

# **J/ $\psi$ production as a function of charged-particle multiplicity in pp collisions at $\sqrt{s} = 13$ TeV with ALICE at the LHC**

***Dhananjaya Thakur***

***(For the ALICE Collaboration)***

***Indian Institute of Technology Indore, India***



**3<sup>rd</sup> Heavy Flavor Meet**

**18<sup>th</sup> -20<sup>th</sup> March, 2019, IIT Indore, India**



# Heavy flavor vs. multiplicity

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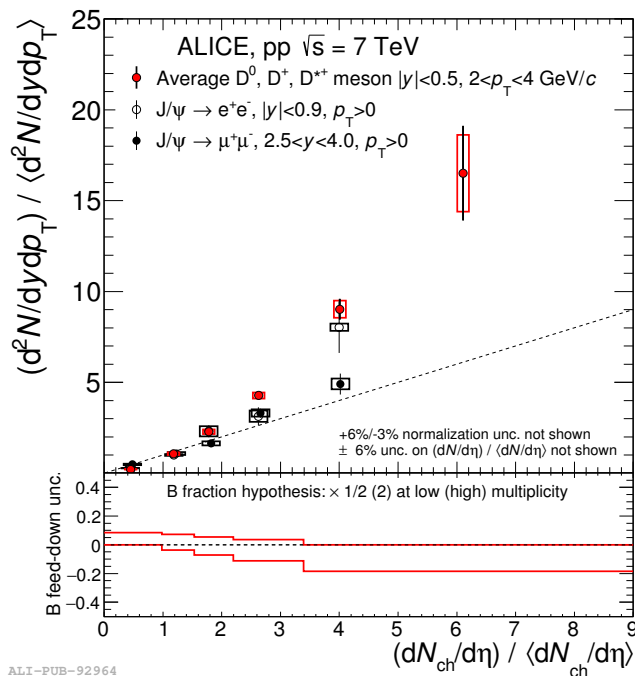
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e.g  $g \rightarrow c\bar{c}$ ,  $q\bar{q} \rightarrow c\bar{c}$  etc.
- ❖ In particular, these measurements help in understanding the interplay between the soft and hard mechanisms.

# J/ψ(D) meson vs. multiplicity in pp collisions



JHEP09(2015)148 (ALICE Collaboration)

❖ ALICE has measured D mesons and J/ψ as a function of multiplicity.

❖ Mid rapidity J/ψ(D) yield ( $|\eta| < 0.9$ ) vs. mid rapidity charged particles ( $|\eta| < 1.0$ )

**Observation:**

➤ A stronger than linear increase towards higher multiplicity

❖ Forward rapidity ( $2.5 < y < 4.0$ ) J/ψ yield vs. mid rapidity charged particles ( $|\eta| < 1.0$ )

**Observation:**

➤ Nearly linear increase

❖ The increase of J/ψ and D meson yield with multiplicity reveals that MPIs are relevant for J/ψ and D meson production.

❖ The higher multiplicity reach for pp collisions at  $\sqrt{s} = 13$  TeV might help to provide further insight.

# The ALICE Detector

## EMCAL

High momentum electron  
Triggering  
PID

## ITS (SPD+SDD+SSD)

Tracking  
Vertexing  
Multiplicity ( $|\eta| < 1.0$ )

## V0

Triggering

Charged-particle multiplicity is measured using the number of SPD tracklets in  $|\eta| < 1$ .

$e^+$   
 $J/\psi \rightarrow e^+e^-$

$e^-$

## TPC

$-0.9 < \eta < 0.9$

Di- $e$ :  $p_T > 0$  GeV/c

Single- $e$ :  $p_T > 0.2$  GeV/c

## ALICE Central barrel

$|\eta| < 0.9$

$J/\psi(\Upsilon) \rightarrow \mu^+\mu^-$

## MUON

$-4.0 < \eta < -2.5$

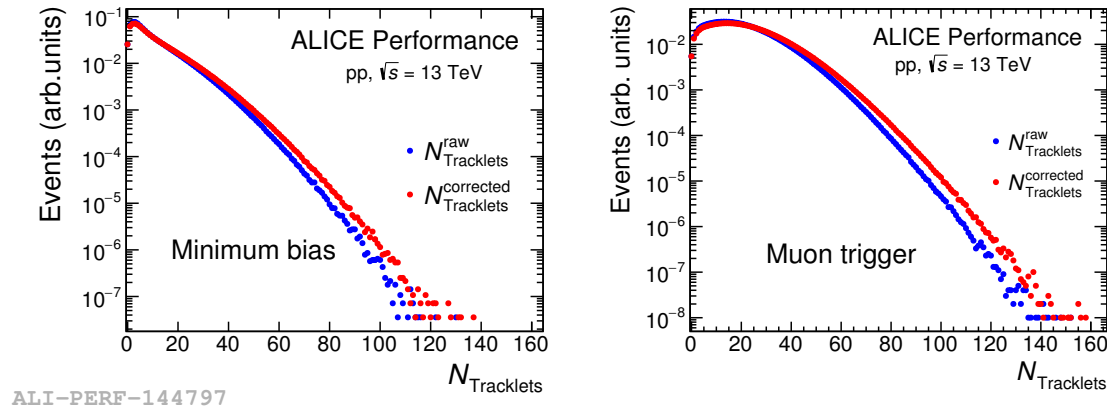
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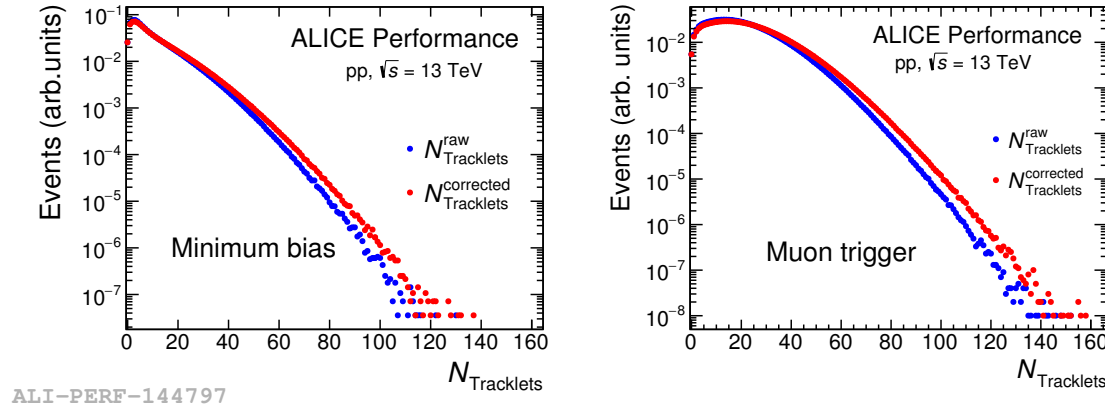
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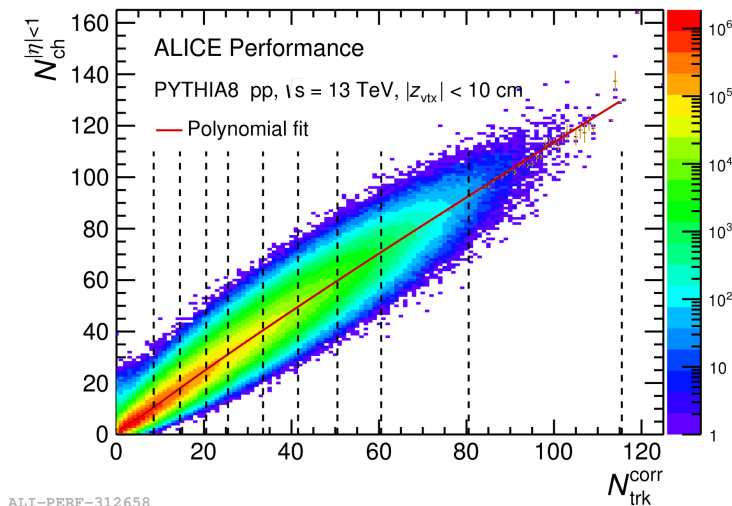


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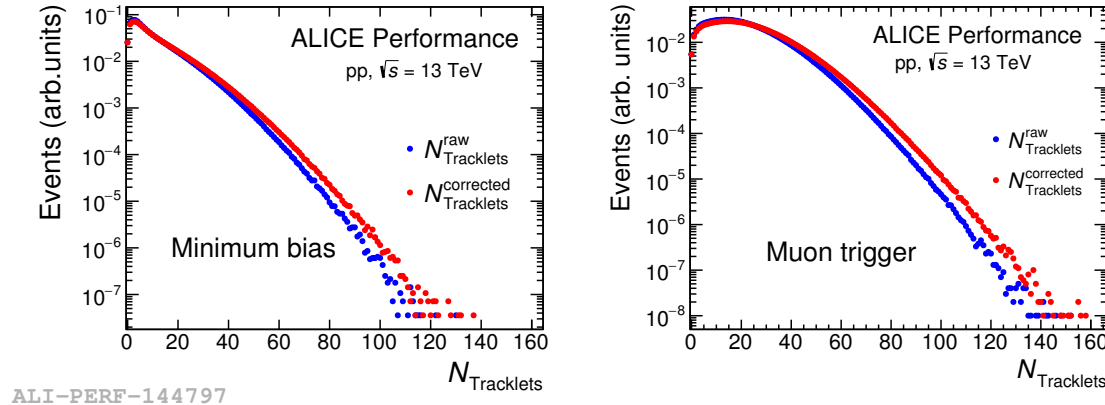


- ❖ The conversion function from corrected- $N_{\text{trk}}$  to  $dN_{\text{ch}}/d\eta$ , “ $f$ ”, is estimated using Monte-Carlo simulations.

