



# Experimental Methods in Particle Physics

---

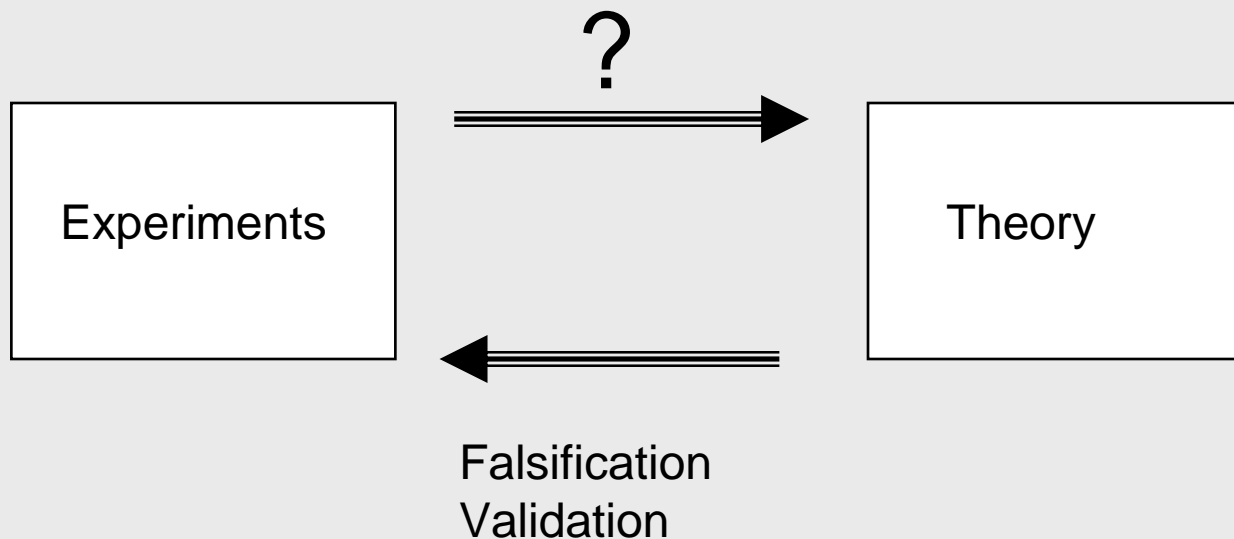
How to get a  
meaningful result out  
of a particle physics  
detector

Emmanuel Tsesmelis  
CERN / University of Oxford  
APPEAL  
3 July 2010

# Design of a Particle Physics Experiment

---

- How physicists go from the basic ideas of measuring some quantities to analyzing physics events from experiments.



# Theory and Experiments

---

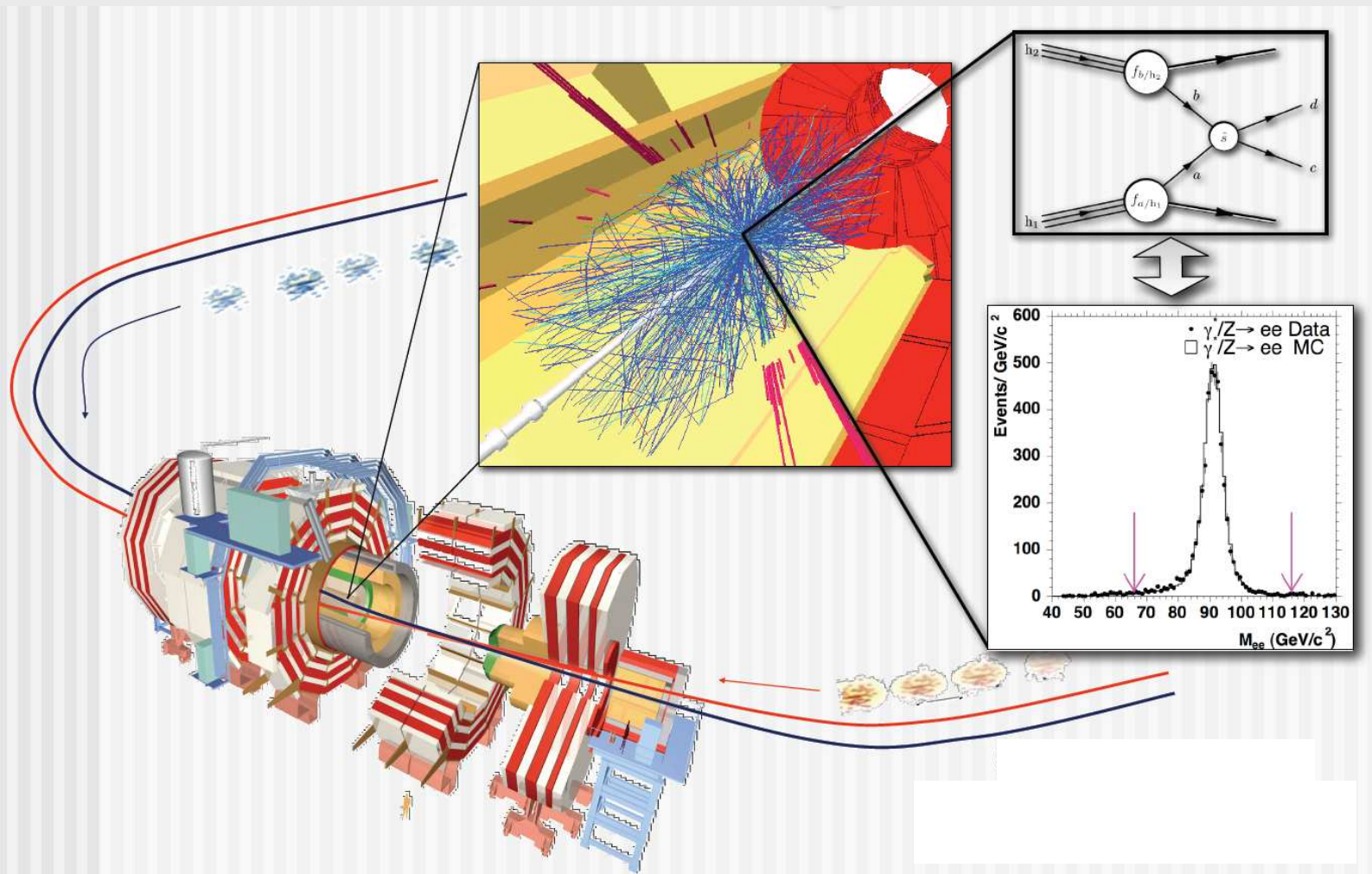
- Expt: Particles have masses: .....why ?
- Theo: Mass is given by the interaction with the Higgs field.
- Expt: Find the Higgs Boson
  
- Expt: There are 3 Forces: .....why ?
- Theo: Supersymmetry unifies the forces
- Expt: Find the signals of Supersymmetry

# Two Classes of Experiments

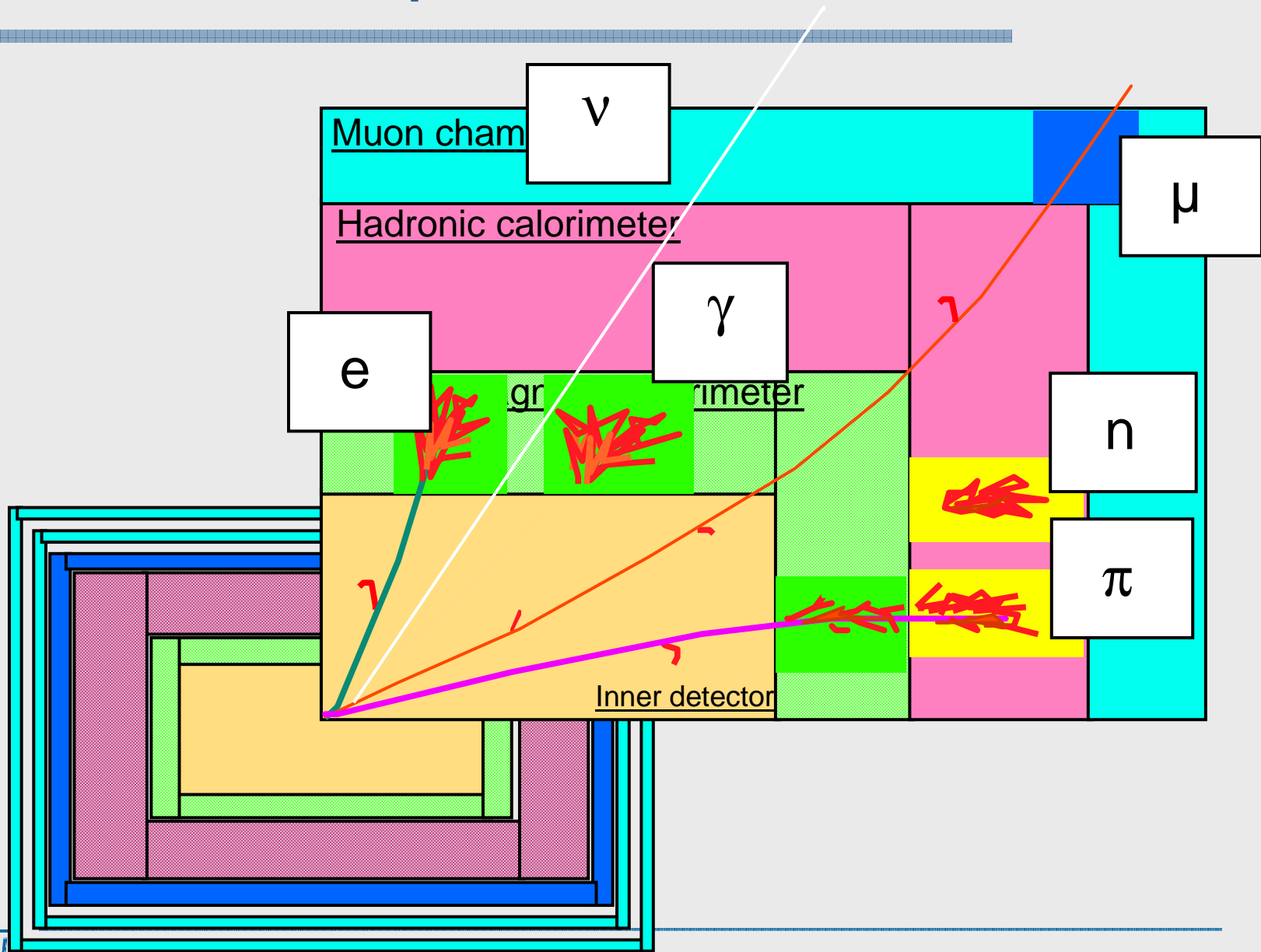
---

1. **E**xperiments designed to measure typically “one quantity”, exploring the “very rare” or “very precise”. *They need high intensity beams and/or very high precision.*
2. **E**xperiments that explore the high energy frontier. *They are typically multi-purpose experiments and they also need high intensity beams.*

# From Raw Data to Physics Results

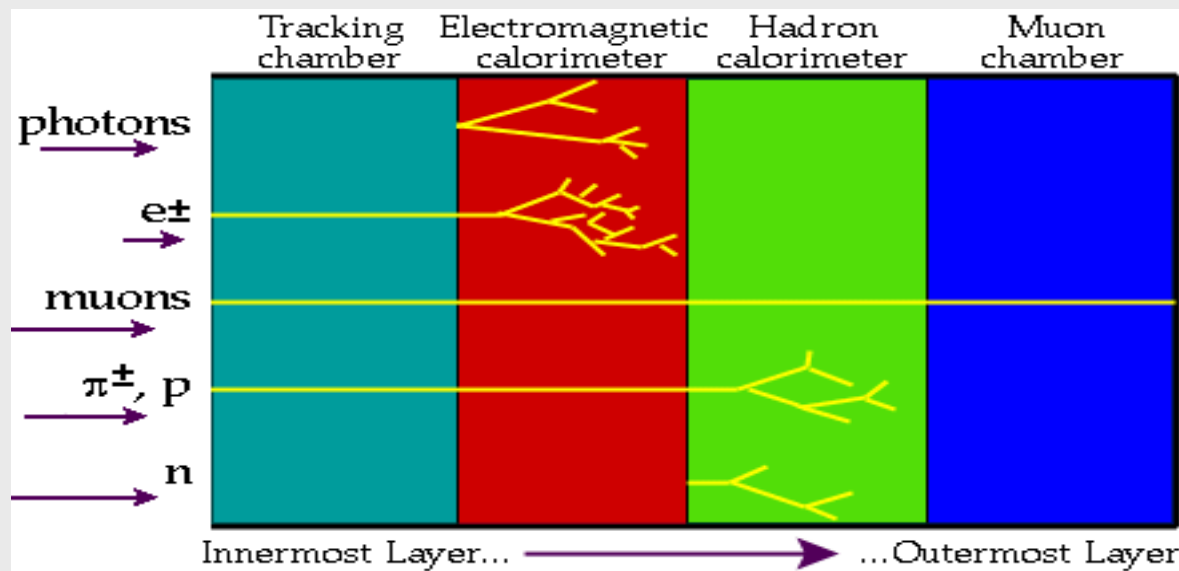


# General Principle



# General Principle

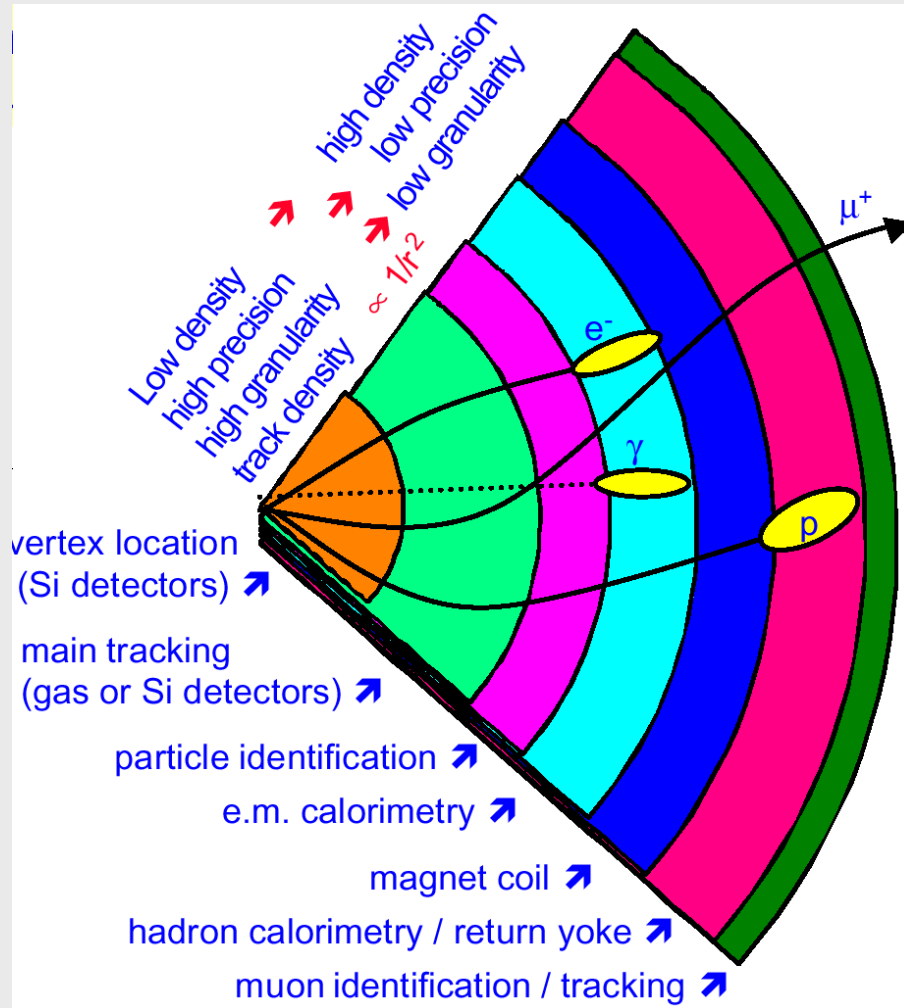
Visible particles are measured by the various sub-detectors and identified from their characteristic pattern.



The parameters of the quarks are reconstructed from the hadronic jets.

The flavor of the quark is determined by reconstructing the hadronic decays of heavy mesons or detecting their detached decay vertex.

# Arrangement of Detectors

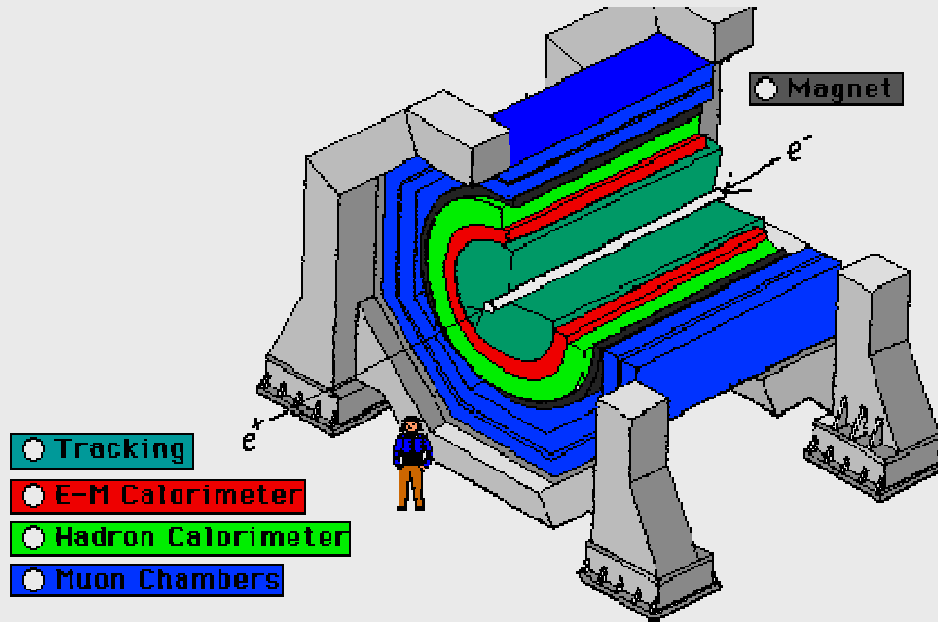


Various detectors and combination of information can provide particle identification.

p versus EM energy for electrons; EM/HAD provides additional information, so do the muon detectors, EM response without tracks indicate a photon; secondary vertices identify b,c, $\tau$ ; isolation cuts help to identify leptons.



# General Principle



Collider detectors are similar since they must perform in sequence the same basic measurements.

The dimensions of the detector are driven by the required resolution.

