Quarkonium Production in Jets at the LHC

Lin Dai Duke University



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Motivations



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NRQCD factorization formalism (LDME extractions)

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- Fragmenting Jet Functions (FJF, analytic tool)

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+ Summary

Some Motivations

 Much of quarkonium production pheno.
 based on NRQCD factorization formalism



[Butenschon, Kniehl, 1105.0820]

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In NRQCD, LDMEs supposed to be universal (however...)

 Quarkonium production in jets provides us new way of studying these issues (cleaner)



Estimate v: $mv^2 \sim \alpha_s/r \sim \alpha_s mv \Rightarrow v \sim \alpha_s$

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$$mv^2 \sim \alpha_s/r \sim \alpha_s mv \Rightarrow v \sim \alpha_s$$

Power count fields:

$$\begin{array}{l} \langle H| \int d^3 x \psi^{\dagger} \psi | H \rangle \sim 1 \\ \int d^3 x \sim 1/(mv)^3 \\ \psi \sim (mv)^{3/2} \end{array}$$

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Power count fields:



Construct NRQCD Lagrangian (and other composite operators) order by order:

Charmonium: $v^2 \sim 0.3$

Operator	Estimate
ψ	$(Mv)^{3/2}$
χ	$(Mv)^{3/2}$
D_0 (acting on ψ or χ)	Mv^2
D	Mv
$g\mathbf{E}$	$M^2 v^3$
$g\mathbf{B}$	$M^2 v^4$
gA_0 (in Coulomb gauge)	Mv^2
$g\mathbf{A}$ (in Coulomb gauge)	Mv^3

NRQCD Factorization

Cross section:

$$\sigma(H) = \sum_{n} d_n \quad (0 | \mathcal{O}_n^H | 0)$$

$$\text{LDMEs (nonperturbative, scale < m_H)}$$

$$n: {}^{1}S_0^{[8]}, {}^{3}S_1^{[1,8]}, {}^{3}P_J^{[8]}, \cdots$$
e.g.
$$\langle 0 | \mathcal{O}_{{}^{3}S_1^{[8]}}^H | 0 \rangle = \sum_{X} \langle 0 | \chi^{\dagger} \sigma^i T^a \psi | HX \rangle \langle HX | \psi^{\dagger} \sigma^i T^a \chi | 0 \rangle$$

NRQCD Factorization

Factorized form of FF:



NRQCD Factorization

Factorized form of FF:

$$D_{q/g}^{H} = \sum_{n} d_{q/g,n} \left(0 | \mathcal{O}_{n}^{H} | 0 \right)$$

$$LDMEs (nonperturbative)$$
The same as those in X-section

e.g.
$$D_{g \to \psi}^{{}^{1}S_{0}^{(8)}}(z, 2m_{c}) = \frac{5\alpha_{s}^{2}(2m_{c})}{96m_{c}^{3}} \langle \mathcal{O}^{\psi}({}^{1}S_{0}^{(8)}) \rangle \left(3z - 2z^{2} + 2(1-z)\log(1-z)\right)$$

The channels we need: ${}^{1}S_{0}^{[8]}$, ${}^{3}S_{1}^{[1,8]}$, ${}^{3}P_{J}^{[8]}$ (power counting up to v^{4})



[Chao et al., 1201.2675]





[Bodwin et al., 1403.3612]





	$\langle \mathcal{O}^{J/\psi}({}^3S_1^{[1]}) \rangle$	$\langle \mathcal{O}^{J/\psi}({}^3S_1^{[8]}) \rangle$	$\langle \mathcal{O}^{J/\psi}({}^1S_0^{[8]}) \rangle$	$\langle \mathcal{O}^{J/\psi}({}^{3}P_{0}^{[8]})\rangle/m_{c}^{2}$
	$\times \text{GeV}^3$	$\times 10^{-2} { m GeV^3}$	$\times 10^{-2} {\rm GeV^3}$	$\times 10^{-2} \text{GeV}^3$
B & K	1.32 ± 0.20	0.224 ± 0.59	4.97 ± 0.44	-0.72 ± 0.88
Chao et al.	1.16 ± 0.20	0.30 ± 0.12	8.9 ± 0.98	0.56 ± 0.21
Bodwin et al.	1.32 ± 0.20	1.1 ± 1.0	9.9 ± 2.2	0.49 ± 0.44



$$\mathcal{G}_{i}^{H}(s, z, \mu) = \sum_{j} \int_{z}^{1} \frac{dz'}{z'} \mathcal{J}_{ij}(s, z', \mu) \mathcal{D}_{j}^{H}(z/z', m_{H}, \mu)$$
[M. Procura, I. Stewart, PRD, 81, 074009 (2010)]
[M. Baumgart, A. Leibovich, T. Mehen, I. Rothstein, Jet scale Hadron scale

Physical meaning:

JHEP, 11, 003 (2014)]





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Frag. at jet scale:





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$$D_q^H(z) \equiv \text{Tr}\langle 0|\frac{\not n}{2}\delta(\frac{p_-}{z} - \mathcal{P}_-)\psi|HX\rangle\langle HX|\psi|0\rangle$$