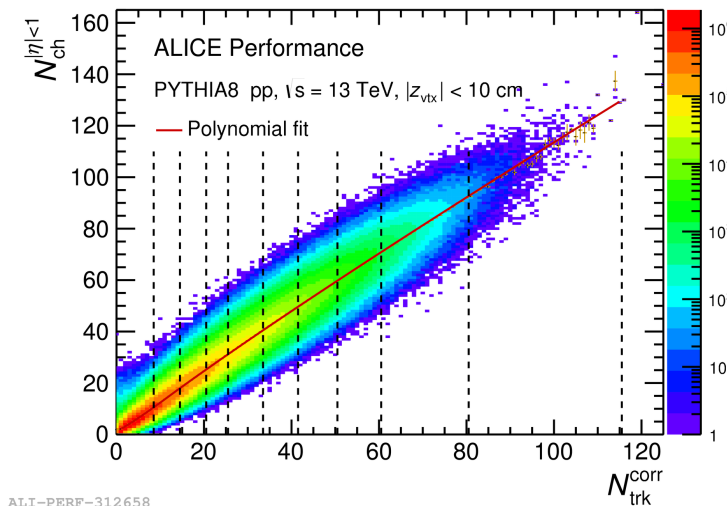


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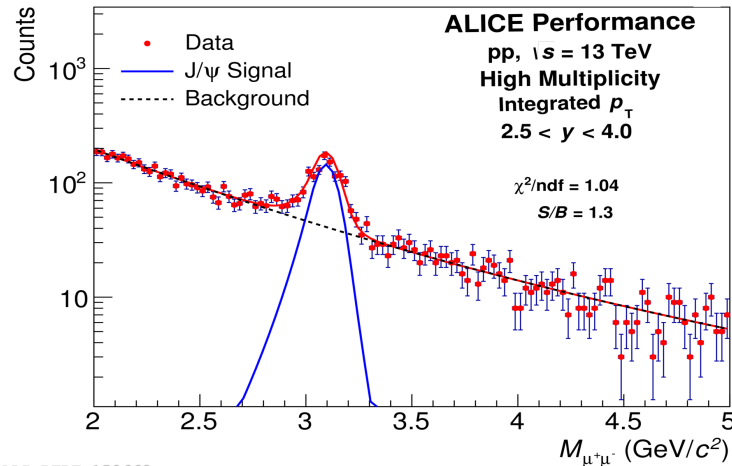
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$$\frac{\langle dN_{\text{ch}}/d\eta \rangle_i}{\langle dN_{\text{ch}}/d\eta \rangle_{\text{INEL} > 0}} = \frac{f(\langle N_{\text{trk}}^{\text{corr}} \rangle_i)}{\langle dN_{\text{ch}}/d\eta \rangle_{\text{INEL} > 0}}$$

# Signal Extraction

## Dimuon, pp at $\sqrt{s} = 13$ TeV

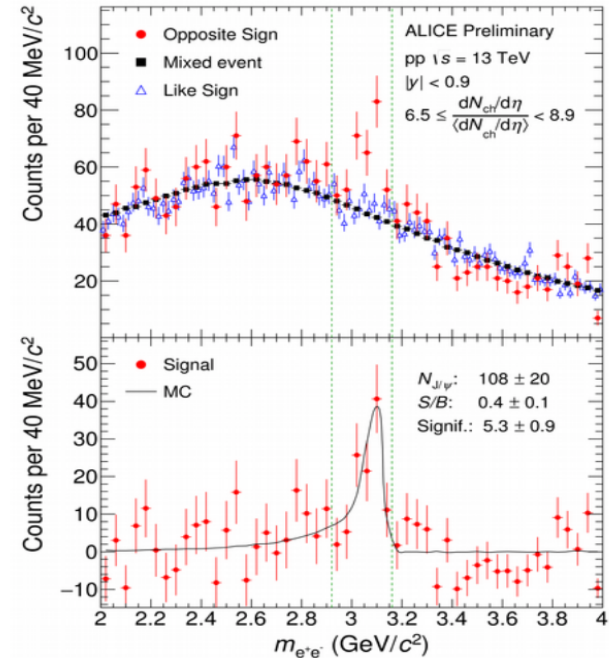


ALI-PERF-150668

### Fitting procedure based on

- **Signal:**  
Extended Crystal Ball function
- **Background:**  
variable-width Gaussian function

## Dielectron, pp at $\sqrt{s} = 13$ TeV



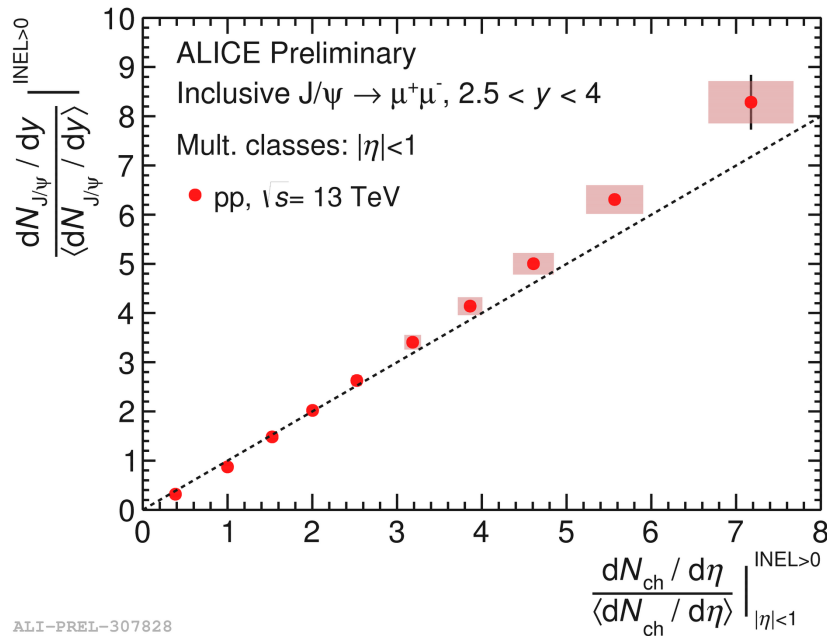
ALI-PREL-118405

- **Signal:**  
bin by bin counting in  $2.92 - 3.16 \text{ GeV}/c^2$
- **Background:**  
Subtracted using normalized  
like-sign pair distribution



