A Beautiful Way of Going Beyond the Standard Model

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March 18, 2019, @Heavy Flavours 3, IIT Indore



A. Kundu (Calcutta U) BSM with flavour 18/03/19

This is T = 0 flavour physics and a sequel to the talk by Nita Sinha

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 - Quantum corrections induced by the heavy fields



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Are there any tensions with the SM?

Yes!!!

Not yet at the 5σ level to claim definite evidence of BSM Still, worth exploring.

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And there we go into the beautiful world of b-hadrons



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B-factories: past, present, and future

BaBar@SLAC :
$$e^+e^-$$
, 429 fb⁻¹, 4.7 \times 10⁸ $B\bar{B}$ pairs

Belle@KEK :
$$e^+e^-$$
, over 1 ab⁻¹, 7.72 × 10⁸ $B\bar{B}$ pairs

LHCb : 6.8 fb⁻¹ till 2017 (3.6 fb⁻¹ at 13 TeV) 7 TeV:
$$\sigma(pp \to b\bar{b}X) = (89.6 \pm 6.4 \pm 15.5)~\mu$$
b, scales linearly with \sqrt{s}

ATLAS and CMS also have dedicated flavour physics programme

LHCb:

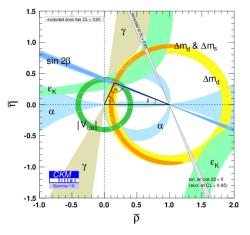
Upgrade I: $\mathcal{L}_{\rm int} > 50 \text{ fb}^{-1}$, $2 \times 10^{33} \text{ cm}^{-2} \text{s}^{-1}$

Phase II with HL-LHC: $\mathcal{L}_{int} > 300 \text{ fb}^{-1}$, $2 \times 10^{34} \text{ cm}^{-2} \text{s}^{-1}$

Belle-II:

 $\mathcal{L}_{\mathrm{int}} = 50~\mathrm{ab^{-1}}$ in 5 years, can go up even higher

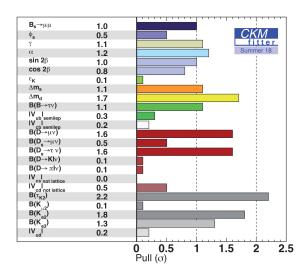




α	$91.6^{+1.7}_{-1.1}$
β direct β indirect β average	$22.14^{+0.69}_{-0.67} \\ 23.9 \pm 1.2 \\ 22.51^{+0.55}_{-0.40}$
γ	$65.81^{+0.99}_{-1.66}$

CKM paradigm rules !!!
NP has to be subdominant



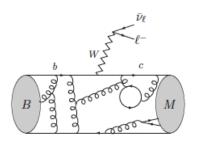




A few interesting anomalies



Experiment	R(D*)	R(D)
	0.332 +/- 0.024+/- 0.018	0.440 +/- 0.058 +/- 0.042
	0.293 +/- 0.038 +/- 0.015	0.375 +/- 0.064 +/- 0.026
	0.302 +/- 0.030 +/- 0.011	-
	0.336 +/- 0.027 +/- 0.030	-
BELLE	0.270 +/- 0.035 ⁺ 0.028 -0.025	-
LHCb	0.291 +/- 0.019 +/- 0.029	-
Average .txt	0.306 +/- 0.013 +/- 0.007	0.407 +/- 0.039 +/- 0.024

