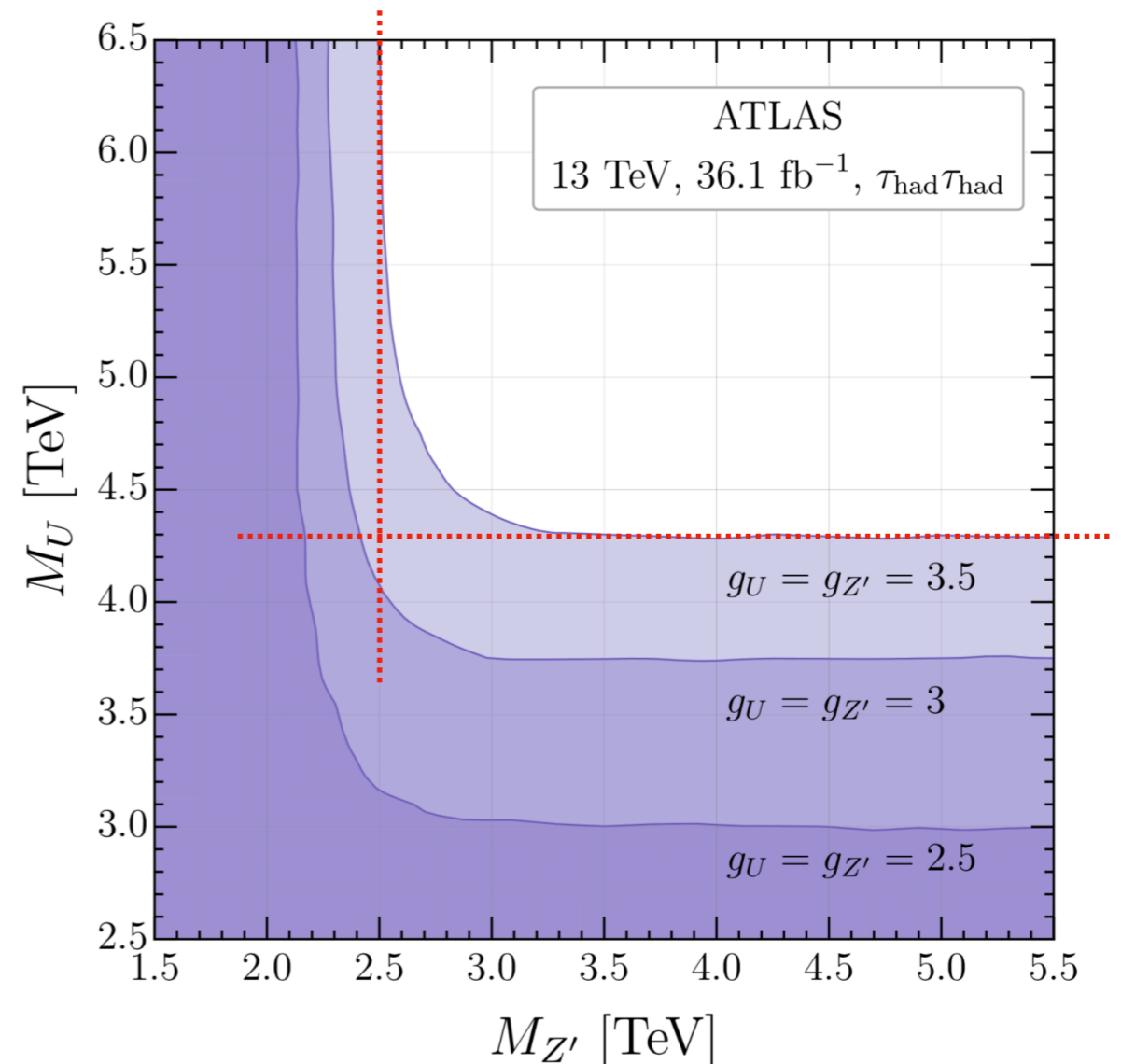
Differential distributions, polarizations,... could also be different from the SM

Recast of the $pp \rightarrow \tau\tau$ data: Z' search and combined

$pp \rightarrow \tau\tau$ tail (in m_T^{tot})