# J/ψ vs. multiplicity in pp at $\sqrt{s} = 13$ TeV

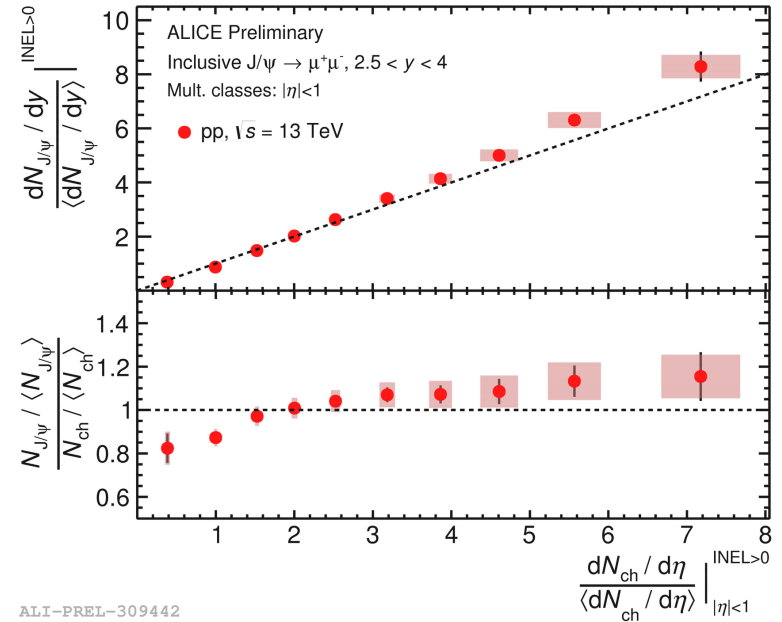
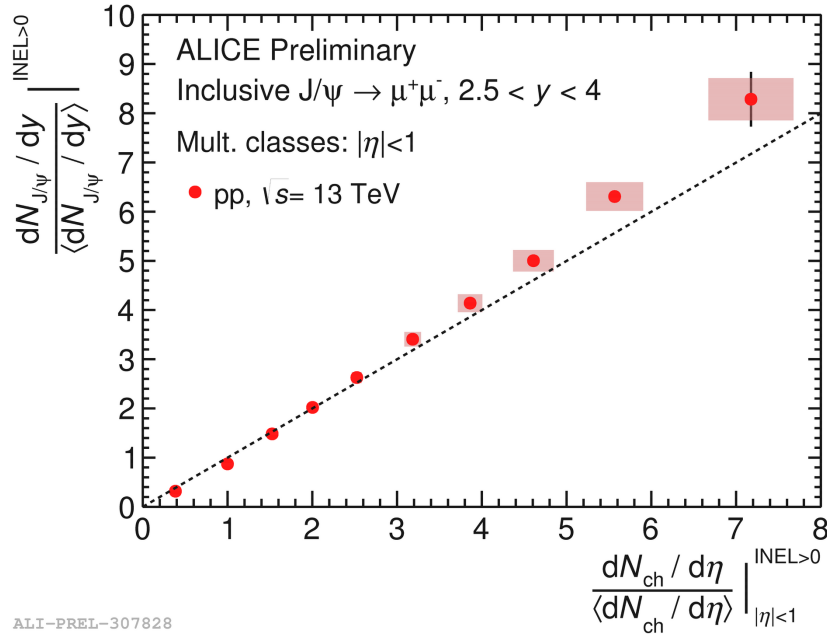
❖ Highest ever multiplicity reached by ALICE in pp collisions



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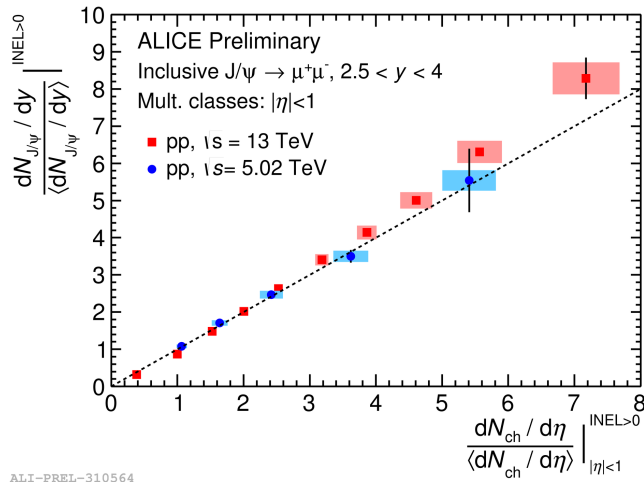
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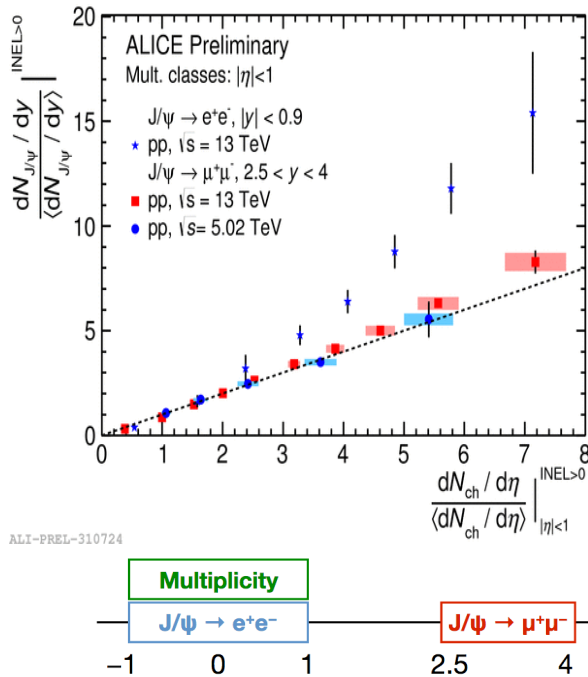
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# $\sqrt{s}$ and rapidity dependence of J/ψ vs. multiplicity



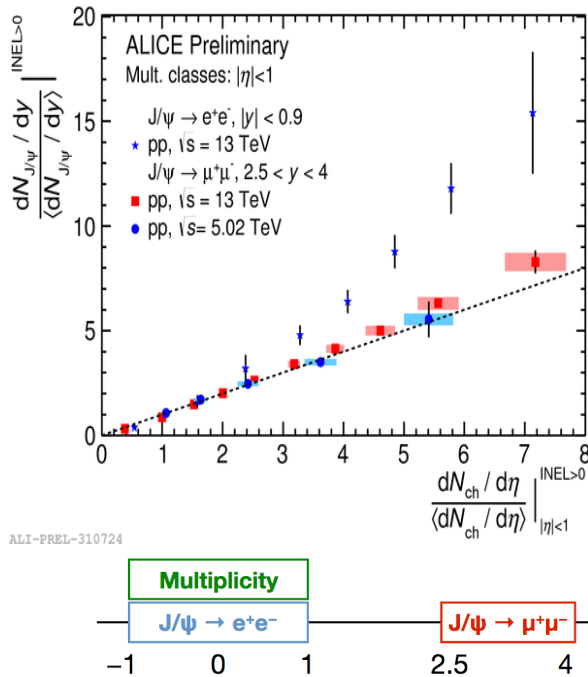
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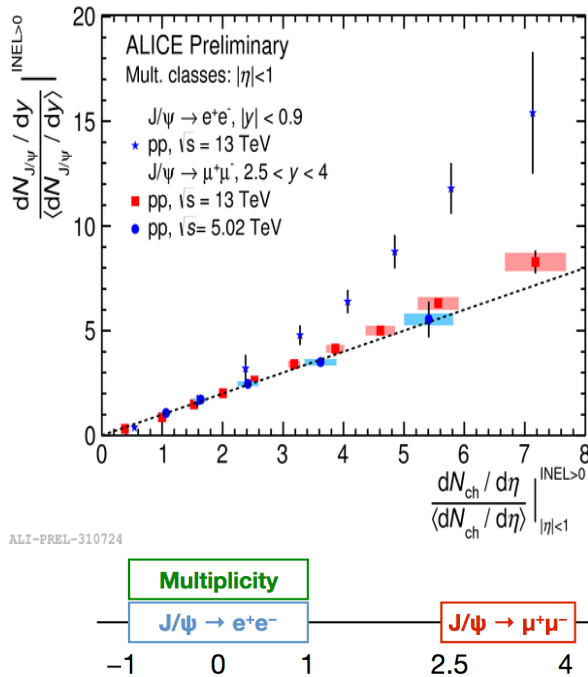
- Similar trends observed for forward-rapidity  $J/\psi$  in pp collisions at  $\sqrt{s} = 5$  and 13 TeV.
- Faster than linear scaling with multiplicity for  $J/\psi$  at mid-rapidity in pp at  $\sqrt{s} = 13$  TeV
  - i.e. w/o rapidity gap between signal and multiplicity estimator

