Identified Particle Multiplicity fluctuations as a measure of deconfinement and hadronization in the QCD phase diagram

R. Bellwied (University of Houston)

Thanks: *Fernando Flor*, Anders Knospe, *Nalinda Kulathunga*, *Jake Martinez, Corey Myers, Paolo Parotto,* Claudia Ratti, Jihye Song, *Jamie Stafford*, Cristina Terrevoli, *Ejiro Umaka*

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The QCD Phase Transition – where do we stand ?

Discovery of the deconfined phase (SPS/RHIC results) by 2005

Signatures: jet quenching (partonic energy loss), quark scaling of large anisotropic flow (hydrodynamics, viscosity limit), photon temperature, J/ψ melting, strangeness enhancement,……..

Characterization measurements (GSI/SPS/RHIC/LHC results) still ongoing.

Two avenues

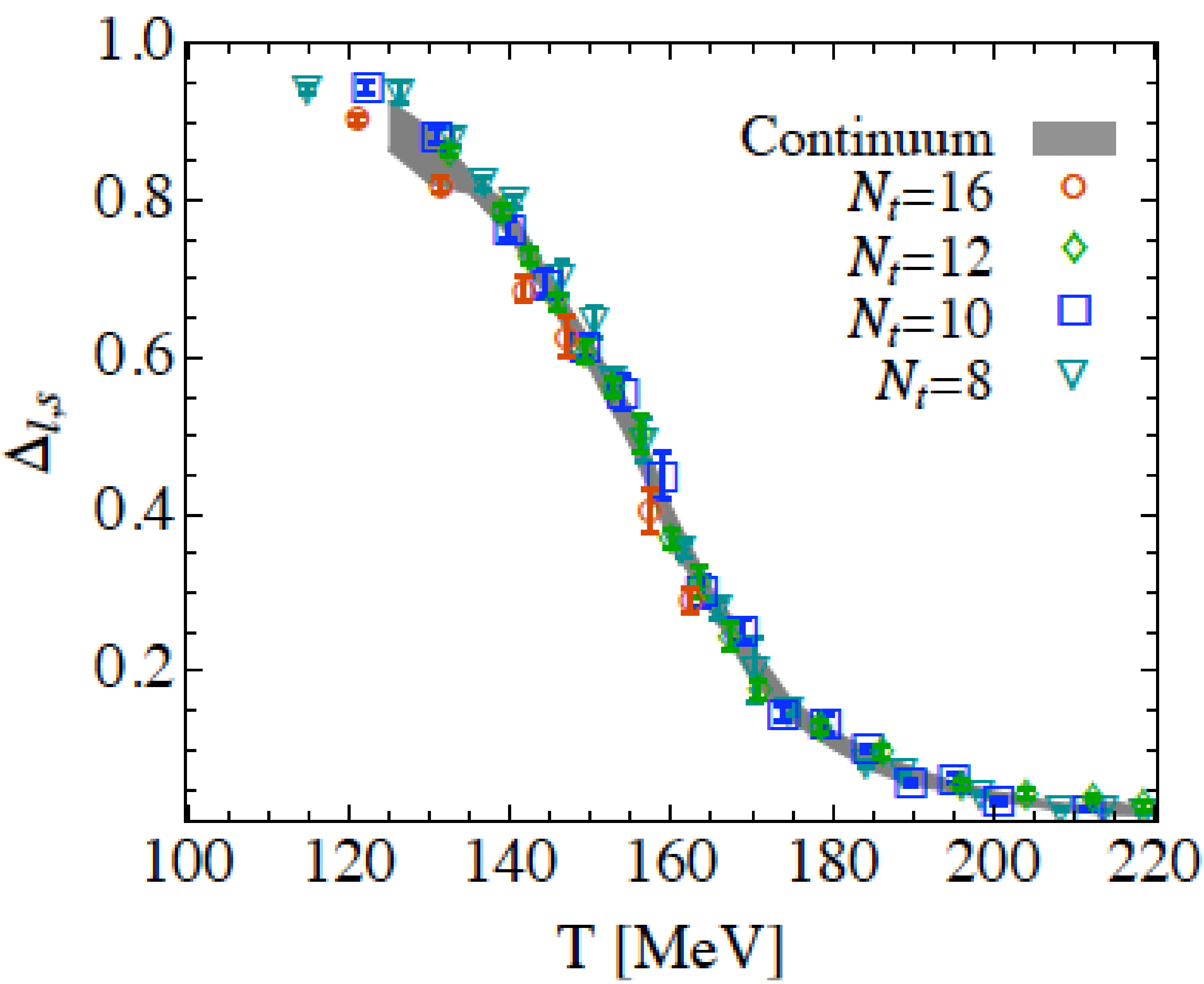
Characterization of the phase:
(phase identified)
- -transport coefficients
- -viscosity, conductivity, etc.
- -vorticity, chiral magnetic effects
- -……

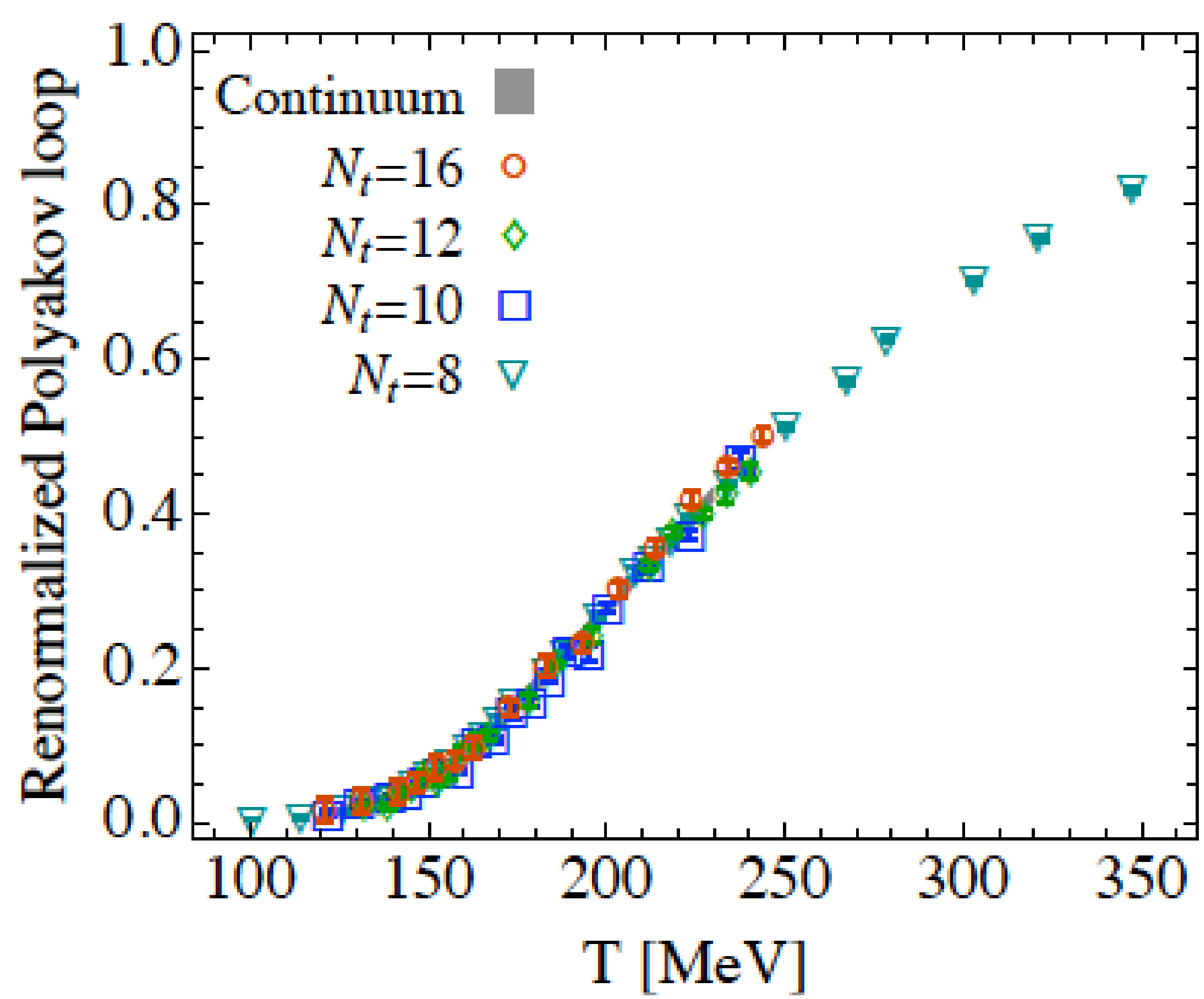
Characterization of the transition:
(particle identified)
- -hadronization
- -chiral symmetry
- -confinement
- -degrees of freedom
- -critical point
- -……

The order parameters of the QCD phase diagram

Chiral restoration vs. Deconfinement, ‘good’ or ‘bad’ ?

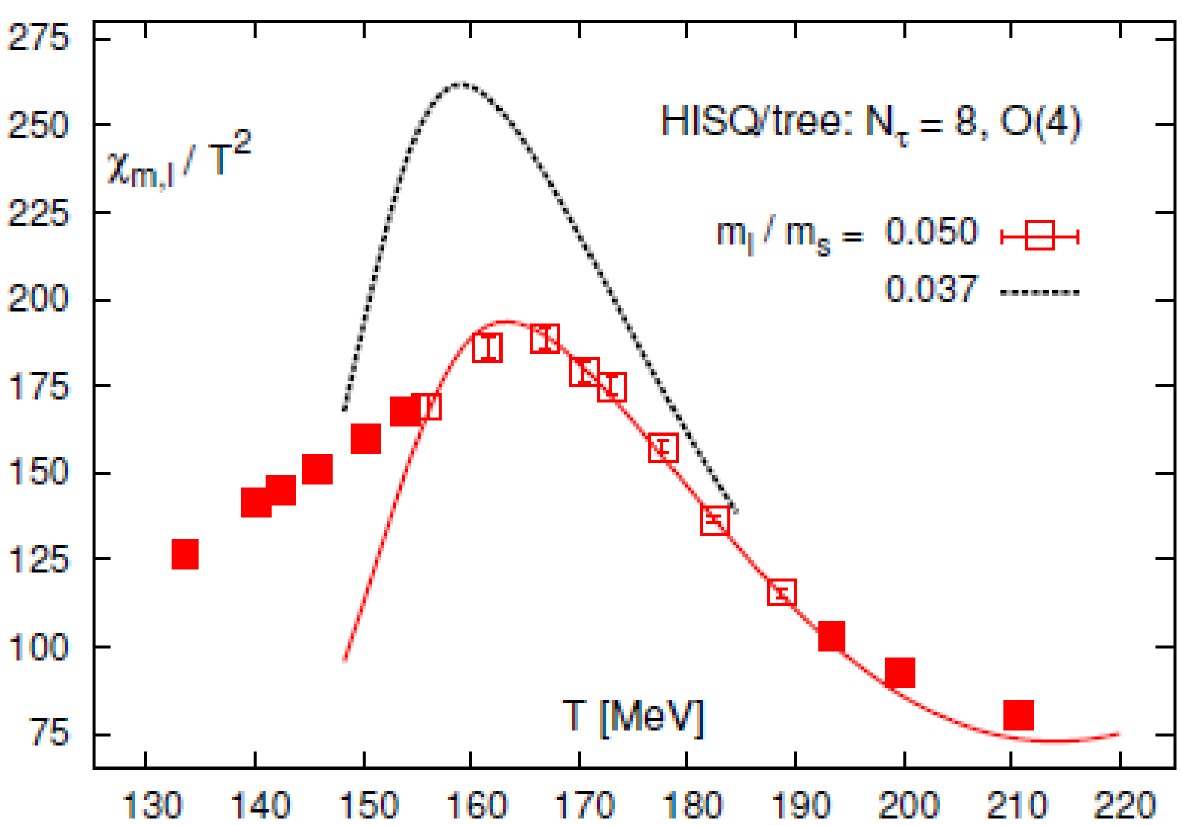
All of them show analytic crossover

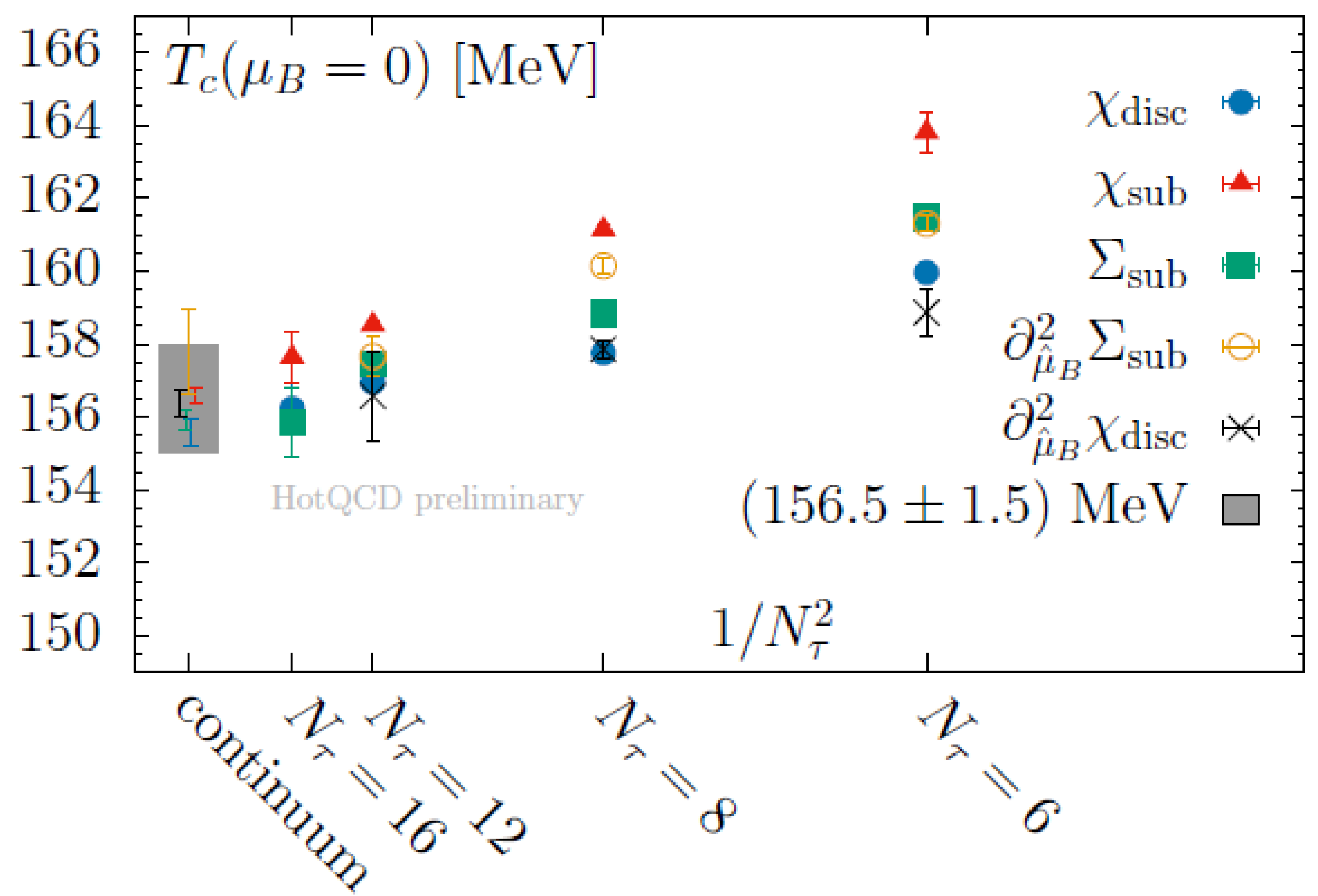




The order parameters of the QCD phase diagram

The chiral transition defines the pseudo-critical T

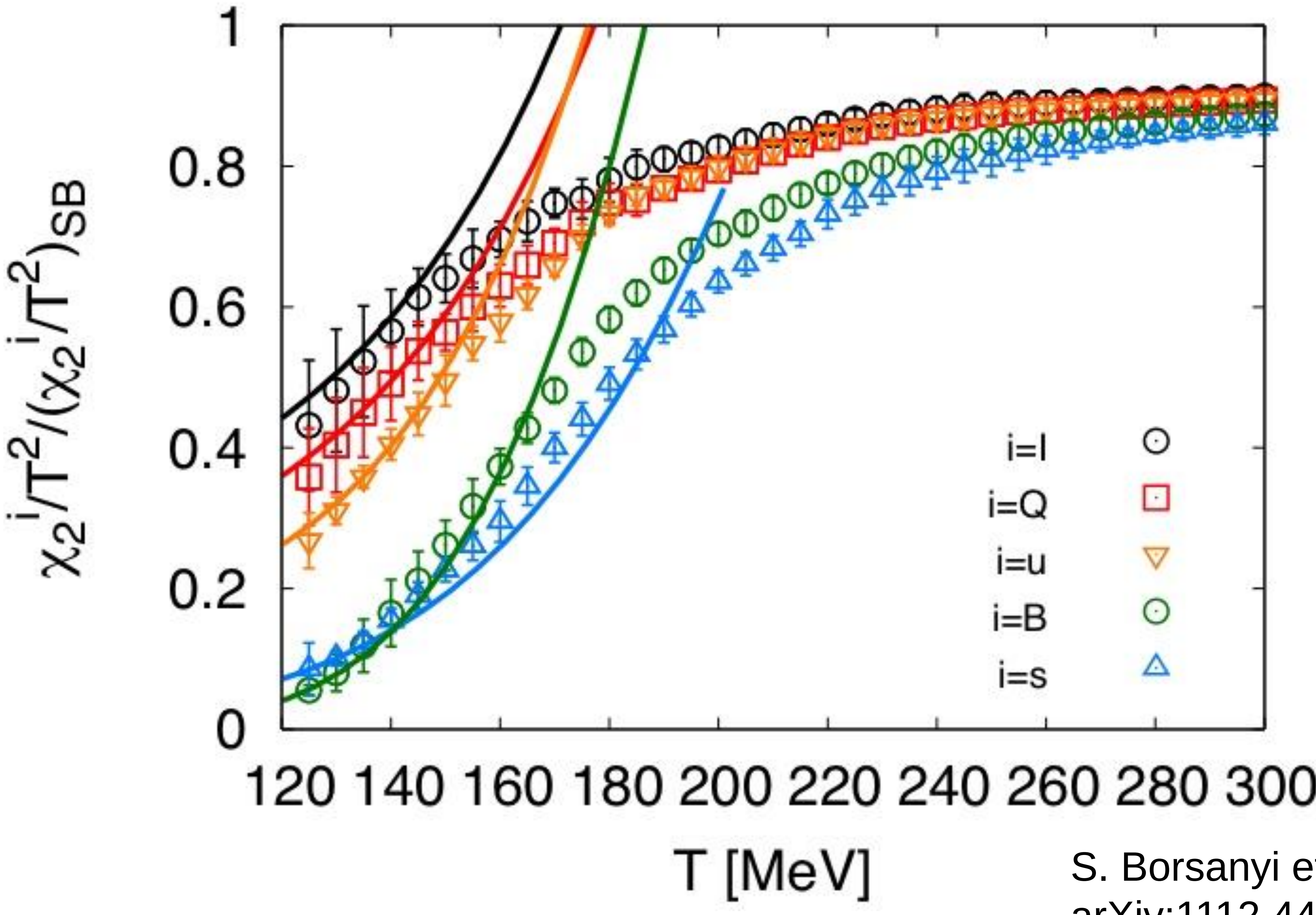




HotQCD (1504.05274)
T_{pc}: 154+- 9 MeV

HotQCD (1807.05607)
T_{pc}: 156.5+- 1.5 MeV

The relevance of conserved charges as order parameters for the phase transition – Understanding hadronization microscopically