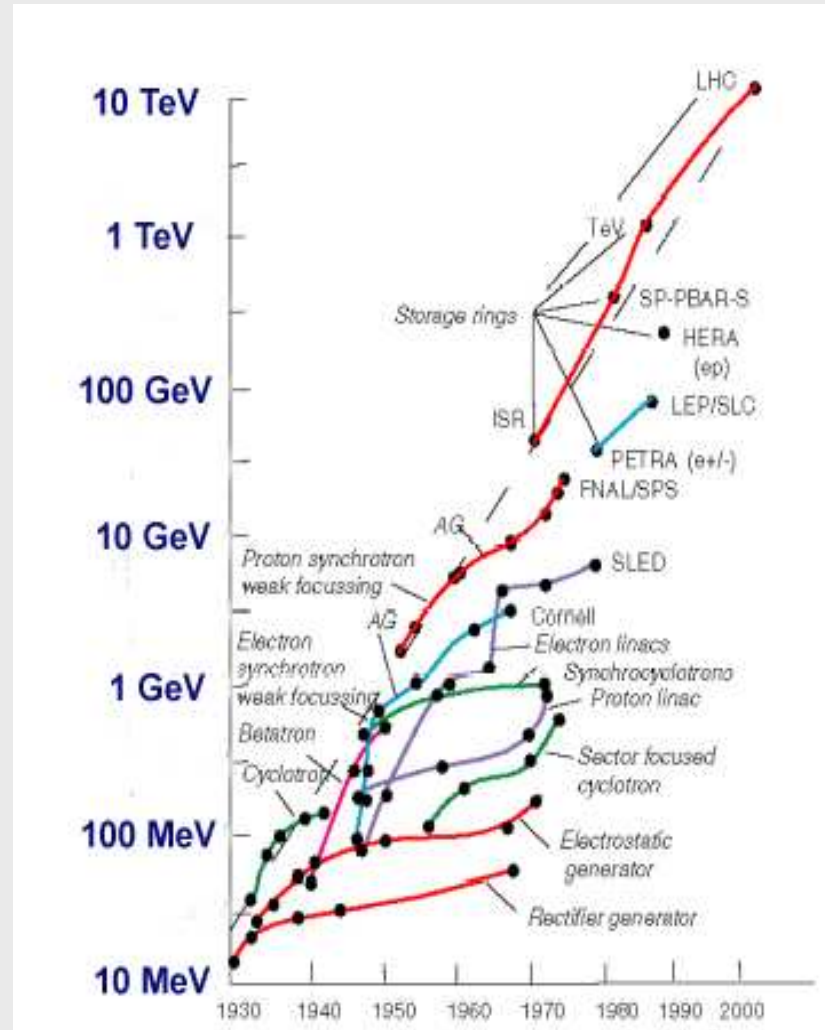
The calorimeter thickness changes only with the logarithm of the energy: for this reason the dimension of the detectors change only slightly with the energy.

# Quest for Energy

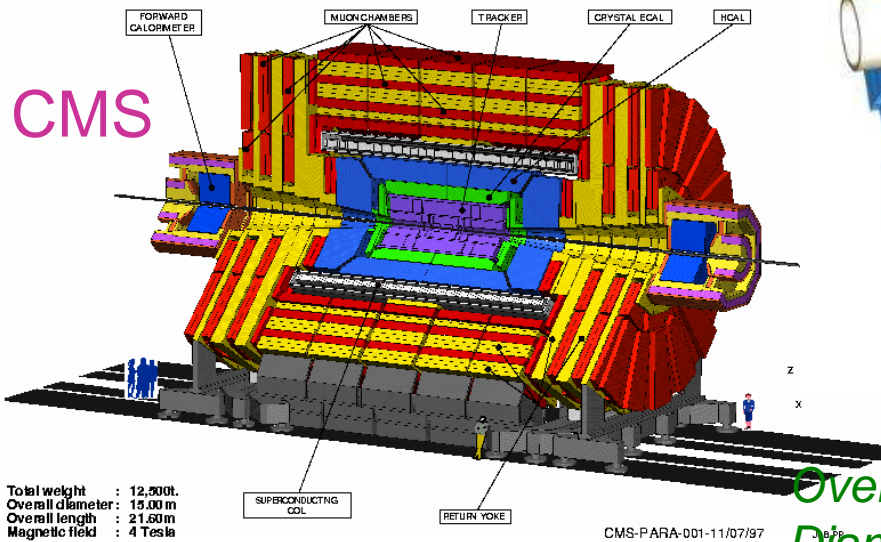
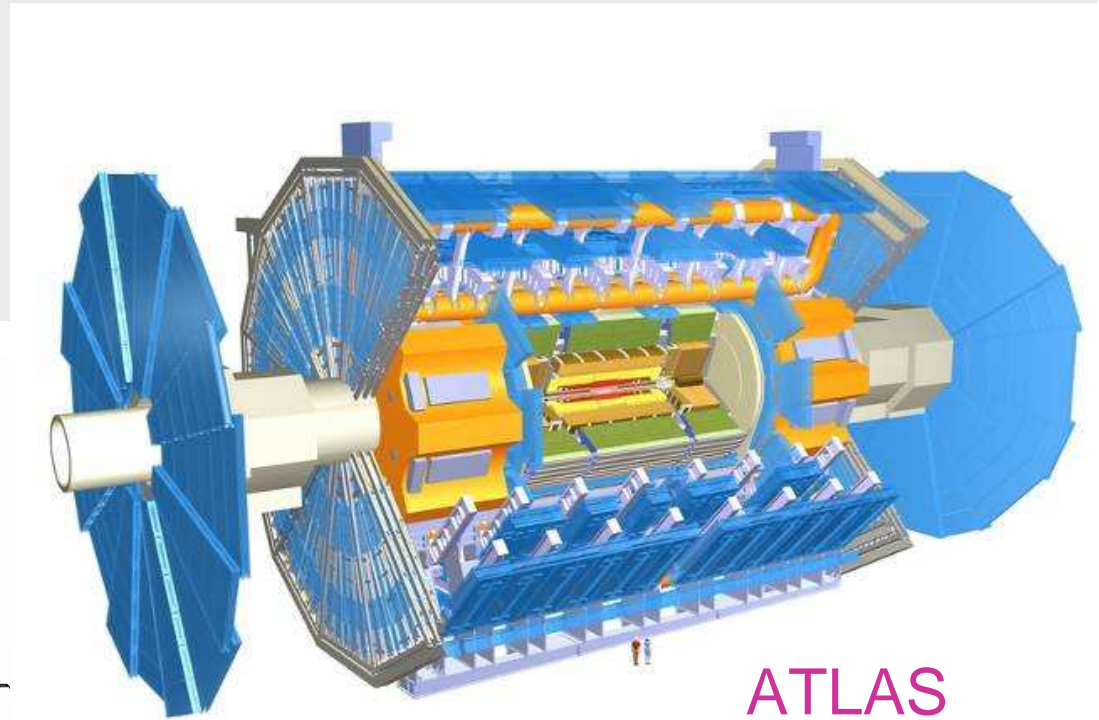
$$\lambda = \frac{hc}{2\pi E}$$

$$\lambda(fm) \approx \frac{0.2}{E(GeV)}$$

Increasing energy allows experiments to probe Nature at smaller distances with the possibility of crossing the threshold of observing new phenomena



# ATLAS and CMS



Total weight : 12,500t  
 Overall diameter : 15.00 m  
 Overall length : 21.60 m  
 Magnetic field : 4 Tesla

CMS-PARA-001-11/07/97

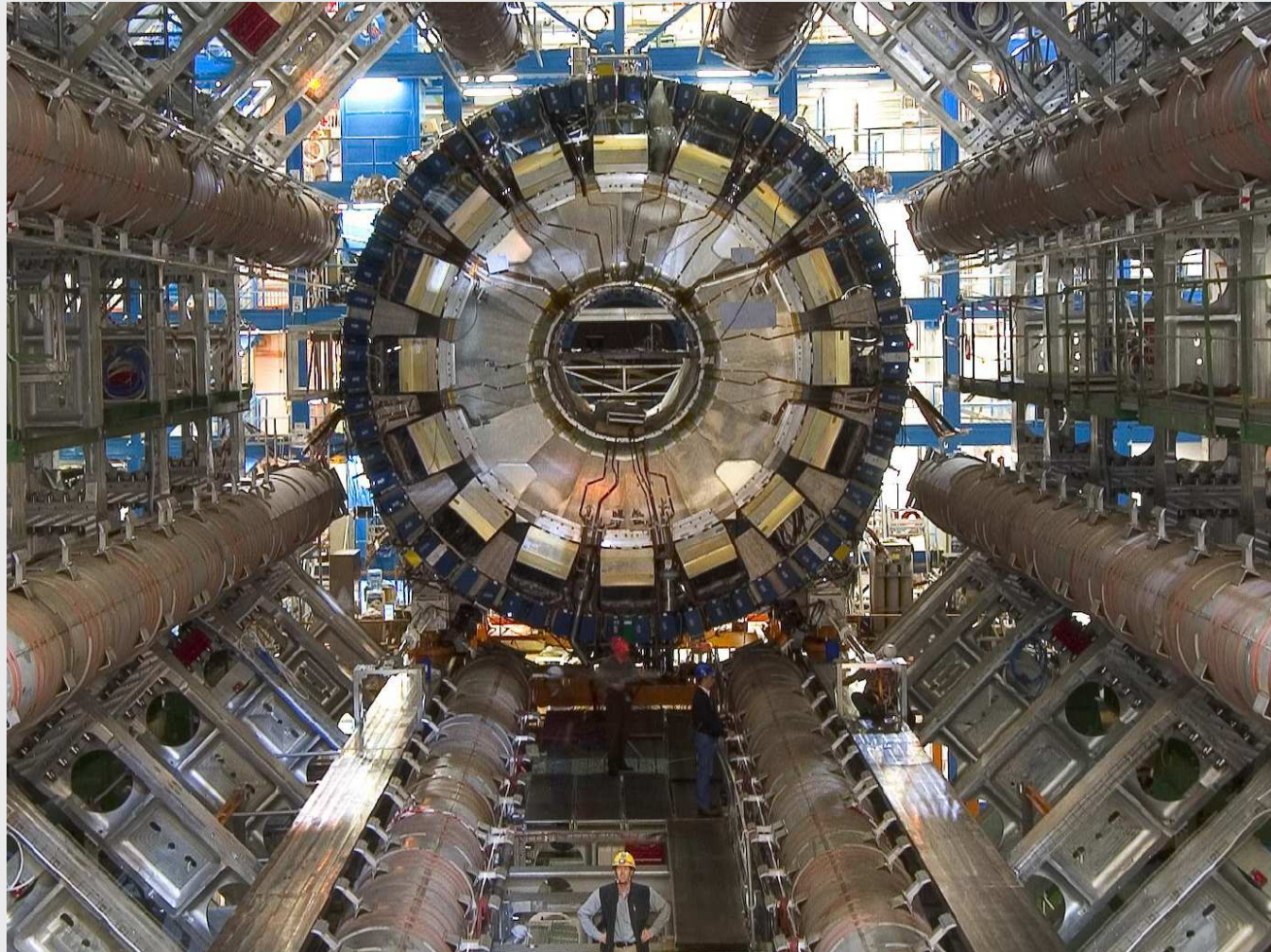
ATLAS

Overall weight (tons)  
 Diameter (m)  
 Length (m)  
 Solenoid field (T)

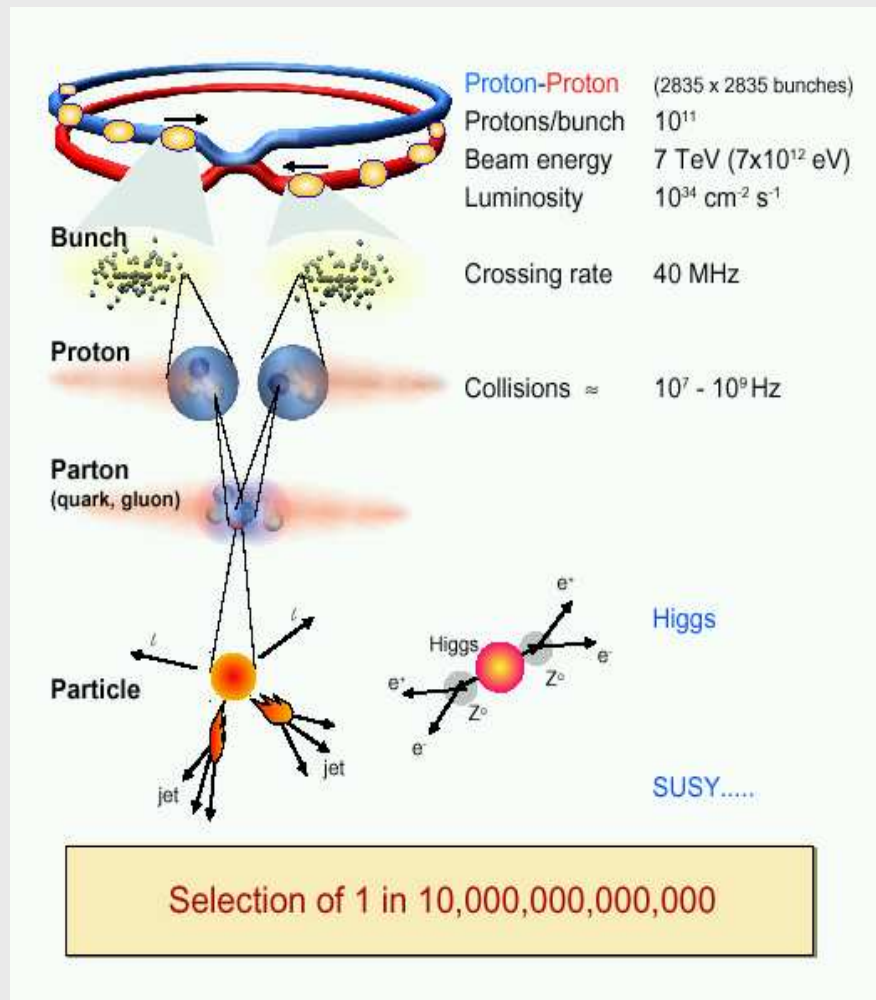
	<u>ATLAS</u>	<u>CMS</u>
Overall weight (tons)	7000	12500
Diameter (m)	22	15
Length (m)	46	22
Solenoid field (T)	2	4



# ATLAS



# Collisions at the LHC



# Higgs at the LHC: The Challenge

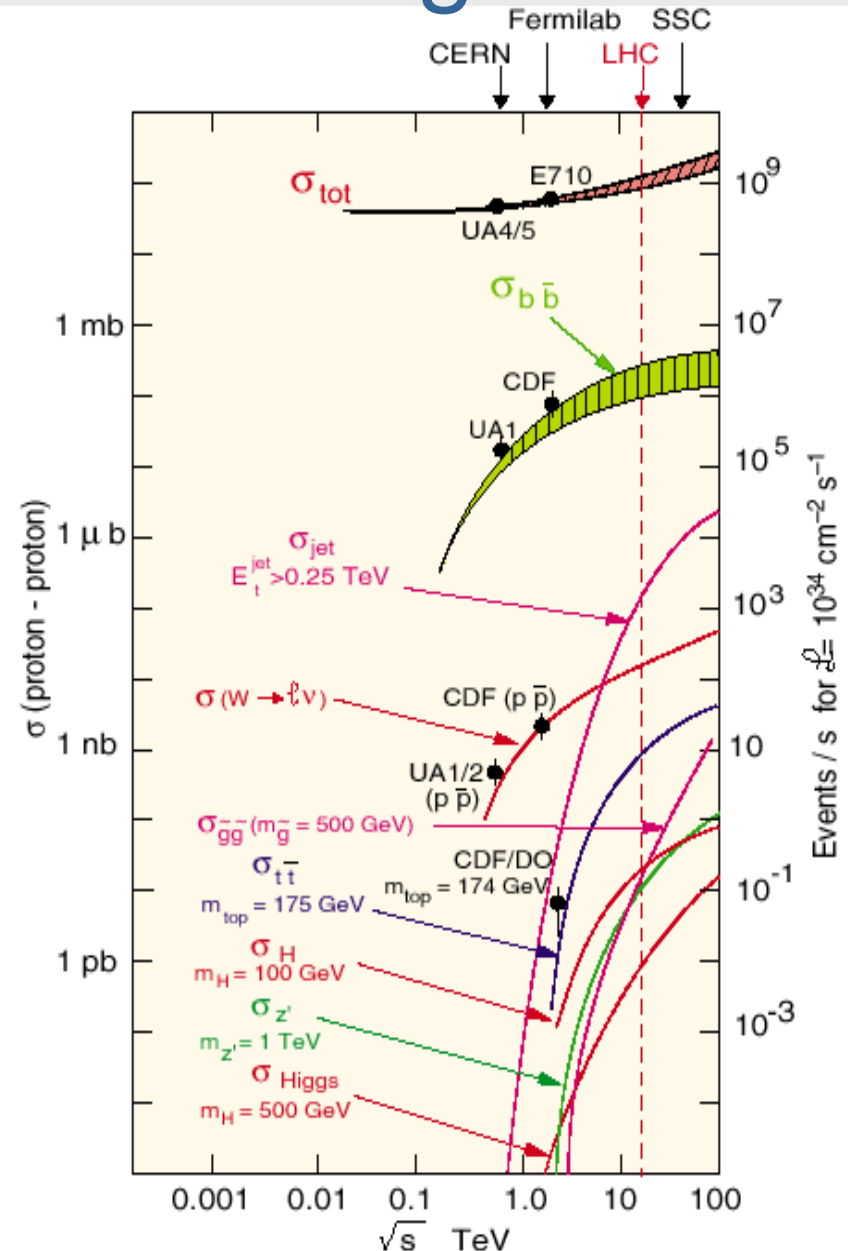
Small cross-sections  
need highest luminosity

→  $L = 10^{34-35} \text{ cm}^{-2}\text{s}^{-1}$

Event rates for various physics channels:

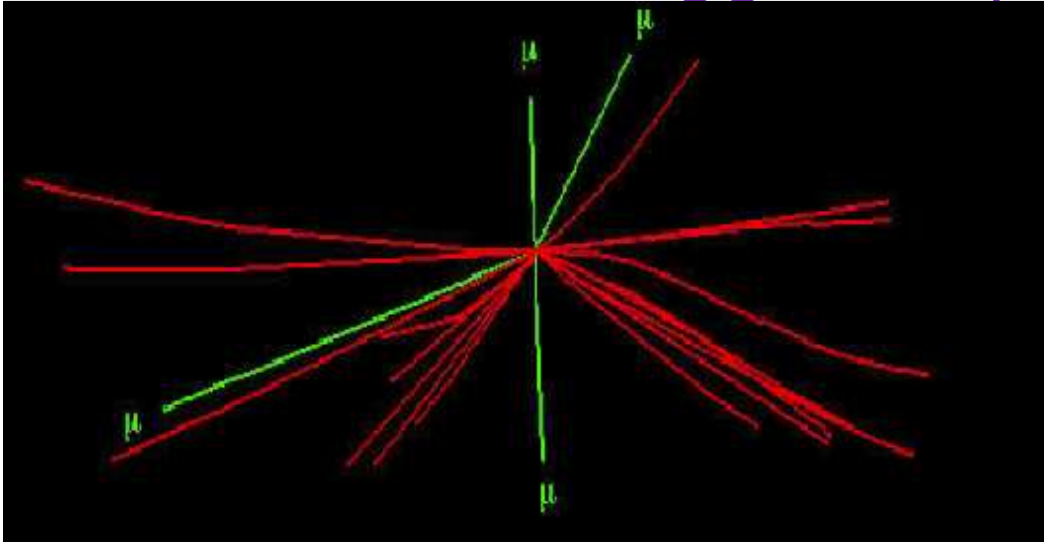
- Inelastic :  $10^9 \text{ Hz}$
  - $W \rightarrow \ell\nu$  :  $10^2 \text{ Hz}$
  - $t\bar{t}$  production :  $10^1 \text{ Hz}$
  - Higgs ( $m=100 \text{ GeV}$ ) :  $10^{-1} \text{ Hz}$
  - Higgs ( $m=600 \text{ GeV}$ ) :  $10^{-2} \text{ Hz}$
- (including branching ratios:  $\sim 10^{-3}$ )

Selection power for  
Higgs discovery  $\approx 10^{14-15}$

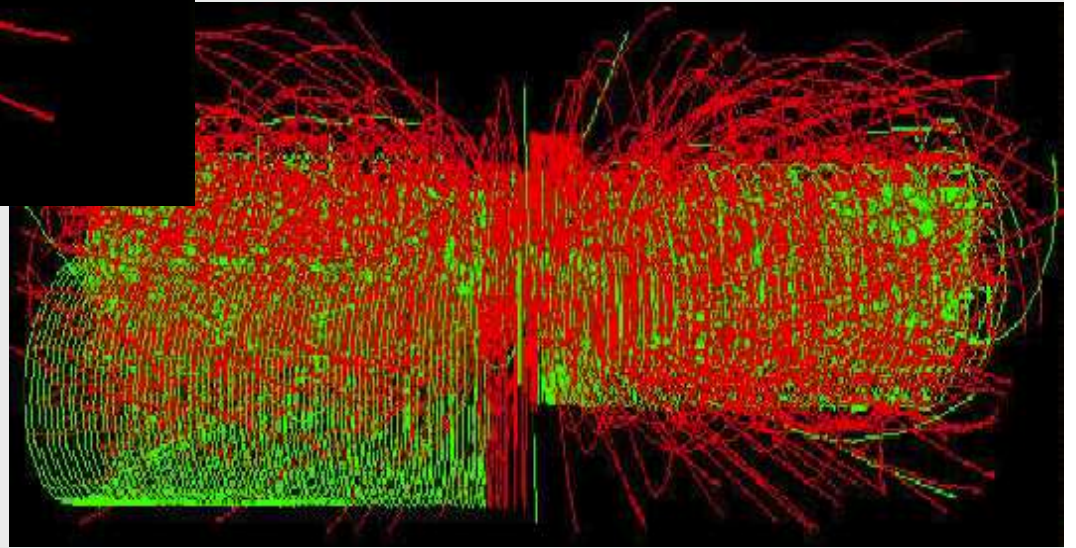


# Higgs at the LHC: The Challenge

How to extract Higgs  $\rightarrow 4\mu$



Within 20 overlapping events

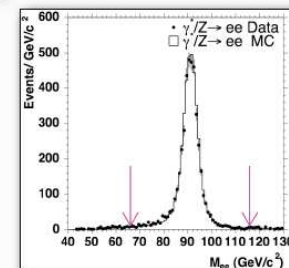


Without knowing really where to look for!