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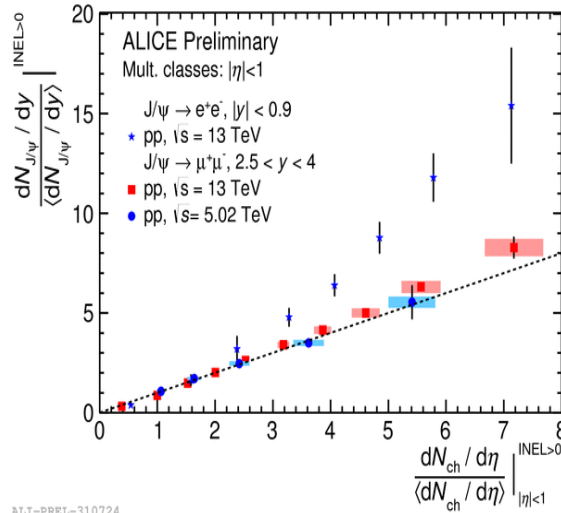
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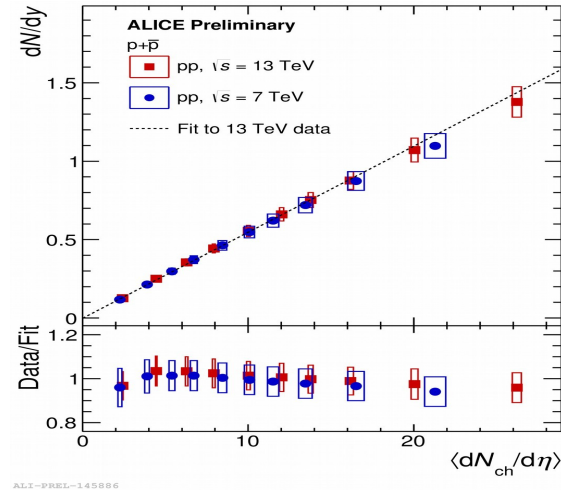
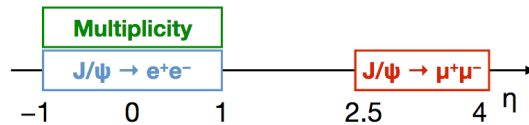
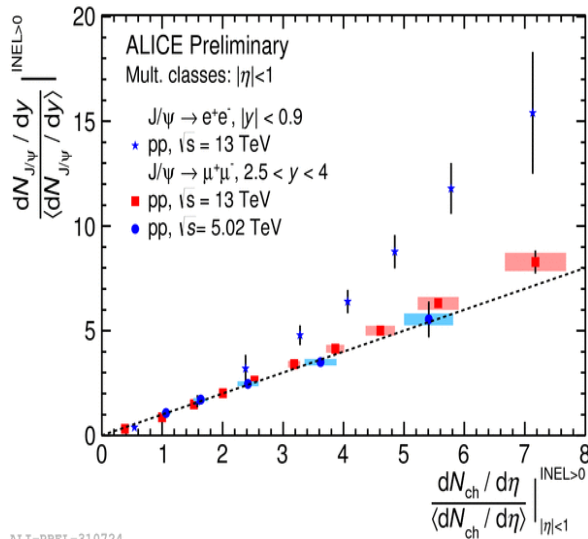


ALI-PREL-310724



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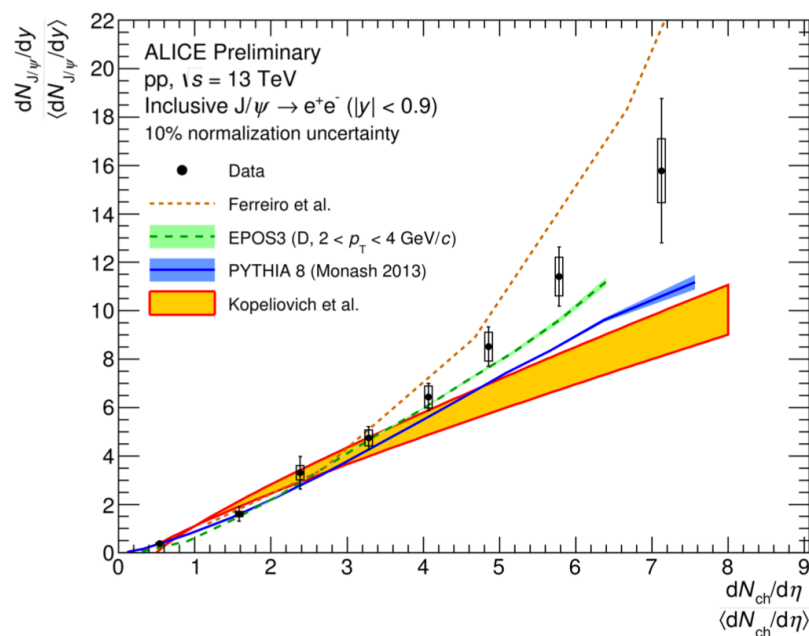
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- Physics scenario such as saturation would influence differently HF production at mid or forward rapidity
- Linear trend is also observed for light flavors, where there is rapidity gap in the measurement of multiplicity (by forward rapidity detector V0M) and signal (by mid-rapidity detector ).
  - ▶ Event activity associated with production of light flavor and heavy flavor is nearly same.



# Model Study: Mid-rapidity J/ψ yield vs. multiplicity in pp at $\sqrt{s} = 13$ TeV



- ❖ Stronger than linear increase of J/ψ yield is observed towards higher multiplicity

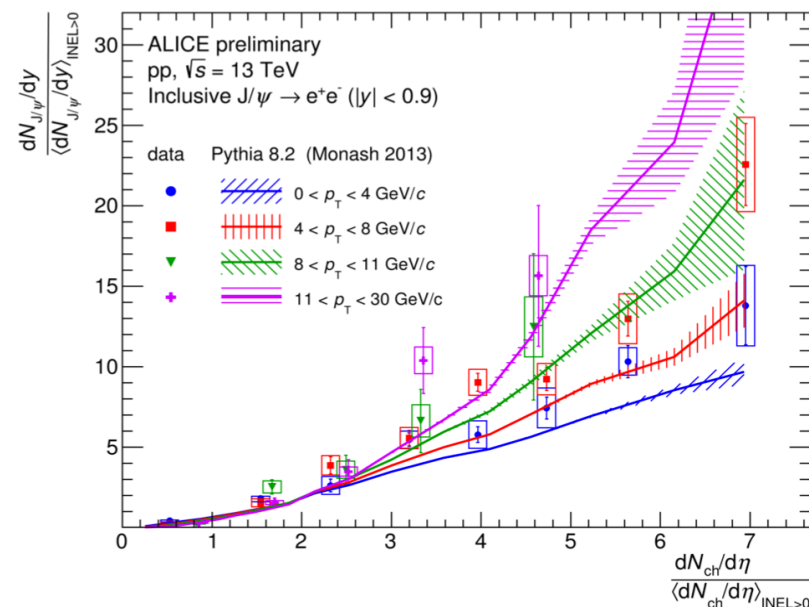
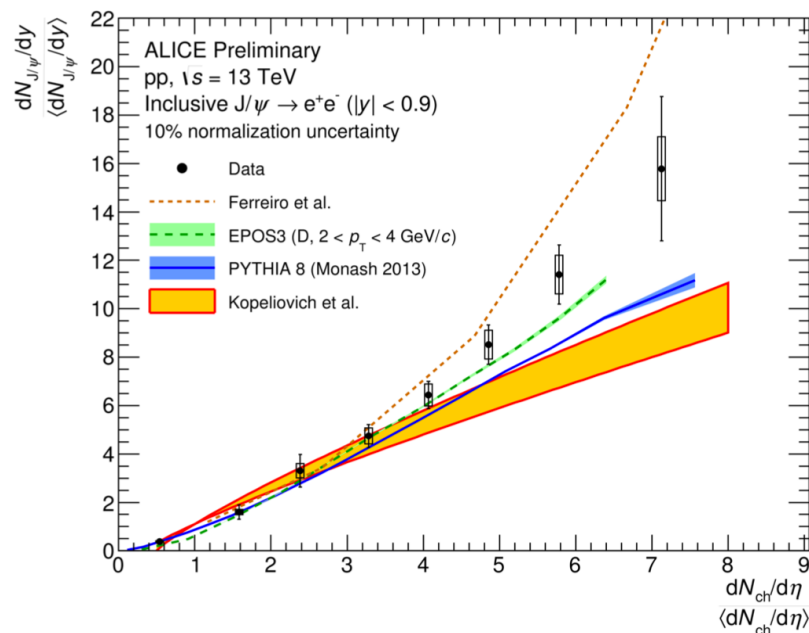
## Theoretical models

- String percolation
- Hydro dynamical evolution (EPOS3)
- Multiple parton interaction (PYTHIA8)
- Contributions of higher Fock states



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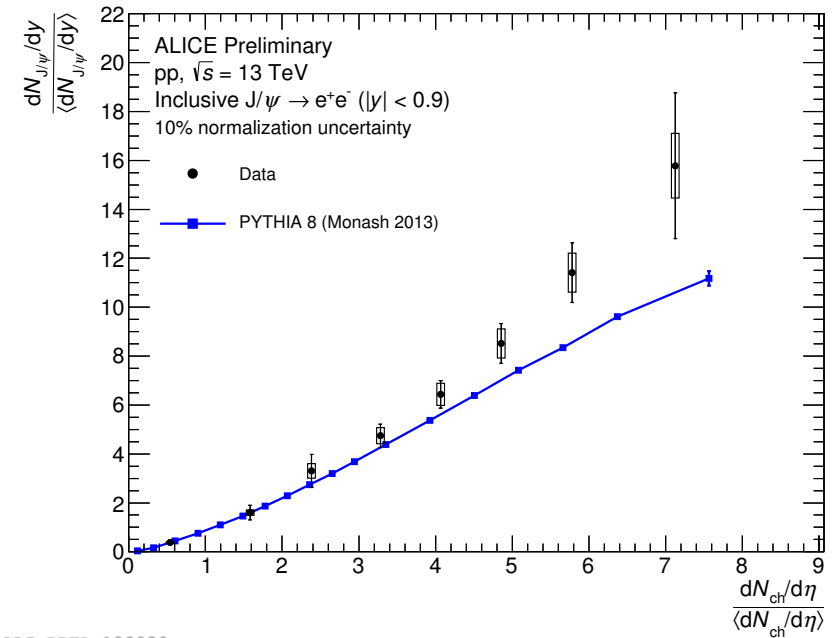


- ❖ The high  $p_T$  analysis is based on EMCAL triggered data
- ❖ The increase of J/ψ production as a function of multiplicity seems steeper at higher transverse momenta
- ❖ The  $p_T$  dependence behavior is explained at least qualitatively by PYTHIA8, which includes MPI processes.

