

Simulations for wire BBLR compensation in HL-LHC

Stephane Fartoukh, Nikos Karastathis, Yannis Papaphilippou, Dario Pellegrini, Axel Poyet, <u>Kyriacos Skoufaris</u> and Guido Sterbini

CERN, Geneva



8th HL-LHC Collaboration Meeting, October 17, 2018

Contents

- Quantification and solution of the problem generated from the BBLR interactions.
- BBLR compensation with wire in HL-LHC v1.3.
- Conclusions



Quantification of the BBLR problem



Large DA (lifetime) degradation, at least 3σ , in the present of the beam-beam long range interaction.



Treatment of the perturbation generated by the BBLR interactions (I)

The integrated electromagnetic field (4D) that is generated by the BBLR encounters (assuming a round beam $\sigma_x = \sigma_y$) is given by:

$$\int_{-\infty}^{\infty} B_{\theta} ds = \frac{N_{p} q c \mu_{0} \beta_{st}}{2\pi} \frac{1 - Exp\left(-\frac{r^{2}}{2\sigma^{2}}\right)}{r}$$

This field is similar to the integrated magnetic filed from an "infinite" current currying wire.

$$\int_{-\infty}^{\infty} B_{\theta} \ ds = \frac{N_p \ q \ c \ \mu_0}{2\pi} \ \frac{1}{r}$$



Treatment of the perturbation generated by the **BBLR** interactions (II)

The wire is calibrated such as to **compensate the non-linear RDT** that are driven by the long-range beam-beam interactions.



s (m)

Configuration for the simulated machine

HL-LHC v1.3 configuration table						
Attributes	Symbol	Value [units]				
Energy	E	7000 [GeV]				
Bunch population (end of leveling)	Np	1.2x10 ¹¹ [1]				
Normalized emittance	٤n	2.5 [µm rad]				
Horizontal tune	Qx	62.31 or 62.315 [1]				
Vertical tune	Qy	60.32 [1]				
Horizontal chromaticity	ξx	15 [1]				
Vertical chromaticity	ξy	15 [1]				
Beta function at IP1 & IP5	β*	15 [cm]				
Half crossing angle at IP1 & IP5	Ф/2	210 – 250 [µrad]				
Octupole current	lo	-300 or 0 [A]				
Wires longitudinal position from the IP	Sw	+/- 195 [m]				
Number of BBLR kicks per IP per sited	NBBLR	25 [1]				
Number of wires per IP per sited	Nw	1 [1]				

In order to find the best BBLR compensation for different lattice configurations, a set of DA scans for **different wire current (lw)** and **wire transvers position (D)** are performed.



Lattice tunes	lo	Half crossing angle	Normalized crossing angle	Luminosity
62.315 ; 60.32	0 [A]	250 [µrad]	10.6	5x10 ³⁴ [cm ⁻² s ⁻¹]
HL-LH β	łC_V1.3 ; Q = (6 * = 15 [<i>cm</i>] ; I ₀ =	2.315, 60.320) ; $\xi = (15, 15)$; $\varepsilon_n =$ = 0 [A] ; $\frac{\phi_{15}}{2} = 250 [\mu rad]$; $N_p = 1.2$	2.5 [μm] ; 2 × 10 ¹¹	
0.324		Resonant HO:NO; L	R:15; WC:NO	spread is
0.322		HO:NO; L	R:NO; WC:NO generated lattice sext	from the upoles.
0.320 රි			Destructive (wings form generated	e tune spread nation) is from the long
0.318			range bear interactions	n beam S.
0.316			Using the v spread fror	vire the tune n the BBLR can
0.310	0.312	0.314 0.316 Q_X	5 0.318 De compen compressio	sated (wings on).





 $3.10 \downarrow 3.06$

3.08

3.10

3.12

 Q_X

3.14

3.16

HO15; LR15; WC15

3.18

3.20

1e-1



A single value that describe the min DA and correspond to a single trajectory in the phase space it is not enough to describe the effect of the wire on the different particles (different phase space trajectories). Thus, a more detailed DA analysis is performed.



Between 39° and 72° where the largest DA degradation occurred there are wire configurations that improve the average DA (effective mean DA) more than 1σ .