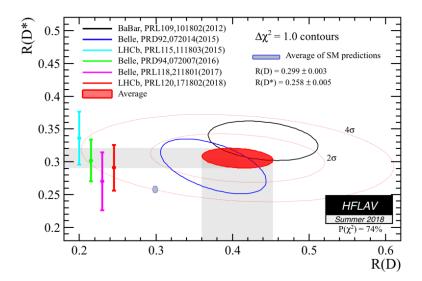


$$R(D^{(*)}) = \frac{BR(B \to D^{(*)} \tau \nu)}{BR(B \to D^{(*)} \ell \nu)}$$

	R(D)	R(D*)
D.Bigi, P.Gambino, Phys.Rev. D94 (2016) no.9, 094008 [arXiv:1606.08030 [hep-ph]]	0.299 +- 0.003	
[F.Bernlochner, Z.Ligeti, M.Papucci, D.Robinson, Phys.Rev. D95 (2017) no.11, 115008 [arXiv:1703.05330 [hep-ph]]	0.299 +- 0.003	0.257 +- 0.003
D.Bigi, P.Gambino, S.Schacht, JHEP 1711 (2017) 061 [arXiv:1707.09509 [hep-ph]]		0.260 +- 0.008
S.Jaiswal, S.Nandi, S.K.Patra, JHEP 1712 (2017) 060 [arXiv:1707.09977 [hep-ph]]	0.299 +- 0.004	0.257 +- 0.005
Arithmetic average	0.299 +- 0.003	0.258 +- 0.005

 2.3σ for R(D), 3.0σ for $R(D^*)$, 3.78σ combined with corr.







Longitudinal polarization fraction for $B o D^* au
u$

$$F_L = 0.457 \pm 0.010 \text{ (SM)}, \quad 0.60 \pm 0.09 \text{ (Belle 1903.03102)}$$

While we are talking about b o c au
u

$$R_{J/\psi} = \frac{\mathrm{BR}(B_c \to J/\psi \, \tau \nu)}{\mathrm{BR}(B_c \to J/\psi \, \ell \nu)}$$

= 0.71 ± 0.17 ± 0.18 (exp), 0.283 ± 0.048 (SM)

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And the neutral current $b o s \ell^+ \ell^-$

$$R_{K(K^*)} = \frac{\mathrm{BR}(B \to K(K^*)\mu^+\mu^-)}{\mathrm{BR}(B \to K(K^*)e^+e^-)}$$



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e or μ ? $B_s \to \phi \mu^+ \mu^-$ is also interesting \cdots



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$$R_K = 0.745^{+0.090}_{-0.074} \pm 0.036$$
 $q^2 \in [1:6] \text{ GeV}^2$,
 $R_{K^*}^{\text{low}} = 0.66^{+0.11}_{-0.07} \pm 0.03$ $q^2 \in [0.045:1.1] \text{ GeV}^2$,
 $R_{K^*}^{\text{central}} = 0.69^{+0.11}_{-0.07} \pm 0.05$ $q^2 \in [1.1:6] \text{ GeV}^2$.

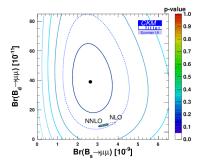
$$\frac{d}{dq^2} BR(B_s \to \phi \mu \mu) \Big|_{q^2 \in [1:6] \text{ GeV}^2}$$

$$= \begin{cases} \left(2.58^{+0.33}_{-0.31} \pm 0.08 \pm 0.19\right) \times 10^{-8} \text{ GeV}^{-2} & \text{(exp.)} \\ (4.81 \pm 0.56) \times 10^{-8} \text{ GeV}^{-2} & \text{(SM)}, \end{cases}$$

Is there some pattern?



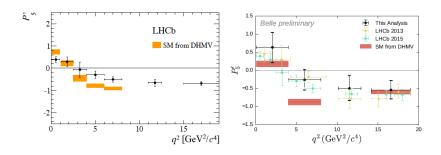
But $B_s/B_d \to \mu\mu$ is consistent with the SM (Only theory errors are from f_{B/B_s} and CKM. NLO EW, NNLO QCD, soft photon, large $\Delta\Gamma_s$ effects taken into account)



while $B \to K^* \mu \mu$ observable P_5' shows a deviation



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LHCb: two bins deviating by 2.8σ and 3.0σ Belle confirms with larger uncertainty CMS and ATLAS: Consistent with both LHCb/Belle and SM, large uncertainties



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Effective theory approach

$$\mathcal{H}_{ ext{eff}} = (\textit{CKM}) \sum_{i} \textit{C}_{i}\textit{O}_{i}$$

Main source of uncertainty: FF in $\langle M|\mathcal{H}_{\mathrm{eff}}|B\rangle$ Ratios are relatively insensitive