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Practical: resum $\log(\frac{m_H}{s})$

Factorize Everything

 $\sigma(H) = \sum \mathcal{H} \otimes \mathcal{S} \otimes \mathcal{G}_i^H \otimes \mathcal{J}_1 \otimes \mathcal{J}_2 \otimes \cdots$ i

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Factorize Everything

 $\sigma(H) = \sum \mathcal{H} \otimes \mathcal{S} \otimes \mathcal{G}_i^H \otimes \mathcal{J}_1 \otimes \mathcal{J}_2 \otimes \cdots$ *i* $\mathcal{G}_i^H(s,z,\mu) = \sum_j \int_z^1 \frac{dz'}{z'} \mathcal{J}_{ij}(s,z',\mu) D_j^H(z/z',m_H,\mu)$ **B** meson, or J/ψ FF.



B Fragmentation Function

Power Model FF $D_b^B(z) = N z^{\alpha} (1-z)^{\beta}$

•
$$\alpha = 16.87, \beta = 2.02$$

2 parameters,
N fixed by:
 $\int_{0}^{1} dzzD(z) = 1$
• Fit at m_B scale
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[Kniehl and Kramer, 0705.4392]

Analytic & Monte Carlo (B in jet)



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$$z \equiv \frac{E_B}{E_J} \qquad \tau_0 \propto m_J$$

FF: Pythia8 & Analytic: Power Model Herwig: Cluster hadronization

[R. Bain, L. Dai, A. Hornig, A. Leibovich, Y. Makris, T. Mehen, JHEP, 06, 121 (2016)]



J/psi Fragmentation Functions



J/ ψ Production with Pythia

Pythia default:

 $Q\bar{Q}({}^3S_1^{[8]},\cdots)$ Produced in hard process

Pythia Shower with $2P_{qq}$



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Gluon Fragmentation Improved Pythia (GFIP):

Initiating gluon produced in hard process

Pythia Shower (to *2m_c*, Hadronization off)

Manual convolve with FFs







[R. Bain, L. Dai, A. Hornig, A. Leibovich, Y. Makris, T. Mehen, JHEP, 06, 121 (2016)]

Analytic is Eqiv. to GFIP

FJF: (Fragmenting Jet Function) $G_g^{J/\psi} = \sum_i \mathcal{J}_{gi} \otimes D_i^{J/\psi}$	GFIP: (Gluon FF Improve Pythis)
$\frac{d\mathcal{J}(\mu, m_J)}{d\log\mu} = \gamma_J \otimes \mathcal{J}(\mu, m_J)$ $\frac{dD_i(\mu, m_{J/\psi})}{d\log\mu} = P_{ij} \otimes D_j(\mu, m_{J/\psi})$	Pythia Shower (to 2m _c , Hadronization off) Manual convolve with FFs

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RG Running	SHOWER And J/W Constant J/W Constant J/W



 J/ψ in jets @ LHCb

Cuts:

$$p_T(ext{jet}) > 20 \ ext{GeV}$$

 $2 > \eta(ext{jet}) > 2.5$
 $p_T(\mu) > 0.5 \ ext{GeV}$
 $p(\mu) > 5 \ ext{GeV}$
 $4.5 > \eta(\mu) > 2$

Similar to e⁻e⁺ \rightarrow J/ ψ : J/ ψ too hard in jets



[LHCb, Phys. Rev. Lett. 118, 192001 (2017)]

Analytic & GFIP (J/psi in jets @ LHCb)



Comparison with LHCb Data



[R. Bain, L. Dai, A. Leibovich, Y. Makris, T. Mehen, Phys. Rev. Lett. 119, 032002 (2017)]

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[R. Bain, L. Dai, A. Leibovich, Y. Makris, T. Mehen, Phys. Rev. Lett.119, 032002 (2017)]

- Error band only from LDMEs (B&K narrower, used more data)
- Both FJF & GFIP work better than default Pythia;

- High pT LDMEs give better description of data than the global fit ones

Preliminary

 $25 \text{GeV} < p_{T,\text{jet}} < 35 \text{GeV}$ $3 \text{GeV} < p_{T,J/\psi} < 35 \text{GeV}$





[L. Dai, A. Leibovich, Y. Makris, T. Mehen, work in progress]

CMS-PAS-HIN-18-012 25

Preliminary







[L. Dai, A. Leibovich, Y. Makris, T. Mehen, work in progress]

CMS-PAS-HIN-18-012 *26*

E(jet) distribution

Normalized cross sections for the LHC (*i* denotes production channels):





[L. Dai, P. Shrivastava, PRD 96, 036020 (2017)] 28



[L. Dai, P. Shrivastava, PRD 96, 036020 (2017)] 28



Decreasing at z > 0.5 (high pT), consistent with the dominance of ${}^{1}S_{0}{}^{[8]}$ at high pT



[L. Dai, P. Shrivastava, PRD 96, 036020 (2017)] 29



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Generalize to pp collision

 Compare with LHC data, high pT extractions give better description of data than the global fit. And both are better than the default PYTHIA

 Study E(jet) distributions (the potential to distinguish production mechanisms and LDME extractions)

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