# A Beautiful Way of Going Beyond the Standard Model 

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# This is $\mathbf{T}=0$ flavour physics and a sequel to the talk by Nita Sinha 

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## This is $\mathbf{T}=0$ flavour physics and a sequel to the talk by Nita Sinha

- SM passed all experimental tests, last missing piece discovered in 2014
- We all know this is not the ultimate theory
- DM, DE, BAU, $m_{\nu}$, hierarchy
- No luck so far on direct search front
- Have to look for indirect effects
- Quantum corrections induced by the heavy fields

Are there any tensions with the SM?

## Yes!!! <br> Not yet at the $5 \sigma$ level to claim definite evidence of BSM Still, worth exploring.

Circumstantial evidence is occasionally very convincing, as when you find
a trout in the milk.

- Arthur Conan Doyle

And there we go into the beautiful world of $b$-hadrons

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## B-factories: past, present, and future

BaBar@SLAC : $e^{+} e^{-}, 429 \mathrm{fb}^{-1}, 4.7 \times 10^{8} B \bar{B}$ pairs
Belle@KEK : $e^{+} e^{-}$, over $1 \mathrm{ab}^{-1}, 7.72 \times 10^{8} B \bar{B}$ pairs

## LHCb : $6.8 \mathrm{fb}^{-1}$ till 2017 ( $3.6 \mathrm{fb}^{-1}$ at 13 TeV )

$7 \mathrm{TeV}: \sigma(p p \rightarrow b \bar{b} X)=(89.6 \pm 6.4 \pm 15.5) \mu$ b, scales linearly with $\sqrt{s}$
ATLAS and CMS also have dedicated flavour physics programme
LHCb:
Upgrade I: $\mathcal{L}_{\text {int }}>50 \mathrm{fb}^{-1}, 2 \times 10^{33} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$
Phase II with HL-LHC: $\mathcal{L}_{\text {int }}>300 \mathrm{fb}^{-1}, 2 \times 10^{34} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$
Belle-II:
$\mathcal{L}_{\text {int }}=50 \mathrm{ab}^{-1}$ in 5 years, can go up even higher


| $\alpha$ | $91.6_{-1.1}^{+1.7}$ |
| :--- | :--- |
| $\beta$ direct | $22.14_{-0.67}^{+0.69}$ |
| $\beta$ indirect | $23.9 \pm 1.2$ |
| $\beta$ average | $22.51_{-0.40}^{+0.55}$ |
| $\gamma$ | $65.81_{-1.66}^{+0.99}$ |

CKM paradigm rules !!! NP has to be subdominant


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## A few interesting anomalies

| Experiment | R(D*) | R(D) |
| :--- | :--- | :--- |
| BaBar | $0.332+/-0.024+/-$ <br> 0.018 | $0.440+/-0.058$ <br> $+/-0.042$ |
| BELLE | $0.293+/-0.038+/-$ <br> 0.015 | $0.375+/-0.064$ <br> $+/-0.026$ |
| BELLE | $0.302+/-0.030+/-$ <br> 0.011 | - |
| LHCb | $0.336+/-0.027+/-$ <br> 0.030 | - |
| BELLE | $0.270+/-0.035+$ <br> 0.028 <br> -0.025 | - |
| LHCb | $0.291+/-0.019+/-$ <br> 0.029 | - |
| Average | $\mathbf{0 . 3 0 6}+/-\mathbf{0 . 0 1 3}+/-$ <br> txt | $\mathbf{0 . 0 0 7}$ |



$$
R\left(D^{(*)}\right)=\frac{\operatorname{BR}\left(B \rightarrow D^{(*)} \tau \nu\right)}{\operatorname{BR}\left(B \rightarrow D^{(*)} \ell \nu\right)}
$$

|  |  | R(D) |
| :--- | :--- | :--- |
| D.Bigi, P.Gambino, Phys.Rev. D94 (2016) no.9,094008 [arXiv:1606.08030 [hep-ph]] | $0.299+-0.003$ |  |
| F.Bernlochner, Z.Ligeti, M.Papucci, D.Robinson, Phys.Rev. D95 (2017) no.11, 115008 [arXiv:1703.05330 [hep-ph]] | $0.299+-0.003$ | $0.257+-0.003$ |
| D.Bigi, P.Gambino, S.Schacht, JHEP 1711 (2017) 061 [arXiv:1707.09509 [hep-ph]] |  |  |
| S.Jaiswal, S.Nandi, S.K.Patra, JHEP 1712 (2017) 060 [arXiv:1707.09977 [hep-ph]] | $0.260+-0.008$ |  |
| Arithmetic average | $0.299+-0.004$ | $0.257+-0.005$ |