Example: $b \rightarrow s \mu^+ \mu^-$

$$\mathcal{H}_{\text{eff}}^{\text{SM}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i C_i(\mu) O_i(\mu)$$

with the relevant operators

$$O_{7} = \frac{e}{16\pi^{2}} m_{b} (\bar{s} \sigma_{\mu\nu} P_{R} b) F^{\mu\nu} , \quad C_{7} = -0.304$$

$$O_{9} = \frac{e^{2}}{16\pi^{2}} (\bar{s} \gamma^{\mu} P_{L} b) (\bar{\mu} \gamma_{\mu} \mu) , \quad C_{9} = 4.211$$

$$O_{10} = \frac{e^{2}}{16\pi^{2}} (\bar{s} \gamma^{\mu} P_{L} b) (\bar{\mu} \gamma_{\mu} \gamma_{5} \mu) , \quad C_{10} = -4.103$$



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Top-down:

UV complete theory \to Get C_i at high scale with proper matching \to Run down to $m_b \to$ Check consistency with data

Examples: leptoquarks, extra Z'

Bottom-up

Fit data with set of chosen operators \rightarrow Get the corresponding C_i



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How reliable are the form factors?

- $B \to K, D$: Only two FF, f_0 and f_1 , determined over the entire q^2 -range from lattice
- $B \to K^*, D^*$: Four FF, V, A_0, A_1, A_2 , lattice not yet complete, HQET is helpful, higher-order corrections can be estimated
- There can be more FF with BSM operators (like tensor)

Are there other pitfalls?

 D^* is detected as $D\pi$, take finite decay width into consideration

Reduces tension to 2.2σ

Chavez-Saab and Toledo, 1806.06997]

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- Tension for CC with $\ell= au$, comparable with SM tree (\sim 15% enhancement in amplitude)
- Tension for NC with $\ell=\mu$, comparable with SM loop only. Destructive interference needed



- \bullet Tension for CC with $\ell=\tau,$ comparable with SM tree (\sim 15% enhancement in amplitude)
- \bullet Tension for NC with $\ell=\mu,$ comparable with SM loop only. Destructive interference needed
- Consider a new operator involving au. Rotate the leptonic (au,μ) basis to (au',μ')

$$au = au' \cos heta + \mu' \sin heta \, , \quad
u_ au' =
u_ au \cos heta +
u_\mu \sin heta$$

• If the mixing angle θ is small, $\sin^2 \theta$ suppression makes the BSM tree comparable with SM loop



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$$\tau = \tau' \cos \theta + \mu' \sin \theta \,, \quad \nu_\tau' = \nu_\tau \cos \theta + \nu_\mu \sin \theta \,$$

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$$\mathcal{O}_{I} = \sqrt{3} A_{1} (\bar{Q}_{2L} \gamma^{\mu} L_{3L})_{3} (\bar{L}_{3L} \gamma_{\mu} Q_{3L})_{3} -2 A_{2} (\bar{Q}_{2L} \gamma^{\mu} L_{3L})_{1} (\bar{L}_{3L} \gamma_{\mu} Q_{3L})_{1}$$

- Only 3rd gen leptons, but can rotate to get muons
- Can give a good fit to R(D), $R(D^*)$, R_K , R_{K^*} , $R_{J/\psi}$, $\mathrm{BR}(B_s \to \phi \mu \mu)$, $\mathrm{BR}(B_s \to \mu \mu)$ and within limits for $b \to s+$ invisible and $B \to K^{(*)} \mu \tau$
- Much improved χ^2 compared to the SM

$$\chi^2 = \sum_{i=1}^8 \frac{\left(\mathcal{O}_i^{\text{exp}} - \mathcal{O}_i^{\text{th}}\right)^2}{\left(\Delta \mathcal{O}_i^{\text{exp}}\right)^2 + \left(\Delta \mathcal{O}_i^{\text{th}}\right)^2}$$

• $\chi^2/d.o.f.=1.5$ (this model), 6.1 (SM), with $A_1=0.028/{\rm TeV^2}$, $A_2=-2.90/{\rm TeV^2}$, $|\sin\theta|=0.018$, $C_9^{\rm NP}=-C_{10}^{\rm NP}=-0.61$