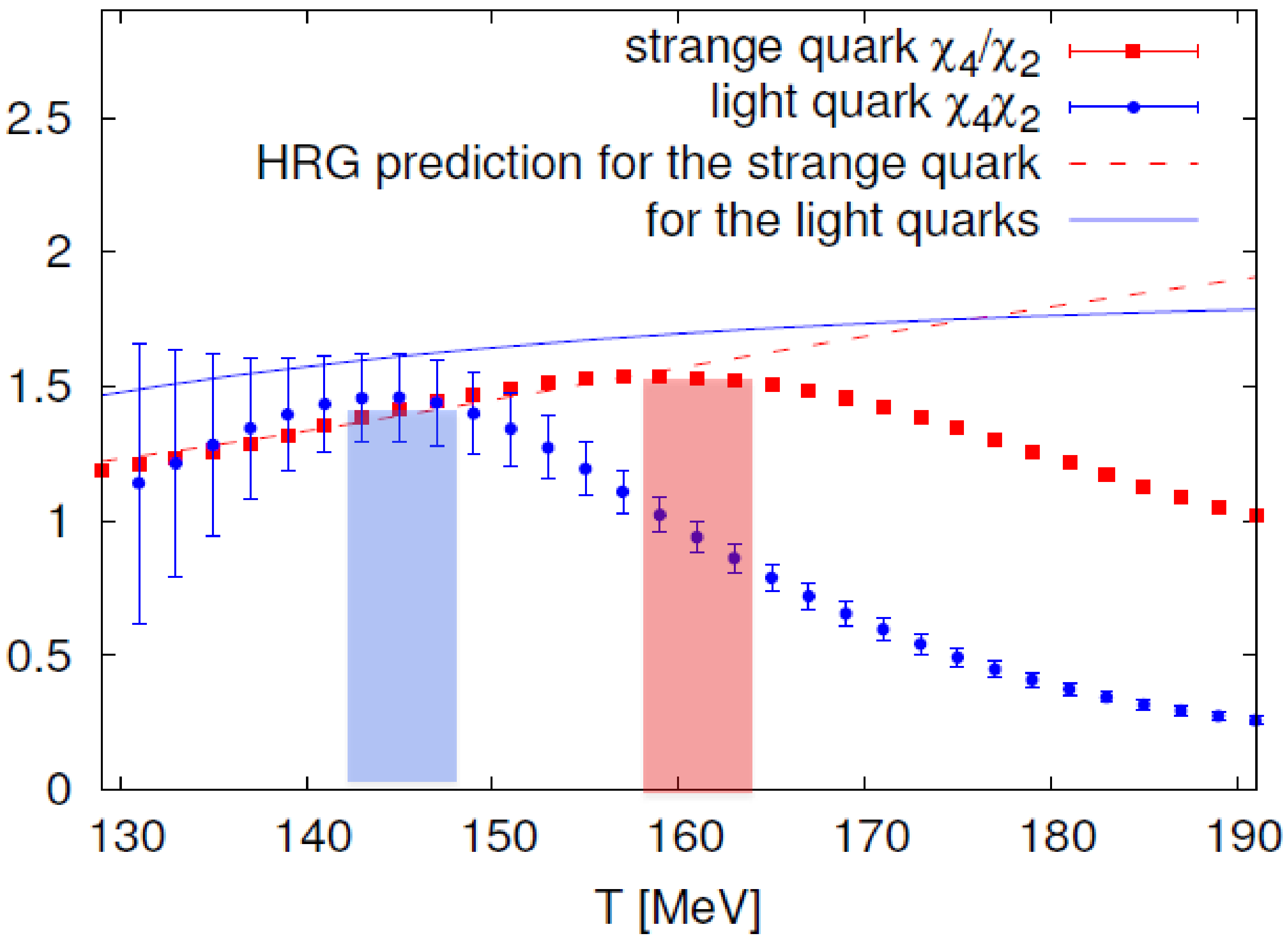


S. Borsanyi et al., arXiv:1112.4416

Direct determination of freeze-out parameters from first principles (lattice QCD)

$$\kappa_B \sigma_B^2 \equiv \frac{\chi_{4,\mu}^B}{\chi_{2,\mu}^B} = \frac{\chi_4^B(T)}{\chi_2^B(T)} \left[\frac{1 + \frac{1}{2}\frac{\chi_6^B(T)}{\chi_4^B(T)}(\mu_B/T)^2 + \dots}{1 + \frac{1}{2}\frac{\chi_4^B(T)}{\chi_2^B(T)}(\mu_B/T)^2 + \dots} \right]$$

Susceptibility ratios are a model independent measure of the chemical freeze-out temperature near μ=0. (Karsch, arXiv:1202.4173)



R. Bellwied & WB Collab., PRL (2013), arXiv:1305.6297

<u>Indication of sequential hadronization</u>

$$\chi_{lmn}^{BSQ} = \frac{\partial^{l+m+n}(p/T^4)}{(\partial\mu_B/T)^l(\partial\mu_S/T)^m(\partial\mu_Q/T)^n}.$$

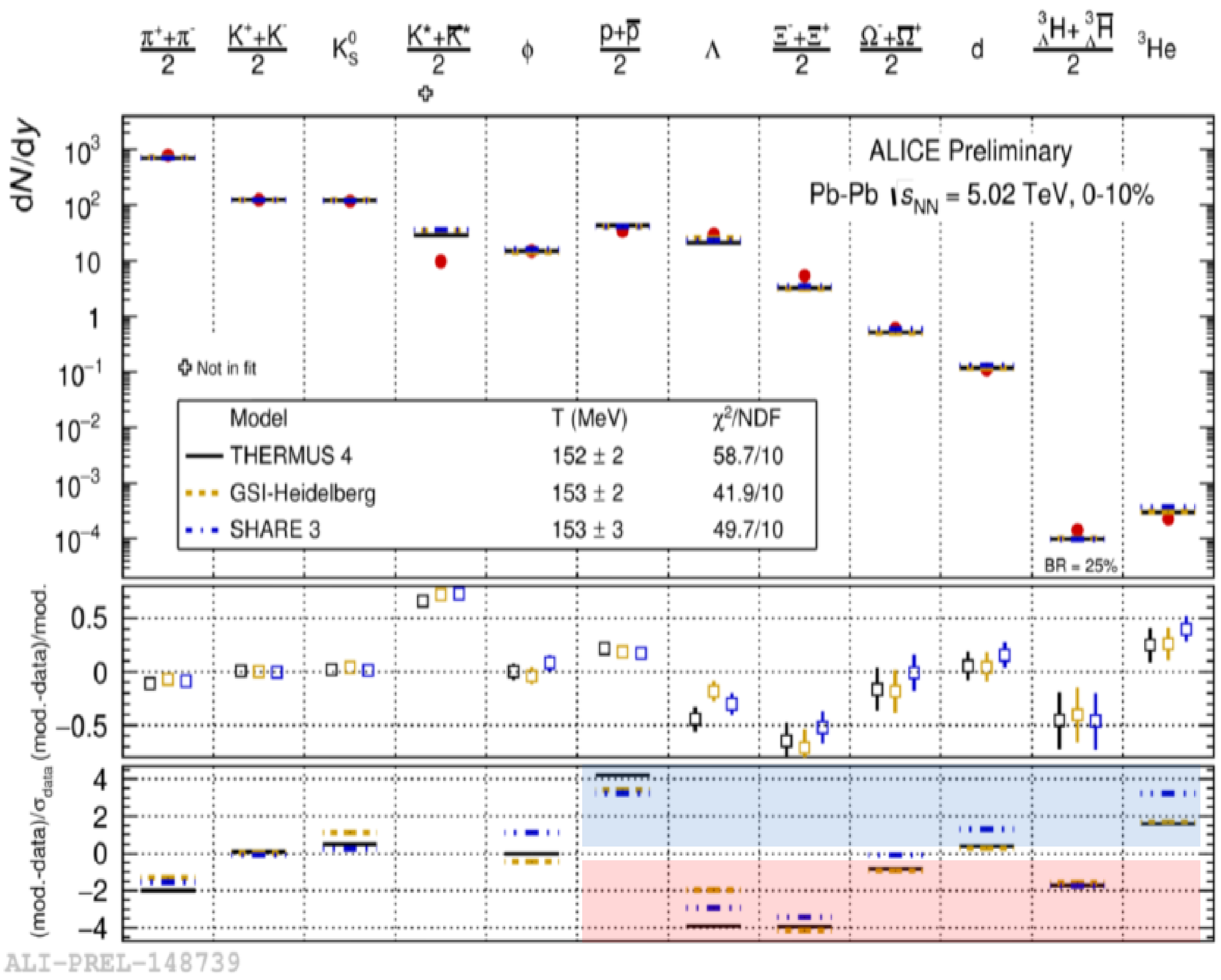
<u>Needs experimental verification</u>

<u>Fluctuations of e-by-e multiplicity distributions (mean (M), variance (σ), skewness (S), kurtosis (κ)</u>

$$M/\sigma^2 = \chi_1/\chi_2 \qquad S\sigma = \chi_3/\chi_2$$

$$S\sigma^3/M = \chi_3/\chi_1 \qquad \kappa\sigma^2 = \chi_4/\chi_2.$$

Experimental evidence: HRG model comparisons to ALICE yields

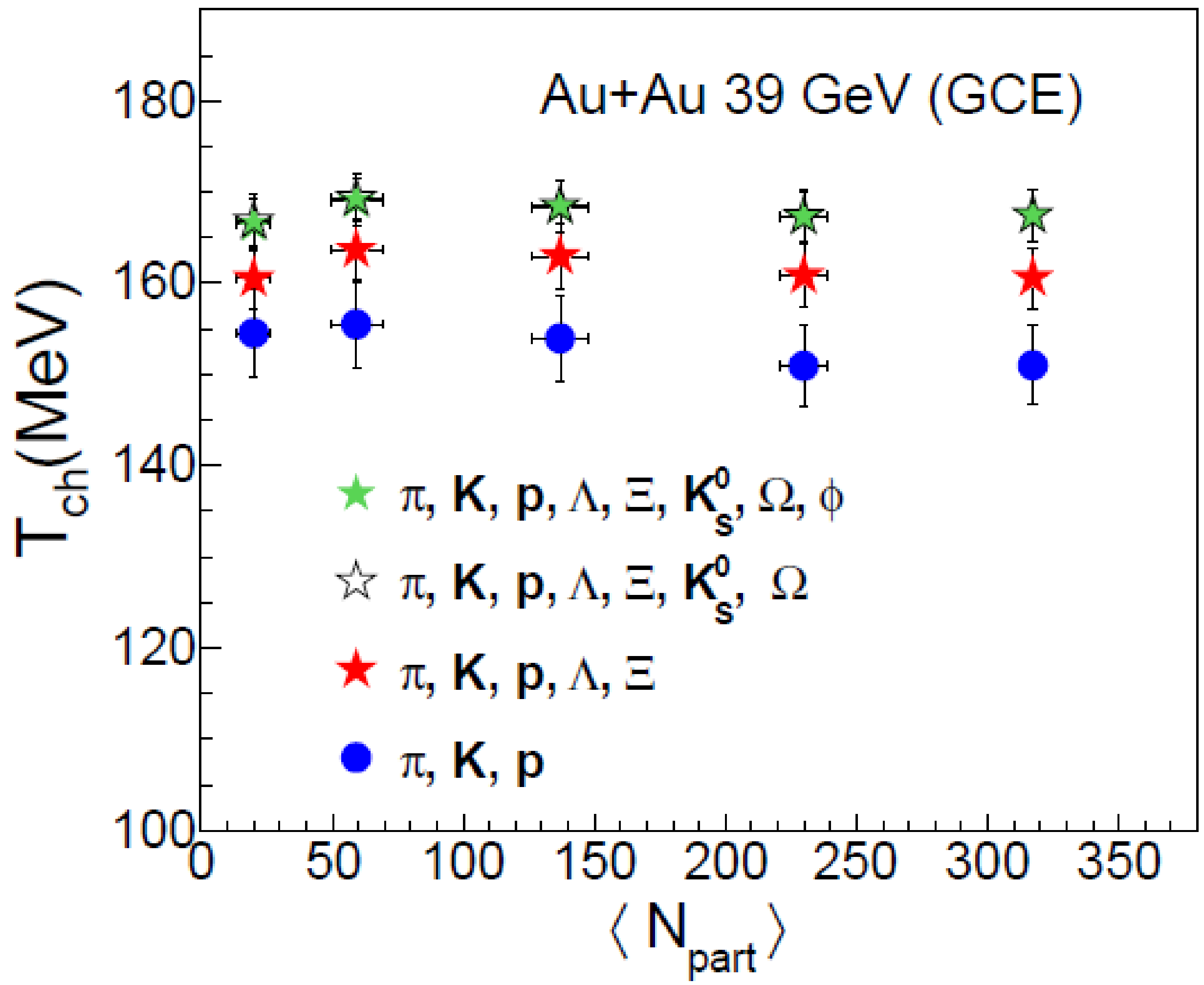


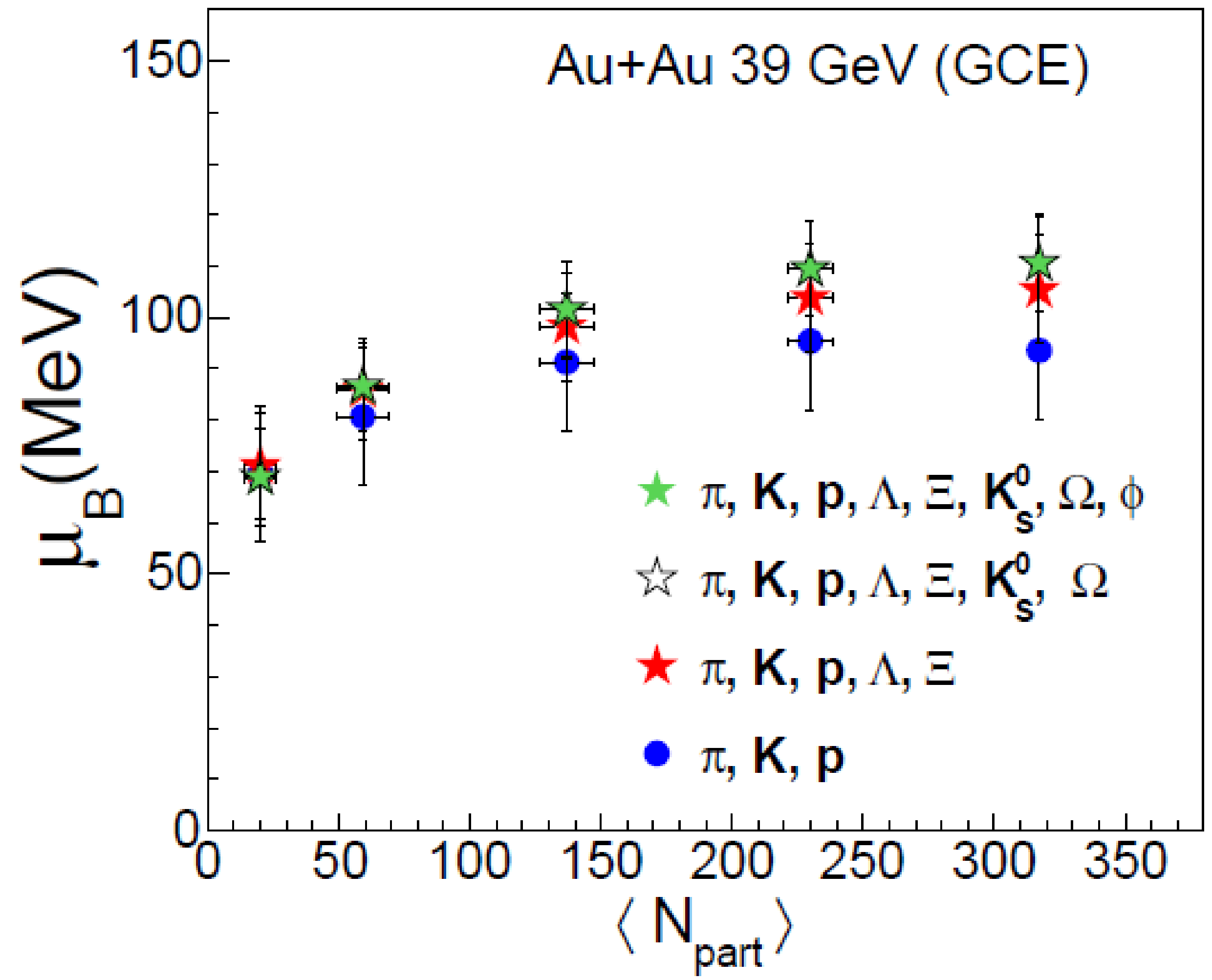
The new ALICE 5.02 TeV Pb-Pb data show a more pronounced and more precise tension between strange and non-strange particles in the baryonic sector (+4σ effect in protons vs. -4σ effect in Ξ baryons)