# Data Analysis Chain



- Have to collect data from many channels on many sub-detectors (millions)
- Decide to read out everything or throw event away (Trigger)
- Build the event (put info together)
- Store the data
- Analyze them
  - reconstruction, user analysis algorithms, data volume reduction
- do the same with a simulation
  - correct data for detector effects
- Compare data and theory

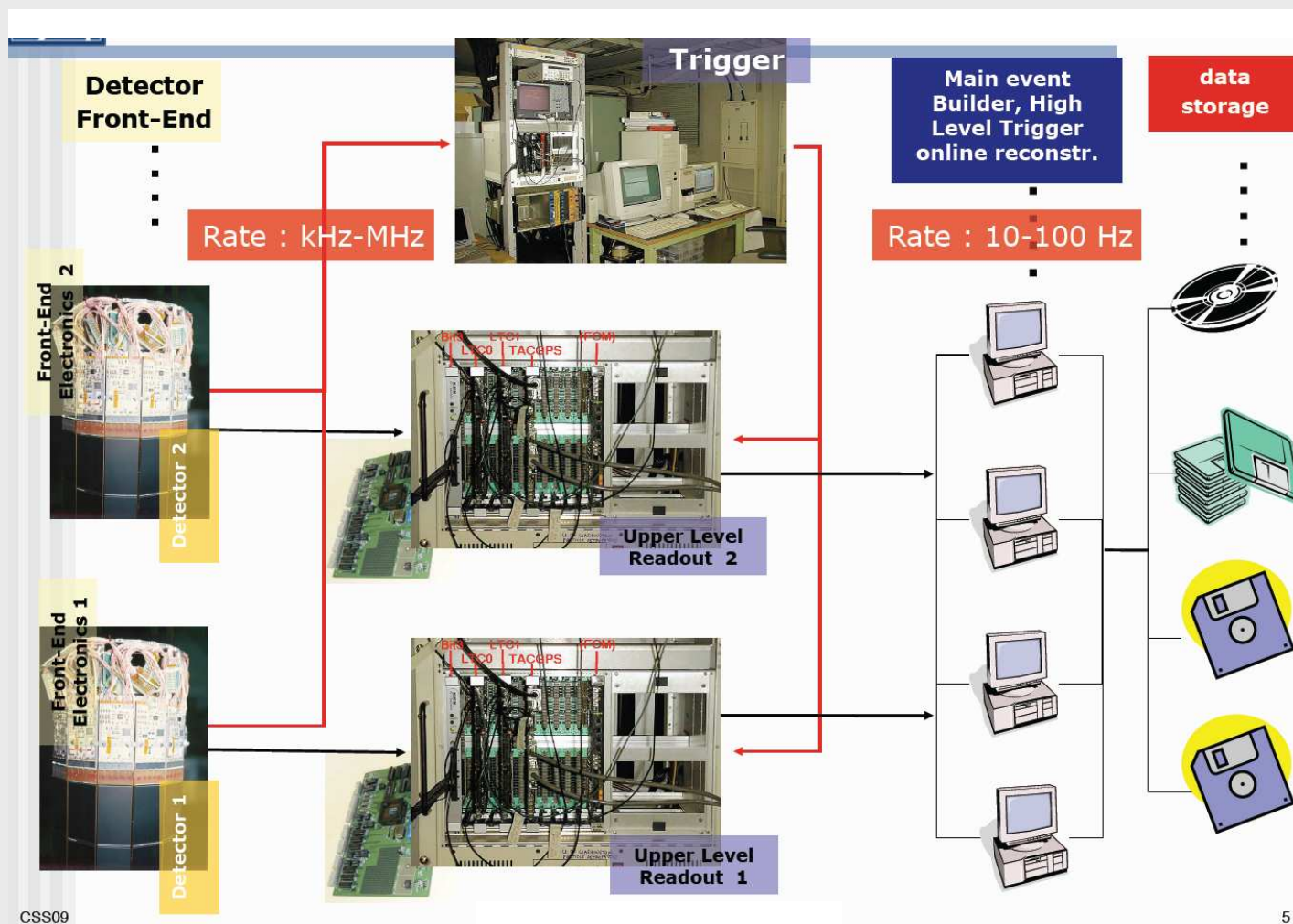


CSS09

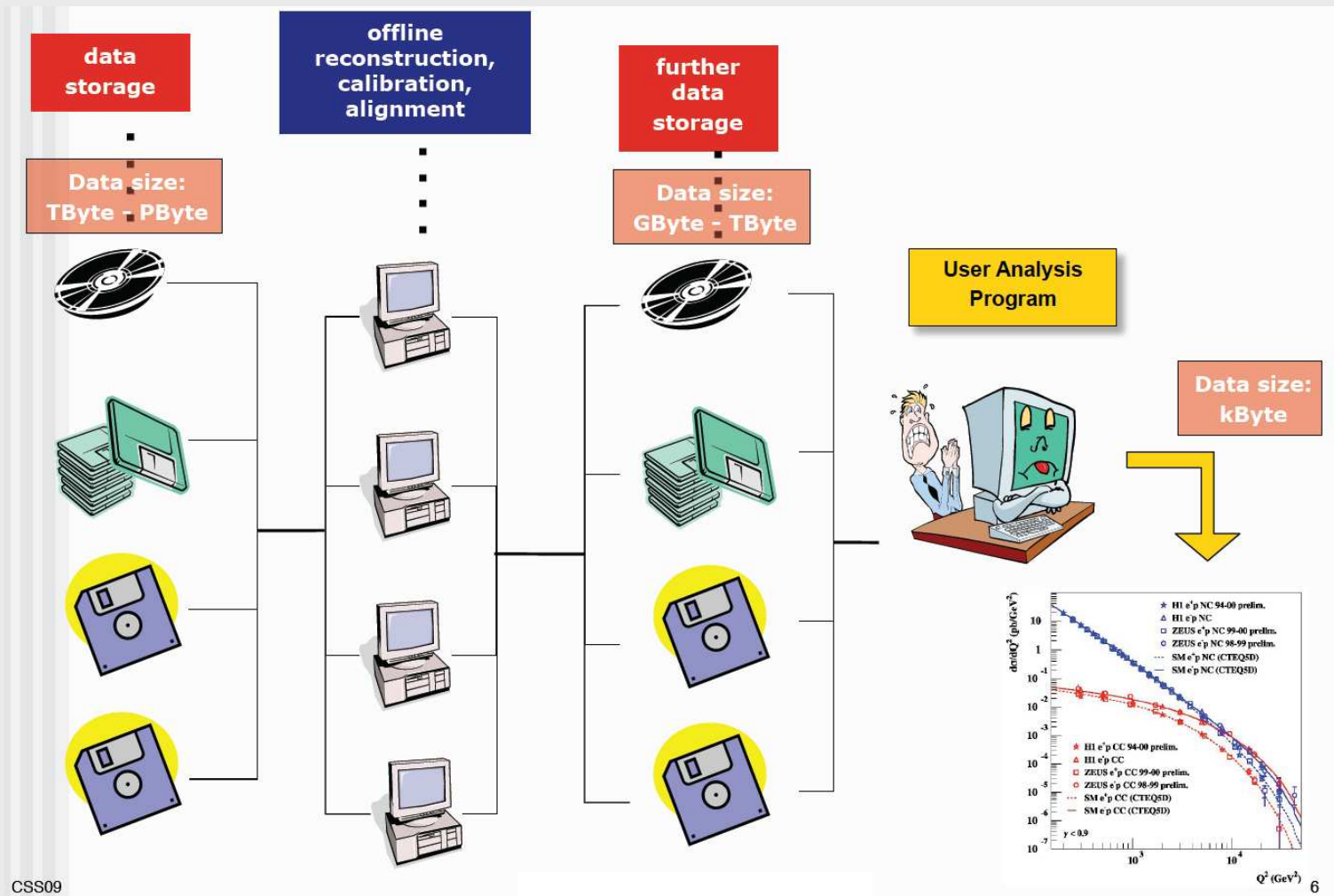
4



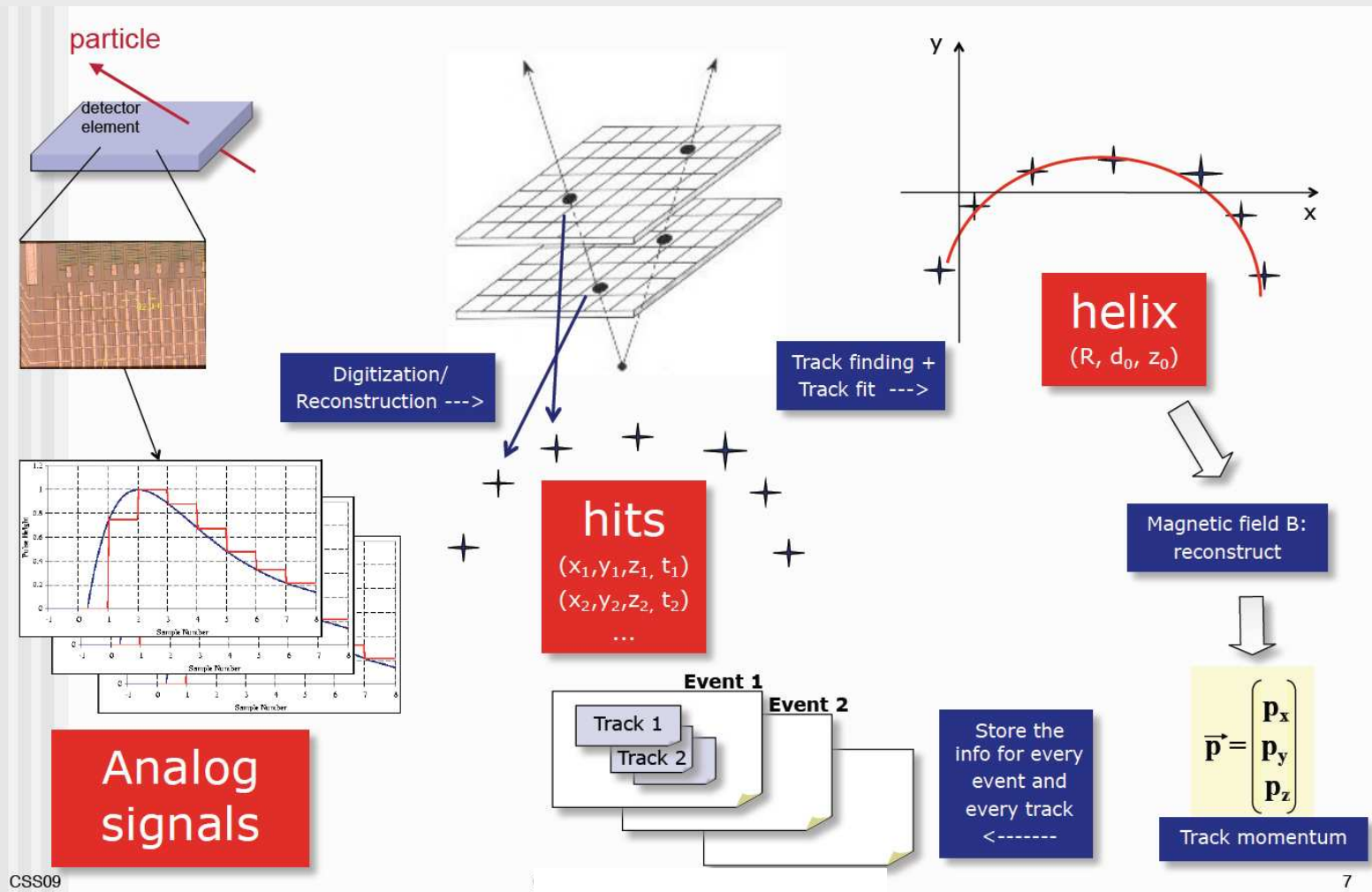
# Data Chain



# Offline Analysis Chain



# Data Reduction



CSS09

7

# High-level Storage

- Data are stored sequentially in files...

## Event 1

**Nch** (charged tracks) :  
2

**Pcha**  
(Momentum of each track):

```
{ {"-7.65698", "42.9725", "14.3404"},  
  {" 7.54101", "-42.1729", "-14.0108"} }
```

px py pz

**Qcha**  
(Charge of each track):  
{-1,1}

## Event 2

**Nch** (charged tracks) :  
3

**Pcha**  
(Momentum of each track):

```
{ {"-12.9305", "12.2713", "40.5615"},  
  {" 12.2469", "-11.606", "-38.7182"},  
  {"0.143435", "-0.143435", "-0.497444"} }
```