- ullet For these models $C_9^{
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 m NP}$: only LH currents
- $B_s o au^+ au^-$ gets sizable contribution from C_{10} , not C_9
- R_K and R_{K^*} need at least one of C_9 and C_{10} to be significant
- This is ruled out by $B_s \to \tau^+ \tau^-$ (as well as by ΔM_s)
- We need to break $C_0 = -C_{10}$ introduce RH currents

$$\mathcal{O}_{II} = \sqrt{3} A_{1} \left[-(Q_{2L}, Q_{3L})_{3} (L_{3L}, L_{3L})_{3} + \frac{1}{2} (Q_{2L}, L_{3L})_{3} (L_{3L}, Q_{3L})_{3} \right]$$

$$+ \sqrt{2} A_{5} (Q_{2L}, Q_{3L})_{1} \{ \tau_{R}, \tau_{R} \}$$

$$= \frac{3 A_{1}}{4} (c, b) (\tau, \nu_{\tau}) + \frac{3 A_{1}}{4} (s, b) (\tau, \tau) + A_{5} (s, b) \{ \tau, \tau \}$$

$$+ \frac{3 A_{1}}{4} (s, t) (\nu_{\tau}, \tau) + A_{5} (c, t) \{ \tau, \tau \} + \frac{3 A_{1}}{4} (c, t) (\nu_{\tau}, \nu_{\tau})$$

with $\{x,y\} \equiv \bar{x}_R \gamma^\mu y_R$, $(x,y) \equiv \bar{x}_L \gamma^\mu y_L \quad \forall \quad x,y$



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with
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, $(x,y) \equiv \bar{x}_L \gamma^{\mu} y_L \quad \forall \quad x,y$

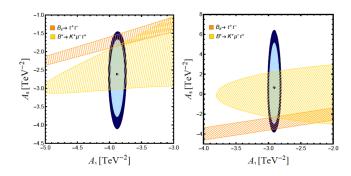


Can also play the same game with

$$\mathcal{O}_{\text{III}} = -\sqrt{3} A_1 (Q_{2L}, Q_{3L})_3 (L_{3L}, L_{3L})_3 + A_1 (Q_{2L}, Q_{3L})_1 (L_{3L}, L_{3L})_1 + \sqrt{2} A_5 (Q_{2L}, Q_{3L})_1 \{\tau_R, \tau_R\} = A_1 (c, b) (\tau, \nu_\tau) + A_1 (s, b) (\tau, \tau) + A_5 (s, b) \{\tau, \tau\} + A_1 (s, t) (\nu_\tau, \tau) + A_1 (c, t) (\nu_\tau, \nu_\tau) + A_5 (c, t) \{\tau, \tau\}$$

Best fit points	Model II	Model III
$ {\sf sin} heta $	0.016	0.016
A_1 in TeV $^{-2}$	-3.88	-2.91
A_5 in TeV $^{-2}$	-2.61	0.66





[Slightly different fit taking all ~ 160 observables into account. Also, Model I seems to be allowed. (Bhattacharya, Biswas, Calcuttawala, Patra, 1902.02796)]

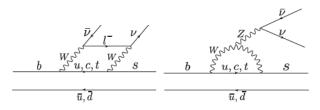


Something futuristic: $b \rightarrow s + \text{invisibles}$ at Belle-II

[Calcuttawala, AK, Nandi, Patra 2016]



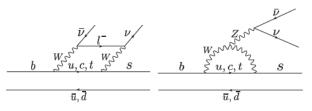
• SM: $b \rightarrow s \nu \bar{\nu}$, only penguin and box



- Not always related to $b \to s \ell^+ \ell^-$:
 - Leptons can be R with no neutrino counterpart

 - The invisibles can be something different!

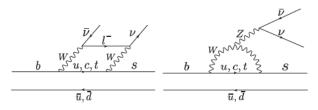
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 - $\bullet \epsilon_{ab} \bar{L}_L^a \gamma^\mu Q_L^b : b \to \nu, t \to \ell$
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- Observables: BR, $d\Gamma/dq^2$, $F'_T(q^2)$ (neutrinos), $F'_I(q^2)$ (light scalars