Overall there seems to be a light vs. strange particle trend

Experimental evidence: from varying the input particles into the chemical fit

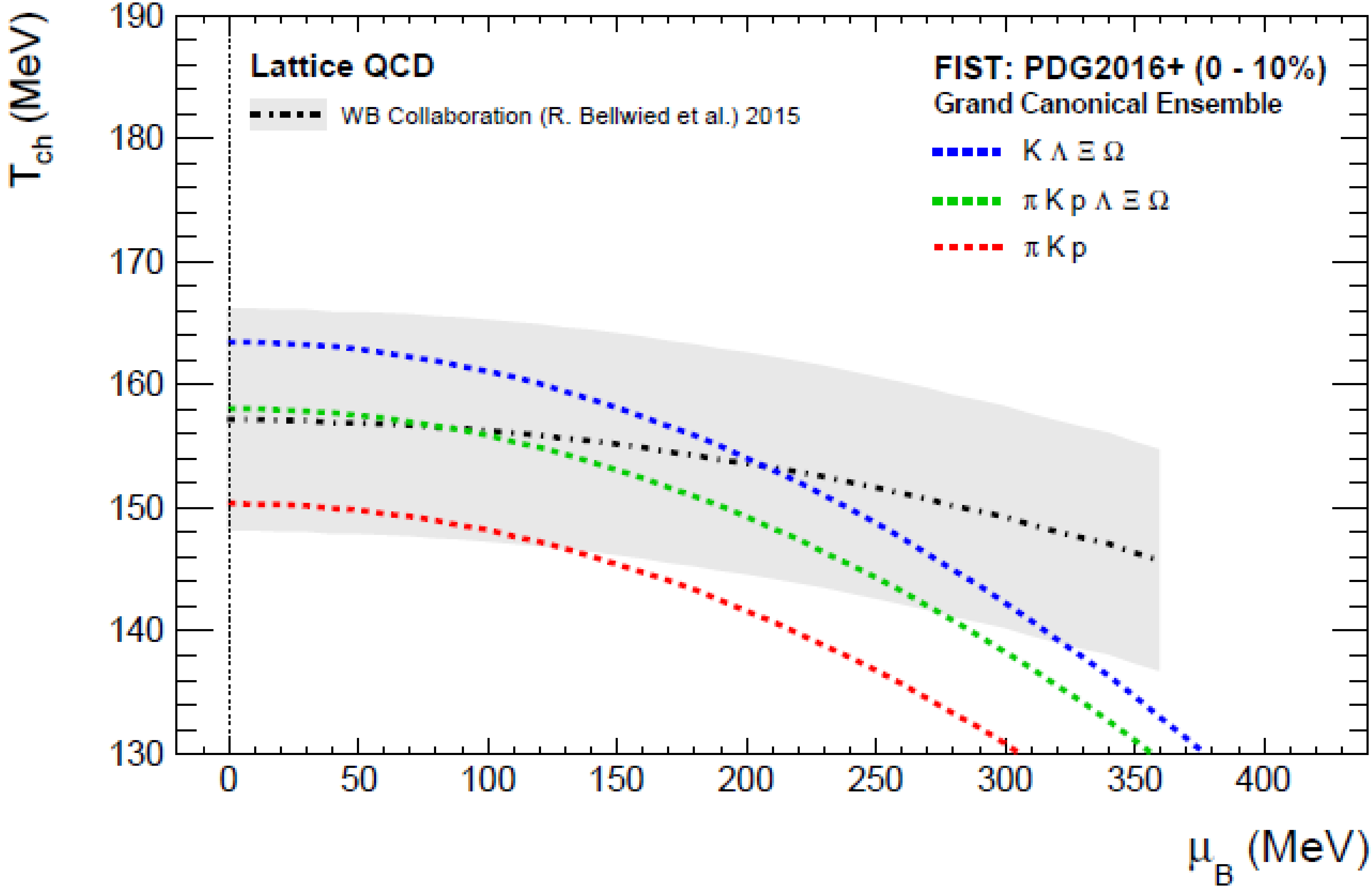
Latest example: Beam Energy Scan data from STAR (arXiv:1701.07065)





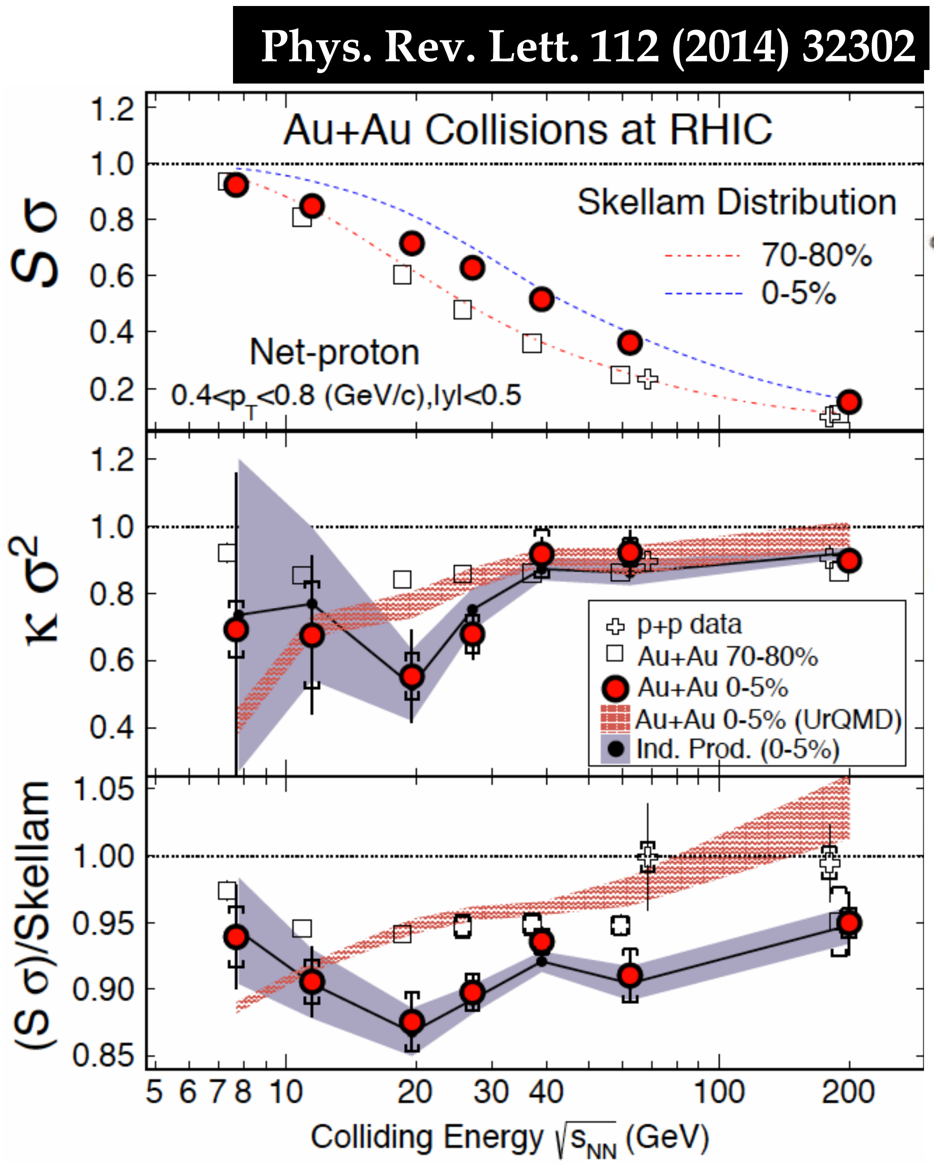
This is a long known fact in SHM, always argued as 'the more states the better', but all additional states (to π,k,p) are strange states

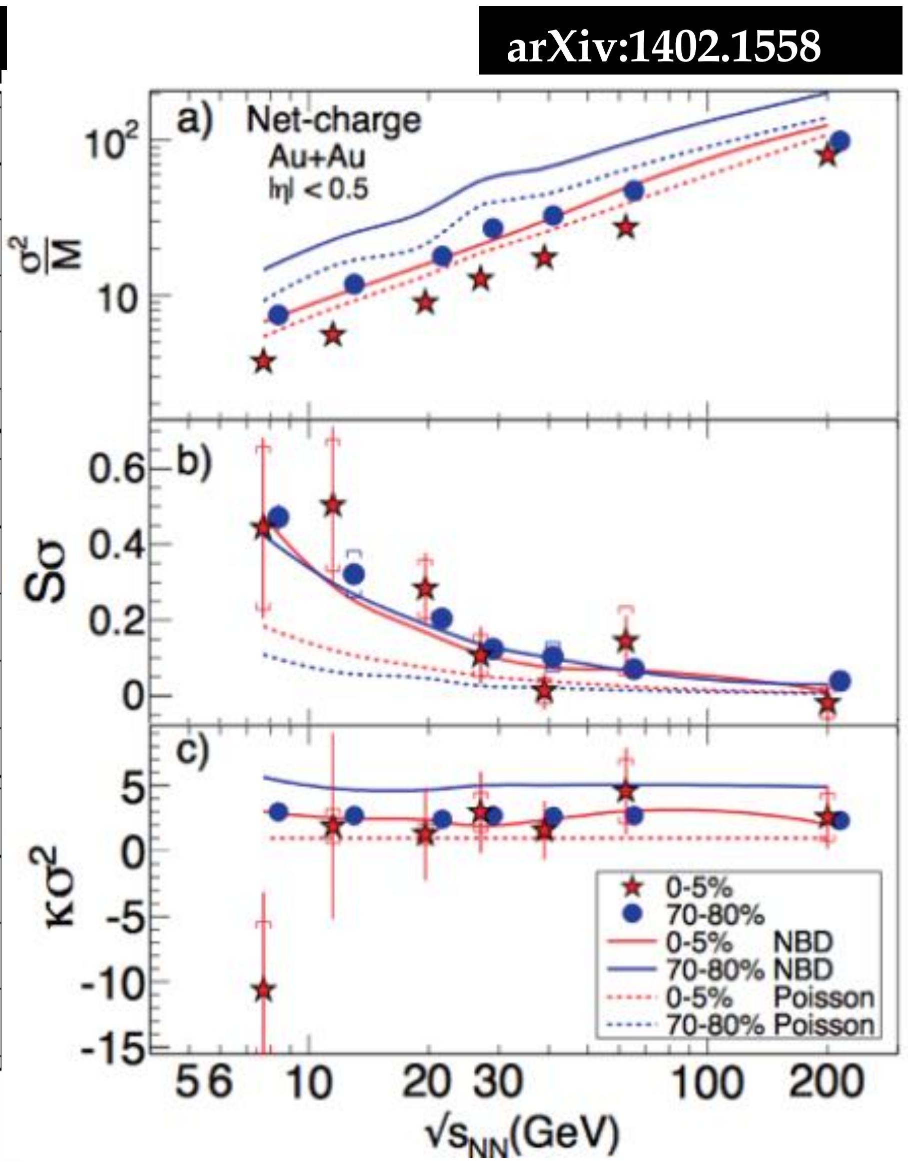
Latest study using FIST fits to ALICE and STAR data (F.Flor, G. Olinger (UH))

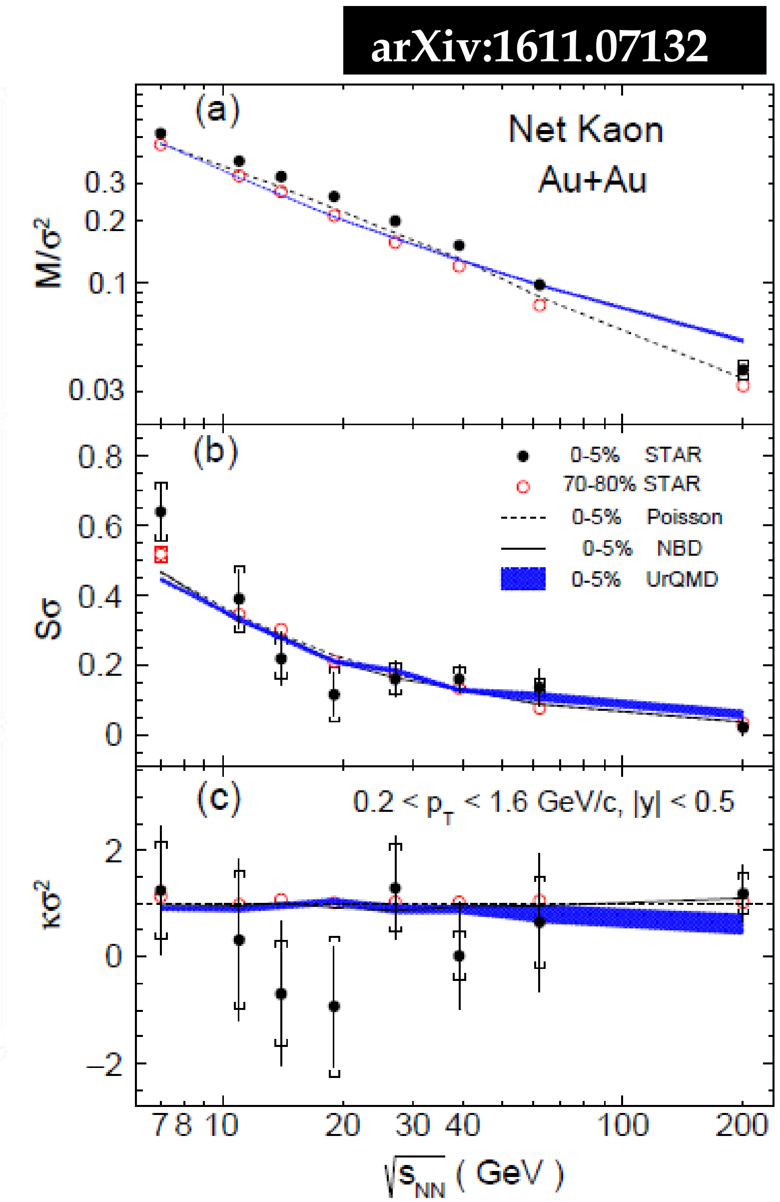


Data: ALICE PbPb 2.76 TeV, STAR BES

Higher moment ratios of e-by-e multiplicity distributions for net-proton, net-charge & net-Kaon distributions from STAR