$2.3 \sigma$ for $R(D), 3.0 \sigma$ for $R\left(D^{*}\right), 3.78 \sigma$ combined with corr.
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Longitudinal polarization fraction for $B \rightarrow D^{*} \tau \nu$

$$
F_{L}=0.457 \pm 0.010(\mathrm{SM}), \quad 0.60 \pm 0.09 \text { (Belle 1903.03102) }
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While we are talking about $b \rightarrow c \tau \nu$

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\begin{aligned}
R_{J / \psi} & =\frac{\operatorname{BR}\left(B_{c} \rightarrow J / \psi \tau \nu\right)}{\operatorname{BR}\left(B_{c} \rightarrow J / \psi \ell \nu\right)} \\
& =0.71 \pm 0.17 \pm 0.18(\exp ), \quad 0.283 \pm 0.048(\mathrm{SM})
\end{aligned}
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And the neutral current $b \rightarrow s \ell^{+} \ell^{-}$

$$
R_{K\left(K^{*}\right)}=\frac{\operatorname{BR}\left(B \rightarrow K\left(K^{*}\right) \mu^{+} \mu^{-}\right)}{\operatorname{BR}\left(B \rightarrow K\left(K^{*}\right) e^{+} e^{-}\right)}
$$

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$$

$e$ or $\mu$ ? $B_{s} \rightarrow \phi \mu^{+} \mu^{-}$is also interesting $\cdots$

$$
\left.\begin{array}{rl}
R_{K}=0.745_{-0.074}^{+0.090} \pm 0.036 & q^{2} \in[1: 6] \mathrm{GeV}^{2}, \\
R_{K^{*}}^{\text {low }}=0.66_{-0.07}^{+0.11} \pm 0.03 & q^{2} \in[0.045: 1.1] \mathrm{GeV}^{2}, \\
R_{K^{*}}^{\text {central }}=0.69_{-0.07}^{+0.11} \pm 0.05 & q^{2} \in[1.1: 6] \mathrm{GeV}^{2} .
\end{array}\right\} \begin{aligned}
& \left.\frac{d}{d q^{2}} \mathrm{BR}\left(B_{s} \rightarrow \phi \mu \mu\right)\right|_{q^{2} \in[1: 6] \mathrm{GeV}^{2}} \\
& = \begin{cases}\left(2.58_{-0.31}^{+0.33} \pm 0.08 \pm 0.19\right) \times 10^{-8} \mathrm{GeV}^{-2} & \text { (exp.) } \\
(4.81 \pm 0.56) \times 10^{-8} \mathrm{GeV}^{-2} & \text { (SM), }\end{cases}
\end{aligned}
$$

Is there some pattern?

But $B_{s} / B_{d} \rightarrow \mu \mu$ is consistent with the SM
(Only theory errors are from $f_{B / B_{s}}$ and CKM. NLO EW, NNLO QCD, soft photon, large $\Delta \Gamma_{s}$ effects taken into account)

while $B \rightarrow K^{*} \mu \mu$ observable $P_{5}^{\prime}$ shows a deviation


LHCb: two bins deviating by $2.8 \sigma$ and $3.0 \sigma$
Belle confirms with larger uncertainty
CMS and ATLAS: Consistent with both LHCb/Belle and SM, large uncertainties

Effective theory approach

$$
\mathcal{H}_{\mathrm{eff}}=(C K M) \sum_{i} C_{i} O_{i}
$$

Main source of uncertainty: FF in $\langle M| \mathcal{H}_{\text {eff }}|B\rangle$ Ratios are relatively insensitive

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Example: $b \rightarrow s \mu^{+} \mu^{-}$

$$
\mathcal{H}_{\mathrm{eff}}^{\mathrm{SM}}=-\frac{4 G_{F}}{\sqrt{2}} V_{t b} V_{t s}^{*} \sum_{i} C_{i}(\mu) O_{i}(\mu)
$$

with the relevant operators

$$
\begin{aligned}
O_{7} & =\frac{e}{16 \pi^{2}} m_{b}\left(\bar{s} \sigma_{\mu \nu} P_{R} b\right) F^{\mu \nu}, \quad C_{7}=-0.304 \\
O_{9} & =\frac{e^{2}}{16 \pi^{2}}\left(\bar{s} \gamma^{\mu} P_{L} b\right)\left(\bar{\mu} \gamma_{\mu} \mu\right), \quad C_{9}=4.211 \\
O_{10} & =\frac{e^{2}}{16 \pi^{2}}\left(\bar{s} \gamma^{\mu} P_{L} b\right)\left(\bar{\mu} \gamma_{\mu} \gamma_{5} \mu\right), \quad C_{10}=-4.103
\end{aligned}
$$