Distortions in the tails of the m_T^{tot} distribution, as in the U_1 case

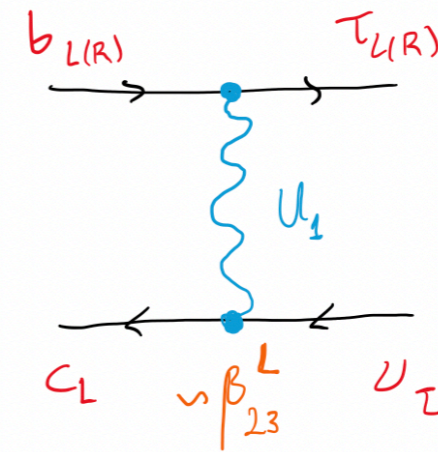
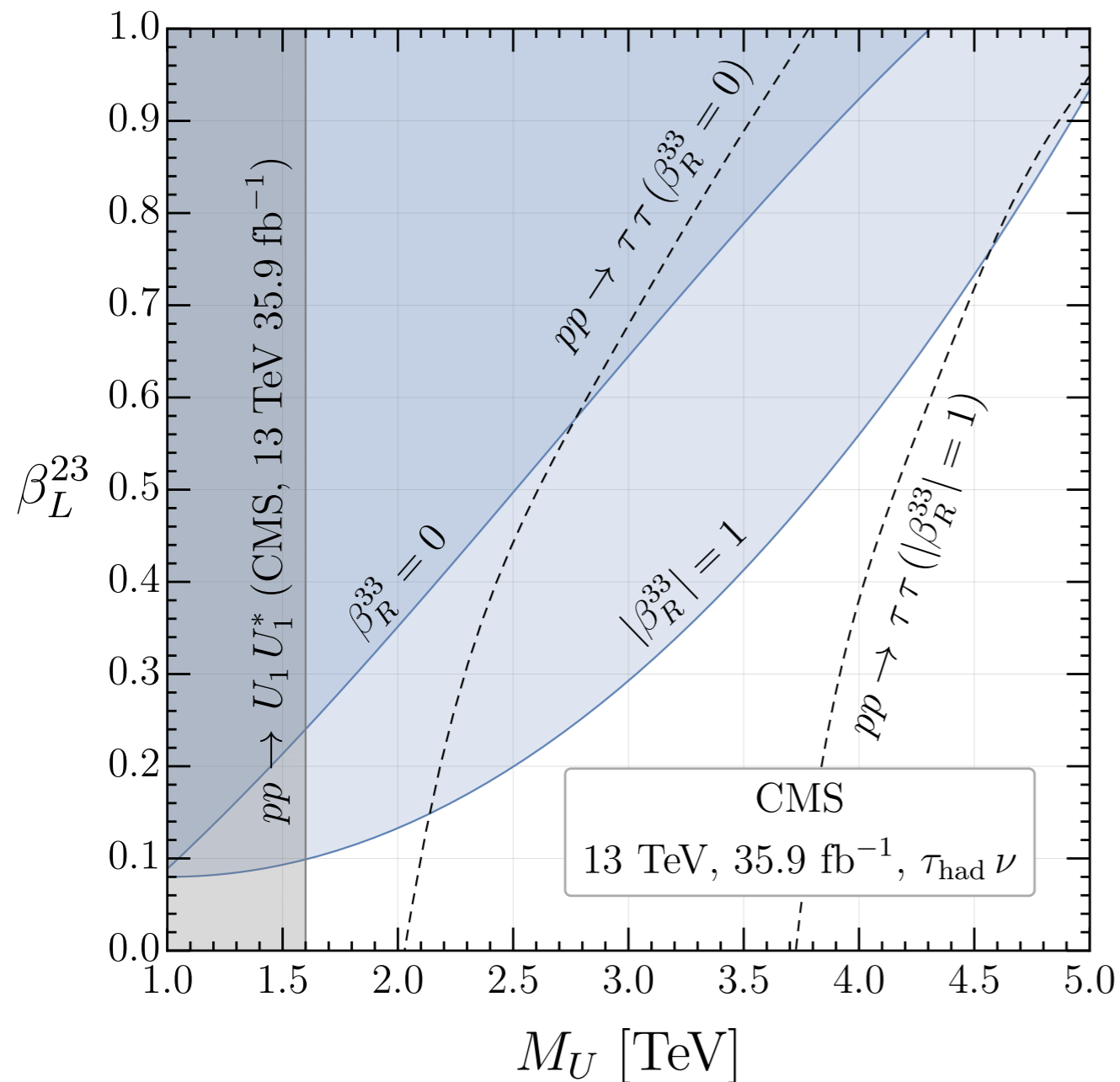


The mass region where the two particles overlap is small

Recast of the $pp \rightarrow \tau\nu$ data

[Baker, JFM, Isidori, König, 1901.10480]

$$\mathcal{L} \supset \frac{3}{\sqrt{2}} U_\mu (\beta_{23}^L \bar{Q}_2 \gamma^\mu L_3 + \bar{Q}_3 \gamma^\mu L_3 - \beta_R \bar{b}_R \gamma^\mu \tau_R)$$



Contrary to other NP scenarios for $R(D^{(*)})$, here $pp \rightarrow \tau\tau$ gives stronger limits ($\beta_{23}^L \sim 0.1$ preferred by $R(D^{(*)})$ fit)

Interesting prospects. Good complementarity between $pp \rightarrow \tau\tau$ and $pp \rightarrow \tau\nu$