# What does PYTHIA8 tell us ?

❖ J/ψ production has contributions from dedicated processes in PYTHIA8:

- ✓ Initial  $c$  or  $b$  quarks originate via first hardest 2->2 partonic interactions
- ✓ Has finite production probability from the subsequent hard processes in MPI
- ✓ Heavy quarks from gluon splitting
- ✓ Gluons from initial/final state radiations
- ✓ Color reconnection (at the hadronization stage)



❖ The events with a small number of MPI contribute to the low multiplicity interval, while high multiplicity events are dominated by a large number of MPI

❖ Monash 2013 tuned PYTHIA8 describes well the data in the low multiplicity region

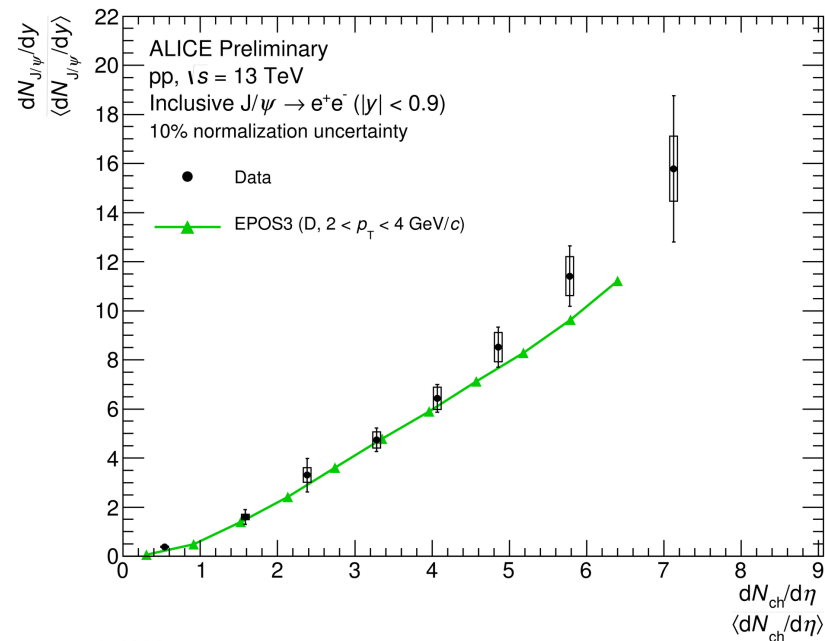


# What does EPOS3 tell us ?

❖ EPOS3 imposes the same theoretical scheme in pp, pA and AA systems

✓ Initial conditions followed by a hydrodynamical evolution

✓ Initial conditions based on “Gribov-Regge” formalism.  
Multiple interaction occurs in parallel



❖ The EPOS calculation is for D meson, which should be a very good proxy for J/ψ.

❖ The good description of the data with EPOS3 model shows that the energy density reached in pp collisions at the LHC is high enough to apply hydrodynamical evolution

❖ Result of EPOS version 3.1 and 3.2 differ significantly

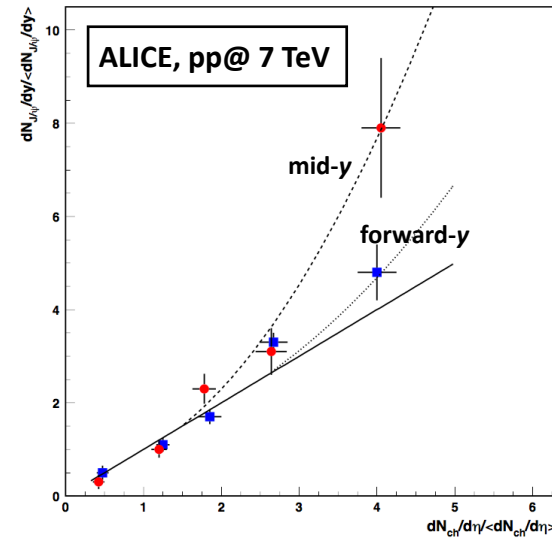
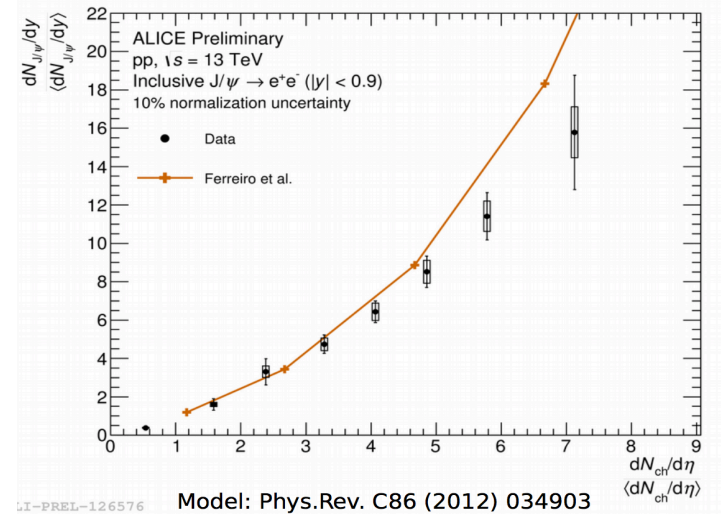


# What does Percolation tell us ?

- ❖ High-energy hadronic collisions are driven by the exchange of color sources (strings) between the projectile and the target
- ❖ The number of parton-parton collisions is reflected as the number of produced strings ( $N_s$ )

✓  $J/\psi$  multiplicity  $\propto N_s$

✓ Charged particle multiplicity  $\propto \sqrt{N_s}$



Ferreiro et al. Phys. Rev. C86 (2012) 034903



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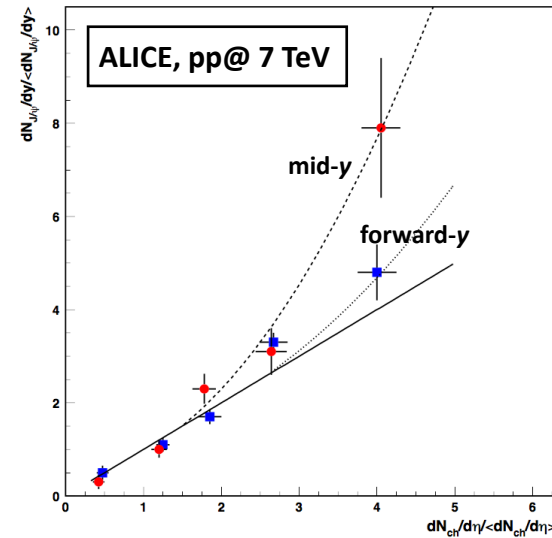
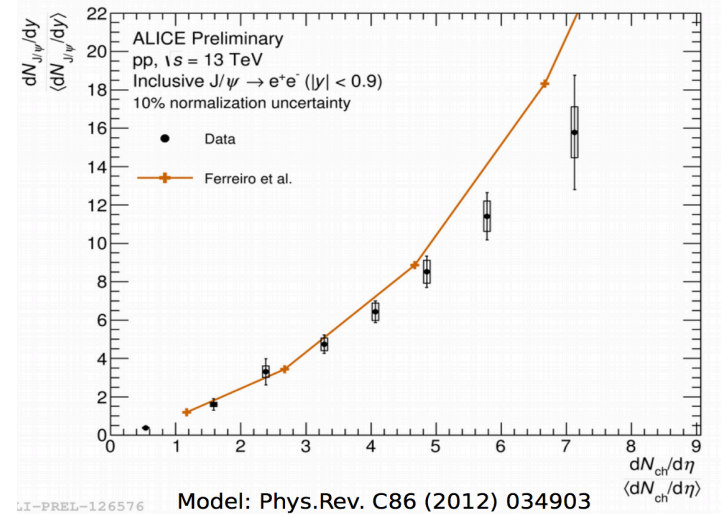
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At Low multiplicity

$$\frac{n_{J/\psi}}{\langle n_{J/\psi} \rangle} = \frac{\frac{dN}{d\eta}}{\langle \frac{dN}{d\eta} \rangle}$$