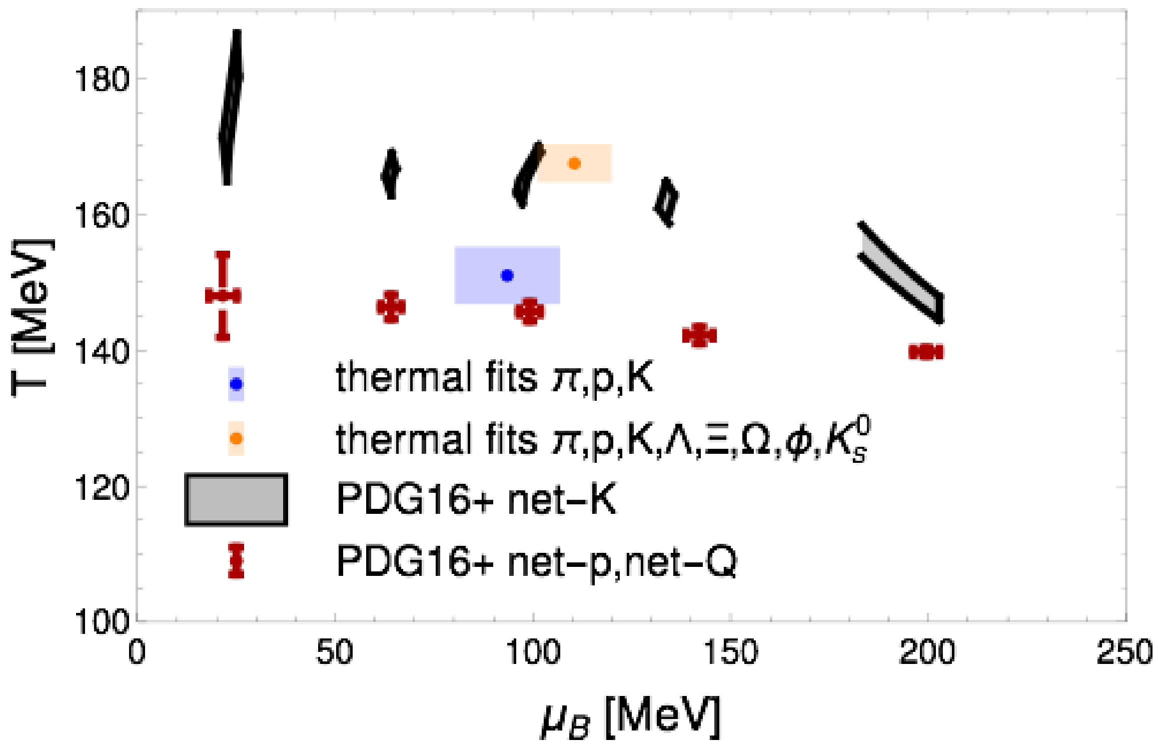


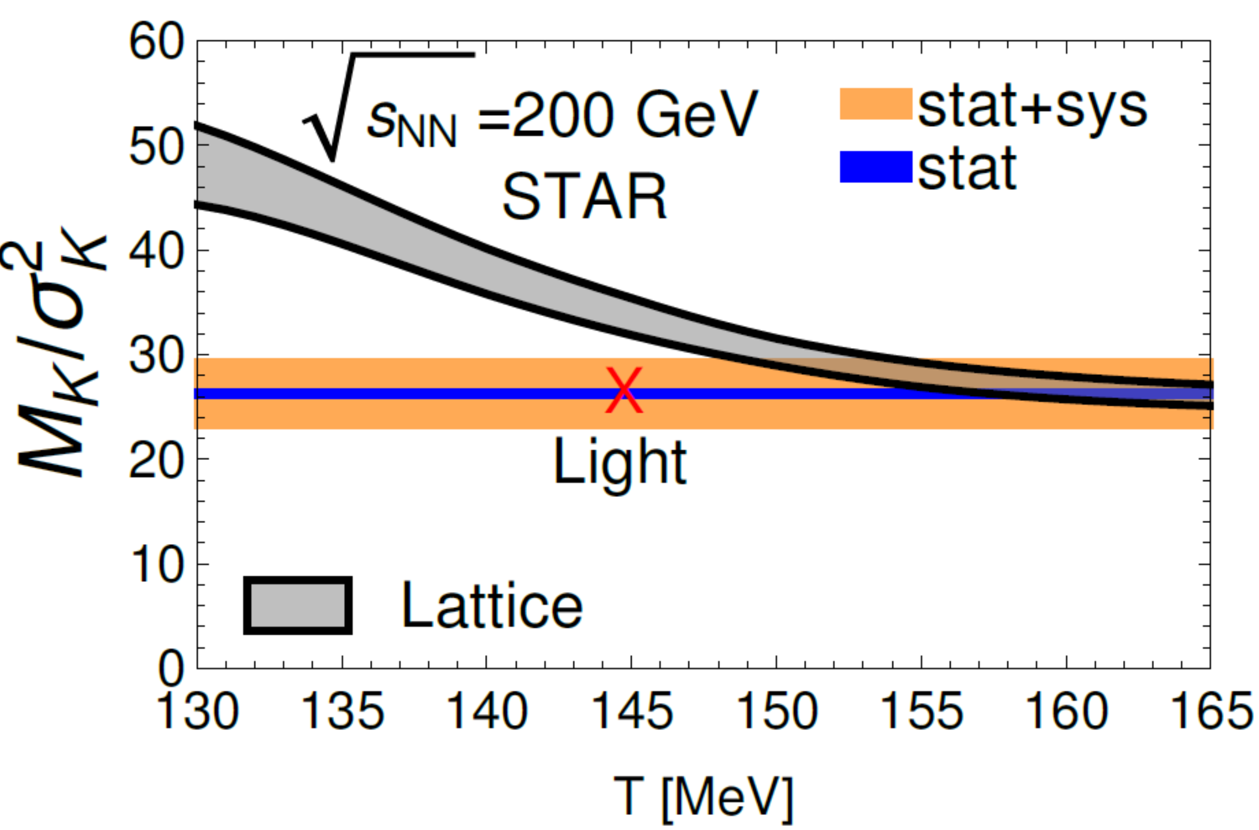


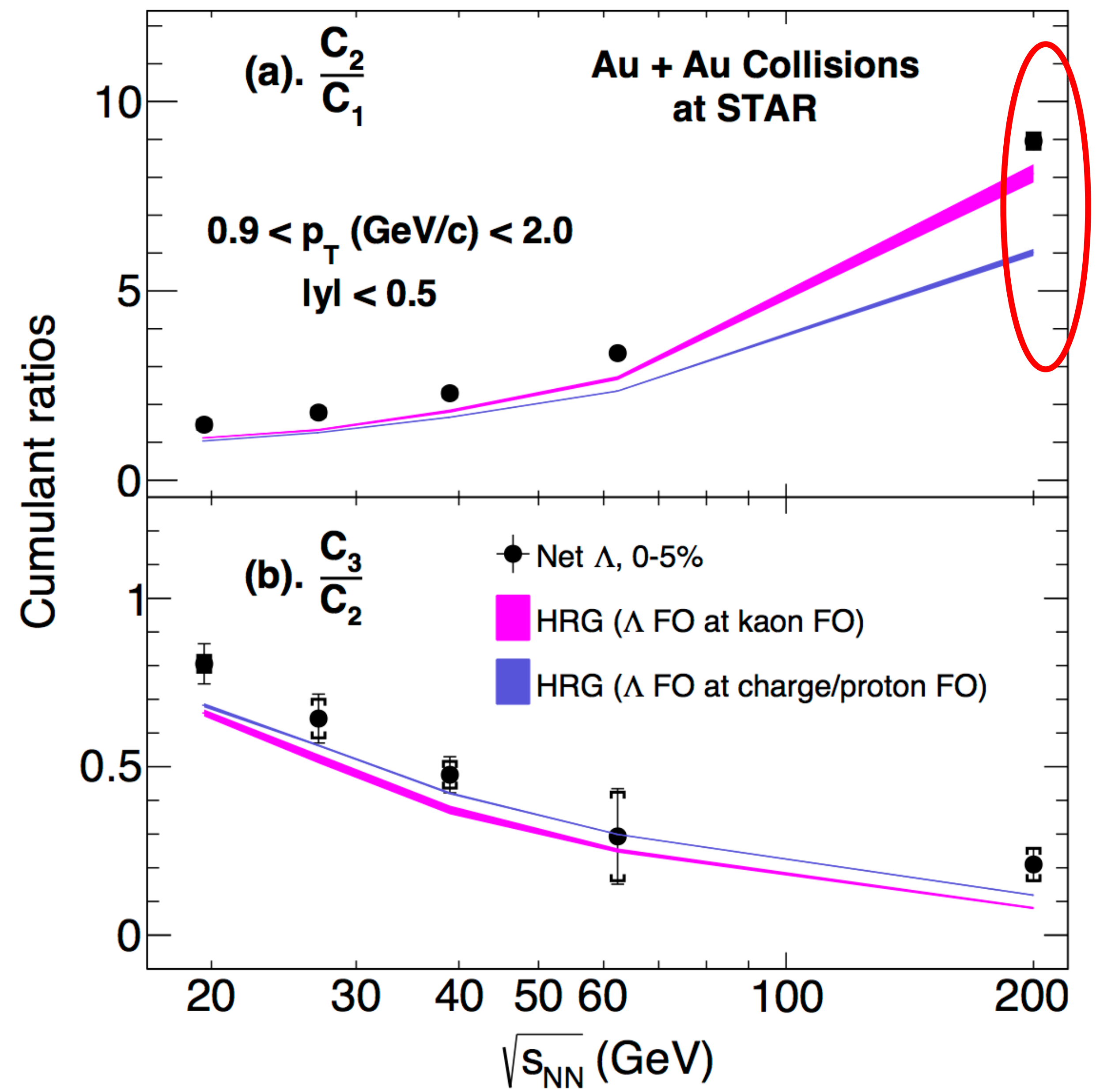


Fluctuations are more sensitive to chemical freeze-out as simple yields. They can be directly compared to susceptibilities on the lattice (P.Alba et al., PRC, (arXiv:1504.03262))

HRG fits to σ^2/M for net-protons, net-charge, net-kaons and net-Lambda STAR data





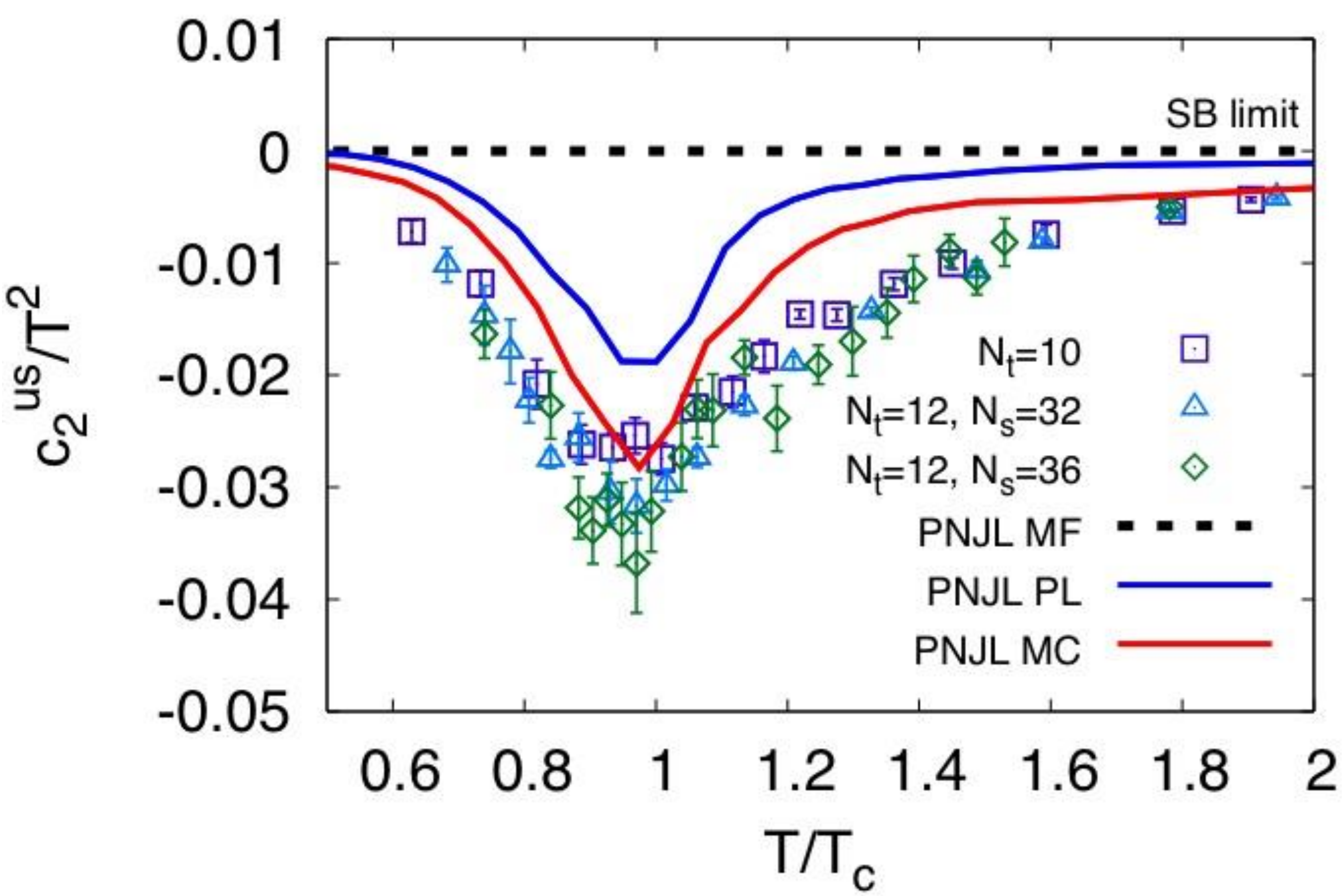


Net-Λ follow the net-kaon freeze-out (R. Bellwied, SQM 2019 based on N.Kulathunga, STAR/UH Ph.D. thesis 2018)

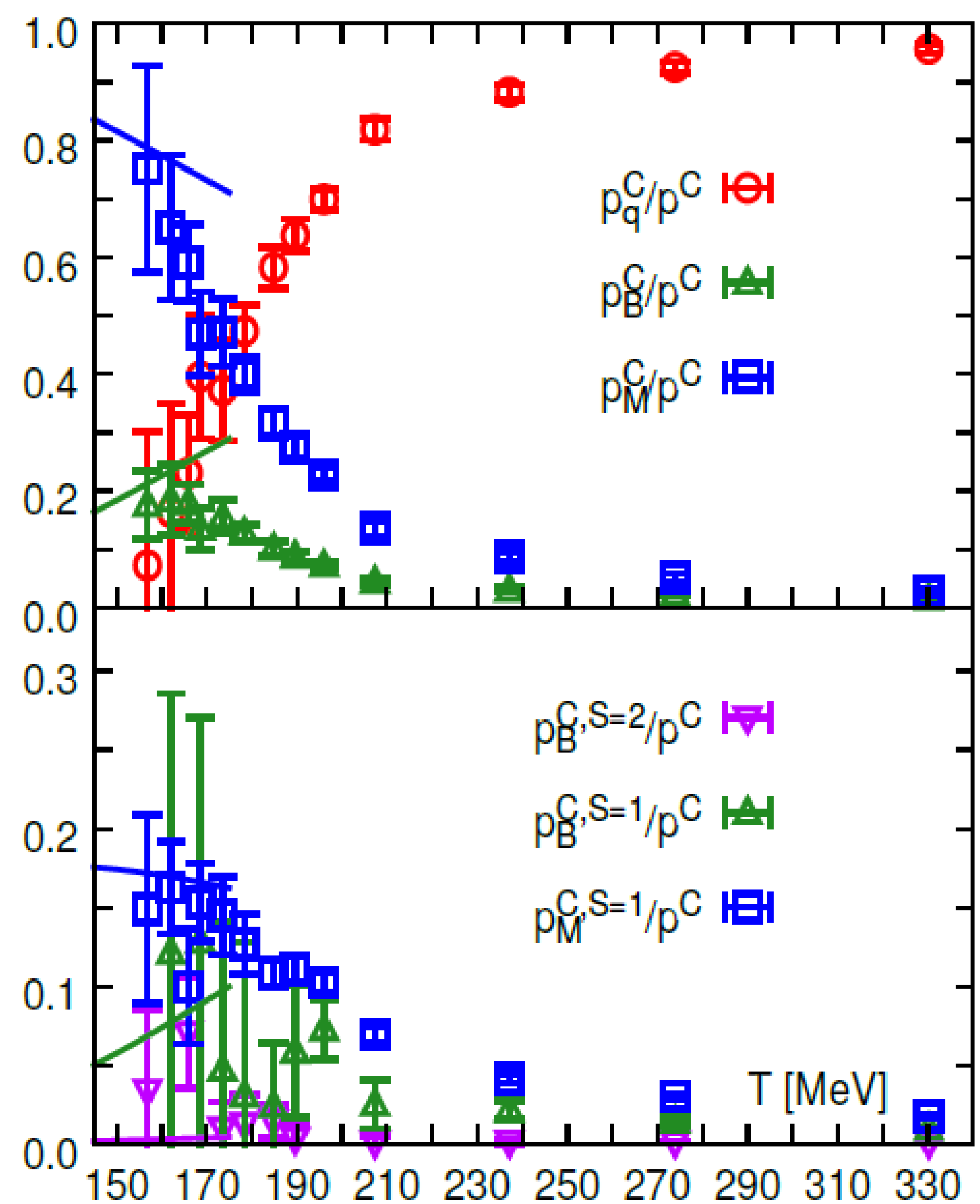
C. Ratti et al., arXiv:1607.02527 and R. Bellwied et al. arXiv: 1805.00888

Is there evidence from other lattice studies for a flavor dependence ?

Bound states in the strange sector (C. Ratti et al., PRD 85 (2012)) through BS Correlator



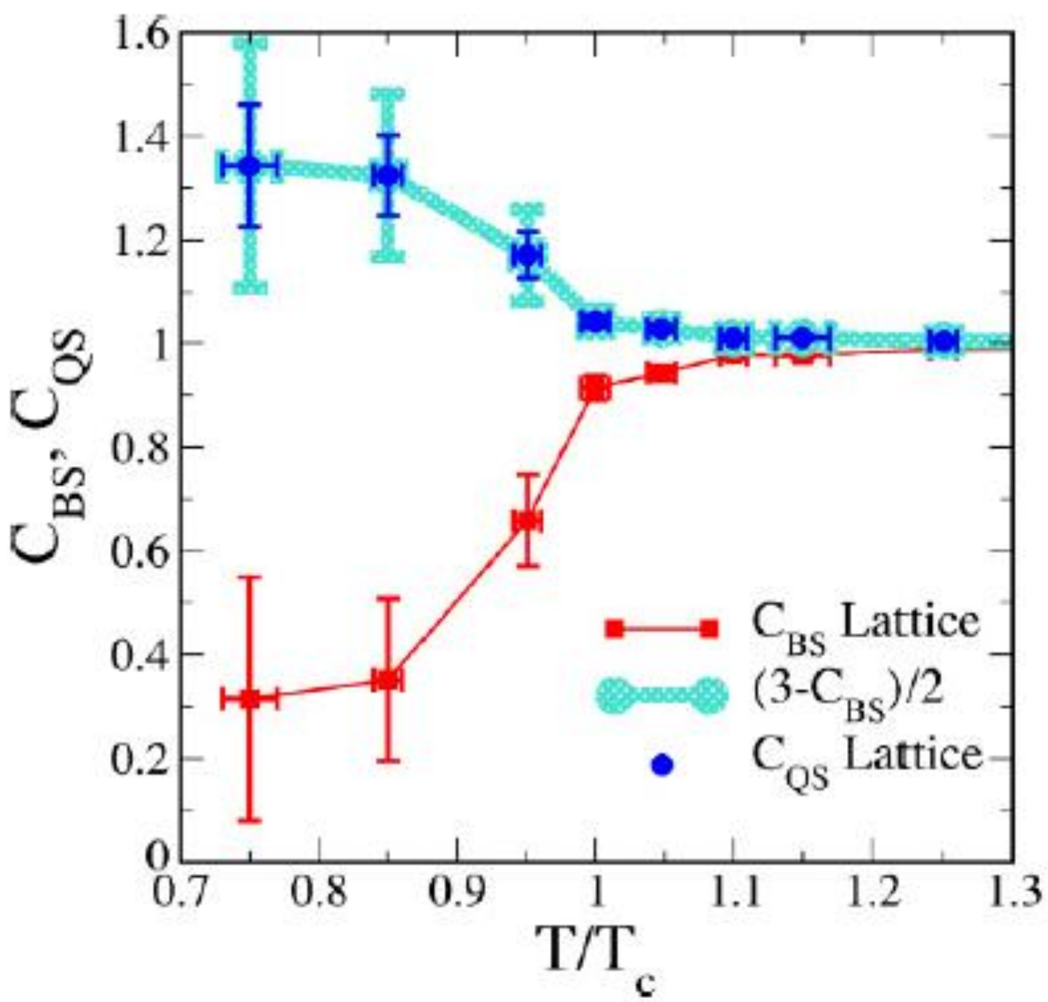
Bound states in the charm sector (S. Mukherjee et al., PRD 93 (2016))



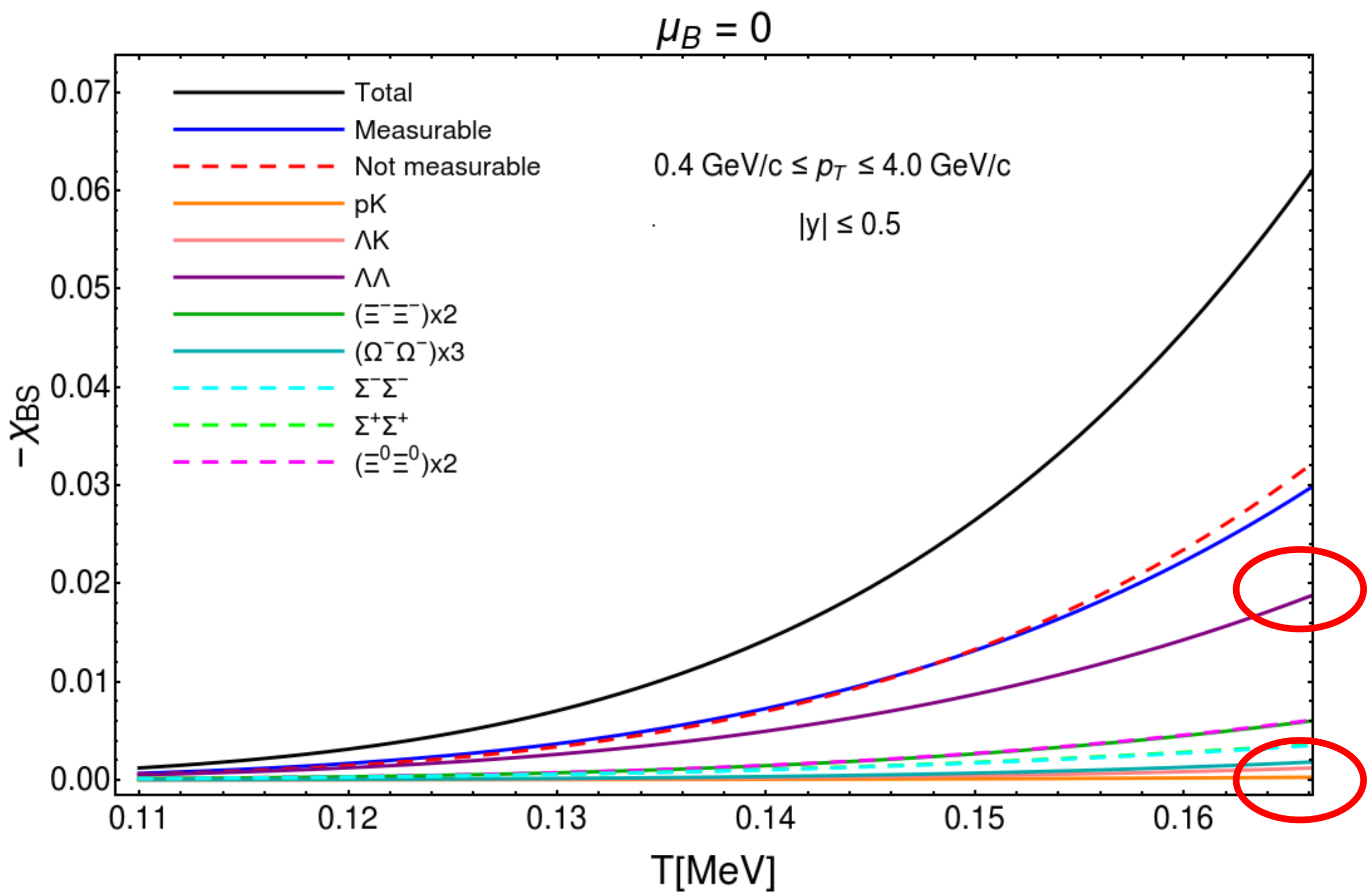
What can we learn from cross-correlators (specifically BS-correlator)?