px py pz

**Qcha**  
(Charge of each track):  
{-1,1,-1}

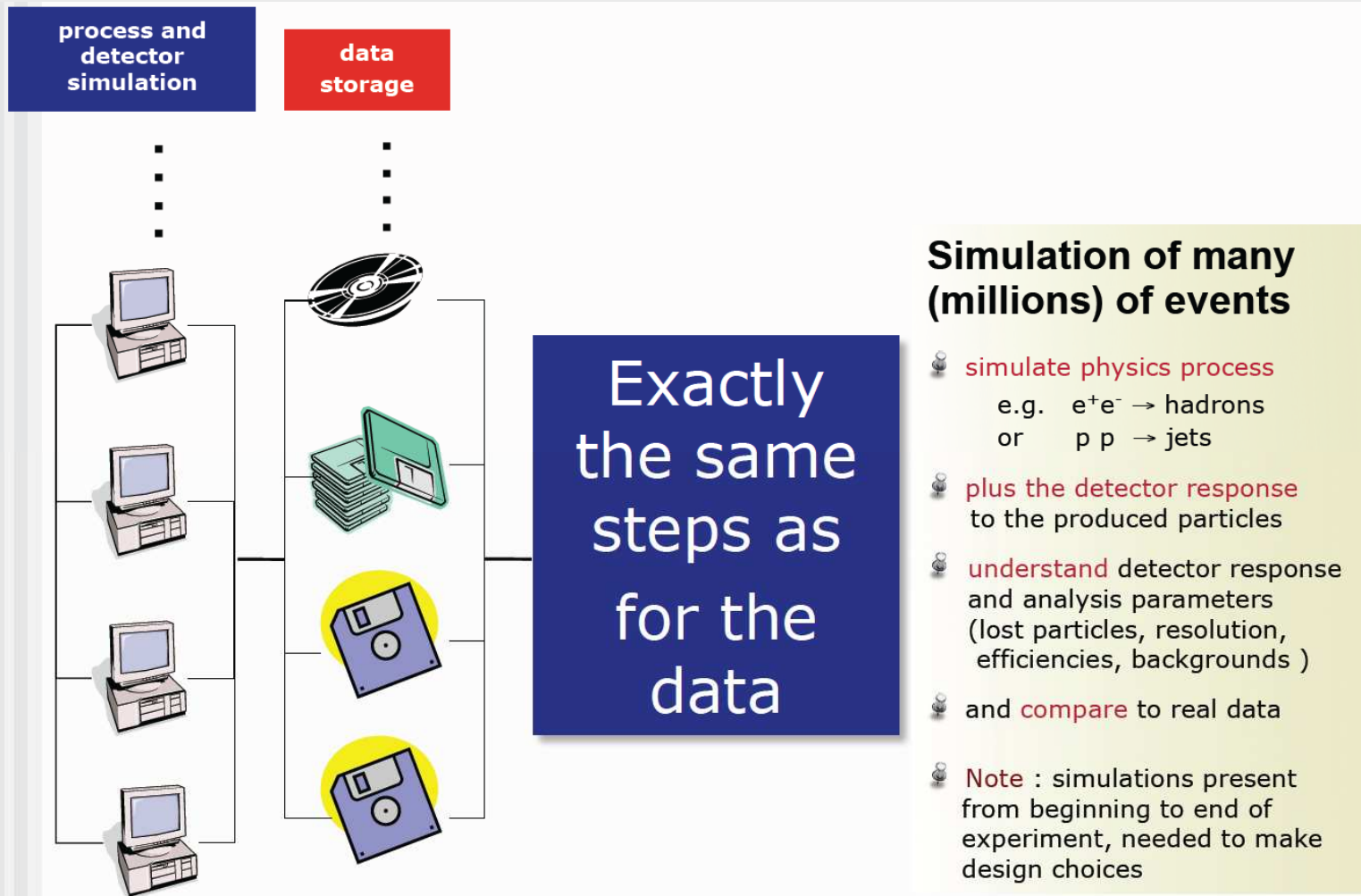
File A

CSS09

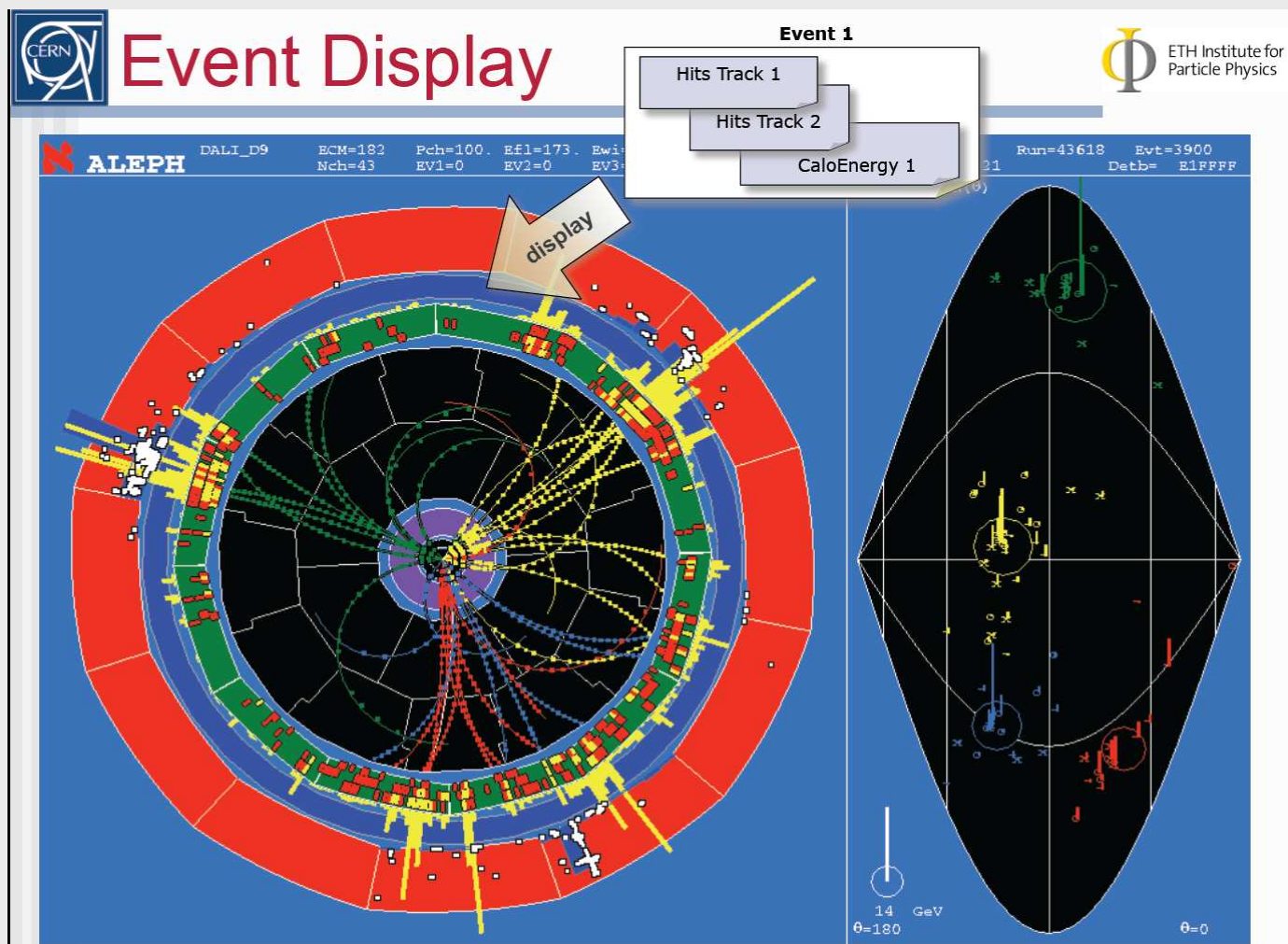
8



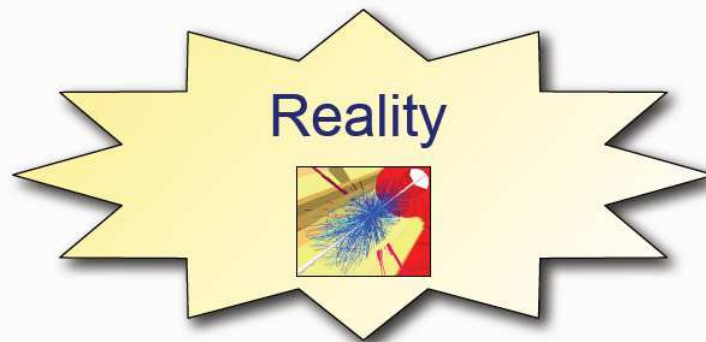
# Simulation



# Event Display



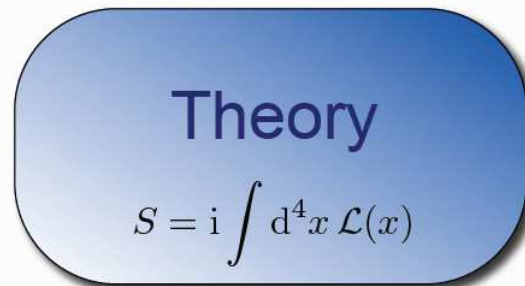
# Our Task



We use experiments to inquire about what “reality” (nature) does



We intend to fill this gap



*The goal is to understand in the most general; that's usually also the simplest.*

- A. Eddington

# Theory

$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{4} \mathbf{W}_{\mu\nu} \cdot \mathbf{W}^{\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} & \left\{ \begin{array}{l} W^\pm, Z, \gamma \text{ kinetic} \\ \text{energies and} \\ \text{self-interactions} \end{array} \right. \\
 & + \bar{L} \gamma^\mu (i\partial_\mu - g \frac{1}{2} \boldsymbol{\tau} \cdot \mathbf{W}_\mu - g' \frac{Y}{2} B_\mu) L & \left\{ \begin{array}{l} \text{lepton and quark} \\ \text{kinetic energies} \\ \text{and their} \\ \text{interactions with} \\ W^\pm, Z, \gamma \end{array} \right. \\
 & + \bar{R} \gamma^\mu (i\partial_\mu - g' \frac{Y}{2} B_\mu) R & \\
 & + \left| (i\partial_\mu - g \frac{1}{2} \boldsymbol{\tau} \cdot \mathbf{W}_\mu - g' \frac{Y}{2} B_\mu) \phi \right|^2 & \left\{ \begin{array}{l} W^\pm, Z, \gamma \text{ and} \\ \text{Higgs masses} \\ \text{and couplings} \end{array} \right. \\
 & - V(\phi) & \\
 & - (G_1 \bar{L} \phi R + G_2 \bar{L} \phi_c R + h.c.) & \left\{ \begin{array}{l} \text{lepton and quark} \\ \text{masses and} \\ \text{coupling to Higgs} \end{array} \right.
 \end{aligned}$$

$L$  ... left-handed fermion ( $l$  or  $q$ ) doublet  
 $R$  ... right-handed fermion singlet

$\mathcal{L}$  from QCD:

$$\mathcal{L} = \underbrace{\bar{q} (i\gamma^\mu \partial_\mu - m) q}_{E_{\text{kin}}(q)} - g \underbrace{(\bar{q} \gamma^\mu T_a q) G_\mu^a}_{\text{Interaction } q, g} - \frac{1}{4} \underbrace{G_{\mu\nu}^a G_a^{\mu\nu}}_{E_{\text{kin}}(g) \text{ includes self-interaction between gluons}}$$

the Standard Model

has parameters

coupling constants

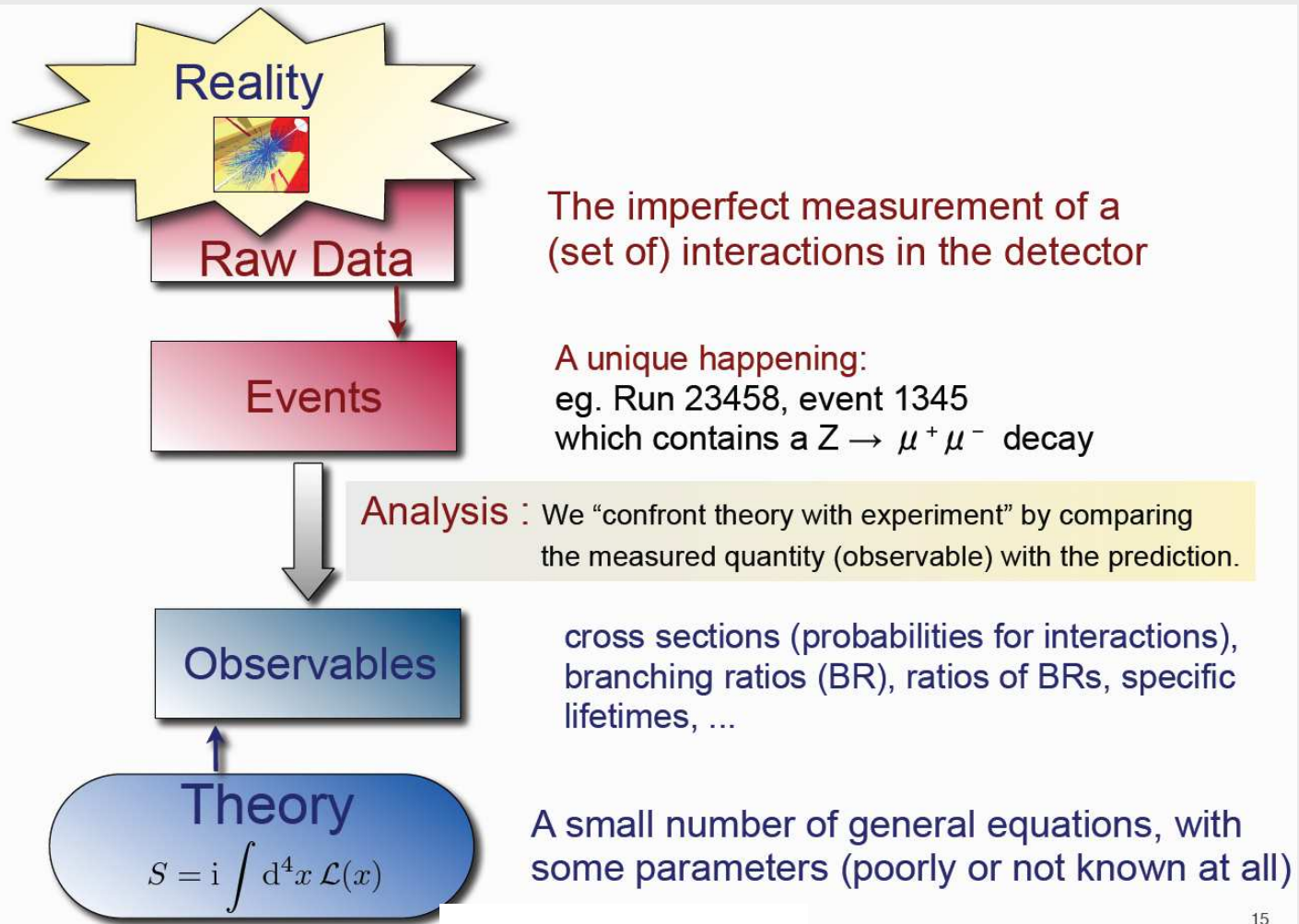
masses

predicts:  
cross sections,  
branching ratios,  
lifetimes, ...





# Making the Connection

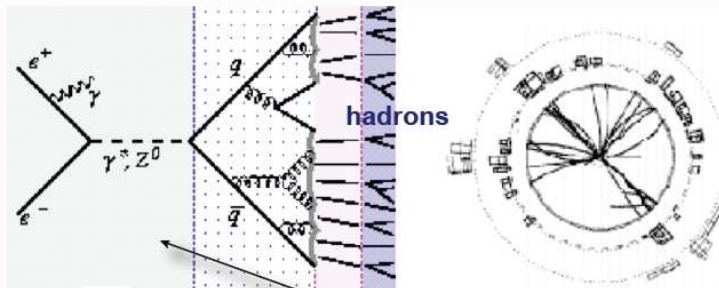


CSS09

15

# A Simple Counting Experiment

## Measurement of $e^+e^-$ annihilation into hadrons and muons:



### Hadronic final state

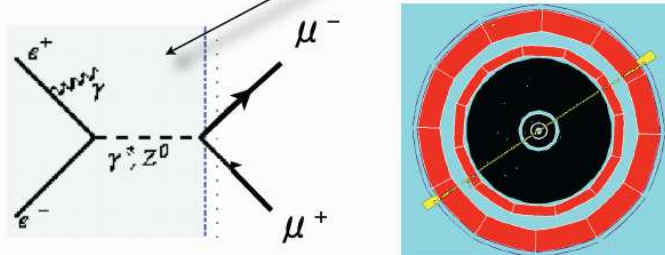
- many charged tracks (  $> \sim 10$  )
- sum of energy deposits in calorimeters not too far from centre-of-mass energy

sum over all quark flavours, which can be produced at a certain  $e^+e^-$  centre-of-mass energy  $E_{CM}$ , e.g. d, u, s, c, b, t

$$R := \frac{\sigma(e^+e^- \rightarrow q_f \bar{q}_f)}{\sigma(e^+e^- \rightarrow \mu^+ \mu^-)} = N_c \sum_f z_f^2$$

Number of colours

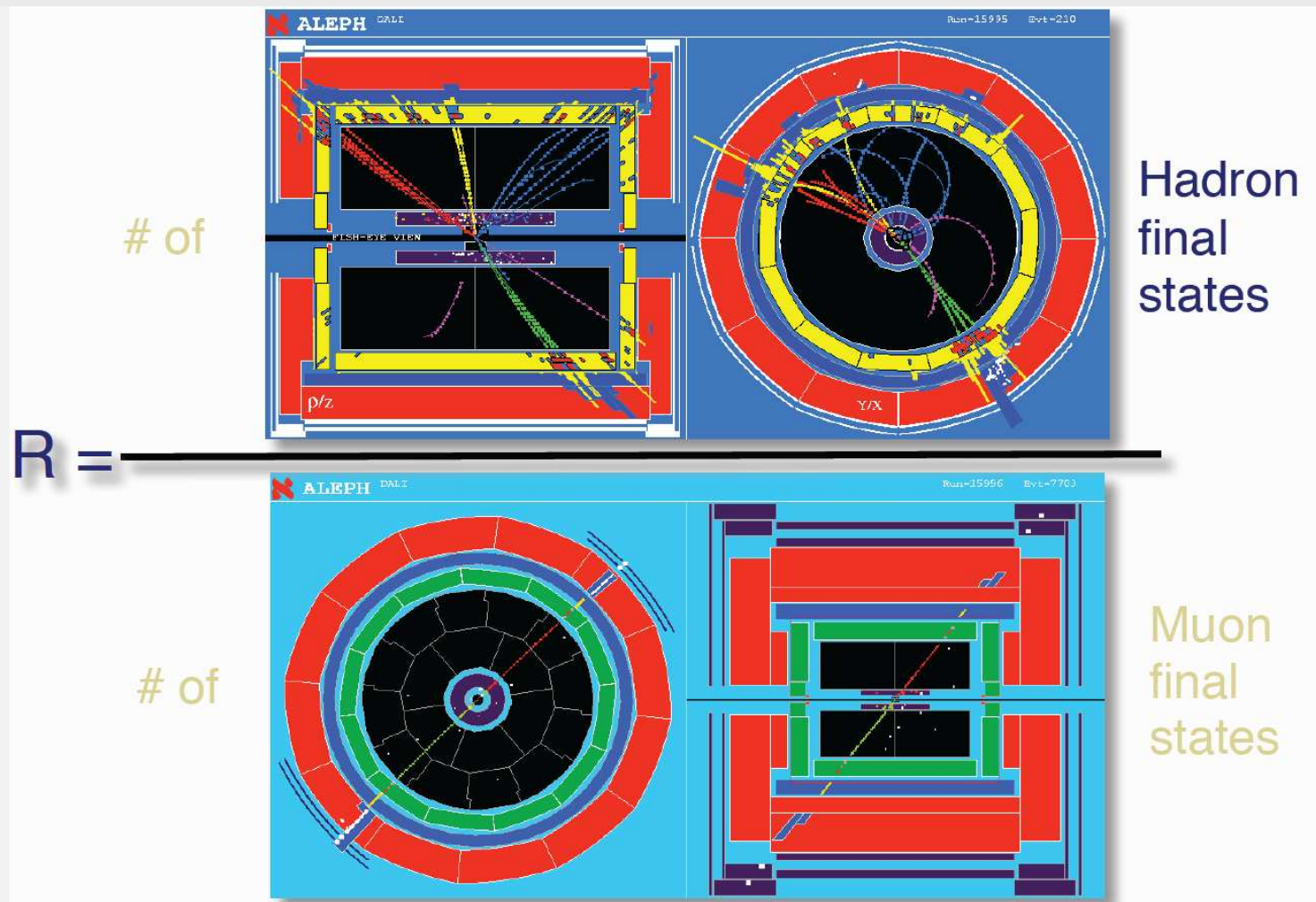
electric charges of quarks, in units of electron charge