Plotting only the wire configurations with positive effective mean DA it is clear the beneficial effect of the wire (blue curve with rhombus).

A more detailed analysis of the strongest resonances vs angle is needed.



Without smoothing

Gaussian weighted average

Smoothed using

Wire configuration





Over all the angles the wire configurations with the best positive effective mean DA is slightly better than the case without wire.







The use of the wires not only improve the problem around the 45° but also preserve the very good DA for the rest of the angles.

Without smoothing

Smoothed using Gaussian weighted average



By fixing the transvers position of the wires at $D\sim10\sigma$ and adjusting the wire current of the left and right wire independently, it is again possible to obtain the same min DA as in the case of the optimized tune.







The wire configuration with the best effective mean DA is the one resulting in the preservation of a good DA for the different angles.

Without smoothing

Smoothed using Gaussian weighted average





Lattice tunes	lo	Half crossing angle	Normalized crossing angle	Luminosity
62.315 ; 60.32	-300 & 0 [A]	230 [µrad]	9.7	5.4x10 ³⁴ [cm ⁻² s ⁻¹]

Pushing the crossing angle at 230 μ rad, the use of the wire guaranty a min DA close to 6 σ and a boosted luminosity independent of the octupoles current.



The effective mean DA is increased more than 0.6σ in the case with negative octupoles and more than 1σ without octupoles.





Without smoothing

Smoothed using Gaussian weighted average



Conclusions

- In all the studies a good min or effective mean DA is found at transvers wire distance larger than 10σ and with wire current lower than 130 Am.
- The DA degradation seen when a non optimized tune is used can be overcome if the left and right wire current or the wire current and its transvers position are adjusted appropriately.
- The wire good performances are independent from the octupole current.
- A deeper understanding of the impact of the different resonances is needed through non-linear dynamics analysis (FMAs,...)
- A working point scan for optimized wire currents is on-going.
- The impact of the wire on lifetime will be assessed.





Thank you for your time!



Backup slides



Bassetti-Erskine formulas

▶ 4D treatment of the beam-beam long range interaction (Bassetti-Erskine)

$$B_{\theta} = -\frac{\beta_{st}}{c} \ E_r \ \rightarrow \ F_{\perp} = q \ E_r \ (1 + \beta_{we}\beta_{st}) = q \ E_{reff} \quad \text{and for } \sigma_x > \sigma_y:$$

$$\int_{-\infty}^{\infty} E_{xeff} \ ds = \frac{N_p \ q \ (1 + \beta_{we}\beta_{st})}{2\epsilon_0 \sqrt{2\pi(\sigma_x^2 - \sigma_y^2)}} \Im \left[\mathcal{F}\left(\frac{x + iy}{\sqrt{2(\sigma_x^2 - \sigma_y^2)}}\right) - E_{xp}\left(-\frac{x^2}{2\sigma_x^2} - \frac{y^2}{2\sigma_y^2}\right) \mathcal{F}\left(\frac{x\sigma_y^2 + iy\sigma_x^2}{\sigma_x\sigma_y\sqrt{2(\sigma_x^2 - \sigma_y^2)}}\right) \right]$$

$$\int_{-\infty}^{\infty} E_{yeff} \ ds = \frac{N_p \ q \ (1 + \beta_{we}\beta_{st})}{2\epsilon_0 \sqrt{2\pi(\sigma_x^2 - \sigma_y^2)}} \Re \left[\mathcal{F}\left(\frac{x + iy}{\sqrt{2(\sigma_x^2 - \sigma_y^2)}}\right) - E_{xp}\left(-\frac{x^2}{2\sigma_x^2} - \frac{y^2}{2\sigma_y^2}\right) \mathcal{F}\left(\frac{x\sigma_y^2 + iy\sigma_x^2}{\sigma_x\sigma_y\sqrt{2(\sigma_x^2 - \sigma_y^2)}}\right) \right]$$

All the quantities are measured from the center of the strong beam in the lab rest frame



Beam-beam field vs wire-like beam-beam field vs wire field



Skoufaris Kyriacos





Min DA HL-LHC v1.3, I = 1.2×10^{11} ppb, β^*_{IP1} =0.15m $(Q_X, Q_Y) = (62.315, 60.320), \phi/2=250\mu rad, \epsilon=2.5\mu m$