Determined by the ratio of off-diagonal to diagonal cumulants: $\kappa^{1,1}_{B,S}/\kappa^{2}_{S}$

The related susceptibility ratio: $C_{B,S} = -3\chi^{1,1}_{B,S}/\chi^{2}_{S}$
(-3 is just a normalization factor so that the asymptotic value = +1)



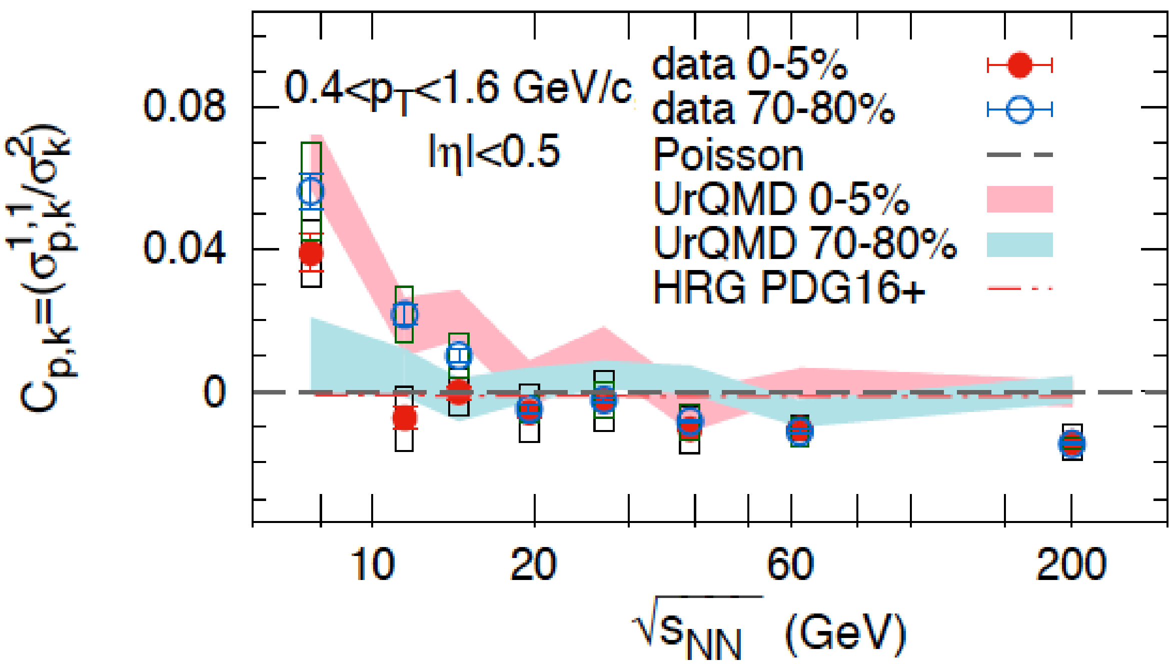
Koch, Majumder, Randrup (2005), Mueller, Majumder (2006)



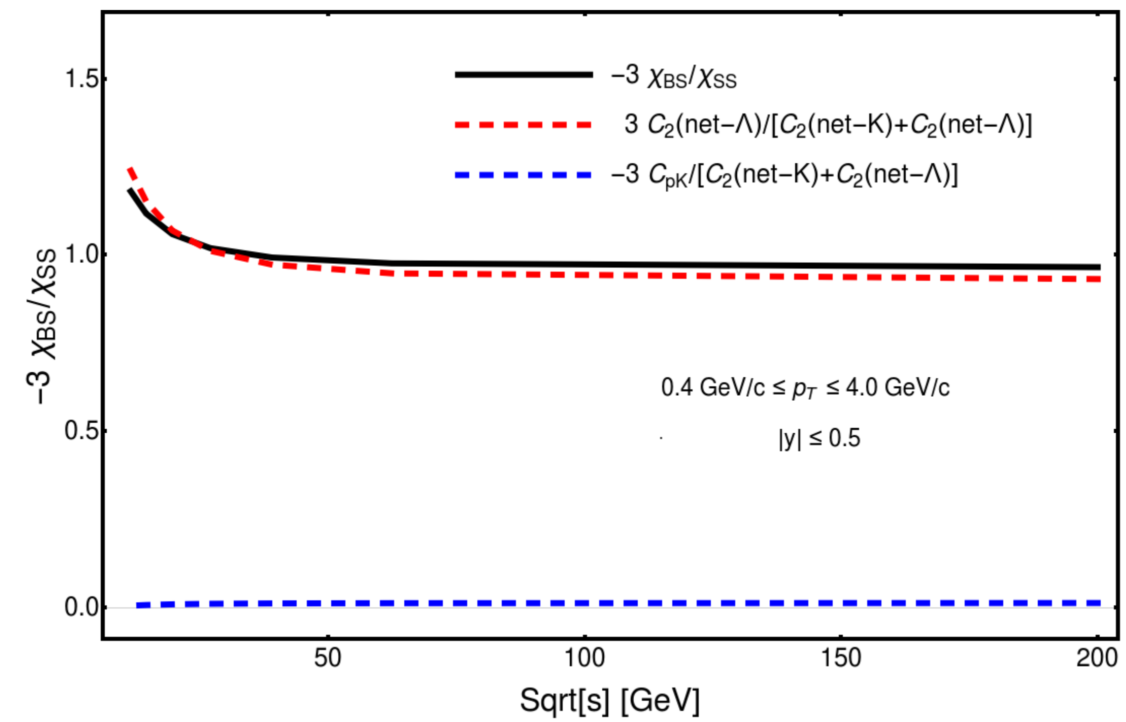
Parotto, Ratti, Stafford (SQM 2019)

Contributions to the 2nd-order off-diagonal BS cumulant

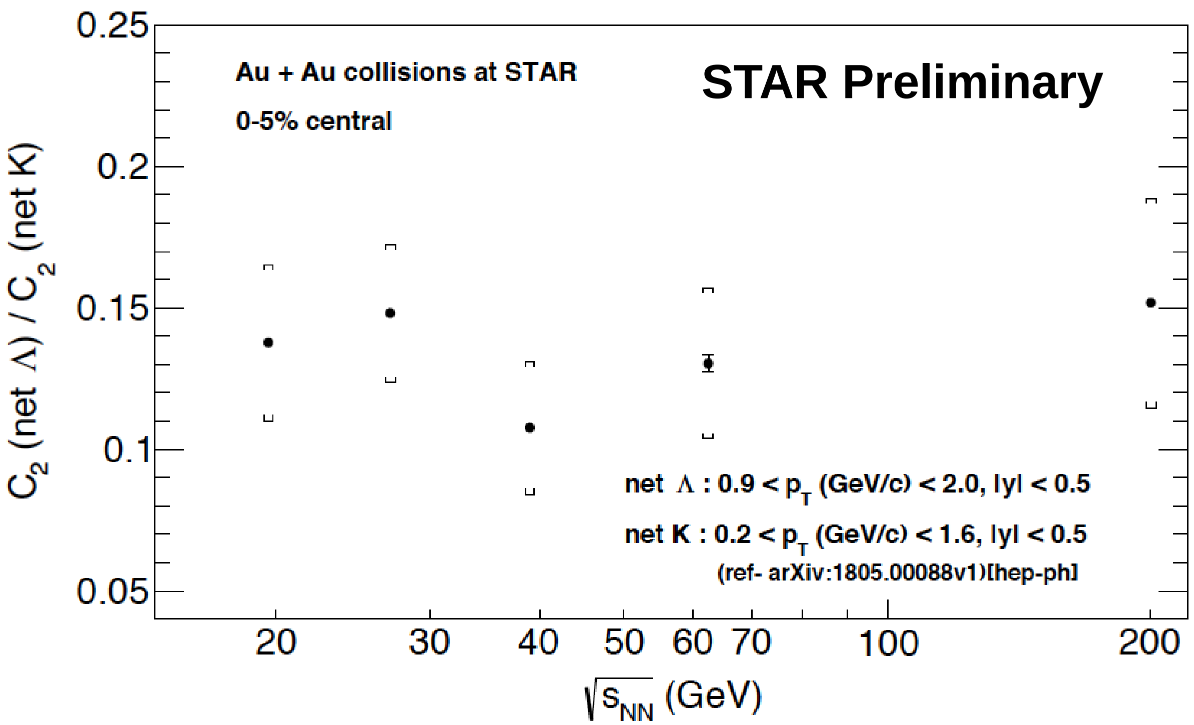
STAR measurements:
pK contribution, arXiv:1903.05370

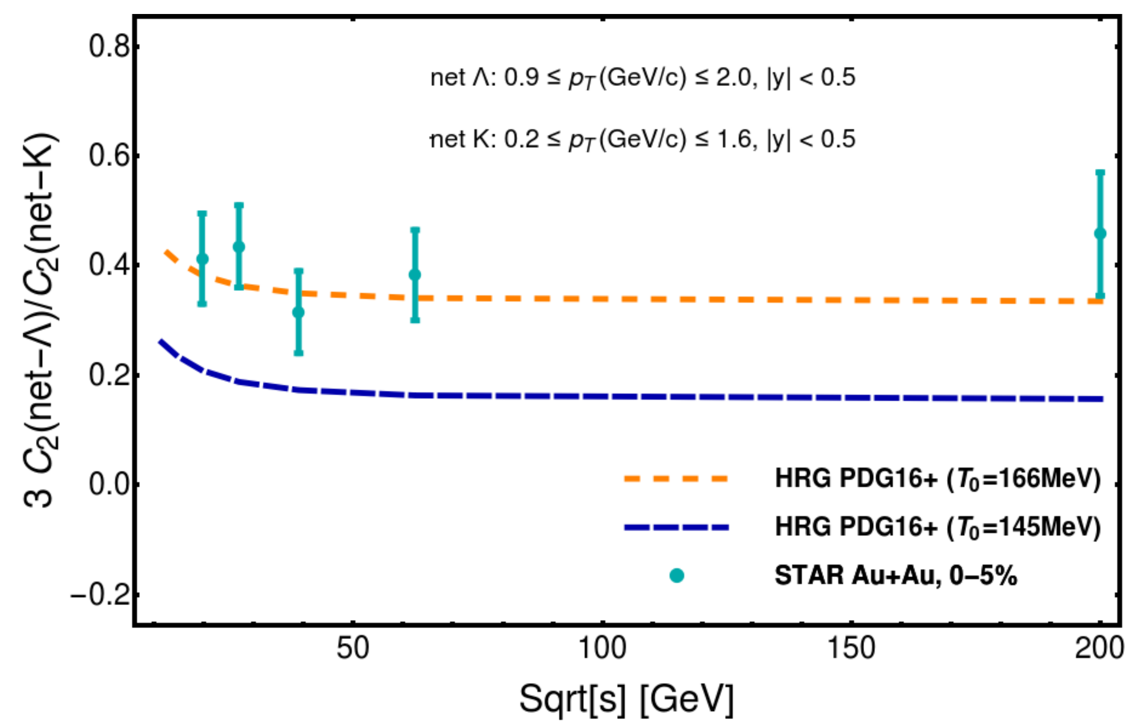


HRG predictions:
Parotto, Ratti, Stafford, SQM 2019



Λ contribution, R. Bellwied (SQM 2019)





Freeze-out is driven by flavor, not baryon number ?

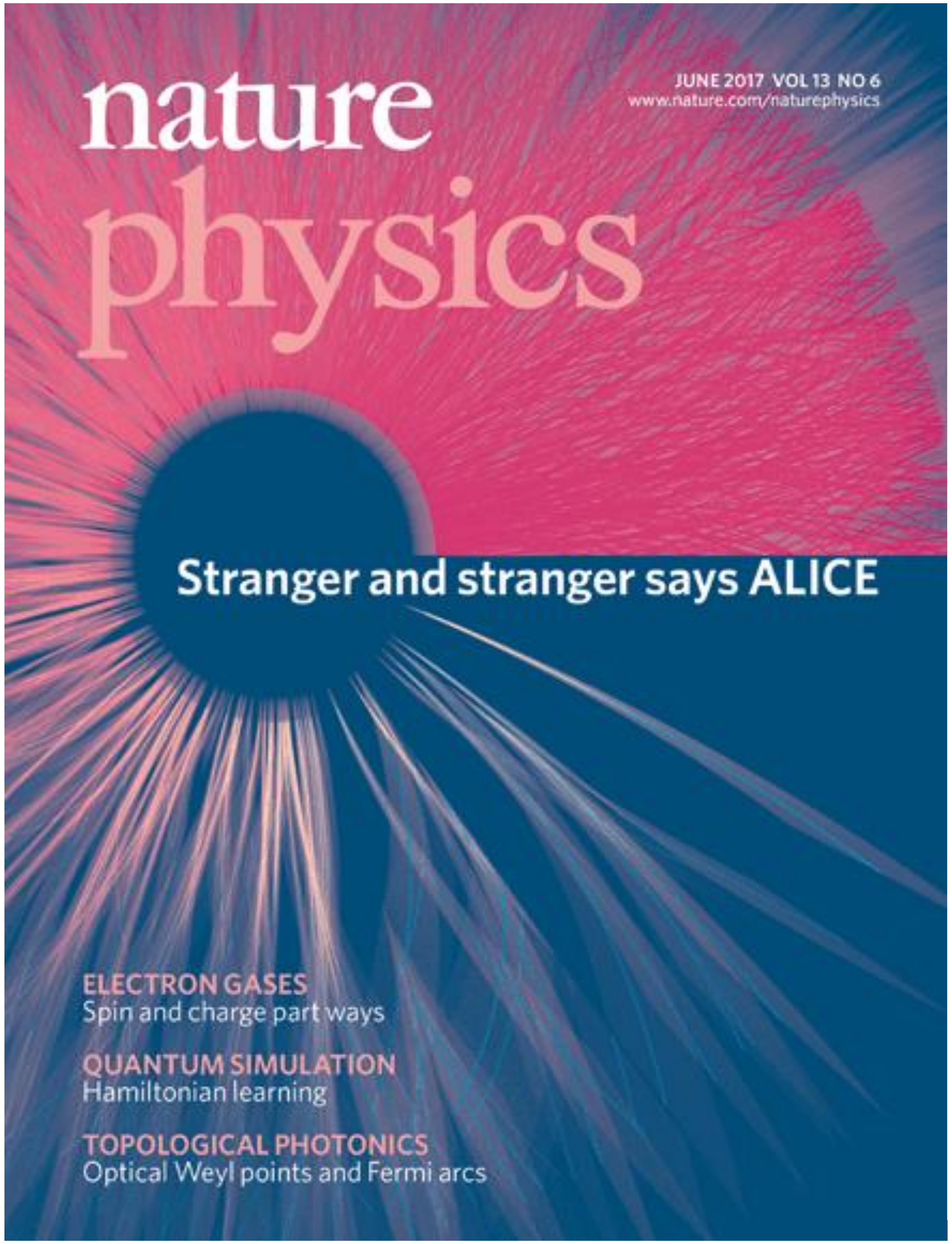
Potential impact of higher temperature freeze-out hypersurface for strangeness

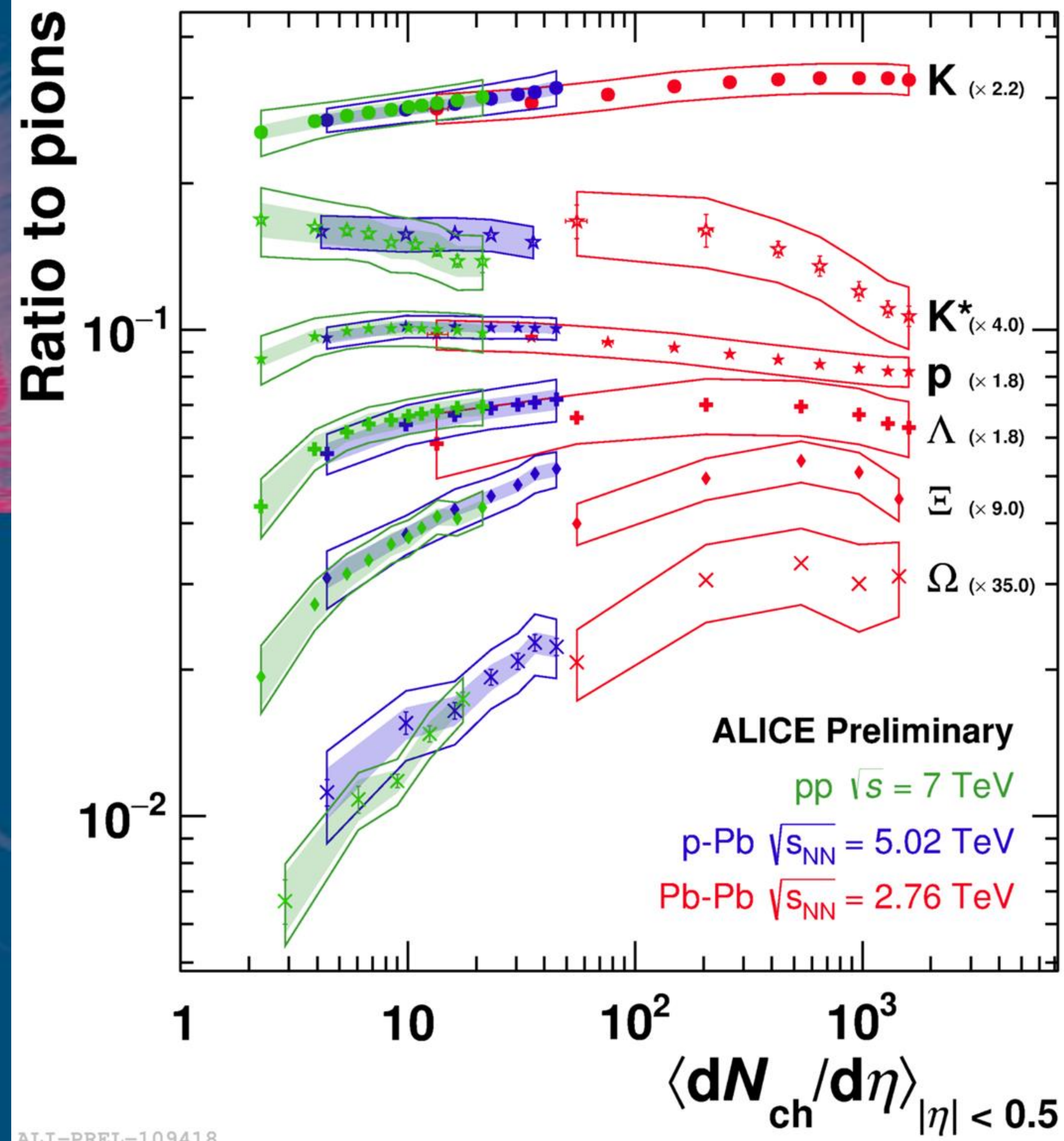
1.) strangeness enhancement

2.) chiral restoration at different T

3.) exotica

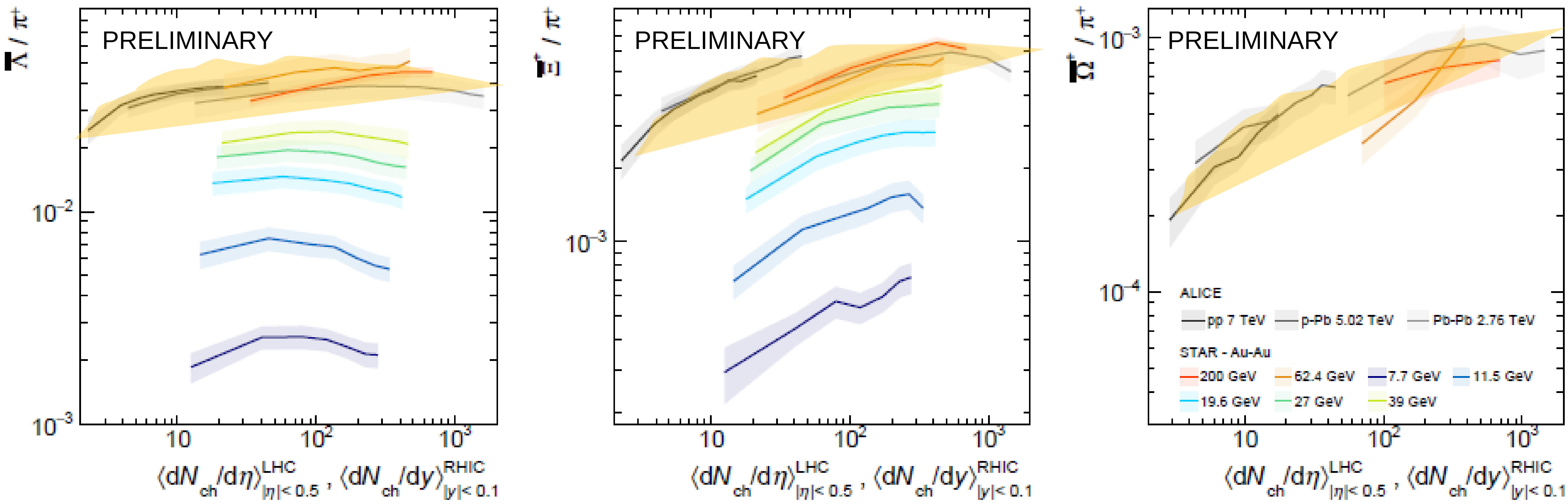
Stranger and stranger from small to large systems (ALICE, arXiv:1606.07424)