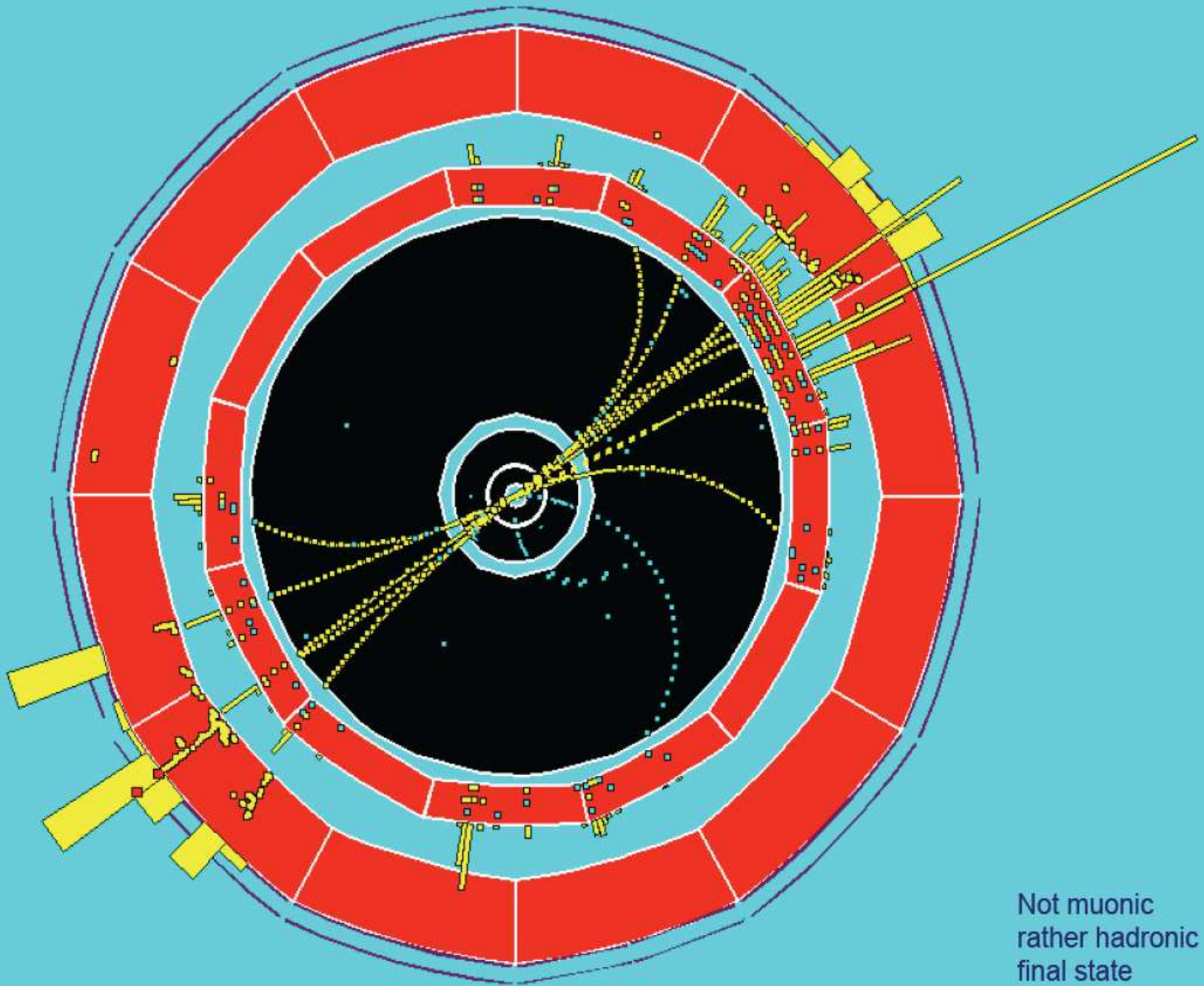


### Muonic final state

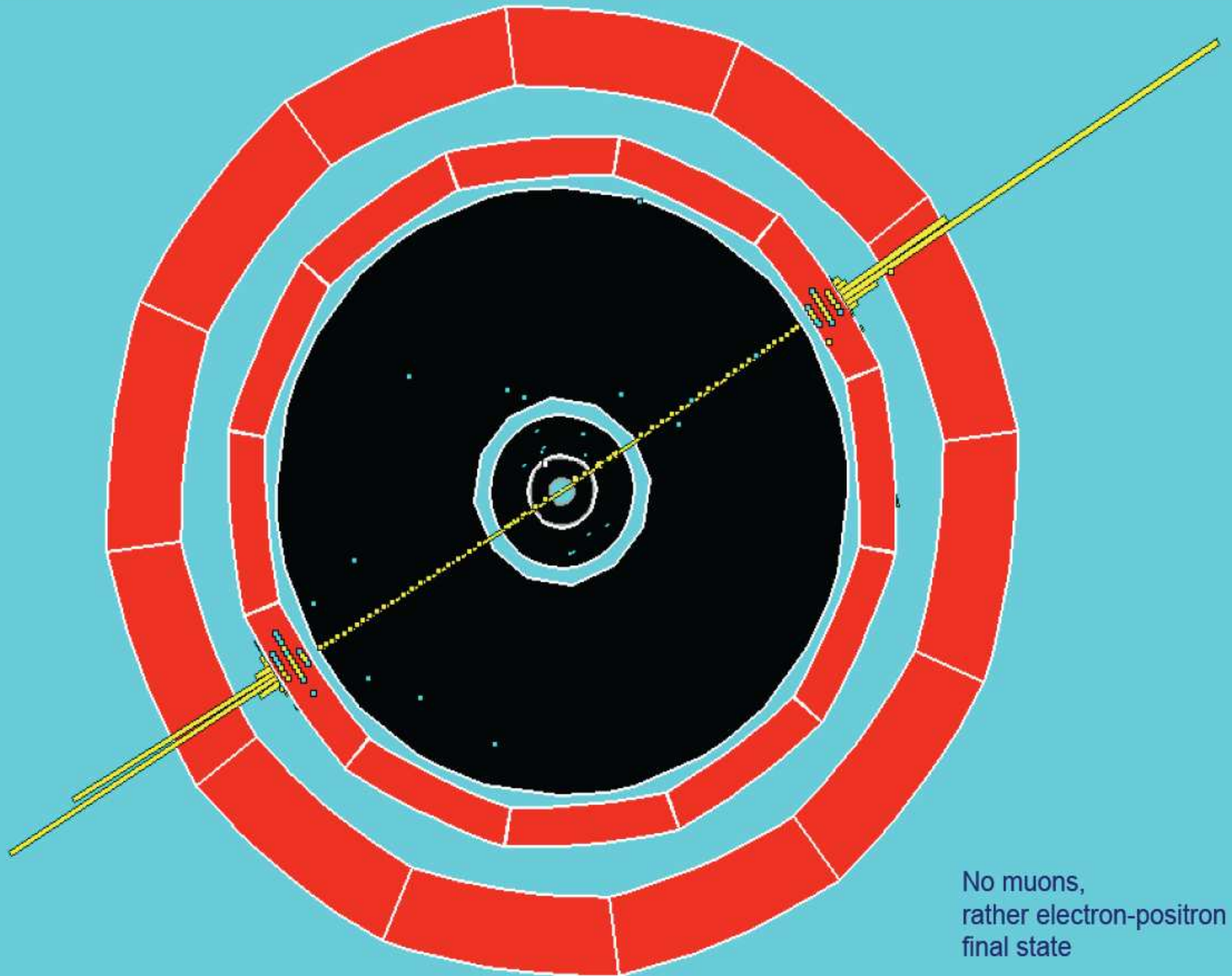
- two charged tracks, approx. back-to-back, with expected momentum (  $\sim 1/2 E_{CM}$  )
- right number of muon hits in outer layers (muons very penetrating, traverse whole detector)
- expected energy in calorimeter (electrons deposit all their energy, muons leave little)

# A Simple Counting Experiment

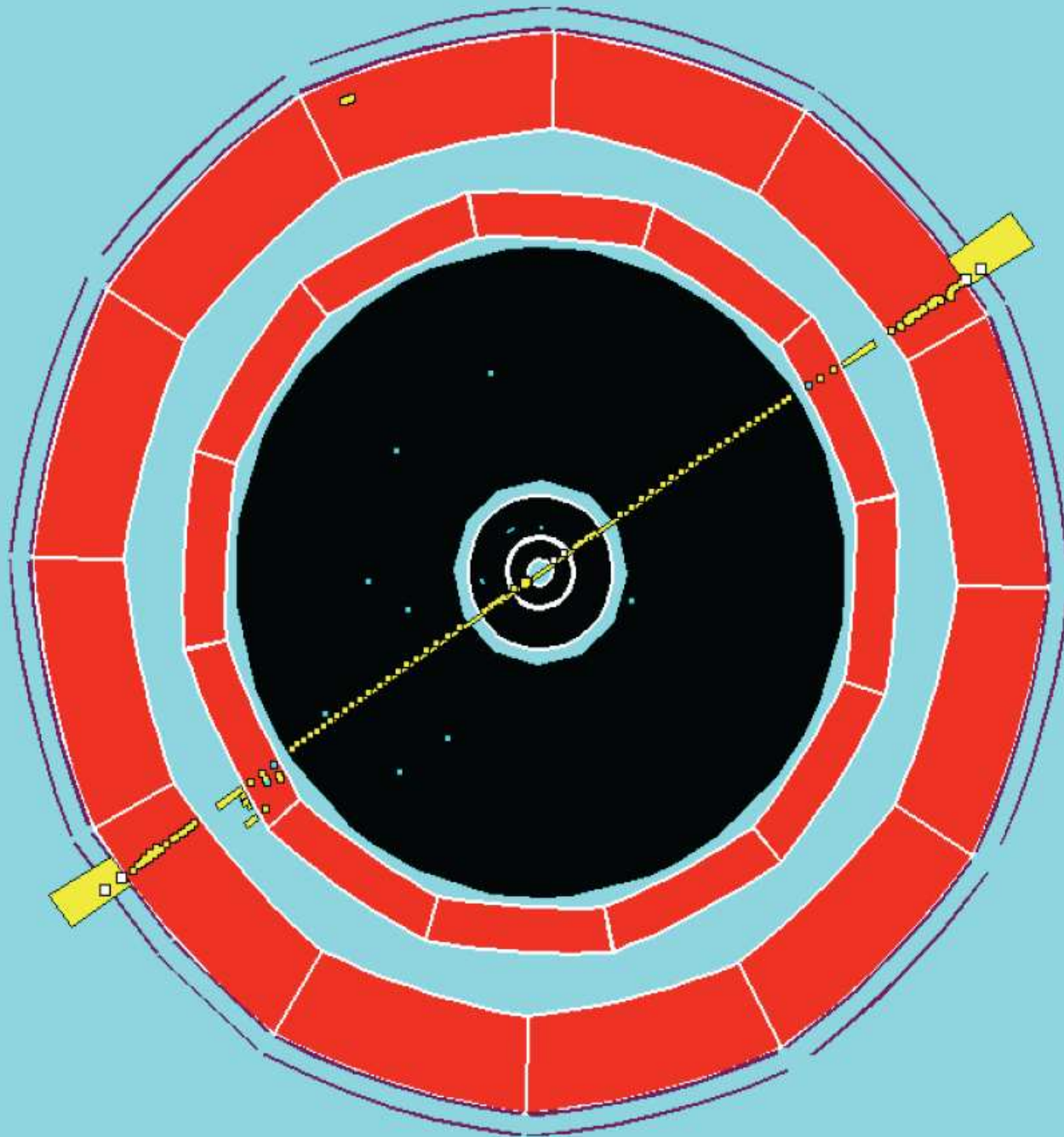


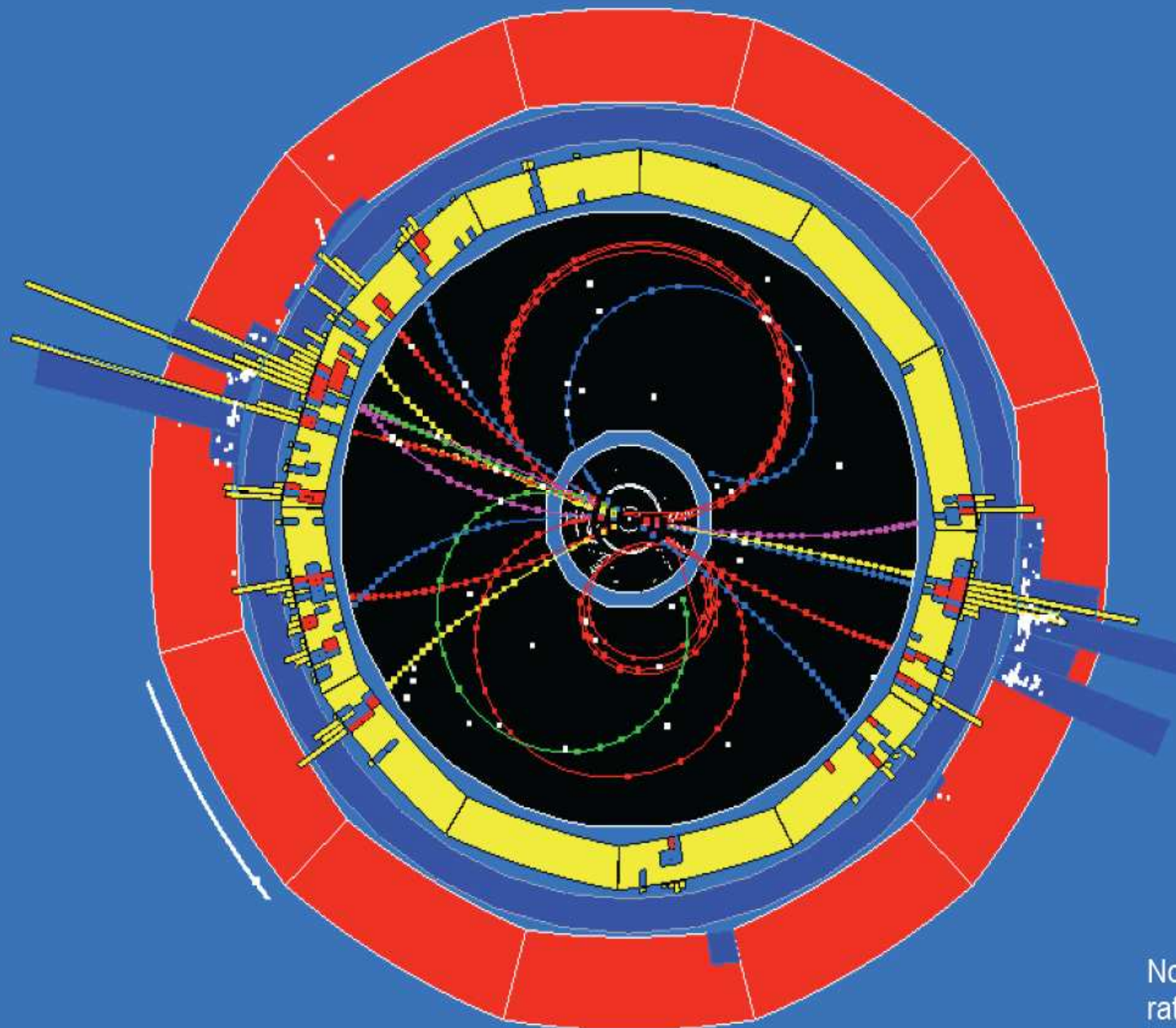


Not muonic  
rather hadronic  
final state



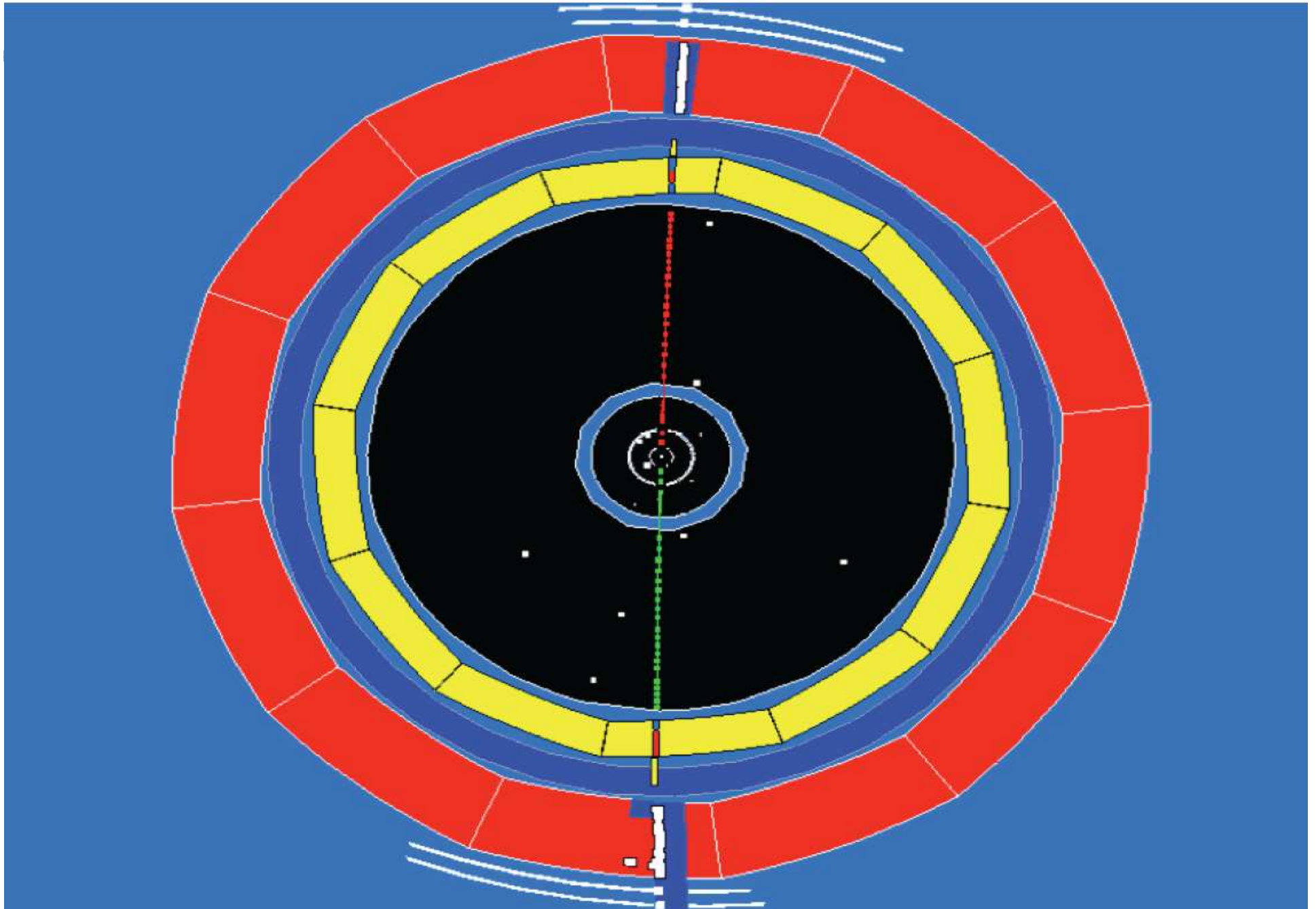
No muons,  
rather electron-positron  
final state





Not muonic  
rather hadronic  
final state

$$Z \rightarrow \mu^+ \mu^-$$

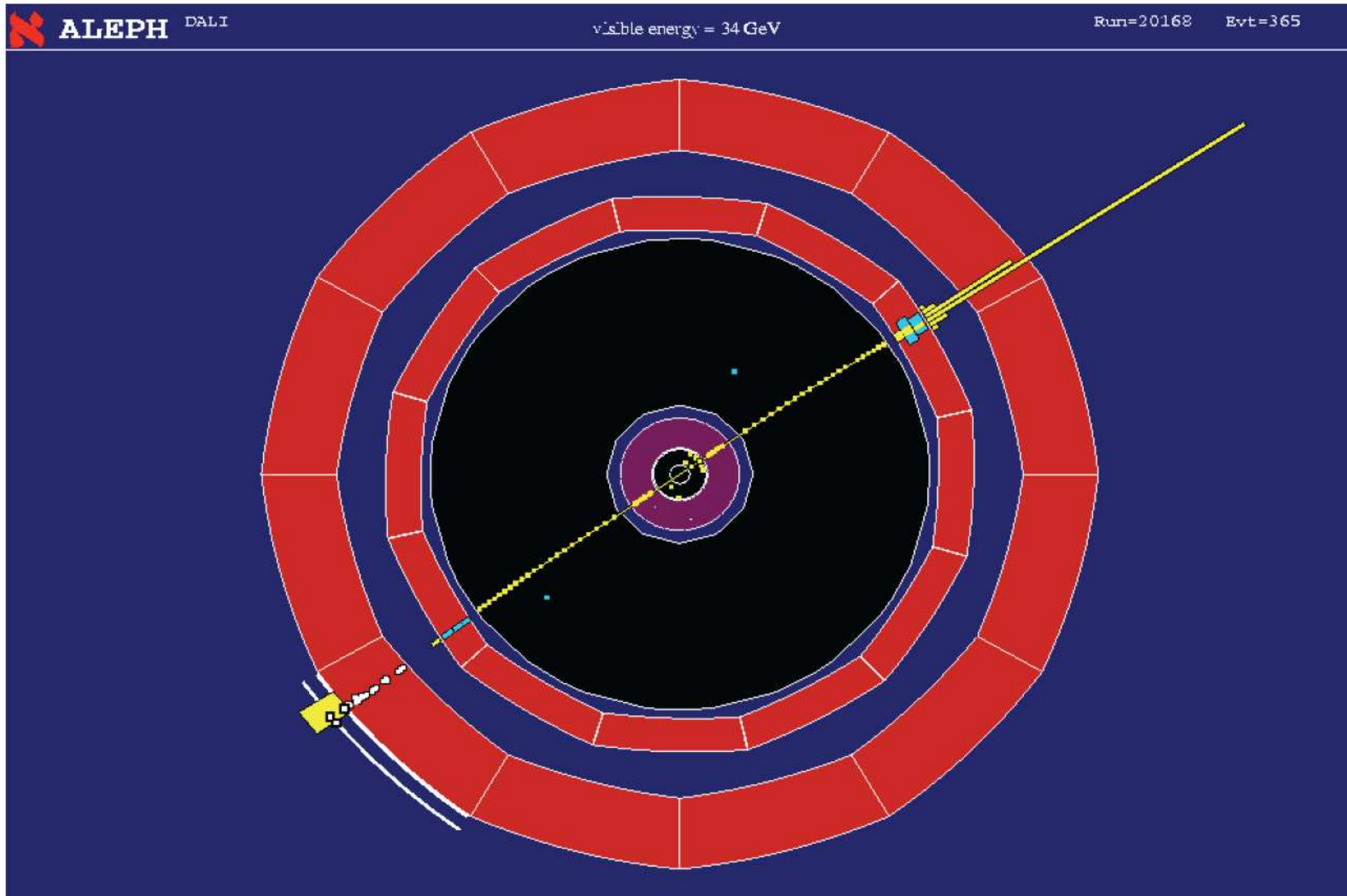






not  $Z \rightarrow \mu^+ \mu^-$

rather Z decay to  $\tau^+ \tau^-$ ,  
one tau decayed to electron + 2 neutrinos  
the other tau decayed to muon + 2 neutrinos