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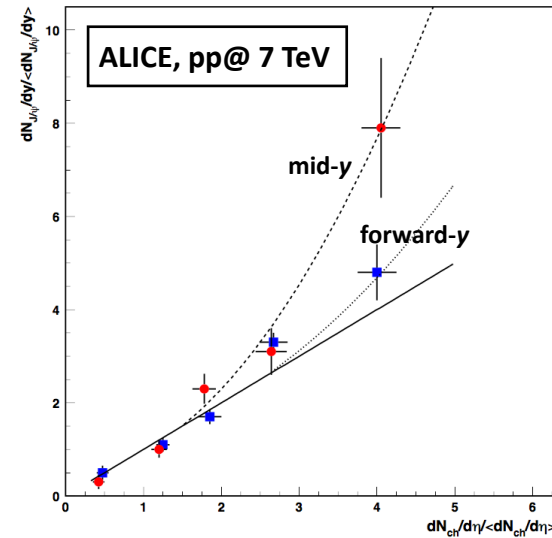
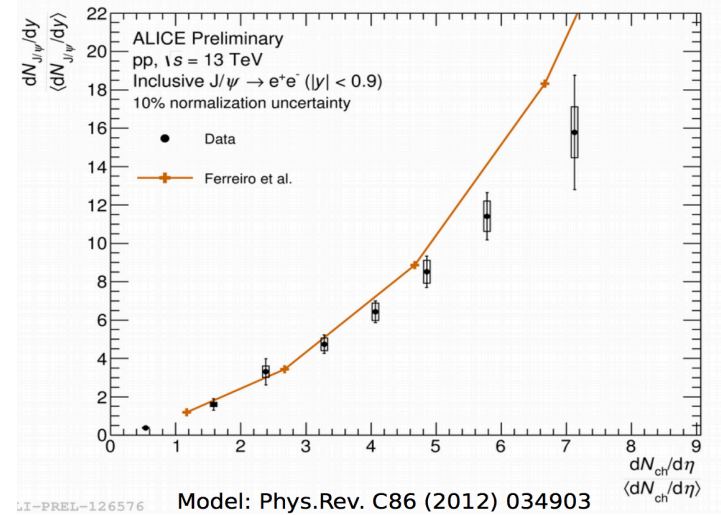
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At High multiplicity

$$\frac{n_{J/\psi}}{\langle n_{J/\psi} \rangle} = \langle \rho \rangle \left( \frac{\frac{dN}{d\eta}}{\langle \frac{dN}{d\eta} \rangle} \right)^2$$



Ferreiro et al. Phys. Rev. C86 (2012) 034903



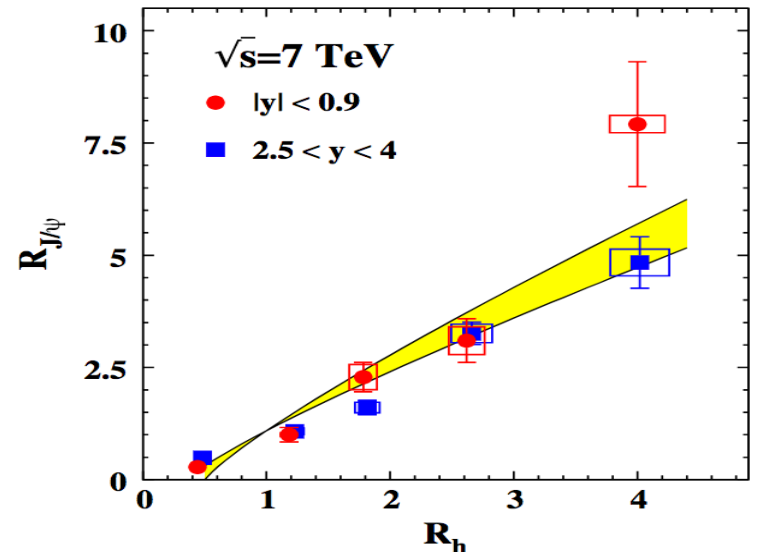
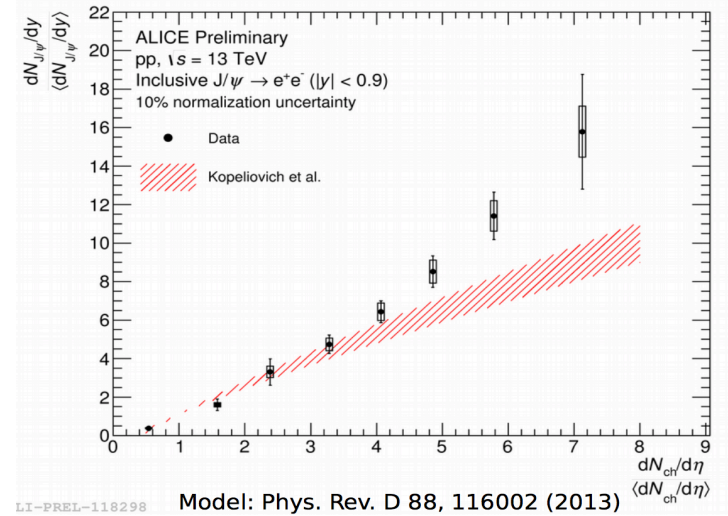
# What does “higher Fock states” model tell us ?

- ❖ **Higher Fock component:** In high energy nuclei, gluons at small- $x$  overlap longitudinally, act as a single source of gluons
- ❖ The inelastic collisions of the Fock components lead to high hadron multiplicity
- ❖ The relative production of  $J/\psi$  is enhanced in such gluon-rich collisions

$$R_h^{pp} \equiv \frac{dN_h^{pp}/dy}{\langle dN_h^{pp}/dy \rangle},$$

$$R_{J/\psi}^{pp} \equiv \frac{dN_{J/\psi}^{pp}/dy}{\langle dN_{J/\psi}^{pp}/dy \rangle}.$$

- ❖ More gluons participating in collisions with  $R_h^{pp} > 1$ , explains why  $R_{pp}^{J/\psi}$  rises with increasing  $R_h$





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## Thank you!!

# Multiplicity determination

- ❖ Charged-particle multiplicity is measured using the number of SPD tracklets in  $|\eta| < 1$ . The variation of the SPD efficiency with the  $z$  position of the primary vertex ( $z_{\text{vertex}}$ ) is corrected using a data-driven method.

$$\Delta N = \frac{\langle N_{\text{trk}} \rangle(z_v^0) - \langle N_{\text{trk}} \rangle(z_v)}{\langle N_{\text{trk}} \rangle(z_v)}$$

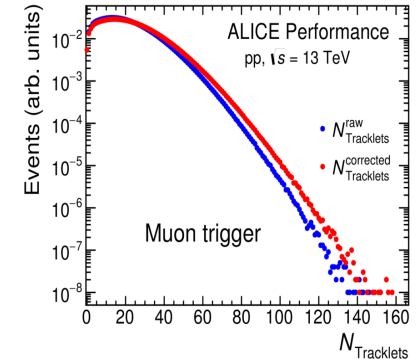
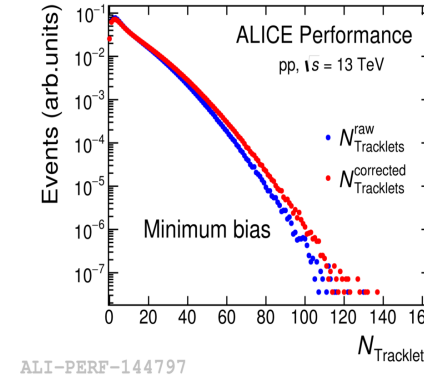
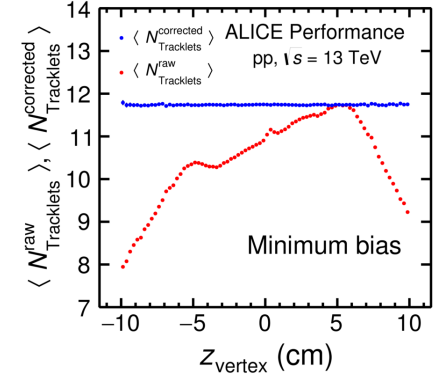
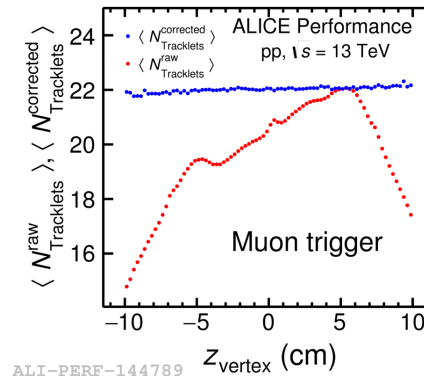
$$N_{\text{trk}}^{\text{corr}}(z_v) = N_{\text{trk}}(z_v) + \Delta N_{\text{rand}}$$

- ❖ Here,  $\Delta N_{\text{rand}}$  follows a Poissonian distribution centered around  $\Delta N$ .
- ❖  $z_v^0$  corresponds to  $z_{\text{vertex}}$  position where  $\langle N_{\text{trk}} \rangle$  is maximum.