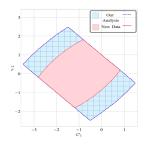
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$$\mathcal{H}_{ ext{eff}} = rac{4 \textit{G}_{\textit{F}}}{\sqrt{2}} \textit{V}_{\textit{tb}} \textit{V}_{\textit{ts}}^* \textit{C}_{\textit{SM}} \left[\textit{O}_{\textit{SM}} + \textit{C}_1' \textit{O}_{\textit{V}_1} + \textit{C}_2' \textit{O}_{\textit{V}_2}
ight] \, ,$$

$$\begin{split} \textit{O}_{\textit{SM}} = \textit{O}_{\textit{V}_1} &= \left(\bar{s}_{\textit{L}} \gamma^{\mu} \textit{b}_{\textit{L}} \right) \left(\bar{\nu}_{\textit{iL}} \gamma_{\mu} \nu_{\textit{iL}} \right) \,, \\ \textit{O}_{\textit{V}_2} &= \left(\bar{s}_{\textit{R}} \gamma^{\mu} \textit{b}_{\textit{R}} \right) \left(\bar{\nu}_{\textit{iL}} \gamma_{\mu} \nu_{\textit{iL}} \right) \,. \end{split}$$



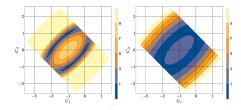
$${
m Br}(B o K(K^*) \nu \bar{\nu}) < 1.6(2.7) imes 10^{-5}$$

Detection efficiencies are small (Belle, 1303.3719)

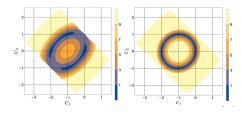
Mode	$N_{ m tot}$	$N_{ m sig}$	Significance	$\epsilon, 10^{-4}$	Upper limit
$B^+ \to K^+ \nu \bar{\nu}$	43	$13.3^{+7.4}_{-6.6}(\mathrm{stat}) \pm 2.3(\mathrm{syst})$	2.0σ	5.68	$< 5.5 \times 10^{-1}$
$B^0 \rightarrow K_s^0 \nu \bar{\nu}$	4	$1.8^{+3.3}_{-2.4}(\mathrm{stat}) \pm 1.0(\mathrm{syst})$	0.7σ	0.84	$< 9.7 \times 10^{-1}$
$B^+ \rightarrow K^{*+} \nu \bar{\nu}$	21	$-1.7^{+1.7}_{-1.1}(\mathrm{stat}) \pm 1.5(\mathrm{syst})$	-	1.47	$<4.0\times10^{-8}$
$B^0 \rightarrow K^{*0} \nu \bar{\nu}$	10	$-2.3^{+10.2}_{-3.5}(stat) \pm 0.9(syst)$	_	1.44	$< 5.5 \times 10^{-5}$



$B \rightarrow K^* \nu \bar{\nu}$ (50 and 2 ab⁻¹)



 F_T , $B o X_s
u ar{
u}$ (50 ab⁻¹)

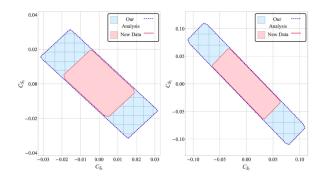




It can also be light invisible scalars (DM?)

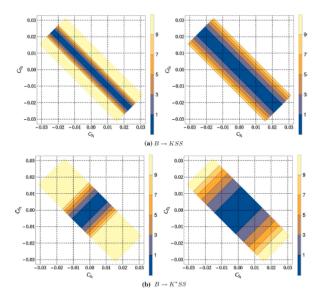
$$\mathcal{L}_{b\to sSS} = C_{S_1} m_b \bar{s}_L b_R S^2 + C_{S_2} m_b \bar{b}_L s_R S^2 + \text{H.c.}$$
 (1)

Higgs portal DM – $\langle S \rangle =$ 0, hSS coupling small to evade LHC limits





B o K and $B o K^*$ for $m_S = 0.5$ (1.8) GeV, $\mathcal{L}_{\mathrm{int}} = 50~\mathrm{ab}^{-1}$





To conclude:

- The CKM paradigm works quite well. BSM CPV needed to explain the baryon asymmetry, but it has to be subleading at least in the B sector (also in K and probably D)
- Flavour physics is the only tool to probe BSM if the scale is beyond the direct reach of LHC
- There are some intriguing anomalies. The pattern is not yet clear but LFU violation is indicated
- The third generation may be the window to BSM
- Watch out for LHCb and Belle-II for new results, confirmatory tests, and possible surprises!



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