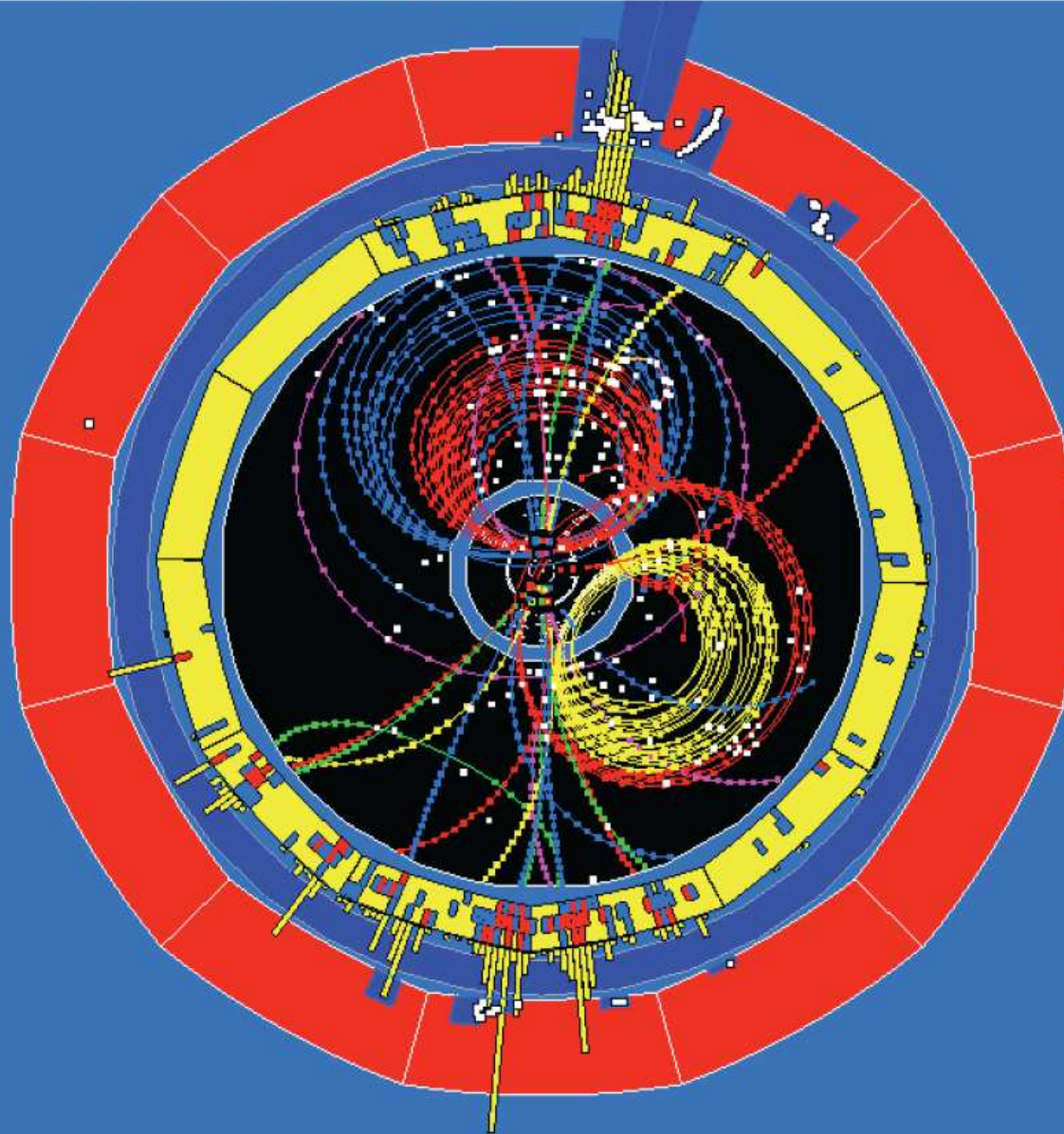
Not muonic, rather hadronic final state

**ALEPH** DALI

$e^+ e^- \rightarrow q \bar{q} \rightarrow \text{hadrons}$

Run=15995

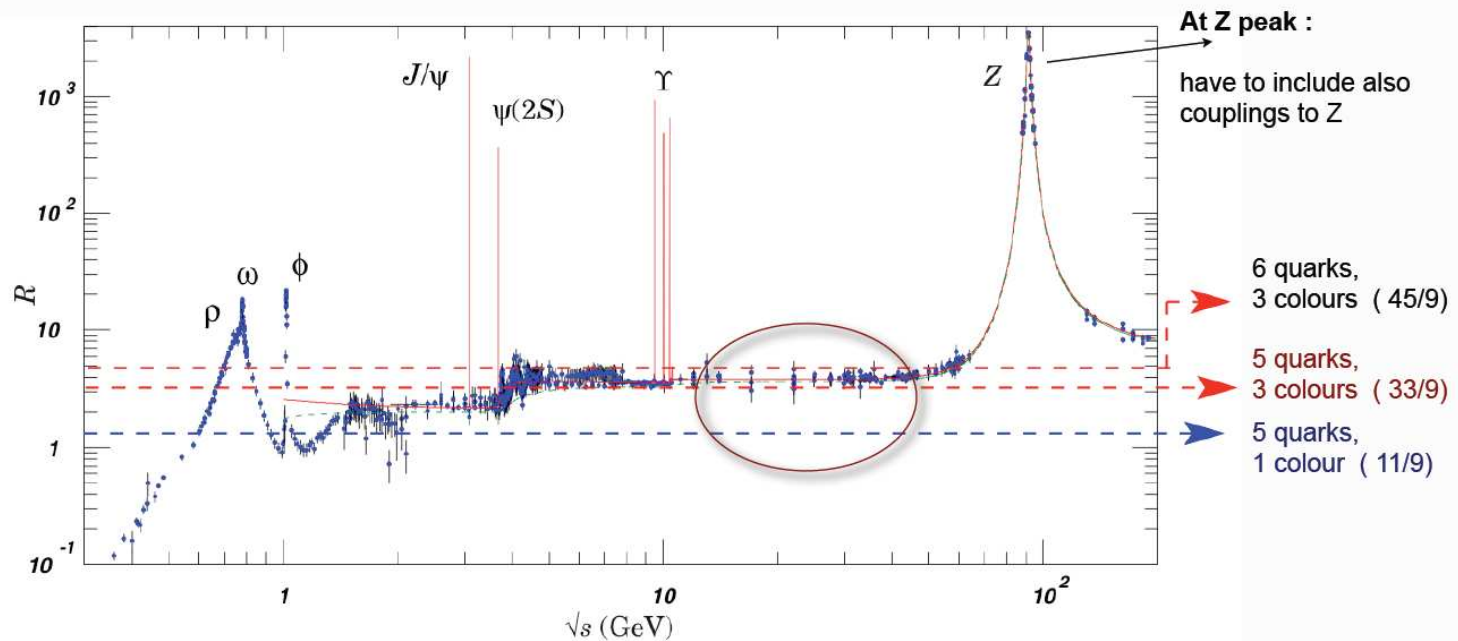
Evt=349



# Result

For  $E_{CM}$  below the Z peak and above the  $\Upsilon$  resonance we expect:

$$R = N_c \sum_f z_f^2 = N_c \cdot \left[ \left(\frac{2}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 + \left(\frac{2}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 \right] = N_c \cdot \frac{11}{9}$$



🕒 Confirmation of : Number of colours = 3 !

Note : small remaining difference : because of QCD correction (gluon radiation) =  $1 + \alpha_s/\pi$

# Uncertainties

Just having a “counting result” is not all, there’s lot more to do!

## Statistical error

- We saw 2 muon events, could easily have been 1 or 3
- Those fluctuations go like the square-root of the number of events

$$BR(Z^0 \rightarrow \mu^+ \mu^-) = \frac{N_{\mu\mu}}{N_{total}} \pm \frac{\sqrt{N_{\mu\mu}}}{N_{total}}$$

- To reduce this uncertainty, you need to record lots (millions) of events in the detector, and process them

## Systematic error

What if you only see 50% of the  $\mu^+ \mu^-$  events?

“efficiency”

$$N_{\mu\mu\text{seen}} = \varepsilon N_{\mu\mu}$$

- because of event selection (cut), detector imperfections, poor understanding, etc.

$$BR(Z^0 \rightarrow \mu^+ \mu^-) = \frac{N_{\text{seen}}/\varepsilon}{N_{total}}$$

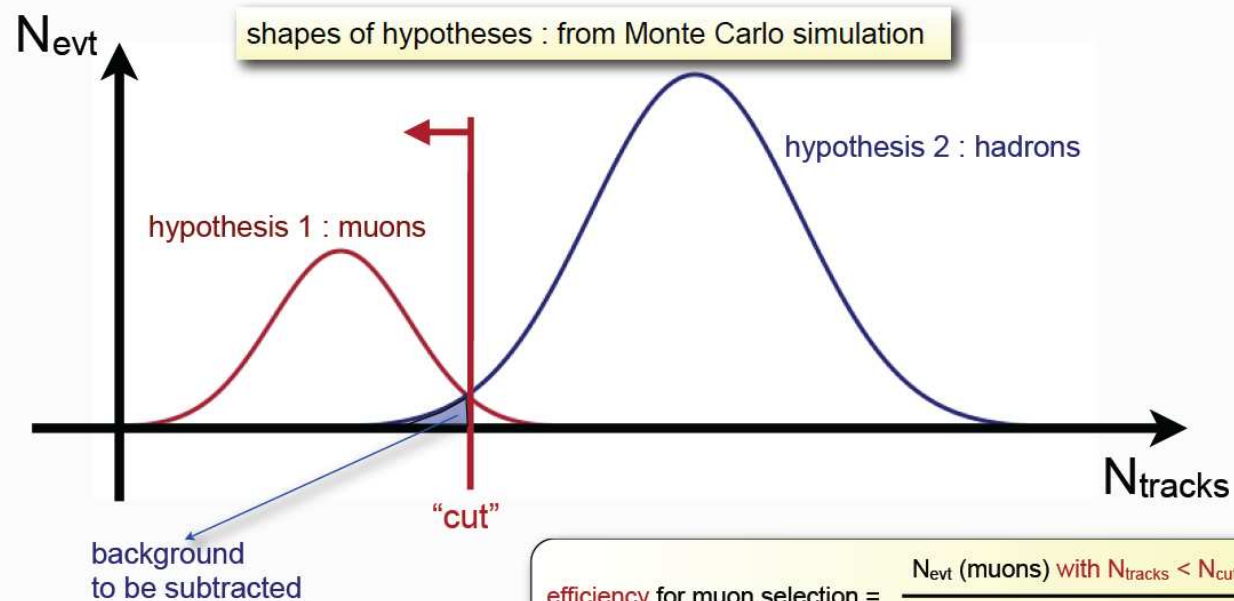
$$\varepsilon = 0.50 \pm 0.05$$

- from statistical error of detector simulation
- imperfect modeling of geometry in simulation
- model of muon interactions in simulation, etc

# Event Selection

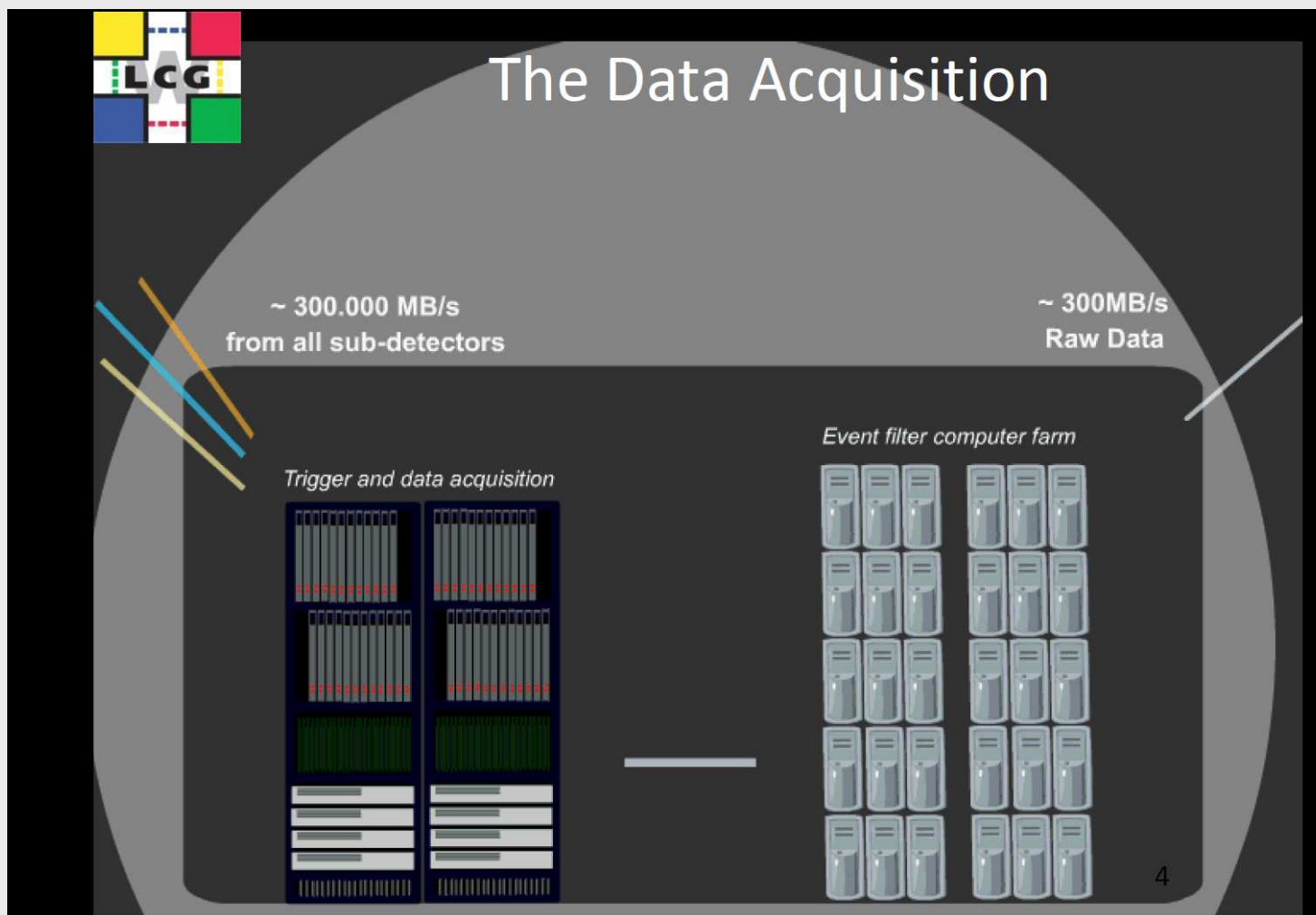
## Event per event have to decide how to categorize it

- eg. do we call it a muon event, or a hadronic event?
- how do we estimate the efficiency?
- Define an **event selection**, eg. “cut-based”
- see statistics lectures, *hypothesis testing* etc...



$$\text{efficiency for muon selection} = \frac{N_{\text{evt}}(\text{muons}) \text{ with } N_{\text{tracks}} < N_{\text{cut}}}{\text{all } N_{\text{evt}}(\text{muons})}$$

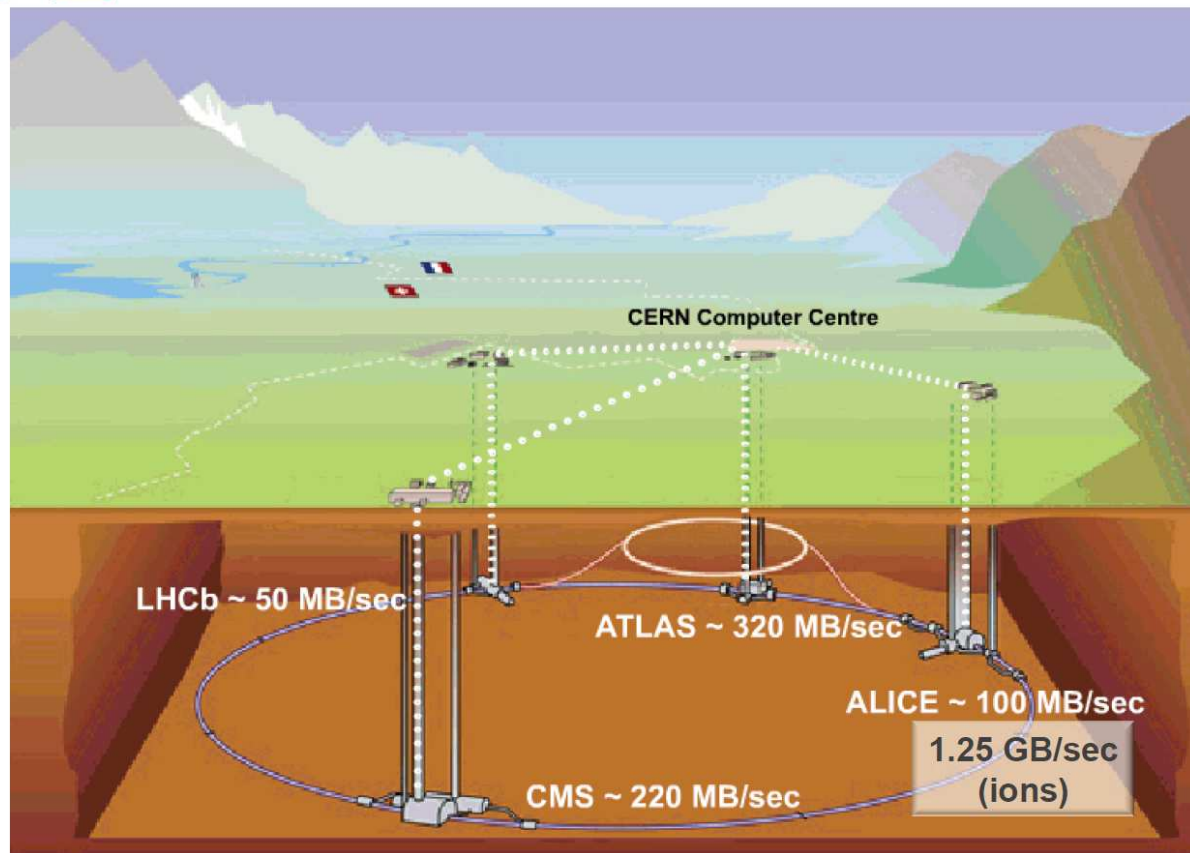
# The Data Acquisition



# Grid Computing



Tier 0 at CERN: Acquisition, First pass processing  
Storage & Distribution



# Grid Computing



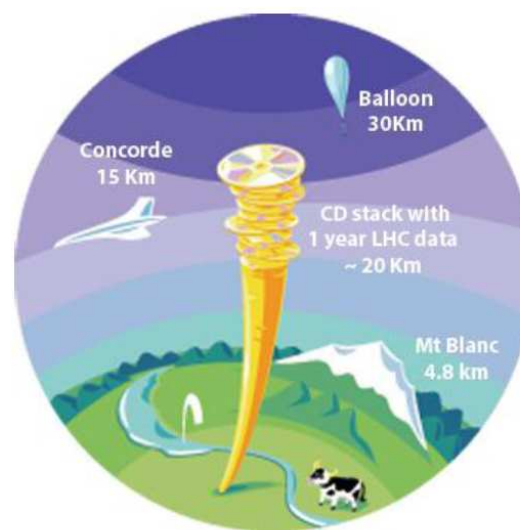
CERN IT Department  
CH-1211 Genève 23  
Switzerland