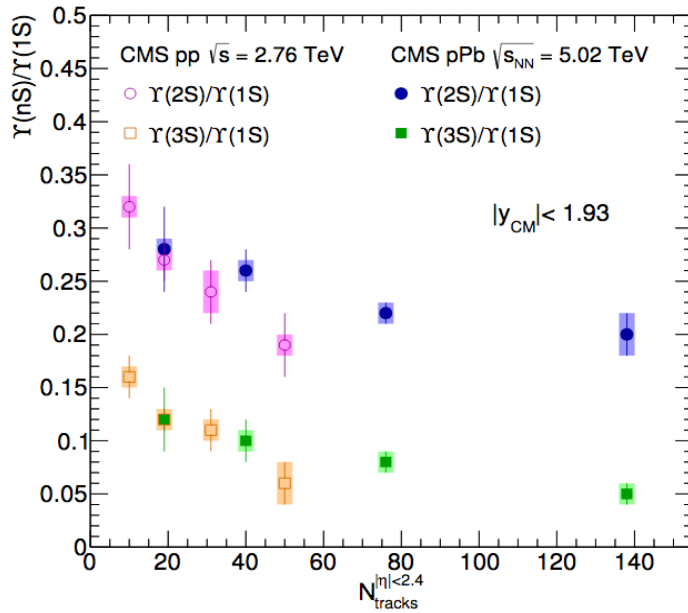
- ❖ The efficiency loss at  $z_0$  and other track-to-particle-corrections need to be taken into account to evaluate the actual charged-particle value.

$$\frac{\langle dN_{\text{ch}}/d\eta \rangle_i}{\langle dN_{\text{ch}}/d\eta \rangle} = \frac{f(\langle N_{\text{trk}}^{\text{corr}} \rangle_i)}{\langle dN_{\text{ch}}/d\eta \rangle_{\text{INEL}} > 0}$$

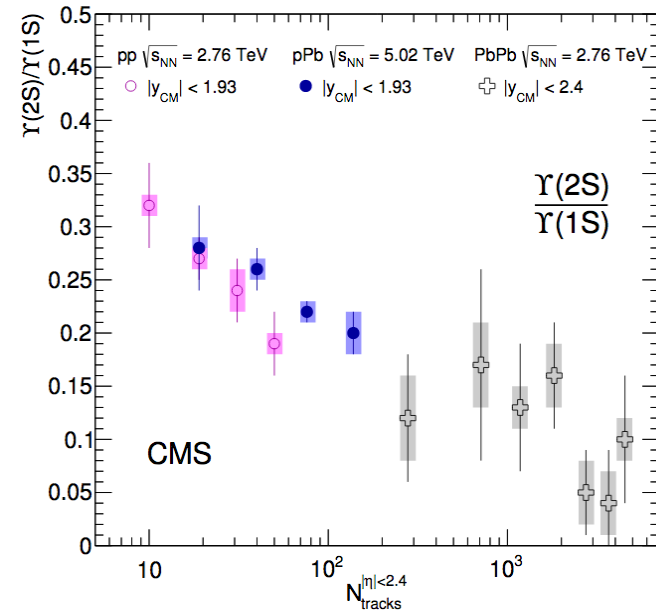
- ❖ The conversion function “f” is estimated using Monte Carlo simulations.



# Y Vs. multiplicity



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- ❖ The excited-to-ground-states ratios,  $Y(nS)/Y(1S)$ , are found to decrease with increasing charged-particle multiplicity.
- ❖ The tightest bound state,  $Y(1S)$ , was observed to be less suppressed than the more loosely bound excited states,  $Y(2S)$  and  $Y(3S)$ .
- ❖ Global behavior of double ratio has been observed irrespective of the collisions system; pp, p-Pb, and Pb-Pb.



- **Silicon Pixel Detector (SPD) is used for charged particle and vertex determination**

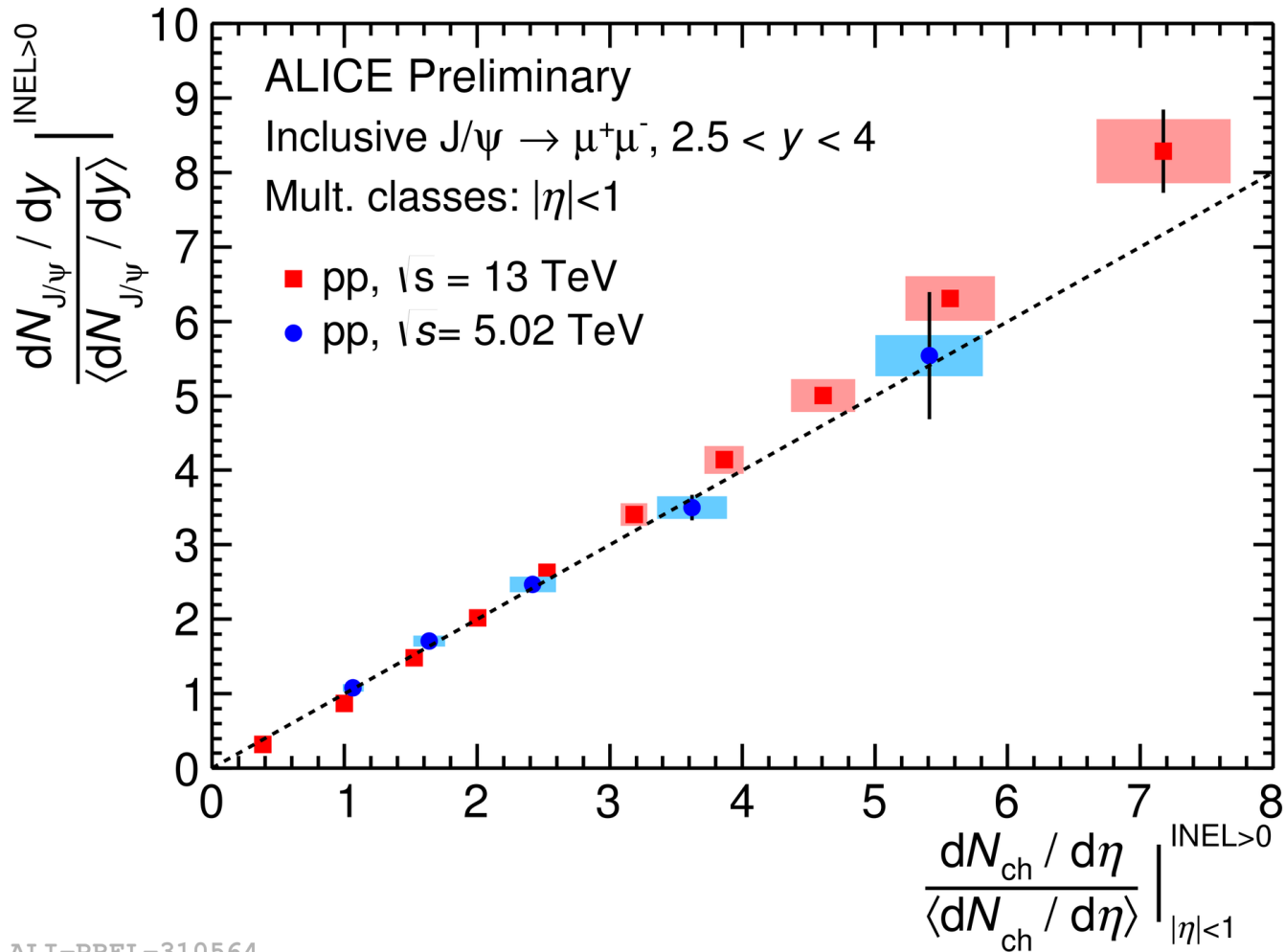
$$J/\psi \rightarrow e^+e^-$$

- **MB trigger**
- **$-0.9 < \eta < 0.9$**
- **Track quality cuts**
- **Rejection of tracks from photon conversion**
- **TPC electron identification**

$$J/\psi \rightarrow \mu^+\mu^-$$

- **Dimuon trigger: MB and two opposite sign muon tracks**
- **$-4.0 < \eta < -2.5$**
- **$17.6 \text{ cm} < R_{\text{abs}} < 89.5 \text{ cm}$**   
( $R_{\text{abs}}$  = Radial position of the track at the end of the absorber)