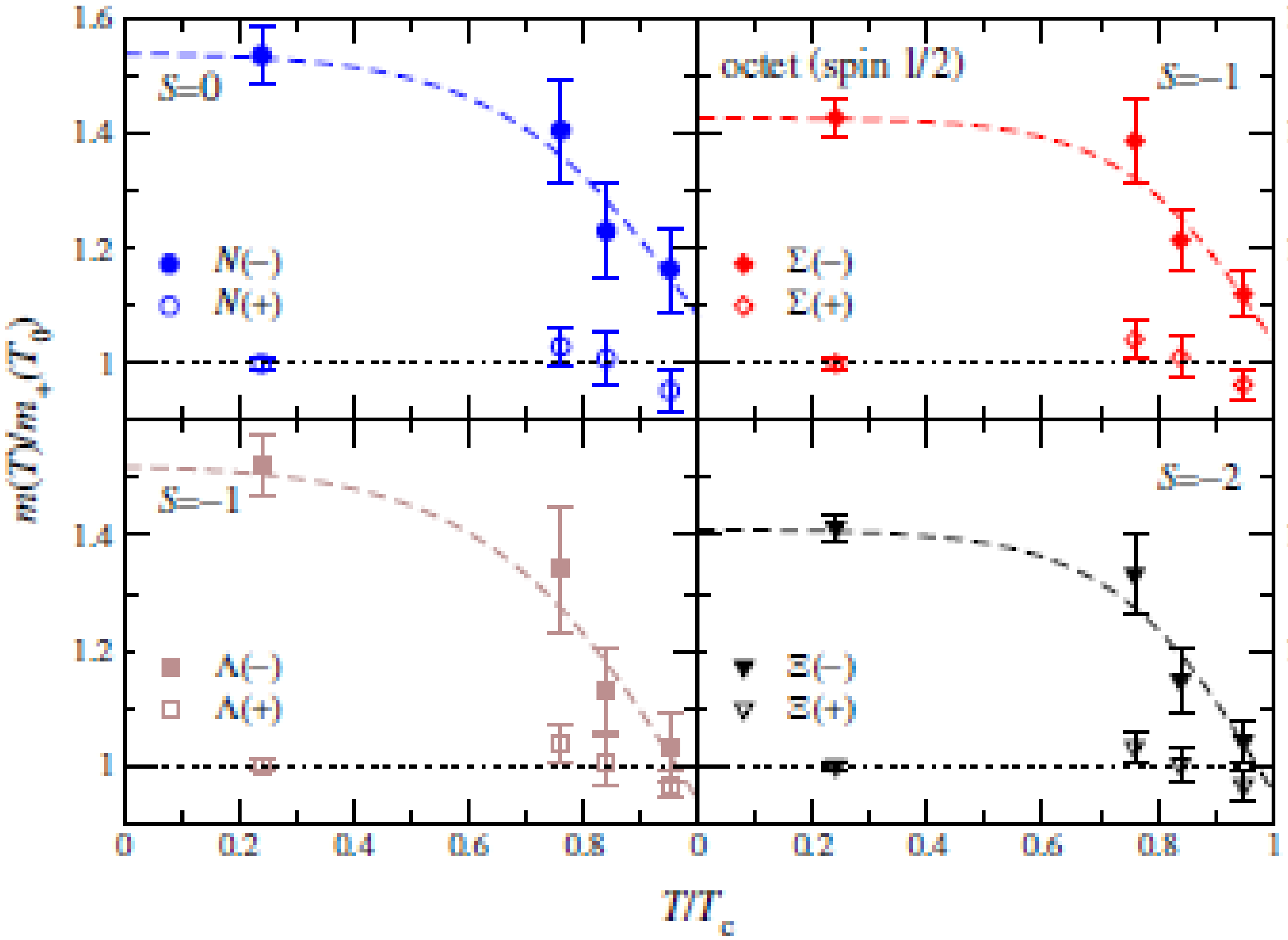
Is it time to re-evaluate strangeness suppression/enhancement ?

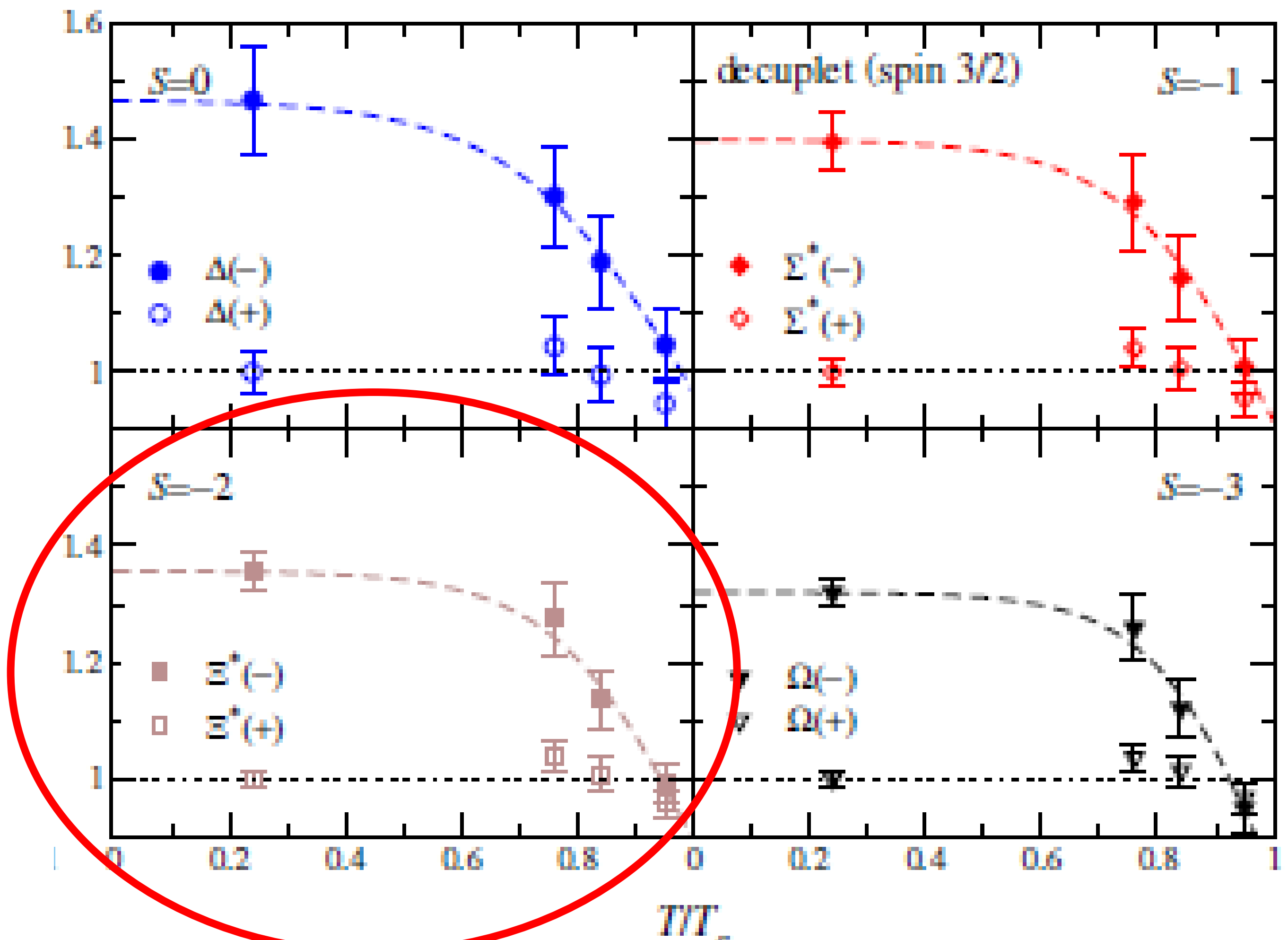
Canonical suppression reduces as a function of energy and as a function of system size (Tounsi, Redlich (2001)). Is suppression over at LHC energies ? Do we only see enhancement ? Can we distinguish ?



Above 39 GeV the curves seem to fall together, no more energy dependence. The volume dependence is still there (γ_s dependence ?). A higher T freeze-out surface in PbPb will lead to actual strangeness enhancement

Effects studied in STAR (arXiv.1906.03732) and ALICE



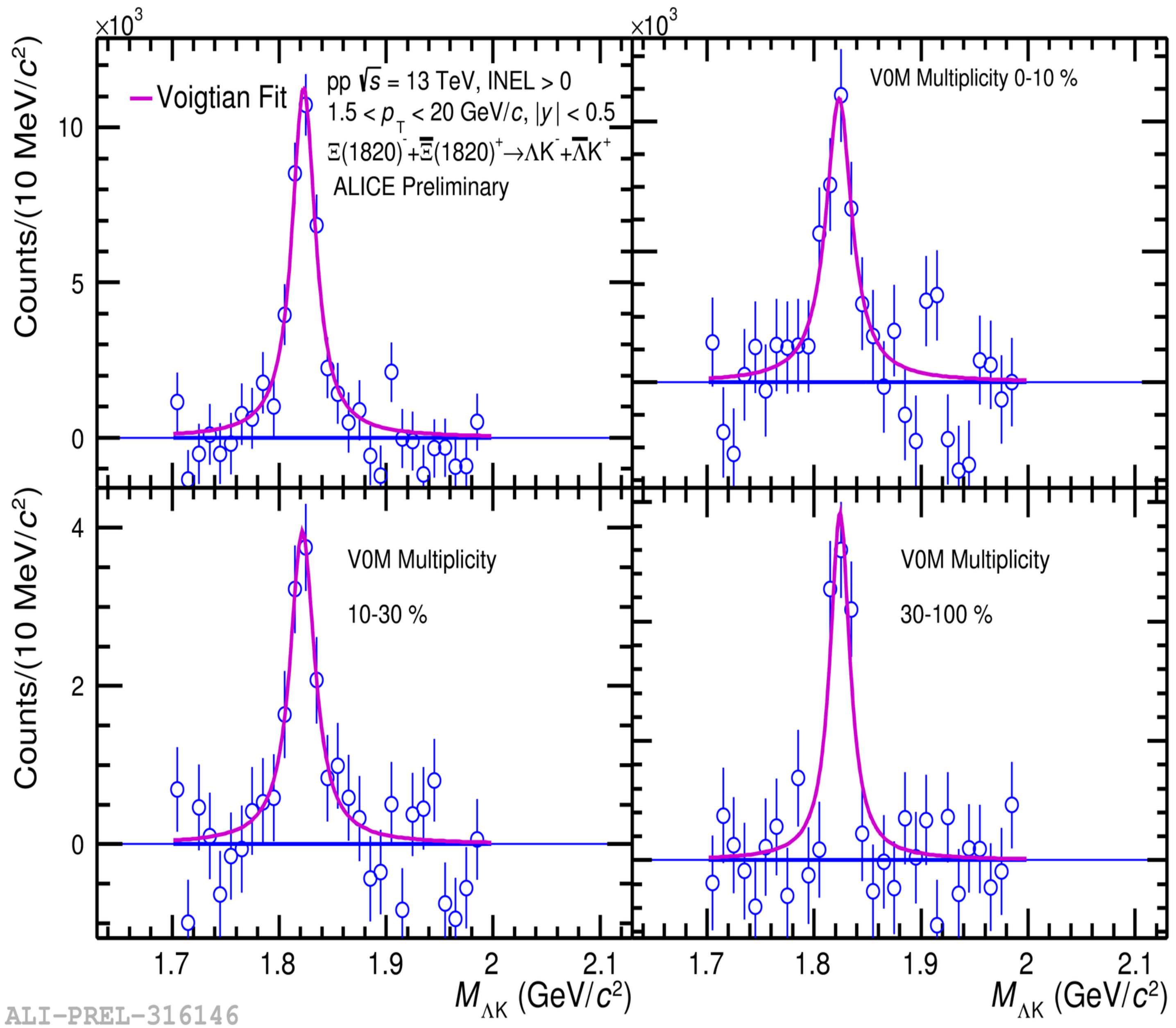


Chiral restoration in the strange baryonic resonance sector

FASTSUM Collaboration:
Baryon spectral functions
JHEP 06 (2017) 034
(arXiv:1703.09246)

- Emerging degeneracy around T_c for chiral partners
- Positive parity masses nearly temperature independent
- Negative parity masses drop as temperature increases
- Experiment: find appropriate chiral partners.
- Study of Ξ(1530) and Ξ(1820) ongoing in ALICE

Parity doubling study
First measurement of Ξ(1820) in 30 years

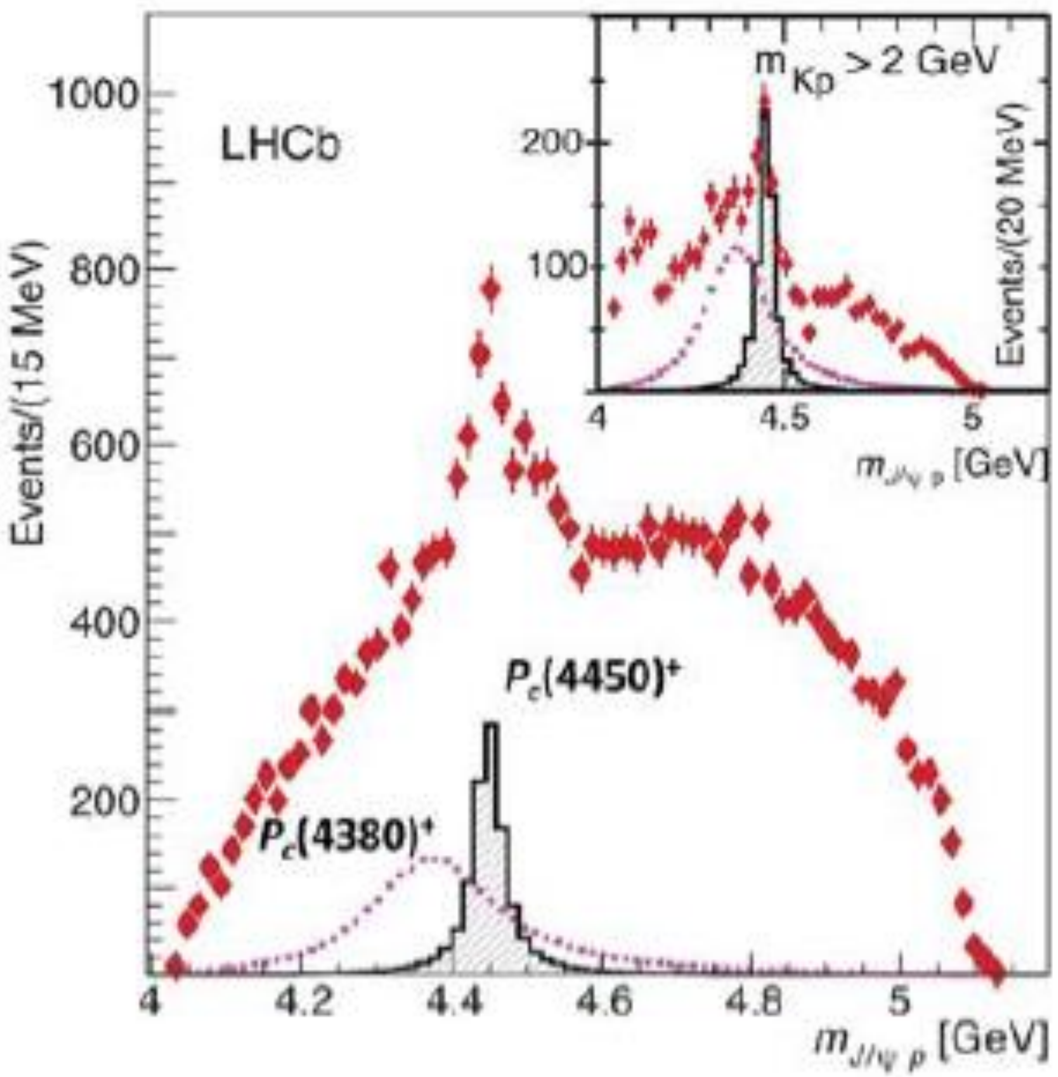


C. J. Myers for ALICE Coll. (SQM 2019)

Exotica:
Penta- and Tetra-quarks from LHCb

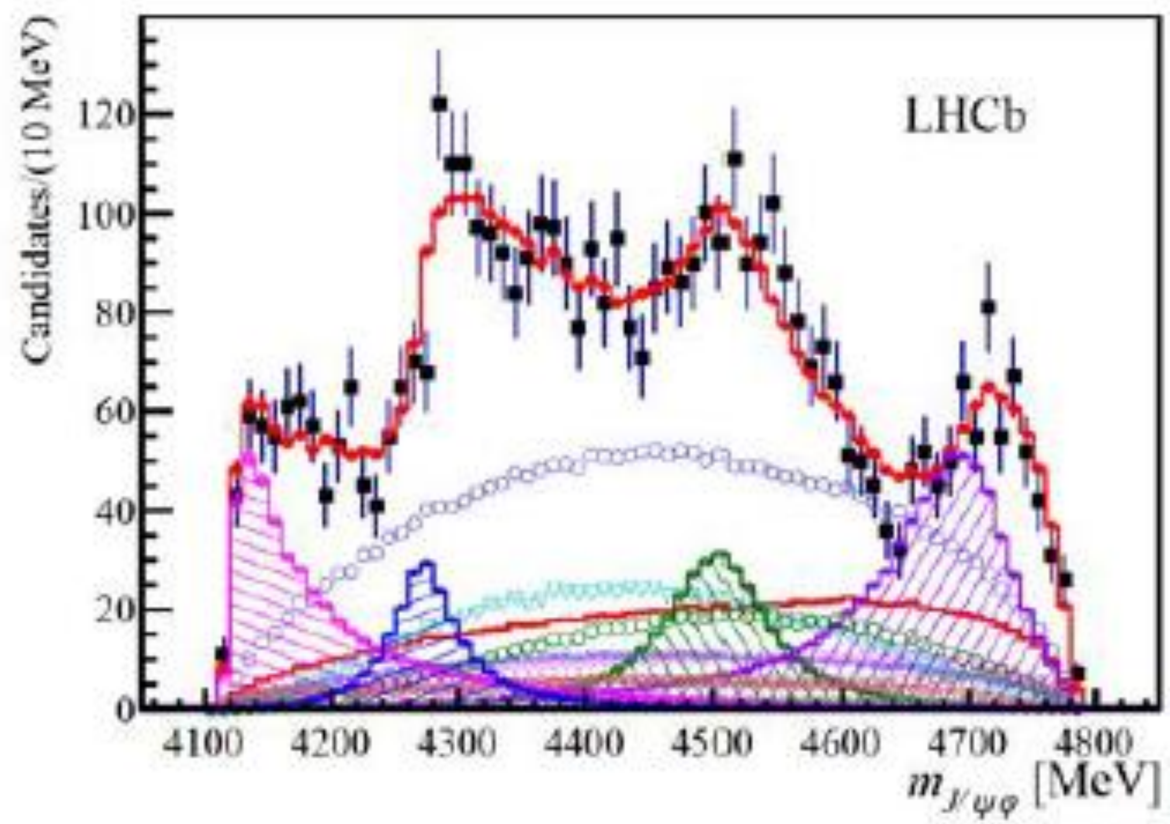
Penta-quark in 2015, 9σ evidence by 2016