## The LHC Data Challenge

CERN IT  
Department

### The LHC will have a lifetime of ~20 years

- Experiments will produce about **15 Million Gigabytes** of data each year (about 20 million CDs!)
- LHC data analysis requires a computing power equivalent to **~100,000 of today's fastest PC processors**
- Requires many cooperating computer centres, as CERN can **only provide ~20% of the capacity**

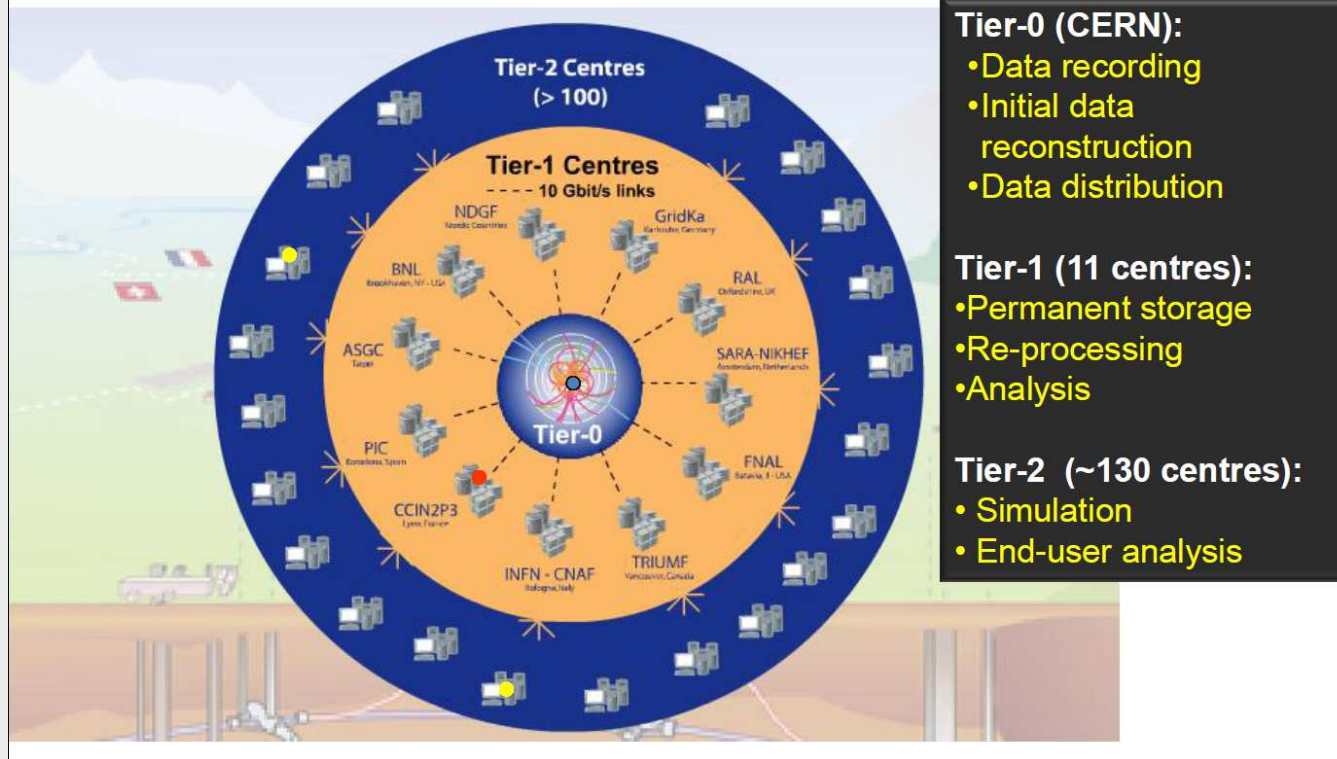




# Grid Computing

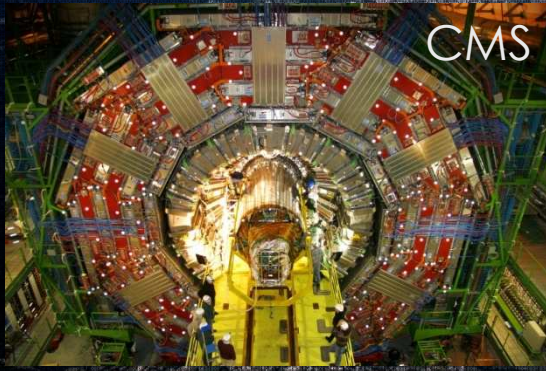


## Tier 0 – Tier 1 – Tier 2



# Enter a New Era in Fundamental Science

Start-up of the Large Hadron Collider (LHC), one of the largest and truly global scientific projects ever, is the most exciting turning point in particle physics.

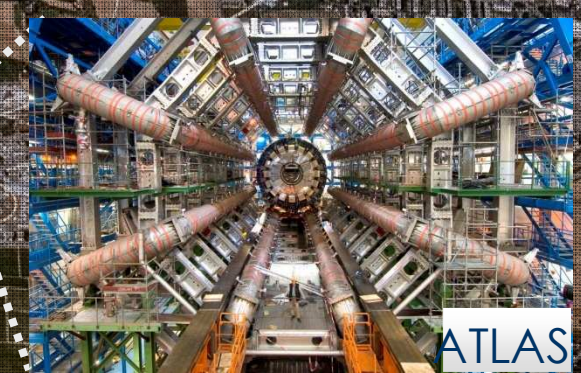


Exploration of a new energy frontier

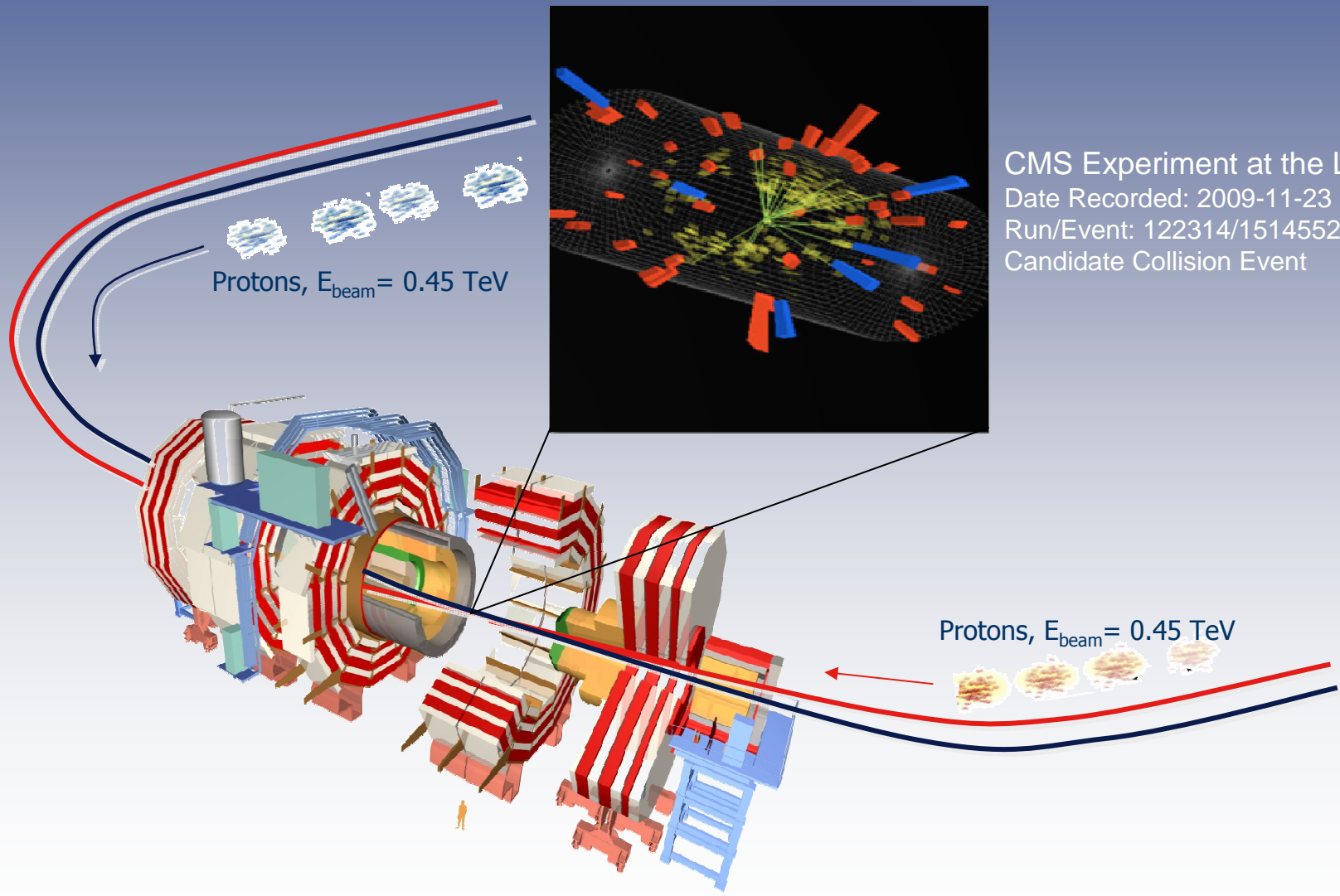


plus  
three smaller experiments

TOTEM  
LHCf  
MoEDAL

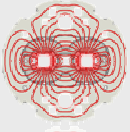


# First Collisions at LHC on 23 November 2009 at $E_{CM} = 900 \text{ GeV}$

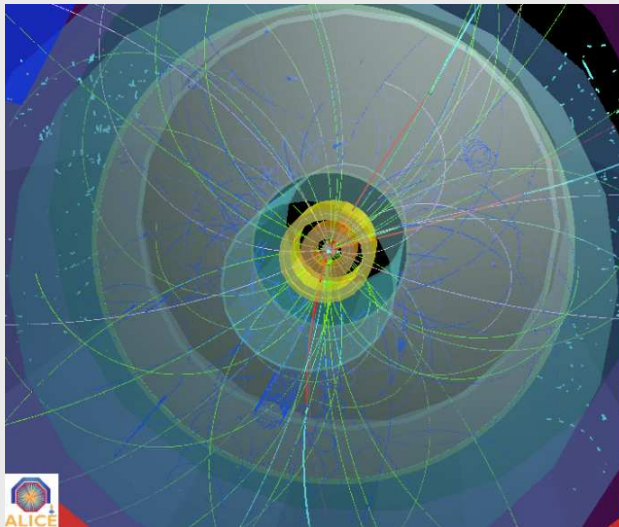
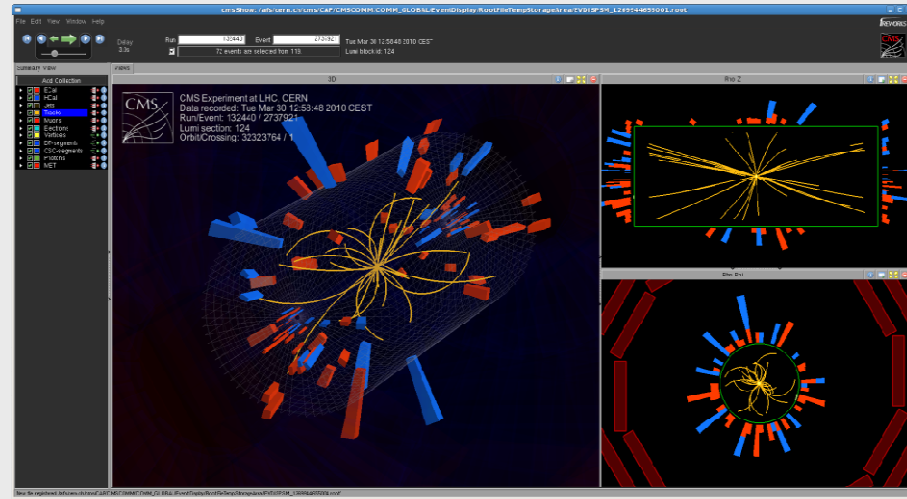
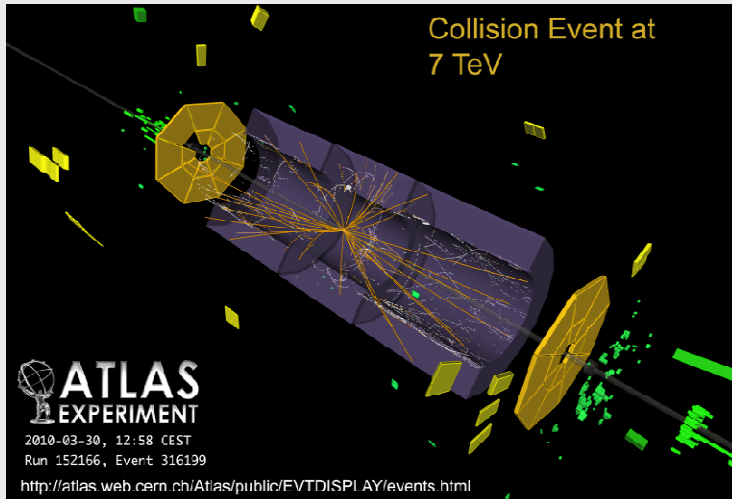


CMS Experiment at the LHC  
Date Recorded: 2009-11-23 19:21 CET  
Run/Event: 122314/1514552  
Candidate Collision Event

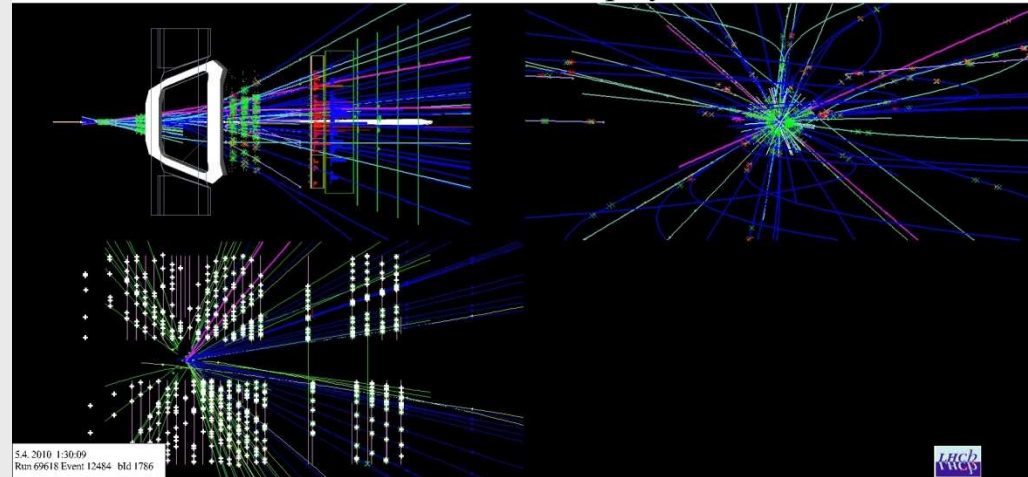
Protons,  $E_{\text{beam}} = 0.45 \text{ TeV}$



# LHC Started 7-TeV Collisions on 30 March 2010



LHCb Event Display



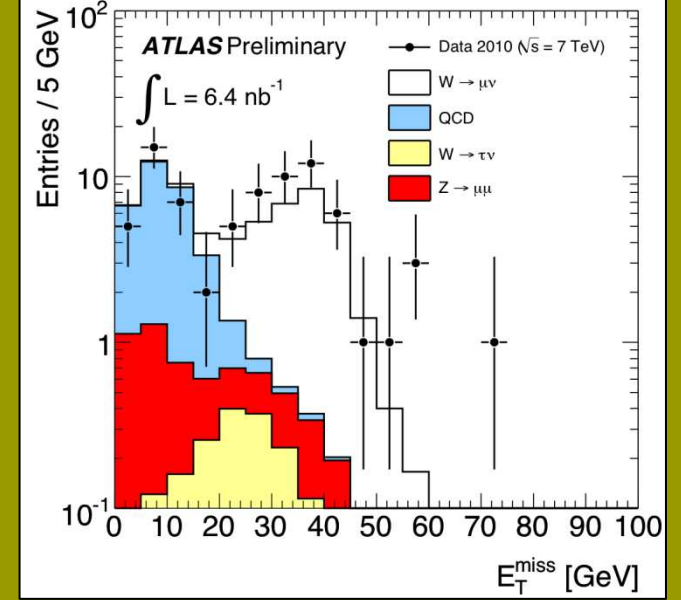
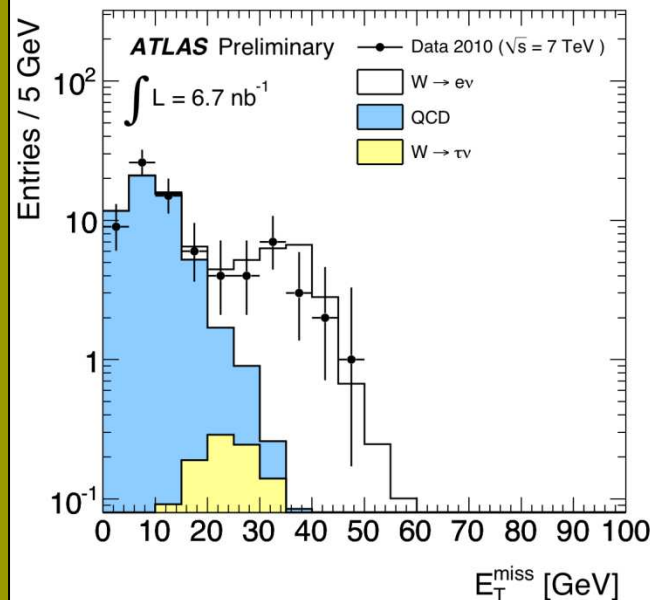
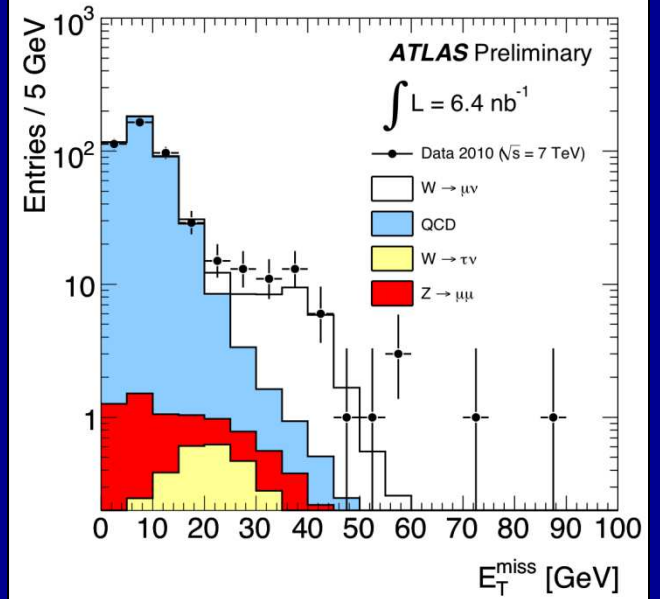
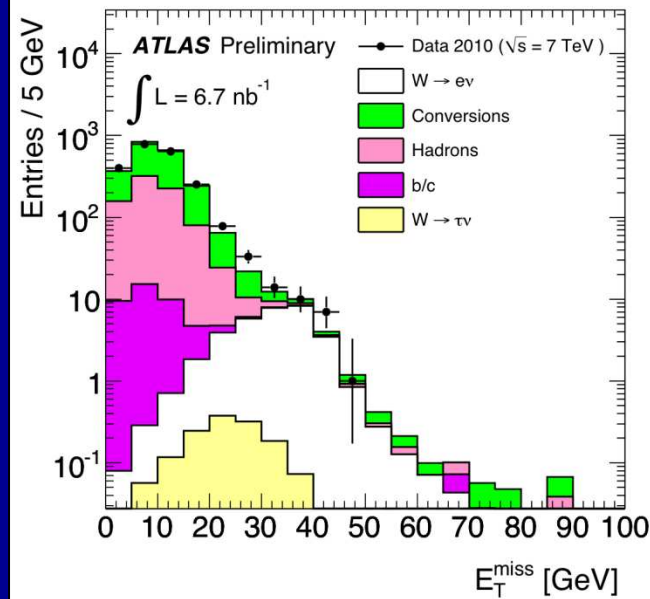
## After pre-selection:

- $W \rightarrow e\nu$ :  
loose  $e^\pm$ ,  $E_T > 20 \text{ GeV}$
- $W \rightarrow \mu\nu$ :  
 $p_T(\mu) > 15 \text{ GeV}$   
 $|\Delta p_T(\text{ID-MS})| < 15 \text{ GeV}$   
 $|Z_\mu - Z_{\nu\tau}| < 1 \text{ cm}$

MC: normalised to data  
(total number of events)

Observed events: 57

After all cuts  
but  $E_T^{\text{miss}}$  and  $m_T$



# Z $\rightarrow$ e<sup>+</sup>e<sup>-</sup> Observation

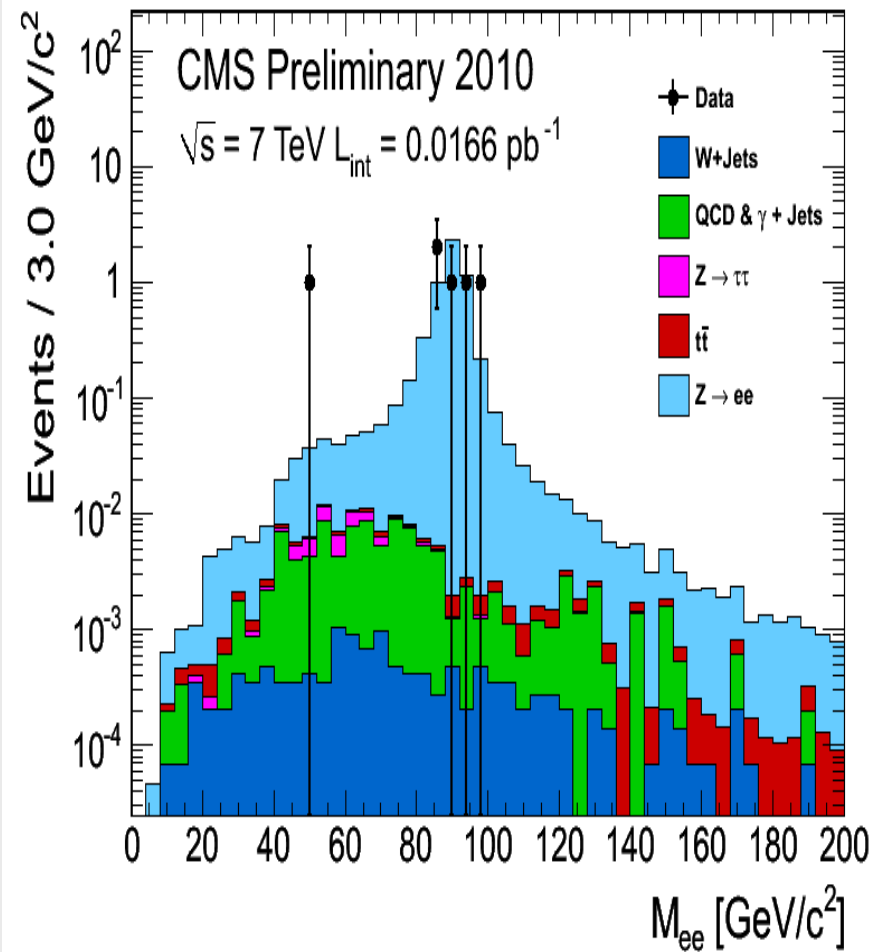
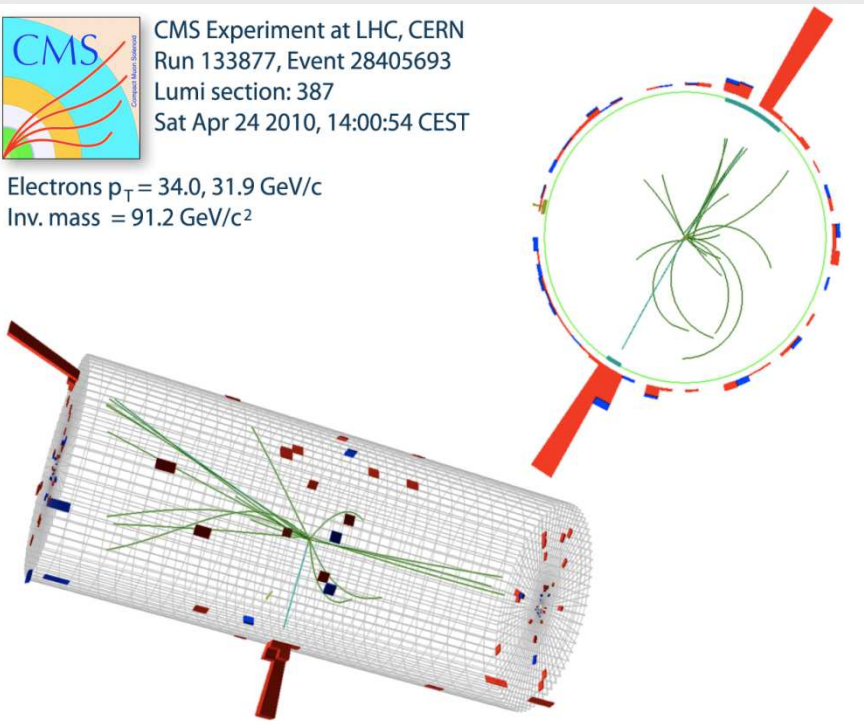
Event selection: *both electrons with a SuperCluster with  $E_t > 20$  GeV*

Monte Carlo : *cross section normalized to 17 nb<sup>-1</sup> integrated luminosity*



CMS Experiment at LHC, CERN  
Run 133877, Event 28405693  
Lumi section: 387  
Sat Apr 24 2010, 14:00:54 CEST

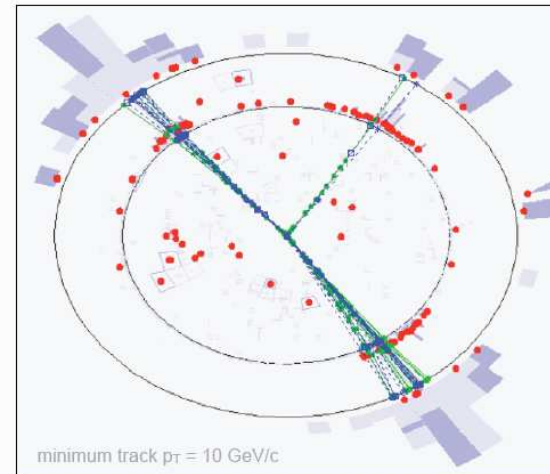
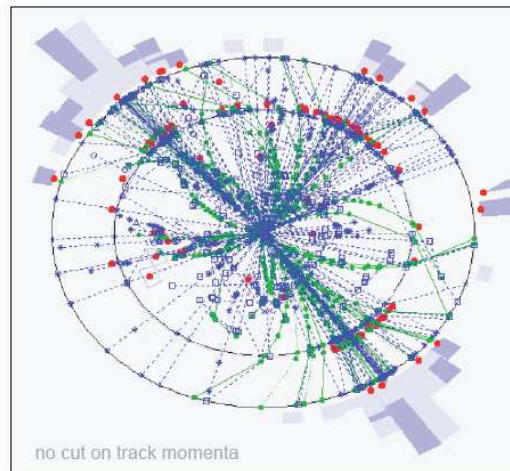
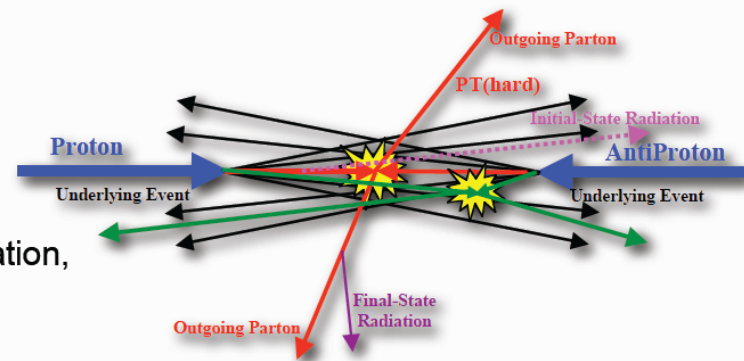
Electrons  $p_T = 34.0, 31.9$  GeV/c  
Inv. mass = 91.2 GeV/c<sup>2</sup>



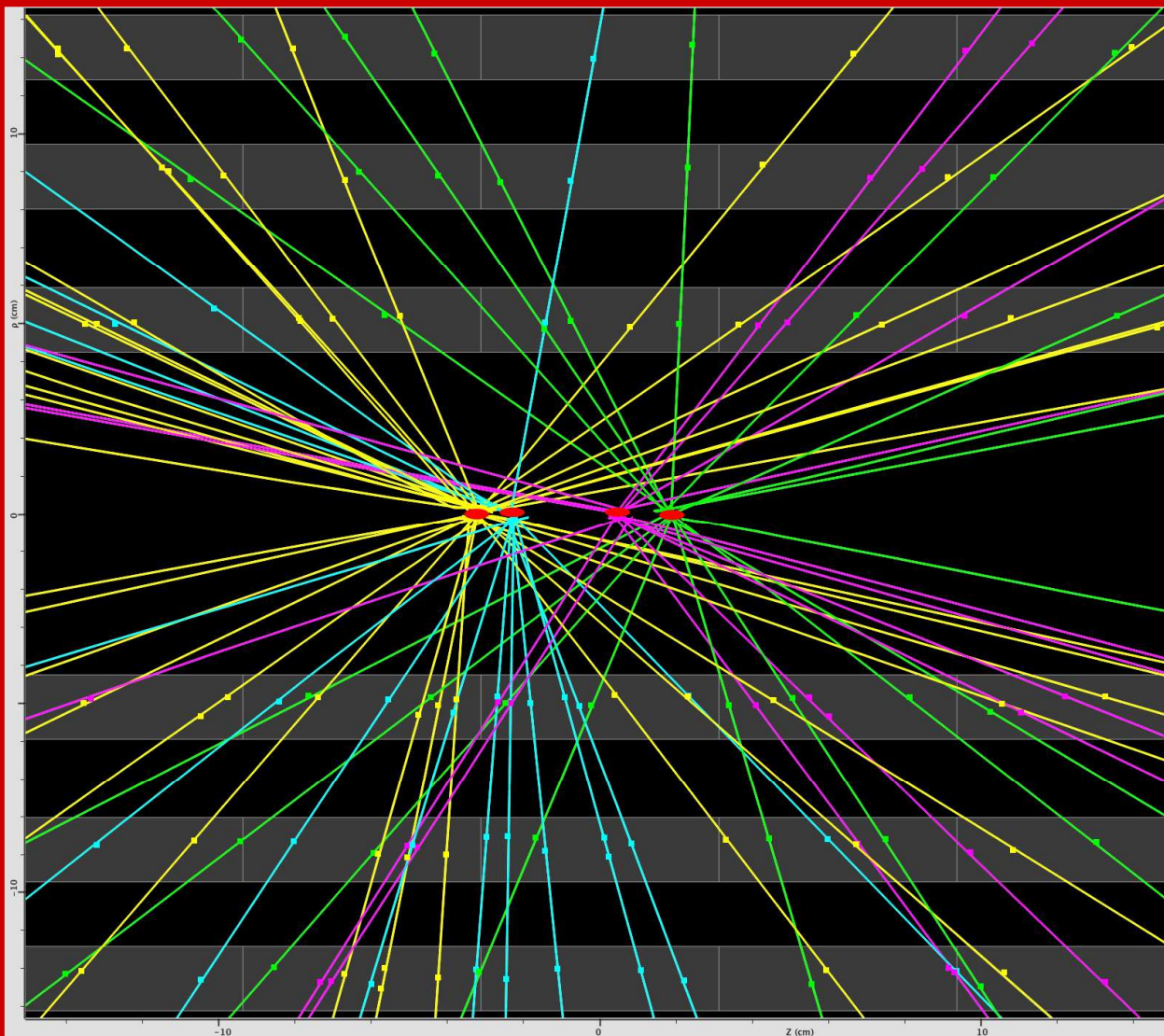
**5 Z  $\rightarrow$  e<sup>+</sup>e<sup>-</sup> candidates**

# Further Difficulties

- **Pile Up** : many additional soft proton-proton interactions
  - up to 20 at highest LHC luminosity
- **Underlying event**
  - beam-beam remnants, initial state radiation, multiple parton interactions
  - gives additional energy in the event
- All this additional energy has nothing to do with jet energies
  - **have to subtract it**



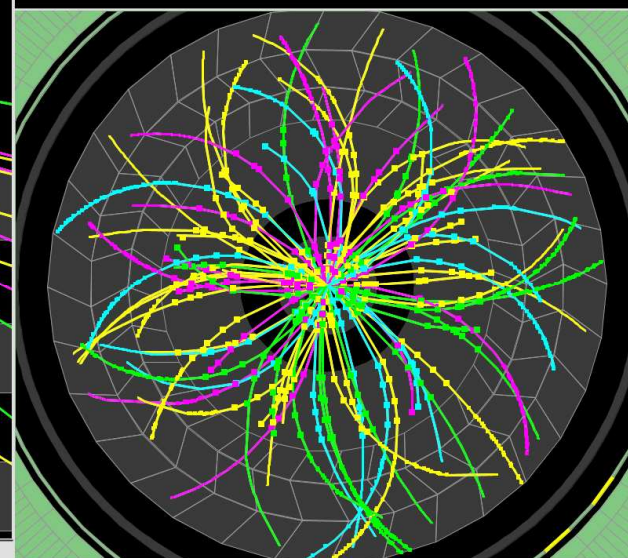
Preparing for the future : pile-up reconstruction  
4 pp interactions in the same bunch-crossing



Run Number: 153565, Event Number: 4487360

Date: 2010-04-24 04:18:53 CEST

Event with 4 Pileup Vertices  
in 7 TeV Collisions



~ 10-45 tracks with  $p_T > 150$  MeV per vertex  
Vertex z-positions : -3.2, -2.3, 0.5, 1.9 cm (vertex resolution better than ~200  $\mu$ m)



# Conclusions

---

- Experimental methods in Particle Physics is an inter-disciplinary study between
  - Experimental physics
  - Theoretical physics
  - Detector physics (and technology)
  - High-performance computing
  - Accelerator physics
  - Statistical analysis
  - .....

***An excellent training ground for young scientists***



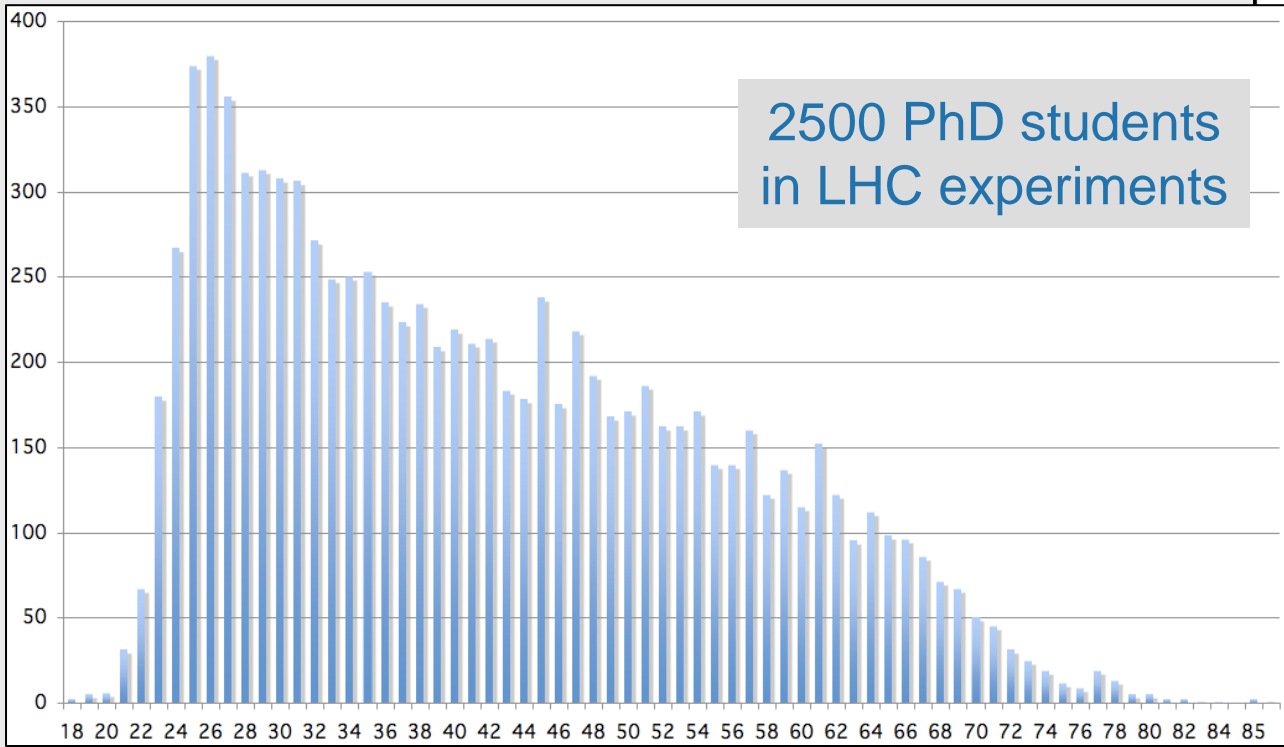


# Age Distribution of Scientists

- and where they go afterwards

Survey in March 2009

2500 PhD students in LHC experiments



They do not all stay: where do they go?