## Top-down:

UV complete theory $\rightarrow$ Get $C_{i}$ at high scale with proper matching $\rightarrow$ Run down to $m_{b} \rightarrow$ Check consistency with data

Examples: leptoquarks, extra $Z^{\prime}$

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UV complete theory $\rightarrow$ Get $C_{i}$ at high scale with proper matching $\rightarrow$ Run down to $m_{b} \rightarrow$ Check consistency with data

Examples: leptoquarks, extra $Z^{\prime}$

## Bottom-up:

Fit data with set of chosen operators $\rightarrow$ Get the corresponding $C_{i}$

How reliable are the form factors?

- $B \rightarrow K, D$ : Only two FF, $f_{0}$ and $f_{1}$, determined over the entire $q^{2}$-range from lattice
- $B \rightarrow K^{*}, D^{*}$ : Four FF, $V, A_{0}, A_{1}, A_{2}$, lattice not yet complete, HQET is helpful, higher-order corrections can be estimated


## Are there other pitfalls? <br> $D^{*}$ is detected as $D \pi$, take finite decay width into consideration <br> Reduces tension to $2.2 \sigma$ <br> For $B \rightarrow K^{(*)}$, no estimate for charmonium-dominated bins, have to be removed

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- There can be more FF with BSM operators (like tensor)

Are there other pitfalls?
$D^{*}$ is detected as $D \pi$, take finite decay width into consideration
Reduces tension to $2.2 \sigma$
[Chavez-Saab and Toledo, 1806.06997]
For $B \rightarrow K^{(*)}$, no estimate for charmonium-dominated bins, have to be removed

- Tension for CC with $\ell=\tau$, comparable with SM tree ( $\sim 15 \%$ enhancement in amplitude)
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- Tension for NC with $\ell=\mu$, comparable with SM loop only. Destructive interference needed
- Tension for CC with $\ell=\tau$, comparable with SM tree ( $\sim 15 \%$ enhancement in amplitude)
- Tension for NC with $\ell=\mu$, comparable with SM loop only. Destructive interference needed
- Consider a new operator involving $\tau$. Rotate the leptonic ( $\tau, \mu$ ) basis to $\left(\tau^{\prime}, \mu^{\prime}\right)$

$$
\tau=\tau^{\prime} \cos \theta+\mu^{\prime} \sin \theta, \quad \nu_{\tau}^{\prime}=\nu_{\tau} \cos \theta+\nu_{\mu} \sin \theta
$$

- If the mixing angle $\theta$ is small, $\sin ^{2} \theta$ suppression makes the BSM tree comparable with SM loop

$$
\begin{aligned}
\mathcal{O}_{\mathrm{I}}= & \sqrt{3} A_{1}\left(\bar{Q}_{2 L} \gamma^{\mu} L_{3 L}\right)_{3}\left(\bar{L}_{3 L} \gamma_{\mu} Q_{3 L}\right)_{3} \\
& -2 A_{2}\left(\bar{Q}_{2 L} \gamma^{\mu} L_{3 L}\right)_{1}\left(\bar{L}_{3 L} \gamma_{\mu} Q_{3 L}\right)_{1}
\end{aligned}
$$

- Only 3rd gen leptons, but can rotate to get muons
- Can give a good fit to $R(D), R\left(D^{*}\right), R_{K}, R_{K^{*}}, R_{J / \psi}, \operatorname{BR}\left(B_{s} \rightarrow \phi \mu \mu\right)$, $\operatorname{BR}\left(B_{s} \rightarrow \mu \mu\right)$ and within limits for $b \rightarrow s+$ invisible and $B \rightarrow K^{(*)} \mu \tau$
- Much improved $\chi^{2}$ compared to the SM

$$
\chi^{2}=\sum_{i=1}^{8} \frac{\left(\mathcal{O}_{i}^{\text {exp }}-\mathcal{O}_{i}^{\text {th }}\right)^{2}}{\left(\Delta \mathcal{O}_{i}^{\exp }\right)^{2}+\left(\Delta \mathcal{O}_{i}^{\text {th }}\right)^{2}}
$$