# Energy dependence of J/ψ yield vs. multiplicity

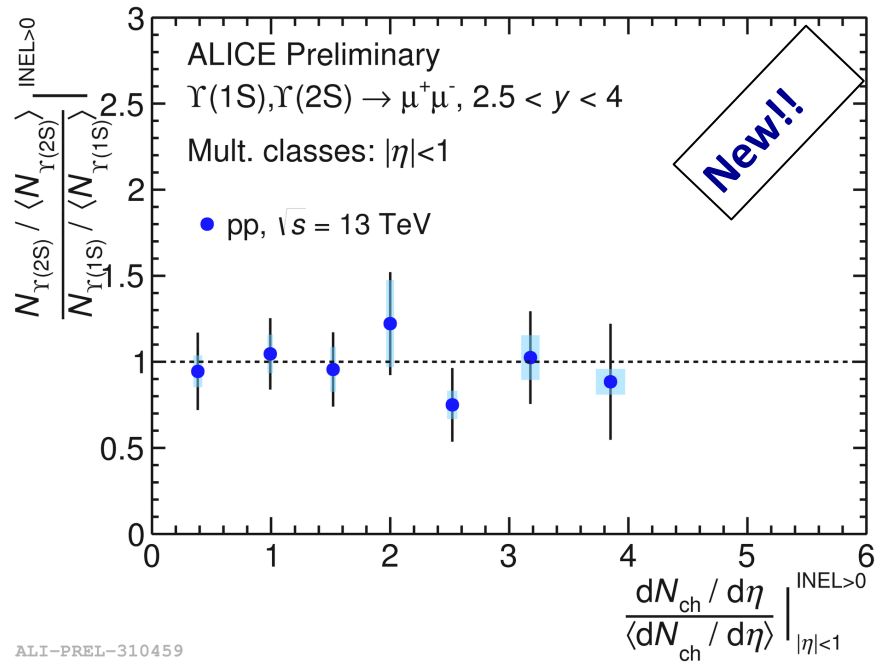
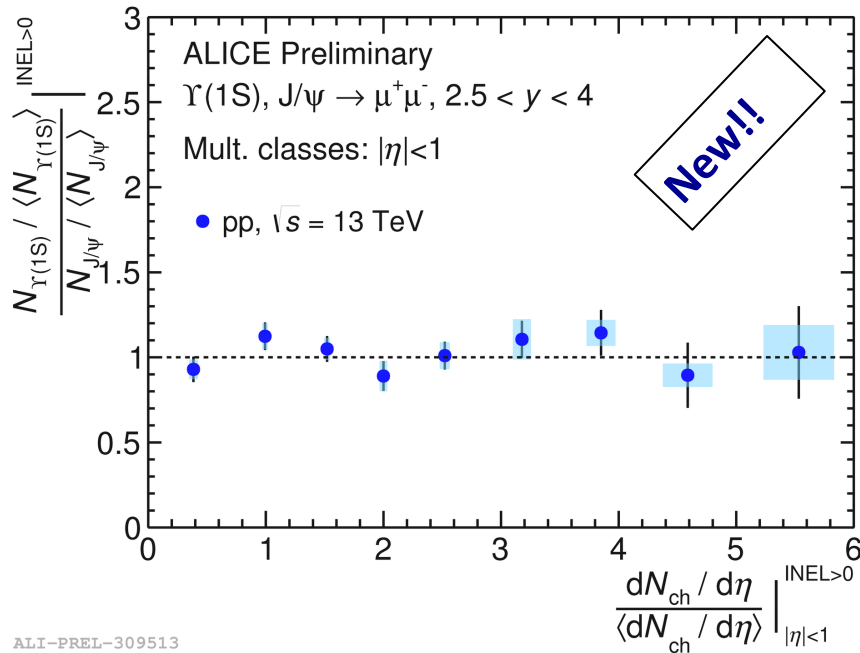


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❖ Trend is independent of colliding energy.



# Double ratio

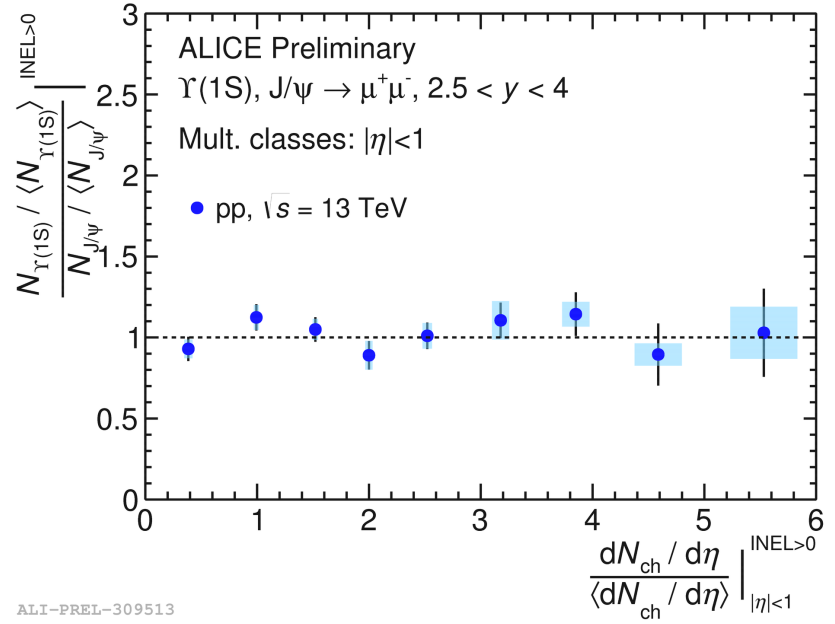
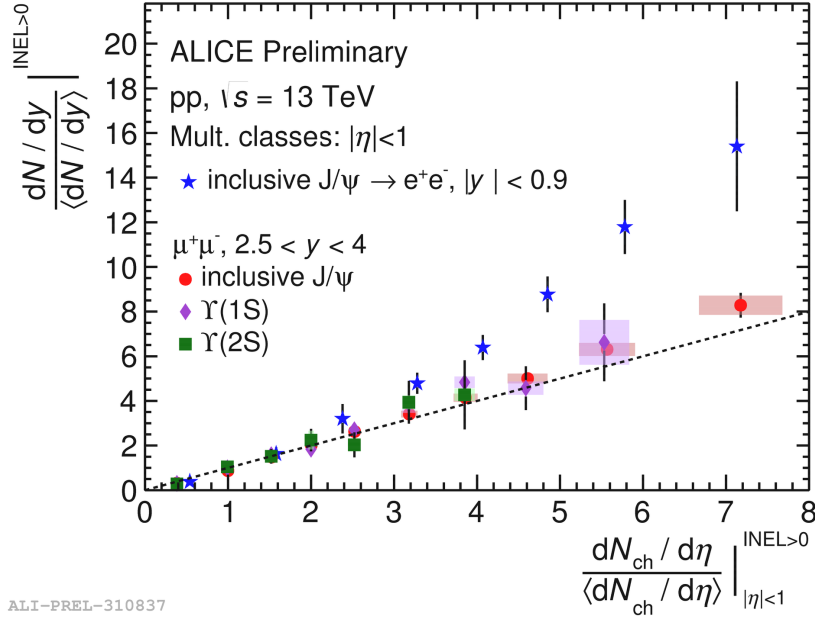


❖ The double ratio  $\Upsilon(1S)/J/\psi$  and  $\Upsilon(2S)/\Upsilon(1S)$  has been studied for pp@ 13 TeV

- The ratio is found to be unity irrespective of charged-particle multiplicity.
- The multiplicity dependence production is same for  $J/\psi$ ,  $\Upsilon(1S)$  and  $\Upsilon(2S)$ .
- $\Upsilon(nS)/\Upsilon(1S)$ , are found to decrease with increasing charged-particle multiplicity, when  $\Upsilon$  and charged particles are measured in mid-rapidity. (CMS Collaboration, JHEP04(2014)103)



# Quarkonium vs. multiplicity



- The multiplicity dependence production is same for  $J/\psi$ ,  $\Upsilon(1S)$  and  $\Upsilon(2S)$ .
  - The event activity associated with the production of heavy-flavor is also independent of quark content of the particle