In the charm sector: J/ψ p resonance
In Λ_b decays to J/ψ p K^-



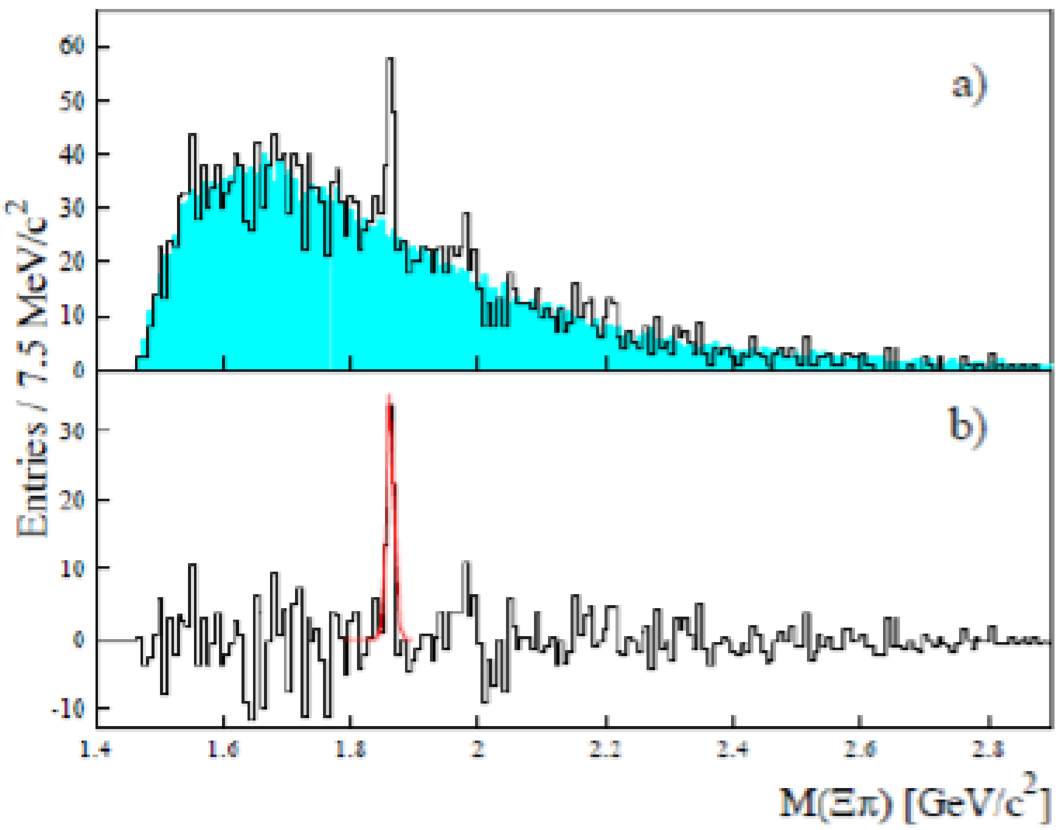
Tetra-quarks in 2016

In the charm sector: J/ψ ϕ resonance
In B^+ decays to J/ψ ϕ K^+

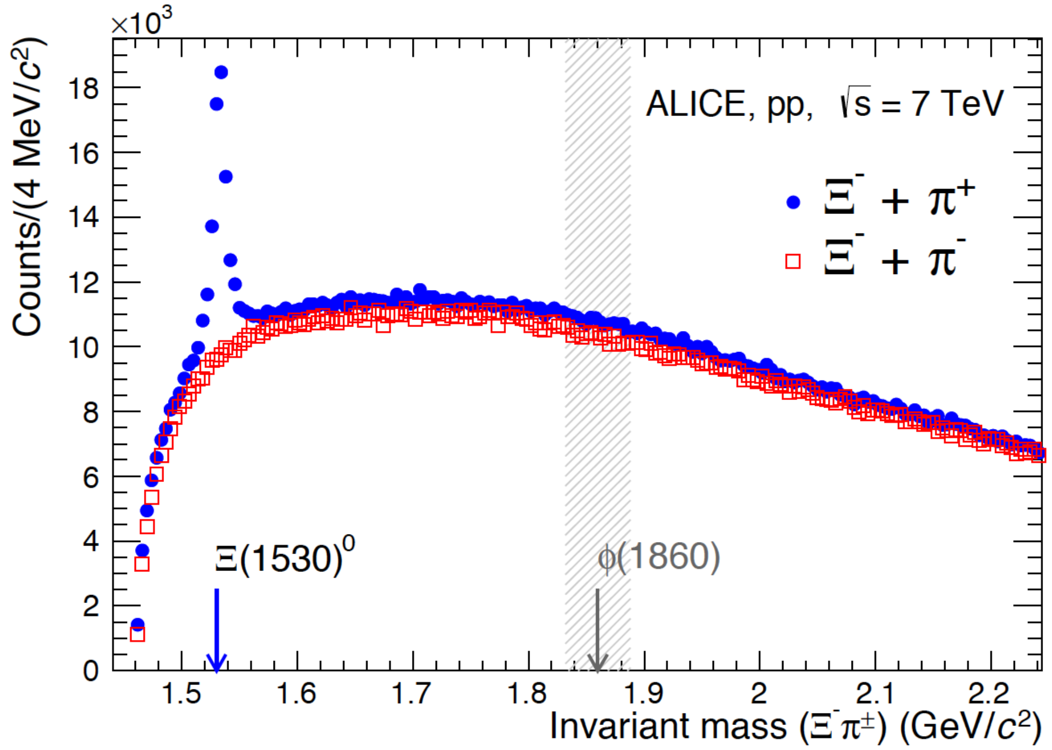


Why nothing in the strange sector ?

Exotica in strange sector ?



Famous pentaquark candidate from NA49 (2008) in Ξπ channels (ϕ(1860)) (dsdsubar)
Never retracted, never confirmed



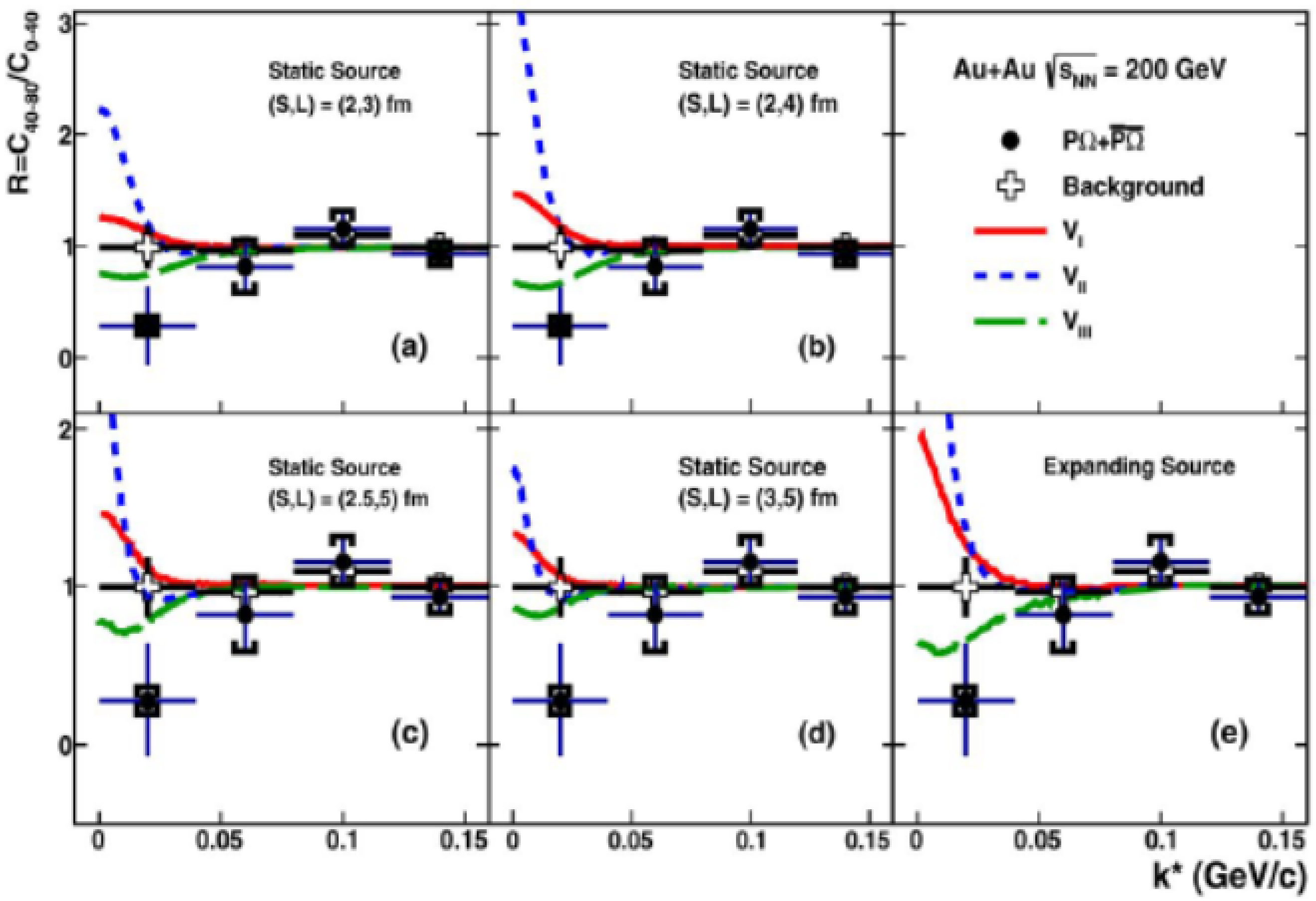
No evidence for H-dibaryon or ϕ(1860) in ALICE data.

Maybe we are looking in the wrong channels. In the charm sector all tetra- and penta-quarks seem to require closed charm components.

Keep looking !!

Another avenue: Femtoscopic studies of di-baryon strong interactions (STAR & ALICE)

- Study of the p-Ω^- correlation function in Au-Au collisions at $\sqrt{s_{NN}}$ = 200GeV STAR Collaboration. Phys. Lett. B790 (2019) 490-497
- Observable: ratio of the correlation function peripheral/central collisions.
- Comparison with Lattice QCD calculations (with large masses)



- Test different fits to Lattice QCD data (delivering **three different binding energies of the NΩ**):

Binding energy ($\mathbf{E_b}$), scattering length ($\mathbf{a_0}$) and effective range ($\mathbf{r_{eff}}$) for the Spin-2 proton–Ω potentials [24].

Spin-2 pΩ potentials	V_I	V_{II}	V_{III}
$\mathbf{E_b}$ (MeV)	–	6.3	26.9
$\mathbf{a_0}$ (fm)	−1.12	5.79	1.29
$\mathbf{r_{eff}}$ (fm)	1.16	0.96	0.65

[24] K. Morita, A. Ohnishi, F. Etminan, T. Hatsuda, Phys. Rev. C 94 (2016), 031901

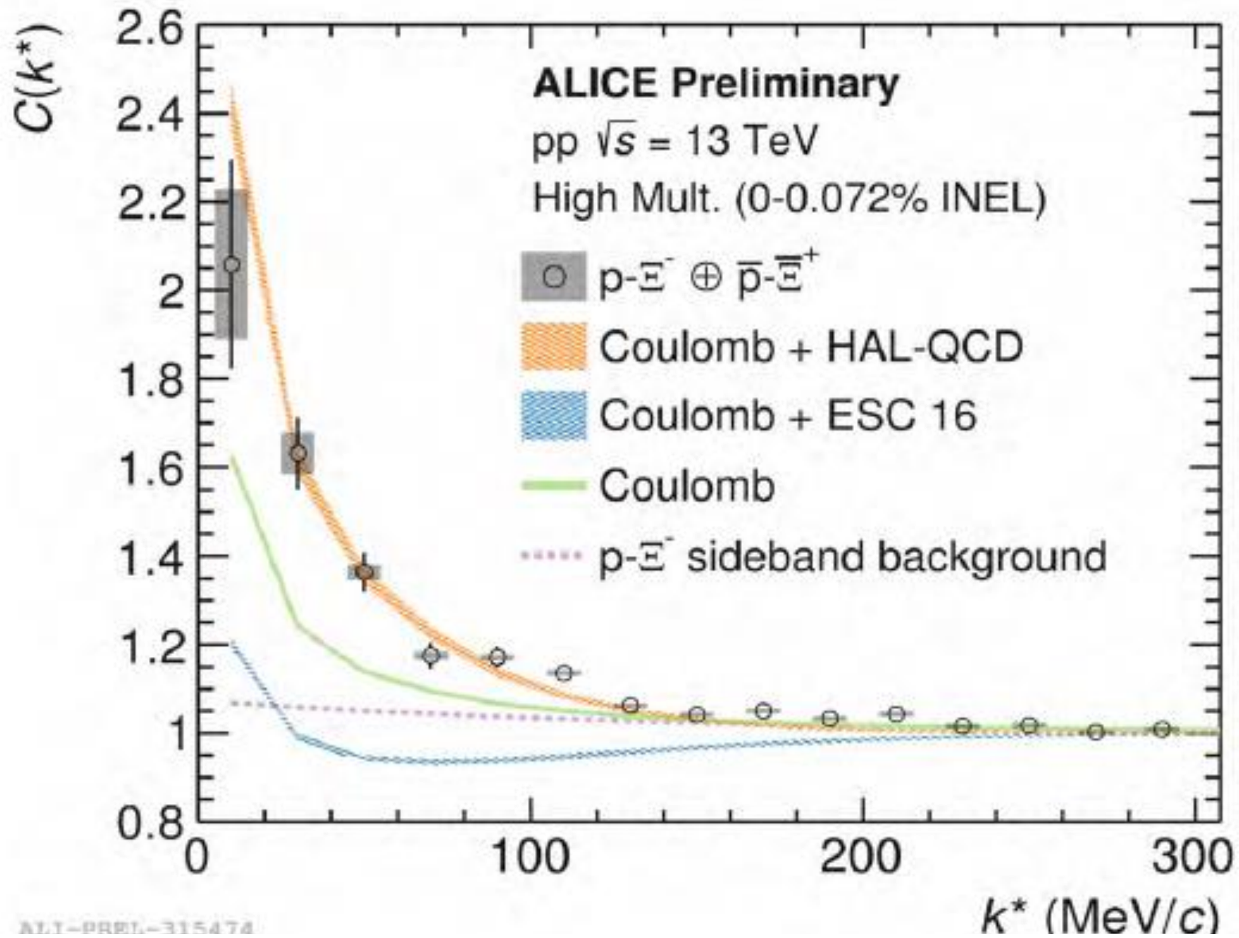
STAR data favor V_{III}, with $\mathbf{E_b}$ **= 27 MeV**

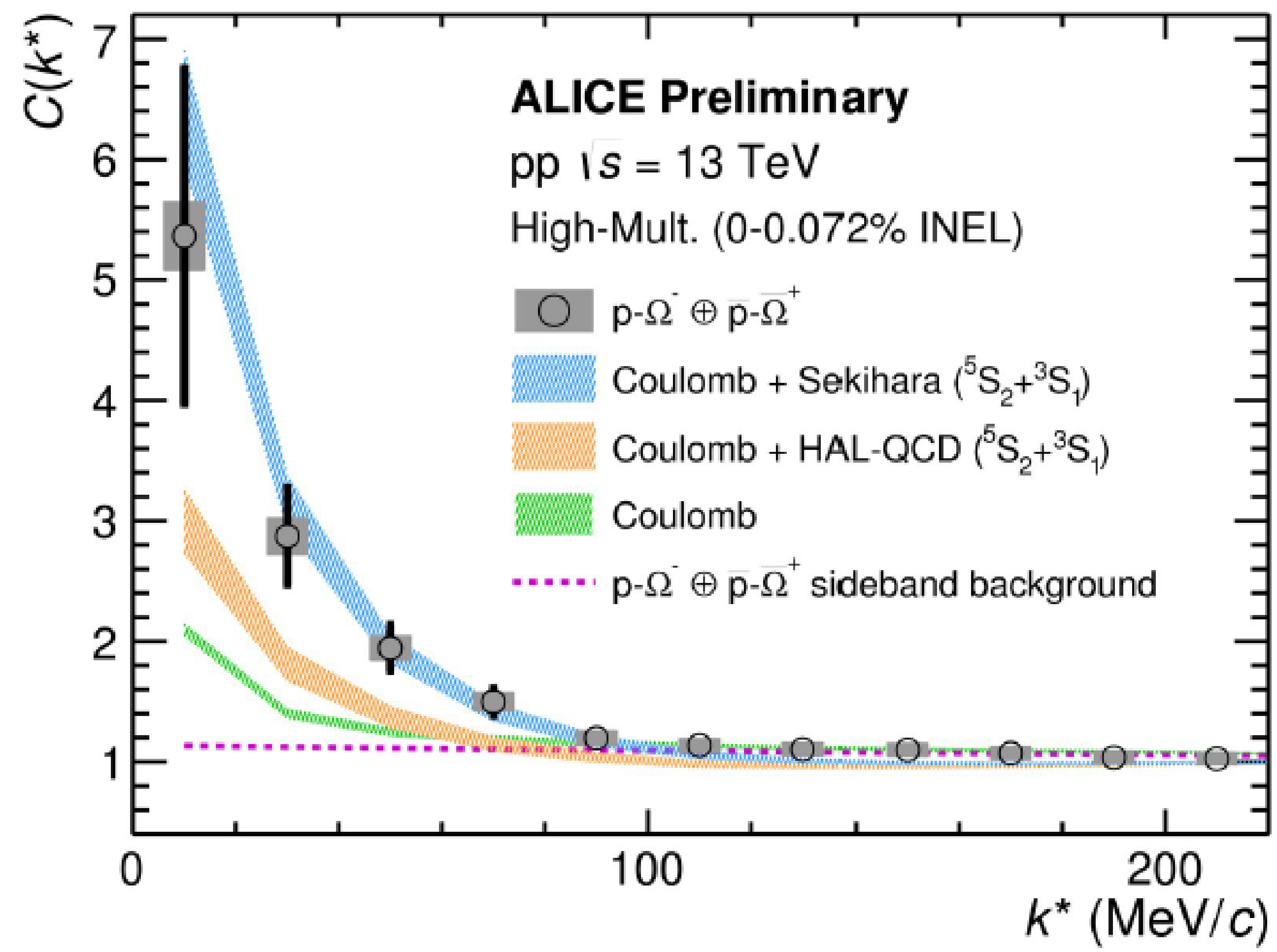
3

STAR concluded attractive potential and large binding energy For a di-baryon state (based on comparison to lattice QCD (with wrong hadron masses).