- $\chi^{2} /$ d.o.f. $=1.5$ (this model), $6.1(\mathrm{SM})$, with $A_{1}=0.028 / \mathrm{TeV}^{2}$, $A_{2}=-2.90 / \mathrm{TeV}^{2},|\sin \theta|=0.018, C_{9}^{\mathrm{NP}}=-C_{10}^{\mathrm{NP}}=-0.61$
- For these models $C_{9}^{\mathrm{NP}}=-C_{10}^{\mathrm{NP}}$ : only LH currents
- $B_{s} \rightarrow \tau^{+} \tau^{-}$gets sizable contribution from $C_{10}$, not $C_{9}$
- $R_{K}$ and $R_{K^{*}}$ need at least one of $C_{9}$ and $C_{10}$ to be significant - We need to break $C_{0}=-C_{10}$ - introduce RH currents
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- $R_{K}$ and $R_{K^{*}}$ need at least one of $C_{9}$ and $C_{10}$ to be significant
- This is ruled out by $B_{s} \rightarrow \tau^{+} \tau^{-}$(as well as by $\Delta M_{s}$ )
- We need to break $C_{0}=-C_{10}$ - introduce RH currents

$$
\begin{aligned}
\mathcal{O}_{\mathrm{II}} & =\sqrt{3} A_{1}\left[-\left(Q_{2 L}, Q_{3 L}\right)_{3}\left(L_{3 L}, L_{3 L}\right)_{3}+\frac{1}{2}\left(Q_{2 L}, L_{3 L}\right)_{3}\left(L_{3 L}, Q_{3 L}\right)_{3}\right] \\
& +\sqrt{2} A_{5}\left(Q_{2 L}, Q_{3 L}\right)_{1}\left\{\tau_{R}, \tau_{R}\right\} \\
& =\frac{3 A_{1}}{4}(c, b)\left(\tau, \nu_{\tau}\right)+\frac{3 A_{1}}{4}(s, b)(\tau, \tau)+A_{5}(s, b)\{\tau, \tau\} \\
& +\frac{3 A_{1}}{4}(s, t)\left(\nu_{\tau}, \tau\right)+A_{5}(c, t)\{\tau, \tau\}+\frac{3 A_{1}}{4}(c, t)\left(\nu_{\tau}, \nu_{\tau}\right)
\end{aligned}
$$

with $\{x, y\} \equiv \bar{x}_{R} \gamma^{\mu} y_{R}, \quad(x, y) \equiv \bar{x}_{L} \gamma^{\mu} y_{L} \quad \forall x, y$

Can also play the same game with

$$
\begin{aligned}
\mathcal{O}_{\text {III }} & =-\sqrt{3} A_{1}\left(Q_{2 L}, Q_{3 L}\right)_{3}\left(L_{3 L}, L_{3 L}\right)_{3}+A_{1}\left(Q_{2 L}, Q_{3 L}\right)_{1}\left(L_{3 L}, L_{3 L}\right)_{1} \\
& +\sqrt{2} A_{5}\left(Q_{2 L}, Q_{3 L}\right)_{1}\left\{\tau_{R}, \tau_{R}\right\} \\
& =A_{1}(c, b)\left(\tau, \nu_{\tau}\right)+A_{1}(s, b)(\tau, \tau)+A_{5}(s, b)\{\tau, \tau\} \\
& +A_{1}(s, t)\left(\nu_{\tau}, \tau\right)+A_{1}(c, t)\left(\nu_{\tau}, \nu_{\tau}\right)+A_{5}(c, t)\{\tau, \tau\}
\end{aligned}
$$

| Best fit points | Model II | Model III |
| :---: | :---: | :---: |
| $\|\sin \theta\|$ | 0.016 | 0.016 |
| $A_{1}$ in $\mathrm{TeV}^{-2}$ | -3.88 | -2.91 |
| $A_{5}$ in $\mathrm{TeV}^{-2}$ | -2.61 | 0.66 |


[Slightly different fit taking all $\sim 160$ observables
into account. Also, Model I seems to be allowed (Bhattacharya, Biswas, Calcuttawala, Patra, 1902.02796)]

## Something futuristic: $b \rightarrow s+$ invisibles at Belle-II

[Calcuttawala, AK, Nandi, Patra 2016]

- SM: $b \rightarrow s \nu \bar{\nu}$, only penguin and box

(1) Leptons can be R with no neutrino counterpart
(-) The invisibles can be something different!
- SM: $b \rightarrow s \nu \bar{\nu}$, only penguin and box

- Not always related to $b \rightarrow s \ell^{+} \ell^{-}$:
(1) Leptons can be R with no neutrino counterpart
(2) $\epsilon_{a b} \bar{L}_{L}^{a} \gamma^{\mu} Q_{L}^{b}: b \rightarrow \nu, t \rightarrow \ell$
(3) The invisibles can be something different!
- SM: $b \rightarrow s \nu \bar{\nu}$, only penguin and box

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(2) $\epsilon_{a b} \bar{L}_{L}^{a} \gamma^{\mu} Q_{L}^{b}: b \rightarrow \nu, t \rightarrow \ell$
(3) The invisibles can be something different!
- Observables:
$\mathrm{BR}, d \Gamma / d q^{2}, F_{T}^{\prime}\left(q^{2}\right)$ (neutrinos), $F_{L}^{\prime}\left(q^{2}\right)$ (light scalars)

$$
\mathcal{H}_{\mathrm{eff}}=\frac{4 G_{F}}{\sqrt{2}} V_{t b} V_{t s}^{*} C_{S M}\left[O_{S M}+C_{1}^{\prime} O_{V_{1}}+C_{2}^{\prime} O_{V_{2}}\right]
$$

$$
\begin{aligned}
& O_{S M}=O_{V_{1}}=\left(\bar{s}_{L} \gamma^{\mu} b_{L}\right)\left(\bar{\nu}_{i L} \gamma_{\mu} \nu_{i L}\right), \\
& \\
& O_{V_{2}}=\left(\bar{s}_{R} \gamma^{\mu} b_{R}\right)\left(\bar{\nu}_{i L} \gamma_{\mu} \nu_{i L}\right) . \\
& \operatorname{Br}\left(B \rightarrow K\left(K^{*}\right) \nu \bar{\nu}\right)<1.6(2.7) \times 10^{-5}
\end{aligned}
$$



Detection efficiencies are small (Belle, 1303.3719)

| Mode | $N_{\text {tot }}$ | $N_{\text {sig }}$ | Significance | $\epsilon, 10^{-4}$ | Upper limit |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $B^{+} \rightarrow K^{+} \nu \bar{\nu}$ | 43 | $13.3_{-6.6}^{+7.4}$ (stat) $\pm 2.3$ (syst) | $2.0 \sigma$ | 5.68 | $<5.5 \times 10^{-5}$ |
| $B^{0} \rightarrow K_{s}^{0} \nu \bar{\nu}$ | 4 | $1.8_{-2.4}^{+3.3}$ (stat) $\pm 1.0($ syst $)$ | $0.7 \sigma$ | 0.84 | $<9.7 \times 10^{-5}$ |
| $B^{+} \rightarrow K^{*+} \nu \bar{\nu}$ | 21 | $-1.7_{-1.1}^{+1.7}($ stat $) \pm 1.5($ syst $)$ | - | 1.47 | $<4.0 \times 10^{-5}$ |
| $B^{0} \rightarrow K^{* 0} \nu \bar{\nu}$ | 10 | $-2.3_{-3.5}^{+10.2}$ (stat) $\pm 0.9$ (syst) | - | 1.44 | $<5.5 \times 10^{-5}$ |

$B \rightarrow K^{*} \nu \bar{\nu}\left(50\right.$ and $\left.2 \mathrm{ab}^{-1}\right)$

$F_{T}, B \rightarrow X_{s} \nu \bar{\nu}\left(50 \mathrm{ab}^{-1}\right)$


It can also be light invisible scalars (DM?)

$$
\begin{equation*}
\mathcal{L}_{b \rightarrow s S S}=C_{S_{1}} m_{b} \bar{s}_{L} b_{R} S^{2}+C_{S_{2}} m_{b} \bar{b}_{L} s_{R} S^{2}+\text { H.c. } \tag{1}
\end{equation*}
$$

Higgs portal DM $-\langle S\rangle=0$, hSS coupling small to evade LHC limits

$B \rightarrow K$ and $B \rightarrow K^{*}$ for $m_{S}=0.5(1.8) \mathrm{GeV}, \mathcal{L}_{\mathrm{int}}=50 \mathrm{ab}^{-1}$


## To conclude:

- The CKM paradigm works quite well. BSM CPV needed to explain the baryon asymmetry, but it has to be subleading at least in the $B$ sector (also in $K$ and probably $D$ )
- Flavour physics is the only tool to probe BSM if the scale is beyond the direct reach of LHC


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- There are some intriguing anomalies. The pattern is not yet clear but LFU violation is indicated
- The third generation may be the window to BSM.
- Watch out for LHCb and Belle-II for new results, confirmatory tests, and possible surprises!


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> Thank you!