Another avenue: Femtoscopic studies of di-baryon strong interactions (STAR & ALICE)

ALICE, High multiplicity pp collisions, pΞ and pΩ correlations





ALICE concluded attractive potential and but very small binding energy for a di-baryon state (based on comparison to lattice QCD (with correct hadron masses) and meson exchange model.

Conclusions / Outlook

- High precision (continuum limit) lattice QCD susceptibility ratios indicate *flavor separation in the crossover from the partonic to the hadronic matter.*
- There are hints, when comparing to hadron resonance gas and PNJL calculations, that this could lead to a short phase during the crossover in which strange particle formation is dominant.

- If the abundance of strange quarks is sufficiently high (LHC) this could lead to *enhancements in the strange hadron yields (evidence from ALICE)* and it could lead to *strangeness clustering (exotic states: dibaryons, strangelets)* or *higher mass strange Hagedorn states* (as predicted by Quark Models).
- Dynamic quantities that evolve during the deconfined phase will be affected as long as the hadronization temperature plays a significant role, i.e. quark phase is shortened for heavier flavors, which could explain flavor effects in R_{AA} if energy loss builds up near Tc.

- Ongoing project (BEST Collaboration): The phases can be linked in a hydrodynamic calculation by using a mixed EOS from lattice and HRG with varying flavor-dependent switching temperatures.

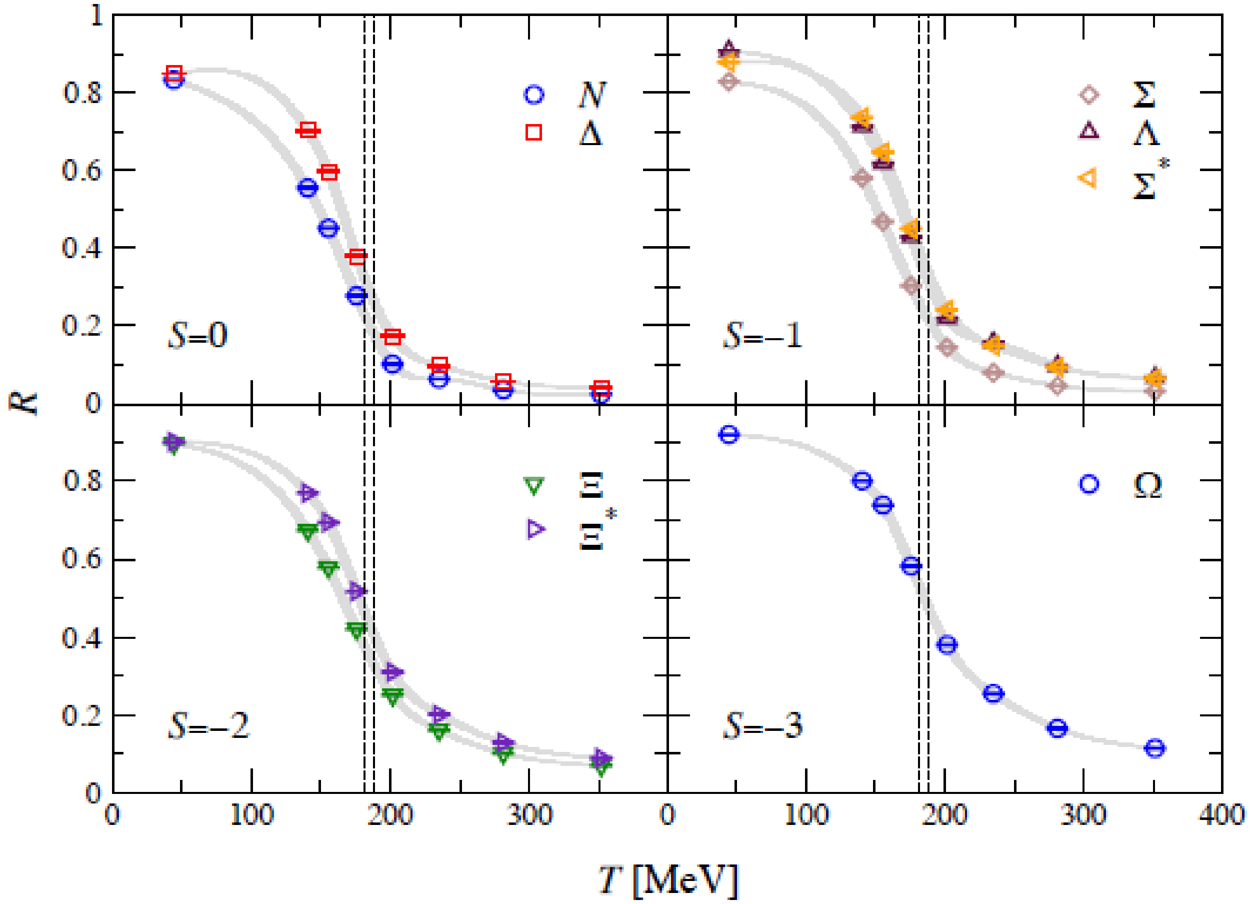
Some speculation, conclusions

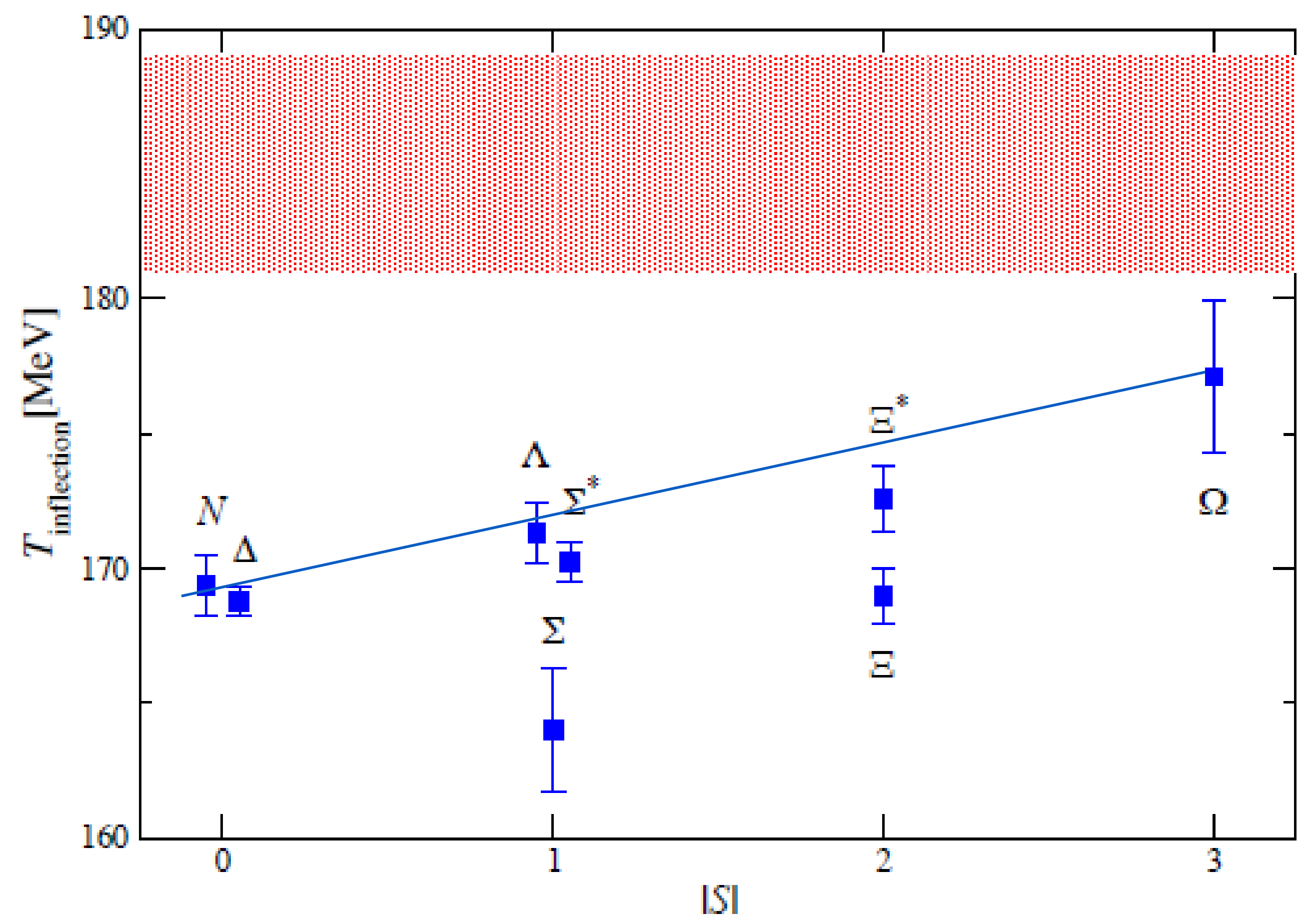
- We are starting to learn about the intricate hadronization mechanism in the QCD crossover region.
- There are plenty of ideas of the dynamic system ranging from quark clustering into Hagedorn states (Greiner, Noronha-Hostler) over interacting hadron states (Vovchenko, Stoecker) to colored and color neutral quasiparticles (Bratkovskaya, Cassing) to constituent quarks embedded in gluon clouds (Stock).
- Lattice seems to indicate quantum number dependencies in the crossover region. Flavor (thus quark mass) seems to play more of a role than baryon number or charge.
- By studying identified particle production features in terms of quantum number fluctuations we can learn detailed features of the hadronization process not only from following the flavor dependencies (up to charm), but also the charge, isospin and baryon number dependencies.

Backup

Flavor dependent 'melting' temperature for baryons

$$R \equiv \frac{\sum_{n=1}^{N_\tau/2-1} R(\tau_n)/\sigma^2(\tau_n)}{\sum_{n=1}^{N_\tau/2-1} 1/\sigma^2(\tau_n)}$$





The thermal charm opportunities / challenges

Ultimately the question of hadronization in an analytic crossover region can be reduced to:

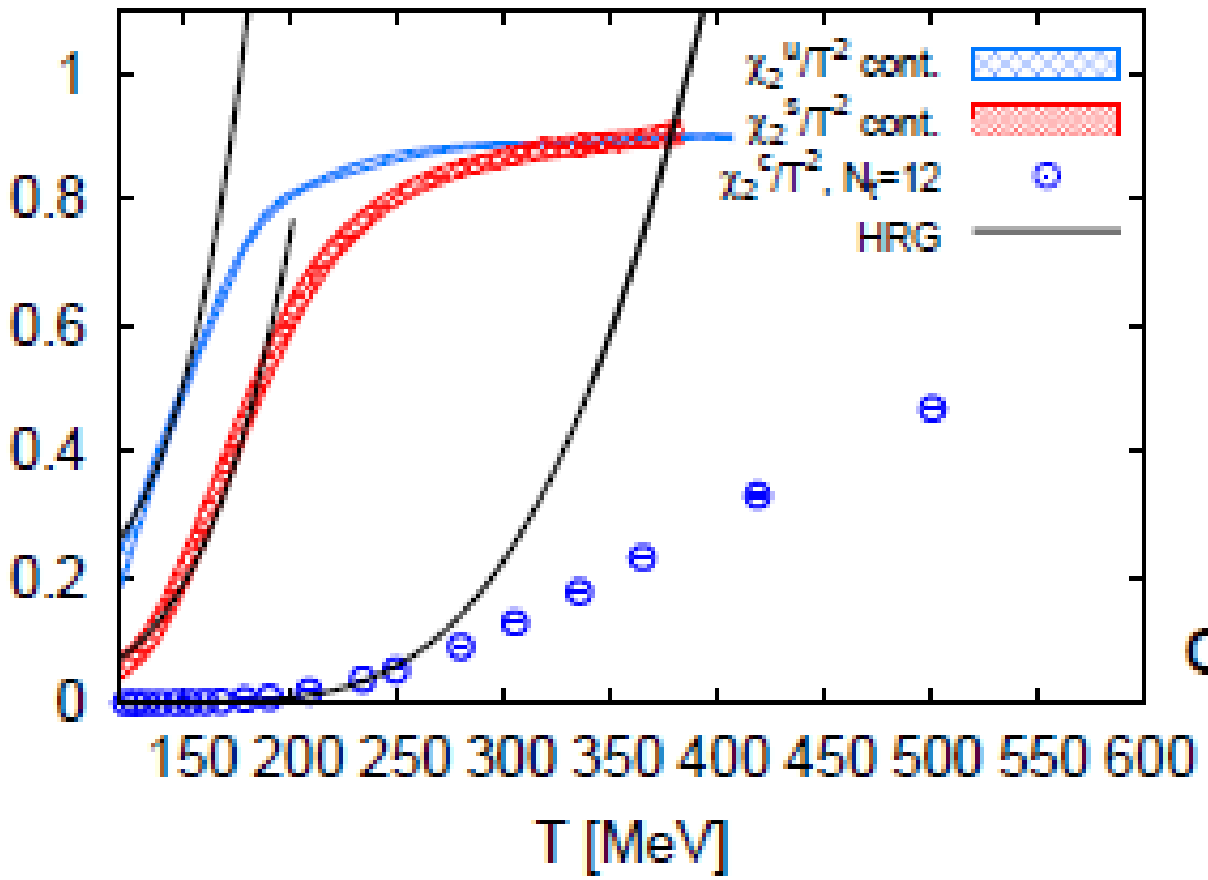
- Is there a flavor (quark mass) or other quantum number dependence when looking at lattice results from the QGP side (WB et al.)?
- Is there a hadron mass dependence when looking at HRG results from the hadronic side (Vovchenko et al.)?

At high collision energies charm can be thermally produced (C.M. Ko et al.) and/or equilibrate during the cooling of the deconfined phase.

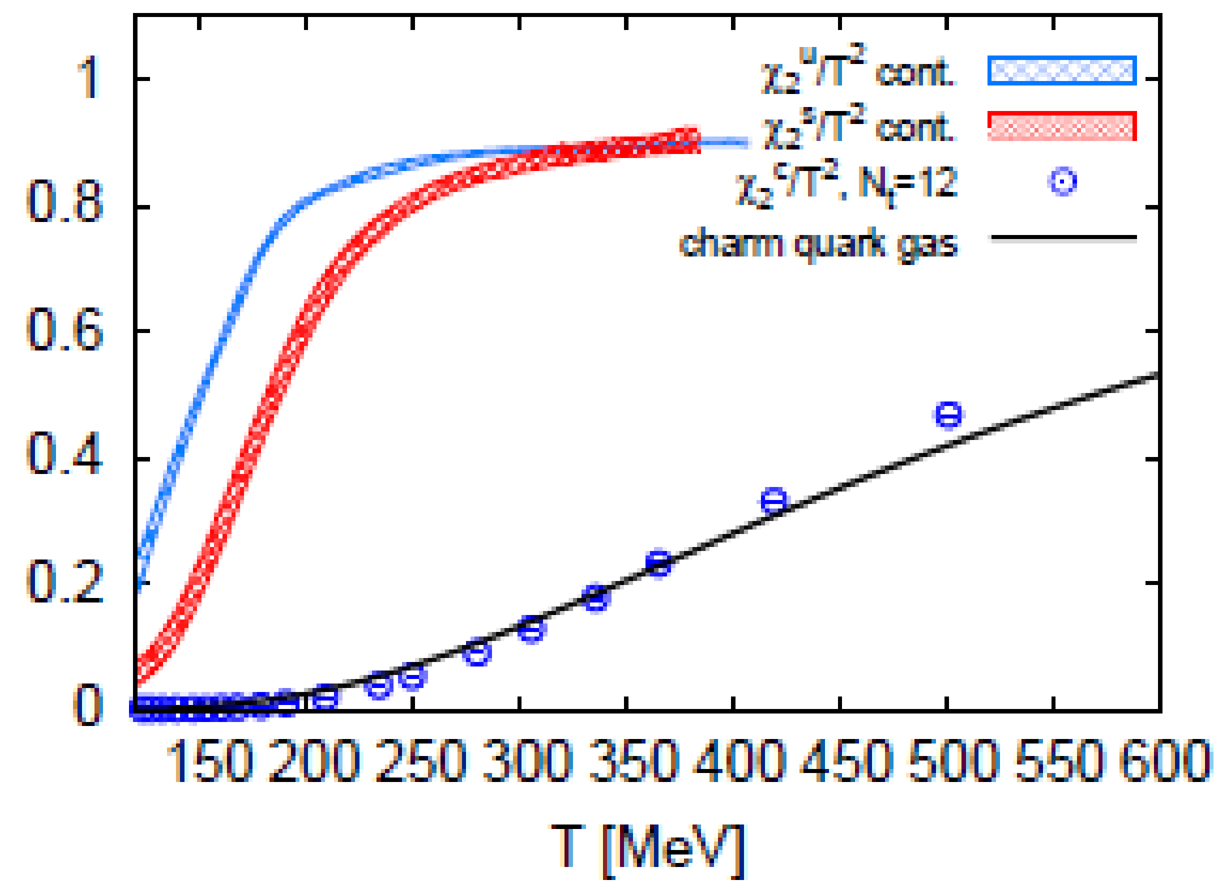
Charm fluctuation measurements could be used to explore a flavor dependent decoupling temperature for produced particles (Graf et al. (1802.07908)).

What does lattice have to say about a flavor/quark mass dependence of the freeze-out based on open charm ?

Lattice QCD: WB results (1507.04627)



or



- survival of open charm hadrons up to $T \simeq 2T_c$?
- HRG results agree with the lattice up to the inflection point in the data

- thermal excitation of charm quarks takes place at larger temperatures
- ideal gas of charm quarks agrees with lattice

need for non-diagonal quark number susceptibilities